

March 1, 2017 in Toyama

ILC, Project Status and Physics with Focus on Electroweak Symmetry Breaking

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KEK

arXiv: 1506.05992 (ILC Physics Case)

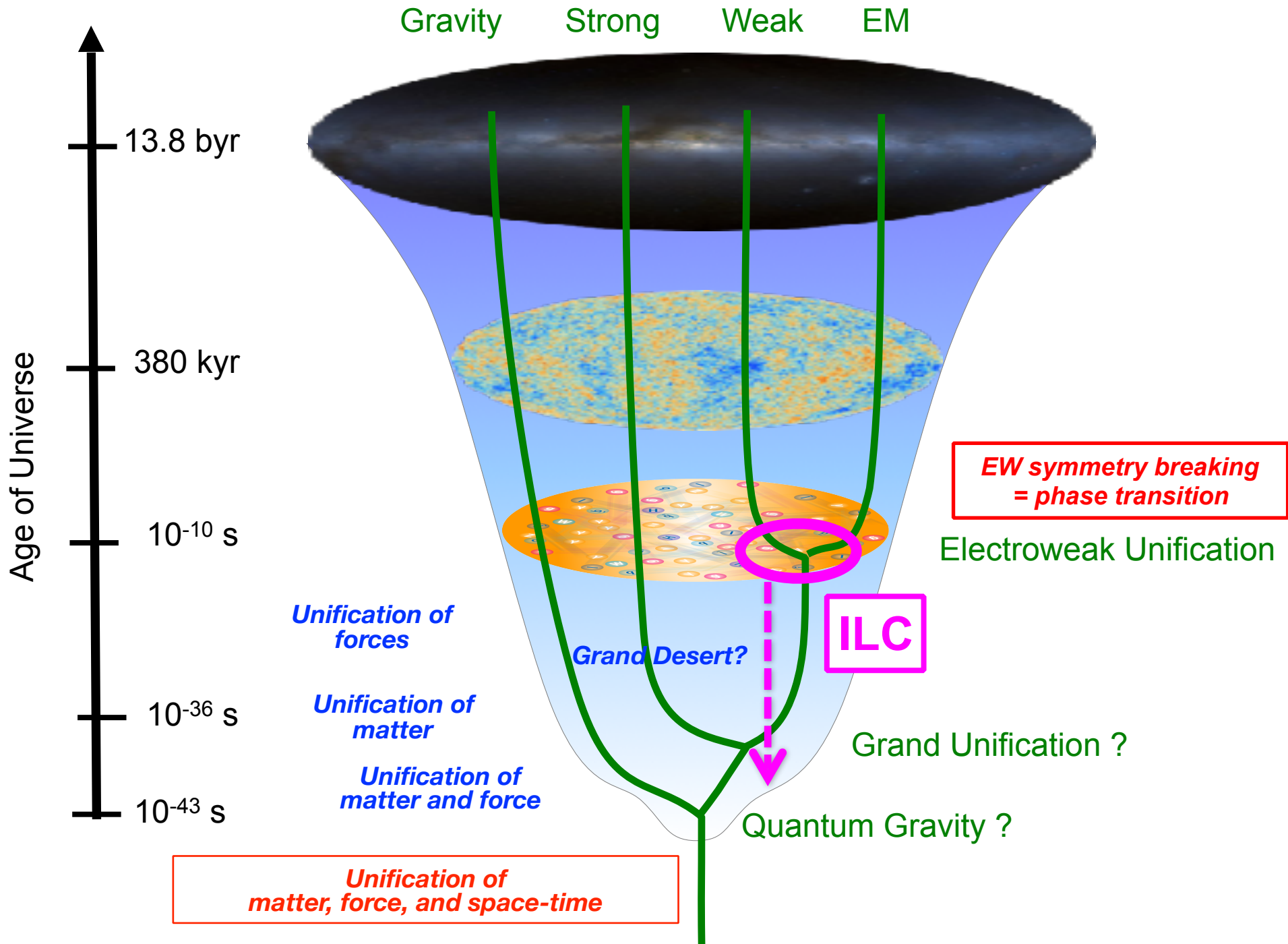
arXiv: 1506.07830 (ILC Run Scenarios)

arXiv: 1306.6352 (ILC TDR: Physics)

EPJC (2015) 75:371 (LC Physics)

arXiv: 1702.05333 (ILC New Particles)

Towards ultimate unification

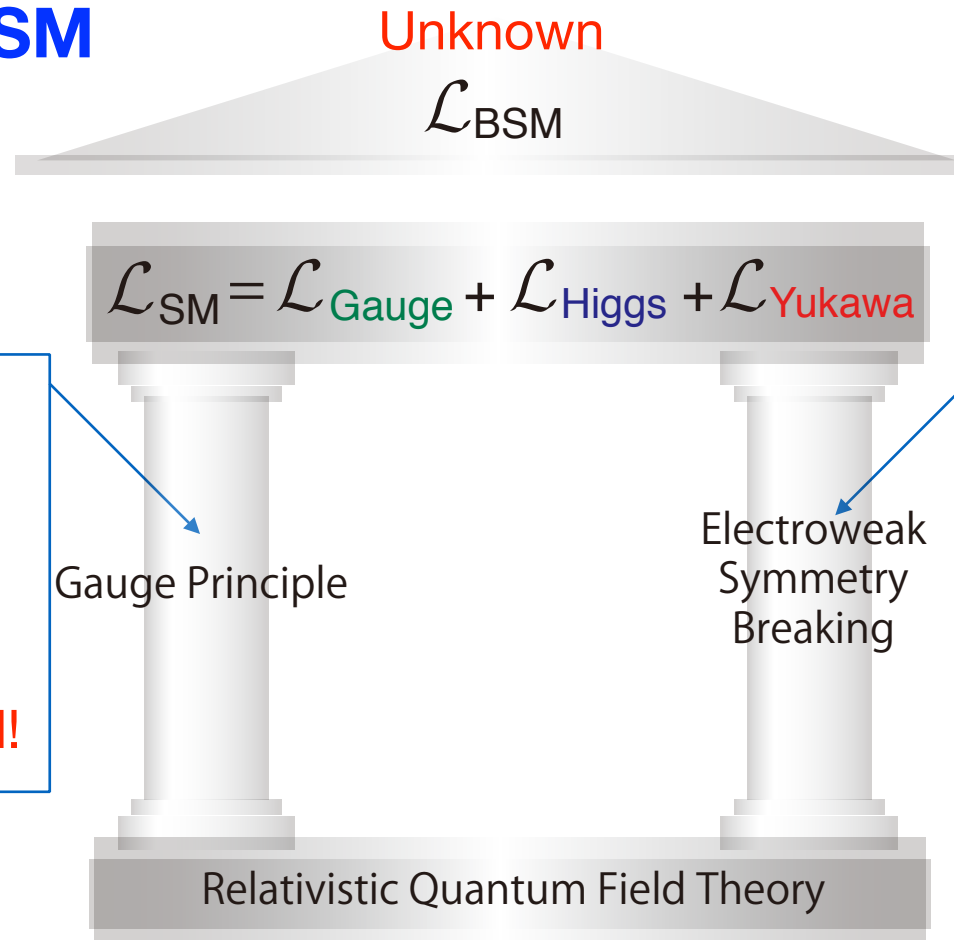


**Why is the EW scale
so important ?**

Why is the EW scale so important?

Mystery of something in the vacuum

2 Pillars of SM



Success of SM
= success of
gauge theory
(left pillar)

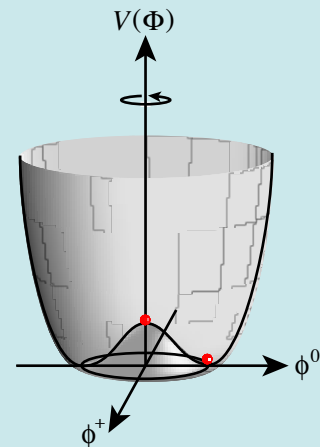
Precisely tested!

Vacuum filled with weak
charge (evidence: H125)

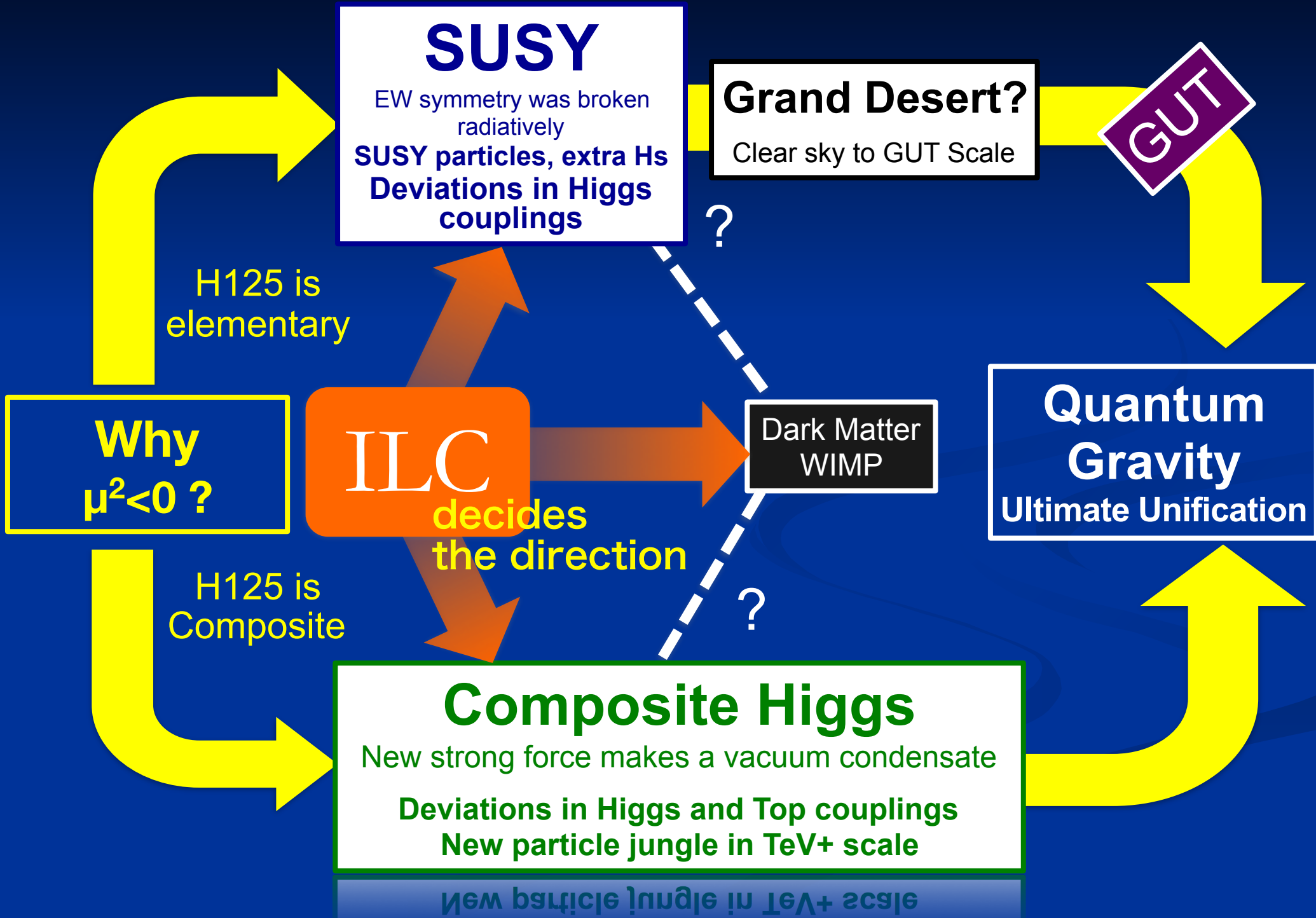
The nature of the
Higgs field - its
multiplet structure &
dynamics behind it -
is all unknown!

The SM does not explain **why the Higgs field developed a vacuum expectation value** (*Why $\mu^2 < 0$?*)! The answer forks depending on whether **H125 is elementary or composite!**

$$V(\phi) = \mu^2 |\phi|^2 + \lambda |\phi|^4$$



Big Branching Point at the EW Scale

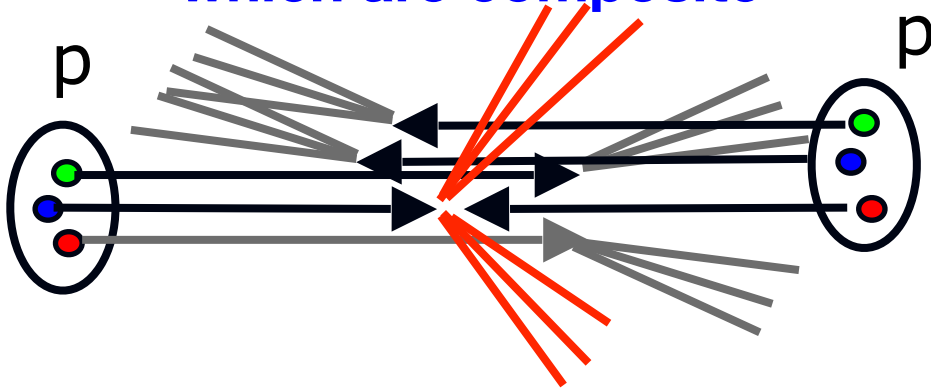


**The 3 major probes
for BSM at ILC:**

Higgs, Top, and
search for
New Particles

3 Powerful Tools

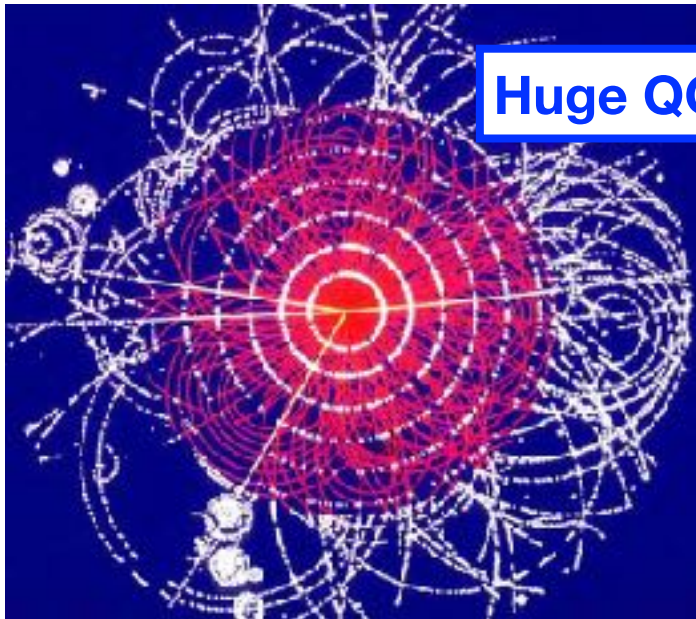
**LHC: Collision of protons
which are composite**



E_{cm} 7-14 TeV

Pileup

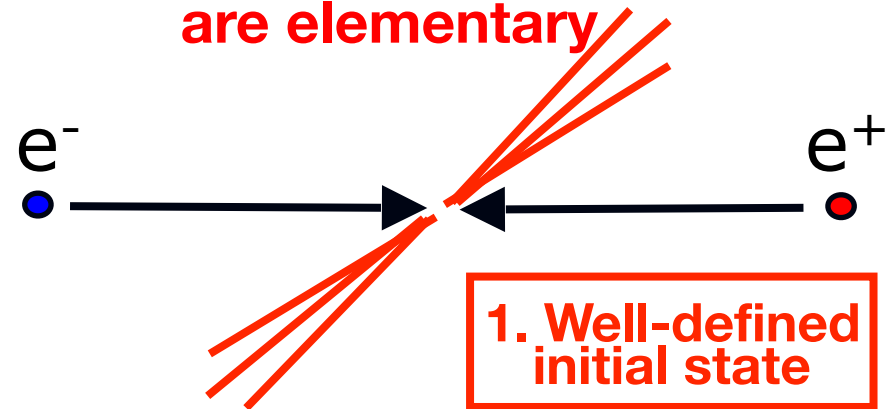
Initial state not very well defined



Huge QCD BG

proton is composite \Rightarrow events are complicated but
maximum reachable energy is high!

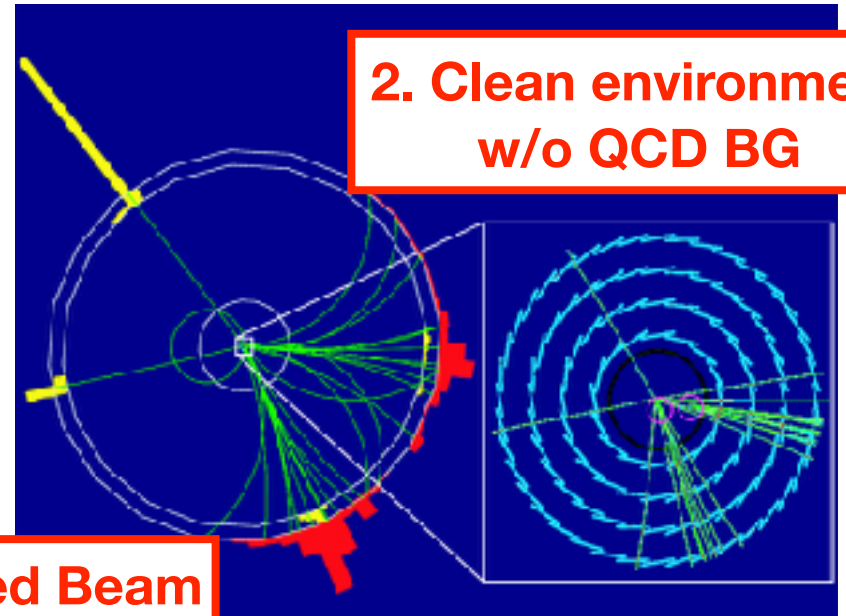
**ILC: Collision of e^+e^- which
are elementary**



E_{cm} 0.25-1 TeV

Lab. frame = CM frame

**2. Clean environment
w/o QCD BG**

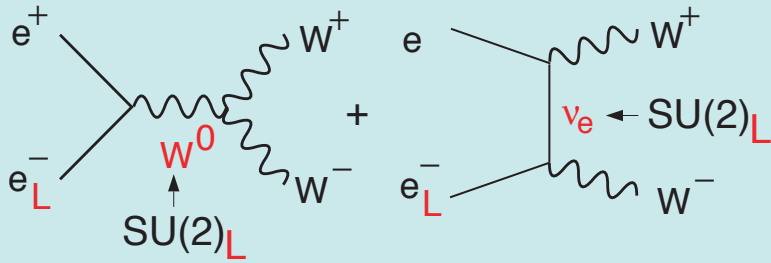


3. Polarized Beam

clean and and able to detect everything produced!

Power of Beam Polarization

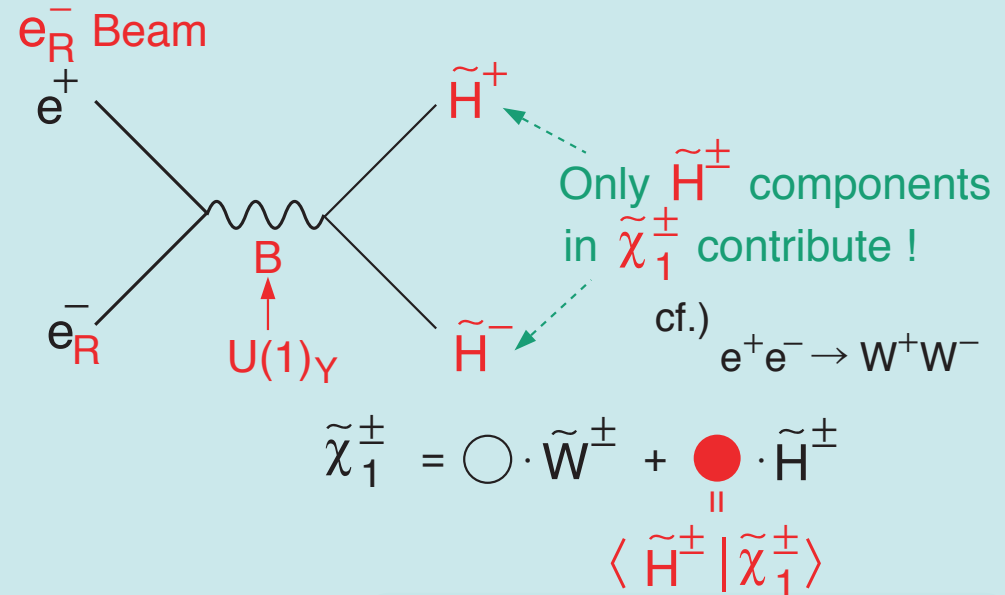
$W^+ W^-$ (Largest SM BG in SUSY searches)



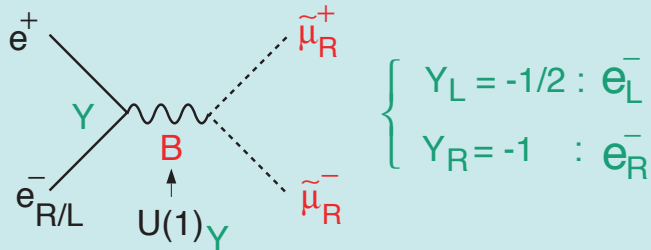
In the symmetry limit, $\sigma_{WW} \rightarrow 0$ for e_R^- !

BG Suppression

Chargino Pair



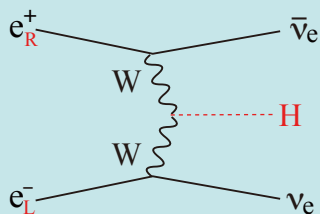
Slepton Pair



In the symmetry limit, $\sigma_R = 4 \sigma_L$!

Decomposition

WW-fusion Higgs Prod.



	ILC
Pol (e ⁻)	-0.8
Pol (e ⁺)	+0.3
$(\sigma/\sigma_0)_{\nu H}$	$1.8 \times 1.3 = 2.34$

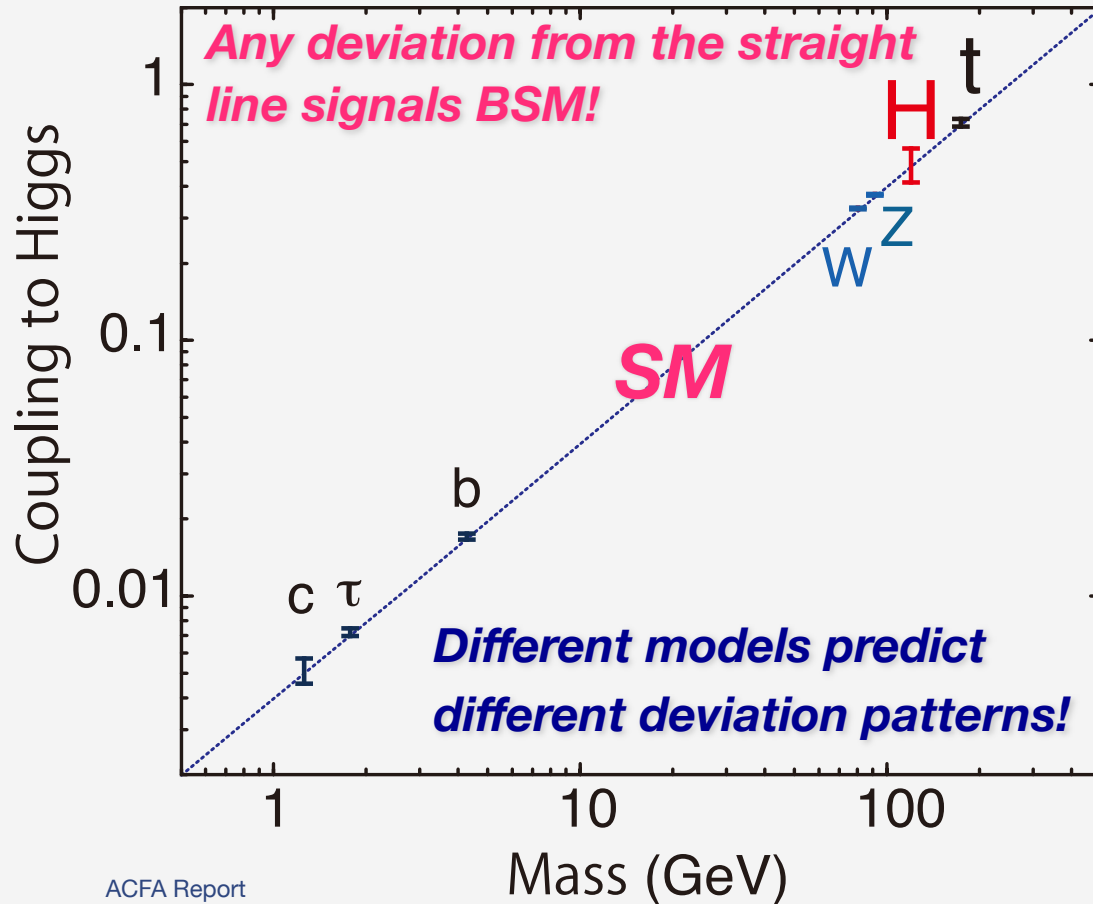
Signal Enhancement

Among three probes, today, we will focus on

Higgs

Deviation in Higgs Couplings

Mass-coupling relation



The size of the deviation depends on the new physics scale (Λ)!

Decoupling Theorem:
 $\Lambda \uparrow \rightarrow SM$

example 1: **Minimal SUSY**

(MSSM : $\tan\beta=5$, radiative correction factor ≈ 1)

$$\frac{g_{hbb}}{g_{h_{SM}bb}} = \frac{g_{h\tau\tau}}{g_{h_{SM}\tau\tau}} \simeq 1 + 1.7\% \left(\frac{1 \text{ TeV}}{m_A} \right)^2$$

heavy Higgs mass

example 2: **Minimal Composite Higgs Model**

$$\frac{g_{hVV}}{g_{h_{SM}VV}} \simeq 1 - 8.3\% \left(\frac{1 \text{ TeV}}{f} \right)^2$$

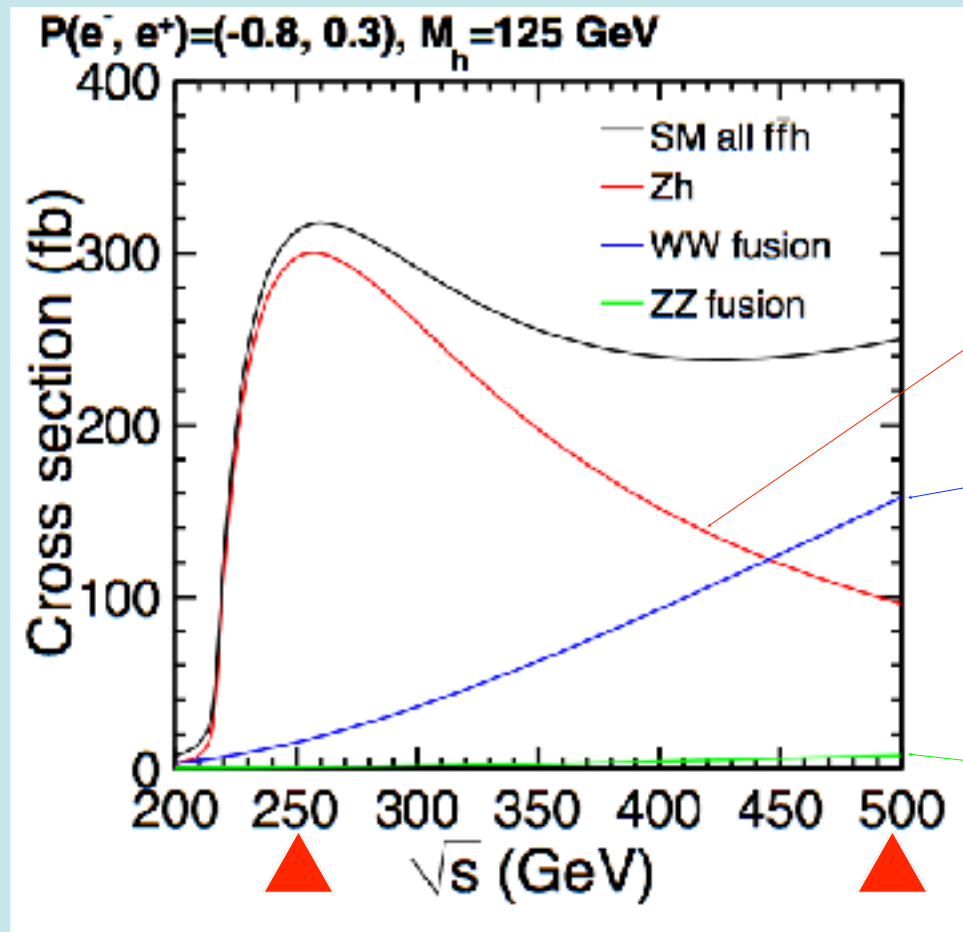
composite scale

New physics at 1 TeV \rightarrow deviation is at most $\sim 10\%$
 We need a %-level precision \rightarrow ILC

Main Production Processes

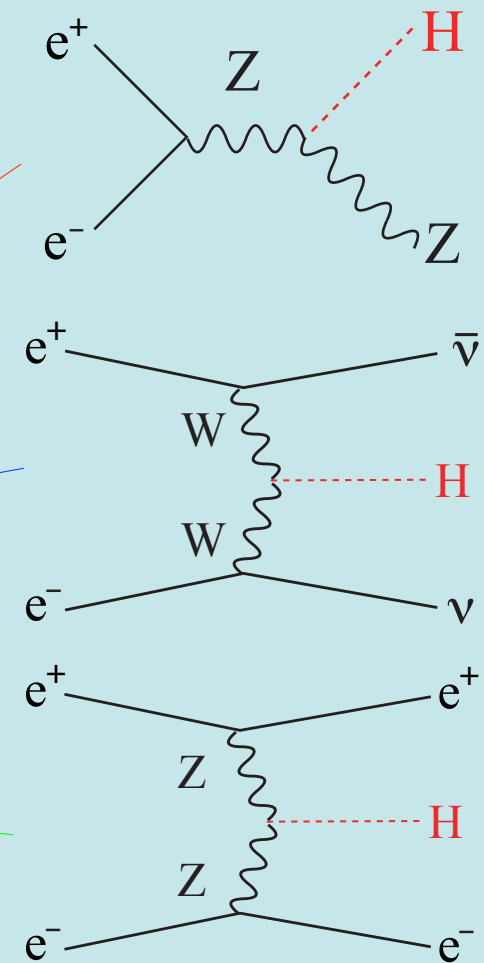
Single Higgs Production

Production cross section



ZH dominates at 250 GeV
(~80k ev: 250 fb⁻¹)

vvH takes over at 500 GeV
(~125k ev: 500 fb⁻¹)

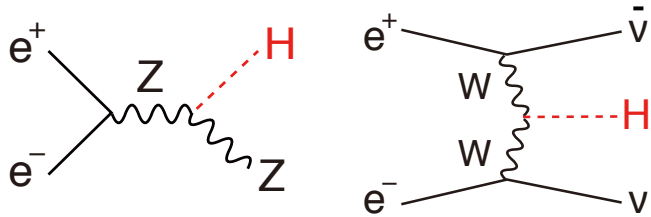


200k w/ TDR baseline, eventually >1M Higgs events!

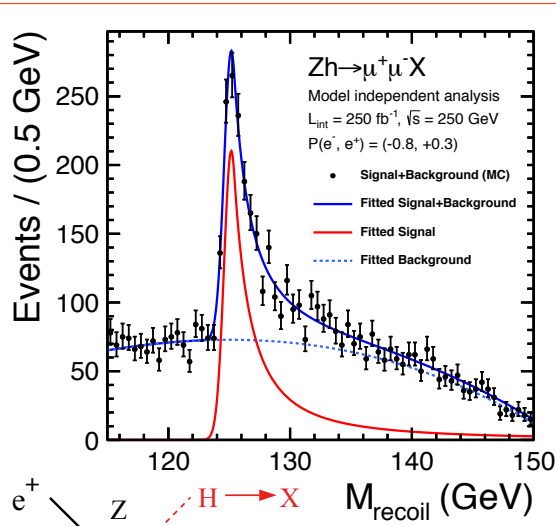
Key Point

At LHC all the measurements are $\sigma \times BR$ measurements.

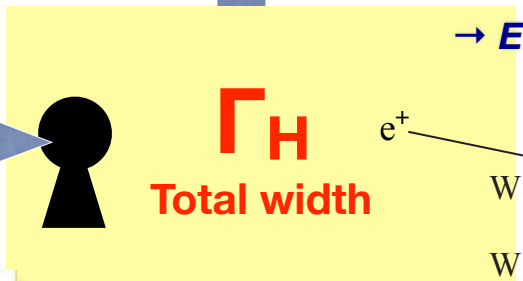
At ILC all but **the σ measurement using recoil mass technique** is $\sigma \times BR$ measurements.



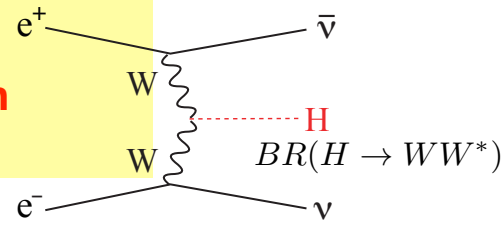
$$g_{HAA}^2 \propto \Gamma(H \rightarrow AA) = \Gamma_H \cdot BR(H \rightarrow AA)$$



The Key



WW-fusion is crucial for precision total width measurement
 → $E_{cm} > 350 \text{ GeV}$



Can detect even if Higgs decays invisibly!

$$M_X^2 = (p_{CM} - (p_{\mu^+} + p_{\mu^-}))^2$$

Higgs Couplings

Model-independent coupling fit, impossible at LHC

H20 Scenario

arXiv: 1506.05992
arXiv: 1506.07830

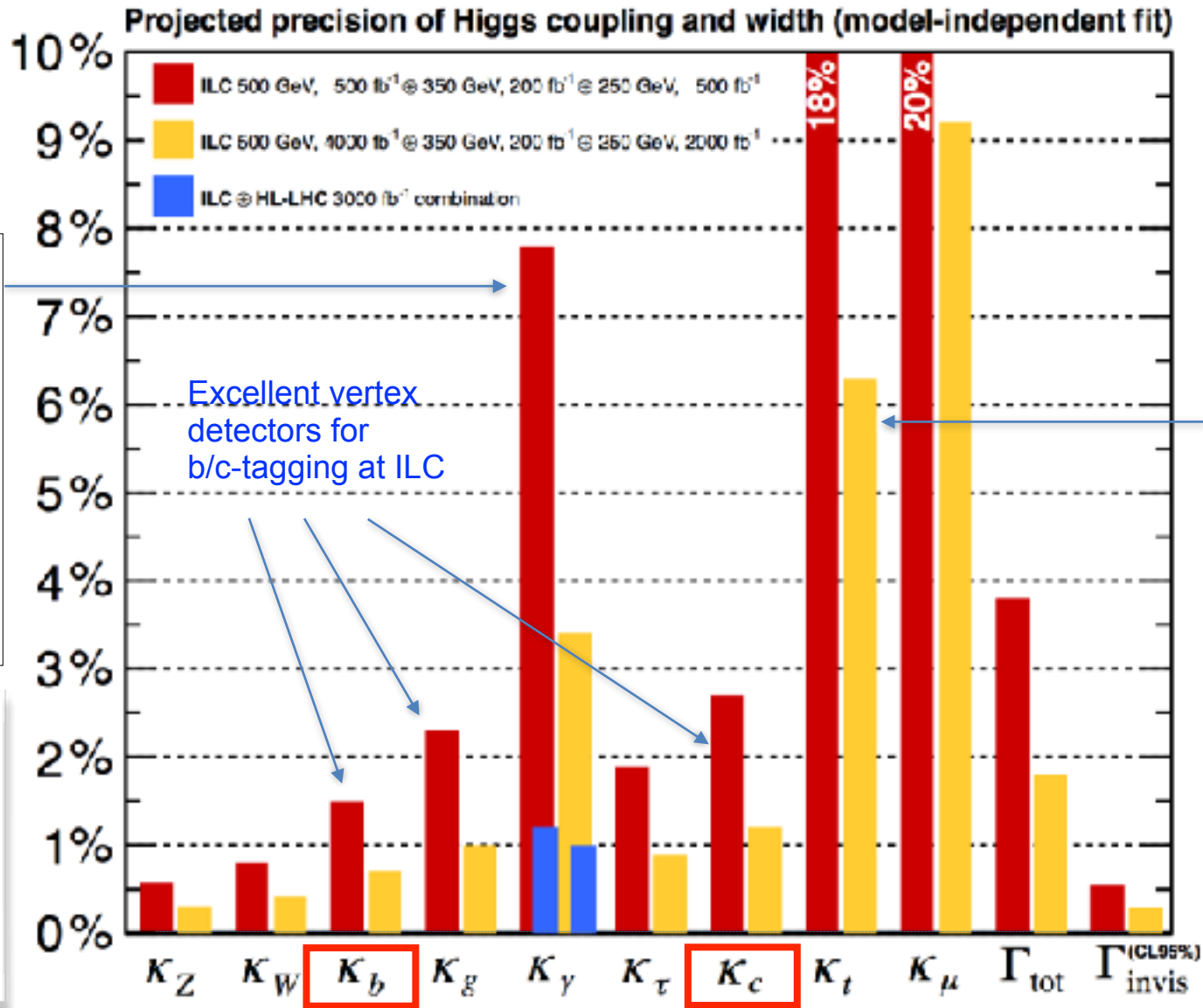
Better hyy with LHC/ILC synergy

LHC can precisely measure

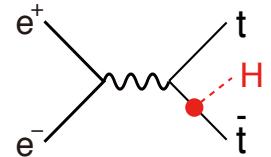
$$BR(h \rightarrow \gamma\gamma) / BR(h \rightarrow ZZ^*) = (K_\gamma / K_Z)^2$$

ILC can precisely measure K_Z

All of major Higgs decay modes accessible at ILC with 250-500GeV!



Top Yukawa improves by going to 550 GeV

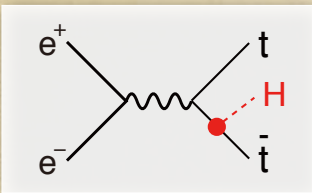


Near threshold → a factor of 4 enhancement of σ_{th} by going from 500GeV to 550 GeV

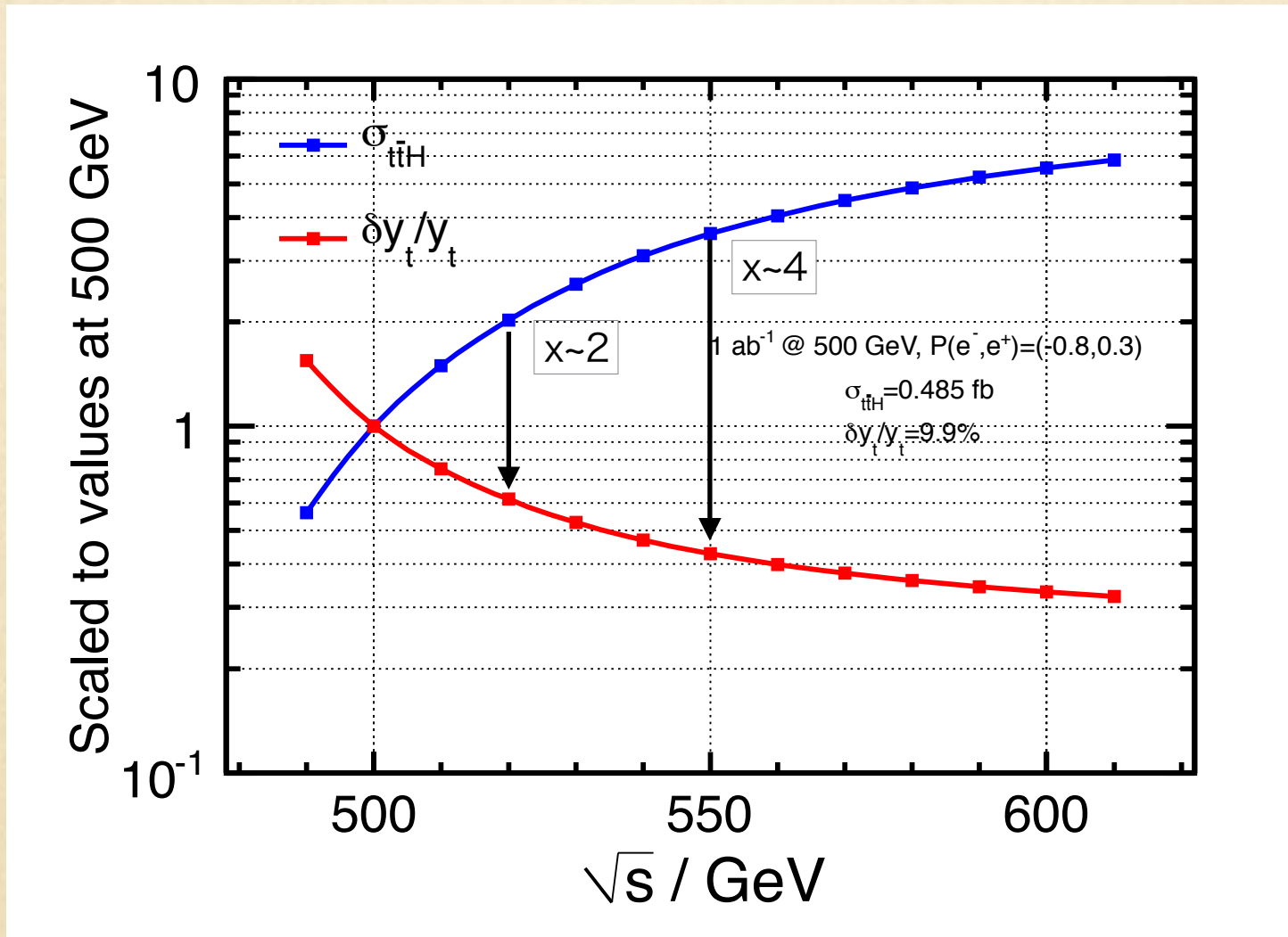
→ 3%

500 GeV already excellent except for K_t , K_μ , and K_γ

~1% or better for most couplings!



Top Yukawa coupling



Y. Sudo

Slight increase of E_{max} is very beneficial!

Model-dependent coupling fit (LHC-style 7-parameter fit)

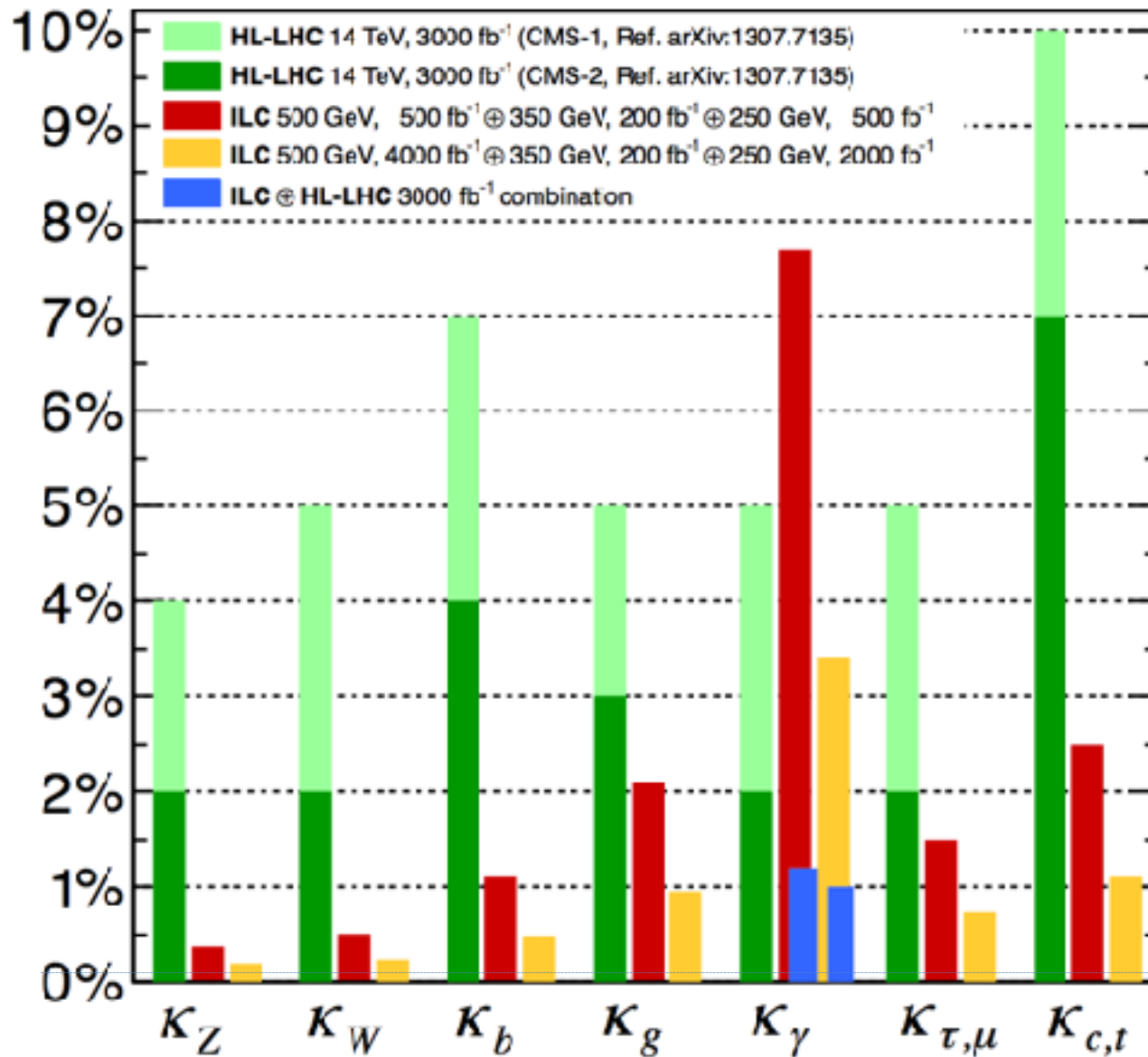
H20 Scenario

arXiv: 1506.05992

arXiv: 1506.07830

Projected Higgs coupling precision (7-parameter fit)

$\Sigma_{SM} BR = 1$

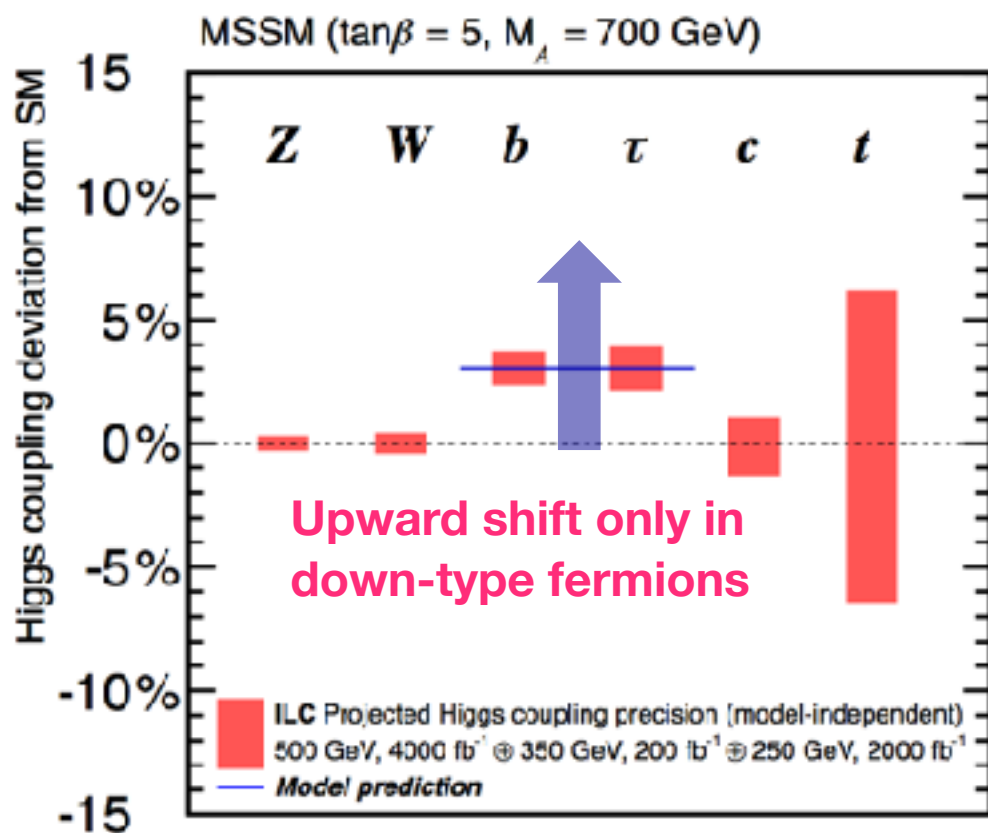


Possible to achieve precision far exceeding LHC!

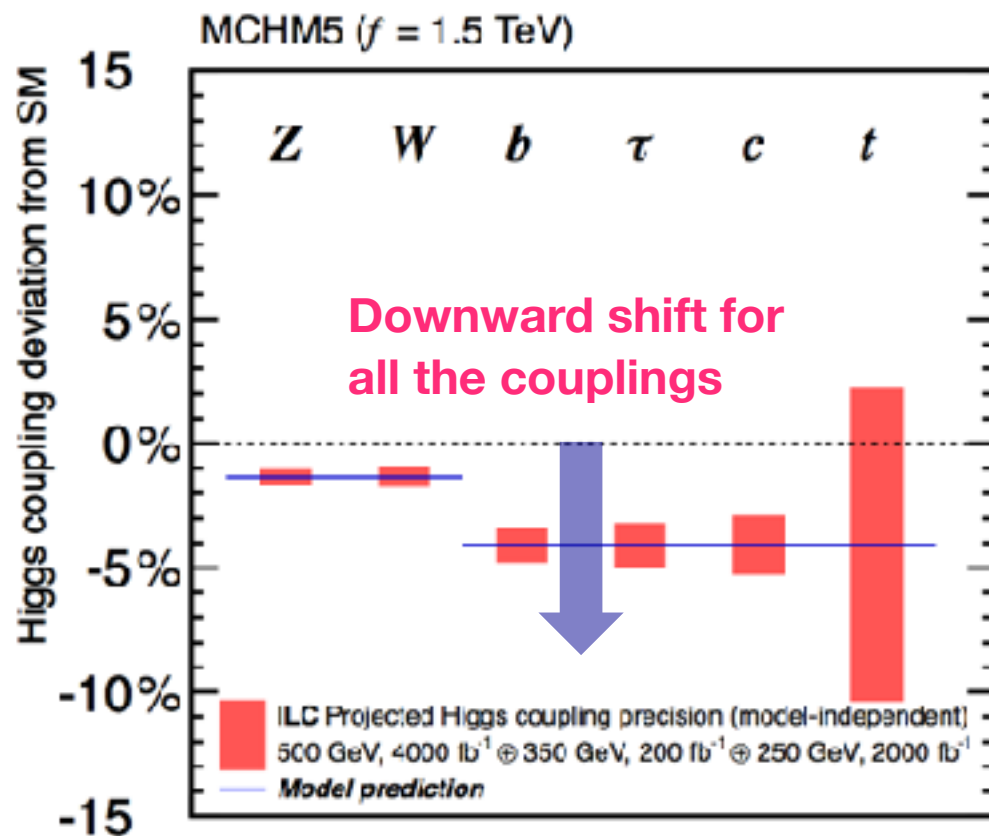
Fingerprinting

Elementary v.s. Composite?

Supersymmetry (MSSM)



Composite Higgs (MCHM5)



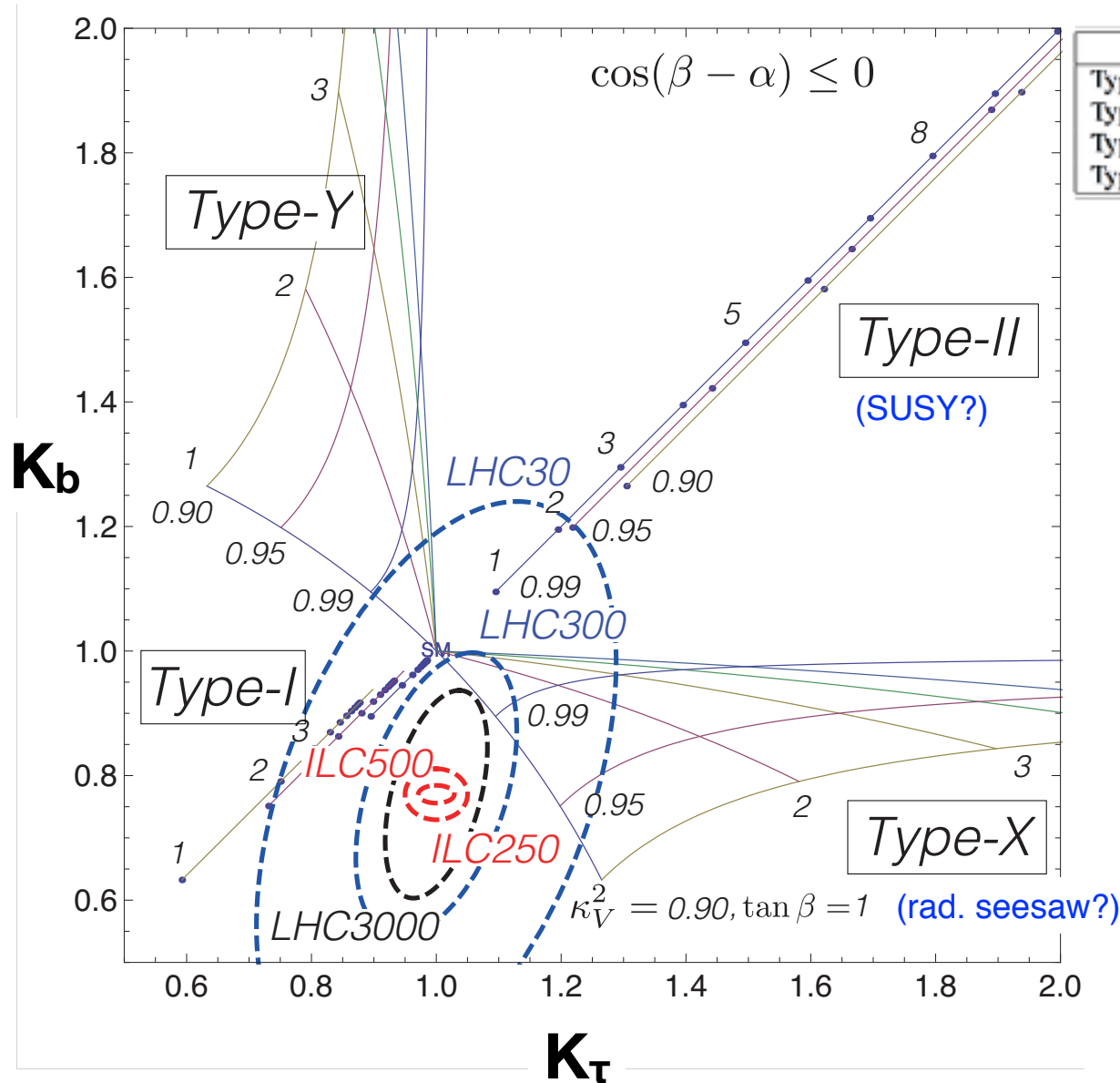
ILC 250+500 LumiUP

Complementary to direct searches at LHC: Depending on parameters, ILC's sensitivity far exceeds that of LHC!

Fingerprinting

2HDM

Multiplet Structure



	Φ_1	Φ_2	ν_R	d_R	ℓ_R	Q_L, L_L
Type I	+	-	-	-	-	+
Type II (SUSY)	+	-	-	+	+	+
Type X (Lepton-specific)	+	-	-	-	+	+
Type Y (Flipped)	+	-	-	+	-	+

4 Possible Z_2 Charge Assignments that forbids tree-level Higgs-induced FCNC

$$\kappa_V^2 = \sin(\beta - \alpha)^2 = 1 \Leftrightarrow \text{SM}$$

Given a deviation of the Higgs to Z coupling: $\Delta \kappa_V^2 = 1 - \kappa_V^2 = 0.01$ we will be able to **discriminate the 4 models!**

Model-dependent
7-parameter fit
ILC: Baseline lumi.

ILC TDR

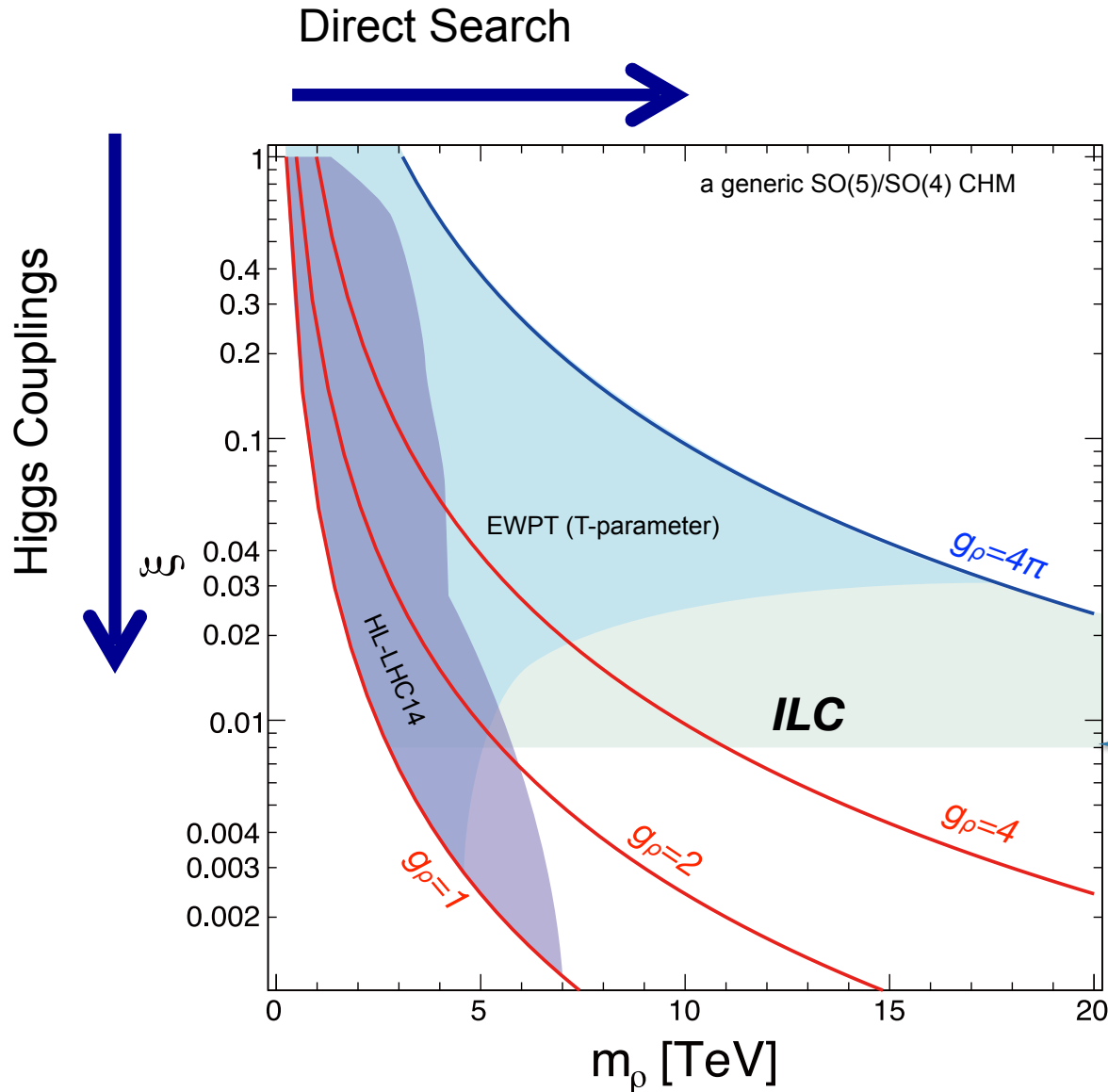
Snowmass ILC Higgs White Paper (arXiv: 1310.0763)

Kanemura et al (arXiv: 1406.3294)

Composite Higgs: Reach

Complementary approaches to probe composite Higgs models

- Direct search for heavy resonances at the LHC
 - Indirect search via Higgs couplings at the ILC
- Comparison depends on the coupling strength (g_*)



Based on Contino, et al, JHEP 1402 (2014) 006
 Torre, Thamm, Wulzer 2014
 Grojean @ LCWS 2014

$$\xi = \frac{g_\rho^2}{m_\rho^2} v^2 = \frac{v^2}{f^2}$$

$$\frac{g_{hVV}}{g_{h_{SM}VV}} = \sqrt{1 - \xi}$$

ILC (250+500 LumiUP)

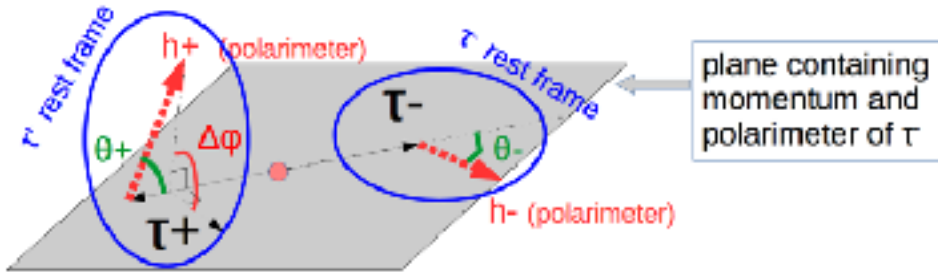
$$\Delta \frac{g_{hVV}}{g_{h_{SM}VV}} = 0.4\%$$

CP Mixing

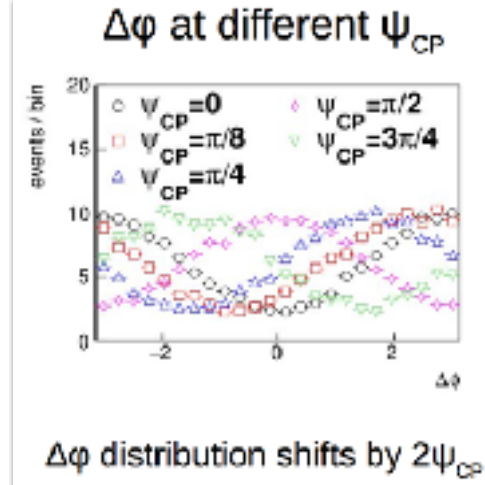
$$\mathcal{L}_{h\tau\tau} = g\bar{\tau} (\cos \Psi_{CP} + i\gamma_5 \sin \Psi_{CP}) \tau h$$

250GeV

CP from polarimeters : taus from spin 0 parent



- $\theta_{\pm}, \varphi_{\pm}$ direction of h_{\pm} with respect to τ - boost in τ_{\pm} rest frame
- $\Delta\varphi$ angle between polarimeter planes
- Ψ_{CP} CP mixing angle we want to measure



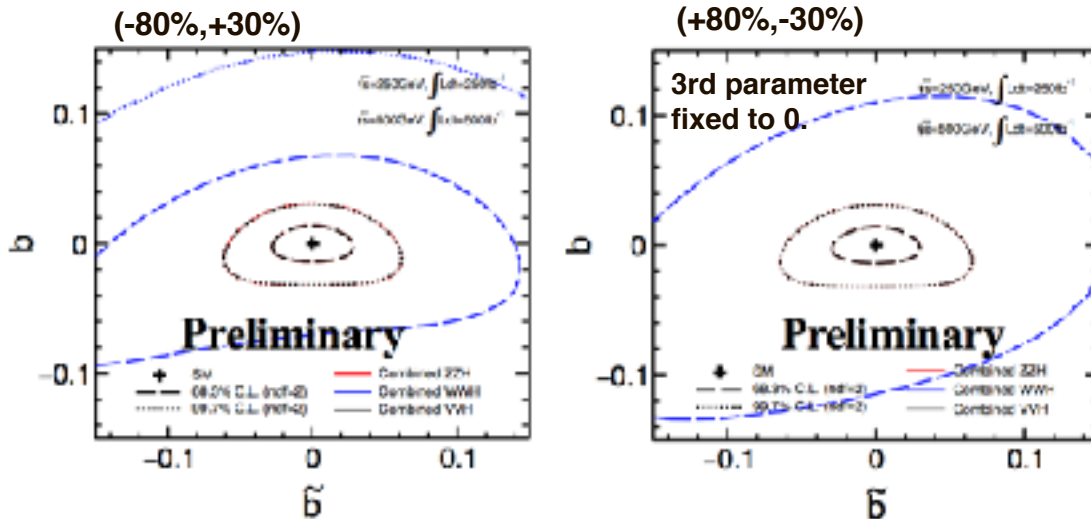
2ab⁻¹ @ 250 GeV

$\langle \delta \Psi_{CP} \rangle \simeq 3.8^\circ$
(preliminary)

$$\mathcal{L}_{VVh} = 2M_V^2 \frac{1}{\Lambda} \left(\frac{\Lambda}{v} + a \right) h V_\mu V^\mu + C_V \frac{b}{\Lambda} h V_{\mu\nu} V^{\mu\nu} + C_V \frac{\tilde{b}}{\Lambda} h V_{\mu\nu} \tilde{V}^{\mu\nu} \quad (C_W, C_Z) = (1, 1/2)$$

250+500GeV

Angular shape for ZZh (7 processes) + WW h (5 processes) + σ : 250fb⁻¹@250 GeV and 500fb⁻¹@500 GeV



Translation to Snowmass convention

$$f_{CP\tilde{b}} = (f_{\tilde{b}}/f_{a_4}) f_{CP}$$

$$\rightarrow f_{CP\tilde{b}} \simeq 1.5 \times 10^{-5}$$

($\tilde{b} \simeq 0.07(3\sigma)$ @ 500GeV)

A factor of 2.5 better than Snowmass fast simulation results (arXiv: 1309.4819).

Further improvement expected for H20!

EW Phase Transition

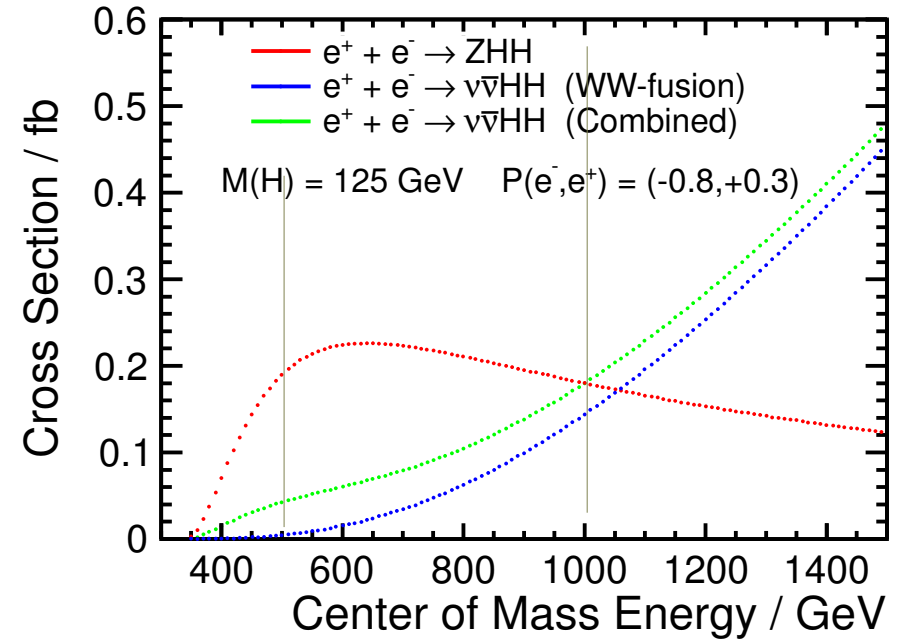
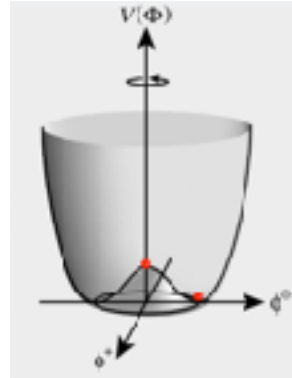
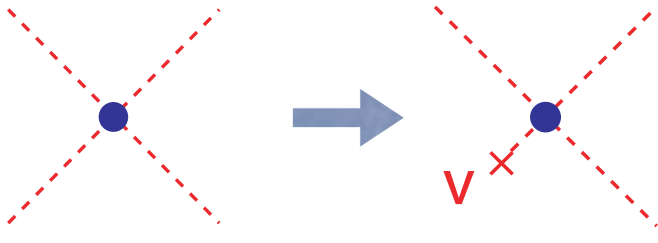
1st order

or

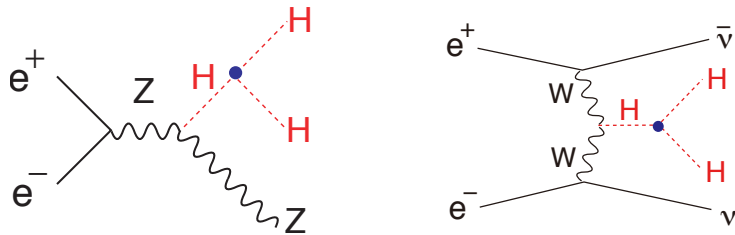
2nd order ?

Higgs Self-Coupling

The **Higgs 3-point self-coupling** is at the heart of EWSB!



There are **two ways to measure it** at ILC



arXiv:1310.0763	ILC500	ILC500-up	ILC1000	ILC1000-up
\sqrt{s} (GeV)	500	500	500/1000	500/1000
$\int \mathcal{L} dt$ (fb ⁻¹)	500	1600 [†]	500+1000	1600+2500 [†]
$P(e^-, e^+)$	(-0.8, 0.3)	(-0.8, 0.3)	(-0.8, 0.3/0.2)	(-0.8, 0.3/0.2)
$\sigma(ZHH)$	42.7%		42.7%	23.7%
$\sigma(\nu\bar{\nu}HH)$	-	-	26.3%	16.7%
λ	83%	46%	21%	13%

26% (H20)

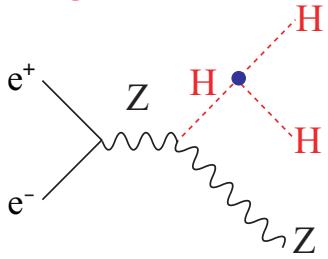
Challenging even at ILC because of

- Small cross section
- **Presence of irreducible BG diagrams**

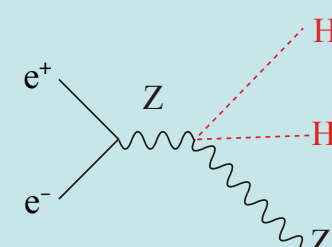
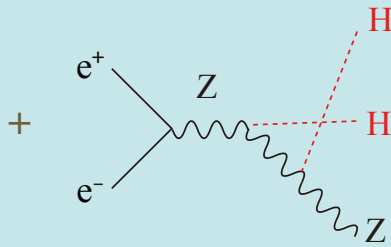
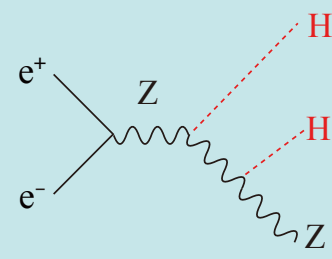
Ongoing analysis improvements **towards O(10)% measurement**

The Problem : BG diagrams dilute self-coupling contribution

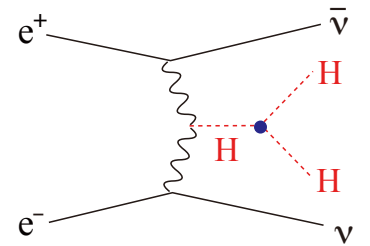
Signal diagram



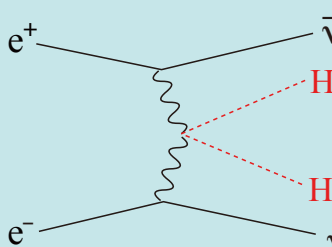
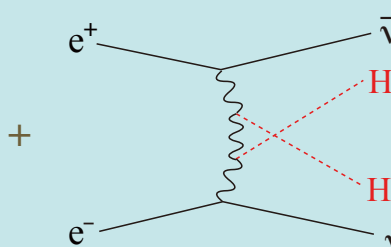
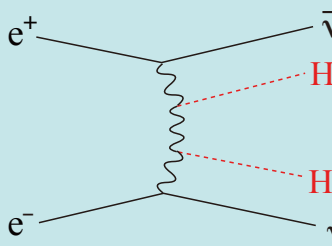
Irreducible BG diagrams



Signal diagram



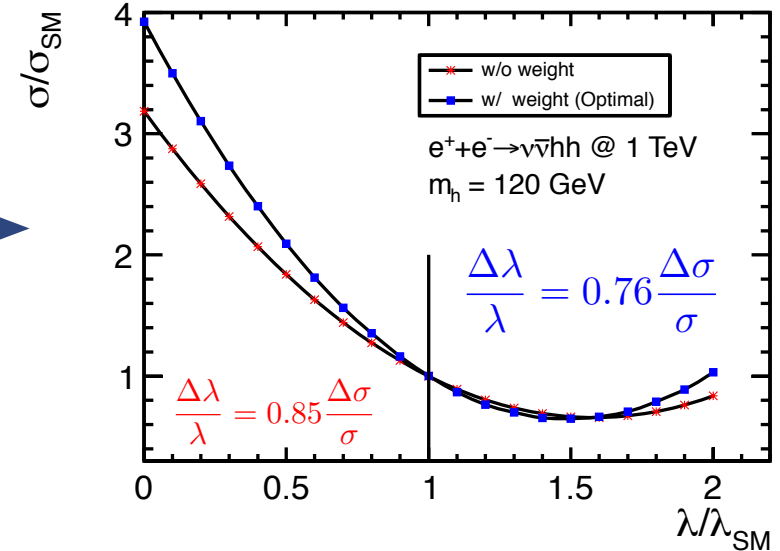
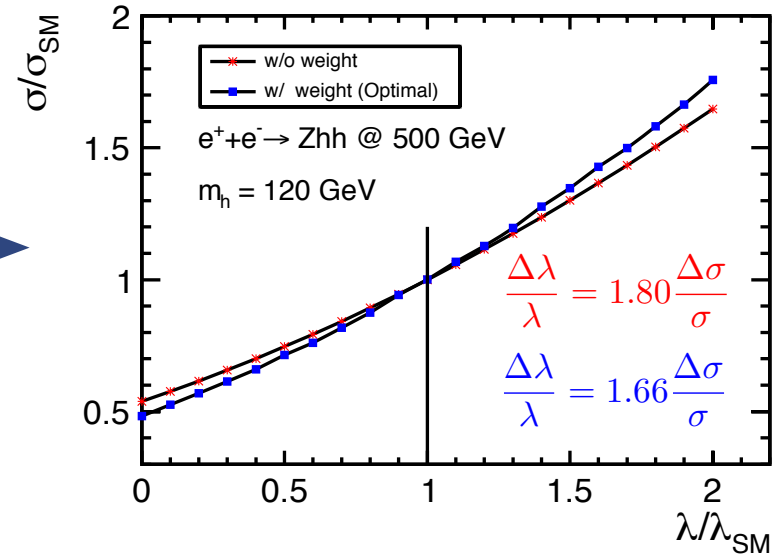
Irreducible BG diagrams



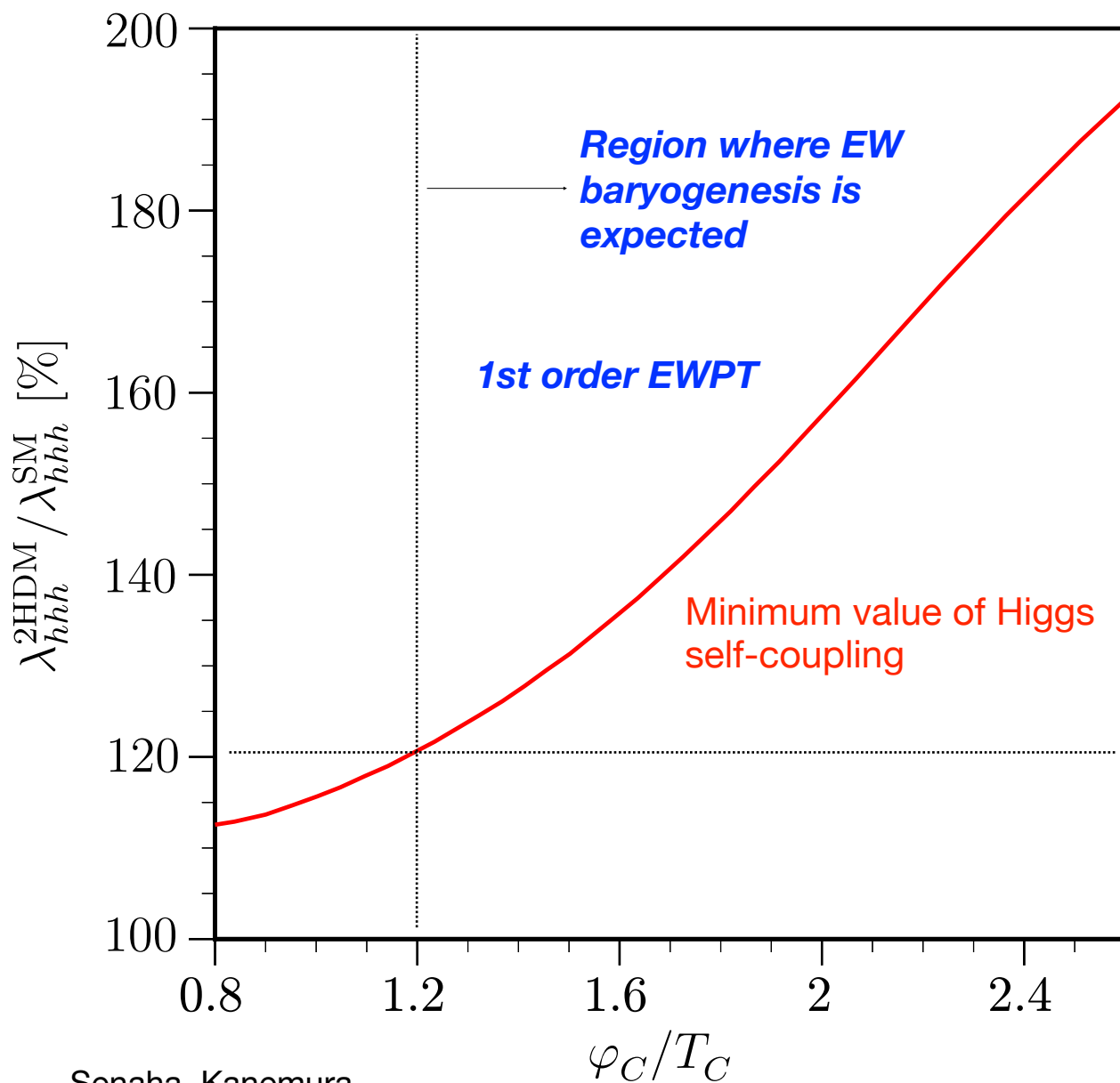
$$\sigma = \lambda^2 S + \lambda I + B$$

$$\frac{\Delta\lambda}{\lambda} = F \cdot \frac{\Delta\sigma}{\sigma}$$

$F=0.5$ if no BG diagrams



Electroweak Baryogenesis



Example:

Electroweak baryogenesis in a **Two Higgs Doublet Model**

Large deviations in Higgs self-coupling

→ **1st order EW phase transition**

→ **Out of equilibrium**

+ **CPV in Higgs sector**

→ **EW baryogenesis possible**

Constructive interference between signal and BG diagrams:

→ **if +100% deviation, then 14% precision expected on λ at 500GeV.**

ILC can address the idea of **baryogenesis occurring at the electroweak scale.**

Summary of Physics Part

- The primary goal for the next decades is ***to uncover the secret of the EW symmetry breaking.*** The discovery of H(125) completed the SM particle spectrum and taught us how the EW symmetry was broken. However, it does not tell us why it was broken. ***Why $\mu^2 < 0$?*** To answer this question we need to go beyond the SM.
- There is a big fork concerning the question: ***Is H(125) elementary or composite?*** There are ***two powerful probes*** in hand: ***H(125) itself and the top quark.*** Different models predict different deviation patterns in Higgs and top couplings. ***ILC will measure these couplings with unprecedented precision.***
- This will open up ***a window to BSM*** and ***fingerprint BSM models***, otherwise it will ***set the energy scale for the E-frontier machine that will follow LHC and ILC.***
- ***Cubic self-coupling measurement*** will decide whether the EWSB was strong 1st order phase transition or not. If it was, it will provide us the possibility of understanding ***baryogenesis at the EW scale.***
- ***The ILC is an ideal machine to address these questions*** (regardless of BSM scenarios) and we can do this ***model-independently.***
- Though I could not cover it today, It is also very important to stress that ***ILC, too, is an energy frontier machine.*** It will access the energy region never explored with any lepton collider. It is not a tiny corner of the parameter space that will be left after LHC. ***There is a wide and interesting region for ILC to explore.***
- Once a new particle is discovered, we can precisely determine its properties. In the SUSY case, for instance, we might even hope to probe GUT scale physics through RGE.
- ***In this way, ILC will pave the way to the moment of creation.***

ILC Project Status

ILC Promotion: Recent History

- Oct. 2013: Japanese HEP community proposed to host the ILC in Japan as a global project. “A Proposal for a Phased Execution of the ILC”.
- Statements on ILC hosted in Japan
 - The European Strategy for Particle Physics Update 2013
 - ACFA/AsiaHEP Statement on the ILC (Sep. 2013)
 - P5 Report (May 2014)
 - ICFA statements (Jan. & July 2014)
- Sep. 2013: Science Council of Japan (SCJ) sent report on the ILC project to MEXT.
- May 2014: MEXT set up ILC Advisory Panel.

MEXT's ILC Review

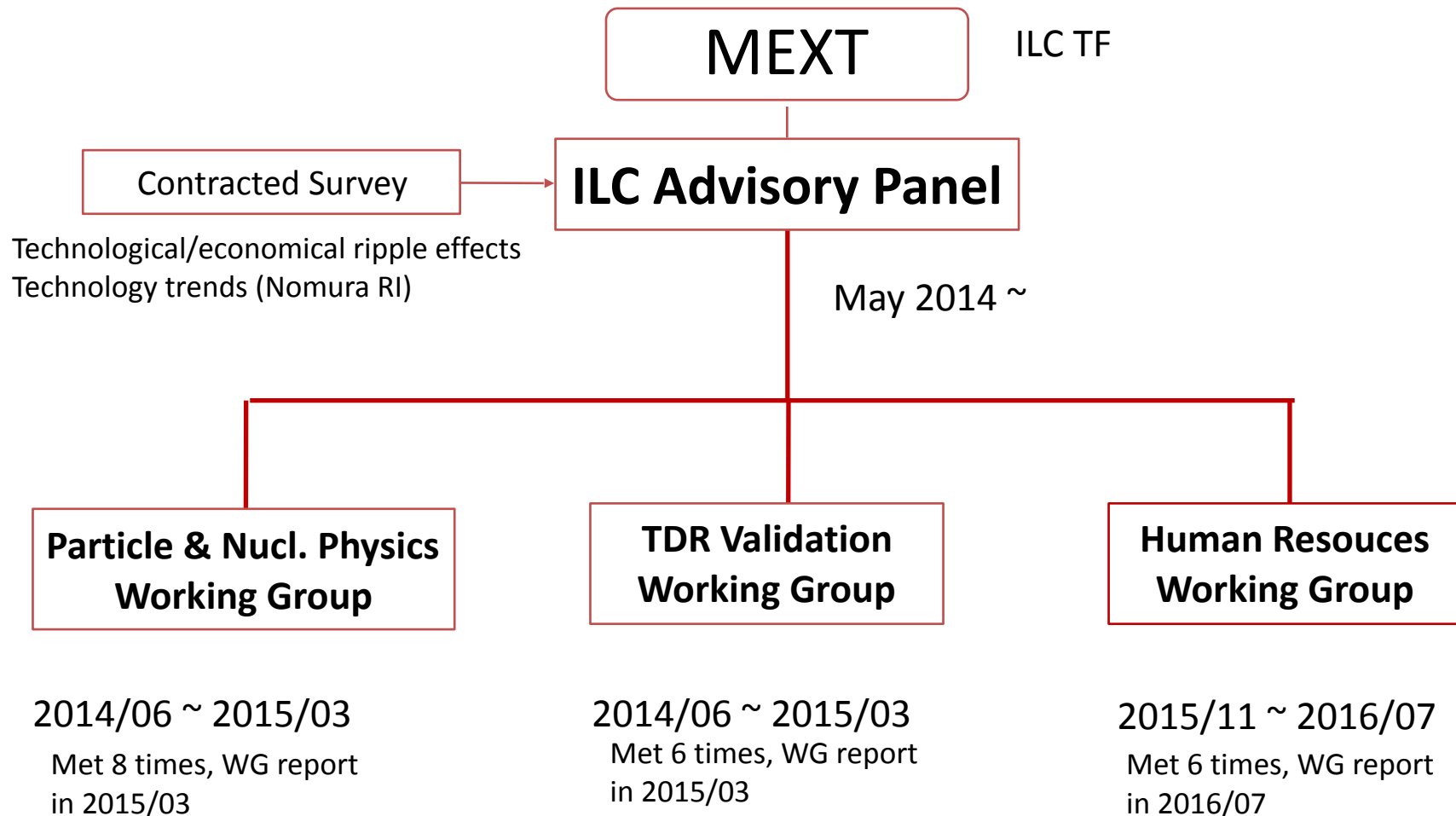
MEXT

=

Japan's
**Ministry of
Education,
Culture, Sports, Science and
Technology**

ILC Advisory Panel

Set up in May 2014 under MEXT ILC Task Force to investigate various issues concerning the possibility of hosting the ILC in Japan



Interim Summary

http://www.mext.go.jp/b_menu/shingi/chousa/shinkou/038/gaiyou/1360593.htm

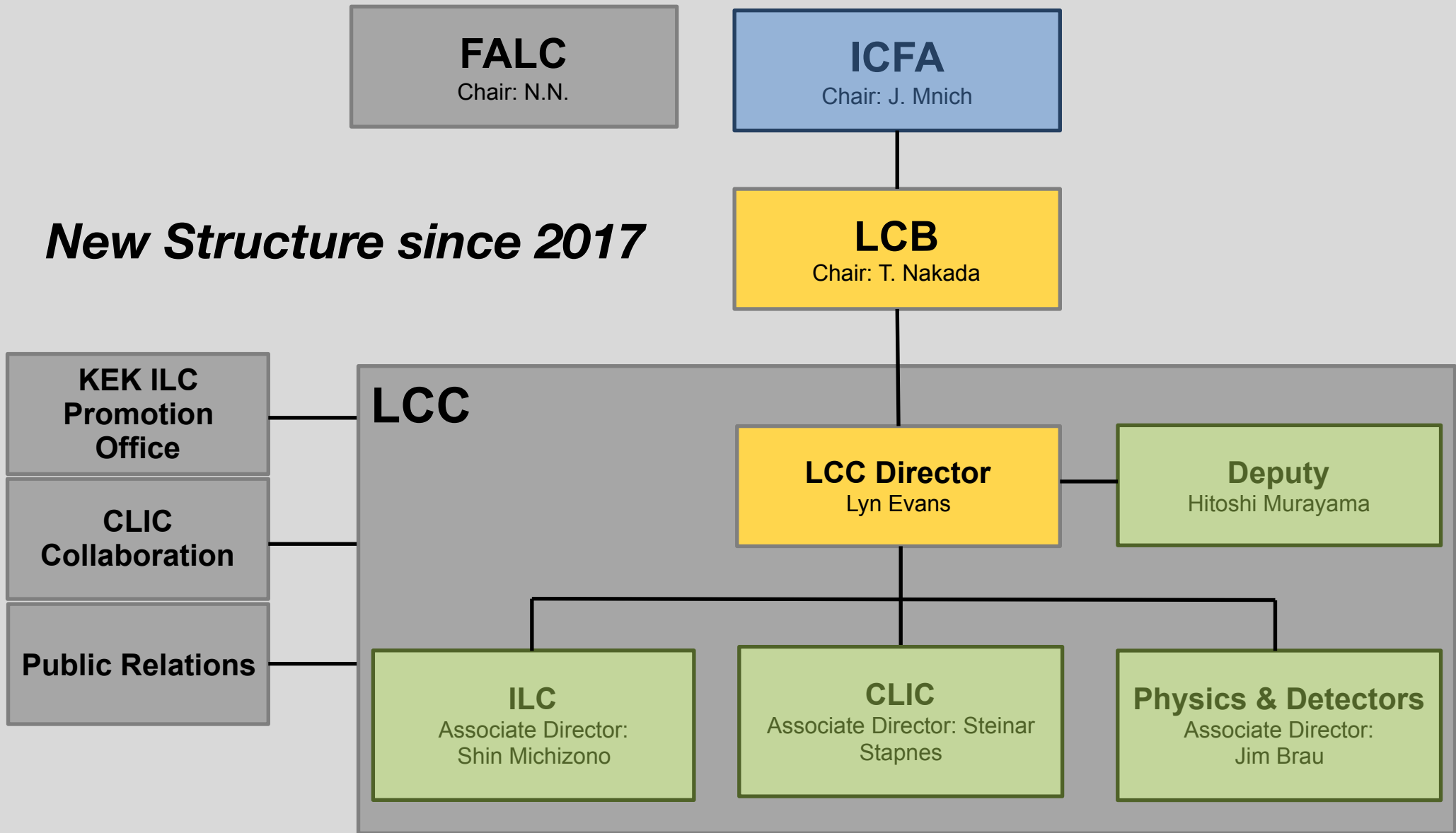
- ILC Advisory Panel published an interim summary of their discussions based on the reports from the two working groups (Particle & Nuclear Physics WG and TDR Validation WG).
- The interim summary pointed out the following issues
 - Obtain clear vision for international cost sharing
 - **Make clear scientific merits (not only precision studies of Higgs and top but also possibilities of new particle discoveries) that match the investment**
 - **Monitor, analyze, and examine the development of LHC experiments.**
 - Solve remaining technological issues and **mitigate cost risk.**
 - Get **understanding from the general public** and other scientific communities.

Linear Collider Collaboration (LCC)

**international collaboration
to address these issues**



New Structure since 2017



ILC Brochure



ILC communicators with consultation by LCC Physics WG

Essentially completed → to be publicized soon!

And did the homework following MEXT's recommendation

ILC-NOTE-2016-067
DESY 16-146, IPMU16-0108
KEK Preprint 2016-9, LAL 16-185
MPP-2016-174, SLAC-PUB-16751

July, 2016

Implications of the 750 GeV $\gamma\gamma$ Resonance as a Case Study for the International Linear Collider

ILC PHYSICS WORKING GROUP

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TAKEO MOROI¹⁶, FRANÇOIS RICHARD¹⁰, JUNPING TIAN¹², MARCEL VOŠ¹⁷,
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ABSTRACT

If the $\gamma\gamma$ resonance at 750 GeV suggested by 2015 LHC data turns out to be a real effect, what are the implications for the physics case and upgrade path of the International Linear Collider? Whether or not the resonance is confirmed, this question provides an interesting case study testing the robustness of the ILC physics case. In this note, we address this question with two points: (1) Almost all models proposed for the new 750 GeV particle require additional new particles with electroweak couplings. The key elements of the 500 GeV ILC physics program—precision measurements of the Higgs boson, the top quark, and 4-fermion interactions—will powerfully discriminate among these models. This information will be important in conjunction with new LHC data, or alone, if the new particles accompanying the 750 GeV resonance are beyond the mass reach of the LHC. (2) Over a longer term, the energy upgrade of the ILC to 1 TeV already discussed in the ILC TDR will enable experiments in $\gamma\gamma$ and e^+e^- collisions to directly produce and study the 750 GeV particle from these unique initial states.

arXiv:1607.03829v2 [hep-ph] 31 Jul 2016

The Potential of the ILC for Discovering New Particles

Document Supporting the ICFA Response Letter to the ILC Advisory Panel

LCC PHYSICS WORKING GROUP

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LIST², MIHOKO NOJIRI^{1,9}, MAXIM PERELSTEIN¹⁰, ROMAN PÖSCHL¹¹, JÜRGEN
REUTER², FRANK SIMON¹², TOMOHIKO TANABE¹³, JAMES D. WELLS¹⁴, JAEHOON
YU¹⁵; HOWARD BAER¹⁶, MIKAEL BERGGREN², SVEN HEINEMEYER¹⁷, SUVI-LEENA
LEHTINEN², JUNPING TIAN¹³, GRAHAM WILSON¹⁸, JACQUELINE YAN¹; HITOSHI
MURAYAMA^{9,19,20}, JAMES BRAU²¹

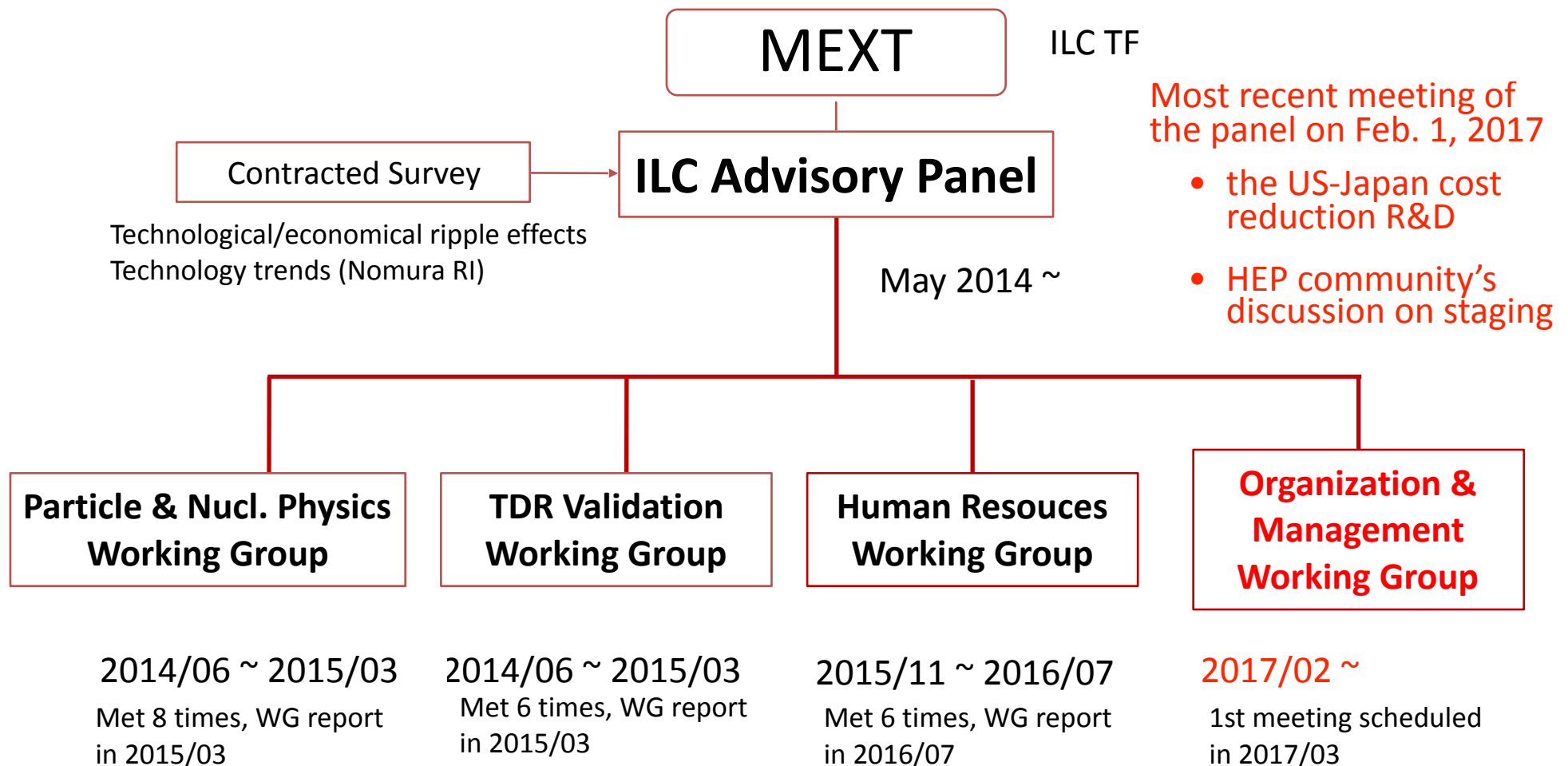
Abstract

This paper addresses the question of whether the International Linear Collider has the capability of discovering new particles that have not already been discovered at the CERN Large Hadron Collider. We summarize the various paths to discovery offered by the ILC, and discuss them in the context of three different scenarios: 1. LHC does not discover any new particles, 2. LHC discovers some new low mass states and 3. LHC discovers new heavy particles. We will show that in each case, ILC plays a critical role in discovery of new phenomena and in pushing forward the frontiers of high-energy physics as well as our understanding of the universe in a manner which is highly complementary to that of LHC.

For the busy reader, a two-page executive summary is provided at the beginning of the document.

ILC Advisory Panel

Set up in May 2014 under MEXT ILC Task Force to investigate various issues concerning the possibility of hosting the ILC in Japan



Support from Diet Members and Industrial Sector in Japan

- ***Federation of Diet Members for the ILC*** (since 2008 with >150 members)
- ***Advanced Accelerator Association Promoting Science & Technology (AAA)*** (since 2008 with 100 companies and 40 universities and research institutions)
- Event in Washington DC on Feb. 2016 coordinated by Hudson Institute and AAA. 4th visit to Washington by Diet members with MEXT officials.



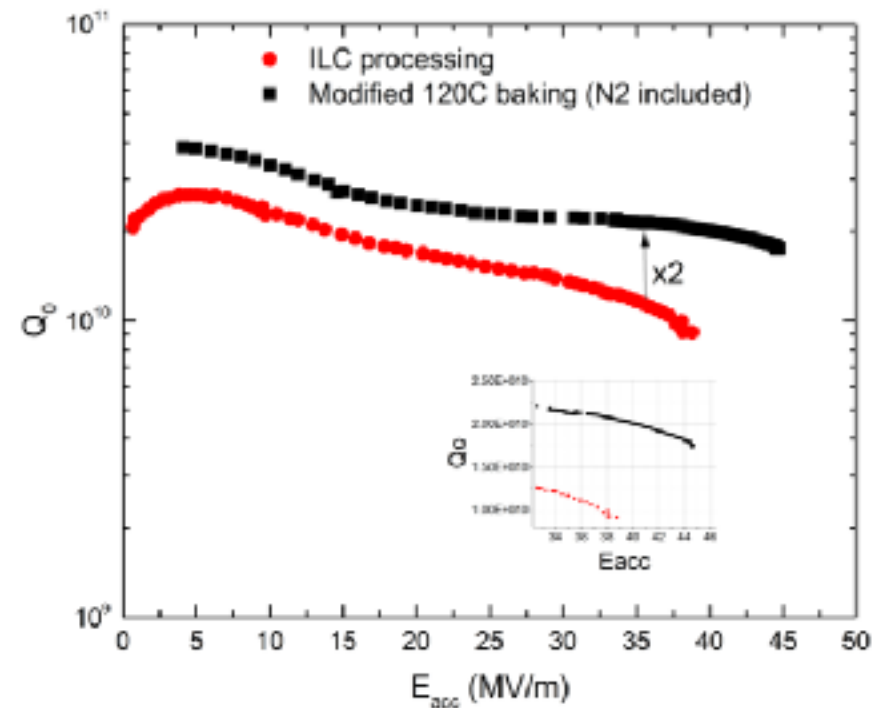
From LC NEWSLINE

<http://newsline.linearcollider.org/2016/03/03/us-japan-symposium/>

Hon. Shionoya is recommending the Kasoku Kids cartoon book to the roundtable discussion chaired by Dr. William Schneider, Jr. (Hudson Institute)

MET-DOE Discussion Group

- High level officers from MEXT visited their DOE counter part at the end of May and it was agreed to start a US-Japan discussion group co-chaired by Director of Office of Science of DOE and a corresponding level officer in MEXT. They decided to meet every 2-3 months.
- In their Oct. 2016 meeting, it was agreed to start ***US-Japan collaborative research for ILC cost reduction***: aiming at 10-12% cost reduction of the ILC machine construction cost.
 - Cost reduction in Nb material preparation
 - High-Q high-gradient SCRF cavity using nitrogen infusion



Science First with the ILC- Keynote speech by Takeo Kawamura *from LC NEWSLINE*

<http://newsline.linearcollider.org/2016/12/08/science-first/>



In his keynote at LCWS2016, former MEXT Minister Takeo Kawamura stressed that while fundamental research may have application in the long run, it's the science that is most important.

Hon. Takeo Kawamura giving a keynote speech at the LCWS2016 in Morioka, Japan. Image: LCWS2015 LOC



LCWS2016 in Morioka, Japan.



Standing ovations for Hon. Takeo Kawamura's speech by LCWS2016 participants. Image: LCWS2016 LOC

Staging Discussion

- In LCWS 2016, Nov. in Morioka, it was agreed to start seriously considering a staging scenario of the ILC ***to significantly reduce the initial construction cost.***
 - 1st stage as a Higgs factory
 - and later stages taking advantage of flexible energy expandability of a linear collider.
- LCB/LCC started working on possible staging scenarios to build consensus among the worldwide HEP community.

Summary

- **MEXT is seriously investigating various issues to be solved to host the ILC in Japan.**
- **KEK/JHEP is taking various actions together with the LCC to address issues pointed out by the MEXT ILC Advisory Panel.**
- **MEXT-DOE joint discussion group started.**
- **US-Japan collaborative research on cost reduction started.**
- **There are important political interactions happening also in Europe and Asia.**
- **Serious discussions on staging started.**
- **As Hon. Kawamura said in LCWS 2016, 2017-2018 will be a very important time for the ILC.**

Backup

The Current Official Operation Scenario: H20

<i>ILC Physics Goals</i>	500 GeV	350 GeV	250 GeV
• precision Higgs couplings	✓	✓	✓
• gHWW and overall normalization of Higgs couplings	✓	✓	
• search for invisible and exotic Higgs decay modes	✓	✓	✓
• Higgs couplings to top	✓		
• Higgs self-coupling	✓		
• search for extended Higgs states	✓		
• precision electroweak couplings of the top quark	✓		
• precision W couplings	✓	✓	
• precision search for Z'	✓		
• search for supersymmetry	✓		
• search for Dark Matter	✓		
• top quark mass from threshold scan		✓	
• precision Higgs mass			✓

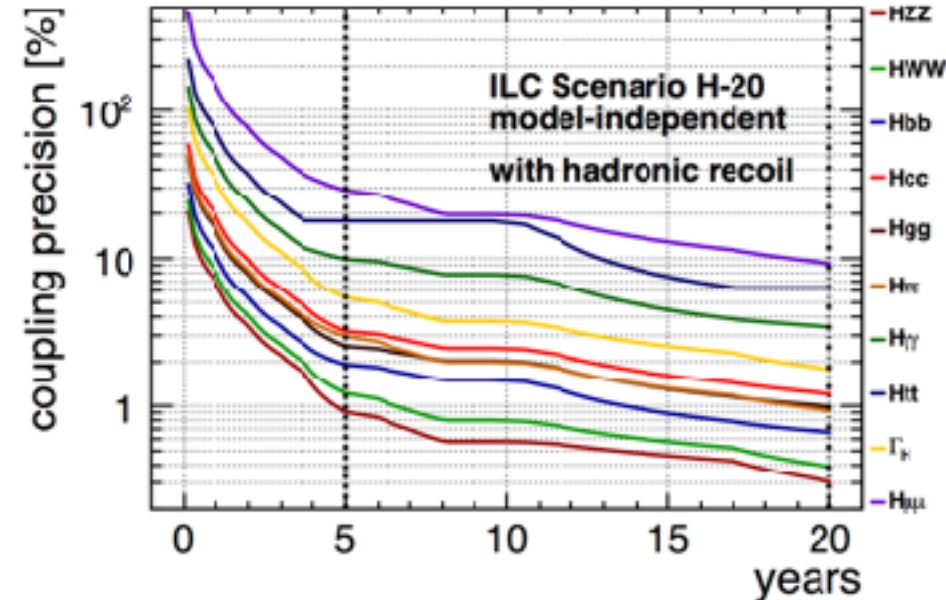
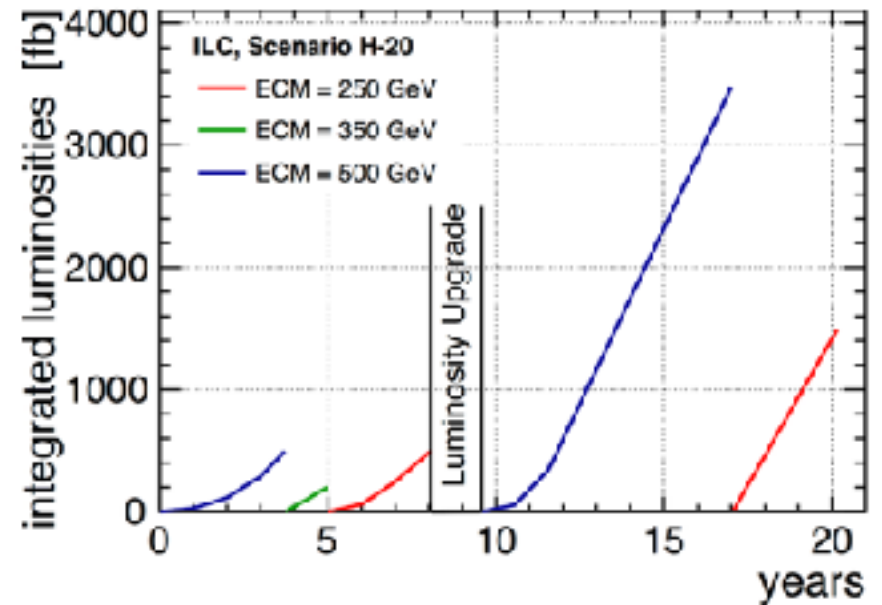
Figure 1: ILC Physics Goals.

$$\Delta m_h = 20 \text{ MeV}$$

→ 0.2% coupling uncertainty
for hWW and hZZ

High luminosity 250GeV run will be needed anyway, and the 250GeV stage alone can produce significant physics outputs, but, of course, the full program needs higher energy running.

Integrated Luminosities [fb]



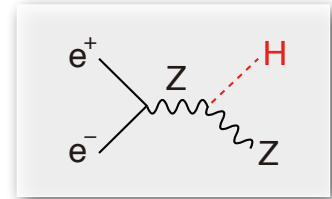
Higgs-related Physics at $E_{cm} \approx 500 \text{ GeV}$

Three well know thresholds

ZH @ 250 GeV ($\sim M_Z + M_H + 20 \text{ GeV}$) :

- Higgs mass, width, J^{PC}
- Gauge quantum numbers
- Absolute measurement of HZZ coupling (**recoil mass**)
- BR($h \rightarrow VV, qq, ll, \text{invisible}$) : $V=W/Z(\text{direct}), g, \gamma(\text{loop})$

\rightarrow Higgs couplings (other than top)

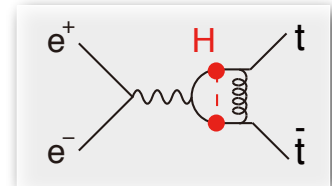


$t\bar{t}$ @ 340-350 GeV ($\sim 2m_t$) : ZH meas. Is also possible

- Threshold scan \rightarrow theoretically clean m_t measurement:
 - \rightarrow test stability of the SM vacuum
 - \rightarrow indirect meas. of top Yukawa coupling
- A_{FB} , Top momentum measurements
- Form factor measurements

$$\Delta m_t(\overline{MS}) \simeq 100 \text{ MeV}$$

$\gamma\gamma \rightarrow HH$ @ 350 GeV possibility



$v\bar{v}H$ @ 350 - 500 GeV :

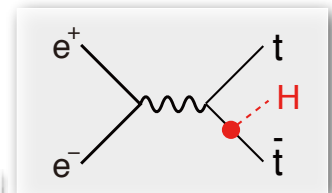
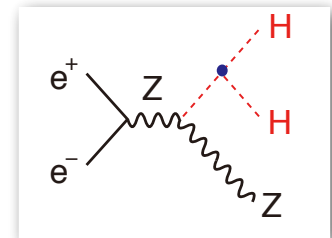
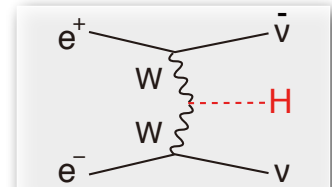
- **HWW coupling** \rightarrow total width \rightarrow absolute normalization of Higgs couplings

ZHH @ 500 GeV ($\sim M_Z + 2M_H + 170 \text{ GeV}$) :

- Prod. cross section attains its maximum at around 500 GeV \rightarrow Higgs self-coupling

$t\bar{t}H$ @ 500 GeV ($\sim 2m_t + M_H + 30 \text{ GeV}$) :

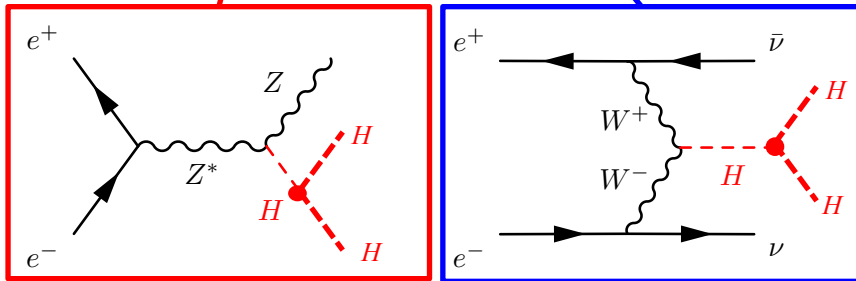
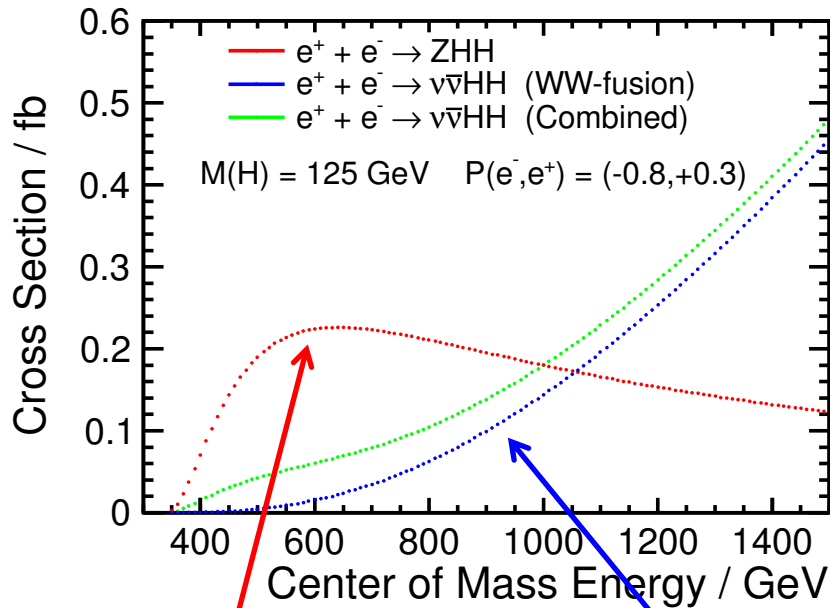
- Prod. cross section becomes maximum at around 800 GeV.
- QCD threshold correction enhances the cross section \rightarrow top Yukawa measurable at 500 GeV concurrently with the self-coupling



We can access all the relevant Higgs couplings at $\sim 500 \text{ GeV}$ for the mass-coupling plot!

Direct Measurement

Cross section vs CM energy (e^+e^-)



Diagrams with triple-Higgs coupling

Expected precision based on **full detector simulation** studies:

ILC
500 GeV, 4 ab⁻¹
 $\delta\lambda = 27\%$

ILC
500 GeV, 4 ab⁻¹
& 1 TeV, 8 ab⁻¹
 $\delta\lambda = 10\%$

References:

J. Tian, LC-REP-2013-003

M. Kurata, LC-REP-2014-025

C. Duerig, Ph.D. thesis at DESY, 2016

HH \rightarrow bbbb, bbWW* combination

CLIC
1.4 GeV, 1.5 ab⁻¹
 $\delta\lambda = 21\%$

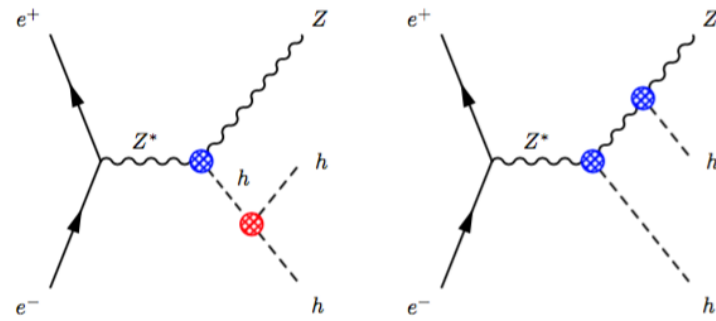
CLIC
1.4 TeV, 1.5 ab⁻¹
& 3 TeV, 2 ab⁻¹
 $\delta\lambda = 10\%$

References:

arXiv: 1307.5288

HH \rightarrow bbbb only, upgrade in progress including bbWW*

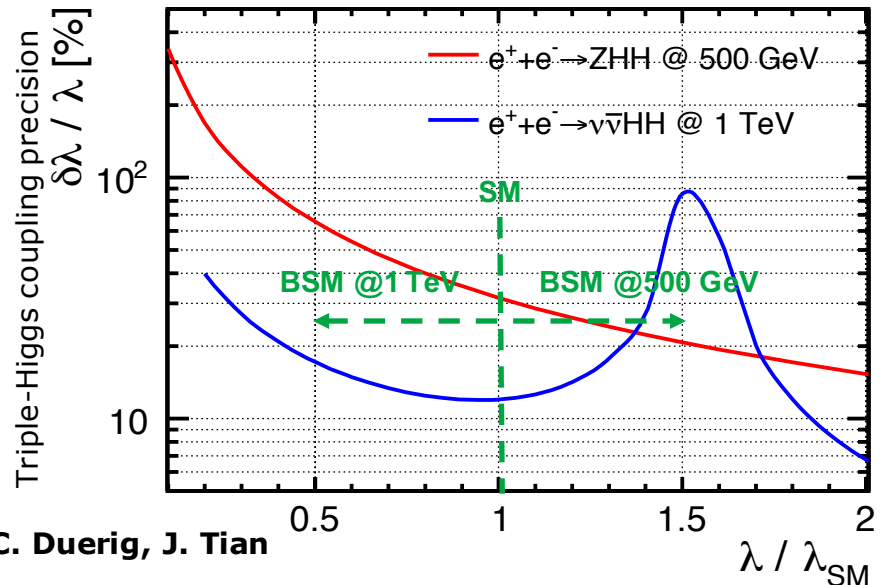
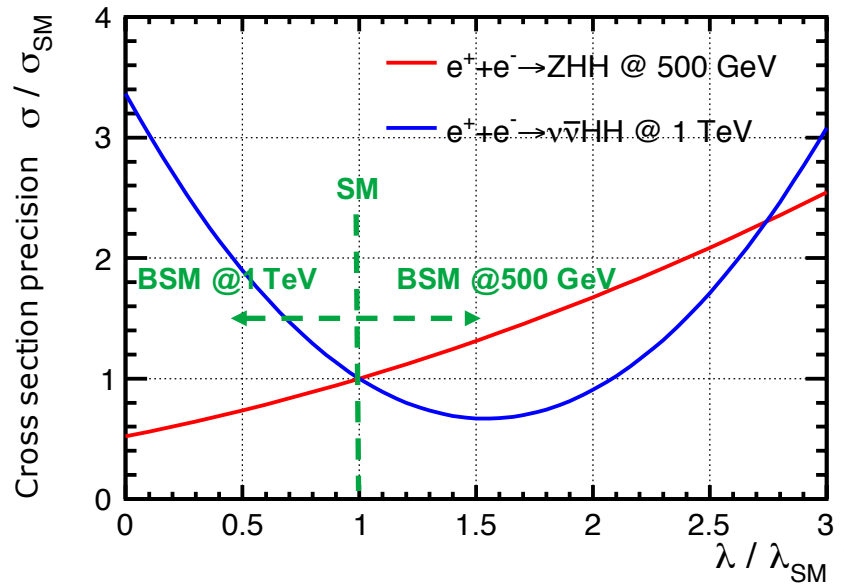
Triple-Higgs coupling with BSM



BSM can modify the triple-Higgs coupling. What effect does it have on the total cross section?

At 500 GeV, the cross section **increases** with increasing λ .

At 1 TeV, the cross section **decreases** with increasing λ . [Same as LHC]

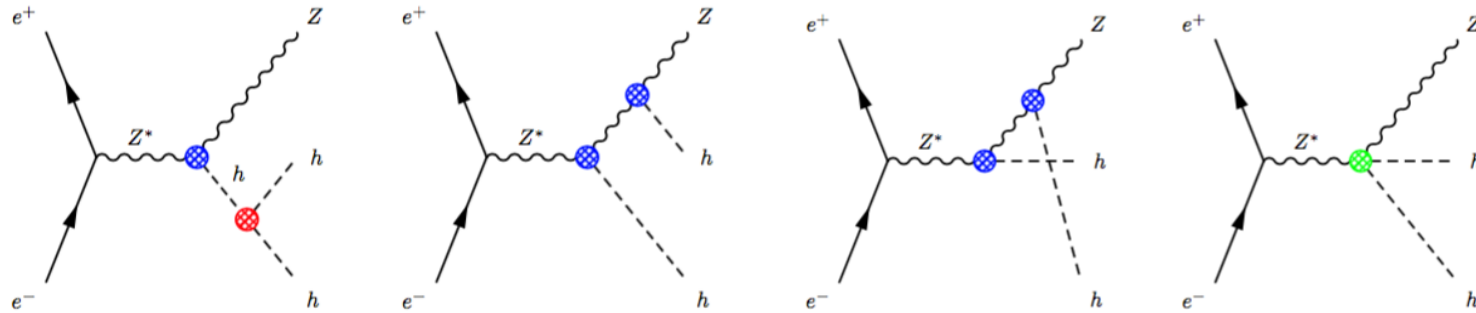


C. Duerig, J. Tian

500 GeV has unique sensitivity to larger triple-Higgs coupling

Model-independent extraction with EFT (1)

[Barklow, Fujii, Jung, Peskin, Tian]



- The authors ask the question: **If there is a deviation from the SM cross section, how to interpret this as a shift of the Higgs potential?**
- EFT analysis with dimension 6 operators: goal is to extract the parameter c_6 in

$$\Delta\mathcal{L} = -c_6 \frac{\lambda}{v^2} (\Phi^\dagger \Phi)^3$$

- **9 additional dimension 6 coefficients** contribute to the double Higgs production. Of which:
 - **3** are determined by **precision electroweak** data
 - **3** are determined by measurement of **$e^+e^- \rightarrow W+W^-$**
 - **1** combination is constrained by the **small size of $h \rightarrow \gamma\gamma$**
- **Need 2 more constraints \rightarrow will be provided by $e^+e^- \rightarrow Zh$**

Model-independent extraction with EFT (2)

[Barklow, Fujii, Jung, Peskin, Tian]

- Estimates of $e^+e^- \rightarrow Zh$ for the ILC program [Ogawa, Fujii, Tian]:
 - $\delta(c_H) = 1\%$ $\delta(16c_{WW}) = 0.25\%$ (highly correlated)
- The effect of these parameters on the double Higgs cross section is

$$\frac{\sigma(e^+e^- \rightarrow Zhh)}{\sigma_{SM}} = 1 - \underline{3.6 c_H} + \underline{7.4 (16c_{WW})} + 0.56 c_6$$

Higgs wavefunction renormalization
& new vertex $\Delta\mathcal{L} = \frac{c_H}{v_0} h\partial_\mu h\partial^\mu h$

dim-6 vertices enhanced by (s/m_Z^2)

*These are issues for any double Higgs production process.

For $e^+e^- \rightarrow Zhh$, precision measurement of single-Higgs process brings these effects under control at the **10% level in c_6** . Within this uncertainty, the extraction of c_6 is completely **model-independent**.

Conclusions

on triple-Higgs coupling at future e+e- colliders

- **Indirect measurement** at the **30%-level** may be possible with high luminosity (e.g. circular colliders)
 - Provides complementary (but model-dependent) information about the loop process.
- For **direct measurement**, expected precision is **$\delta\lambda/\lambda=27\%$ at 500 GeV**, and **$\delta\lambda/\lambda=10\%$ combining 1 TeV**. These numbers are supported by studies with full detector simulation.
 - **500 GeV** has unique sensitivity when **$\lambda > \lambda_{SM}$** .
 - This condition is theoretically well-motivated.
 - This is **complementary to 1 TeV and LHC**, which are sensitive for $\lambda < \lambda_{SM}$.
- Based on EFT analysis, a **model-independent** extraction of the coefficient **c_6** is possible at the **10% level**.

Top

Top

Reference

arXiv:1604.08122

Report from workshop at IFIC Valencia (July 2015)

An up-to-date summary with an extensive bibliography

Top physics at high-energy lepton colliders

Summary of TopLC15, IFIC Valencia, 30th June - 2nd July, 2015

M. Vos (IFIC, editor)

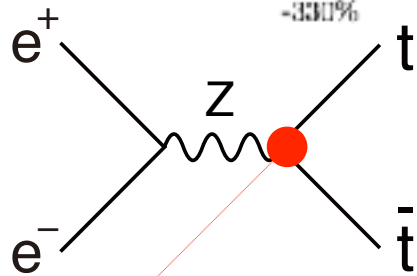
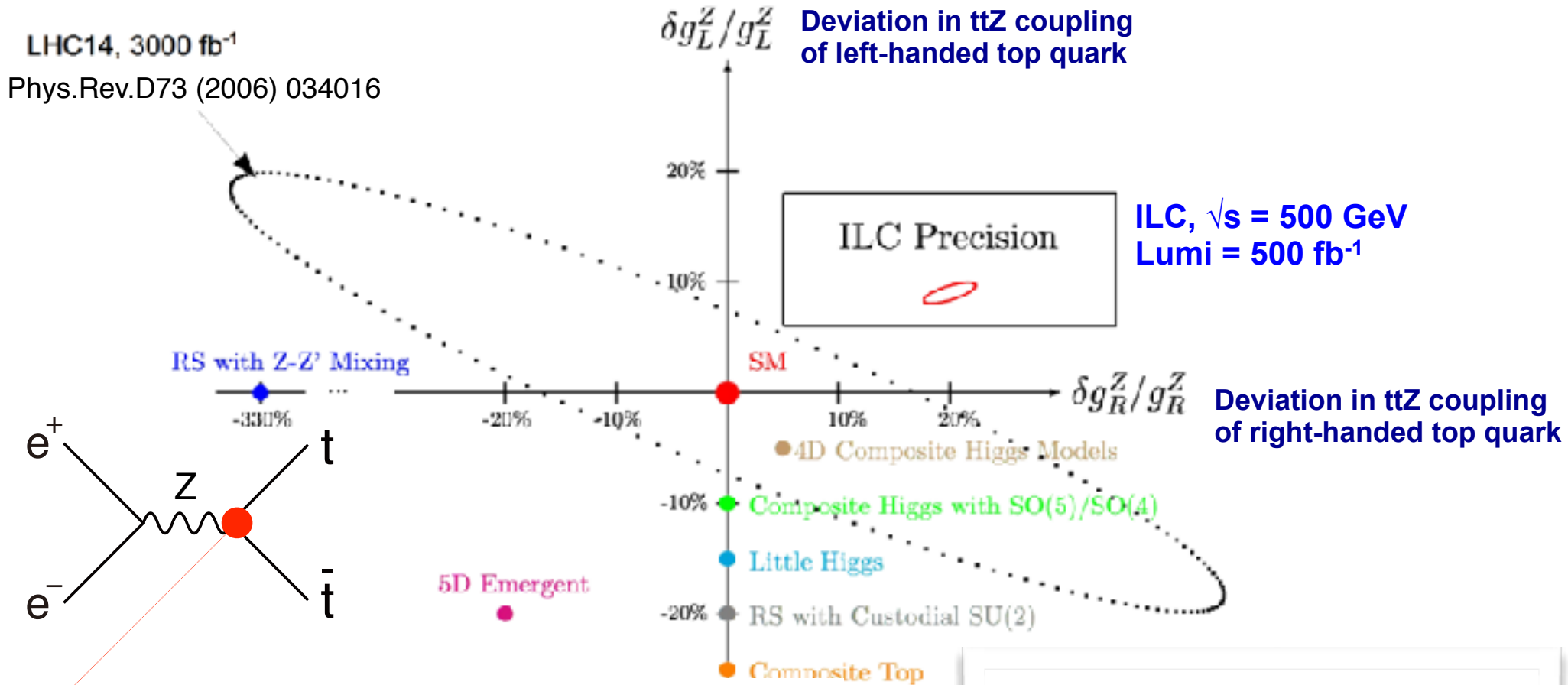
Attendees of the workshop:

*G. Abbas (IFIC), M. Beneke (TUM), S. Bilokin (LAL), M.J. Costa (IFIC),
S. de Curtis (U. & INFN Firenze), K. Fujii (KEK), J. Fuster (IFIC),
I. Garcia Garcia (IFIC), P. Gomez (IFIC), A. Hoang (U. Vienna), A. Irlin
(DESY), Y. Kiyo (Yuntendo), M. Kurata (Tokyo), I. Linszen (CERN), J. List
(DESY), M. Nebot (Lisboa), M. Perello (IFIC), R. Pöschl (LAL), N. Quach
(KEK), J. Reuter (DESY), F. Richard (LAL), G. Rodrigo (IFIC), Ph. Roloff
(CERN), E. Ros (IFIC), F. Simon (MPI Munich), J. Tian (KEK), A.F. Żarnecki
(Univ. of Warsaw)*

*Corresponding author: M. Vos, IFIC (UVEG/CSIC), Edificios de Investigacion,
c./ Catedratico Jose Beltran 2, E-46980 Paterna, Valencia, SPAIN (marcel.vos@ific.uv.es)*

Search for Anomalous tZ Couplings

Top: Heaviest in SM \rightarrow Must couple strongly to EW breaking sector (source of $\mu^2 < 0$)!
 \rightarrow **Specific deviation pattern** expected in ttZ form factors depending on new physics.
 \rightarrow **Beam polarization essential** to separate L- and R-couplings (Strength of ILC)



$$\Gamma_{\mu}^{ttX}(k^2, q_1, \bar{q}) = ie \left\{ \gamma_{\mu} \left(\tilde{F}_{1V}^X(k^2) + \gamma_5 \tilde{F}_{1A}^X(k^2) \right) + \frac{(q - \bar{q})_{\mu}}{2m_t} \left(\tilde{F}_{2V}^X(k^2) + \gamma_5 \tilde{F}_{2A}^X(k^2) \right) \right\}$$

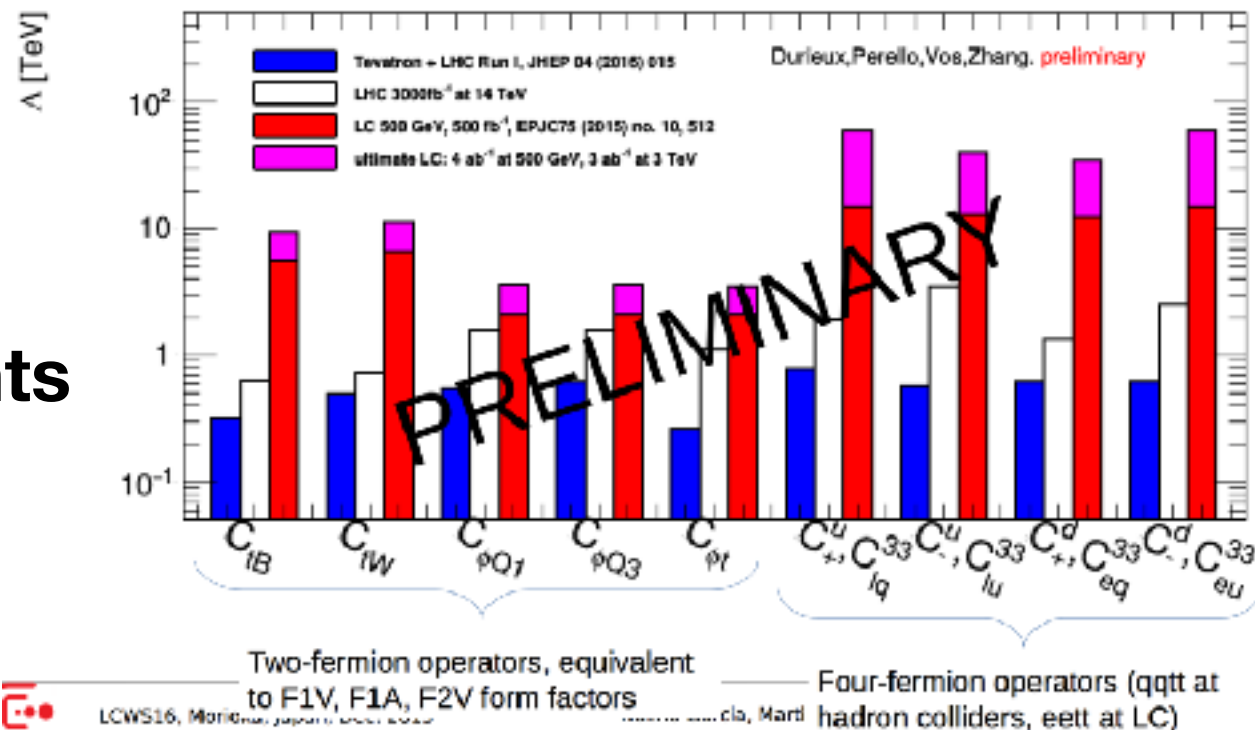
Deviation expected for various new physics models (new physics scale ~ 1 TeV)
 arXiv:1505.06020

ILC is sensitive to M_{KK} up to ~ 25 TeV for typical RS scenarios (even up to ~ 80 TeV in extreme cases)!

Top/QCD Talks from ILD at LCWS 2016

1. $e^+e^- \rightarrow tt$: semi-leptonic (Sviatslav Bilokin)
2. $e^+e^- \rightarrow bb$ (Sviatslav Bilokin) → Dec.6 (R.Poeschl)
3. $e^+e^- \rightarrow tt : bb\mu+\mu-\nu\nu$: MEM (Yo Sato) → Dec.7
4. mt reconstruction at 1TeV or higher (Nacho Garcia, Martin Perello, Philipp Roloff, Rickard Strom) with CLICdp → Dec.8 (R.Strom)
5. mt using radiative return to threshold (Marça Boronat and Pablo Gomis) → Dec.8 (M.Vos)
6. Global fit with D6 EFT (Martin Perello, et al.) → Dec.6 (M.Vos)

Form factors
→ EFT coefficients



**What if no deviation from
the SM would be seen?**

Clarify the Range of Validity of SM

Stability of SM Vacuum

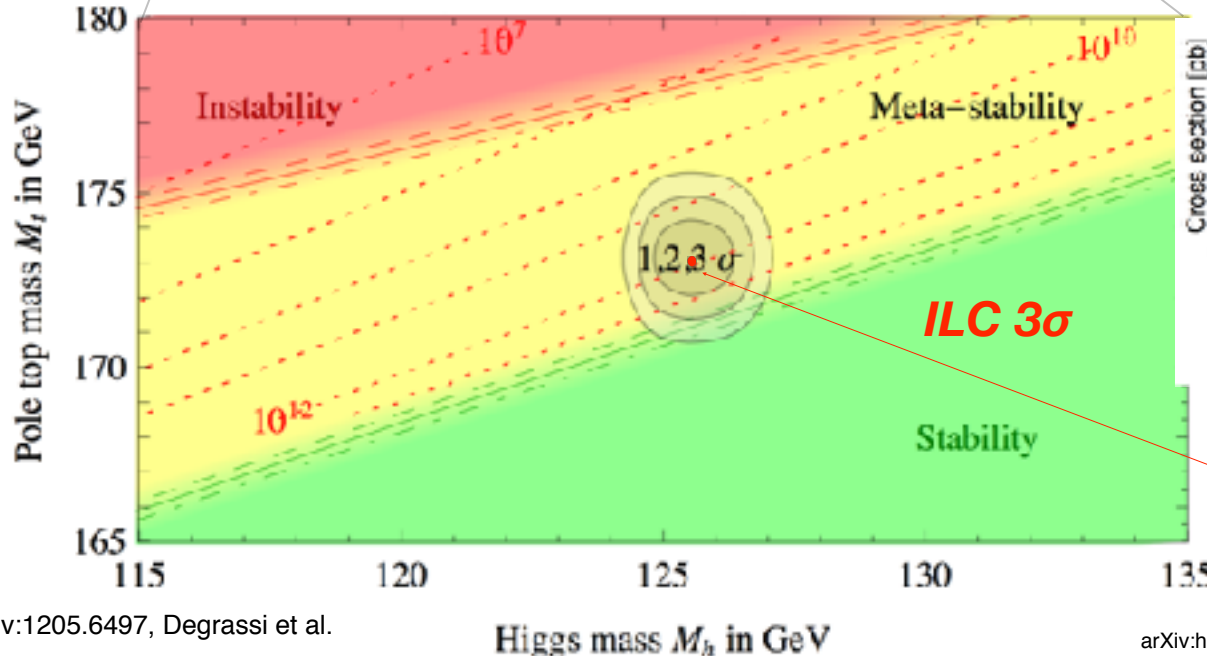
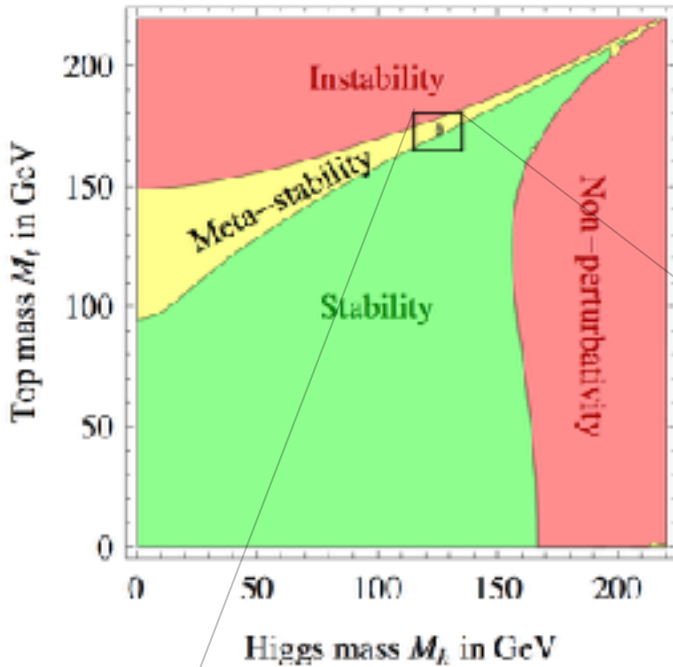
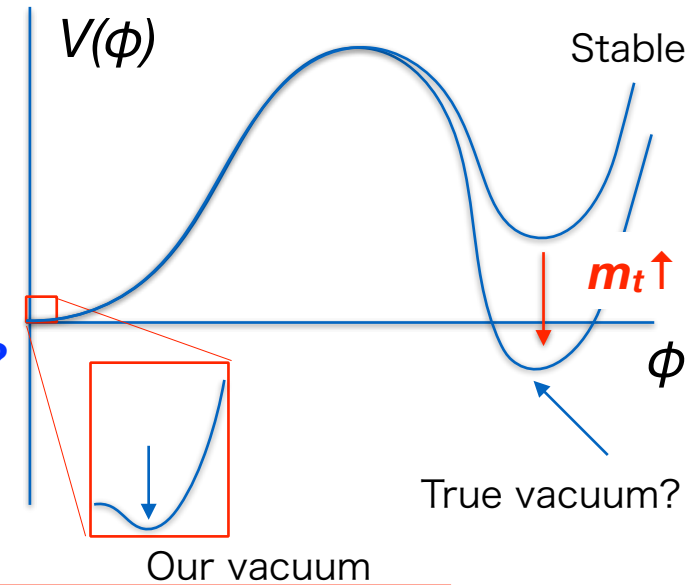
Top Yukawa coupling drives the 4-point Higgs coupling (λ) to negative!

The current values of m_t and m_h :
Subtle point of meta-stability!

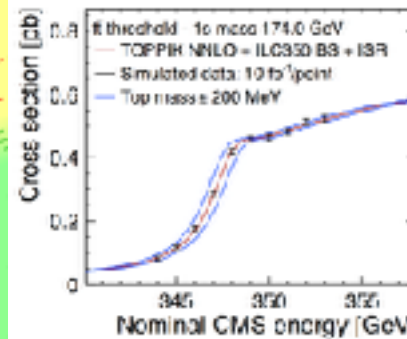
Does λ go to negative below Λ_P ?
 or $\lambda(\Lambda_P) = 0$?

To answer this, we need
precision m_t measurement!

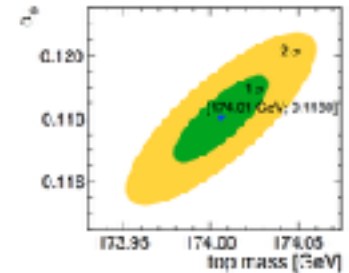
At LHC, theory error limits the precision to $\sim 500\text{MeV}$.



$T\bar{T}$ Threshold Scan @ ILC



Theoretically very clean measurement of m_t



$\Delta m_t(\overline{MS}) \lesssim 50 \text{ MeV}$
 $\Delta m_H = 30 \text{ MeV}$
ILC pinpoints the vacuum location

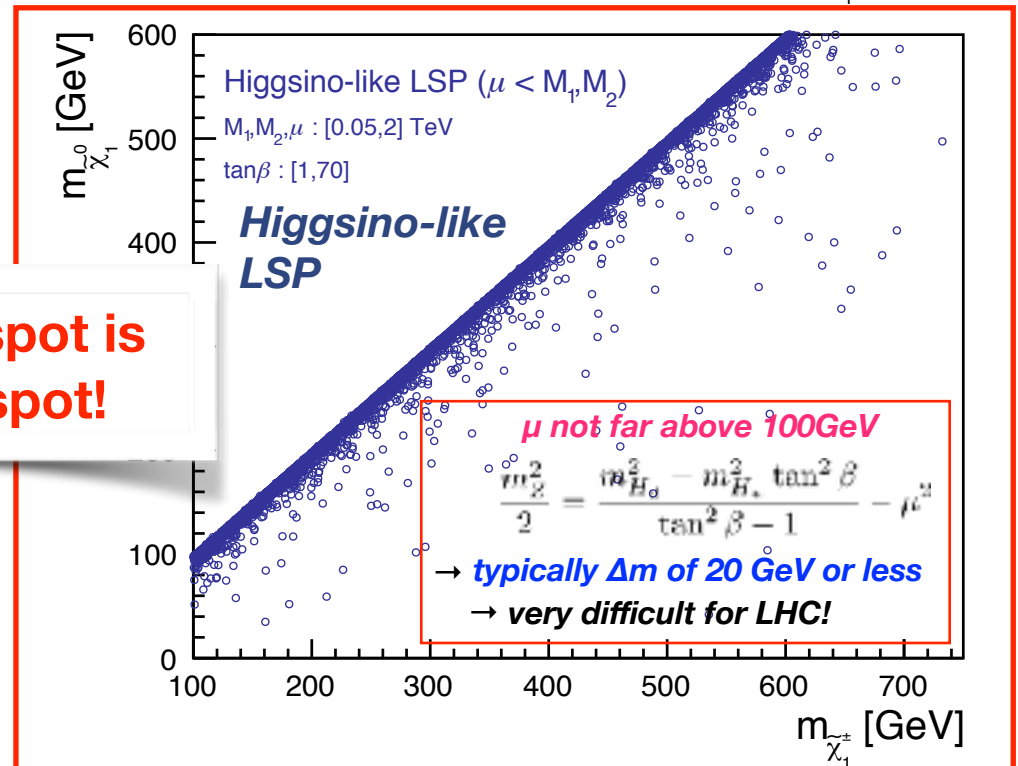
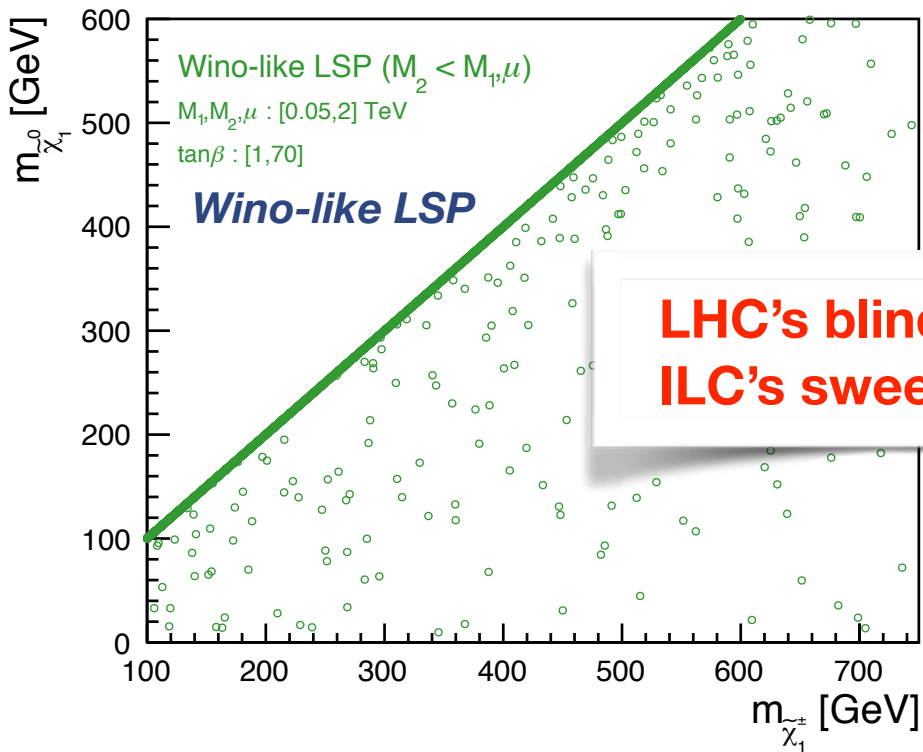
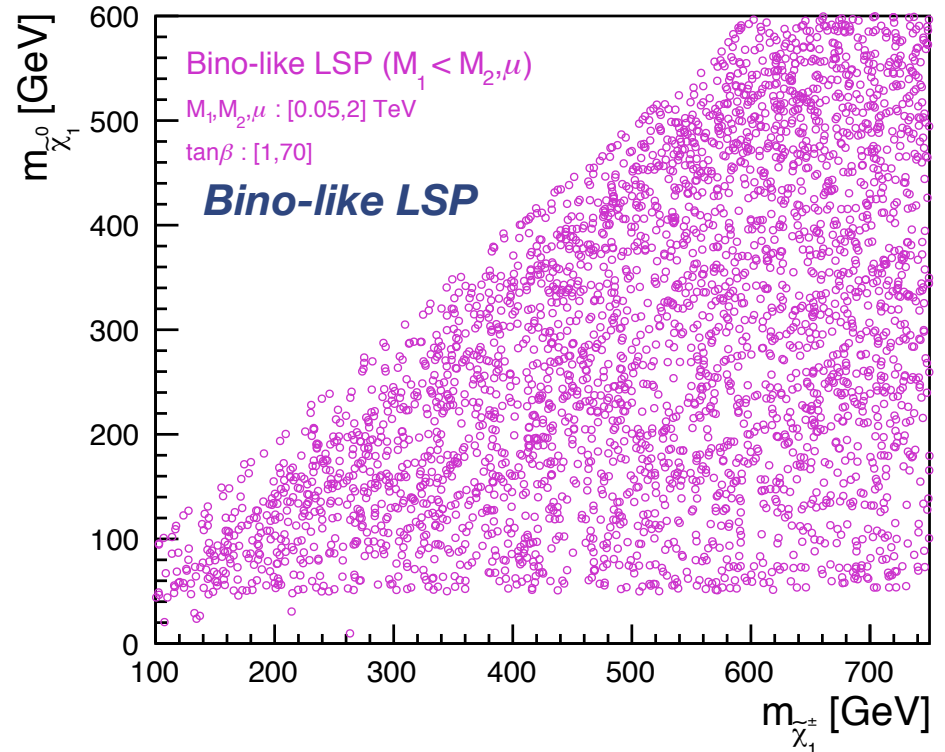
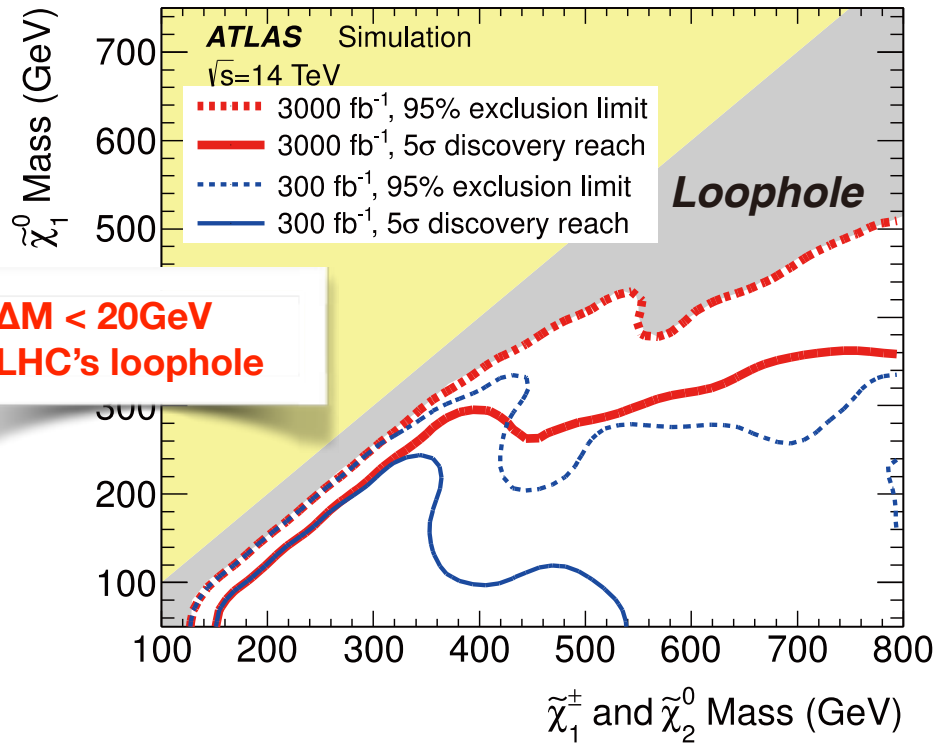
Direct Searches
for
New Particles

ILC, too, is an energy frontier machine!

*It will enter **uncharted waters of e^+e^- collisions***

Thanks to well-defined initial states,
clean environment w/o QCD BG, and polarized beams
ILC can cover blind spots of LHC

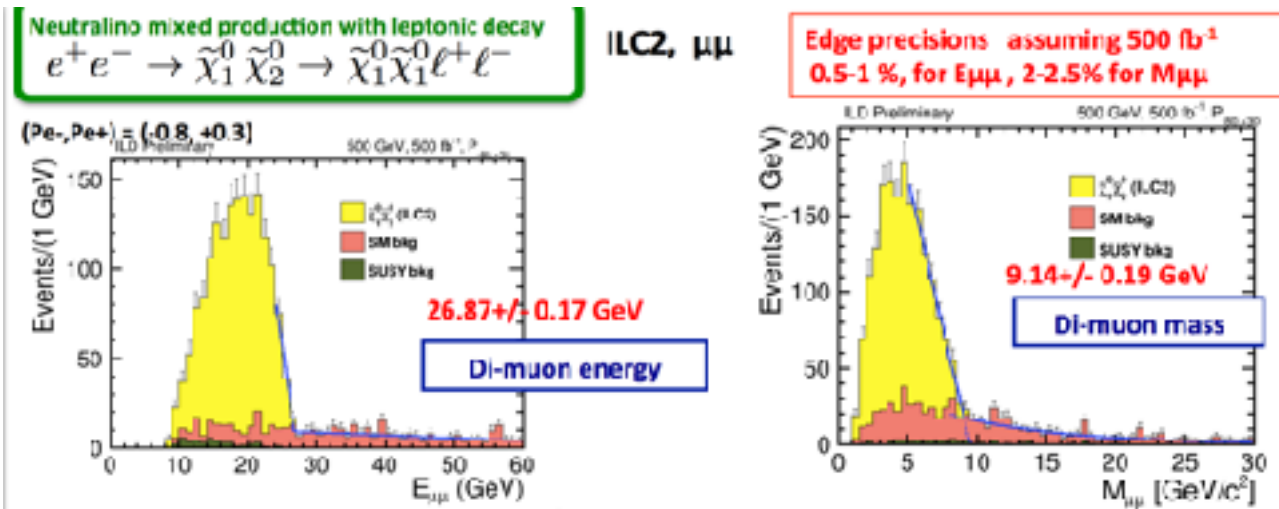
Chargino Search



LHC's blind spot is ILC's sweet spot!

BSM Talks from ILD at LCWS 2016

1. Generic WIMP searches (Moritz Habermehl) → Dec.8 (T. Tanabe)
2. SUSY co-annihilation (Mikael Berggren) → Dec.8
3. Higgsinos (Jacqueline Yan) → Dec.8
4. SUSY parameters from Higgsinos (Suvi-Leena Lehtinen) → Dec.8



J. Yan : LCWS2016

End points → M_X

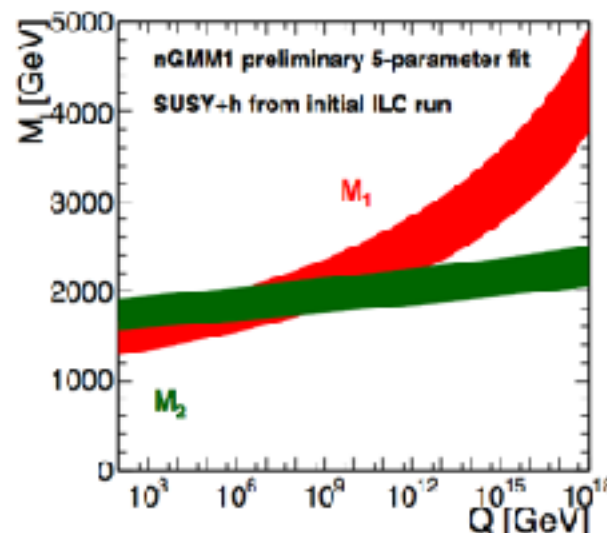
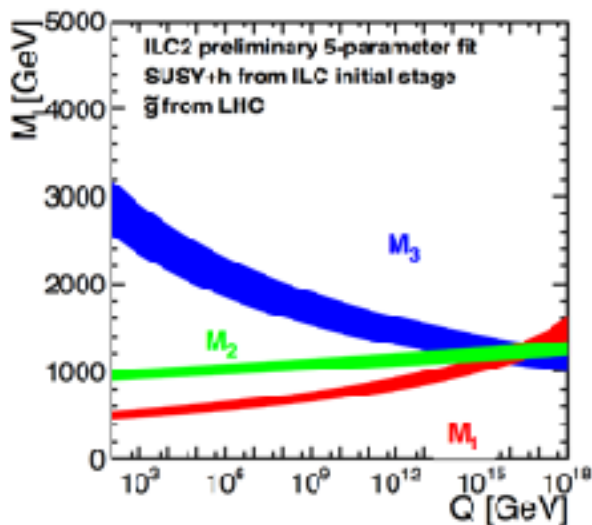
“ILC2 benchmark”: $\Delta M \sim 10$ GeV
 $\sigma_M / M < 1\%$ (H20)

500GeV

ILC1: 250GeV
 ILC2: 350GeV

Mass [GeV]	1.1	1.2	1.3
M_1	102.7	140.1	151.4
M_2	1019	1010	1000
M_3	1173	1065	1053
M_4	16.5	10.1	7.9

Probing very high scale physics



S. Lehtinen : LCWS2016

Left: Test of gaugino mass unification

Right: Select SUSY breaking models (gravity mediated SUSY breaking vs mirage unification)

WIMP Dark Matter Search @ ILC

Weakly Interacting Massive Particle

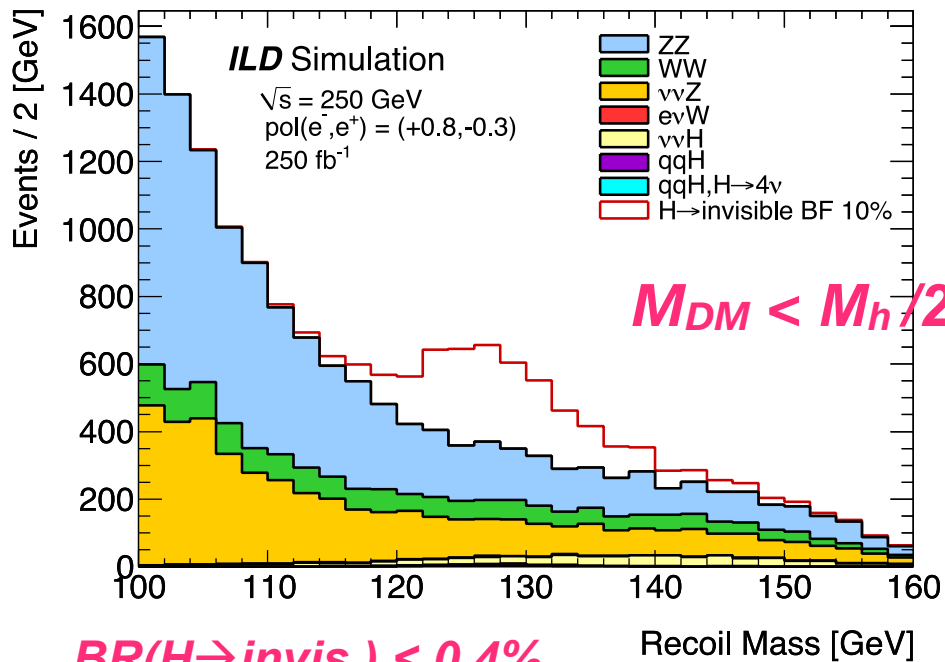
Decay of a new particle to Dark Matter (DM)

DM has a charged partner in many new physics models.

SUSY: The Lightest SUSY Particle (LSP) = DM → Its partner decays to a DM.

- Events with missing Pt (example: light chargino: see the previous page)

Higgs Invisible Decay

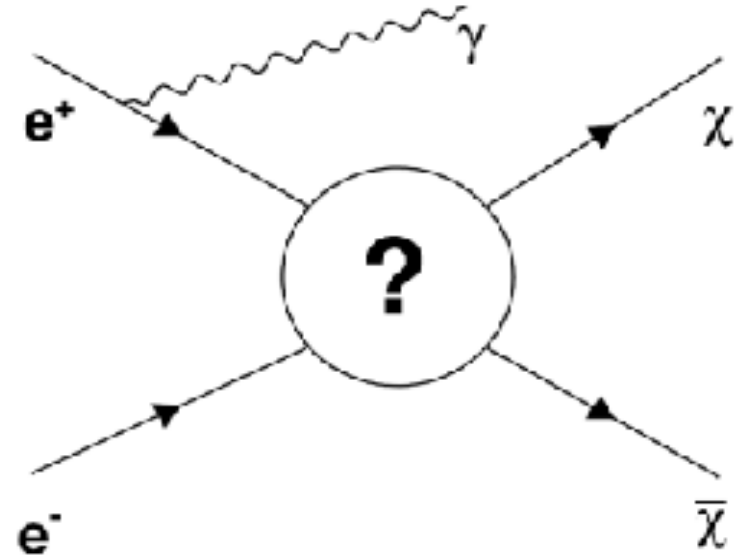


$BR(H \rightarrow \text{invis.}) < 0.4\%$

at 250 GeV, 1150 fb⁻¹ (<0.3% at 95%CL: H20)

Possible to access BR_{inv} to 0.4%!

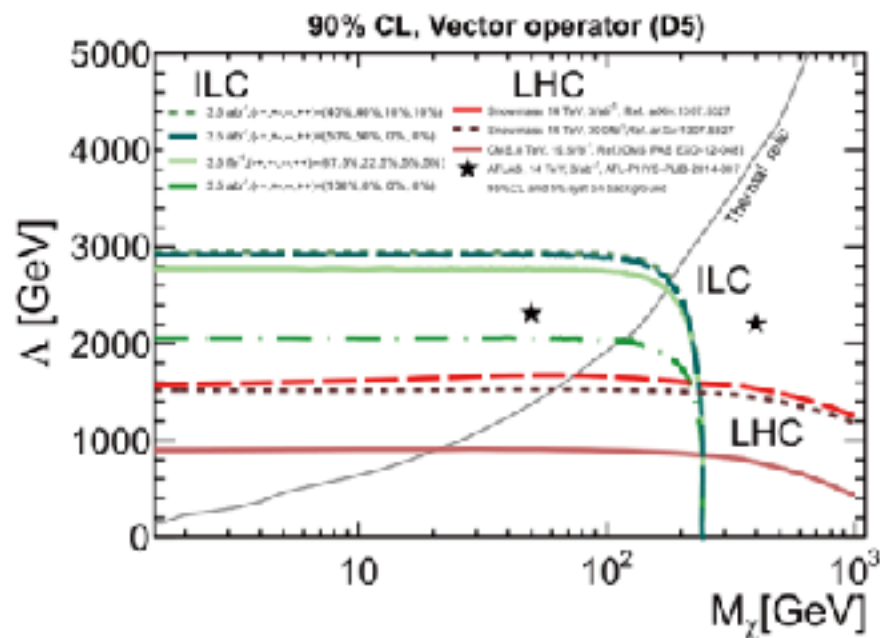
Mono-photon Search



→ $M_{DM} \text{ reach} \sim E_{cm} / 2$

Possible to access DM to $\sim E_{cm} / 2!$

DM: WIMP Searches



Previous result

LHC-ILC Comparison [A. Chau]

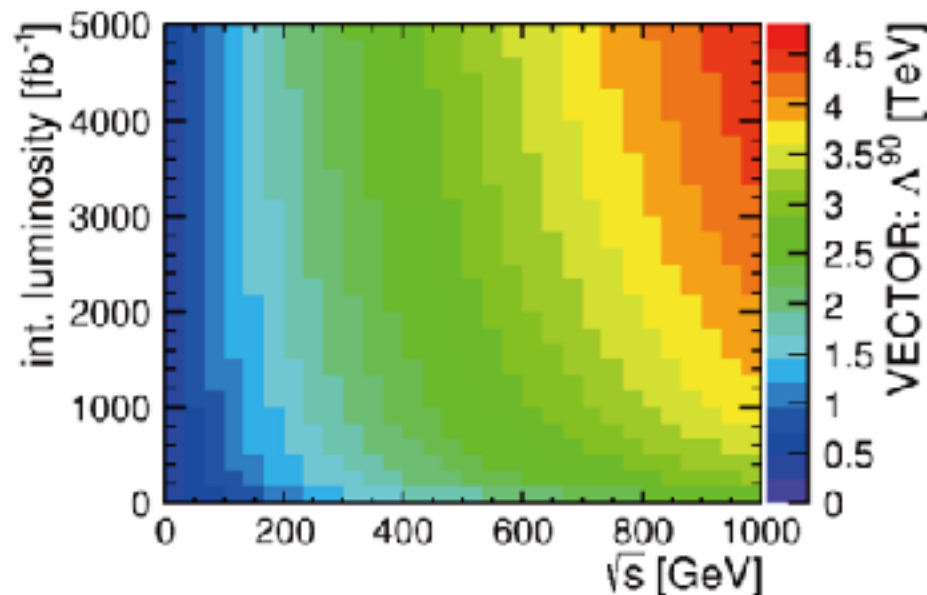
Example: Vector operator

- LHC sensitive to higher mass
- ILC sensitive to higher Λ

Recent result

Extrapolation to other \sqrt{s} [M. Habermehl]

- ILC reach of Λ at different CM energies and integrated luminosities
- for small M_χ (< 100 GeV)
- Allows study of run scenarios



ILC's H20 run scenario allows us to access Λ up to 3 ~ 4 TeV

Additional Slides

Higgs

Why 500 GeV?

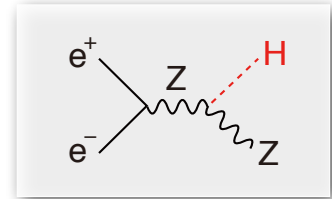
Higgs-related Physics at $E_{cm} \approx 500 \text{ GeV}$

Three well know thresholds

ZH @ 250 GeV ($\sim M_Z + M_H + 20 \text{ GeV}$) :

- Higgs mass, width, J^{PC}
- Gauge quantum numbers
- Absolute measurement of HZZ coupling (**recoil mass**)
- BR($h \rightarrow VV, qq, ll, \text{invisible}$) : $V=W/Z(\text{direct}), g, \gamma(\text{loop})$

\rightarrow Higgs couplings (other than top)

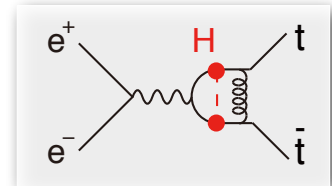


$t\bar{t}$ @ 340-350 GeV ($\sim 2m_t$) : ZH meas. Is also possible

- Threshold scan \rightarrow theoretically clean m_t measurement:
 - \rightarrow test stability of the SM vacuum
 - \rightarrow indirect meas. of top Yukawa coupling
- A_{FB} , Top momentum measurements
- Form factor measurements

$$\Delta m_t(\overline{MS}) \simeq 100 \text{ MeV}$$

$\gamma\gamma \rightarrow HH$ @ 350 GeV possibility



$v\bar{v}H$ @ 350 - 500 GeV :

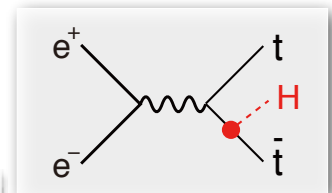
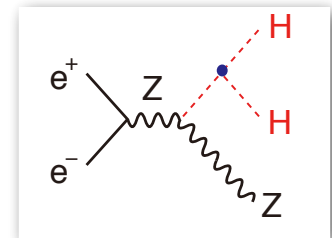
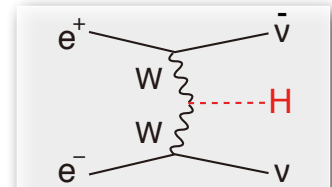
- **HWW coupling** \rightarrow total width \rightarrow absolute normalization of Higgs couplings

ZHH @ 500 GeV ($\sim M_Z + 2M_H + 170 \text{ GeV}$) :

- Prod. cross section attains its maximum at around 500 GeV \rightarrow Higgs self-coupling

$t\bar{t}H$ @ 500 GeV ($\sim 2m_t + M_H + 30 \text{ GeV}$) :

- Prod. cross section becomes maximum at around 800 GeV.
- QCD threshold correction enhances the cross section \rightarrow top Yukawa measurable at 500 GeV concurrently with the self-coupling



We can access all the relevant Higgs couplings at $\sim 500 \text{ GeV}$ for the mass-coupling plot!

Higgs Physics at Higher Energy

Self-coupling with WBF, top Yukawa at xsection max., other higgses, ...

vvH @ $\sqrt{s} > 1\text{TeV}$: $> 1\text{ab}^{-1}$ (pol e^+, e^-)=(+0.2,-0.8)

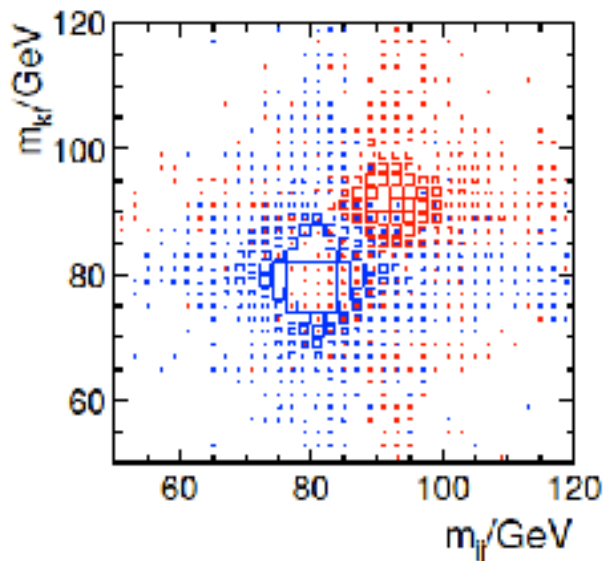
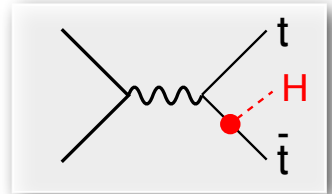
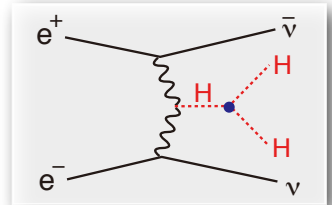
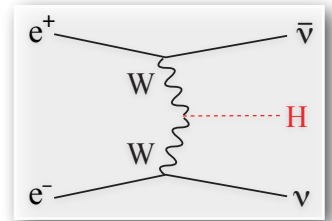
- allows us to measure rare decays such as $H \rightarrow \mu^+ \mu^-$, ...
- further improvements of coupling measurements

vvHH @ 1TeV or higher : 2ab^{-1} (pol e^+, e^-)=(+0.2,-0.8)

- cross section increases with E_{cm} , which compensates the dominance of the background diagrams at higher energies, thereby giving a better precision for the self-coupling.
- If possible, we want to see the running of the self-coupling (very very challenging).

ttbarH @ 1TeV : 1ab^{-1}

- Prod. cross section becomes maximum at around 800GeV.
- CP mixing of Higgs can be unambiguously studied.



Obvious but most important advantage of higher energies in terms of Higgs physics is, however, its **higher mass reach to other Higgs bosons** expected in extended Higgs sectors and **higher sensitivity to $W_L W_L$ scattering** to decide whether the Higgs sector is strongly interacting or not.

In any case we can improve the mass-coupling plot by including the data at 1TeV!

Total Width and Coupling Extraction

One of the major advantages of the LC

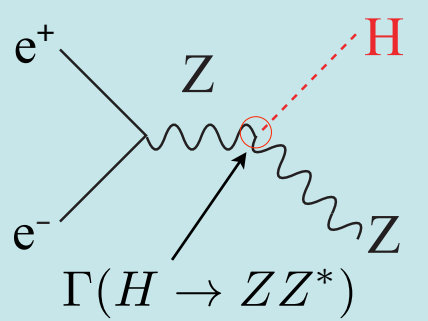
To extract couplings from BRs, we need the total width:

$$g_{HAA}^2 \propto \Gamma(H \rightarrow AA) = \Gamma_H \cdot BR(H \rightarrow AA)$$

To determine the total width, we need at least one partial width and corresponding BR:

$$\Gamma_H = \Gamma(H \rightarrow AA) / BR(H \rightarrow AA)$$

In principle, we can use $A=Z$, or W for which we can measure both the BRs and the couplings:

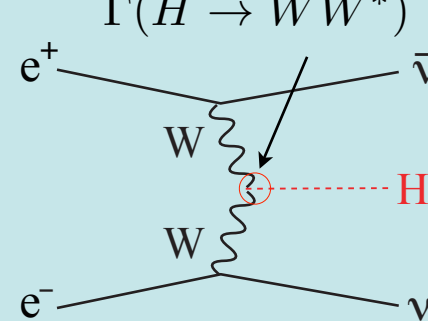


$BR(H \rightarrow ZZ^*)$

$\Gamma(H \rightarrow ZZ^*)$

BR=O(1%): precision limited by low stat. for H->ZZ* events

$250 \text{ fb}^{-1} @ 250 \text{ GeV}$
 $\Delta\Gamma_H / \Gamma_H \simeq 20\%$



$\Gamma(H \rightarrow WW^*)$

$BR(H \rightarrow WW^*)$

More advantageous but not easy at low E

$250 \text{ fb}^{-1} @ 250 \text{ GeV}$
 $\Delta\Gamma_H / \Gamma_H \simeq 11\%$

C.F.Durig, Helmholtz Alliance 6th WS,
Dec. 2012

Model-independent Global Fit for Couplings

33 σ_{BR} measurements (Y_i) and σ_{ZH} ($Y_{34,35}$)

$$\chi^2 = \sum_{i=1}^{35} \left(\frac{Y_i - Y'_i}{\Delta Y_i} \right)^2$$

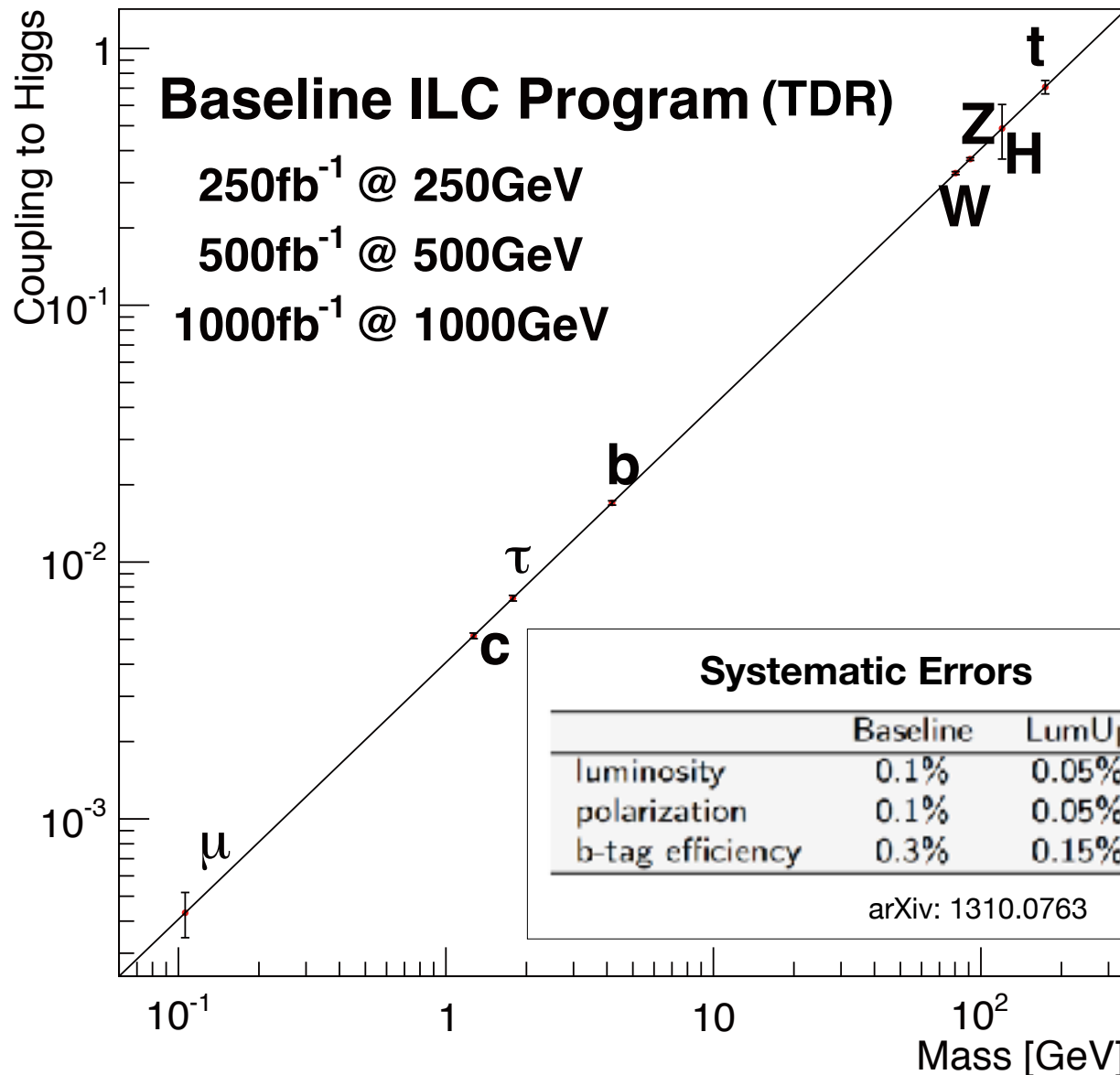
$$Y'_i = F_i \cdot \frac{g_{HA_i A_i}^2 \cdot g_{HB_i B_i}^2}{\Gamma_0}$$

($i = 1, \dots, 33$)
 ($A_i = Z, W, t$)
 ($B_i = b, c, \tau, \mu, g, \gamma, Z, W$: decay)

$$F_i = S_i G_i$$

$G_i = \left(\frac{\Gamma_i}{g_i^2} \right)$

$$S_i = \left(\frac{\sigma_{ZH}}{g_{HZZ}^2} \right), \left(\frac{\sigma_{\nu\bar{\nu}H}}{g_{HWW}^2} \right), \text{ or } \left(\frac{\sigma_{t\bar{t}H}}{g_{Htt}^2} \right)$$



ILC's precisions will eventually reach sub-% level!

Independent Higgs Measurements at ILC

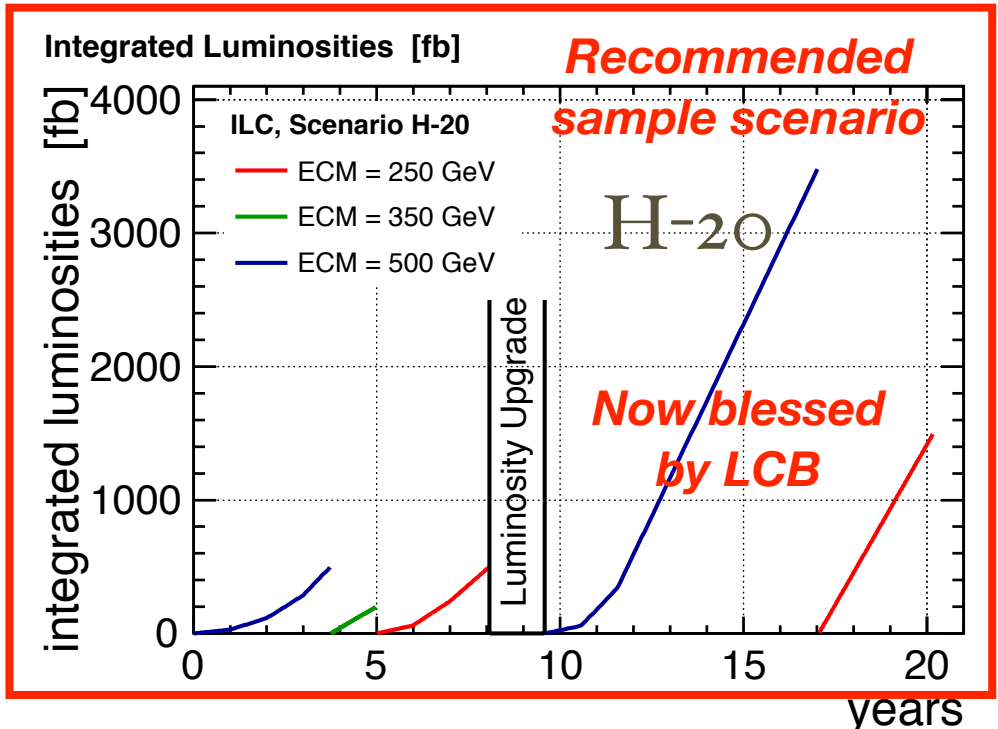
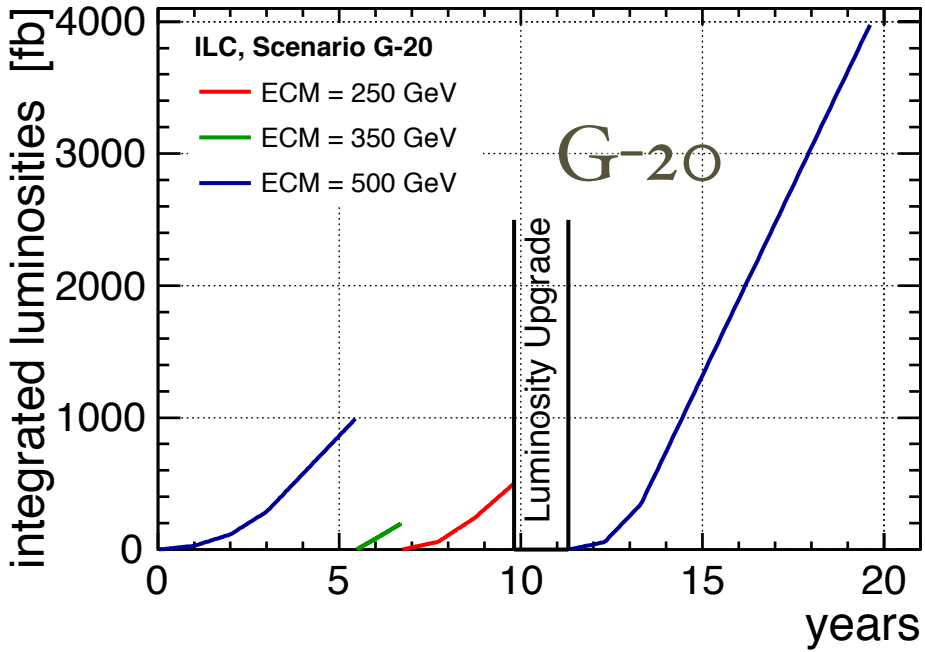
Baseline (=TDR) ILC program

250 GeV: 250 fb⁻¹
 500 GeV: 500 fb⁻¹
 1 TeV: 1000 fb⁻¹

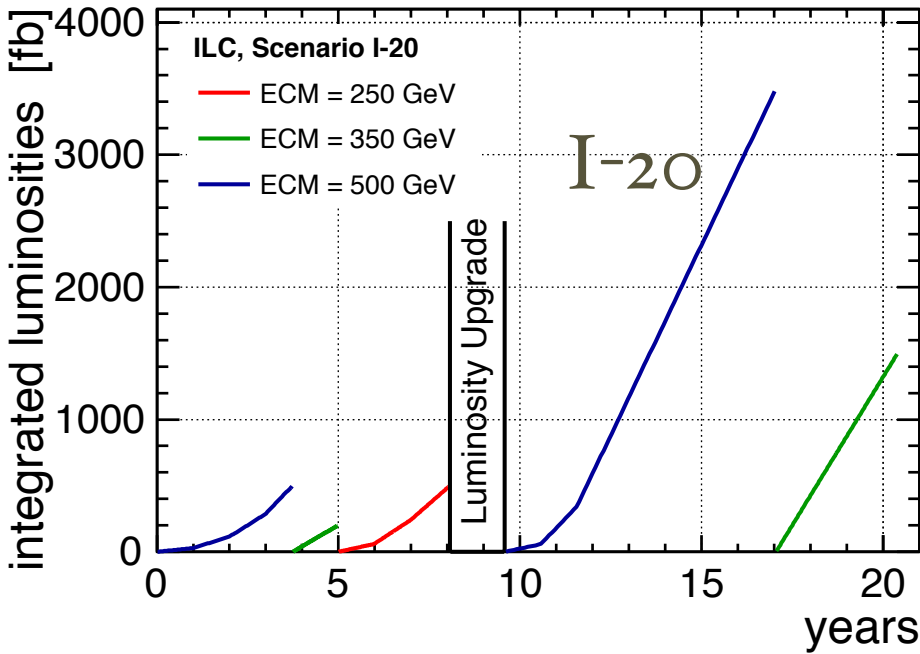
(M_H = 125 GeV)

Ecm	250 GeV		500 GeV		1 TeV
luminosity [fb ⁻¹]	250		500		1000
polarization (e ⁻ ,e ⁺)	(-0.8, +0.3)		(-0.8, +0.3)		(-0.8, +0.2)
process	ZH	vvH(fusion)	ZH	vvH(fusion)	vvH(fusion)
cross section	2.6%	-	3%	-	
	σ·Br	σ·Br	σ·Br	σ·Br	σ·Br
H→bb	1.2%	10.5%	1.8%	0.66%	0.32%
<u>H→cc</u>	8.3%		13%	6.2%	3.1%
H→gg	7%		11%	4.1%	2.3%
H→WW*	6.4%		9.2%	2.4%	1.6%
H→ττ	3.2%		5.4%	9%	3.1%
H→ZZ*	19%		25%	8.2%	4.1%
H→γγ	34%		34%	19%	7.4%
H→μμ	72%	-	88%	72%	31%
tth/H→bb	-		28% (12%@550GeV)		6.2%

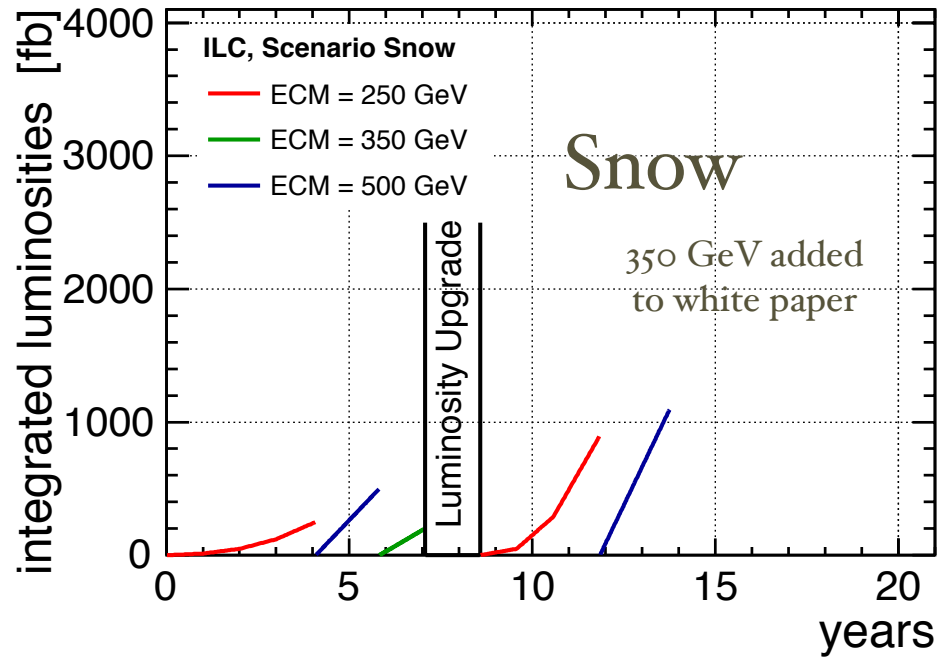
Integrated Luminosities [fb]



Integrated Luminosities [fb]



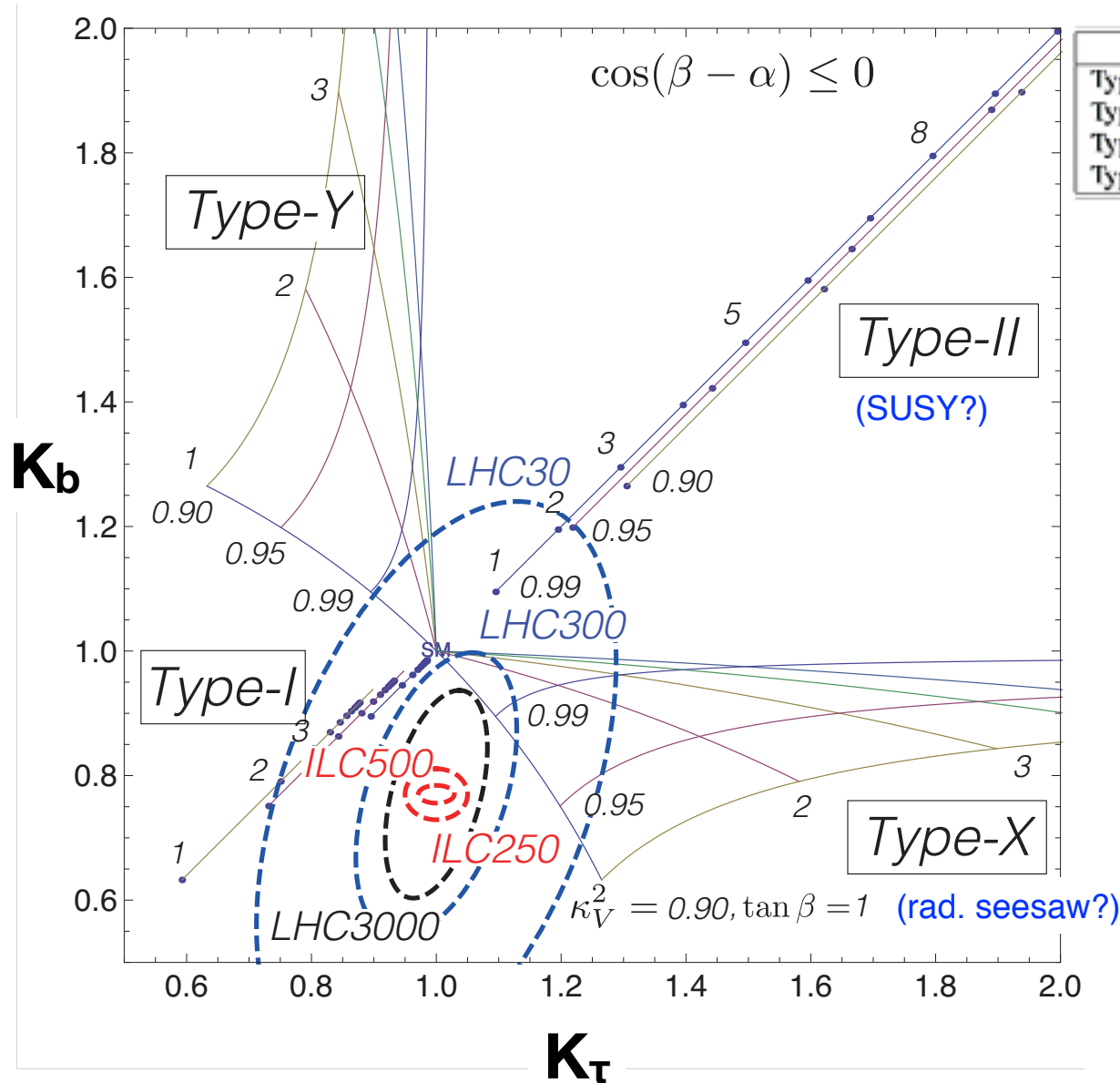
Integrated Luminosities [fb]



Fingerprinting

2HDM

Multiplet Structure



	Φ_1	Φ_2	ν_R	d_R	ℓ_R	Q_L, L_L
Type I	+	-	-	-	-	+
Type II (SUSY)	+	-	-	+	+	+
Type X (Lepton-specific)	+	-	-	-	+	+
Type Y (Flipped)	+	-	-	+	-	+

4 Possible Z_2 Charge Assignments that forbids tree-level Higgs-induced FCNC

$$\kappa_V^2 = \sin(\beta - \alpha)^2 = 1 \Leftrightarrow \text{SM}$$

Given a deviation of the Higgs to Z coupling: $\Delta \kappa_V^2 = 1 - \kappa_V^2 = 0.01$ we will be able to **discriminate the 4 models!**

Model-dependent
7-parameter fit
ILC: Baseline lumi.

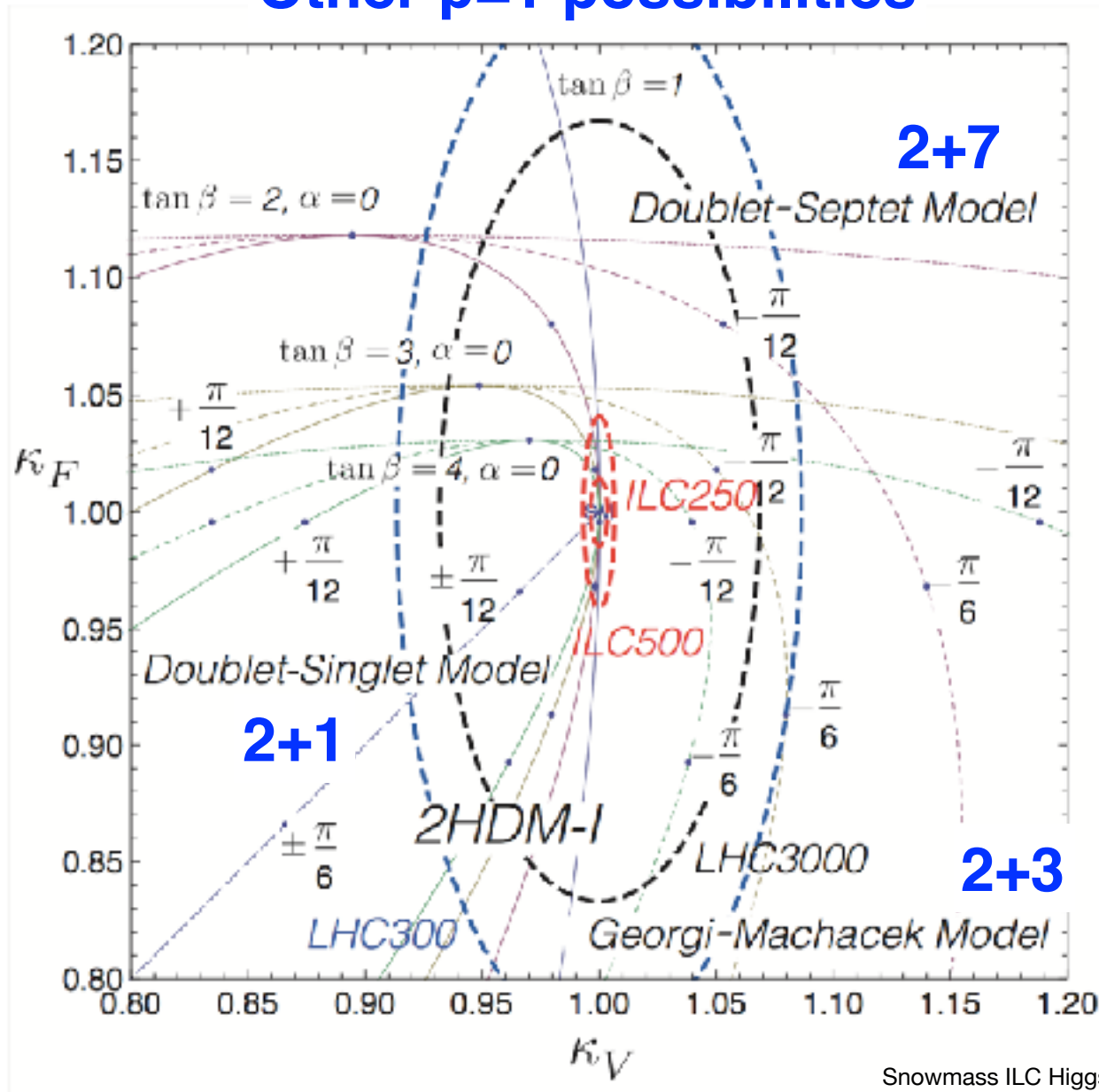
ILC TDR

Snowmass ILC Higgs White Paper (arXiv: 1310.0763)

Kanemura et al (arXiv: 1406.3294)

Multiplet Structure

Other $\rho=1$ possibilities



Snowmass ILC Higgs White Paper (arXiv: 1310.0763)

Kanemura et al (arXiv: 1406.3294)

Figure 1.18. The scaling factors in models with universal Yukawa coupling constants.

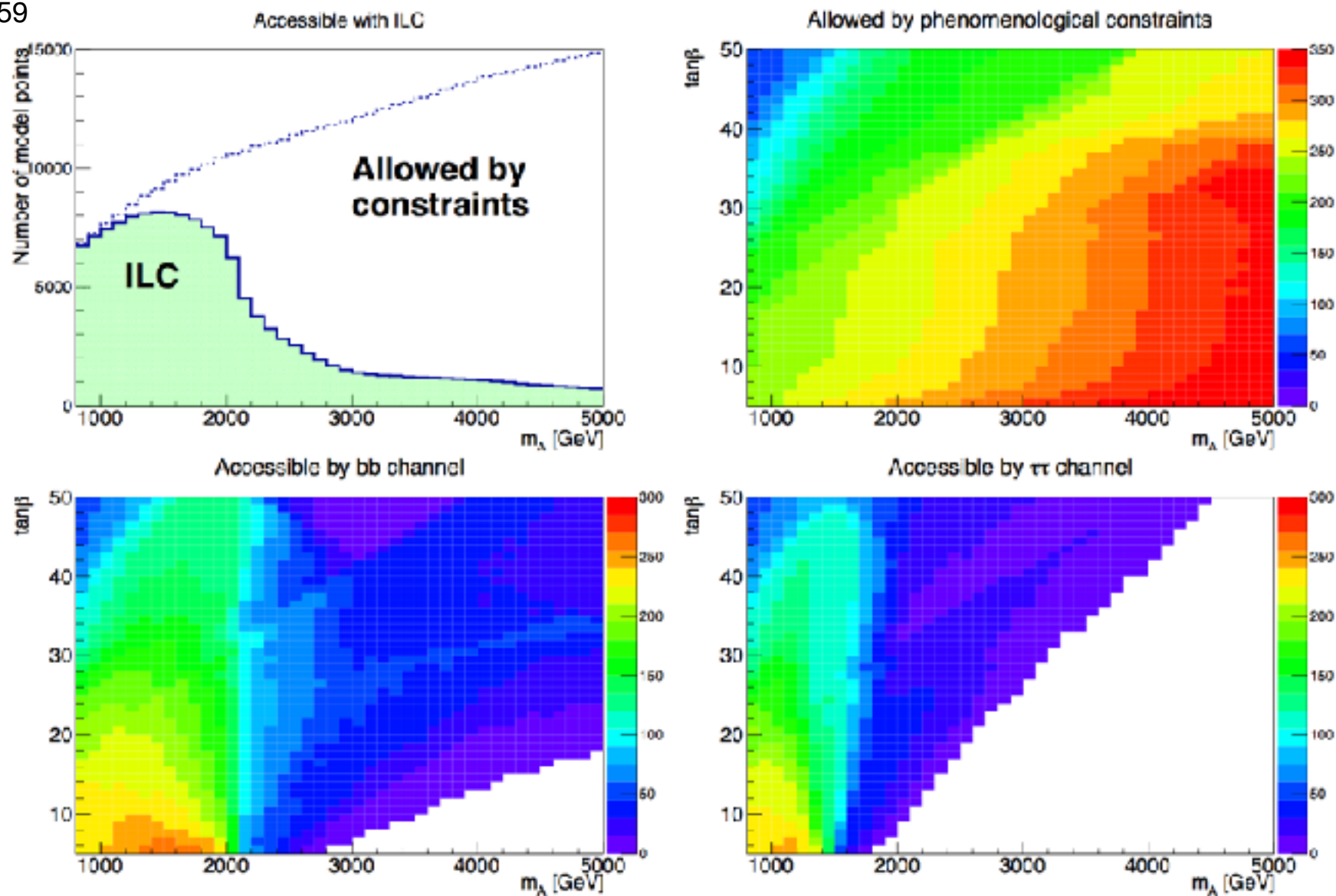
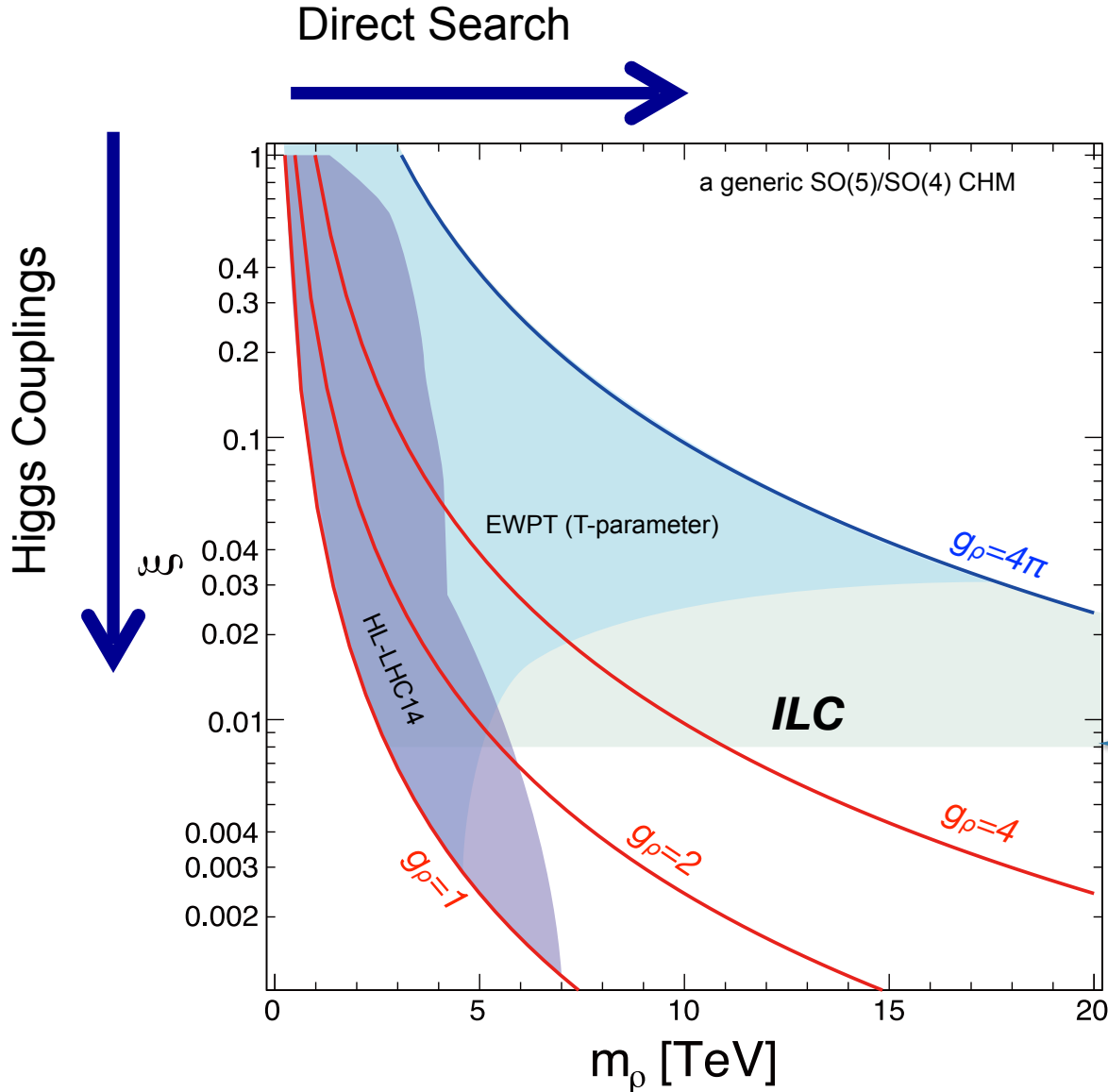
Motoi Endo^(a,b), Takeo Moroi^(a,b), and Mihoko M. Nojiri^(b,c,d)

Figure 8: Upper-left: The number of model points accessible with ILC by at least one decay mode of h as a function of m_A (green histogram), as well as that of model points allowed by the phenomenological constraints (dotted histogram). Upper-right: The number of model points allowed by the phenomenological constraints on m_A vs. $\tan\beta$ plane. Lower-left: The number of model points accessible with ILC by $h \rightarrow bb$. Lower-right: The number of model points accessible with ILC by $h \rightarrow \bar{\tau}\tau$.

Composite Higgs: Reach

Complementary approaches to probe composite Higgs models

- Direct search for heavy resonances at the LHC
 - Indirect search via Higgs couplings at the ILC
- Comparison depends on the coupling strength (g_*)



Based on Contino, et al, JHEP 1402 (2014) 006
Torre, Thamm, Wulzer 2014
Grojean @ LCWS 2014

$$\xi = \frac{g_\rho^2}{m_\rho^2} v^2 = \frac{v^2}{f^2}$$

$$\frac{g_{hVV}}{g_{h_{SM}VV}} = \sqrt{1 - \xi}$$

ILC (250+500 LumiUP)

$$\Delta \frac{g_{hVV}}{g_{hVV}} = 0.4\%$$

New resonance scale and fingerprint identification in minimal composite Higgs models

Shinya Kanemura,¹ Kunio Kaneda,² Naoki Machida,¹ and Tetsuo Shindou³

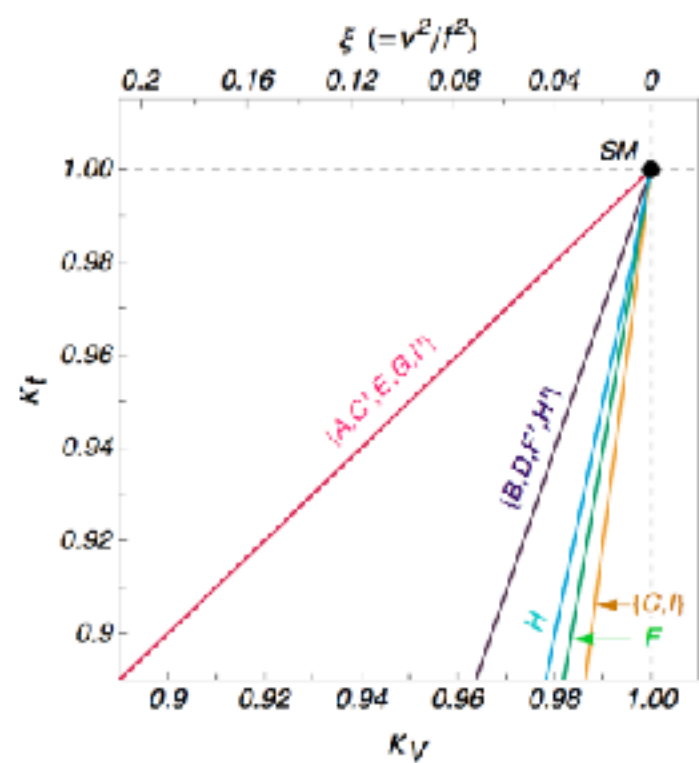
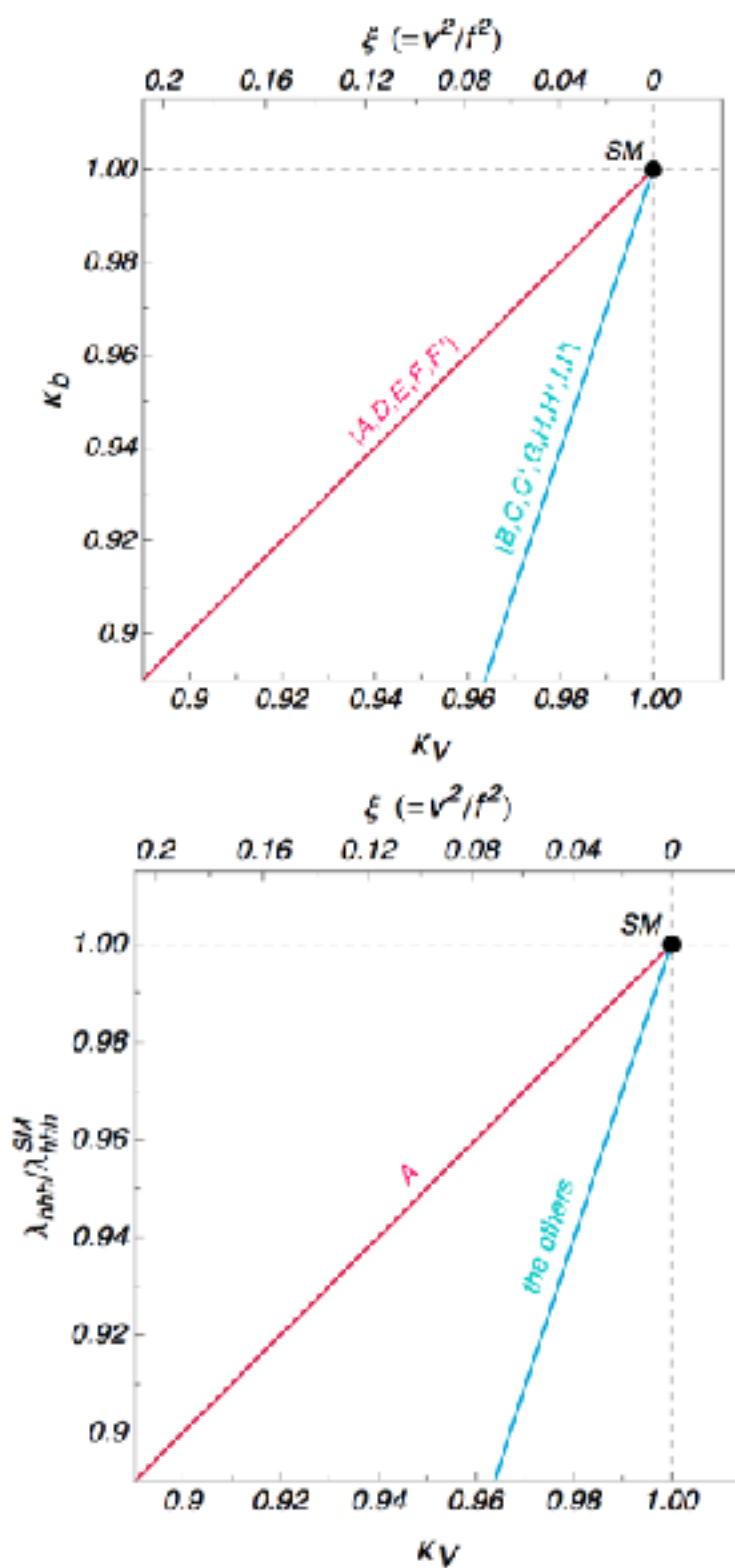
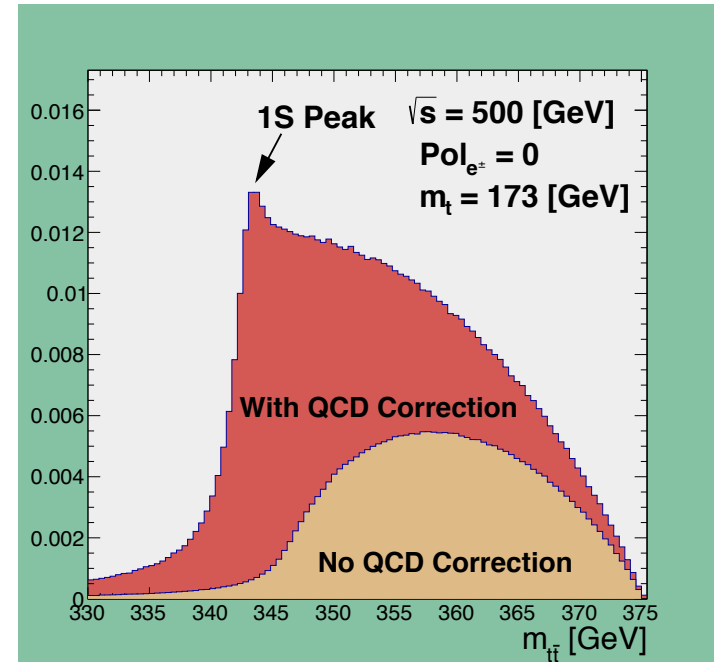
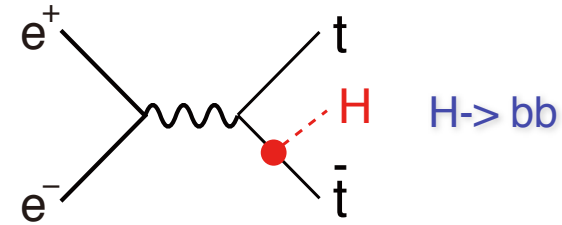
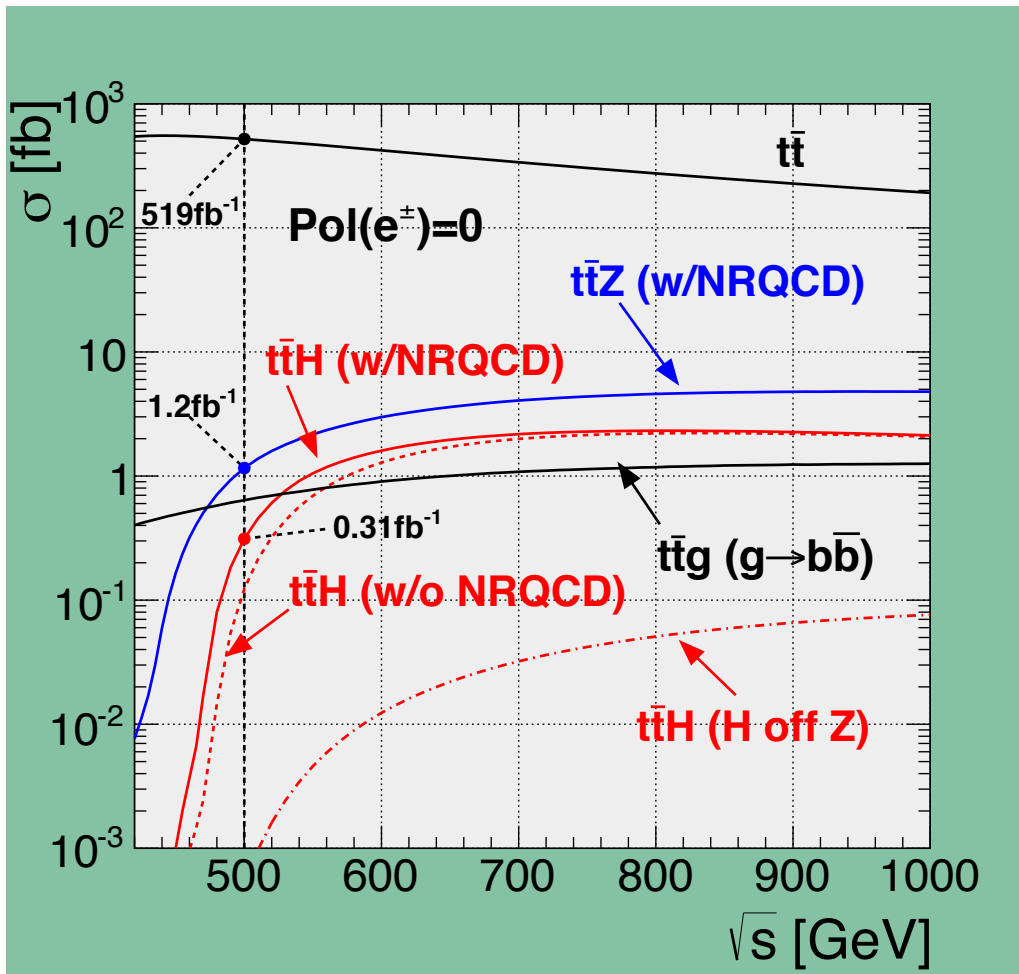


TABLE I: Scale factors for MCHMs with various matter representations. The labels are used in Fig. 7, where C, E and I are the case of $M_1^2 \rightarrow 0$, and C', H' and I' are the case of $M_2^2 \rightarrow 0$.

Label	Model	κ_V	$\kappa_{\Delta\Delta V}$	$\kappa_{\Delta\Delta\Delta}$	$\kappa_{\Delta\Delta\Delta}$	κ_t	κ_b	$\kappa_{\Delta\Delta}$	$\kappa_{\Delta\Delta}$
A	MCHM ₁	$\sqrt{1-\xi}$	$1-2\xi$	$\sqrt{1-\xi}$	$1-\frac{1}{2}\xi$	$\sqrt{1-\xi}$	$\sqrt{1-\xi}$	$-\xi$	$-\xi$
B	MCHM ₂	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(2+2\xi^2)/3}{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
B	MCHM ₁₀	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(2+2\xi^2)/3}{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
C, C'	MCHM ₁₆	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_2	F_3	$\frac{1-2\xi}{\sqrt{1-\xi}}$	F_3	4ξ
D	MCHM ₅₊₂₀	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(2+2\xi^2)/3}{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\sqrt{1-\xi}$	-4ξ	$-\xi$
E	MCHM ₁₊₁₆₊₁₆	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(2+2\xi^2)/3}{1-\xi}$	$\sqrt{1-\xi}$	$\sqrt{1-\xi}$	$-\xi$	$-\xi$
F, F'	MCHM ₃₊₁₆₊₁₆	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_2	F_3	$\sqrt{1-\xi}$	F_3	$-\xi$
G	MCHM ₁₀₊₁₀	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(2+2\xi^2)/3}{1-\xi}$	$\sqrt{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$-\xi$	-4ξ
B	MCHM ₁₀₊₁₆₊₁₆	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_3	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
B	MCHM ₁₀₊₁₆₊₁₆	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(2+2\xi^2)/3}{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
H, H'	MCHM ₁₊₁₊₁₆	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_3	F_3	$\frac{1-2\xi}{\sqrt{1-\xi}}$	F_3	-4ξ
B	MCHM ₁₀₊₁₆₊₁₆	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_2	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
I, I'	MCHM ₁₀₊₁₆₊₁₆	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_3	F_3	$\frac{1-2\xi}{\sqrt{1-\xi}}$	F_3	-4ξ

Top Yukawa Coupling

The largest among matter fermions, but not yet directly observed



A factor of 2 enhancement from QCD bound-state effects

Cross section maximum at around $E_{cm} = 800 \text{ GeV}$

Philipp Roloff, LCWS12

Tony Price, LCWS12

DBD Full Simulation

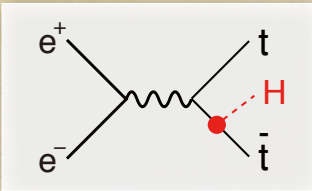
$$1 \text{ ab}^{-1} @ 500 \text{ GeV} \quad m_H = 125 \text{ GeV}$$

$$\Delta g_Y(t) / g_Y(t) = 9.9\%$$

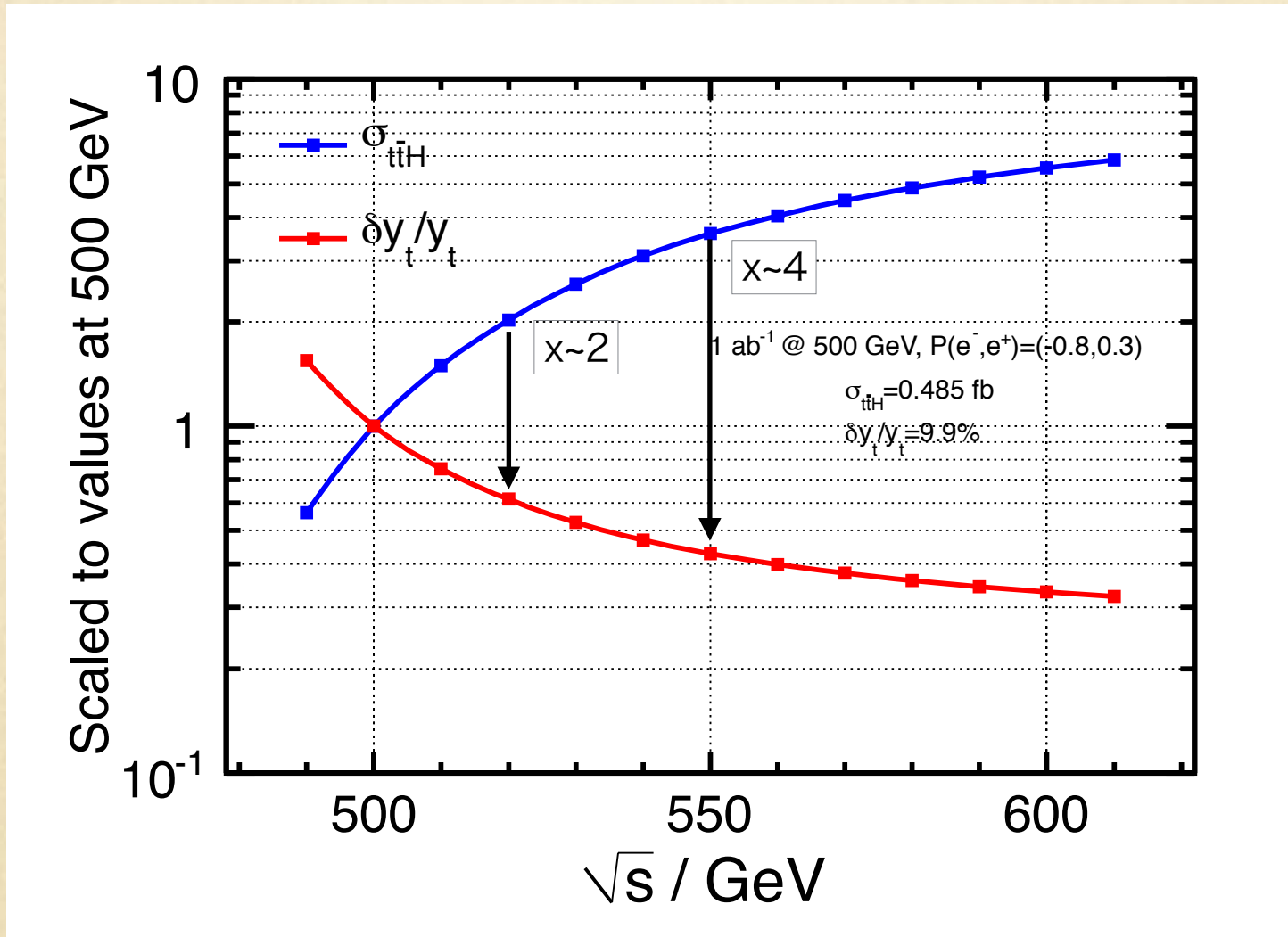
Tony Price, LCWS12

scaled from $m_H = 120 \text{ GeV}$

Notice $\sigma(500+20 \text{ GeV}) / \sigma(500 \text{ GeV}) \sim 2$
Moving up a little bit helps significantly!



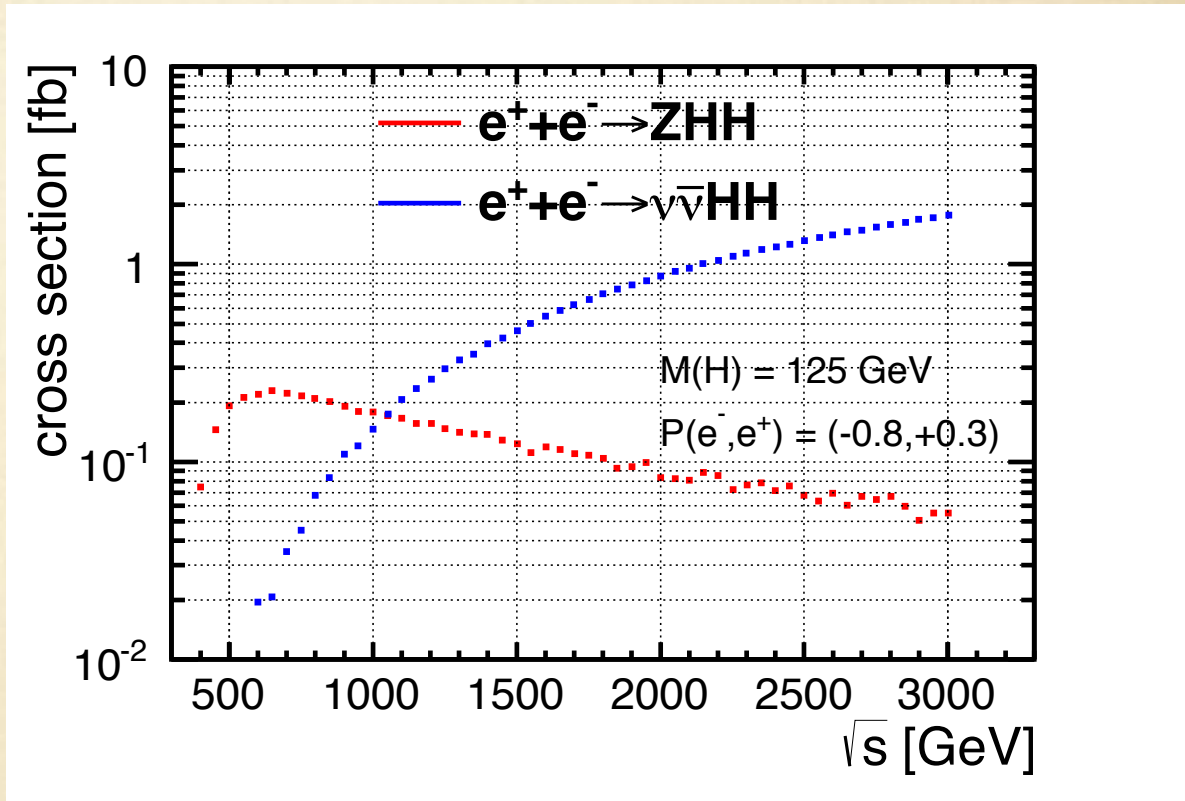
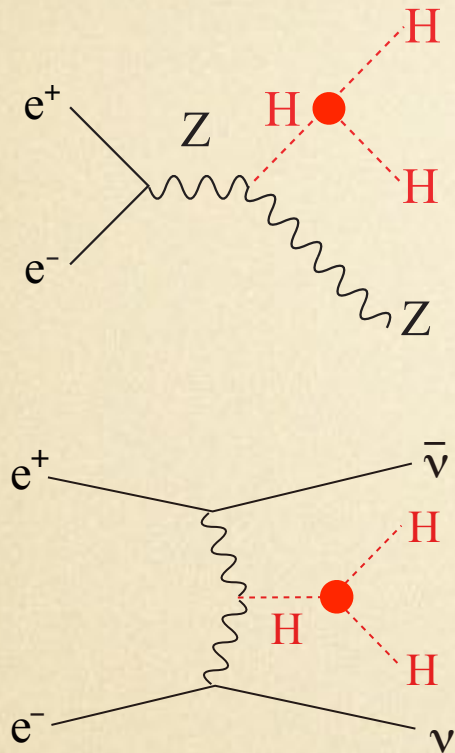
Top Yukawa coupling



Y. Sudo

Slight increase of E_{max} is very beneficial!

prospects of Higgs self-coupling @ linear colliders



prospects from full simulation studies:

ILC	$\Delta\lambda_{HHH}/\lambda_{HHH}$	500 GeV	+ 1 TeV
	Snowmass	46%	13%
	H20	27%	10%

(ref. H20 arXiv: 1506.07870)

J. Tian, LC-REP-2013-003

CLIC	1.4 TeV (1.5 ab^{-1})	+3 TeV (2 ab^{-1})
	21%	10%

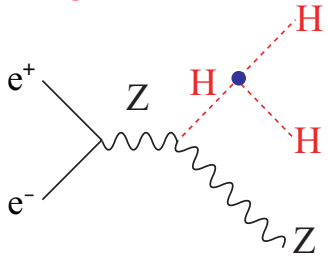
(arXiv: 1307.5288)

M. Kurata, LC-REP-2014-025

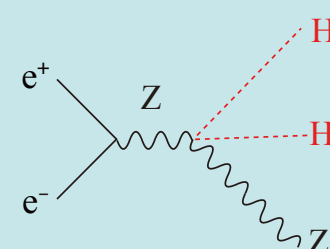
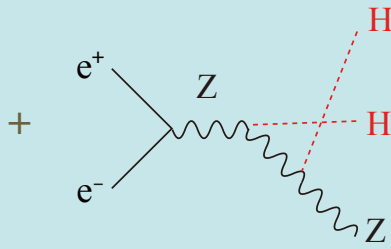
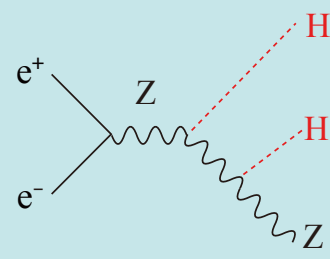
C. Dürig @ ALCW15

The Problem : BG diagrams dilute self-coupling contribution

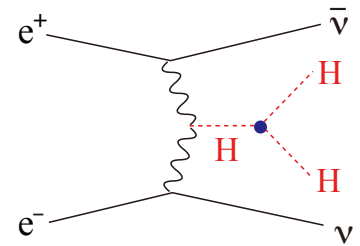
Signal diagram



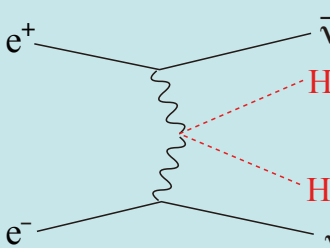
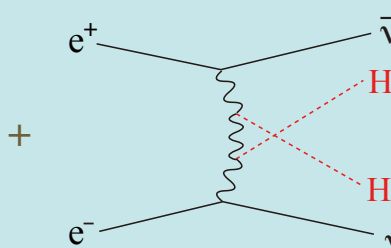
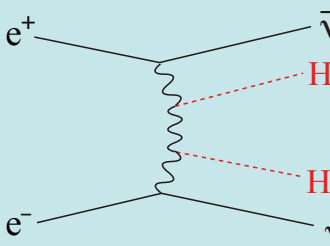
Irreducible BG diagrams



Signal diagram



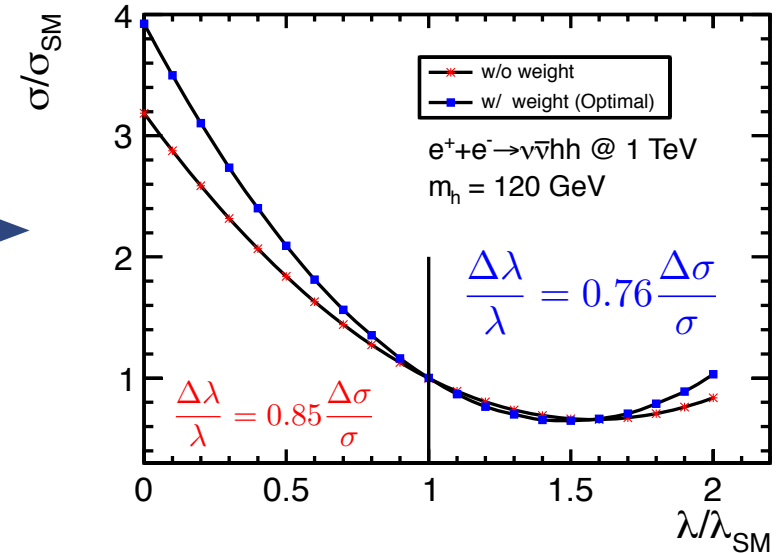
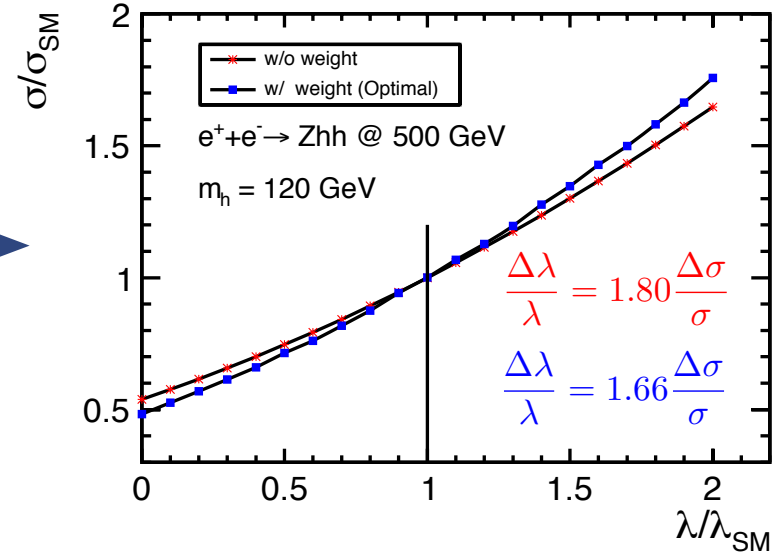
Irreducible BG diagrams



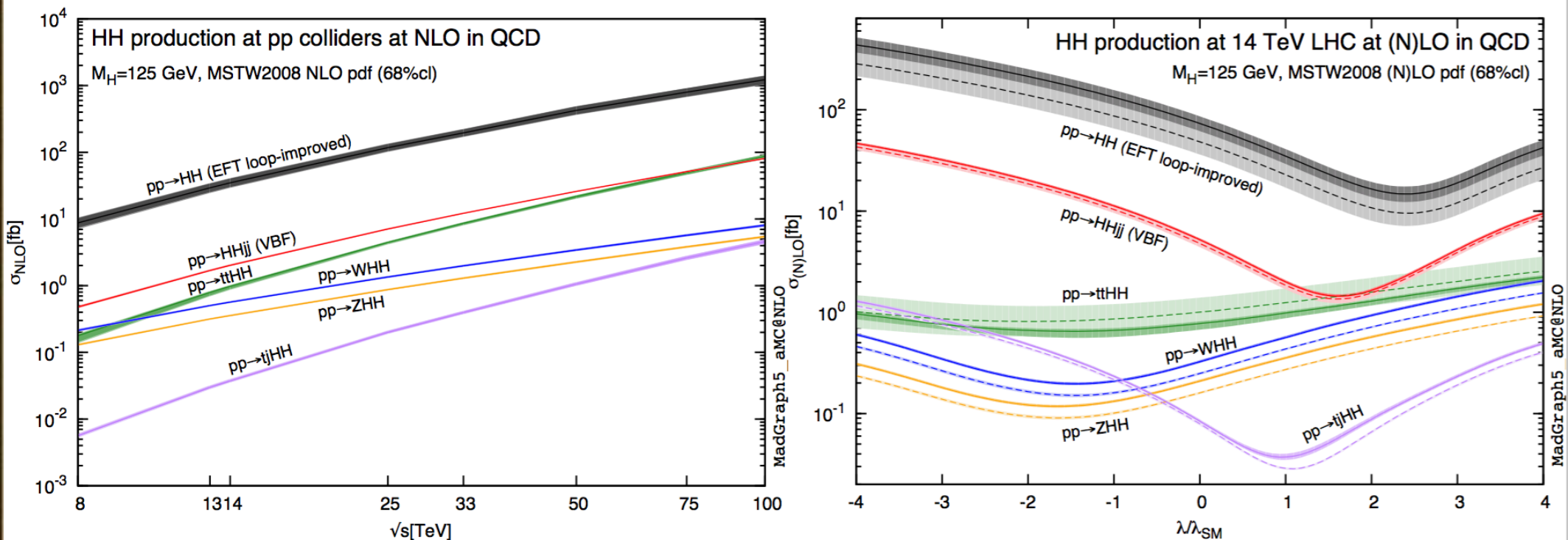
$$\sigma = \lambda^2 S + \lambda I + B$$

$$\frac{\Delta\lambda}{\lambda} = F \cdot \frac{\Delta\sigma}{\sigma}$$

$F=0.5$ if no BG diagrams



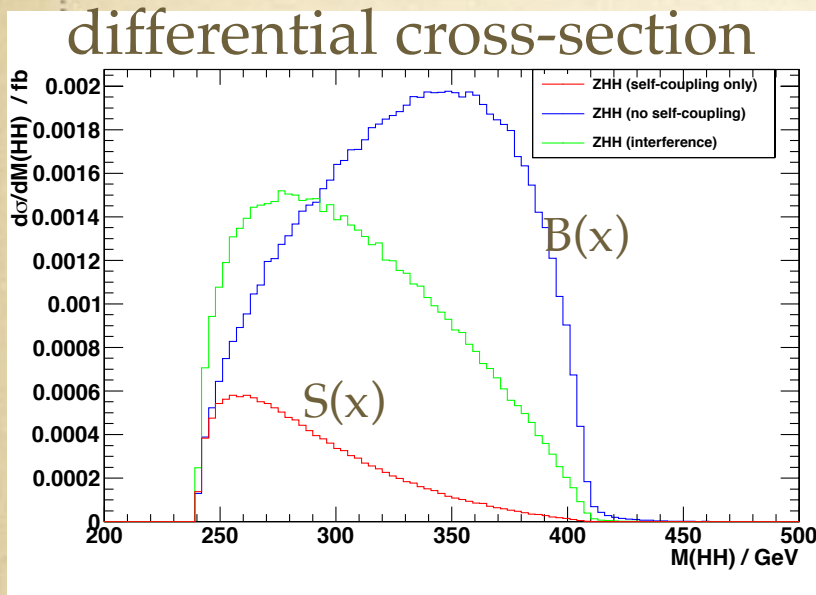
What if $\lambda \neq \lambda_{SM}$? @ LHC



arXiv:1401.7304

- interference is destructive, σ minimum at $\lambda \sim 2.5\lambda_{SM}$; if λ is enhanced, it's going to be very difficult (from snowmass study by 3000 fb⁻¹ @ 14 TeV, significance of double Higgs production is only $\sim 2\sigma$, if cross section decreases by a fact of 2~3, very challenging to observe pp→HH)

a new general method to improve the sensitivity of λ

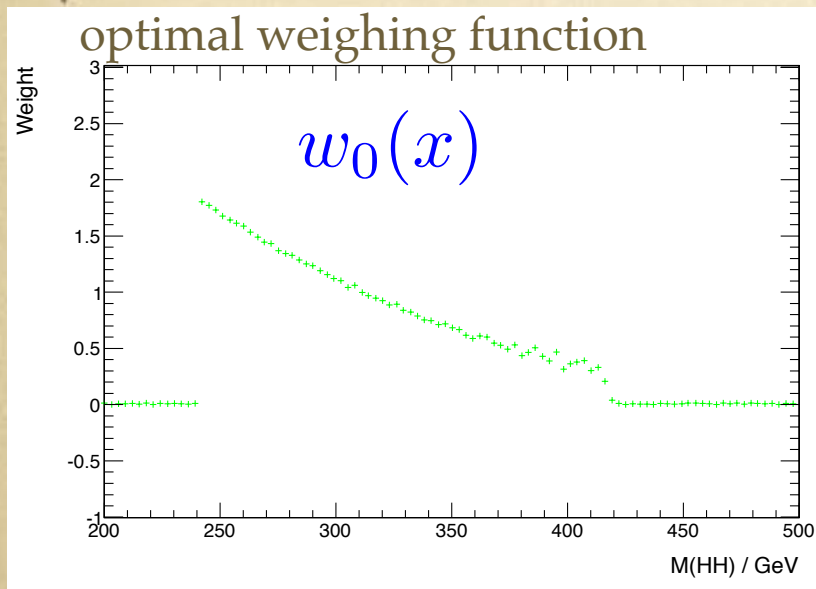


$$\frac{d\sigma}{dx} = B(x) + \lambda I(x) + \lambda^2 S(x)$$

irreducible
interference
self-coupling

observable: weighted cross-section

$$\sigma_w = \int \frac{d\sigma}{dx} w(x) dx$$



equation of the optimal $w(x)$ (variance principle):

$$\sigma(x)w_0(x) \int (I(x) + 2S(x))w_0(x)dx = (I(x) + 2S(x)) \int \sigma(x)w_0^2(x)dx$$

general solution:

$$w_0(x) = c \cdot \frac{I(x) + 2S(x)}{\sigma(x)}$$

c: arbitrary normalization factor

σ_{Zh} in EFT \rightarrow Composite Scale

The size comes from the scale of an EFT operator:

$$\mathcal{L} \supset \left(\frac{c_H}{\Lambda^2}\right) \frac{1}{2} (\partial_\mu |H|^2)^2$$

$$\rightarrow \left(\frac{2C_H v^2}{\Lambda^2}\right) \frac{1}{2} (\partial_\mu h)^2$$

250GeV

ILC direct Zh
(Yan et al. 1604.07524)

\sqrt{s}	250 GeV		350 GeV		500 GeV	
	$\int \mathcal{L} dt$	$\Delta\sigma_{Zh}/\sigma_{Zh}$	$\int \mathcal{L} dt$	$\Delta\sigma_{Zh}/\sigma_{Zh}$	$\int \mathcal{L} dt$	$\Delta\sigma_{Zh}/\sigma_{Zh}$
e^+e^-	1300 fb ⁻¹	1.1%	115 fb ⁻¹	5.0%	1600 fb ⁻¹	2.6%
e^+e^-	450 fb ⁻¹	2.2%	45 fb ⁻¹	9.8%	1600 fb ⁻¹	3.1%

$c_H \frac{v^2}{\Lambda^2} < 0.0044$
 $\Lambda > 2.6 \text{ TeV} \quad (c_H = 1)$
 $r_H < 0.076 \text{ am}$

My naive ILC combo: $\delta\sigma_{Zh}/\sigma_{Zh} = 0.88\%$

This requires the absolute value, not ratio. \rightarrow recoil mass technique essential \rightarrow e⁺e⁻ colliders.

The current state of the art

Kinematic fit, optimized event selection, leading to 10% relative improvement.

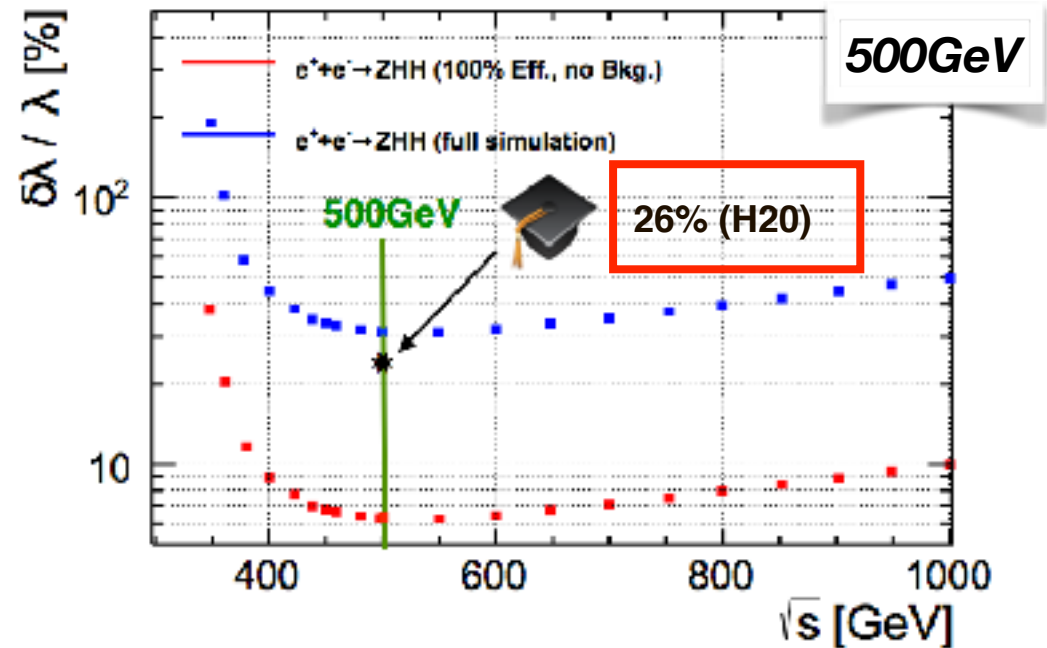
\rightarrow for HH \rightarrow bbbb

$$\frac{\Delta\sigma(\text{ZHH})}{\sigma(\text{ZHH})} = 21.1\% \rightarrow 5.9\sigma \text{ discovery}$$

\rightarrow combined with HH \rightarrow bbWW*

$$\frac{\Delta\sigma(\text{ZHH})}{\sigma(\text{ZHH})} = 16.8\% \rightarrow 8.0\sigma \text{ discovery}$$

\rightarrow results in 26.6% precision on λ_{SM}



Note: $\delta\lambda/\lambda = 13\%$ (H20) if $\lambda = 2 \lambda_{SM}$