

# Particle Physics

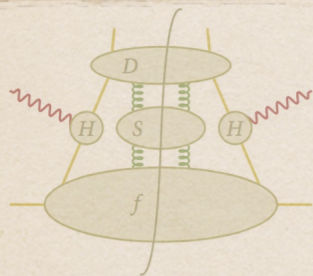


Figure 8.11: Factorisation in SIDIS: the bull diagram. All IR divergences are absorbed in the soft factor S, that hence only interacts with the TMD and FF. Note that there is no real radiation coming from the hard vertex.

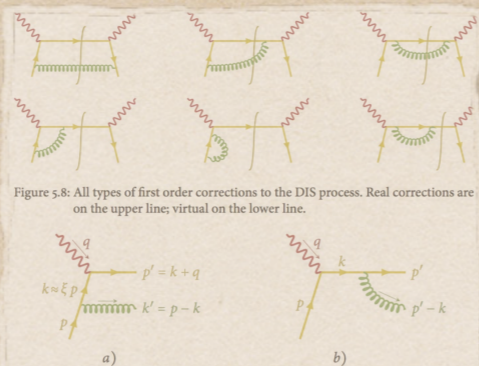


Figure 5.8: All types of first order corrections to the DIS process. Real corrections are on the upper line; virtual on the lower line.

Figure 5.9: a) Initial state gluon radiation. b) Final state gluon radiation.

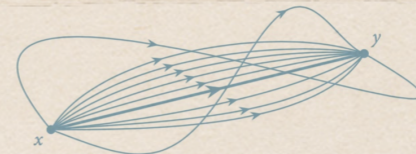
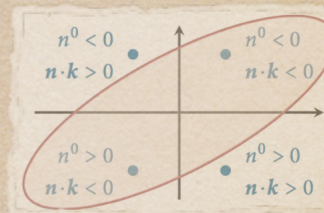


Figure 2.1: As a parallel transporter transforms in function of its path endpoints only, all paths shown will give rise to equivalent  $U_{(y,x)}$ 's, shifting a field at x to a field at y.



The Mikowskian loop integrals are then the same as the Euclidean ones, up to a possible sign difference:

$$\int \frac{d^\omega k}{(2\pi)^\omega} \frac{1}{(k^2 - \Delta)^n} = i \frac{(-)^n}{(4\pi)^{\frac{\omega}{2}}} \frac{\Gamma(n - \frac{\omega}{2})}{\Gamma(n)} \Delta^{\frac{\omega}{2} - n}, \quad (B.25a)$$

$$\left( \begin{array}{l} d \geq 2n \\ d \text{ even} \end{array} \right) = i \frac{\Delta^{\frac{d}{2} - n}}{(4\pi)^{\frac{d}{2}}} \frac{(-)^{\frac{d}{2}}}{(n-1)! (\frac{d}{2} - n)!} \left( \frac{1}{\epsilon} - \gamma_E + \sum_j \frac{1}{j} + \ln 4\pi - \ln \Delta \right),$$

$$\int \frac{d^\omega k}{(2\pi)^\omega} \frac{k^2}{(k^2 - \Delta)^n} = i \frac{(-)^{n+1}}{(4\pi)^{\frac{\omega}{2}}} \frac{\omega \Gamma(n - \frac{\omega}{2} - 1)}{2 \Gamma(n)} \Delta^{\frac{\omega}{2} + 1 - n}, \quad (B.25b)$$

$$\left( \begin{array}{l} d \geq 2n - 2 \\ d \text{ even} \end{array} \right) = i \frac{\Delta^{\frac{d}{2} + 1 - n}}{(4\pi)^{\frac{d}{2}}} \frac{\omega}{2} \frac{(-)^{\frac{d}{2}}}{(n-1)! (\frac{d}{2} + 1 - n)!} \left( \frac{1}{\epsilon} - \gamma_E + \sum_j \frac{1}{j} + \ln 4\pi - \ln \Delta \right)$$

$$\int \frac{d^\omega k}{(2\pi)^\omega} \frac{k^4}{(k^2 - \Delta)^n} = i \frac{(-)^n}{(4\pi)^{\frac{\omega}{2}}} \frac{\omega(\omega+2) \Gamma(n - \frac{\omega}{2} - 2)}{4 \Gamma(n)} \Delta^{\frac{\omega}{2} + 2 - n}, \quad (B.25c)$$

$$\left( \begin{array}{l} d \geq 2n - 4 \\ d \text{ even} \end{array} \right) = i \frac{\Delta^{\frac{d}{2} + 2 - n}}{(4\pi)^{\frac{d}{2}}} \frac{\omega(\omega+2)}{4} \frac{(-)^{\frac{d}{2}}}{(n-1)! (\frac{d}{2} + 2 - n)!} \left( \frac{1}{\epsilon} - \gamma_E + \sum_j \frac{1}{j} + \ln 4\pi - \ln \Delta \right)$$

We list some other common Minkowskian integrals:

$$\int \frac{d^\omega k}{(2\pi)^\omega} \ln(k^2 - a) = -\frac{i}{(4\pi)^{\frac{\omega}{2}}} \Gamma\left(-\frac{\omega}{2}\right) a^{\frac{\omega}{2}}, \quad (B.26a)$$

$$\int \frac{d^\omega k}{(2\pi)^\omega} e^{ak^2 - ib \cdot k} = \frac{i}{(4\pi)^{\frac{\omega}{2}}} a^{-\frac{\omega}{2}} e^{\frac{b^2}{4a}}, \quad (B.26b)$$

$$\int \frac{d^\omega k}{(2\pi)^\omega} \frac{1}{(-k^2)^\alpha} e^{-ib \cdot k} = \frac{i}{4^\alpha \pi^{\frac{\omega}{2}}} \frac{\Gamma(\frac{\omega}{2} - \alpha)}{\Gamma(\alpha)} \frac{1}{(-b^2)^{\frac{\omega}{2} - \alpha}}. \quad (B.26c)$$

## IV. Lagrangian

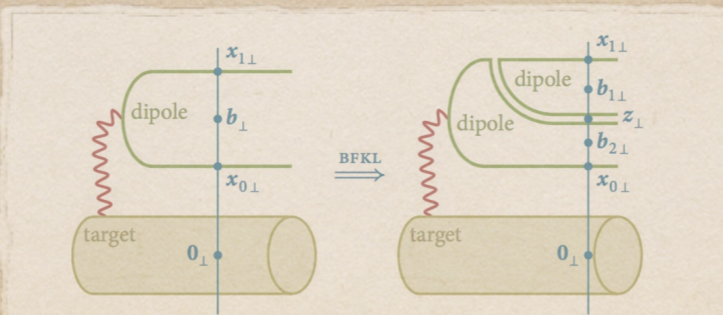
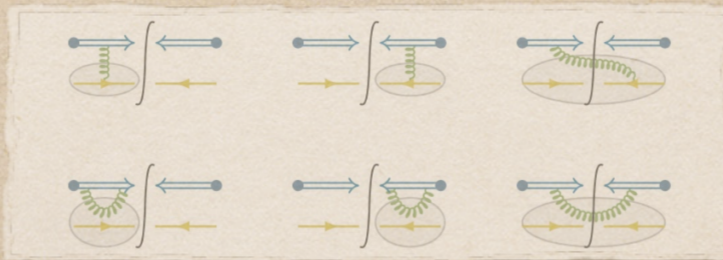


Figure 9.6: In the dipole picture, the BFKL evolution is an evolution in dipoles, i.e. new dipoles are created during the evolution. A gluon that is radiated from the dipole can be represented as two fundamental lines (see Equation 10.13). This essentially splits the dipole in two at the point  $z_{1\perp}$ , as is illustrated in the second diagram.

$$\oint_C dx \cdot A = \int_\Sigma d\sigma \cdot (\partial \wedge A)$$



$$\begin{aligned} \text{tr}(t^a t^x t^b t^x) &= -\frac{1}{4N_c} \delta^{ab}, \\ \text{tr}(t^b t^x t^y) f^{axy} &= -i \frac{N_c}{4} \delta^{ab}, \\ \text{tr}(t^y t^z) f^{axy} f^{bzx} &= -\frac{N_c}{2} \delta^{ab}, \\ f^{xay} f^{ycz} f^{zbw} f^{wcx} &= \frac{N_c^2}{2} \delta^{ab}, \\ f^{avw} f^{xbv} f^{yvw} f^{zvx} &= \frac{N_c^2}{2} \delta^{ab}, \\ f^{awv} f^{bwz} f^{xzy} f^{yvx} &= N_c^2 \delta^{ab}, \\ f^{xay} f^{ycz} f^{zbw} f^{wcx} &= \frac{N_c^2}{2} \delta^{ab}, \\ f^{vaw} f^{wbz} f^{xzy} f^{yvx} &= N_c^2 \delta^{ab}, \end{aligned}$$

and similarly for the seven remaining diagrams.

# Goals

- ¿ Lagrangian ?
- ¿ Symmetries ?
- ¿ Feynman diagrams ?
- ¿ Standard Model ?

# Goals

**This is not a regular lecture**

**→ feel free to interrupt  
& ask questions**



# Symmetries:

## Building the Lagrangian

# Basics

- basic components are **fields**
  - just mathematical tools
  - will give rise to particles
- principal quantity is the **action**, which is the integral of the Lagrangian:

$$S = \int d^4x \mathcal{L}(x, \phi, \partial\phi)$$



- all paths possible (simultaneous), but path with **least action** is favoured
- minimising action leads to **equations of motion**

# Lagrangian

- is kinetic energy minus potential energy

$$\mathcal{L} = T - V$$

- classical example: spring



$$\mathcal{L} = \frac{1}{2}m \left( \frac{dx}{dt} \right)^2 - \frac{1}{2}kx^2 \quad \Rightarrow \quad x = x_0 \cos \sqrt{\frac{k}{m}}t$$

- field example: free electron field

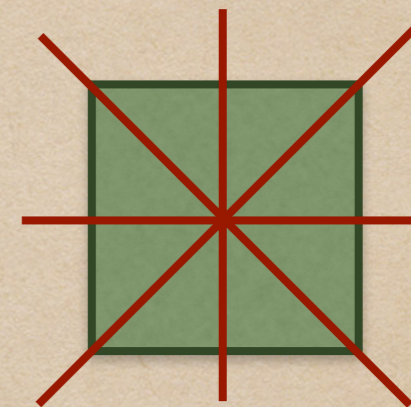
$$\mathcal{L} = i\bar{\psi}\not{\partial}\psi - m\bar{\psi}\psi \quad \Rightarrow \quad (i\not{\partial} - m)\psi = 0 \quad (\text{QM})$$

# Lagrangian

- kinetic terms are **quadratic** and have **derivatives**  $\bar{\psi}\not{\partial}\psi$ ,  $\partial_{\mu}\phi\partial^{\mu}\phi$ , ...
- potential terms are what is left
  - special type: **mass terms**:  $m\bar{\psi}\psi$ ,  $m^2|\phi|^2$ , ...  
quadratic without derivatives
  - others are **interaction terms**  $\bar{\psi}A\psi$ , ...

# Symmetries

- leave theory unchanged
- **symmetry**  $\Rightarrow$  **conservation**



Emmy Noether

- homogeneity of space
  - $\Rightarrow$  translational invariance
  - $\Rightarrow$  momentum conservation
- isotropy of space
  - $\Rightarrow$  rotational invariance
  - $\Rightarrow$  angular momentum conservation





# Symmetries

- in quantum mechanics,  $\psi$  is an amplitude
  - not physical
  - $|\psi|^2$  is probability, physical
- **phase** is undetermined, because we can scale  $\psi \rightarrow e^{i\alpha}\psi$ , then  $\bar{\psi} \rightarrow e^{-i\alpha}\bar{\psi}$ , such that  $|\psi|^2 \rightarrow |\psi|^2$   
  
 $\Rightarrow$  invariant!

# Symmetries

- similar in quantum field theories:

$$i\bar{\psi}\not{\partial}\psi \rightarrow i\bar{\psi}\not{\partial}\psi \quad m\bar{\psi}\psi \rightarrow m\bar{\psi}\psi$$

in other words  $\mathcal{L} = \mathcal{L}'$

- free electron field

=> conservation of electric charge

- BUT... phase  $\alpha$  can depend on spacetime

coordinates:  $\alpha = \alpha(x)$

$$i\bar{\psi}\not{\partial}\psi \rightarrow i\bar{\psi}\not{\partial}\psi - \bar{\psi}(\not{\partial}\alpha)\psi \quad \mathcal{L}' = \mathcal{L} - \bar{\psi}(\not{\partial}\alpha)\psi$$

=> no longer invariant!

# Lagrangian

- add term  $g\bar{\psi}A\psi$  to the Lagrangian
  - with property  $A_\mu \rightarrow A_\mu + 1/g \partial_\mu a$
  - because then  $g\bar{\psi}A\psi \rightarrow g\bar{\psi}A\psi + \bar{\psi}(\not{\partial}a)\psi$
  - invariant !
- but new field also needs kinetic terms
  - symmetry  $\Rightarrow$  conserved tensor:
$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$
  - its square will be kinetic term:
$$-\frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

# QED:

From the Lagrangian  
to a full theory

# QED

Full Lagrangian for Quantum Electro Dynamics:

$$\mathcal{L}_{\text{QED}} = i\bar{\psi}\not{\partial}\psi - m\bar{\psi}\psi + g\bar{\psi}\not{A}\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

electron kinetic term

electron mass term

electron-photon interaction term

photon kinetic term

# QED

Full Lagrangian for Quantum Electro Dynamics:

$$\mathcal{L}_{\text{QED}} = i\bar{\psi}\not{\partial}\psi - m\bar{\psi}\psi + g\bar{\psi}\not{A}\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

'free' theory

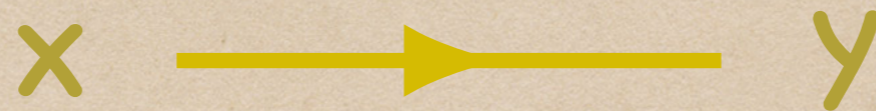
'interaction' theory

$$\mathcal{L}_{\text{QED}} = \mathcal{L}_0 + \mathcal{L}_I$$

Keep  $\mathcal{L}_0$  exact, but expand  $\mathcal{L}_I \Rightarrow$  perturbation theory

# Free Theory

Propagation of an electron from  $x$  to  $y$ :



$$\langle 0 | \psi(y) \bar{\psi}(x) | 0 \rangle = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \psi(y) \bar{\psi}(x) e^{iS_0}$$

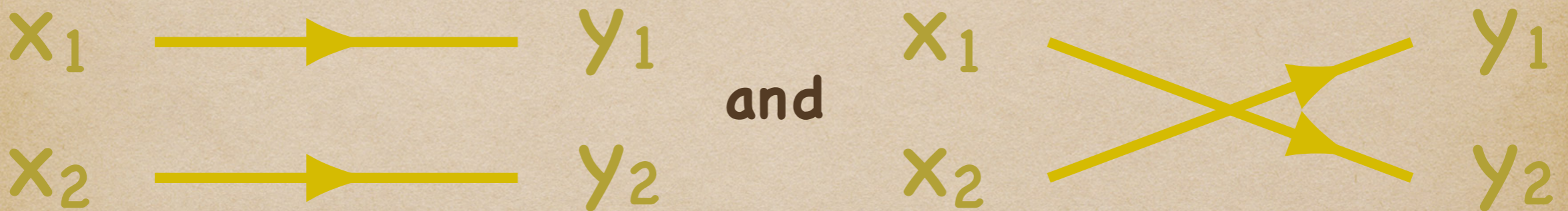
Propagation of a photon from  $x$  to  $y$ :



$$\langle 0 | A_\mu(x) A_\mu(y) | 0 \rangle = \int \mathcal{D}A_\mu A_\mu(x) A_\mu(y) e^{iS_0}$$

# Free Theory

Easily generalised to more points:



$$\langle 0 | \psi_{y_1} \psi_{y_2} \bar{\psi}_{x_1} \bar{\psi}_{x_2} | 0 \rangle = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \psi_{y_1} \psi_{y_2} \bar{\psi}_{x_1} \bar{\psi}_{x_2} e^{S_0}$$



# Interaction

- Add interaction part from action (Lagrangian) to the exponential  $e^{S_0 + S_I}$
- Equations are not solvable anymore

=> expand interaction part:

$$e^{S_I} \approx 1 + S_I + \frac{1}{2} S_I^2 + \dots$$

- Propagation of electron from  $x$  to  $y$  is now:

$$\langle \psi(y) \bar{\psi}(x) \rangle = \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \mathcal{D}A_\mu \psi(y) \bar{\psi}(x) e^{S_0} \left( 1 + S_I + \frac{1}{2} S_I^2 + \dots \right)$$

# Interaction

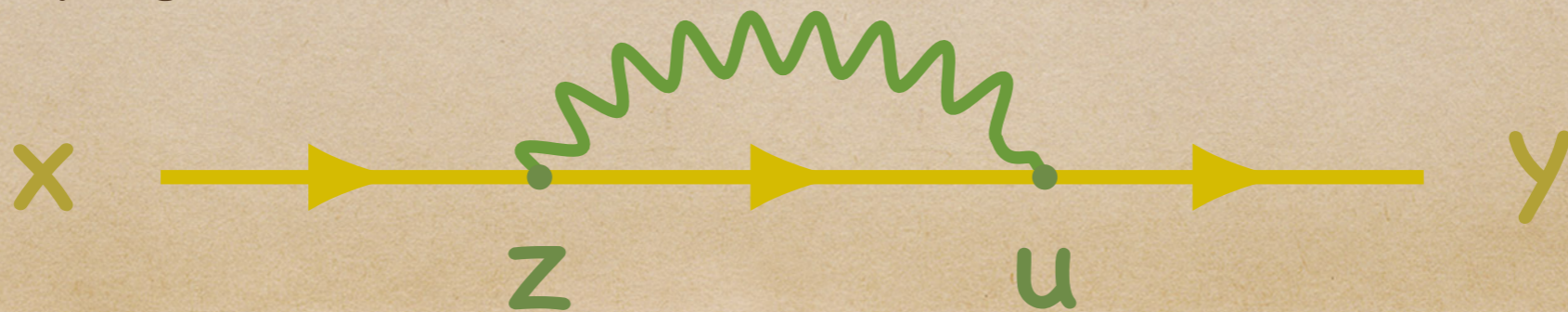
- Propagation of electron from  $x$  to  $y$  is now:

$$\langle \Psi(y) \bar{\Psi}(x) \rangle = \int \mathcal{D}\bar{\Psi} \mathcal{D}\Psi \mathcal{D}A_\mu \Psi(y) \bar{\Psi}(x) e^{S_0} \left( 1 + S_I + \frac{1}{2} S_I^2 + \dots \right)$$

- Take the second order as example:

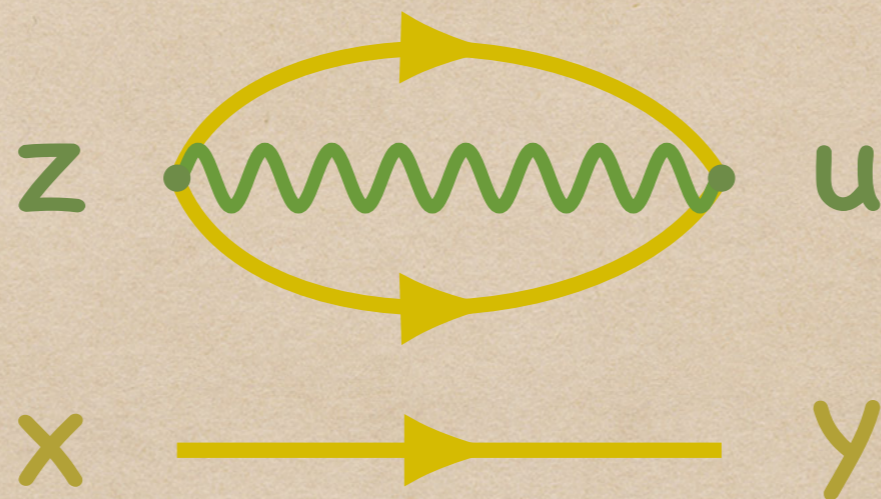
$$S_I^2 = g^2 \int dz du (\bar{\Psi} \not{A} \Psi)_z (\bar{\Psi} \not{A} \Psi)_u$$

- So we have electron propagation from  $x$  to  $z$ , from  $z$  to  $u$ , and from  $u$  to  $y$ . We also have photon propagation from  $z$  to  $u$ . Schematically:



# Interaction

- Other possibility:



- Vacuum diagrams

=> These are unwanted and have to be cancelled:

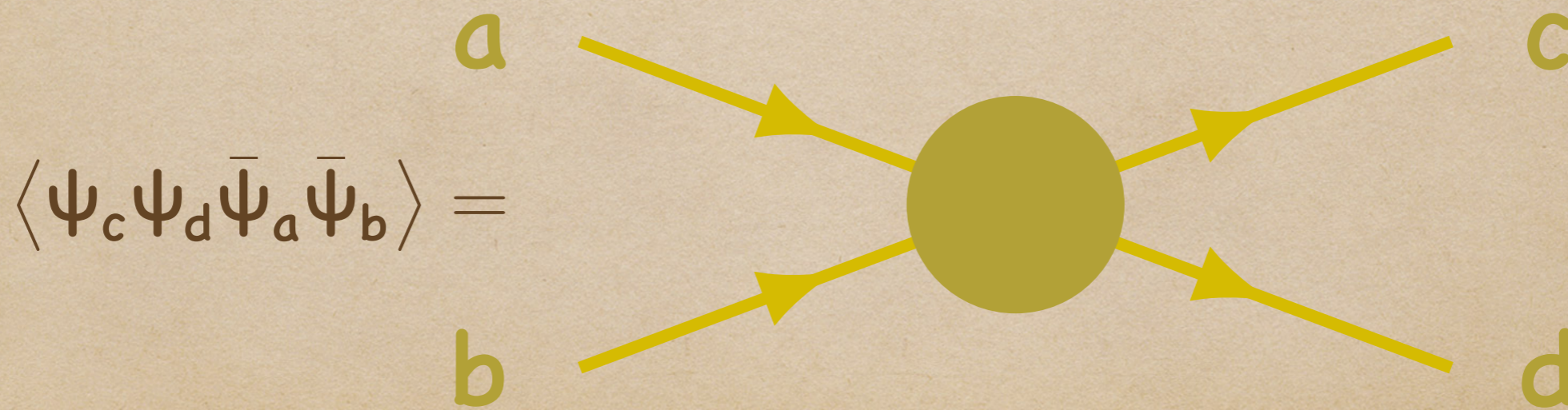
$$\langle \psi(y) \bar{\psi}(x) \rangle = \frac{\int \mathcal{D}\bar{\psi} \mathcal{D}\psi \psi(y) \bar{\psi}(x) e^{\mathcal{S}}}{\int \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{\mathcal{S}}}$$

# Interaction

- So the full propagator is:



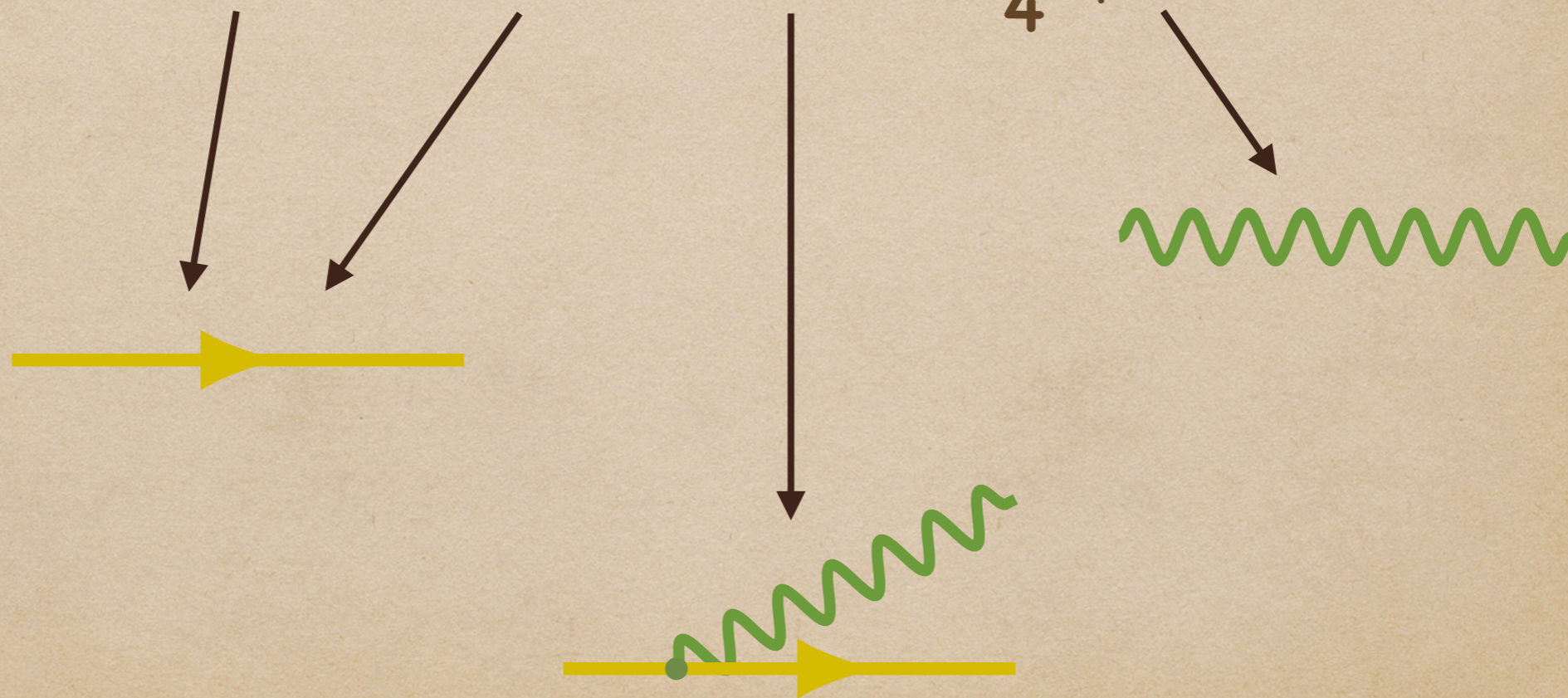
- This can be extended to any number of particles, i.e. electron-electron collision:



# Feynman Rules

Full Lagrangian for Quantum Electro Dynamics:

$$\mathcal{L}_{\text{QED}} = i\bar{\psi}\not{\partial}\psi - m\bar{\psi}\psi + g\bar{\psi}\not{A}\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$





**SM:**


**The Standard Model  
of particle physics**

# Lagrangian

- We can easily extend our theory by adding new parts to our Lagrangian:

$$\mathcal{L} = \mathcal{L}_{\text{QED}} + \mathcal{L}_{\text{QFD}} + \mathcal{L}_{\text{QCD}} + \dots$$

- QFD (weak force) and QCD (strong force) are very similar to QED. They only add two different types of interaction terms:

$$A_\mu A_\nu \partial^\mu A^\nu$$


$$A_\mu A_\nu A^\mu A^\nu$$

