

Phenomenological signatures of additional scalar bosons at the LHC

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Outline

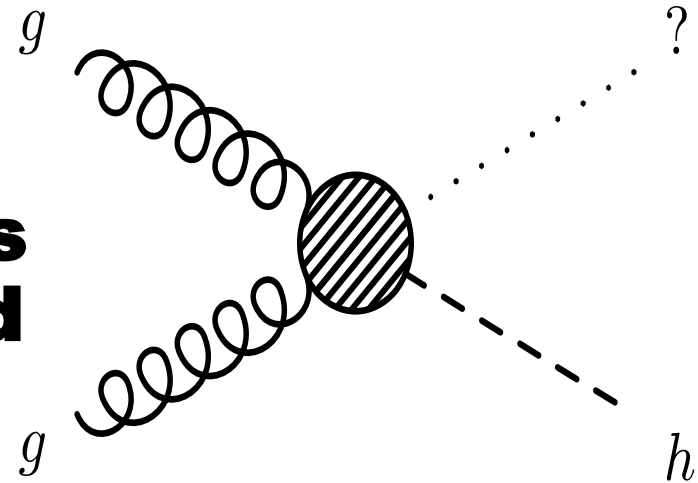


- **The Effective Lagrangian**
- **Study with Run I data**
- **Formulation of the hypothesis**
- **Prediction of signatures at the LHC**
 - **Three same sign leptons**
 - **4 W production**

Views expressed here are of the authors only

Bottom-up approach: What if?

- Initially were interested in investigating the Higgs boson transverse momentum
- What if the Higgs boson is being also being produced in association with something else...
- What can we fill the blob with?



The Lagrangian

arXiv:1506.00612

arXiv:1603.01208

arXiv:1606.01674

Introduce H and X fields with the interactions listed below

$$\mathcal{L} = \mathcal{L}_{SM} + \mathcal{L}_{BSM}$$

$$\mathcal{L}_{BSM} = \mathcal{L}_K + \mathcal{L}_T + \mathcal{L}_Q + \mathcal{L}_{Hgg} + \mathcal{L}_{HVV}$$

$$\mathcal{L}_K = \frac{1}{2} \partial_\mu X \partial^\mu X + \frac{1}{2} \partial_\mu H \partial^\mu H - \frac{1}{2} M_X^2 X^2 - \frac{1}{2} M_H^2 H^2$$

$$\mathcal{L}_T = -\frac{1}{2} \mu_1 h^2 H - \frac{1}{2} \mu_2 X^2 h - \frac{1}{2} \mu_3 X^2 H$$

$$\mathcal{L}_Q = -\frac{1}{4} \lambda_1 H^2 h^2 - \frac{1}{4} \lambda_2 X^2 h^2 - \frac{1}{4} \lambda_3 H^2 X^2 - \frac{1}{2} \lambda_4 H h X^2$$

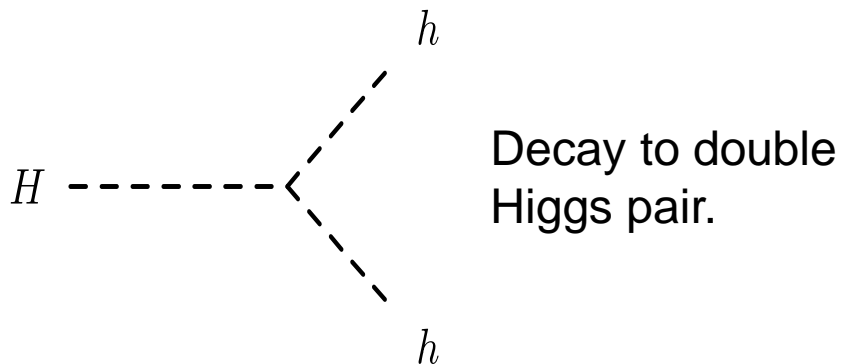
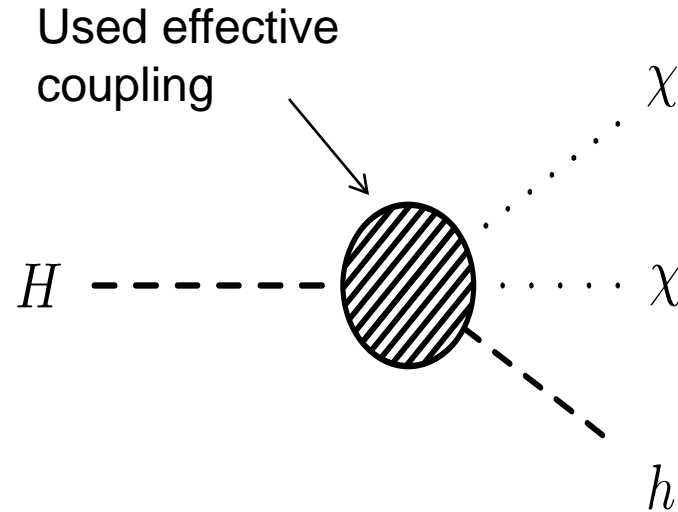
$$\mathcal{L}_{Hgg} = -\frac{1}{4} \beta_g \kappa_{hgg}^{SM} G_{\mu\nu} G^{\mu\nu} H$$

$$\mathcal{L}_{HVV} = \frac{2M_W^2}{v} \beta_W W_\mu W^\mu H + \frac{M_Z^2}{v} \beta_Z Z_\mu Z^\mu H$$

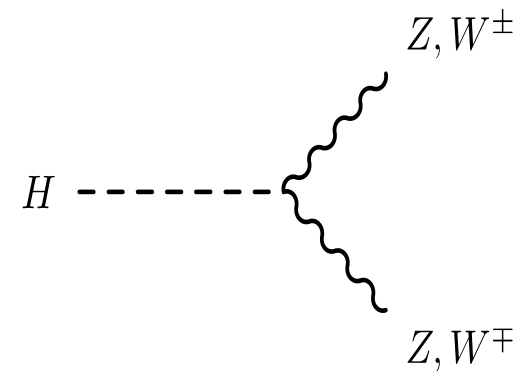
Main decay modes of H

Decay to single Higgs and a dark matter (DM) candidate

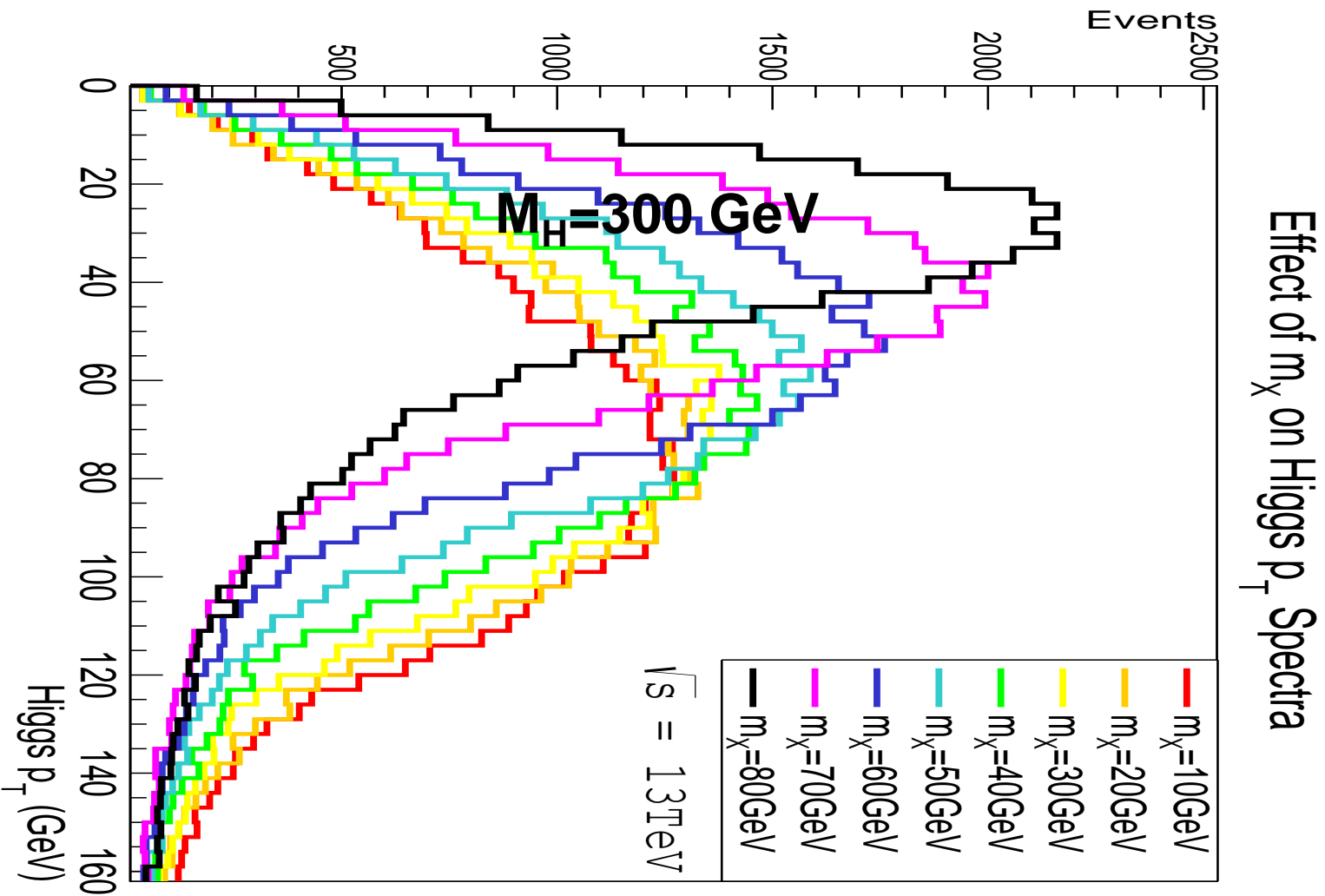
- DM is assumed scalar for simplicity
- This was our strategy, but we can infer different physics in the blob



Decay to vector boson pairs.



Higgs boson p_T Spectra

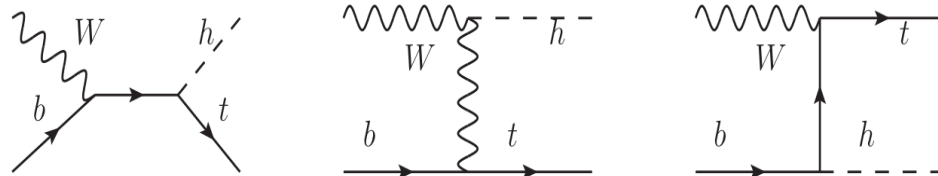


Enhancement of tH production

□ In experiment, top associated Higgs production is measured as a sum of single top and double top cross sections

□ In the SM, we find that $\sigma_{th} \ll \sigma_{tth}$

$$A = \frac{g}{\sqrt{2}} \left[(c_F - c_V) \frac{m_t \sqrt{s}}{m_W v} A \left(\frac{t}{s}, \varphi; \xi_t, \xi_b \right) + \left(c_V \frac{2m_W s}{v} \frac{1}{t} + (2c_F - c_V) \frac{m_t^2}{m_W v} \right) B \left(\frac{t}{s}, \varphi; \xi_t, \xi_b \right) \right]$$



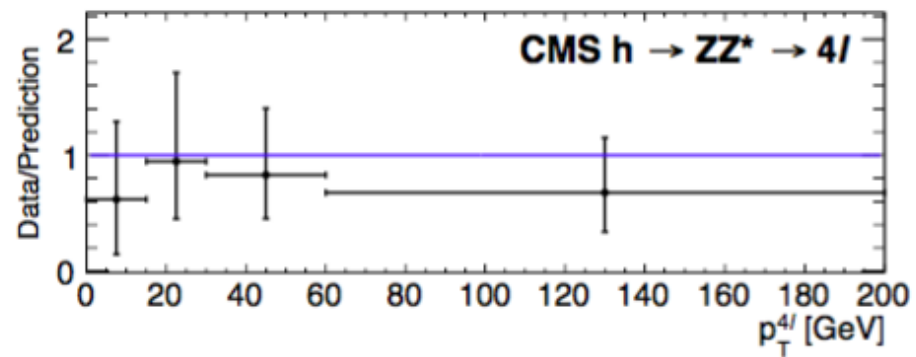
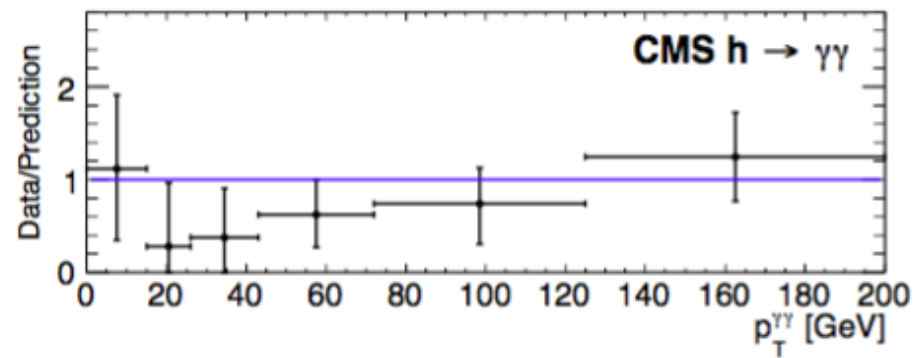
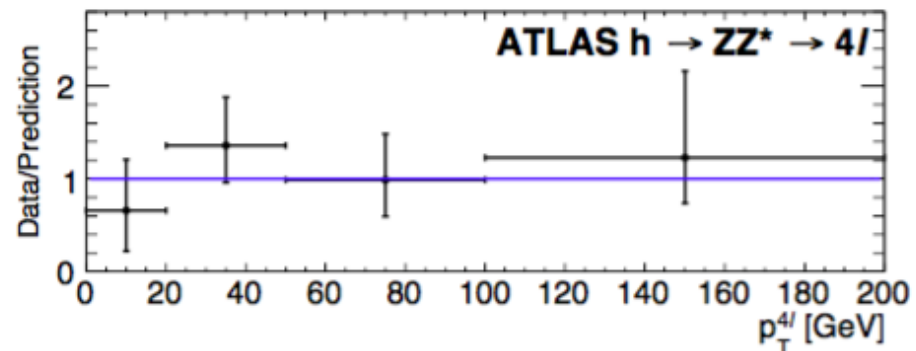
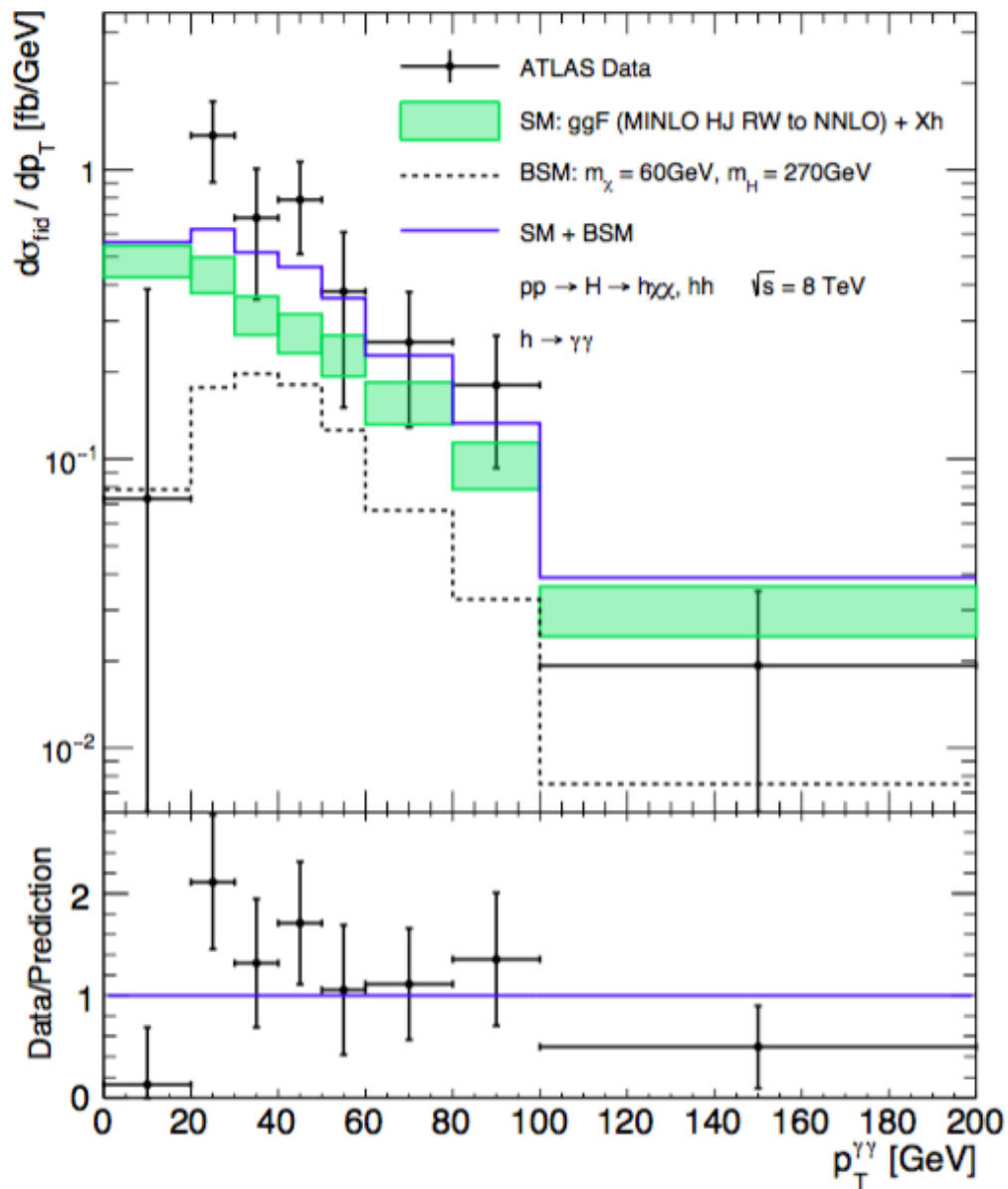
□ For the heavy scalar considered here, $c_V \ll c_F$

□ We expect a sizeable cross section to come from top associated heavy scalar production ($\sigma_{tH} \simeq \sigma_{ttH}$)

Results incorporated in the fit

Four groups of final states
received consideration

Category	Experiment	Result
Higgs p_T spectra	ATLAS	$h \rightarrow \gamma\gamma$ and $h \rightarrow ZZ$
	CMS	$h \rightarrow \gamma\gamma$ and $h \rightarrow ZZ$
Di-Higgs resonance searches	ATLAS	Limits on $H \rightarrow hh \rightarrow bb\tau\tau, \gamma\gamma WW, \gamma\gamma bb,$ and $bbbb$
	CMS	Limits on $H \rightarrow hh \rightarrow bb\tau\tau, \gamma\gamma bb,$ and multi-lepton
Top associated Higgs production	ATLAS	Limits on $h \rightarrow \gamma\gamma$ Measurements on $h \rightarrow bb,$ and multi-lepton
	CMS	Measurements on $h \rightarrow \gamma\gamma, h \rightarrow bb,$ and multi-lepton
Decays to weak vector bosons	ATLAS	Limits on $H \rightarrow ZZ$ and WW
	CMS	Limits on $H \rightarrow ZZ$ and WW



Satisfactory goodness of the global fit, including Higgs p_T

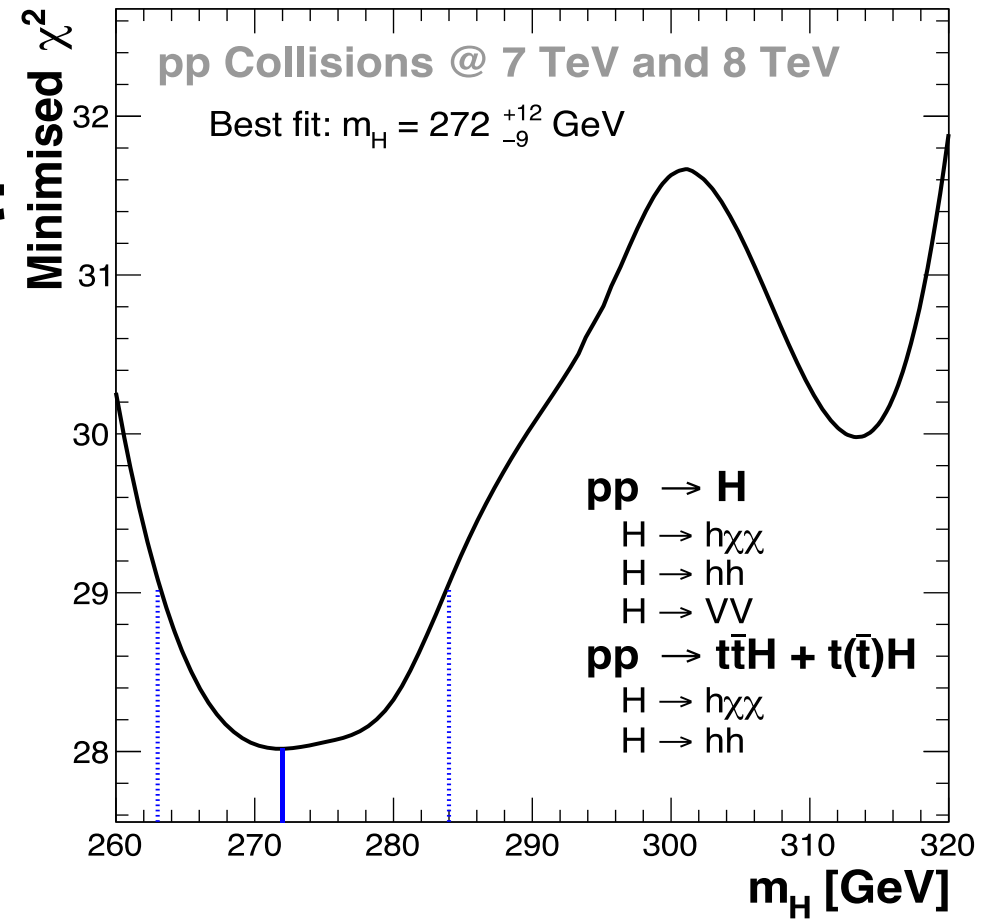
The combined result

□ **Combining all of the results produces a best fit at $m_H = 272$ GeV**

□ **The errors are +12 GeV and -9 GeV, which are one sigma deviations from the best fit point**

□ **At this point:**

$$\beta_g = 1.5 \pm 0.6$$



arXiv:1506.00612

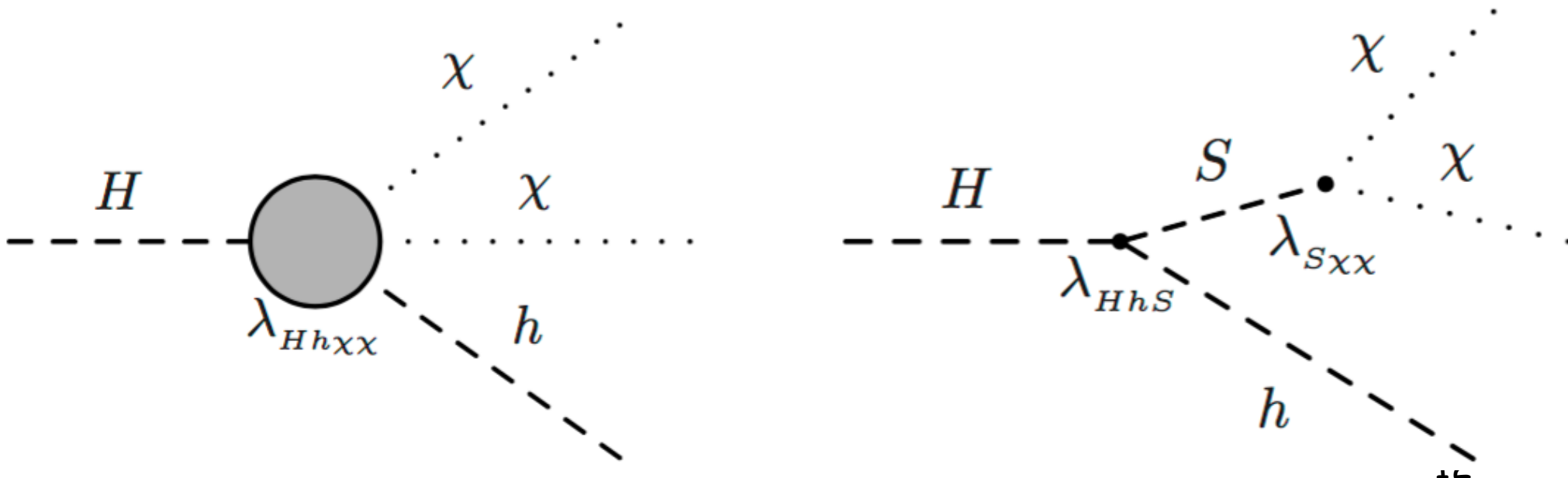
arXiv:1603.01208

arXiv:1606.01674

Formulation of the Hypothesis

The Hypothesis

1. The starting point of the hypothesis is the existence of **a boson, H** , that contains Higgs-like interactions, with a mass **in the range 250-295 GeV**
2. In order to avoid large quartic couplings and to incorporate **a mediator with Dark Matter a real scalar, S** , is introduced:

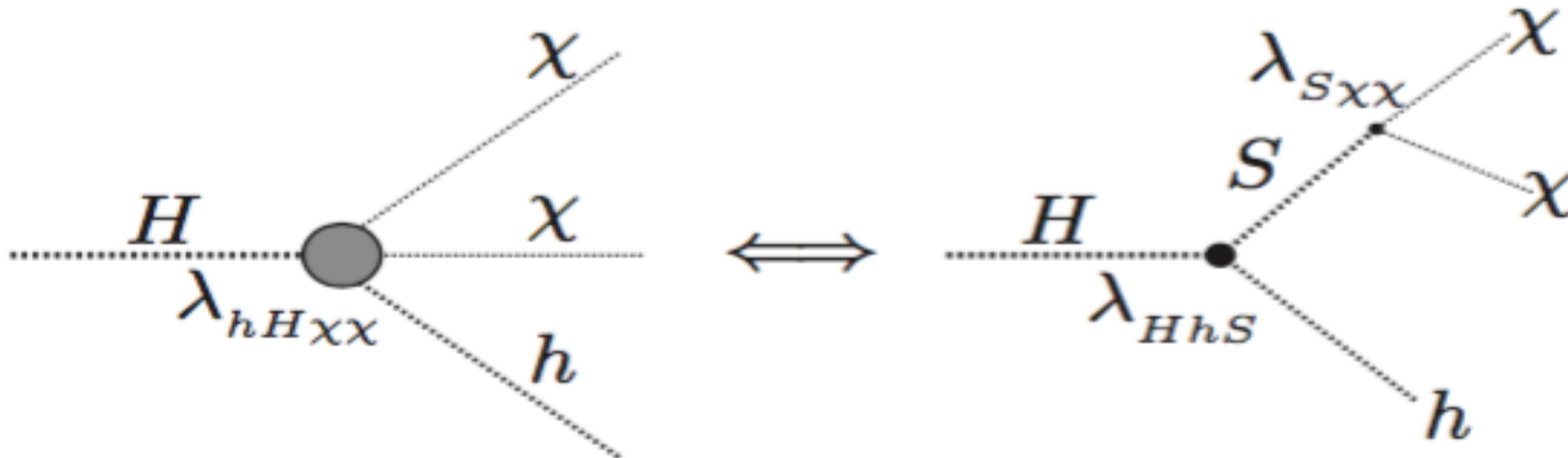


The intermediate scalar, S

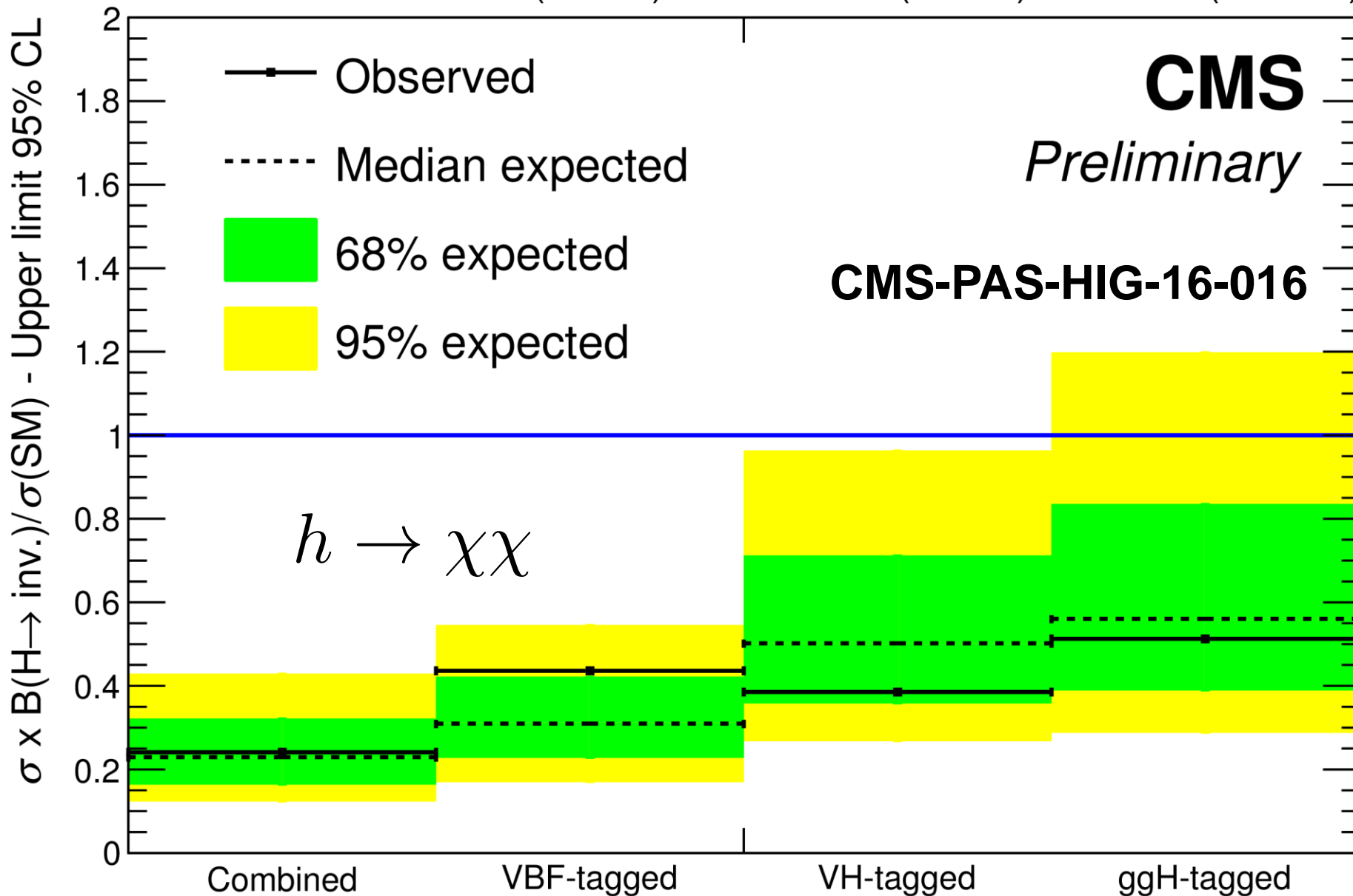
- Dark Matter is introduced in the form of a scalar and the decay $H \rightarrow h \chi \chi$ via effective quartic couplings

$$\mathcal{L}_Q = -\frac{1}{2}\lambda_{Hh\chi\chi}Hh\chi\chi - \frac{1}{4}\lambda_{HHhh}HHhh - \frac{1}{4}\lambda_{hh\chi\chi}hh\chi\chi - \frac{1}{4}\lambda_{HH\chi\chi}HH\chi\chi$$

- Due to gauge invariance we encounter an awkward situation where a three body decay may be larger or comparable to a two body decay. This can be naturally explained by introducing an intermediate real scalar S



4.9 fb⁻¹ (7 TeV) + ≤ 19.7 fb⁻¹ (8 TeV) + 2.3 fb⁻¹ (13 TeV)



Room for invisible decays of the Higgs is narrowing down.

Seems reasonable to introduce a mediator S.

The Lagrangian

$$\mathcal{L}_K = \frac{1}{2} \partial_\mu S \partial^\mu S - \frac{1}{2} m_S^2 S S,$$

$$\begin{aligned} \mathcal{L}_{SVV'} = & \frac{1}{4} \kappa_{Sgg} \frac{\alpha_s}{12\pi v} S G^{a\mu\nu} G_{\mu\nu}^a + \frac{1}{4} \kappa_{S\gamma\gamma} \frac{\alpha}{\pi v} S F^{\mu\nu} F_{\mu\nu} + \frac{1}{4} \kappa_{SZZ} \frac{\alpha}{\pi v} S Z^{\mu\nu} Z_{\mu\nu} \\ & + \frac{1}{4} \kappa_{SZ\gamma} \frac{\alpha}{\pi v} S Z^{\mu\nu} F_{\mu\nu} + \frac{1}{4} \kappa_{SWW} \frac{2\alpha}{\pi s_w^2 v} S W^{+\mu\nu} W_{\mu\nu}^-, \end{aligned}$$

$$\mathcal{L}_{Sf\bar{f}} = - \sum_f \kappa_{Sf} \frac{m_f}{v} S \bar{f} f,$$

$$\mathcal{L}_{HhS} = - \frac{1}{2} v \left[\lambda_{h h S} h h S + \lambda_{h S S} h S S + \lambda_{H H S} H H S + \lambda_{H S S} H S S + \lambda_{H h S} H h S \right],$$

$$\mathcal{L}_{S\chi} = - \frac{1}{2} v \lambda_{S\chi\chi} S \chi\chi - \frac{1}{2} \lambda_{SS\chi\chi} S S \chi\chi.$$

$$\mathcal{L}_S = \mathcal{L}_K + \mathcal{L}_{SVV'} + \mathcal{L}_{Sf\bar{f}} + \mathcal{L}_{hHS} + \mathcal{L}_{S\chi}$$

Note that some of the effective quartic couplings shown earlier appear here as trilinear. What was formerly a three body decay is now a two body decay (see below).

The Decays of H

- In the general case, H can have couplings as those displayed by a Higgs boson in addition to decays involving the intermediate scalar and Dark Matter

$$H \rightarrow WW, ZZ, q\bar{q}, gg, Z\gamma, \gamma\gamma, \chi\chi$$
$$+ H \rightarrow SS, Sh, hh$$

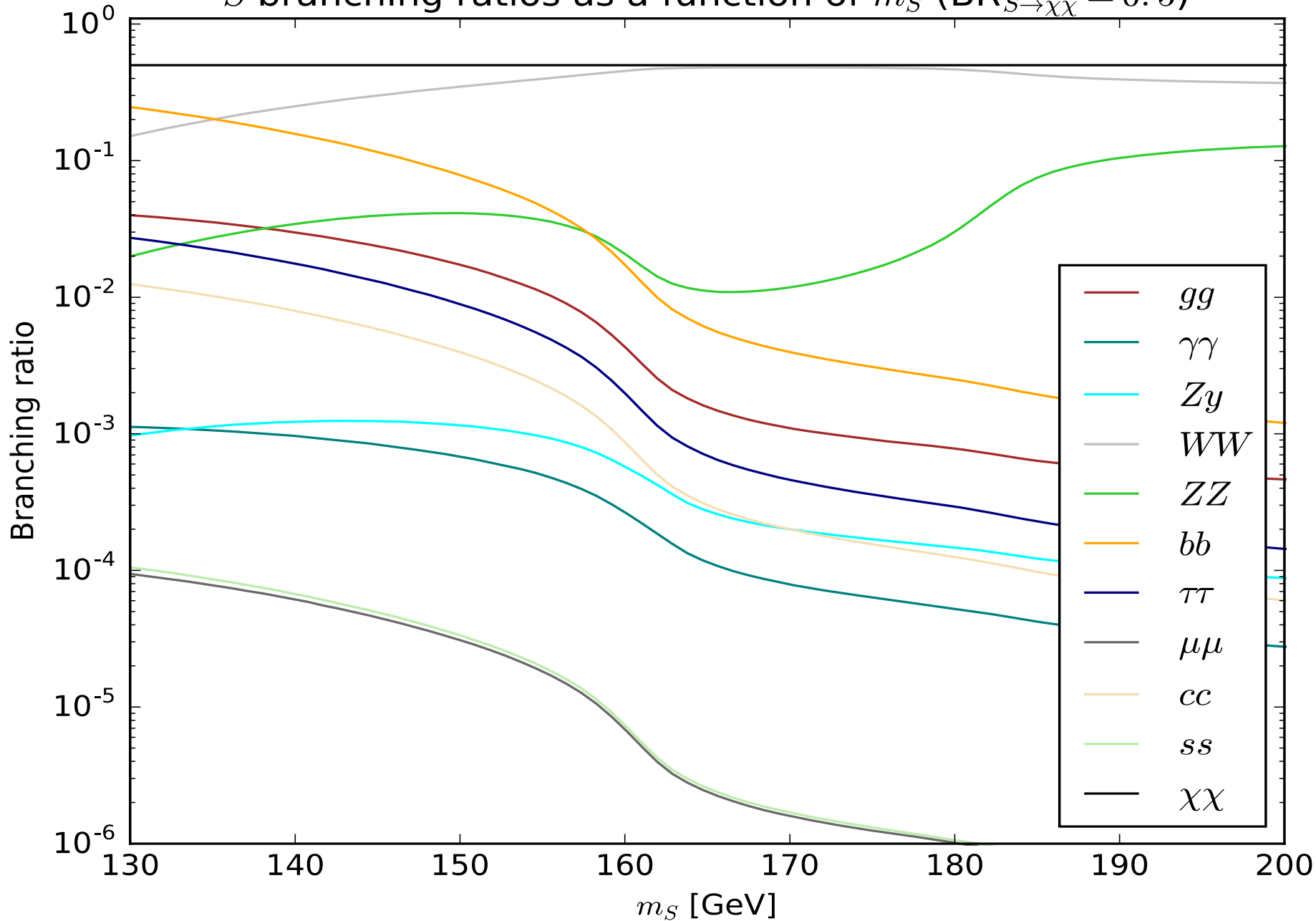
Dominant decays

Diboson decay

$$H \rightarrow h(+X), S(+X)$$

In a simplified model treat S as Higgs-like

S branching ratios as a function of m_S ($\text{BR}_{S \rightarrow \chi\chi} = 0.5$)



A simplified model (cont.)

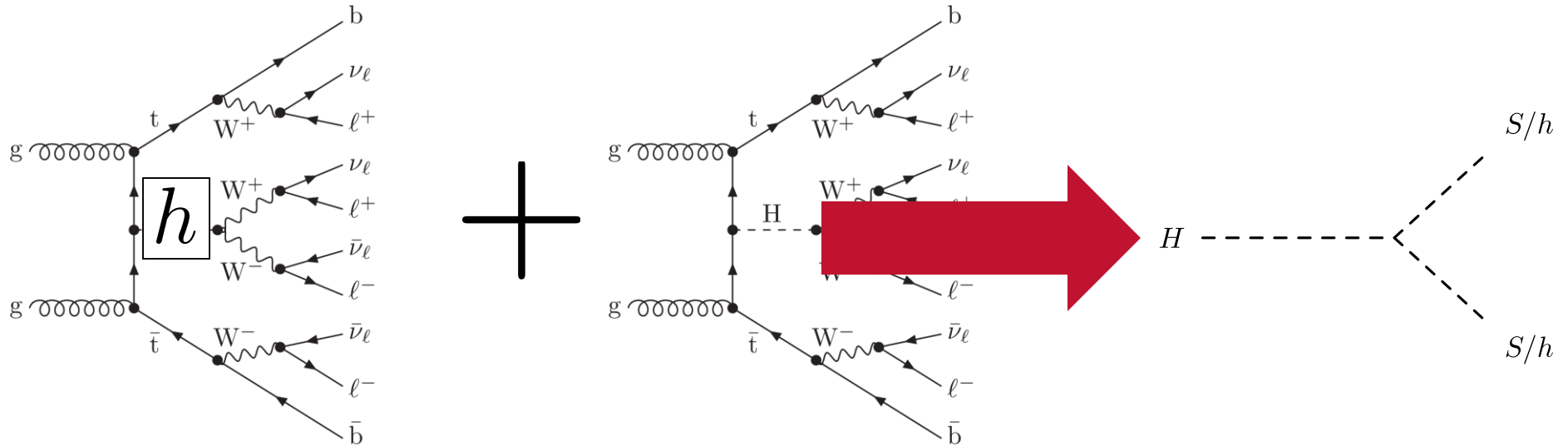
□ The following parameters are considered:

$$m_h < m_S < m_H - m_h$$

$$\frac{\Gamma(H \rightarrow SS)}{\Gamma(H \rightarrow Sh)}$$

$$Br(S \rightarrow \chi\chi)$$

Top associated Higgs production



Reduced cross-section of $ttH+tH$ is compensated by di-boson, (SS , Sh) decay and potentially large $Br(S \rightarrow WW)$. Production of same sign leptons, three leptons is enhanced

$$S, h \rightarrow WW, \tau\tau, ZZ$$

“The philosophers have only interpreted the world, in various ways; the point is to change it”



The prediction of signatures at the LHC

arXiv:1603.01208

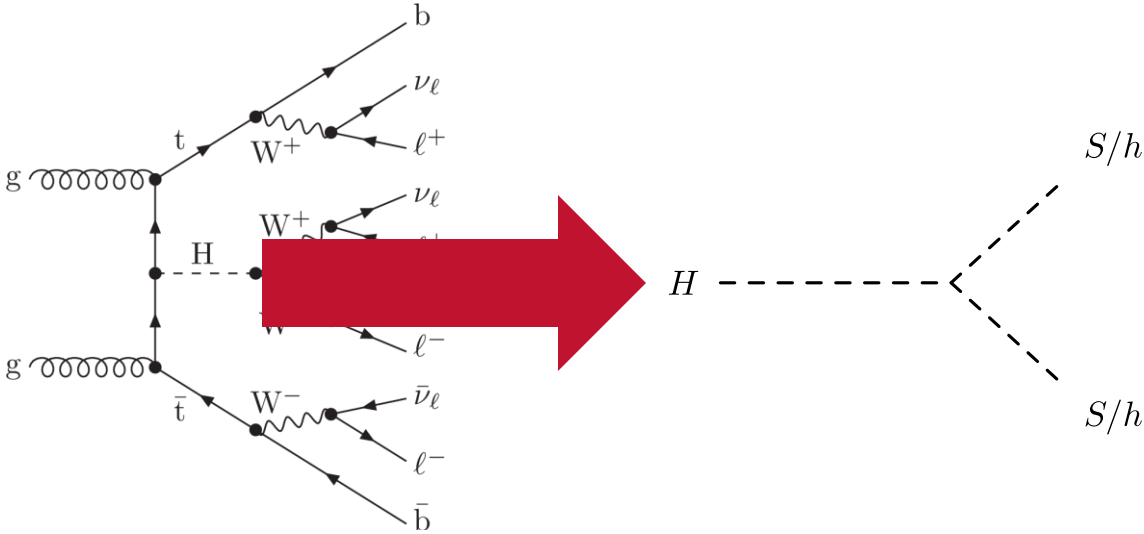
arXiv:1606.01674

The hypothesis leads to rich phenomenology

S. No.	Scalars	Decay modes
D.1	h	$b\bar{b}, \tau^+\tau^-, \mu^+\mu^-, s\bar{s}, c\bar{c}, gg, \gamma\gamma, Z\gamma, W^+W^-, ZZ$
D.2	H	D.1, hh, SS, Sh
D.3	A	D.1, $t\bar{t}, Zh, ZH, ZS, W^\pm H^\mp$
D.4	H^\pm	$W^\pm h, W^\pm H, W^\pm S$
D.5	S	D.1, $\chi\chi$

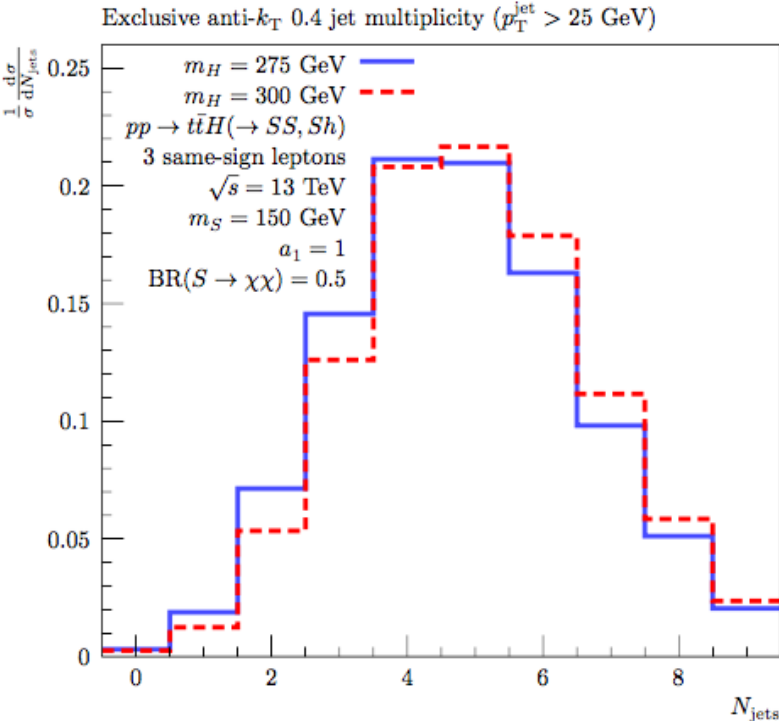
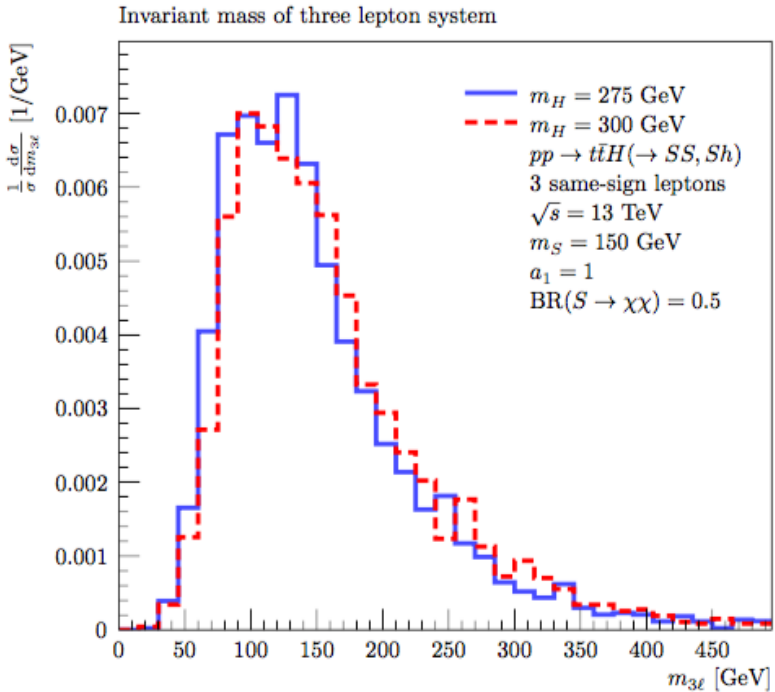
Scalar	Production mode	Search channels
H	$gg \rightarrow H, Hjj$ (ggF and VBF)	Direct SM decays as in Table 1 $\rightarrow SS/Sh \rightarrow 4W \rightarrow 4\ell + E_T^{\text{miss}}$ $\rightarrow hh \rightarrow \gamma\gamma b\bar{b}, b\bar{b}\tau\tau, 4b, \gamma\gamma WW$ etc. $\rightarrow Sh$ where $S \rightarrow \chi\chi \implies \gamma\gamma, b\bar{b}, 4\ell + E_T^{\text{miss}}$
	$pp \rightarrow Z(W^\pm)H$ ($H \rightarrow SS/Sh$)	$\rightarrow 6(5)l + E_T^{\text{miss}}$ $\rightarrow 4(3)l + 2j + E_T^{\text{miss}}$ $\rightarrow 2(1)l + 4j + E_T^{\text{miss}}$
	$pp \rightarrow t\bar{t}H, (t+\bar{t})H$ ($H \rightarrow SS/Sh$)	$\rightarrow 2W + 2Z + E_T^{\text{miss}}$ and b -jets $\rightarrow 6W \rightarrow 3$ same sign leptons + jets and E_T^{miss}
H^\pm	$pp \rightarrow tH^\pm$ ($H^\pm \rightarrow W^\pm H$)	$\rightarrow 6W \rightarrow 3$ same sign leptons + jets and E_T^{miss}
	$pp \rightarrow tbH^\pm$ ($H^\pm \rightarrow W^\pm H$)	Same as above with extra b -jet
	$pp \rightarrow H^\pm H^\mp$ ($H^\pm \rightarrow HW^\pm$)	$\rightarrow 6W \rightarrow 3$ same sign leptons + jets and E_T^{miss}
	$pp \rightarrow H^\pm W^\pm$ ($H^\pm \rightarrow HW^\pm$)	$\rightarrow 6W \rightarrow 3$ same sign leptons + jets and E_T^{miss}
A	$gg \rightarrow A$ (ggF)	$\rightarrow t\bar{t}$ $\rightarrow \gamma\gamma$
	$gg \rightarrow A \rightarrow ZH$ ($H \rightarrow SS/Sh$)	Same as $pp \rightarrow ZH$ above, but with resonance structure over final state objects
	$gg \rightarrow A \rightarrow W^\pm H^\mp$ ($H^\mp \rightarrow W^\mp H$)	$6W$ signature with resonance structure over final state objects

Production of three same sign leptons

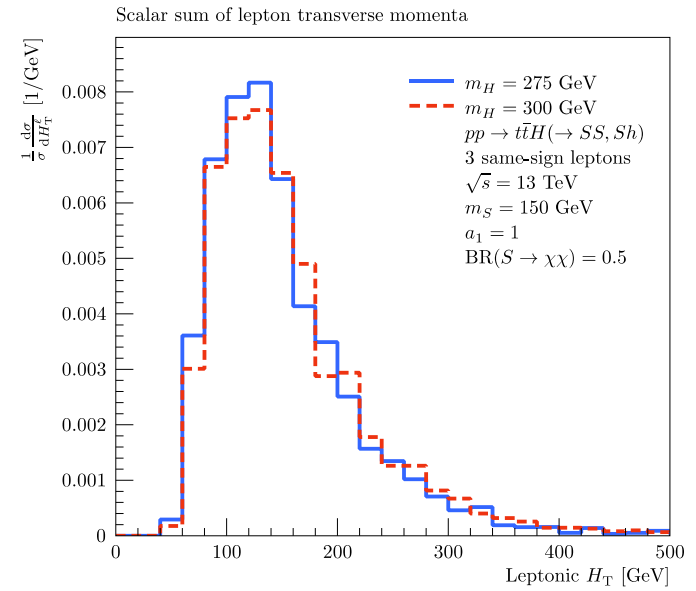
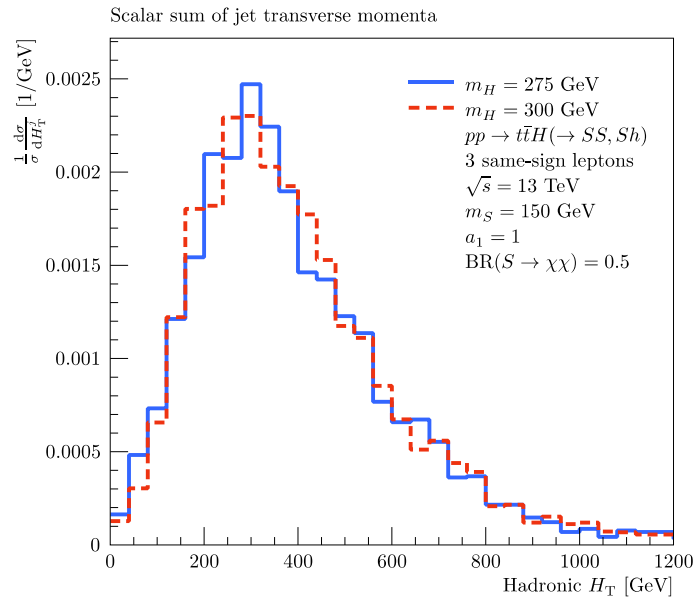
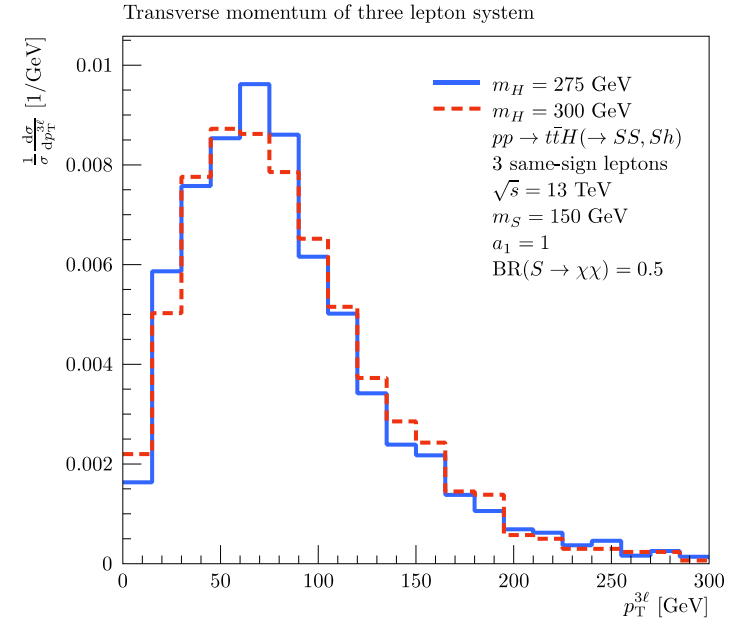
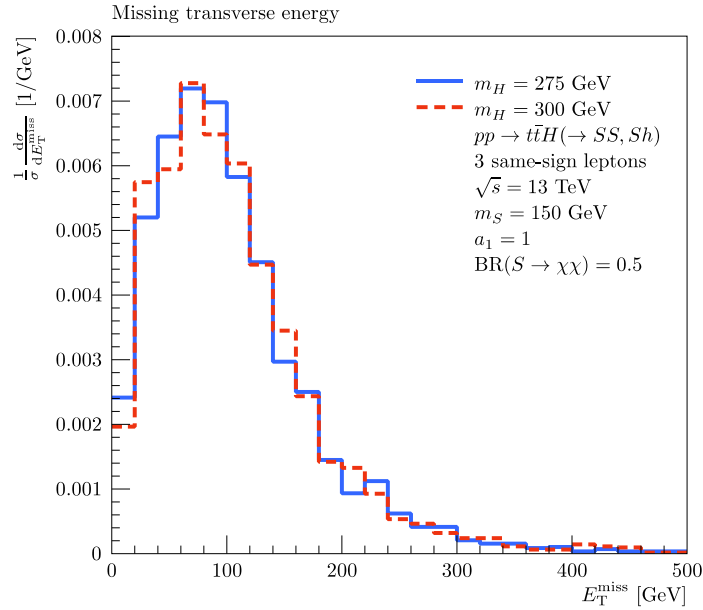


$$t(t)H \rightarrow 6W + X$$

$$\rightarrow l^\pm l^\pm l^\pm + X$$

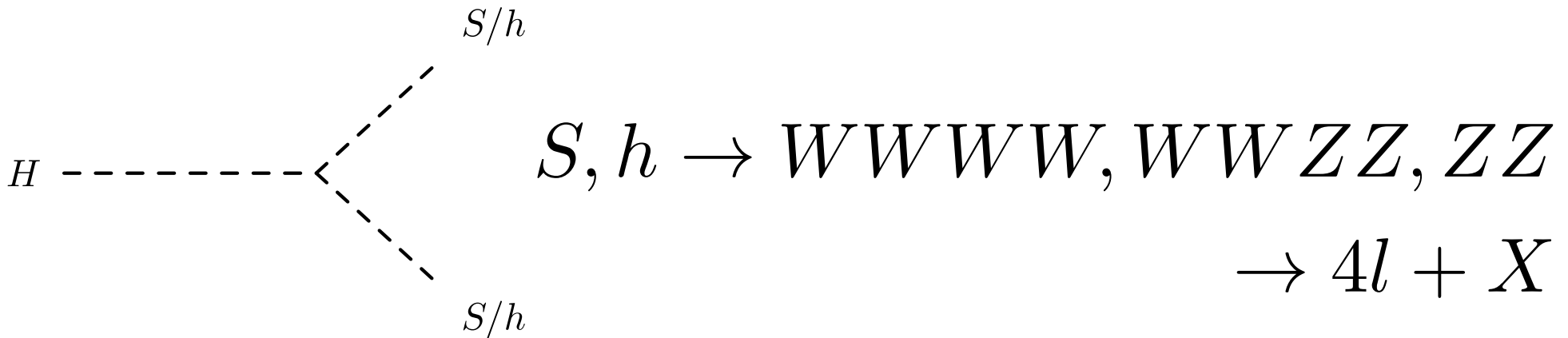


Production of three same sign leptons



Production of 4 isolated leptons

Coming predominantly from production of 4W

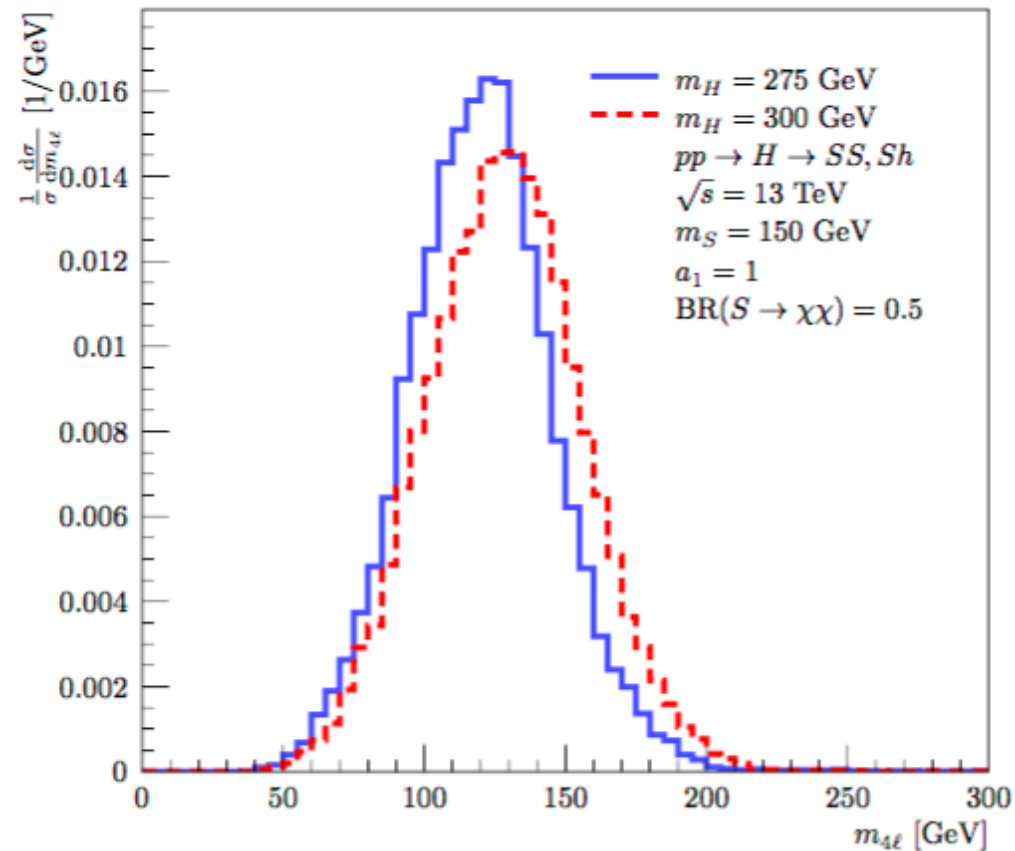


Features:

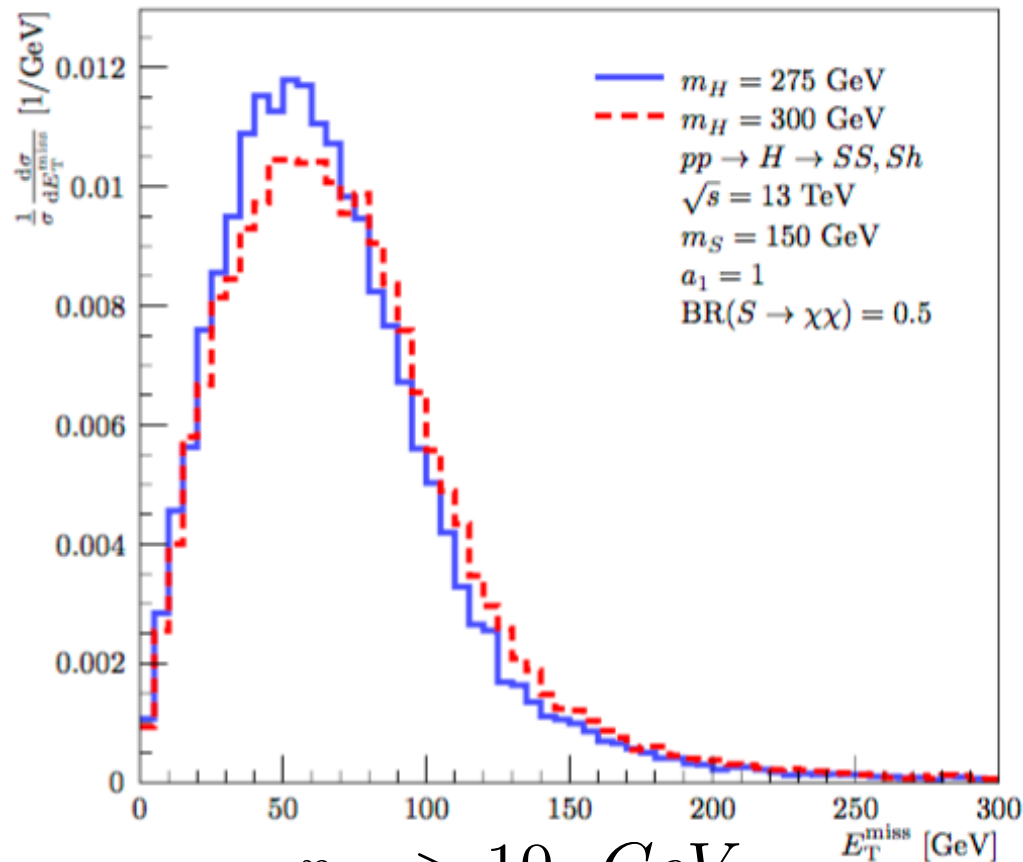
1. **Low backgrounds -> excellent S/B**
2. **Clean signature with fake leptons under control**
3. **Unique signature of the hypothesis**
4. **Sensitive to the mass of H**

The production of 4W from a resonance is a unique signature leading to the production of 4 isolated charged leptons and missing energy. The LHC experiments have not reported on this signature to date

Invariant mass of four lepton system



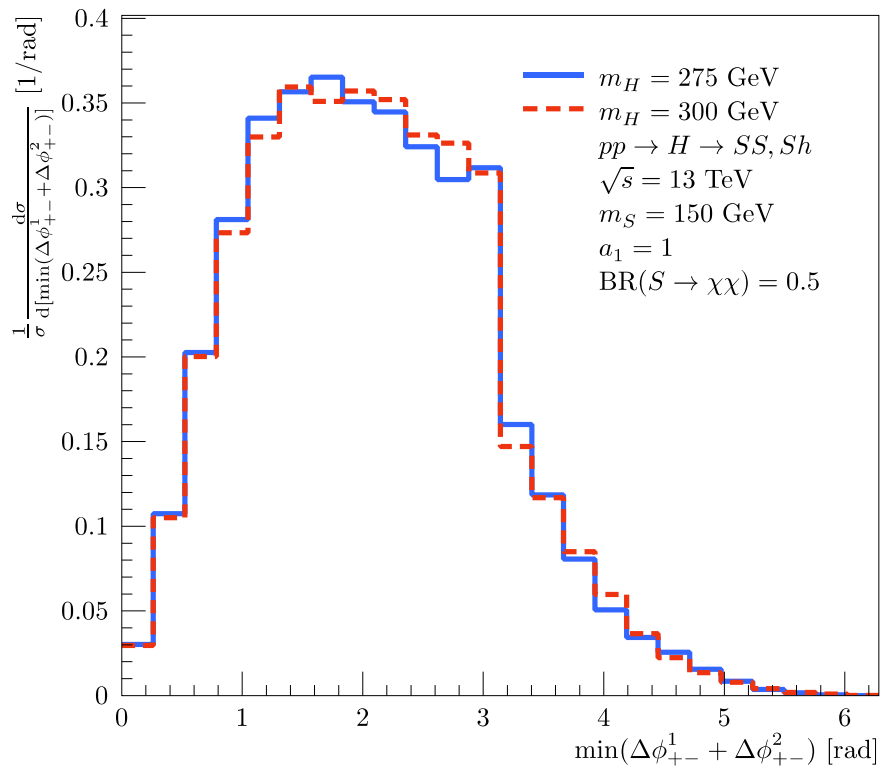
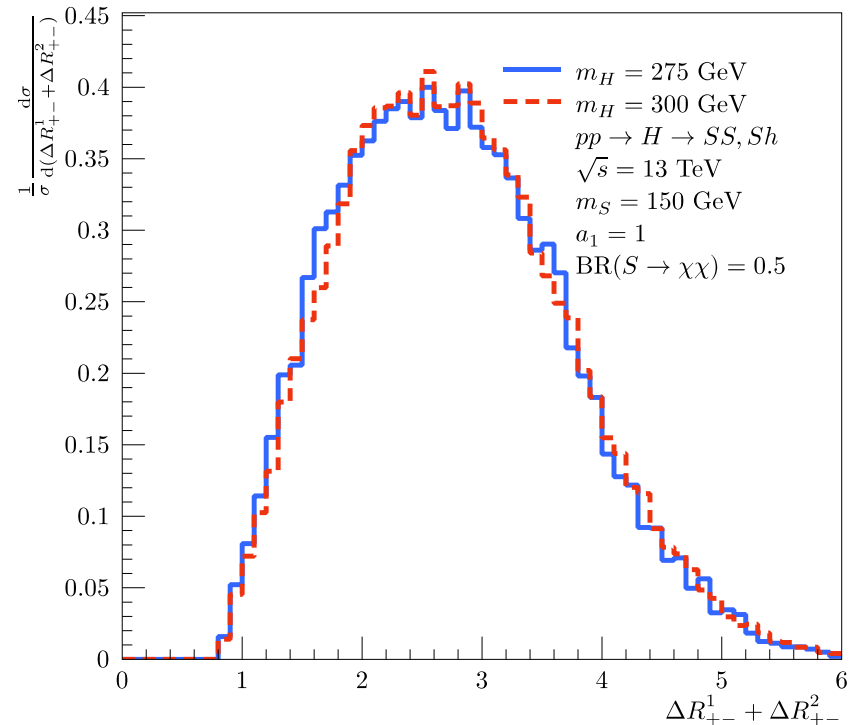
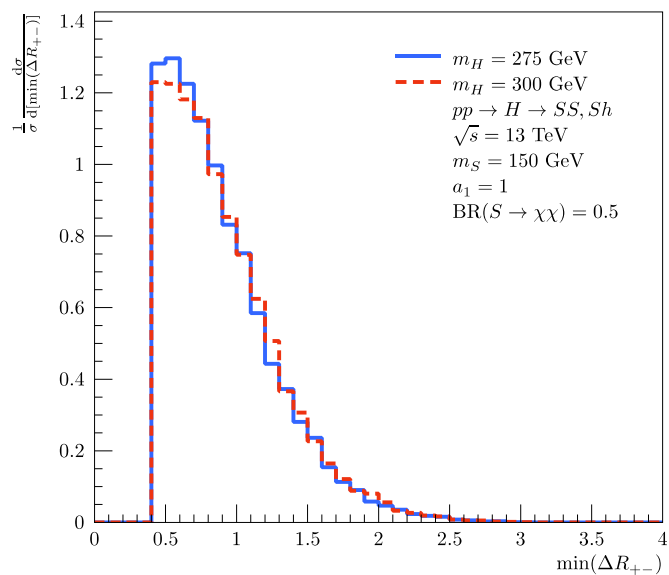
Missing transverse energy



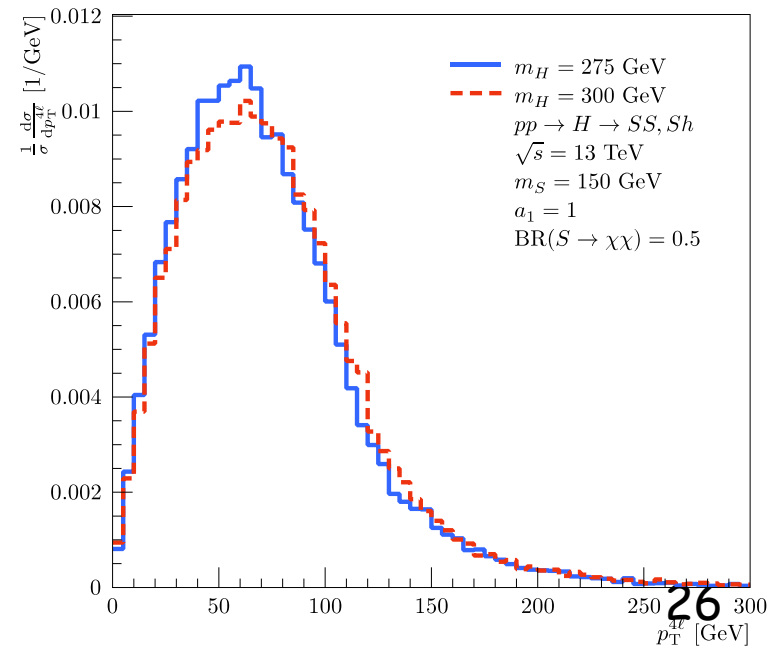
$$p_{Tl} > 10 \text{ GeV}$$

Predict ~1.5 fb of fiducial cross-section

$$|\eta_l| < 2.5 \quad 25$$

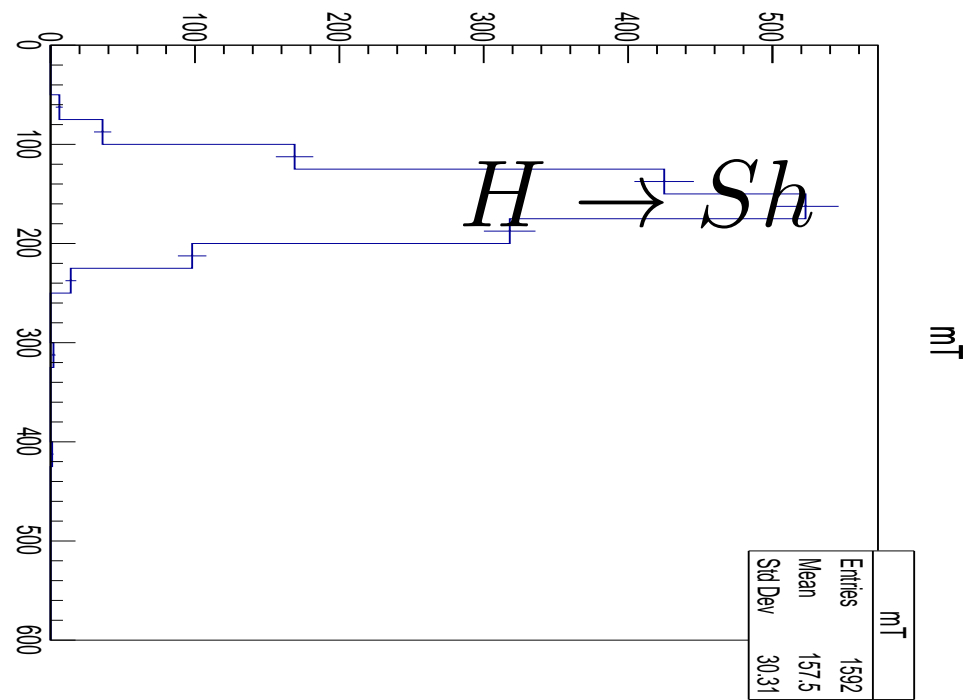
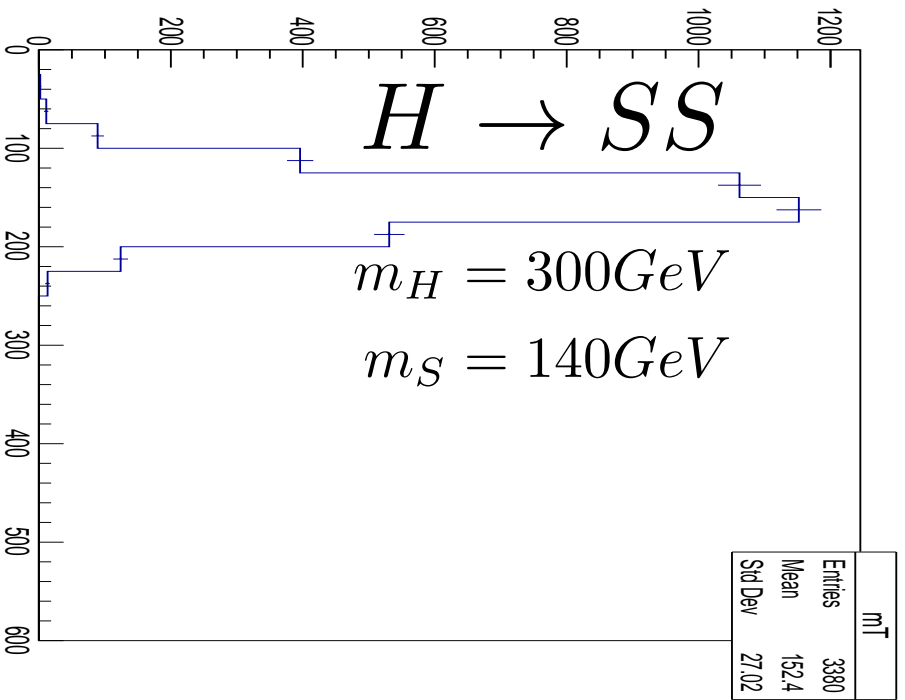
Minimised sum of combinatoric opposite sign lepton $\Delta\phi$

 Sum of ΔR of $\Delta\phi$ matched lepton pairs

 Minimised ΔR of opposite sign lepton pairs


Transverse momentum of four lepton system

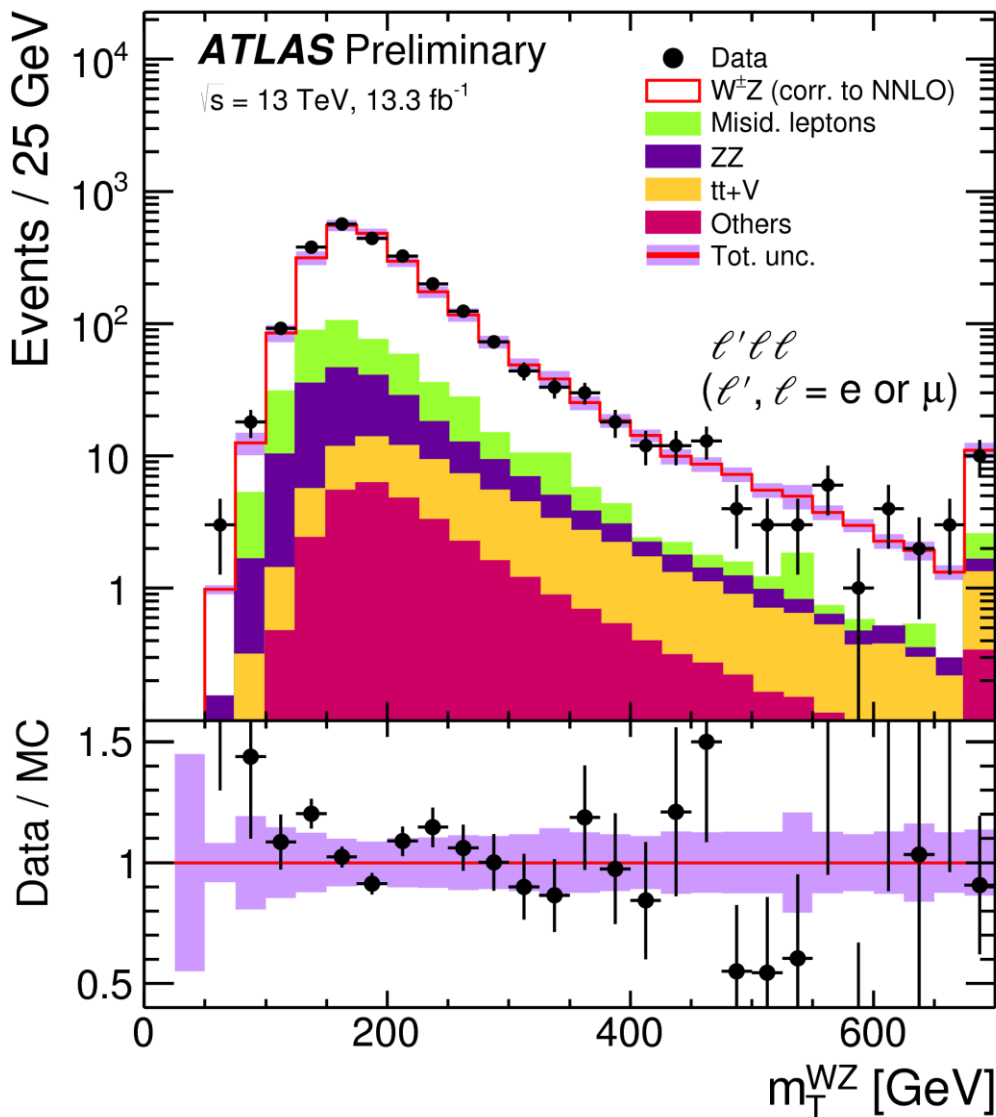
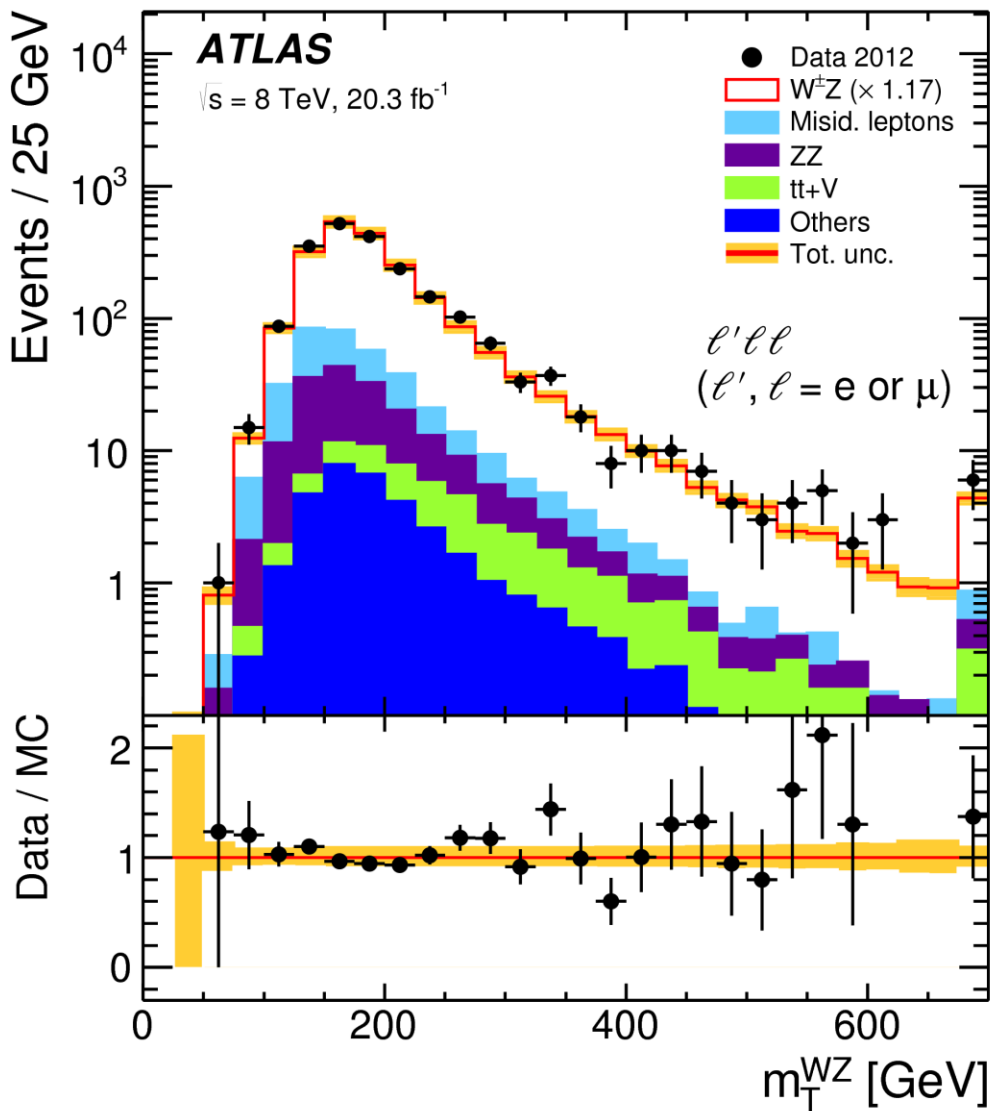


It is very important to check if the production of 3 leptons, as predicted by the model is not excluded by the existing data. Below is the transverse mass (parton level, no detector smearing applied) after the application of requirements of leptons outlined in ATLAS-CONF-2016-043

$$m_T^{WZ} = \sqrt{\left(\sum_{\ell=1}^3 p_T^\ell + E_T^{\text{miss}}\right)^2 - \left[\left(\sum_{\ell=1}^3 p_x^\ell + E_x^{\text{miss}}\right)^2 + \left(\sum_{\ell=1}^3 p_y^\ell + E_y^{\text{miss}}\right)^2\right]}$$



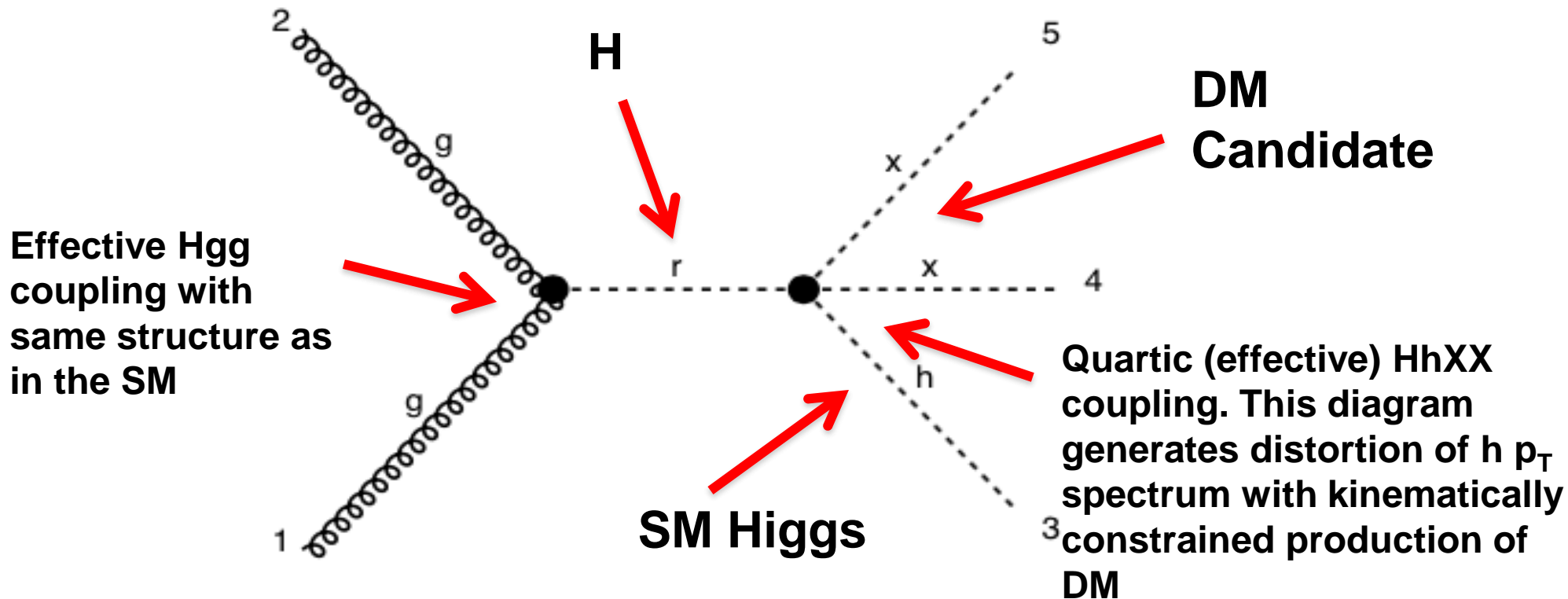
The data reported with Run I and Run II by ATLAS overshoots the MC with $M_T < 200$ GeV. The 4W prediction is not excluded with the current results.

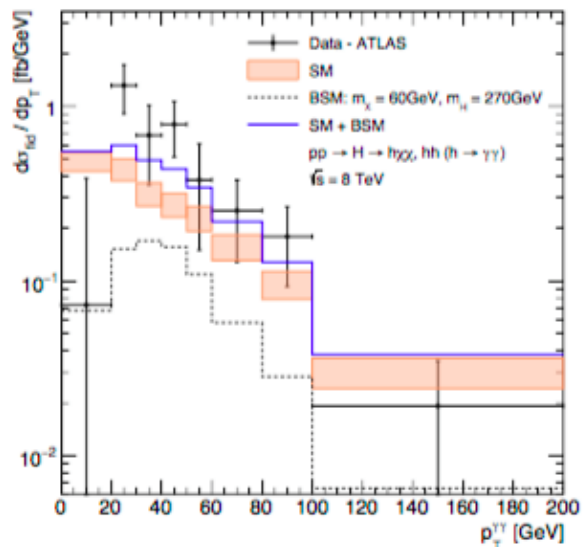


Additional Slides

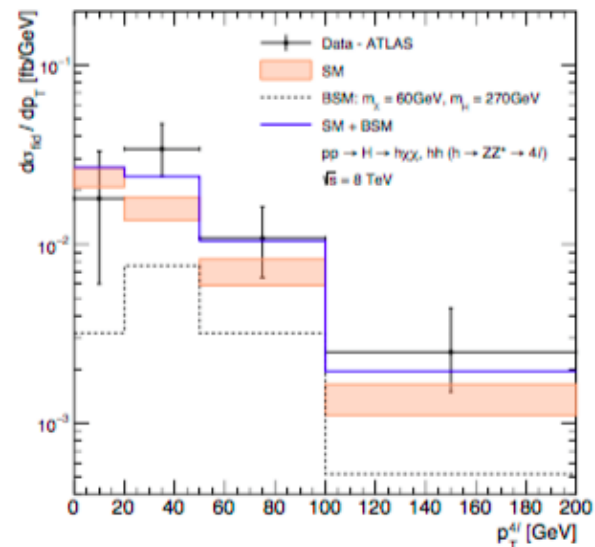
Event Generation

Generated complete Gauge invariant set of diagrams.
Suppressed hXX and $hhXX$ couplings to study diagram below

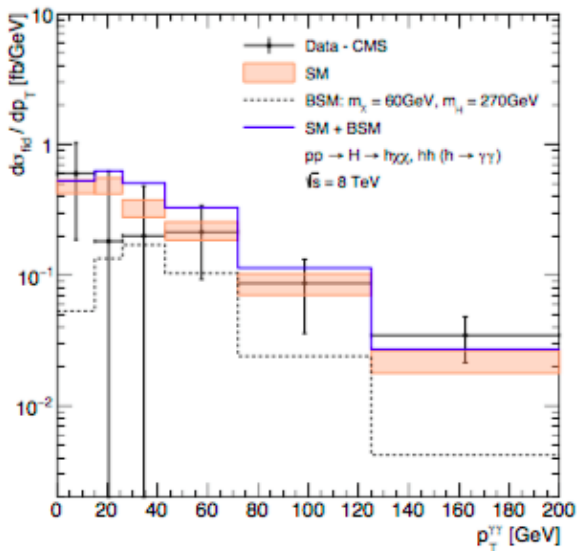




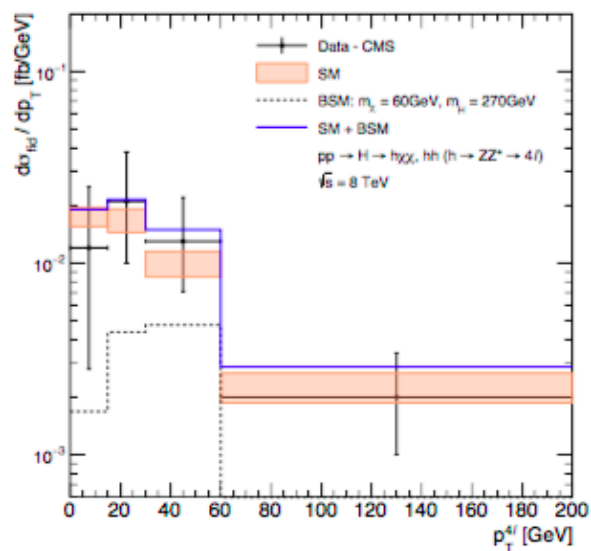
(a) ATLAS $h \rightarrow \gamma\gamma$



(b) ATLAS $h \rightarrow ZZ^* \rightarrow 4\ell$



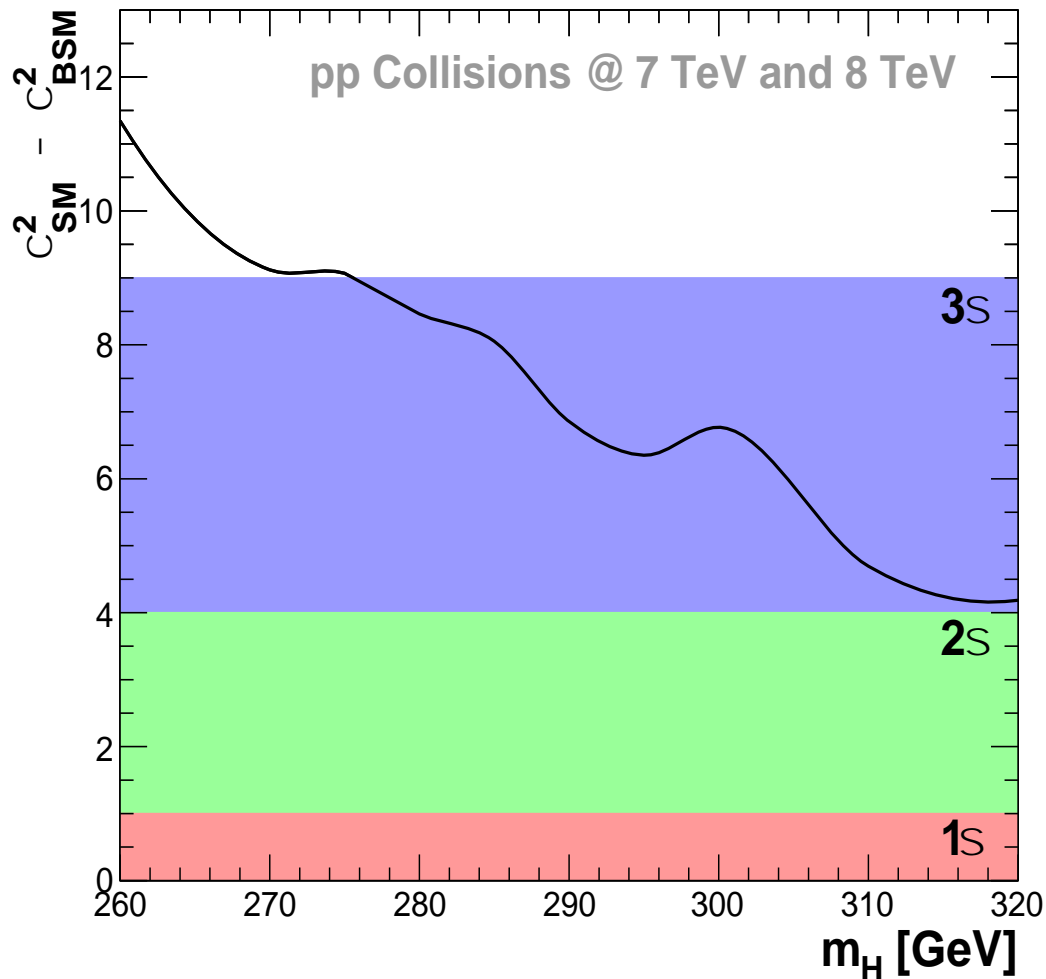
(c) CMS $h \rightarrow \gamma\gamma$



(d) CMS $h \rightarrow ZZ^* \rightarrow 4\ell$

Figure 3: Fits to the ATLAS and CMS Run I differential p_T spectra using the point $m_H = 270\text{ GeV}$ and $m_\chi = 60\text{ GeV}$.

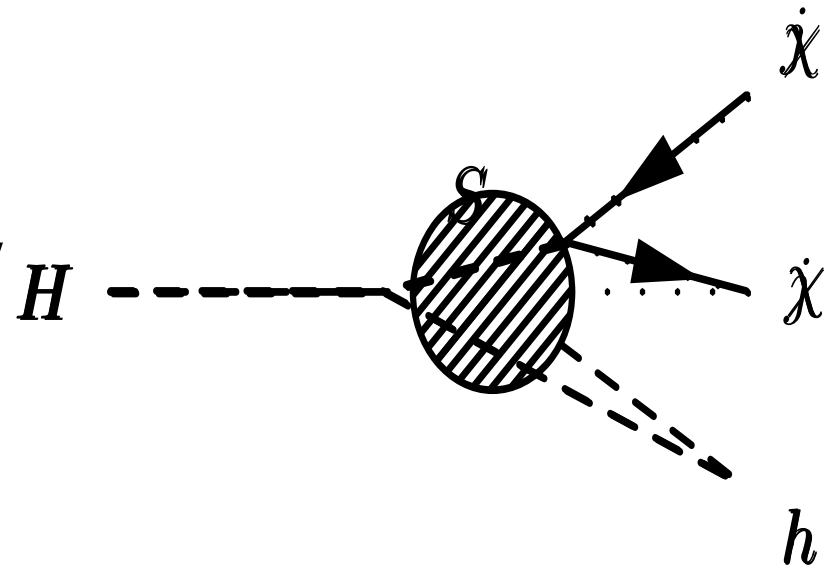
In terms of significance



- To see how significant the result is, we use a test statistic: $\chi_{SM}^2 - \chi_{BSM}^2$
- This gives an improvement on the null hypothesis (the Standard Model) in units of sigma
- For one degree of freedom, the best fit point has a 3 sigma improvement. **This does not mean evidence yet.**

What is in the blob?

- Recall the effective interaction of $H \rightarrow h\chi\chi$
- Theoretical issues make this decay unnatural
- Propose the idea that H decays into an inert scalar S and the Higgs boson h
- S can also decay into jets



Calculating χ^2

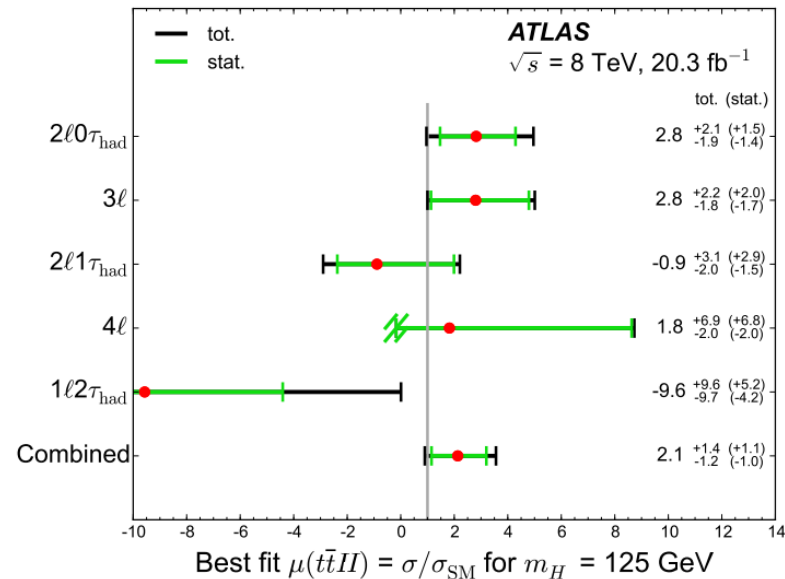
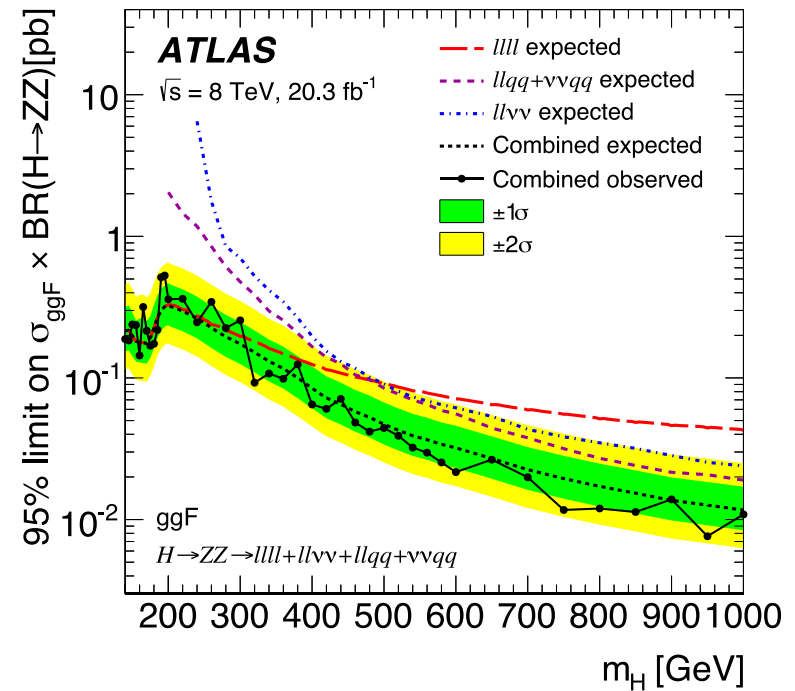
Two types of results:

Measurements

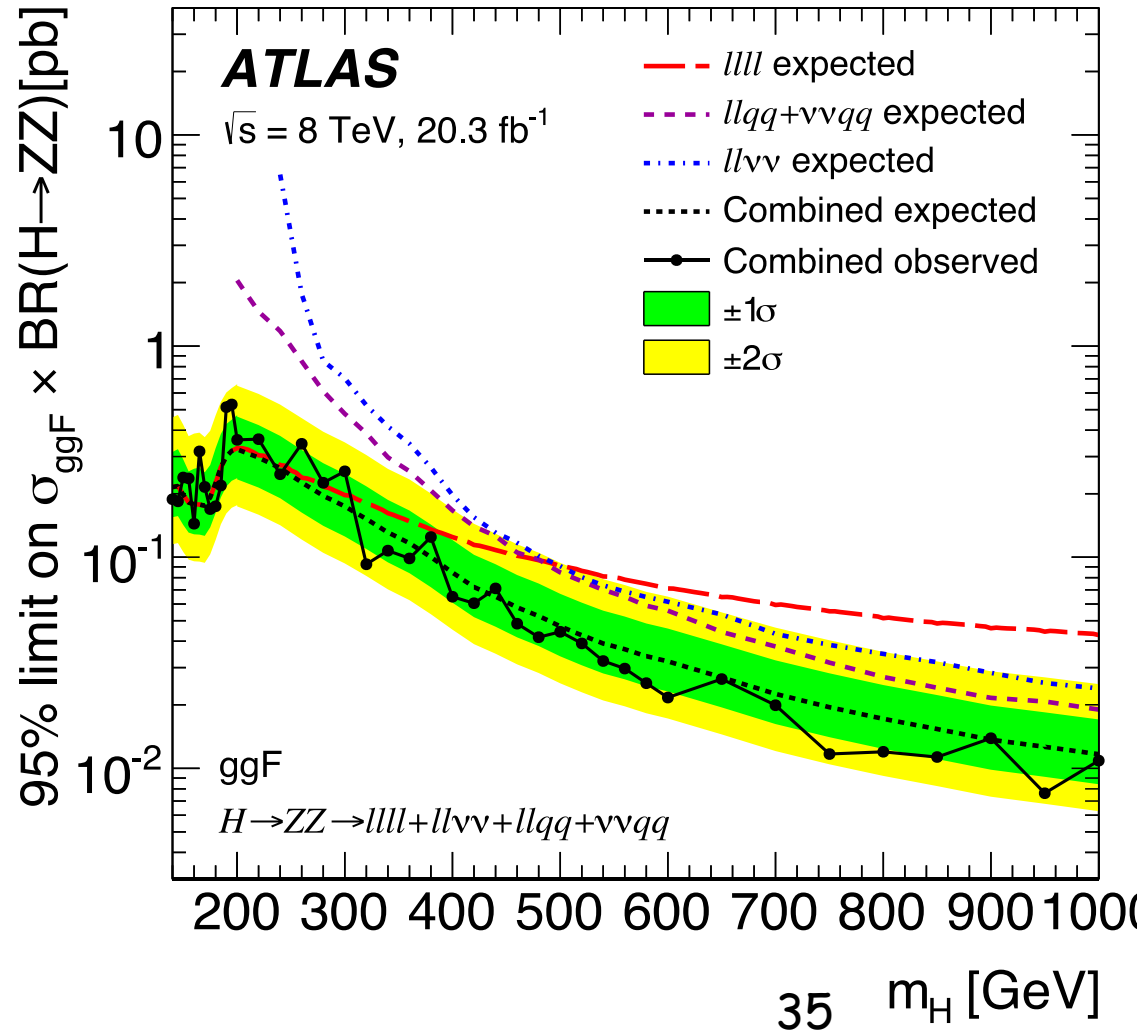
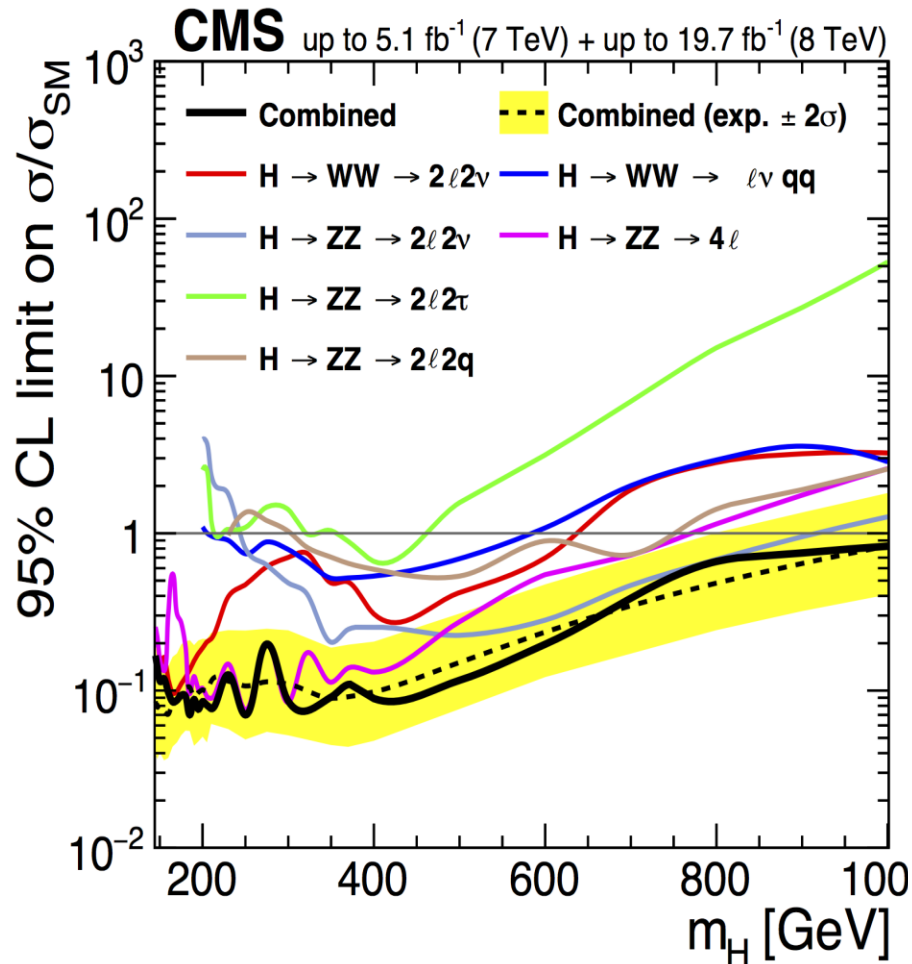
$$\chi^2 = \frac{(\mu - \mu^{\text{th}})^2}{(\Delta\mu)^2 + (\Delta\mu^{\text{th}})^2}$$

Limits

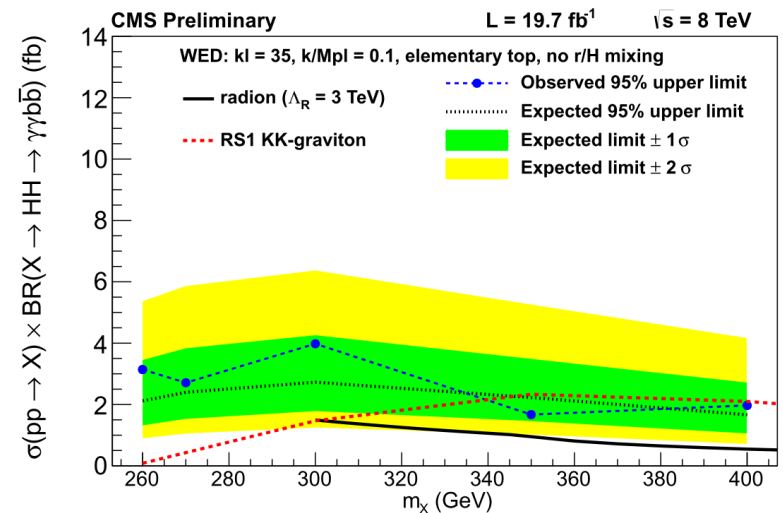
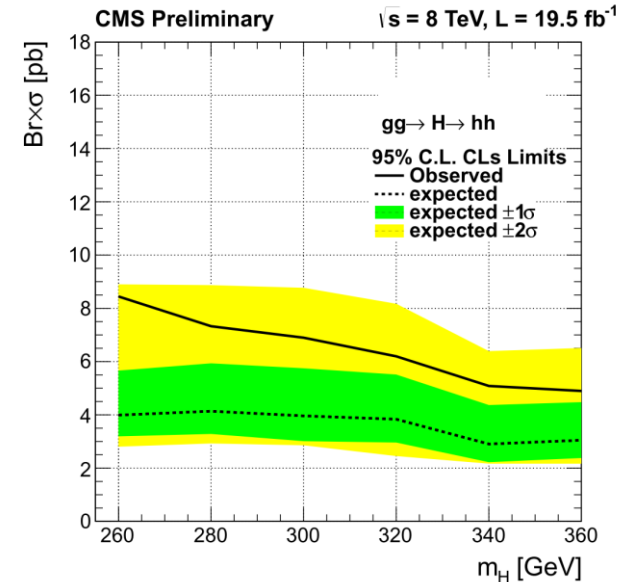
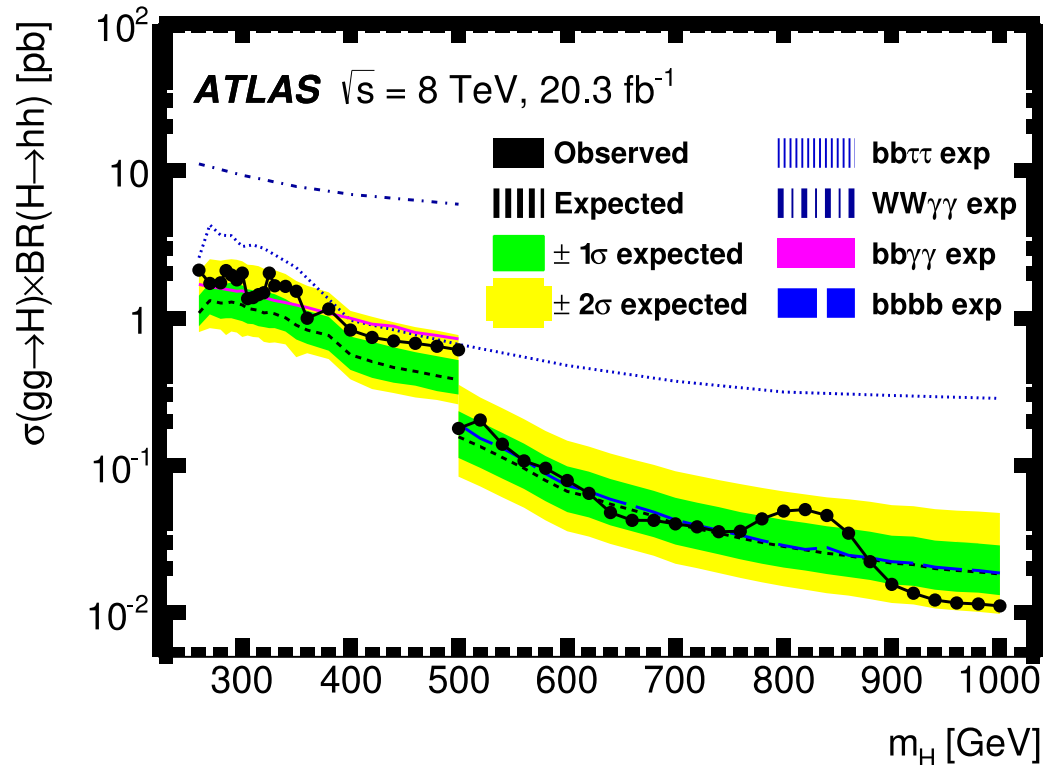
$$\chi^2 = \frac{(\mu^{\text{obs}} - \mu^{\text{exp}} - \mu^{\text{th}})^2}{(\mu^{\text{exp}} / 1.96)^2}$$



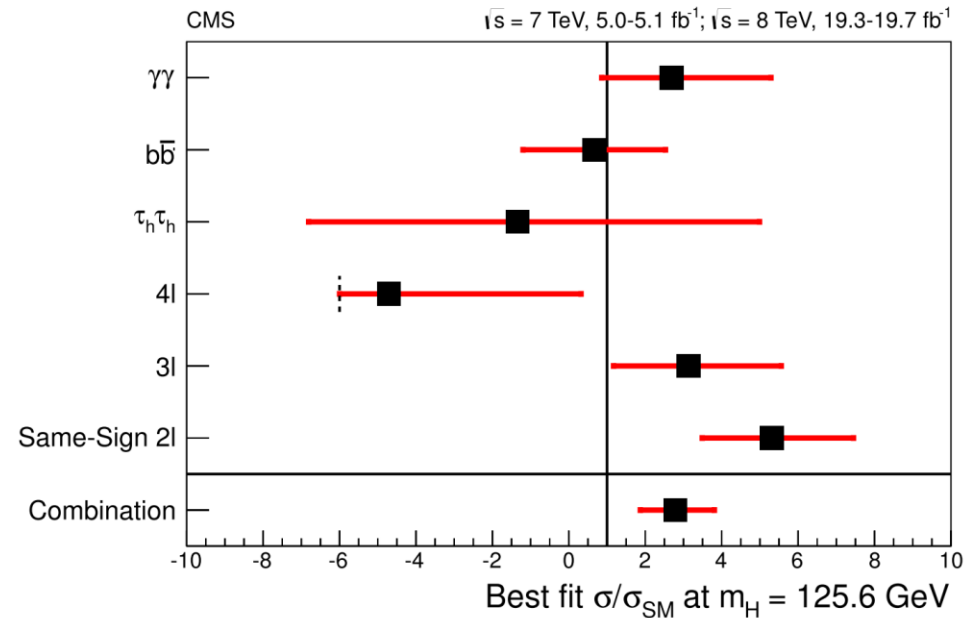
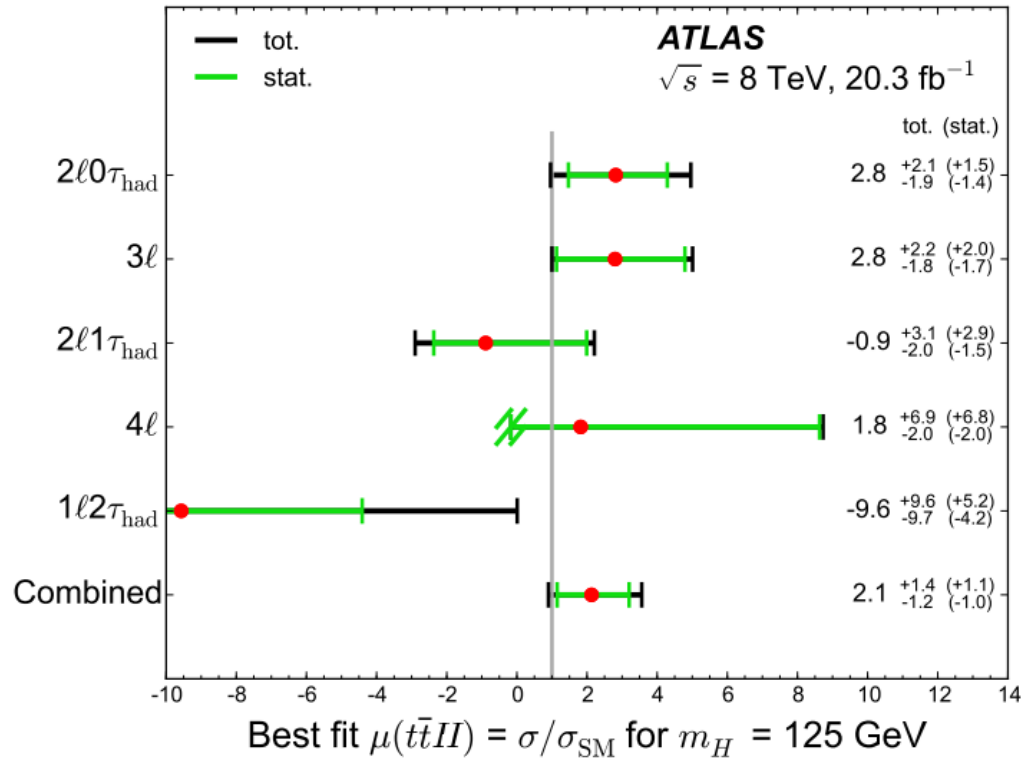
Inputs: Higgs decays to weak vector bosons



Inputs: Di-Higgs searches



Inputs: Top associated Higgs production



Step-by-step results

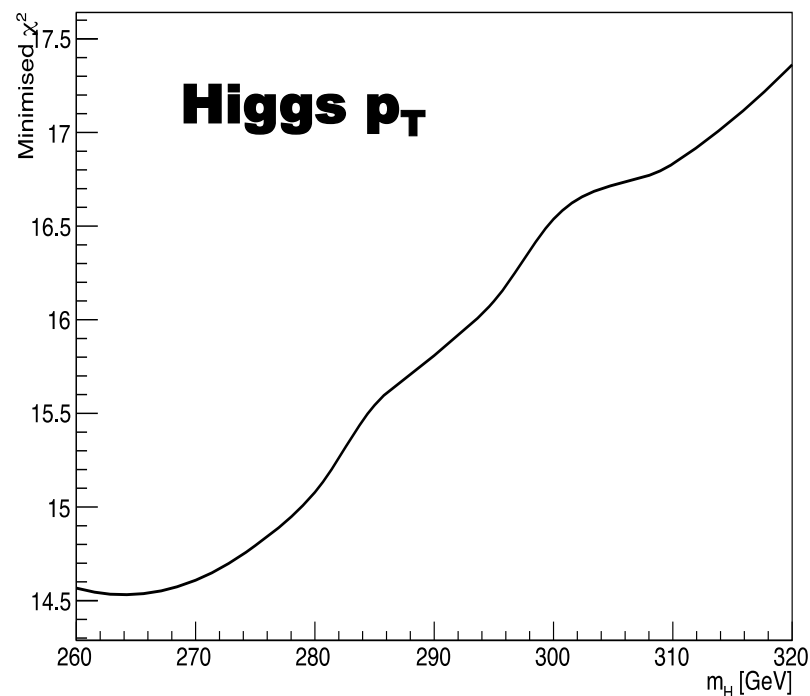
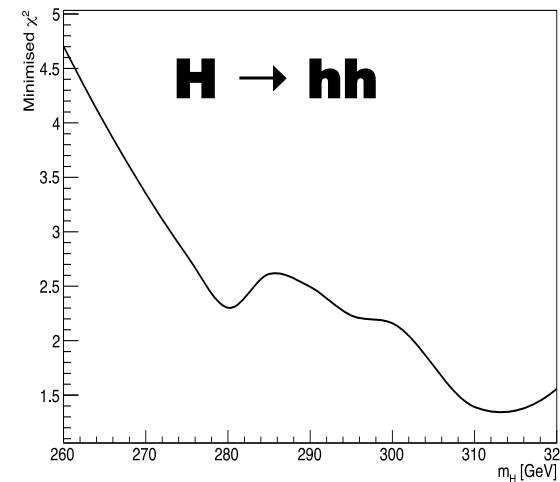
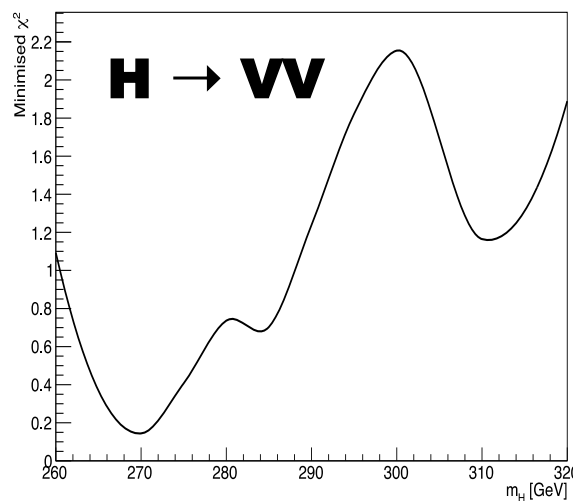
□ The global χ^2 was calculated as a function of m_H and by marginalising β_g

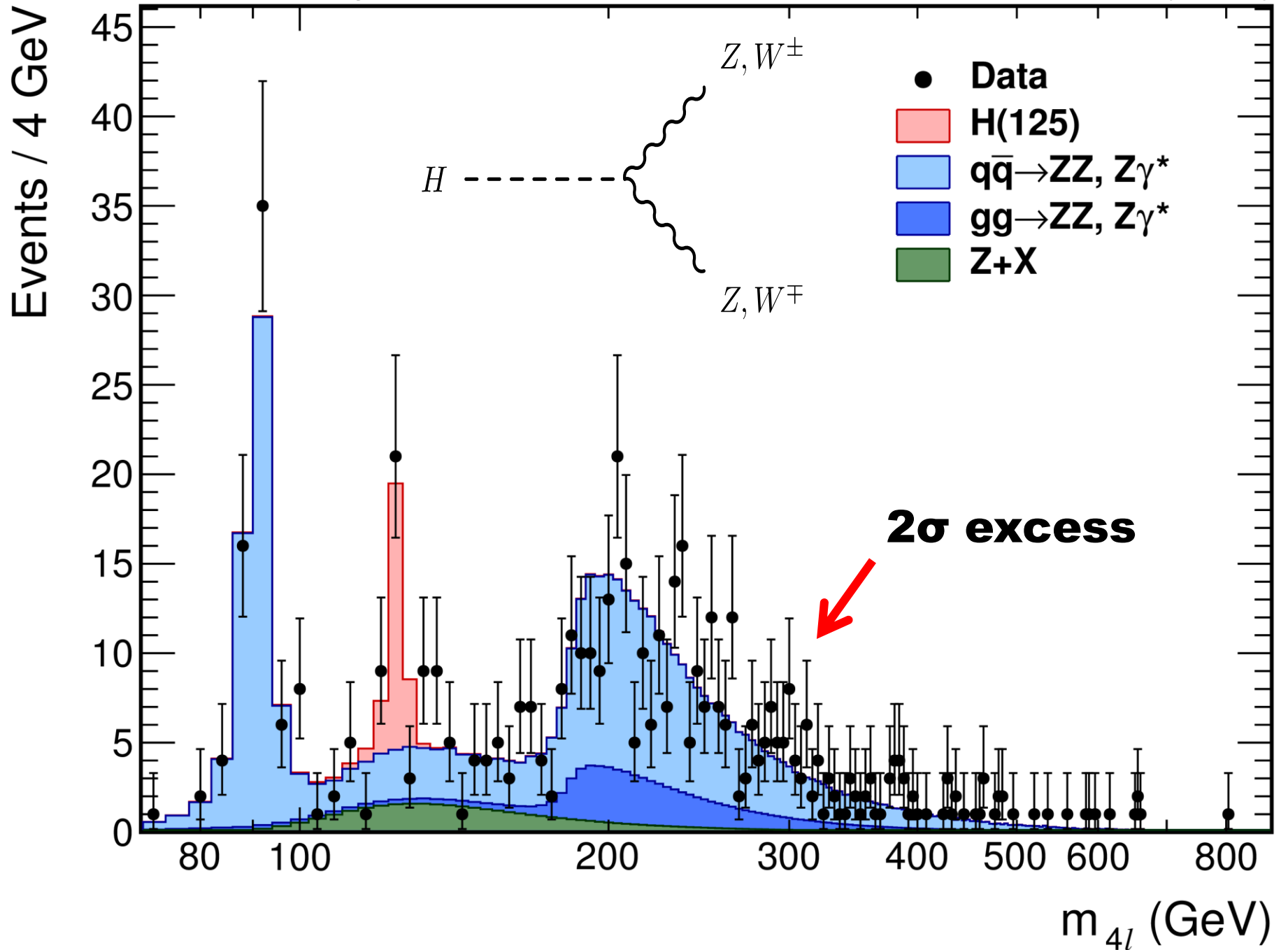
□ p_T spectra

□ $H \rightarrow VV$ decays

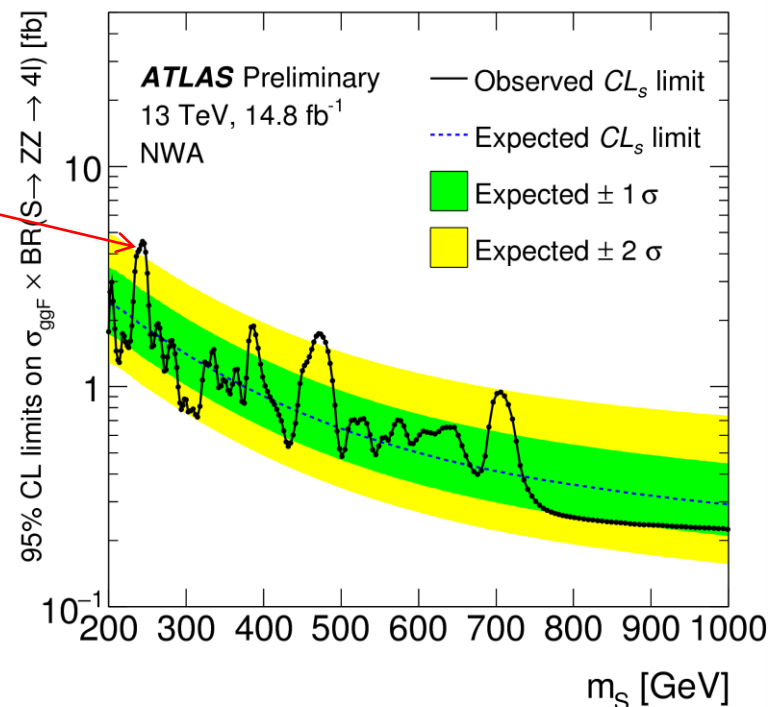
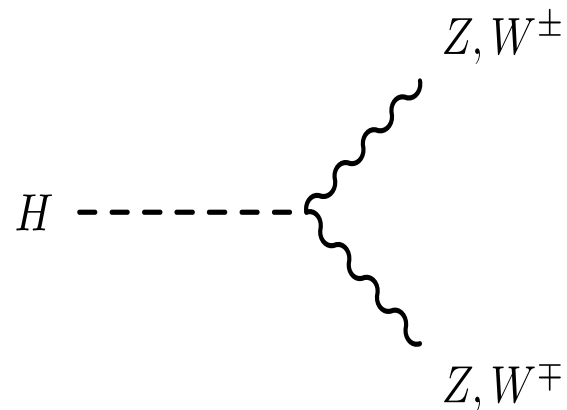
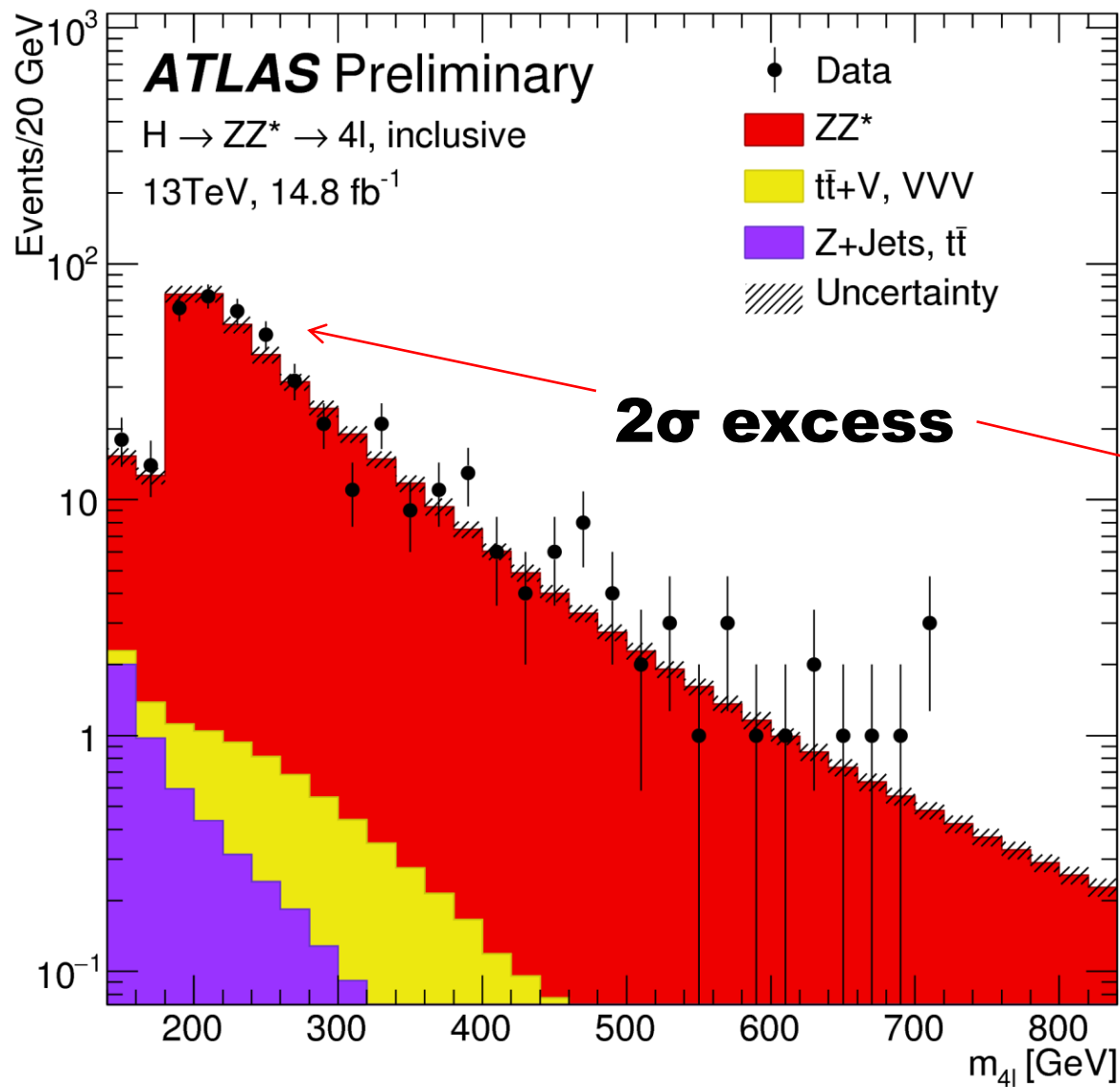
□ $H \rightarrow hh$ decays

□ These results are not significant on their own



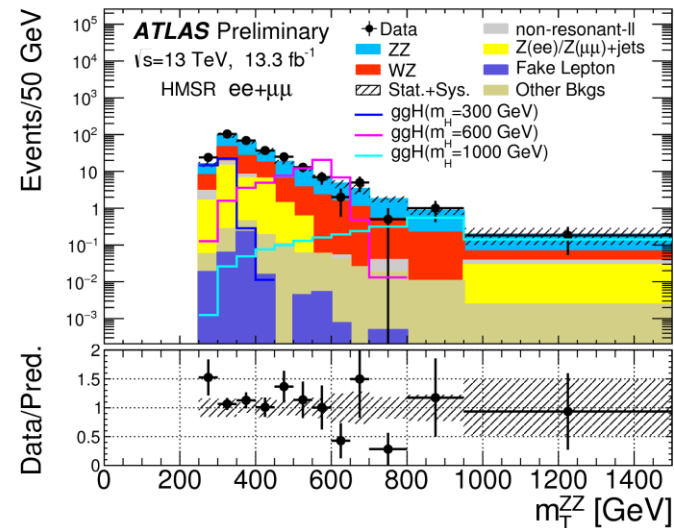
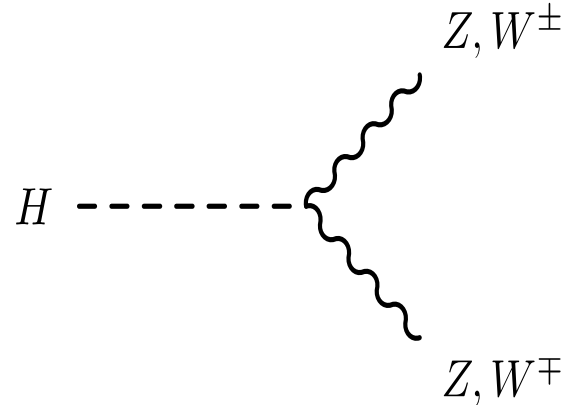
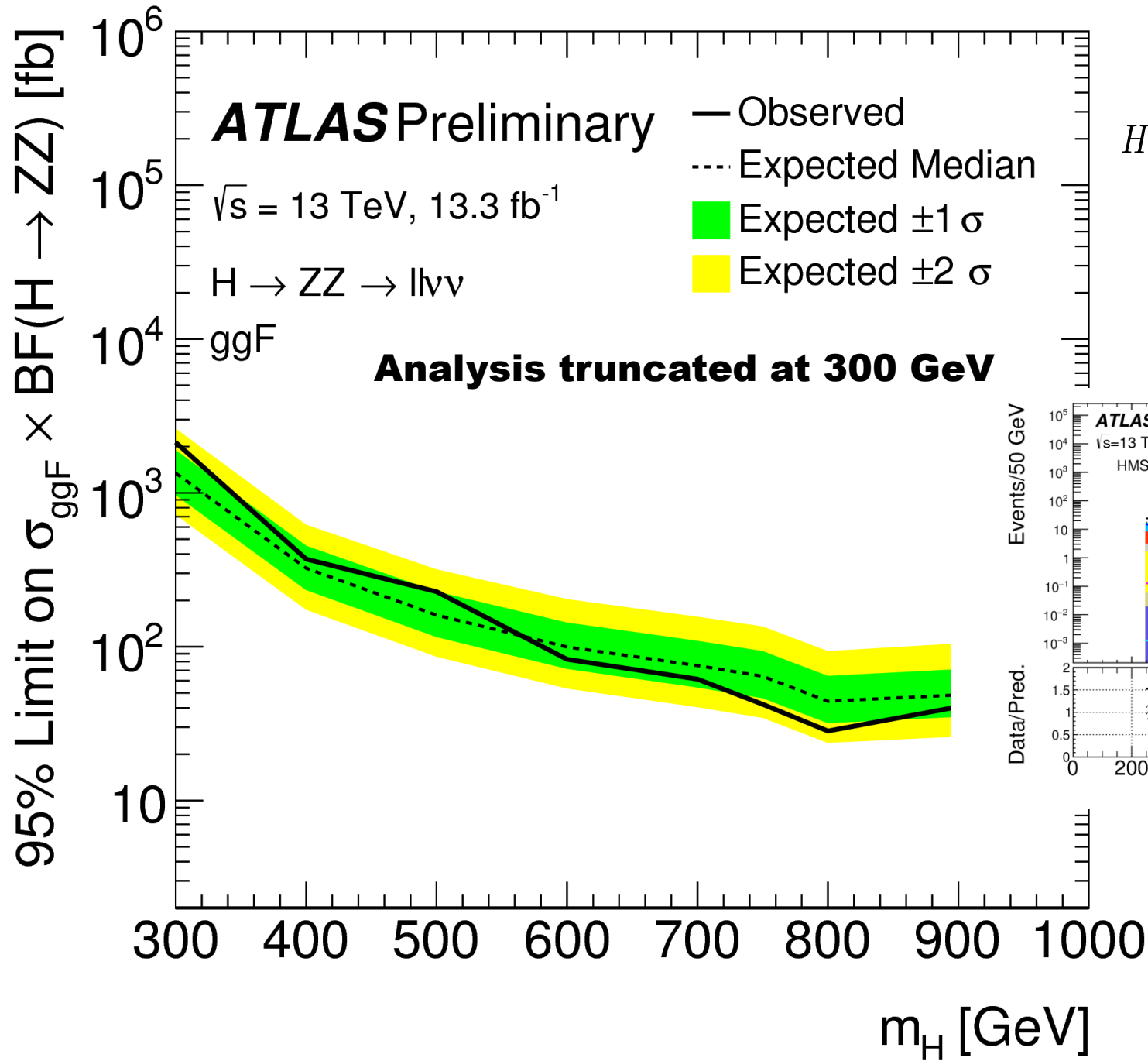


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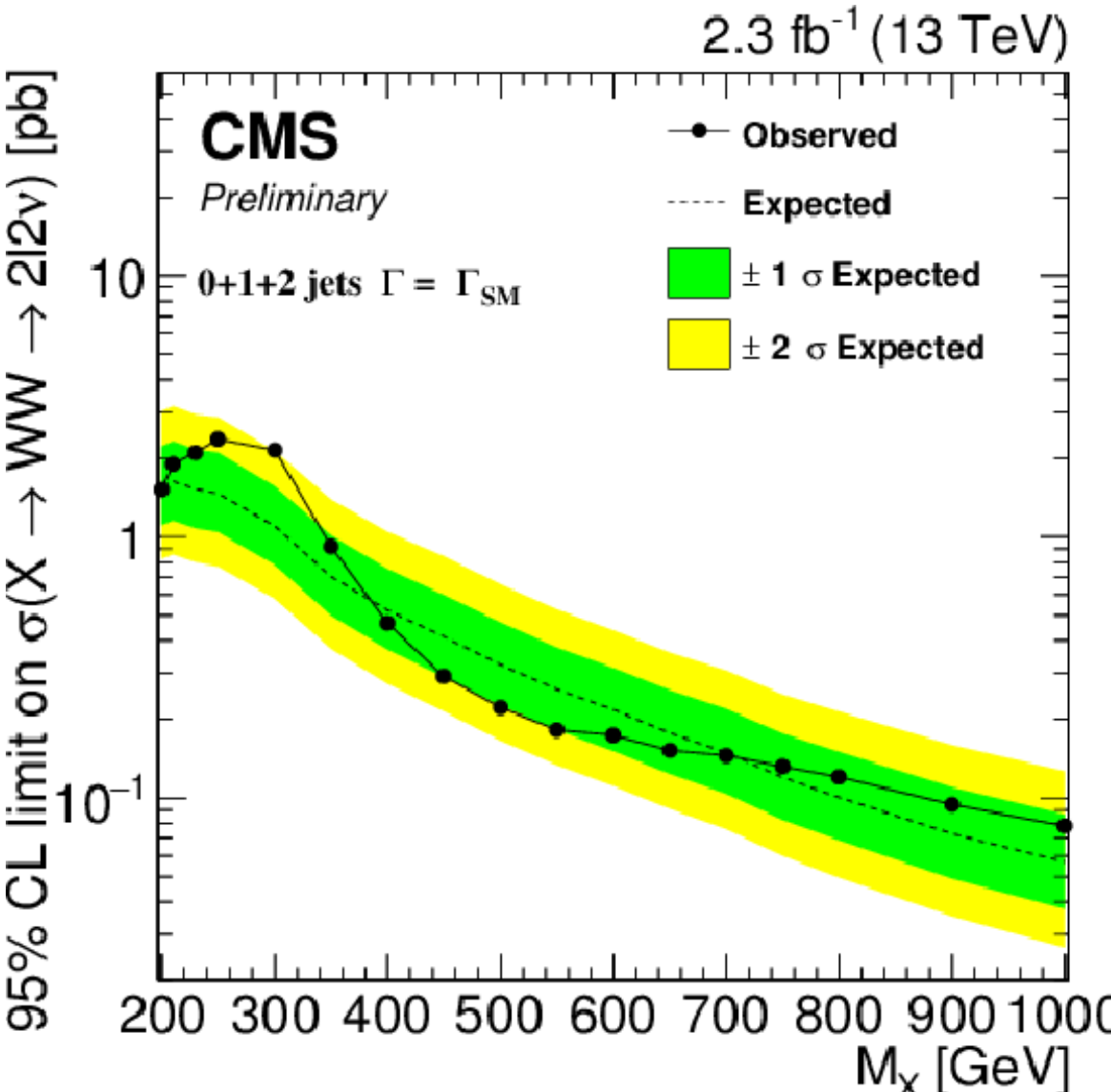
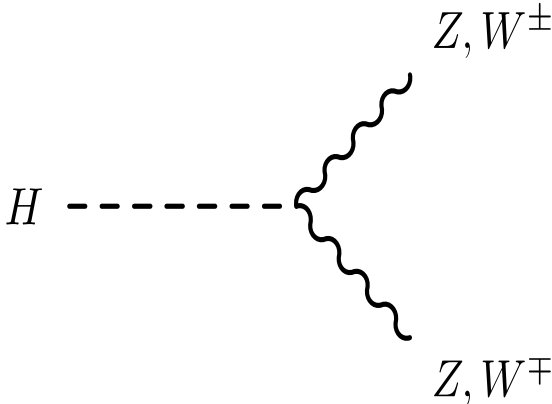


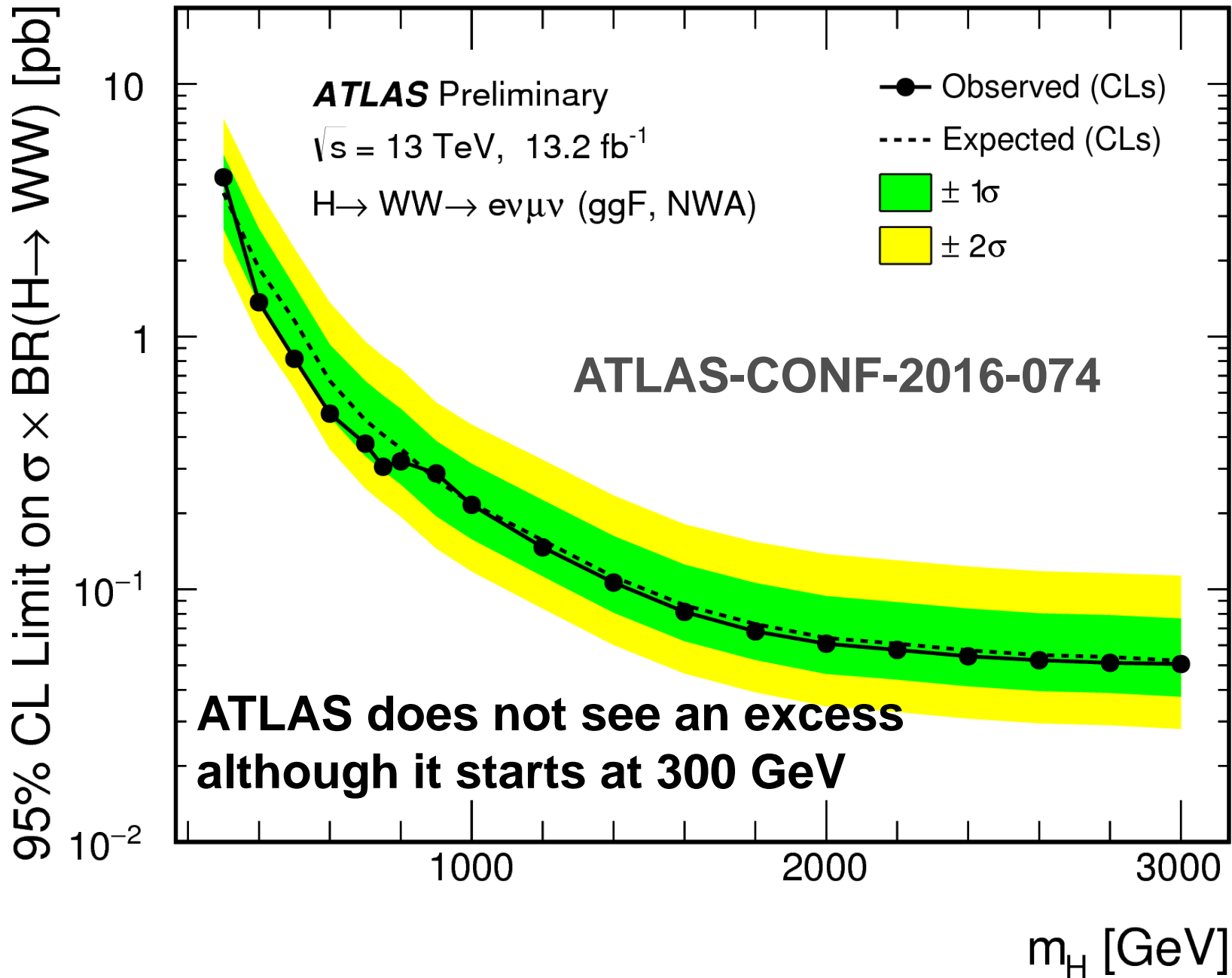
**Excess right below 250,
could consistent with Run I
measurement**

40



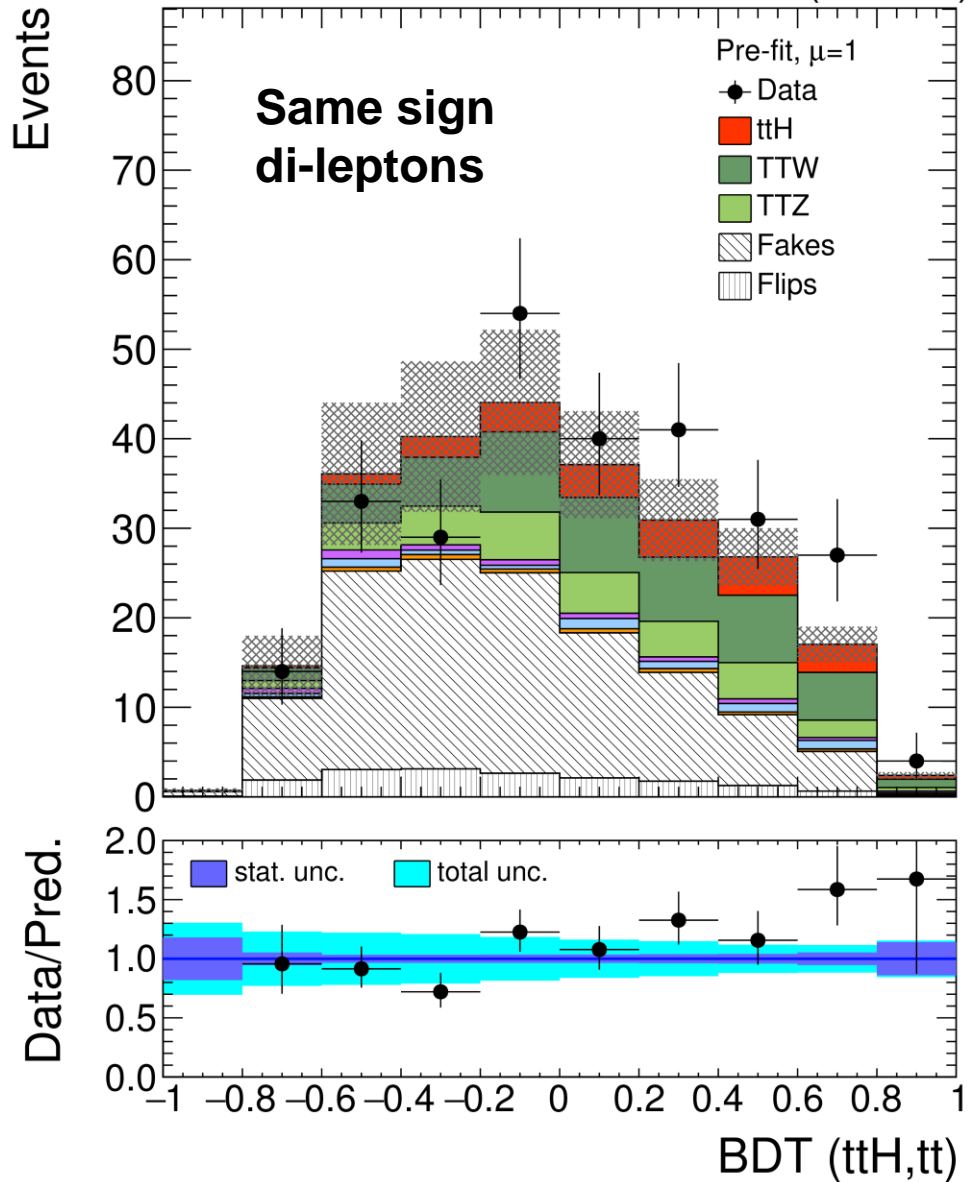
Search for high mass Higgs to WW with fully leptonic decays using 2015 data



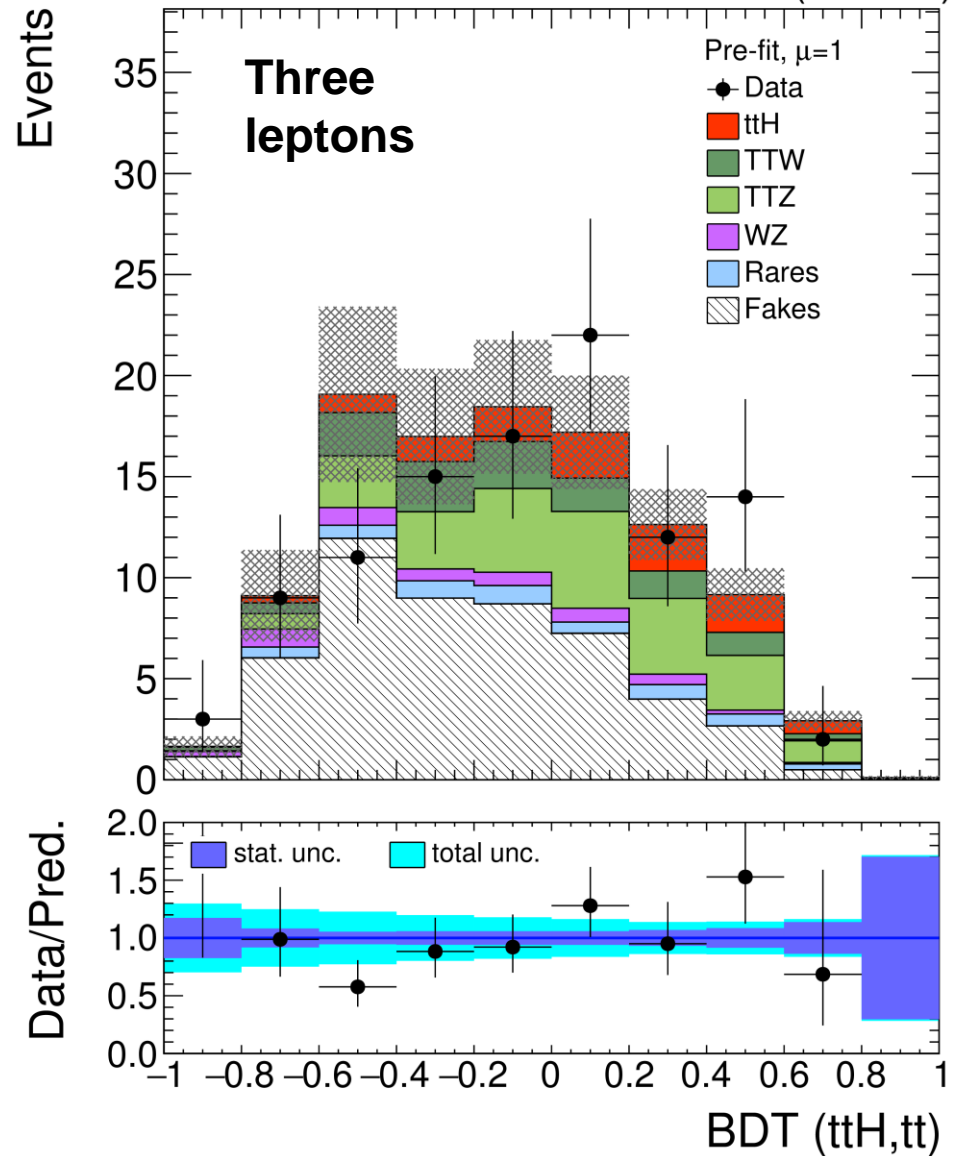


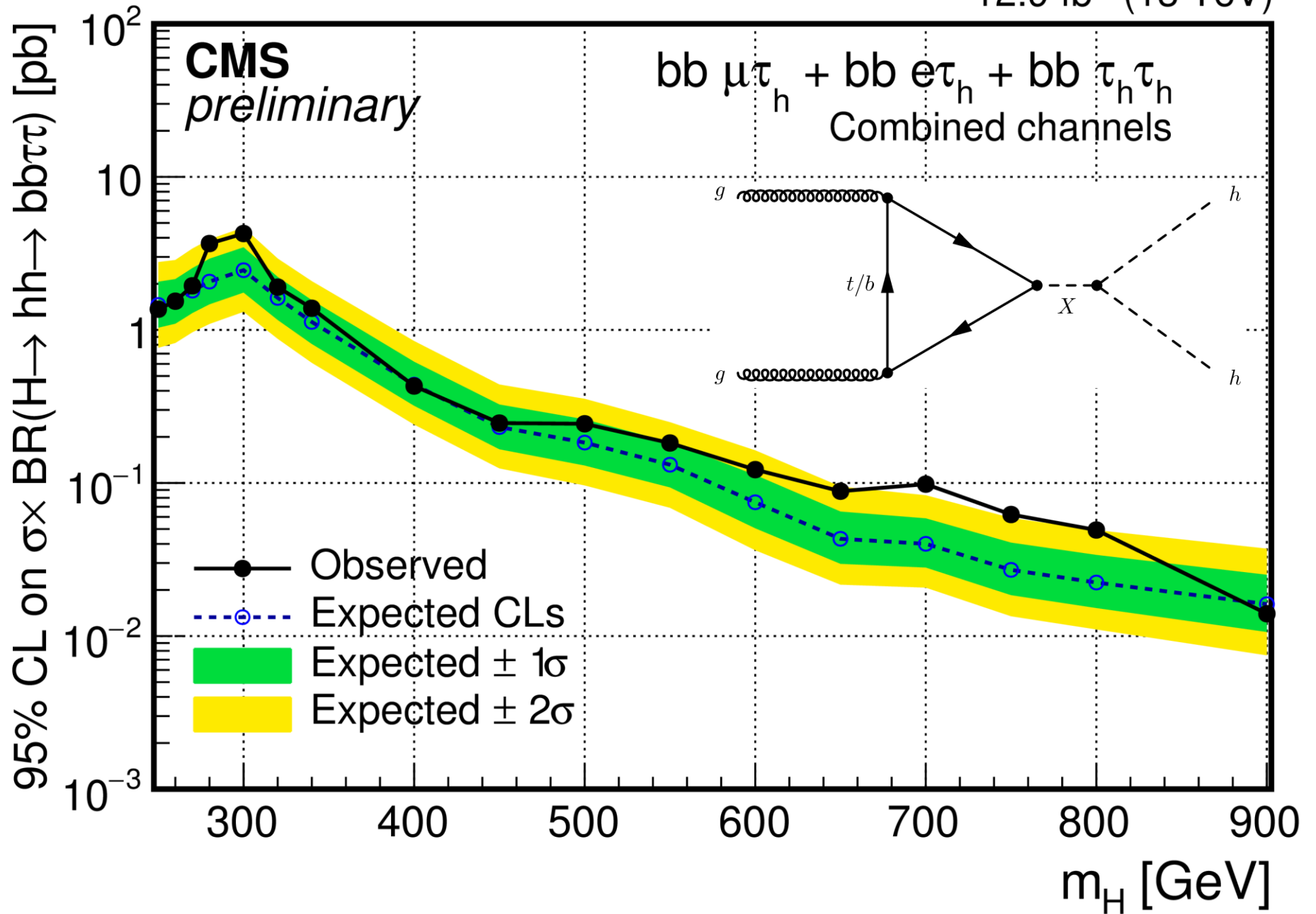
CMS-PAS-HIG-16-022

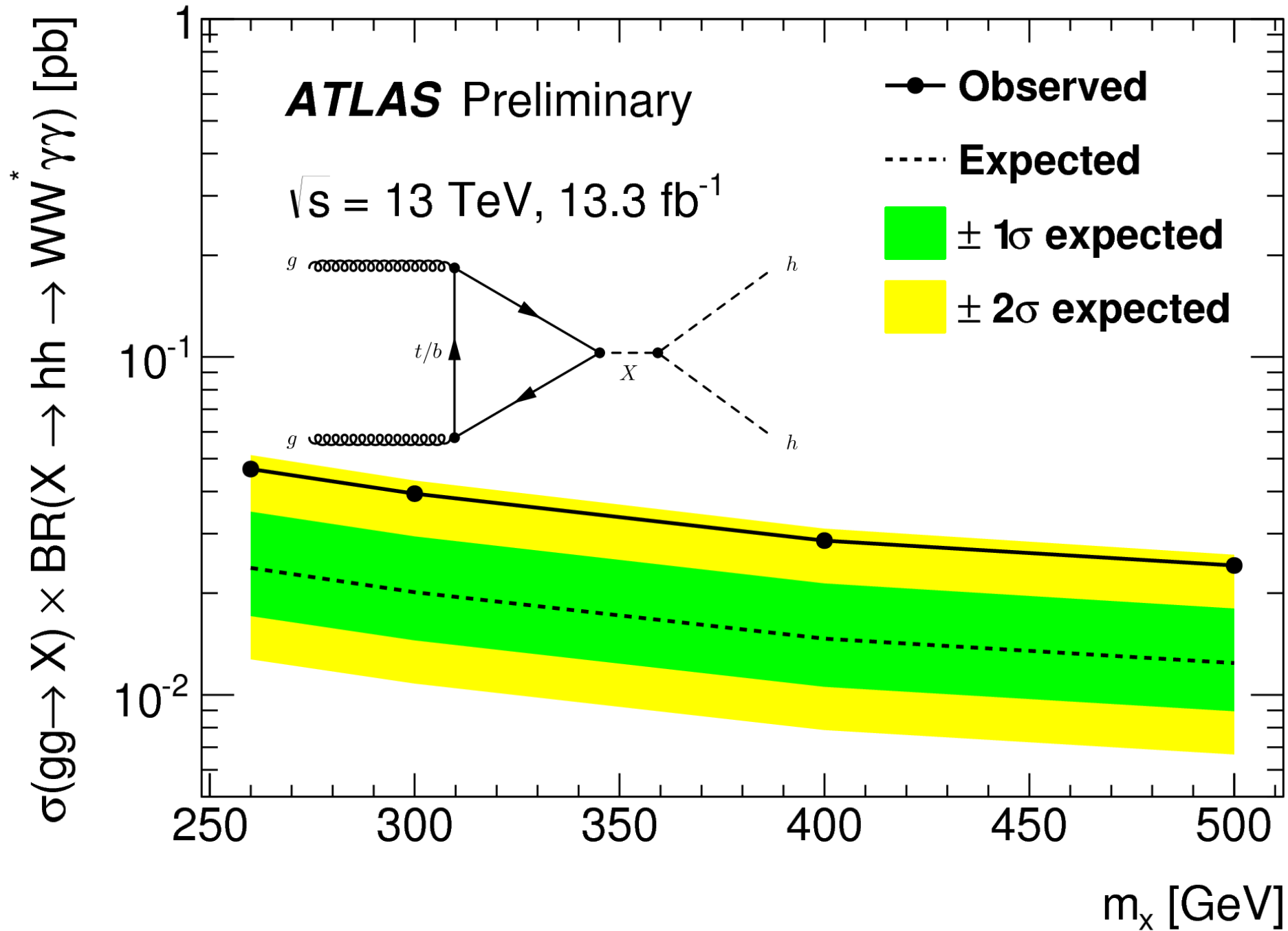
CMS Preliminary 12.9 fb⁻¹ (13 TeV)

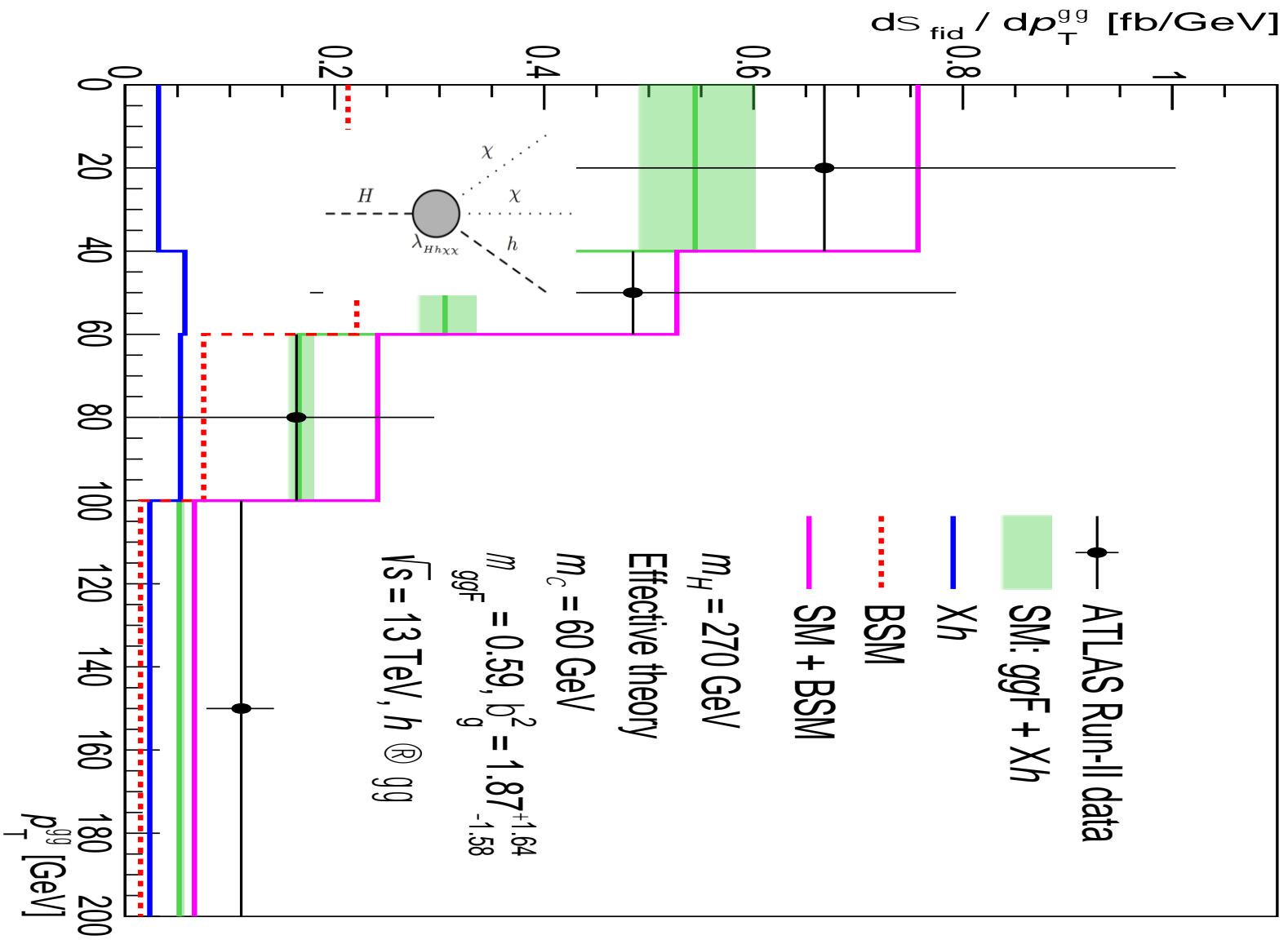


CMS Preliminary 12.9 fb⁻¹ (13 TeV)









Type of final state	Final state	Significance
Decays to weak vector bosons	ZZ->4l, ZZ->llvv, WW->2llvv	2σ, 2σ, 2σ, 1σ, 0σ
Top associated Higgs production	Leptons, $\gamma\gamma$, bbar	1σ
Di-Higgs resonance searches	bb$\tau\tau$, WW$\gamma\gamma$	2σ, 2σ
Higgs pT spectrum	$\gamma\gamma$	1σ

**Overall, positive fluctuations. One can argue that the new data does not contradict the original hypothesis
Combination underway (cannot add up in quadrature!)**

Top associated Higgs production (Run II)

CMS Preliminary $2.3+12.9 \text{ fb}^{-1}$ (13 TeV)

$m_H = 125 \text{ GeV}$

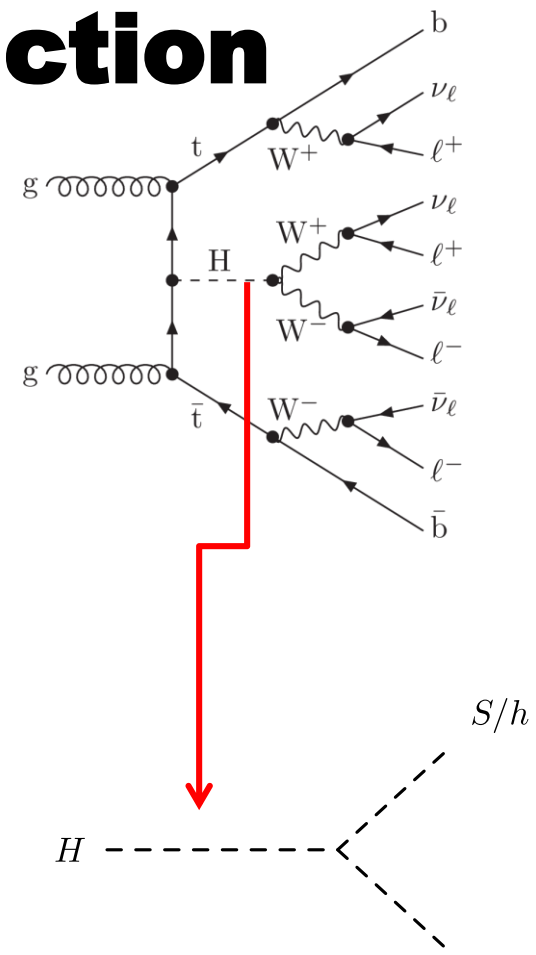
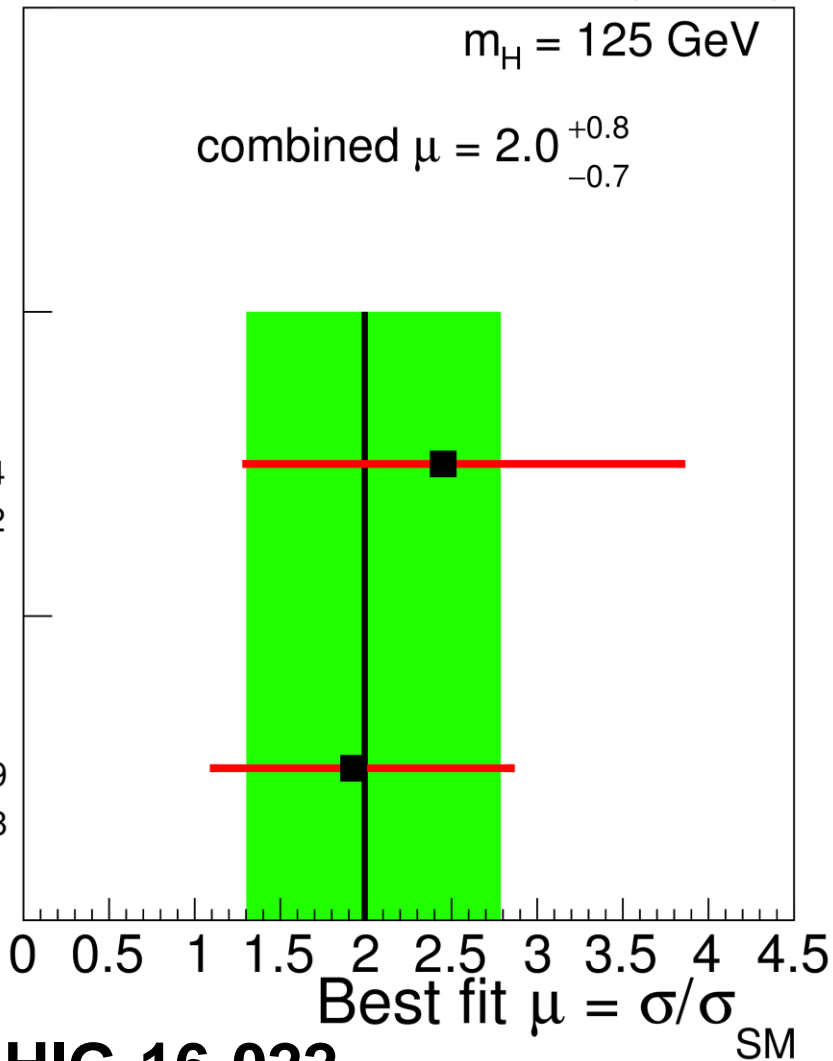
combined $\mu = 2.0^{+0.8}_{-0.7}$

trilepton

$\mu = 2.5^{+1.4}_{-1.2}$

dilepton

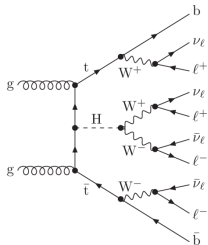
$\mu = 1.9^{+0.9}_{-0.8}$



$$S, h \rightarrow WW, \tau\tau ZZ$$

$$\rightarrow l^\pm l^\pm, 3l + X$$

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ATLAS Preliminary

$\sqrt{s}=13$ TeV, 13.2-13.3 fb⁻¹

— total

— stat.

(tot.) (stat. , syst.)

$ttH(H \rightarrow \gamma\gamma)$
(13 TeV 13.3 fb⁻¹)

-0.3 ^{+1.2}_{-1.0} (^{+1.2} ^{+0.2}
_{-1.0} , _{-0.2})

$ttH(H \rightarrow WW/\tau\tau/ZZ)$
(13 TeV 13.2 fb⁻¹)

2.5 ^{+1.3}_{-1.1} (^{+0.7} ^{+1.1}
_{-0.7} , _{-0.9})

$ttH(H \rightarrow b\bar{b})$
(13 TeV 13.2 fb⁻¹)

2.1 ^{+1.0}_{-0.9} (^{+0.5} ^{+0.9}
_{-0.5} , _{-0.7})

ttH combination
(13 TeV)

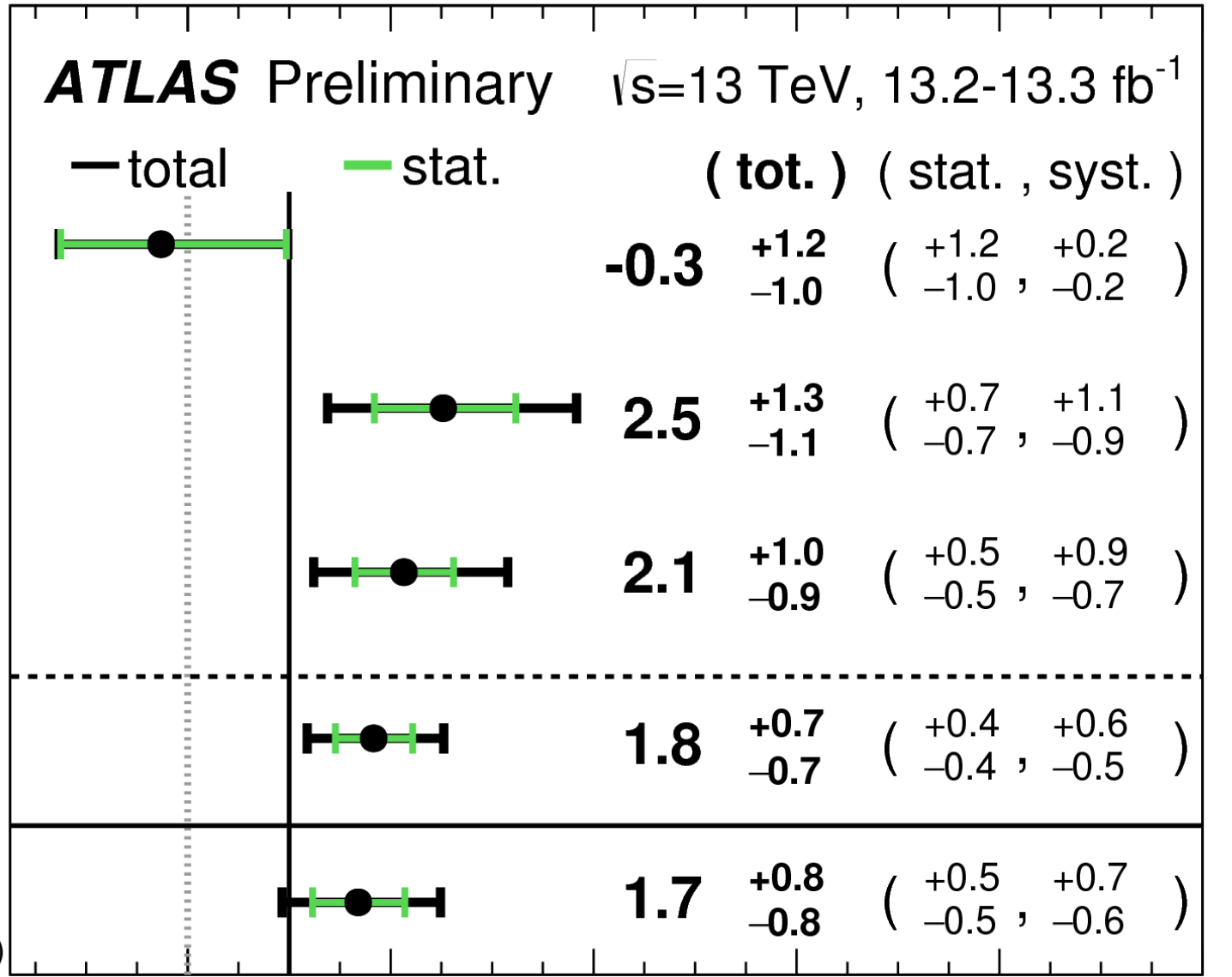
1.8 ^{+0.7}_{-0.7} (^{+0.4} ^{+0.6}
_{-0.4} , _{-0.5})

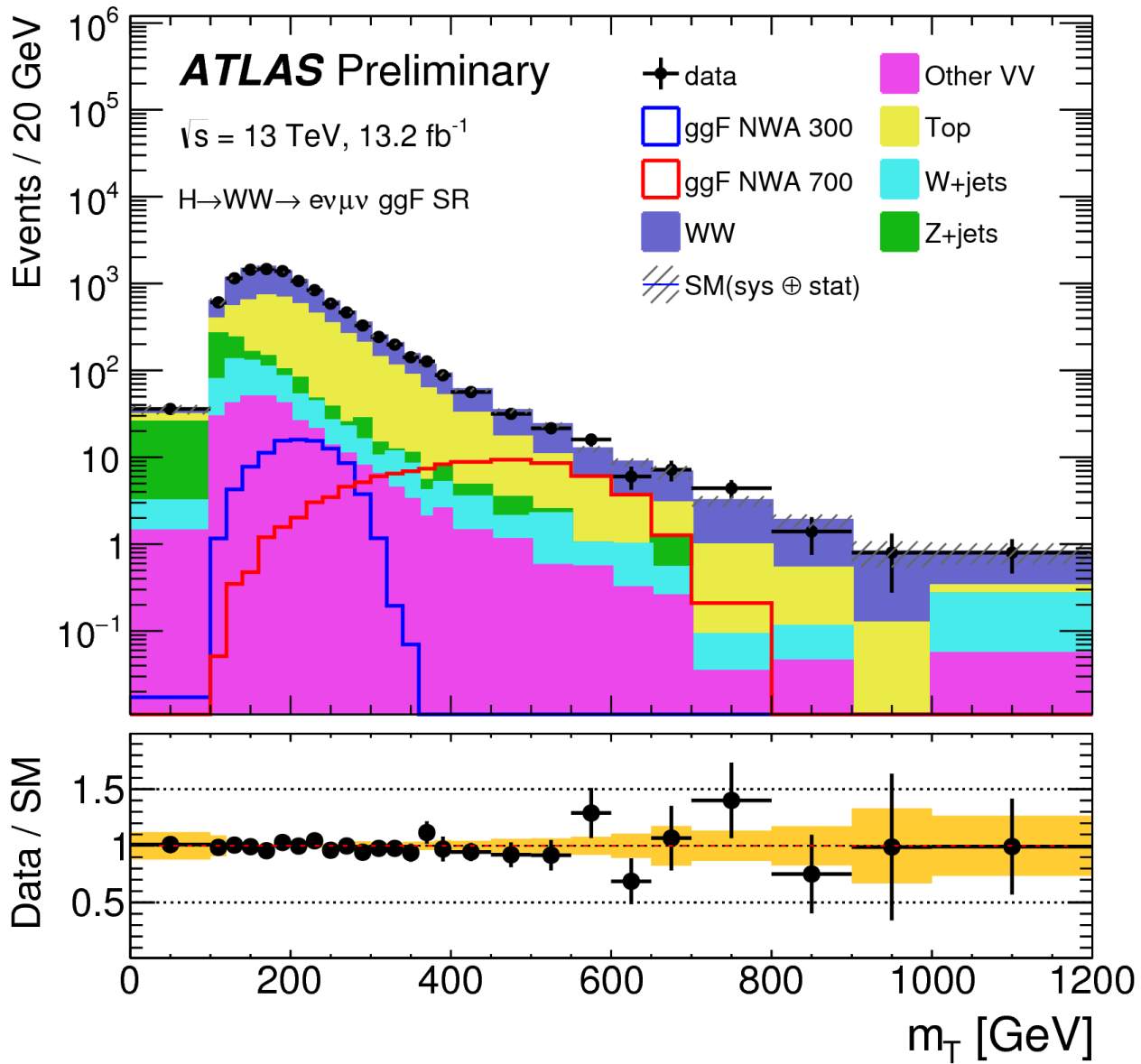
ttH combination
(7-8TeV, 4.5-20.3 fb⁻¹)

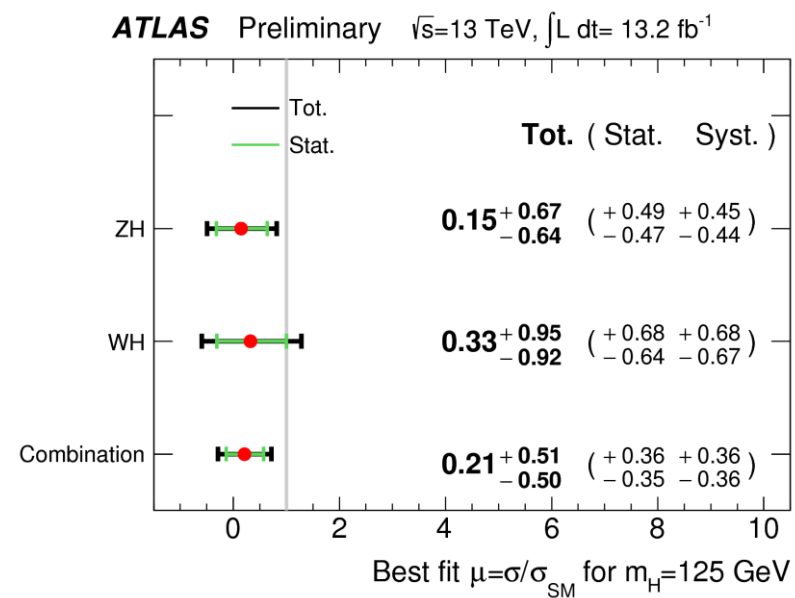
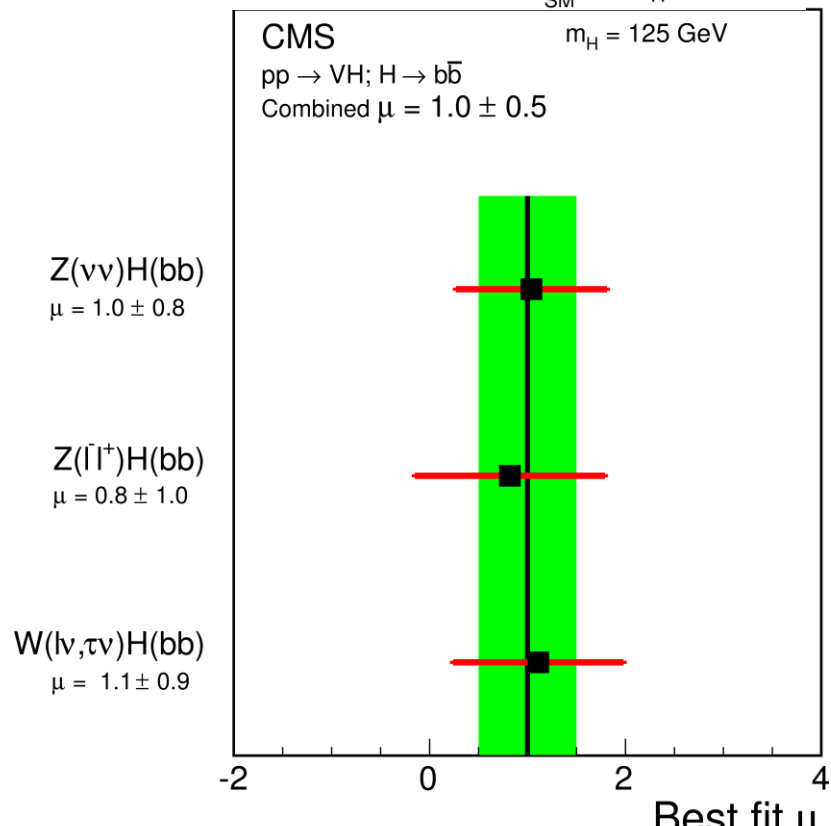
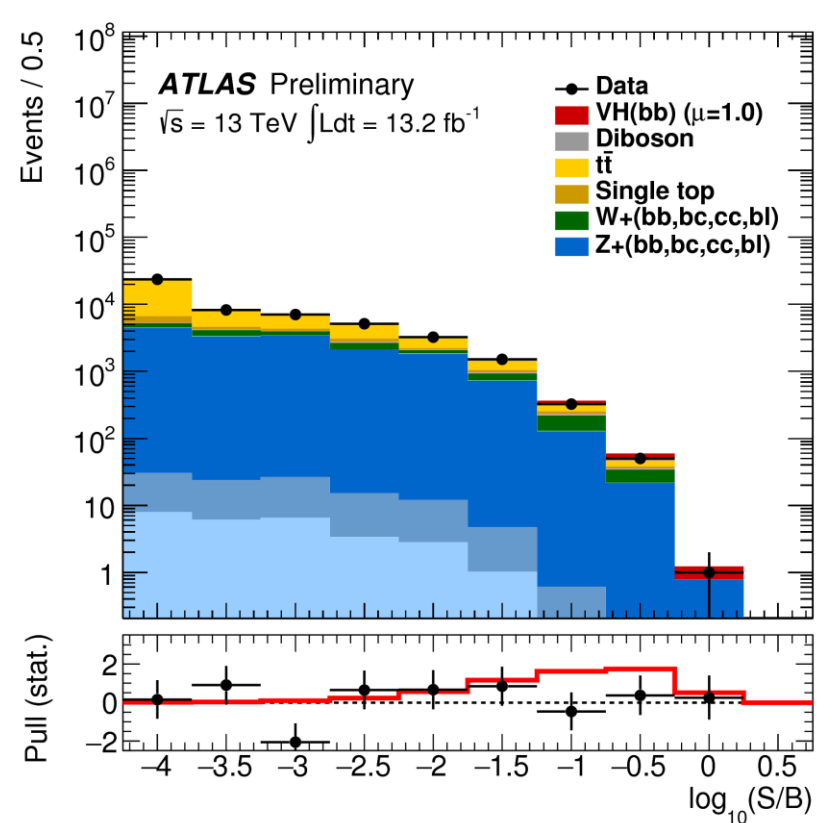
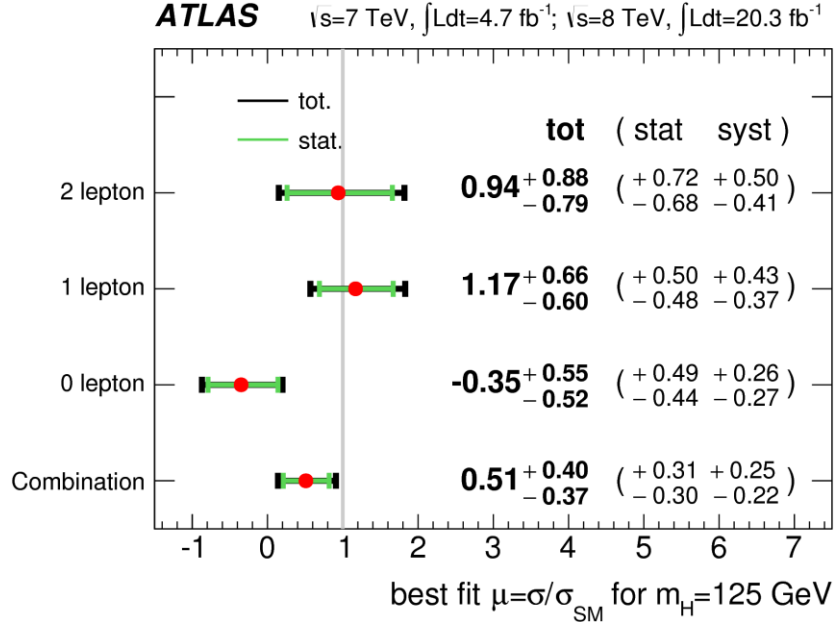
1.7 ^{+0.8}_{-0.8} (^{+0.5} ^{+0.7}
_{-0.5} , _{-0.6})

$S, h \rightarrow WW, \tau\tau ZZ$
 $\rightarrow l^\pm l^\pm, 3l + X$

best fit μ_{ttH} for $m_H=125$ GeV







Using the full theory including the S boson

