

Insights into sea quark asymmetries from global PDF analysis

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CTEQ-Jefferson Lab (CJ) collaboration

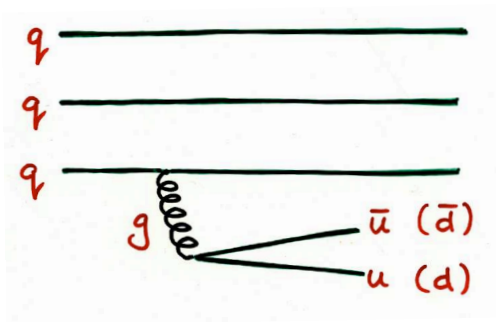
<http://www.jlab.org/CJ>

Outline

- Asymmetries between various PDFs ($\bar{d} - \bar{u}$, $s - \bar{s}$, $\Delta\bar{u} - \Delta\bar{d} \dots$) reveal nonperturbative structure of the nucleon
- $\bar{d} - \bar{u}$ asymmetry and chiral symmetry in QCD: a brief history
- Evidence of sign change at large x ?
- New (preliminary) SeaQuest data
 - constraint on $\bar{d} - \bar{u}$ at large x
 - implications for pion cloud models and pion structure function extraction
- Outlook

Light quark sea

- From perturbative QCD expect symmetric $q\bar{q}$ sea generated by gluon radiation into $q\bar{q}$ pairs (if quark masses are the same)

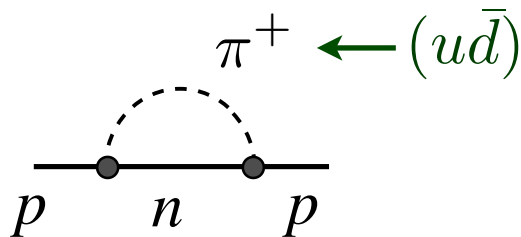


→ since u and d quarks nearly degenerate, expect flavor-symmetric light-quark sea

$$\bar{d} \approx \bar{u}$$

Ross, Sachrajda (1979)

- Thomas suggested that chiral symmetry of QCD (important at low energy) should have consequences for antiquark PDFs in nucleon (measured at high energy)

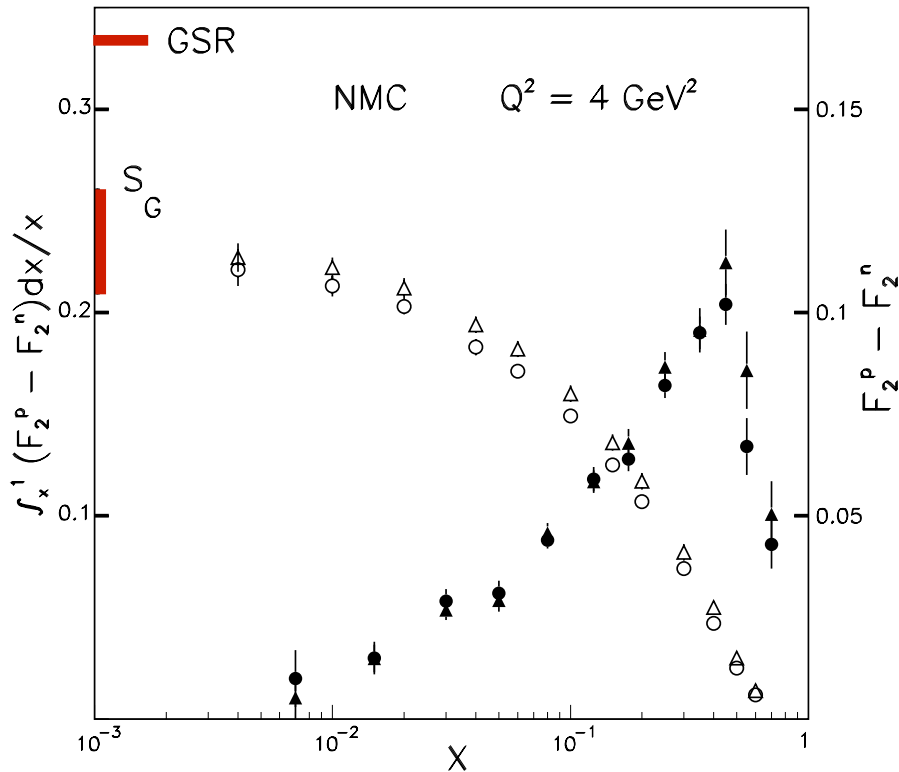


$$\rightarrow \bar{d} > \bar{u}$$

A.W. Thomas (1984)

Light quark sea

- First clear experimental support for $\bar{d} \neq \bar{u}$ came from violation of Gottfried sum rule observed by NMC



$$\frac{1}{x}(F_2^p - F_2^n) = \frac{1}{3}(u^+ - d^+) + \frac{2}{3}(\bar{u} - \bar{d})$$

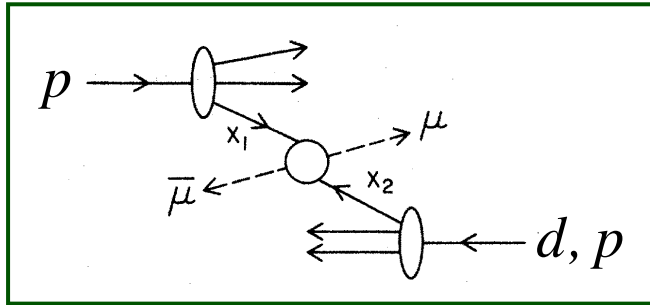
$$S_G \equiv \int_0^1 \frac{dx}{x} (F_2^p - F_2^n) = 0.235 \pm 0.026$$

NMC (1991, 1994)

→ clear evidence for $\bar{d} - \bar{u} > 0$ (or at least integrated value)

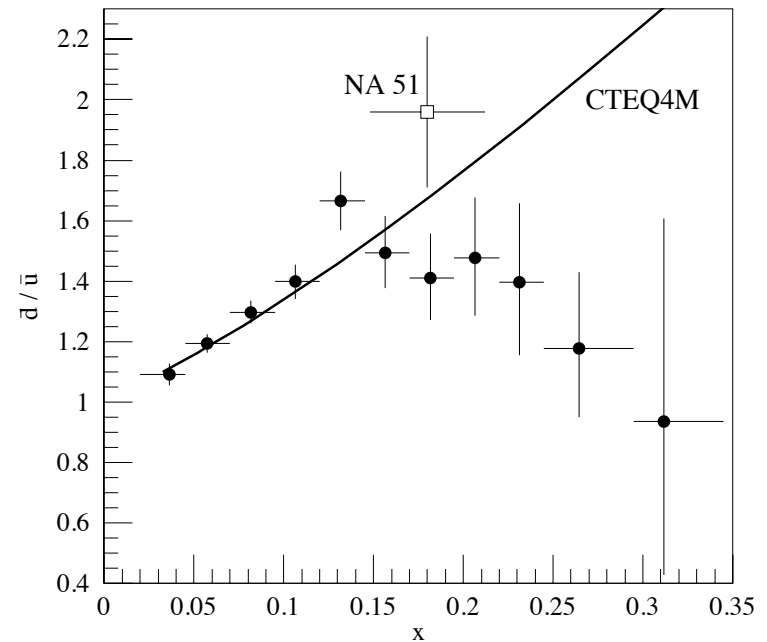
Light quark sea

- x dependence of $\bar{d} - \bar{u}$ asymmetry established in Fermilab E866 pp/pd Drell-Yan experiment



$$\frac{d\sigma}{dx_1 dx_2} \sim \sum_q e_q^2 q(x_1) \bar{q}(x_2) + (x_1 \leftrightarrow x_2)$$

$$\frac{\sigma^{pd}}{\sigma^{pp}} \approx 1 + \frac{\bar{d}(x_2)}{\bar{u}(x_2)} \quad \text{for } x_1 \gg x_2$$

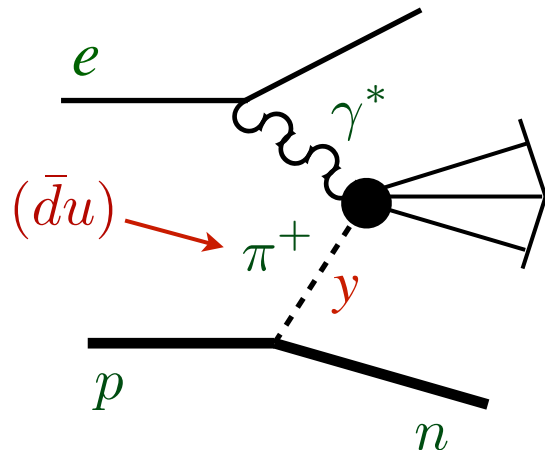


E866 (2001)

- strong enhancement of \bar{d} at $x \sim 0.1 - 0.2$, but intriguing behavior at large x hinting at possible sign change of $\bar{d} - \bar{u}$

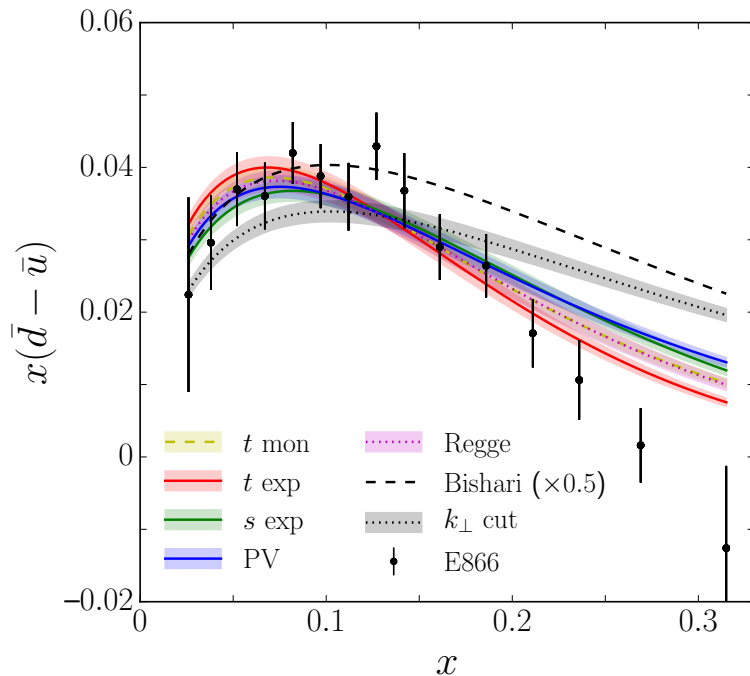
Light quark sea

■ General agreement with pion loop models



$$(\bar{d} - \bar{u})(x) = \int_x^1 \frac{dy}{y} f_{\pi^+ n}(y) \bar{q}_v^\pi(x/y)$$

$p \rightarrow \pi^+ n$ splitting function



→ shape qualitatively reproduced by most models (except at high x), but is there a direct connection with QCD?

Chiral EFT

- Recently rigorous connection with QCD established via chiral effective field theory

$$\mathcal{L}_{\text{eff}} = \frac{g_A}{2f_\pi} \bar{\psi}_N \gamma^\mu \gamma_5 \vec{\tau} \cdot \partial_\mu \vec{\pi} \psi_N - \frac{1}{(2f_\pi)^2} \bar{\psi}_N \gamma^\mu \vec{\tau} \cdot (\vec{\pi} \times \partial_\mu \vec{\pi}) \psi_N$$

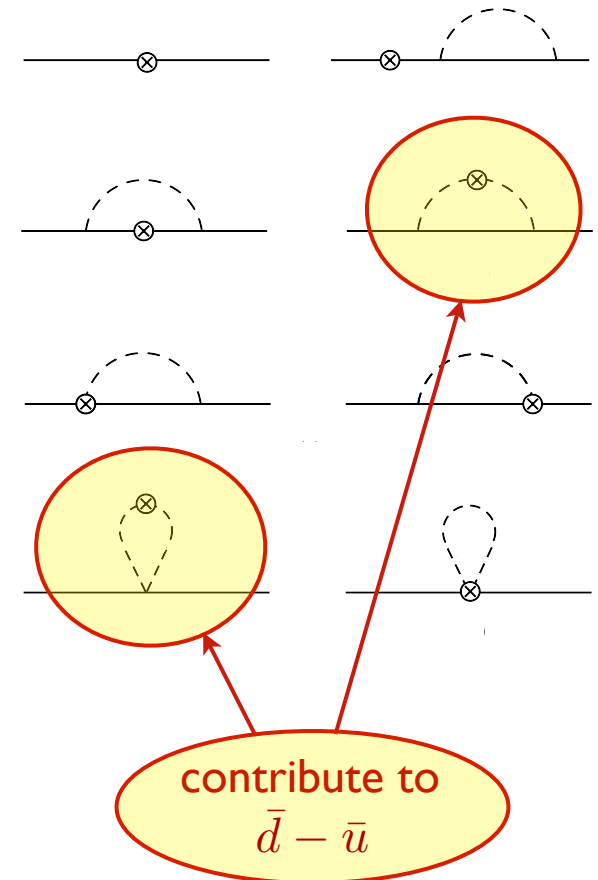
Weinberg (1967)

- lowest order πN interaction includes pion rainbow and tadpole contributions
- matching quark- and hadron-level operators

$$\mathcal{O}_q^{\mu_1 \cdots \mu_n} = \sum_h c_{q/h}^{(n)} \mathcal{O}_h^{\mu_1 \cdots \mu_n}$$

yields convolution representation

$$q(x) = \sum_h \int_x^1 \frac{dy}{y} f_h(y) q_v^h(x/y)$$



Ji, WM, Thomas (2013)

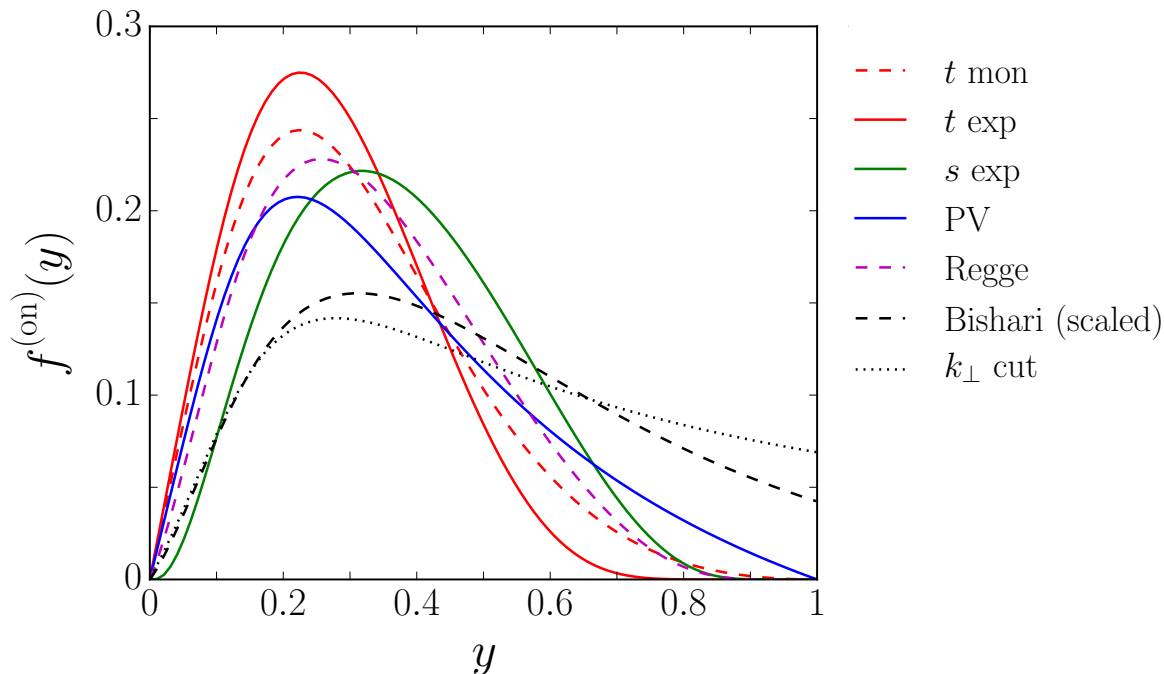
Chiral EFT

- Recently rigorous connection with QCD established via chiral effective field theory
- Splitting function for pion rainbow diagram has on-shell and δ -function contributions

$$f_{\pi}(y) = f^{(\text{on})}(y) + f^{(\delta)}(y)$$

$$f^{(\text{on})}(y) = \frac{g_A^2 M^2}{(4\pi f_{\pi})^2} \int dk_{\perp}^2 \frac{y(k_{\perp}^2 + y^2 M^2)}{[k_{\perp}^2 + y^2 M^2 + (1-y)m_{\pi}^2]^2} \mathcal{F}^2$$

$$f^{(\delta)}(y) = \frac{g_A^2}{4(4\pi f_{\pi})^2} \int dk_{\perp}^2 \log\left(\frac{k_{\perp}^2 + m_{\pi}^2}{\mu^2}\right) \delta(y) \mathcal{F}^2$$



UV regulator

Chiral EFT

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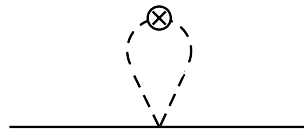
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- Bubble diagram contributes only at $y=0$ (hence $x=0$)

$$f^{(\text{bub})}(y) = \frac{8}{g_A^2} f^{(\delta)}(y)$$



Salamu, Ji, WM, Wang (2015)

→ contributes to lowest moment, but not at $x > 0$

Chiral EFT

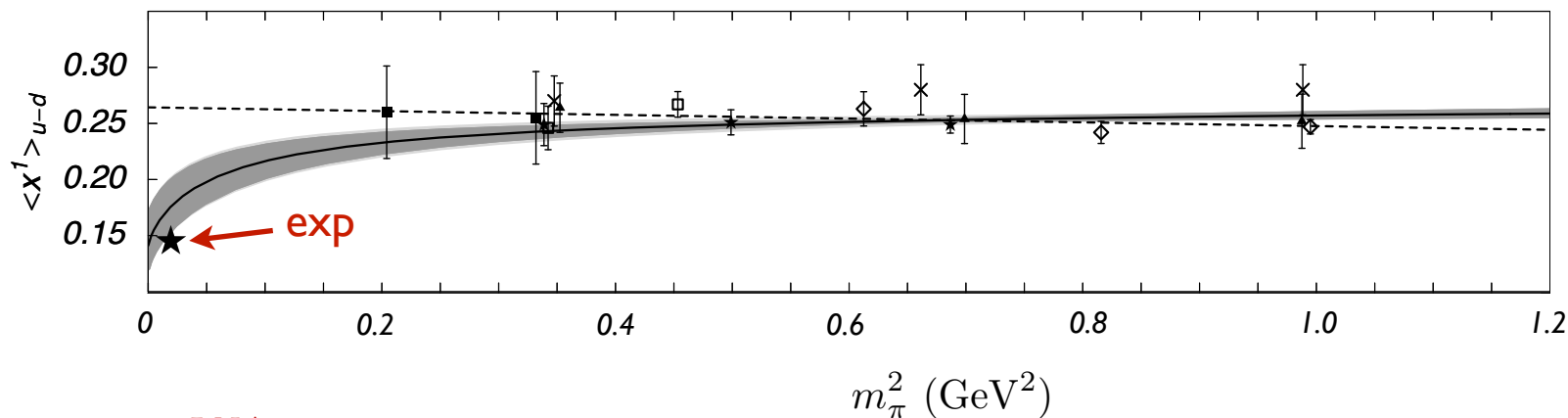
■ Expand moments of PDFs in powers of m_π

- coefficients of leading nonanalytic (LNA) terms, reflecting infrared behavior, are model-independent!
- QCD therefore *predicts* a nonzero asymmetry from π loops

$$\int_0^1 dx (\bar{d} - \bar{u}) = \frac{(3g_A^2 - 1)}{(4\pi f_\pi)^2} m_\pi^2 \log(m_\pi^2/\mu^2) + \text{analytic in } m_\pi^2$$

Thomas, WM, Steffens (2000)

- nonanalytic behavior vital for chiral extrapolation of lattice data on PDF moments from large m_π



$$\langle x \rangle_{u-d}^{\text{LNA}} \sim m_\pi^2 \log m_\pi^2$$

Detmold et al. (2001)

Chiral EFT

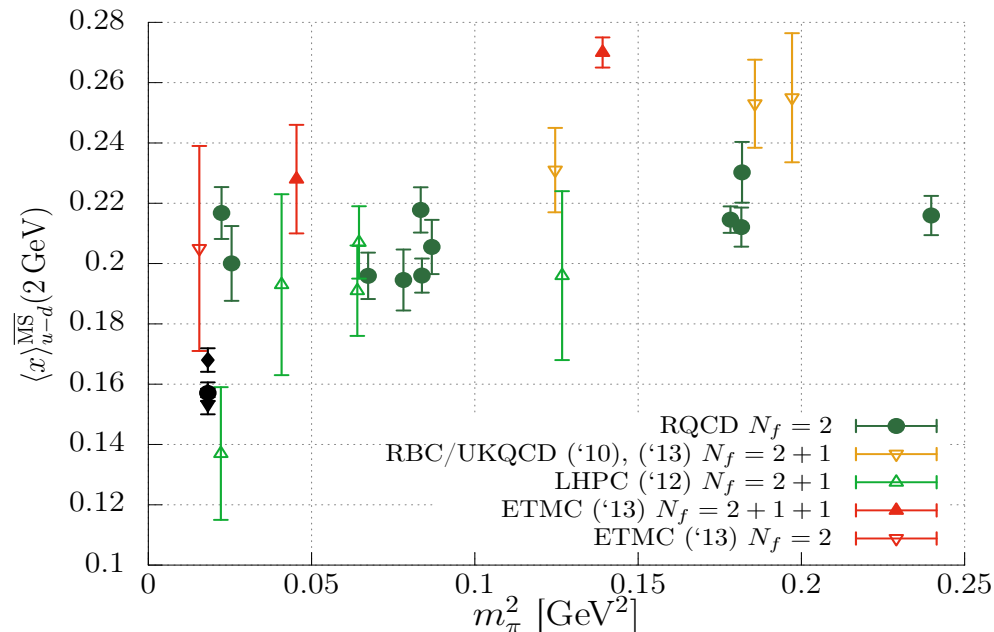
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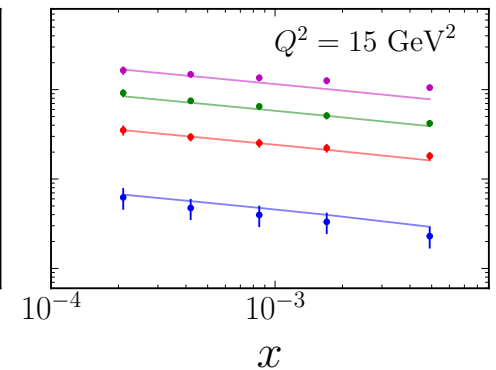
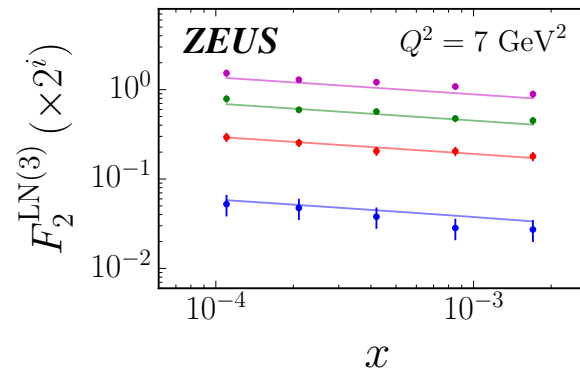
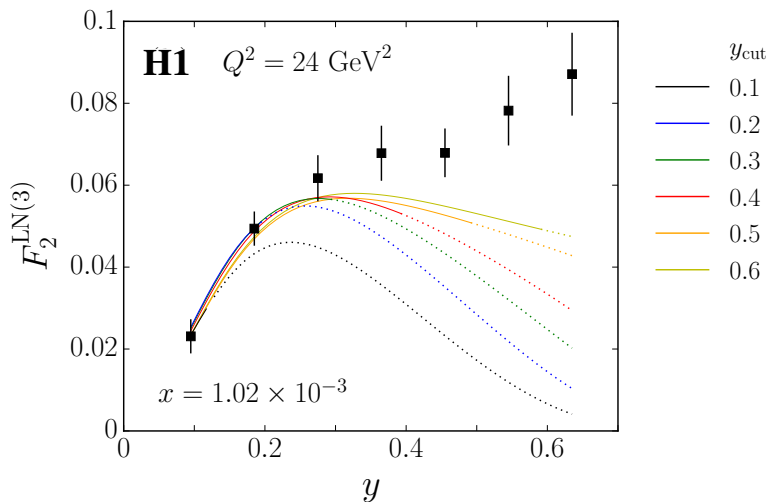
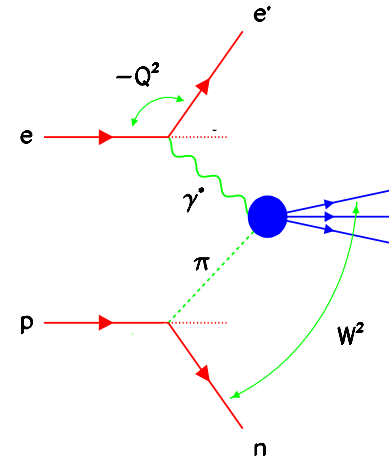
Bali et al. (2014)

Constraints from leading neutrons

- Drell-Yan $\bar{d} - \bar{u}$ data can be described with range of UV regulators (shapes of pion splitting functions)

→ semi-inclusive production of “leading neutrons” at HERA can discriminate between different shapes

$$\frac{d^3\sigma^{\text{LN}}}{dx dQ^2 dy} \sim F_2^{\text{LN}(3)}(x, Q^2, y) = f_\pi(y) F_2^\pi(x/y, Q^2)$$



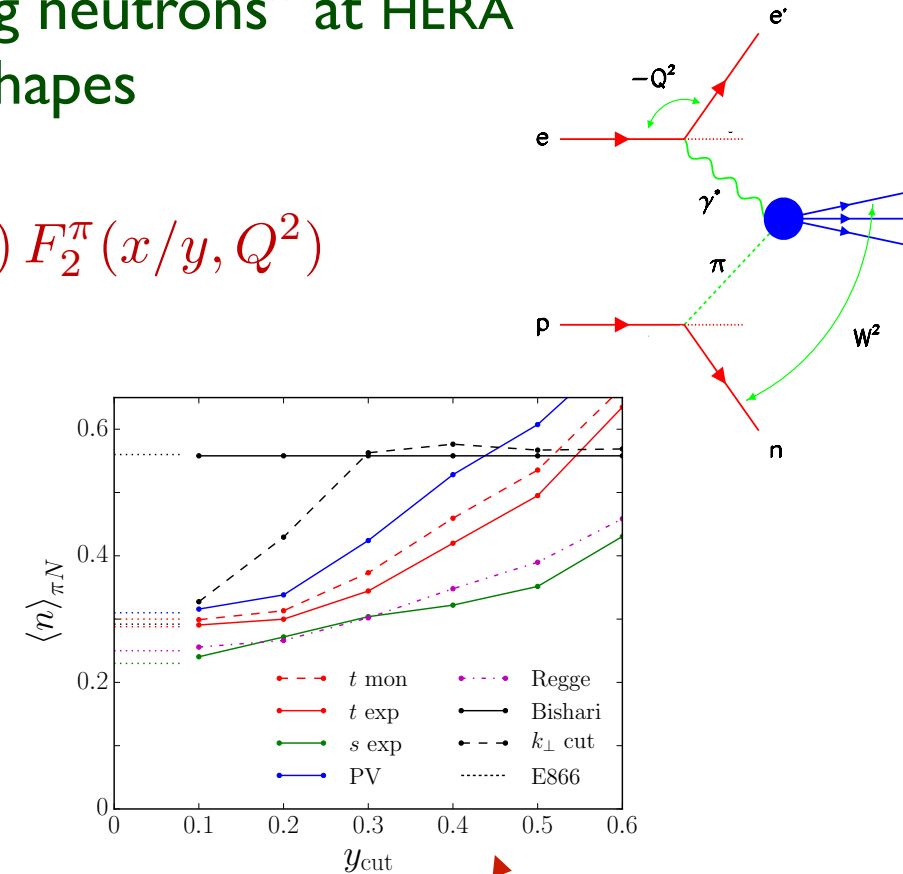
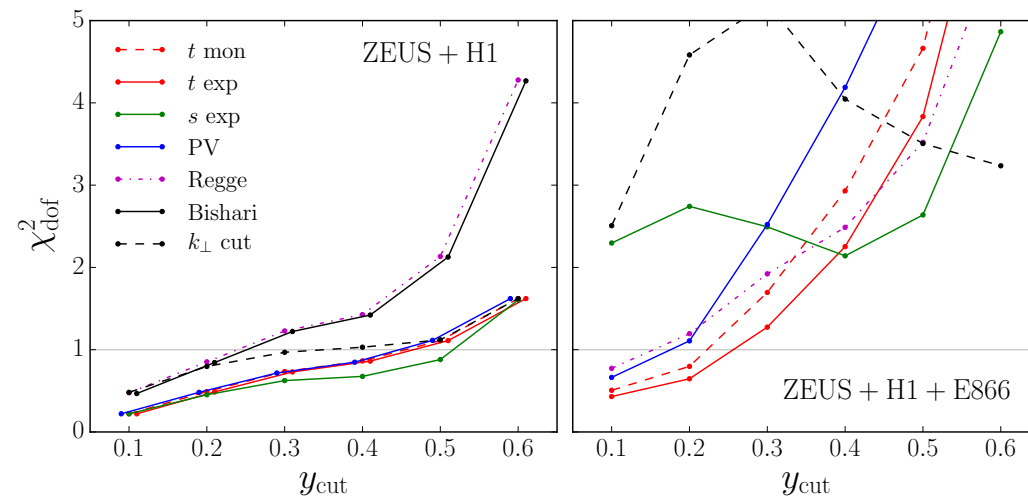
McKinney, Ji, WM, Sato (2016)

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→ best fit to combined ZEUS/H1 & DY data for t -dependent exponential regulator

average pion “multiplicity”

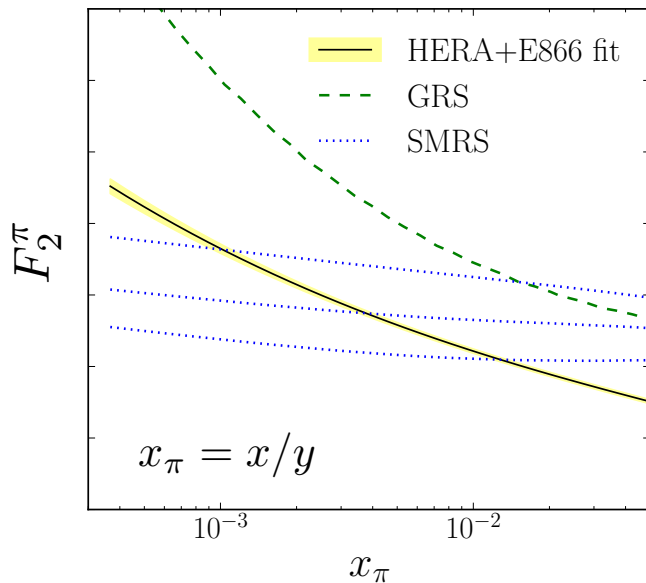
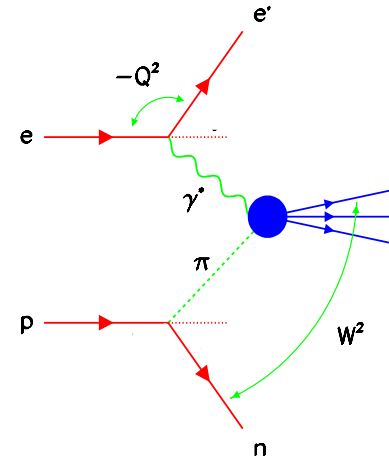
$$\langle n \rangle_{\pi N} = 3 \int_0^1 dy f_N^{(\text{on})}(y) \sim 0.3$$

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→ constrain shape of F_2^π at $10^{-4} \lesssim x_\pi \lesssim 0.03$ from combined HERA + Drell-Yan fit

→ global analysis under way of HERA LN, Drell-Yan $\pi N + pd/pp$ (+ upcoming JLab TDIS data) to determine pion PDFs at all x

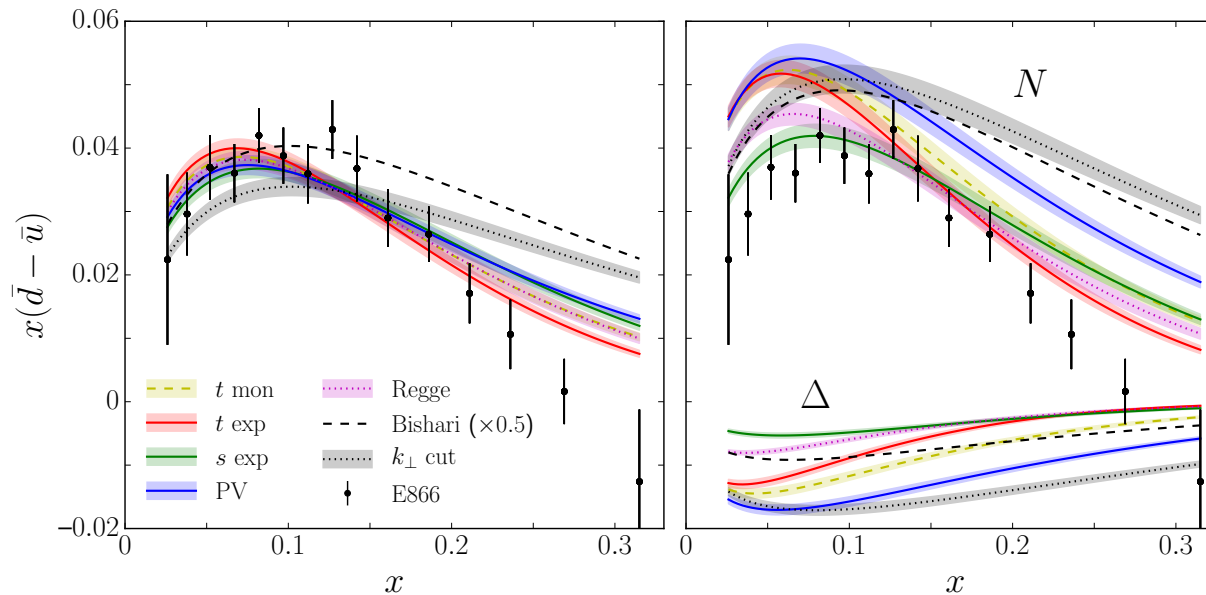


C. Keppel (Thu 9:40 am – WG1)



Sign change at large x ?

- E866 data has driven successful phenomenology through interplay of PDFs and chiral physics



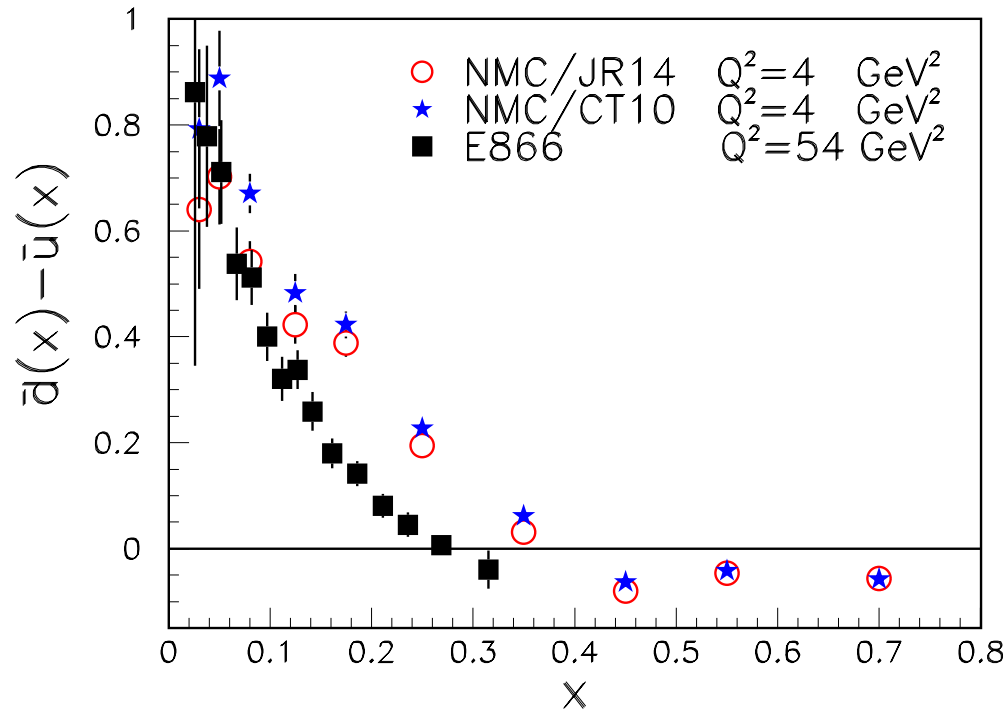
... but lingering question of possible sign change of $\bar{d} - \bar{u}$ at high x

- sign change cannot be accommodated within chiral EFT framework since (negative) Δ contribution \ll (positive) N contribution
- evidence for other mechanisms?

Sign change at large x ?

- “Independent evidence for $\bar{d} - \bar{u}$ sign change at $x \sim 0.3$ ” from NMC

Peng et al., PLB 736 (2014) 411



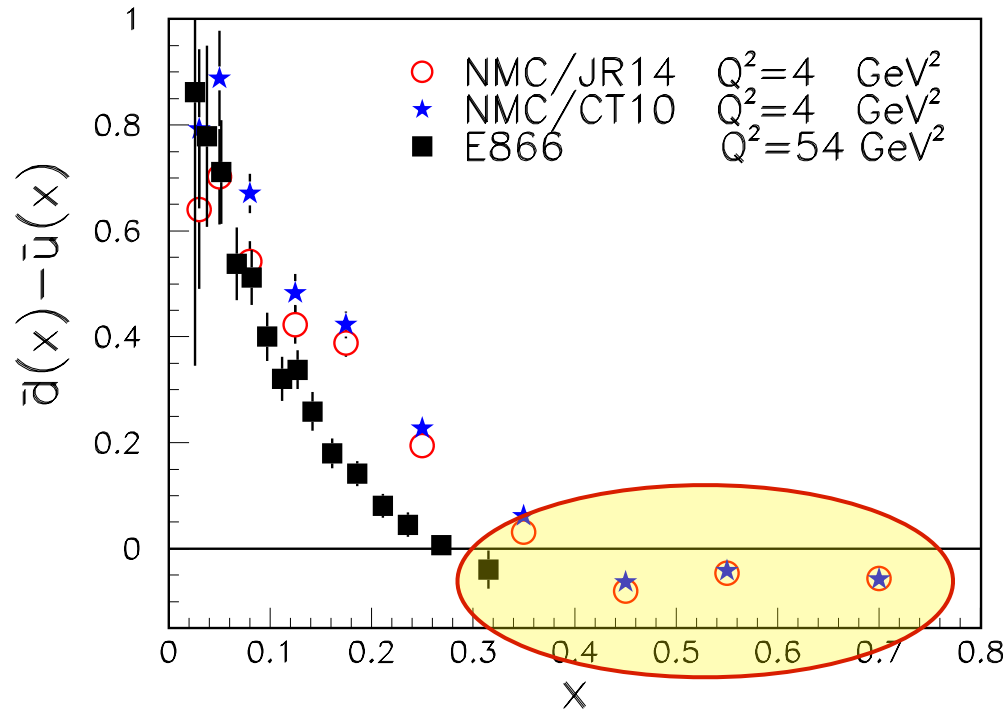
$$\bar{d} - \bar{u} \equiv \frac{1}{2}(u_v - d_v) - \frac{3}{2x}(F_2^p - F_2^n)$$

→ conclusions based on LO analysis ... how robust?

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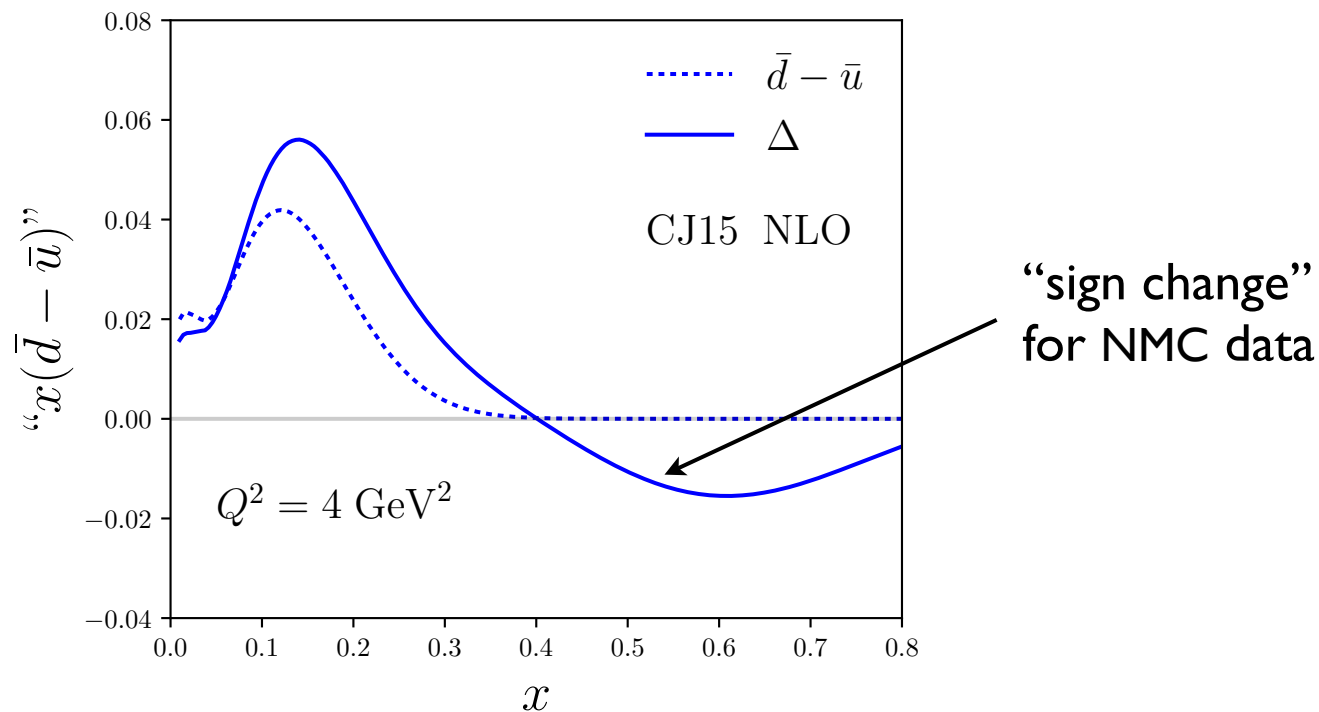
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Sign change at large x ?

- At higher order can easily generate zero crossing in

$$\Delta \equiv \frac{1}{2}(u_v - d_v) - \frac{3}{2x}(F_2^p - F_2^n)$$

with no $\bar{d} - \bar{u}$ asymmetry!



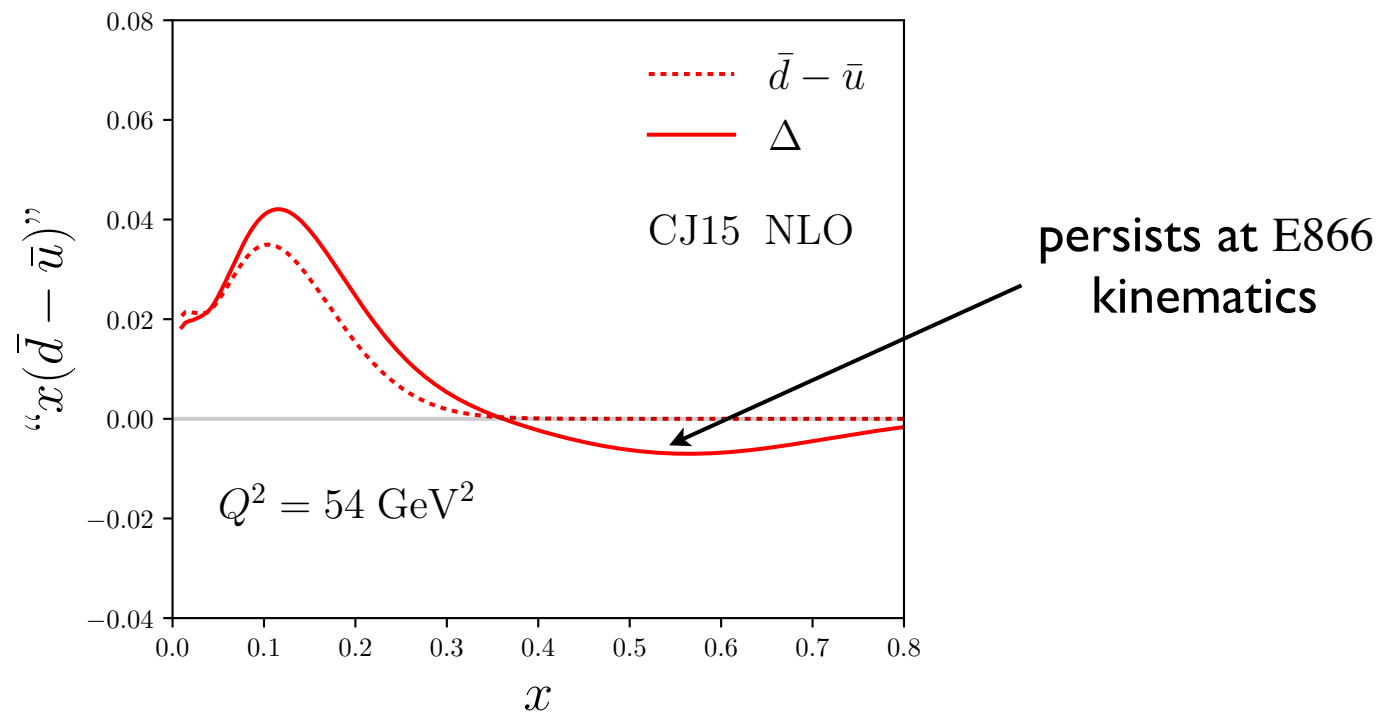
→ no evidence of sign change from DIS data!

Sign change at large x ?

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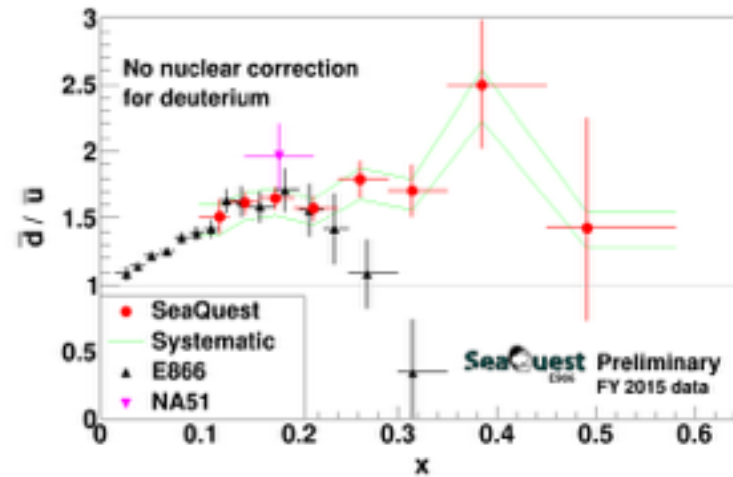
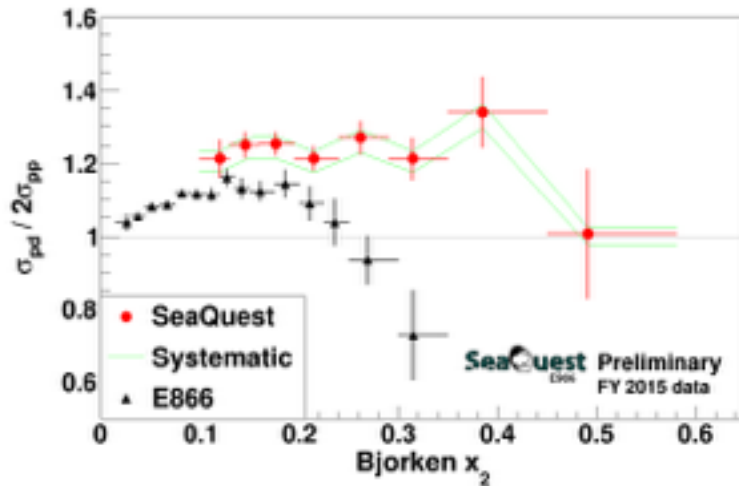
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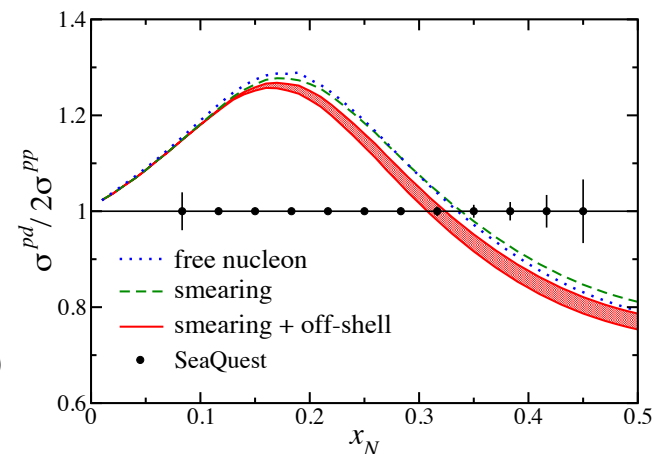
Sign change at large x ?

- Preliminary data from SeaQuest (E906) Drell-Yan experiment at Fermilab shows no evidence for sign change



P. Reimer (2016)

- SeaQuest data consistent with E866 data up to $x \sim 0.2$, remains above unity up to $x \sim 0.5$
- Results not significantly affected if include nuclear corrections

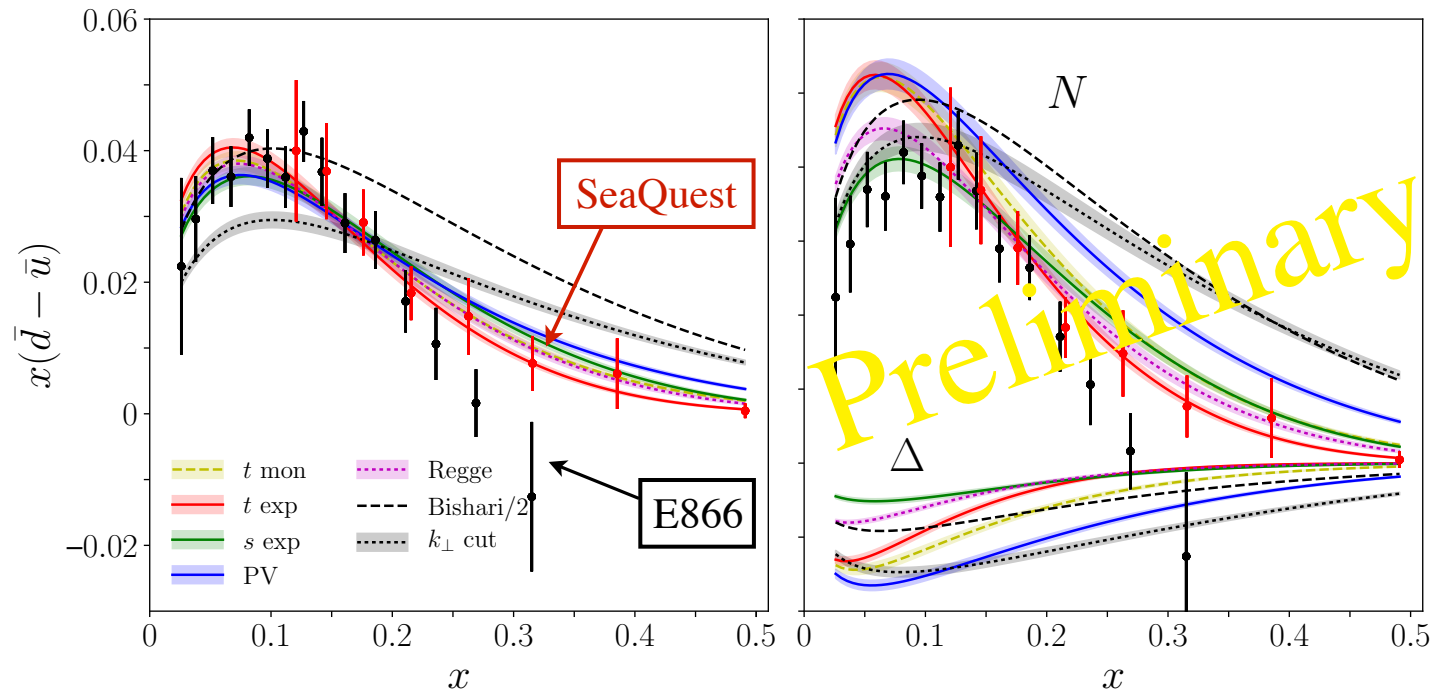


Ehlers, Accardi, Brady, WM (2014)

Sign change at large x ?

Consequences for chiral analysis?

→ including both E866 and preliminary SeaQuest data in fit ...



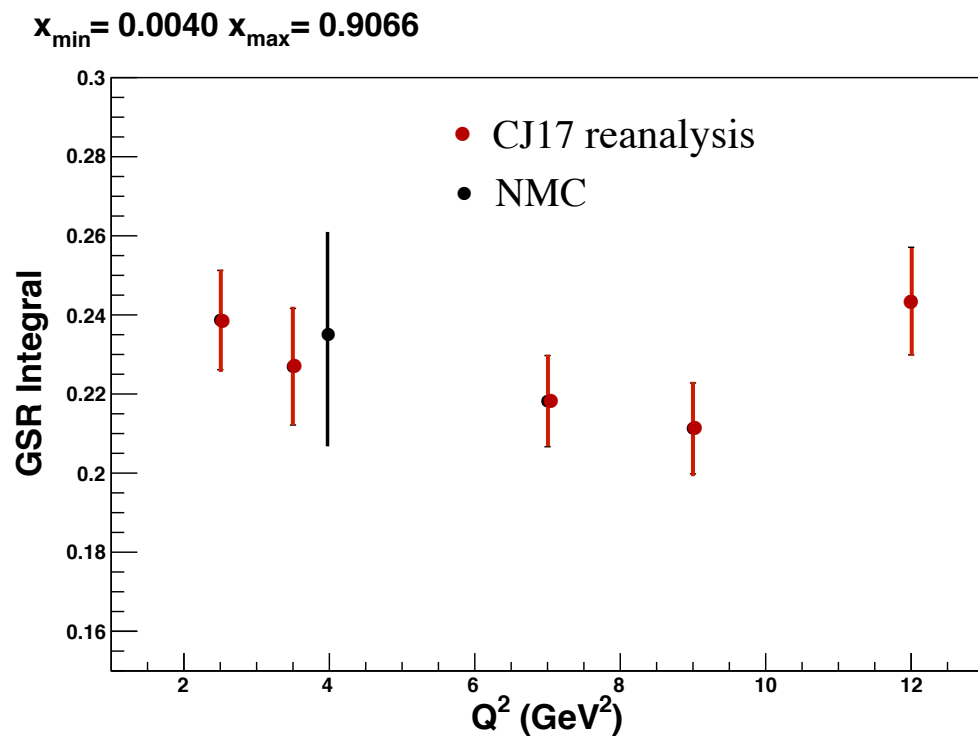
Barry, Sato, WM et al. (2017)

	E866	E866 + SeaQuest
χ^2_{dof}	1.2	0.8
Λ (GeV)	0.85	0.85
$\langle n \rangle_{\pi N}$	0.29	0.29

→ Improved fit, but almost no effect on shape!

Outlook

- Eagerly await final SeaQuest data!
 - settle question of sign change in $\bar{d} - \bar{u}$ at high x
- Combine “leading neutron” analysis with πN Drell-Yan data to constrain pion PDFs at low and high x
 - upcoming “tagged DIS” experiment at JLab → *C. Keppel (Thu 9:40 am – WG1)*
- Extend to strangeness sector to analyze $s - \bar{s}$ asymmetry within chiral SU(3) EFT framework



Niculescu et al. (2016)