

# Test of proton intrinsic charm in $\gamma(Z)+c(b)$ production at LHC



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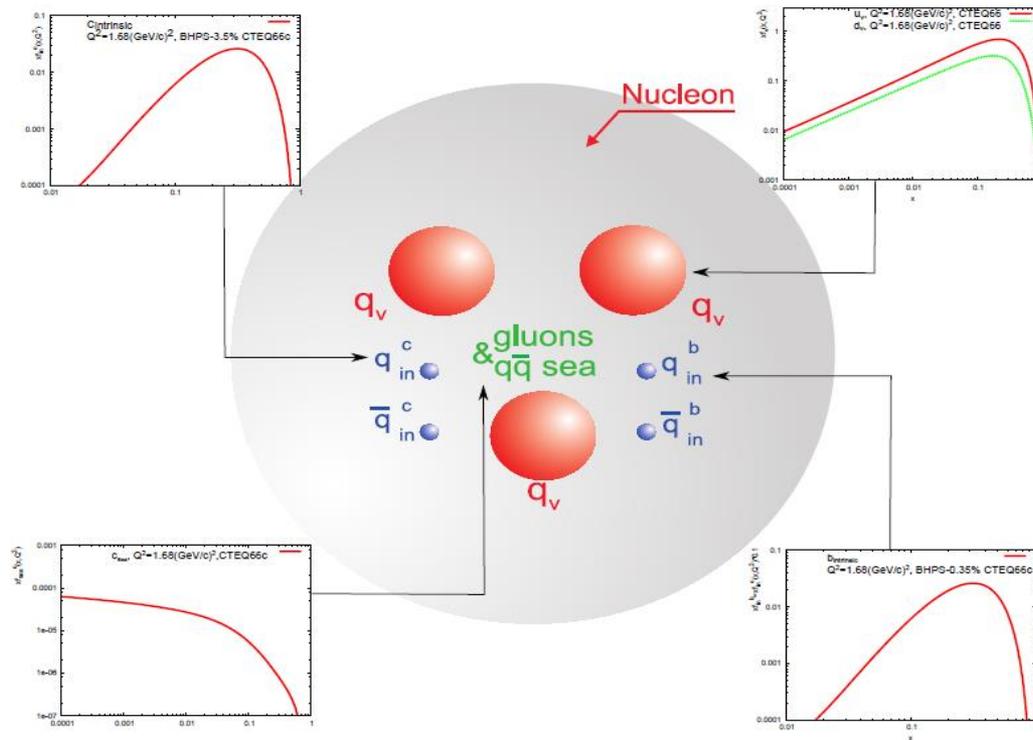


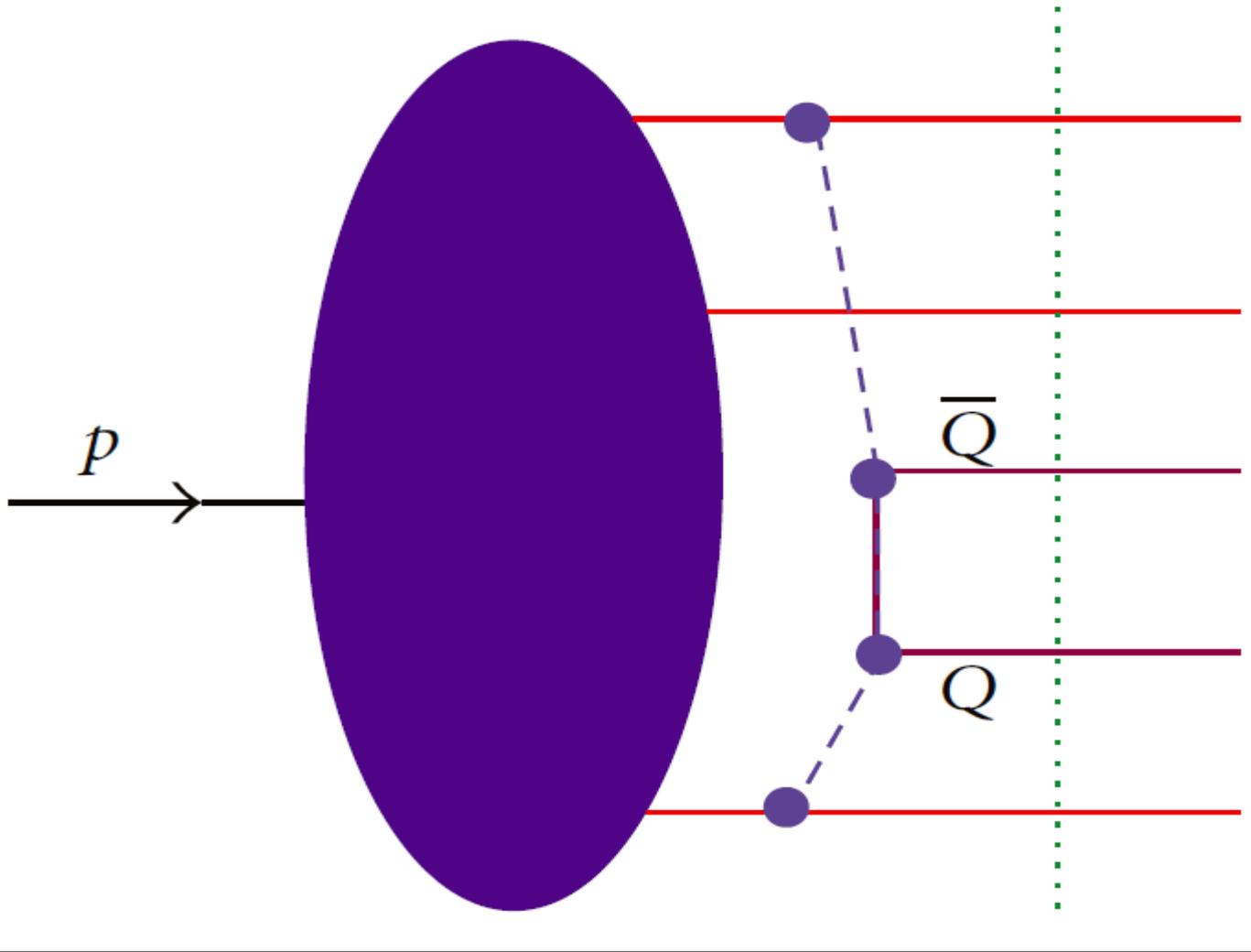
# ***OUTLINE***

- 1. Intrinsic heavy flavours in proton**
- 2. Main goal of our study**
- 3. Search for intrinsic charm (*IC*) in  $p_T$  -spectra  
the  $\gamma/Z/W+c(b)$  production in **p-p** collision**
- 4. Ratio  $\sigma(\gamma/Z/W+c)/\sigma(\gamma/Z/W+c)$  as a *significant*  
signal of the *IC contribution in proton***
- 5. Summaruy**

**BHPS** model: S.J. Brodsky, P. Hoyer, C. Peterson and N. Sakai, Phys.Lett.B9(1980) 451; S.J. Brodsky, S.J. Peterson and N. Sakai, Phys.Rev. D23 (1981) 2745.

## Intrinsic $Q\bar{Q}$ in proton





**Cut gluon-gluon scattering box diagram  $gg \rightarrow Q \bar{Q}$  inserted into the proton self-energy**

$$dP = N \prod_{j=1}^5 \frac{dx_j}{x_j} \delta \left( 1 - \sum_{j=1}^5 x_j \right) \prod_{j=1}^5 d^2 p_{Tj} \delta^{(2)} \left( \sum_{j=1}^5 p_{Tj} \right) \frac{F^2(s)}{(s - m_N^2)^2},$$

where

$$s = \sum_{j=1}^5 \frac{p_{Tj}^2 + m_j^2}{x_j},$$

$$P(x) = \frac{N x^2}{6(1 - \alpha)^5} \left( \phi_1(x) + \phi_2(x) [\ln(x) - \ln[1 - c(1 - x)x]] \right),$$

where  $x = x_5$ ,  $c = m_N^2/m_c^2$ ,

$$\phi_1(x) = (1 - x)(1 - \alpha) \left[ 1 + x \left[ 10 + x - c(1 - x) \left( x(10 - c(1 - x)) + 2 \right) \right] \right],$$

and

$$\phi_2(x) = 6x [1 + x(1 - c(1 - x))] [1 - c(1 - x)x].$$

Here  $N$  is found from the normalization equation:

$$\int_0^1 P(x) dx = w,$$

where  $w$  is the integral fraction of the intrinsic charm. Setting  $c \rightarrow 0$  leads to the BHPS result |

$$P(x) = 600wx^2 \left[ 6x(1 + x) \ln x + (1 - x)(1 + 10x + x^2) \right].$$

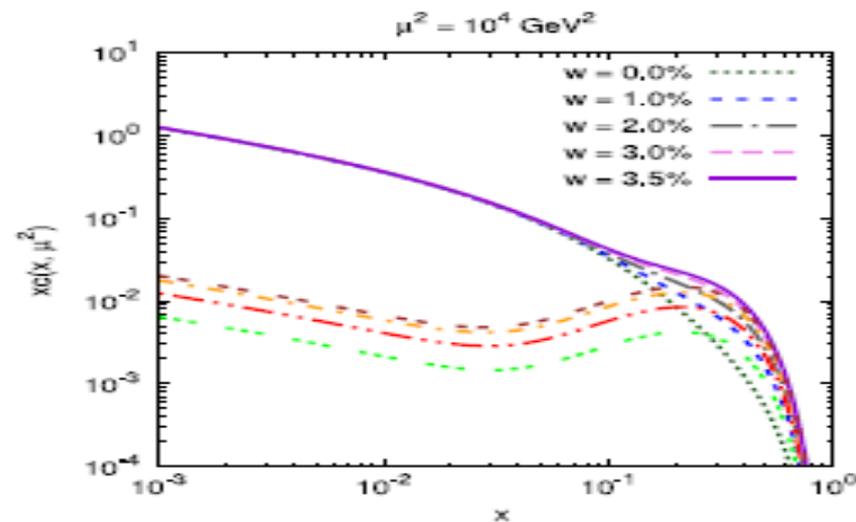
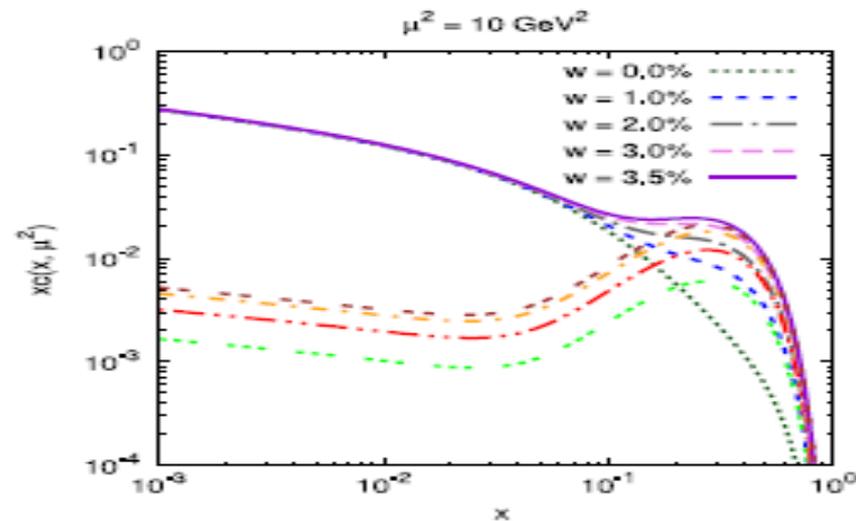
$$F^2 = \exp \left( \frac{-(s - m_N^2)}{\Lambda^2} \right)$$

or the *power-law suppression* factor:

$$F^2 = \frac{1}{(s + \Lambda^2)^n}.$$

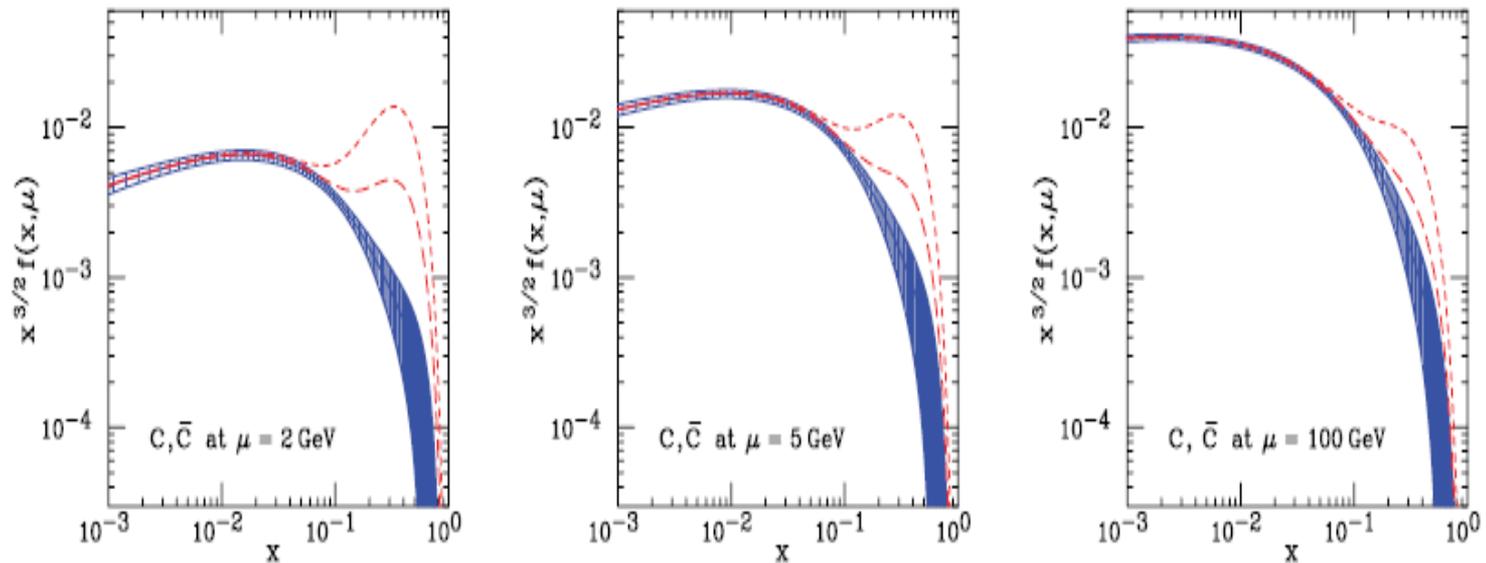
# Intrinsic charm density in a proton as a function of IC probability $w$

$$XC(x, \mu_0^2) = XC_{\text{ex}}(x, \mu_0^2) + XC_{\text{in}}(x, \mu_0^2).$$



$$XC_{\text{in}}(x, \mu^2) = \frac{w}{w_{\text{max}}} XC_{\text{in}}(x, \mu^2) \Big|_{w=0}^{w=w_{\text{max}}}$$

## CHARM QUARK DISTRIBUTIONS IN PROTON



Charm quark distributions within the BHPS model. The three panels correspond to the renormalization scales  $\mu=2, 5, 100$  GeV respectively. The long-dashed and the short-dashed curves correspond to  $\langle x_{c\bar{c}} \rangle = 0.5\% / 2\%$  respectively using the PDF CTEQ66c. The solid curve and shaded region show the central value and uncertainty from CTEQ6.5, which contains no **IC**.

**There is an enhancement at  $x > 0.1$  due to the IC contribution**

*Main goal: searching for the signal of the intrinsic charm (IC) contribution in proton from the analysis of the prompt photon or Z/W boson production in p-p collision accompanied by heavy c(b)-jet.*

**We have predictions on IC**

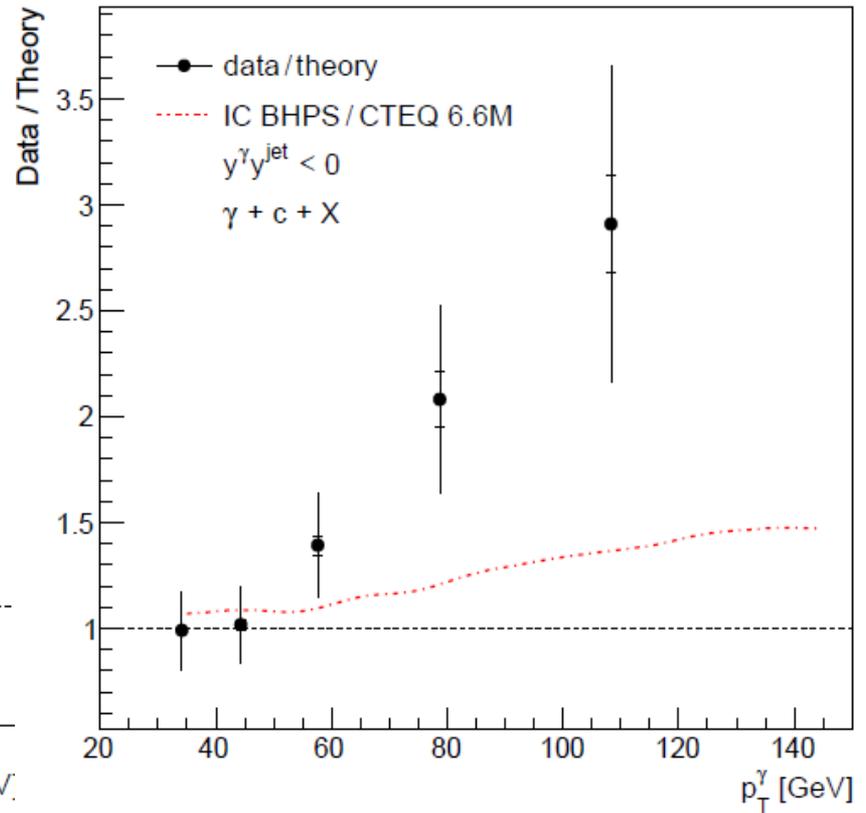
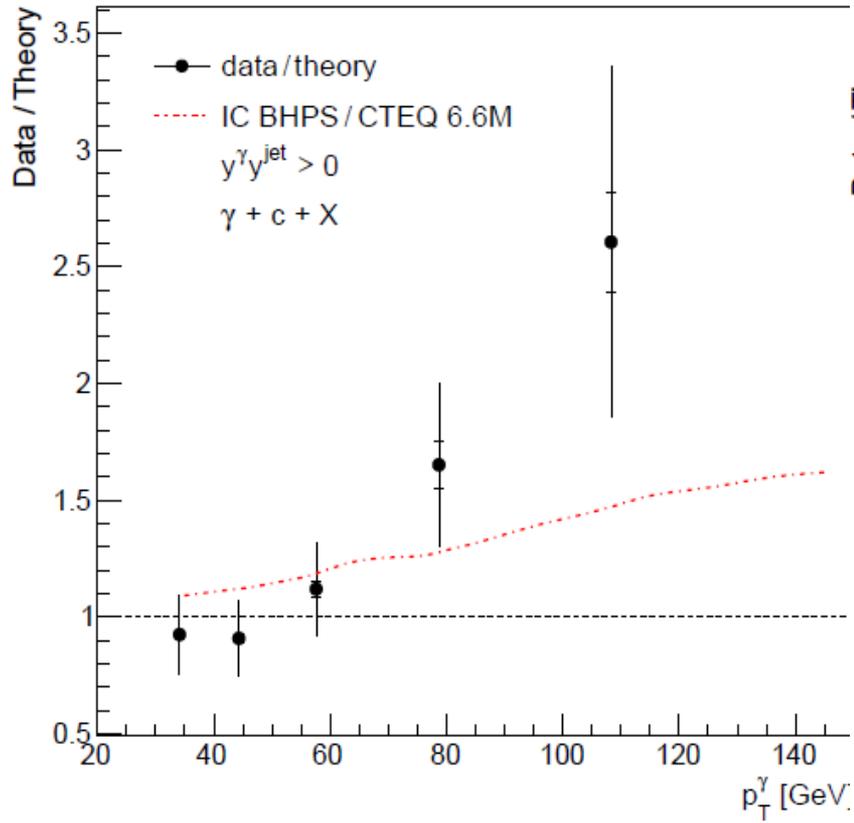
*PP- $\rightarrow\gamma + c + X$  : V.A.Bednyakov, M.A.Demichev, G.L., T.Stavreva, M.Stockton, Phys.Lett. B728, 602 (1914)*

*PP- $\rightarrow Z/W + c(b) + X$  : H.Beauchemin, V.A.Bednyakov, G.L., Yu. Yu. Stepanenko, Phys.Rev.D92, 034014 (2015)*

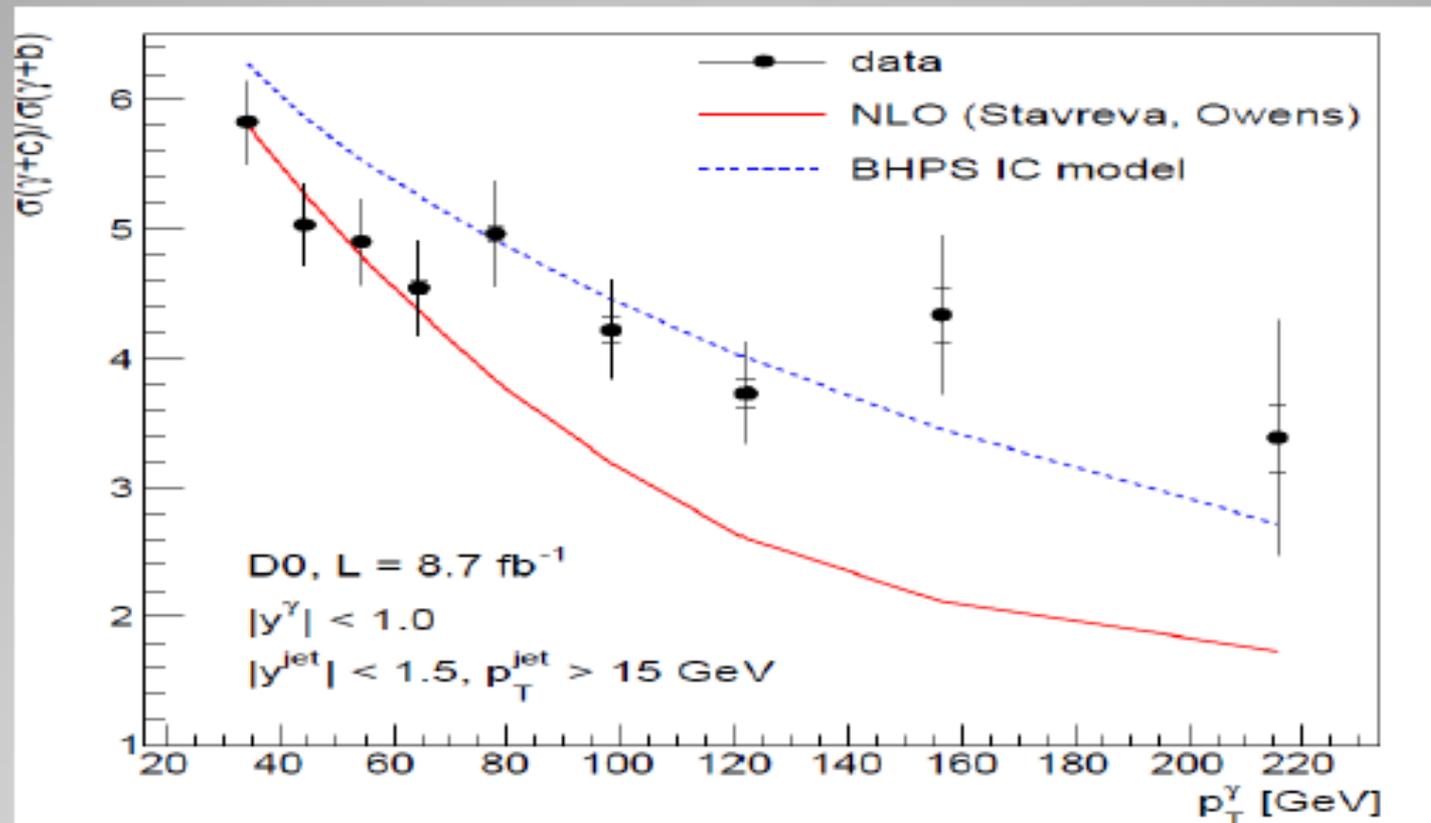
*PP- $\rightarrow\gamma/Z + c(b) + X$  : A. Lipatov, G.L., Yu. Stepanenko V.A.Bednyakov, Phys.Rev.D94, 05301 (2016)*

*Collider tests of heavy PDF: S.J.Brodsky, V.A.Bednyakov, G.L., S.Tokar, J.Smiesko, Prog. In Part. Nucl.Phys., 93, 108, (2017).*

$pp \rightarrow \gamma + c(b) + X$  D0 experiment at Tevatron  $s^{1/2} = 1.96 \text{ TeV}$

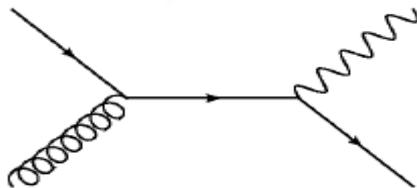


$R = \sigma(\gamma + c) / \sigma(\gamma + b)$  for  $p \bar{p} \rightarrow \gamma + Q$  at  $s^{1/2} = 1.98$  TeV

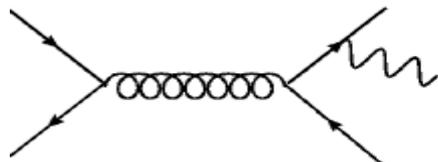


$p_T$  -distribution of  $R$ , points are the D0 data; red solid line is NLO without *IC* ; short dash line is BHPS with *IC* probability about 1 %

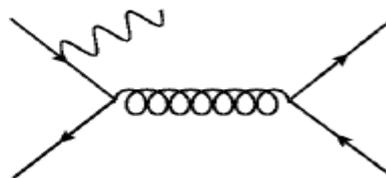
$$pp \rightarrow \gamma + Q + X, Q = c, b$$



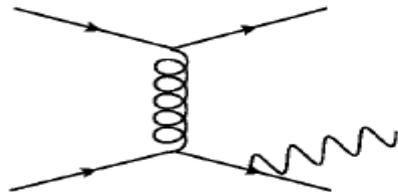
a)



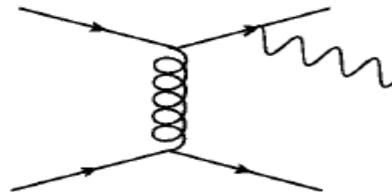
b)



c)



d)

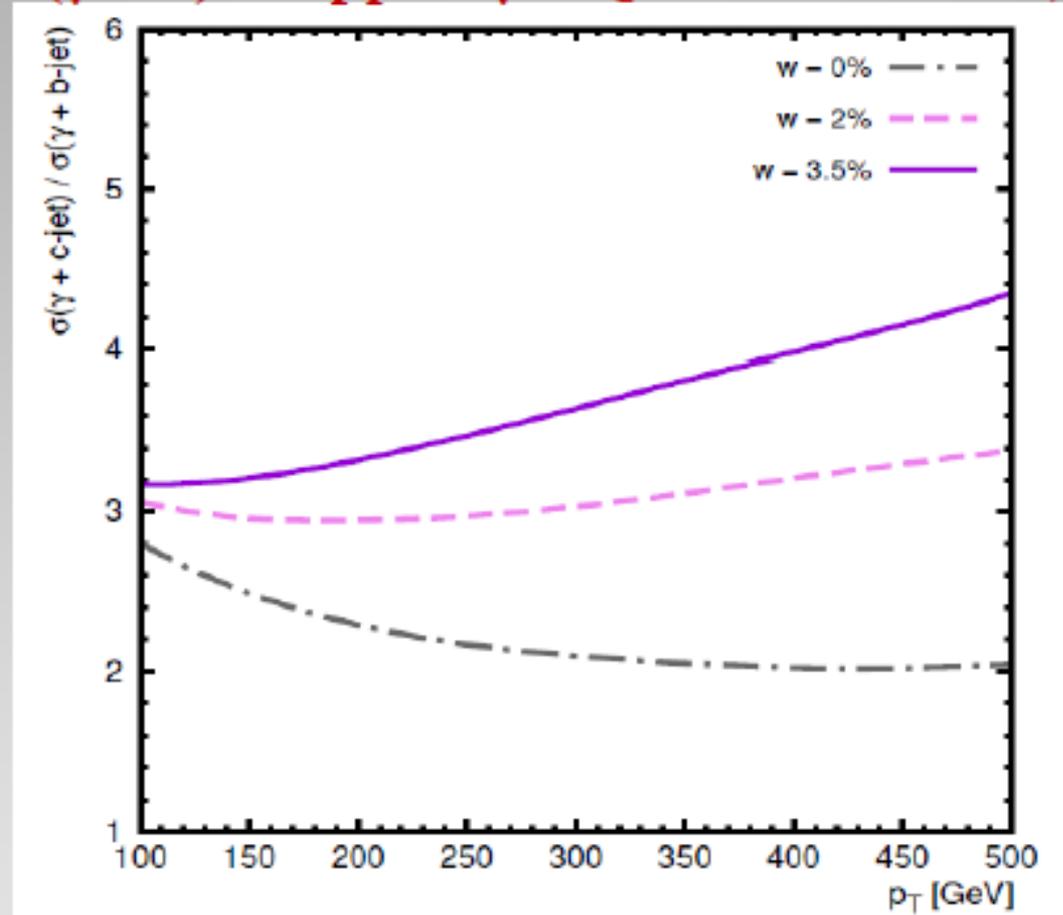


e)

**Feynman QCD diagrams: a):  $Q + g \rightarrow \gamma + Q$ ;**

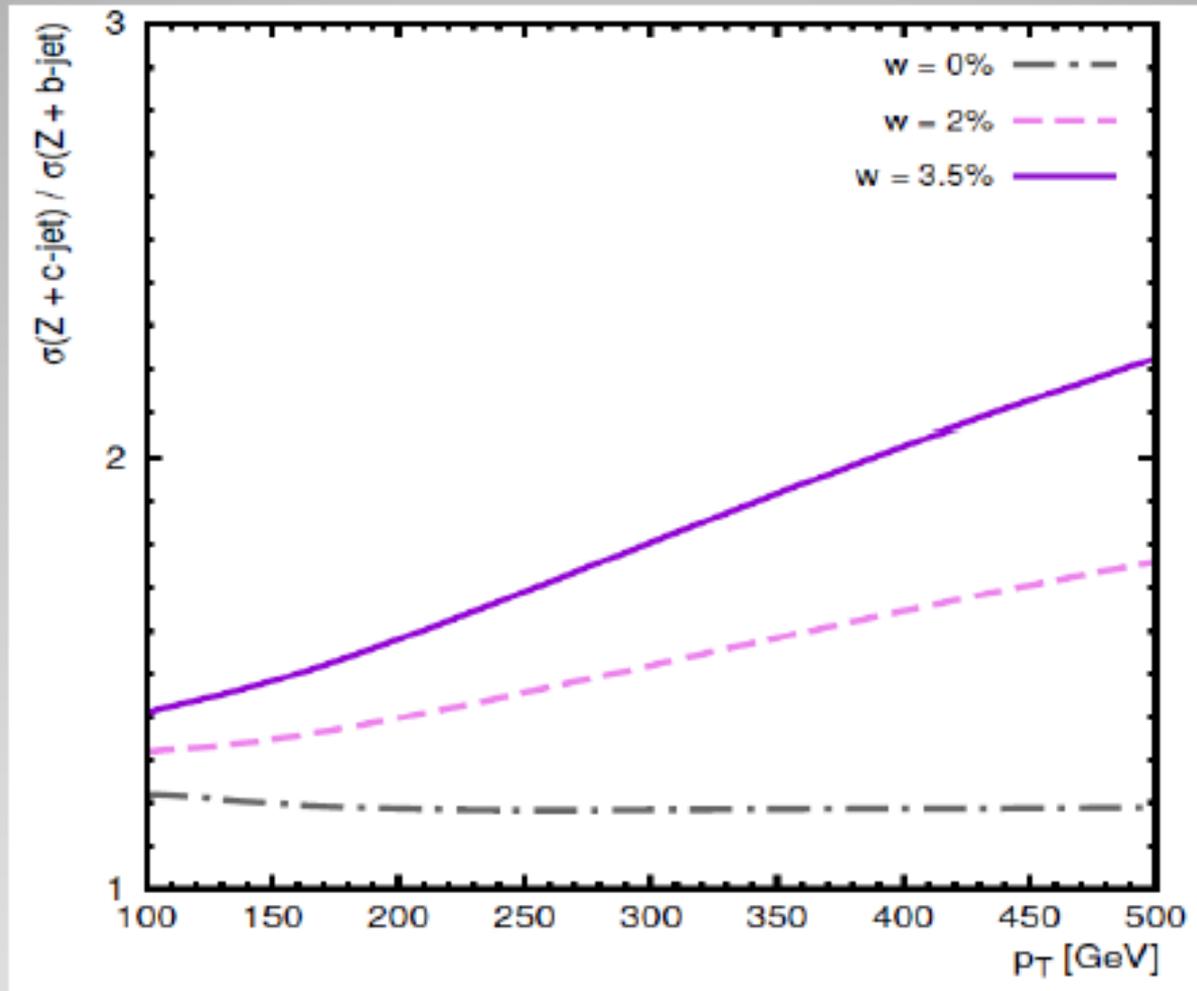
**b-c):  $Q + \bar{Q} \rightarrow Q + \bar{Q} + \gamma$ ; d-e):  $Q(q) + q(Q) \rightarrow Q(q) + q(Q) + \gamma$**

$R = \sigma(\gamma + c) / \sigma(\gamma + b)$  for  $pp \rightarrow \gamma + Q$  at  $s^{1/2} = 8 \text{ TeV}$ ;  $1.5 < \eta < 2.4$



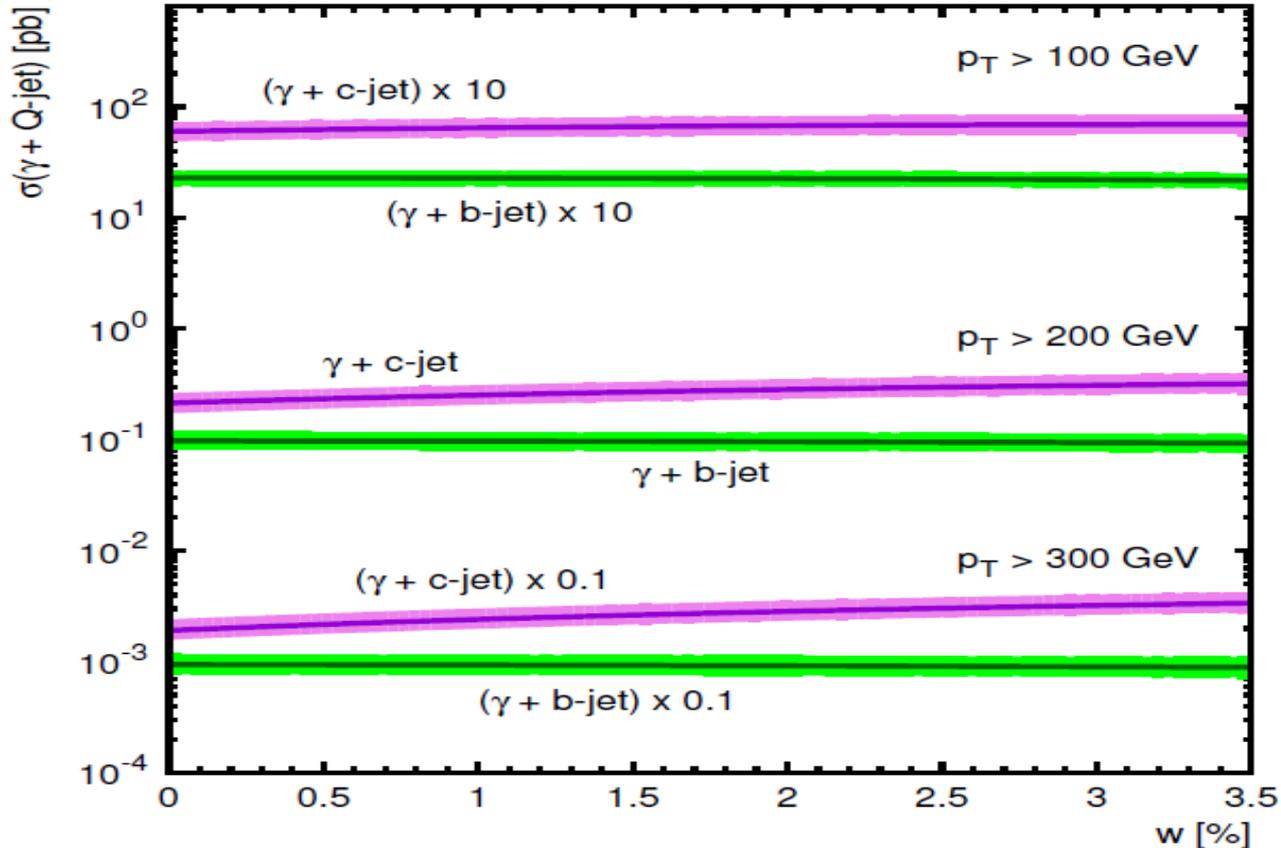
$p_T$  -distribution of  $R$  at different IC probability  $w$

$R = \sigma(Z + c) / \sigma(Z + b)$  for  $pp \rightarrow Z + Q$  at  $s^{1/2} = 8 \text{ TeV}$ ;  $1.5 < \eta < 2.4$



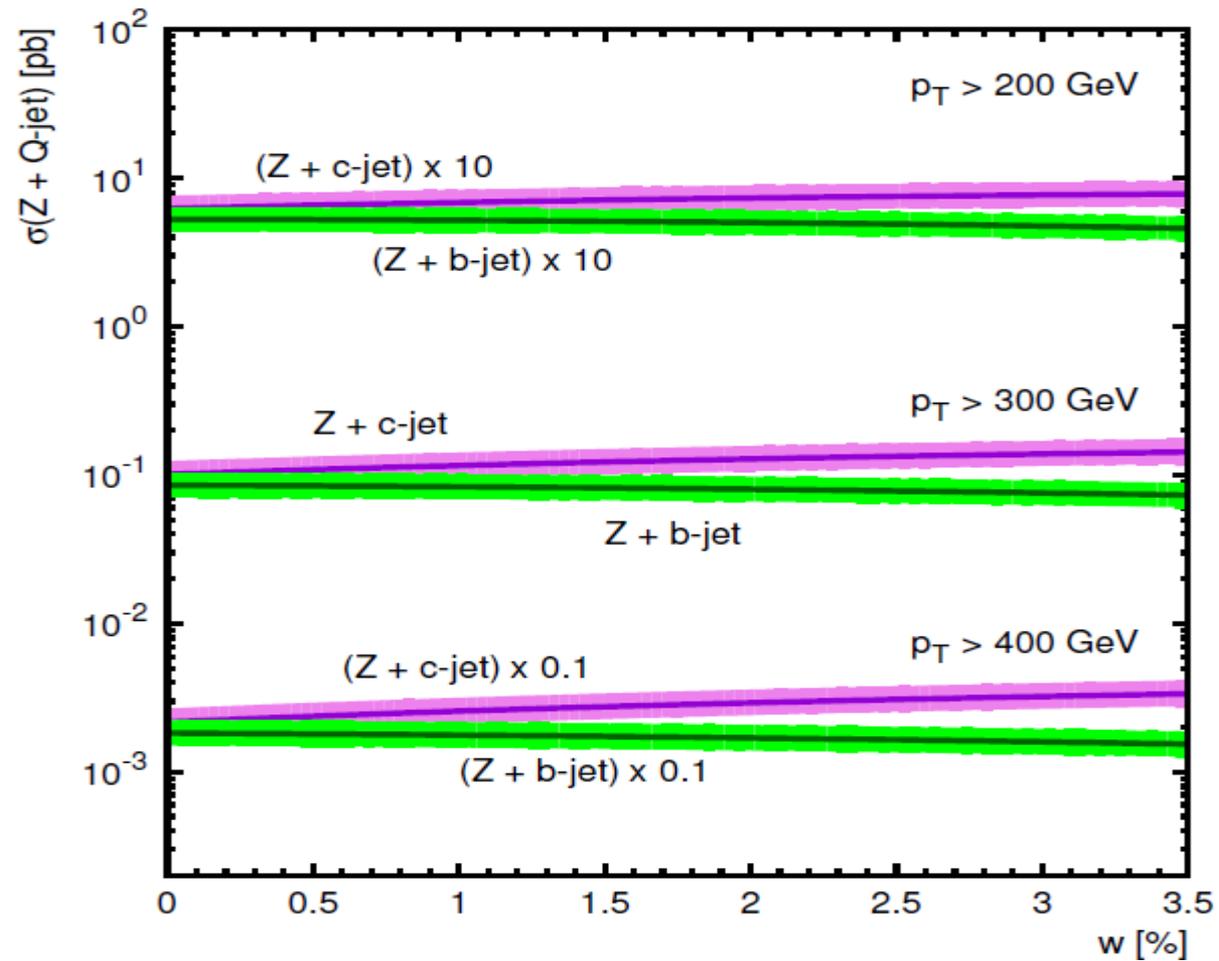
$p_T$  –distribution of  $R$  at different IC probability  $w$

$\sigma(\gamma + Q)$  at  $s^{1/2} = 8 \text{ TeV}$  ;  $1.5 < \eta < 2.4$ ;  $Q = c, b$



$p_T^\gamma$  – spectrum integrated over  $p_T^\gamma$ , i.e.,  $\sigma(\gamma+c)$  and  $\sigma(\gamma+b)$  at  $p_T^\gamma > 100 \text{ GeV}$  or  $200 \text{ GeV}$ , or  $300 \text{ GeV}$ , vs. *IC* probability  $w$

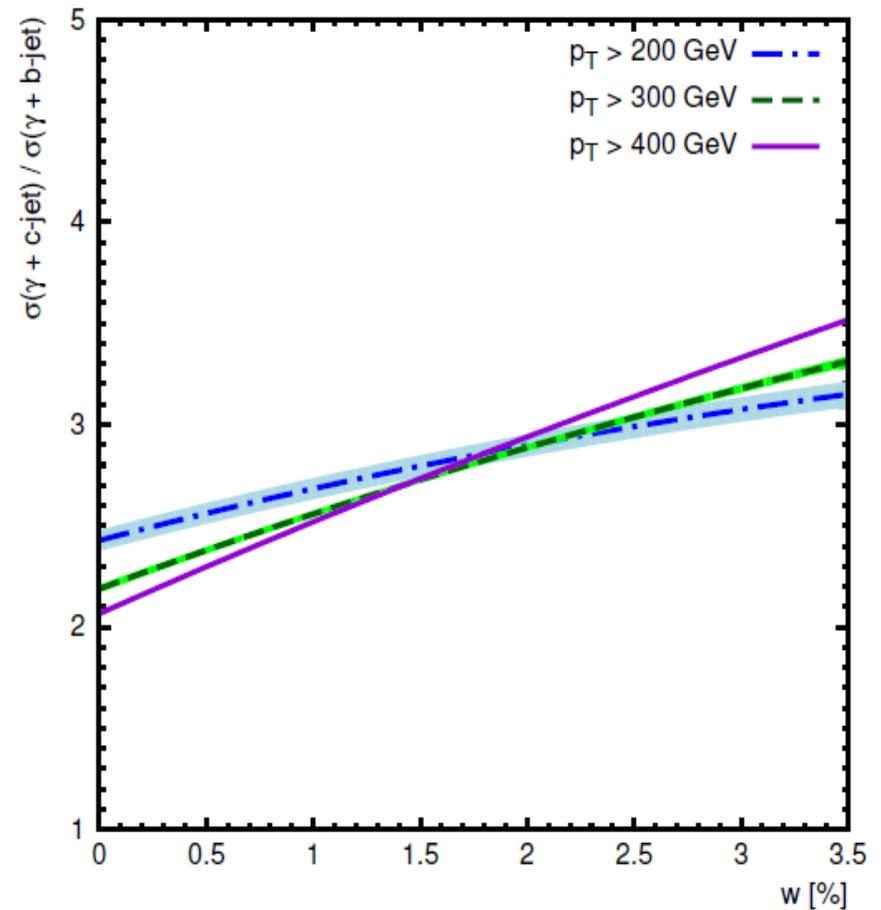
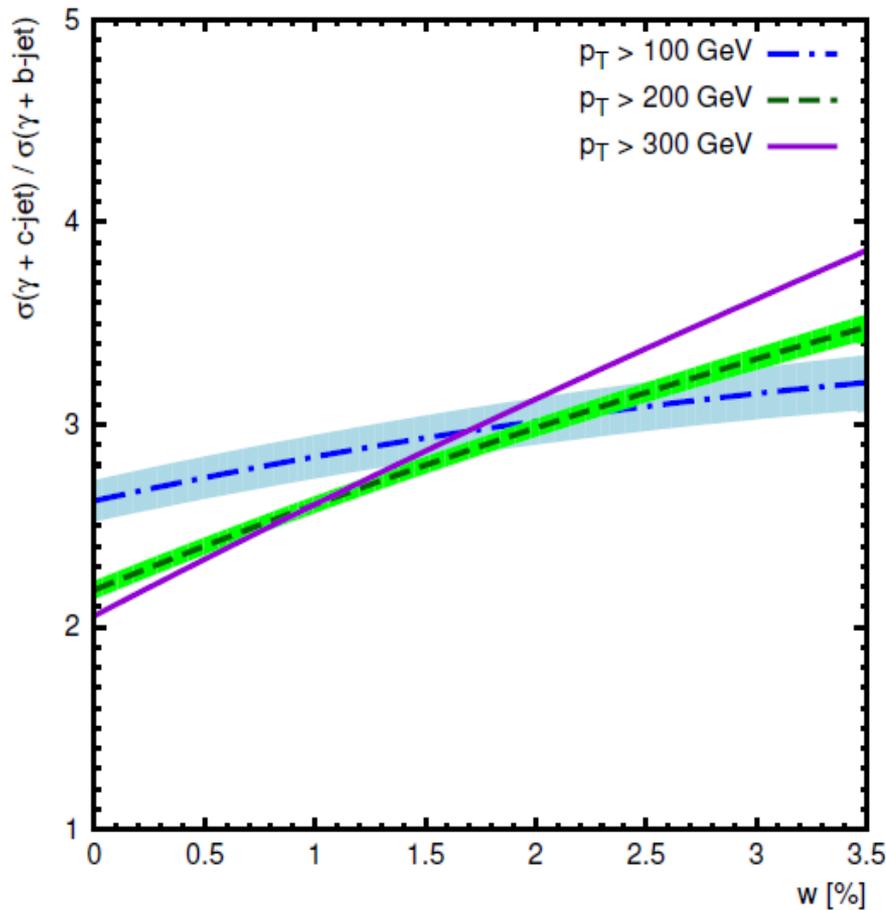
$\sigma(Z + Q)$  at  $s^{1/2} = 13 \text{ TeV}$  ;  $1.5 < \eta < 2.4$ ;  $Q = c, b$



$s^{1/2} = 8 \text{ TeV}$

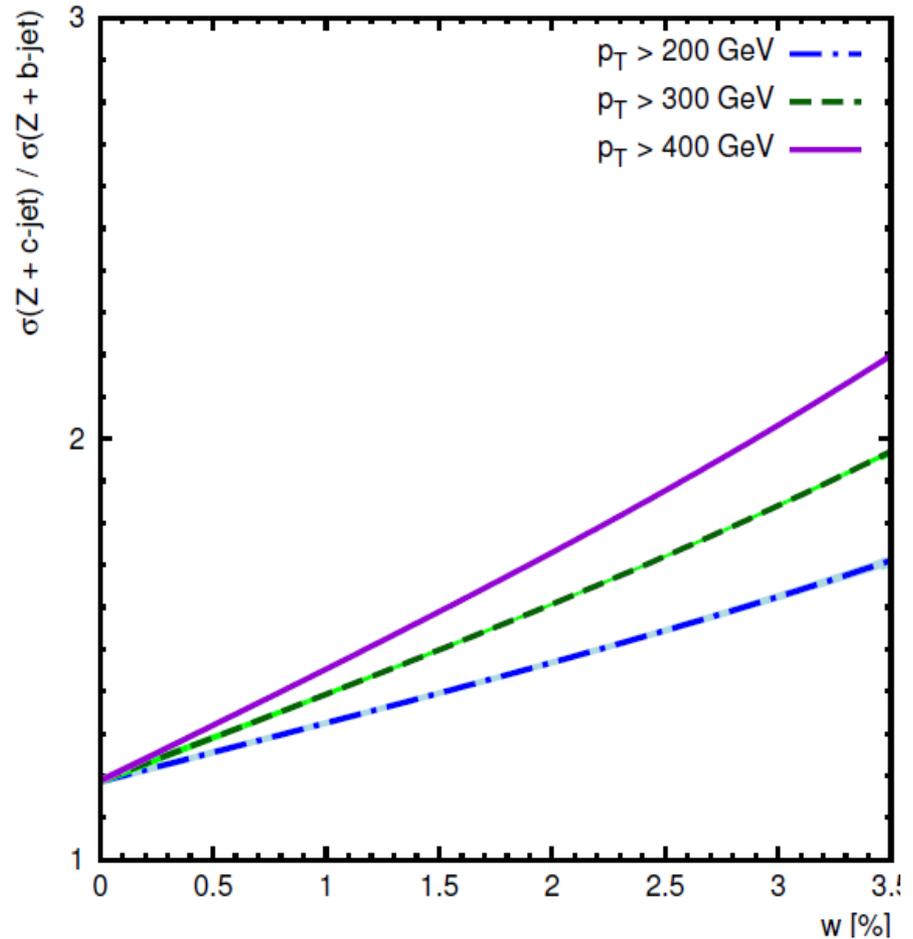
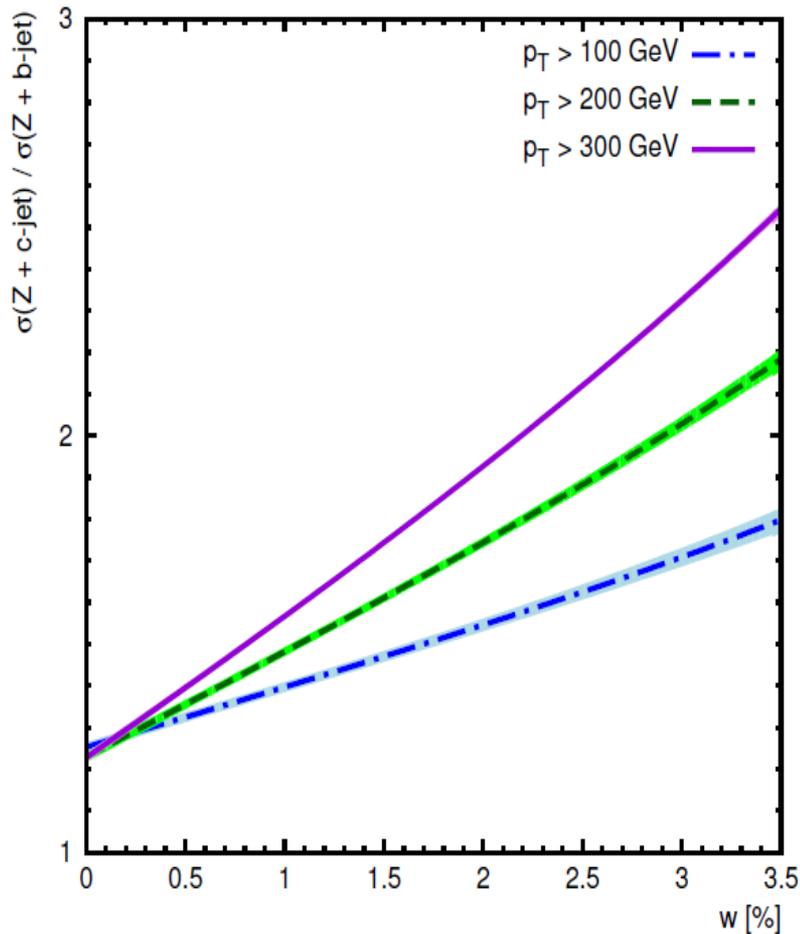
$PP \rightarrow \gamma + Q + X, Q = c, b$

$s^{1/2} = 13 \text{ TeV}$



Ratio between the x-sections of  $\gamma + c$  and  $\gamma + b$  production integrated over  $p_T$ . Bands mean the QCD scale uncertainty

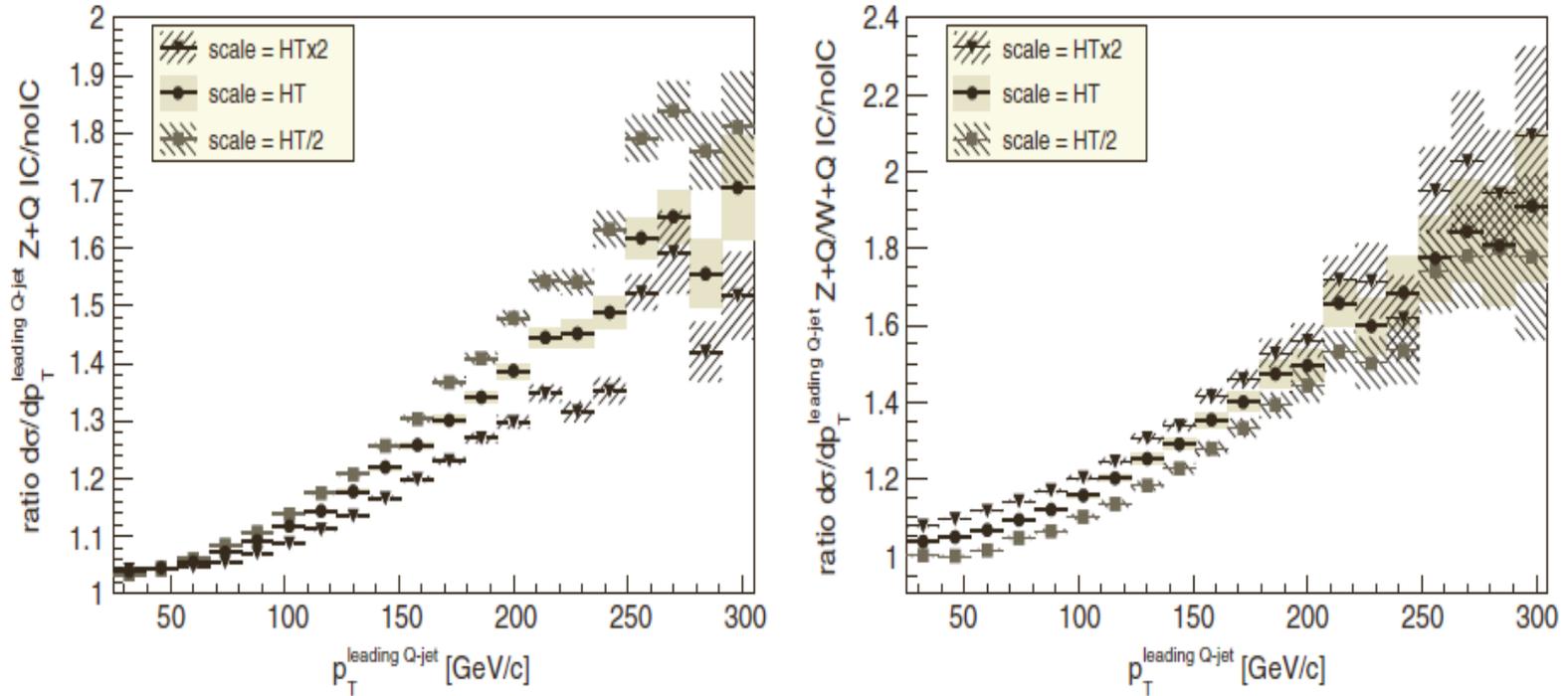
**Ratio  $\sigma(Z+c)/\sigma(Z+b)$**   
 **$S^{1/2} = 8$  TeV (left) and  $S^{1/2} = 13$  TeV (right)**



**Ratio between the x-sections of  $\gamma + c$  and  $\gamma + b$  production**  
**integrated over  $p_T$ . Bands mean the QCD scale uncertainty**

**A.V.Lipatov, G.L., Yu.Yu.Stepanenko, V.A.Bednyakov, Phys.Rev. D94 , 053011 (2016) ;**  
**S.J.Brodsky, V.A.Bednyakov, G.L., J.Smiesko, S.Tokar, arXiv:1612.01351 ,**  
**Prog. Part.Nucl.Phys.,v. 93, p.108 (2017)**

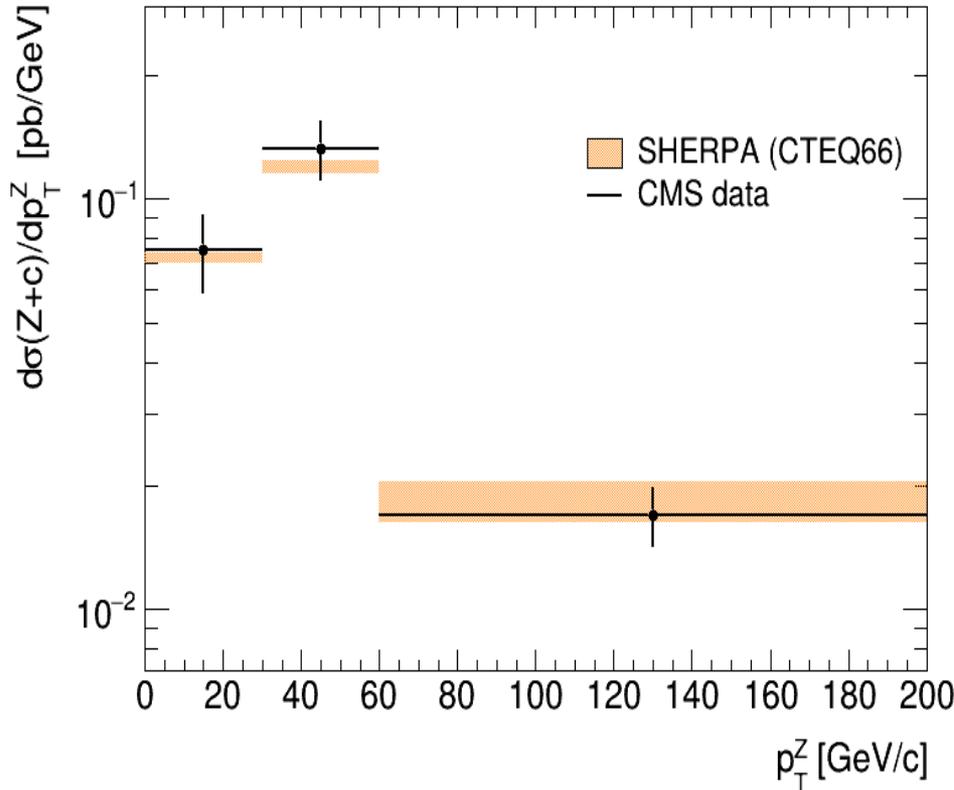
## Scale uncertainty for Z+Q and Z+Q/W+Q



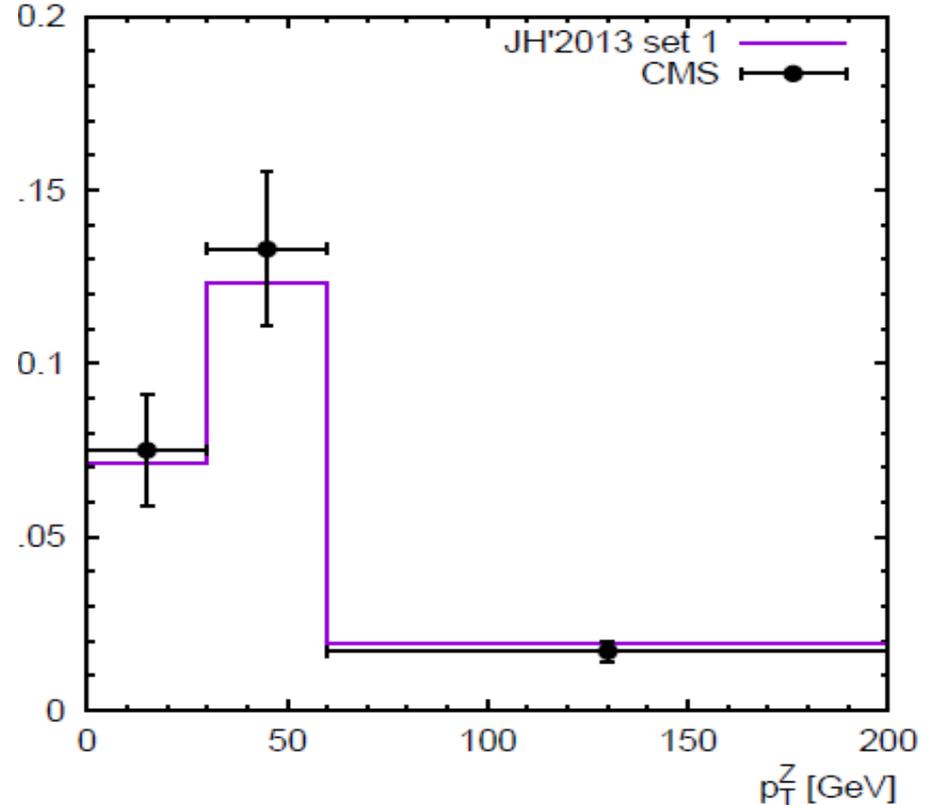
**Left: Z+Q with IC and without IC at different scales**

**Right: Z+Q/W+Q with IC and without IC at different scales**

**$pp \rightarrow Z+c\text{-jet} + X$  at  $s^{1/2} = 8$  TeV**  
 **$p_T^{\text{jet}} > 25$  GeV/c;  $|\eta^{\text{jet}}| < 2.5$**



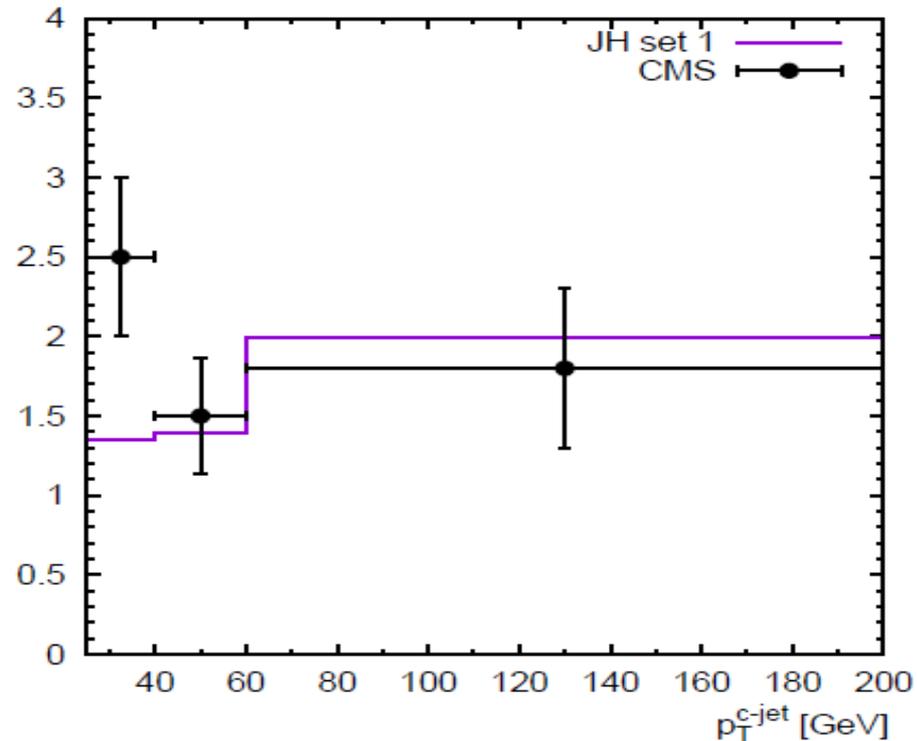
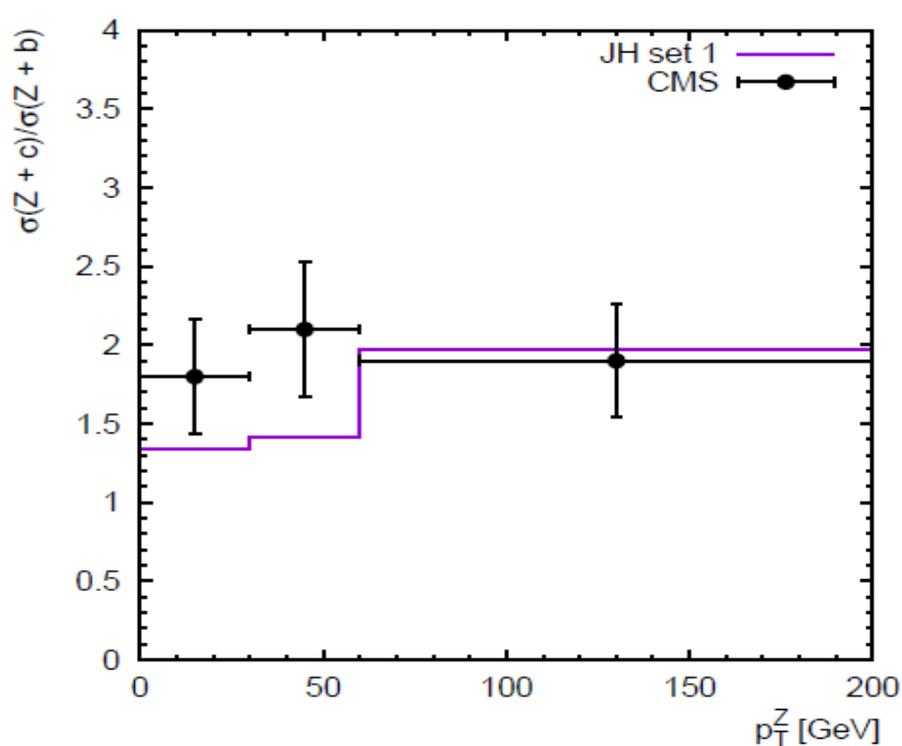
**Left: SHERPA results**



**Right: QCD calculation done by A.V.Lipatov, M.A.Malyshev, S.P. Baranov**

**Experimental data are taken from CMS PAS SMP-15-009, CMS-SMP-14-10**

**$pp \rightarrow Z+c(b)\text{-jet} + X$  at  $s^{1/2} = 8$  TeV**  
 **$p_T^{\text{jet}} > 25$  GeV/c;  $|\eta^{\text{jet}}| < 2.5$**



**Ratio  $\sigma(Z+c)/\sigma(Z+b)$  as a function of  $p_T^Z$**

**QCD calculation done by A.Lipatov, S.Baranov, M.Malyshev**

**Experimental data are taken from CMS PAS SMP-15-009,**

**CMS-SMP-14-10**

# SUMMARY

1. QCD predicts two sources of heavy quarks: the standard small- $x$  **extrinsic contribution** at from gluon splitting  $g \rightarrow Q \bar{Q}$  and **intrinsic contribution at large  $x$** , which arises from the cut gluon-gluon scattering box.
2. The hypothesis of *intrinsic* quark components at large  $x$  was motivated by possible explanation of the large cross section for the forward open charm production in p-p at ISR.
3. However, the accuracy of such experimental data at large  $x$  does not provide precise constraints on the *IC* probability.
4. The production of prompt photons or gauge bosons accompanied by heavy jets (c,b) can provide an ideal method to verify the *IC* probability in proton.
5. The increase of  $p_T$  – spectrum of  $\gamma/Z/W$  or *c/b*-jets produced at large  $p_T$  and the forward rapidity region of ATLAS or CMS due the the *IC* enhancement in the PDF is predicted.
6. The ratio between the cross sections of the  $\gamma/Z + c$  production and  $\gamma/Z + b$  one can give a significant *IC* signal.

**THANK YOU VERY MUCH  
FOR YOUR ATTENTION !**

**BACK UP**

## PRODUCTION OF HEAVY FLAVOURS IN HARD P-P COLLISIONS

$$E \frac{d\sigma}{d^3p} = \sum_{i,j} \int d^2k_{iT} \int d^2k_{jT} \int_{x_i^{\min}}^1 dx_i \int_{x_j^{\min}}^1 dx_j f_i(x_i, k_{iT}) f_j(x_j, k_{jT}) \frac{d\sigma_{ij}(\hat{s}, \hat{t})}{d\hat{t}} \frac{D_{i,j}^h(z_h)}{\pi z_h}$$

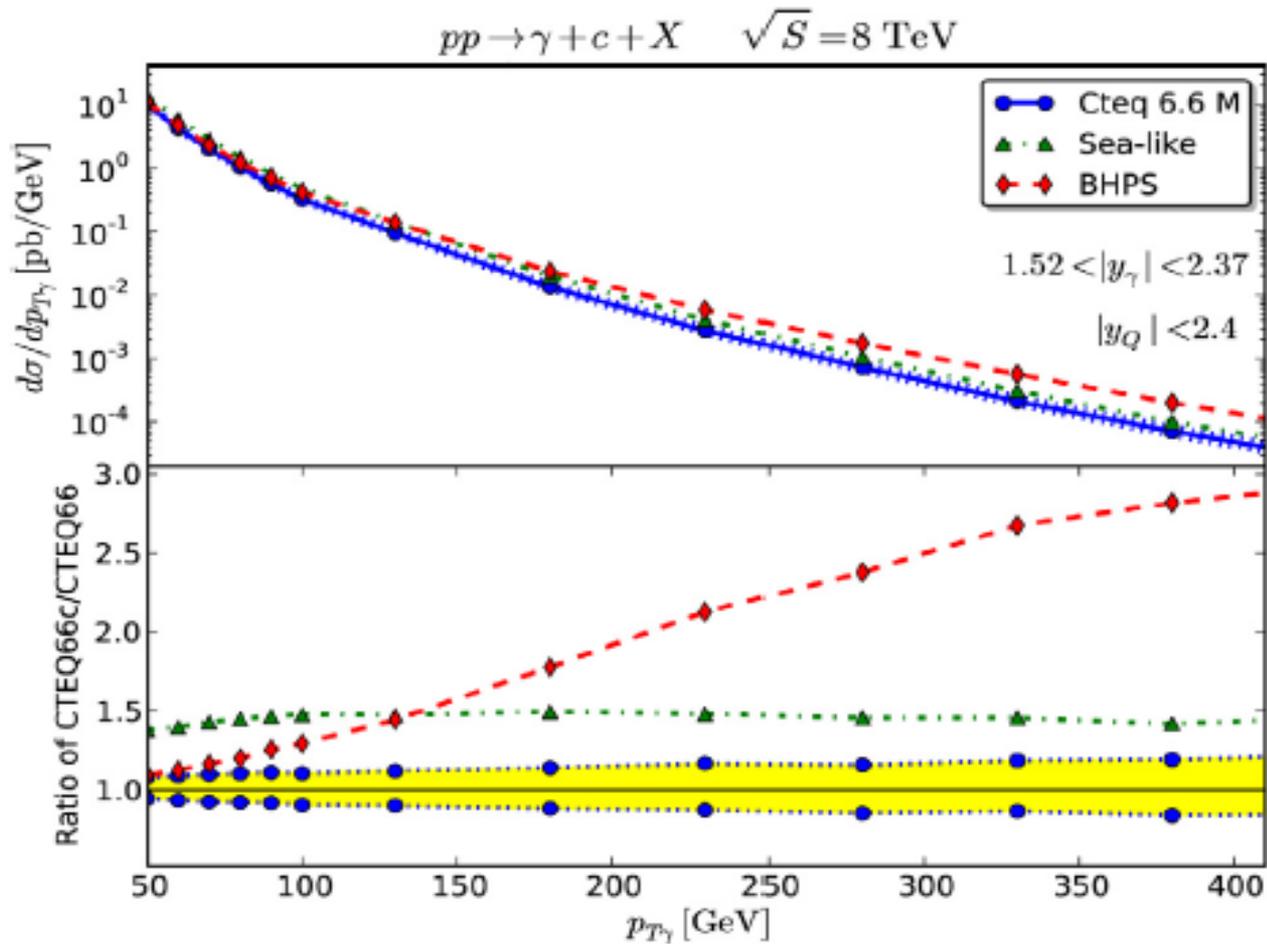
$$x_i^{\min} = \frac{x_T \cot(\frac{\theta}{2})}{2 - x_T \tan(\frac{\theta}{2})} \quad x_F \equiv \frac{2p_z}{\sqrt{s}} = \frac{2p_T}{\sqrt{s}} \frac{1}{\tan \theta} = \frac{2p_T}{\sqrt{s}} \sinh(\eta)$$

$$x_i^{\min} = \frac{x_R + x_F}{2 - (x_R - x_F)} \quad x_R = 2p/\sqrt{s}$$

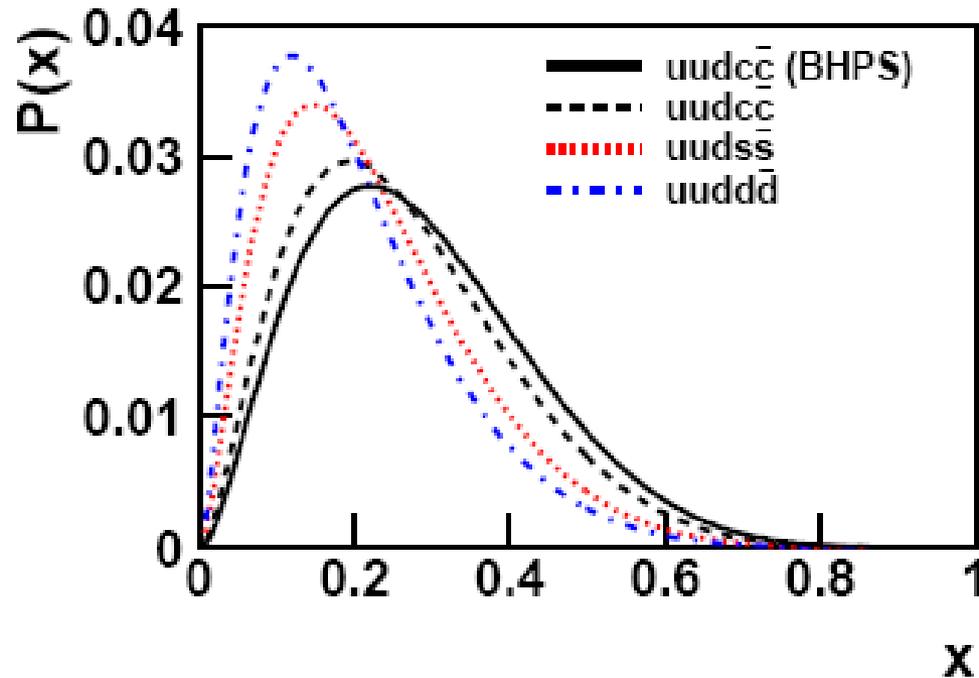
One can see that  $x_i \geq x_F$ . If  $x_F > 0.1$  then,  $x_i > 0.1$  and the **conventional sea** heavy quark (extrinsic) contributions are suppressed in comparison to the **intrinsic** ones.

$x_F$  is related to  $p_T$  and  $\eta$ . So, at certain values of these variables, in fact, there is **no conventional sea** heavy quark (**extrinsic**) contribution. And we can study the **IQ contributions** in hard processes at the **certain** kinematical region.

# Predictions for ATLAS and SMC

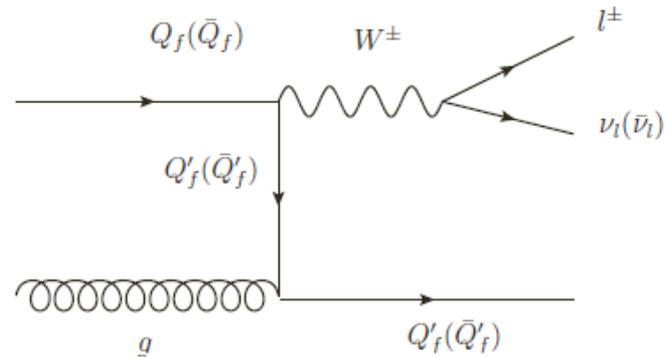


V.A. Bednyakov, M.A. Demichev, G.L. Stavreva, M. Stockton,  
Phys.Lett. B728, 602 (2014).

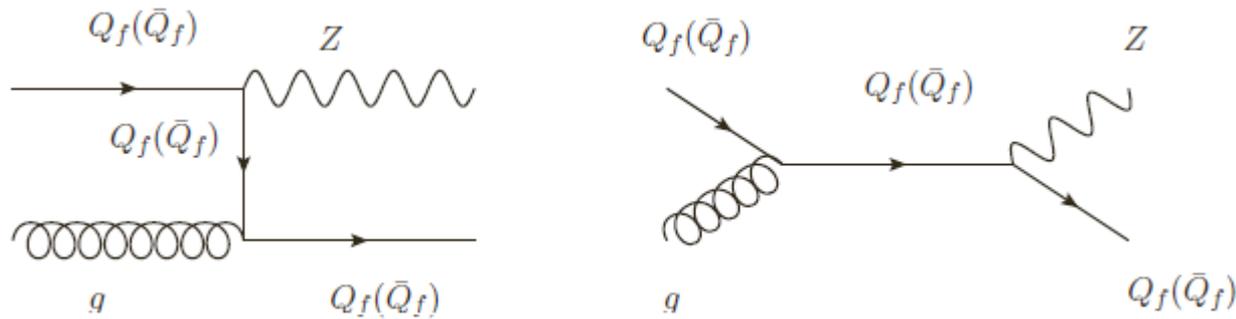


The  $x$ -distribution of the intrinsic  $Q$  calculated within the BHPS model. **There is an enhancement at  $x > 0.1$**   
 Jen-Chieh Peng & We-Chen Chang, hep-ph/1207.2193.

## pp $\rightarrow$ W/Z+heavy flavour jets



The LO Feynman diagrams for the process  $Q_f(\bar{Q}_f)g \rightarrow W^\pm Q'_f(\bar{Q}'_f)$ , where  $Q_f = c, b$  and  $Q'_f = b, c$  respectively.



Feynman diagram for the process  $Q_f(\bar{Q}_f)g \rightarrow ZQ_f(\bar{Q}_f)$

# PHOTON (DI-LEPTON) AND c(b)-JETS PRODUCTION IN P-P

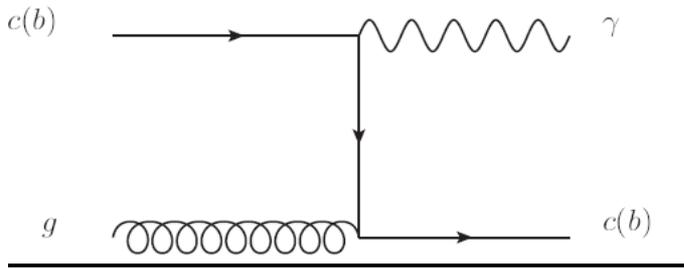


Fig.a. Feynman diagram for the process  $c(b) + g \rightarrow \gamma + c(b)$

$$x_F = \frac{2p_T}{s^{1/2}} \sinh(\eta); p_{T\gamma} = -p_{Tc}$$

for Fig.a

$$x_c \geq x_F > 0.1$$

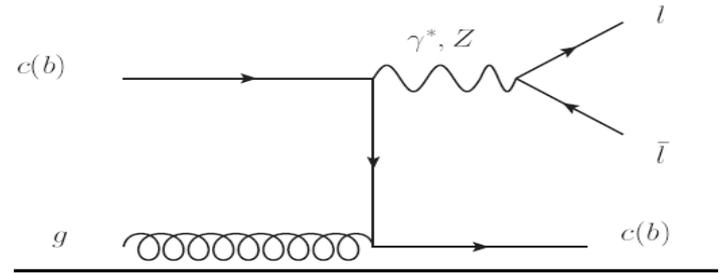


Fig.b. Feynman graph for the process  $c(b) + g \rightarrow \gamma/Z^0 + c(b)$

$$x_{c(b)} = \frac{m_{l+l^-}^2}{x_g s} + x_{c(b)}^f$$

To observe the IC

for Fig.b

$$x_{c(b)} = \frac{m_{l+l^-}^2}{x_g s} + x_{c(b)}^f > 0.$$