STUFF about the Two Higgs Doublet Model Pedro Ferreira

ISEL and CFTC, UL

Higgs



Higgs Days - Santander, 20/09/2016



The STUFF

(random selection of things Pedro's interested in)

- Symmetries of the 2HDM: *what a waste*!
- One-loop vacuum structure of 2HDM.
- 750 GeV anomaly and 2HDM: a tale of how a model *nearly* died.
- tth how LHC is looking at it sideways.
- Run II amazing data: the importance of "h" precision studies.

• LHC discovered a new particle with mass ~125 GeV.

• Up to now, all is compatible with the Standard Model (SM) Higgs particle.

BORING!

Two-Higgs Dublet model, 2HDM (Lee, 1973) : one of the easiest extensions of the SM, with a richer scalar sector. Can help explain the matter-antimatter asymmetry of the universe, provide dark matter candidates, ...

> G.C. Branco, P.M. Ferreira, L. Lavoura, M. Rebelo, M. Sher, J.P Silva, Physics Reports 716, 1 (2012)

The Two-Higgs Doublet potential

Most general SU(2) \times U(1) scalar potential:

$$V_{-} = m_{11}^{2} \Phi_{1}^{\dagger} \Phi_{1} + m_{22}^{2} \Phi_{2}^{\dagger} \Phi_{2} - [m_{12}^{2} \Phi_{1}^{\dagger} \Phi_{2} + \text{H.c.}] + \frac{1}{2} \lambda_{1} (\Phi_{1}^{\dagger} \Phi_{1})^{2} + \frac{1}{2} \lambda_{2} (\Phi_{2}^{\dagger} \Phi_{2})^{2} + \lambda_{3} (\Phi_{1}^{\dagger} \Phi_{1}) \times (\Phi_{2}^{\dagger} \Phi_{2}) + \lambda_{4} (\Phi_{1}^{\dagger} \Phi_{2}) (\Phi_{2}^{\dagger} \Phi_{1}) + [\frac{1}{2} \lambda_{5} (\Phi_{1}^{\dagger} \Phi_{2})^{2} + \lambda_{6} (\Phi_{1}^{\dagger} \Phi_{1}) (\Phi_{1}^{\dagger} \Phi_{2}) + \lambda_{7} (\Phi_{2}^{\dagger} \Phi_{2}) (\Phi_{1}^{\dagger} \Phi_{2}) + \text{H.c.}]_{e}$$

 m_{12}^2 , λ_5 , λ_6 and λ_7 complex - seemingly 14 independent real parameters

Most frequently studied model: softly broken theory with a Z₂ symmetry,

MODEL I: Only Φ_2 couples to fermions.

MODEL II: Φ_2 couples to up-quarks, Φ_1 to down quarks and leptons.

Symmetries of the potential of 2HDM
Higgs Family
Symmetries:

$$\begin{array}{c}
\mathbf{Z}_{2}: \quad \Phi_{1} \rightarrow \Phi_{1} \quad , \quad \Phi_{2} \rightarrow -\Phi_{2} \\
\mathbf{U}(1): \quad \Phi_{1} \rightarrow \Phi_{1} \quad , \quad \Phi_{2} \rightarrow e^{i\theta} \Phi_{2} \qquad \theta \neq \{0, \pi\} \\
\mathbf{U}(2): \begin{pmatrix} \Phi_{1} \\ \Phi_{2} \end{pmatrix} \rightarrow U \begin{pmatrix} \Phi_{1} \\ \Phi_{2} \end{pmatrix} \qquad \forall_{U \in U(2)} \\
\begin{array}{c}
\mathbf{CP1}: \quad \Phi_{1} \rightarrow \Phi_{1}^{*} \quad , \quad \Phi_{2} \rightarrow \Phi_{2}^{*} \\
\mathbf{CP2}: \quad \Phi_{1} \rightarrow \Phi_{1}^{*} \quad , \quad \Phi_{2} \rightarrow \Phi_{2}^{*} \\
\mathbf{CP2}: \quad \Phi_{1} \rightarrow \Phi_{2}^{*} \quad , \quad \Phi_{2} \rightarrow -\Phi_{1}^{*} \\
\mathbf{CP3}: \begin{pmatrix} \Phi_{1} \\ \Phi_{2} \end{pmatrix} \rightarrow \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \Phi_{1}^{*} \\ \Phi_{2}^{*} \end{pmatrix} \quad 0 < \theta < \pi/2
\end{array}$$

Symmetries of the potential of 2HDM

symmetry m_{11}^2	m_{22}^2	$m_{12}^2 \lambda_1$	$\lambda_2 \ \lambda_3$	λ_4	λ_5	λ_6	λ_7			
Z_2		0			real	0	0	7		
U(1)		0			0	0	0	6		
U(2)	m_{11}^2	0	λ_1	$\lambda_1 - \lambda_3$	0	0	0	3		
CP1		real			real	real	λ_6	10		
CP2	m_{11}^2	0	λ_1		real	0	0	5		
CP3	m_{11}^2	0	λ_1		$\lambda_1 - \lambda_3 - \lambda_4$ (real)	0	0	4		
$V = V = m_{11}^{2}(\varphi_{1}^{\dagger}\varphi_{1}) + m_{11}^{2}(\varphi_{2}^{\dagger}\varphi_{2}) + \frac{1}{2}\lambda_{1}(\varphi_{1}^{\dagger}\varphi_{1})^{2} + \frac{1}{2}\lambda_{1}(\varphi_{2}^{\dagger}\varphi_{2})^{2} + \lambda_{3}(\varphi_{1}^{\dagger}\varphi_{1})(\varphi_{2}^{\dagger}\varphi_{2}) + (\lambda_{1}-\lambda_{3})(\varphi_{1}^{\dagger}\varphi_{2})(\varphi_{2}^{\dagger}\varphi_{1}) $										

Symmetries of the LAGRANGIAN of 2HDM

symmetry	m_{11}^2	m_{22}^2	m_{12}^2	λ_1	λ_2	λ_3	λ_4	λ_5	λ_6	λ_7	
Z_2			0					real	0	0	7
U(1)			0					0	0	0	6
U(2)	1	m_{11}^2	0		λ_1		$\lambda_1 - \lambda_3$	0	0	0	3
CP1			real					real	real	λ_6	10
CP2	1	m_{11}^2	0		λ_1			real	0	0	5
CP3		m_{11}^2	0		λ_1			$\lambda_1 - \lambda_3 - \lambda_4$ (real)	0	0	4

Three generations of massive fermions: CP1, Z₂, U(1) and CP3 (but bad CKM!)

Plus absence of tree-level FCNC: Z_2 , U(1)

Remainder USELESS...? Not necessarily so...

One-loop contributions to inert minima in 2HDM

Z₂-symmetric model:

$$V = m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - m_{12}^2 \left(\Phi_1^{\dagger} \Phi_2 + h.c. \right)$$

$$+ \frac{1}{2} \lambda_1 |\Phi_1|^4 + \frac{1}{2} \lambda_2 |\Phi_2|^4 + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 |\Phi_1^{\dagger} \Phi_2|^2 + \frac{1}{2} \lambda_5 \left[\left(\Phi_1^{\dagger} \Phi_2 \right)^2 + h.c. \right]$$

Tree-level vacuum solutions:

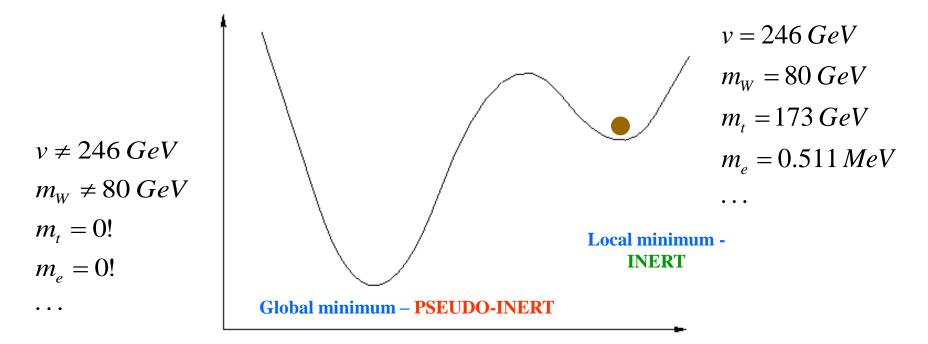
INERT:
$$v_1^2 = -\frac{2m_{11}^2}{\lambda_1}$$
, provided $m_{11}^2 < 0$.

FERMIONS MASSIVE – SM-LIKE PHENOMENOLOGY

INERT-LIKE:
$$v_2^2 = -\frac{2 m_{22}^2}{\lambda_2}$$
, provided $m_{22}^2 < 0$.

FERMIONS MASSLESS – UNPHYSICAL VACUUM

These minima can coexist in the potential, which raises a troubling possibility...



Tree-Level Conclusions:

Inert and inert-like minima can coexist in the potential if $m_{11}^2 < 0$ and $m_{22}^2 < 0$.

$$V_{I} - V_{IL} = \frac{1}{2} \left(\frac{m_{22}^{4}}{\lambda_{2}} - \frac{m_{11}^{4}}{\lambda_{1}} \right)$$
$$= \frac{1}{4} \left[\left(\frac{m_{H^{\pm}}^{2}}{v_{2}^{2}} \right)_{IL} - \left(\frac{m_{H^{\pm}}^{2}}{v_{1}^{2}} \right)_{I} \right] v_{1}^{2} v_{2}^{2}$$

Tree-level to one-loop...

 $V = V_0 + V_1,$

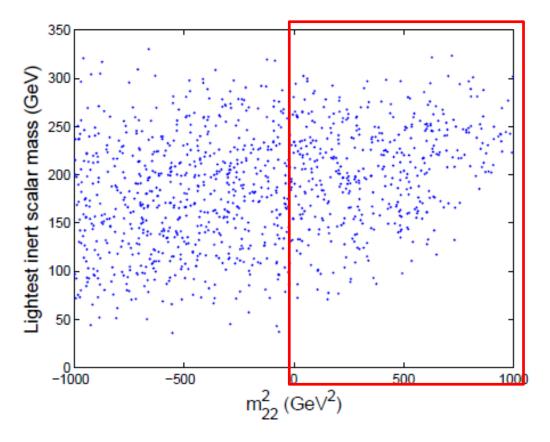
$$V_{1} = \frac{1}{64\pi^{2}} \sum_{\alpha} n_{\alpha} m_{\alpha}^{4}(\varphi_{i}) \left[\log\left(\frac{m_{\alpha}^{2}(\varphi_{i})}{\mu^{2}}\right) - \frac{3}{2} \right]$$
$$\frac{\partial V}{\partial \varphi_{i}} = \frac{\partial V_{0}}{\partial \varphi_{i}} + \frac{1}{32\pi^{2}} \sum_{\alpha} n_{\alpha} m_{\alpha}^{2} \frac{\partial m_{\alpha}^{2}}{\partial \varphi_{i}} \left[\log\left(\frac{m_{\alpha}^{2}}{\mu^{2}}\right) - 1 \right]$$
$$m_{h}^{2} = m_{h_{0}}^{2} + \frac{1}{32\pi^{2}} m_{h_{1}}^{2},$$

 $m_{h_{1,S}}^{2} = \lambda_{1}A(m_{G_{0}}) + 2\lambda_{1}A(m_{G_{0}^{\pm}}) + 3\lambda_{1}A(m_{h_{0}}) + \lambda_{345}A(m_{H_{0}}) + \bar{\lambda}_{345}A(m_{A_{0}}) + 2\lambda_{3}A(m_{H_{0}^{\pm}})$ $+ \lambda_{1}^{2}v_{1}^{2}B(m_{G_{0}}, m_{G_{0}}, p^{2}) + 2\lambda_{1}^{2}v_{1}^{2}B(m_{G_{0}^{\pm}}, m_{G_{0}^{\pm}}, p^{2}) + 9\lambda_{1}^{2}v_{1}^{2}B(m_{h_{0}}, m_{h_{0}}, p^{2})$ $+ \lambda_{345}^{2}v_{1}^{2}B(m_{H_{0}}, m_{H_{0}}, p^{2}) + \bar{\lambda}_{345}^{2}v_{1}^{2}B(m_{A_{0}}, m_{A_{0}}, p^{2}) + 2\lambda_{3}^{2}v_{1}^{2}B(m_{H_{0}^{\pm}}, m_{H_{0}^{\pm}}, p^{2})$

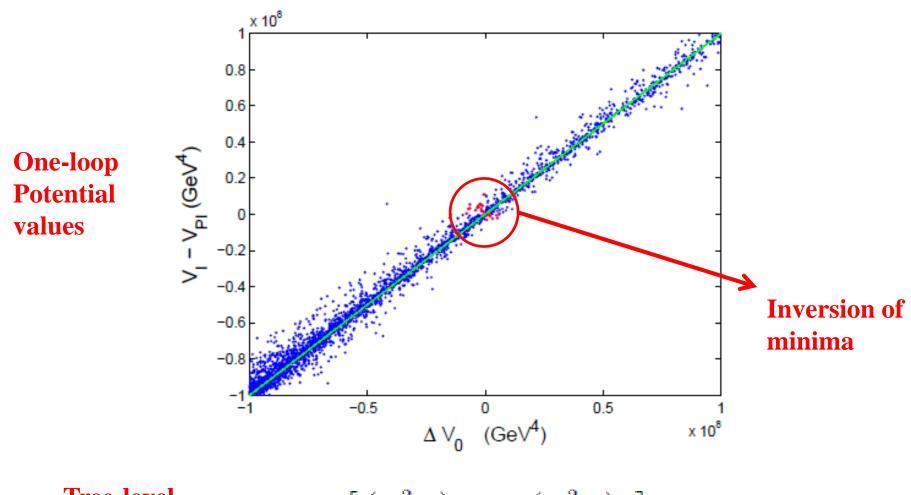
Tree-level results:

Inert and inert-like minima can coexist in the potential if $m_{11}^2 < 0$ and $m_{22}^2 < 0$.

One-loop results:

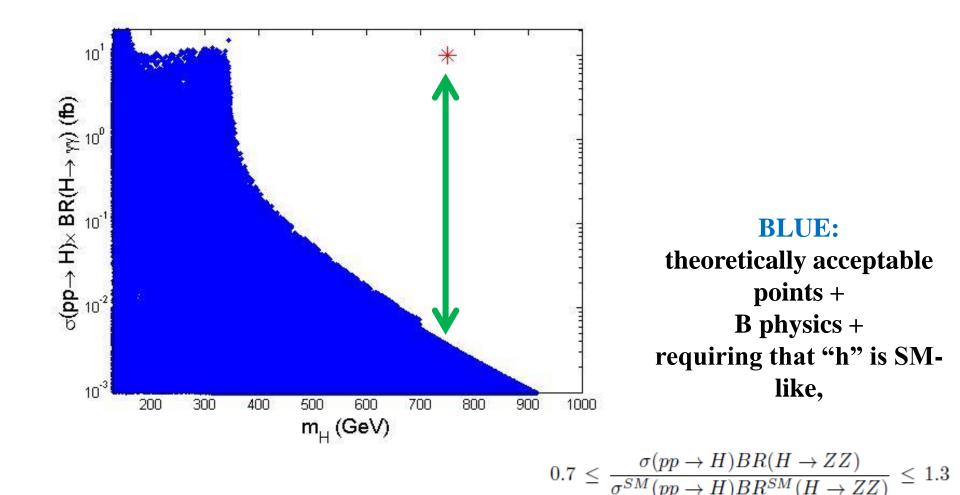


IMPOSSIBLE TO HAVE SIMULTANEOUS MINIMA IN THIS REGION AT TREE-LEVEL!

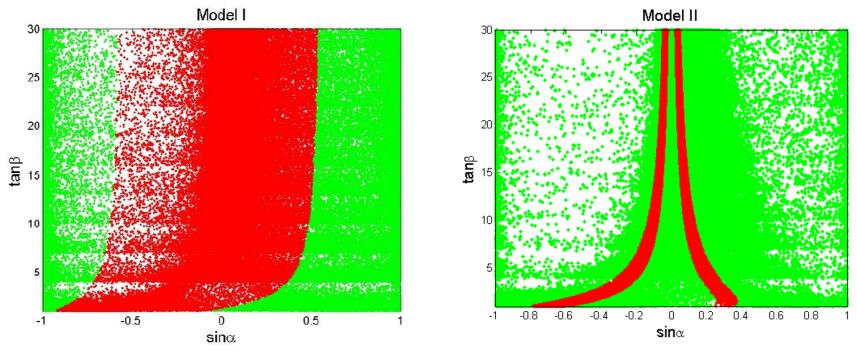


Tree-level Potential $\Delta V_0 = \frac{1}{4} \left[\left(\frac{m_{H^{\pm}}^2}{v_2^2} \right)_{IL} - \left(\frac{m_{H^{\pm}}^2}{v_1^2} \right)_I \right] v_1^2 v_2^2$ values

The 750 GeV WHATEVER...



Run-I parameter space restrictions



Model I

$$k_t = \frac{\cos \alpha}{\sin \beta} > 0$$
$$k_b = \frac{\cos \alpha}{\sin \beta} > 0$$

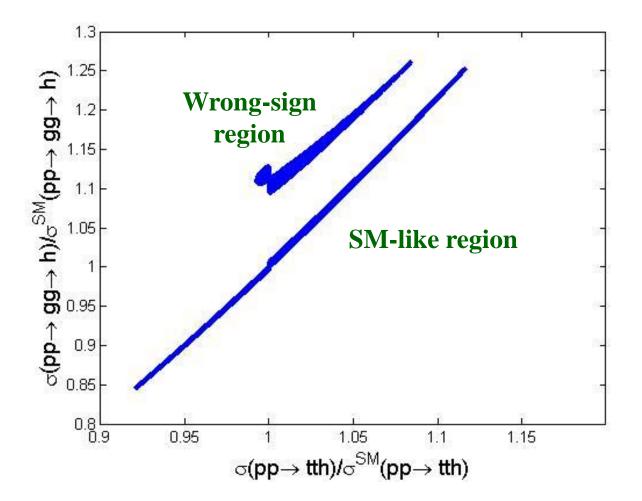
Model II

$$k_t = \frac{\cos \alpha}{\sin \beta} > 0$$

$$k_b = \frac{\sin \alpha}{\cos \beta} > 0 \text{ or } < 0$$

Wrong-Sign Limit

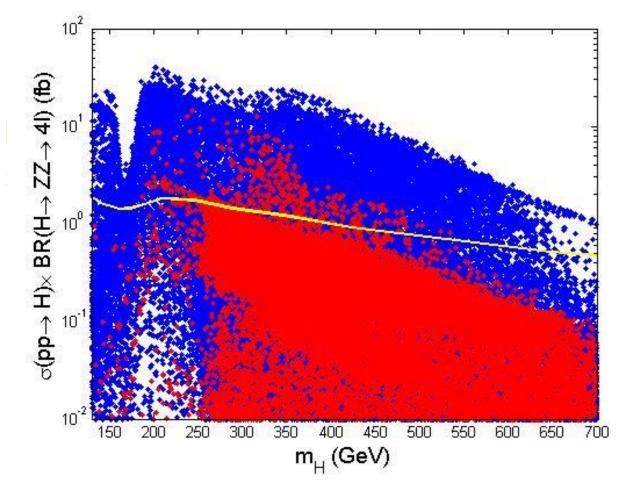
σ_{tth} versus σ_{ggh} in model II



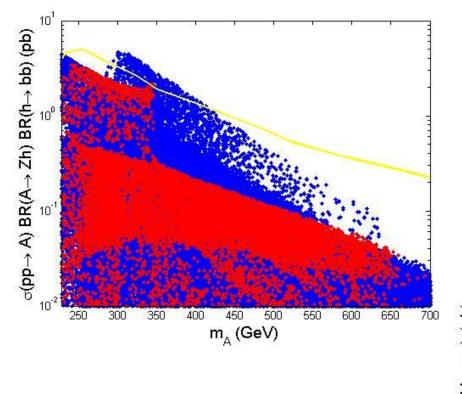
LHC – YOU HAVE THE WRONG RESULT!

The Importance of Being Earnest h

Run II has limits on high mass resonances in the 4 lepton channel...

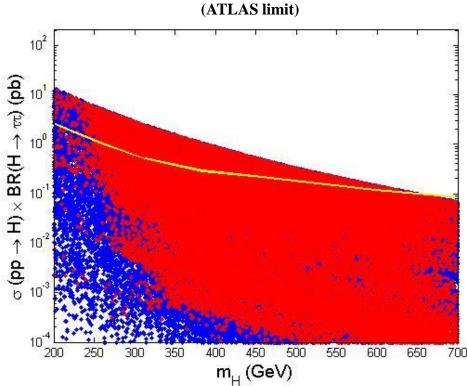


(yellow line upper bound on non-observation from CMS PAS HIG-16-033) (red points are what remains after demanding "h" rates are within 30% of SM values) (ATLAS limit)



... Though not for ALL observables

Demanding "h" behaviour being SMlike complies with latest high-mass exclusions...



NOT AWESOME CONCLUSIONS

• STILL NO EXTRA SCALARS!

SO-SO CONCLUSIONS

• 750 GeV a bust... But we learned quite a lot from it.

AWESOME CONCLUSIONS

- Run II already providing amazing data.
- Already excluding significant regions of 2HDM parameter space.

NOT-SO-AWESOME CONCLUSIONS

• All of this is tree-level. Possibility of RG-improving? Need to use 1-Loop effective potential?

• No cosmological considerations undertaken – that calculation is arguable...

• Discriminant exists for more complicated versions of the model (no Z_2 symmetry) but it can no longer be cast in a nice analytical expression.

