Parton shower and finite top mass effects in HH production



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In collaboration with

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Motivation

Dominant production mechanism for HH production is gluon fusion



Sensitive to HHH coupling

Test of Higgs potential & EW symmetry breaking

Most calculations are done in $m_t \rightarrow \infty$ limit (Higgs EFT)



HEFT valid for $\sqrt{s} \ll 2 \, m_T$

Higgs pair production:

 $2m_H < \sqrt{s}$

full top quark mass dependence required for accurate predictions

- 1. LO, including full m_T dependence Glover, van der Bij `88
- 2. NLO, (Born-improved) HEFT K≈2 Dawson, Dittmaier, Spira `98
 - including full m_T dependence in real radiation (FT approx.)
 -10% Maltoni, Vryonidou, Zaro `14
 - including $1/m_T$ expansion $\pm 10\%$ Grigo, Hoff, Melnikov, Steinhauser `13; Grigo, Hoff, Steinhauser `15 Degrassi, Giardino, Gröber `16
- 3. NLO, including full m_T dependence Borowka, Greiner, Heinrich, Jones, MK, Schlenk, Schubert, Zirke `16
 - NLO matched to parton shower Heinrich, Jones, Luisoni, MK, Vryonidou `17 this talk
 - transverse momentum NLL+NLO Ferrera, Pires `16

- 4. NNLO (HEFT) de Florian, Mazzitelli `13
 - including all matching coefficients Grigo, Melnikov, Steinhauser `14
 - including $1/m_T$ expansion Grigo, Hoff, Steinhauser `15
 - NNLL soft gluon resummation Shao, Li, Li, Wang `13
 - NNLL + NNLO matching de Florian, Mazzitelli `15
 - fully differential de Florian, Grazzini, Hanga, Kallweit, Lindert, Maierhöfer, Mazzitelli, Rathlev `16

+20%

Two Loop Diagrams lagrams



→ numeric calculation required

NLO Calculation

1. Reduction to master integrals using Reduze von Manteuffel, Studerus`12

- full reduction achieved only for planar integrals (non-planar integrals evaluated directly)
- use finite basis von Manteuffel, Panzer, Schabinger `15
- 2. Numerical evaluation of 2-loop integrals using SecDec Borowka, Heinrich, Jahn, Jones, MK, Schlenk, Zirke
 - using Quasi-Monte-Carlo integration with $O(n^{-1})$ scaling Li, Wang, Yan, Zhao `15; Review: Dick, Kuo, Sloan
 - dynamically set number of sampling points for each integral
 - parallelization on gpu
- 3. Use unweighted events (based on LO) for phase space integration of virtuals

4. Real radiation amplitudes using GoSam Cullen, van Deurzen, Greiner, Heinrich, Luisoni, Mastrolia, Mirabella, Ossola, Peraro, Schlenk, von Soden-Fraunhofen, Tramontano

5. Dipole subtraction Catani Seymour to deal with IR singularities

NLO Results — Invariant Mass



- basic HEFT leads to wrong shape
- B.I. HEFT overestimates by 16% / 30%
- FT approx closer to full result (difference increasing with m_{hh})

	$14 { m TeV}$	$100 { m TeV}$
LO	$19.85^{+27.6\%}_{-20.5\%}$	$731.3^{+20.9\%}_{-15.9\%}$
B.i. HEFT	$38.32^{+18.1\%}_{-14.9\%}$	$1511^{+16.0\%}_{-13.0\%}$
FT approx	$34.26^{+14.7\%}_{-13.2\%}$	$1220^{+11.9\%}_{-10.7\%}$
NLO full	$32.91^{+13.6\%}_{-12.6\%}$	$1149^{+10.8\%}_{-10.0\%}$

NLO Results – Higgs Momentum



- basic HEFT leads to wrong shape
- B.I. HEFT overestimates by 16% / 30%
- FT approx closer to full result (difference increasing with $p_{T,h}$)

top mass effects important, in particular at $\sqrt{s} = 100 \,\mathrm{TeV}$

Parton Shower Interface

2-loop amplitude too slow (median 2h on gpu) for direct interface to PS \rightarrow construct grid for interpolation of virtual amplitude (3741 2-loop results used)

input parameters (\hat{s}, \hat{t}) transformed to

$$x = f(\beta(\hat{s})), \quad c_{\theta} = |\cos \theta| = \left| \frac{\hat{s} + 2\hat{t} - 2m_{H}^{2}}{\hat{s}\beta(\hat{s})} \right|, \quad \beta = \left(1 - \frac{4m_{H}^{2}}{\hat{s}} \right)^{\frac{1}{2}}$$

 \rightarrow nearly unify distribution of phase space points in $(x, c_{\theta}) \in [0, 1]^2$ if $f(\beta)$ chosen according to cumulative distribution of points in original calculation

 (x, c_{θ})

Grid validation:



LHE Events in HEFT & comparison with NNLO

Les Houches Event Level:

Sudakov factor included, but no parton shower

Powheg allows to split real radiation into (exponentiated) singular and regular part



• $h = \infty$: LHE level results close to NNLO

• h = 250: LHE level approaches NLO in tail of p_T^{hh} distribution

Results including Parton Shower

Powheg + Pythia8



only small parton shower effects on NLO accurate observables

Results including Parton Shower

Powheg

Parton shower effects large for observables sensitive to real radiation, e.g. p_T^{hh}

MadGraph5_aMC@NLO



- parton shower enhances tail p_T^{hh} distribution by factor of ~2
- difference of matching schemes of ~20%
- small difference between full NLO and FT approx. result

Results including Parton Shower

Parton shower effects large for observables sensitive to real radiation, e.g. ΔR^{hh}



MadGraph5 aMC@NLO



$\Delta R^{hh} < \pi$

- filled by real radiation only LO accurate
- parton shower corrections up to factor of ~2.5
- differences due to matching method visible

- NLO accurate
- small dependence on parton shower / matching

Higgs pair production at NLO

- retaining full m_t dependence
- numeric evaluation of 2-loop amplitudes
- reduces cross section by 14% compared to Born-improved HEFT
- corrections not uniform over phase space

Matching to Parton Showers

- up to ~20% differences for NLO accurate observables only small dependence on matching method
- effects can be large for LO accurate observables

Outlook

- comparison with Herwig and Sherpa parton shower
- combination of NLO in full theory with NNLO HEFT

Backup

Two Lockample it gg $\Rightarrow hh$

• tensor structure Glover, van der Bij `88

MAX-PLANCK-GESELLSCHAFT

$$\mathcal{M} = \epsilon_{\mu}(p_1, n_1)\epsilon_{\nu}(p_2, n_2) \mathcal{M}^{\mu\nu}$$

 $\mathcal{M}^{\mu\nu} = A_1(s, t, m_H^2, m_t^2, D) T_1^{\mu\nu} + A_2(s, t, m_H^2, m_t^2, D) T_2^{\mu\nu}$

with

$$T_{1}^{\mu\nu} = g^{\mu\nu} - \frac{p_{1}^{\nu} p_{2}^{\mu}}{p_{1} \cdot p_{2}}$$

$$T_{2}^{\mu\nu} = g^{\mu\nu} + \frac{1}{p_{T}^{2} (p_{1} \cdot p_{2})} \left\{ m_{H}^{2} p_{1}^{\nu} p_{2}^{\mu} - 2 \left(p_{1} \cdot p_{3} \right) p_{5}^{\nu} p_{1}^{\mu} - 2 \left(p_{2} \cdot p_{3} \right) p_{5}^{\nu} p_{1}^{\mu} + 2 \left(p_{1} \cdot p_{2} \right) p_{5}^{\nu} p_{5}^{\mu} \right\}$$
triangle diagrams $gg \rightarrow H \rightarrow HH$
only contribute to A_{1}

$$M^{++} = M^{-+} = -A_{2}$$

$$T_{2}^{\mu\nu} = g^{\mu\nu} + \frac{1}{p_{T}^{2} (p_{1} \cdot p_{2})} \left\{ m_{H}^{2} p_{1}^{\nu} p_{2}^{\mu} - 2 \left(p_{1} \cdot p_{3} \right) p_{5}^{\nu} p_{1}^{\mu} + 2 \left(p_{1} \cdot p_{2} \right) p_{5}^{\nu} p_{5}^{\mu} \right\}$$

projectors

construct projectors $P_j^{\mu\nu}$ such the $P_1^{\mu\nu}\mathcal{M}_{\mu\nu} = A_1(s,t,m_H^2,m_t^2,D)$ **Construct** $P_i^{\mu\nu} = \sum c_{ij} T_j^{\mu\nu}$ such that

$$P_2^{\mu\nu}\mathcal{M}_{\mu\nu} = A_2(s,t,m_H^2,m_t^2,D)$$





MAX-PLANCK-GESELLSCH.

Amplitude Structure

rewrite loop integrals with r propagators and s inverse propagators as

$$I_{r,s}(s,t,m_h^2,m_t^2) = (M^2)^{-L\epsilon} (M^2)^{2L-r+s} I_{r,s} \left(\frac{s}{M^2},\frac{t}{M^2},\frac{m_h^2}{M^2},\frac{m_t^2}{M^2}\right)$$

arbitrary scale

and write renormalized form factors as

$$F^{\text{virt}} = aF^{(1)} + a^{2} \left(\frac{n_{g}}{2} \,\delta Z_{A} + \delta Z_{a}\right) F^{(1)} + a^{2} \delta m_{t}^{2} F^{ct,(1)} + a^{2} F^{(2)} + \mathcal{O}(a^{3})$$

$$F^{(1)} = \left(\frac{\mu_{R}^{2}}{M^{2}}\right)^{\varepsilon} \left[b_{0}^{(1)} + b_{1}^{(1)}\varepsilon + b_{2}^{(1)}\varepsilon^{2} + \mathcal{O}(\varepsilon^{3})\right], \qquad \text{(1-loop)}$$

$$F^{ct,(1)} = \left(\frac{\mu_{R}^{2}}{M^{2}}\right)^{\varepsilon} \left[c_{0}^{(1)} + c_{1}^{(1)}\varepsilon + \mathcal{O}(\varepsilon^{2})\right], \qquad \text{(mass counter-term)}$$

$$F^{(2)} = \left(\frac{\mu_{R}^{2}}{M^{2}}\right)^{2\varepsilon} \left[\frac{b_{-2}^{(2)}}{\varepsilon^{2}} + \frac{b_{-1}^{(2)}}{\varepsilon} + b_{0}^{(2)} + \mathcal{O}(\varepsilon)\right], \qquad \text{(2-loop)}$$

 \rightarrow scale variations do not require re-computation of $b_i^{(n)}, c_i^{(n)}$

Amplitude Evaluation — Example

$\sqrt{s} = 327.25 \,\text{GeV}, \, \sqrt{-t} = 170.05 \,\text{GeV}, \, M^2 = s/4$

contributing integrals:



Results - Combination with NNLOHEFT



 $^{\prime}\mathrm{d}m_{hh}$

900

1000

modified Higgs self-interactions



modified Higgs self-interactions



Calculation of σ^{v}

Importance sampling:



 σ^V with 2.5% accuracy using

~1000 phase-space points

- Accuracy goal: 3% for form factor F₁
 - 5-20% for form factor F_2 (depending on F_2/F_1)

 Run time: (gpu time)

- 80 min 2 d (=wall-clock limit)
- median: 2h



Grigo, Hoff, Steinhauser `15



Grigo, Hoff, Steinhauser `15

Differential Cross Section





NNLO and NNLL results



3-point, 1 off-shell leg



Spira, Djouadi et al. `93, `95 Bonciani, Mastrolia `03, `04 Anastasiou, Beerli et al. `06

$$ightarrow \mathsf{HPLs}$$

3-point, 2 off-shell leg



Gehrmann, Guns, Kara `15 → generalized HPLs, 12 letters

Amplitude Structure (II)

Form factors are sums of rational functions multiplied by integrals that depend on ratios of the scales s, t, m_h^2, m_t^2 and the arbitrary scale M^2

$$\begin{split} F^{(L)} &= \sum_{i} \left[\left(\sum_{j} C_{i,j}^{(L)} \epsilon^{j} \right) \cdot \left(\sum_{k} I_{i,k}^{(L)} \epsilon^{k} \right) \right] \\ &= \epsilon^{-2} \left[C_{1,-2}^{(L)} \cdot I_{1,0}^{(L)} + C_{1,-1}^{(L)} \cdot I_{1,-1}^{(L)} + \dots \right] \\ &+ \epsilon^{-1} \left[C_{1,-1}^{(L)} \cdot I_{1,0}^{(L)} + \dots \right] + \dots \\ &\quad \text{compute only once} \end{split}$$

Additionally, all *L*-loop form factors are computed simultaneously without re-evaluating common integrals

Note: $gg \rightarrow HH$ is a loop induced process, real subtraction and mass factorisation contained in $\mathbf{I}, \mathbf{P}, \mathbf{K}$ operators (not discussed here) Catani, Seymour 96

Slide: Stephen Jones — L&L 2016

Phase-Space Sampling Ce Sampling

Phase-space implemented by hand

limited to 2-3 w/ 2 massive particles Events for virtual:

1) VEGAS algorithm applied to LO matrix element $\mathcal{O}(100k)$ events computed

2) Using LO events unweighted events generated using accept/reject method $\mathcal{O}(30k)$ events remain

3) Randomly select 666 Events (woops), compute at NLO, exclude 1

Note: No grids used either for integrals or phase-space

Slide: Stephen Jones — L&L 2016

