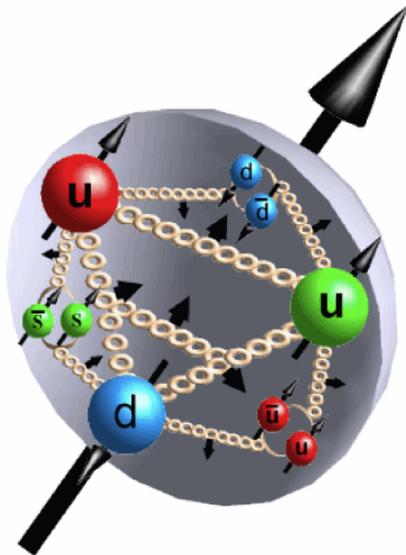


# Study of TMDs with SoLID at 12-GeV Jefferson Lab

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Kalyan Allada  
Massachusetts Institute of Technology

DIS Workshop 2017, University of Birmingham, UK  
4<sup>th</sup> April 2017

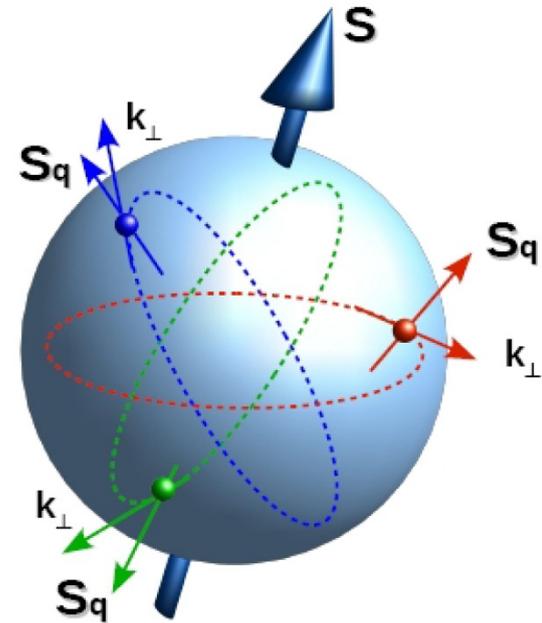


# Spin and Intrinsic Transverse Momentum

We cannot learn about the spin structure of the nucleon without *rigorously* taking into account the transverse motion of the partons inside it !

Distribution and fragmentation functions depends on:

- lightcone momentum fraction ( $x$  for the distributions and  $z$  for the fragmentation)
- $Q^2$  (pQCD evolution)
- intrinsic transverse momentum of partons ( $\mathbf{k}_T$  for the distributions and  $\mathbf{p}_T$  for the fragmentation)



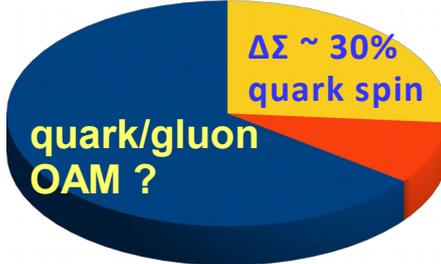
## Open questions:

- How do TMD's depend on the intrinsic transverse momentum ?
  - Gaussian behaviour in the central region ...
  - Power law decrease at large transverse momentum...
- Does the partonic intrinsic transverse momentum  $\mathbf{k}_T(\mathbf{p}_T)$  depend on  $x$  ( $z$ ) ?

# Nucleon Spin Puzzle

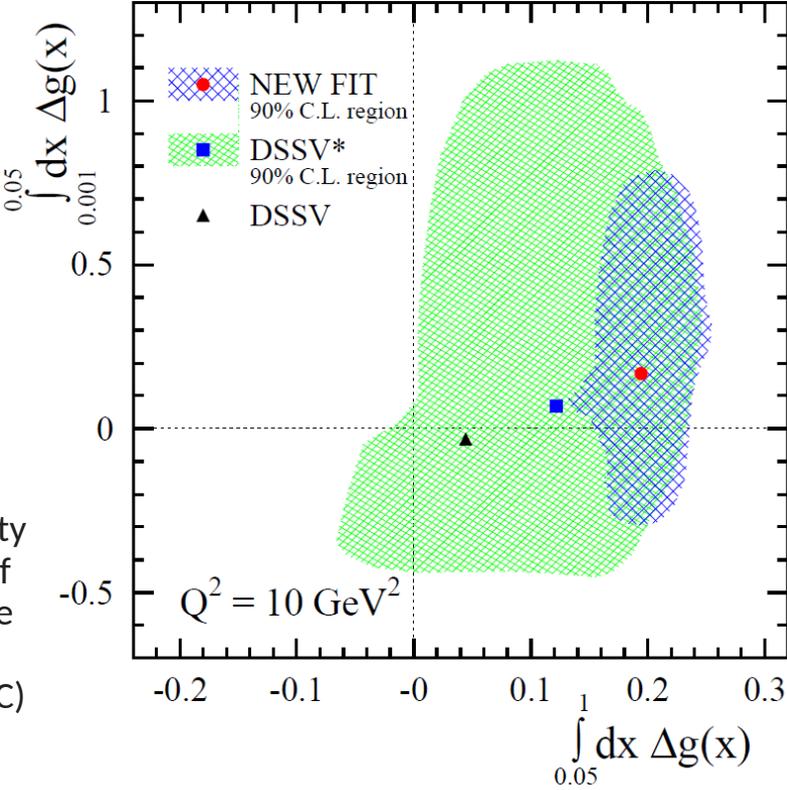
Nucleon spin:  $\frac{1}{2} = \frac{1}{2} \Delta \Sigma + \Delta G + L_q + L_g$

$L_q + L_g$   
orbital angular momentum



$\Delta G \sim$  not small?

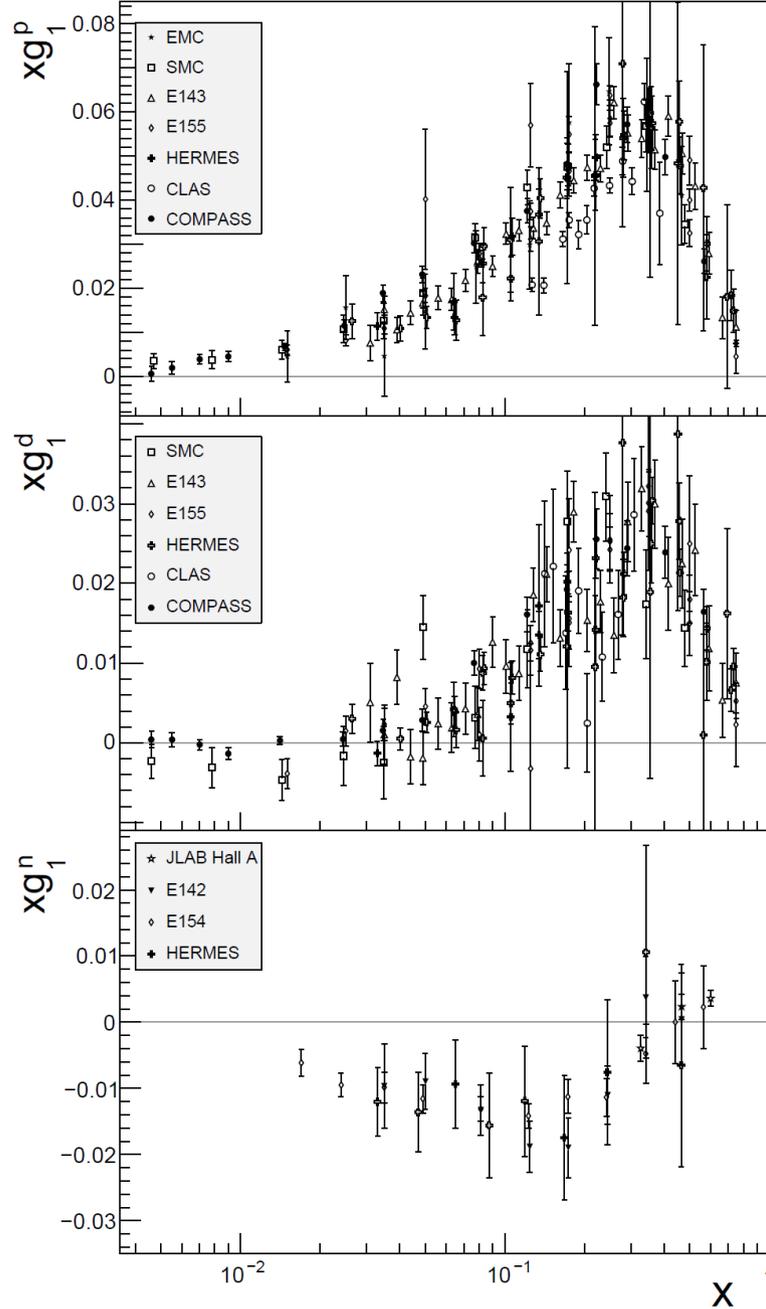
## Gluon polarization



deFlorian, et al.,  
PRL 113, 012001 (2014)

Global analysis of helicity PDFs - clear evidence of gluon polarization in the intermediate momentum scales (RHIC)

## Polarized DIS data



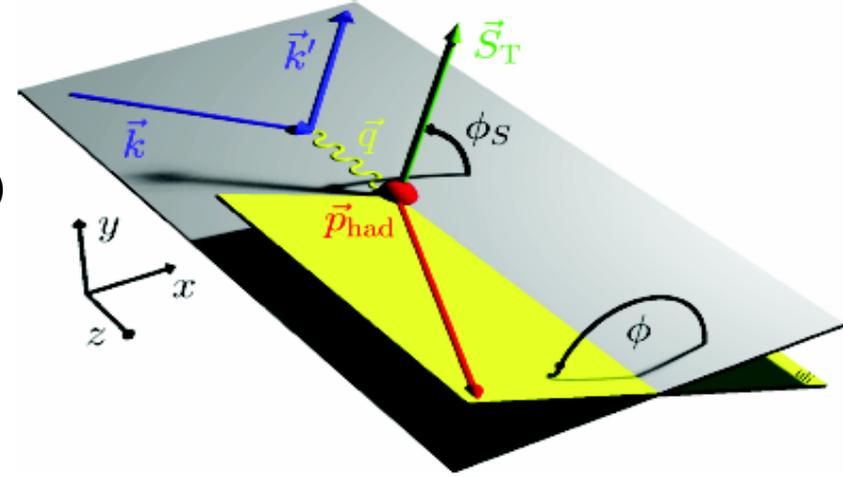
# Structure functions from SIDIS

SIDIS differential cross section

18 structure functions  $F(x, z, Q^2, P_T)$ ,

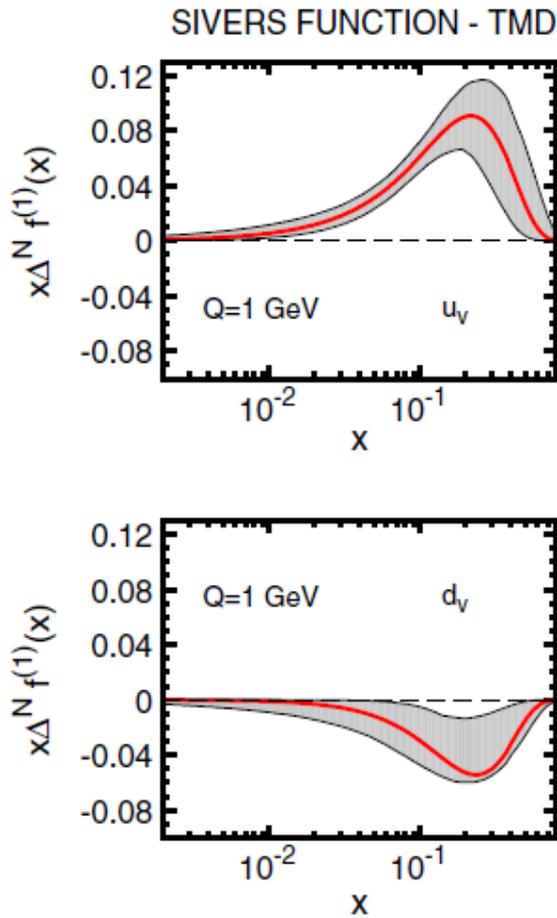
model independent (one photon exchange approximation)

$$\begin{aligned}
 & \frac{d\sigma}{dx dy dz dP_T^2 d\phi_h d\phi_S} \\
 &= \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left( 1 + \frac{\gamma^2}{2x} \right) \\
 & \times \left\{ F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h + \lambda_e \sqrt{2\epsilon(1-\epsilon)} F_{LU}^{\sin\phi_h} \sin\phi_h \right. \\
 & + S_L \left[ \sqrt{2\epsilon(1+\epsilon)} F_{UL}^{\sin\phi_h} \sin\phi_h + \epsilon F_{UL}^{\sin 2\phi_h} \sin 2\phi_h \right] + \lambda_e S_L \left[ \sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} F_{LL}^{\cos\phi_h} \cos\phi_h \right] \\
 & + S_T \left[ (F_{UT,T}^{\sin(\phi_h-\phi_S)} + \epsilon F_{UT,L}^{\sin(\phi_h-\phi_S)}) \sin(\phi_h - \phi_S) + \epsilon F_{UT}^{\sin(\phi_h+\phi_S)} \sin(\phi_h + \phi_S) + \epsilon F_{UT}^{\sin(3\phi_h-\phi_S)} \sin(3\phi_h - \phi_S) \right. \\
 & \quad \left. + \sqrt{2\epsilon(1+\epsilon)} F_{UT}^{\sin\phi_S} \sin\phi_S + \sqrt{2\epsilon(1+\epsilon)} F_{UT}^{\sin(2\phi_h-\phi_S)} \sin(2\phi_h - \phi_S) \right] \\
 & + \lambda_e S_T \left[ \sqrt{1-\epsilon^2} F_{LT}^{\cos(\phi_h-\phi_S)} \cos(\phi_h - \phi_S) \right. \\
 & \quad \left. + \sqrt{2\epsilon(1-\epsilon)} F_{LT}^{\cos\phi_S} \cos\phi_S + \sqrt{2\epsilon(1-\epsilon)} F_{LT}^{\cos(2\phi_h-\phi_S)} \cos(2\phi_h - \phi_S) \right] \left. \right\}
 \end{aligned}$$

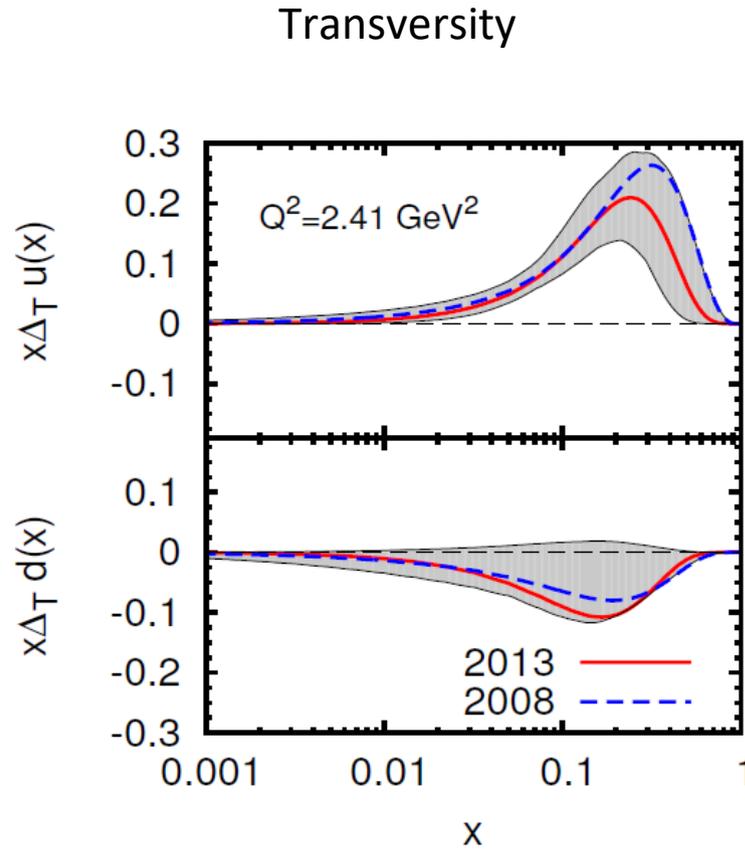


In parton model,  $F(x, z, Q^2, P_T)$ s are expressed as the convolution of TMDs.

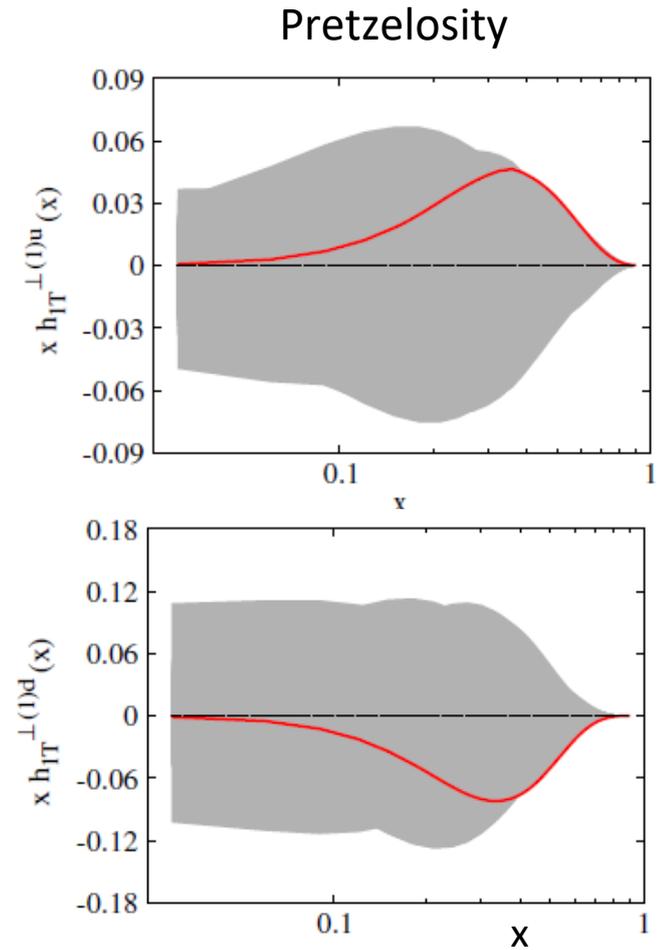
# Present Status of TMD Extractions



Anselmino et al.,  
PRD 86, 014028 (2012)



Anselmino et al.,  
Phys. Rev. D 87, 094019



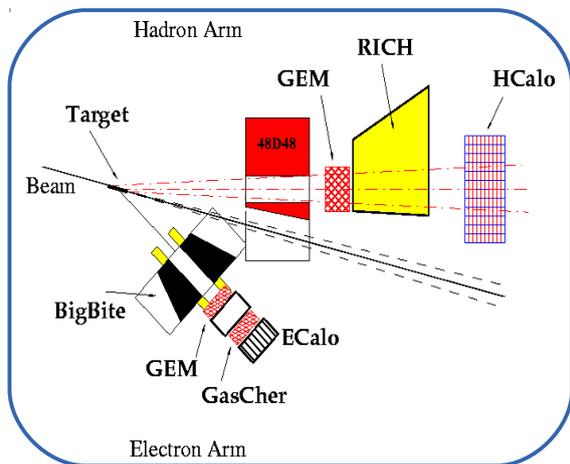
C. Lefky, A. Prokudin  
PRD 91, 034010 (2015)

Transverse SSA data available from HERMES, COMPASS and JLab Hall A

# Multi-Hall SIDIS Program at JLab

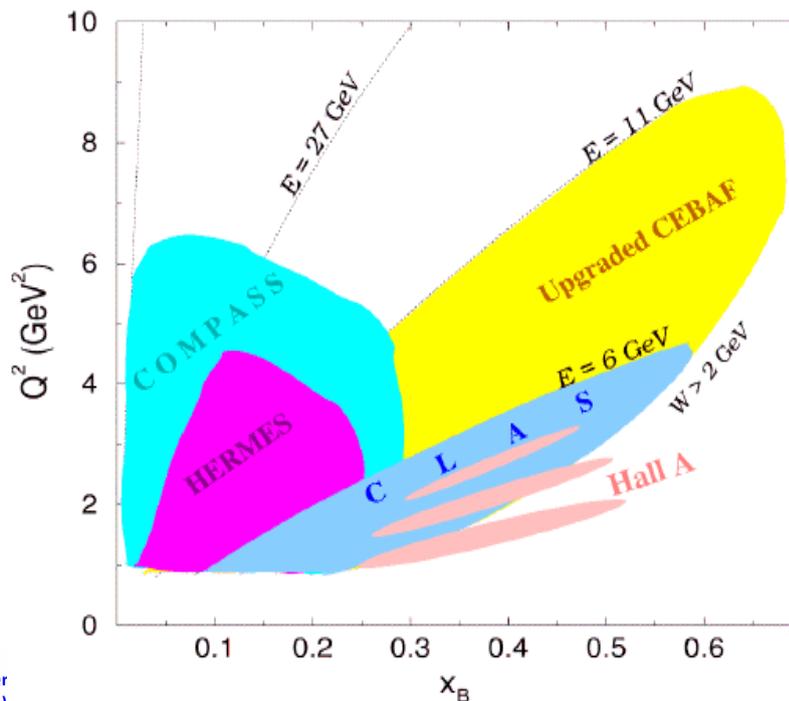
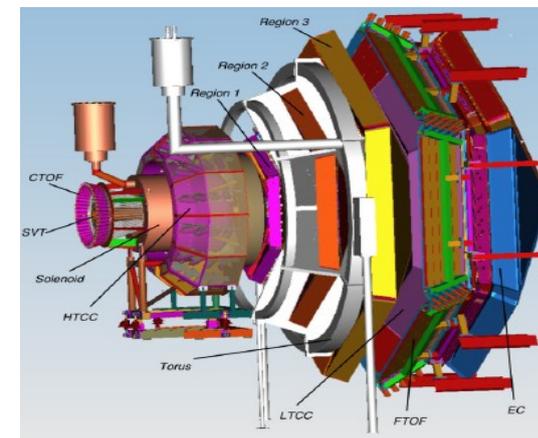
## Hall A: Super BigBite

SIDIS with  $^3\text{He}$ , high  $x$ ,  $Q^2$



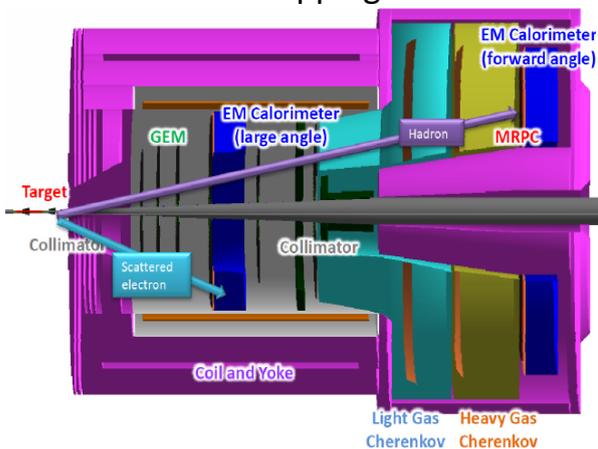
## Hall B: CLAS12

SIDIS with polarized H/D  
Comprehensive SIDIS program



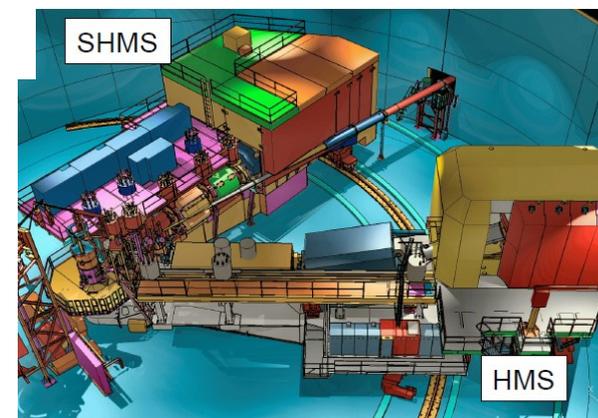
## Hall A: SOLID

SIDIS with polarized  $^3\text{He}/\text{NH}_3$   
Precision 4D mapping



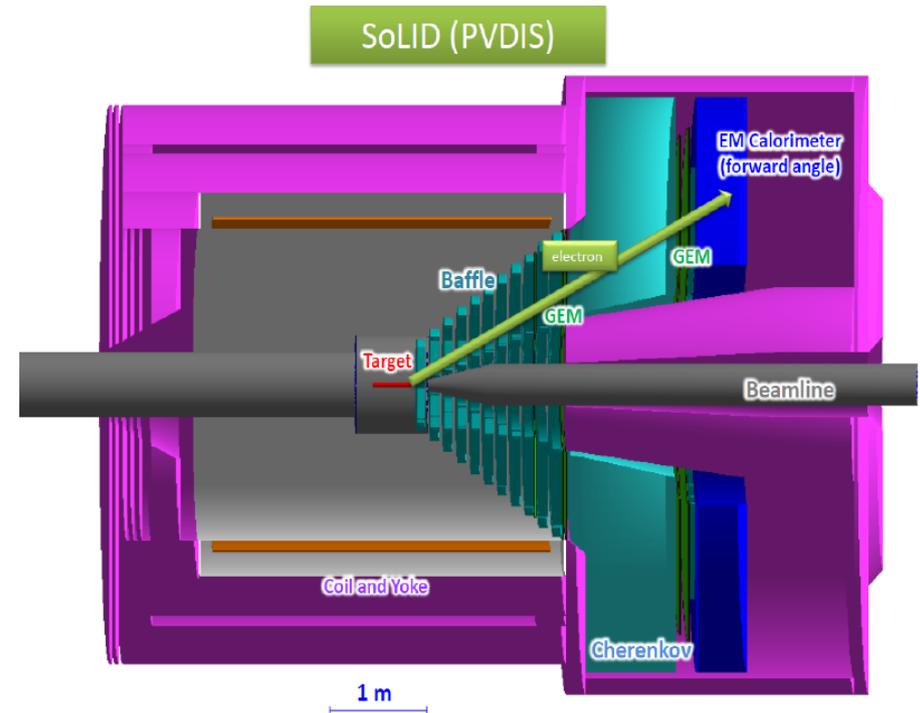
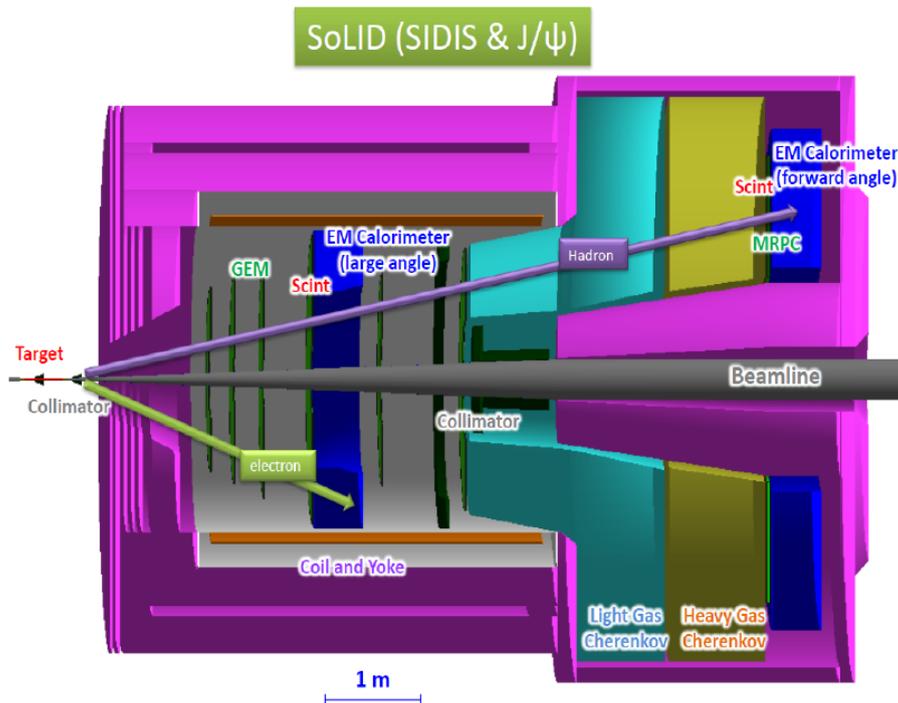
## Hall C: SHMS

SIDIS with unpolarized H/D



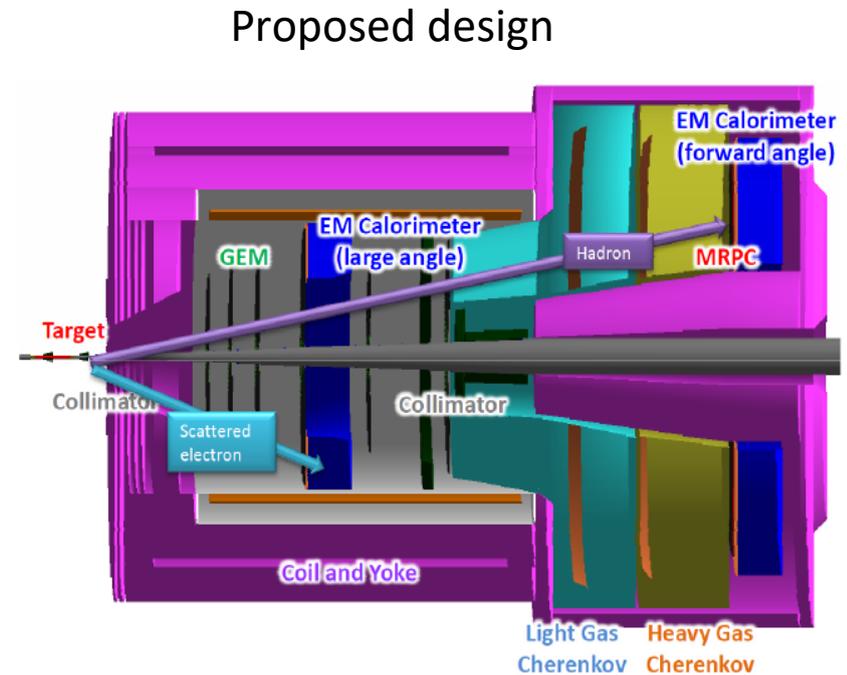
# Overview of SoLID

- Proposed to take full advantage of 12 GeV upgrade at JLab
  - A **large acceptance** detector which can handle **high luminosity** ( $10^{37}$  -  $10^{39}$   $\text{cm}^{-2}\text{s}^{-1}$ )
  - Modified CLEO-II magnet, modern detectors, latest data-acquisition technologies
- Five highly rated approved experiments:
  - Three SIDIS experiments, one PVDIS and one  $J/\psi$  production
- Large collaboration (250+ collaborators for 70+ institutes in 13 countries)
  - Major collaboration with Chinese institutes



# SoLID Spectrometer for SIDIS

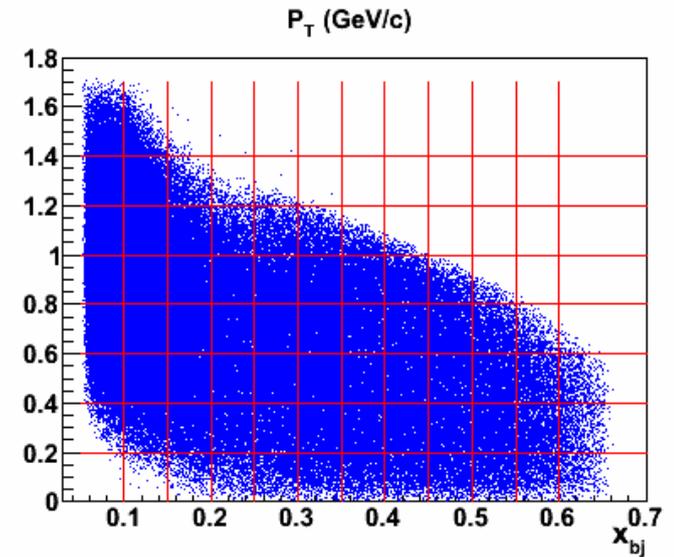
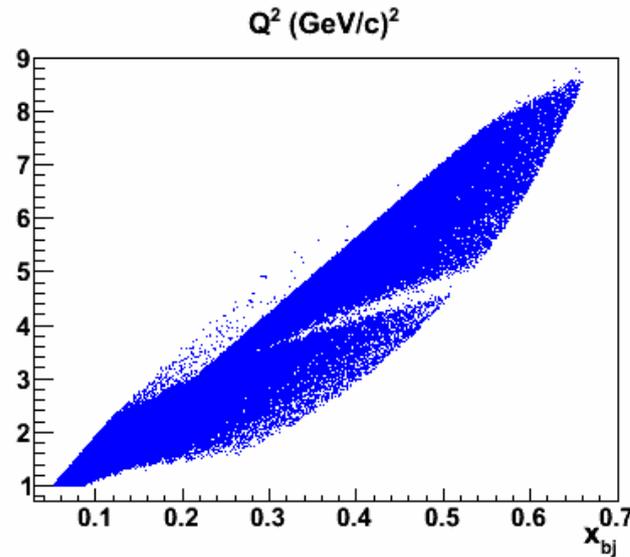
- Longitudinal and transverse polarized targets
  - proton ( $\text{NH}_3$ ) and neutron ( $^3\text{He}$ ) targets
- Three approved Semi-inclusive DIS proposals:
  - SSA/DSA on longitudinal neutron ( $^3\text{He}$ ) - E12-11-007
  - SSA on transverse neutron ( $^3\text{He}$ ) - E12-10-006
  - SSA on transversely proton ( $\text{NH}_3$ ) - E12-11-108
- Precision 4-D ( $x, Q^2, p_T, z$ ) mapping of SSA/DSA
  - high luminosity ( $10^{36} \text{cm}^{-2} \text{s}^{-1}$ ) and larger acceptance
- Tensor charge, TMDs (Collins, Sivers, Pretzelosity...)
- Two run group experiments:
  - SSA in dihadron production and  $A_y$  (two photon-exchange)



# Kinematic Coverage with SoLID

Wide phase space coverage – essential for 4D binning of SSA/DSA

- $x_B = 0.05 - 0.68$
- $Q^2 = 1.0 - 9.0 \text{ GeV}^2$
- $P_T = 0 - 1.5 \text{ GeV}^2$
- $Z = 0.3 - 0.7$
- $W > 2.3 \text{ GeV}$

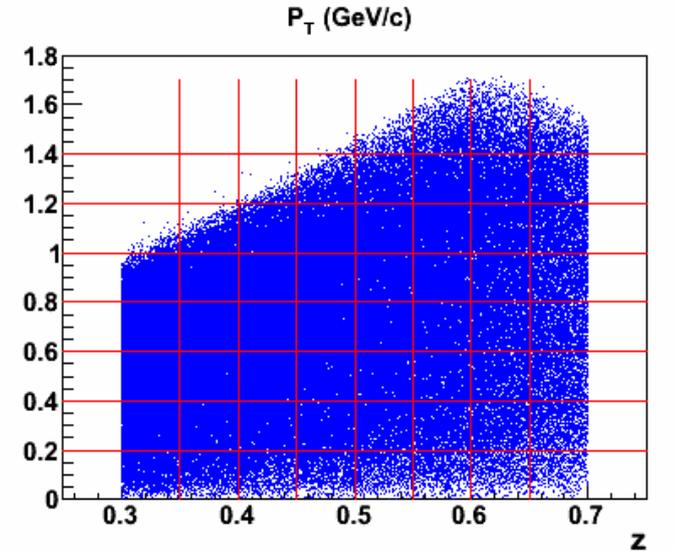
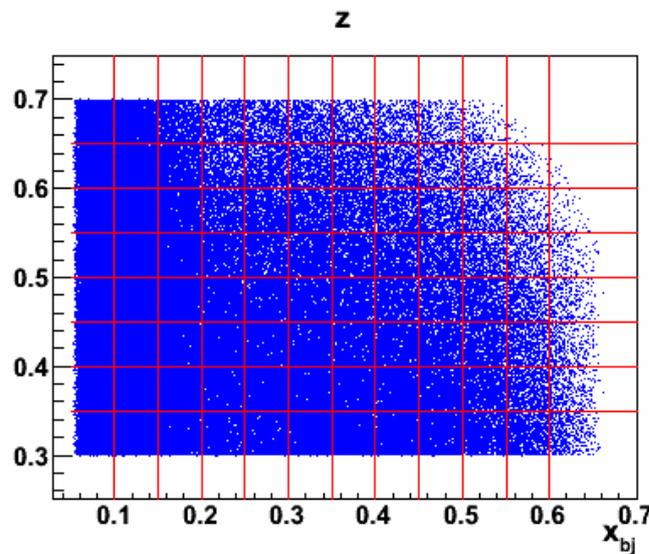


**Forward acceptance:**

- $\phi = 2\pi$
- $\theta = 8^\circ - 14^\circ$
- $P : 1.0 - 7.0 \text{ GeV/c}$

**Large angle acceptance:**

- $\phi = 2\pi$
- $\theta = 16^\circ - 24^\circ$
- $P : 3.5 - 7.0 \text{ GeV/c}$

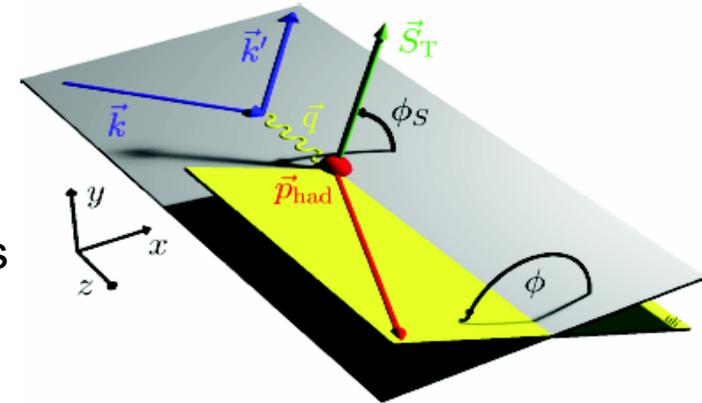


# Target SSA Measurement using SoLID

**Target SSA (used in 6 GeV experiment):**

$$A_{UT}^h(\phi_h, \phi_S) = \frac{1}{(p_T)} \cdot \frac{N_1(\phi_h, \phi_S) - N_2(\phi_h, \phi_S + \pi)}{N_1(\phi_h, \phi_S) + N_2(\phi_h, \phi_S + \pi)}$$

- Relative luminosity monitored using single particle triggers
- Target spin flip (every 10-20 mins for  $^3\text{He}$ )



**Target SSA definition for SoLID experiments:**

$$A_{UT}^h(\phi_h, \phi_S) = \frac{2}{P_T^1 + P_T^2} \cdot \frac{\sqrt{N_1(\phi_h, \phi_S)N_2(\phi_h, \phi_S + \pi)} - \sqrt{N_1(\phi_h, \phi_S + \pi)N_2(\phi_h, \phi_S)}}{\sqrt{N_1(\phi_h, \phi_S)N_2(\phi_h, \phi_S + \pi)} + \sqrt{N_1(\phi_h, \phi_S + \pi)N_2(\phi_h, \phi_S)}}$$

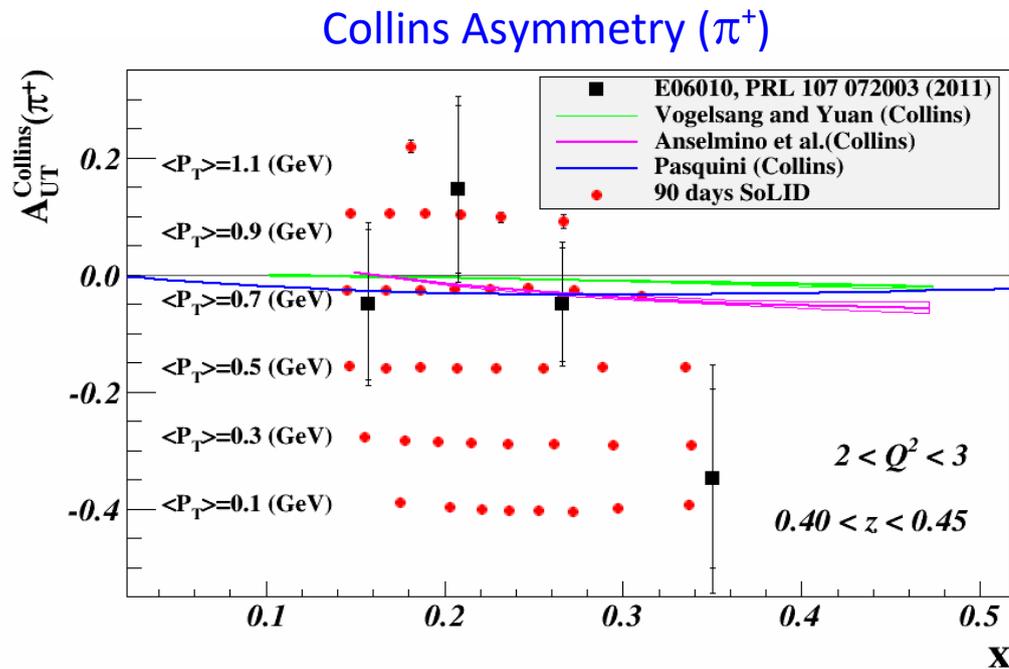
- $N_1(\phi_h, \phi_S)$  and  $N_1(\phi_h, \phi_S + \pi)$  taken at the same time
- Luminosity at different times cancels
- In first order, acceptance and detection efficiency cancels
- **double cancellation due to spin-flip and full azimuthal coverage**

**Unique advantage in reducing the systematics associated with detection efficiency, luminosity and acceptance effects**

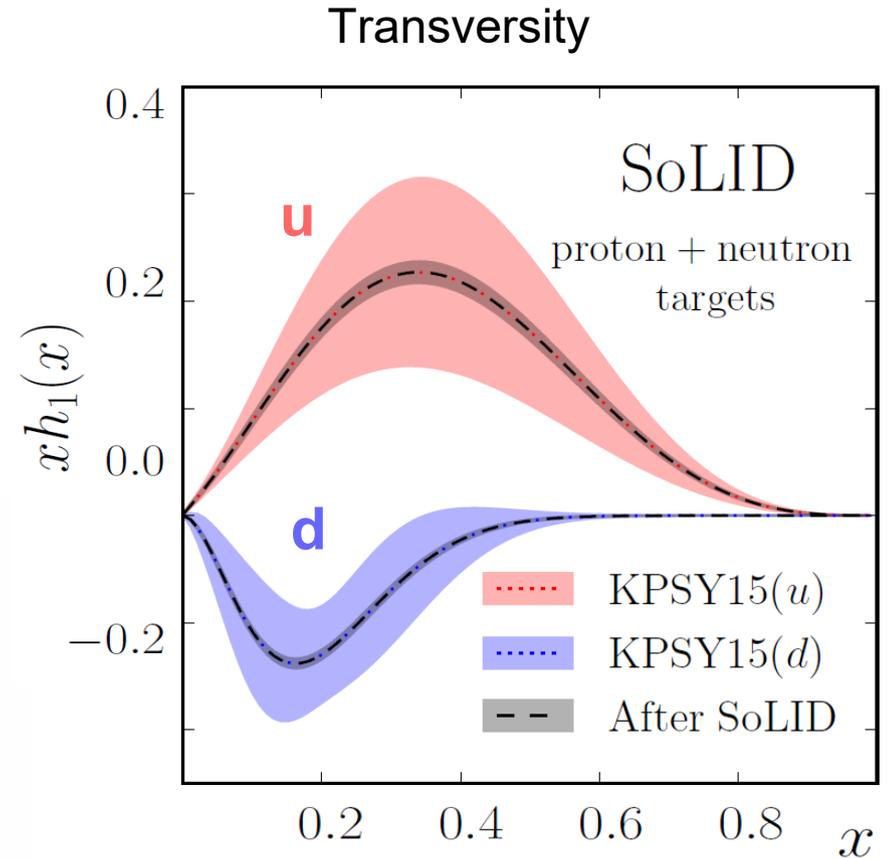
# Expected Impact on Transversity

$$\sigma_{UT}^{SIDIS} \propto \sin(\phi_h + \phi_S) h_1 \otimes H_1^\perp$$

- Global fit to SIDIS and e+e- experimental data
- Covers wide region in  $x$  – important for transversity and tensor charge extraction
- Relatively large  $P_T$  range – important for testing TMD factorization approach



$P_T$  vs.  $x$  for one ( $Q^2, z$ ) bin  
 Total > 1400 data points

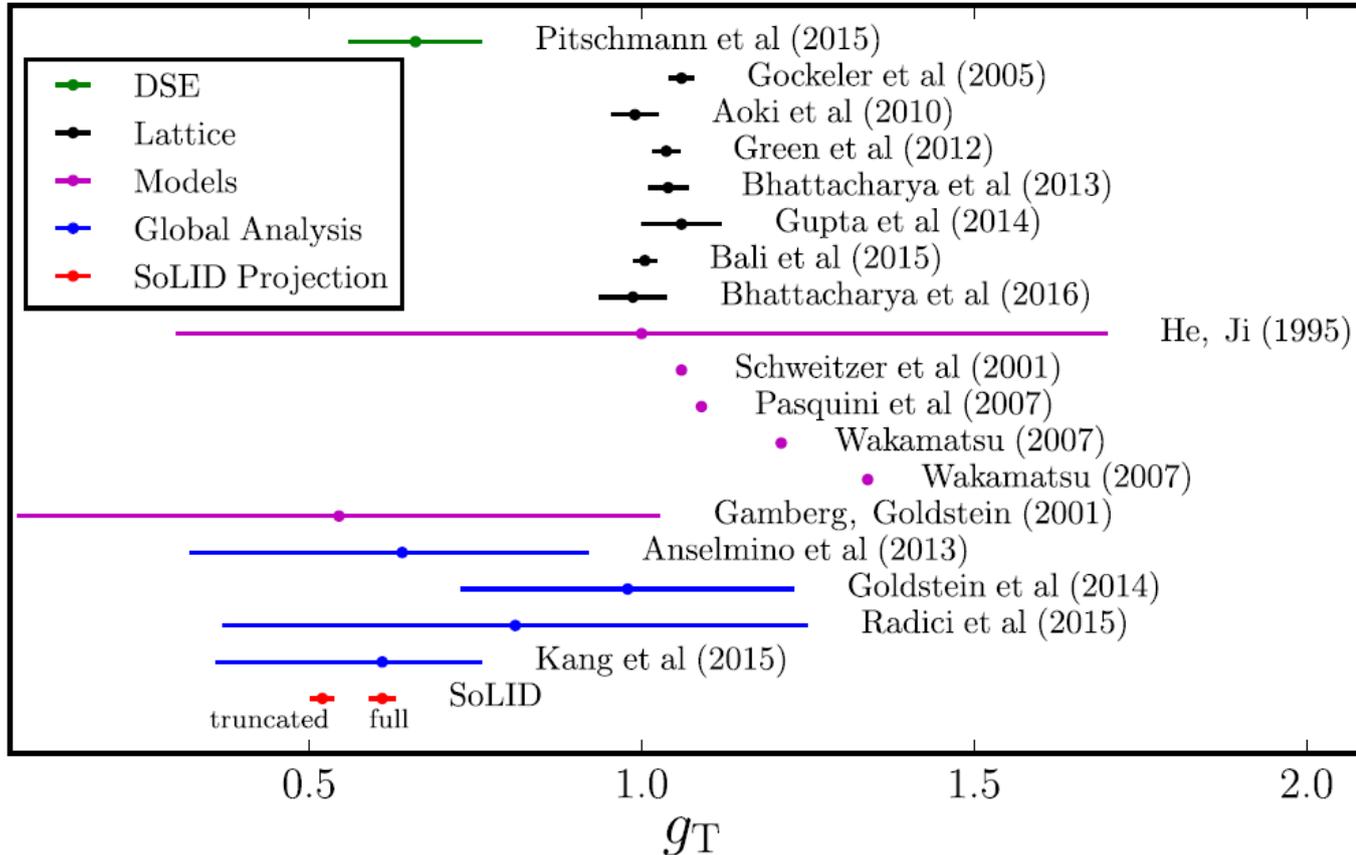


Z. Ye, N. Sato et al.,  
 Phys. Lett. B 767C (2017)

# Tensor Charge from SoLID

Isovector nucleon tensor charge  $g_T = \delta u - \delta d$ ,

$$\delta q = \int_0^1 [h_1^q(x) - h_1^{\bar{q}}(x)] dx$$



Z. Ye, N. Sato et al.,  
Phys. Lett. B 767C (2017)

- truncated range of  $x$ : 0.05 – 0.6
- assumption of a negligible contribution from sea quarks
- model dependent assumptions on the shape of underlying TMD distributions
- $Q^2$  evolution included

# Tensor Charge and Neutron EDM

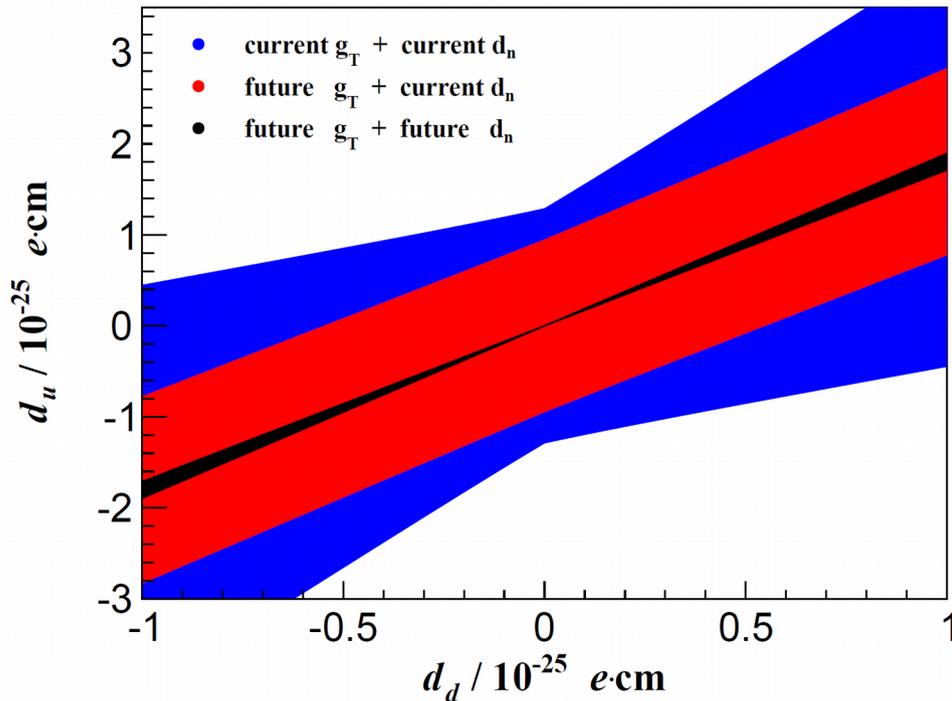
quark EDM ( $d_u, d_d$ ) contribution to the nucleon EDM ( $d_p, d_n$ )

$$d_p = g_T^u d_u + g_T^d d_d,$$

$$d_n = g_T^d d_u + g_T^u d_d,$$

Bounds on  $d_u$  and  $d_d$ , assuming negligible strange quark contribution ( $\delta_T^s = 0$ )

Constraints on quark EDMs with the upper limits on nucleon EDMs and the tensor charge extraction



■ Kang et al., with  $|d_n| < 2.1 \times 10^{-26}$  e.cm  
■ SoLID with  $|d_n| < 2.1 \times 10^{-26}$  e.cm  
■ SoLID with  $|d_n| < 2.1 \times 10^{-28}$  e.cm

current knowledge of tensor charge comes from lattice QCD

SoLID can provide precision comparable to lattice QCD

Current neutron EDM limit  $|d_n| < 2.1 \times 10^{-26}$  e.cm

H. Gao, T. Liu et al.,  
(arXiv:1704.00113)

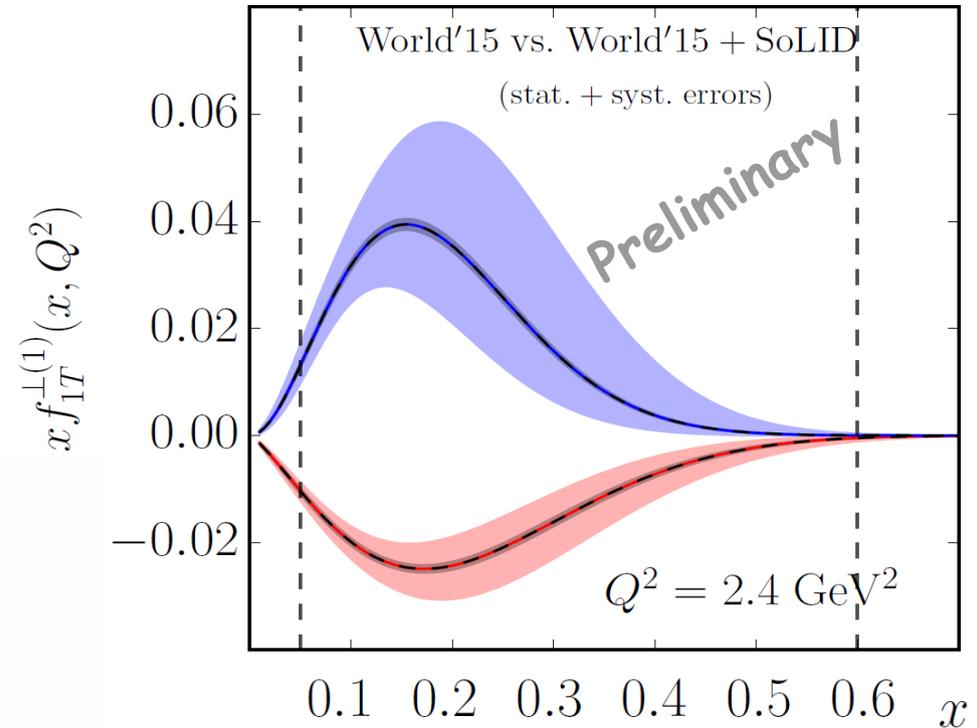
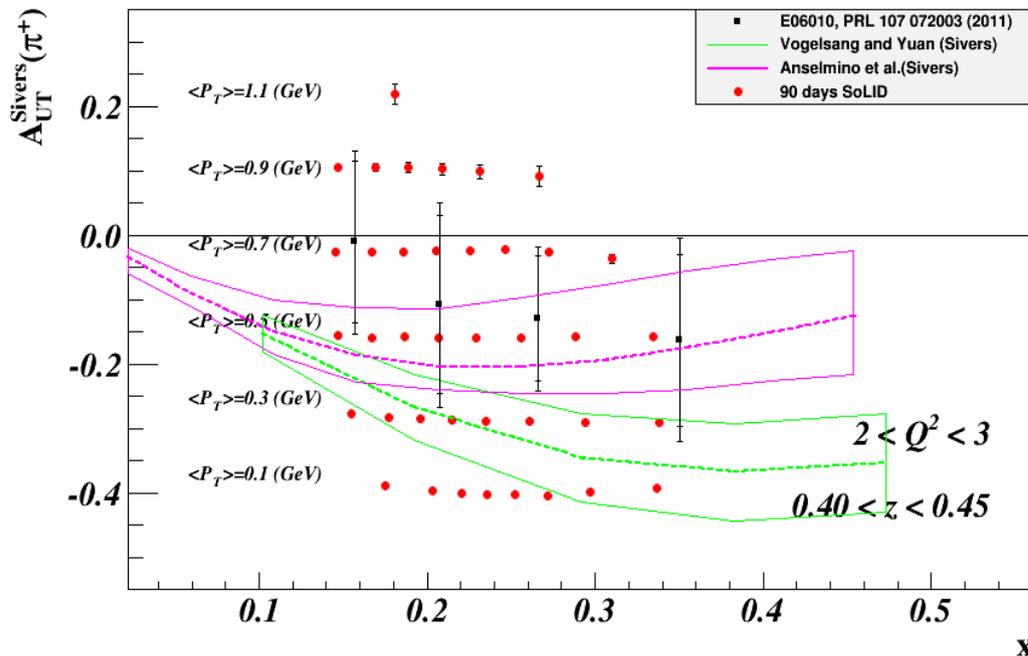
# Sivers Asymmetry and Impact of SoLID

$$f_{q/p\uparrow}(x, \mathbf{k}_\perp) = f_1^q(x, k_\perp) - f_{1T}^\perp(x, k_\perp) \frac{\hat{\mathbf{P}} \times \mathbf{k}_\perp \cdot \mathbf{S}}{M}$$

## Sivers Asymmetry in SIDIS

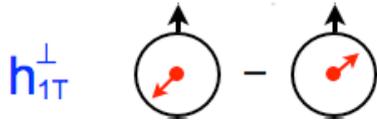
$$A_{UT}^{\sin(\phi_h - \phi_S)} \sim f_{1T}^\perp(x, k_\perp) \otimes D_1(z, p_\perp)$$

### Sivers Asymmetry ( $\pi^+$ )



N. Sato, T. Liu, A. Prokudin et al.,  
(In progress)

# Pretzelosity

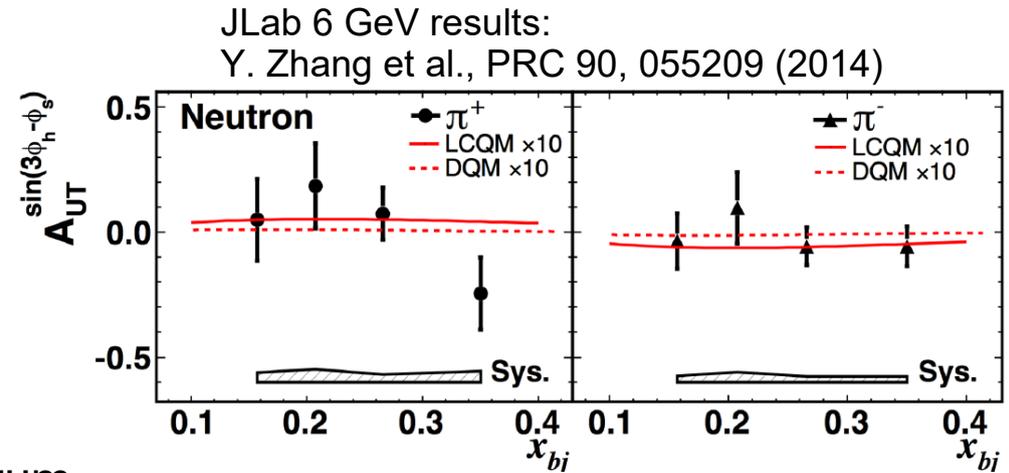


$$\sigma_{UT}^{SIDIS} \propto \sin(3\phi_h - \phi_s) h_{1T}^\perp \otimes H_1^\perp$$

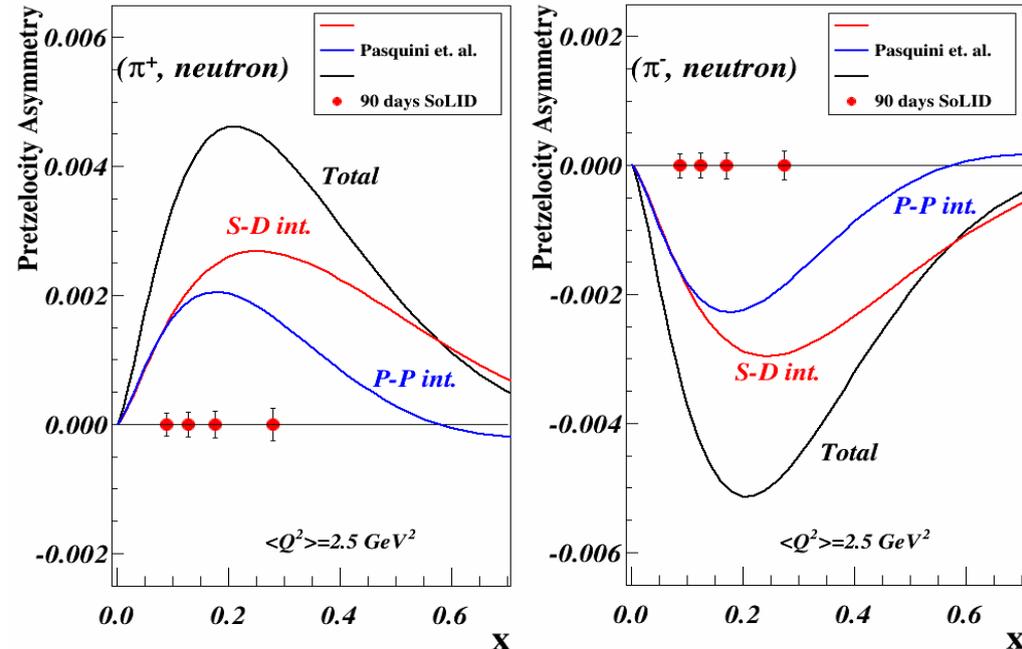
- Correlation between quark transverse momentum and transverse polarization of the nucleon
- Chiral odd, no gluon analogy
- Direct probe of relativistic effects
- Access to quark OAM through models

$$L_z^q = - \int dx d^2\mathbf{k}_\perp \frac{\mathbf{k}_\perp^2}{2M^2} h_{1T}^{\perp q}(x, \mathbf{k}_\perp) = - \int dx h_{1T}^{\perp(1)q}(x)$$

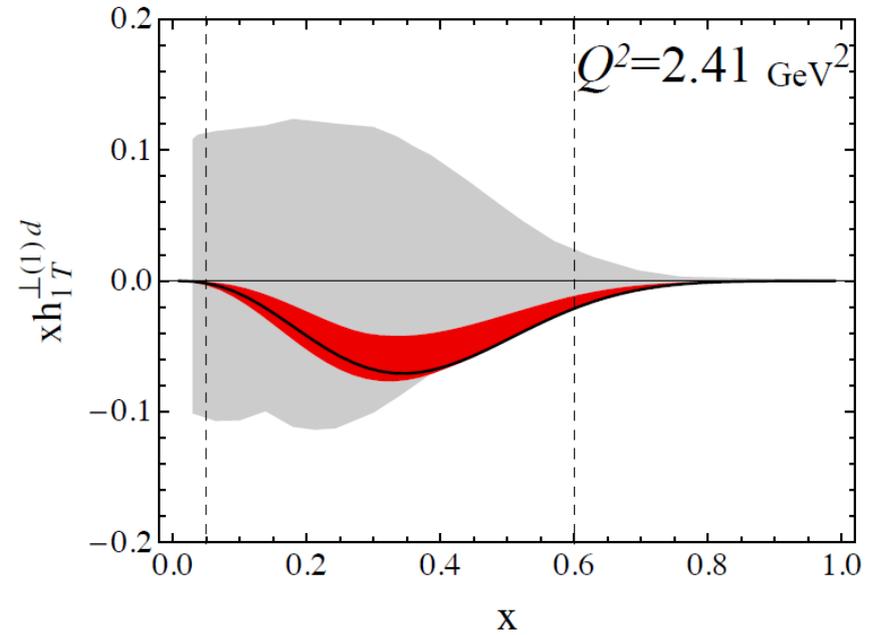
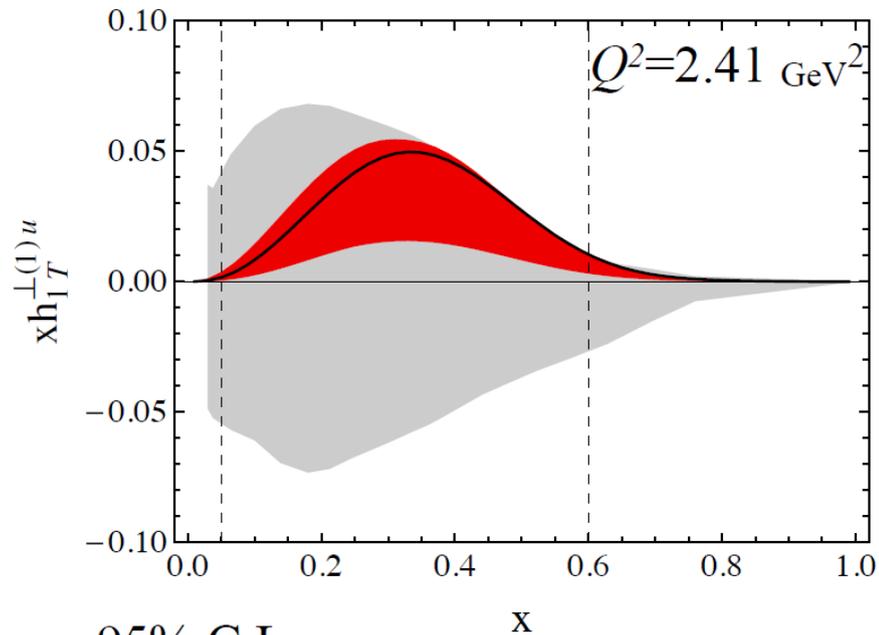
(model dependent)



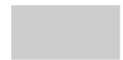
## SoLID SIDIS projections



# Impact of SoLID on Pretzelosity



95% C.L.



parametrization by C. Lefky *et al.*, PR D 91, 034010 (2015).



SoLID projection with transversely polarized neutron and proton data.

C. Lefky, A. Prokudin PRD 91, 034010 (2015)

- A global fit to 175 data from COMPASS, HERMES, and JLab

Tianbo Liu (Duke)

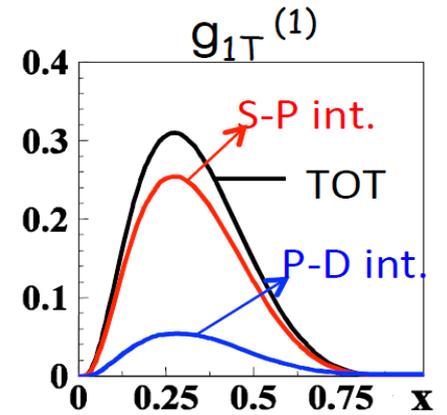
# Summary

- TMDs provide very rich information of parton dynamics in the nucleon beyond the collinear picture
- High precision data in wide kinematic region is essential to study TMDs in SIDIS:
  - Clean extraction of TMDs,  $Q^2$  evolution,  $P_T$  dependence, higher twist effects, current/target fragmentation regions, etc...
- Plans for high precision SIDIS measurements in Hall-A at 11 GeV using SoLID
  - Three fully approved SIDIS experiments
  - Transverse SSA on proton and neutron
- SoLID SIDIS data combined with high precision cross-section data from Hall C HMS+SHMS will allow us to perform a rigorous study of TMDs

# Spare Slides

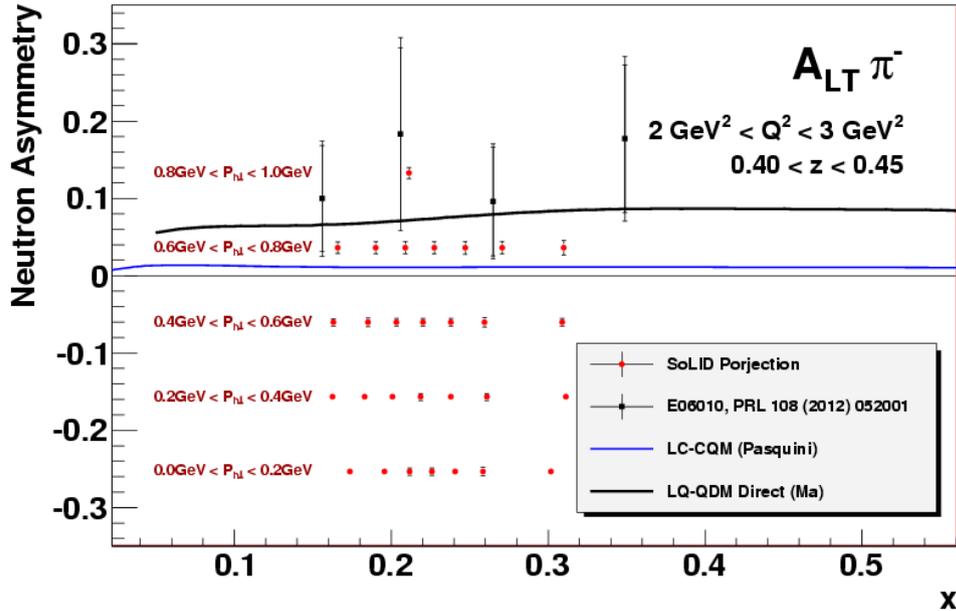
# Worm-Gear Functions

- G1T and h1L
- Dominated by real part of interference between L=0 (S) and L=1 (P) states



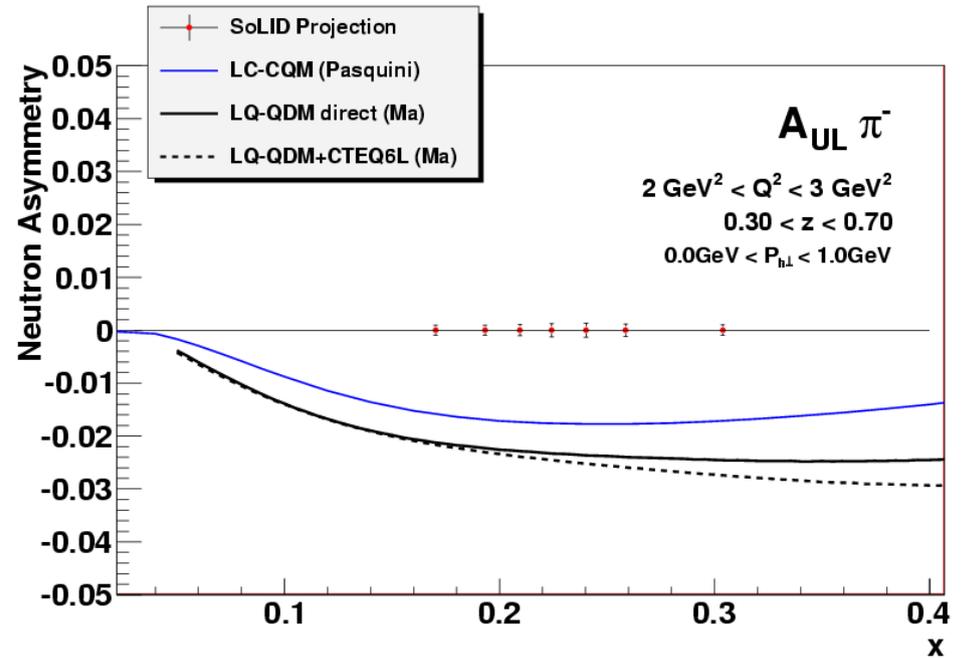
Light-Cone CQM by B. Pasquini  
B.P., Cazzaniga, Boffi, PRD78, 2008

## SoLID projections from neutron



$$A_{LT} \sim g_{1T}(x)D_1(z)$$

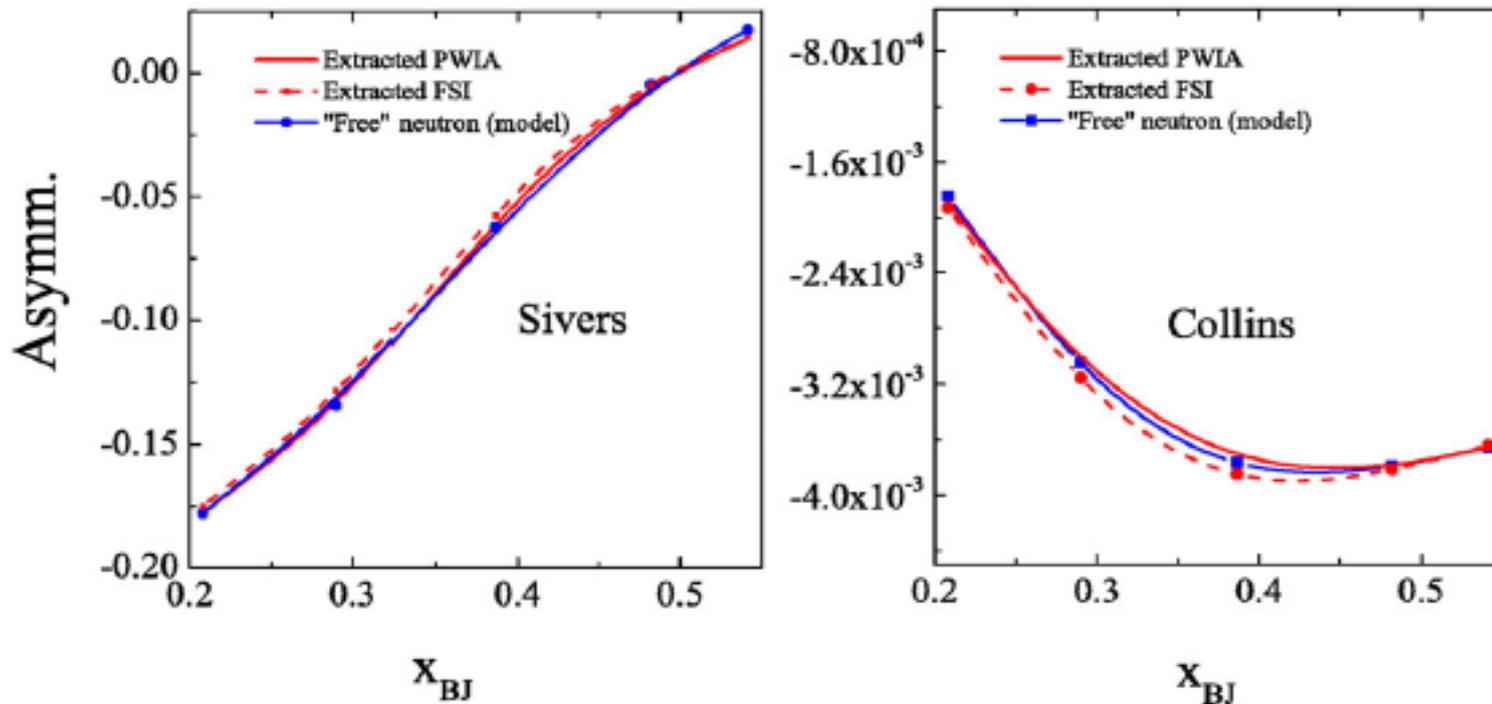
## SoLID projections from neutron



$$A_{UL} \sim h_{1L}^\perp(x) \otimes H_1^\perp(z)$$

# Extraction of neutron SSA from $^3\text{He}$ – effect of FSI

$$A_n \simeq \frac{1}{p_n f_n} (A_3^{\text{exp}} - 2p_p f_p A_p^{\text{exp}})$$



**Fig. 13.** Check of the extraction procedure, eq. (17), with and without FSI taken into account, for the Siverts (left) and Collins (right) SSAs, in the kinematics of [65].

Raphael Dupre and Sergio Scopetta  
 Eur. Phys. J. A (2016) 52: 159

# Hall C E00-108 Experiment: SIDIS Cross Section



$$\frac{\frac{d\sigma}{d\Omega_e dE_e' dz dp_T^2 d\phi}}{\frac{d\sigma}{d\Omega_e dE_e'}} = \frac{dN}{dz} b e^{-b p_T^2} \frac{1 + A \cos(\phi) + B \cos(2\phi)}{2\pi},$$

$$\frac{dN}{dz} \sim \sum_q e_q^2 q(x, Q^2) D_{q \rightarrow \pi}(z, Q^2)$$

Low energy semi-inclusive cross sections consistent with calculation using high energy parameters of frag. functions and CTEQ PDFs (for  $z < 0.7$  and  $M_x^2 > \sim 2.5$ )

