Top quark pair property measurements and $t\bar{t}+X$, using the ATLAS detector at the LHC

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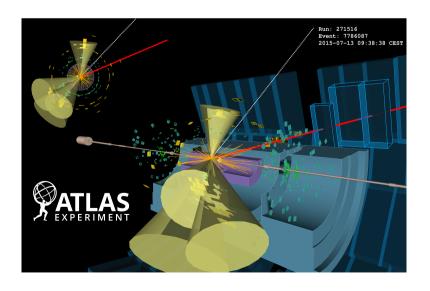




Top quark properties at ATLAS



- \sim 40 millions of $t\bar{t}$ events produced at the ATLAS detector
- Era of precision measurements for top quark physics
- Measurements of top quark properties as tests of Standard Model
- Possibility of probing new observables and extending existing analysis approaches



Publications since last DIS



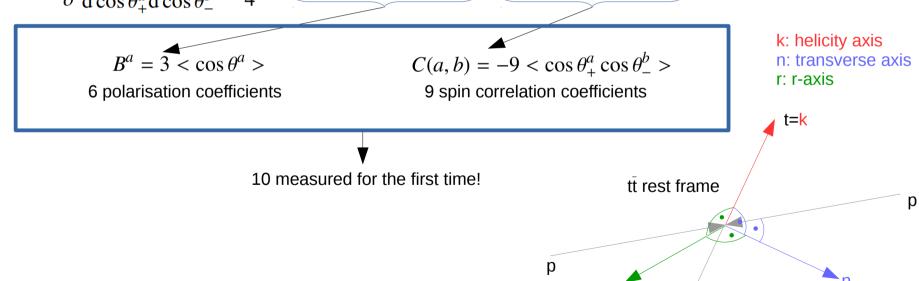
- 1) Measurements of **top quark spin observables** in *tt* events using **dilepton** final states in 8 TeV *pp* collisions with the ATLAS detector: [JHEP 03 (2017) 113]
- 2) Measurements of the **charge asymmetry** in top-quark pair production in the **dilepton** final state at 8 TeV with the ATLAS detector: [Phys. Rev. D 94, 032006]
- 3) Measurements of **charge and CP asymmetries in** *b***-hadron decays** using top-quark events collected by the ATLAS detector in *pp* collisions at 8 TeV: [JHEP02(2017)071]
- 4) Measurement of the **W boson polarisation** in *tt* events from pp collisions at 8 TeV in the **lepton+jets** channel with ATLAS: *submitted to EPJC* [arXiv:1612.02577]
- 5) Measurement of the *ttZ* and *ttW* production cross sections in multilepton final states using 3.2 fb⁻¹ of *pp* collisions at 13 TeV at the LHC: [Eur. Phys. J. C (2017) 77: 40.]

Top Quark Spin Observables (1)



- 8 TeV, 20.2 fb⁻¹, dilepton channel of $t\bar{t}$ decays
- Top quark spin is correlated, strength quantified by quantisation axis and production process
- Spin information is transferred to decay products
 - → use angular observables of decay products

$$\frac{1}{\sigma} \frac{\mathrm{d}^2 \sigma}{\mathrm{d} \cos \theta_+^a \mathrm{d} \cos \theta_-^b} = \frac{1}{4} (1 + B_+^a \cos \theta_+^a + B_-^b \cos \theta_-^b - C(a, b) \cos \theta_+^a \cos \theta_-^b)$$

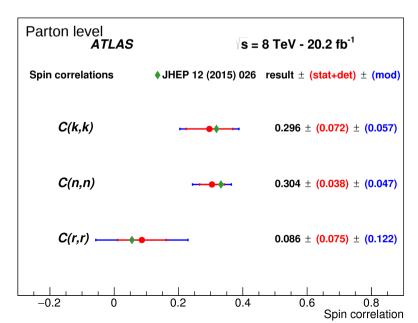


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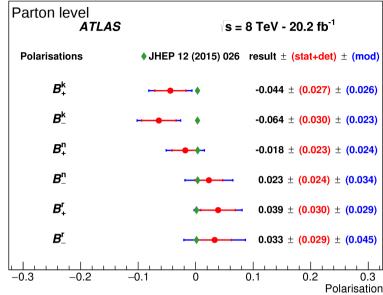
Top Quark Spin Observables (2)

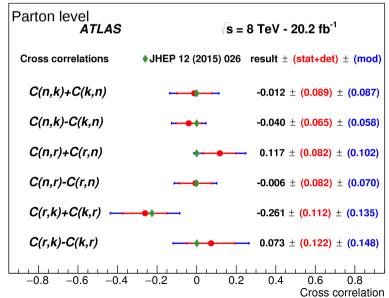


- 2 different analyses:
 - parton level (full phase space)
 - stable particle level (fiducial space)
- *tt* reconstruction using neutrino weighting technique
- Fully Bayesian unfolding to deal with distortions due to cuts and detector resolution
- No significant deviation from SM
- Observation of C(n,n) with 5.1 σ



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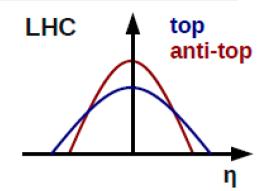




Charge Asymmetry in dilepton (1)



- 8 TeV, 20.3 fb⁻¹, dilepton channel of $t\bar{t}$ decays
- Asymmetries expected from valence quark sea antiquark fusion
 → antitop more central than top



Leptonic asymmetries:

$$A_{
m C}^{\ell\ell}=rac{N(\Delta|\eta|>0)-N(\Delta|\eta|<0)}{N(\Delta|\eta|>0)+N(\Delta|\eta|<0)}$$
 , $\Delta|\eta|=|\eta_{\ell^+}|-|\eta_{\ell^-}|$

tt asymmetries:

$$A_{
m C}^{tar t}=rac{N(\Delta|y|>0)-N(\Delta|y|<0)}{N(\Delta|y|>0)+N(\Delta|y|<0)}$$
 , $\Delta|y|=|y_t|-|y_{ar t}|$

3 different measurements of both observables:

- inclusive measurements on parton level in the full pase space
- inclusive measurements on particle level in the fiducial region
- differential measurements: inv. mass $(m_{t\bar{t}})$, $p_{T,t\bar{t}}$, and longitudinal boost $(\beta_{z,t\bar{t}})$ of $t\bar{t}$ system in the fiducial regions and the full phase space

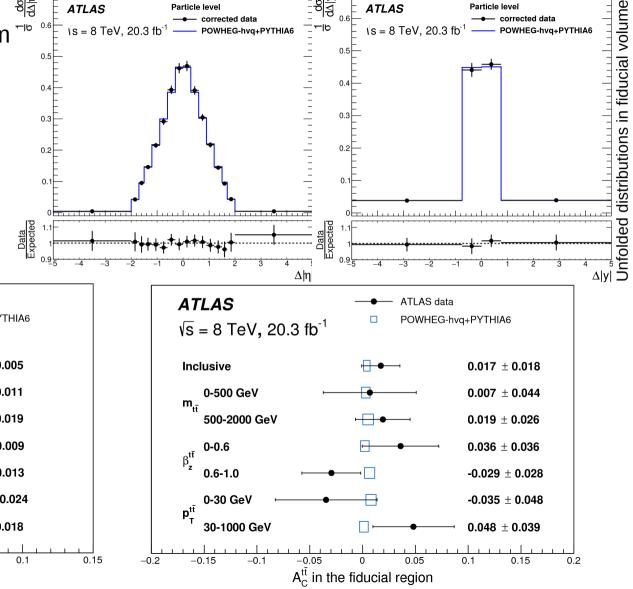
Charge Asymmetry in dilepton (2)

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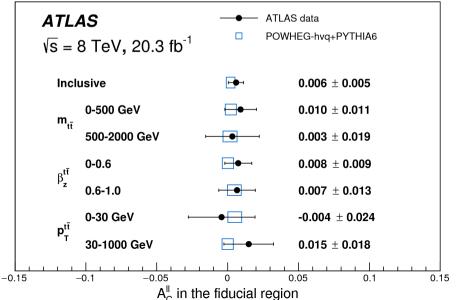
ATLAS



- Kinematic reconstruction of $t\bar{t}$ system
- Fully Bayesean unfolding
- Result: $A_C^{\ell\ell} = 0.008 \pm 0.006$ $A_C^{t\bar{t}} = 0.021 \pm 0.016$
- SM prediction: $A_{\rm C}^{\ell\ell} = 0.0064 \pm 0.0003$ $A_C^{t\bar{t}} = 0.0111 \pm 0.0004$

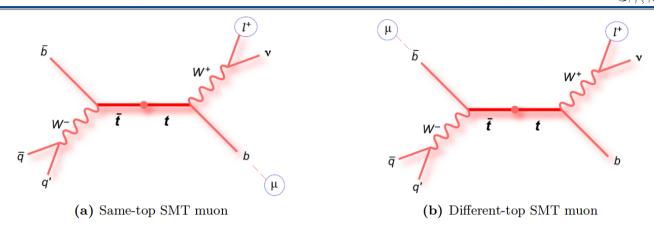


ATLAS



Charge and CP asymmetries in b-decay (1)





- 8 TeV, \(\ell\)+jets, \(b\) decaying semileptonically to a soft muon
- 5 CP asymmetries, 2 charge asymmetries
- charge of lepton from W determines the charge of the produced b-quark
- charge of soft lepton determines the charge of b-quark at decay

$$A^{\text{ss}} = \frac{P(b \to \ell^+) - P(b \to \ell^-)}{P(b \to \ell^+) + P(\overline{b} \to \ell^-)}$$
$$A^{\text{os}} = \frac{P(b \to \ell^-) - P(\overline{b} \to \ell^+)}{P(b \to \ell^-) + P(\overline{b} \to \ell^+)}$$

CP asymmetries: $B_q - \overline{B}_q$ mixing

$$A^{\text{ss}} = \frac{P(b \to \ell^{+}) - P(\bar{b} \to \ell^{-})}{P(b \to \ell^{+}) + P(\bar{b} \to \ell^{-})} \begin{vmatrix} A^{b\ell}_{\text{mix}} = \frac{\Gamma(b \to \bar{b} \to \ell^{+}X) - \Gamma(\bar{b} \to b \to \ell^{-}X)}{\Gamma(b \to \bar{b} \to \ell^{+}X) + \Gamma(\bar{b} \to b \to \ell^{-}X)} \end{vmatrix} A^{b\ell}_{\text{dir}} = \frac{\Gamma(b \to \ell^{-}X) - \Gamma(\bar{b} \to \ell^{+}X)}{\Gamma(b \to \ell^{-}X) + \Gamma(\bar{b} \to \ell^{-}X)} A^{\text{os}} = \frac{P(b \to \ell^{-}) - P(\bar{b} \to \ell^{+})}{P(b \to \ell^{-}) + P(\bar{b} \to \ell^{+})} A^{bc}_{\text{mix}} = \frac{\Gamma(b \to \bar{b} \to \bar{c}X) - \Gamma(\bar{b} \to b \to cX)}{\Gamma(b \to \bar{b} \to \bar{c}X) + \Gamma(\bar{b} \to b \to cX)} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) - \Gamma(\bar{c} \to \ell^{+}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{+}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) - \Gamma(\bar{c} \to \ell^{+}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{+}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) - \Gamma(\bar{c} \to \ell^{+}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{+}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) - \Gamma(\bar{c} \to \ell^{+}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{+}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) - \Gamma(\bar{c} \to \ell^{+}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{+}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) - \Gamma(\bar{c} \to \ell^{+}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{+}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) - \Gamma(\bar{c} \to \ell^{+}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{+}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) - \Gamma(\bar{c} \to \ell^{+}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{+}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) - \Gamma(\bar{c} \to \ell^{+}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{-}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{-}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{-}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L}) + \Gamma(\bar{c} \to \ell^{-}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L})} A^{c\ell}_{\text{dir}} = \frac{\Gamma(\bar{c} \to \ell^{-}X_{L})}{\Gamma(\bar{c} \to \ell^{-}X_{L})} A^{c\ell}_{\text{dir}}$$

CP asymmetries: direct CP violation

$$A^{\text{ss}} = \frac{P\left(b \to \ell^{+}\right) - P\left(\overline{b} \to \ell^{-}\right)}{P\left(b \to \ell^{+}\right) + P\left(\overline{b} \to \ell^{-}\right)} \quad A^{b\ell}_{\text{mix}} = \frac{\Gamma\left(b \to \overline{b} \to \ell^{+}X\right) - \Gamma\left(\overline{b} \to b \to \ell^{-}X\right)}{\Gamma\left(b \to \overline{b} \to \ell^{+}X\right) + \Gamma\left(\overline{b} \to b \to \ell^{-}X\right)} \quad A^{b\ell}_{\text{dir}} = \frac{\Gamma\left(b \to \ell^{-}X\right) - \Gamma\left(\overline{b} \to \ell^{+}X\right)}{\Gamma\left(b \to \ell^{-}X\right) + \Gamma\left(\overline{b} \to \ell^{+}X\right)} \quad A^{bc}_{\text{dir}} = \frac{\Gamma\left(b \to cX_{L}\right) - \Gamma\left(\overline{b} \to \overline{c}X_{L}\right)}{\Gamma\left(b \to cX_{L}\right) + \Gamma\left(\overline{b} \to \overline{c}X_{L}\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}\right) - P\left(\overline{b} \to \ell^{+}\right)}{P\left(b \to \ell^{-}\right) - P\left(\overline{b} \to \ell^{+}\right)} \quad A^{bc}_{\text{mix}} = \frac{P\left(b \to \ell^{-}X\right) - \Gamma\left(\overline{b} \to b \to cX\right)}{P\left(b \to \ell^{-}X\right) - P\left(\overline{b} \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}\right) - P\left(\overline{b} \to \ell^{+}X\right)}{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)}{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)}{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)}{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)}{P\left(b \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)}{P\left(b \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}X\right)}{P\left(b \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}X\right) - P\left(b \to \ell^{-}X\right)}{P\left(b \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{mix}} = \frac{P\left(b \to \ell^{-}X\right)}{P\left(b \to \ell^{-}X\right)} \quad A^{c\ell}_{\text{$$

Charge and CP asymmetries in b-decay (2)



- Soft muon heavy flavour tagging (SMT muons)
- Data is unfolded to well defined fiducial space
- CP result cannot disprove DØ deviation in dimuon asymmetry
 - → result both compatible with SM and DØ results

		Data	(10^{-2})	MC ((10^{-2})	Existing limits (2σ)	(10^{-2})	SM predic	ction (10^{-2})
	A^{ss}	-0.7	± 0.8	0.05	± 0.23	-		$< 10^{-2}$	[19]
	A^{os}	0.4	± 0.5	-0.03	$\pm \ 0.13$	_		$< 10^{-2}$	[19]
	$A_{ m mix}^b$	-2.5	± 2.8	0.2	± 0.7	< 0.1	[95]	$< 10^{-3}$	[96] [95]
	$A_{ m dir}^{b\ell}$	0.5	± 0.5	-0.03	$\pm \ 0.14$	< 1.2	[94]	$< 10^{-5}$	[19] [94]
	$A_{ m dir}^{c\ell}$	1.0	± 1.0	-0.06	$\pm \ 0.25$	< 6.0	[94]	$< 10^{-9}$	[19] [94]
	$A_{ m dir}^{bc}$	-1.0	± 1.1	0.07	± 0.29	-		$< 10^{-7}$	[97]

First direct measurements of direct CP violation in this context, improve existing limits of indirect measurements

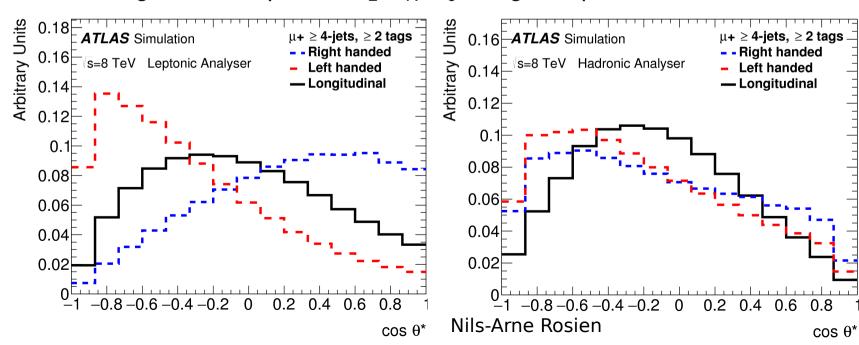
W boson polarisation (1)



- 8 TeV, lepton+jets channel, 20.2 fb⁻¹
- Most precise W boson polarisation measurement to date
- Use orientation of analyser wrt. the b quark (inv. direction) in the W rest frame of top decay

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta^*} = \frac{3}{4} \left(1 - \cos^2\theta^* \right) F_0 + \frac{3}{8} \left(1 - \cos\theta^* \right)^2 F_L + \frac{3}{8} \left(1 + \cos\theta^* \right)^2 F_R$$

- Take both charged lepton or down-type quark from W decay as analyser
- 3 reweighted $t\bar{t}$ samples for F_L , F_R , F_0 + bkg. samples



W boson polarisation (2)



- KLFitter tt reconstruction [Nucl. Instrum. Meth. A 748 (2014) 18–25]
- Perform template fit
- Best fits (=smallest uncertainty):
 - two channel combination (ejets+mujets, ≥2b) for leptonic analyser
 - four channel combination (ejets+mujets, =1b and ≥2b) for hadronic analyser

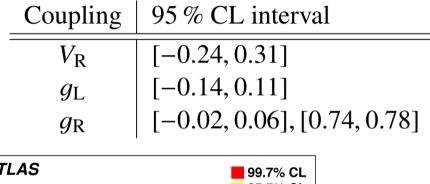
Leptonic analyser (≥2 <i>b</i> -tags)	Hadronic analyser (1 b -tag + \geq 2 b -tags)				
$F_0 = 0.709 \pm 0.012$ (stat.+bkg. norm.) $^{+0.015}_{-0.014}$ (syst.)	$F_0 = 0.659 \pm 0.010$ (stat.+bkg. norm.) $^{+0.052}_{-0.054}$ (syst.)				
$F_{\rm L} = 0.299 \pm 0.008$ (stat.+bkg. norm.) $^{+0.013}_{-0.012}$ (syst.)	$F_{\rm L} = 0.281 \pm 0.021$ (stat.+bkg. norm.) $^{+0.063}_{-0.067}$ (syst.)				
$F_{\rm R} = -0.008 \pm 0.006$ (stat.+bkg. norm.) ± 0.012 (syst.	$F_{\rm R} = 0.061 \pm 0.022$ (stat.+bkg. norm.) $^{+0.101}_{-0.108}$ (syst.)				

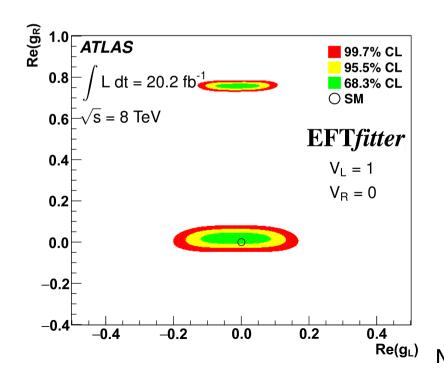
W boson polarisation (3)

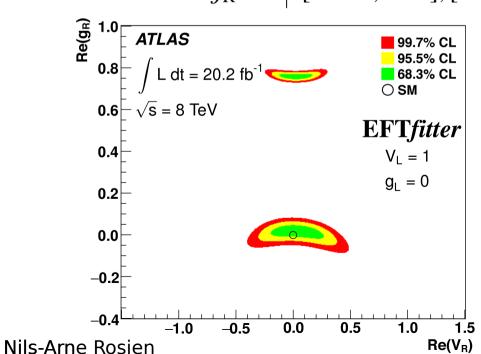


$$\mathcal{L}_{Wtb} = -\frac{g}{\sqrt{2}}\bar{b}\,\gamma^{\mu}\,(\underline{V_{L}}P_{L} + \underline{V_{R}}P_{R})\,t\,W_{\mu}^{-} - \frac{g}{\sqrt{2}}\bar{b}\,\frac{i\sigma^{\mu\nu}q_{\nu}}{m_{W}}\,(\underline{g_{L}}P_{L} + \underline{g_{R}}P_{R})\,t\,W_{\mu}^{-} + \text{h.c.}$$

- F_L , F_R , F_0 can constrain anomalous *Wtb* couplings
- Anomalous couplings are constrained using EFTfitter [Eur. Phys. J. C (2016) 76: 432]

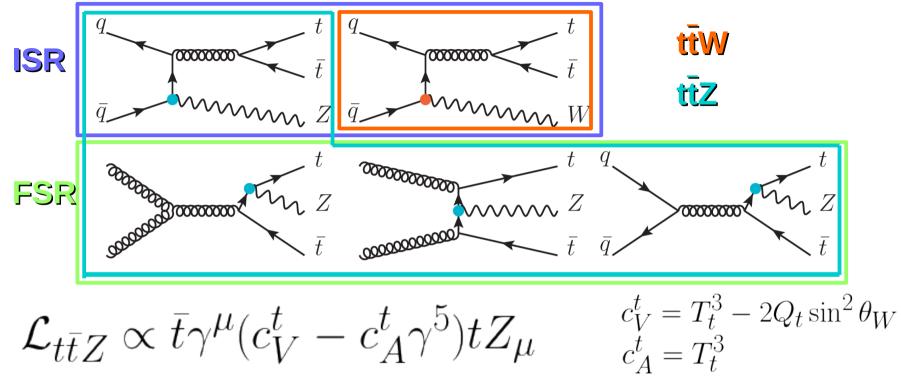






ttZ and ttW cross sections (1)



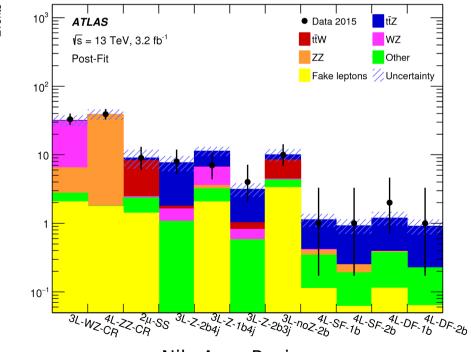


- Access to the the third component of the weak isospin of the top quark (FSR)
- Access to anomalous ttZ Couplings
- Learn about **electroweak symmetry breaking** via interactions of W and Z
- Indicator for strongly coupled Higgs sector, technicolour, heavy top partners
- Important background process for ttH (multilepton channel), SUSY (multilepton, stop pairs) and others
- Possibility to test PDFs via ttW because of ISR

ttZ and ttW cross sections (2)



- Separation into 3 different regions:
 - 2 same sign muons $(2\mu SS)$ \rightarrow sensitive to ttW, dominant bkg: fake leptons
 - $3\ell \rightarrow \text{sensitive to } ttW \text{ and } ttZ, \text{ dominant bkg: } WZ \text{ and fake leptons}$
 - $4\ell \rightarrow \text{sensitive to } t\bar{t}Z$, dominant bkg: ZZ and fake leptons
- 2μ SS and the 3ℓ channel: fakes are estimated using fully data driven matrix method
- 4ℓ channel: fake estimation using FF method using shapes from MC



ttZ and ttW cross sections (3)



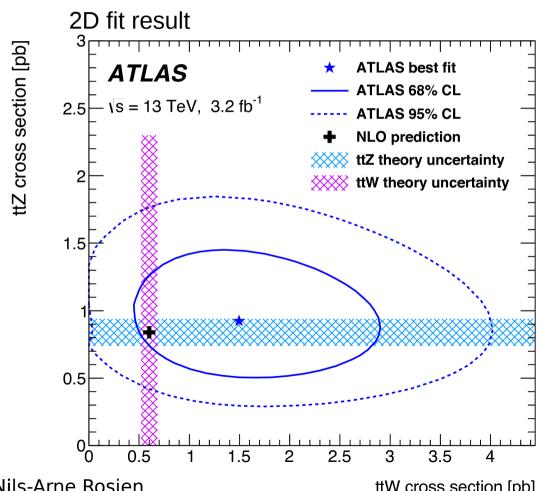
- Statistically limited
- First observation ($>5\sigma$) of this process for ATLAS (parallel to CMS at 8 TeV)

1D fit result

$$\sigma_{t\bar{t}W} = 1.5 \pm 0.8 \text{ pb}$$
 $\sigma_{t\bar{t}Z} = 0.9 \pm 0.3 \text{ pb}$

NLO QCD

$$\sigma_{t\bar{t}W} = 0.60 \pm 0.08 \text{ pb}$$
 $\sigma_{t\bar{t}Z} = 0.84 \pm 0.09 \text{ pb}$



Conclusion



- Top quark properties measurements at ATLAS in spin observables, charge and CP asymmetries, W boson polarisation and $t\bar{t}V$ cross section
- Some spin observables measured for the first time
- $t\bar{t}V$ observed for the first time at 8 TeV (parallel to CMS)
- No disagreements with the SM
- More measurements using 13 TeV data are right around the corner



Thank you very much for your attention!