

# Flavor Structure of the Nucleon Sea from Lattice QCD

Giunn-Wei Chen

National Taiwan U. & MIT

Collaborators: Huey-Wen Lin, Saul D. Cohen,  
Xiangdong Ji, Jianhui Zhang

arXiv: 1402.1462 + 1603.06664

# Parton Distribution Functions (PDF)---an Infinite-Body Problem

- The number of quark anti-quark pairs diverges (manifestation of non-perturbative nature of the problem)

# One Possible Solution...

- Lattice QCD: making the degree of freedom finite by discretizing the space time
- Goal: Computing the x-dependence of PDF's from first principles (QCD).

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# Past Limitation

- Traditional approach: can only calculate lower moments PDFs.
- Still first principle, carried out successfully: close to using physical parameters---highly non-trivial and demanding in computing power.
- However, it also means the community has reached the limit on what one can learn from the lower moments.

# New Hopes

- Smearred sources: Davoudi & Savage
- Current-current correlators: K.-F. Liu & S.-J. Dong; Braun & Müller; Detmold & Lin
- Xiangdong Ji (Phys. Rev. Lett. 110 (2013) 262002): quasi-PDF: computing the x-dependence directly.

# Lattice PDF: from Moments to the Sea

- Quark PDF in a proton:  $(\lambda^2 = 0)$

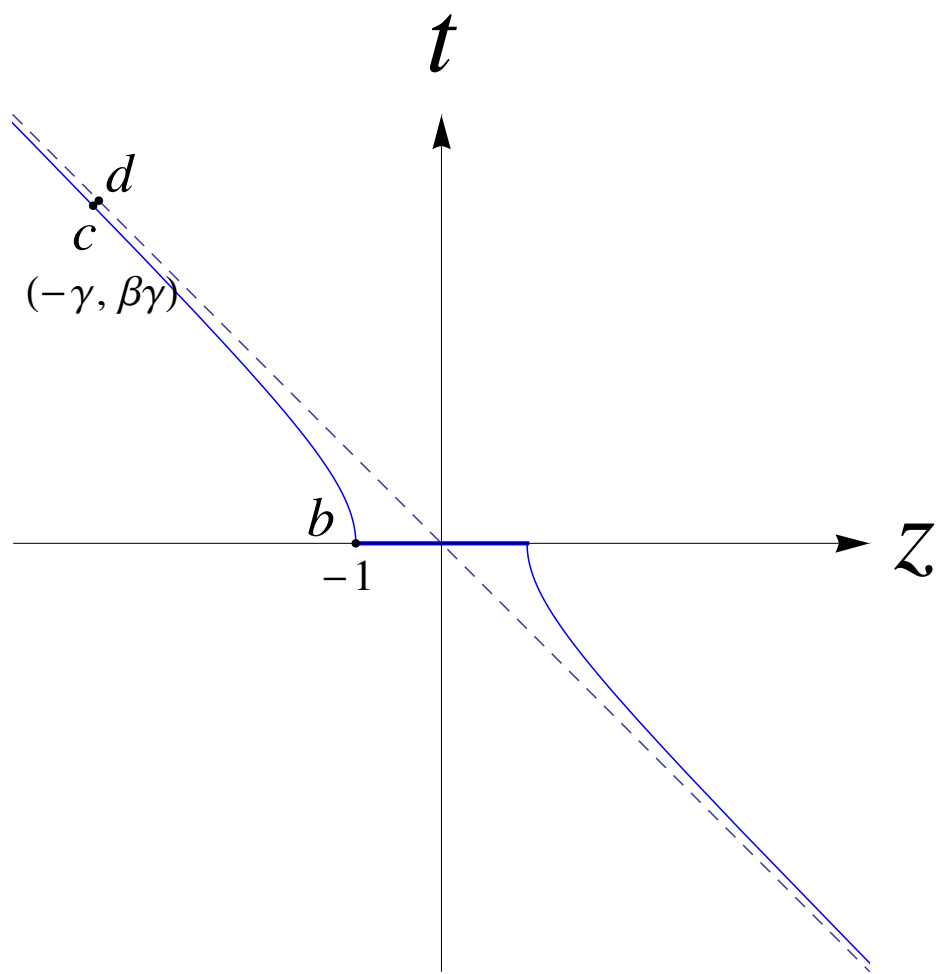
$$q(x, \mu^2) = \int \frac{d\xi^-}{4\pi} e^{ix\xi^- P^+} \langle P | \bar{\psi}(0) \lambda \cdot \gamma \Gamma \psi(\xi^- \lambda) | P \rangle$$

- Euclidean lattice: light cone operators cannot be distinguished from local operators
- Moments of PDF given by local twist-2 operators; LPDF limited to first few moments;  
Sea quarks cannot be isolated

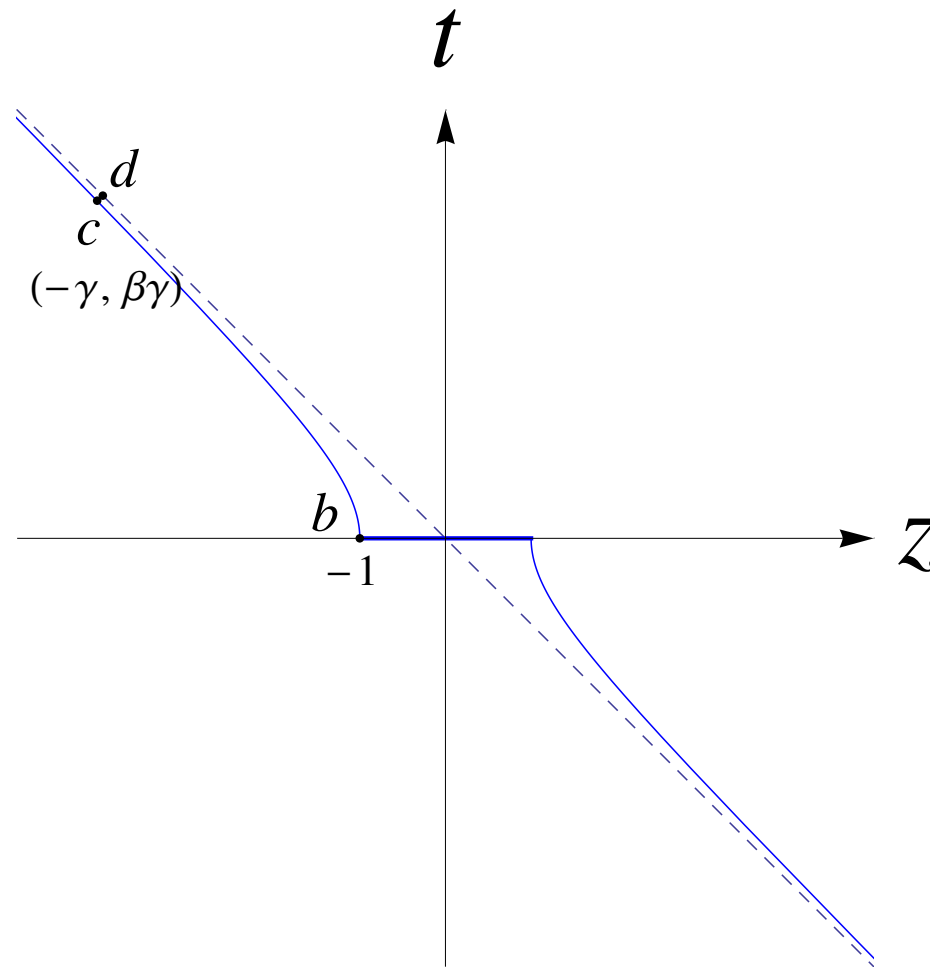
- Quark PDF in a proton:  $(\lambda^2 = 0)$

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- Boost invariant in the z-direction, rest frame OK
- Quark bilinear op. always on the light cone
- What if the quark bilinear is slightly away from the light cone (space-like) in the proton rest frame?







- Then one can find a frame where the quark bilinear is of equal time but the proton is moving.

## Review: Ji's LPDF

$$\begin{aligned}\tilde{q}(x, \mu^2, P^z) &= \int \frac{dz}{4\pi} e^{-ixzP^z} \langle P | \bar{\psi}(0) \lambda \cdot \gamma \Gamma \psi(z\lambda) | P \rangle \\ &\equiv \int \frac{dz}{2\pi} e^{-ixzP^z} h(zP^z) P^z\end{aligned}$$

$$\lambda^\mu = (0, 0, 0, 1)$$

- Taylor expansion yields

$$\bar{\psi} \lambda \cdot \gamma \Gamma (\lambda \cdot D)^n \psi = \lambda_{\mu_1} \lambda_{\mu_2} \cdots \lambda_{\mu_n} O^{\mu_1 \cdots \mu_n}$$

op. symmetric but not traceless

# Review: Ji's LPDF

$$\langle P | O^{(\mu_1 \dots \mu_n)} | P \rangle = 2a_n P^{(\mu_1} \dots P^{\mu_n)}$$

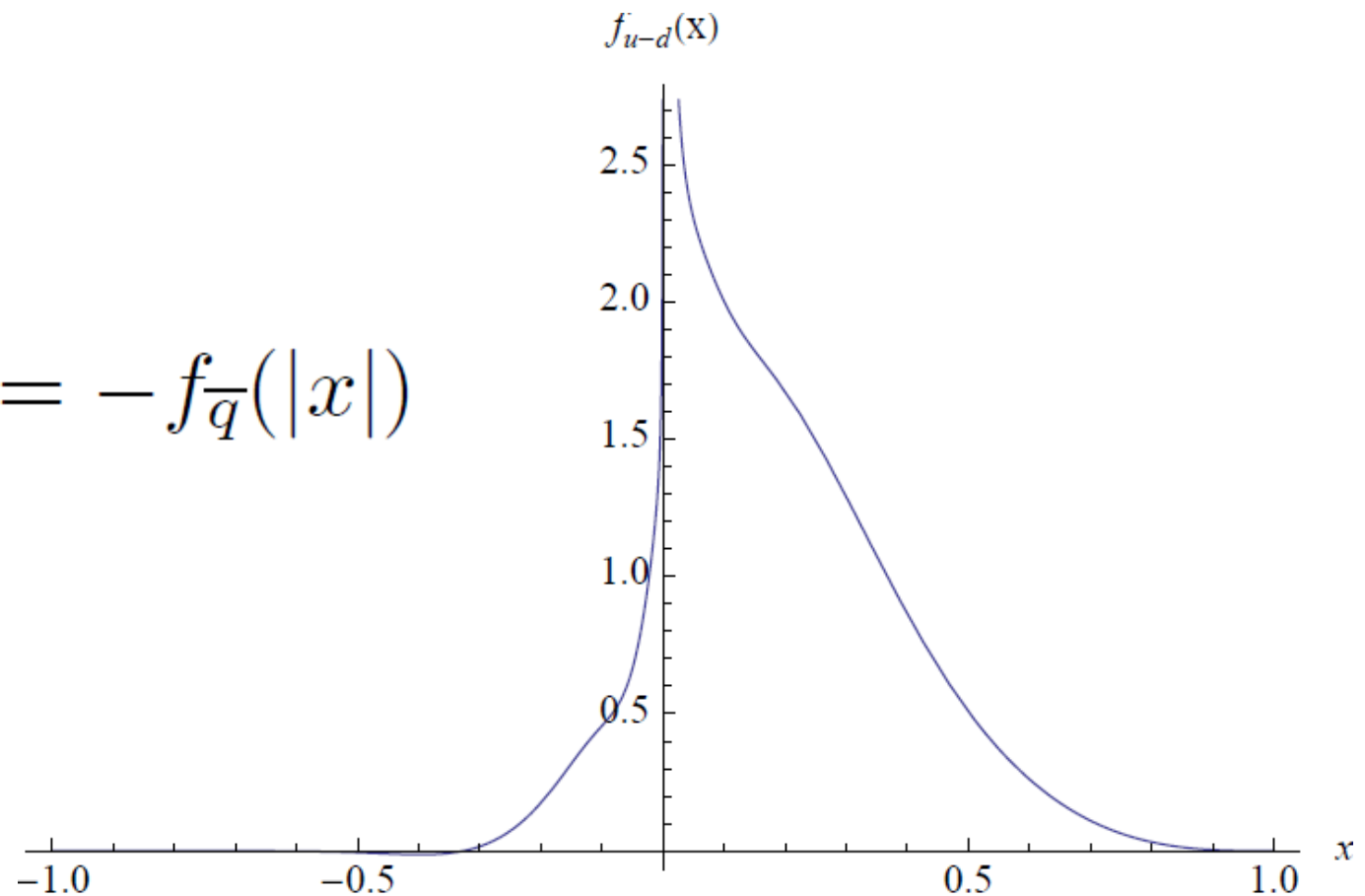
- LHS: trace, twist-4  $\mathcal{O}(\Lambda_{\text{QCD}}^2 / (P^z)^2)$  corrections, parametrized in this work
- RHS: trace  $\mathcal{O}(M^2 / (P^z)^2)$ .
- One loop matching  $\alpha_s \ln P^z$ , OPE

$$\tilde{q}(x, \Lambda, P_z) = \int \frac{dy}{|y|} Z\left(\frac{x}{y}, \frac{\mu}{P_z}, \frac{\Lambda}{P_z}\right) q(y, \mu) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{P_z^2}, \frac{M^2}{P_z^2}\right) + \dots$$

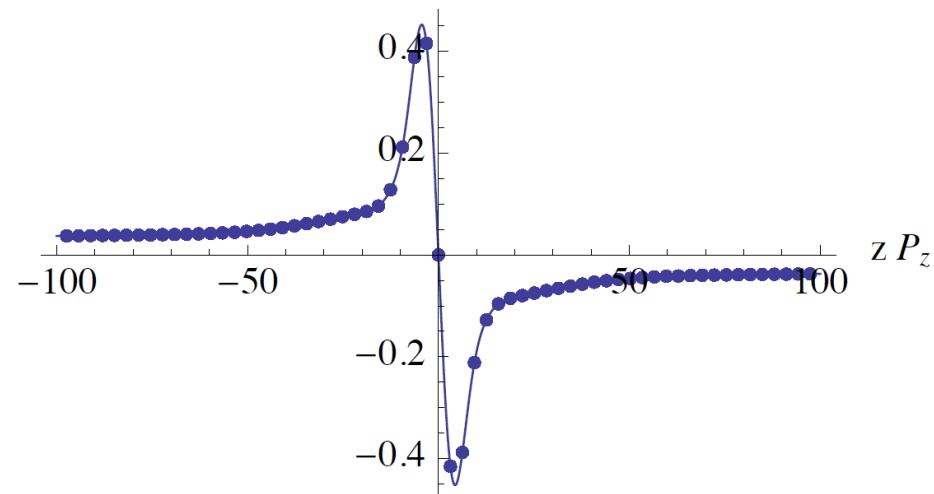
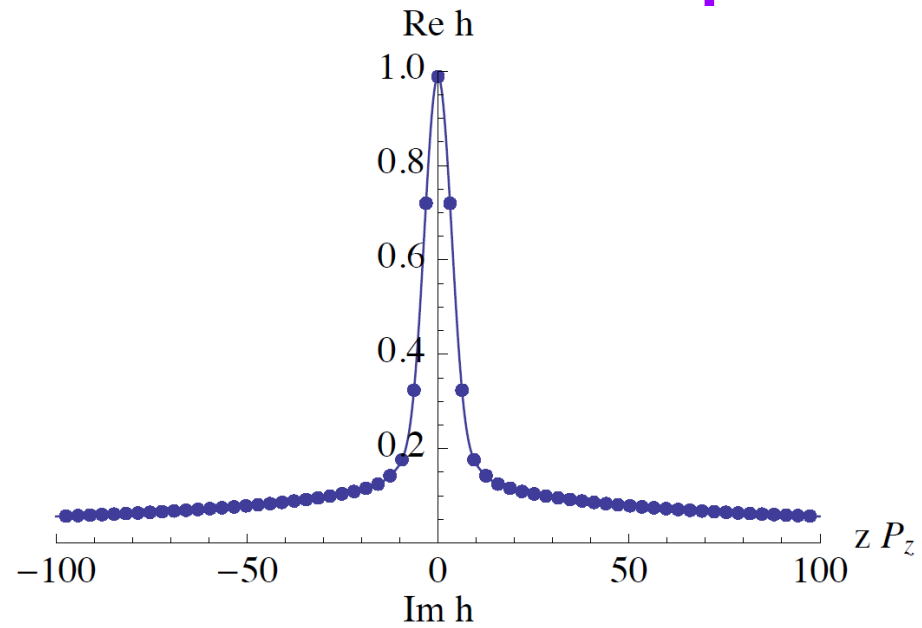
# What do we expect to see on the lattice?

- Suppose LPDF is the CTEQ PDF at  $P^z \rightarrow \infty$

$$f_q(-|x|) = -f_{\bar{q}}(|x|)$$



# in the Fourier Space



# First (isovector) LPDF Computation

- Lattice:  $24^3 \times 64$

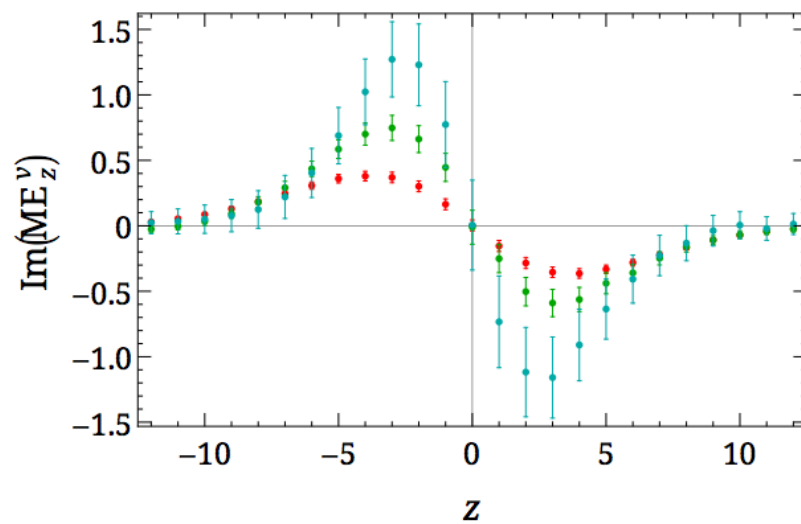
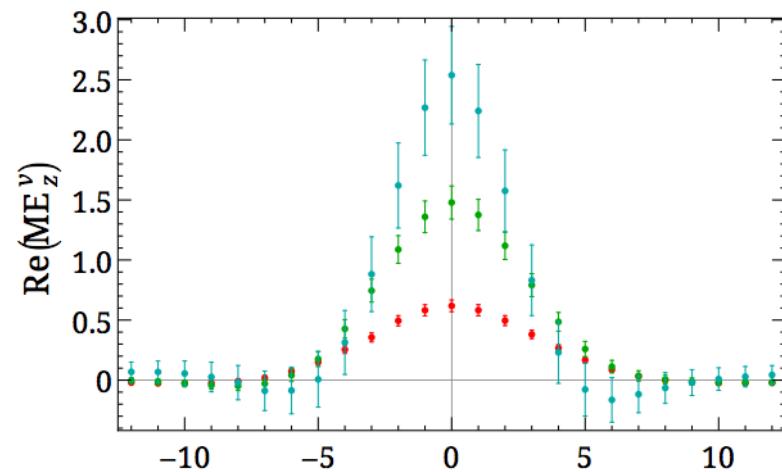
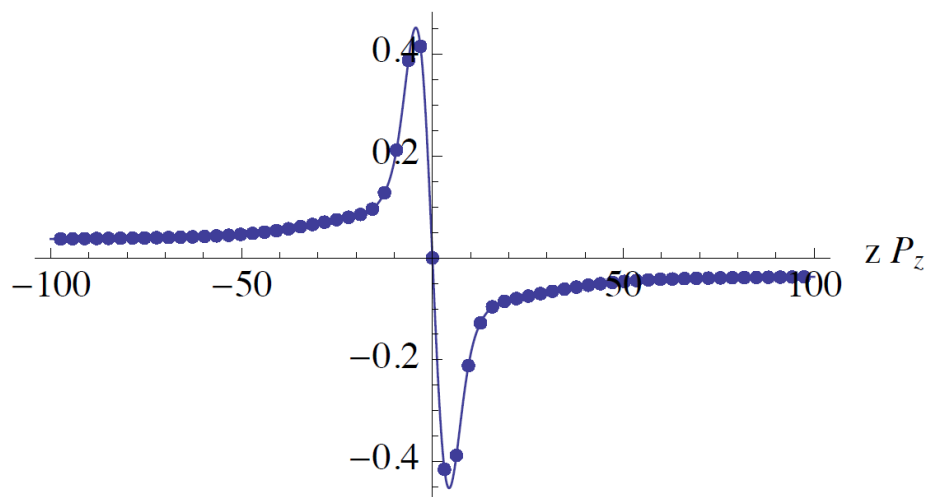
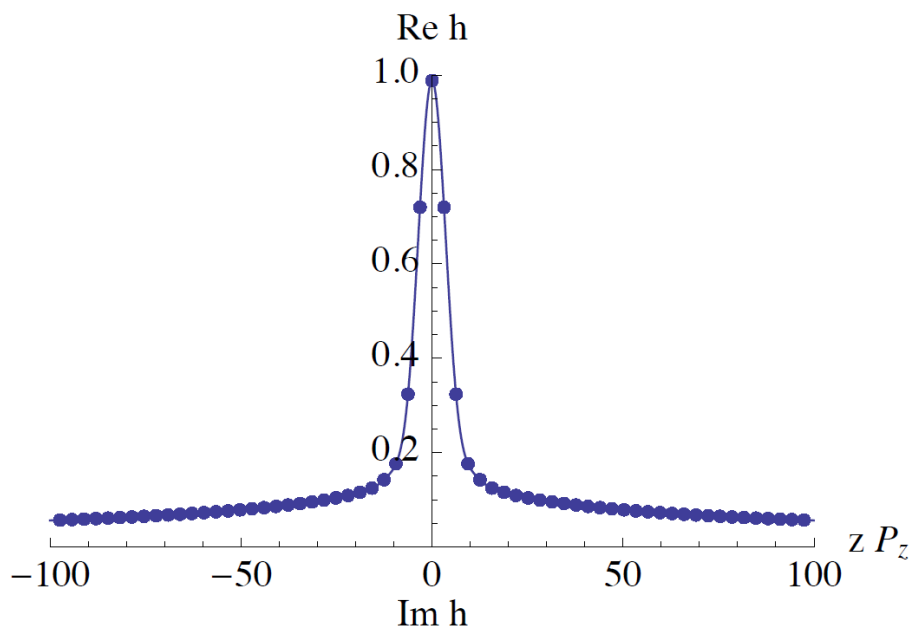
$$a \approx 0.12 \text{ fm} \quad L \approx 3 \text{ fm}$$

- Fermions: MILC highly improved staggered quarks (HISQ) Clover (valence)

$$N_f = 2 + 1 + 1 \quad M_\pi \approx 310 \text{ MeV}$$

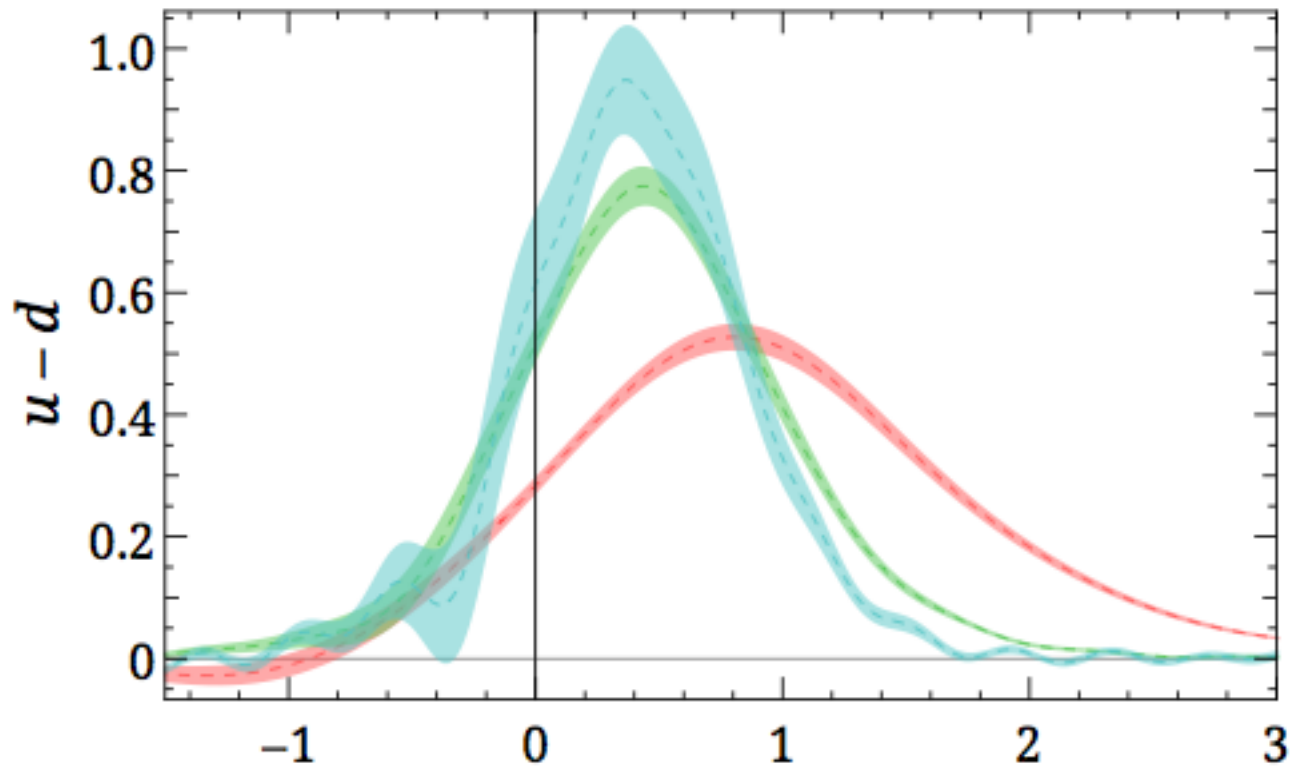
- Gauge fields/links: hypercubic (HYP) smearing, 461 config.

- $P^z = \frac{2\pi}{L}n = n \times 0.43 \text{ GeV} \quad n = 1, 2, 3, \dots$



$$P^z = \frac{0.26n}{a}, \quad n = 1, 2, 3$$

# Quasi-PDF (unpolarized)



$$P^z = \frac{2\pi}{L}n = n \times 0.43 \text{ GeV} \quad n = 1, 2, 3.$$



# RG of Wilson Coefficient

$$\tilde{q}(x, \Lambda, P_z) = \int \frac{dy}{|y|} Z \left( \frac{x}{y}, \frac{\mu}{P_z}, \frac{\Lambda}{P_z} \right) q(y, \mu) + \mathcal{O} \left( \frac{\Lambda_{\text{QCD}}^2}{P_z^2}, \frac{M_N^2}{P_z^2} \right) + \dots$$

Xiong, Ji, Zhang, Zhao

(GPD: Ji, Schfer, Xiong, Zhang; Xiong, Zhang)

factorization, Linear divergence, lattice PT

(Ma, Qiu; Ishikawa, Ma, Qiu, Yoshida; H.N. Li)

# $\mathcal{O}(M^2/(P^z)^2)$ · Corrections

$$P^z = \frac{2\pi}{L}n = n \times 0.43 \text{ GeV}$$

- Computed to all orders in  $\mathcal{O}(M^2/(P^z)^2)$  .

$$q(x) = \sqrt{1+4c} \sum_{n=0}^{\infty} \frac{f_-^n}{f_+^{n+1}} \left[ (1+(-1)^n) \tilde{q}\left(\frac{f_+^{n+1}x}{2f_-^n}\right) + (1-(-1)^n) \tilde{q}\left(\frac{-f_+^{n+1}x}{2f_-^n}\right) \right]$$

$$f_{\pm} = \sqrt{1+4c} \pm 1$$

$$c = M^2/4P_z^2$$

# $\mathcal{O}(\Lambda_{QCD}^2/(P^z)^2)$ Corrections

- Twist-4:

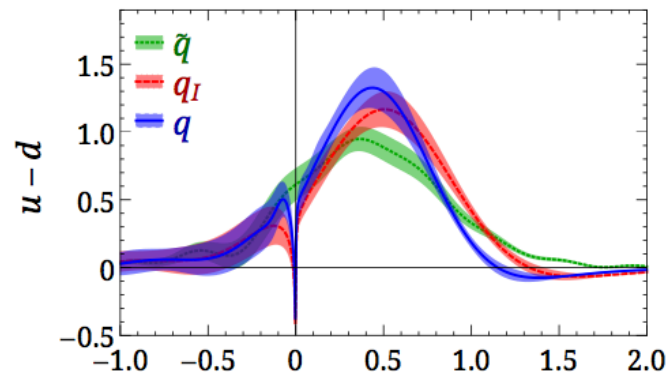
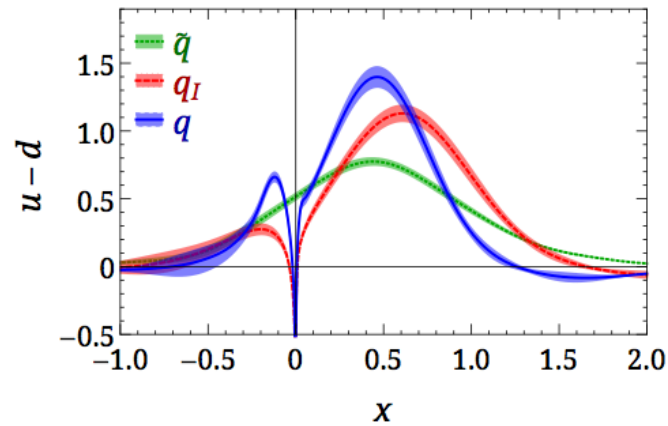
$$q_{tr}(x, \mu^2, P^z) = \frac{\lambda^2}{8\pi} \int_{-\infty}^{\infty} dz \int_0^1 \frac{dt}{t} e^{i\frac{zk^z}{t}} \langle P | \tilde{\mathcal{O}}_{tr}(z) | P \rangle$$

$$\begin{aligned} \tilde{\mathcal{O}}_{tr}(z) = & \int_0^z dz_1 \bar{\psi}(0) [\gamma^\nu \Gamma(0, z_1) D_\nu \Gamma(z_1, z) \\ & + \int_0^{z_1} dz_2 \lambda \cdot \gamma \Gamma(0, z_2) D^\nu \Gamma(z_2, z_1) D_\nu \Gamma(z_1, z)] \psi(z\lambda) \end{aligned}$$



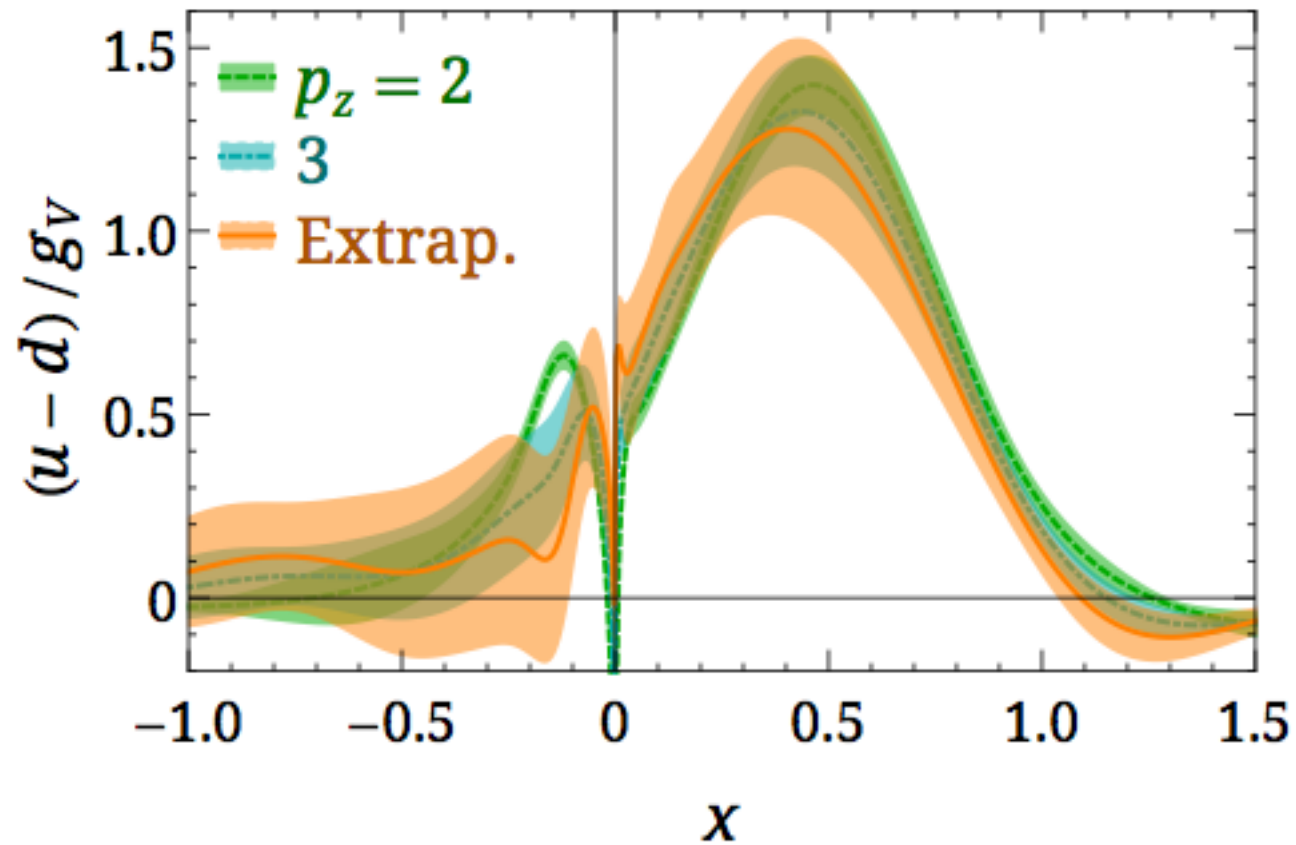
Parameterized

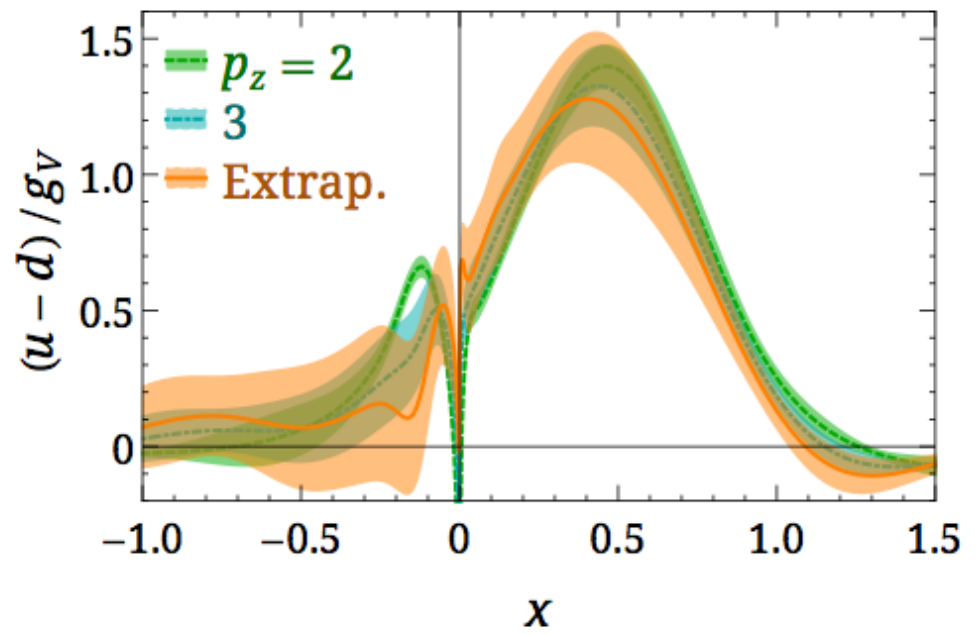
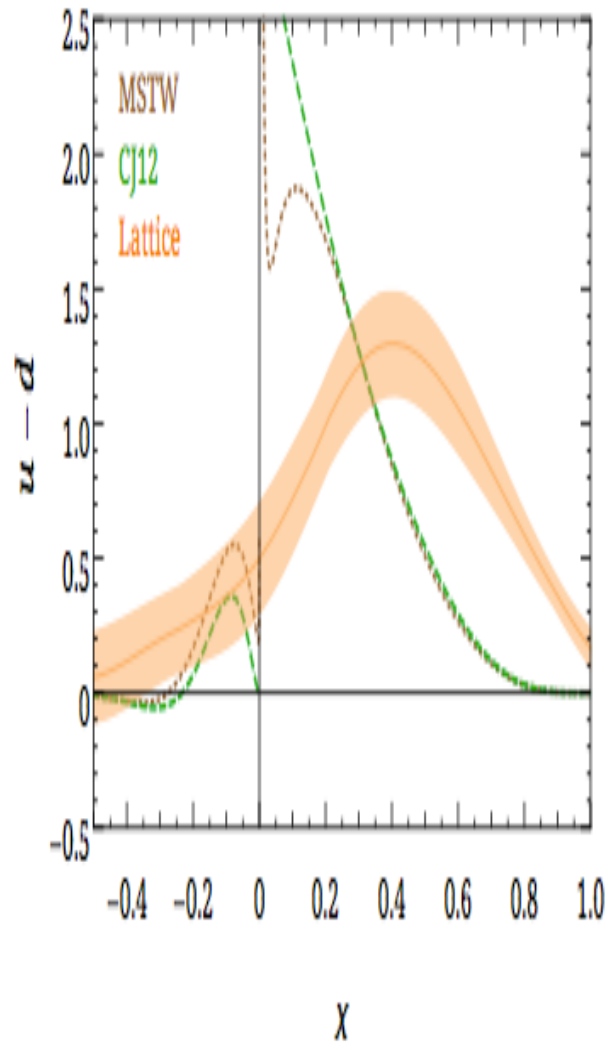
Quasi-PDF (green) w/ loop (red) w/ loop + mass (blue)



$$P^z = \frac{2\pi}{L} n = n \times 0.43 \text{ GeV} \quad n = 2 \text{ (upper) \& } 3$$

# Unpolarized Isovector Proton PDF

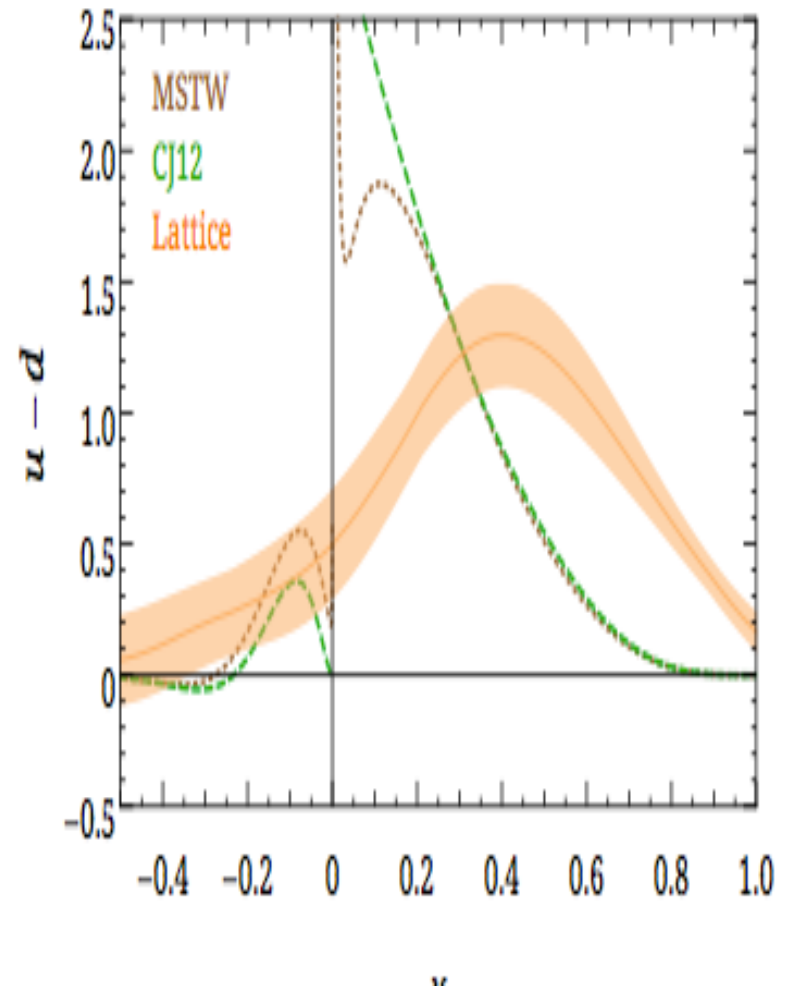
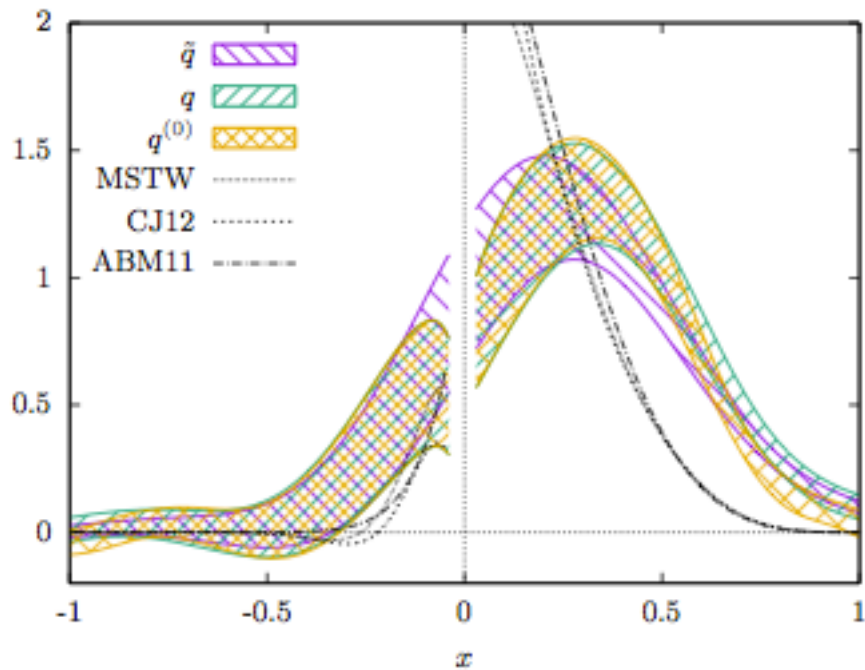




Quark mass effect!

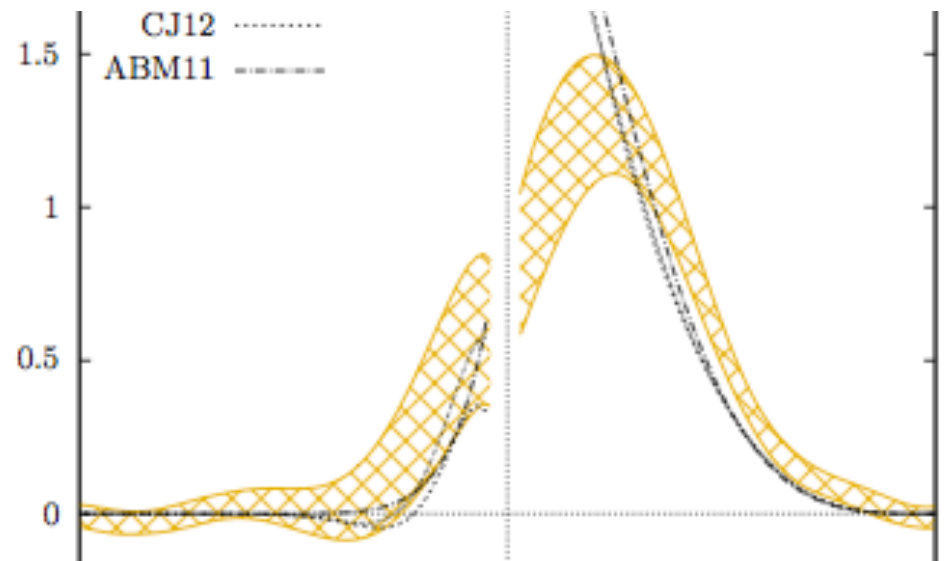
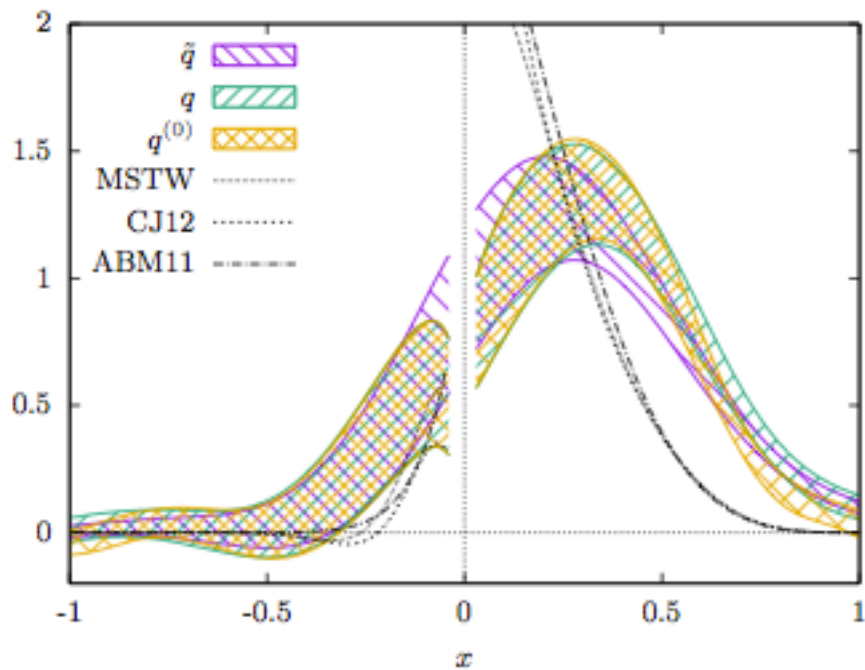
# A follow-up work

(Alexandrou et. al.1504.07455)



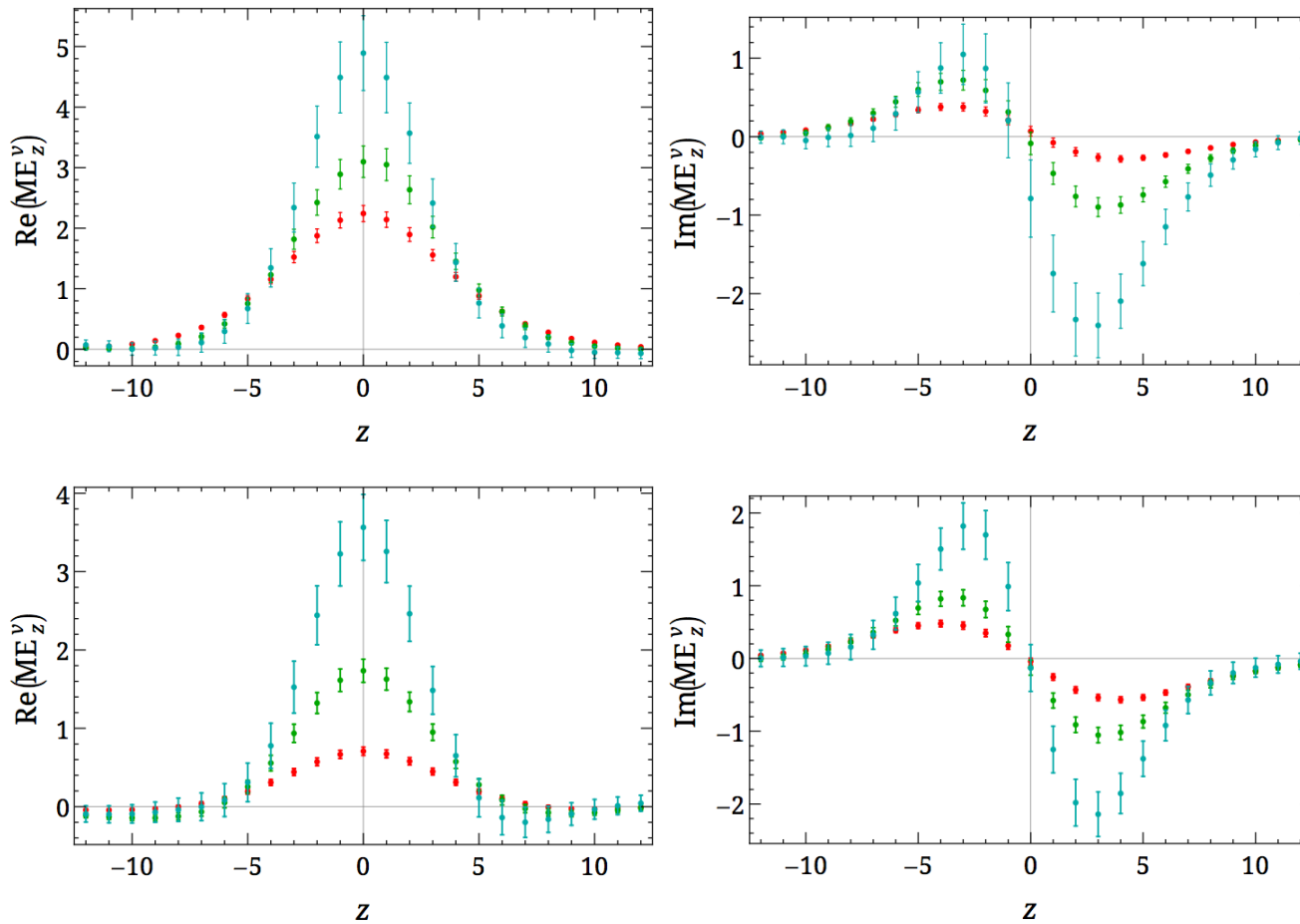
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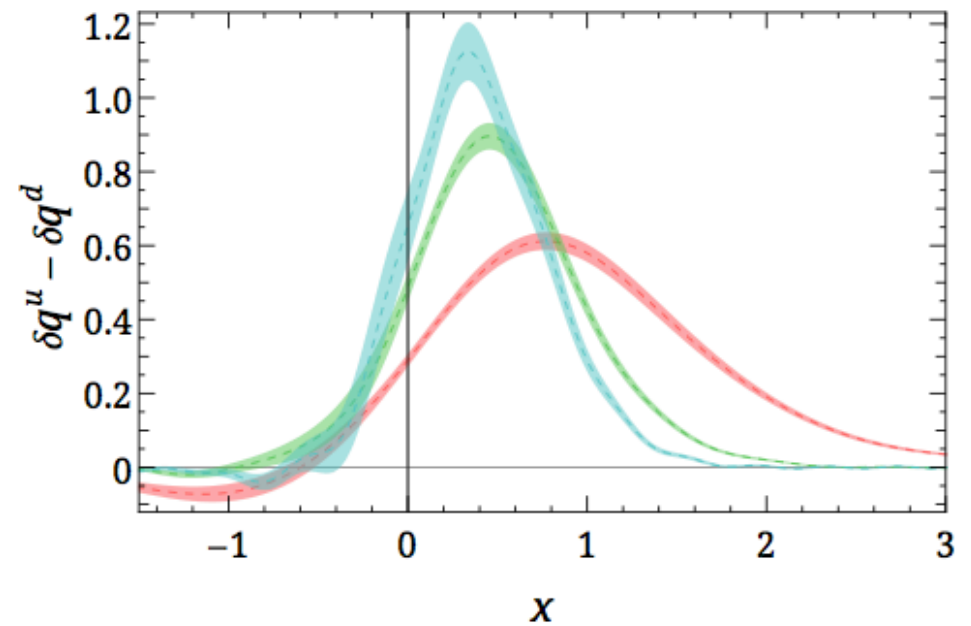
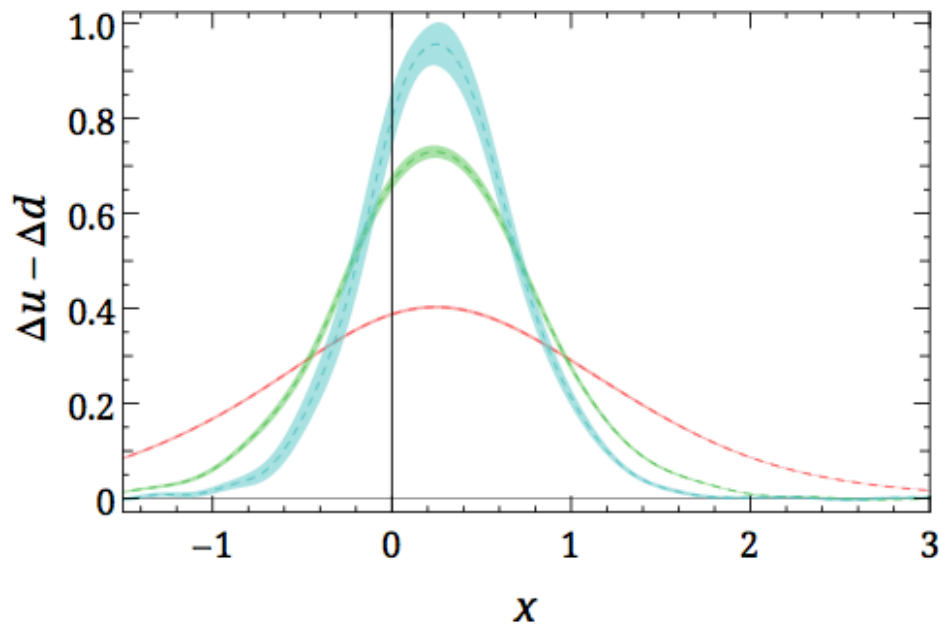




# Helicity and Transversity (isovector)

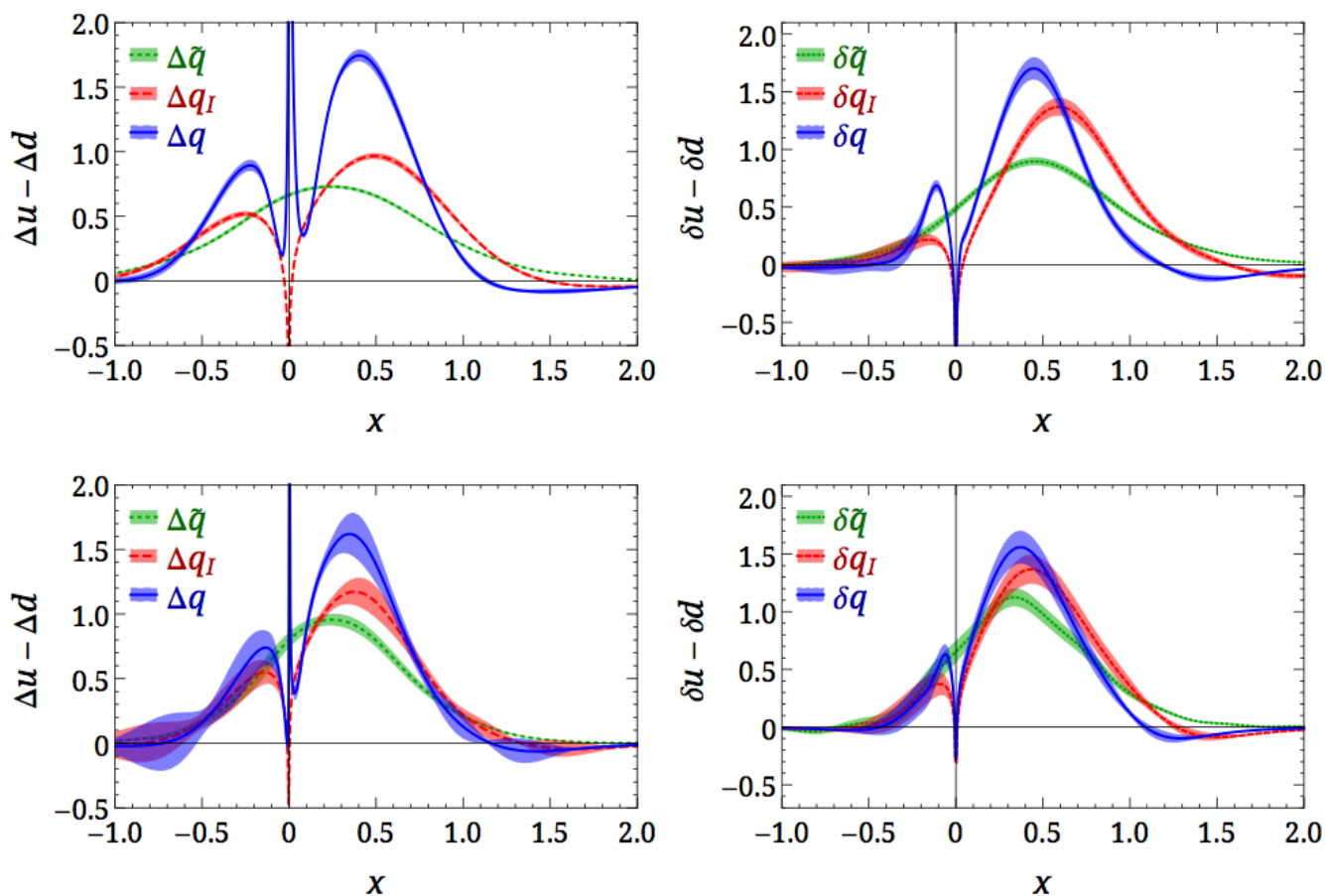


# Quasi-PDF (Helicity and Transversity)



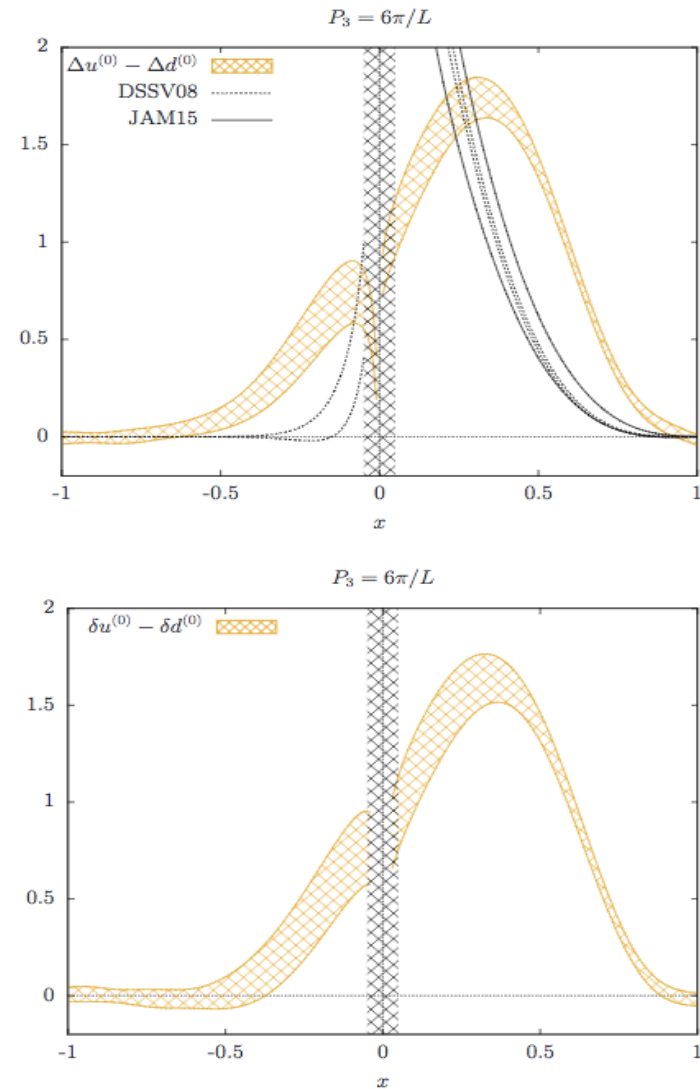
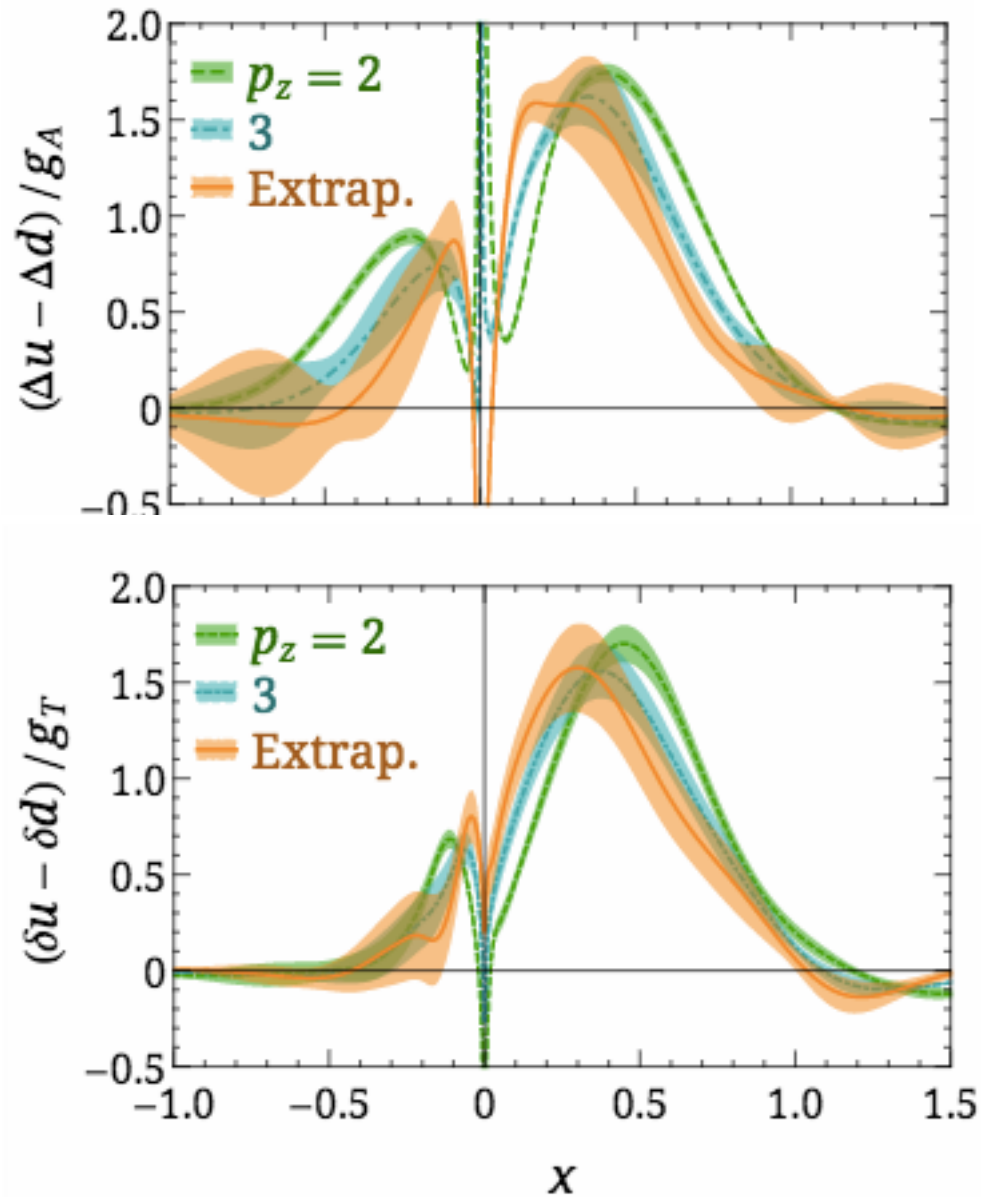
$$P^z = \frac{2\pi}{L} n = n \times 0.43 \text{ GeV} \quad n = 1, 2, 3.$$

Quasi-PDF (green) w/ loop (red) w/ loop + mass (blue)



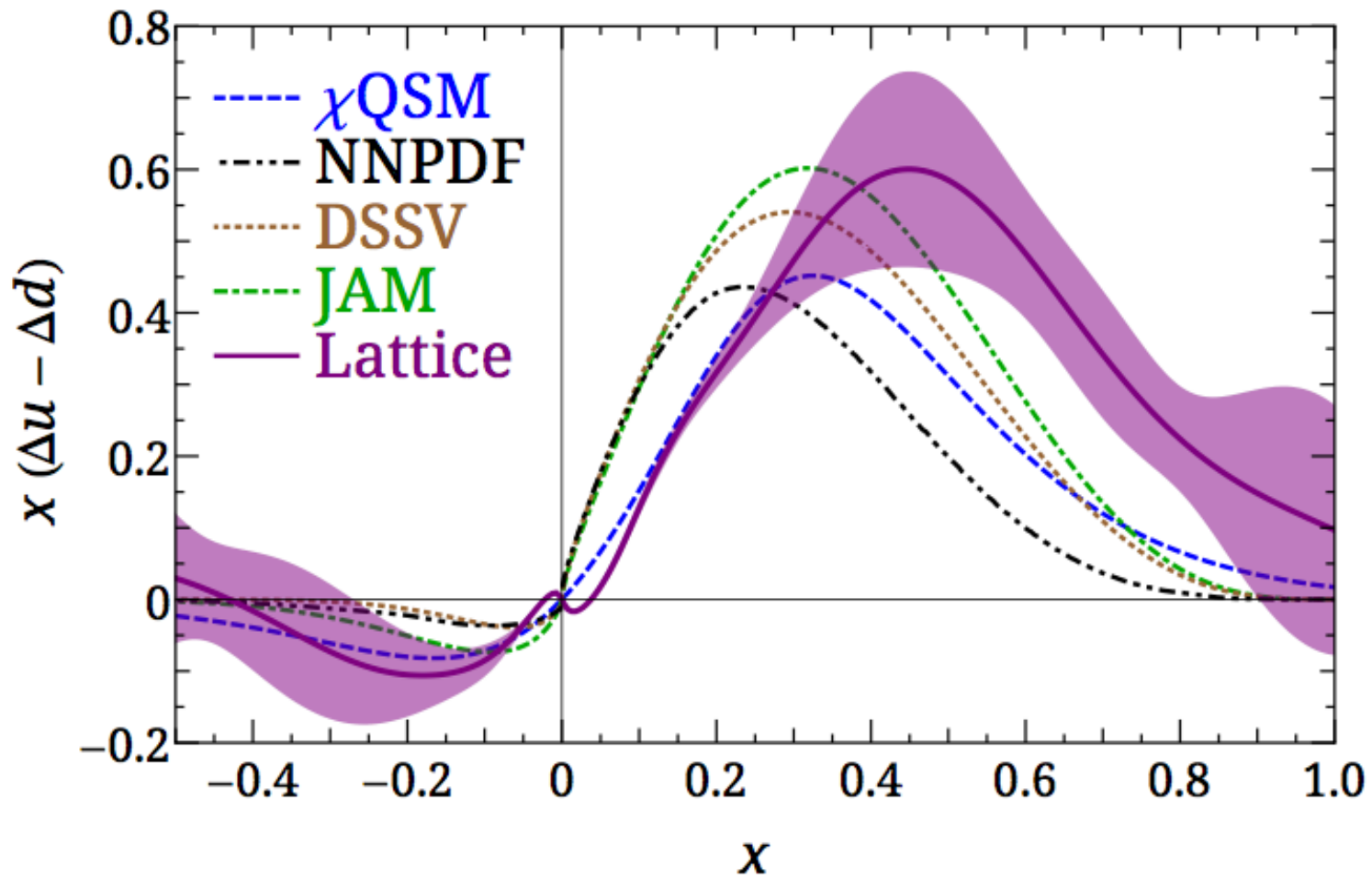
$$P^z = \frac{2\pi}{L} n = n \times 0.43 \text{ GeV} \quad n = 2 \text{ (upper) \& } 3$$

# Isvector Proton Helicity and Transversity

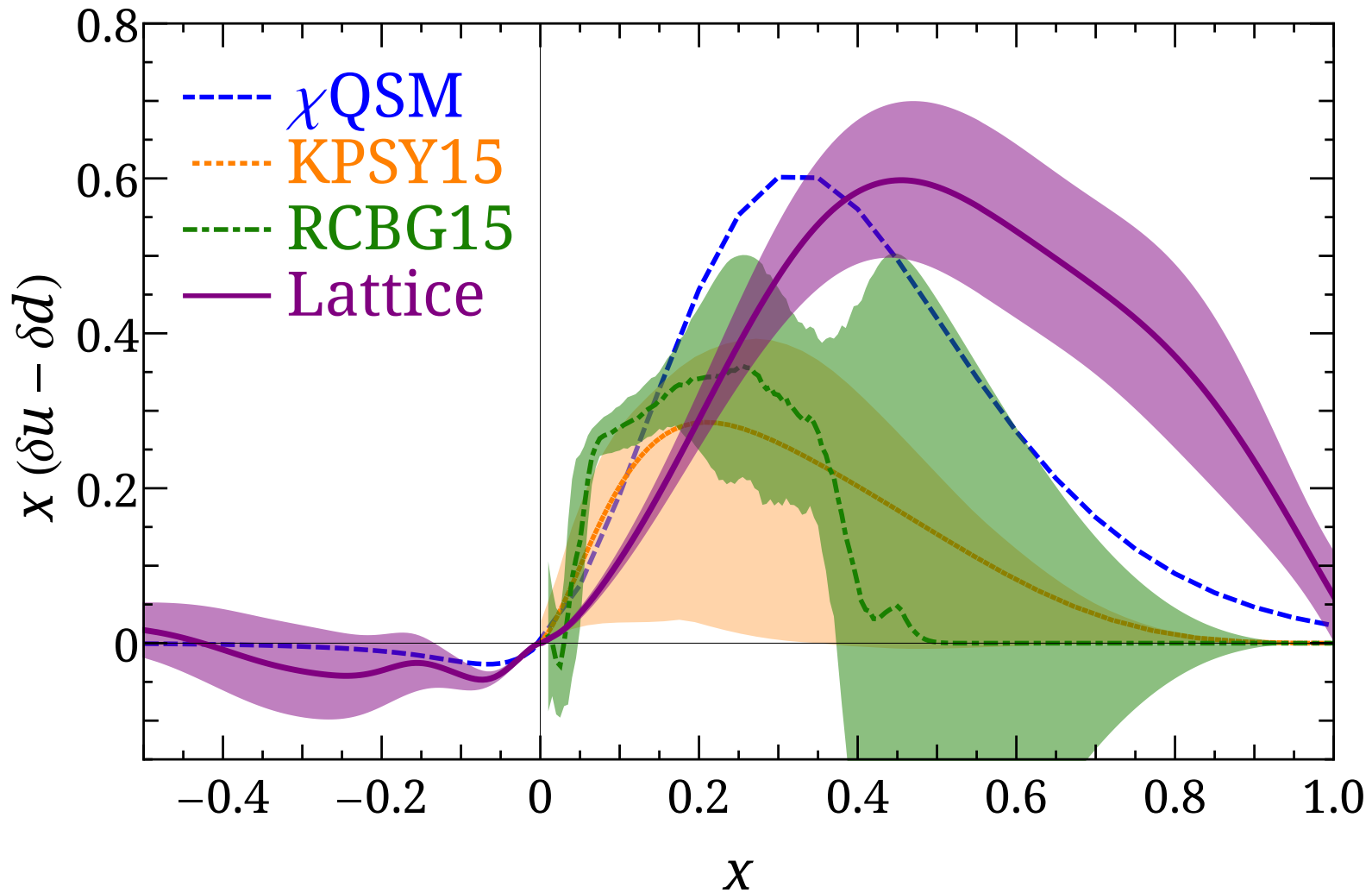


(Alexandrou et al.,  
1609.00172)

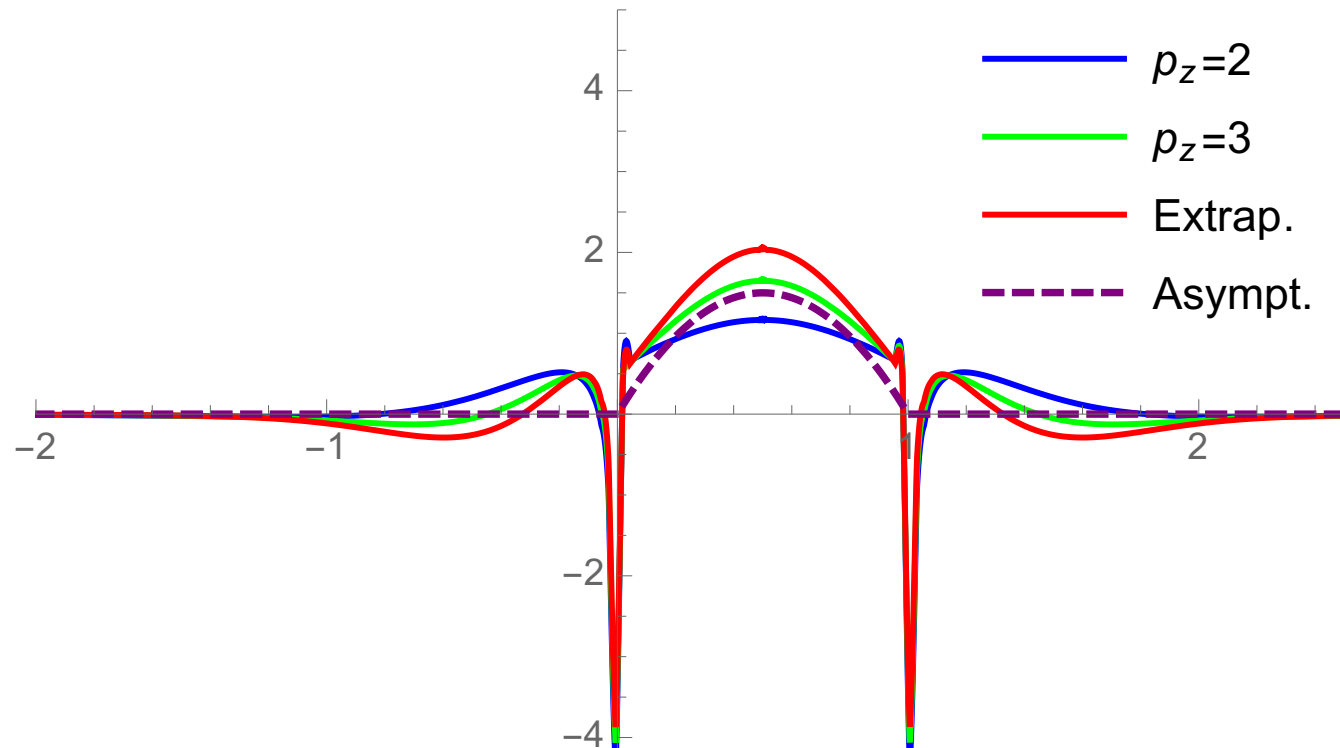
# Isovector Proton Helicity



# Isovector Proton Transversity



# Pion Light Cone Wave Function (preliminary)



$$\begin{aligned}\phi_{\pi}(x, \mu) + 3\phi_{\eta}(x, \mu) &= 2[\phi_{K^+}(x, \mu) + \phi_{K^-}(x, \mu)] \\ &= 2[\phi_{K^0}(x, \mu) + \phi_{\bar{K}^0}(x, \mu)],\end{aligned}$$

JWC, Iain W. Stewart, Phys.Rev.Lett. 92 (2004) 202001

# Outlook

- Wee partons, a good test to this approach---  
smaller quark mass
- Next: linear divergence in the matching kernel,  
lattice perturbation theory, proof of  
factorization
- If it works, lots of things to do: LCW, PDF,  
GPD, TMD ...



# Backup slides

# Linear Divergence in Matching

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