

Experimental Calibration of the Top-Quark Monte-Carlo Mass

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04.04.2017



1) CERN

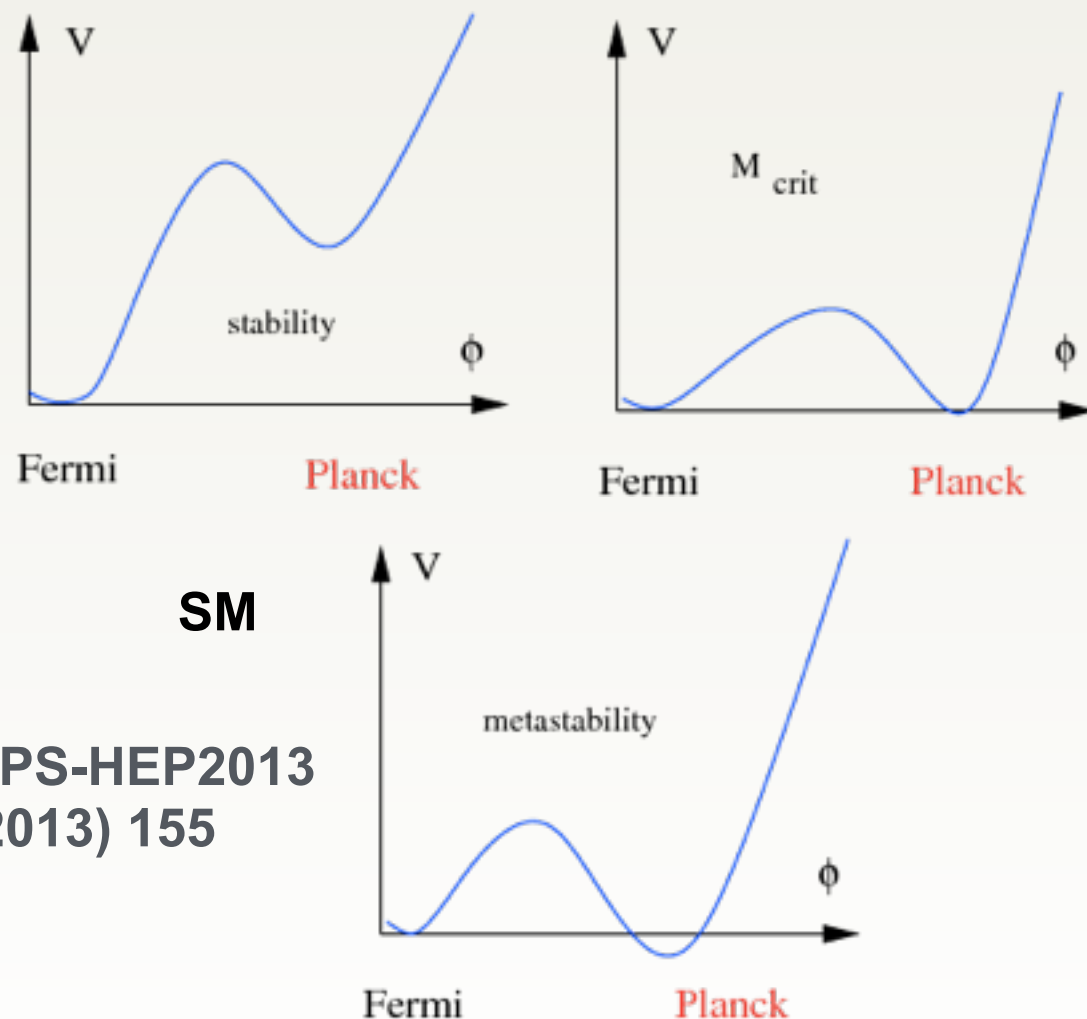
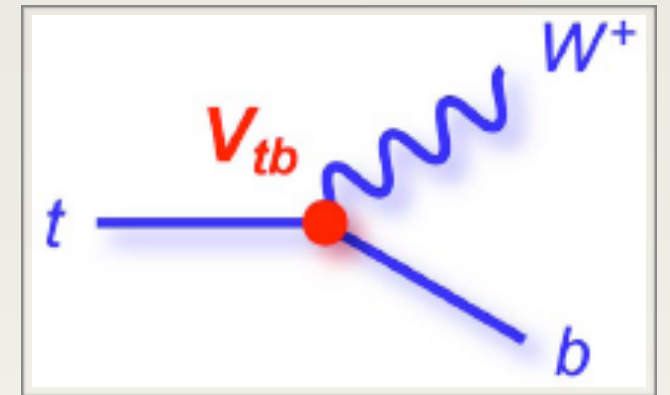
2) DESY

3) Universität Hamburg

The Top Quark and its Mass

⇔ **Heaviest fundamental particle in SM**

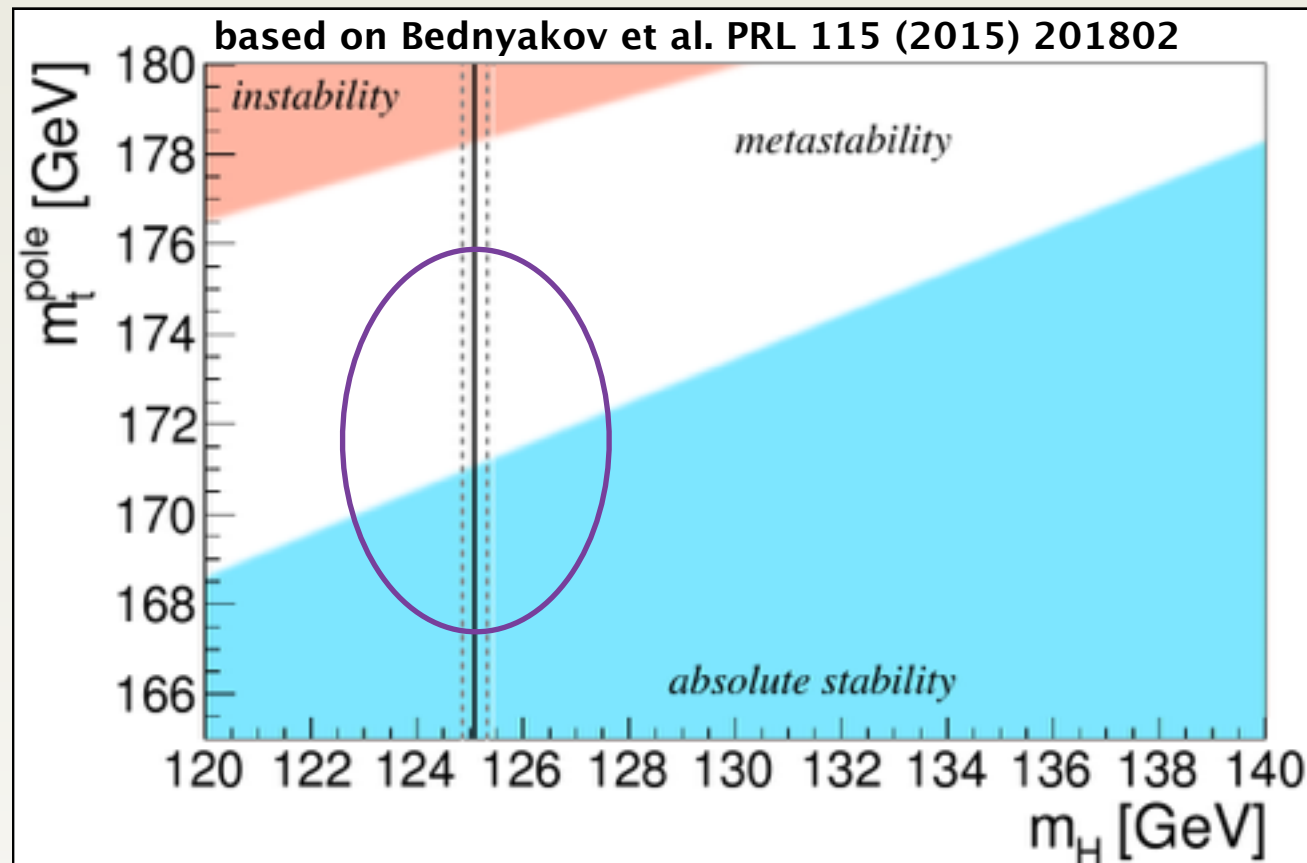
- Possible to study bare-quark properties
- Uniquely strong coupling to Higgs field
- Special role in electroweak symmetry breaking
- New physics may couple preferably to top quarks



- Important role in EWK fits
- EWK vacuum stability critically depends on
 - Higgs-boson mass
 - Top-quark (pole) mass

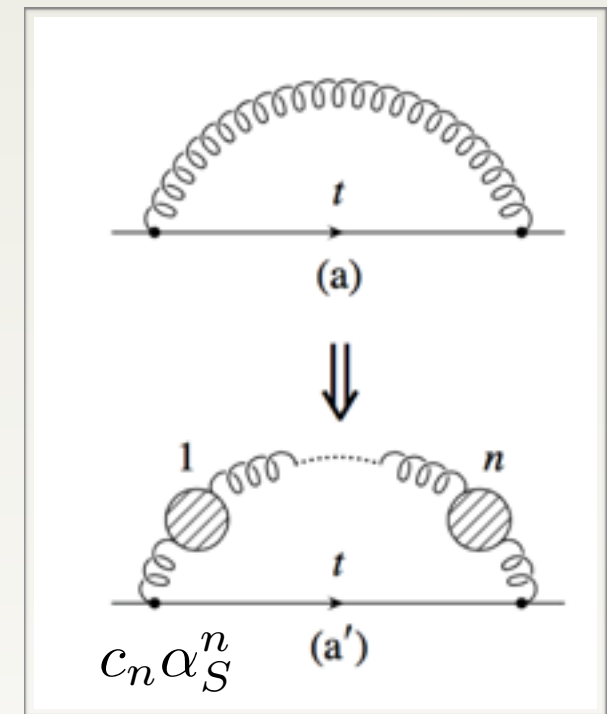
PoS EPS-HEP2013
(2013) 155

Top-Quark Mass In Calculations



$$m_H = 125.09 \pm 0.24 \text{ GeV}$$

ATLAS, CMS Collaborations,
PRL 114 (2015) 191803



- Beyond LO: self-energy corrections
 - Top-quark mass is renormalisation scheme dependent
 - pole mass: $m_{t,\text{pole}} \rightarrow O(\Lambda_{\text{QCD}})$ ambiguity (c_n diverge $\sim n!$)
 - running masses $m(\mu)$, e.g. MSbar mass: $\bar{m}_t(\mu)$
 - ...any many others (see G. Corcella's talk)
- ➔ “well-defined” m_t for calculations

Top-Quark Mass in Monte-Carlo Simulation

- Initial protons
 - Compound objects
 - Described by PDF's

- Hard interaction
 - Calculable in pQCD

- Hard decay

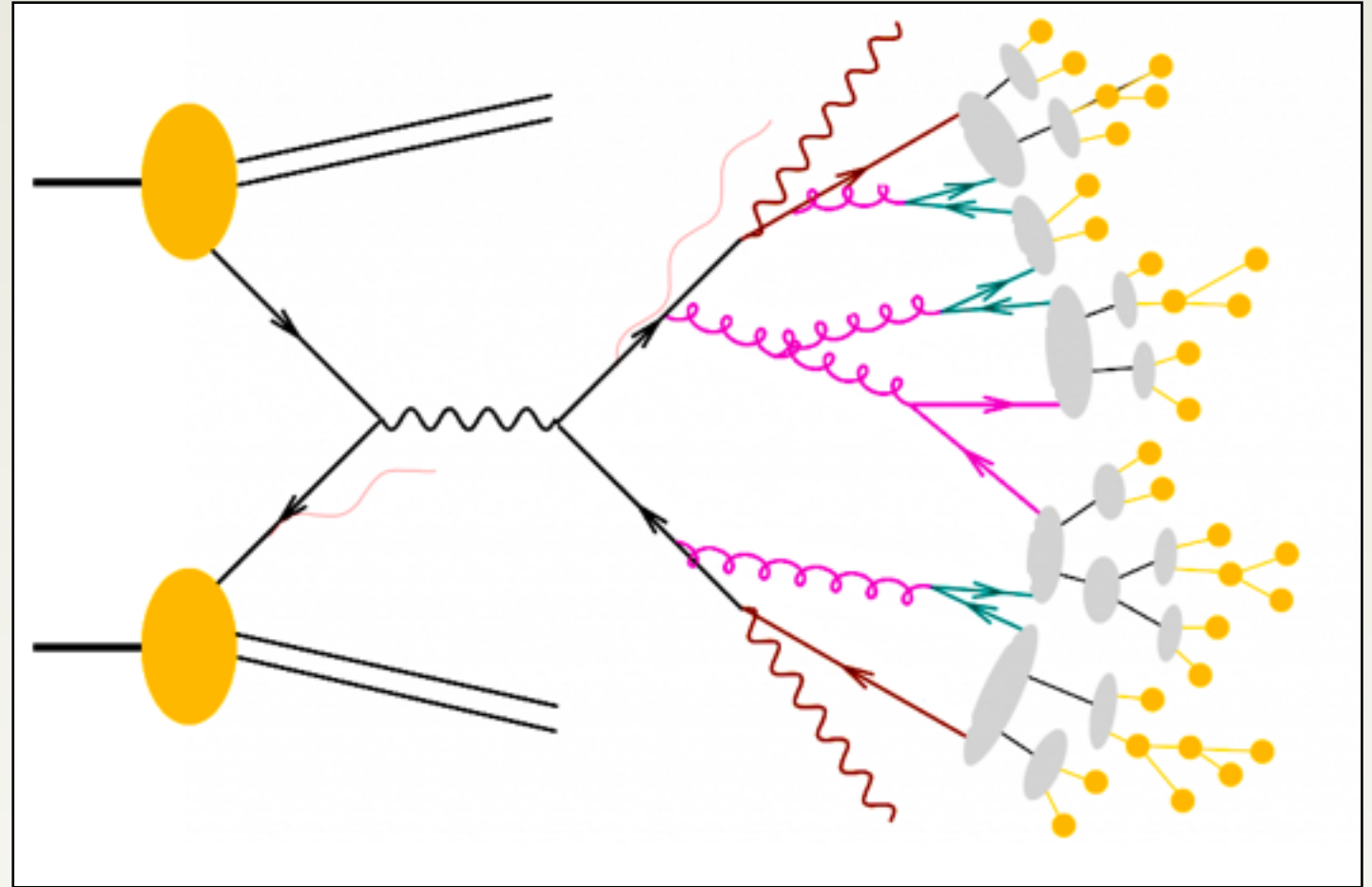
-
- Parton shower

- Hadronization

→
**visible
signature**

production

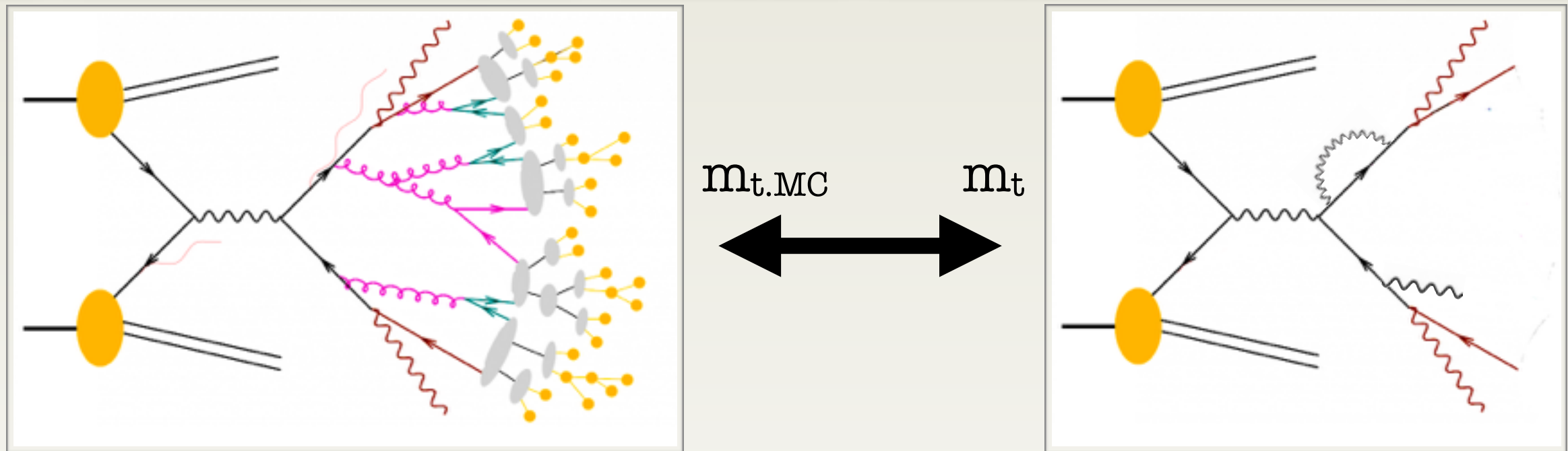
Decay to final state



- Based on heuristic models and approximations with parameters tuned to data
 - Also top-quark mass can be tuned to describe data
- ➡ Mass measurement

→
**visible
signature**

Relation between MC- and well-defined Mass



- Direct top-quark mass measurements

- ▶ Using final states from MC simulation (models)
- ▶ Measure MC parameter $m_{t,MC}$ (in principle depends on generator)

$$m_{t,MC} = 172.44 \pm 0.49 \text{ GeV}$$

CMS Collaboration
PRD 93 (2016) 072004

Hoang, Steward, NPPS 185 (2008)
Butenschoen et al., PRL 117 (2016) 232001

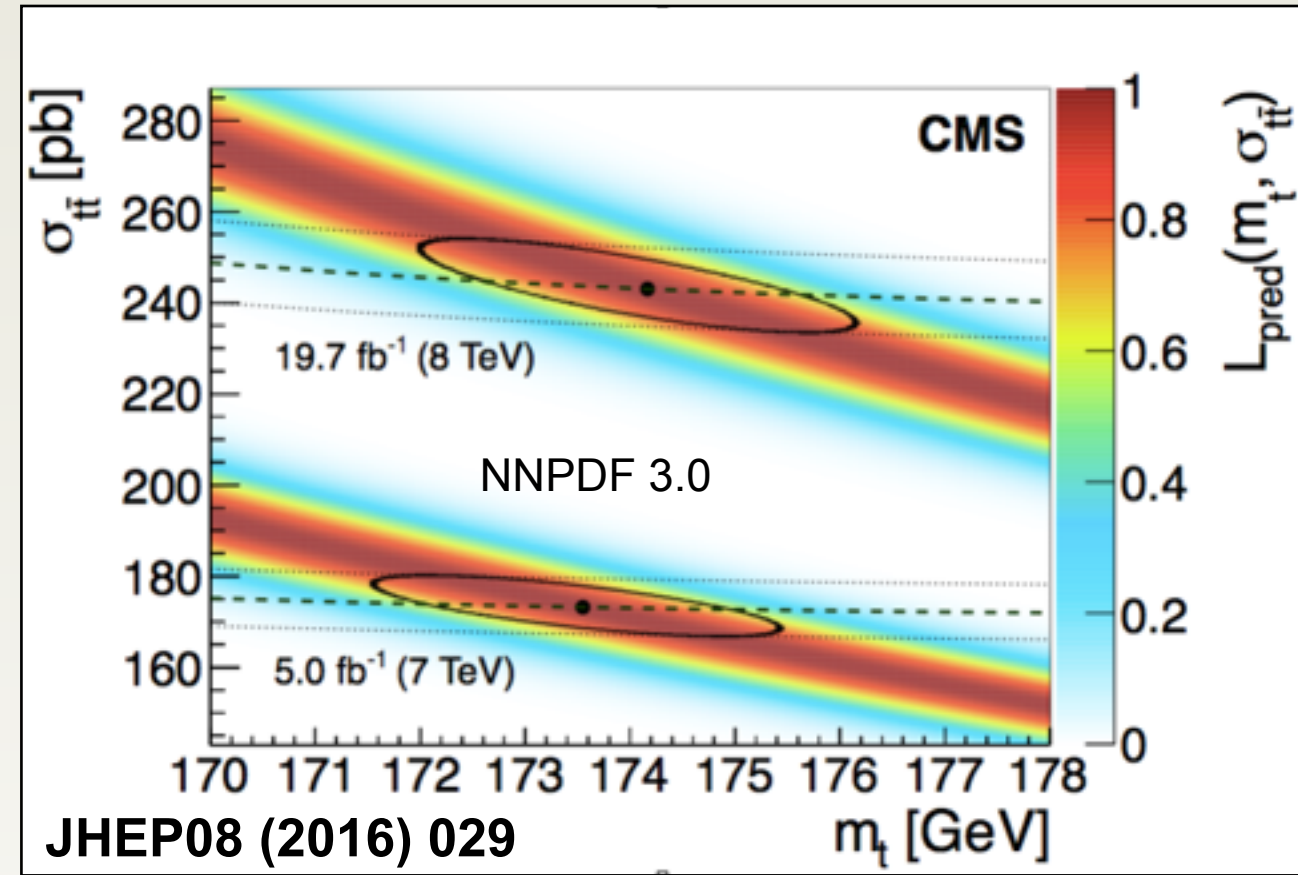
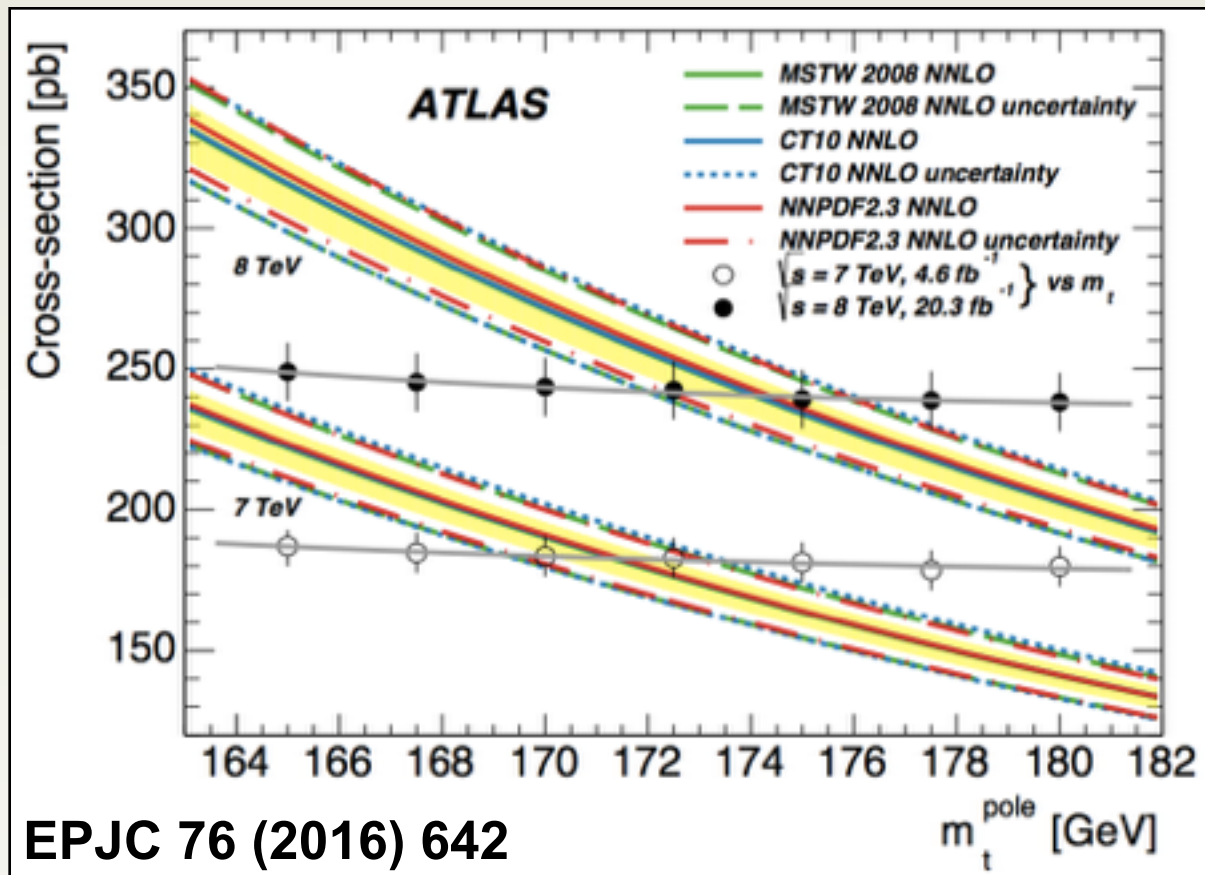
- Exact interpretation of $m_{t,MC}$ in terms of well-defined m_t

- ▶ Uncertainty $\approx 1 \text{ GeV}$ (pp), studies to reduce uncertainty ongoing (see G. Corcellas talk)

Buckley et al, Phys. Rept. 504 (2011)

- ▶ For measurements: often assumed $m_{t,pole} - m_{t,MC} \approx 1 \text{ GeV}$

Determine well-defined Mass directly



- Predicted production cross section depends significantly on m_t
- NNLO predictions using **well-defined m_t** (here pole mass) available

Czakon et al. PRL 110 (2013) 252004

- Measure σ_{tt} precisely (in $e\mu$ channel)
- Dependence of measurement on $m_t(\text{MC})$ mild

Pole-Mass Results

- **Combine** result from 7 and 8 TeV

➔ Most precise single pole mass determination

- (Higher precision can be reached in global PDF fits $O(1 \text{ GeV})$)
arXiv:1701.05838

- Uncertainties from measured and predicted $\sigma_{t\bar{t}}$ contribute equally
- Main difference between ATLAS and CMS: CMS uses more recent PDF sets.

EPJC 76 (2016) 642

ATLAS

$$m_t^{\text{pole}} = 172.9_{-2.6}^{+2.5} \text{ GeV}$$

JHEP08 (2016) 029

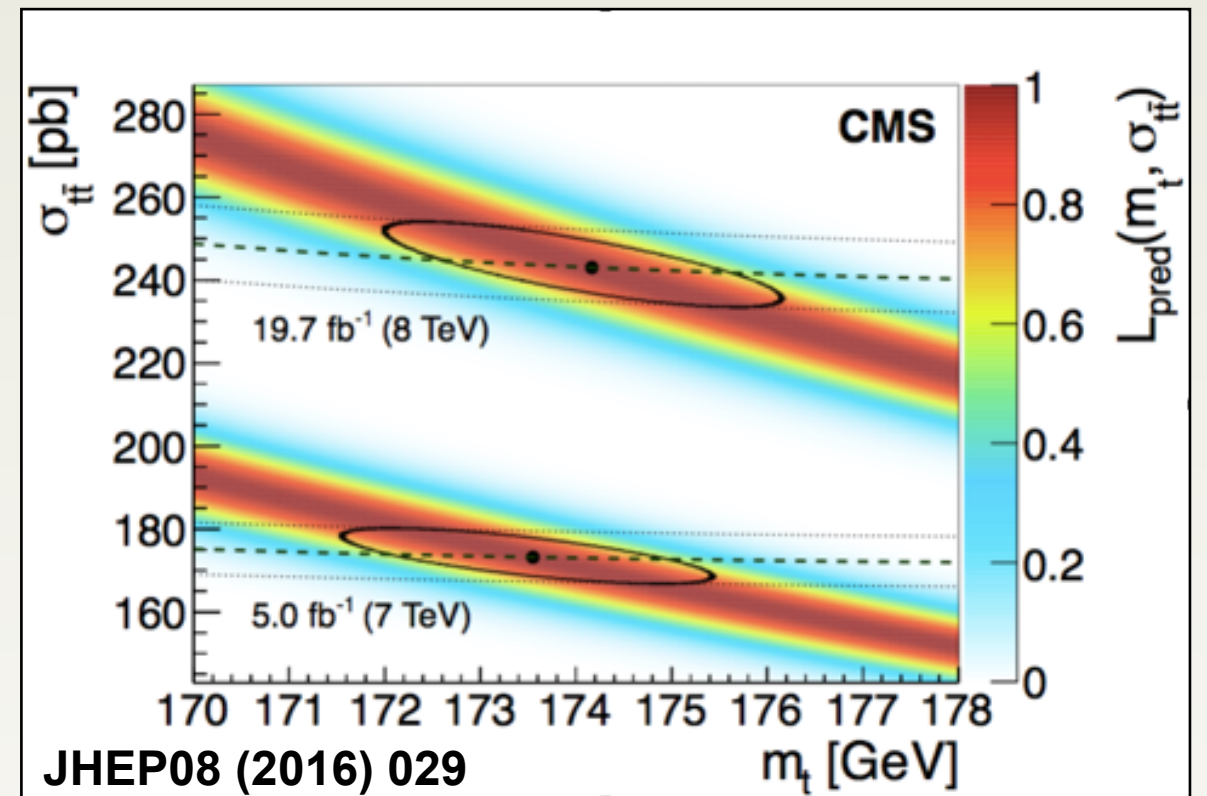
CMS	m_t [GeV]
NNPDF3.0	$173.8_{-1.8}^{+1.7}$
MMHT2014	$174.1_{-2.0}^{+1.8}$
CT14	$174.3_{-2.2}^{+2.1}$

- Working on a combination of results for $\sigma_{t\bar{t}}$ from ATLAS and CMS
 - significant gain in precision expected [1]
 - **paper will include subsequent pole-mass extraction (recent PDF sets)**

[1] JK "Update on t-tbar production results", LHCTopWG open meeting 17.5.2016

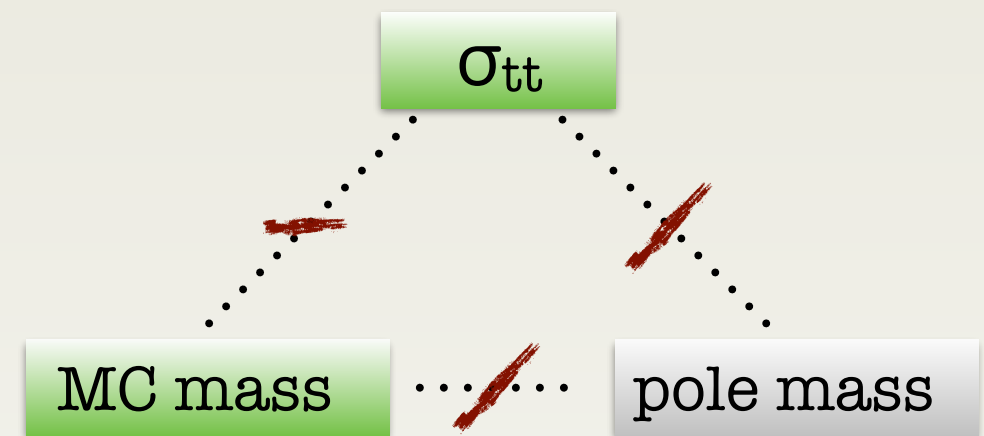
MC Mass in Measurements

- Dependence of experimental result is evaluated using MC mass
- “[...]an additional uncertainty $\Delta m_{t\pm}$ in the obtained cross section dependence is introduced. It is evaluated by shifting the measured dependence by ± 1 GeV [...]”
JHEP08 (2016) 029
- Still: quantitative assumption on relation between MC mass and pole mass (or other well-defined mass) needed
- Assumption can be avoided
 - Use existing measurements of MC mass \rightarrow precise but easily inconsistent (which MC, uncertainties and correlations, ...)
 - Measure MC mass simultaneously

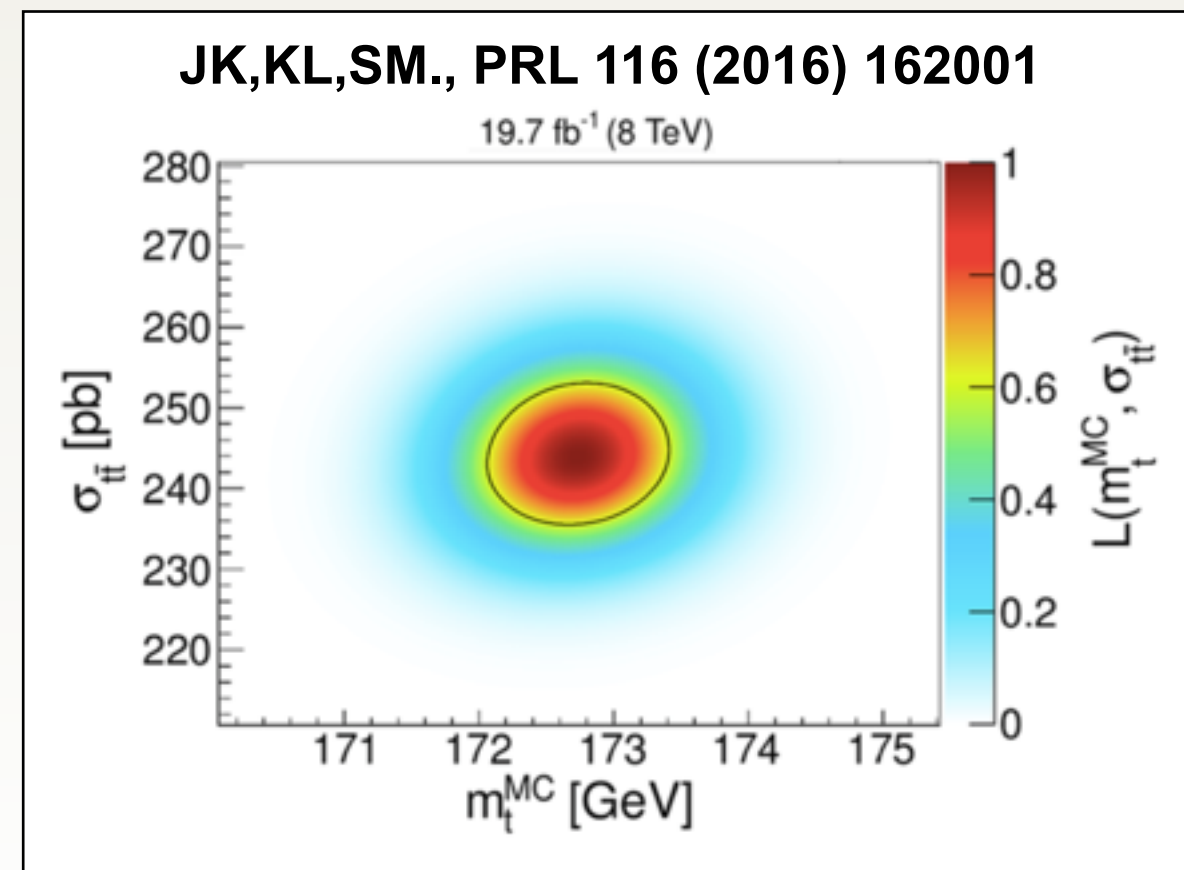


Mitigate dependence on MC mass

- Assume all dependencies to be unknown
- Absorb MC mass dependence in uncertainty on σ_{tt} through simultaneous fit
 - Shape of e.g. m_{lb} : MC mass
 - Normalisation: σ_{tt}
- ➔ **Measurement of σ_{tt} and MC mass**

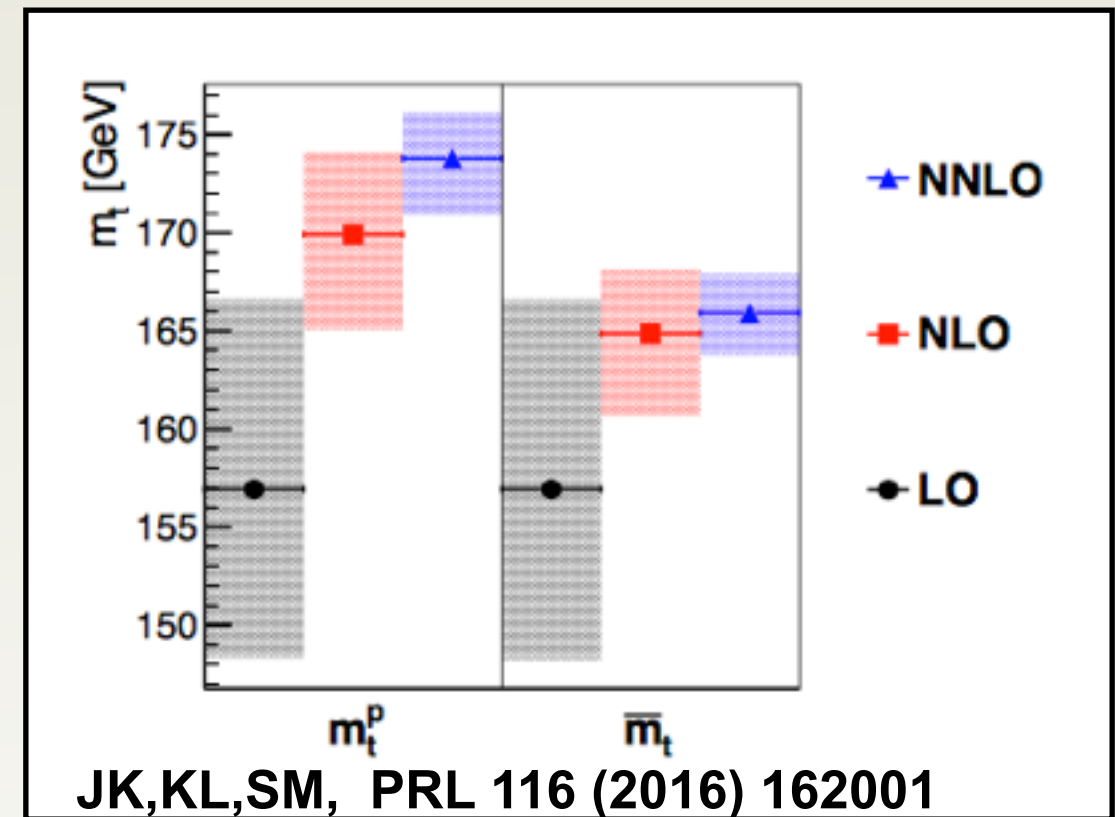


- ➔ Improved physics interpretation of measured $\sigma_{tt}(\cancel{m_t})$
 - Only assumption some weak qualitative relation between MC mass and well-defined mass
- ➔ Determine pole or \overline{MS} mass (or any other) from **direct comparison**
- ➔ **No assumptions** on the relation between MC mass and pole/ \overline{MS} mass needed
- ➔ Difference can be measured



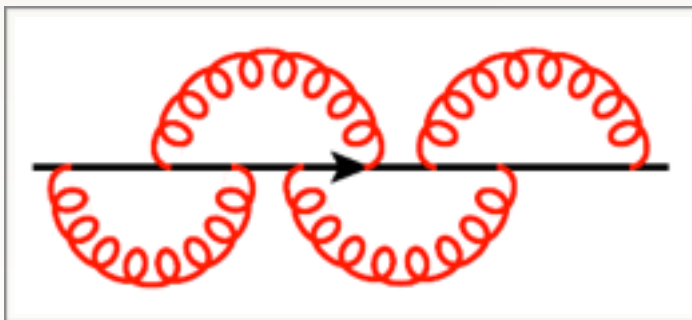
MSbar mass scheme

- Study on extracted top-quark mass
 - Consider measured σ_{tt} independent of m_t
 - Extract m_t by comparison with predicted $\sigma_{tt}(m_t)$
- ➔ Using \bar{m}_t improves perturbative convergence



- Conversion between $\overline{\text{MS}}$ and pole mass known up to 4-loop QCD
 - Indicates the size of higher-order corrections to $m_{t,\text{pole}}$ beyond NNLO (2-loop): about 250 MeV

Marquard et al., PRL 114 (2015) 142002



$$m_t^{\text{pole}}(k) = m_t^{\overline{\text{MS}}}(\mu) \left[1 + \sum_{n=1}^k c_n \left(\frac{\mu}{m_t^{\overline{\text{MS}}}(\mu)} \right) \alpha_S^n(\mu) \right]$$

Resulting m_t

Well-defined m_t :

- **Without assuming** any relation to $m_t(\text{MC})$
- **Higher precision** than accounting for slope (CMS/ATLAS/Tevatron)
- Consistently lower for ABM
- About 1 GeV difference between directly measured and converted pole mass
→ **sizeable corrections beyond NNLO**

	$\alpha_s(M_Z)$	\bar{m}_t [GeV]	m_t^p [GeV]	$m_t^{p,c}$ [GeV]
ABM12	0.113	$158.4 \pm_{1.9}^{1.2}$	$166.6 \pm_{1.9}^{1.6}$	$168.0 \pm_{2.1}^{1.3}$
NNPDF3.0	0.118	$165.2 \pm_{1.7}^{1.1}$	$174.0 \pm_{1.7}^{1.4}$	$175.1 \pm_{1.9}^{1.2}$
MMHT2014	0.118	$165.4 \pm_{1.9}^{1.1}$	$174.3 \pm_{1.8}^{1.4}$	$175.3 \pm_{2.1}^{1.3}$
CT14	0.118	$165.5 \pm_{2.0}^{1.5}$	$174.4 \pm_{2.0}^{1.8}$	$175.4 \pm_{2.2}^{1.7}$

JK,KL,SM, PRL 116 (2016) 162001

	$\bar{\Delta}_m$ [GeV]	Δ_m^p [GeV]	$\Delta_m^{p,c}$ [GeV]
ABM12	$-14.3 \pm_{2.0}^{1.4}$	$-6.1 \pm_{2.0}^{1.7}$	$-4.7 \pm_{2.2}^{1.5}$
NNPDF3.0	$-7.6 \pm_{1.9}^{1.3}$	$1.3 \pm_{1.9}^{1.6}$	$2.4 \pm_{2.0}^{1.5}$
MMHT2014	$-7.3 \pm_{2.1}^{1.3}$	$1.5 \pm_{2.0}^{1.6}$	$2.6 \pm_{2.2}^{1.5}$
CT14	$-7.2 \pm_{2.1}^{1.7}$	$1.6 \pm_{2.1}^{1.9}$	$2.7 \pm_{2.3}^{1.8}$

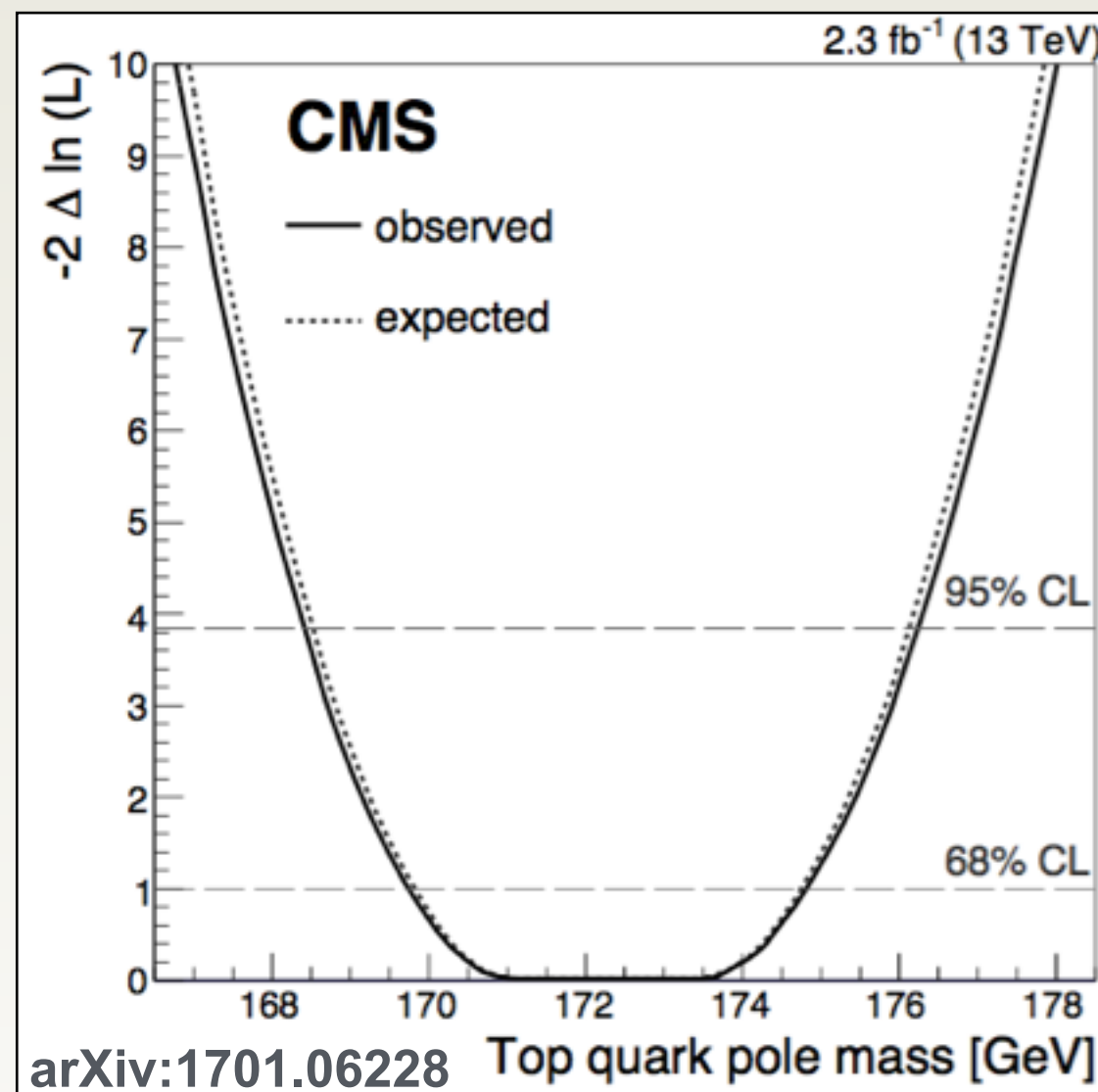
$m_t - m_t(\text{MC})$:

- **Directly** measurable
- **First** consistent experimental calibration
- Precision ~ 2 GeV
- **Consistent** with assumption of $m_t - m_t(\text{MC}) \approx 1$ GeV for most PDF sets

Strategy

- Measure σ_{tt} in l+jets channel
 - Simultaneous nuisance parameter fit of cross sections and MC mass (m_{lb})
- Incorporate likelihood for NNLO prediction
 - Model scale variations with box prior
- Determine m_t from joint likelihood measured \otimes predicted

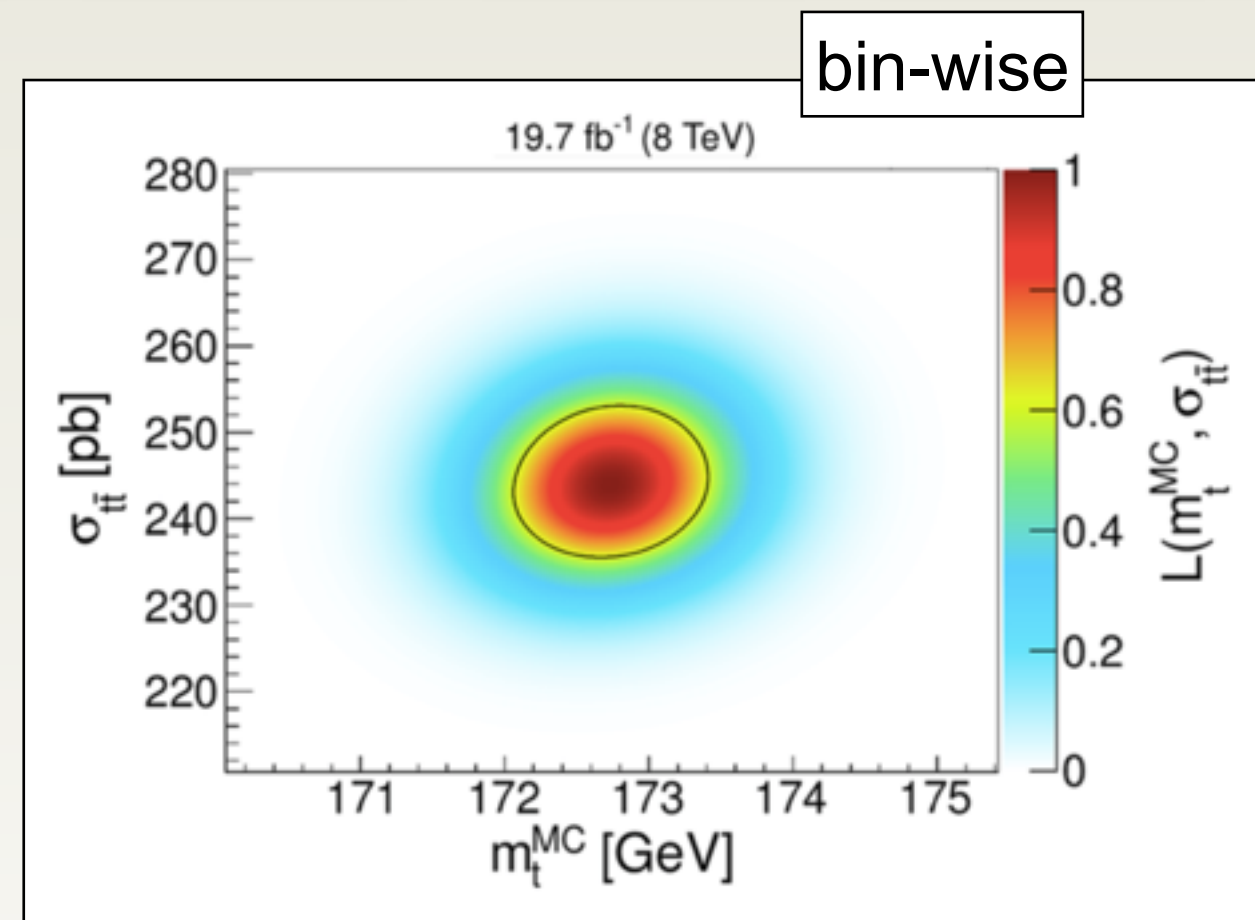
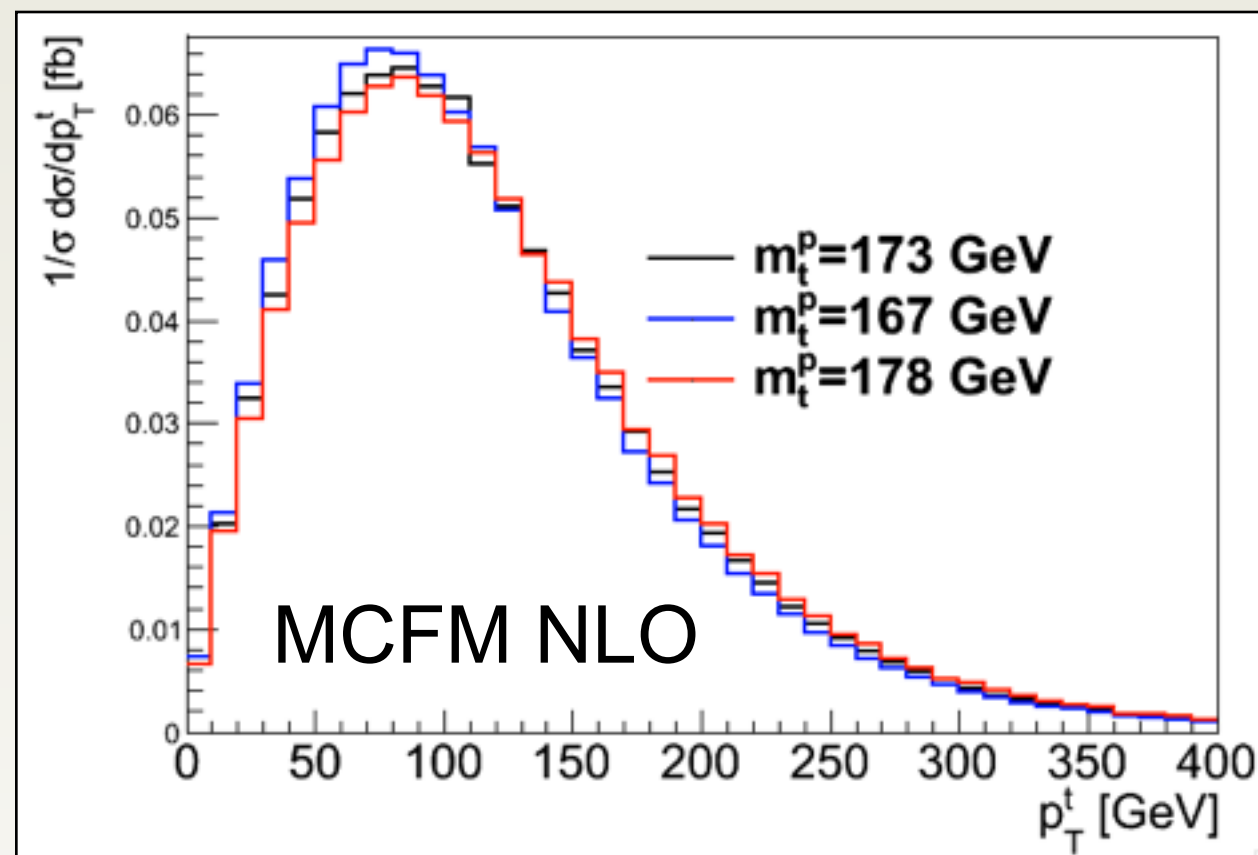
$$m_t(\text{pole}) = 172.7^{+2.4}_{-2.7} \text{ GeV} \\ (\text{CT14}, \alpha_s = 0.118)$$



Source	Δm_t [GeV]
Uncertainties from the fit in the fiducial region	$+2.1 / -2.0$
Extrapolation to the full phase space	$+0.7 / -1.1$
Beam energy	$+0.5 / -0.8$
μ_R / μ_F and PDF+ α_s	$+0.9 / -1.1$
Total	$+2.4 / -2.7$

← 2x larger than for dilepton measurement (jets in final state)

Extend to Differential Distributions



- Residual dependence on MC mass can be absorbed while unfolding similarly as for the inclusive cross section

JK,KL,SM, PRL 116 (2016) 162001

- Measurement of $d\sigma/dX$ (m_t)
- Could be used for **simultaneous parameter extraction** through direct comparison: α_s , pole mass / $\overline{\text{MS}}$ mass, ...
- Will likely provide higher precision

Summary

- Important to clearly define **which top-quark mass** is measured
- More and more precise direct pole-mass measurements
 - Using **inclusive and differential** cross sections
 - @ NNLO: down to ~ 2 GeV uncertainty in a single measurement
- Consistent way of mitigating $m_t(\text{MC})$ dependence in (cross-section) measurements
 - **Improves physics interpretation** of measured quantity
 - Allows to **extract any mass** in a well-defined scheme from direct comparison
 - Offers possibility to **measure relation** between m_t and $m_t(\text{MC})$ fully consistently
 - Precision is likely to increase when extending to differential measurements