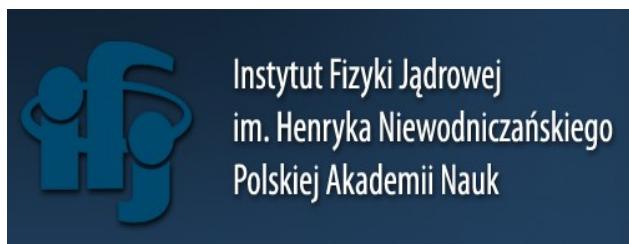


*Supported by Narodowe Centrum Nauki (NCN)
with Sonata BIS grant*



Improved TMD factorization and production of forward jets

Krzysztof Kutak



Based on
Kotko, K.K, Marquet, Petreska, Sapeta, van Hameren,
JHEP 1612 (2016) 034
JHEP 1612 (2016) 034

Why jets?

Jets are manifestations of partonic nature of hadrons which is not completely known

Jets can be used to uncover dynamics of QCD in semi perturbative and perturbative region

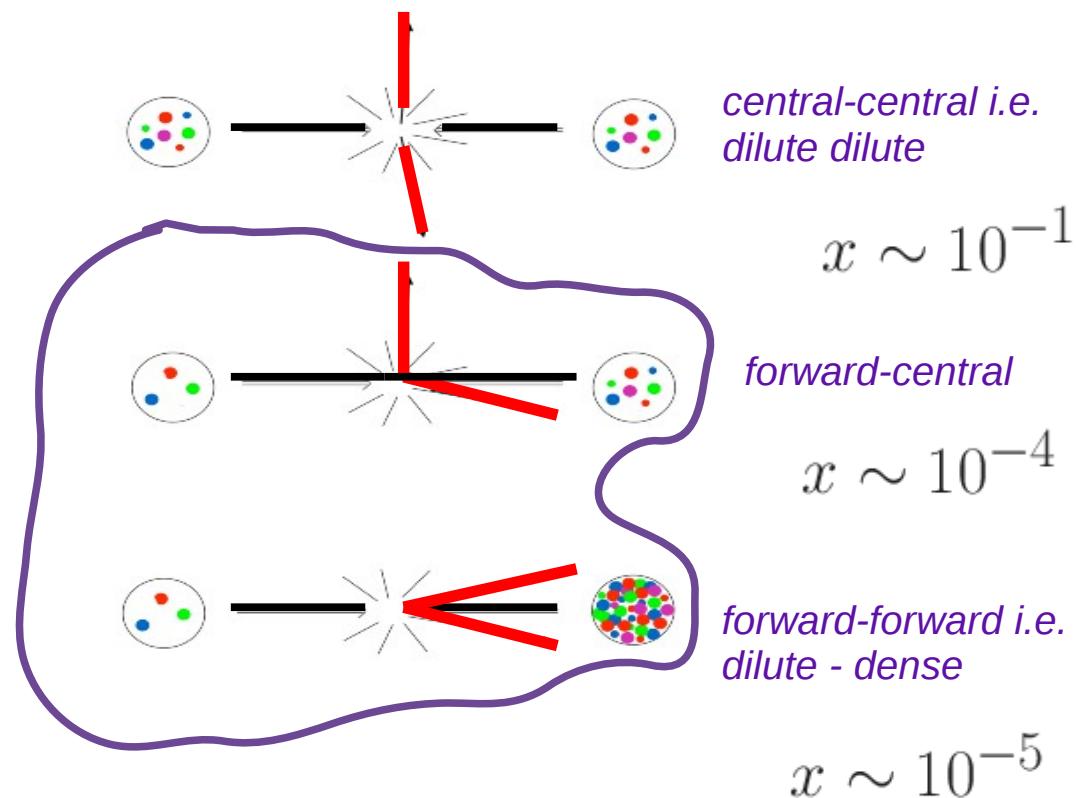
Forward jets can be used to study so called low x phenomena i.e. saturation of gluon density

Jets can be used to perform tomography of QGP

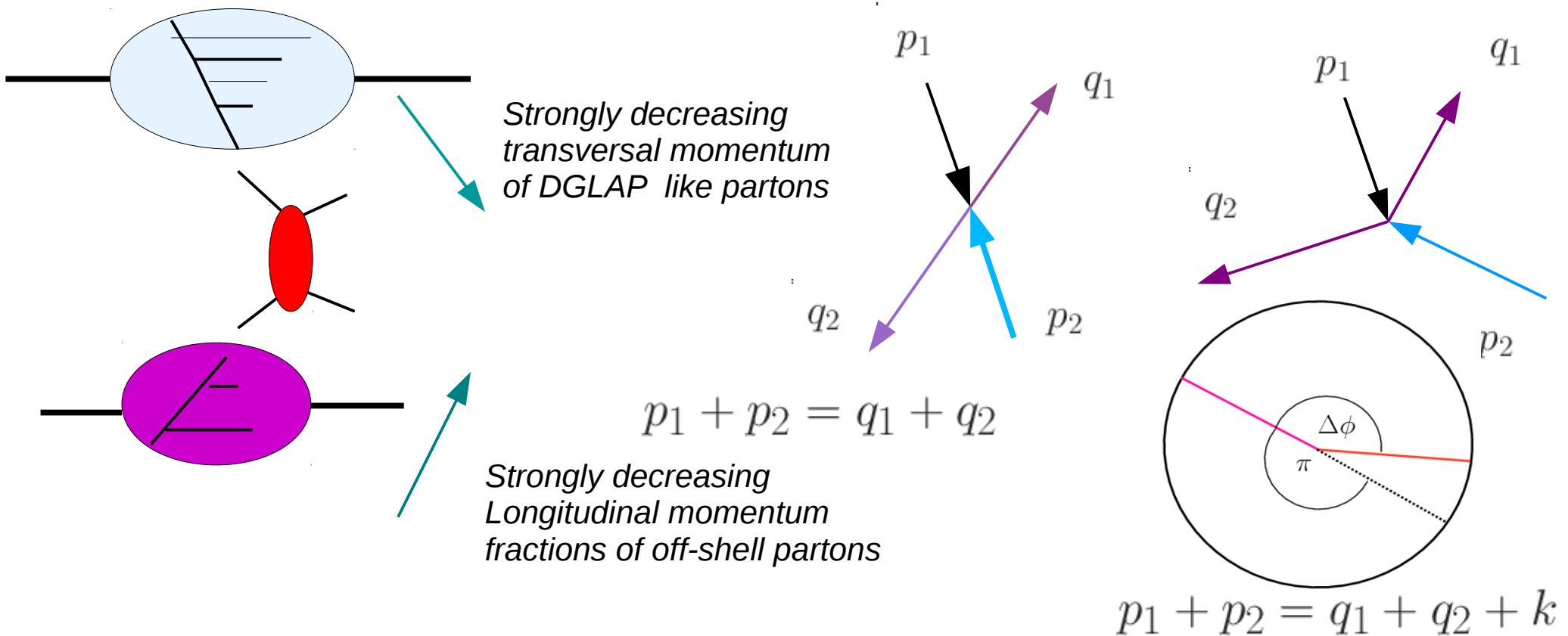
In this talk I am going to address di-jets and single inclusive jet

4-jet production → Mirko Serino

LHC and tomography of partons \rightarrow di-jet example



hybrid High Energy Factorization

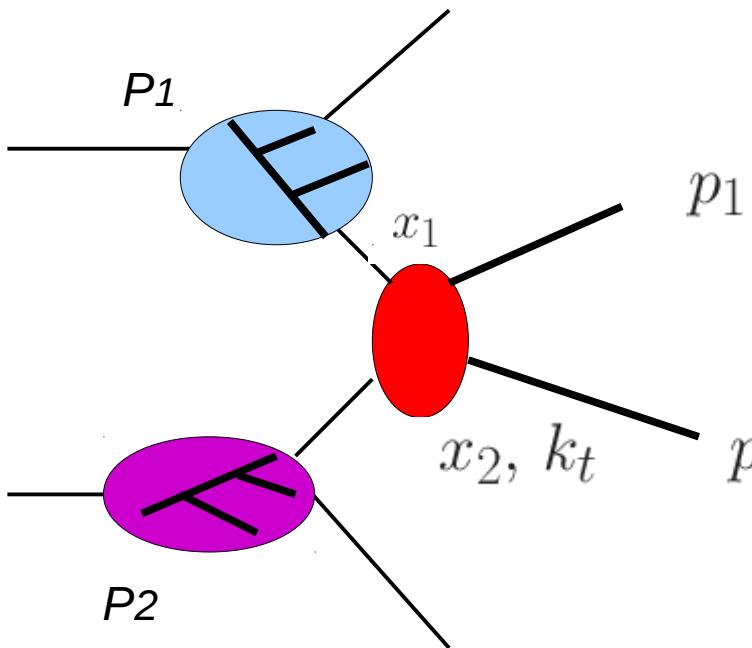


First attempt: hybrid factorization and dijets

$$\frac{d\sigma_{\text{SPS}}^{P_1 P_2 \rightarrow \text{dijets} + X}}{dy_1 dy_2 dp_{1t} dp_{2t} d\Delta\phi} = \frac{p_{1t} p_{2t}}{8\pi^2 (x_1 x_2 s)^2} \sum_{a,c,d} x_1 f_{a/P_1}(x_1, \mu^2) |\overline{\mathcal{M}_{ag^* \rightarrow cd}}|^2 \mathcal{F}_{g/P_2}(x_2, k_t^2) \frac{1}{1 + \delta_{cd}}$$

conjecture

Deak, Jung, Kutak, Hautmann '09



obtained from CGC after neglecting all nonlinearities

$g^*g \rightarrow gg$ Iancu, Laihet

$qg^* \rightarrow qg$ Van Hameren, Kotko, Kutak, Marquet, Petreska, Sapeta

resummation of logs of x

logs of hard scale

knowing well parton densities at large x one can get information about low x physics

$$x_1 = \frac{1}{\sqrt{s}} (|\vec{p}_{1t}| e^{y_1} + |\vec{p}_{2t}| e^{y_2}) \xrightarrow{y_1, y_2 \gg 0} x_1 \sim 1$$

$$x_2 = \frac{1}{\sqrt{s}} (|\vec{p}_{1t}| e^{-y_1} + |\vec{p}_{2t}| e^{-y_2}) \quad x_2 \ll 1$$

Inbalance momentum:

$$|\vec{k}_t|^2 = |\vec{p}_{1t} + \vec{p}_{2t}|^2 = |\vec{p}_{1t}|^2 + |\vec{p}_{2t}|^2 + 2|\vec{p}_{1t}||\vec{p}_{2t}|\cos\Delta\phi$$

Relevant scales and factorization

P_t average transverse momentum of dijets

k_t target gluon's transverse momentum

Q_s scale at which gluon recombination nonlinear effects at the target start to be relevant

$P_t \sim k_t$ High Energy Factorization \rightarrow partons carry some k_t

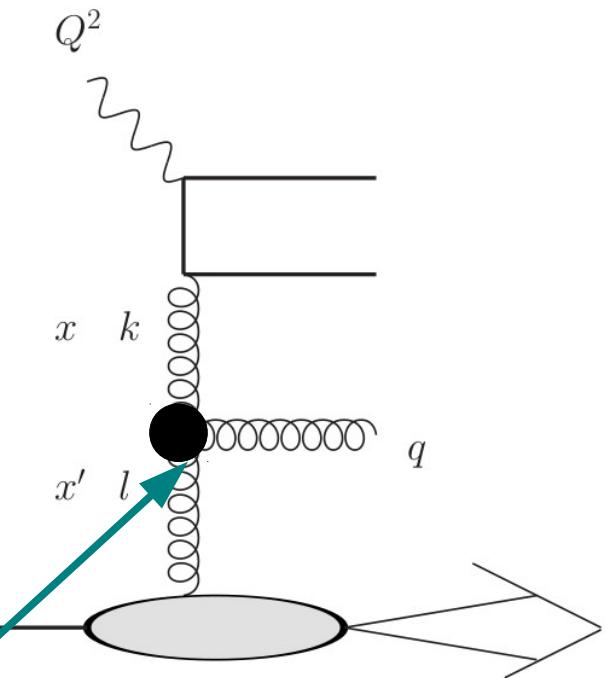
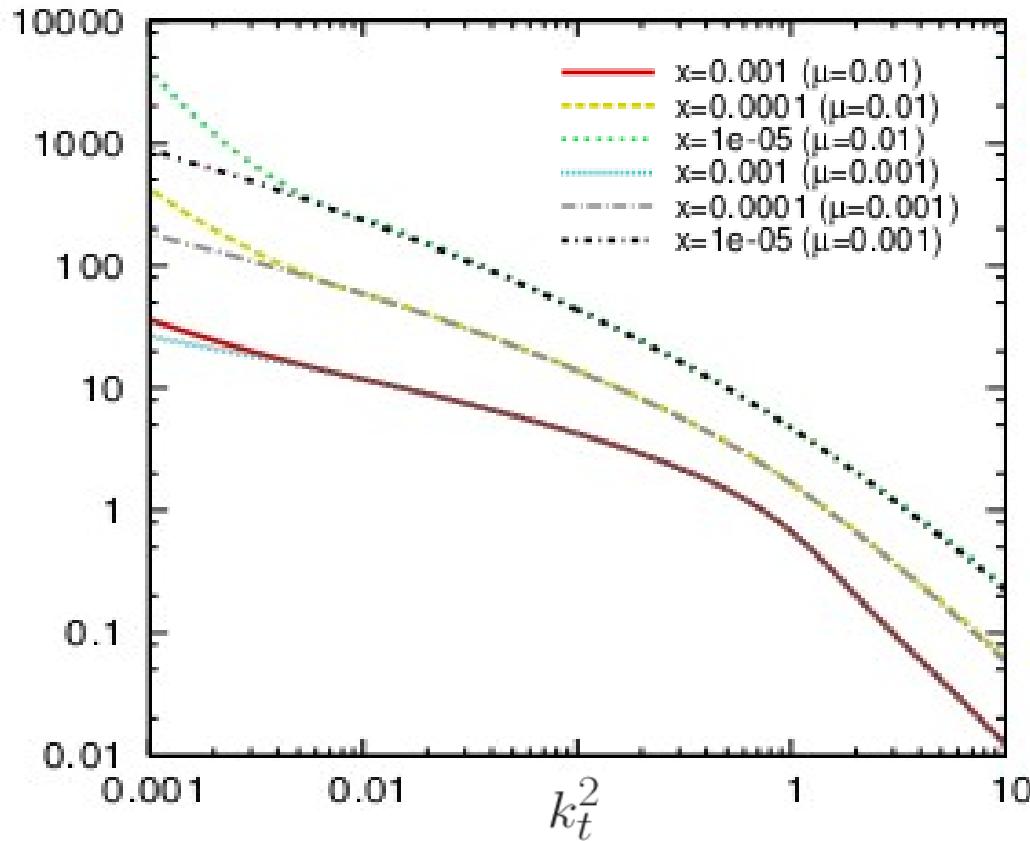
$k_t \ll P_t$ Collinear Factorization \rightarrow partons in one of hadrons are just collinear with hadron k_t is neglected

$Q_s \sim k_t \ll P_t$ generalized Transverse Momentum Dependent Factorization \rightarrow rescatterings formal treatment of nonlinearities but does not allow for calculation of decorrelations

Q_s, k_t, P_t Improved Transverse Momentum Dependent Factorization

The saturation problem: sensitivity to gluons at small k_t

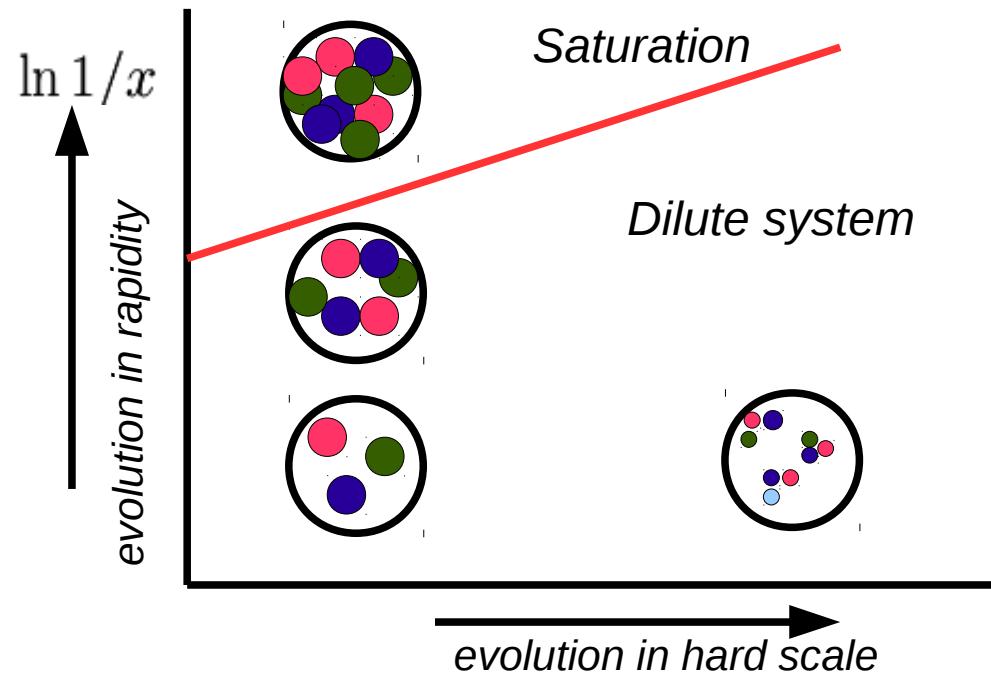
Solution of BFKL equation



$$\mathcal{F} = \mathcal{F}_0 + K \otimes \mathcal{F}$$

High energy factorization and saturation

Saturation – state where number of gluons stops growing due to high occupation number. Way to fulfill unitarity requirements in high energy limit of QCD.



On microscopic level it means that gluon apart splitting recombine

splitting

Linear evolution equation

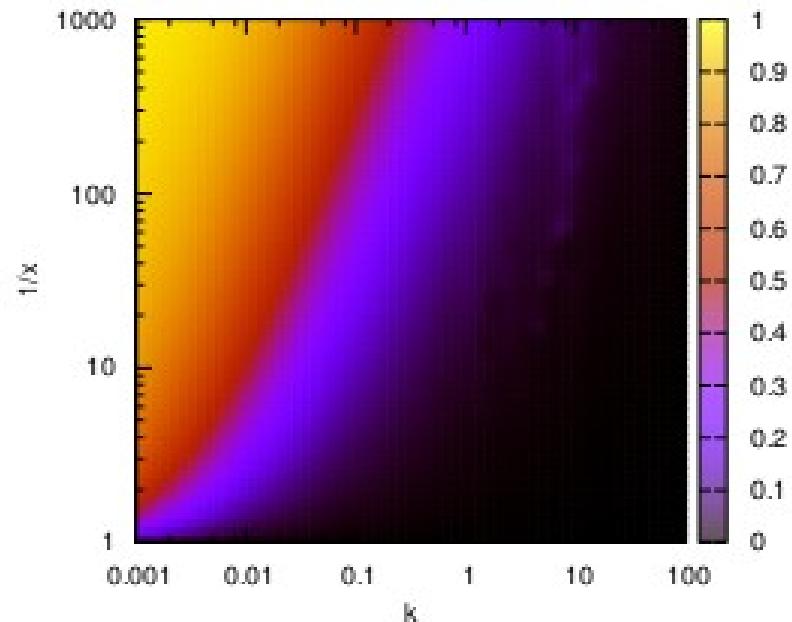
Nonlinear evolution equations

splitting

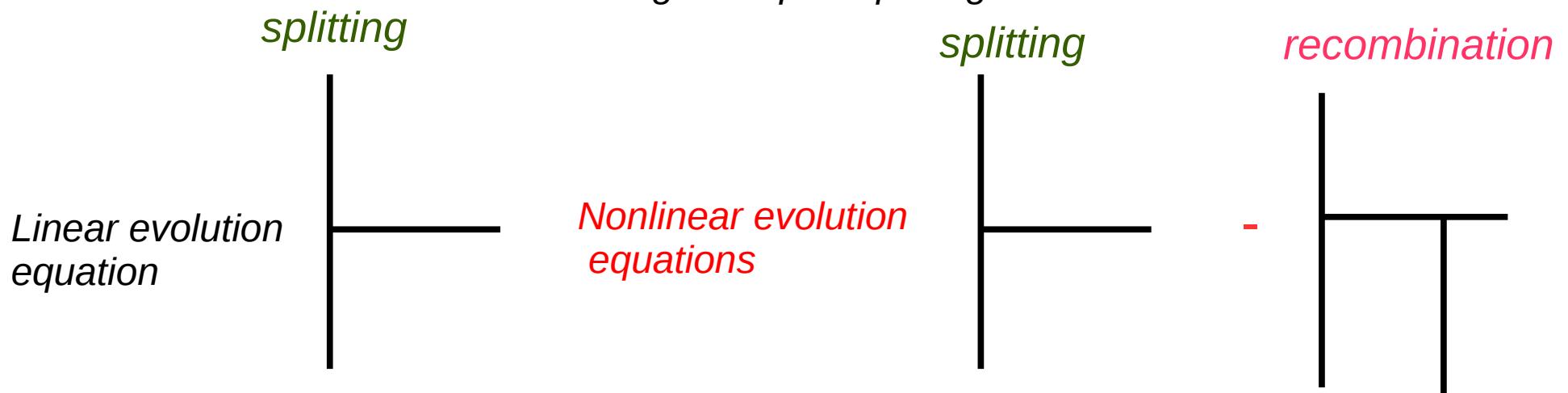
recombination

High energy factorization and saturation

Saturation – state where number of gluons stops growing due to high occupation number. Way to fulfill unitarity requirements in high energy limit of QCD.



On microscopic level it means that gluon apart splitting recombine



The saturation problem: suppressing gluons at small k_t

Originally formulated in coordinate space

Balitsky '96, Kovchegov '99

Now at NLO accuracy

Balitsky, Chirilli '07

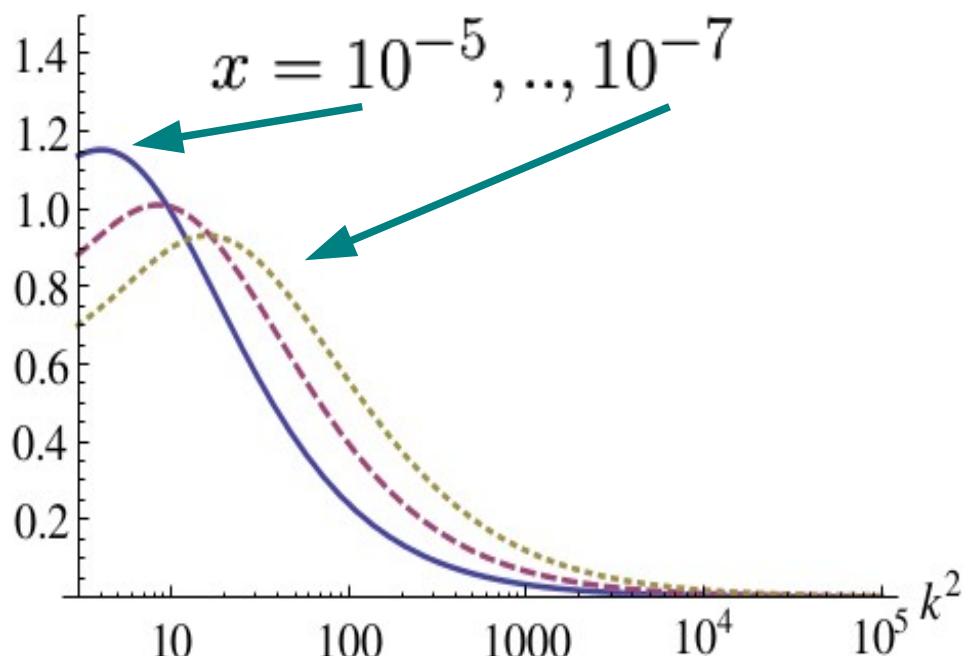
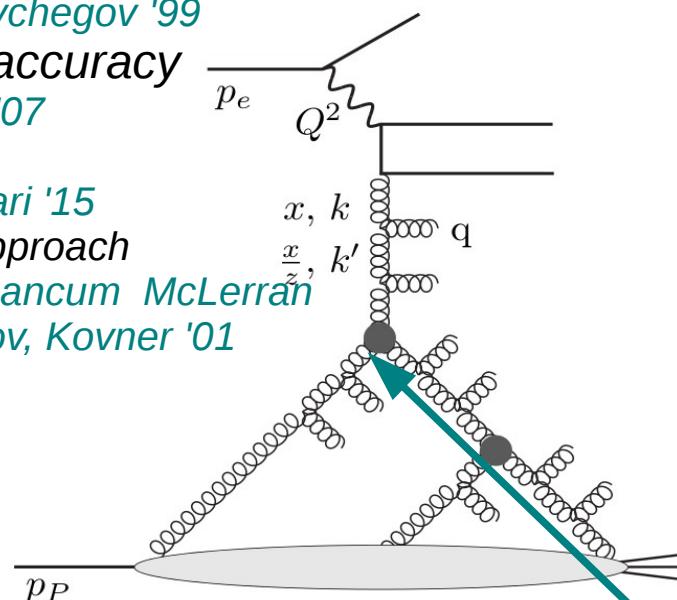
and solved

Lappi, Mantysaari '15

More general approach

Jalilian-Marian, Iancum McLerran

Weigert Leonidov, Kovner '01



Solution of the equation

The BK equation for dipole gluon density

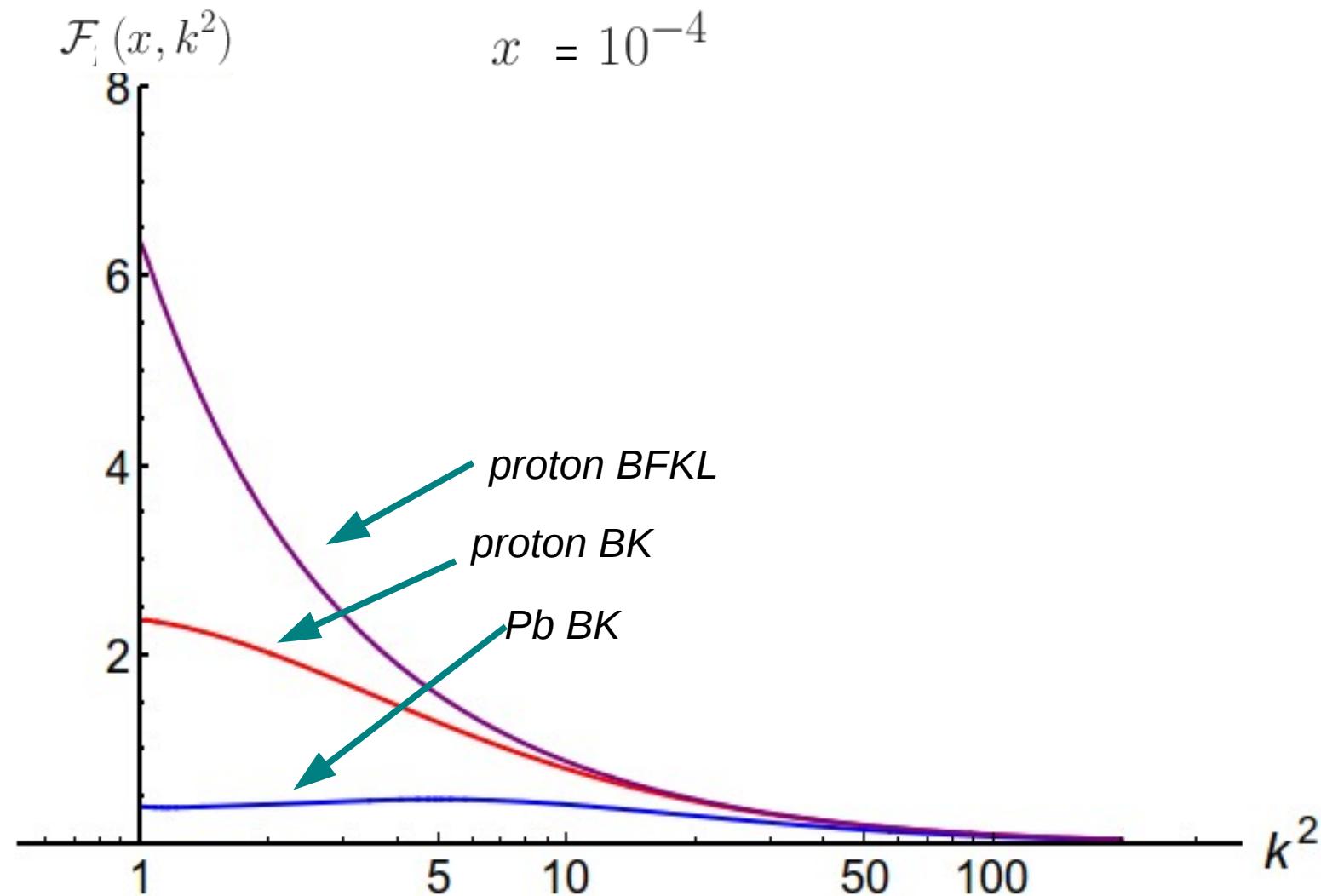
$$\mathcal{F} = \mathcal{F}_0 + K \otimes \mathcal{F} - \frac{1}{R^2} V \otimes \mathcal{F}^2$$

hadron's radius

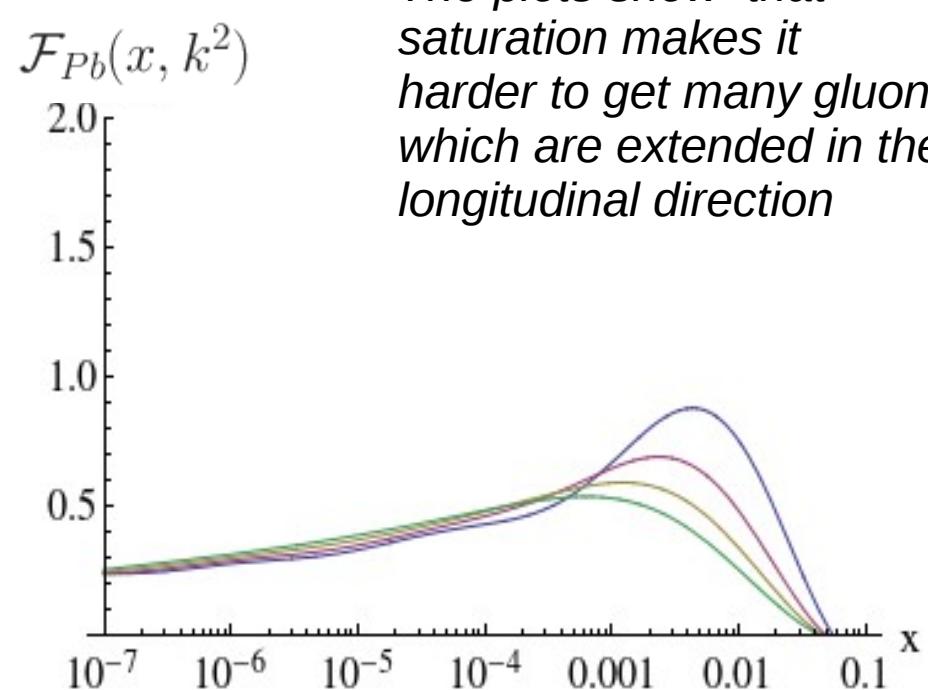
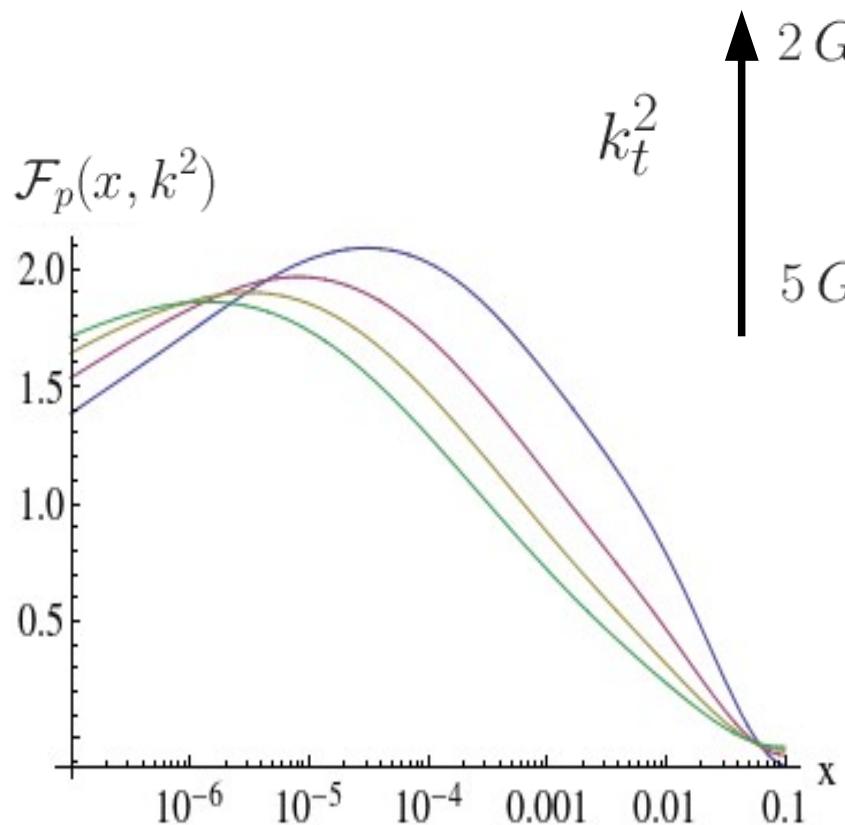
Momentum space

Kwiecinski, Kutak '02
Nikolaev, Schafer '06¹⁰

Glue in p vs. glue in Pb vs. linear - kt dependence



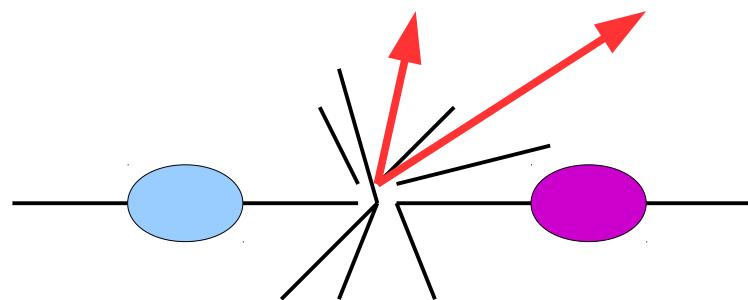
Only nonlinear - glue in p vs. glue in Pb



*In invoking uncertainty principle.
The plots show that
saturation makes it
harder to get many gluons
which are extended in the
longitudinal direction*

Maximum signalize emergence of saturation scale

Central-forward di-jets



PDF we use at present

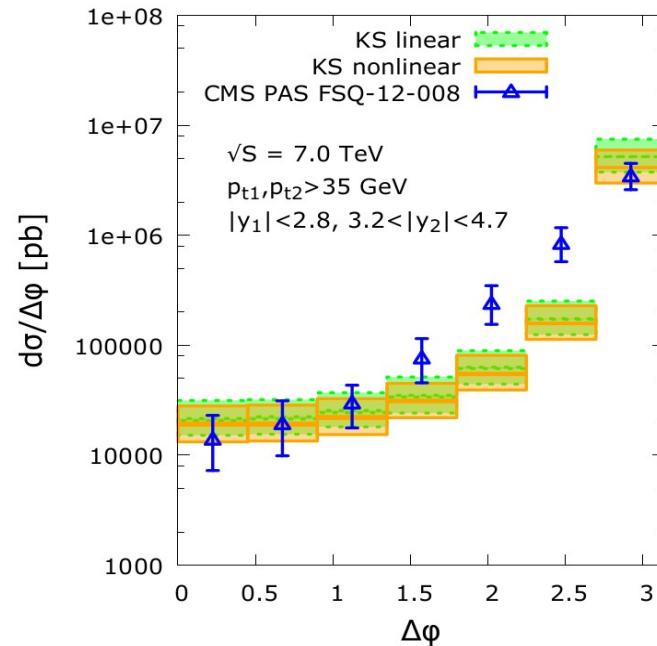
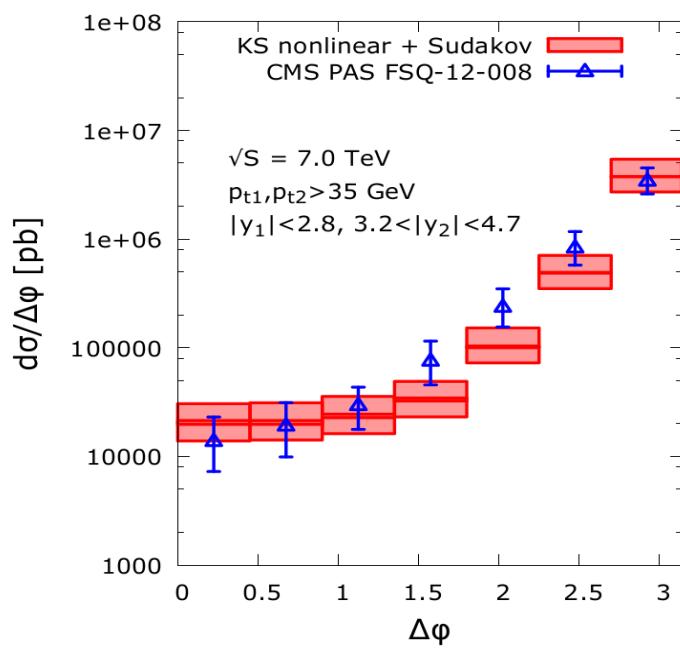
KS (Kutak-Sapeta) nonlinear → gluon density from extension of momentum space version of BK equation to include:

- *kinematical constraint*
- *complete splitting function,*
- *running coupling*
- *quarks*

KK, Kwiecinski '03 fitted to '10 HERA data KK, Sapeta '12, nonlinear extension of unified BFKL+DGLAP Kwiecinski, Martin, Stašto framework '97.

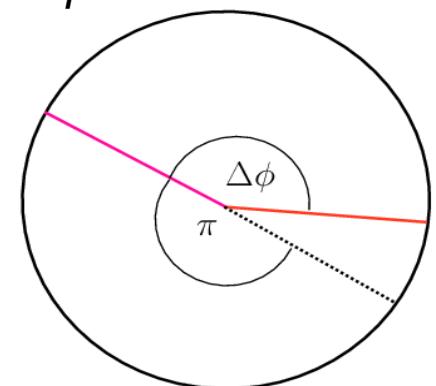
Decorelations inclusive scenario forward-central

Kotko, K.K, Sapeta, van Hameren '14



In DGLAP approach
i.e $2 \rightarrow 2 + \text{pdf}$ one would get delta function

$p_{t1}, p_{t2} > 35 \text{ GeV}$
 $3.2 < |y_2| < 4.7$
 $|y_1| < 2.8$
Leading jets, no further requirement

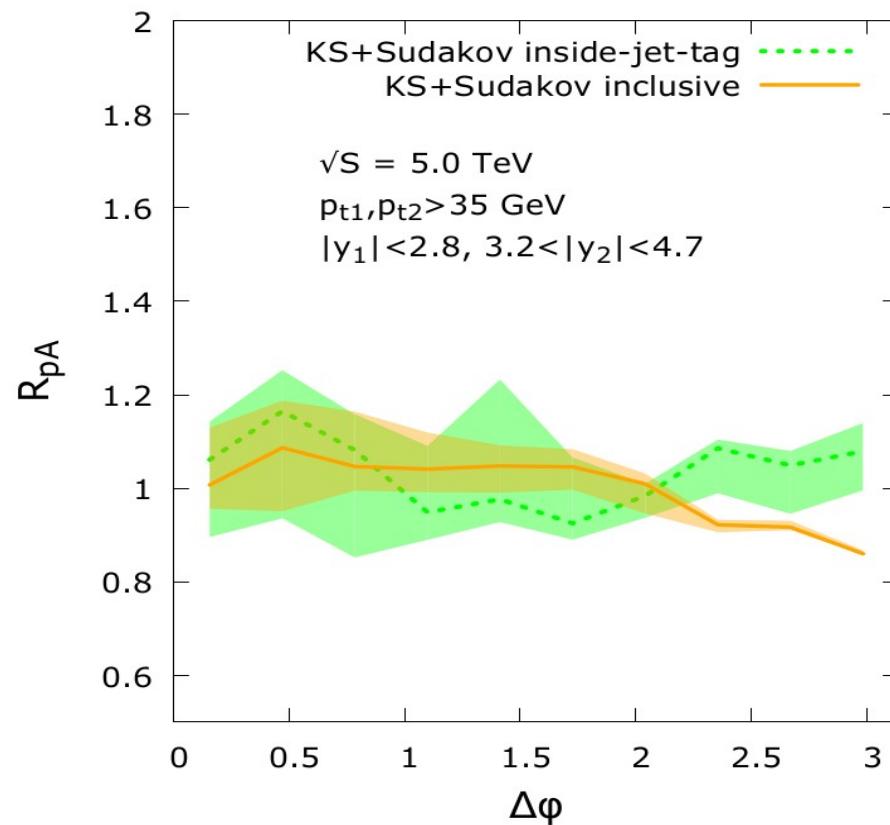


Observable suggested to study BFKL effects
Sabio-Vera, Schwensen '06

Studied also context of RHIC
Albacete, Marquet '10

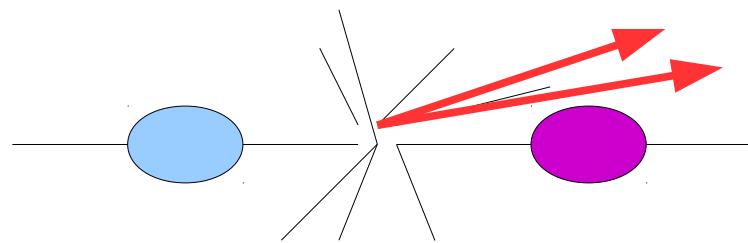
Predictions for p-Pb for forward-central

P.Kotko, KK, S.Sapeta, A. van Hameren '14



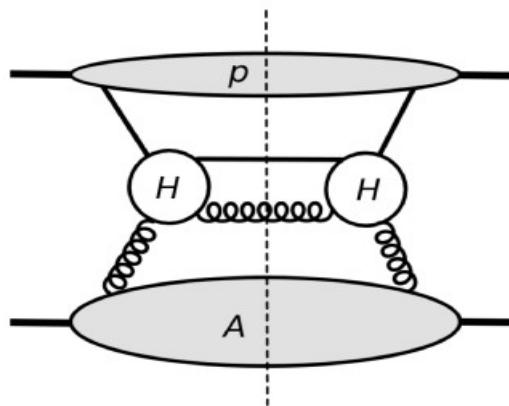
saturation effects are rather weak for forward-central jets

Forward-forward di-jets



Towards TMD for dijets in pA

The used factorization formula for dijets is strictly valid in linear regime and was calculated in a specific gauge. Results for dijets based on it with usage of gluon density coming from nonlinear equation can estimate of strength of saturation. We want to go beyond this



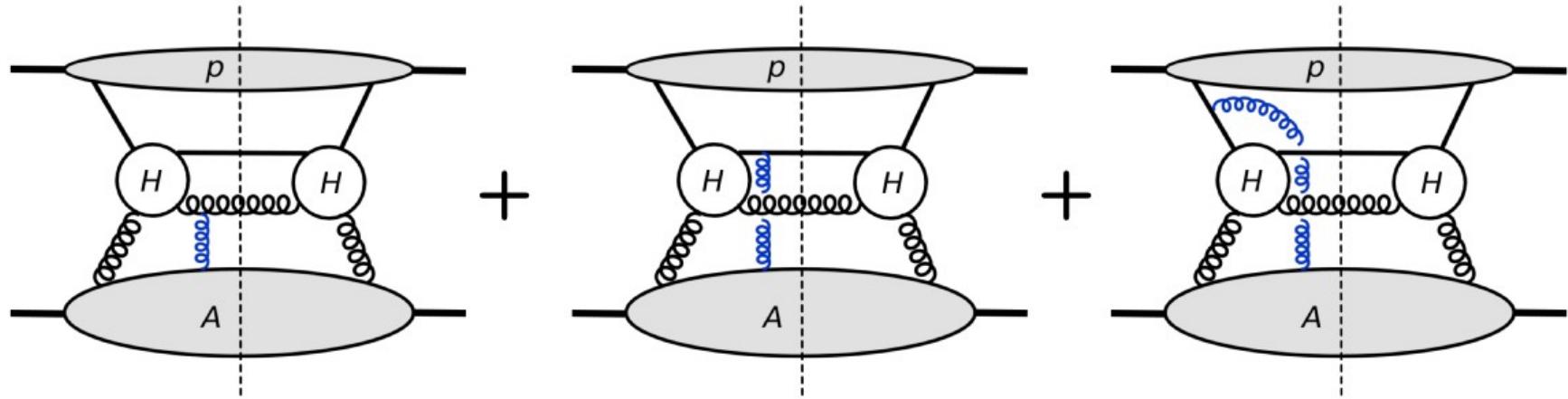
$$\frac{d\sigma^{pA \rightarrow \text{dijets}+X}}{dy_1 dy_2 d^2 p_{1t} d^2 p_{2t}} = \frac{1}{16\pi^3 (x_1 x_2 s)^2} \sum_{a,c,d} x_1 f_{a/p}(x_1, \mu^2) |\mathcal{M}_{ag^* \rightarrow cd}|^2 \mathcal{F}_{g/A}(x_2, k_t) \frac{1}{1 + \delta_{cd}}$$

$$\mathcal{F}_{g/A}(x_2, k_t) \stackrel{\text{naive}}{=} 2 \int \frac{d\xi^+ d^2 \xi_t}{(2\pi)^3 p_A^-} e^{ix_2 p_A^- \xi^+ - ik_t \cdot \xi_t} \langle A | \text{Tr} [F^{i-}(\xi^+, \xi_t) F^{i-}(0)] | A \rangle$$

Taking all complexity into account leads to following generalization of formula above...

Bomhof, Mulders and Pijlman '16.

Towards TMD for dijets in pA – gauge link



+ similar diagrams with 2,3,...gluon exchanges.

All this need to be resummed

Bomhof, Mulders, Pijlman 06

This is achieved via gauge link which
renders the gluon density gauge invariant

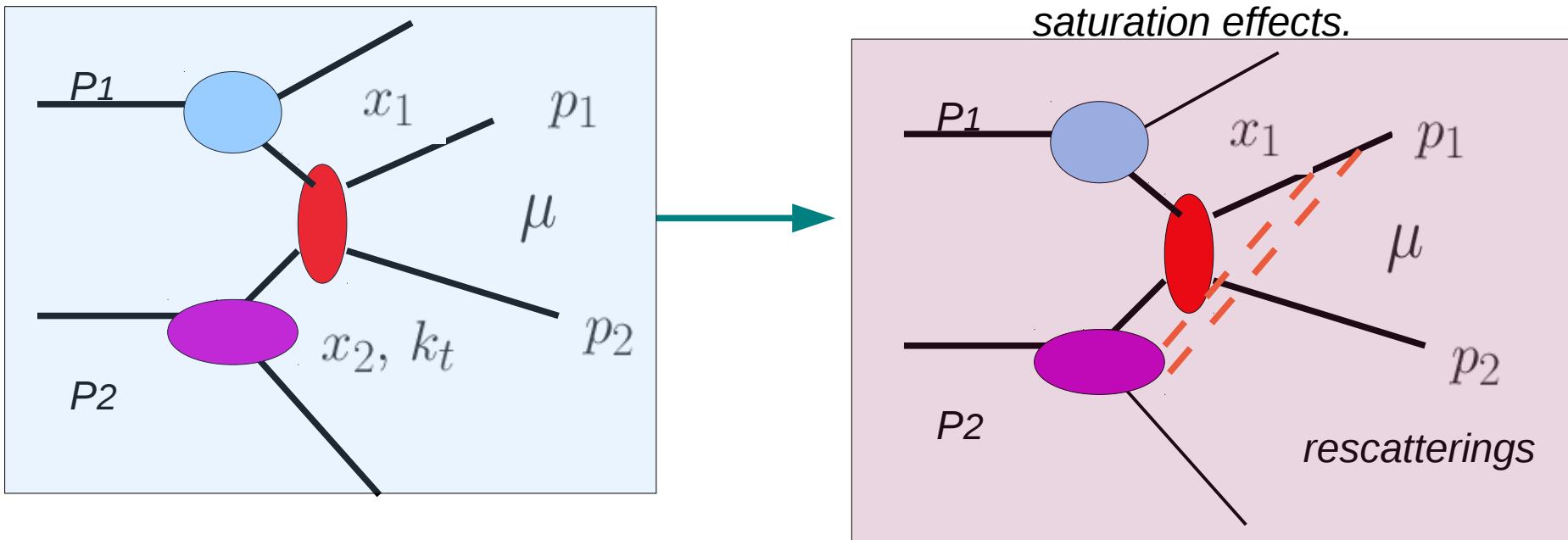
$$\mathcal{P} \exp \left[-ig \int_{\alpha}^{\beta} d\eta^{\mu} A^a(\eta) T^a \right]$$

$$\mathcal{F}_{g/A}(x_2, k_t) = 2 \int \frac{d\xi^+ d^2 \xi_t}{(2\pi)^3 p_A^-} e^{ix_2 p_A^- \xi^+ - ik_t \cdot \xi_t} \langle A | \text{Tr} [F^{i-} (\xi^+, \xi_t) U_{[\xi, 0]} F^{i-} (0)] | A \rangle$$

Improved TMD for dijets

$$\frac{d\sigma_{\text{SPS}}^{P_1 P_2 \rightarrow \text{dijets} + X}}{dy_1 dy_2 dp_{1t} dp_{2t} d\Delta\phi} = \frac{p_{1t} p_{2t}}{8\pi^2 (x_1 x_2 s)^2} \sum_{a,c,d} x_1 f_{a/P_1}(x_1, \mu^2) |\overline{\mathcal{M}_{ag^* \rightarrow cd}}|^2 \mathcal{F}_{g/P_2}(x_2, k_t^2) \frac{1}{1 + \delta_{cd}}$$

can be used for estimates of saturation effects.



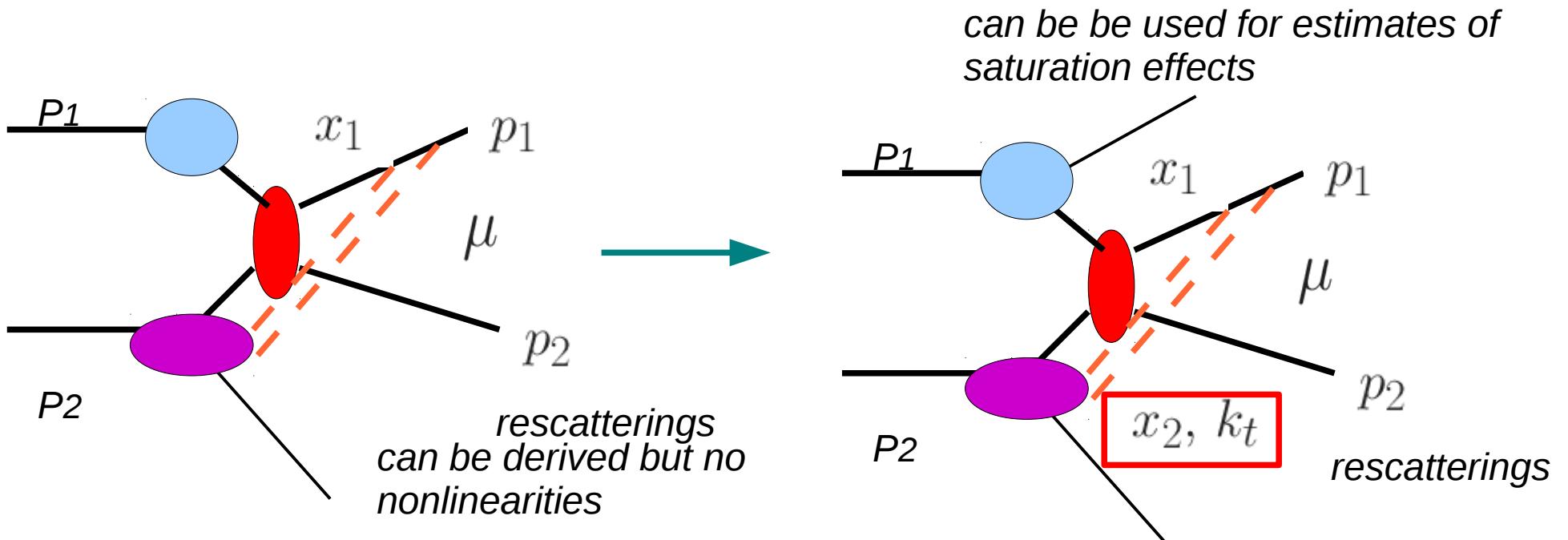
Color structure defines gauge links, different for any sub process

Generalization but **no possibility to calculate decorrelations** since no k_t in ME
Dominguez, Marquet, Xiao, Yuan '11

Application to differential distributions in $d+\text{Au}$
Stasto, Xiao, Yuan '11

$$\frac{d\sigma^{pA \rightarrow cdX}}{d^2 P_t d^2 k_t dy_1 dy_2} = \frac{\alpha_s^2}{(x_1 x_2 s)^2} x_1 f_{q/p}(x_1, \mu^2) \sum_{i=1}^n \mathcal{F}_{ag}^{(i)} H_{ag \rightarrow cd}^{(i)} \frac{1}{1 + \delta_{cd}}$$

Improved TMD for dijets

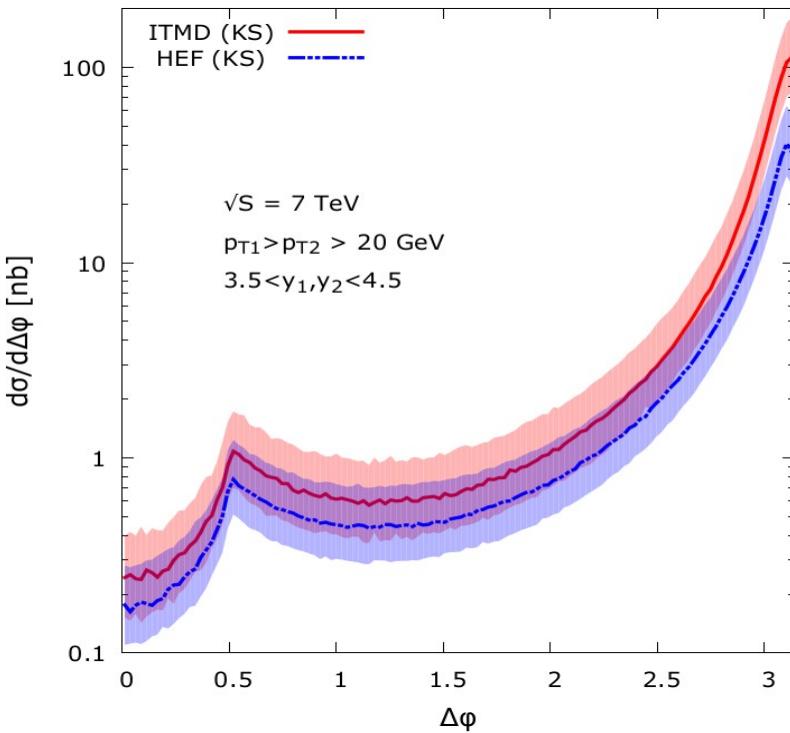
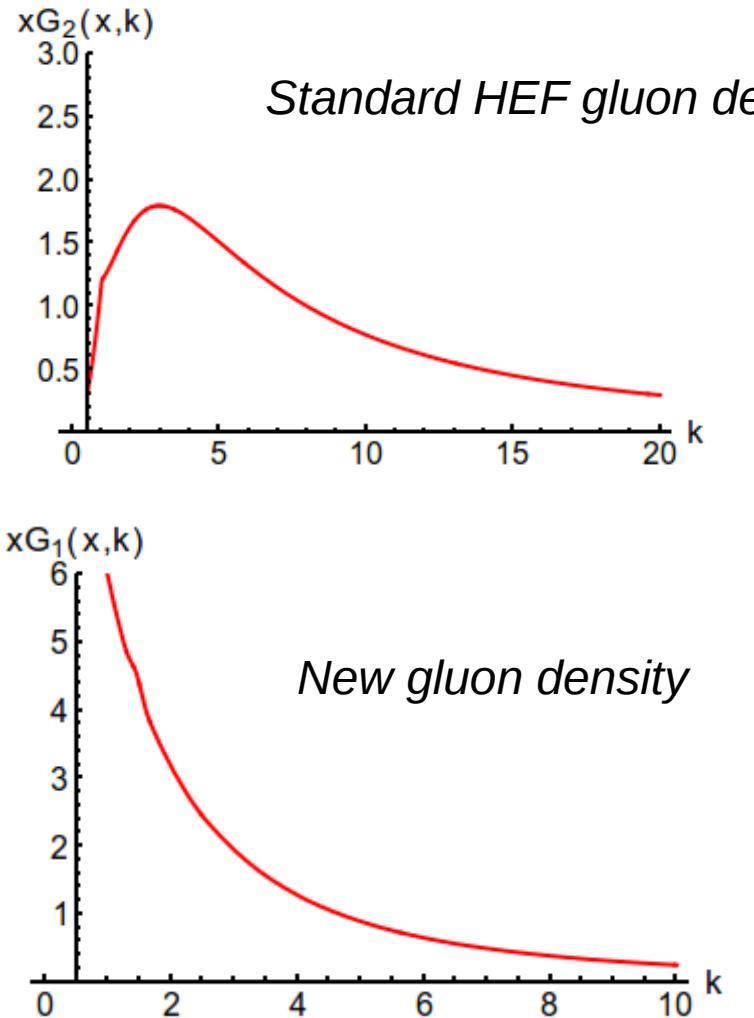


We found a method to include k_t in ME and express the factorization formula in terms of gauge invariant sub amplitudes → more direct relation to two fundamental gluon densities: **dipole gluon density** and **Weizacker-Williams gluon density**
 Kotko, K.K, Marquet, Petreska, Sapeta, van Hameren '15

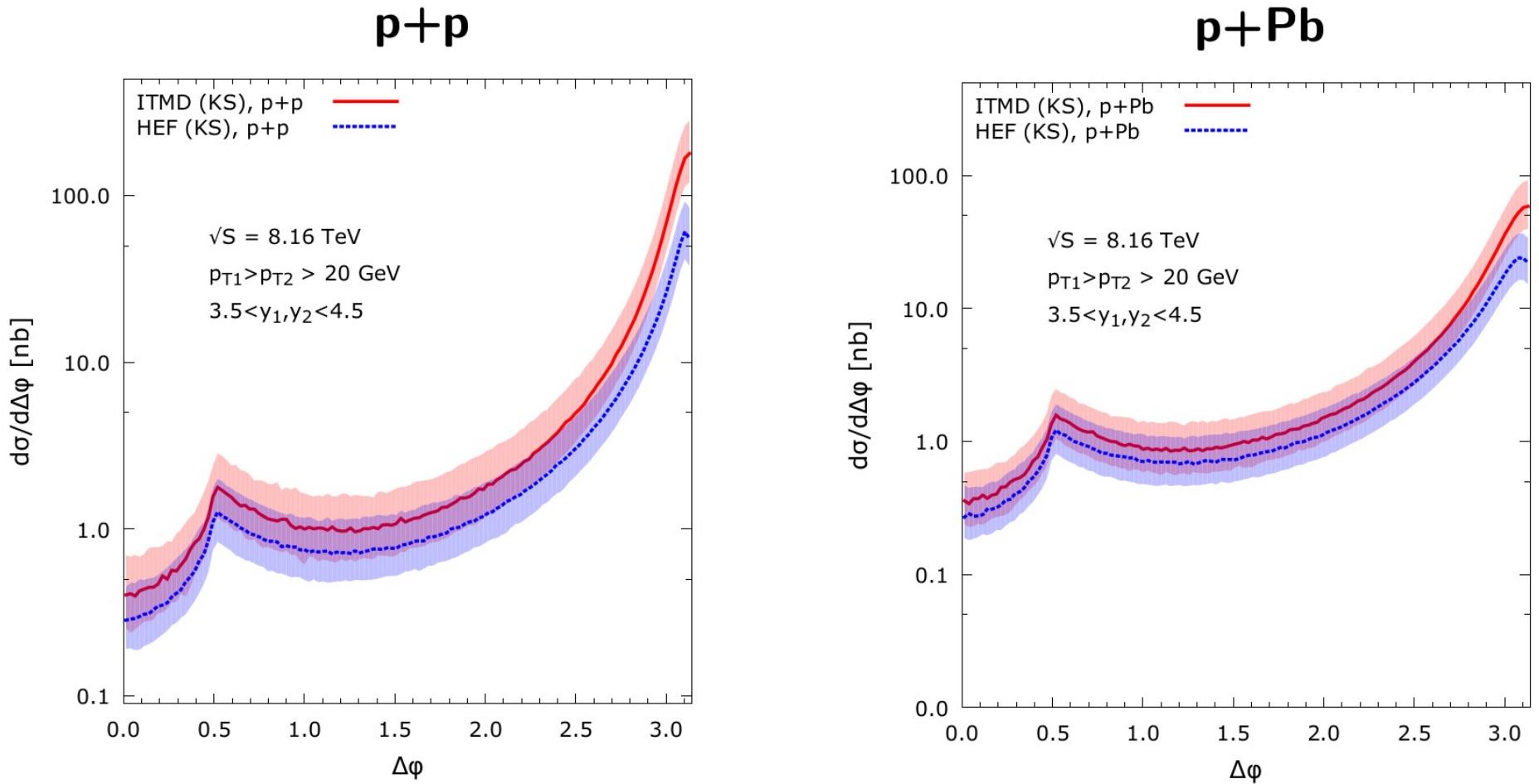
$$\frac{d\sigma^{pA \rightarrow \text{dijets}+X}}{d^2 P_t d^2 k_t dy_1 dy_2} = \frac{\alpha_s^2}{(x_1 x_2 s)^2} \sum_{a,c,d} x_1 f_{a/p}(x_1, \mu^2) \sum_{i=1}^2 K_{ag^* \rightarrow cd}^{(i)} \Phi_{ag \rightarrow cd}^{(i)} \frac{1}{1 + \delta_{cd}} {}^1$$

Glimpse on the first results – HEF vs. ITMD

Kotko, K.K, Marquet, Petreska, Sapeta, van Hameren '16

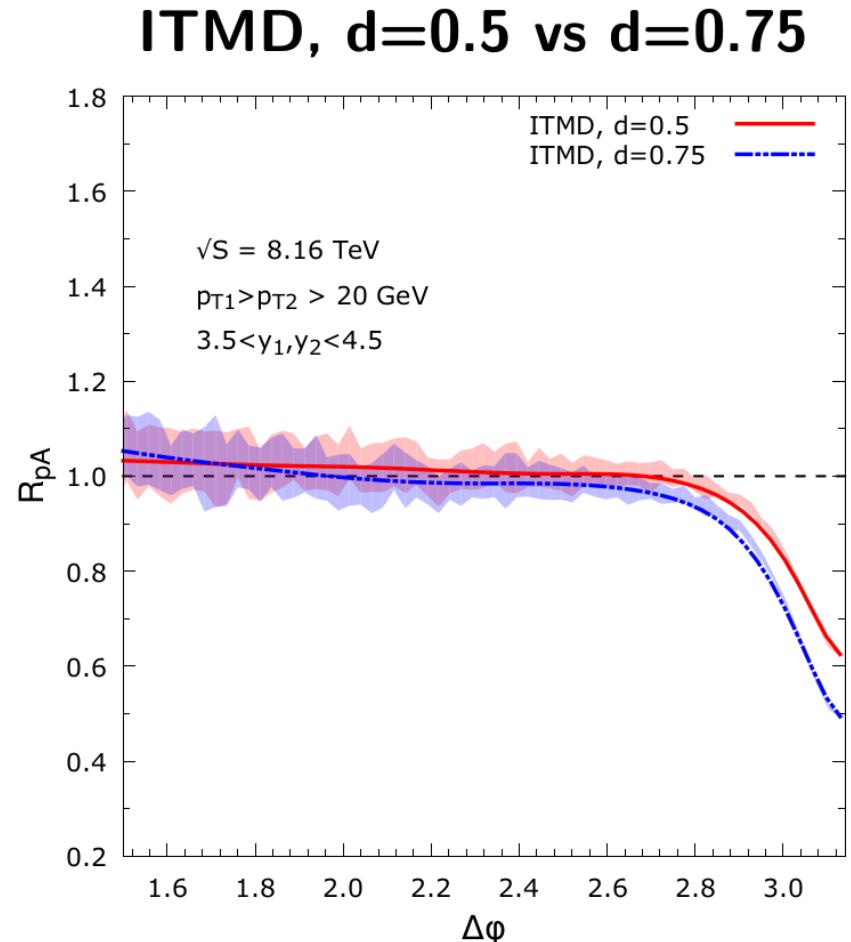
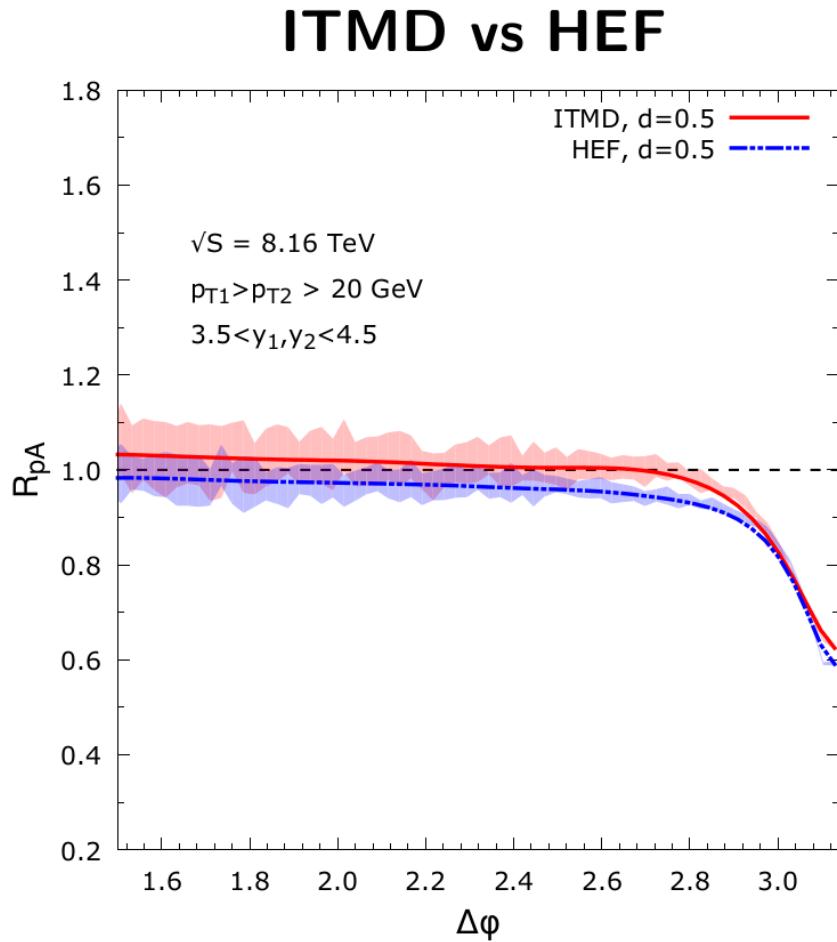


Azimuthal distance between jets



Differences in normalization but altogether similar behavior when switching from HEF to ITMD

R_{pA} : azimuthal distance between jets



R_{pA} comes very similar since differences in the cross section effectively cancel out

Conclusions and outlook

- *New framework ITMD for calculations of forward dijets has been developed*
- *The framework already gives phenomenological predictions for forward-forward dijets*
- *The results for di-jets are consistent with PYTHIA MC simulation. However, the results would agree better after implementing FSR into HEF*
- *Single inclusive jets in p+Pb at very forward rapidities*
- *Update the parton densities*

Backup

Other relevant effects – Sudakov form factor in ISR

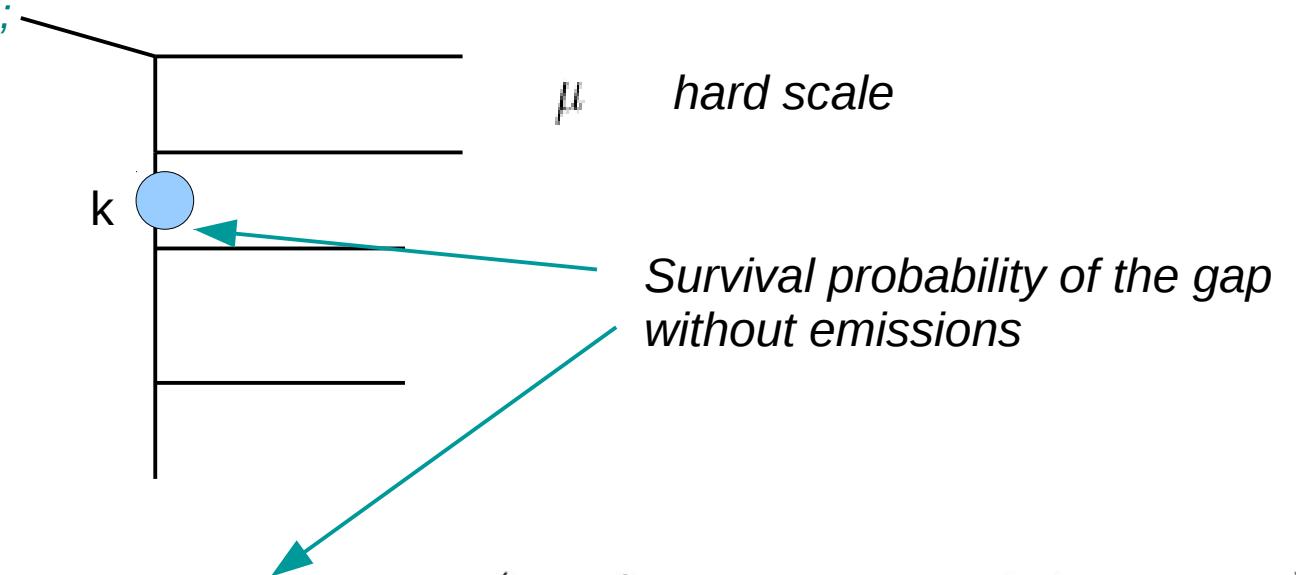
The relevance in low x physics
at linear level recognized by:

Catani, Ciafaloni, Fiorani, Marchesini;

Kimber, Martin, Ryskin;

Collins, Jung

Survival probability
of the gap without
emissions



Kimber, Martin, Ryskin procedure '01: $T_s(\mu^2, k^2) = \exp \left(- \int_{k^2}^{\mu^2} \frac{dk'^2}{k'^2} \frac{\alpha_s(k'^2)}{2\pi} \sum_{a'} \int_0^{1-\Delta} dz' P_{a'a}(z') \right)$

$$\mathcal{F}(x, k^2, \mu^2) \sim \partial_{\lambda^2} (T(\lambda^2, \mu^2) x g(x, \lambda^2))|_{\lambda^2=k^2}$$

$$\Delta = \frac{\mu}{\mu+k}$$

Mueller, Xiao, Yan '12

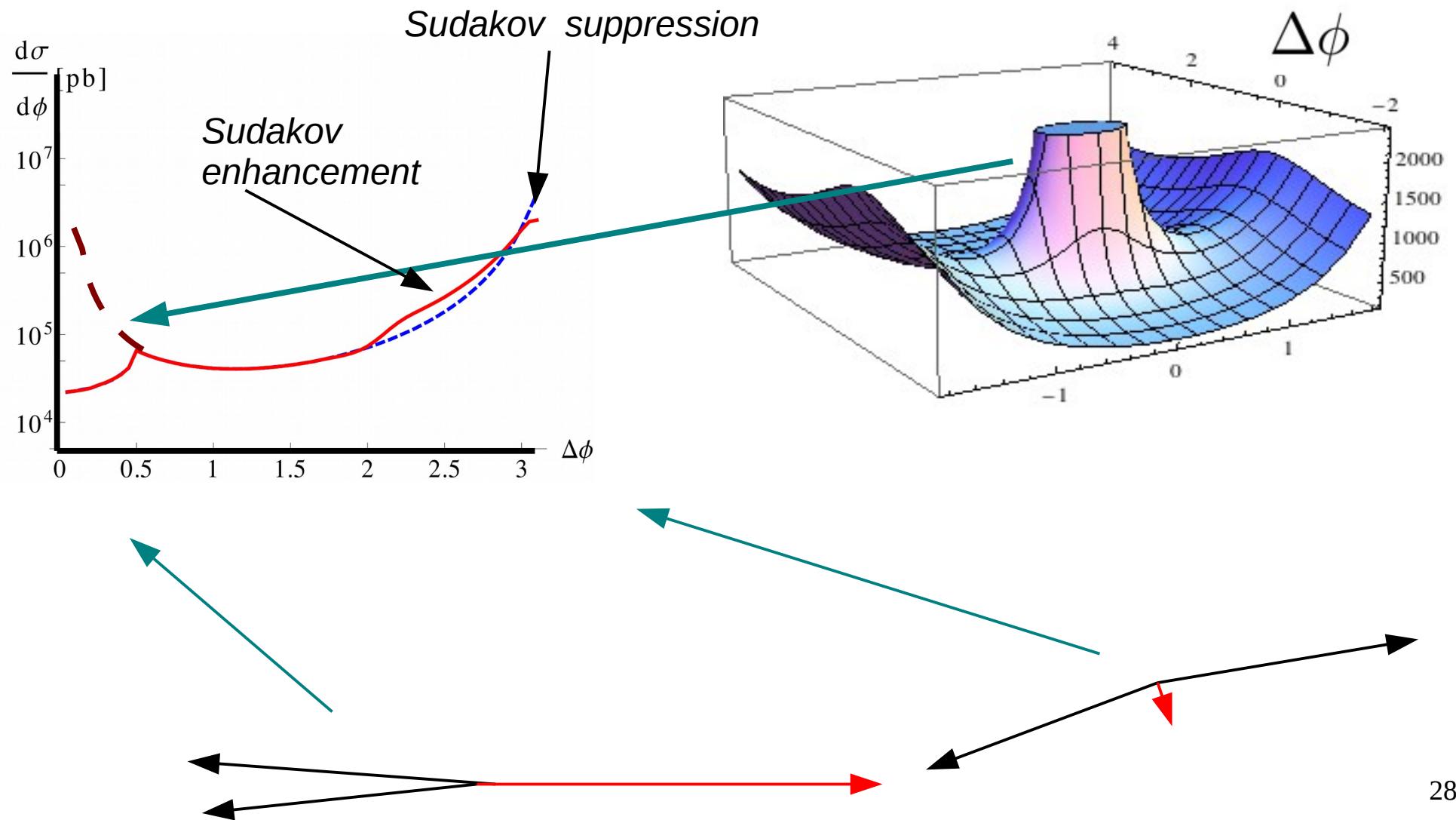
Mueller, Xiao, Yan '13

Kutak '14

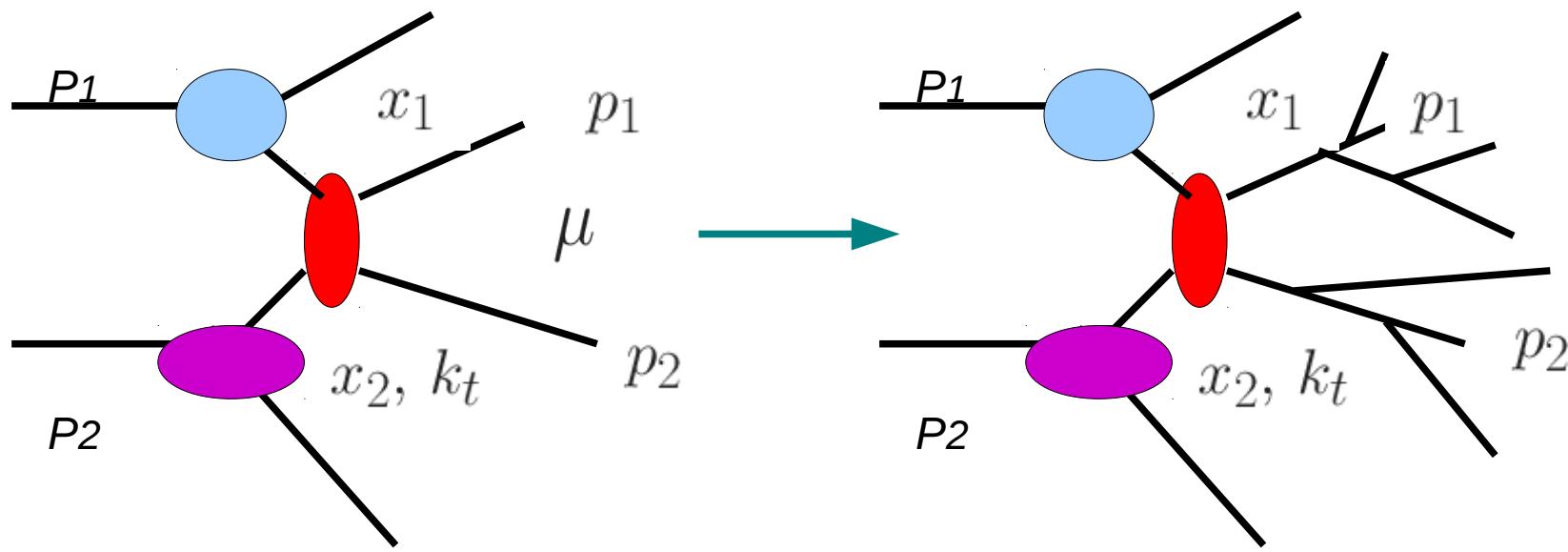
Other relevant effects – Sudakov form factor in ISR

Divergence regularized
by jet algorithm

KK '14



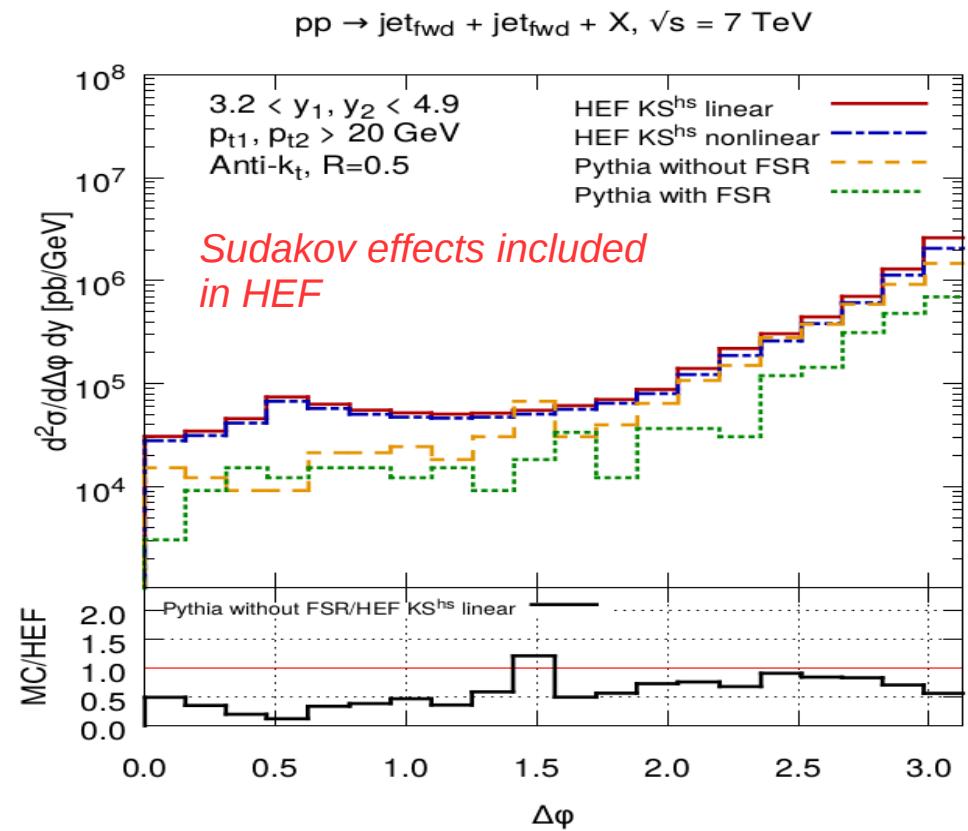
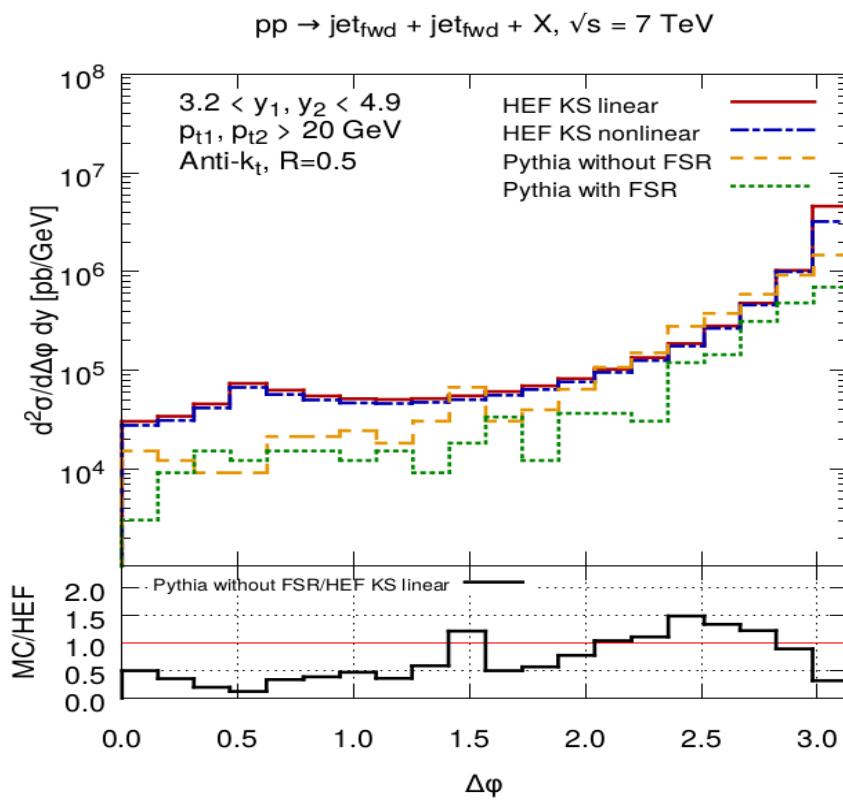
Other relevant effects – Final State Radiation



Wide angle soft emissions lower cross section for hard jets

Other relevant effects – final state parton shower with Pythia

Bury, Deak, Kutak, Sapeta '16



HEF framework with KShardsacle or DLC2016 → compatible with ISR in Pythia at moderate and large angles → importance of Sudakov resummation.