

# Separated $\pi^0$ electroproduction cross sections off the nucleon at Jefferson Lab Hall A

DE LA RECHERCHE À L'INDUSTRIE

cea

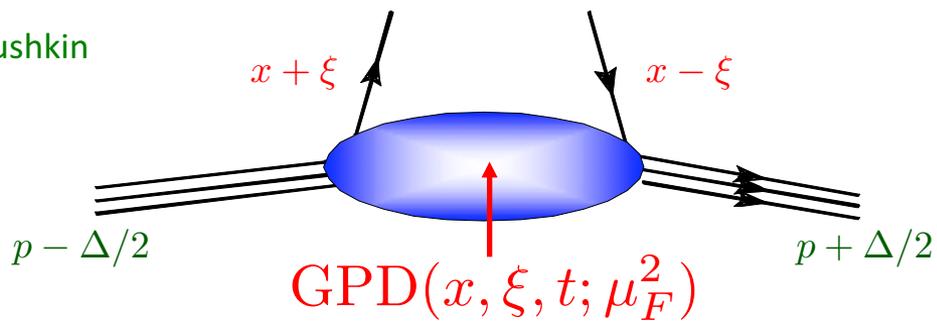
F. Sabatié – CEA Saclay



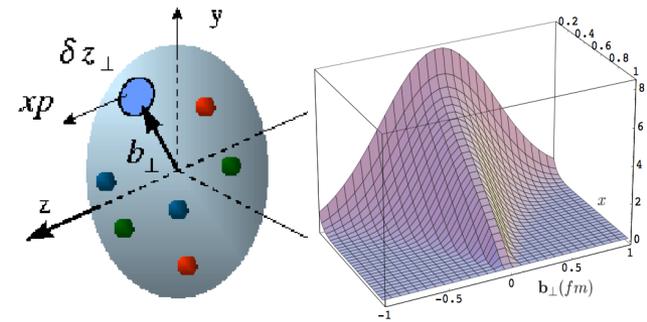
- Introduction
- Generalized Parton Distributions
- $ep \rightarrow ep(\gamma/\pi^0)$  to access GPDs
- $\pi^0$  analysis & proton results  
PRL 117, 262001 (2016)
- $\pi^0$  electroproduction off the **neutron** (first time ever !)  
arXiv:1702.00835
- Summary & Conclusion



Müller  
Ji, Radyushkin



$(x + \xi)$  and  $(x - \xi)$  : longitudinal momentum fractions of quarks



Impact parameter space  
Interpretation  
Burkardt, Diehl, ...

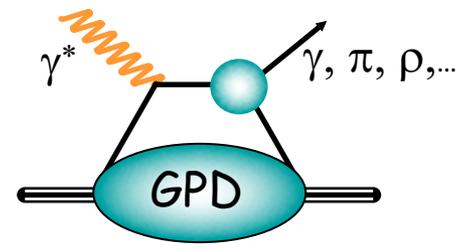
The structure of the nucleon can be described at leading-twist by

4 chiral-even **Generalized Parton Distributions** :

$$H, \tilde{H}, E, \tilde{E}(x, \xi, t; \mu_F^2)$$

4 chiral-odd  $H_T, \tilde{H}_T, E_T, \tilde{E}_T(x, \xi, t; \mu_F^2)$

> They enter the  $\gamma^* p \rightarrow (\gamma \text{ or } M) p$  amplitude  
as **convolution integrals** : no direct access



Deep Exclusive Processes  
**Parton distributions** in  
both coordinate and  
momentum space

> Forward limit ( $t = \xi = 0$ ) of  $H$  and  $\tilde{H}$  : PDFs

> First moment in x : Form Factors

> Second moment of  $(H + E)$  when  $t \rightarrow 0$  : total angular momentum



Belitsky and Mueller'00, Diehl'01

- Quark helicity flip GPDs

$$\begin{aligned}
 & \frac{1}{2} \int \frac{d\lambda}{2\pi} e^{ixP^+\lambda} \langle N(p') | \bar{\Psi}(-\lambda n/2) i\sigma^{+i} \Psi(\lambda n/2) | N(p) \rangle \\
 &= \frac{1}{2P^+} \bar{U}(p') \left[ H_T^q i\sigma^{+i} + \tilde{H}_T^q \frac{P^+ \Delta^i - \Delta^+ P^i}{m^2} \right. \\
 & \quad \left. + E_T^q \frac{\gamma^+ \Delta^i - \Delta^+ \gamma^i}{2m} + \tilde{E}_T^q \frac{\gamma^+ P^i - P^+ \gamma^i}{m} \right] U(p),
 \end{aligned}$$

- Gluon helicity flip GPDs

$$\begin{aligned}
 & \frac{1}{P^+} \int \frac{d\lambda}{2\pi} e^{ixP^+\lambda} \langle p' | \mathbb{S} G^{+i}(-\lambda n/2) G^{j+}(\lambda n/2) | p \rangle \\
 &= \mathbb{S} \frac{1}{2P^+} \frac{P^+ \Delta^j - \Delta^+ P^j}{2mP^+} \bar{U}(p') \left[ H_T^g i\sigma^{+i} + \tilde{H}_T^g \frac{P^+ \Delta^i - \Delta^+ P^i}{m^2} \right. \\
 & \quad \left. + E_T^g \frac{\gamma^+ \Delta^i - \Delta^+ \gamma^i}{2m} + \tilde{E}_T^g \frac{\gamma^+ P^i - P^+ \gamma^i}{m} \right] U(p),
 \end{aligned}$$



## Deeply Virtual Compton Scattering

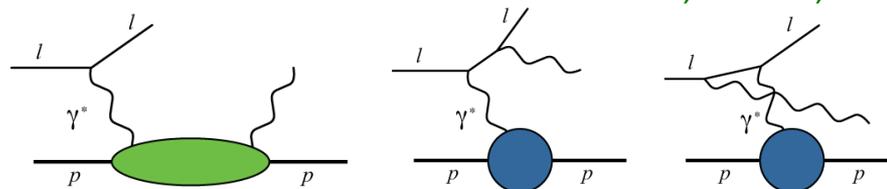
- Theory is under control : up to  $\alpha_S^2$ , twist-3, target mass corrections, etc.
- Sensitive to the quark combination :  $\frac{4}{9}u + \frac{1}{9}d + \frac{1}{9}s$

Müller et al,  
Braun et al, ...

- At Jefferson Lab energies, *mostly* sensitive to valence quarks
- Actually sensitive to *gluon* GPDs at NLO or beyond (even at somewhat large x)
- At LO, direct access the GPDs on the line  $x = \xi$  through Beam Spin Asymmetries sensitive to the interference with known **Bethe-Heitler** process

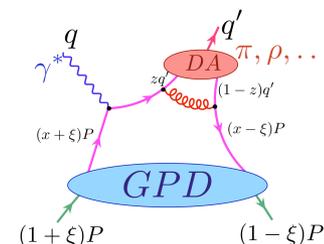
Moutarde, Pire,  
**FS**, Wagner, ...

Diehl, Gousset,  
Pire, Ralston, ...



## Hard Meson Electroproduction

- Many channels available for flavor separation ( $\rho^0$ ,  $\rho^+$ ,  $\pi^0$ ,  $\pi^+$ ,  $\phi$ , ...)
- $J/\Psi$  and  $\phi$  access gluon GPDs. Surprises with the  $\pi^0$  (later in this talk)
- Theory less under control : convolution with (unknown) meson WF,



potentially slow scaling, large power and NLO corrections



$$\frac{d^4\sigma}{dt d\phi dQ^2 dx_B} = \frac{1}{2\pi} \Gamma_{\gamma^*}(Q^2, x_B, E_e) \left[ \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{TL}}{dt} \cos(\phi) + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) \right]$$

$$\Gamma_{\gamma^*}(Q^2, x_B, E_e) = \frac{\alpha}{8\pi} \frac{Q^2}{M^2 E_e^2} \frac{1-x_B}{x_B^3} \frac{1}{1-\epsilon}$$

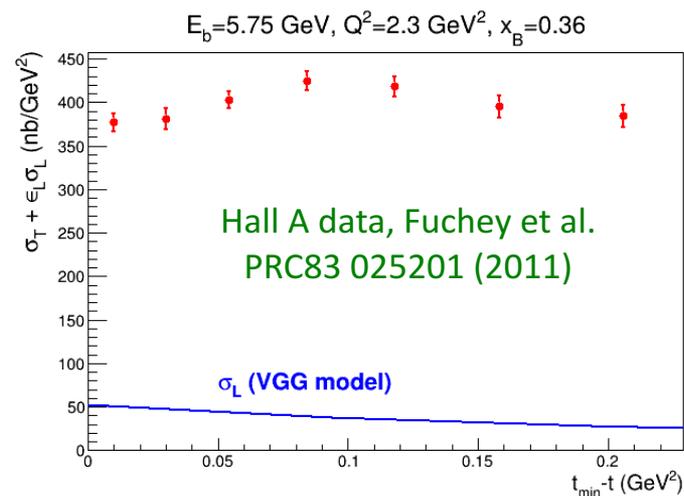
$$\epsilon = \frac{1-y-\frac{Q^2}{4E_e^2}}{1-y+\frac{y^2}{2}+\frac{Q^2}{4E_e^2}}$$

Factorization proven for L cross section only.  $\frac{d\sigma_L}{dt}$  is mostly sensitive to  $\tilde{H}$  :

$$\frac{d\sigma_L}{dt} \propto \frac{1}{Q^6} \left[ (1-\xi^2) \left| \langle \tilde{H} \rangle \right|^2 + \dots \right]$$

## However, exciting results in the last few years

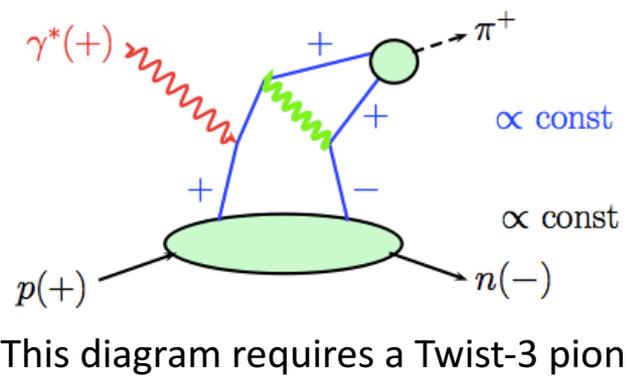
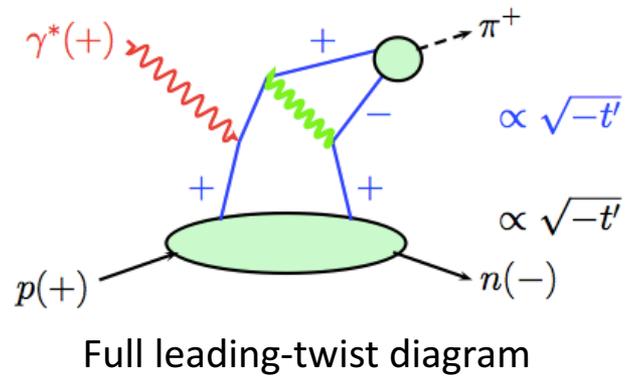
- Hall A and CLAS  $\pi^0$  un-separated cross section data overshoot chiral-even GPD models by a factor  $\sim 10$  (Figure: Hall A data, similar findings from CLAS)
- $Q^2$  dependence is too large for  $\gamma_L\text{-}\pi$  transition and actually compatible with a  $\gamma_T\text{-}\pi$  transition (Hall A)





Latest GK : Kroll, Moutarde, Goloskokov, **FS**, EPJC73, 2278 (2013)  
 Chiral-odd GK: Goloskokov, Kroll, EPJA47, 112 (2011)

Explanation: chiral-odd GPD may contribute to transverse part



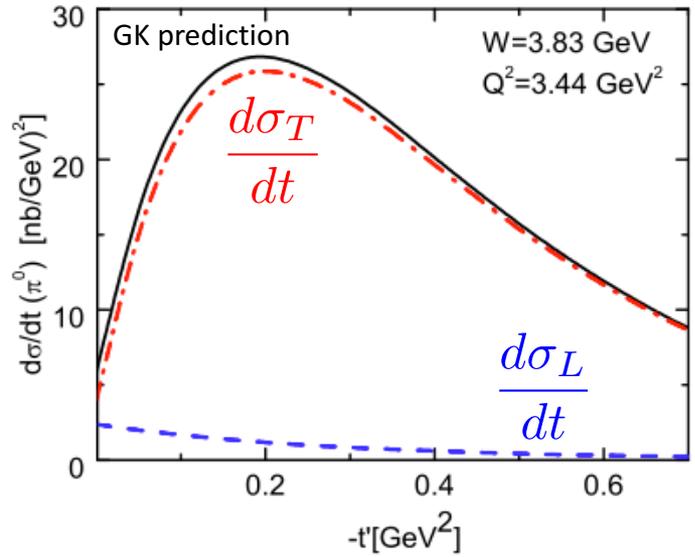
$$\mu_\pi = \frac{m_\pi^2}{m_u + m_d} \simeq 2.5 \text{ GeV}$$

In this approach :

$$\frac{d\sigma_T}{dt} \propto \frac{1}{Q^8} \left[ (1 - \xi^2) |\langle H_T \rangle|^2 - (t'/8m^2) |\langle \bar{E}_T \rangle|^2 \right]$$

with  $\bar{E}_T = 2\tilde{H}_T + E_T$

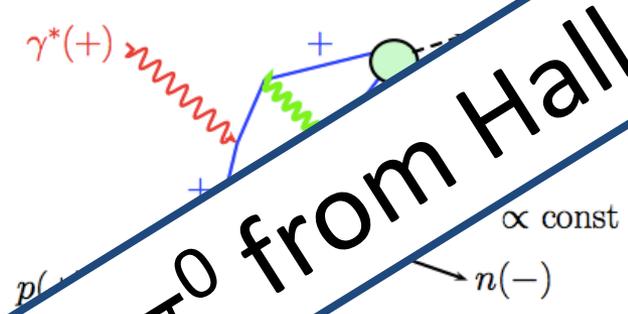
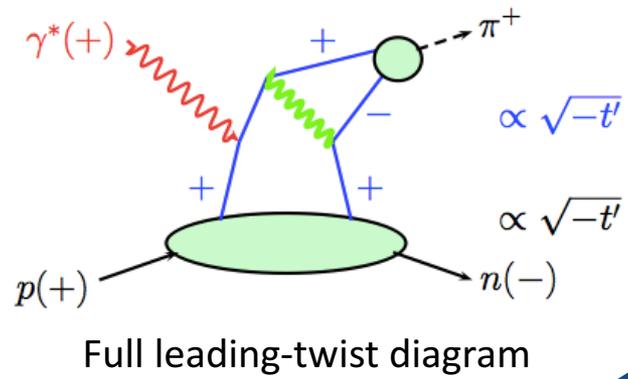
Low- $t'$  behavior results from interplay between  $H_T, \bar{E}_T$



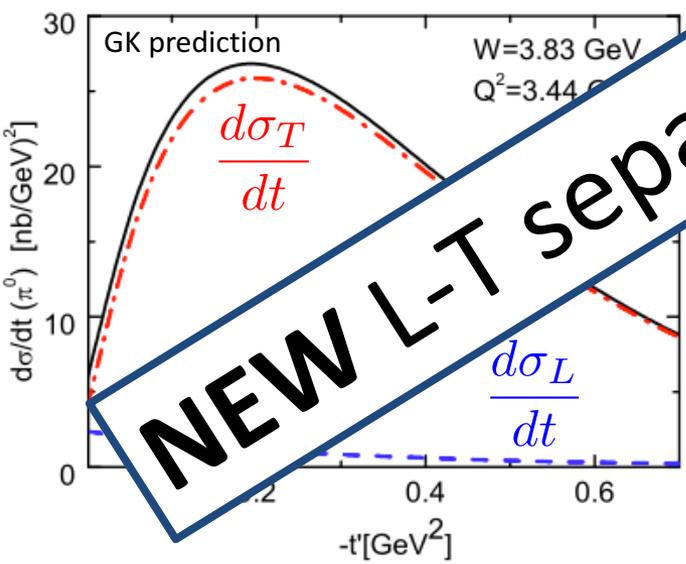


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Explanation: chiral-odd GPD may contribute to transverse part



NEW L-T separation of  $\pi^0$  from Hall A!



requires a Twist-3 pion w.f.  
 by  $\mu_\pi/Q$  wrt leading twist, BUT

$$\mu_\pi = \frac{m_\pi^2}{m_u + m_d} \simeq 2.5 \text{ GeV}$$

in this approach :

$$\frac{d\sigma_T}{dt} \propto \frac{1}{Q^8} \left[ (1 - \xi^2) |\langle H_T \rangle|^2 - (t'/8m^2) |\langle \bar{E}_T \rangle|^2 \right]$$

with  $\bar{E}_T = 2\tilde{H}_T + E_T$

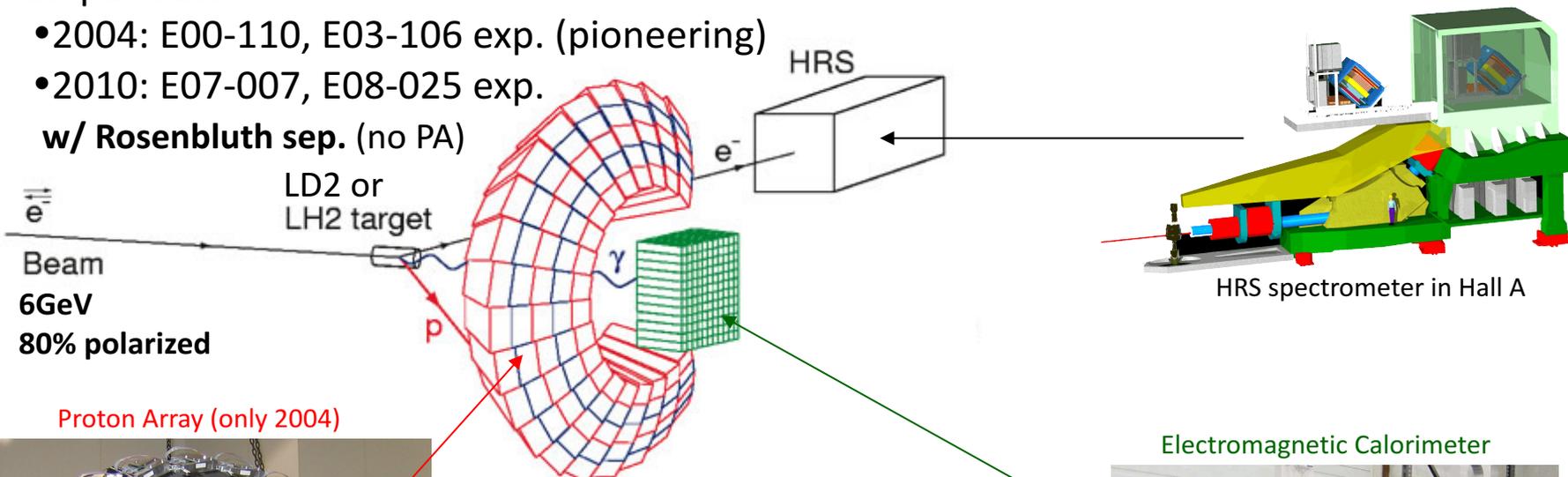
Low- $t'$  behavior results from interplay between  $H_T, \bar{E}_T$



Two run periods:

- 2004: E00-110, E03-106 exp. (pioneering)
- 2010: E07-007, E08-025 exp.

w/ Rosenbluth sep. (no PA)



Beam  
6GeV  
80% polarized

Proton Array (only 2004)



100 scintillator blocks

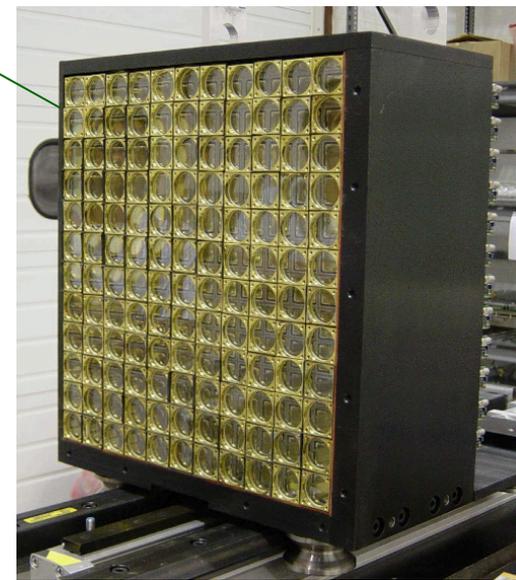
$E_e = 5.75$  GeV

2004		
	0.36	1.5
	0.36	1.9
	0.36	2.3

$E_e = 3.45, 4.55, 5.55$  GeV

2010			
	0.36	1.5	0.52, 0.84
	0.36	1.75	0.65, 0.79
	0.36	2	0.53, 0.72

Electromagnetic Calorimeter

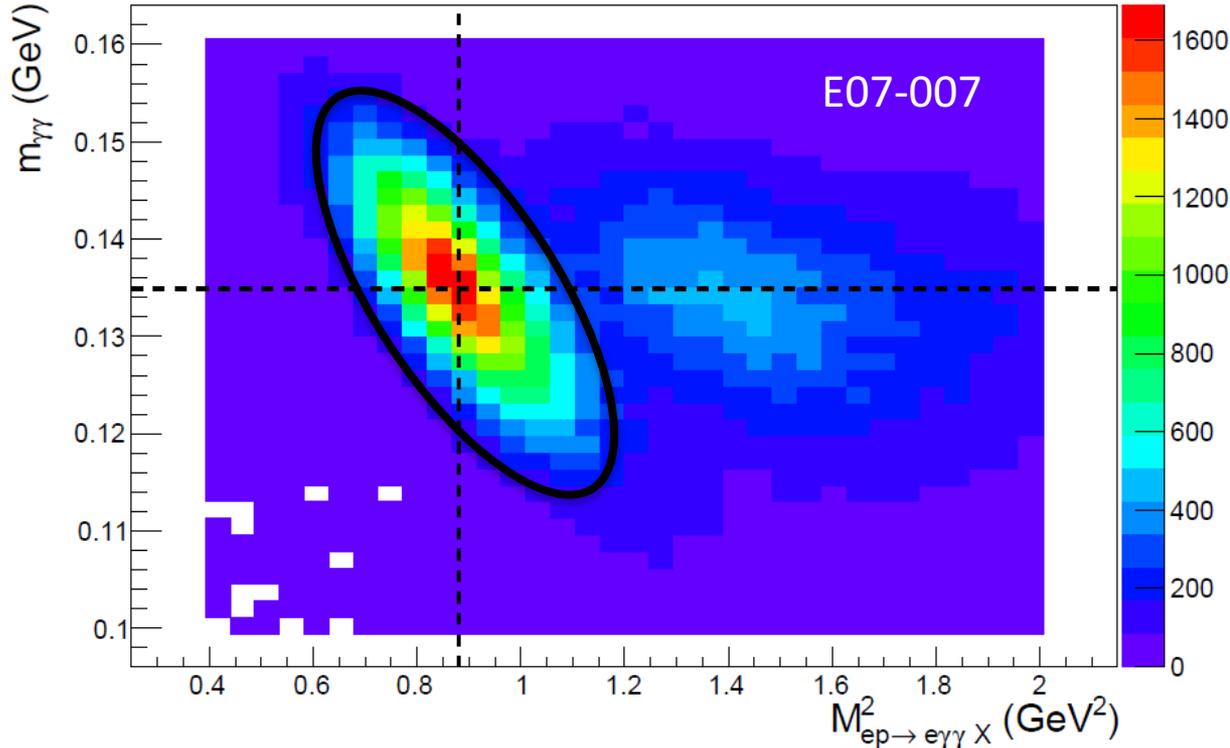


132  $PbF_2$  blocks  
(208 in 2010)



**Inclusive electron trigger** (99% pure) using Hall A Left HRS  
Require **two clusters in the calorimeter**

$ep \rightarrow ep\pi^0$  process cleanly selected by dual cut on  $\gamma\gamma$  invariant mass and  $p$  missing mass



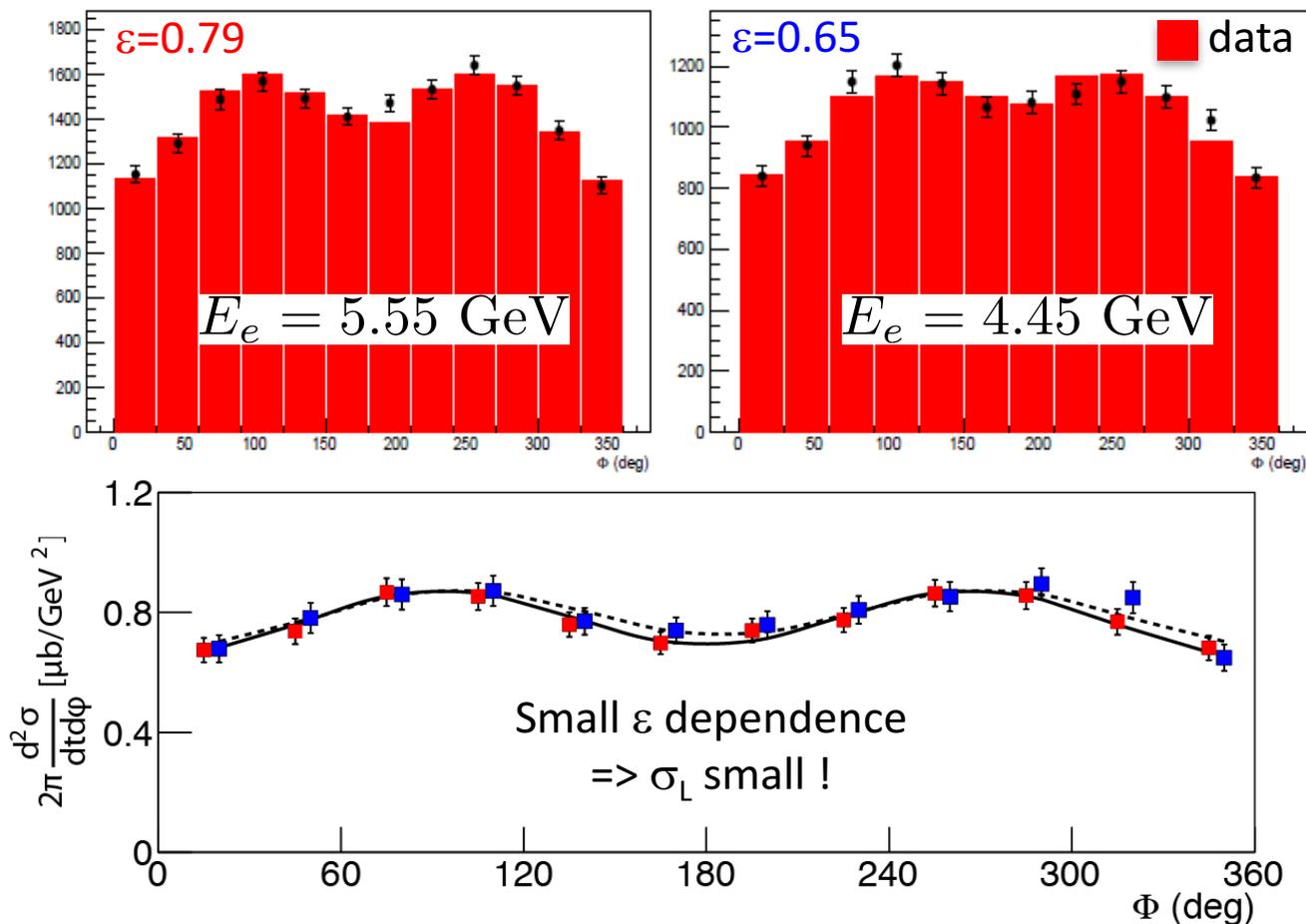
**Matching of Monte-Carlo and data resolutions of utmost importance to reduce systematic uncertainties from the cuts**

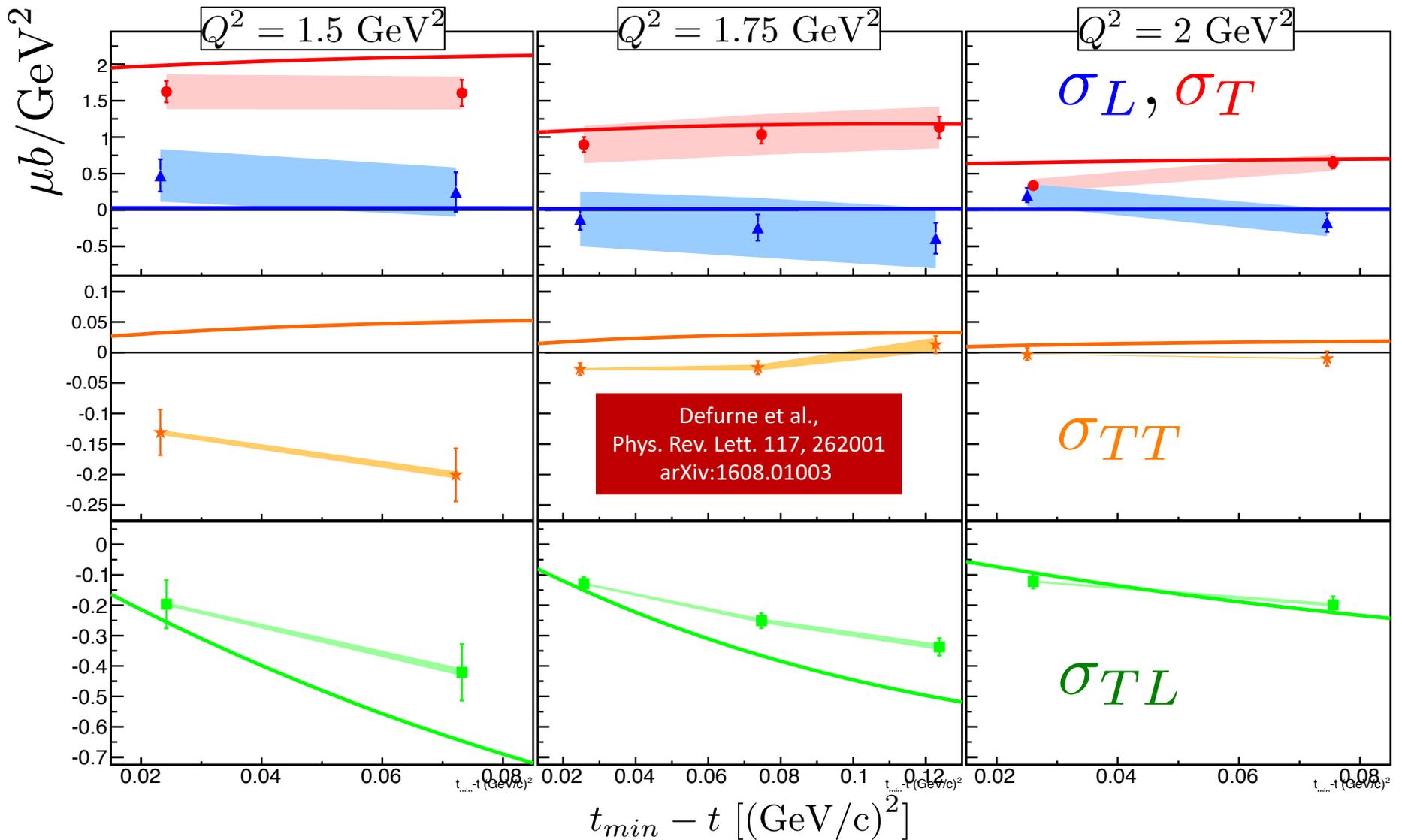


Extraction of all the (unpolarized) responses by a simultaneous fit to the 2 beam energies

$$\frac{d^4\sigma}{dt d\phi dQ^2 dx_B} = \frac{1}{2\pi} \Gamma_{\gamma^*}(Q^2, x_B, E_e) \left[ \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{TL}}{dt} \cos(\phi) + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) \right]$$

$$t = -0.02 \text{ GeV}^2, Q^2 = 1.75 \text{ GeV}^2$$





Shaded area: 2% normalization uncertainty

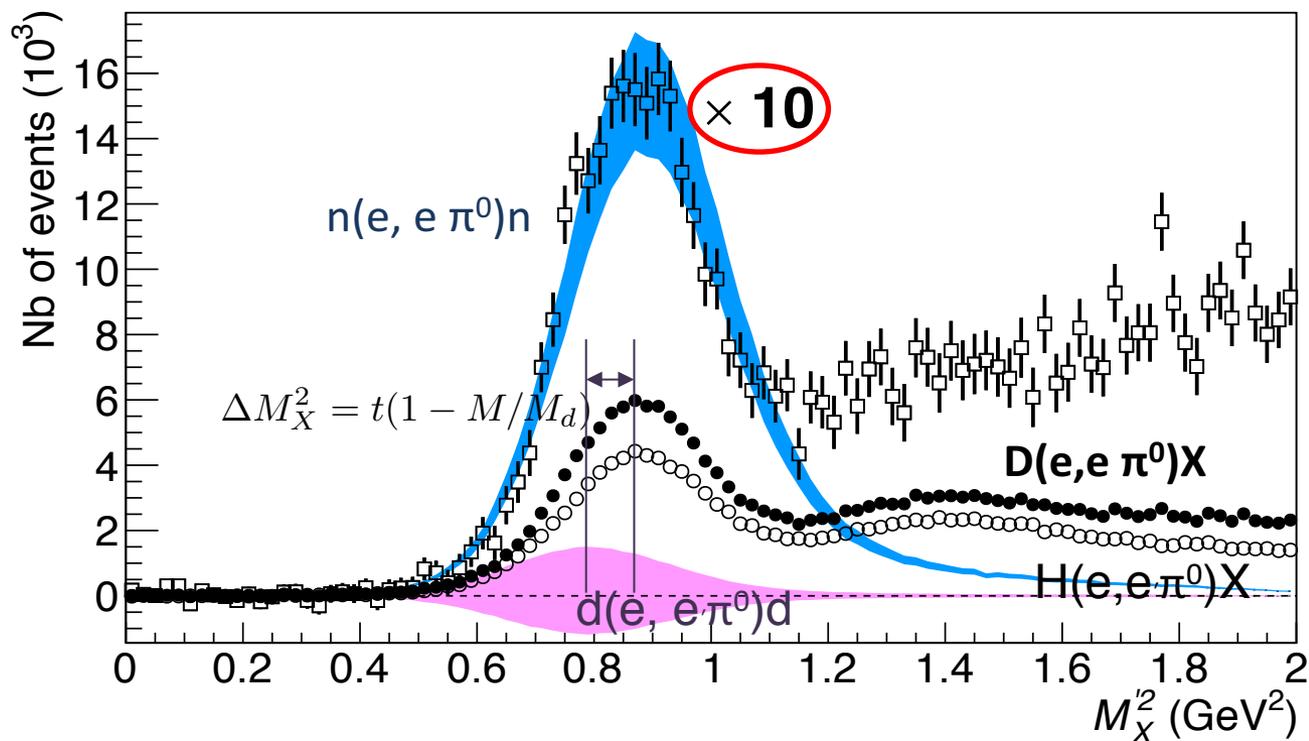
Solid line: GK model

Reasonable agreement with GK model as described before

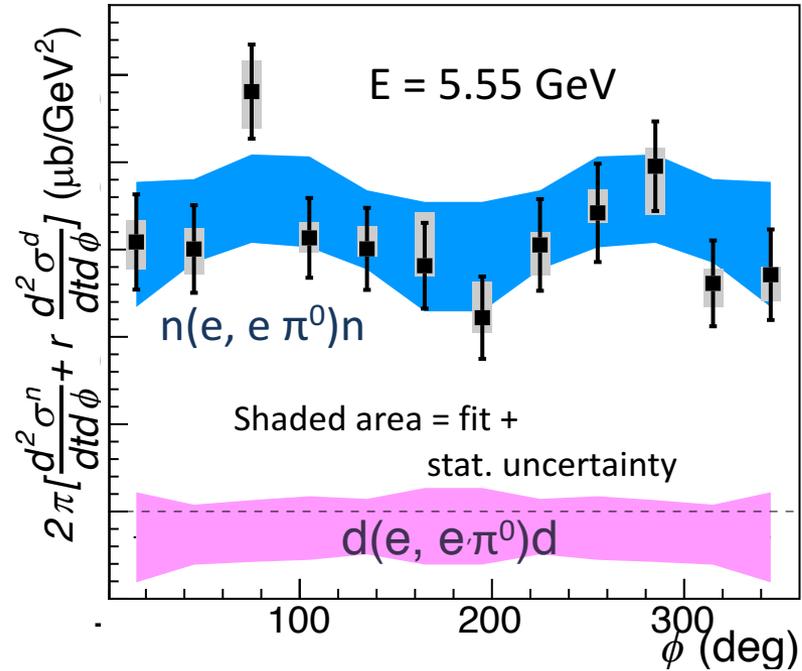


The quasi-free  $\pi^0$  electroproduction events off the proton are **subtracted using the (normalized) data from E07-007** (concurrent data taking, just switching the LH2 target with the LD2 target)

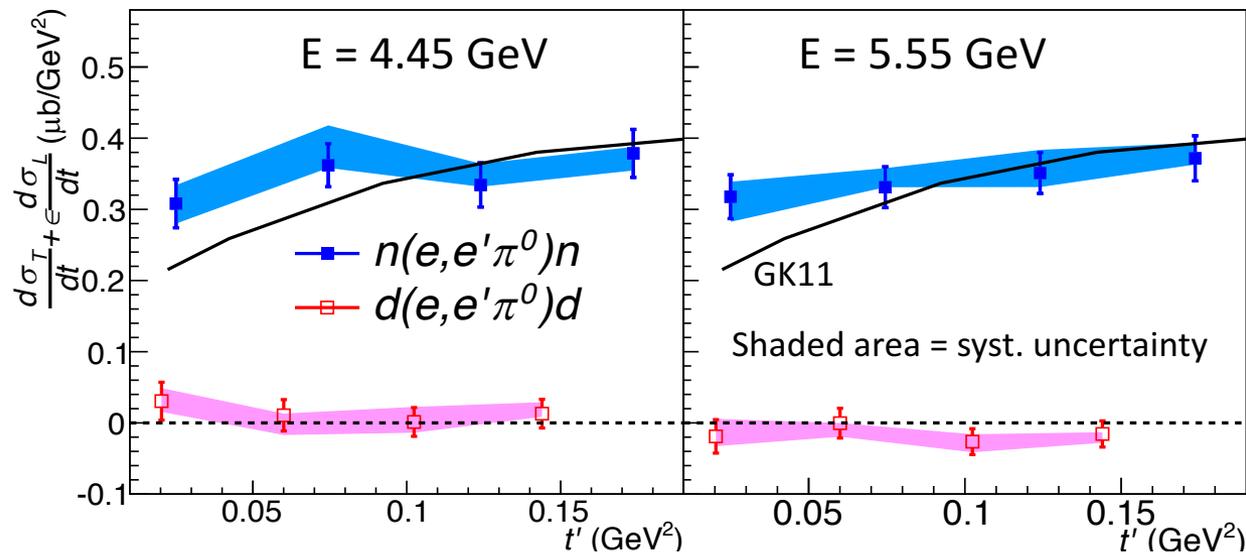
$$D(e, e \pi^0)X - p(e, e \pi^0)p = n(e, e \pi^0)n + d(e, e \pi^0)d$$

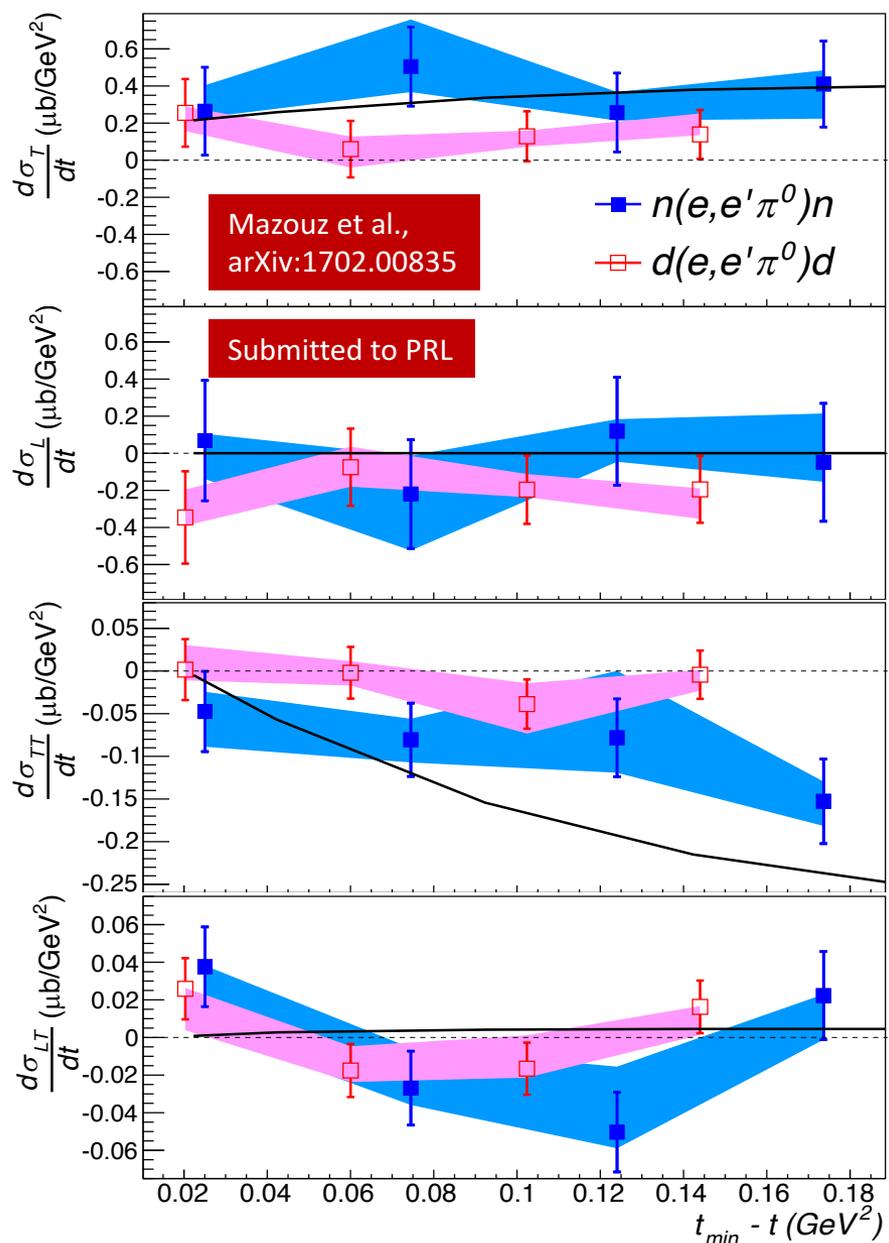


The average momentum transfer to the target is much larger than the np relative momentum, justifying this impulse approximation

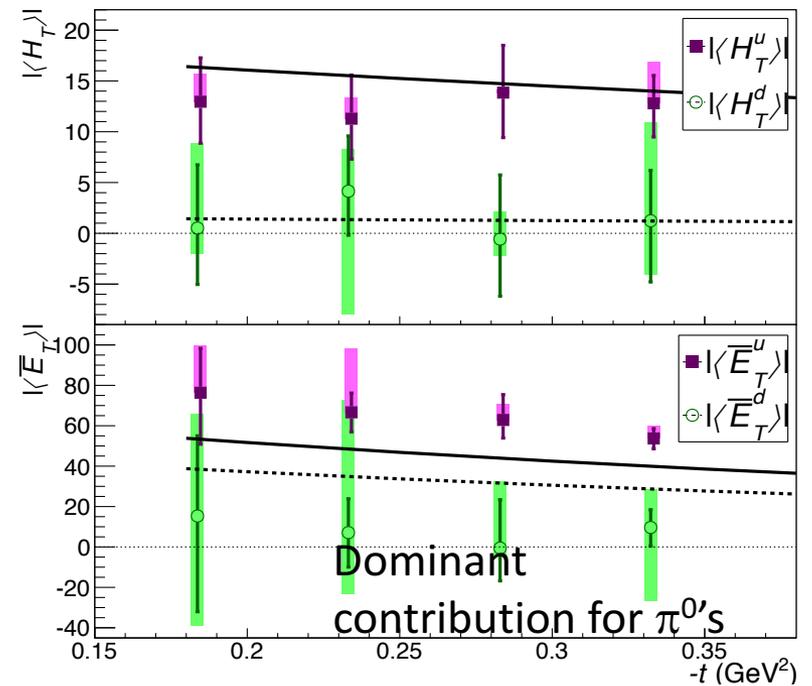


- The 4x2 separated cross sections for (n,d) are fit to the proton-subtracted deuterium data, for the two energy settings at once
- The unpolarized neutron cross section is found to have (again) a **very low dependence on the beam energy**, indicating a small L component
- Unpolarized cross sections are **compared to GK model, with reasonable agreement**



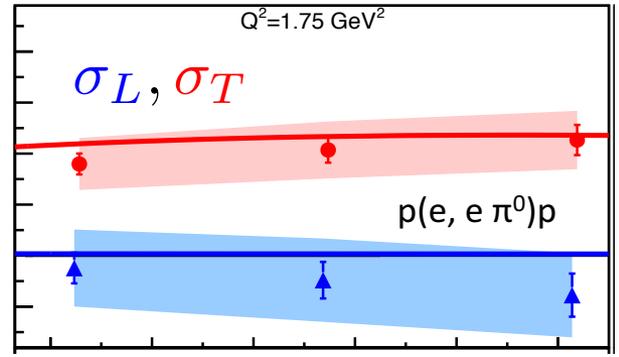
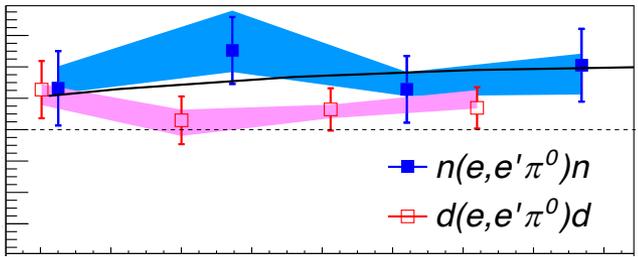


- Challenging measurement for the proton ...**even more challenging for the neutron**
- However, **significant  $\sigma_T$  and  $\sigma_{TT}$**
- And  $\sigma_L$  **still close to 0**, similarly to the proton
- Within GK framework, extraction of flavor dependence of  $|\langle H_T \rangle|$ ,  $|\langle \bar{E}_T \rangle|$  from neutron and proton data
- Data are compared to **GK model**, with **reasonable agreement**



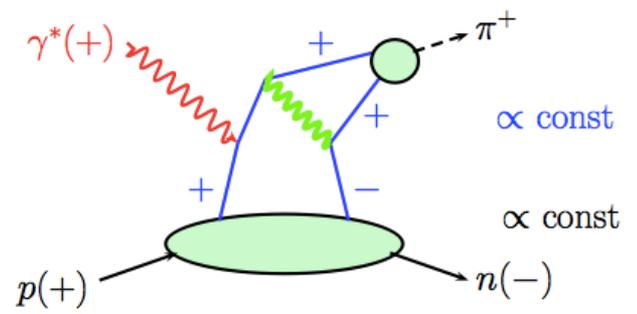


- Finally, **proof** that **transverse  $\pi^0$  cross section is dominant**, for both **proton and neutron targets**



- Cross sections are reasonably well described by the **GK model** using transversity “helicity-flip” GPDs and a *Twist-3* pion wave function

→ **Access to the elusive chiral-odd GPDs**



- An **exciting future** ahead

**12 GeV experiments at JLab** (Hall A already took some data)

COMPASS data at lower x, and then an **exciting program at the EIC !**





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## Experimental Hadron Physics

### [IRFU, SPhN, Saclay](#) - Senior

**Field of Interest:** hep-ex, nucl-ex

**Deadline:** 2017-04-14

**Region:** Europe

**Job description:**

The CEA Paris-Saclay Nucleon Structure Laboratory (LSN) is opening a staff scientist position for an outstanding physicist in the field of experimental hadron physics, with a focus on both the development of the Electron-Ion Collider (EIC) project and the experiments using the recent 12GeV upgrade of the Jefferson Lab (US) electron accelerator.

**DEADLINE : APRIL 14<sup>th</sup>**