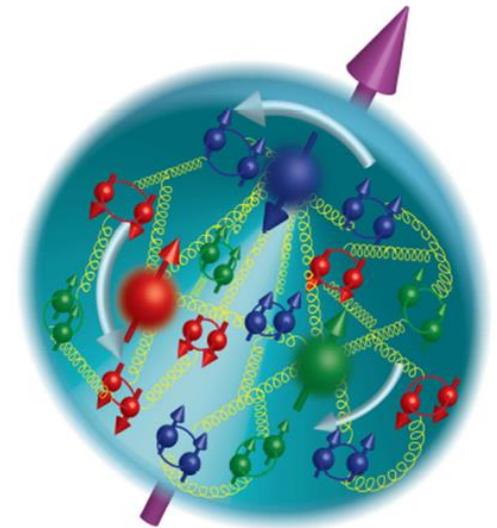


Exploring the Proton Spin with Di-jets at a Future EIC

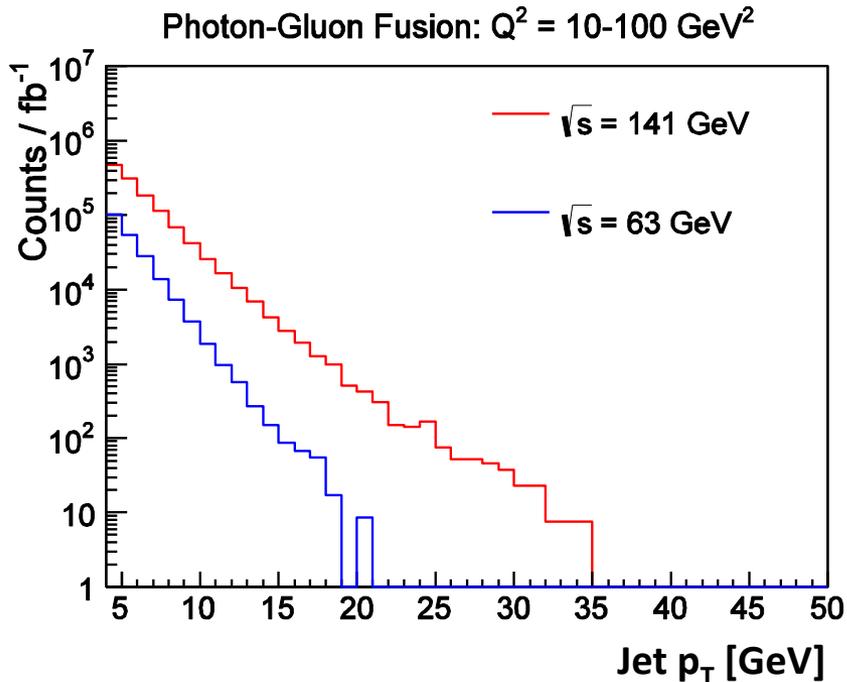
Brian Page
Brookhaven National Laboratory
DIS 2017 - Birmingham



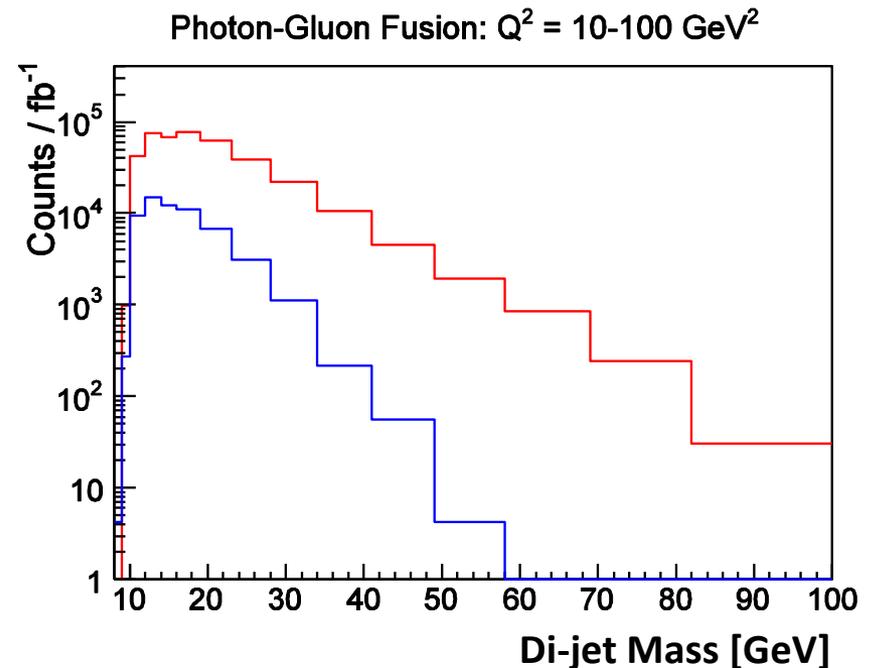
Simulation Details / Particle Cuts

- Electron – Proton events generated using PYTHIA
- Cut on inelasticity: $0.01 \leq y \leq 0.95$
- Jet Algorithm: Anti- k_T ($R = 1.0$)
- Jets found in Breit frame
- Particles used in jet finding:
 - Stable
 - $p_T \geq 250$ MeV
 - $\eta \leq 4.5$
 - Parent cannot originate from scattered electron

Jets at an EIC: Points to Remember

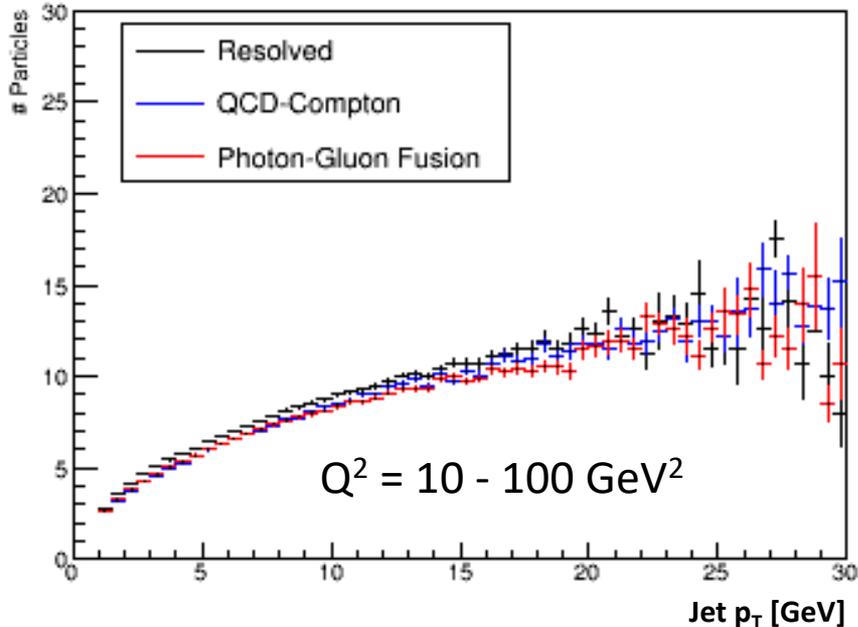


- Lower center of mass energies will lead to lower jet / di-jet yields and more limited p_T / mass reach
- Will need largest available energies and high luminosity to accumulate reasonable statistics at high p_T / mass – use $\sqrt{s} = 141 \text{ GeV}$ for all that follows



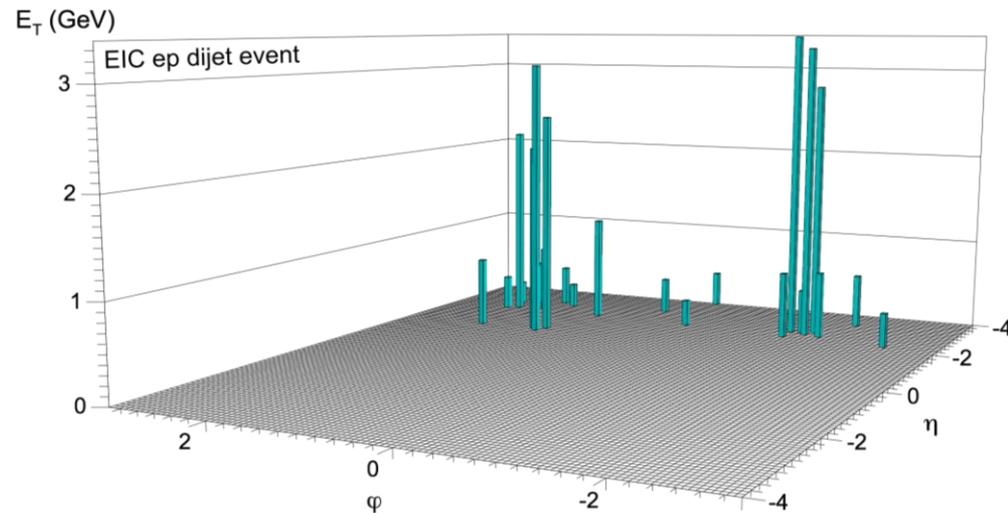
Jets at an EIC: Points to Remember

Number of Particles in Jet Vs Jet Pt



- Jets contain relatively few particles overall
- Events are quite clean with little underlying event activity
- Typical particle p_T is small \rightarrow precision tracking important for reducing jet energy scale uncertainties

- Lower center of mass energies will lead to lower jet / di-jet yields and more limited p_T / mass reach
- Will need largest available energies and high luminosity to accumulate reasonable statistics at high p_T / mass – use $\sqrt{s} = 141$ GeV for all that follows



Jet / Di-jet Applications

- Electron – Nucleus Collisions

- Nuclear PDFs
- Medium Modification / Energy Loss
- Hadronization and Confinement

- Electron – Proton Collisions

- High – x Proton PDFs
- Proton Orbital Angular Momentum / Gluon Wigner Distribution at Low x (Diffractive Di-jets)
- Saturation (Diffractive Di-jets)
- (Un)polarized Photon PDFs
- PGF Tagging for Proton ΔG

Need for ΔG

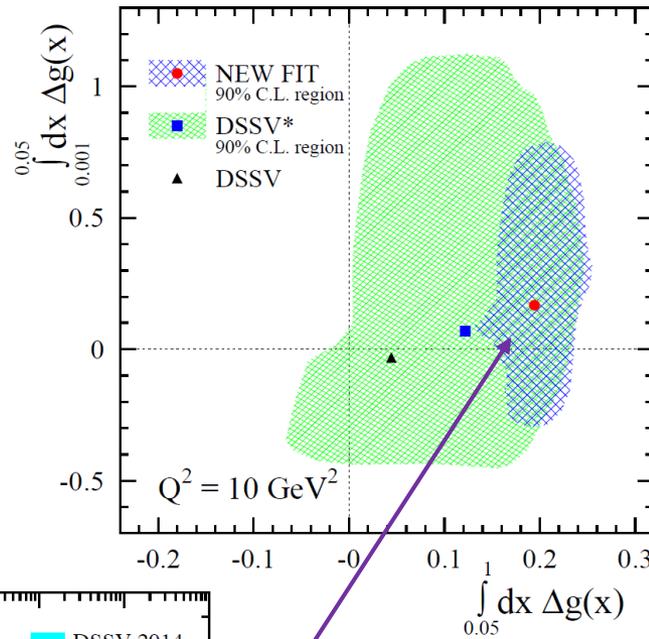
$$S = \frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + L$$

• Current fits to polarized DIS data give $\Delta\Sigma = 0.366^{+0.042}_{-0.062}$ for $10^{-3} < x < 1$

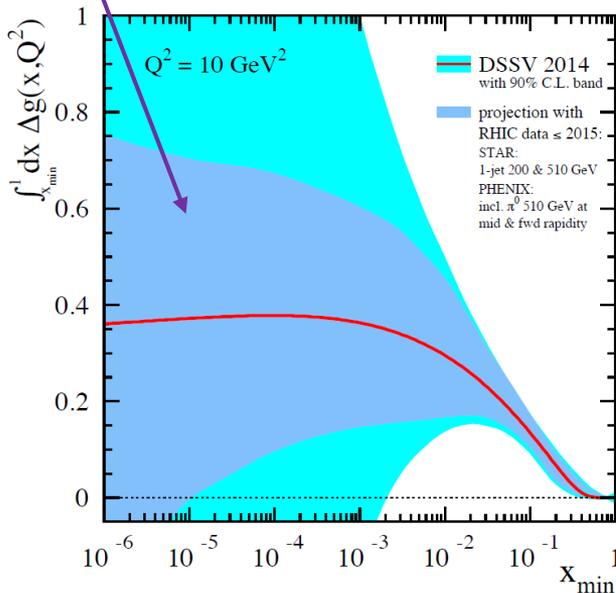
• This leaves over half of the proton spin budget unaccounted for – is any spin carried by the gluons?

• Polarized p+p collision data from RHIC have placed strong constraints on ΔG for $x > 0.05$ and given evidence of non-zero gluon contribution, but large uncertainties at low x will remain

• Need high precision and wide kinematic reach to solve spin puzzle



Uncertainties on low-x region will still be sizable with RHIC data



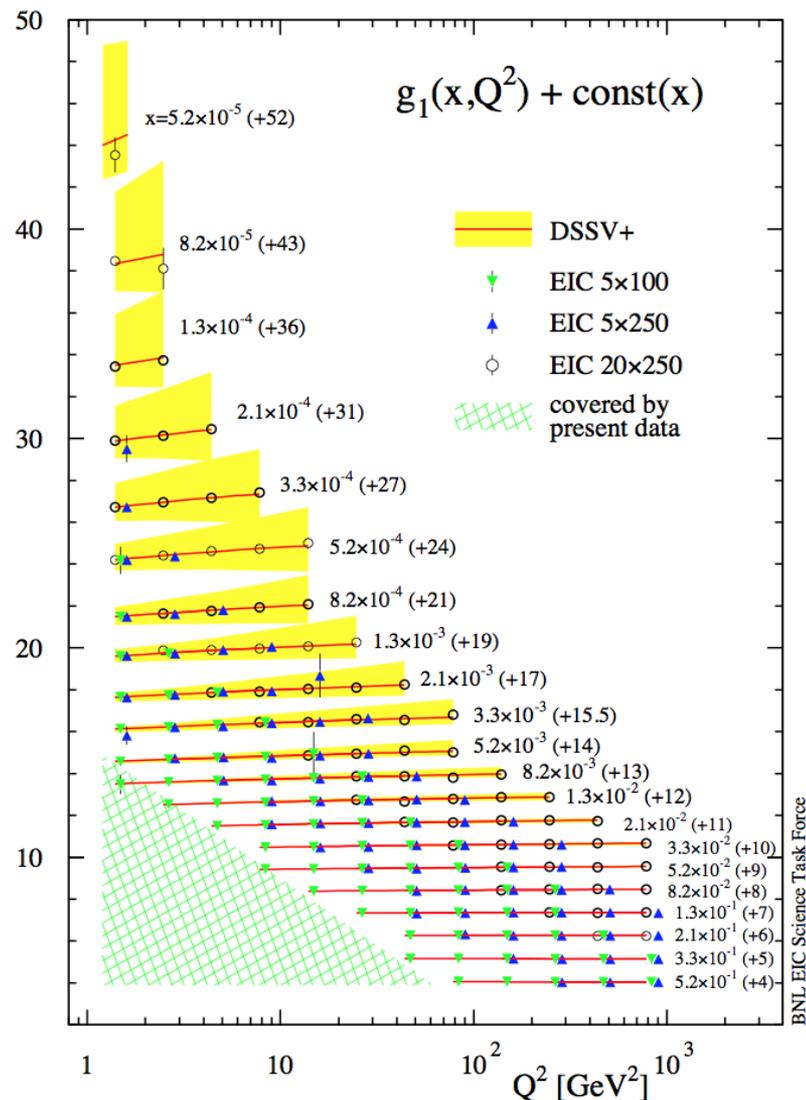
Global analyses with RHIC data through 2009 find a ΔG value of about 0.2 for $x > 0.05$ with small uncertainties

Accessing ΔG in DIS

- Several observables are sensitive to ΔG in DIS but golden measurement at an EIC would be scaling violation of $g_1(x, Q^2)$

$$\frac{dg_1(x, Q^2)}{d\ln(Q^2)} \approx -\Delta g(x, Q^2)$$

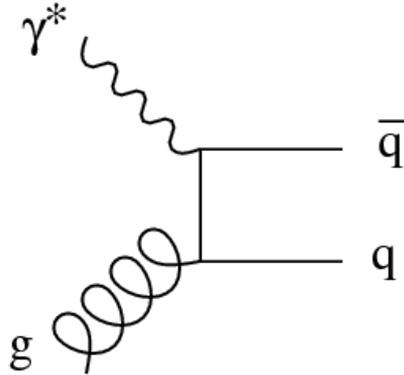
- Current DIS constraints on ΔG hampered by limited x & Q^2 coverage
- EIC would greatly expand kinematic reach and precision of $g_1(x, Q^2)$ measurements!



arXiv:1206.6014

Gluon Polarization with Di-jets

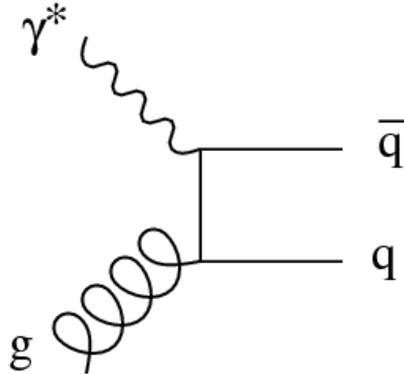
Photon-Gluon Fusion



- Gluons can be also be probed in DIS via the higher-order photon gluon fusion process

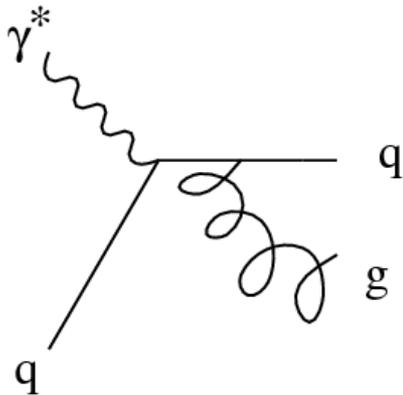
Gluon Polarization with Di-jets

Photon-Gluon Fusion



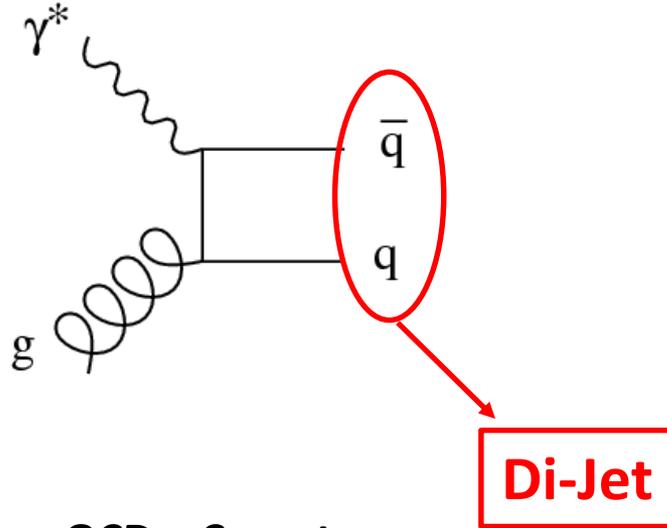
- Gluons can be also be probed in DIS via the higher-order photon gluon fusion process
- Also have the QCD – Compton process which probes quarks at the same order

QCD – Compton

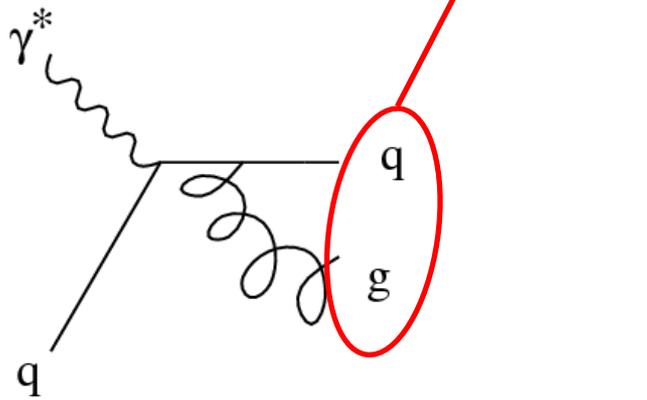


Gluon Polarization with Di-jets

Photon-Gluon Fusion

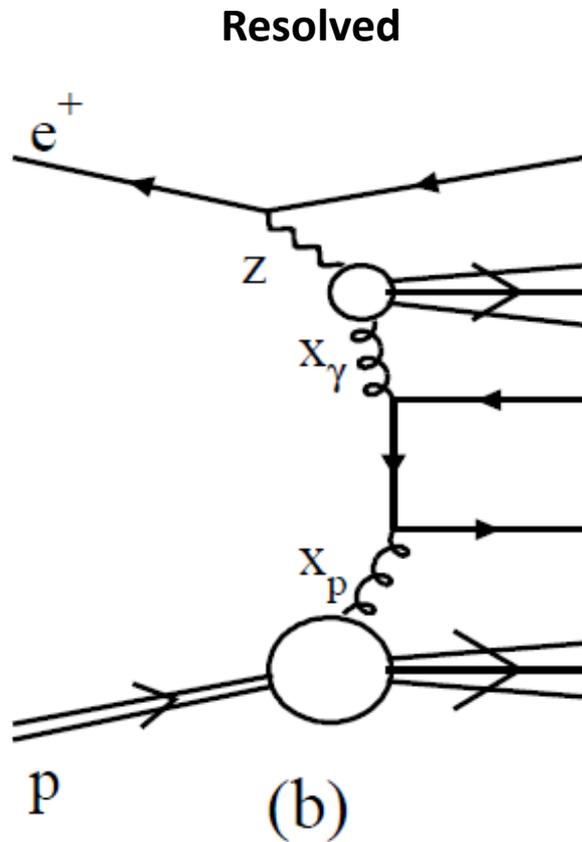


QCD – Compton



- Gluons can be also be probed in DIS via the higher-order photon gluon fusion process
- Also have the QCD – Compton process which probes quarks at the same order
- Both processes produce 2 angularly separated hard partons -> Di-jet

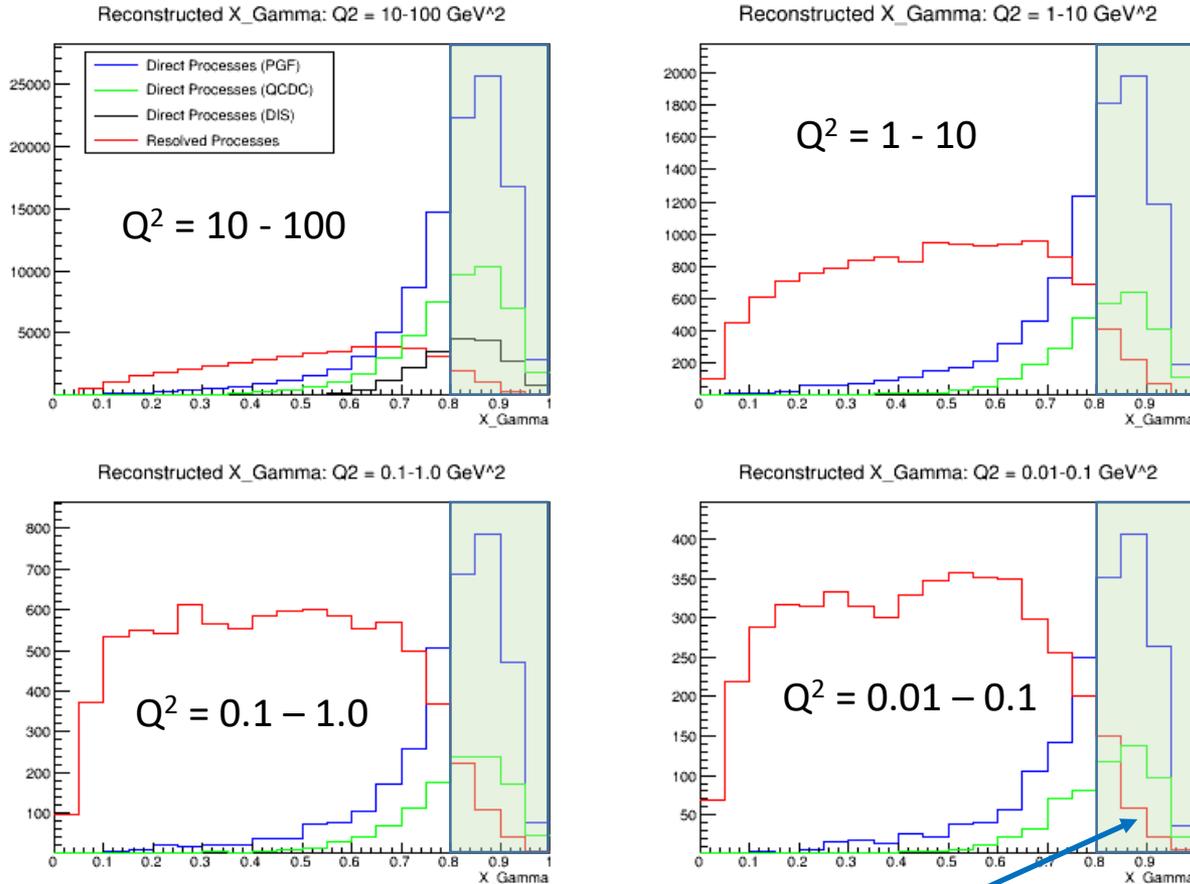
Gluon Polarization with Di-jets



- Gluons can be also be probed in DIS via the higher-order photon gluon fusion process
- Also have the QCD – Compton process which probes quarks at the same order
- Both processes produce 2 angularly separated hard partons -> Di-jet
- At lower Q^2 , resolved processes in which the photon assumes a hadronic structure begin to dominate
- Asymmetry is a convolution of polarized PDF from the proton and polarized photon structure – which is completely unconstrained
- Would like to suppress the resolved component

Direct Vs Resolved Processes

$$X_\gamma = \frac{1}{2E_e y} (m_{T1} e^{-y_1} + m_{T2} e^{-y_2})$$



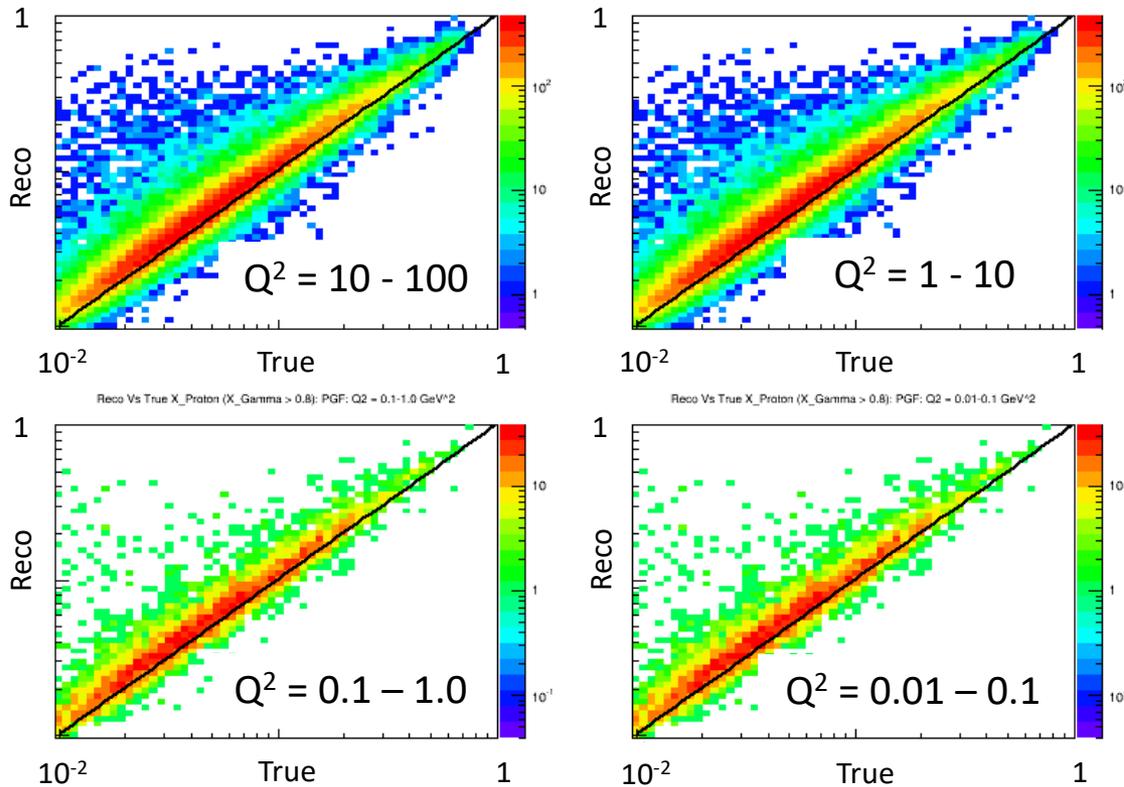
- Plot reconstructed X_γ for direct and resolved processes
- Direct processes should concentrate toward 1 while resolved processes are at lower values
- Direct processes dominate at higher Q^2 while resolved are more prevalent at low Q^2
- Cut of $X_\gamma > 0.8$ enhances the direct fraction at all Q^2

Accepted Region

Proton Partonic Kinematics

Parton Momentum Fraction: Photon Gluon Fusion

- To measure ΔG , need to probe the parton coming from the proton
- Momentum fraction of the parton from proton is well reconstructed



$$X_p = \frac{1}{2E_p} (m_{T1}e^{y_1} + m_{T2}e^{y_2})$$

Proton Partonic Kinematics

$$X_P = x_B \left(1 + \frac{M^2}{Q^2} \right)$$

- To measure ΔG , need to probe the parton coming from the proton
- Momentum fraction of the parton from proton is well reconstructed
- X_p is related to Bjorken- x and Q^2 at leading order

Proton Partonic Kinematics

$$X_P = x_B \left(1 + \frac{M^2}{Q^2} \right)$$

$$Q^2 = syx_B$$

- To measure ΔG , need to probe the parton coming from the proton
- Momentum fraction of the parton from proton is well reconstructed
- X_p is related to Bjorken-x and Q^2 at leading order
- Q^2 and Bjorken-x are also related via the collision energy and inelasticity

Proton Partonic Kinematics

$$X_P = x_B \left(1 + \frac{M^2}{Q^2} \right)$$

$$Q^2 = syx_B$$

$$X_P = x_B + \frac{M^2}{sy}$$

- To measure ΔG , need to probe the parton coming from the proton
- Momentum fraction of the parton from proton is well reconstructed
- X_p is related to Bjorken- x and Q^2 at leading order
- Q^2 and Bjorken- x are also related via the collision energy and inelasticity
- Accessible X_p range basically determined by beam energies

Proton Partonic Kinematics

$$X_P = x_B \left(1 + \frac{M^2}{Q^2} \right)$$

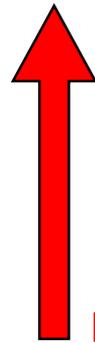
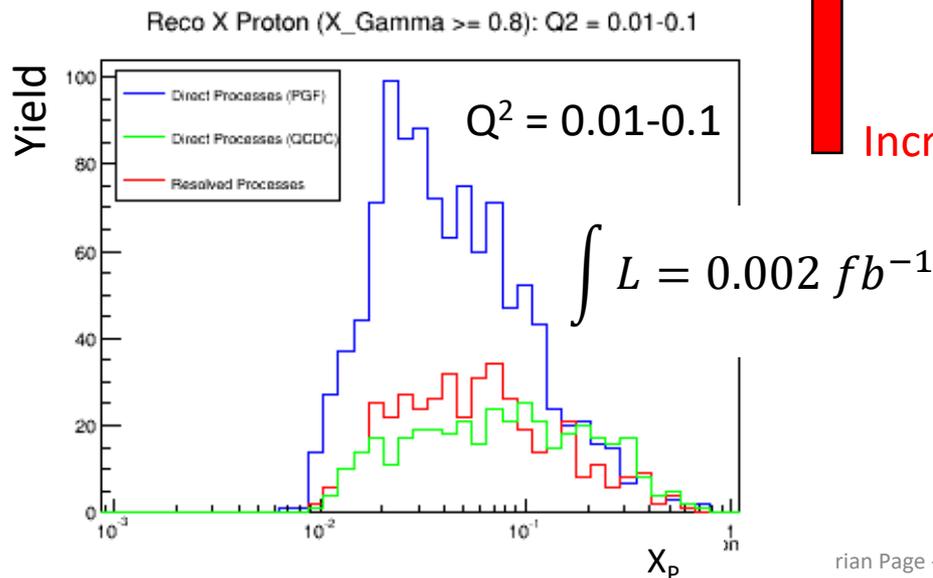
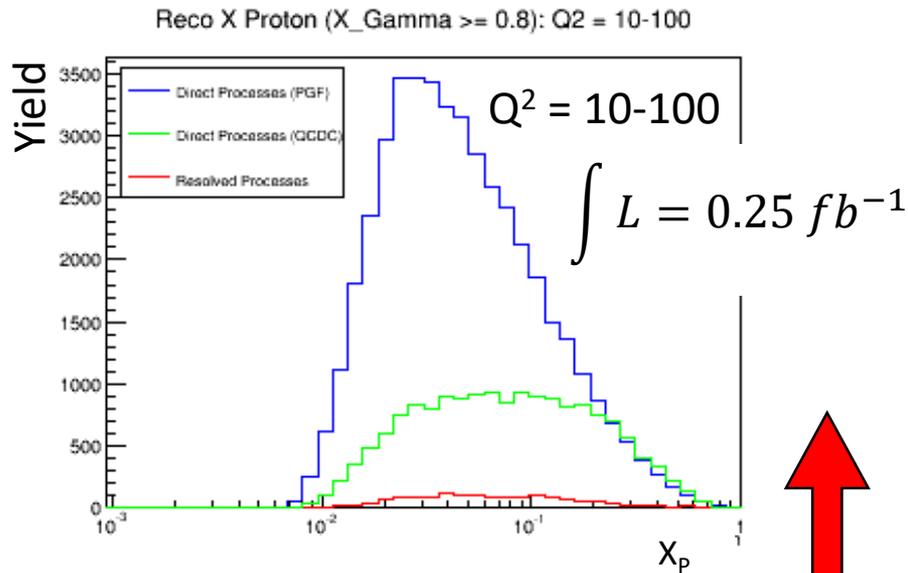
$$Q^2 = syx_B$$

$$X_P = x_B + \frac{M^2}{sy}$$

$$\approx \frac{100}{(20000 \times 0.95)} \approx 0.005$$

- To measure ΔG , need to probe the parton coming from the proton
- Momentum fraction of the parton from proton is well reconstructed
- X_p is related to Bjorken- x and Q^2 at leading order
- Q^2 and Bjorken- x are also related via the collision energy and inelasticity
- Accessible X_p range basically determined by beam energies
- Lowest X_p we can probe is about 0.005

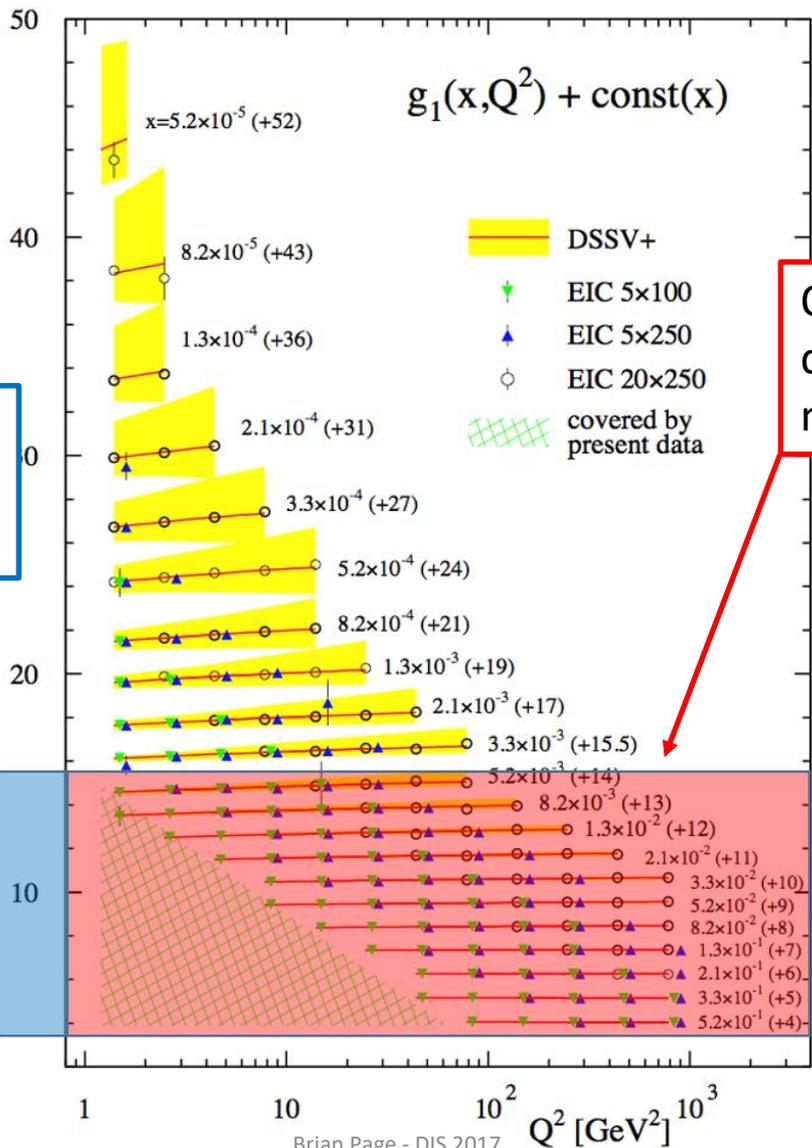
X_p For Different Q^2



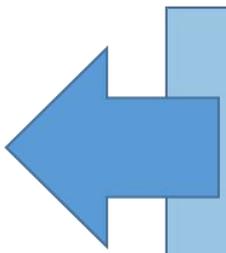
Increasing Q^2

- At lower Q^2 , contribution from resolved process increases while QCD Compton contribution decreases
- For a given di-jet mass range (10 – 20 GeV in this case), same X_p can be reconstructed event-by-event and probed over large range of Q^2
- This will allow for robust tests of the evolution of ΔG

Complementary Coverage



Experimentally possible to extend these measurements to $Q^2 < 1$

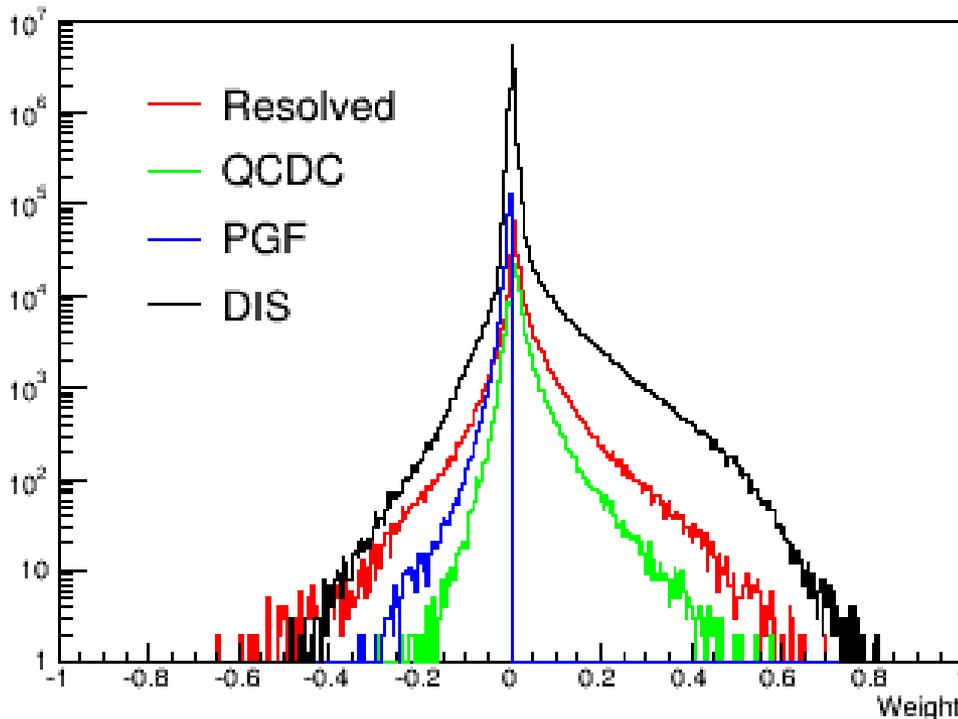


Q^2 and x range covered by di-jet asymmetry measurements

Weighting PYTHIA

$$w = \hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2, Q^2) \cdot \frac{\Delta f_a^{\gamma^*}(x_a, \mu^2)}{f_a^{\gamma^*}(x_a, \mu^2)} \cdot \frac{\Delta f_b^N(x_b, \mu^2)}{f_b^N(x_b, \mu^2)}$$

Total Weight (DSSV14): $Q^2 = 10\text{-}100$ GeV

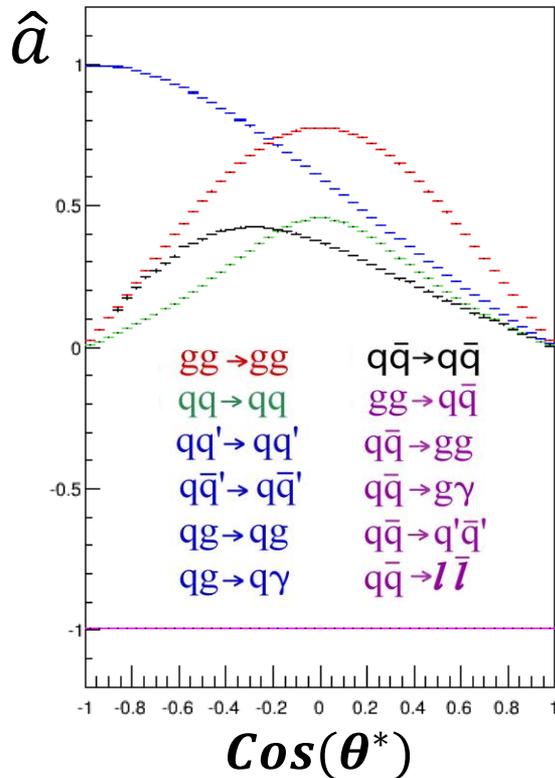


- PYTHIA does not include parton polarization effects, but an asymmetry can be formed by assigning each event a weight depending on the hard-scattering asymmetry and (un)polarized photon and proton PDFs
- Expected asymmetry is then the average over weights
- Weights are sharply spiked near zero -> expect small asymmetries

Weighting PYTHIA

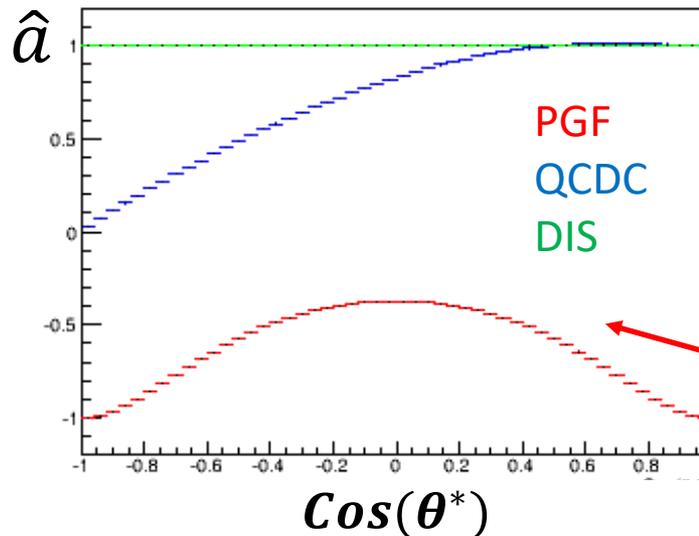
$$w = \hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2, Q^2) \bullet \frac{\Delta f_a^{\gamma^*}(x_a, \mu^2)}{f_a^{\gamma^*}(x_a, \mu^2)} \bullet \frac{\Delta f_b^N(x_b, \mu^2)}{f_b^N(x_b, \mu^2)}$$

Resolved



$$\hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2) = \Delta \hat{\sigma} / (2 \hat{\sigma})$$

Direct



- Process-dependent hard scattering asymmetry is a function of Mandelstam variables ($\text{Cos}(\theta^*)$)

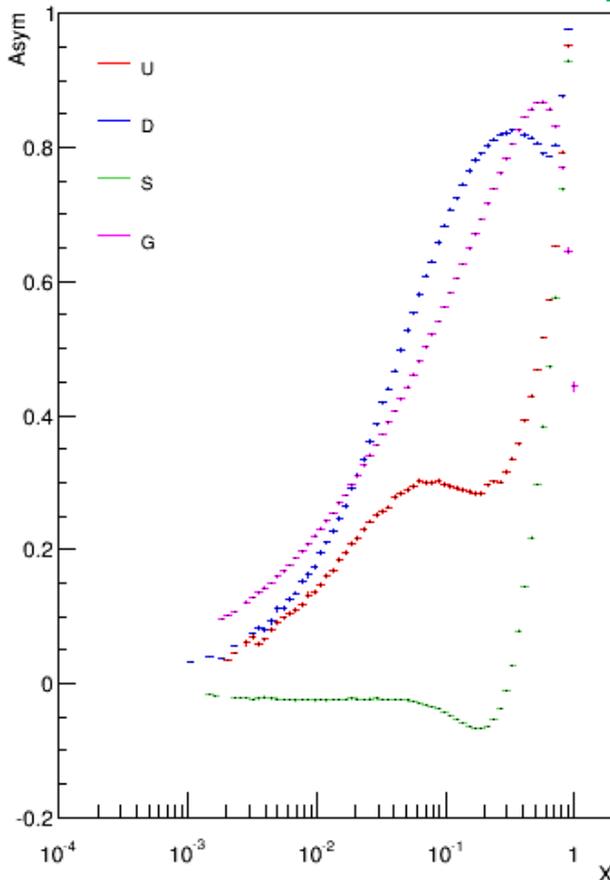
- The direct process distributions will be smeared by the additional depolarization term

- Note that the asymmetry for PGF is negative

Weighting PYTHIA

$$w = \hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2, Q^2) \bullet \frac{\Delta f_a^{\gamma^*}(x_a, \mu^2)}{f_a^{\gamma^*}(x_a, \mu^2)} \bullet \frac{\Delta f_b^N(x_b, \mu^2)}{f_b^N(x_b, \mu^2)}$$

Photon Delta_f/f: Profile: Q2 = 10-100 GeV

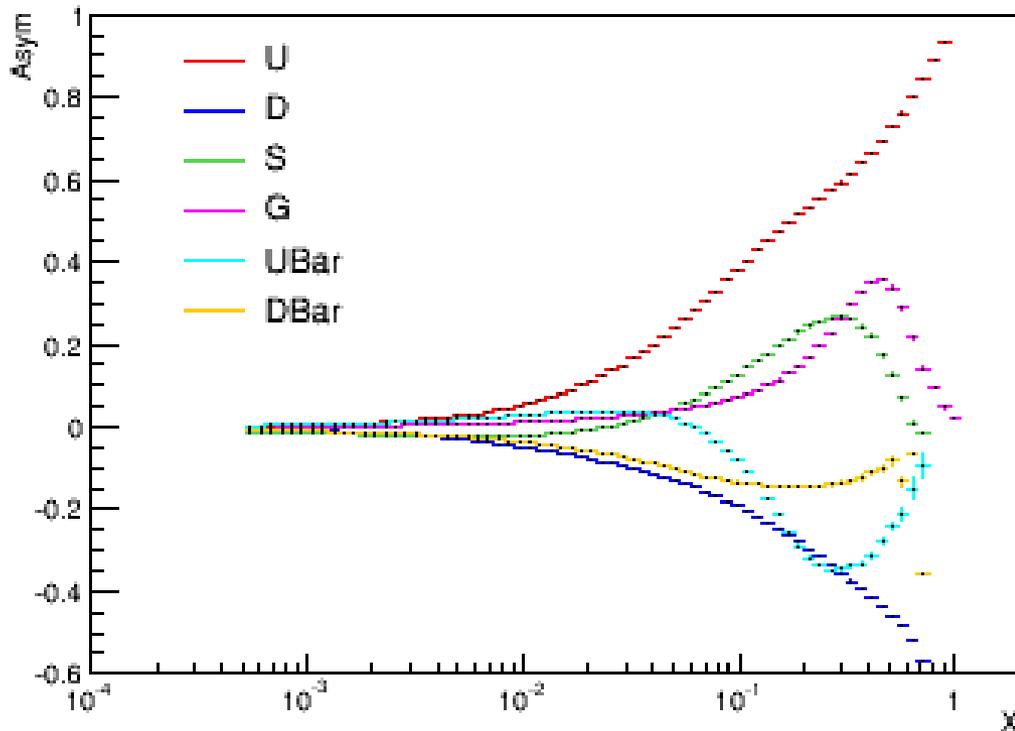


- Second term is the ratio of the polarized to unpolarized photon PDFs
- Use maximal scheme for polarized and GRV-G for unpolarized
- For direct processes such as Photon-Gluon Fusion, this term is identically unity

Weighting PYTHIA

$$w = \hat{a}(\hat{s}, \hat{t}, \hat{\mu}^2, Q^2) \bullet \frac{\Delta f_a^{\gamma^*}(x_a, \mu^2)}{f_a^{\gamma^*}(x_a, \mu^2)} \bullet \frac{\Delta f_b^N(x_b, \mu^2)}{f_b^N(x_b, \mu^2)}$$

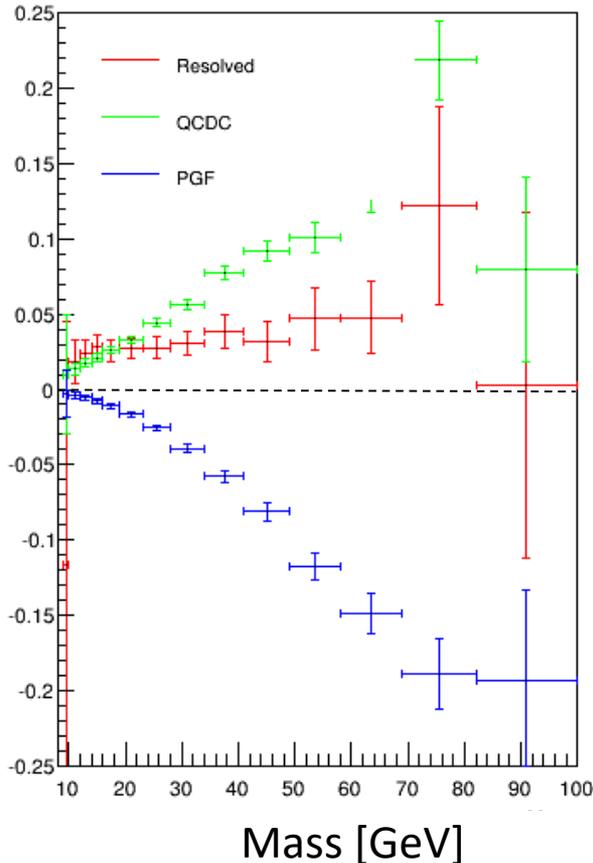
Proton Delta_f/f: Profile (DSSV14): Q2 = 10-100 GeV



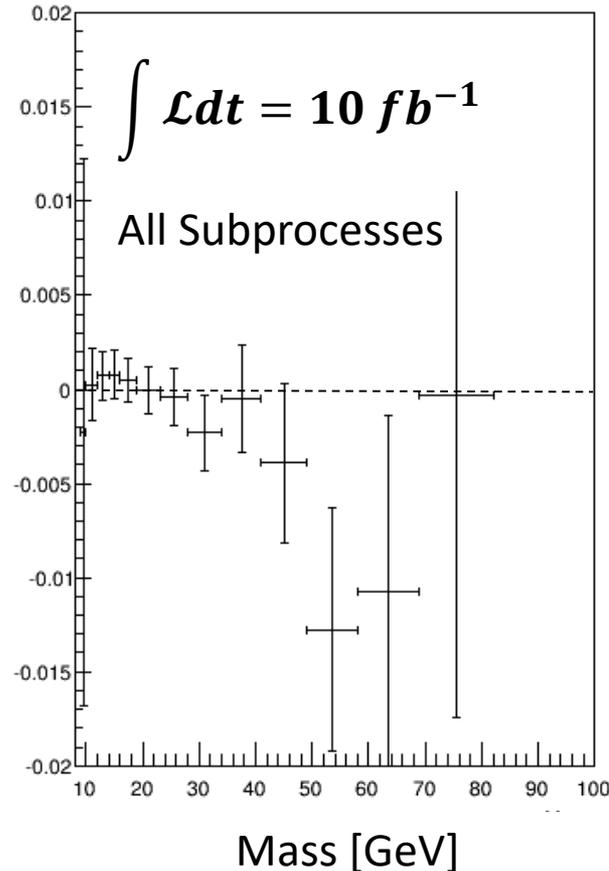
- Last term is the ratio of the polarized to unpolarized proton PDFs
- Use DSSV14 for polarized and CTEQ5M for unpolarized

A_{LL} Vs Di-jet Mass

A_{LL} Vs Di-jet Mass



A_{LL} Vs Di-jet Mass



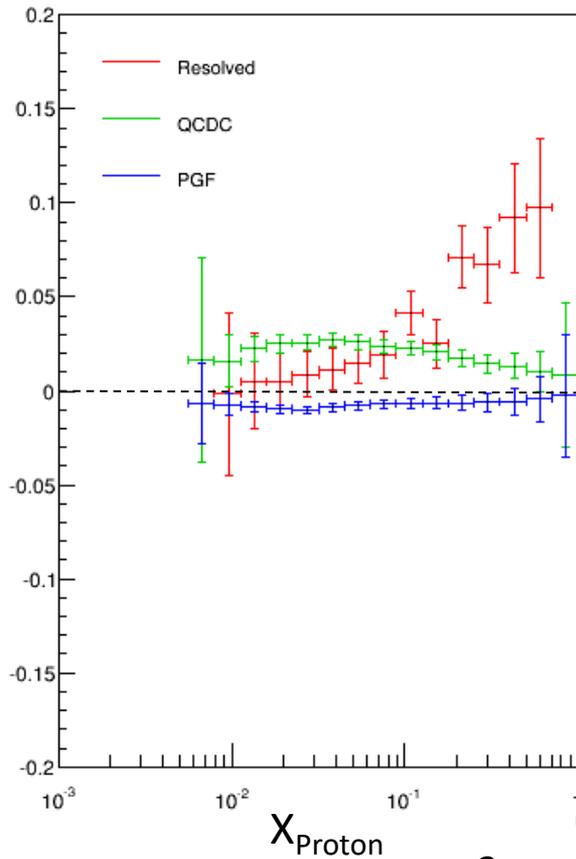
$$Q^2 = 10 - 100 \text{ GeV}^2$$

$$\sigma = \sqrt{\frac{1}{N} - \frac{A^2}{N}}$$

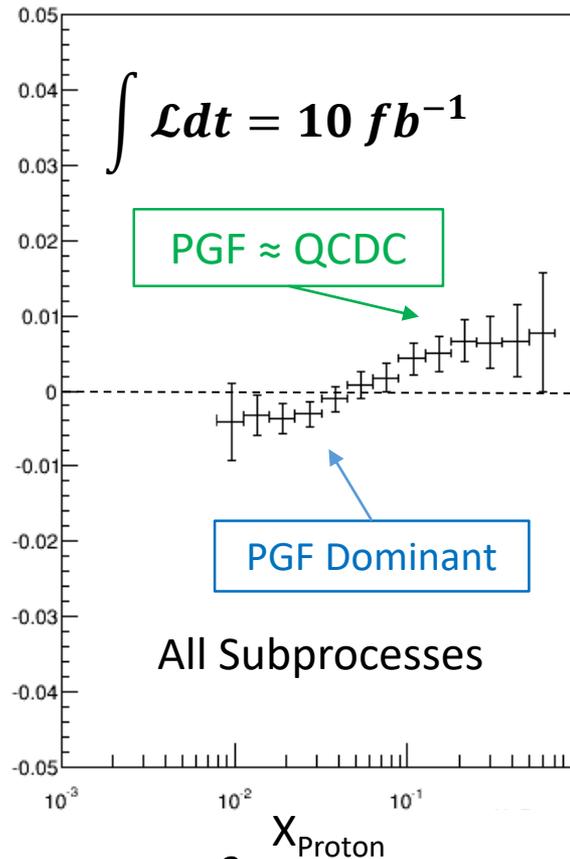
- Plot the expected A_{LL} as a function of di-jet invariant mass for each sub-process separately as well as the combined sample
- PGF asymmetry is nearly canceled out by QCDC asymmetry with opposite sign – would like to reduce QCDC contribution
- Need high integrated luminosity and high energy to probe the high-mass region where asymmetries can be sizable
- Control of systematics will be essential

A_{LL} Vs X_{Proton}

A_{LL} Vs Proton X



A_{LL} Vs Proton X



$$Q^2 = 10 - 100 \text{ GeV}^2$$

$$\sigma = \sqrt{\frac{1}{N} - \frac{A^2}{N}}$$

- Asymmetry is plotted as a function of the momentum fraction of the parton from the proton
- Asymmetry shown for di-jet invariant masses between 10 and 20 GeV/c^2
- Error bars are statistical and scaled to the given integrated luminosity
- Different mass ranges will emphasize different momentum fraction ranges and subprocess mixes

Summary

- Jets at an EIC will contain relatively few total particles and those particles will have low transverse momenta, making tracking essential for reducing jet energy scale uncertainties
- Di-jet measurements can be used to tag photon-gluon fusion events to access ΔG and investigate its evolution
- Combination of QCD-Compton and PGF subprocess asymmetries lead to small overall asymmetries
- Need large integrated luminosity and center-of-mass energy to explore asymmetries at high di-jet mass and good control on systematic effects

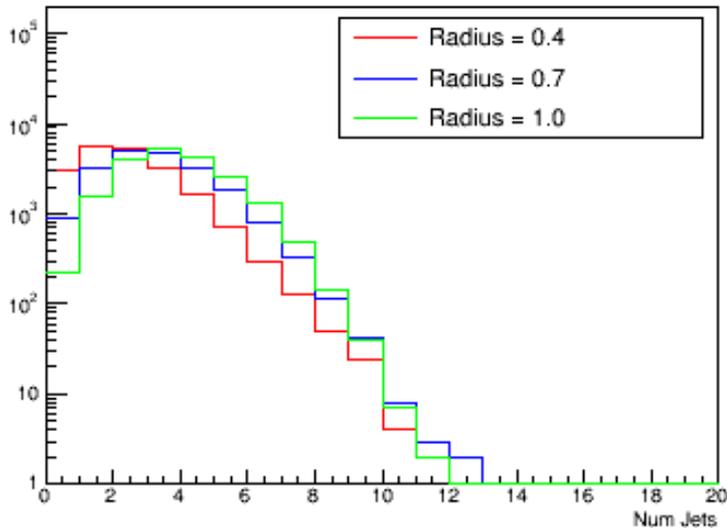
Backup

Jet Basics: Frames

- Can define several useful frames:
 - **Lab:** Detector-based frame
 - **Hadron-Boson:** Beam hadron is at rest, z-direction chosen along virtual photon momentum vector
 - **Breit:** Virtual photon moves in -z direction and boost such that it has zero energy. Separation into target and remnant regions
 - **Center of Mass:** Virtual photon and struck parton have equal and opposite momenta. Can define Feynman-x

Jet Basics: Radius

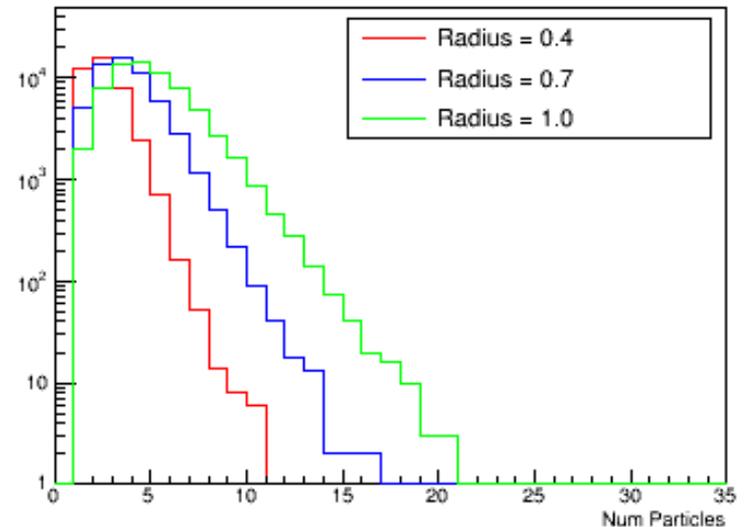
Number of Jets in Event: Anti_kT Hard QCD



- Parameters: Min $p_T = 1.0$ GeV, Resolved processes
- Larger radii result in more found jets as well as more particles in jet

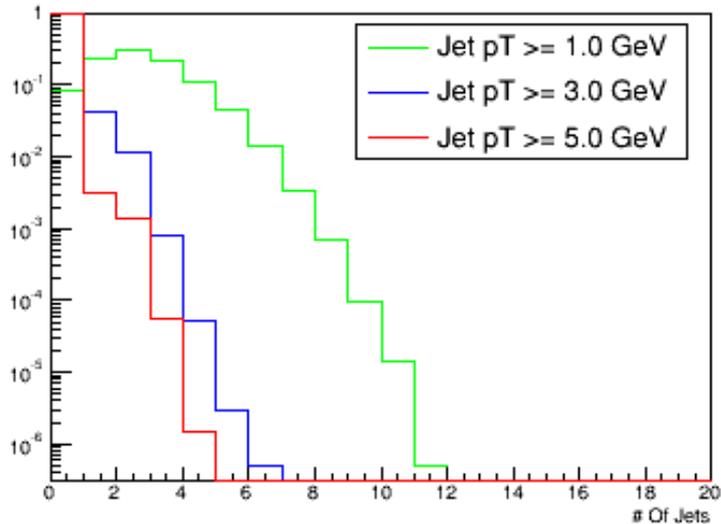
- For anti- k_T algorithm the radius parameter determines the distance at which particles can be grouped together
- Sets the effective size of the jet

Number of Particles in Jet: Anti_kT Hard QCD

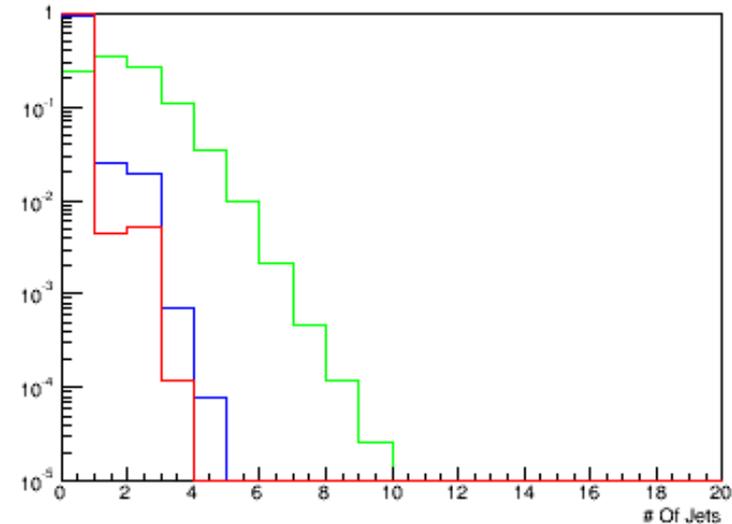


Jet Multiplicity: $Q^2 = 0.01 - 0.1 \text{ GeV}^2$

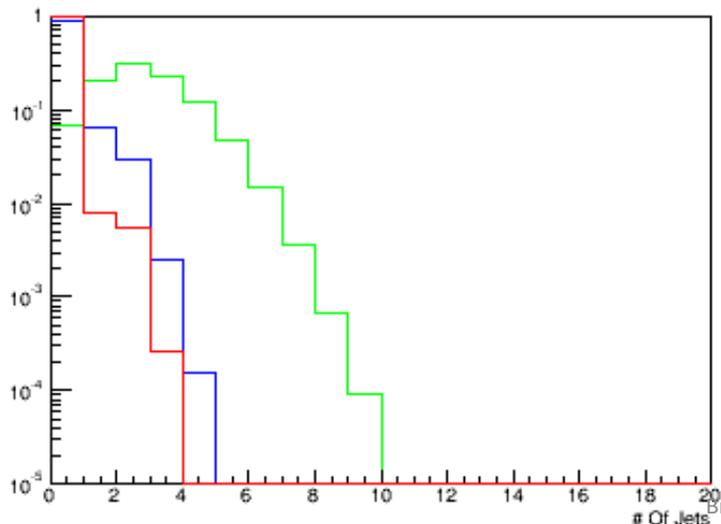
Jets: Resolved Processes



Jets: QCDC



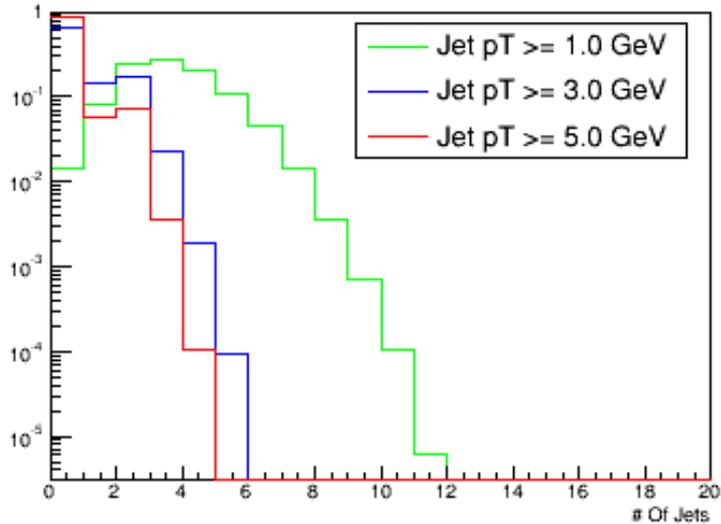
Jets: PGF



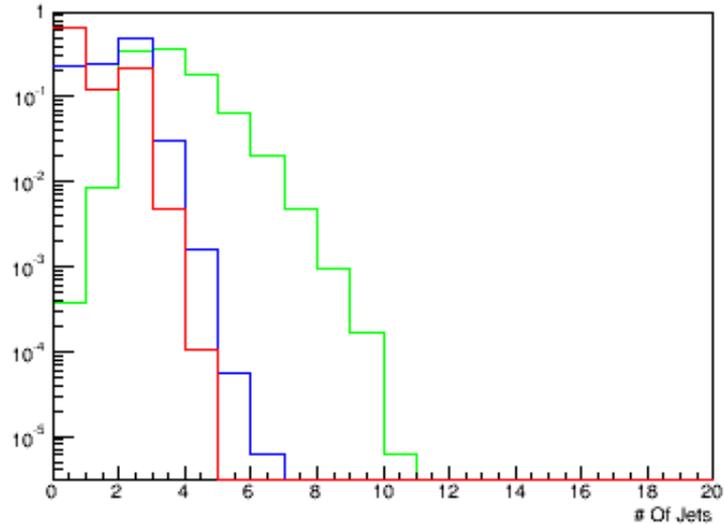
- Percentage of events with a certain number of found jets for different minimum allowed jet p_T s
- See a decrease in number of jets with increasing minimum jet p_T
- Jet p_T of 1 GeV may not be well described theoretically
- Each curve normalized to unity

Jet Multiplicity: $Q^2 = 10 - 100 \text{ GeV}^2$

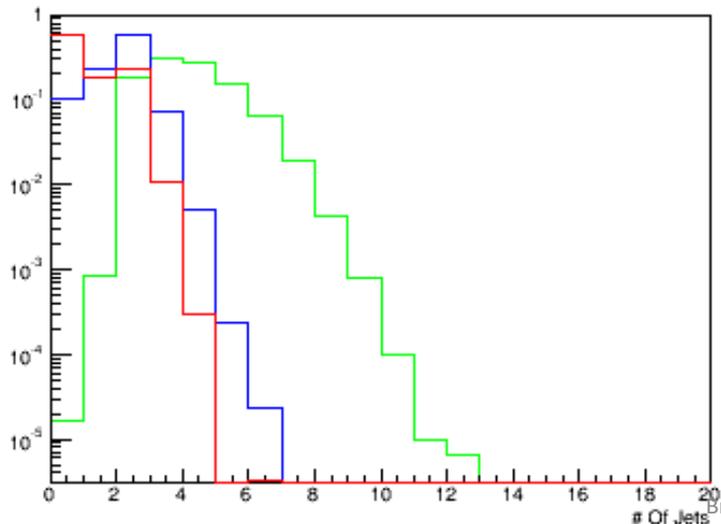
Jets: Resolved Processes



Jets: QCD



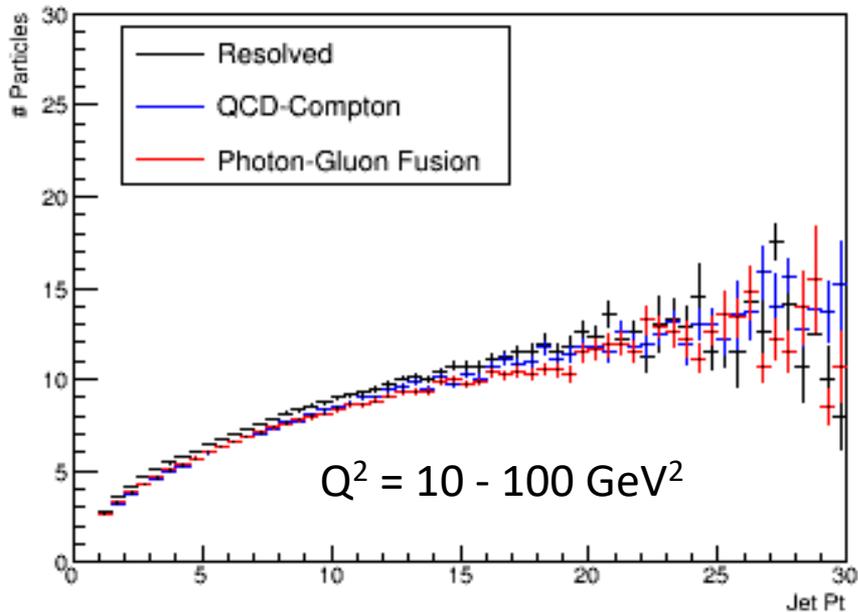
Jets: PGF



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- See a decrease in number of jets with increasing minimum jet p_T
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Jet Particle Multiplicity

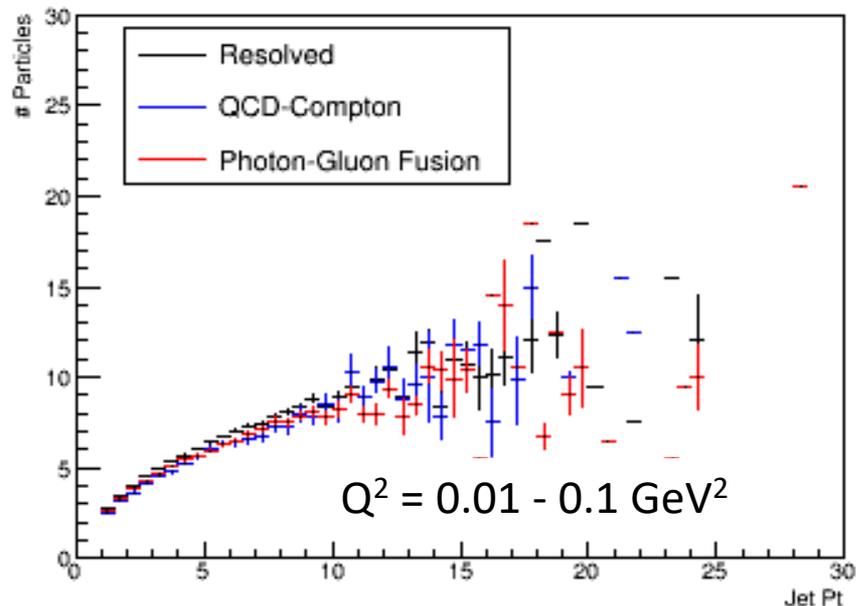
Number of Particles in Jet Vs Jet Pt



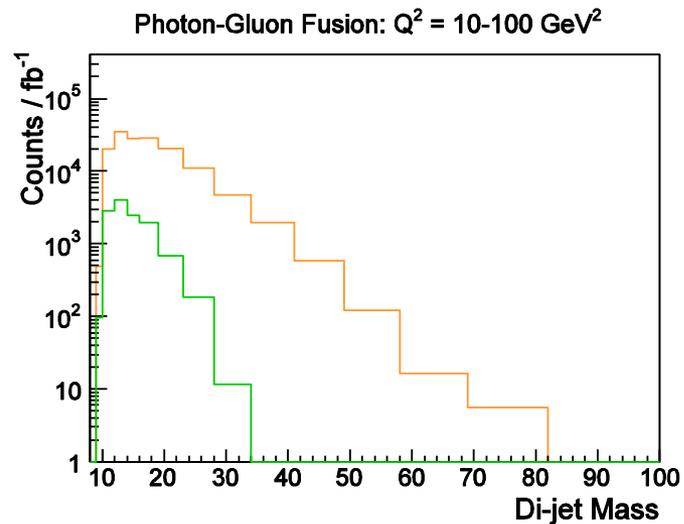
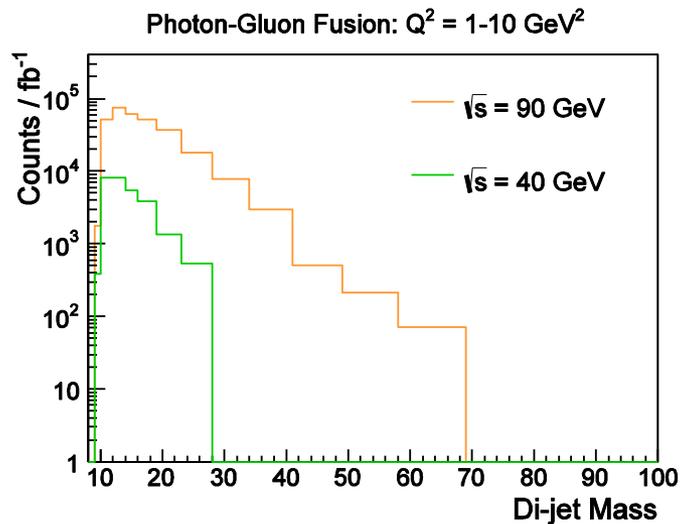
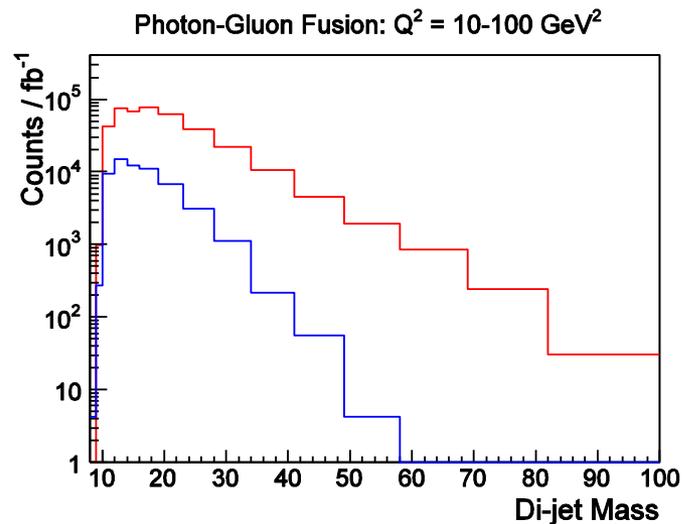
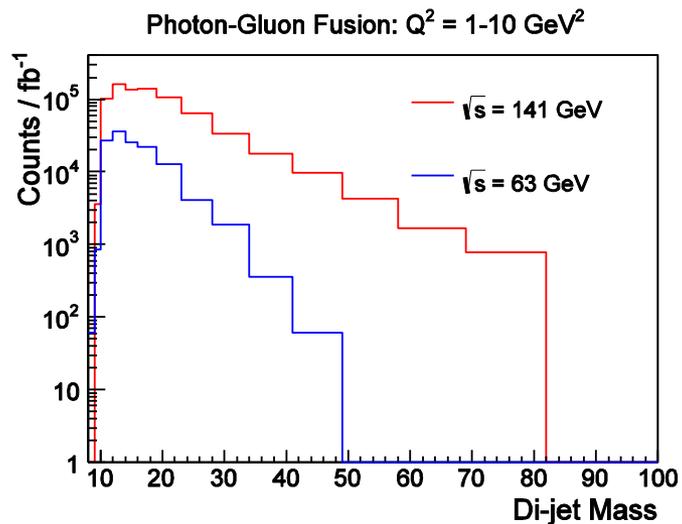
- No dependence on Q^2 or subprocess
- How few particles can be in jet before it doesn't make sense to call the object a jet?

- Look at the average number of particles in jet as a function of jet p_T
- All stable particles (charged and neutral) are counted

Number of Particles in Jet Vs Jet Pt

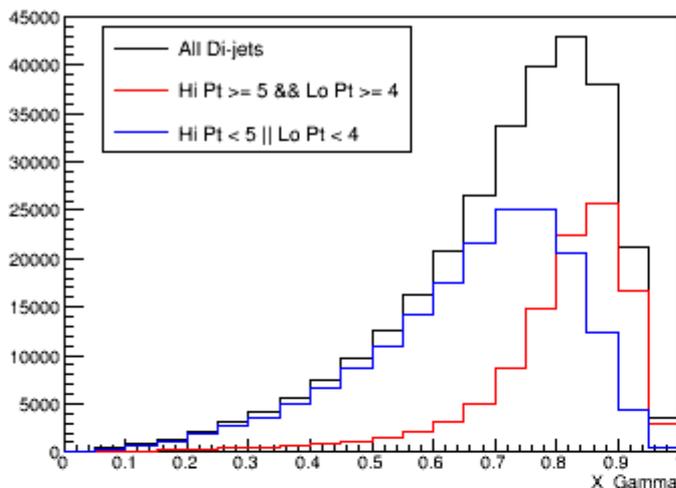


Di-jet Yields

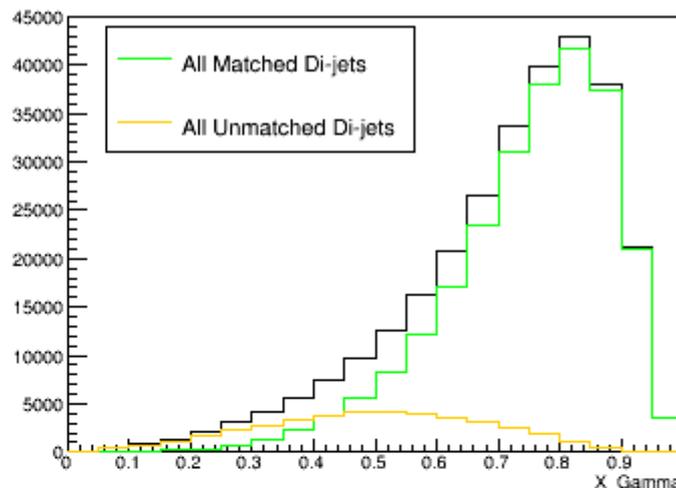


X_γ Reproduction: $Q^2 = 10 - 100 \text{ GeV}^2$

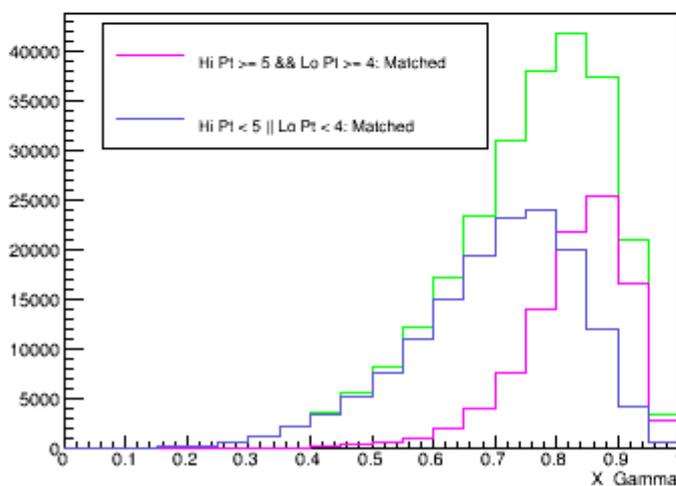
Reconstructed X_γ : Hi / Lo



Reconstructed X_γ : Matched / Unmatched



Reconstructed X_γ : Matched Hi / Lo

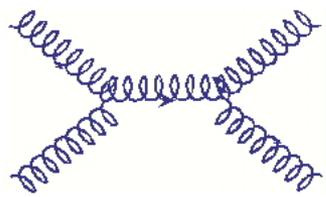


- How does the reproduction of X_γ depend on jet p_T ?
- As expected unmatched events do not reproduce X_γ well
- See that high p_T range is more peaked toward 1 even for matched events

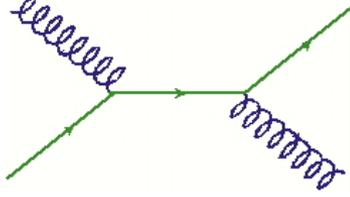
Accessing ΔG at RHIC: A_{LL}

$$A_{LL} = \frac{\sigma^{++} - \sigma^{+-}}{\sigma^{++} + \sigma^{+-}} = \frac{\sum_{a,b,c} \Delta f_a \otimes \Delta f_b \otimes d\hat{\sigma}^{f_a f_b \rightarrow f_c X} \cdot \hat{a}_{LL}^{f_a f_b \rightarrow f_c X} \otimes D_{f_c}^h}{\sum_{a,b,c} f_a \otimes f_b \otimes d\hat{\sigma}^{f_a f_b \rightarrow f_c X} \otimes D_{f_c}^h}$$

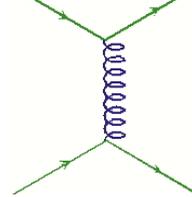
$$\frac{\Delta g}{g} \frac{\Delta g}{g}$$



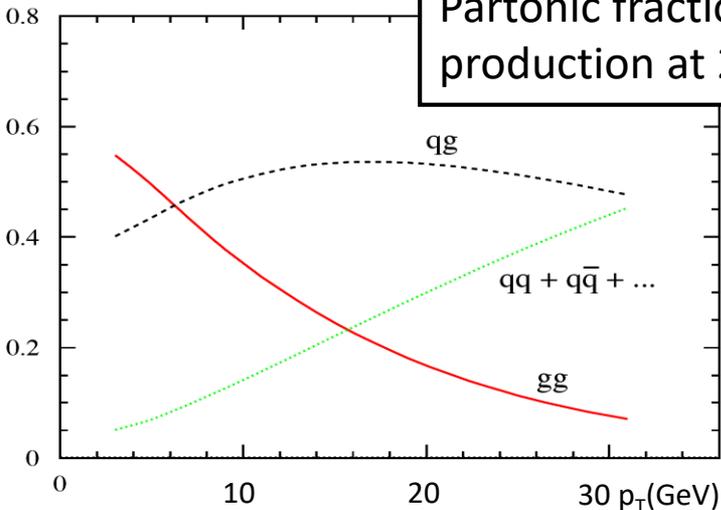
$$\frac{\Delta q}{q} \frac{\Delta g}{g}$$



$$\frac{\Delta q}{q} \frac{\Delta q}{q}$$

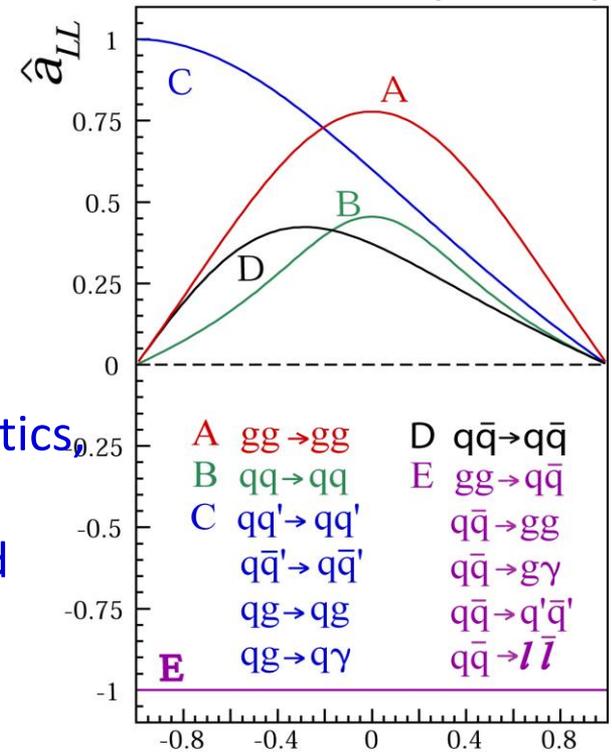


Partonic fractions in jet production at 200 GeV



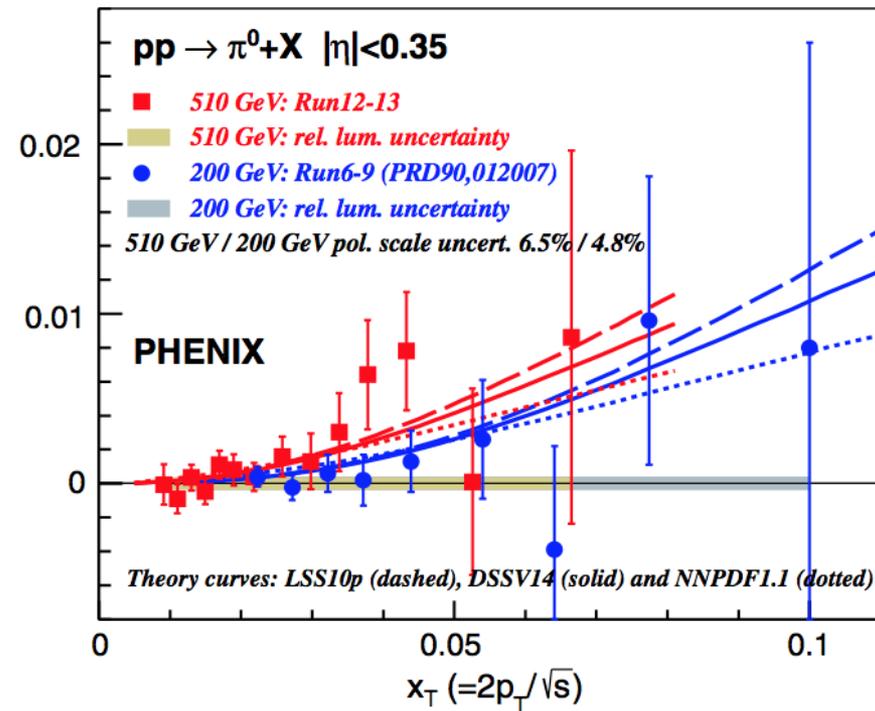
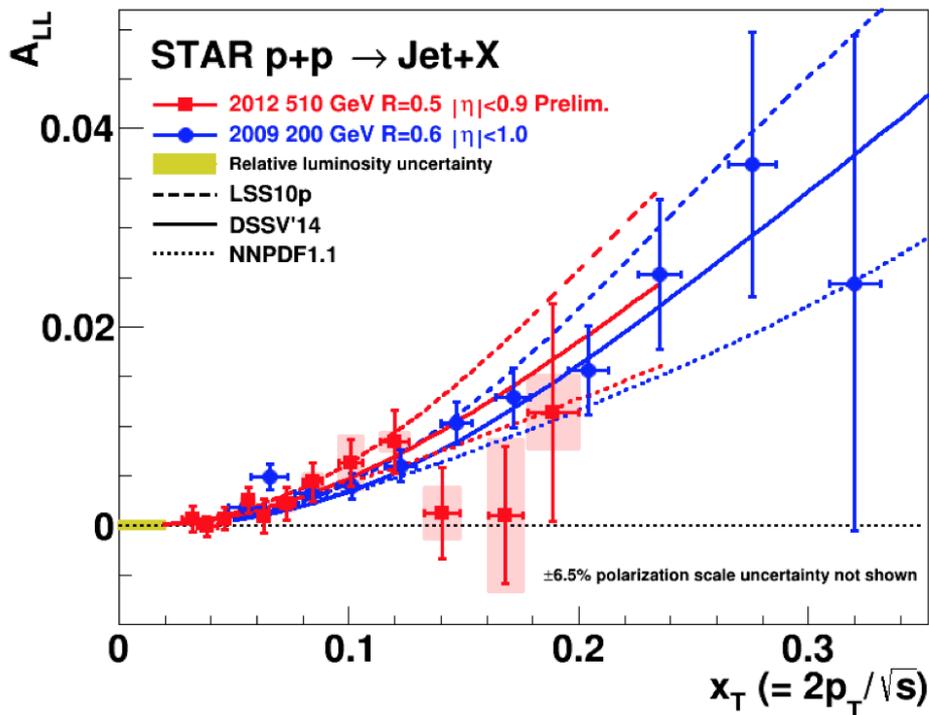
For most RHIC kinematics, **gg** and **qg** dominate, making A_{LL} for jets and hadrons sensitive to **gluon polarization**.

Parton Level Asymmetry



- | | | | |
|---|-----------------------------------|---|-----------------------------------|
| A | $gg \rightarrow gg$ | D | $q\bar{q} \rightarrow q\bar{q}$ |
| B | $qq \rightarrow qq$ | E | $gg \rightarrow q\bar{q}$ |
| C | $q\bar{q}' \rightarrow q\bar{q}'$ | | $q\bar{q} \rightarrow gg$ |
| | $q\bar{q}' \rightarrow q\bar{q}'$ | | $q\bar{q} \rightarrow g\gamma$ |
| | $qg \rightarrow qg$ | | $q\bar{q} \rightarrow q'\bar{q}'$ |
| E | $qg \rightarrow q\gamma$ | | $q\bar{q} \rightarrow l\bar{l}$ |

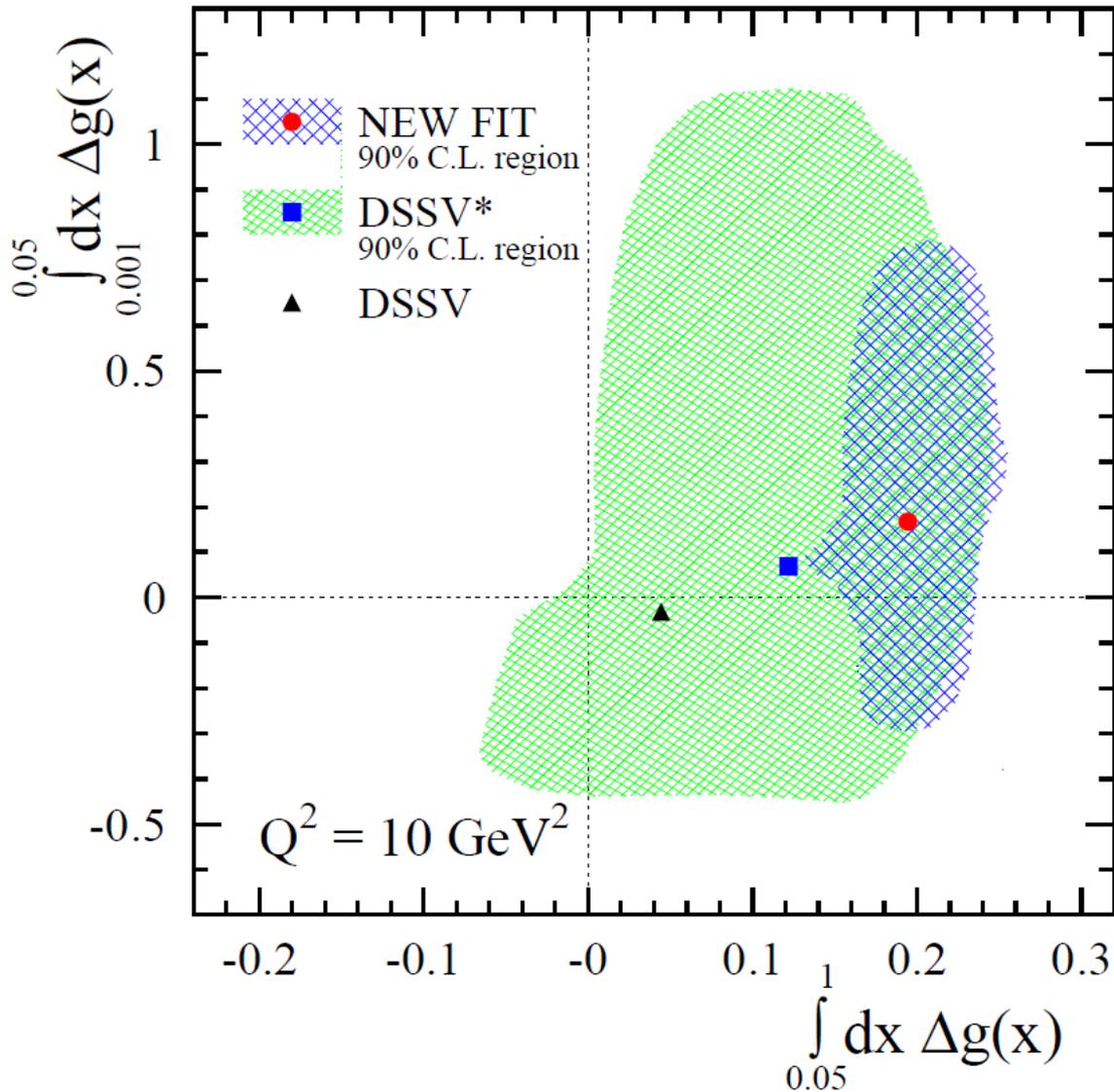
Accessing ΔG at RHIC: A_{LL}



- In polarized pp collisions, access ΔG via the longitudinal double helicity asymmetry A_{LL} which is sensitive to the polarized gluon distribution at leading order
- STAR and PHENIX measure A_{LL} using inclusive jet and π^0 final states, respectively

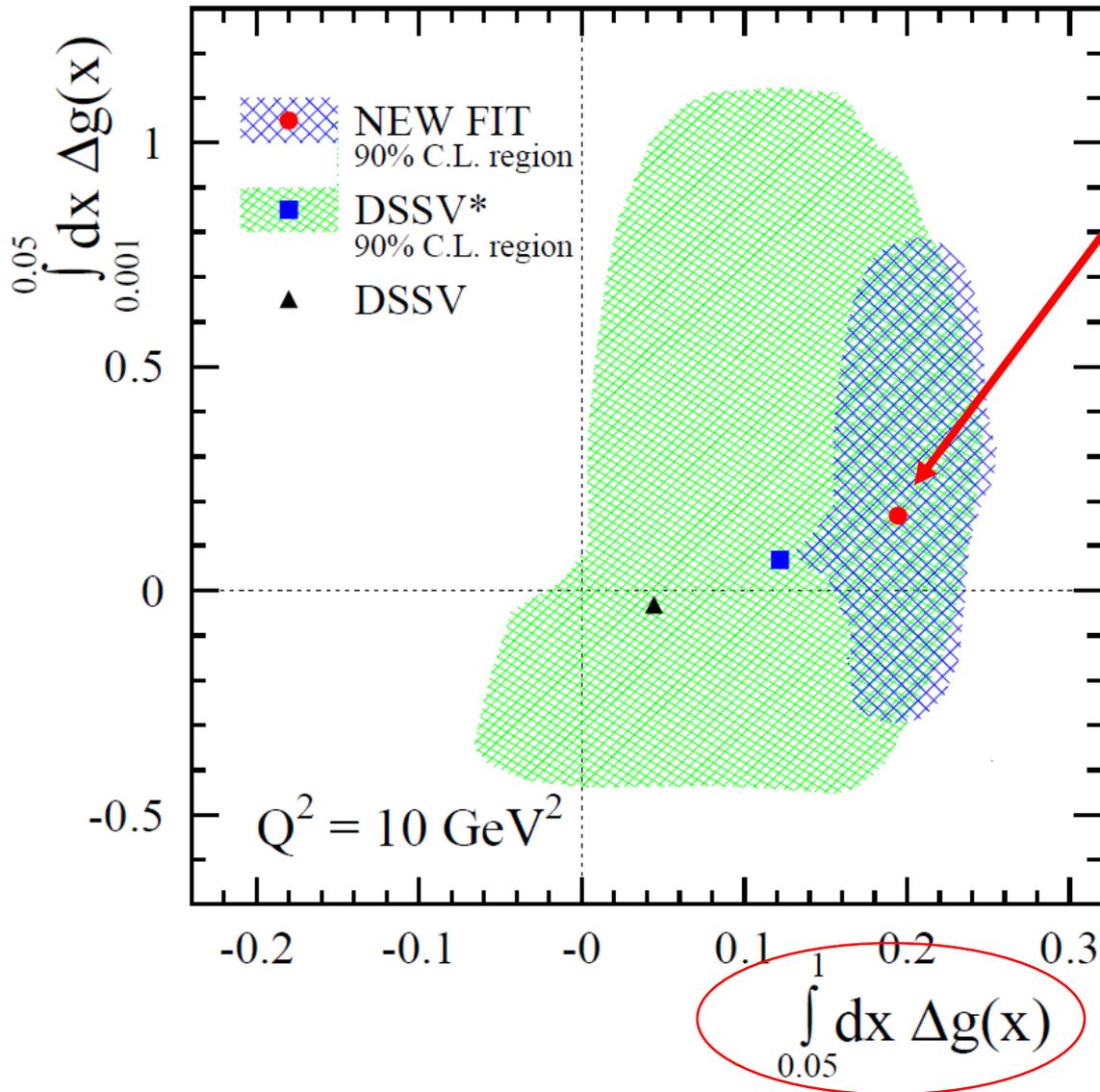
New DSSV Results

Phys. Rev. Lett. 113, 012001 (2014)



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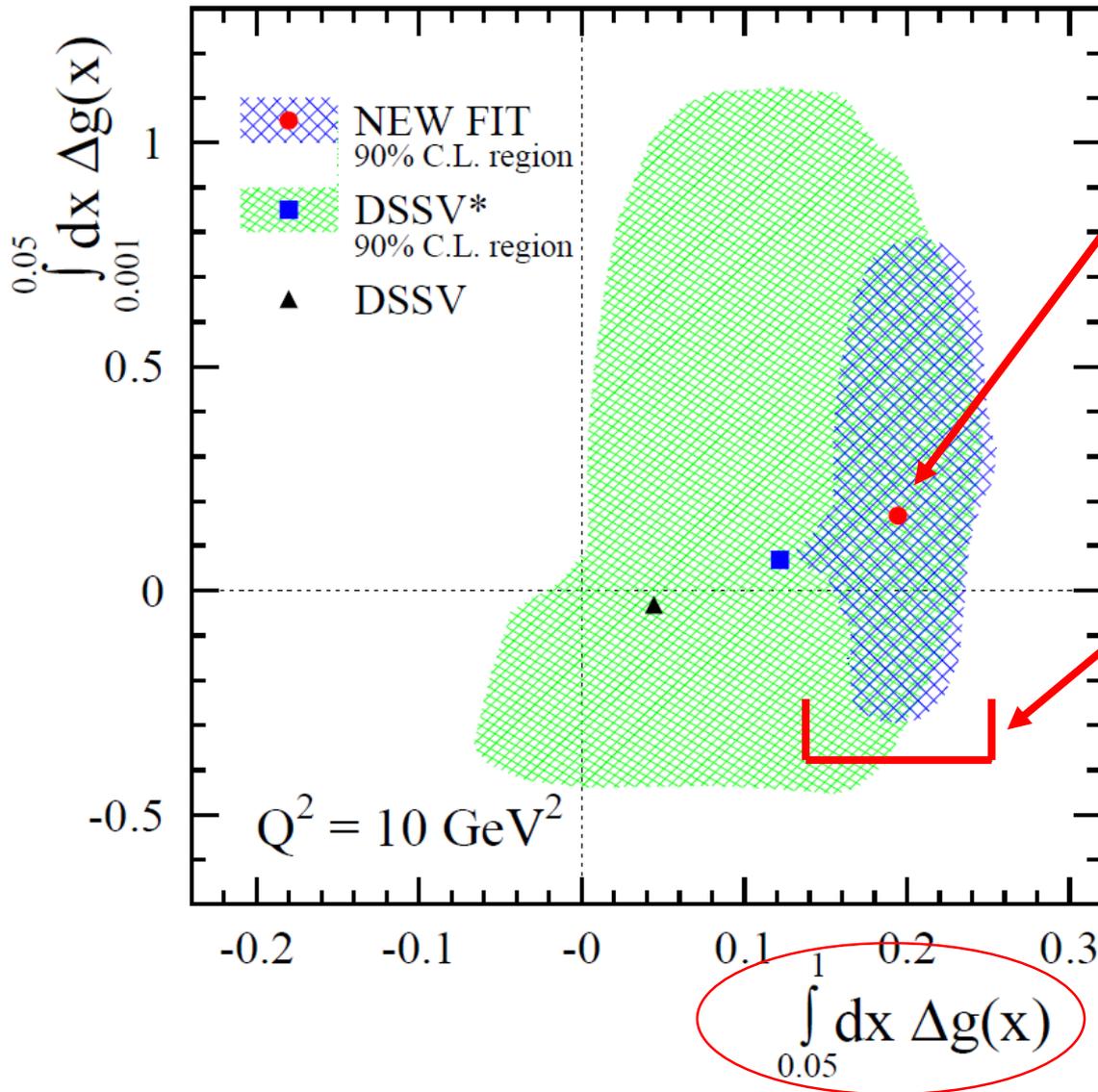
Phys. Rev. Lett. 113, 012001 (2014)



Integral of $\Delta g(x)$ in range $0.05 < x < 1.0$ increases from roughly 0.05 to $0.20^{+0.06}_{-0.07}$. First indication of non-zero gluon polarization!

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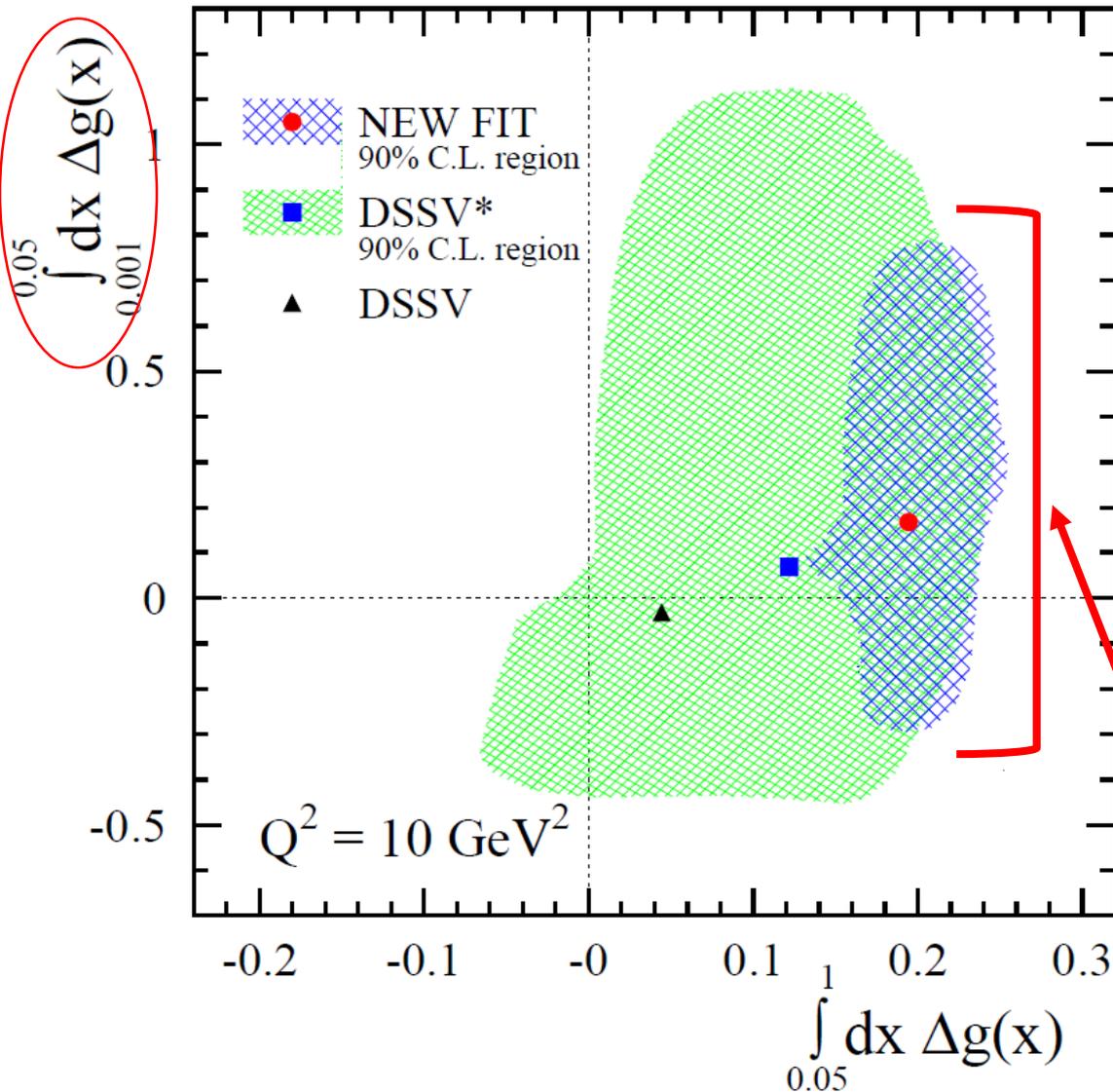
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• Uncertainty shrinks substantially from DSSV* to new DSSV fit

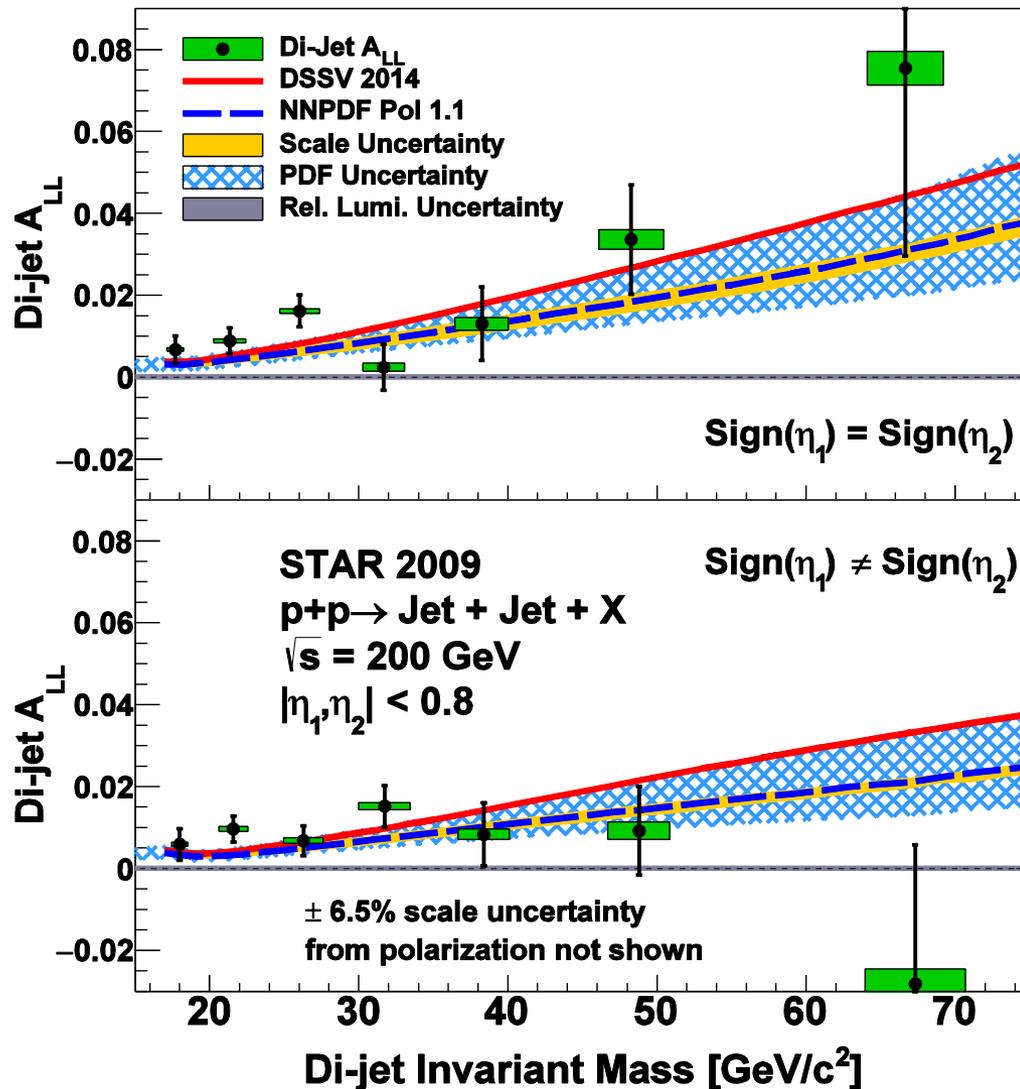
New DSSV Results



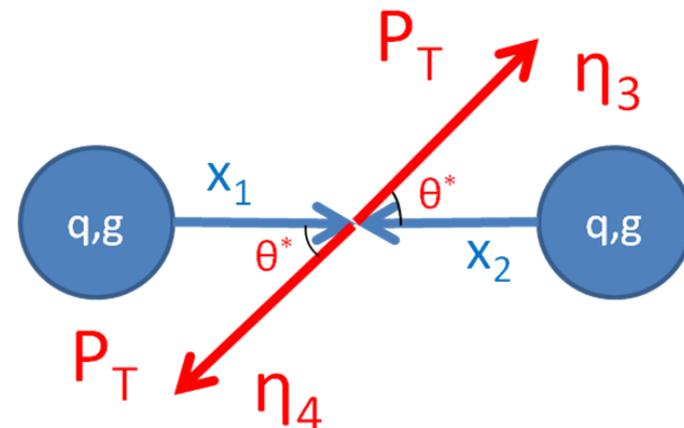
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- Integral of $\Delta g(x)$ in range $0.05 < x < 1.0$ increases from roughly 0.05 to $0.20^{+0.06}_{-0.07}$. First indication of non-zero gluon polarization!
- Uncertainty shrinks substantially from DSSV* to new DSSV fit
- Uncertainty on integral over low x region is still sizable (only $\sqrt{s} = 200$ GeV RHIC data)

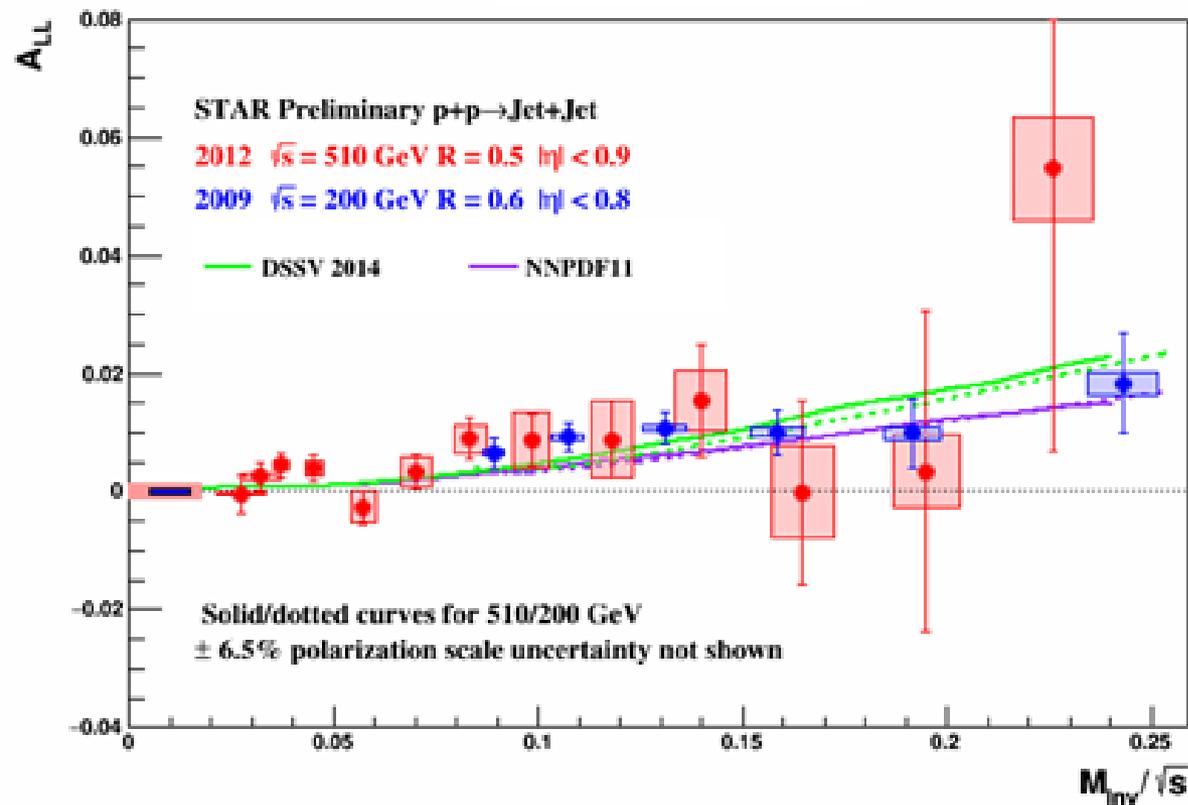
Di-jet A_{LL} (pp)



- Coincidence measurements capture more information about hard scatter and better constrain initial kinematics

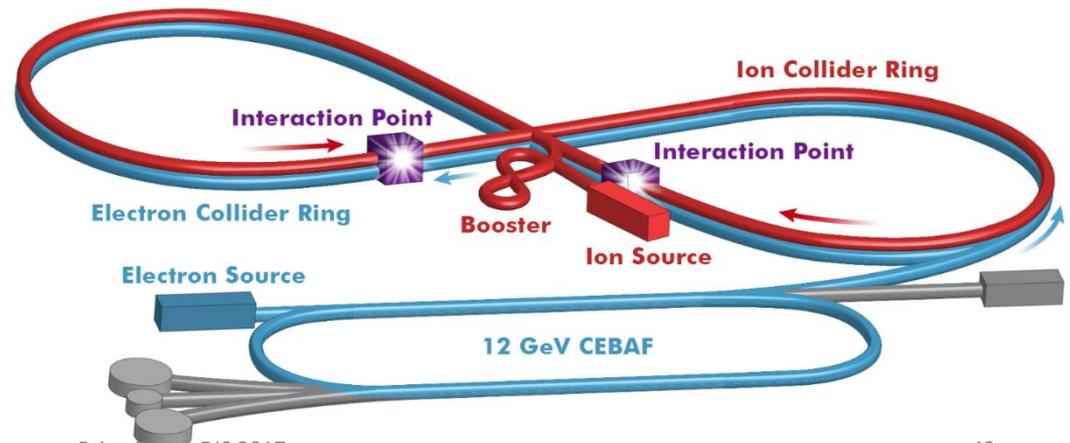
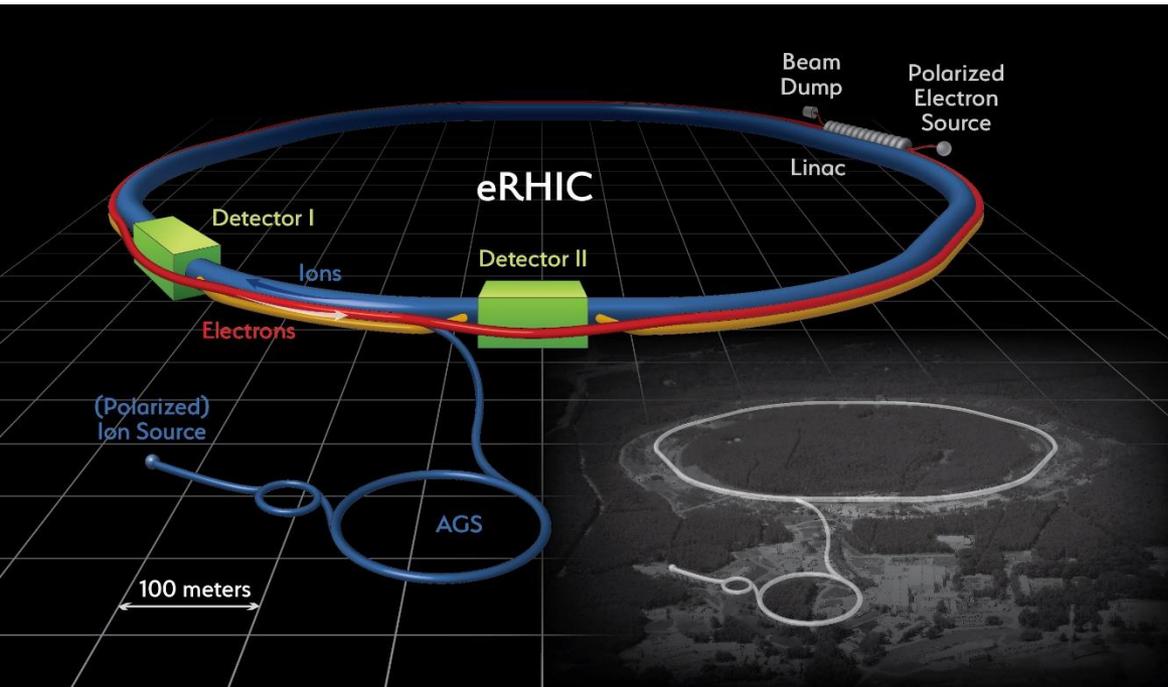


Di-jet A_{LL} (pp)



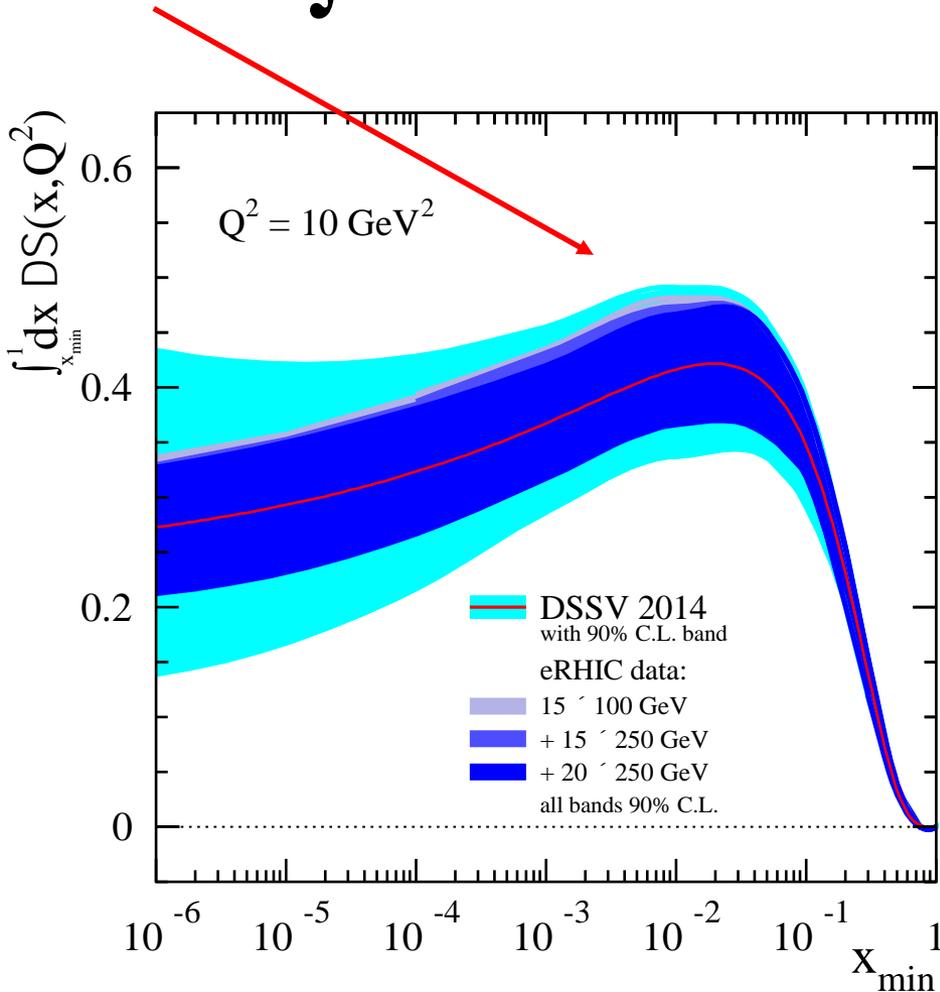
- Coincidence measurements capture more information about hard scatter and better constrain initial kinematics
- Di-jet A_{LL} plotted vs M_{inv}/\sqrt{s} ($\sim \sqrt{x_1 x_2}$ at L.O.) for data taken at $\sqrt{s} = 200$ and 510 GeV
- 510 GeV data extend to lower M_{inv}/\sqrt{s} (lower x) where ΔG not as well constrained while 200 GeV data give better precision at mid to high M_{inv}/\sqrt{s}

EIC: eRHIC or JLEIC



EIC: Impact on Quark Polarizations

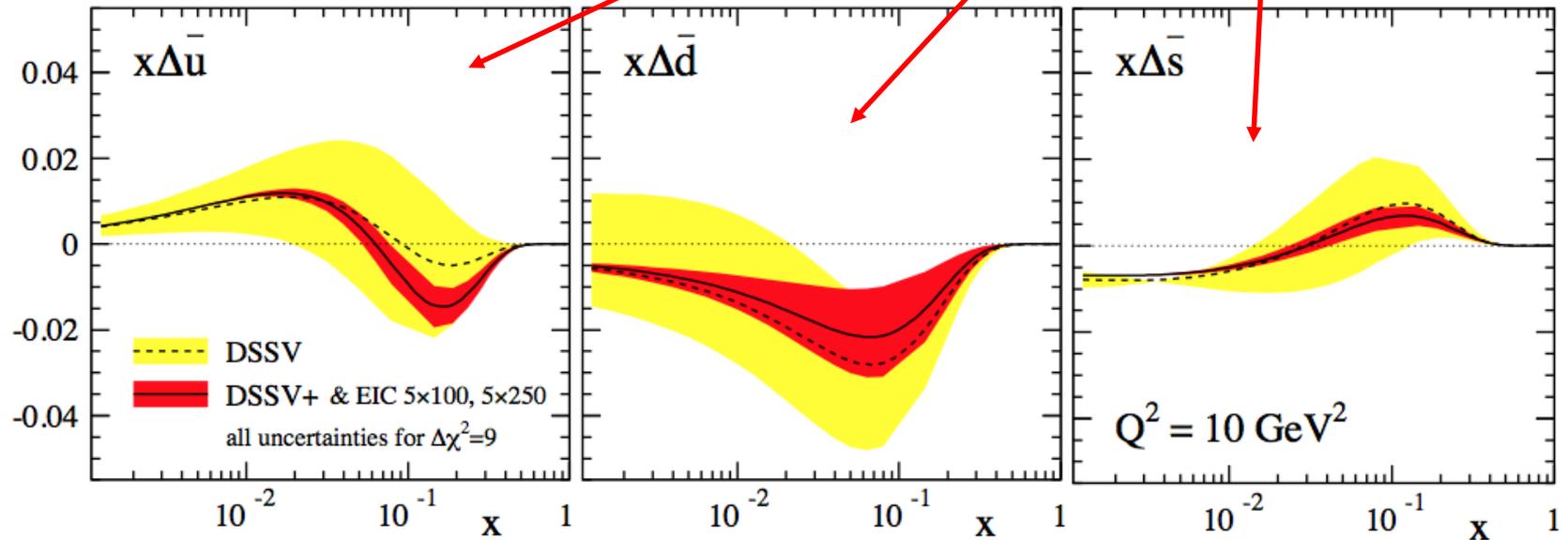
$$\Delta\Sigma = \int (\Delta u + \Delta d + \Delta s + \Delta\bar{u} + \Delta\bar{d} + \Delta\bar{s} + \dots) dx$$



- g_1 is sensitive to the sum of all quark and anti-quark polarized PDFs meaning an EIC will place strong constraints on $\Delta\Sigma$
- An EIC will also be able to constrain the individual quark and anti-quark polarized PDFs via semi-inclusive DIS (SIDIS) measurements
- The polarized anti-quark distributions are of particular interest as they provide information on non-perturbative aspects of proton structure

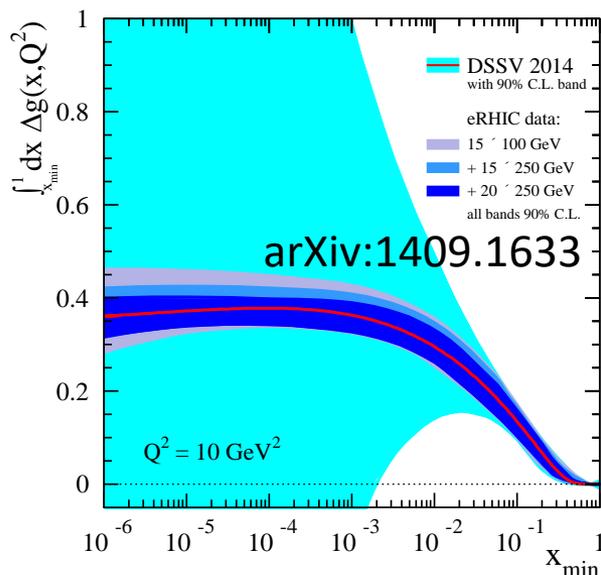
EIC: Impact on Quark Polarizations

$$\Delta\Sigma = \int (\Delta u + \Delta d + \Delta s + \Delta\bar{u} + \Delta\bar{d} + \Delta\bar{s} + \dots) dx$$



- The above plots show the expected reduction in uncertainty for the polarized anti-quark distributions from EIC SIDIS data
- Individual quark and anti-quark distributions can also be measured at an EIC via charged current DIS which access different combinations of PDFs

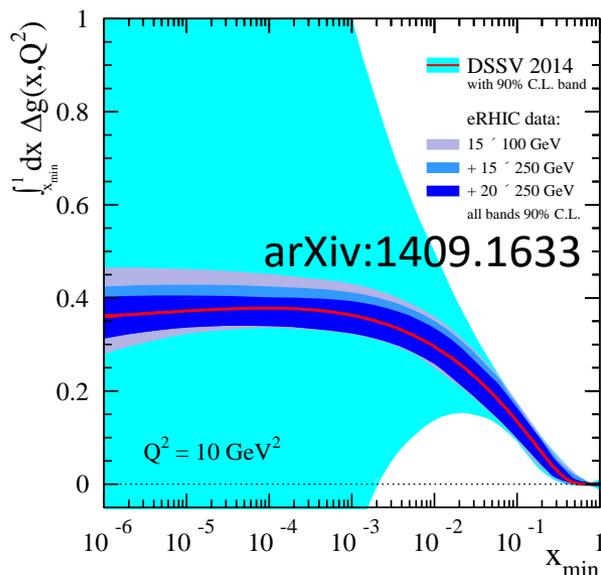
EIC: Solving the Spin Puzzle



Gluon

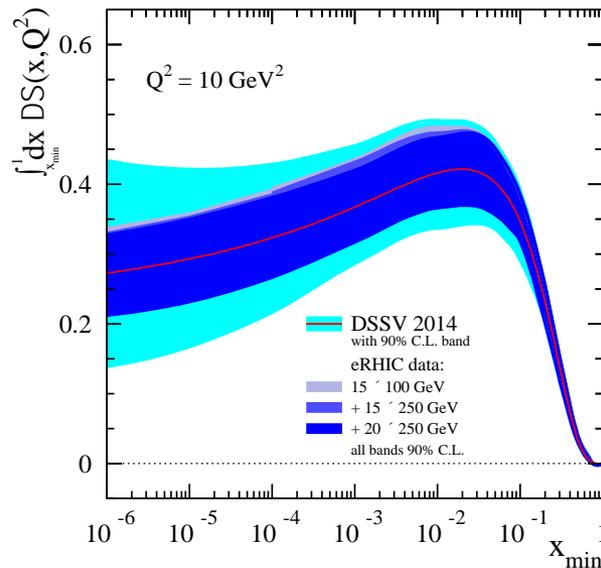
- Above plot shows the running integral of $\Delta g(x, Q^2)$ from x_{\min} to 1 as a function of x_{\min}
- Large reduction in uncertainty on ΔG from EIC can be seen

EIC: Solving the Spin Puzzle



Gluon

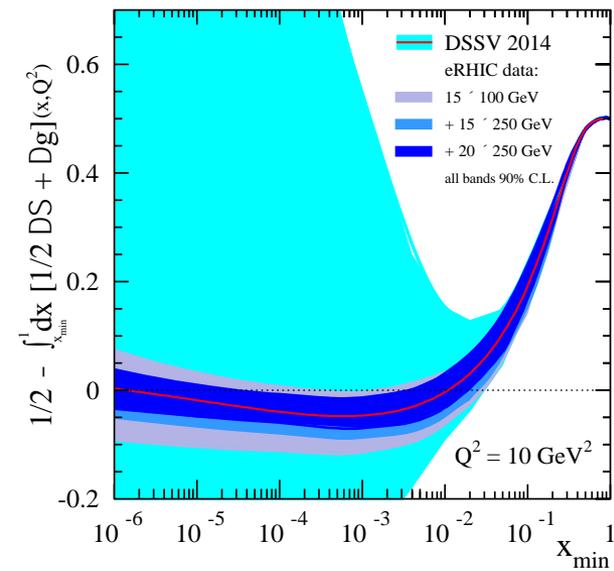
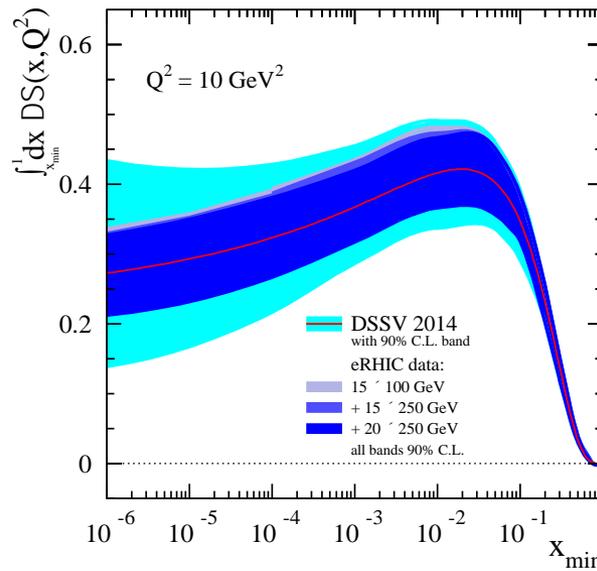
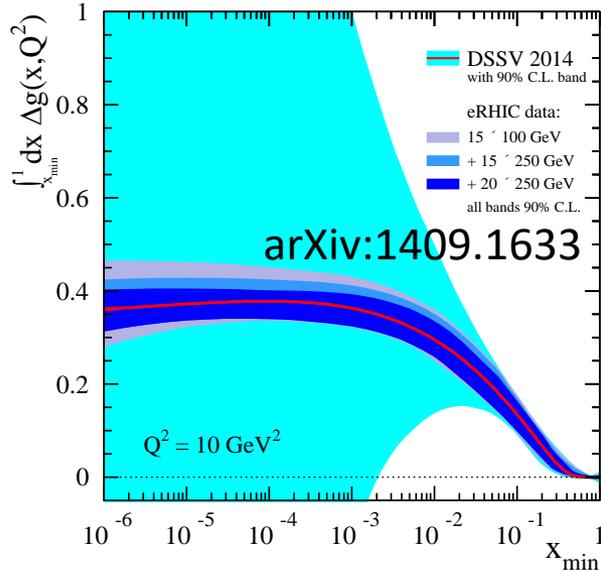
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Quarks

- EIC will also reduce the uncertainty on the quark contribution to the proton spin
- No assumptions about hyperon beta decay in EIC uncertainty

EIC: Solving the Spin Puzzle



$$\frac{1}{2} - \text{Gluon} - \text{Quarks} =$$

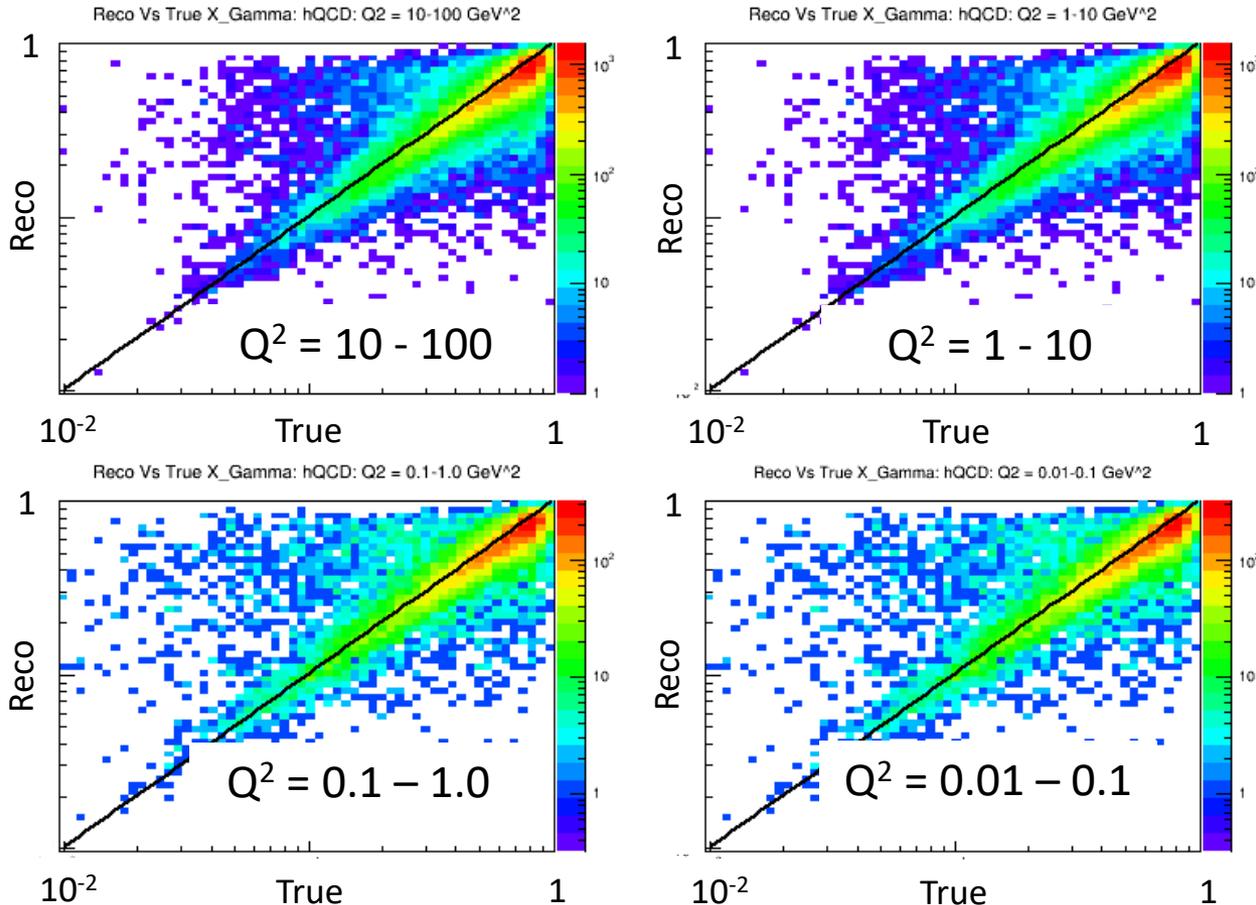
orbital angular momentum

- Above plot shows the running integral of $\Delta g(x, Q^2)$ from x_{min} to 1 as a function of x_{min}
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- EIC will also reduce the uncertainty on the quark contribution to the proton spin
- No assumptions about hyperon beta decay in EIC uncertainty

Constraints on gluon and quark contributions will provide information on the orbital angular momentum component of proton spin

X_γ : Reconstructed Vs True

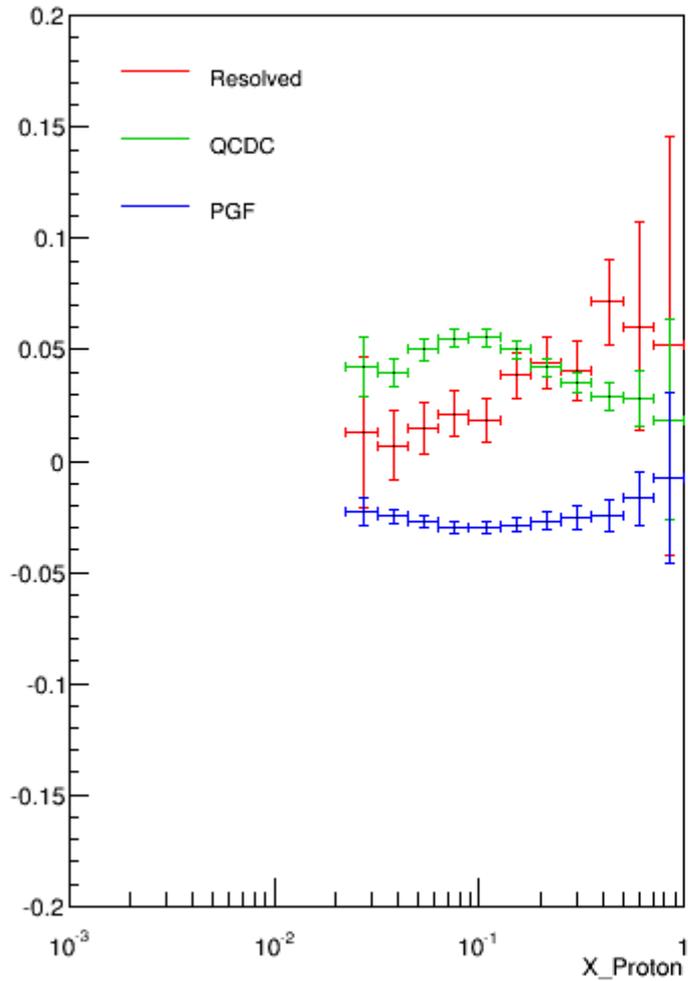


- Will use virtual photon momentum fraction to discriminate between resolved and direct processes
- See good agreement between reconstructed and true X_γ for all Q^2 ranges
- Di-jets found in Breit frame and required one jet with $p_T \geq 5 \text{ GeV}$ and the other with $p_T \geq 4 \text{ GeV}$

$$X_\gamma = \frac{1}{2E_e y} (m_{T1} e^{-y_1} + m_{T2} e^{-y_2})$$

Asymmetry

A_LL Vs Reco X_P: Q2 = 10-100: 20<Mass<40



Total A_LL Vs Reco X_P: Q2 = 10-100: 20<Mass<40

