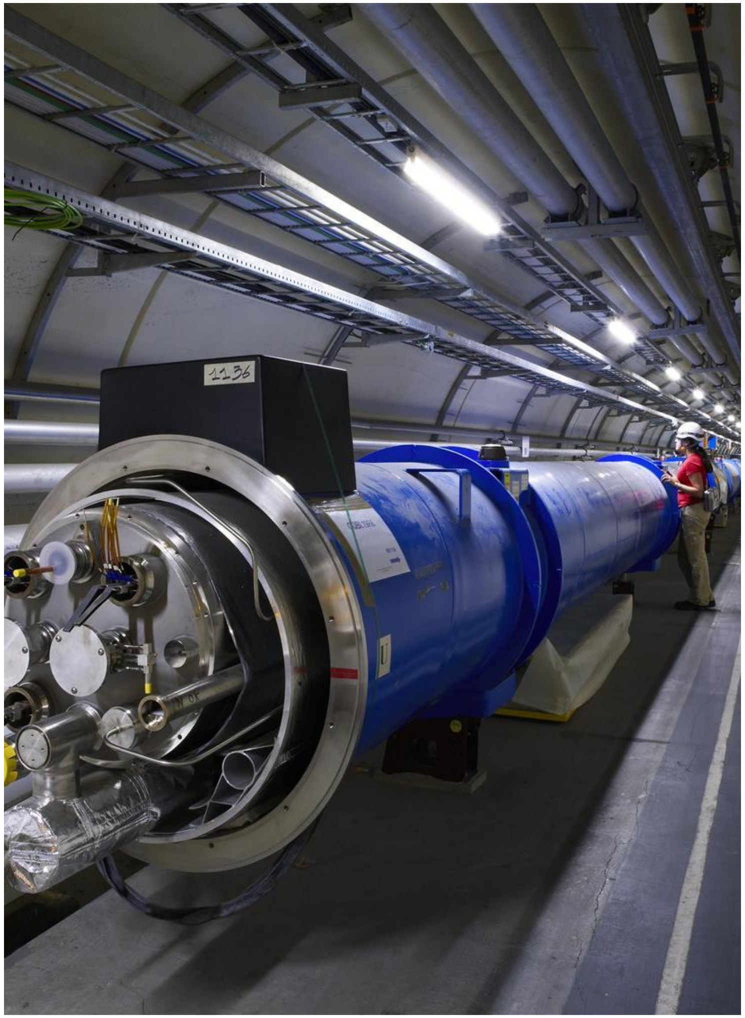
Superconductivity for accelerators – why bother?

Abolish Ohm's Law

- no power consumption
  (although do need refrigeration power)
- high current density ⇒ compact windings, high gradients
- ampere turns are cheap, so don’t need iron ⇒ high fields
  (although often use it for shielding)



Consequences

- lower power bills
- higher magnetic fields mean reduced bend radius
  ⇒ smaller rings
  ⇒ reduced capital cost
  ⇒ new technical possibilities (eg muon collider)
- higher quadrupole gradients
  ⇒ higher luminosity (eg CERN HiLumi)

Superconducting Magnets: plan

1 Introduction to Superconductors

- critical field, temperature & current
- superconductors for magnets
- manufacture of superconducting wires
- high temperature superconductors HTS
- where to find more

2 Magnetization, Cables & AC losses

- superconductors in changing fields, critical state model
- filamentary superconductors & magnetization
- coupling between filaments
  ⇒ magnetization
- flux jumping
- why cables, coupling in cables
- mini tutorial
- AC losses in changing fields

3 Magnetic and mechanical design

Paolo Ferracin

4 Magnet fabrication and assembly

Paolo Ferracin

5 Cryogenics and cryostat design

Philippe Lebrun

6 Quench, quench protection & training

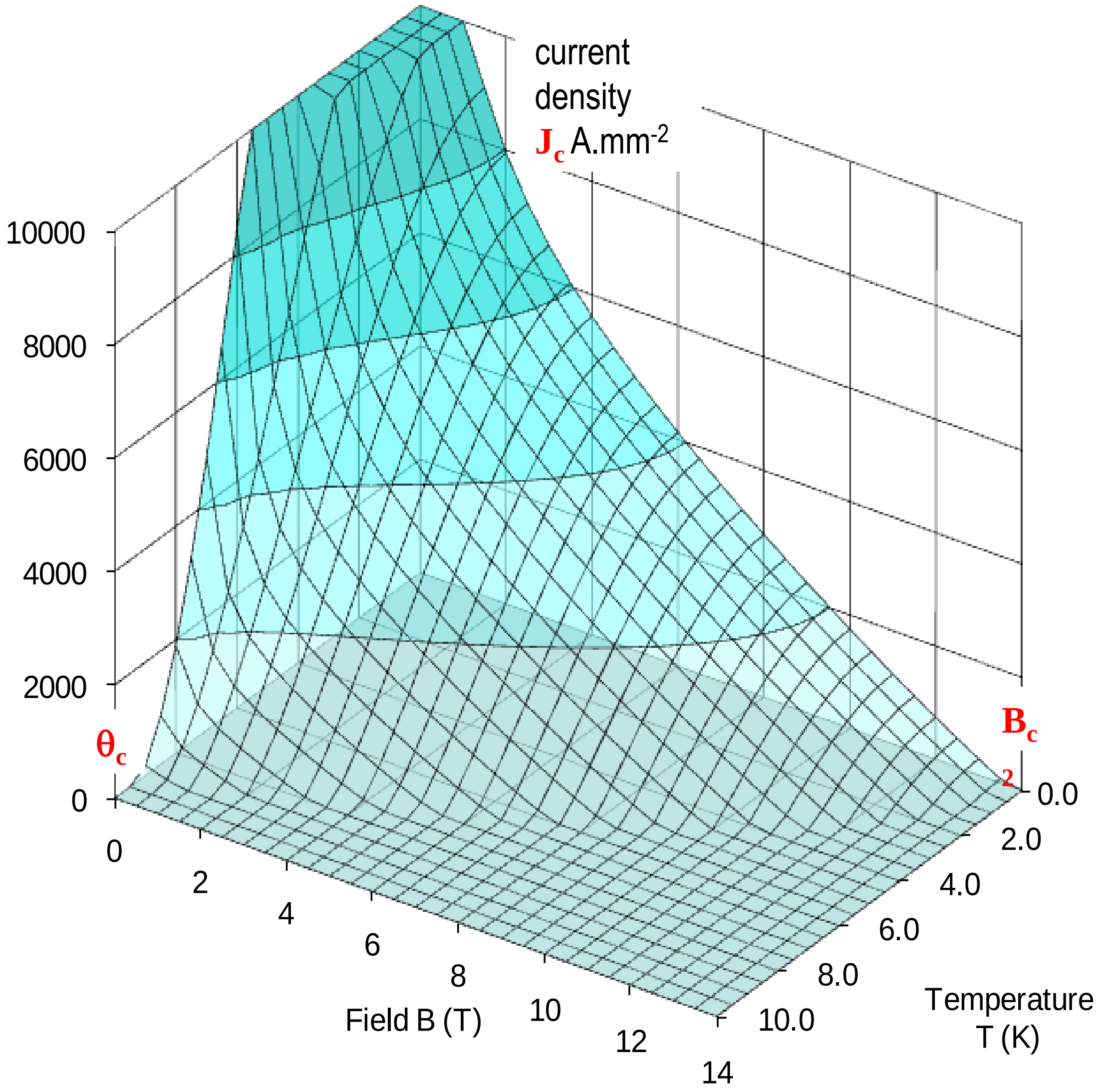
Paolo Ferracin

Mini workshop Superconducting Magnets

Daniel Schoerling Paolo Ferracin
Martin Wilson

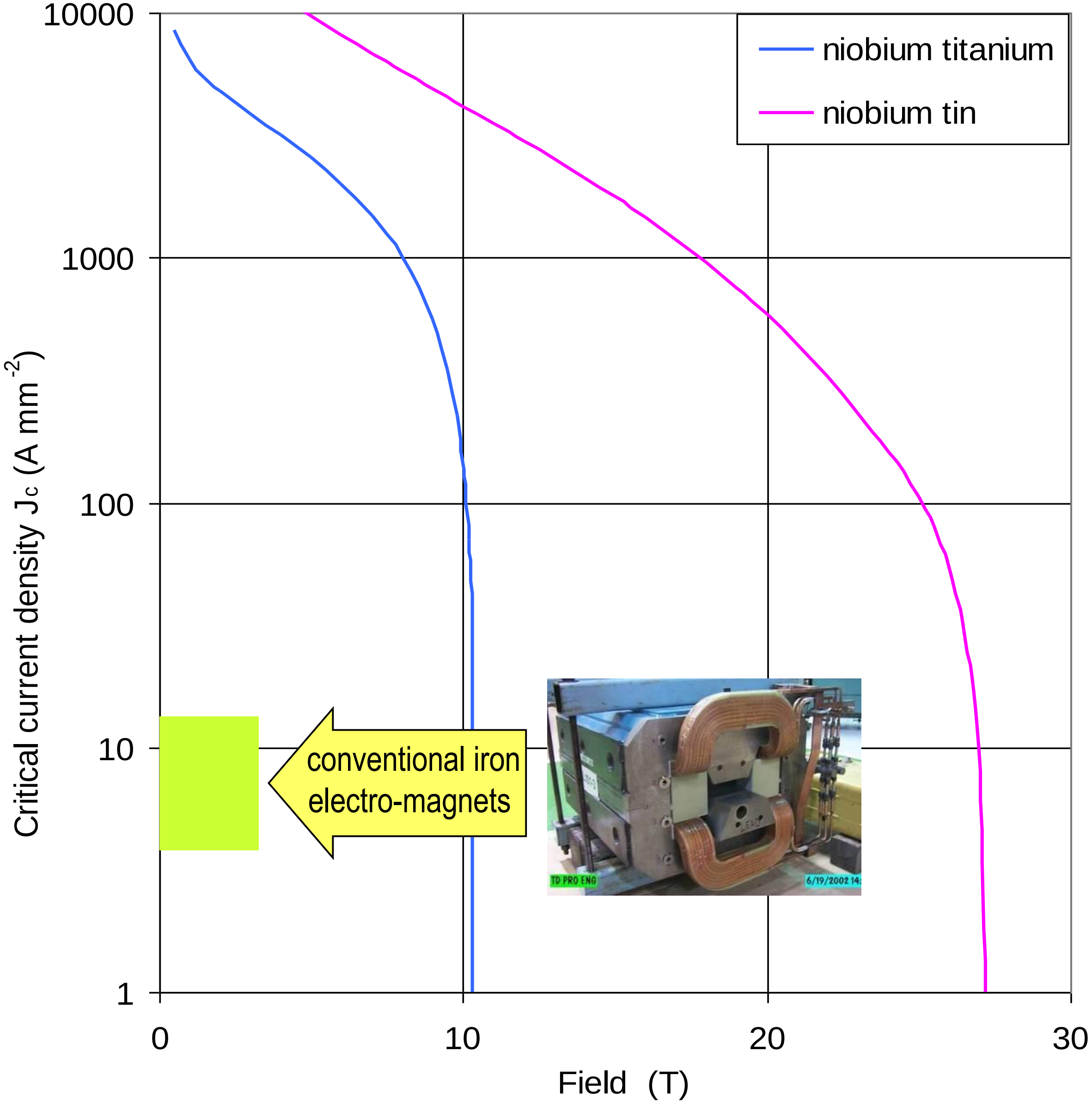
The critical surface of niobium titanium

- **NbTi** is the standard commercial ‘work horse’ of the superconducting magnet business



- ***critical surface*** is the boundary between superconductivity and normal resistivity in ***J, B, θ*** space
- superconductivity prevails everywhere below the surface, resistance everywhere above it
- upper critical field B_{c2} (at zero temperature and current)
- critical temperature θ_c (at zero field and current)
- B_{c2} and θ_c are characteristic of the alloy composition
- critical current density J_c depends on processing
- keep it cold!

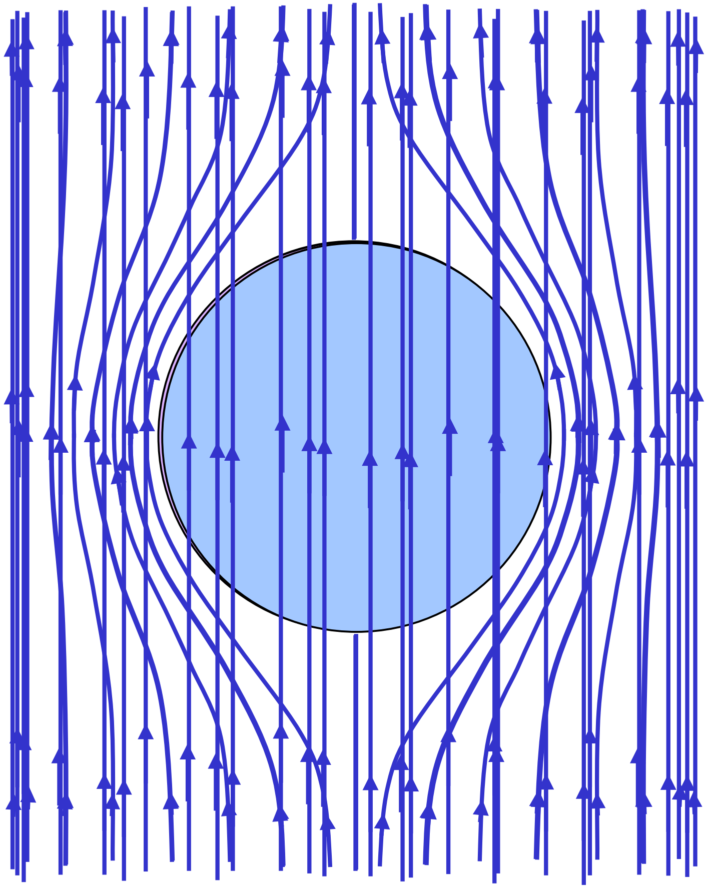
The critical line at 4.2K



- magnets usually work in boiling liquid helium, so the critical surface is often represented by a curve of current versus field at 4.2K
- niobium tin Nb_3Sn has a much higher performance than NbTi
- **<u>but</u>** Nb_3Sn is a brittle intermetallic compound with poor mechanical properties
- both the field and current density of both superconductors are way above the capability of conventional electromagnets

Two kinds of superconductor: type 1

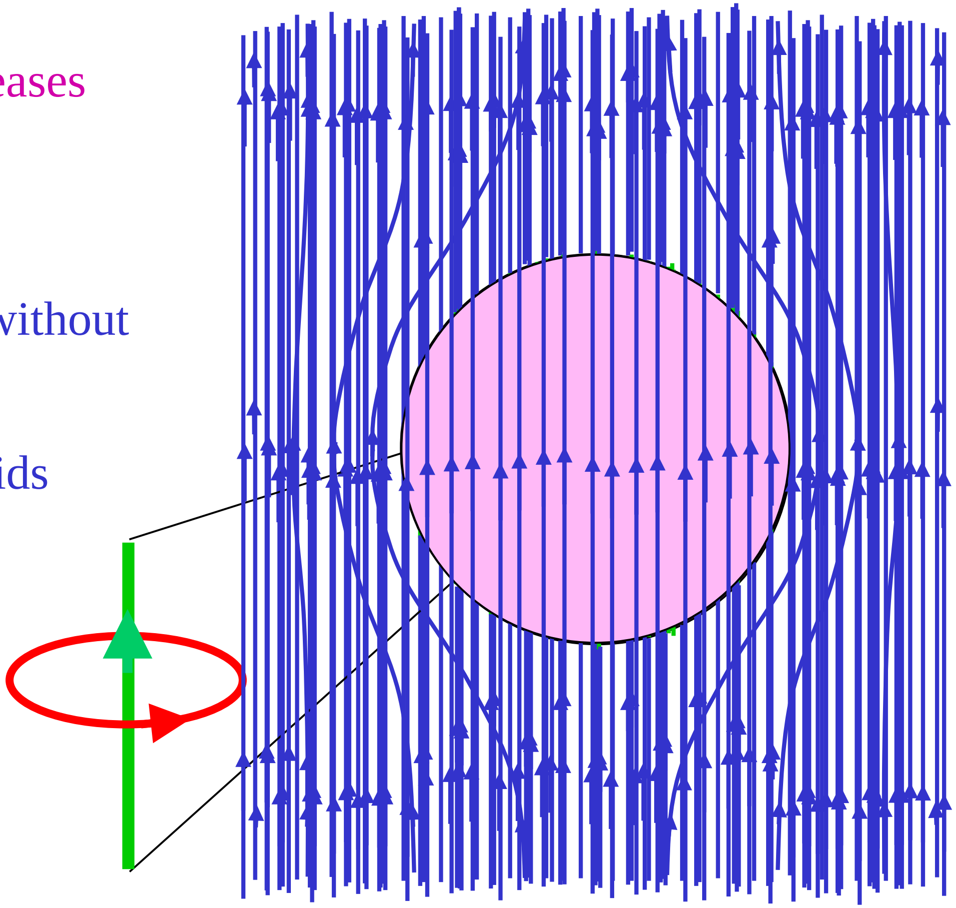
- the materials first discovered by Kammerlingh Onnes in 1911 - soft metals like lead, tin mercury
- sphere of metal at room temperature
- apply magnetic field
- reduce the temperature - resistance decreases
- reduce the temperature some more - resistance decreases some more
- at the critical temperature θ_c the field is pushed out - the **Meissner effect** - superconductivity!
- increase the field - field is kept out
- increase the field some more - superconductivity is extinguished and the field jumps in
- thermodynamic critical field B_c is trade off between **reducing** energy via condensation to superconductivity and **increasing** energy by pushing out field ~ 0.1T



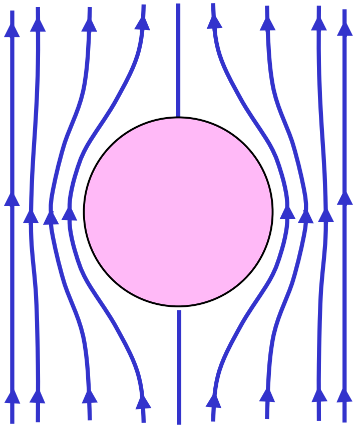
useless for magnets!

Two kinds of superconductor: type 2

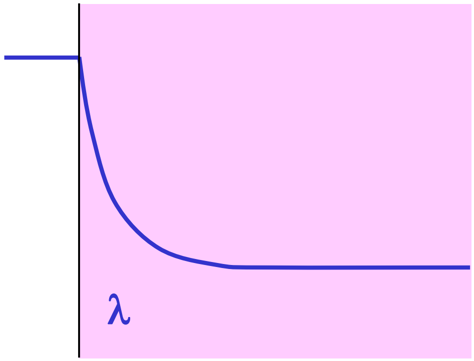
- apply magnetic field
- reduce the temperature - resistance decreases
- at the critical temperature θ_c the field is pushed out
- increase the field - field jumps back in without quenching superconductivity
- it does so in the form of quantized fluxoids
- lower critical field B_{c1}
- supercurrents encircle the resistive core of the fluxoid thereby screening field from the bulk material
- higher field ⇒ closer vortex spacing
- superconductivity is extinguished at the (much higher) upper critical field B_{c2}



Type 1 and type 2 superconductors

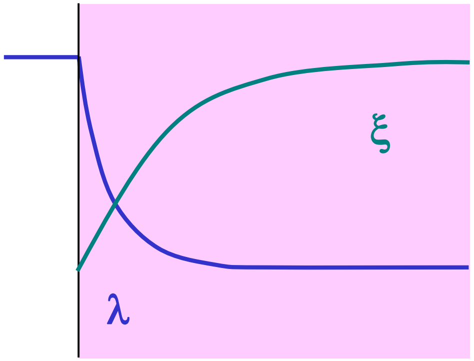


- Meissner effect is not total
- magnetic field penetrates a small distance λ
- the ***London Penetration Depth.***

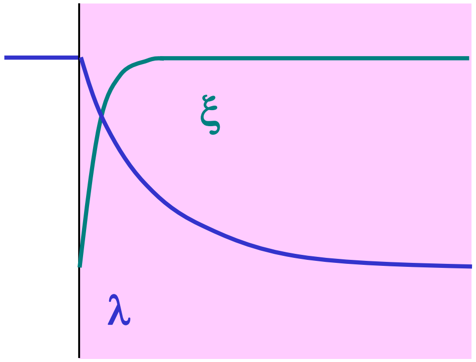


- another characteristic distance is the ***coherence length*** ζ - the minimum distance over which the electronic state can change from superconducting to normal
- theory of Ginsburg, Landau, Abrikosov and Gorkov **GLAG** defines the ratio $\kappa = \lambda / \xi$

if $\kappa < 1/\sqrt{2}$
material is
Type 1



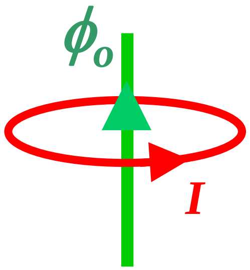
if $k > 1/\sqrt{2}$
material is
Type 2



Critical fields of type 2 superconductors

- recap thermodynamic critical field B_c
- lower critical field $B_{c1} = B_c / \kappa$
- above B_{c1} magnetic field penetrates as discrete quantized ***fluxoids***

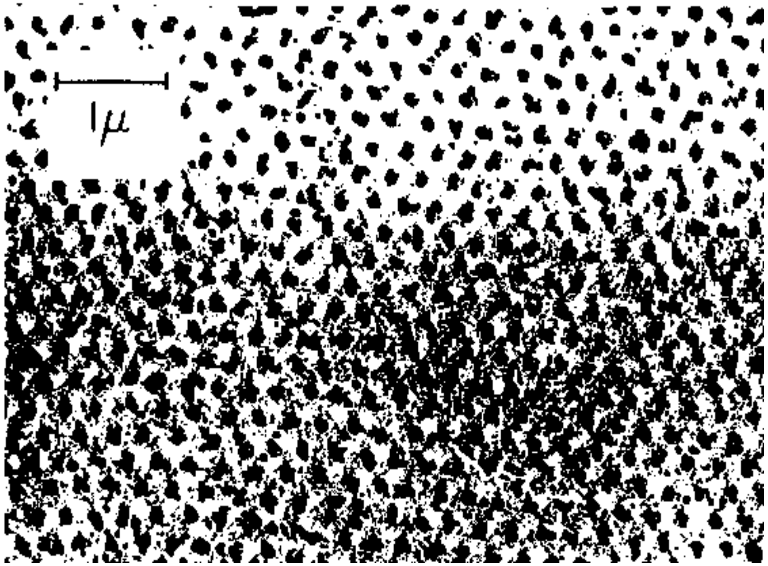
human hair in earth's magnetic field ~ 50ϕ_o



a fluxoid encloses flux

$$\phi_o = \frac{h}{2e} = 2\times10^{-15} Wb$$

h = Planck's constant
e = electronic charge



upper critical field $B_{c2} = \sqrt{2}\kappa B_c$

in the 'dirty limit' $\kappa \approx 2.4\times10^6 \gamma^{\frac{1}{2}} \rho_n$

where ρ_n is the normal state resistivity

- best superconductors are best resistors!

thus the upper critical field $B_{c2} = 3.1\times10^3 \gamma\rho_n\theta_c$

for NbTi: γ ~ 900 J m $^{-3}$ K^{-2} ρ_n ~ 65 x10^{-8} W m θ_c = 9.3 K hence **B_{c2} ~ 18.5 T**

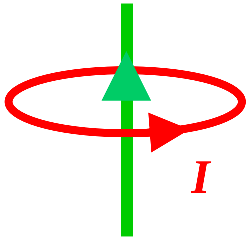
Sommerfeld coefficient of electronic specific heat $C_e = \gamma\theta$

Critical current density: type 2 superconductors

- a single fluxoid encloses flux

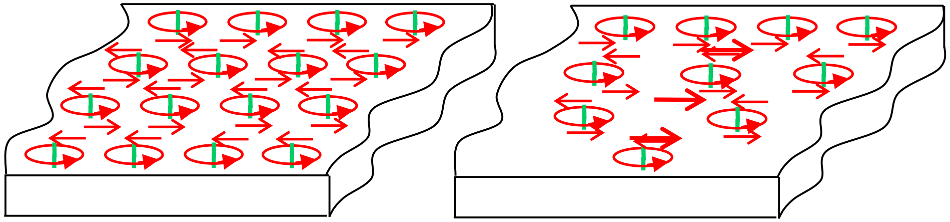
$$\phi_o = \frac{h}{2e} = 2\times10^{-15}\,Webers$$

h = Planck's constant, e = electronic charge



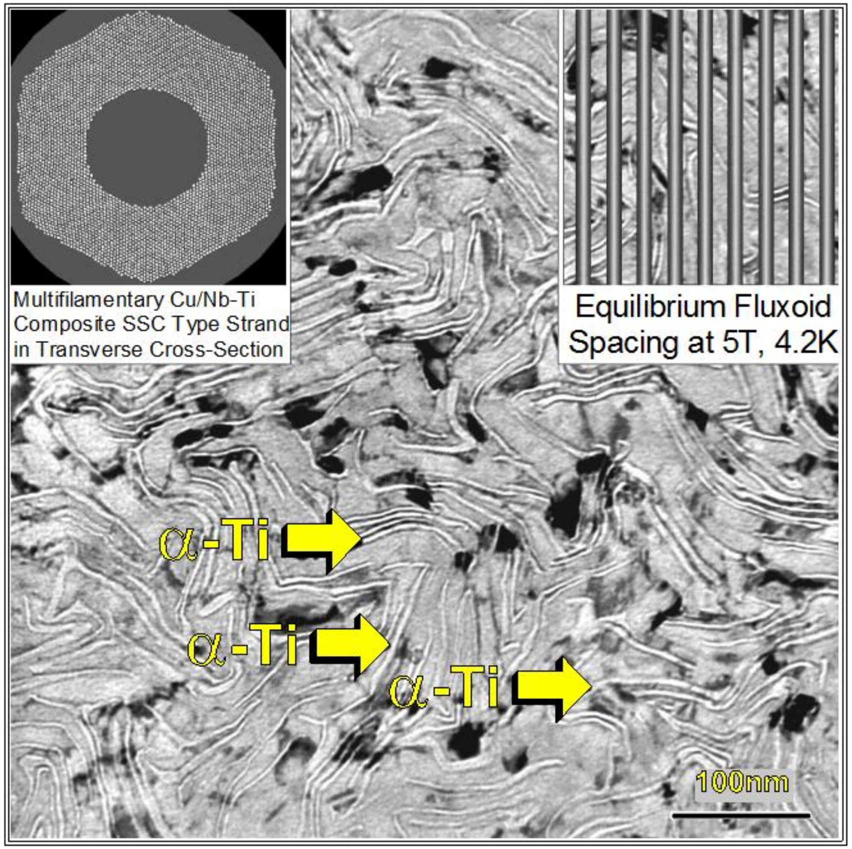
- so density of fluxoids ~ average field
  uniform density ⇒ uniform field
  ⇒ zero J (because $Curl\ B = \mu_o J$)

- to get a current density we must produce a **gradient** in the density of fluxoids



- fluxoids like to distribute uniformly

- must impose a gradient by inhomogeneities in the material, eg dislocations or precipitates



precipitates of α Ti in Nb Ti

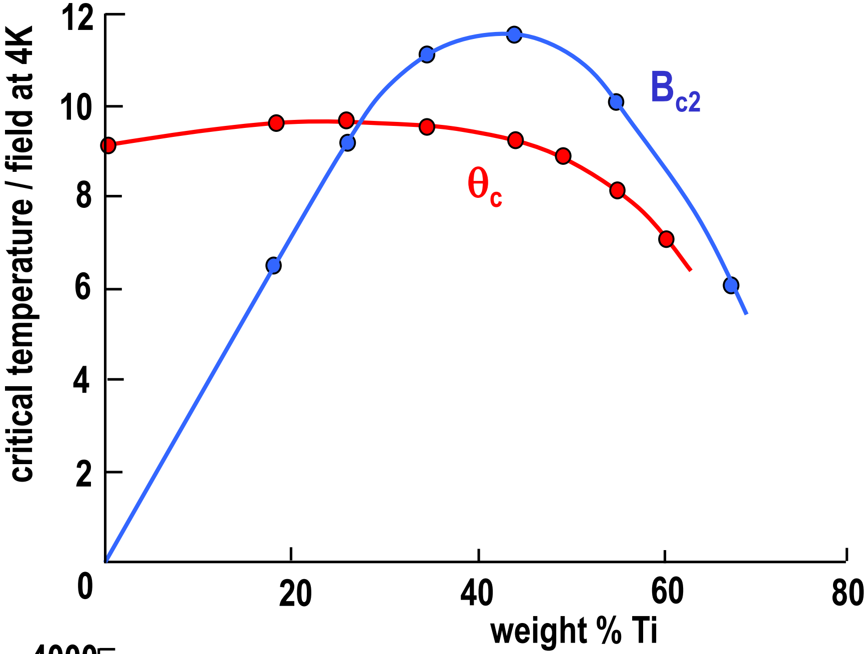
Meingast, P Lee and DC Larbalestier: J. Appl. Phys. 66, 5971

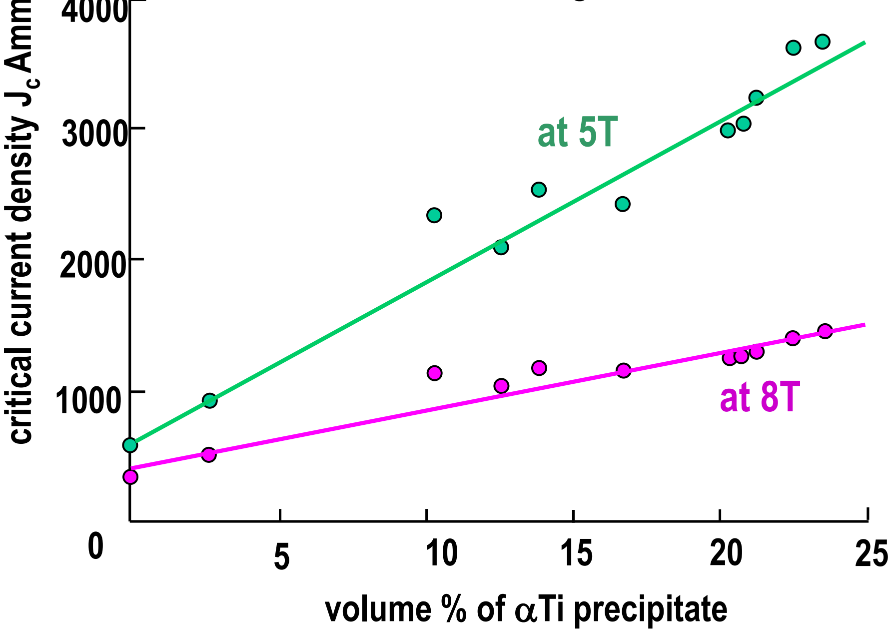
Critical properties

- **Critical temperature θ_c**: choose the right material to have a large energy gap or 'depairing energy'
  property of the material

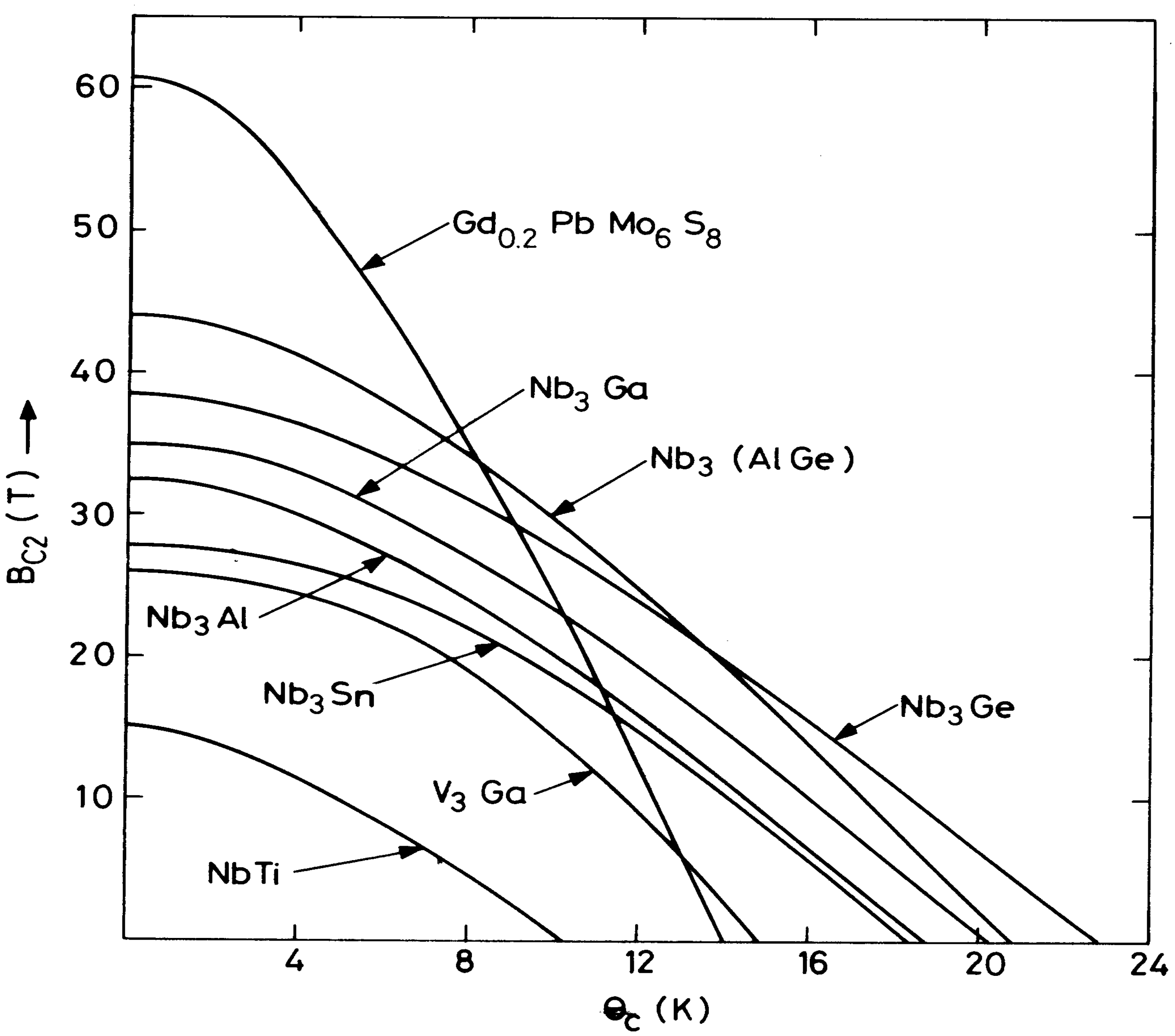
- **Upper Critical field B_{c2}:** choose a Type 2 superconductor with a high critical temperature and a high normal state resistivity
  property of the material

- **Critical current density J_c**: mess up the microstructure by cold working and precipitation heat treatments
  hard work by the producer





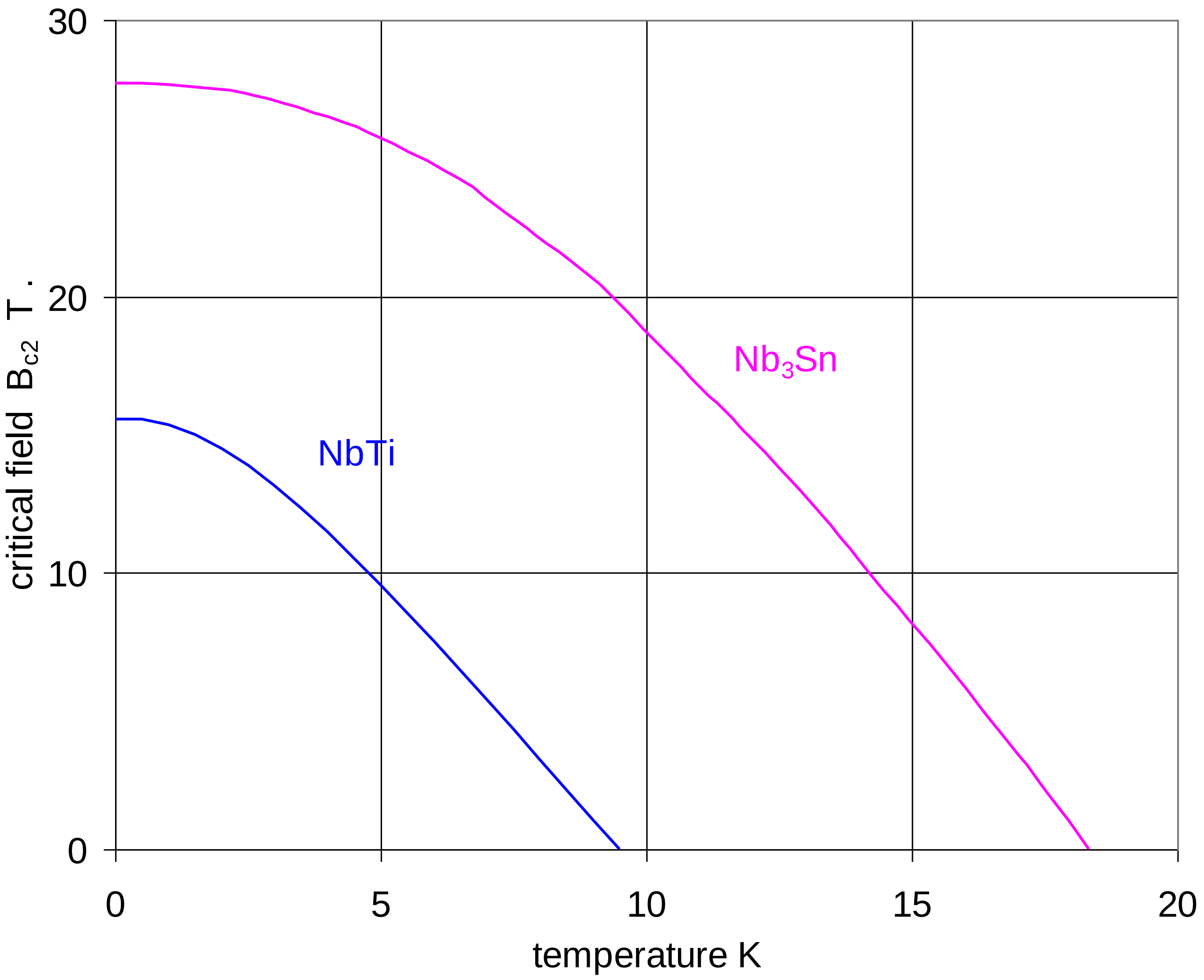
Critical field & temperature of metallic superconductors



Note: of all the metallic superconductors, only NbTi is ductile.

All the rest are brittle intermetallic compounds

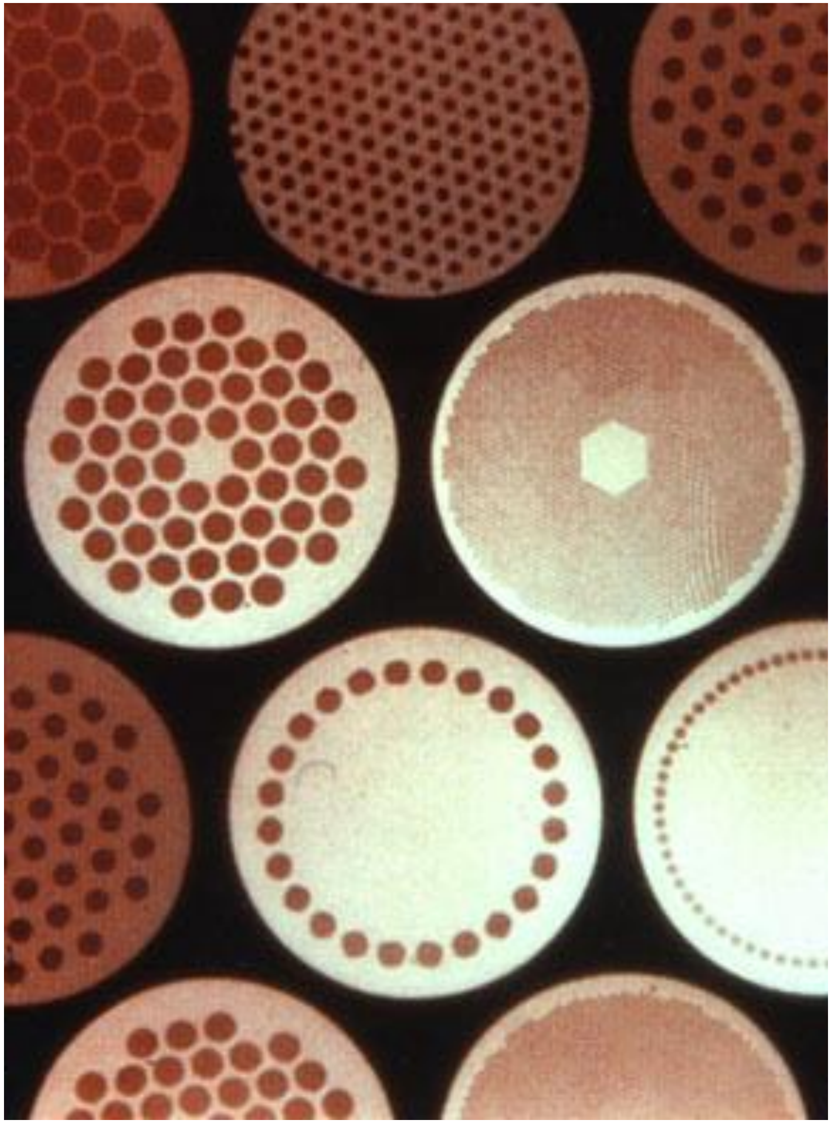
Critical field & temperature of metallic superconductors



To date, all superconducting accelerators have used NbTi.

Of the intermetallics, only Nb_3Sn has found significant use in magnets

Practical wires for magnets



- some 40 years after its development, NbTi is still the most popular magnet conductor, with Nb_3Sn being used for special high field magnets and HTS for some developmental prototypes.

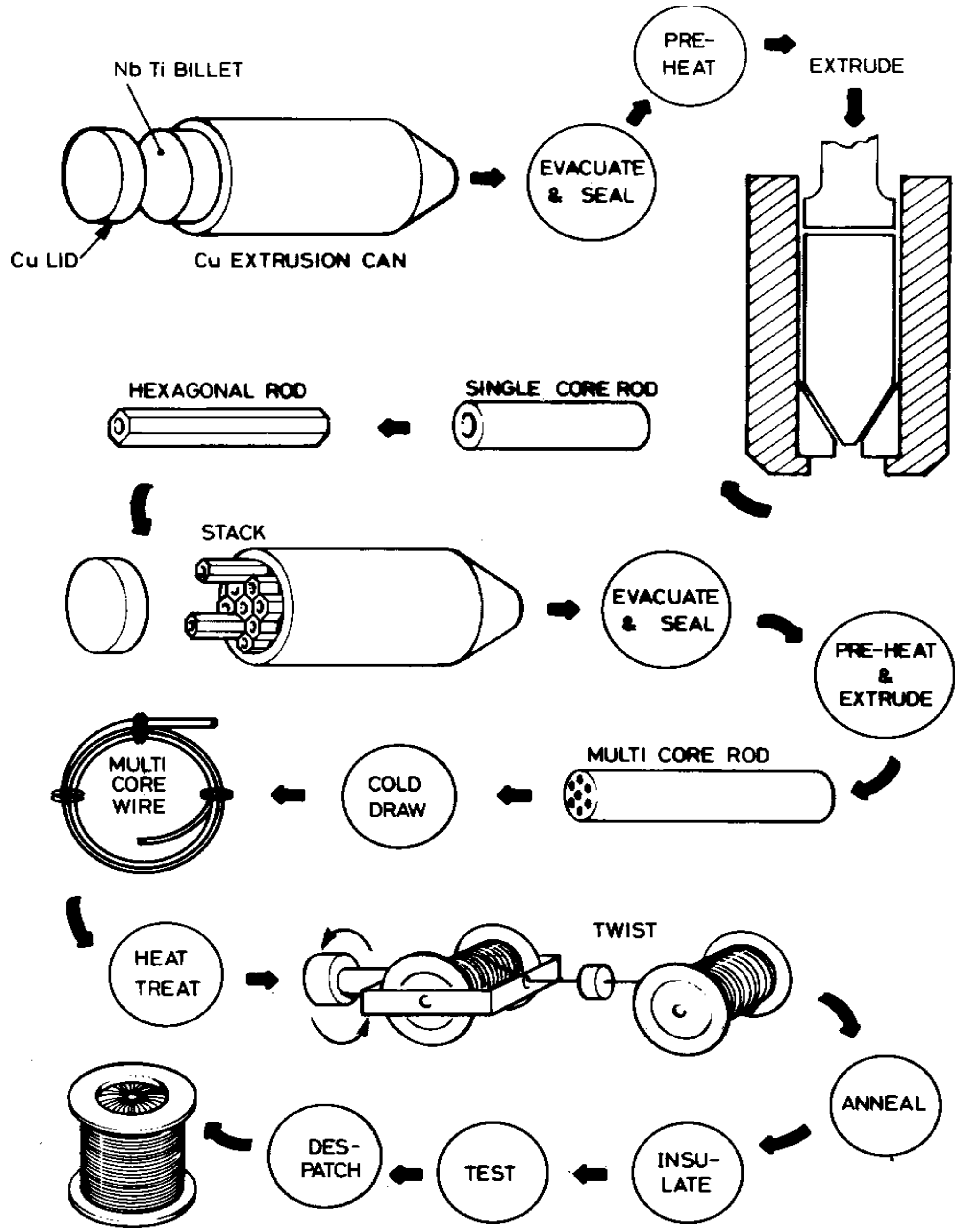
- for reasons that will be described later, superconducting materials are always used in combination with a good normal conductor such as copper

- to ensure intimate mixing between the two, the superconductor is made in the form of fine filaments embedded in a matrix of copper

- typical dimensions are:
  - wire diameter = 0.3 - 1.0mm
  - filament diameter = 5 - 50μm

- for electromagnetic reasons, the composite wires are twisted so that the filaments look like a rope (see Lecture 2)

NbTi manufacture

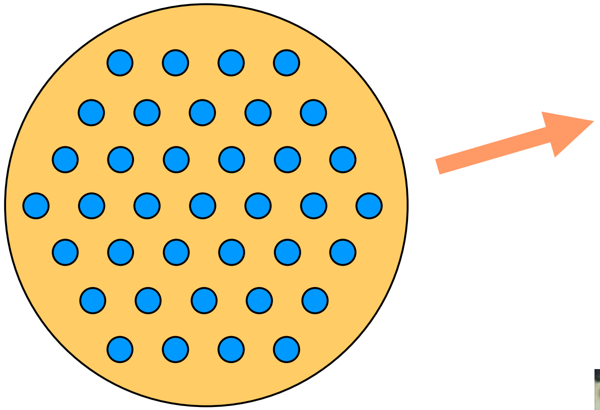


- vacuum melting of NbTi billets
- hot extrusion of the copper NbTi composite
- sequence of cold drawing and intermediate heat treatments to precipitate αTi phases as flux pinning centres
- for very fine filaments, must avoid the formation of brittle CuTi intermetallic compounds during heat treatment
  - usually done by enclosing the NbTi in a thin Nb shell
- twisting to avoid coupling - see lecture 2

Filamentary Nb_3Sn wire via the bronze route

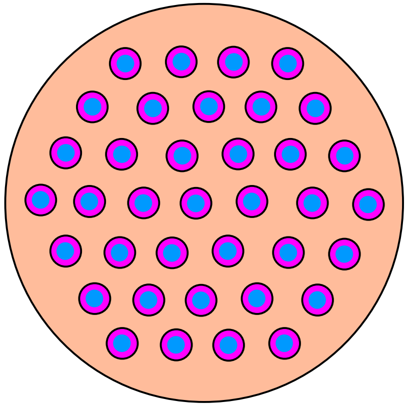
Nb_3Sn is a brittle material and cannot be drawn down.

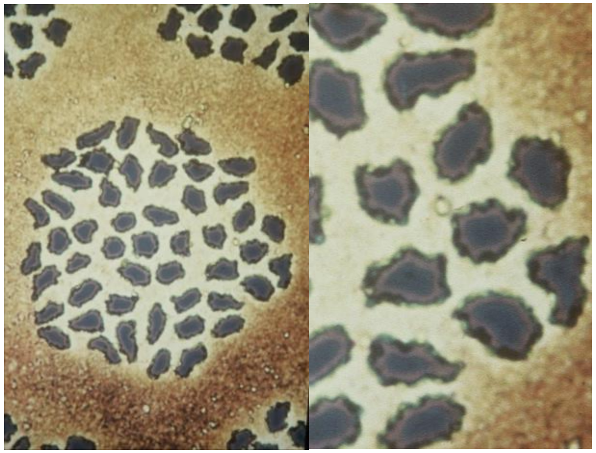
Instead must draw down pure niobium in a matrix of bronze (copper tin)



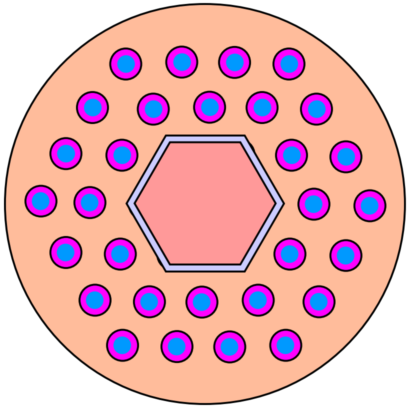
- *BUT maximum ductile bronze is ~13wt% tin,*
- *reaction slows at ~ 3wt%*
- *so low engineering J_c*

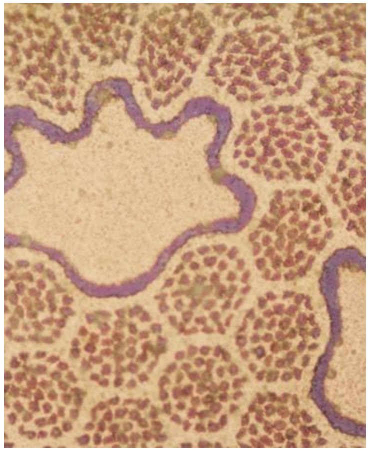
At final size the wire is heated (~700C for some days) tin diffuses through the Cu and reacts with the Nb to form Nb_3Sn



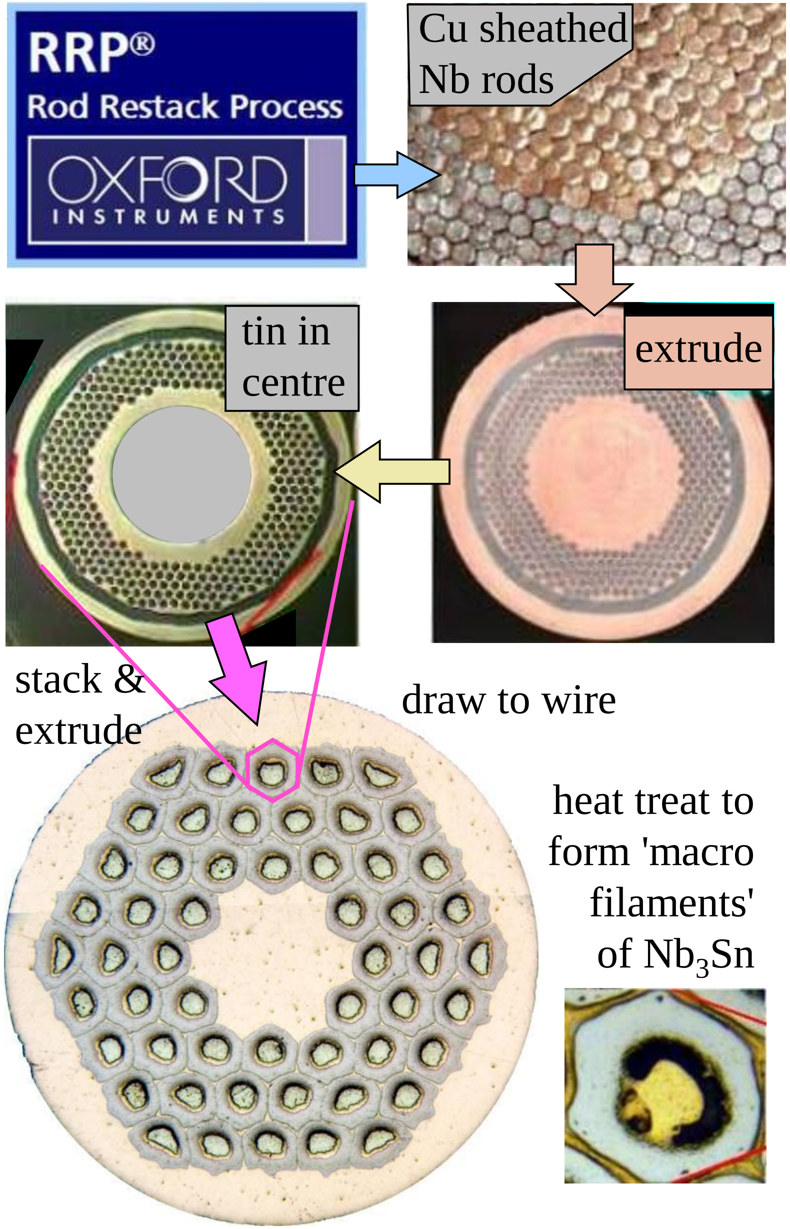


The remaining copper still contains ~ 3wt% tin and has a high resistivity ~ $6\times10^{-8}\Omega$m.
So include 'islands' of pure copper surrounded by a diffusion barrier



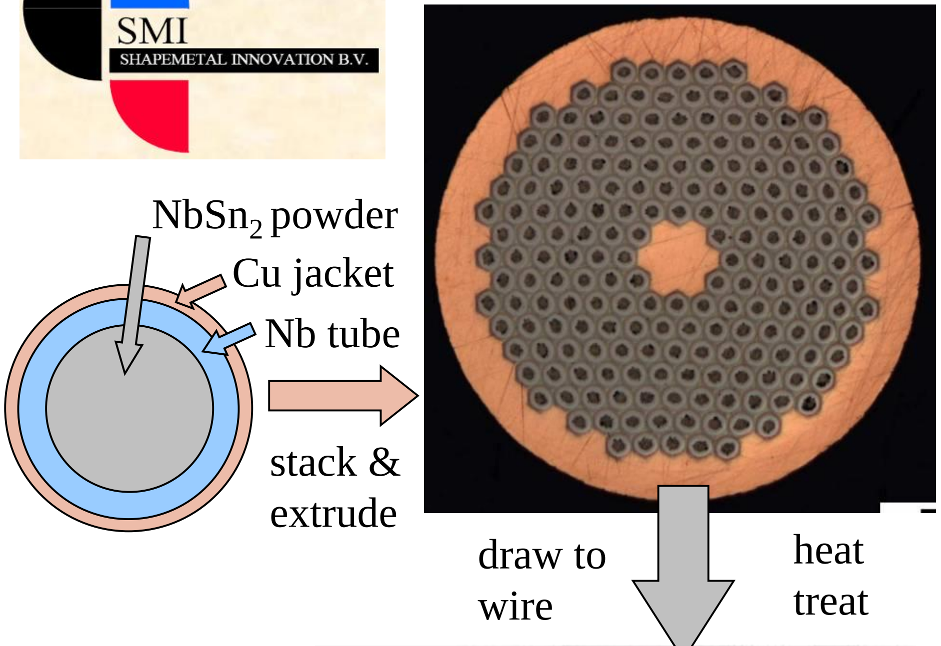


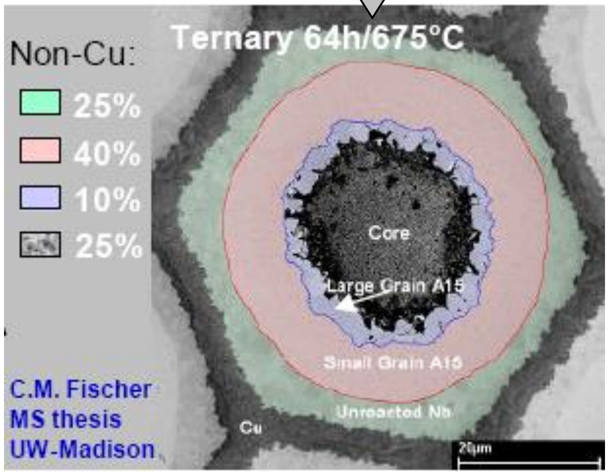
Nb_3Sn with higher engineering J_c



SMI
SHAPEMETAL INNOVATION B.V.

Powder in tube PIT





C.M. Fischer
MS thesis
UW-Madison

both make high J_{eng}

(RRP is the highest)

but large filaments

Measurement of critical current

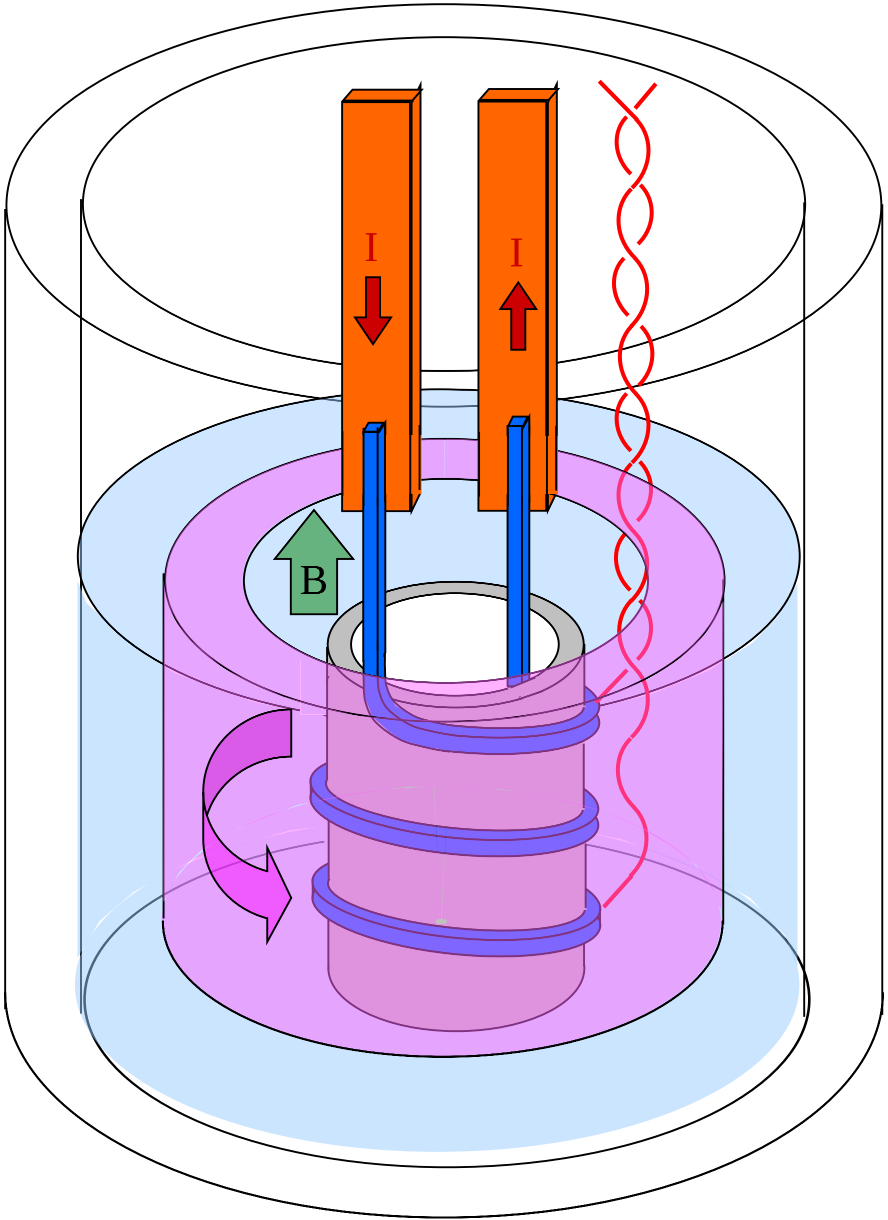
- spiral sample with current leads and voltage taps

- place in the bore of a superconducting solenoid

- put in cryostat

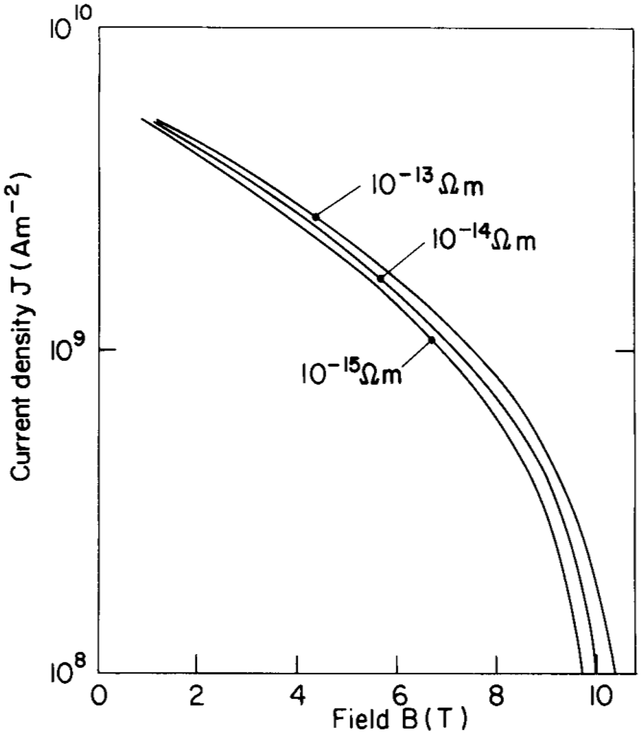
- immerse in liquid helium

- at each field level slowly increase the current is and measure the voltage across the test section

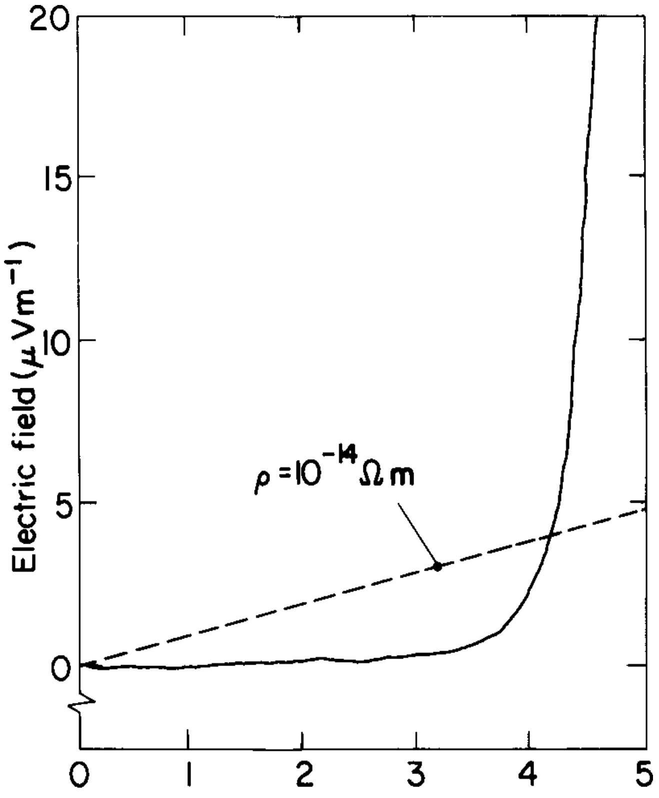


Resistive transition 1

- in practical magnet conductors, the boundary between superconducting and resistive states is not sharp, but slightly blurred.



- measure J_c with voltage taps across sample ⇨ gradual voltage rise
- must define J_c at a given measurement sensitivity
  - electric field
  - effective resistivity.



Overall Current density 10^{8} A.m^{-2}

- usual definitions are **$\rho = 10^{-14}\ \Omega m$ or $E = 1\ \mu V.m^{-1}$**
- *critical current defined at this level is about what you would expect the conductor in a resin impregnated solenoid to achieve. At higher resistivity, self heating would start to raise the internal temperature and reduce the critical current*

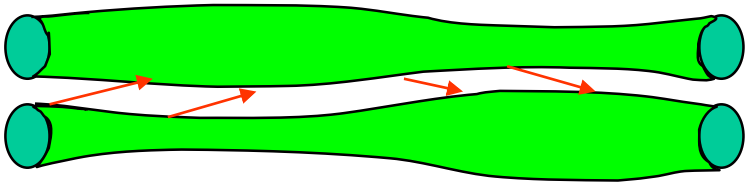
Resistive transition 2

- find empirically that the resistive transition may be represented by a power law

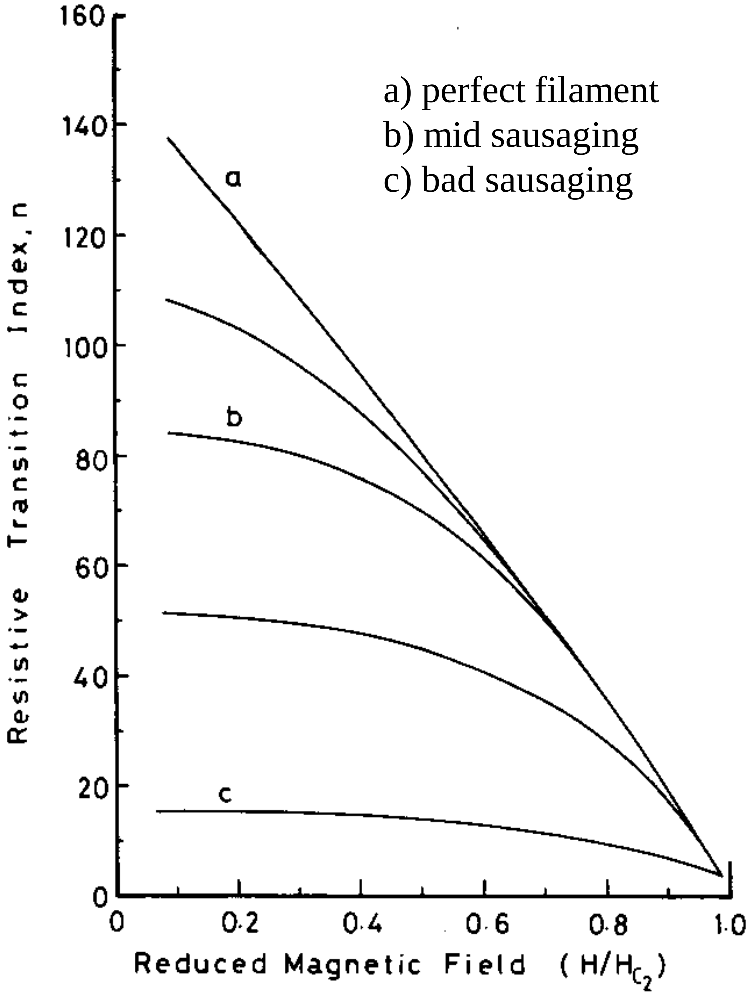
$$\rho(J) = \rho_o \left\{ \frac{J}{J_o} \right\}^n$$

where ***n*** is called the resistive transition index.

- the effect is partly within the filaments (flux flow) and partly between the filaments
- 'sausaging of the filaments, forces current to cross the copper matrix as critical current is approached.

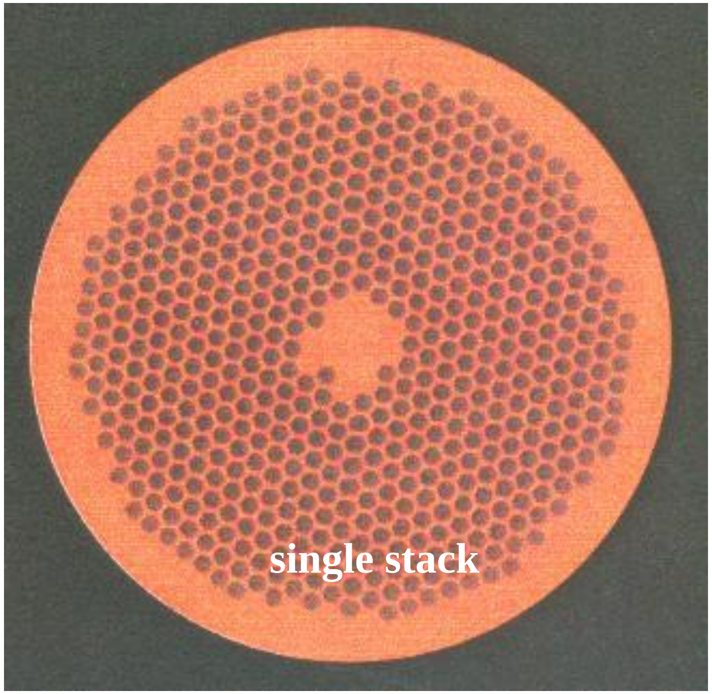


- resistive transition can be the main source of decay in persistent magnets
- 'n' is often taken as a measure of quality - look for n > 50
- HTS conductors so far have low n ~ 5 - 10

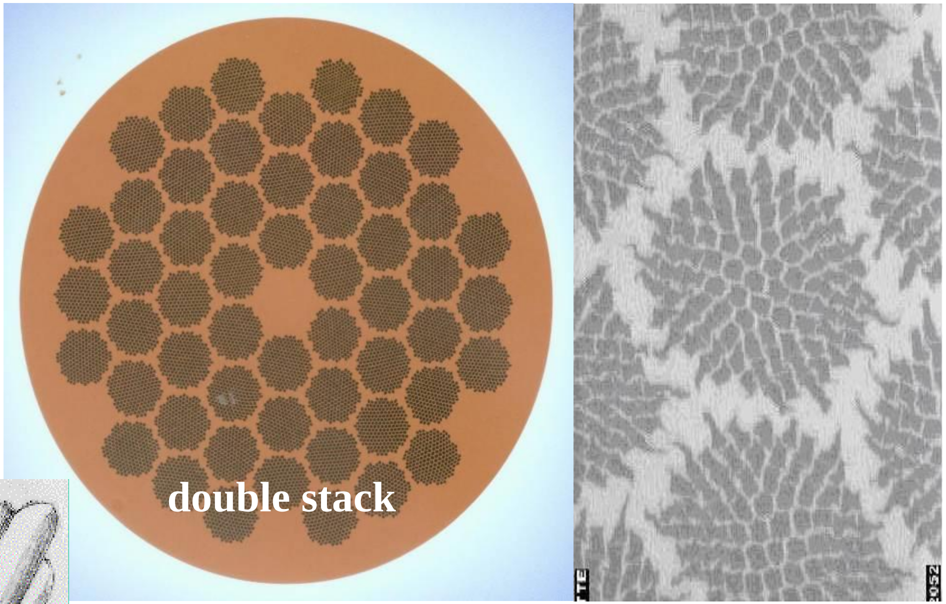


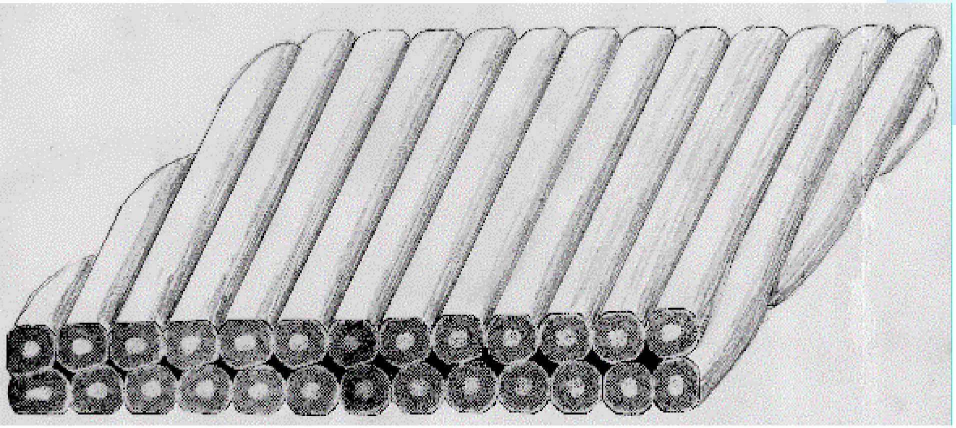
a) perfect filament
b) mid sausaging
c) bad sausaging

Conductors for accelerator magnets



- to date, all superconducting accelerators have used NbTi superconductor.
- to control field errors and ac losses, the filaments must be < 10µm diameter (lectures 2 & 3)

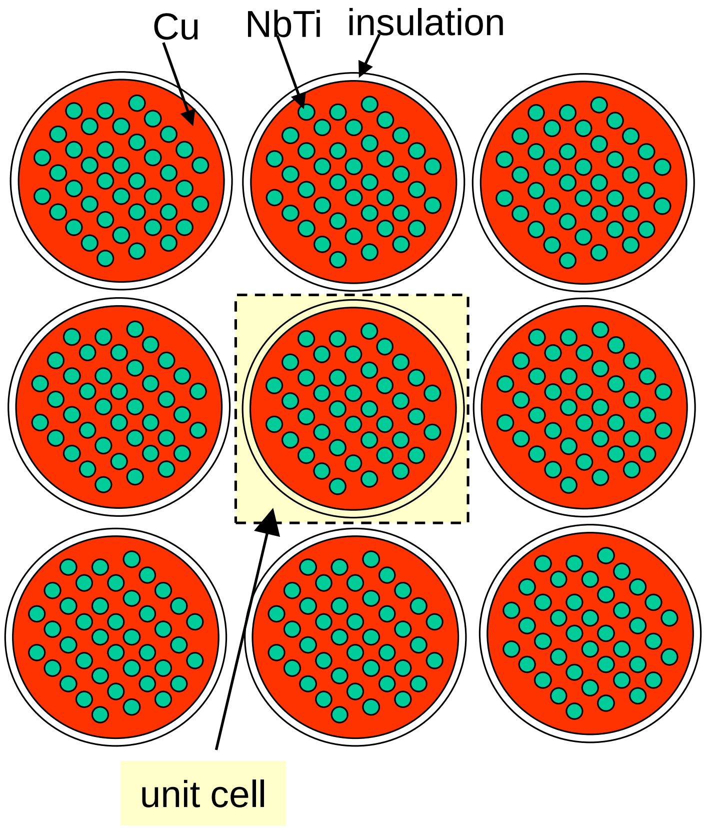




- to get the necessary high operating currents, many wires must be cabled together.

Engineering current density and filling factors

In magnet design, what really matters is the overall 'engineering' current density J_{eng}



$$J_{eng} = current \,/\, unit\ cell\ area = J_{sup} \times \lambda_{su}$$

where λ_{su} = *filling factor of superconductor in unit cell*

filling factor superconductor in the wire $\lambda_{sw} = 1 / (1+mat)$

where *mat* = matrix : superconductor ratio

typically:

for NbTi *mat* = 1.2 to 3.0 ie λ_{sw} = 0.45 to 0.25

for Nb_3Sn *mat* = 2.0 to 4.0 ie λ_{sw} = 0.33 to 0.2

for B2212 *mat* = 3.0 to 4.0 ie λ_{sw} = 0.25 to 0.2

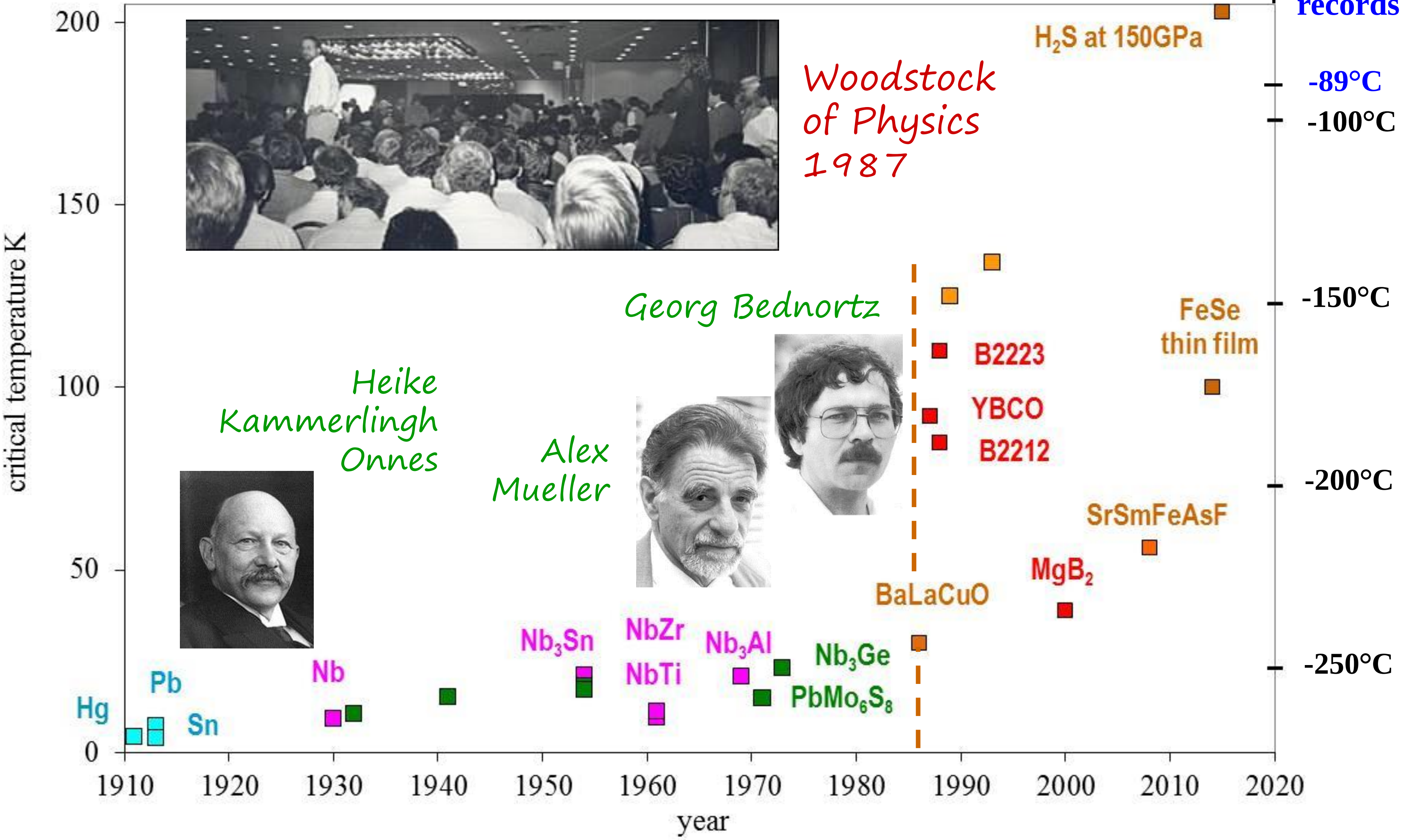
For Nb_3Sn and B2212 the area of superconductor is not well defined, so often define J_{sup} over ‘non matrix’or ‘non Cu’ area, which is greater than superconductor area.

λ_{wu} fraction of wire in unit cell taking account of space occupied by insulation, cooling channels, reinforcement etc: typically λ_{wu}~ 0.7 to 0.8

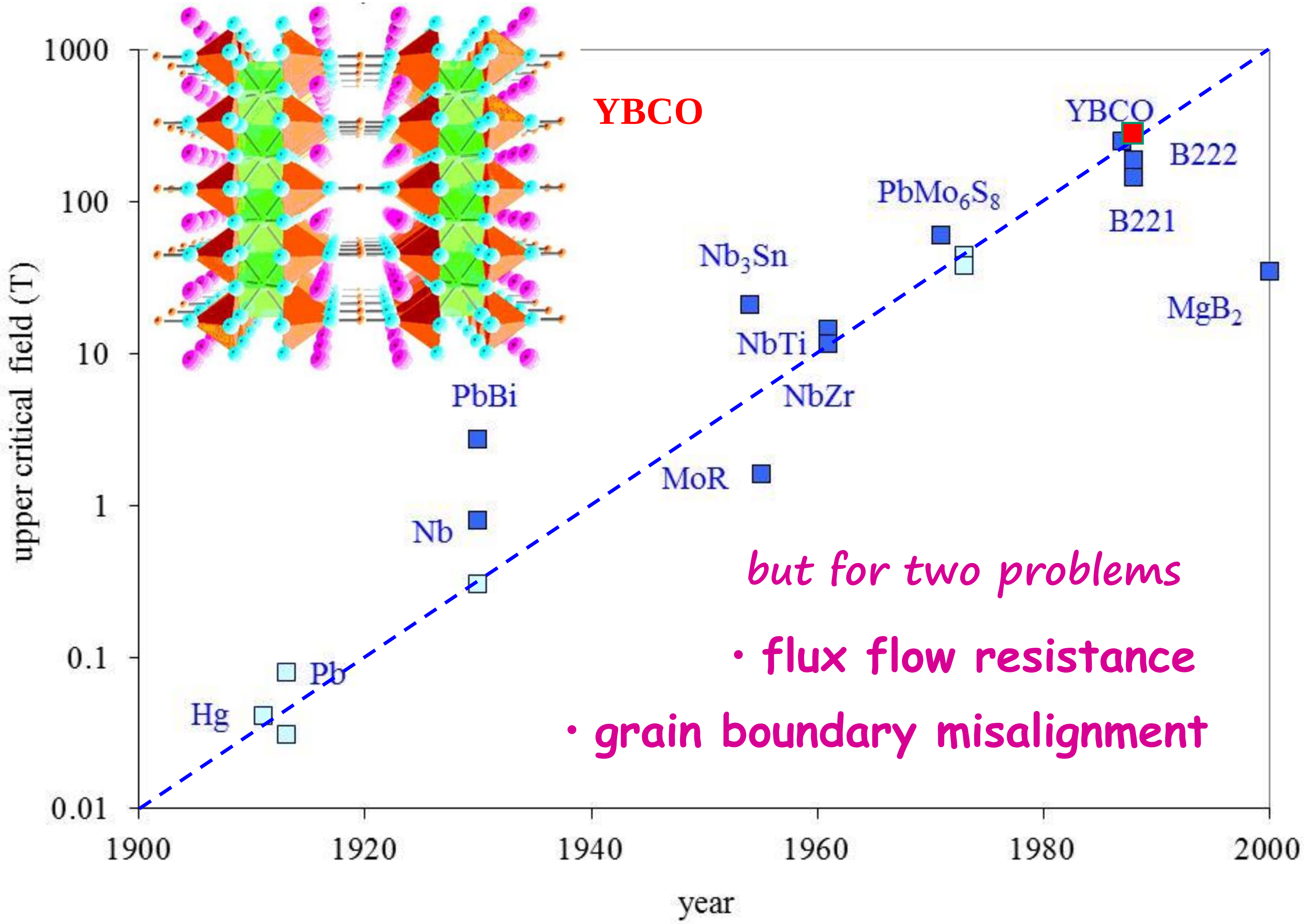
filling factor of superconductor in unit cell $\lambda_{su} = \lambda_{sw} \times \lambda_{wu}$

So typically J_{eng} is only 15% to 30% of $J_{supercon}$

Critical temperatures



Wonderful materials for magnets



but for two problems

- flux flow resistance
- grain boundary misalignment

Yttrium Barium Copper Oxide

$YBa_2Cu_3O_7$

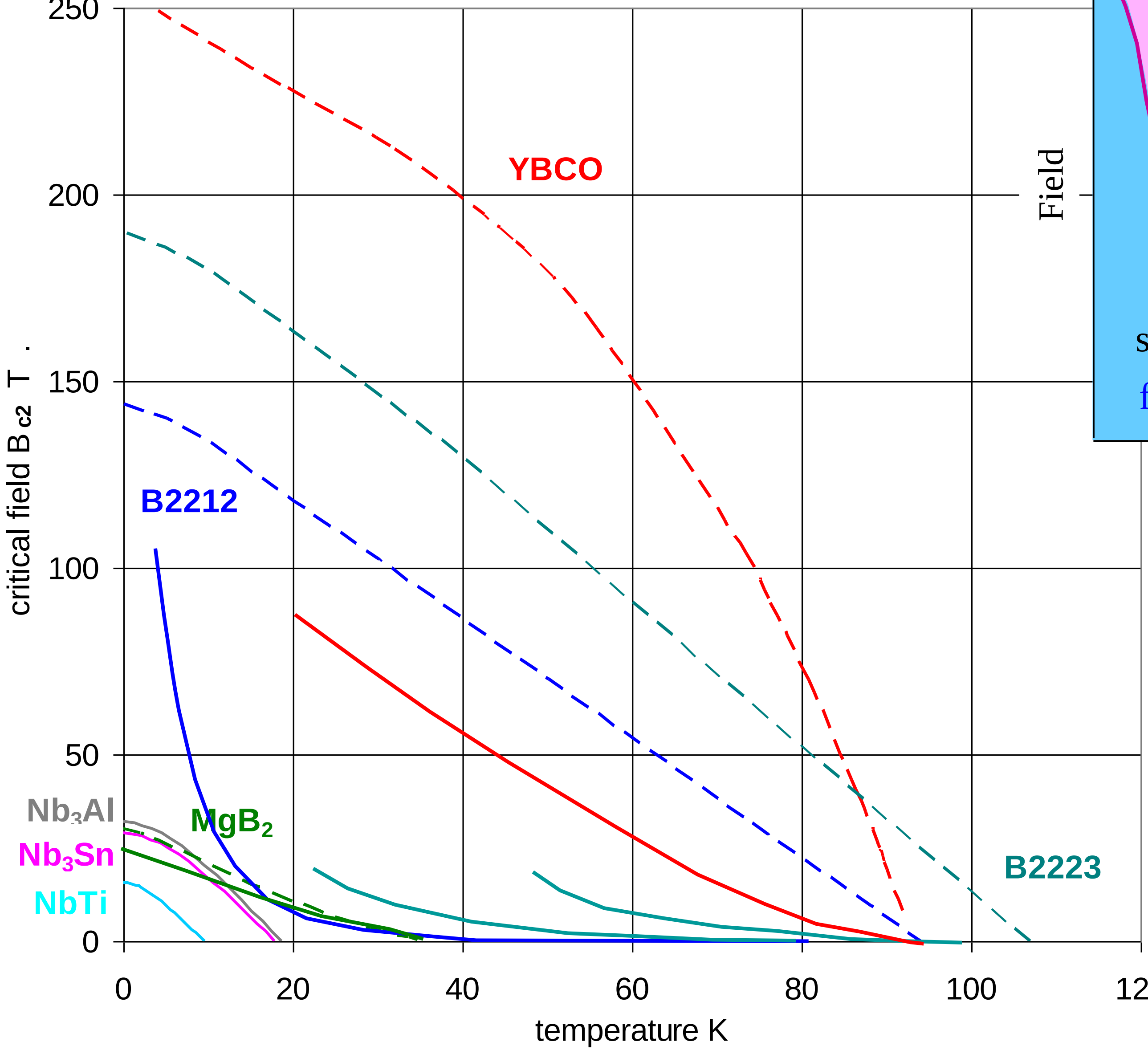
Barium Strontium Calcium Copper Oxide BSCCO

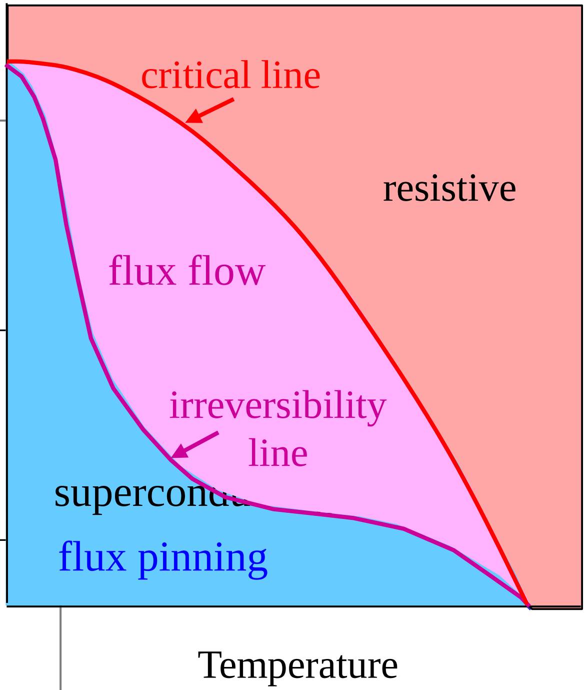
$Bi_2Sr_2CaCu_2O_8$ B2212

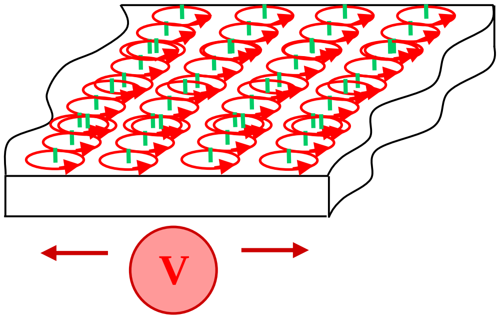
$Bi_2Sr_2Ca_2Cu_3O_6$ B2223

Magnesium Diboride MgB_2

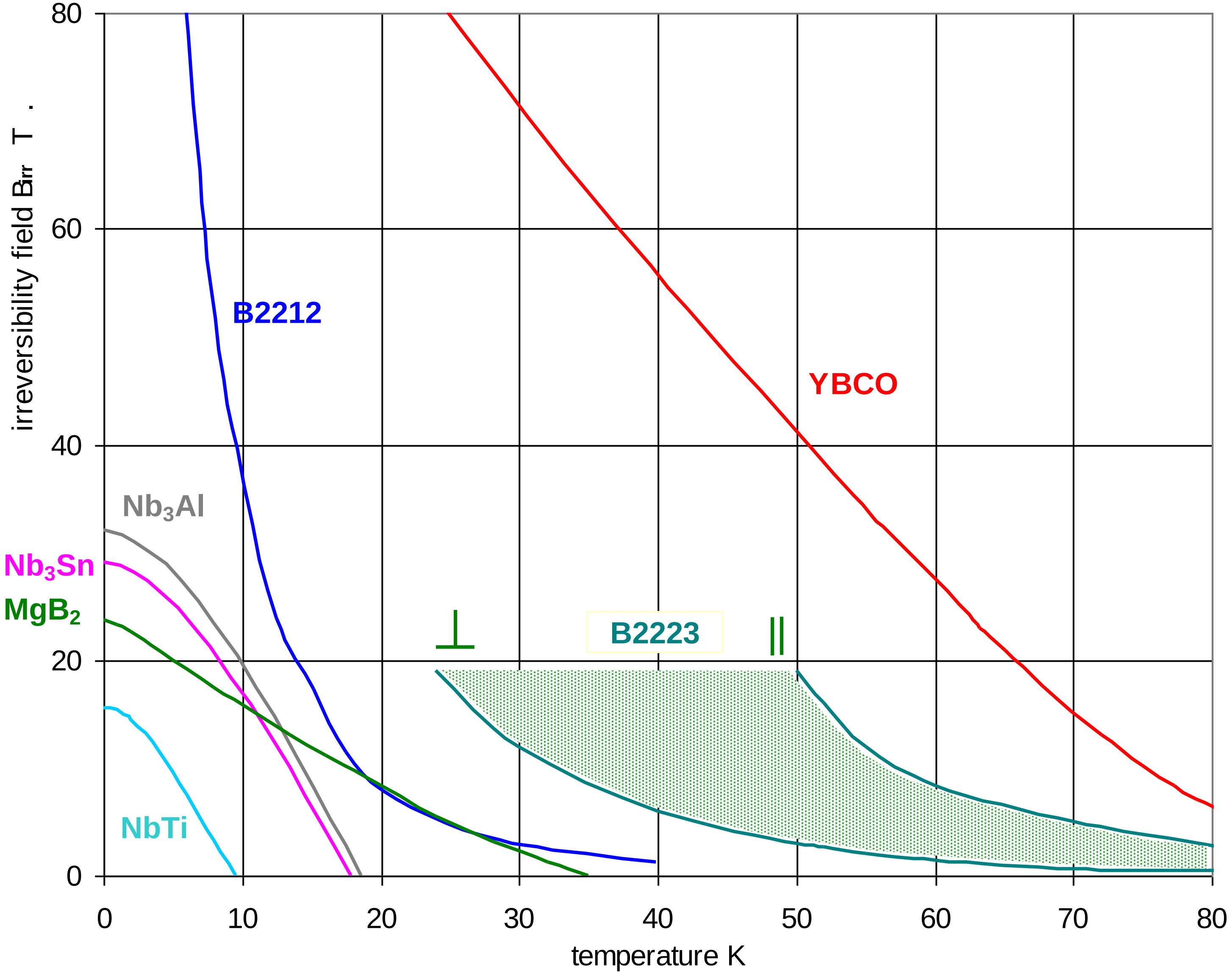
1) Flux flow resistance







Accessible fields for magnets

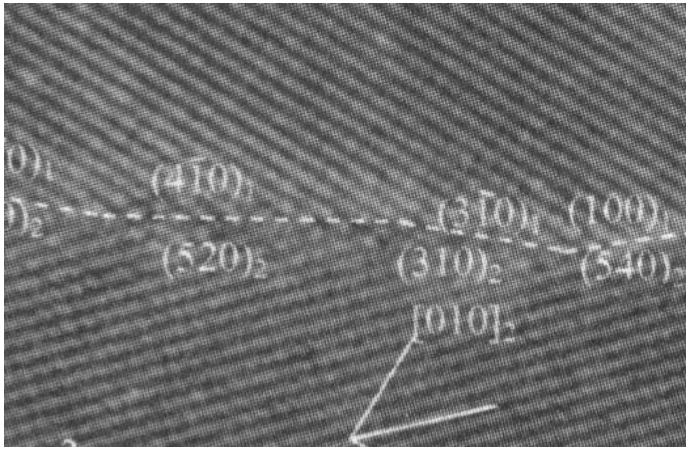


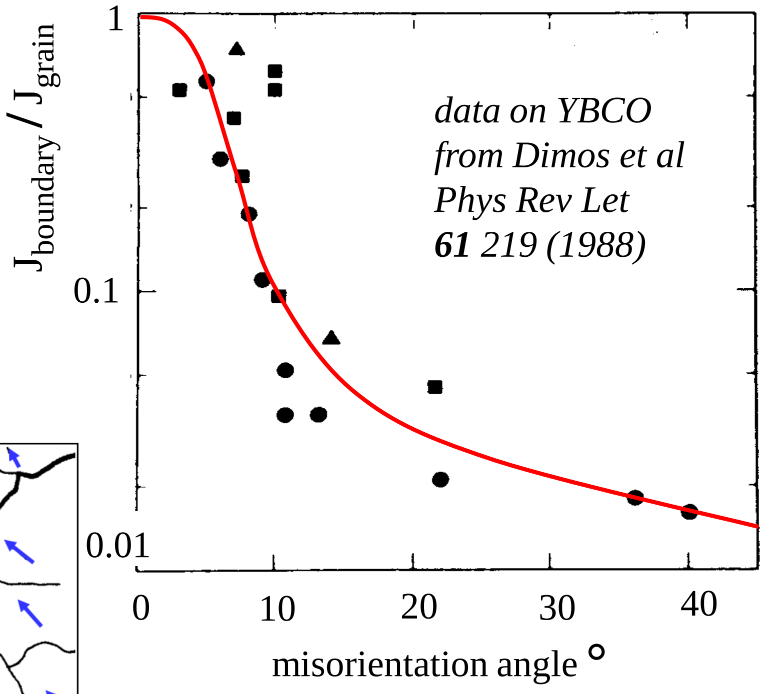
so here's the parameter space for future magnets

and we still need to work on the current density

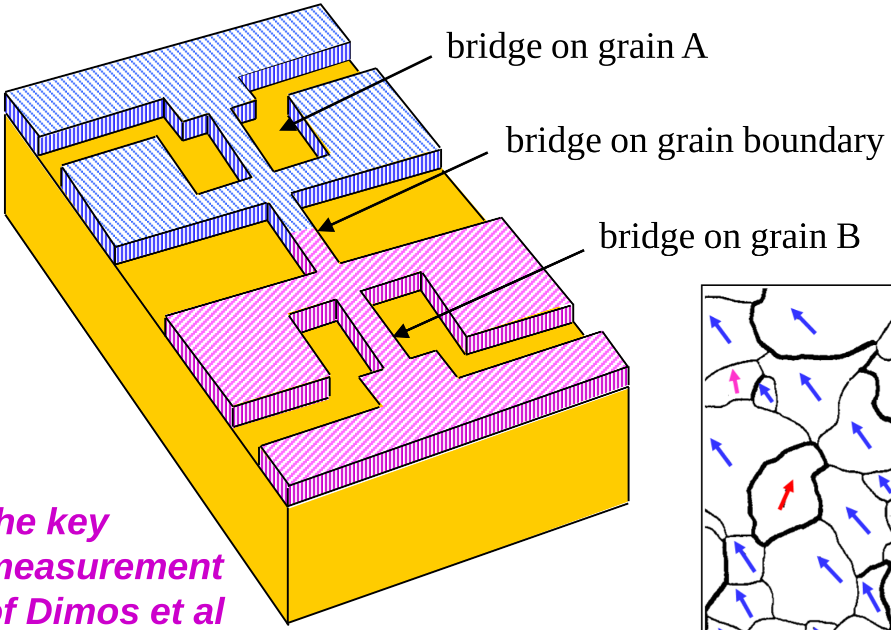
2) Grain boundary mismatch

- crystal planes in grains point in different directions
- critical currents are high within the grains
- ***J*$_c$** across the grain boundary depends on the misorientation angle
- for good ***J*$_c$** must align grains to within a few degrees

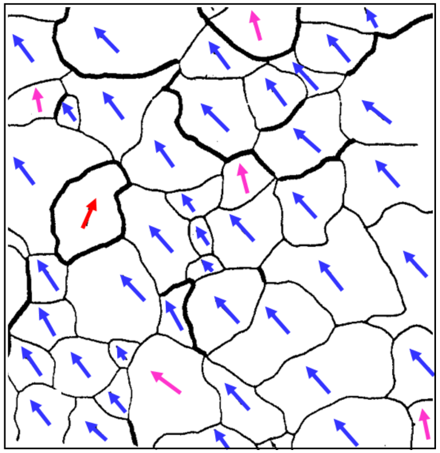




*data on YBCO from Dimos et al Phys Rev Let **61** 219 (1988)*



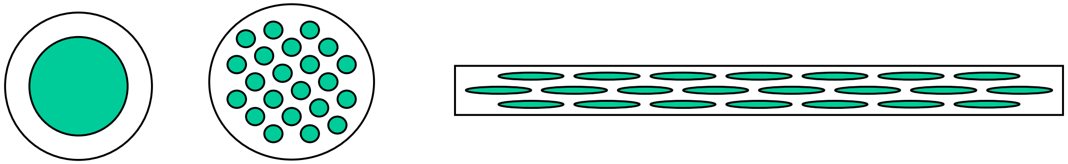
the key measurement of Dimos et al

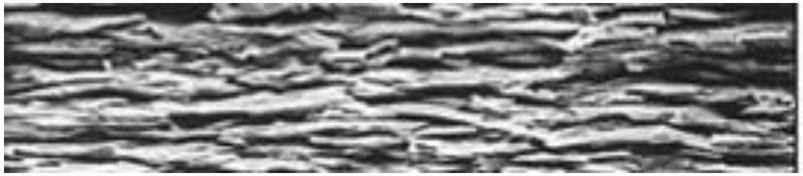


must align (almost) all the grains

BSCCO conductors

- BSCCO powder in silver matrix
- silver is transparent to oxygen
- grains tend to line up during processing in contact with silver
- but low irreversibilty field





$Bi_2Sr_2Ca_2Cu_3O_6$ B2223

- θ_c ~110 K ⇨ 77 K use
- complex thermo-mechanical fabrication route to produce uniaxial texture
- reaction of 2212 + stuff to 2223 occurs primarily in the solid state at ~840°C
- texture requires pressing to a large aspect ratio tape (4 x 0.2 mm)
- anisotropic J_c
- mature conductor – semi commercial production for > 10 years

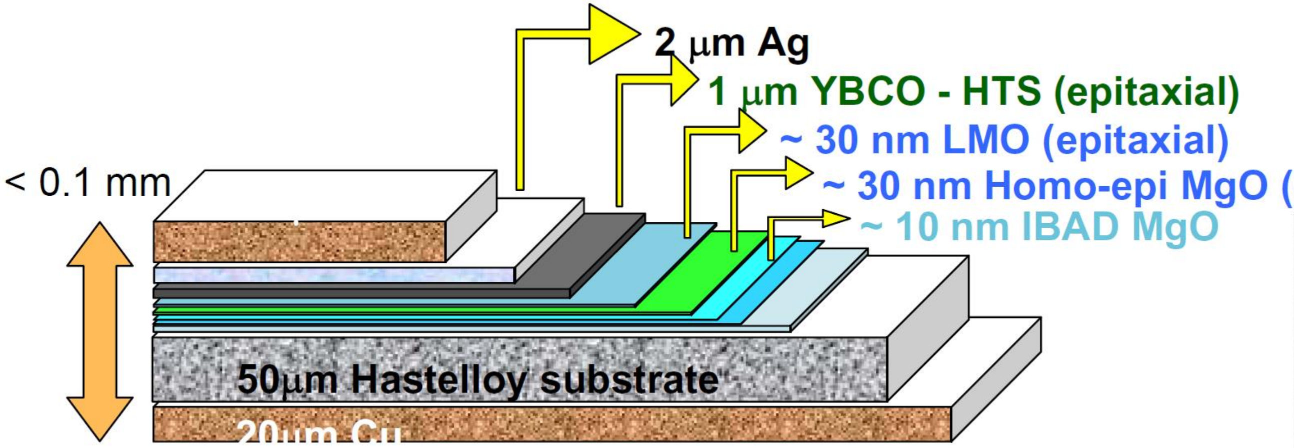
$Bi_2Sr_2CaCu_2O_8$ B2212

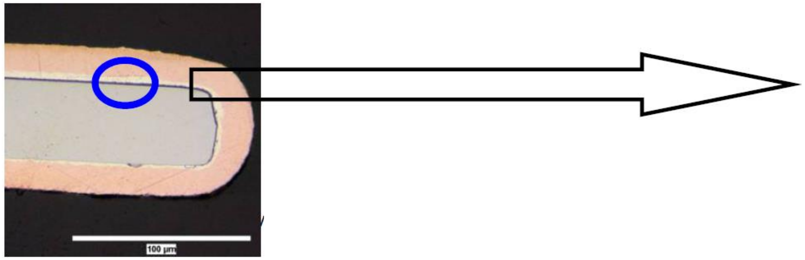
- θ_c ~ 85 K
- produced by melting B2212 powder inside silver matrix
- round filaments in round wires
- isotropic J_c
- high J_c but used to be only in short samples
- recently achieved high J_c in long lengths
- engineers like round wire
- can make Rutherford cables

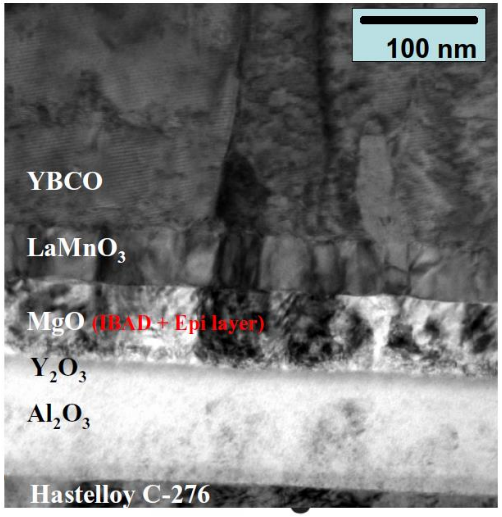
Coated YBCO tape

- YBCO has the best irreversibility field, but it is very sensitive to grain boundary misalignment
- the grains do not line up naturally - they must be persuaded
- deposit YBCO on a substrate where the grains are aligned and the lattice roughly matches YBCO

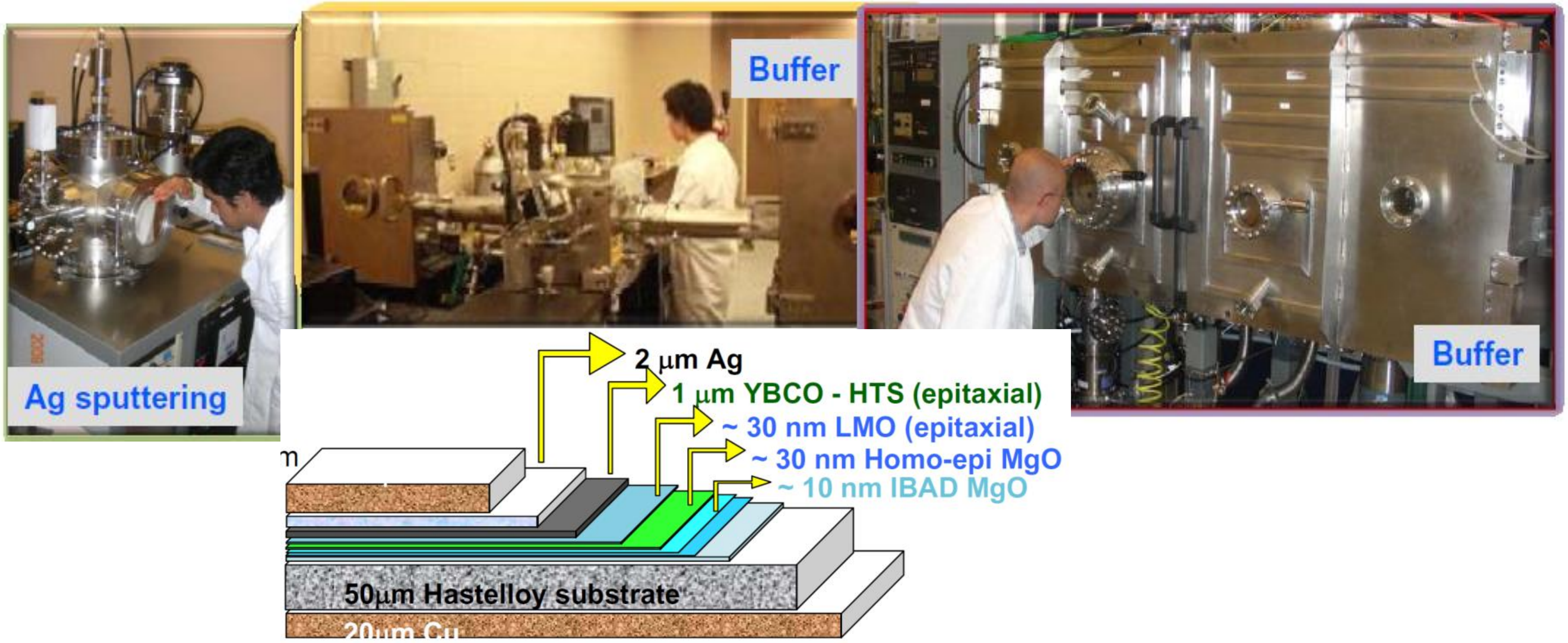
OK high field and high temperature







YBCO coated tape at SuperPower Inc.



2 μm Ag

1 μm YBCO - HTS (epitaxial)

~ 30 nm LMO (epitaxial)

~ 30 nm Homo-epi MgO

~ 10 nm IBAD MgO

50μm Hastelloy substrate

20μm Cu

Lecture 1: concluding remarks

- superconductors allow us to build magnets which burn no power (except refrigeration)
- ampere turns are cheap, so don't need iron
  ⇒ fields higher than iron saturation (but still use iron for shielding)
- performance of all superconductors described by the critical surface in $\boldsymbol{B\ J\ \theta}$ space,
- three kinds of superconductor
  - **type 1**: low temperature, unsuitable for high field
  - **type 2**: low temperature, good for high field - but must create flux pinning to get current density
  - **HTS**: high temperature, high field - but current density is a difficult problem
- NbTi is the most common commercial superconductor - standard production process
- Nb_3Sn has higher critical field & temperature - specialized commercial production
- BSCO high temperature **or** high field, but not both - prototype commercial production
- YBCO high temperature **and** high field, but must align the grains - prototype commercial production
- measure I_c to check specification, the index n indicates quality
- for accelerators, so far it's only been NbTi, usually in Rutherford cables

Some useful references

Superconducting Magnets

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- Case Studies in Superconducting Magnets, Second edition: Y Iwasa, pub Springer (2009), ISBN 978-0-387-09799-2.
- High Field Superconducting Magnets: FM Asner, pub Oxford University Press (1999) ISBN 0 19 851764 5
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- Experimental Techniques for Low-temperature Measurements: J. W. Ekin Pub. Oxford University Press, ISBN 978-0-19-857054-7
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- Materials at Low Temperature: Ed RP Reed & AF Clark, pub Am. Soc. Metals 1983. ISBN 0-87170-146-4
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- High Temperature Superconductors: Processing and Science, A Bourdillon and NX Tan Bourdillon, Academic Press, ISBN 0 12 117680 0
- Superconductivity: A Very Short Introduction by Stephen J. Blundell: Oxford University Press (2009) ISBN978-0-19-954090-7

on the Web

- **Lectures on Superconductivity** http://www.msm.cam.ac.uk/ascg/lectures.
  A series of lectures produced for SCENET by Cambridge University: fundamentals, materials, electronics, applications. Also available as a DVD

- **Superconducting Accelerator Magnets** http://www.mjb-plus.com.
  A course developed from SSC experience, available from website for $20

- www.superconductors.org website run by an enthusiast; gives some basic info and links

- **Superconductivity Course** at the (UK) Open University.
  http://openlearn.open.ac.uk/course/view.php?id=2397 Good coverage of basics.

- **Wikipedia** on Superconductivity http://en.wikipedia.org/wiki/Superconductivity
  Good on basics with lots of references and links.

- **European Society for Applied Superconductivity** http://www.esas.org/
  News, events and people in the area of applied superconductivity

- **CONECTUS** Consortium of European Companies determined to use Superconductivity
  http://www.conectus.org/

- **IEEE Council on Superconductivity** http://www.ewh.ieee.org/tc/csc/
  News, events and people in the area of applied superconductivity (US based)

Materials data on the Web

- Cryogenic properties (1-300 K) of many solids, including thermal conductivity, specific heat, and thermal expansion, have been empirically fitted and the equation parameters are available free on the web at **www.cryogenics.nist.gov**
- Thermodynamic properties of gases (and liquids) available free as a programme which you can interrogate for your own temperature interval etc. **http://webbook.nist.gov/chemistry/fluid/**
- Plots and automated data-look-up using the NIST equations are available on the web for a fee from **www.cpia.jhu.edu**
- Other fee web sites that use their own fitting equations for a number of cryogenic material properties include: **www.cryodata.com** (cryogenic properties of about 100 materials), and **www.jahm.com** (temperature dependent properties of about 1000 materials, many at cryogenic temperatures).
- Commercially supplied room-temperature data are available free online for about 10 to 20 properties of about 24,000 materials at **www.matweb.com**

thanks to Jack Ekin of NIST and Charles Monroe for this information

Cryodata Software Products

GASPAK
properties of pure fluids from the triple point to high temperatures.

HEPAK
properties of helium including superfluid above 0.8 K, up to 1500 K.

STEAMPAK
properties of water from the triple point to 2000 K and 200 MPa.

METALPAK, CPPACK, EXPAK
reference properties of metals and other solids, 1 - 300 K.

CRYOCOMP
properties and thermal design calculations for solid materials, 1 - 300 K.

SUPERMAGNET
four unique engineering design codes for superconducting magnet systems.

KRYOM
numerical modelling calculations on radiation-shielded cryogenic enclosures.