

Cyclotrons

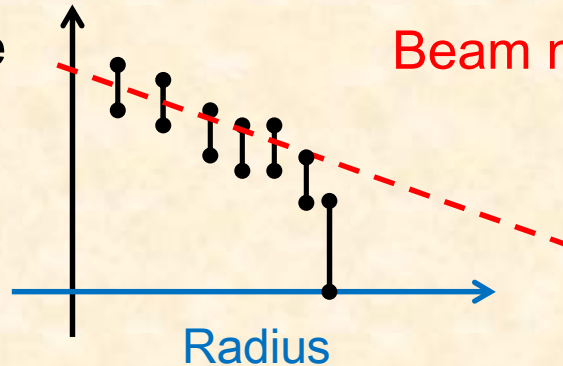
Chapter 4 : theory versus reality

- Isochronism and Phase measurement
- Isochronism 2nd approach & Kf limitation
- Resonances and tune in a cyclotron

- Research applications
- Medical applications

Phase measurement: Isochronous field correction : ΔB

Relative phase
Beam vs RF

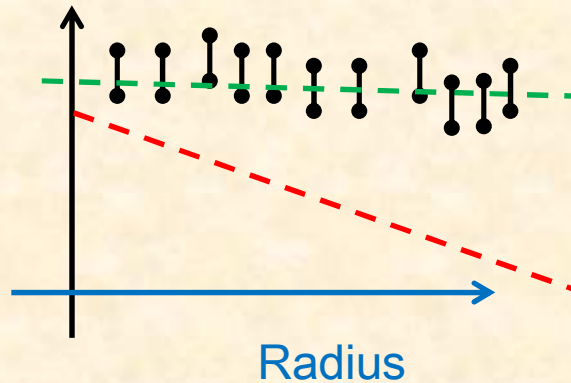


Beam not in phase with RF : B field not OK

$$B_z(R) \neq B_{z0} / \sqrt{1 - (R\omega_{rev})^2 / c^2}$$

Bz(R) not correct

Relative phase
Beam vs RF

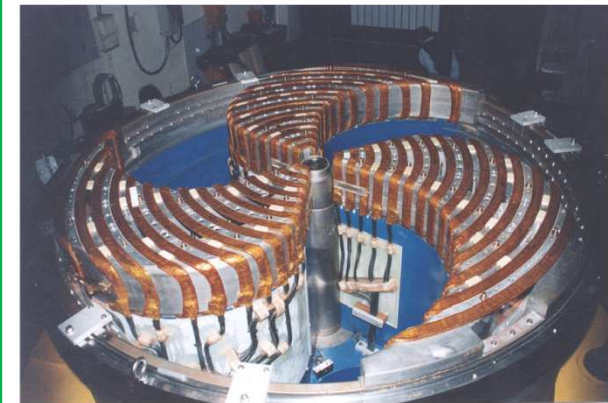


Correction of B : B(/+ ΔI)

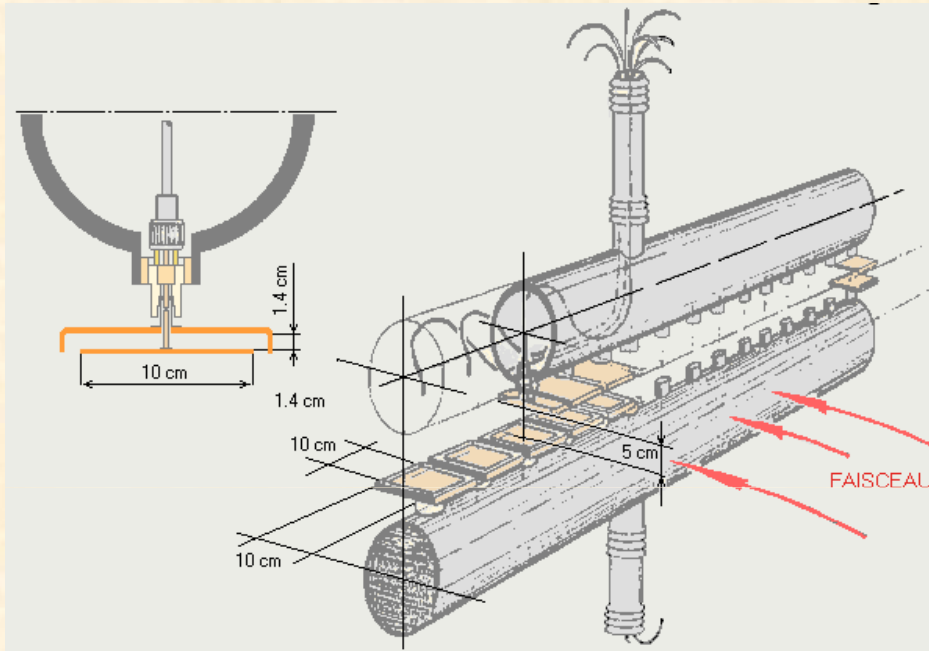
Bz(R) + ΔB_z

Bz(R) correct : Beam in
phase with RF

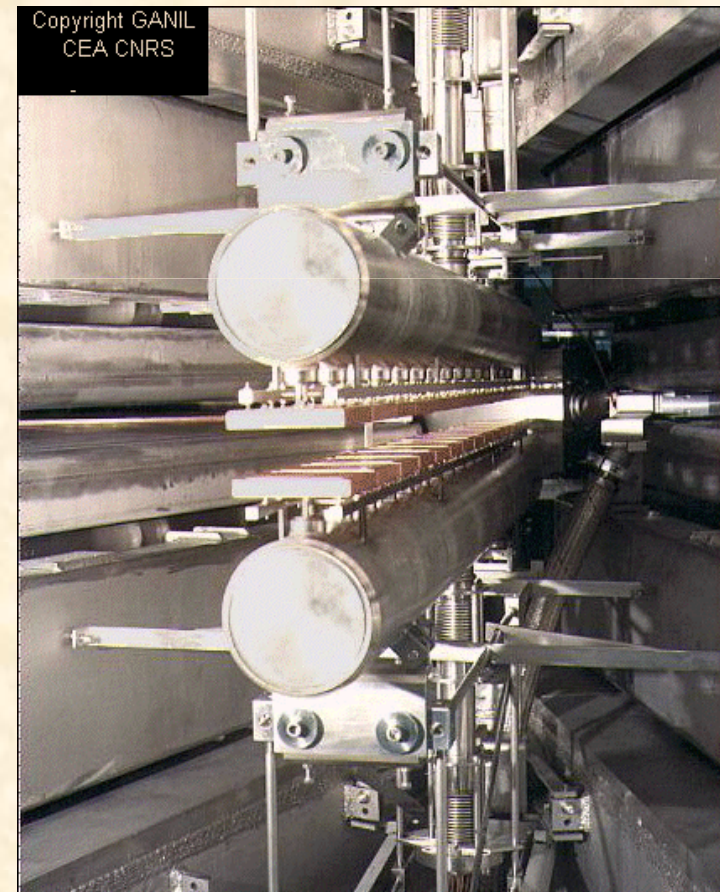
Correction trim coils, AGOR



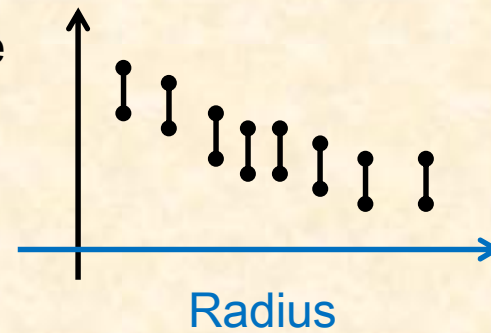
Phase measurement: Isochronism



Ganil, Caen (Fr) : CSS1

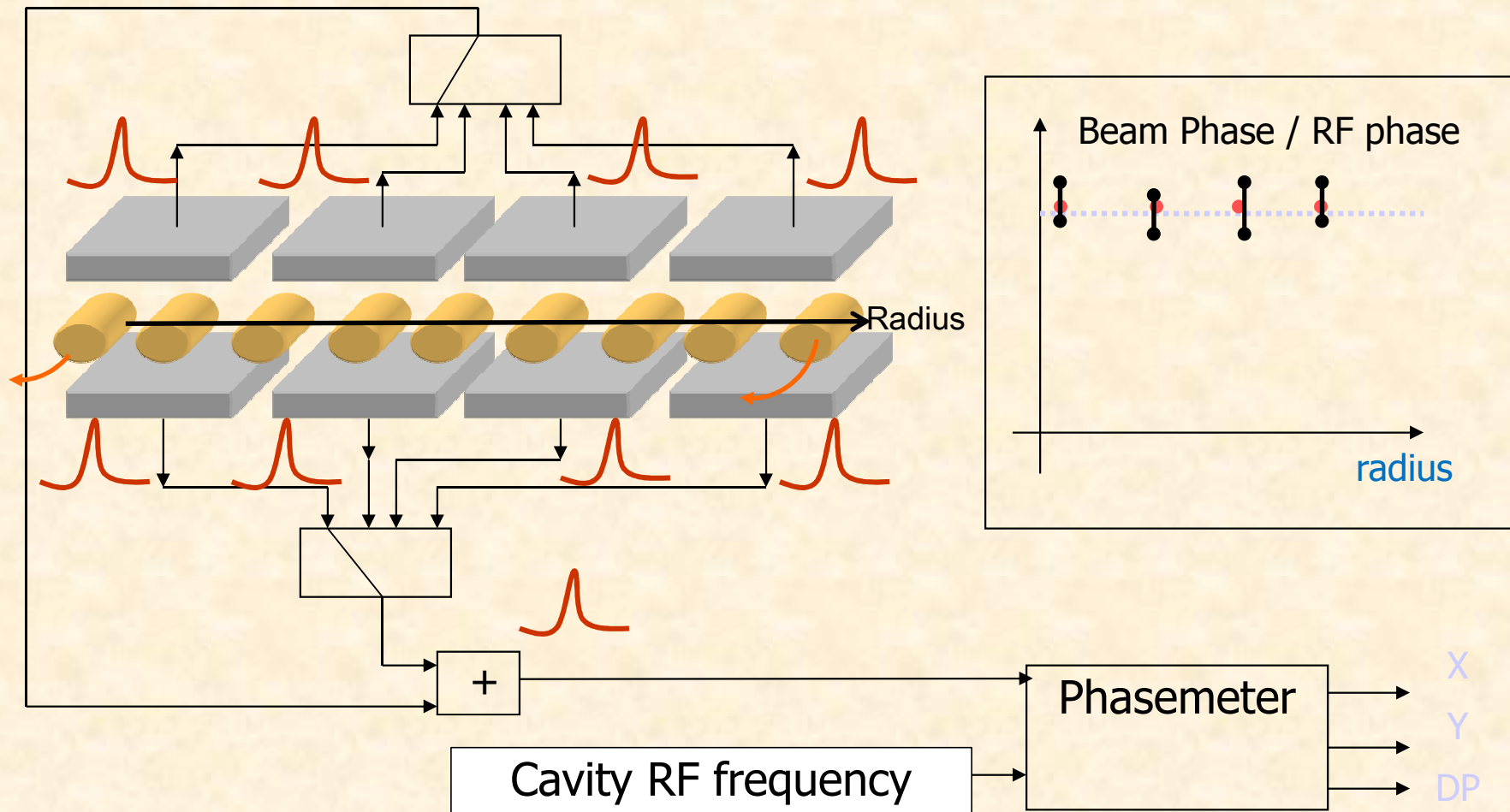


Relative phase
Beam vs RF



Isochronism & Phase measurement

Measuring $B(R)$ or $n(R)$ is difficult, While the $\Phi(R)$ is more sensitive



Isochronous field $B(z)$ = good field index $n(R)$

$$n = - \frac{R}{B_{0z}} \frac{\partial B_z}{\partial R}$$



$$\frac{dB}{B} = - n \frac{dR}{R}$$

$$B \rho = \langle B \rangle \cdot \langle R \rangle = \frac{p}{q}$$

Longitudinal dynamics lecture

$$\frac{dp}{p} = \frac{dB}{B} + \frac{dR}{R} = (1 - n) \frac{dR}{R}$$

$$\frac{dp}{p} = \gamma^2 \frac{d\beta}{\beta} = \gamma^2 \frac{d(\omega_{rev} \cdot R)}{\omega_{rev} \cdot R} = \gamma^2 \frac{dR}{R}$$

$$1 - n = \gamma^2$$

$$n(R) = (1 - \gamma^2)$$

« At high energy » isochronism requires $n \ll 0$

Bz Azimutal modulations are not sufficient
It is a (Focusing) limit for high energy isochronous cyclotron

Max Energy for Superconducting Cyclotrons not limited by (B × Rextraction)

Because of the focusing limitation due to the Flutter dependence on the B field, the **max energy** is not given by $Kb \sim 48 (B.R_{extraction})^2$

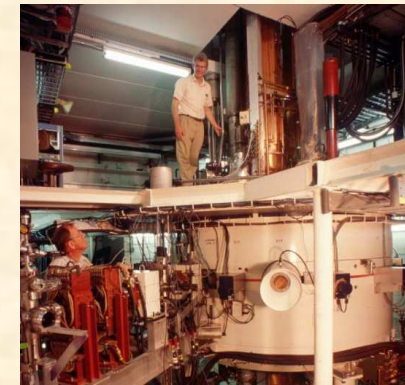
but K_f the so-called “*focusing factor*”:

$$\left[\frac{E}{A} \right]_{\max} \neq Kb \cdot \left\{ \frac{Q}{A} \right\}^2$$

vertical oscillation+isochronous field condition

$$v_z^2 = n + \frac{N^2}{N^2 - 1} F (1 + 2 \tan^2 \epsilon) > 0$$

$$\frac{N^2}{N^2 - 1} F (1 + 2 \tan^2 \epsilon) > -n = \gamma^2 - 1$$



•Focusing limitation (stronger than B limitation)

$$\left[\frac{E}{A} \right]_{\max} = K_f \cdot \left\{ \frac{Q}{A} \right\}^2 < Kb \cdot \left\{ \frac{Q}{A} \right\}^2$$

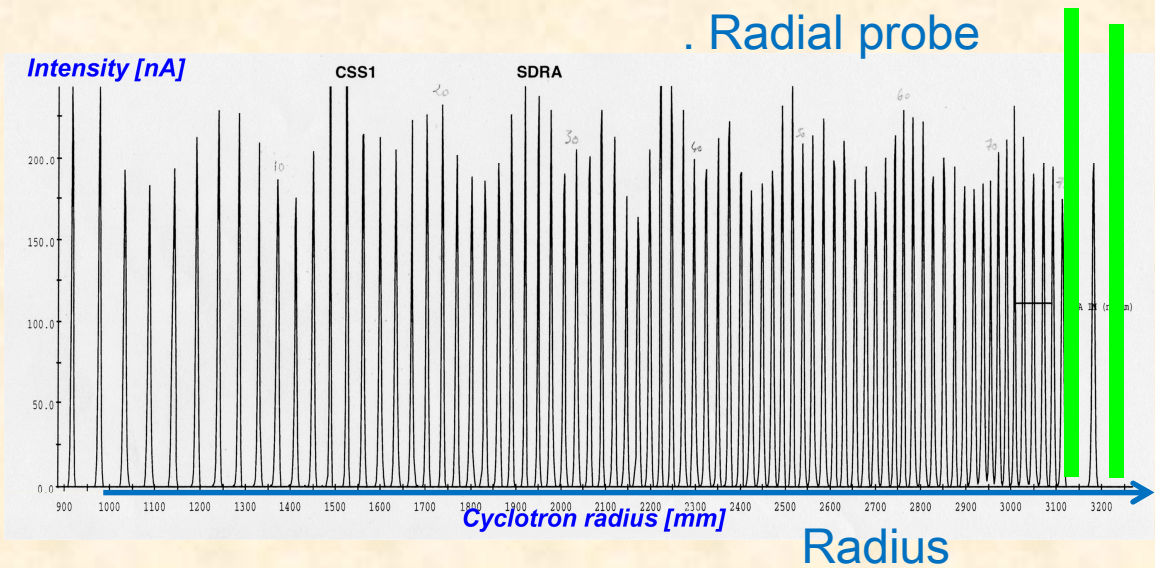
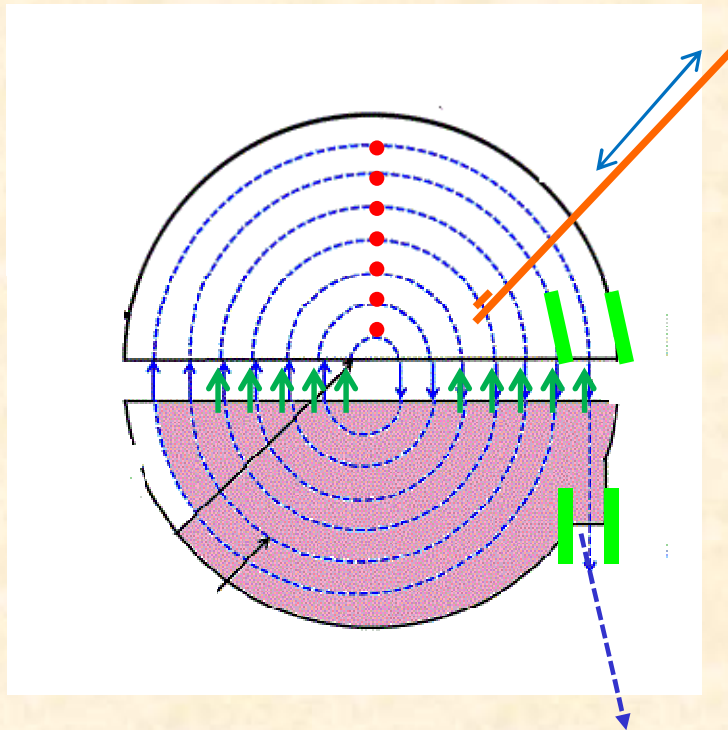
$$Kb \sim 48 (B.R_{extract})^2$$

$$K_f \sim f(\text{FLUTTER})$$

Radial probes

usefull tool for acceleration, precession study

Monitoring turns with a
Radial probes



Radial probe : $I = F(\text{Radius})$

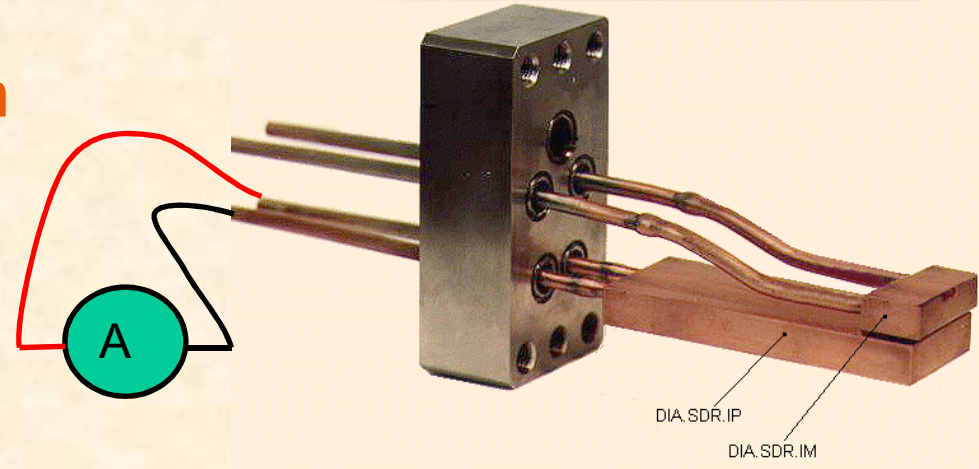
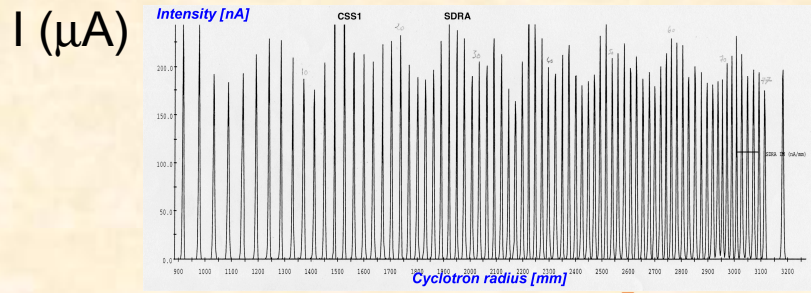
Turn separation : $\delta r = R(\text{turn } N) - R(\text{turn } N-1)$

$\delta r = \text{Acceleration} + \text{Oscillation}$

$\delta r \sim \alpha VRF \cos(\phi) + \text{Oscillation}$

Current measurement: Radial probe

A Faraday cup on mobile Arm



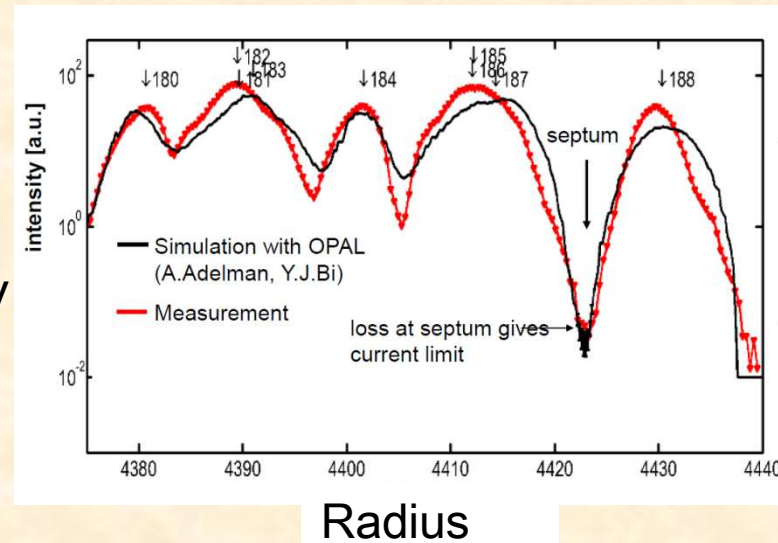
Ganil Csx1 (Fr)
Kb=380 , heavy ions

Radius

Check of the
acceleration & extraction

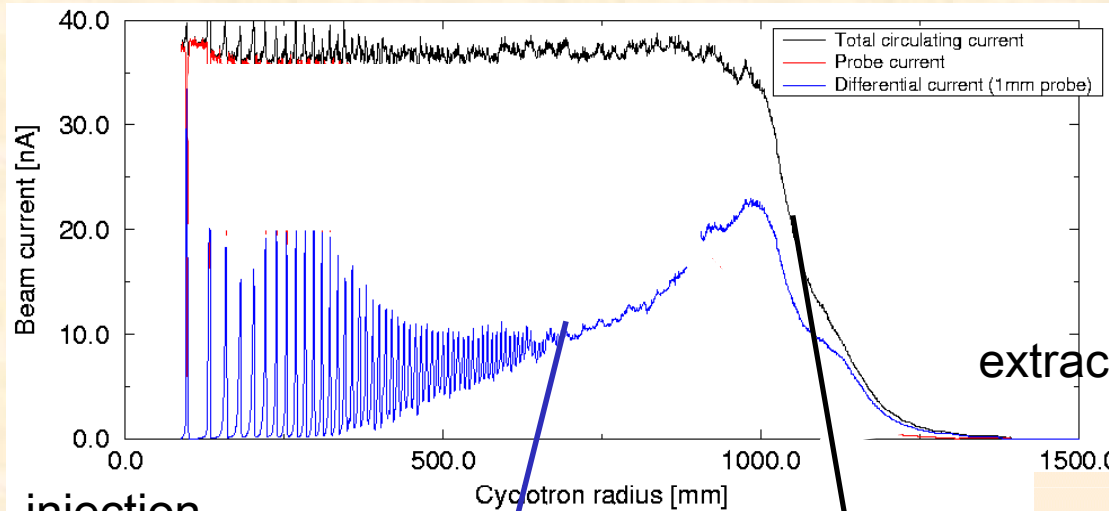
PSI
Kb=590 MeV

Check of the
extraction



Current measurement with a Radial probe

A full check of the dynamics



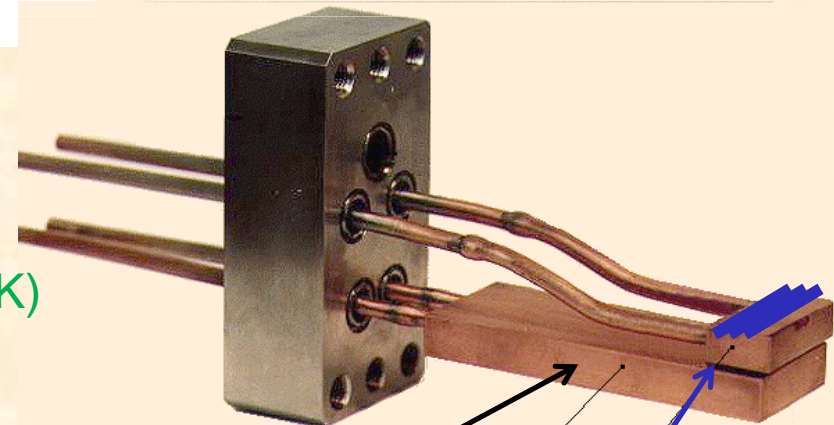
Radial probes with
2 diagnostics :
-1 small finger
-1 large plate

injection

extraction

Beam losses
Before extraction
(B field was not OK)

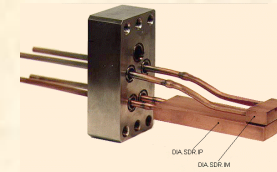
Current on small finger
of the radial probe



Total current
with a large plate

Finger current
Good resolution in r (turn separation)

Tune ν_r measurement with radial probes



:

$$r(t) = \langle r(t) \rangle + X_0 \cos(\nu_r \omega_0 t)$$

$$\nu_r = \sqrt{1 - n}$$

$$\omega_0 t = \text{PHASE}$$

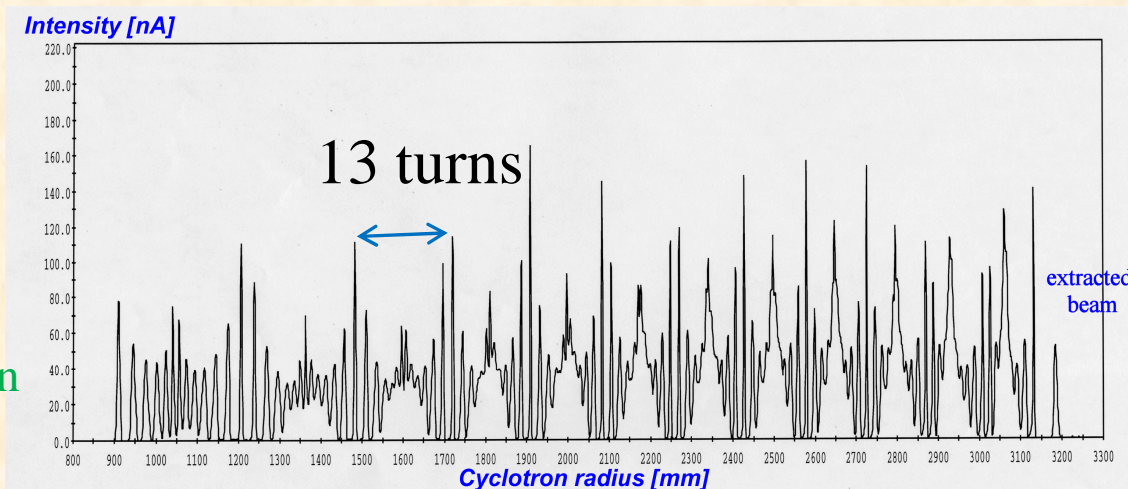
$$\omega_0 \Delta t = 360^\circ = 1 \text{ turn}$$

Precession : $\delta r = \delta r_{\text{acc}} + X_0 \cos(\nu_r \omega_0 t)$

Centering error
At injection

Large X_0

Large Oscillation



1 period for

13 turns \Rightarrow

$$\nu_r = 1/13 = 0.08$$

Back to dynamics and resonances at high energy

During the acceleration, ν_r and ν_z change because $\nu_{r,z} \propto B(r)$

The plot of ν_r vs ν_z is called the working point diagram.

Like any **oscillatory phenomenon**, the amplitude of a betatronic motion can grow uncontrolled whenever an external source excites it with its own frequency.

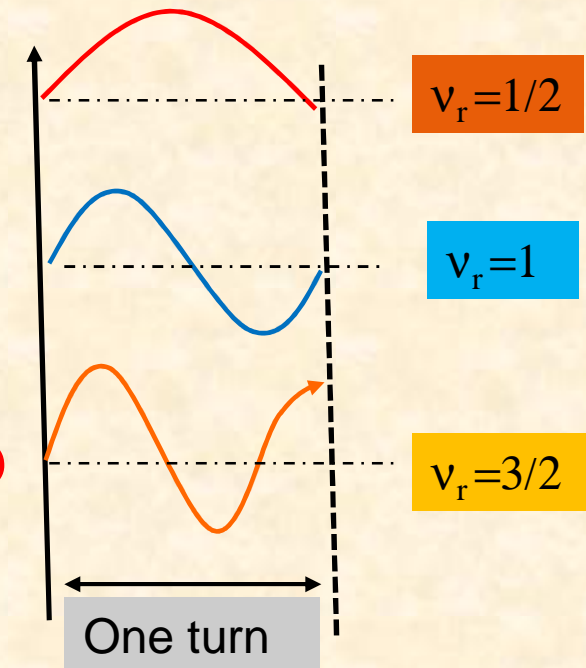
$$\ddot{x} + \nu_r^2 \omega_0^2 x = 0$$

Several kind of radial resonances

$$x(t) = x_0 \cos(\nu_r \omega_0 t) = x_0 \cos(\nu_r 2\pi N \text{turn})$$

can be excited with field defect, injection angle (unwanted)

with field bump, injection angle (desired)



What happen with P field perturbations on 1 turn

$$\ddot{z} + [v_z \omega_0]^2 z = \Delta \cos(P \omega_0 t)$$

Search a particular solution

$$z(t) \propto \cos(P \omega_0 t) ???$$



$$\frac{z(t)}{\Delta} = \frac{\cos(P \omega_0 t)}{[v_z \omega_0]^2 - [P \omega_0]^2}$$

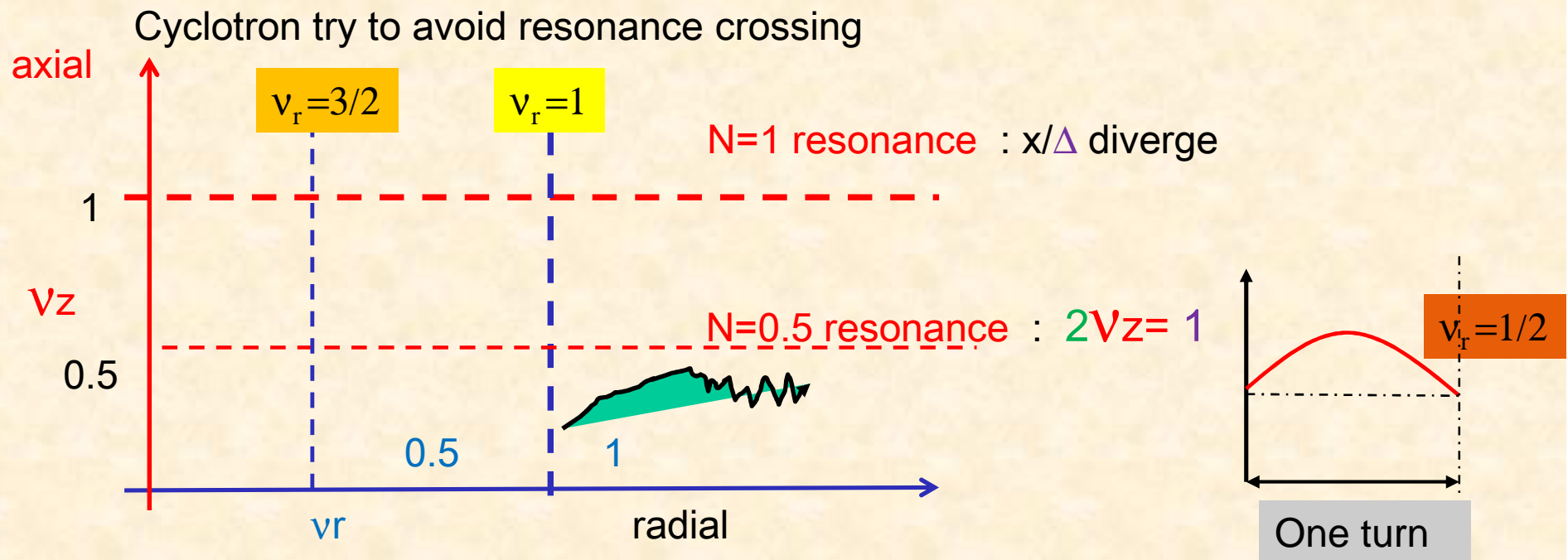
z/Δ diverge at $v_z=P$ (integer)

z(t) is very sensitive to any perturbation Δ

(not good : instabilities = beam losses)

With P field perturbations on L turns

z/Δ diverge at $L \cdot v_z = P$ (integer)



Resonances

Systematic resonances : This resonance occurs as the betatronic frequency is a multiple of the "geometrical frequency" of the cyclotron. In this case, any kick given to the particle because of its particular position will be experienced again and again.

$$(\mathbf{K} \cdot \mathbf{v}_r = \mathbf{P}) \quad // \quad (\mathbf{L} \cdot \mathbf{v}_z = \mathbf{P})$$

Coupling resonances

Under proper circumstances and frequency ratios, the 2 oscillators can be coupled and the energy stored in one motion, transferred to the other.

These are **coupling resonances** ($\mathbf{K} \cdot \mathbf{v}_r + \mathbf{L} \cdot \mathbf{v}_z = \mathbf{P}$).

$$\mathbf{K} \cdot \mathbf{v}_r + \mathbf{L} \cdot \mathbf{v}_z = \mathbf{P} \quad \mathbf{K}, \mathbf{L} \text{ and } \mathbf{P} \text{ integer}$$

The particle's working point curve should avoid or cross as fast as possible those lines.

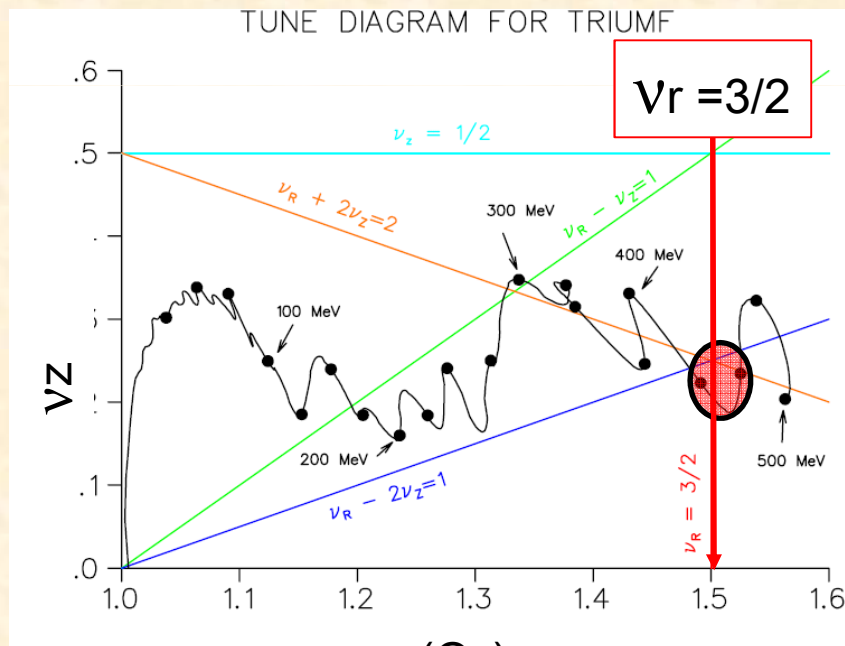
Tunes and resonances at Triumf (Canada) (H⁻ cyclo, K_b=520 MeV, 6 sectors)



$$\mathbf{K} \cdot \mathbf{v}_r + \mathbf{L} \cdot \mathbf{v}_z = \mathbf{P}$$

$|\mathbf{K}| + |\mathbf{L}|$ is called the resonance order (1, 2, 3 ...)

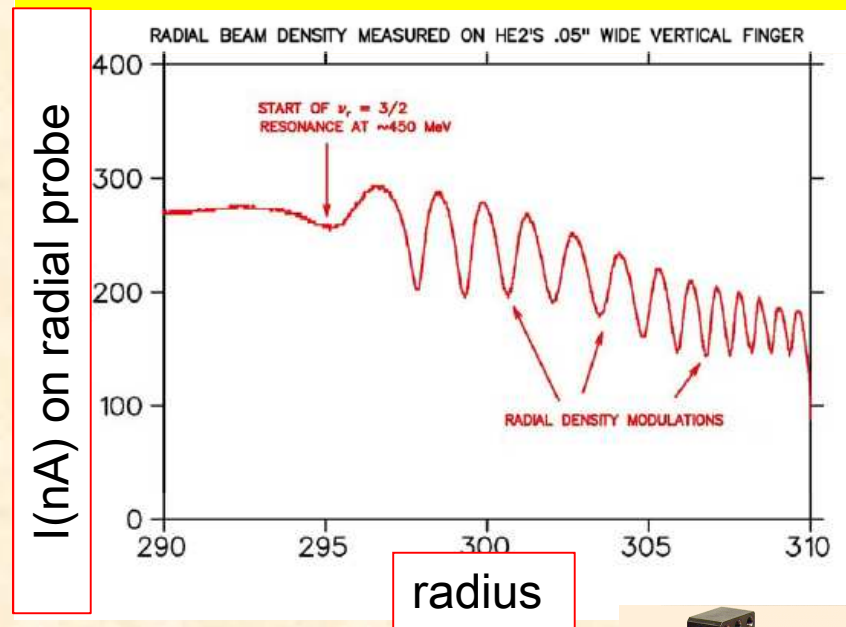
Effect on the crossing of the resonance $2 \cdot \nu_r = 3$ (order 2)



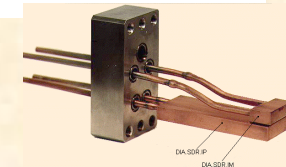
ν_r (Qx)

$$\nu_r^2 = 1 - n + \dots = 1 - (1 - \gamma^2) + \dots \sim \gamma^2$$

Radial Oscillations : $2 \cdot \nu_r = 3$



radius



Cyclotrons in the world

Research laboratories with Cyclotron(s)

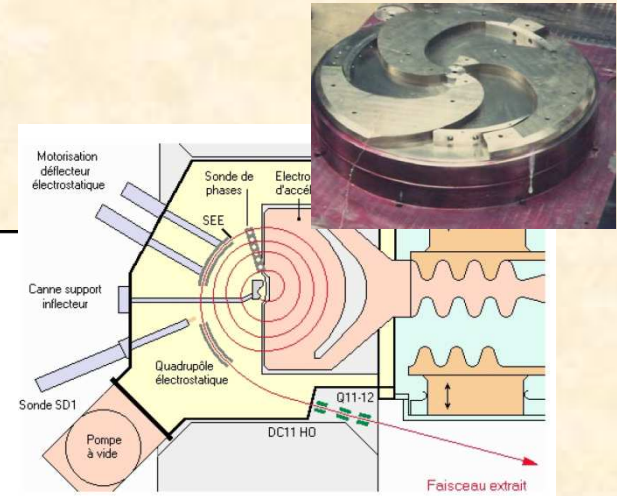
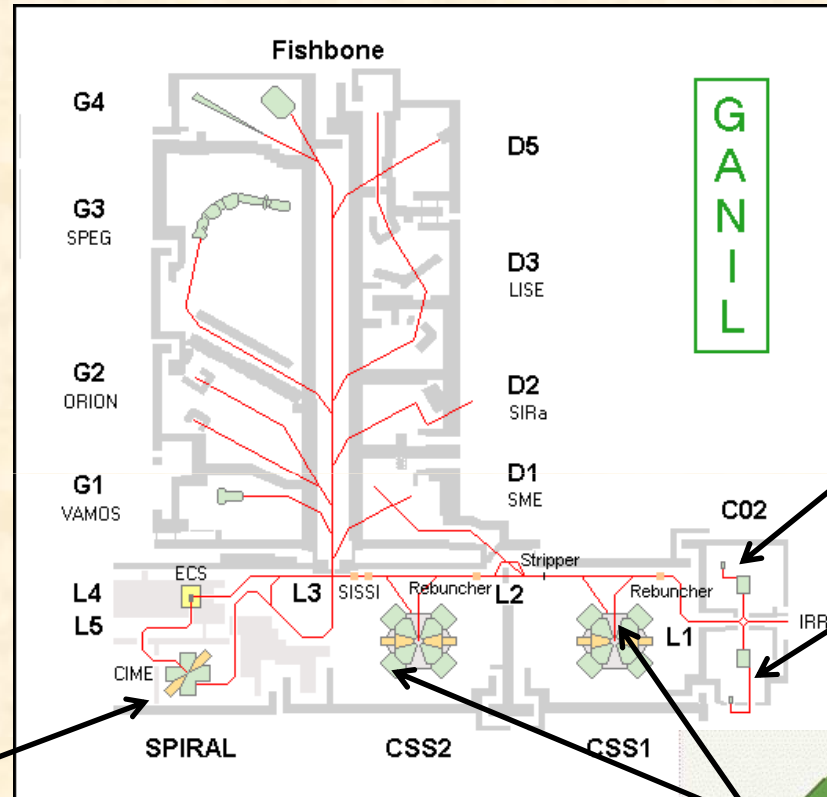


Some of the Research Facility in the world using cyclotrons

GANIL Facility (Caen, Fr)

5 cyclotrons + a new Linac

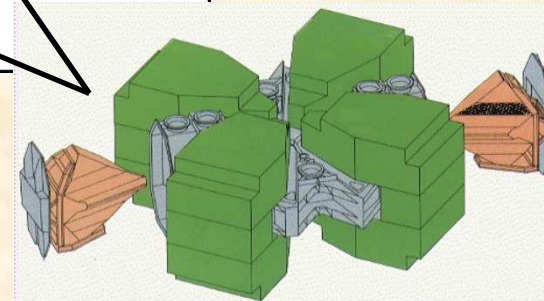
Nuclear physics
Atomic physics
Solid state
Radiobiology



2 compact cyclo
Kb=28 MeV



1 compact cyclo Kb =265
For radioactive ion

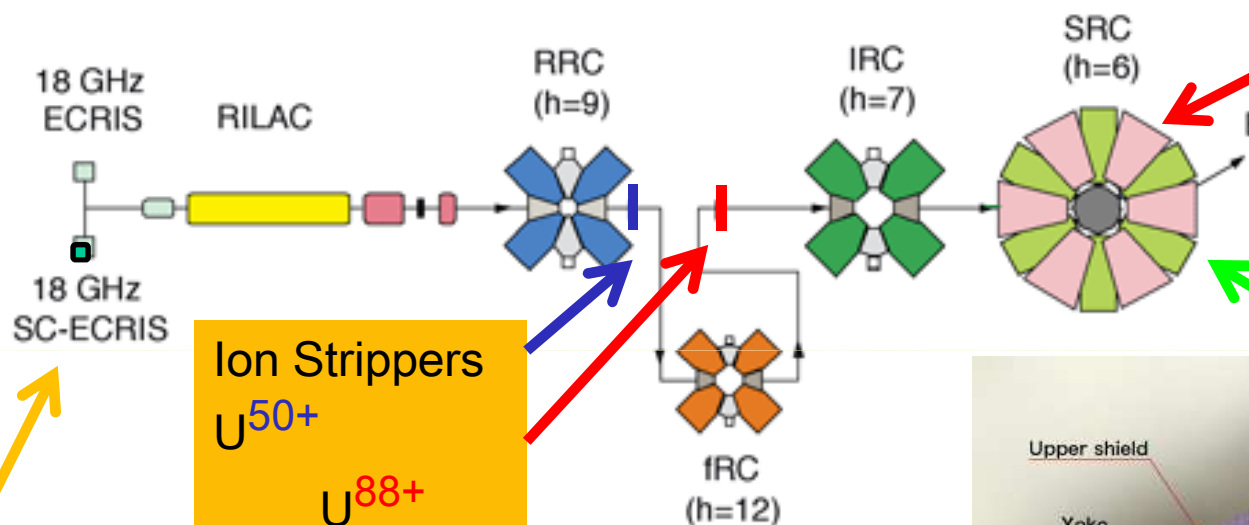


2 Separated Sectors cyclotrons
Kb =380 MeV

RIBF (Tokyo, Japan) :

Uranium beam $^{238}\text{U}^{88+}$ @345 MeV/A cw

Mode (1): RILAC + RRC + (stripper2) + fRC + (stripper3) + IRC + SRC



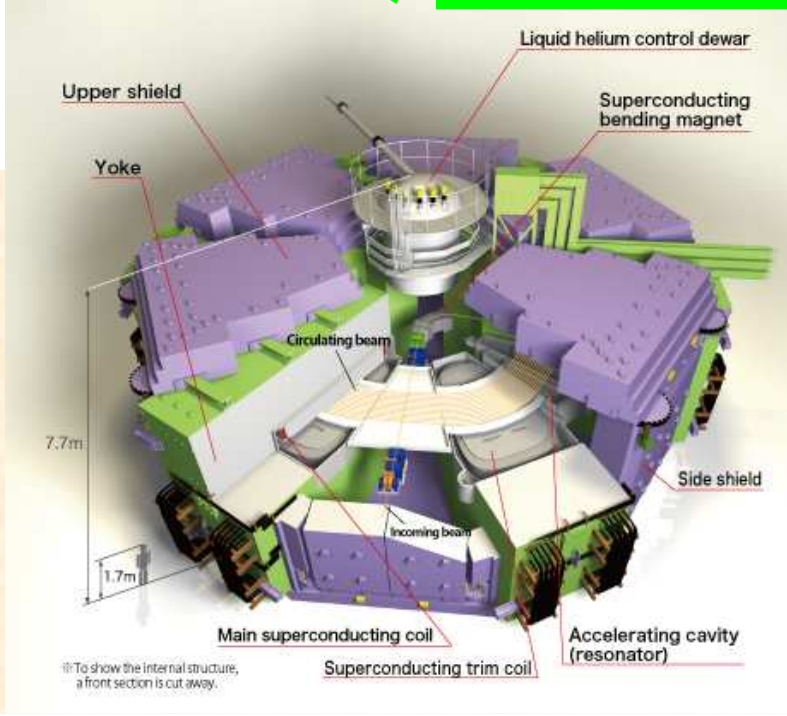
4 RF cavities

Radius=6m
=3.8T

Ion Strippers
 U^{50+}
 U^{88+}

« external »
ECR source
 U^{30+}

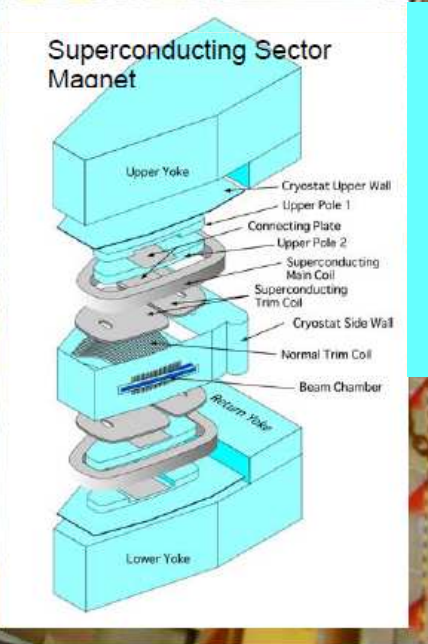
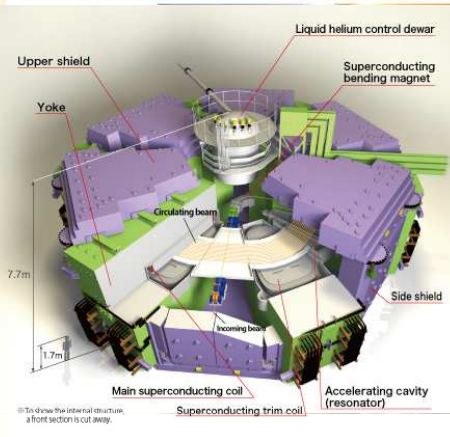
5 accelerators
1 LINAC
+ 4 Cyclotrons



RIBF (Japan) : SRC (K=2600 MeV)

the largest cyclotron in the world

Superconducting Ring Cyclo

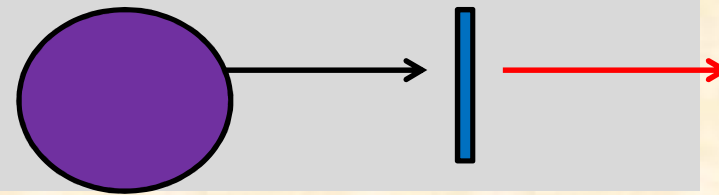


6 sectors :
Iron dominated magnets
With superconducting coils

Completed November 2005 - the 140-ton cold mass cooled to 4.5K.

TREND in Nuclear physics : Physics at RIBF (Tokyo)

Create Unknown "exotic" nuclei
by nuclear fragmentation
or by fission



Heavy ions
Cyclotrons

target

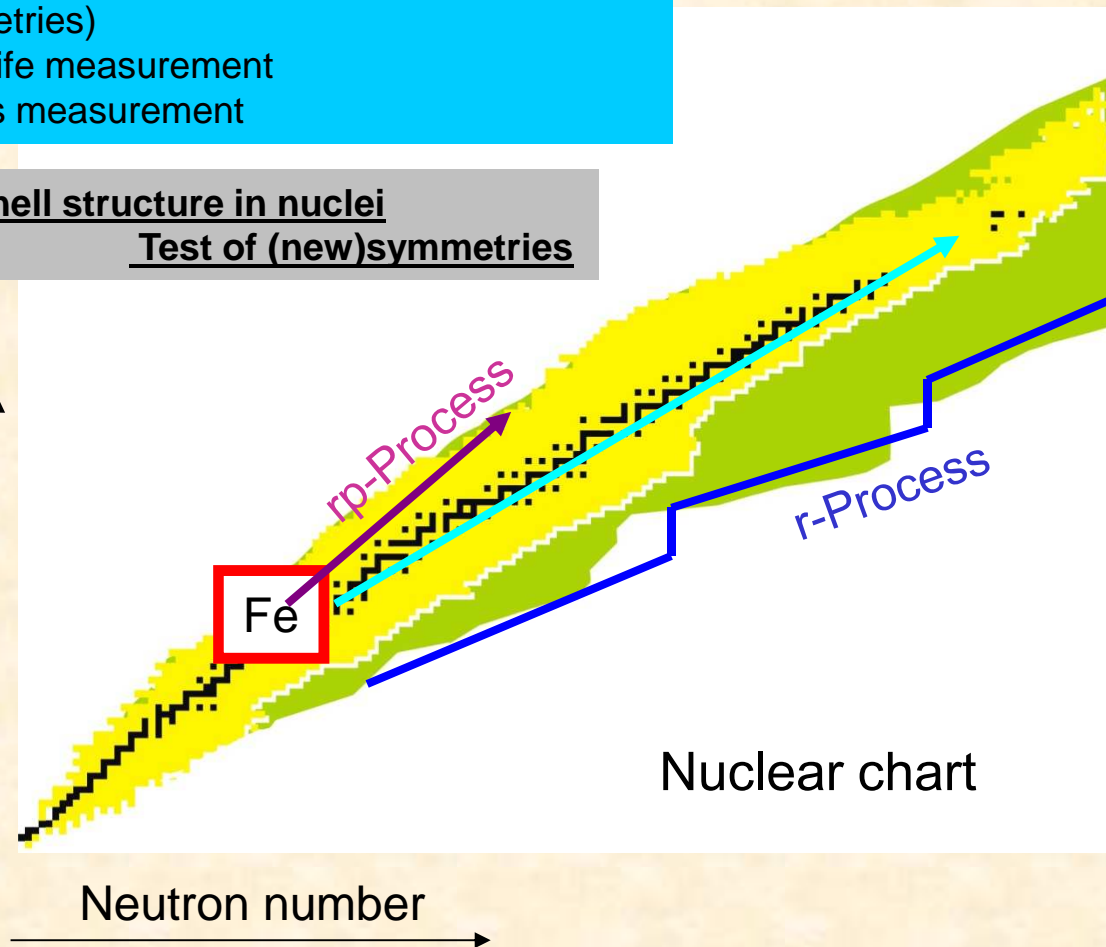
Radioactive
nuclei

Exotic Nuclei (Very large proton-neutron asymmetries)

- half life measurement
- Mass measurement

Shell structure in nuclei
Test of (new)symmetries

Proton number



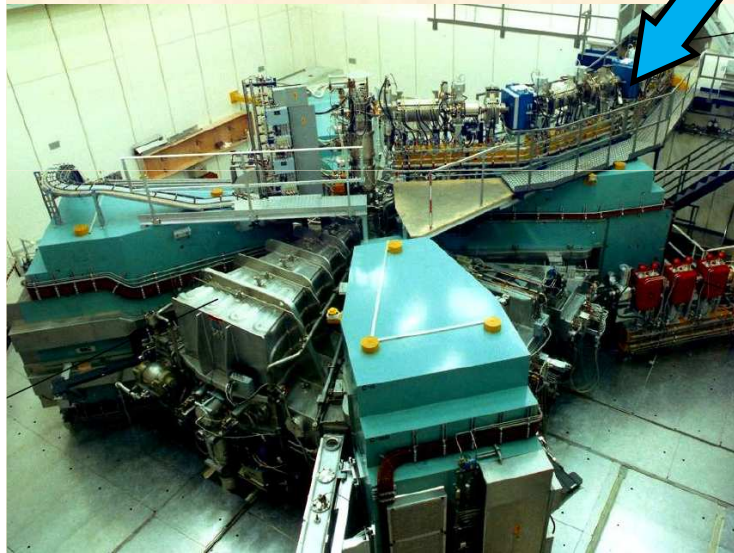
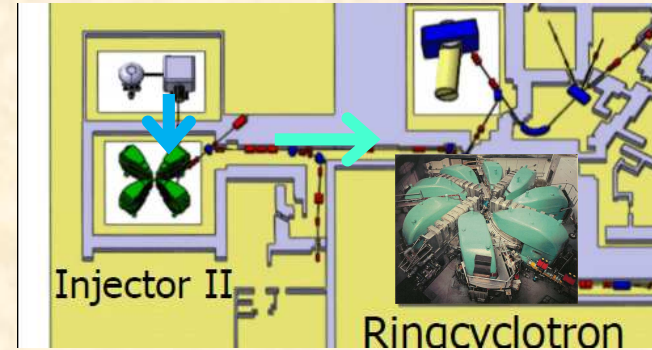
Nuclear chart

	Degrees of Freedom	Energy (MeV)
Physics of Hadrons	 Quarks, Gluons	
	 Constituent Quarks	940 Neutron Mass
Physics of Nuclei	 Baryons, Mesons	140 Pion Mass
	 Protons, Neutrons	8 Proton Separation Energy in Lead
	 Nucleonic Densities and Currents	1.32 Vibrational State in Tin
	 Collective Coordinates	0.043 Rotational State in Uranium

PSI 590Mev proton (Ch)

870 keV

Cockcroft-Walton

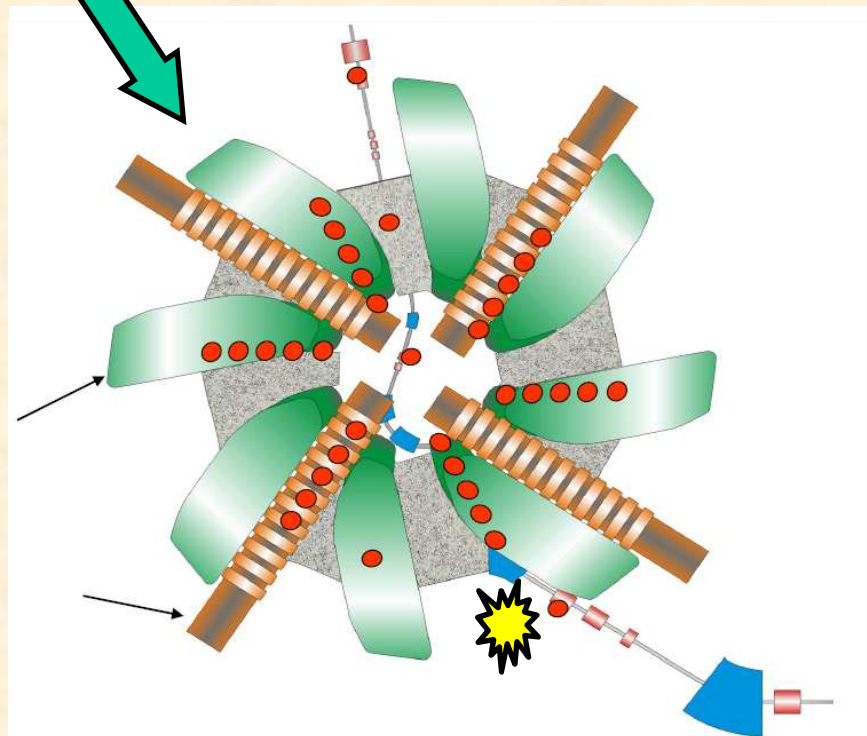


72 MeV injector 2

Coupling of 2 cyclotron :

$$v = \left[\frac{F_{HF} \cdot Re_{jec}}{h} \right]_{cycloA} = \left[\frac{F_{HF} \cdot Rinj}{h} \right]_{cycloB}$$

590 MeV Ring cyclotron

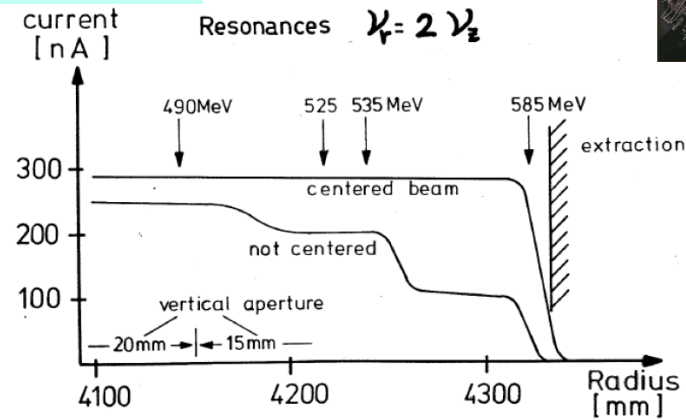


PSI :K= 590 MeV ring cyclotron



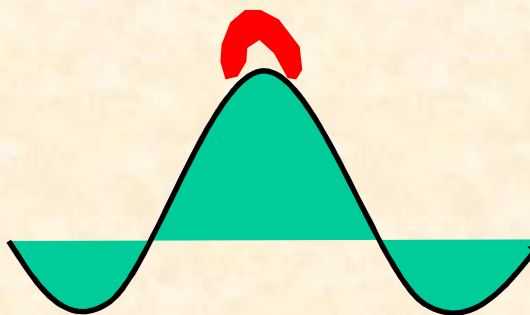
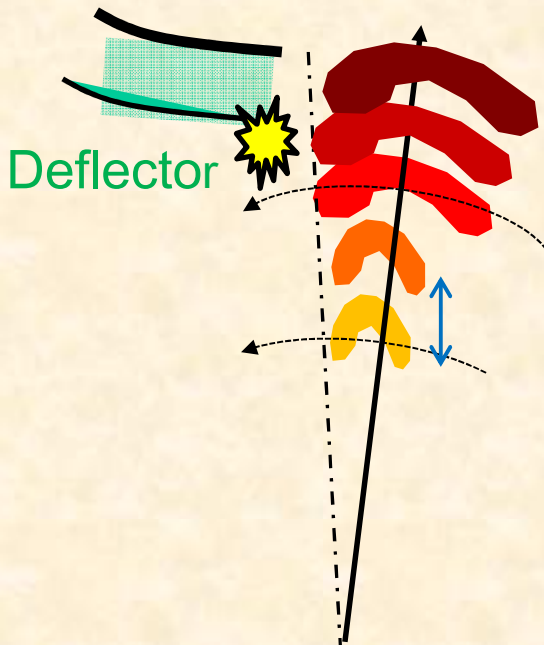
Injection centering :

Watch the resonance $2 \cdot v_z + v_r = 2$



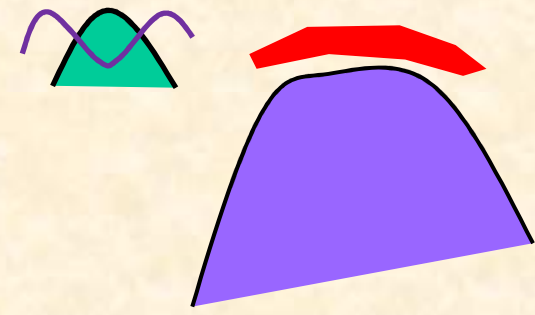
Extraction : Watch the beam losses !

A Flat top RF cavity has been added to reduce losses: $V_{RF} = \cos(\omega t) - \alpha \cdot \cos(3\omega t)$



Standard cavity :

$$V_{RF} = \cos(\omega_{RF} t)$$



Flat top cavity

$$\cos(\omega_{RF} t) - \cos(3\omega t) \quad 22$$

Some Commercial Cyclotrons : manufacturers

IBA (Belgium)

Model	Energy	Particle	Technology	Cost estimate 2008
Cyclone	5/10	proton/D		1 Meuros
Cyclone	9/18	H-/D-		1.5 Meuros
Cyclone	30	H-/D-		???
ProteusOne	250 MeV	p	synchro-cyclo superconducting	
C70		p/D or H-/D-		

Sumitomo HI (Japan)

HM-12	p
HM-18	p

>300 commercial cyclotrons in the World
 - (10-20MeV) protons
 - (230MeV) protons :develloping market

EBCO (Canada)

TR 9/ 18	H-/D-
TR 15/ 30	H-/D-

GE-Scanditronix (USA-Sweden)

MINI TRACE	9/18
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Accel // VARIAN

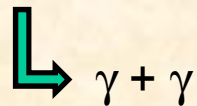
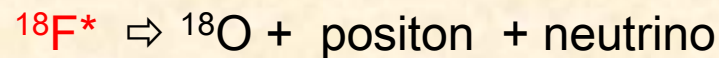
250 MeV	proton	cyclo superconducting
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Commercial Cyclotron

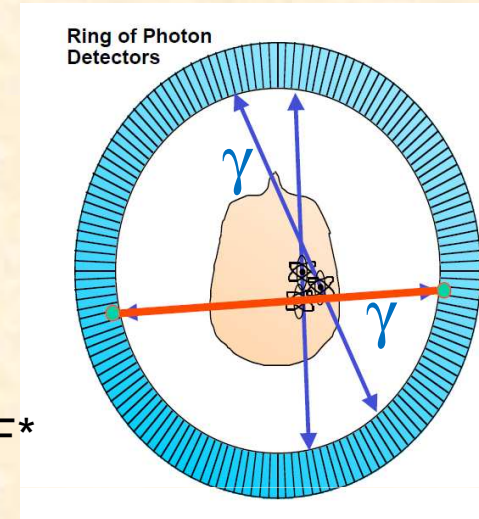
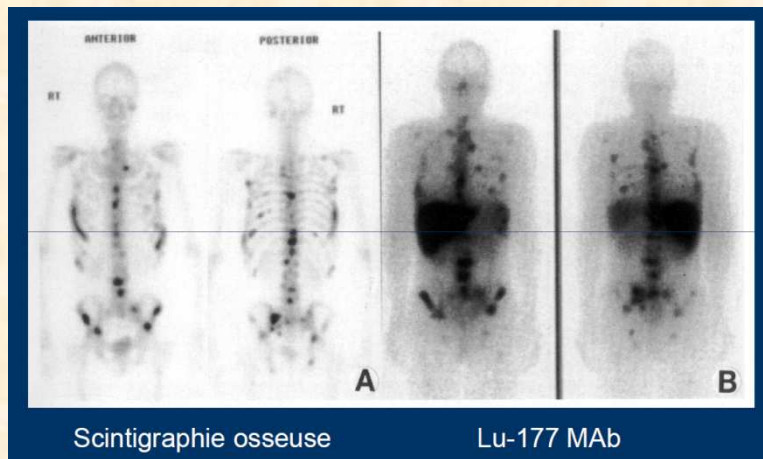
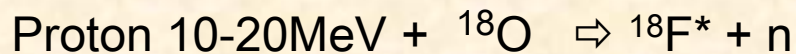
Radioisotope production (5-20 MeV)

Radiotracer $^{18}\text{F}^*$: « Beta+ » emitter
 Fluorine 18

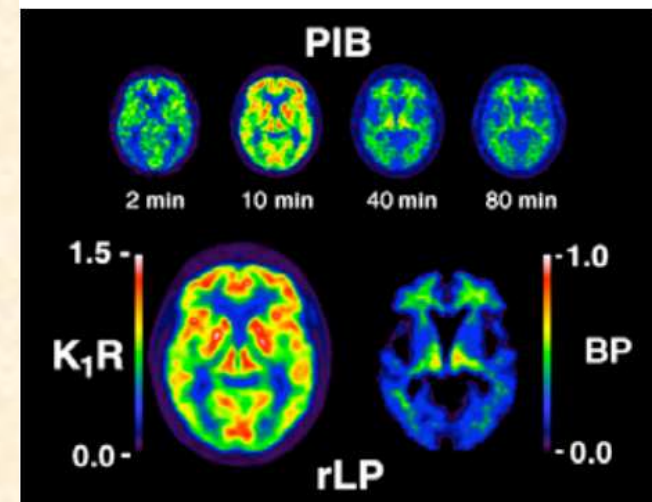
$T_{1/2} = 109.7 \text{ min}$



Production With cyclotrons



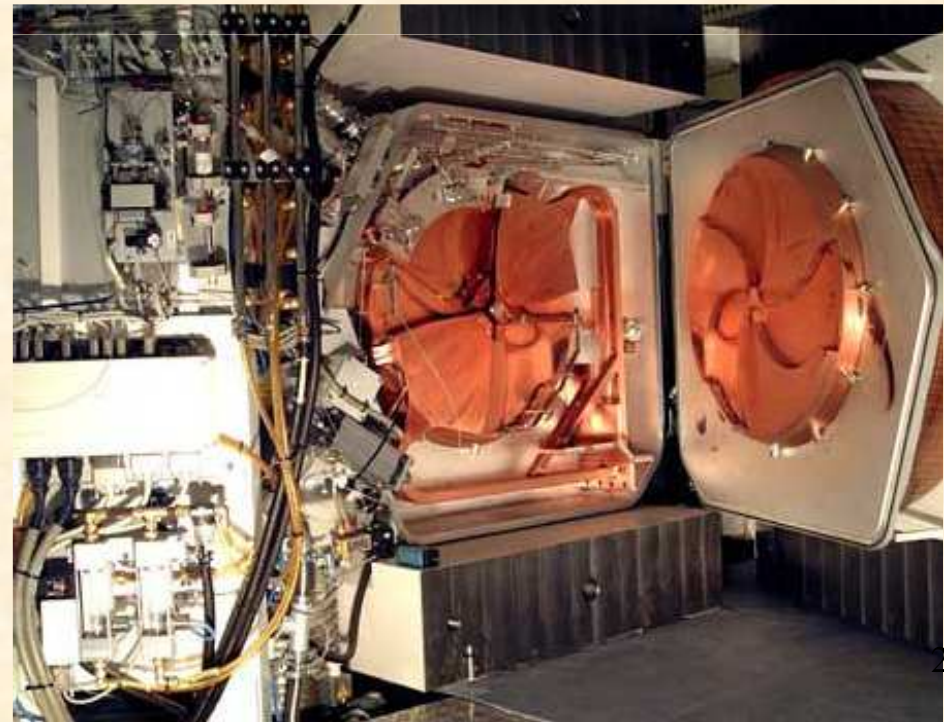
Reconstruction





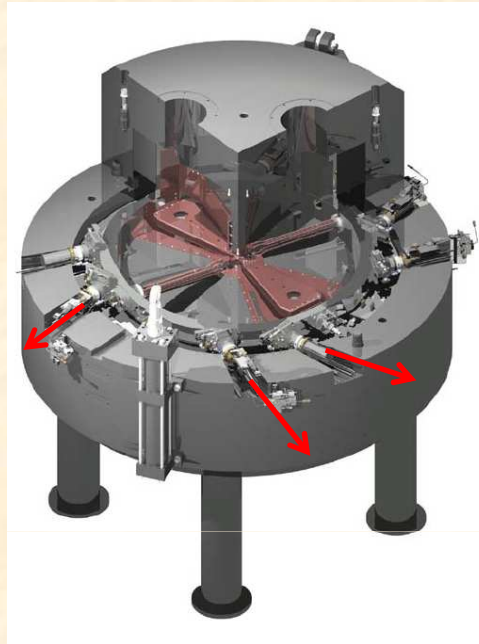
Cyclone 10/5 Mev

cyclone 3D



CYCLONE 18/9 (IBA) : H⁻ 18 MeV

Designed for medical applications (radiotracers production)



Kb=10 MeV

Fixed energy ;

4 straight sector 50°

$\langle B \rangle = 1.35$ Tesla

Hill //valley gap 3cm// 67cm

fixed Frf =42Mhz

2 Dee = 30° , 32 kV

Harmonic h=2(p) ,4 (D)

Internal source

Rextraction=0.46 m

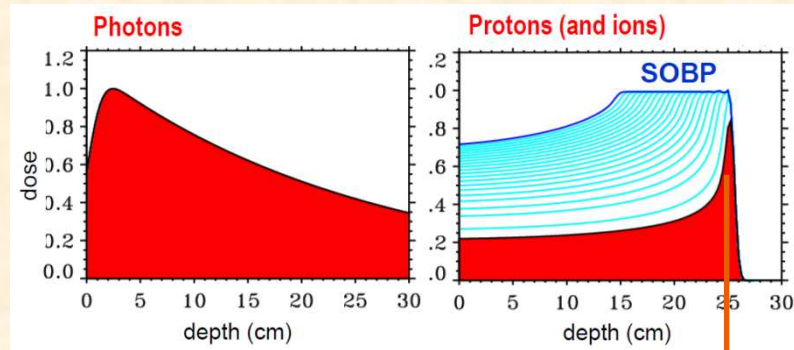
$B_{pmax} = 0.46 \times 1.35 = 0.62$ T.m

Internal PIG source, H⁻ stripping extraction



Commercial Cyclotron : proton therapy (230 MeV)

Photon :
(Radiotherapy)
A Dose in the
whole body

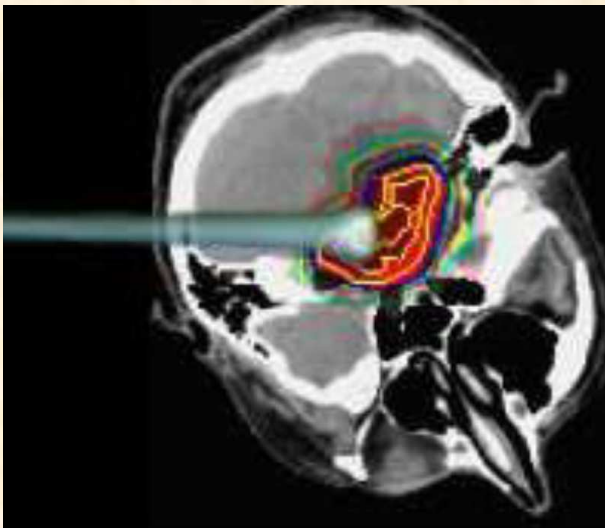


Proton :
Better than
Radiotherapy (photon)

Dose inside the tumor

Bragg Peak

Brain tumor treatment with protons



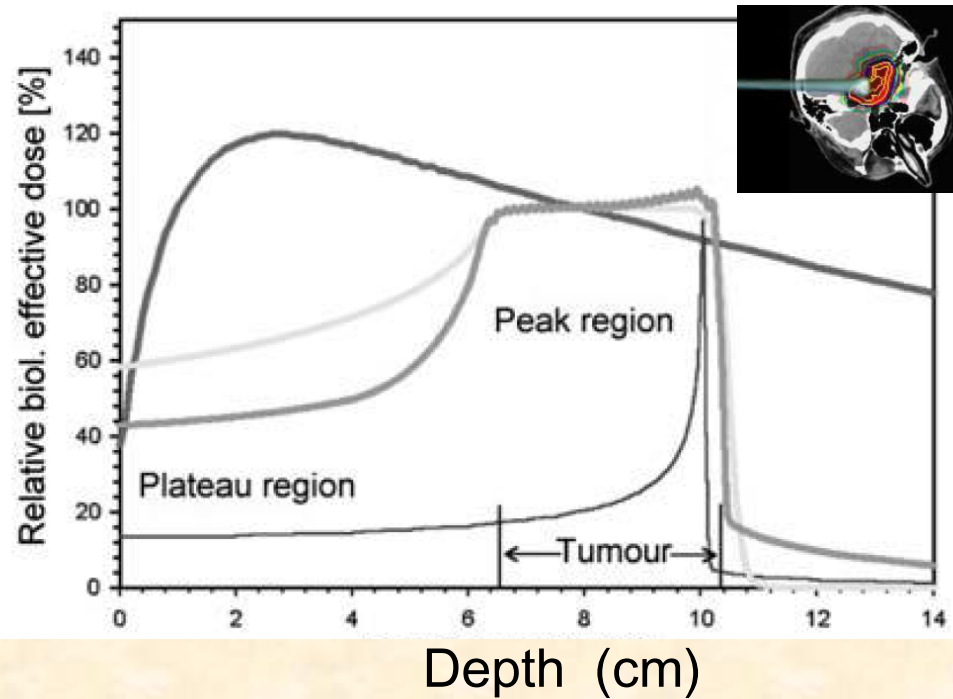
Eyes tumor treatment



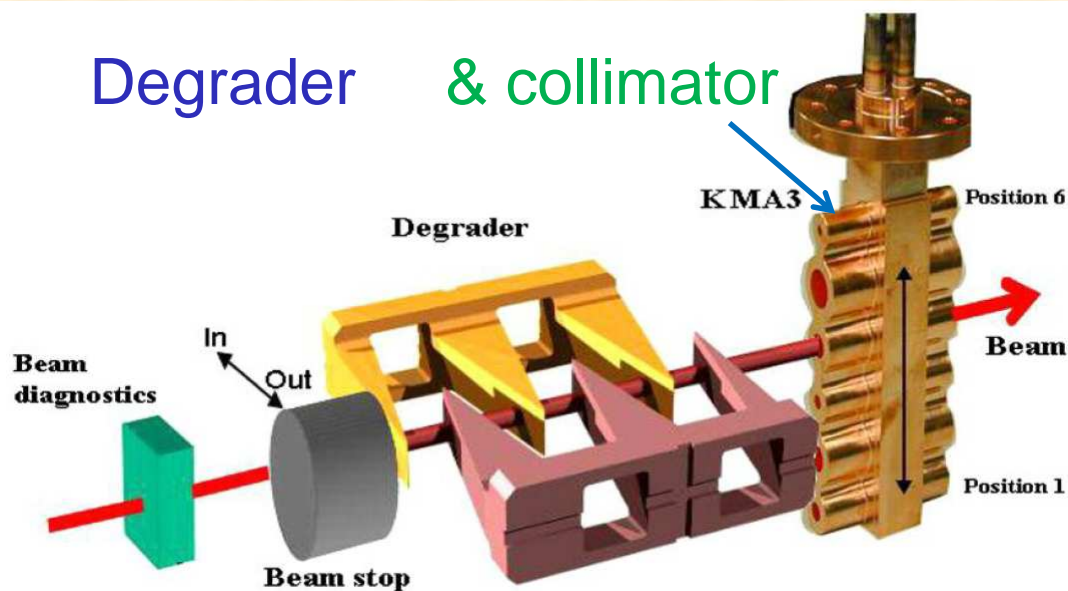
Proton therapy (230 MeV)
 -Energy variation with **degrader**

Scanning the tumors :

With energy variations
 induce by a degrader



Degrader & collimator



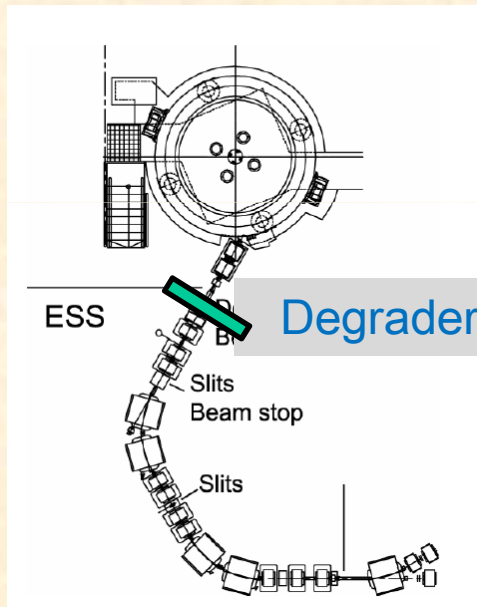
Energy variation
 =
 Range variation
 (tumour scanning)

Proton therapy (230 MeV)

Energy variation with degrader

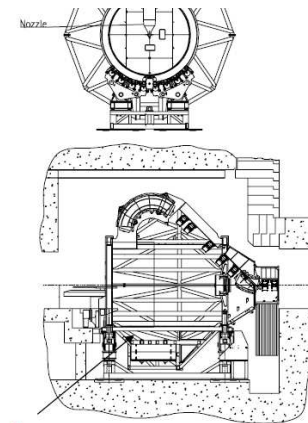
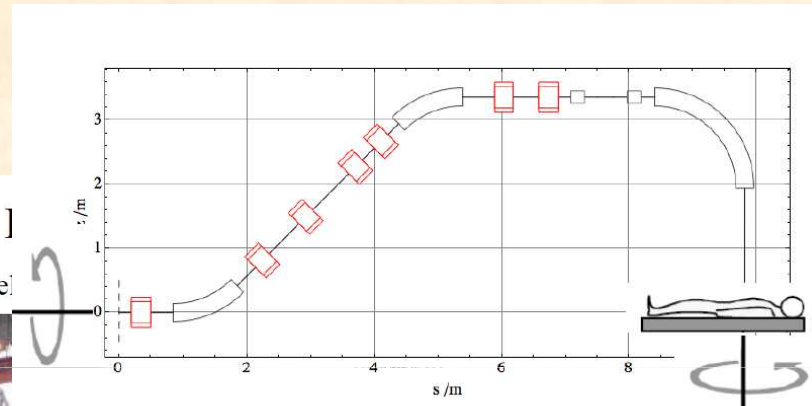
+ Rotating gantry

Optimal dose delivery
Scan the tumors :



Depth variation :
Tumor scanning

Gantry



Polyethylene plates

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The gantry reduces dose in healthy tissues

END

References & Acknowledgements :

F. Chautard, Juas 2015

M. Craddock lecture on Cyclo//FFAG

W.Joho lecture on PSI facility

S. Brandenburg lecture on beam optics

P. Heikinnen, CAS Jyväskylä 1992, CERN 94-01 (1994) “Cyclotrons“ and
“Injection and Extraction“

T. Stammach, CAS La Hulpe, 1994, CERN 96-02 (1996) “Introduction
to Cyclotrons“

The Cyclotron Family

isochronous
cyclotron
(Azimuthally
Varying Field)

$B_z(R) = \text{NOT uniform} = f(r)$

$F_{rev} = \text{Constant}$

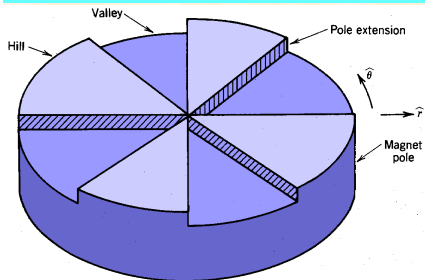
$RF = \text{constant}$

Vertical focusing with $B_z = f(R, \theta)$

Isochronous

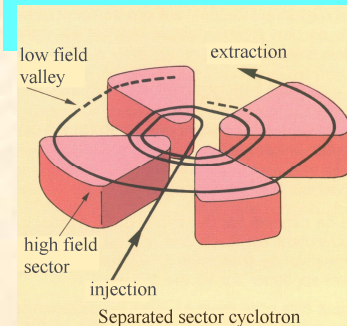
$$\omega_{rev} h = \omega_{RF}$$

Compact cyclotron (with Hills //Valleys)
Straight or Spiralled sectors



Ring cyclotron :
Straight or Spiralled

Separated sectors



Synchrocyclotron

$F_{rev} = \text{NOT Constant}$

$RF = \text{NOT Constant} = \text{beam pulsed}$

$$\omega_{rev}(R) / h = \omega_{RF}(t)$$

Not
Isochronous

Cyclotron vs other RF accelerators

	Radius	F_{rf}	Particles
Cyclotrons	not constant	constant (isochronous)	Ions $\gamma < 2$
Synchro-cyclotrons	not constant	not constant pulsed	Ions
Synchrotrons	constant	not constant pulsed	Ions, electron
Linacs	constant	constant	Ions, electron

Cyclotron Summary

Isochronous cyclotron =
constant revolution frequency
from injection to extraction

$$\omega_{rev} = \frac{qB_z(R)}{\gamma(R) m} = const$$

field index $n < 0$

$$n = - \frac{R}{B_{0z}} \frac{\partial B_z}{\partial x}$$

Vertical stability in isochronous cyclotron
requires Azimuthal field Modulation (N sectors)

$$\nu_z^2 = n + \frac{N^2}{N^2 - 1} F_l (1 + 2 \tan^2 \epsilon) > 0$$



$$z(t) \sim z_0 \exp(-i \nu_z \omega_0 t)$$

: ν_z real for stability