

DETECTOR TECHNOLOGIES

Lecture 2: gaseous detectors

Principle of operation

Proportional counters and beyond

Gaseous detectors : Principle of operation

E. Rutherford and H. Geiger (1908) "An electrical method of counting the number of α particles from radioactive substances," Proceedings of the Royal Society (London)

1. A charged particle is passing through a gaseous medium : loss of energy

$$-\frac{dE}{dx} = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 T_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$

$$T_{\max} = \frac{2 m \beta^2}{1 - \beta^2} \quad \text{Ex : proton } 1 \text{ GeV}/c^2$$

$$T_{\max} = 1.2 \text{ MeV}$$

$$K = 4 \pi N_A r_e^2 m_e = 0.3071$$

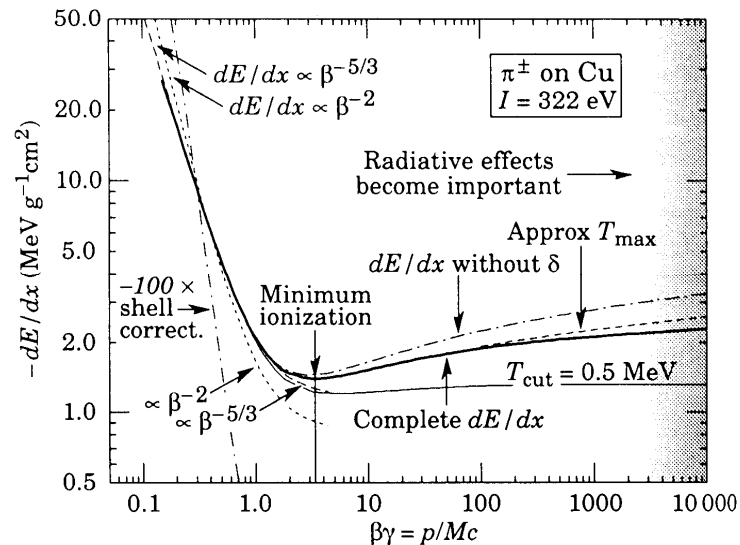
A, Z : atomic mass and number relative to the medium

N_A : Avogadro's number

T_{\max} : maximum possible energy transferred to an electron in the medium

z : charge of the incoming particle

β, γ : relatives to the particle



Gaseous detectors : Principle of operation

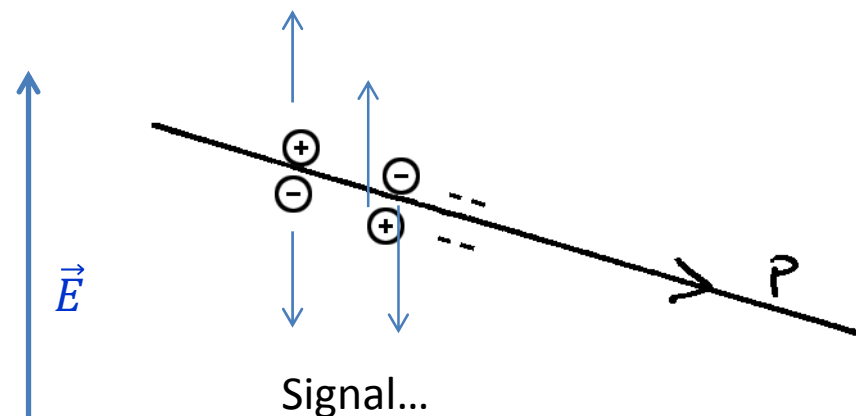
2. E_i : Ionization Energy corresponds to the energy required to remove a single electron from a single atom (or molecule).

Approximation : $E_i \approx 16 Z^{0.9}$

3. If $T_{max} > E_i$ One or more pairs electron – ion is created

Gas	E_i (eV)	$\frac{dE}{dx}$ (MeV)	N pairs /cm
H2	15.4	4.03	5.2
O2	15.2	1.69	22
Ne	21.6	1.68	12
Ar	15.8	1.47	29.4
Xe	12.1	1.23	44
CO2	13.7	1.62	34
CH4	13.1	2.21	16
DME	10.0	1.85	55

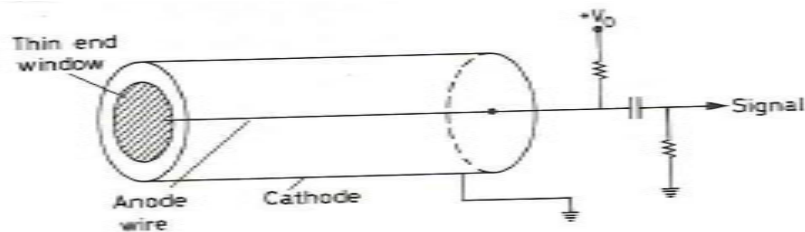
4. If exists an electrical field :
Electrons (and ions) are drifting ...



Gaseous detectors : Principle of operation

First example : Geiger-Muller counter

Idea from Hans Geiger in 1913 – Developpement with Walther Muller in 1928



$$\text{Radial Electrical field : } E(r) = \frac{V_0}{r \ln \frac{r_a}{r}}$$

r_a = anode radius
 r = counter radius

$$\text{Signal collected } V = \frac{N_e}{C}$$

Example :
 $r = 1 \text{ cm}$
Gas : Argon
particle = MIP \rightarrow 120 pairs
 $C = 10 \text{ pF}$
Signal : $2 \mu\text{V}$

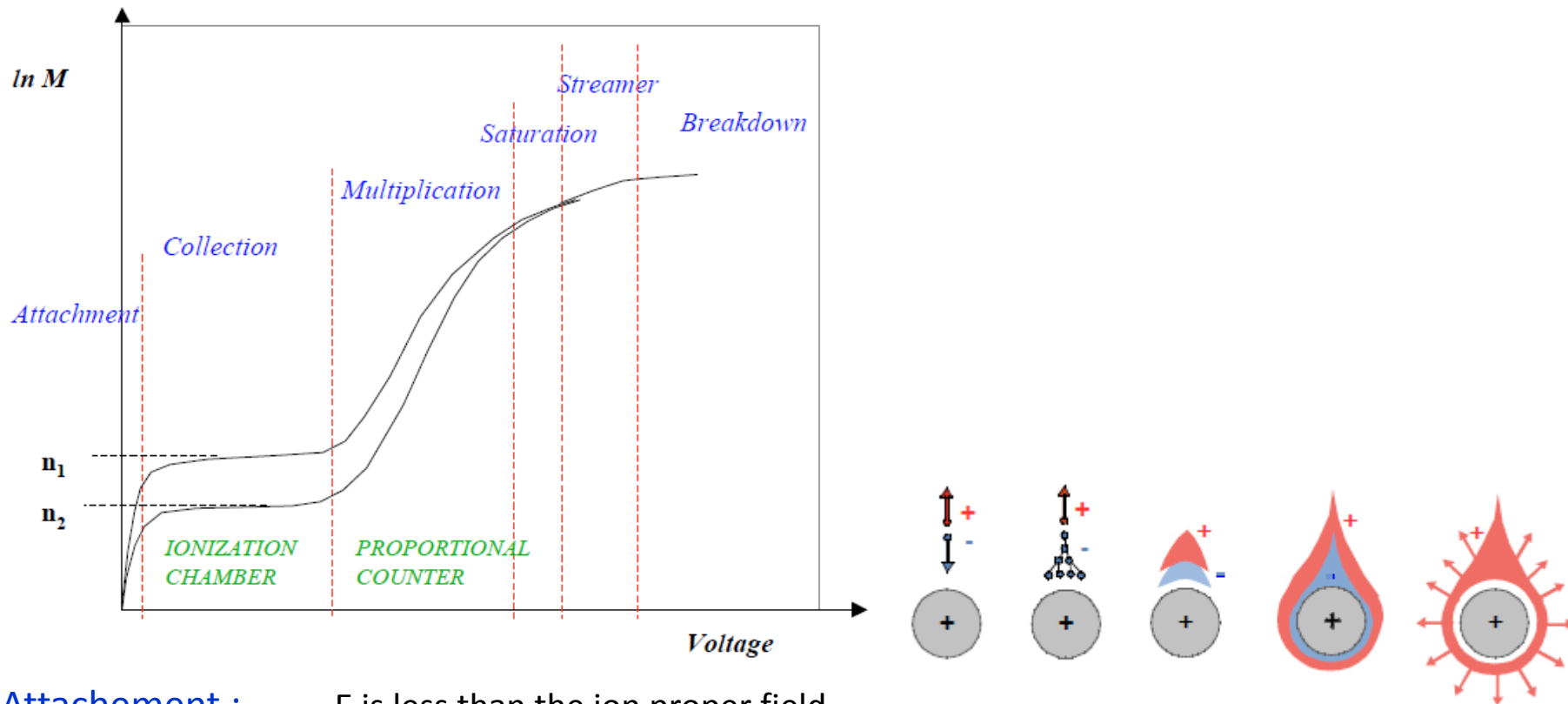
$$\text{Where } C = \frac{2\pi\epsilon}{\ln \frac{r_a}{r}}$$

Extremely weak signal... (One electron = 10^{-9} Coulomb...)

But : what can append to the electrons (and ions) during the drift before collection ?

It depends on the Electrical Field (applied voltage)

Gaseous detectors : Principle of operation

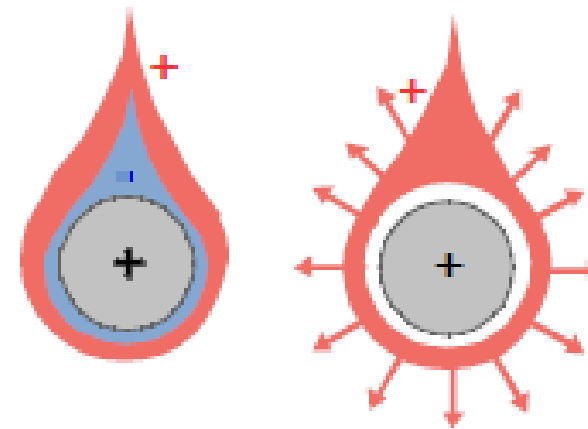
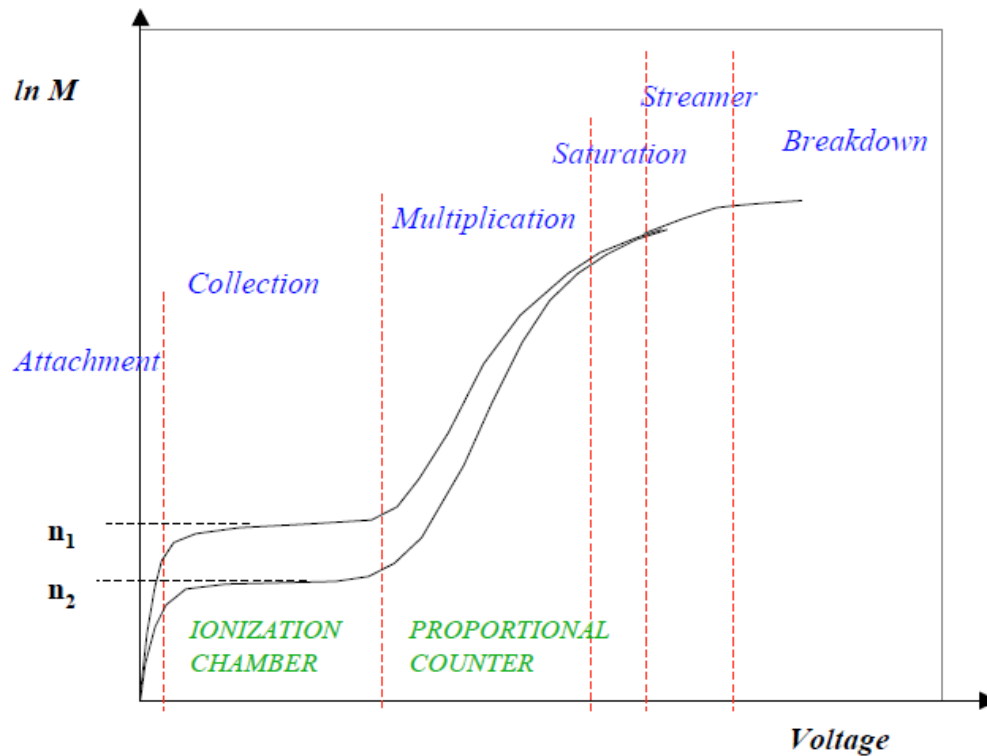


Attachement : E is less than the ion proper field
e⁻ - ion recombination - almost no signal

Collection - Ionization chamber : All e⁻ are drifting towards the anode.
Weak signal (typically 1 e⁻ for 30 eV)

Multiplication - proportionnal regime : E big enough for accelerating e⁻ above E_i
Production of secondary e⁻ ... Avalanche

Gaseous detectors : Principle of operation



Saturation and Streamer mode - Geiger-Muller regime :

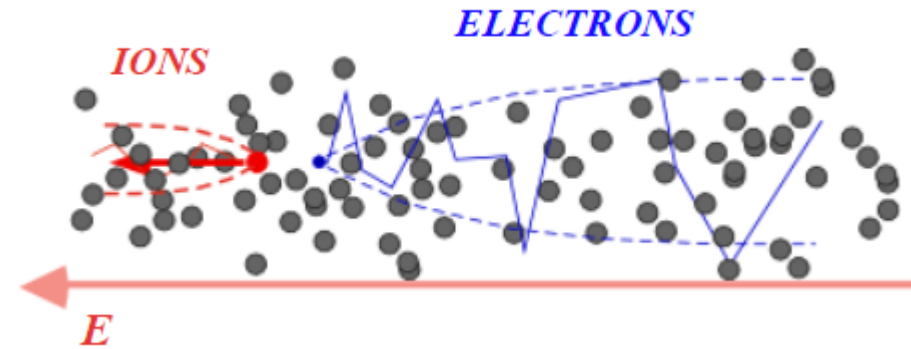
Electronic avalanche amplified by desexcitation of ions
 Saturation of the signal.
 Loss of proportionality

Breakdown : Continuous discharges between anode and cathode... Ultimate destruction...

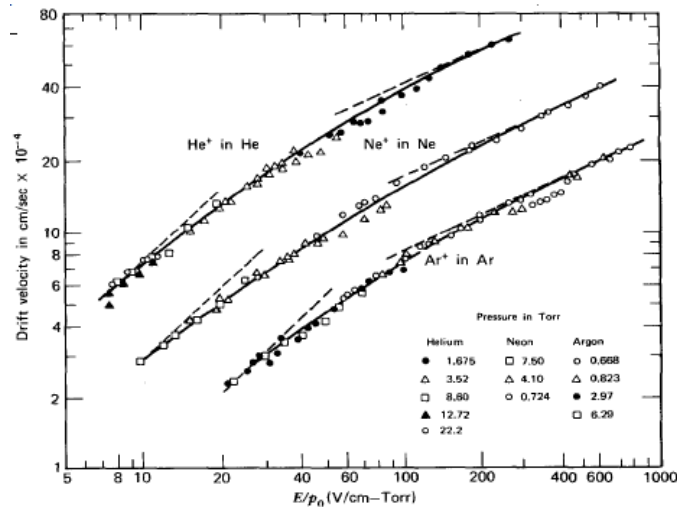
Gaseous detectors : Principle of operation

Transport of electrons and ions in the gas :

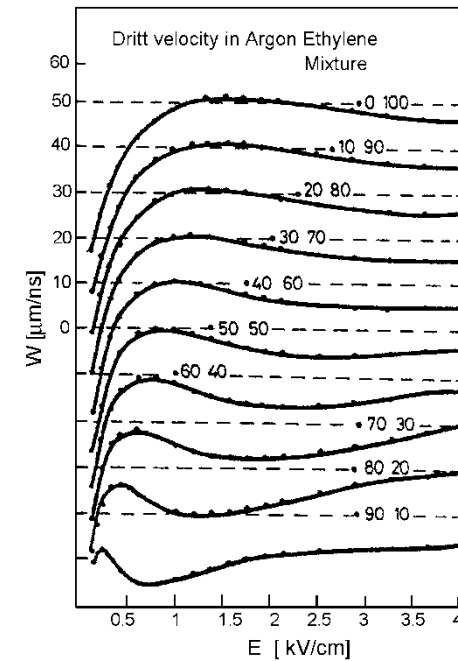
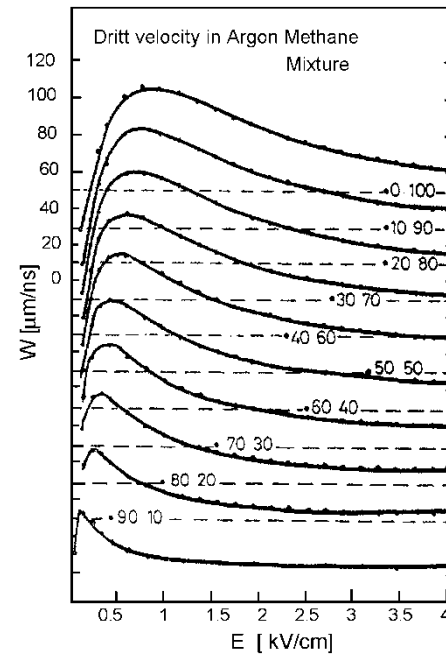
With an electric field, electrons and ions are accelerated along the field lines. Their movement is interrupted by collisions (mean free path...) which limit the maximum average velocity.



This drift velocity is low compare to the thermal velocity.



Ions



Electrons

The ion collection time is determinant !

Gaseous detectors : Principle of operation

Choice of gas : Of course, ionization exists in any possible gas.

- Maximum gain (primary electrons)
- Applied voltage as low as possible
- Avalanche with a good proportionality
- Drift velocity as high as possible

A good compromise : Noble gas (Ar, Xe, Ne...) :

For example : Argon : 30 primary electrons

Possible gain $10^3 - 10^4$

e – drift velocity : 100 $\mu\text{m}/\text{nsec.}$ at $E = 1\text{kV}/\text{cm}$

Limitation : Noble gas have an high excitation energy (typically 10-12 eV). Excited atoms formed in the avalanche desexcite giving photons which can ionize, causing further avalanche... .. Possible discharges.

Solution : the quencher

Gas	E_i (eV)	$\frac{dE}{dx}$ (MeV)	N pairs /cm
H2	15.4	4.03	5.2
O2	15.2	1.69	22
Ne	21.6	1.68	12
Ar	15.8	1.47	29.4
Xe	12.1	1.23	44
CO2	13.7	1.62	34
CH4	13.1	2.21	16
DME	10.0	1.85	55

Gaseous detectors : Principle of operation

Quencher : one has to add a polyatomic gas in order to absorb the photons created either by multiple collisions or molecule dissociation
Usually CH₄, CO₂, CF₃, C₂H₄

With a mixture of Noble gas – Quencher, one can achieve gains up to 10⁶ - 10⁷

Magic Gas : 70% Ar, isobutane 29.6%, Fréon 0.4% .

Problem : after dissociation, the organic molecules will polymerize on the anode.

- Loss of efficiency
- Need gas circulation

One has to add another agent....
(alcohol...)

One of the BEST possible choice :
DME : Dymethylether CH₃OCH₃
No polymerization
Good gain (10⁶)
But it is a solvent !

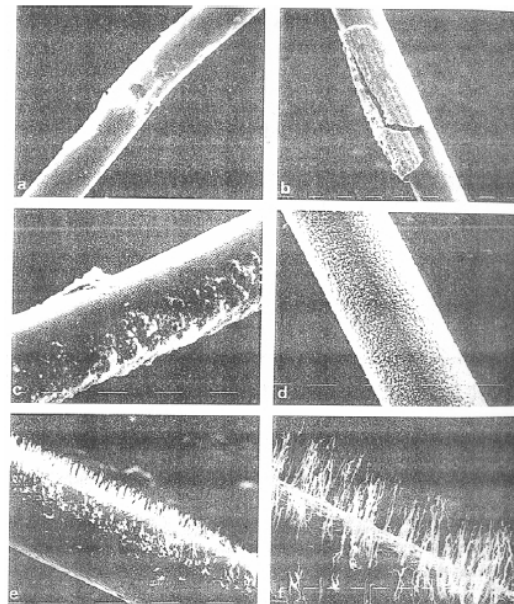


Fig. 4.31. Deposits on anode wires: (a) – Ar + C₂H₆; (b) – Ar + C₂H₆ + methylal; (c) – Ar + CO₂; (d) – perspex chamber, (e, f) – chambers with G10 fiber-glass and a cold trap (Adam 1983)

Gaseous detectors : Principle of operation

Basic requirement for a gas detector : determination of particle trajectories

Recipe for a Gas detector

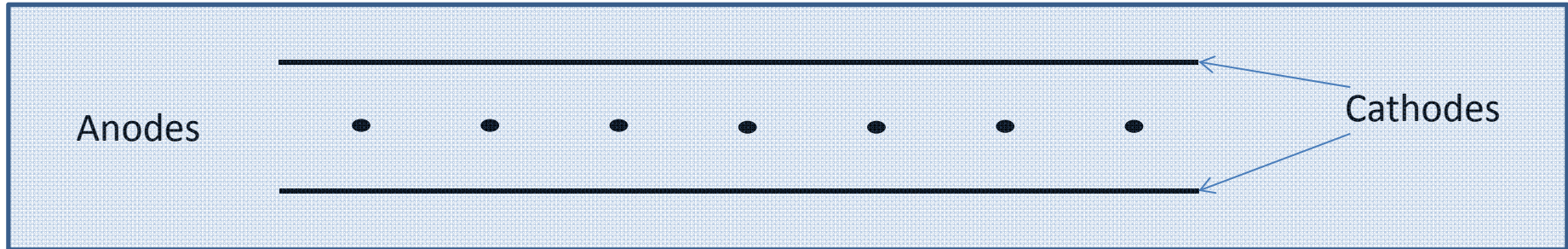
- **Thin** (minimization of dE/dx , does not perturb the particle trajectory)
- **Maximum gain** (choice of gas)
- **Stability** (High Voltage, choice of quencher)
- **Choice of material** (to avoid polymerization)
- **Precision**

Evolution from **GM counter** to **Multiwire Proportional Chamber (MWPC)**

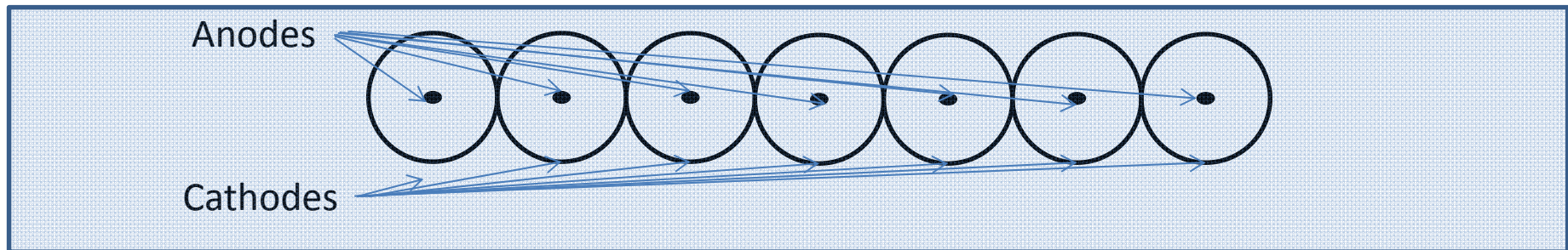
G. Charpak and all., 1968

Gaseous detectors : MWPC

An array of closely spaced anode wires in the same volume

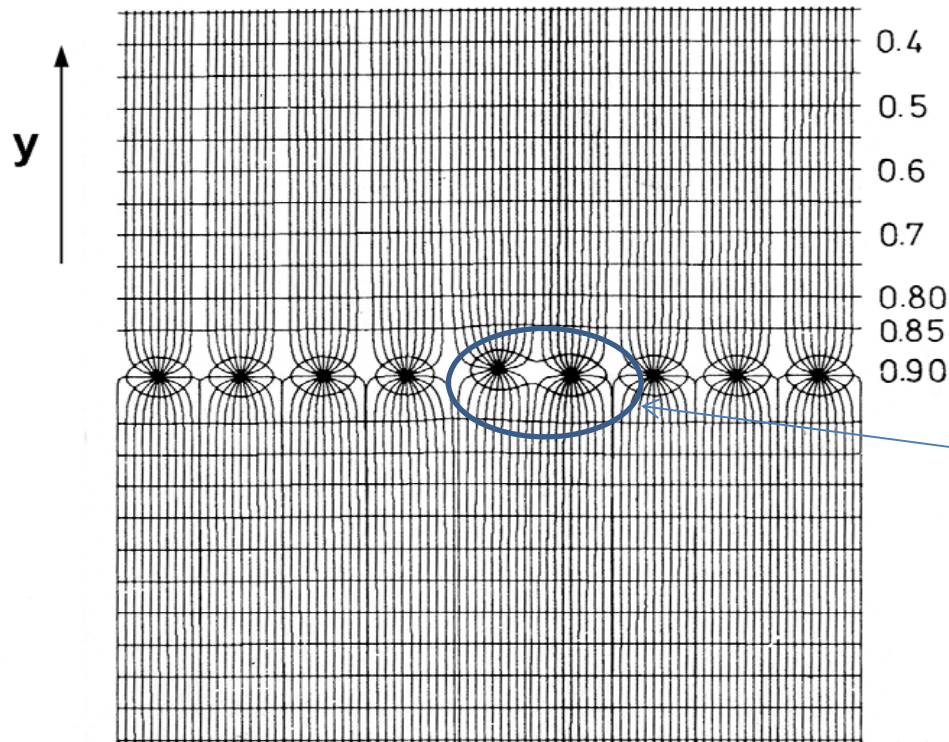
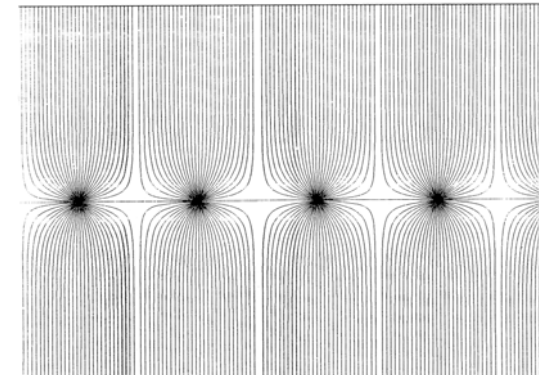
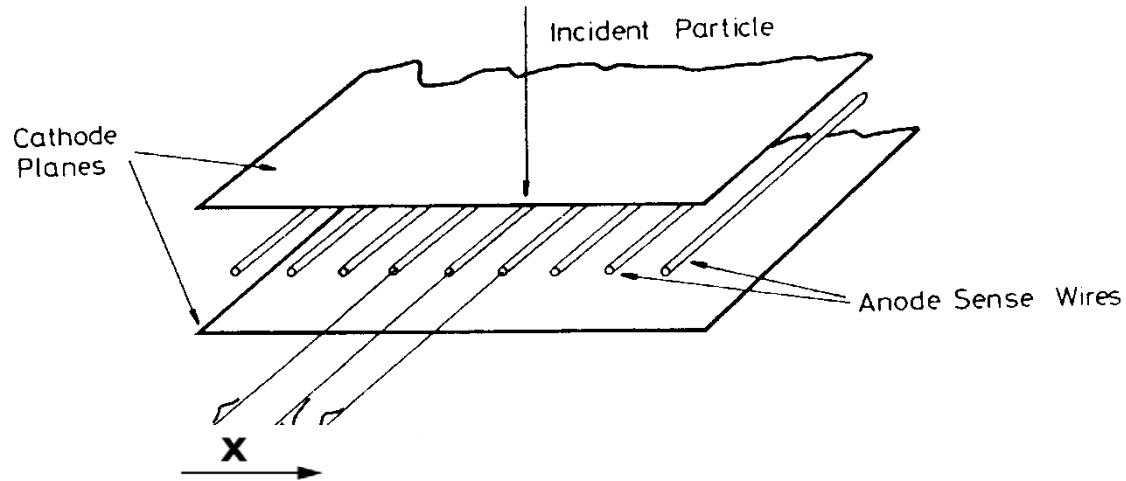


Is equivalent to



An array of proportional counters tubes

Gaseous detectors : MWPC



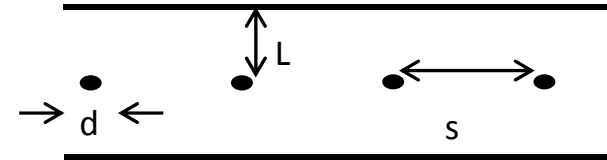
Classical figure of electrical field disturbed by a misplaced anode wire

Gaseous detectors : MWPC

Signal : as seen for the GM proportional counter : $\text{Signal } V = \frac{N_e}{C}$ $C = \frac{2\pi\epsilon}{\ln \frac{r_a}{r}}$

For an MWPC chamber
($L \gg s \gg d$)

$$C = \frac{2\pi\epsilon}{\frac{\pi L}{s} - \ln \frac{\pi d}{s}}$$



Spatial résolution :

The charges due to the particle passing in the gas are distributed over more than one anode. The spatial resolution of a MWPC is the variance of this distribution .

$$\sigma = \frac{a}{\sqrt{12}} \quad \text{Typically } \approx 200 \mu\text{m}$$

Signal formation time : depends on the drift time for electrons (typically 50 nsec)

Dead time : depends on the drift time for the ions (typically 200 nsec)

Gaseous detectors : MWPC limitations

Limitation 1 : Typically : anode spacing of the order of 1-1.5 mm

(Résolution $\approx 200 - 500 \mu\text{m}$)

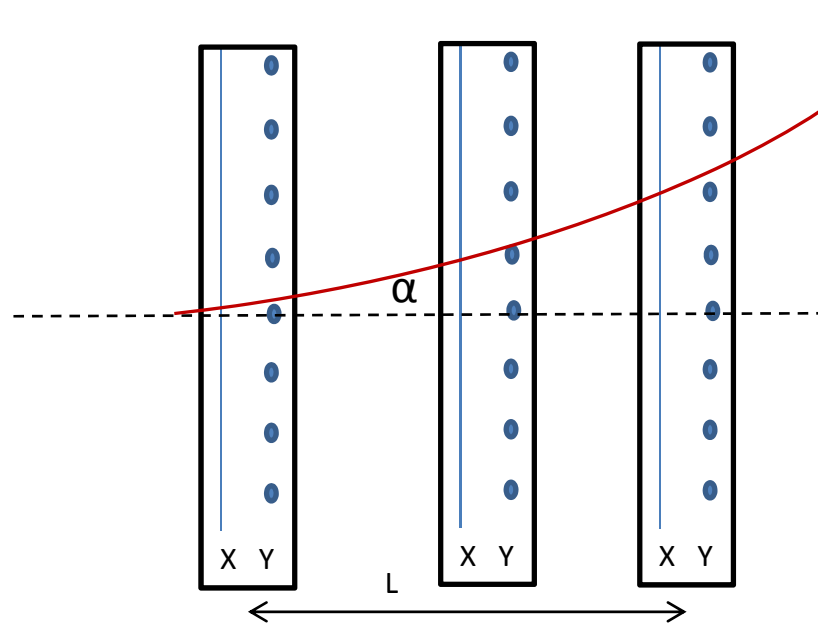
In order to improve the spatial resolution : closer anodes ?

Does not work. Instabilities due to electrostatic forces anode-anode.

Limitation 2 : MWPC can measure only one coordinate.

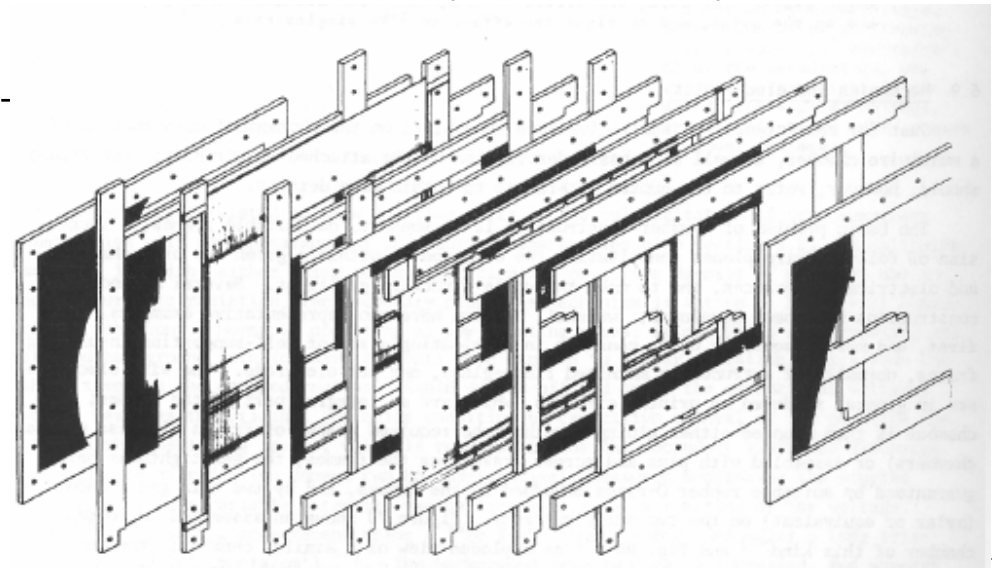
A second MWPC ?

X-Y coordinate with a second anode row



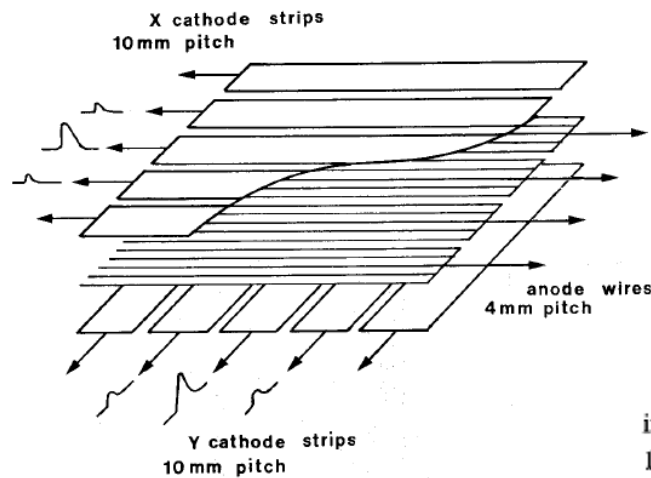
$$\text{Precision : } \Delta \alpha \approx \frac{1}{L} \sigma$$

One can reduce the wire spacing...
Or increase L (dimensions...)



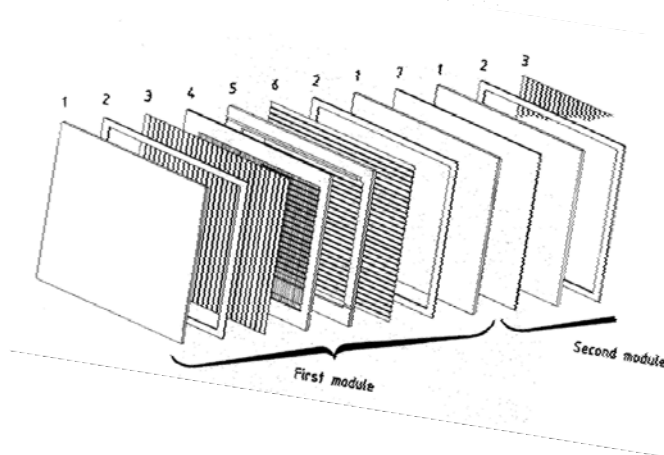
Gaseous detectors : MWPC evolution

Cathode read-out chamber : One anode plan
 Segmented cathode plan
 Analog read-out



The features of one chamber module of the final stack were the following:
 half gap: 5 ± 0.02 mm
 anode pitch: 4 mm
 anode wire diameter: $20 \mu\text{m}$
 cathode strip pitch: 10 mm
 cathode strip width: 9 mm
 sensitive area: $1100 \times 970 \text{ mm}^2$
 gas composition: Ar 80% + CO₂ 20% (vol).
 The total stack thickness was 1 radiation length.

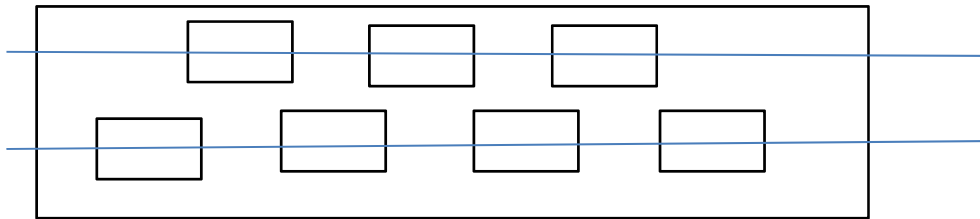
Spatial localization was investigated by comparing the analog chamber results with the information given by the set of digital MWPCs. A value of $\sigma = 2.4$ mm (98% of events inside 10 mm) was obtained for 4 GeV shower electrons after 4 radiation lengths (fig. 8). The resolution deteriorated quickly for lower energies ($\sigma = 5.9$ mm, 91% of events inside 10 mm at 2 GeV) and for wider strips ($\sigma = 5.3$ mm, 86% of events inside 10 mm for 4 GeV electrons, when going from 8 mm to 16 mm strip width).



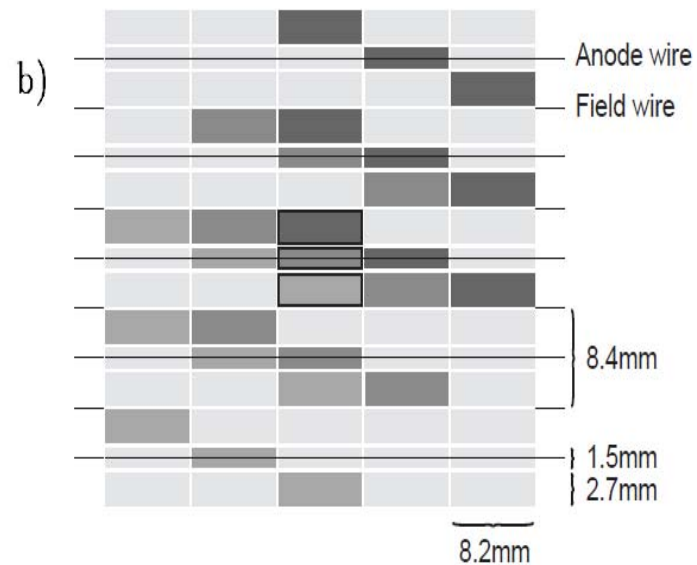
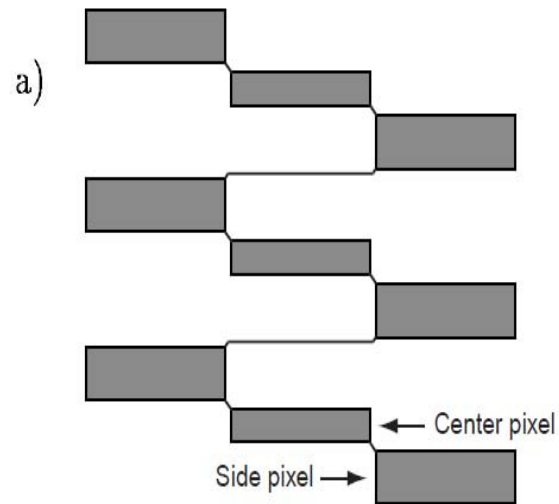
Experiment R704 (CERN) 1981 - 1985

Gaseous detectors : Pad chambers

Direct 2-D detector : Pad chambers
cathode segmented in pads



Needs a lot of electronics

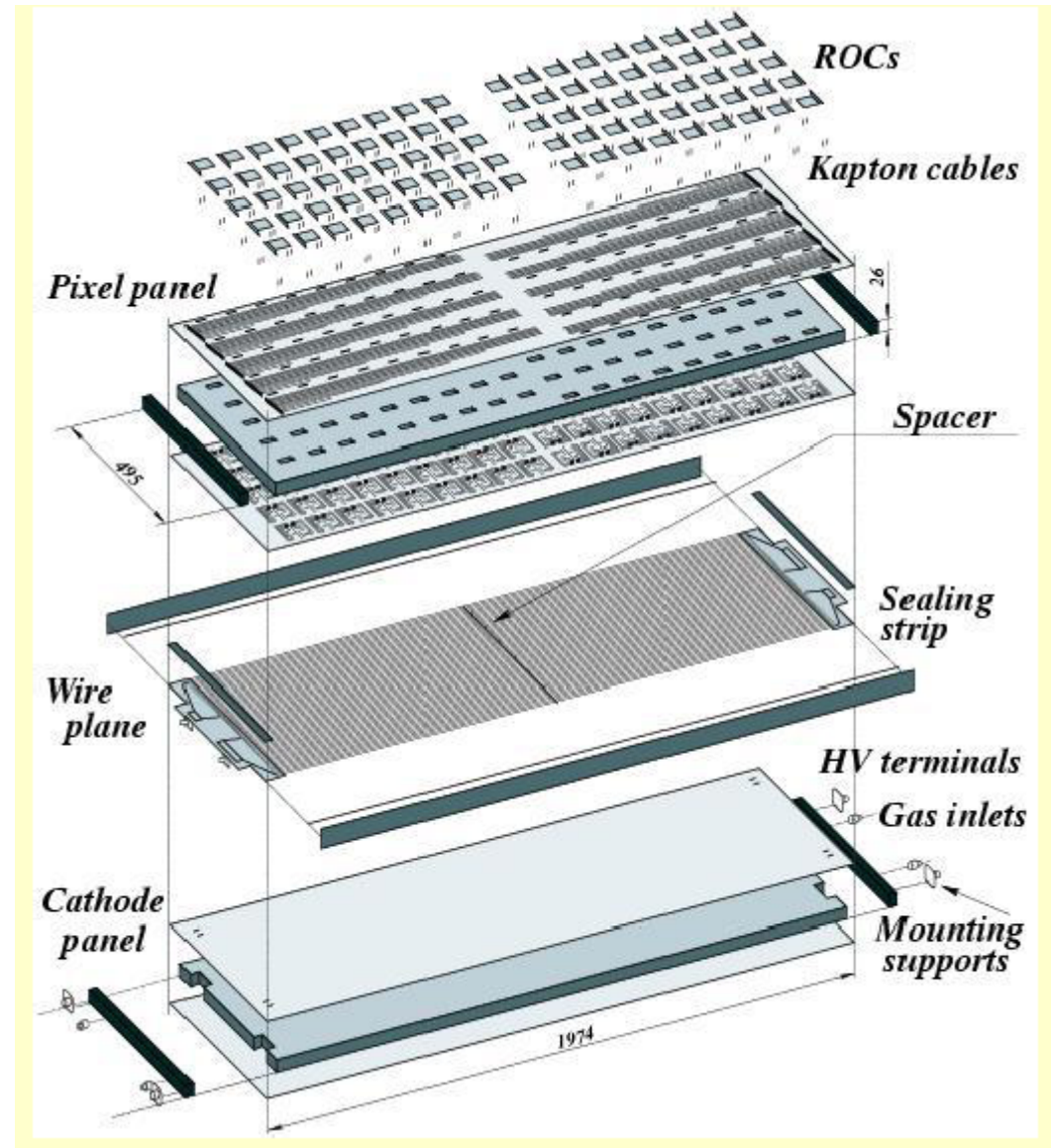
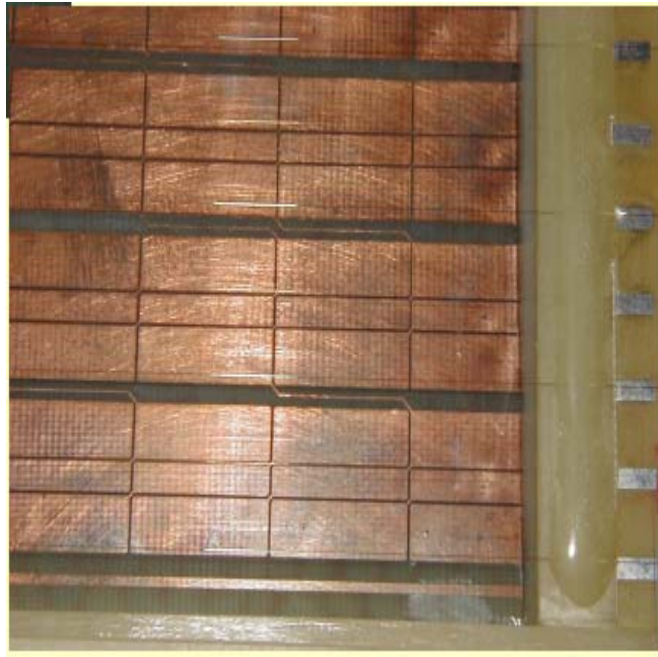


Pads regrouped in cells

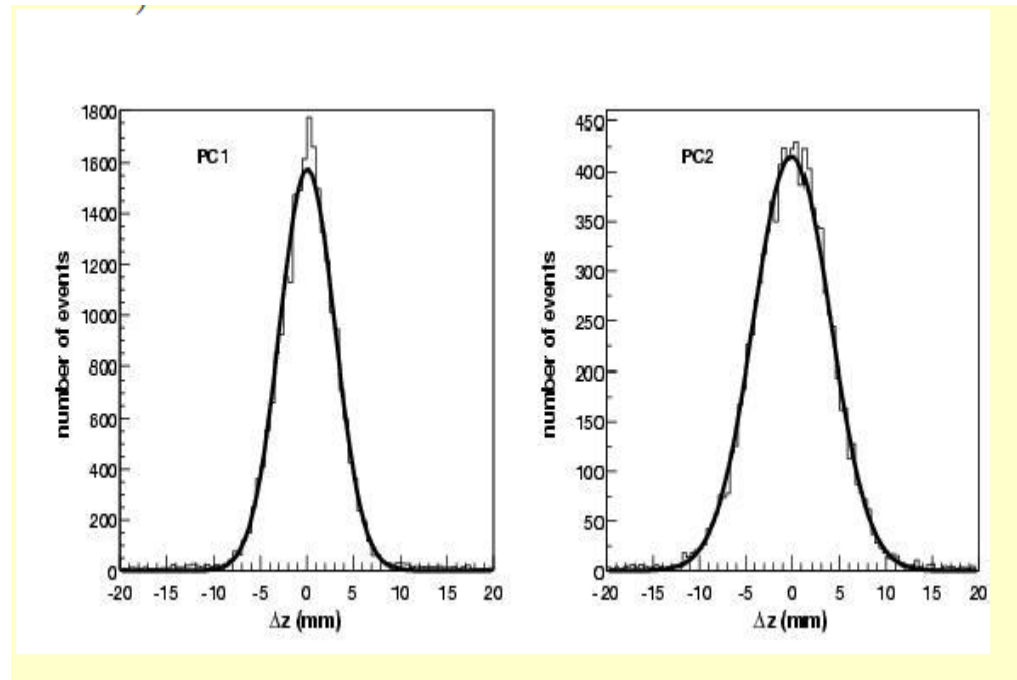
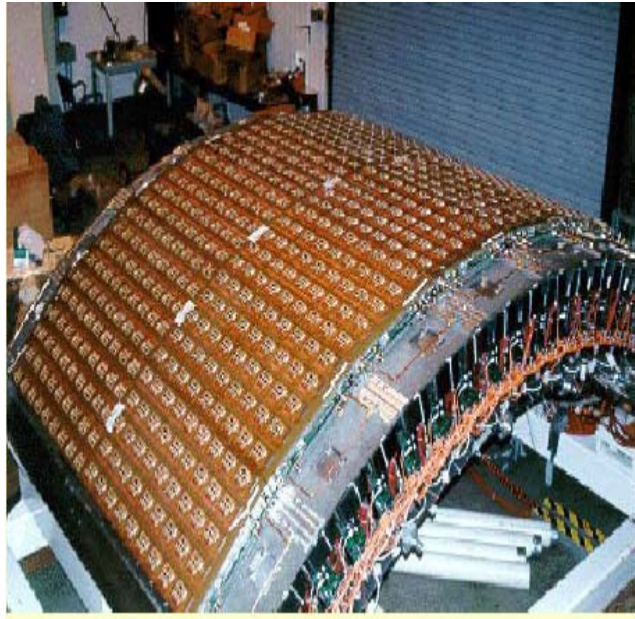
PHENIX

Gaseous detectors : Pad chambers

PHENIX at RHIC



Gaseous detectors : Pad chambers



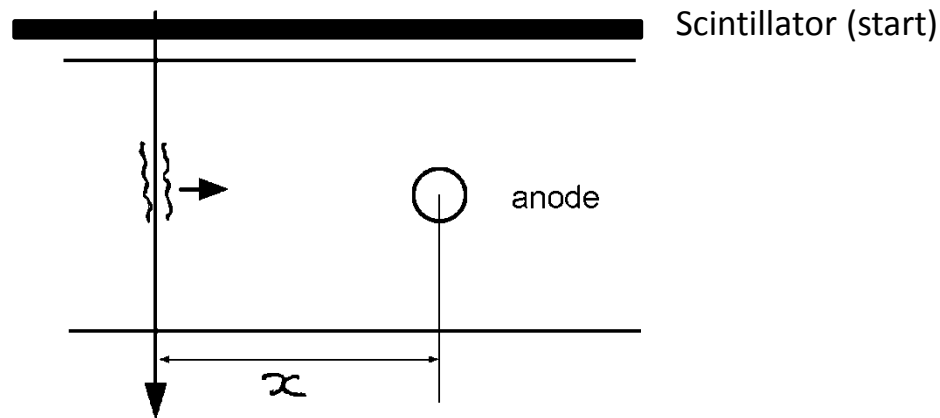
chamber	Wire dist (mm)	Z-resol. (mm)	Perp res (mm)	Rad. Thickn.
PC1	8.4	1.7	2.5	1.2%
PC2	13.6	3.1	3.9	2.4%
PC3	16.0	3.6	4.6	2.4%

Gaseous detectors : Drift Chambers

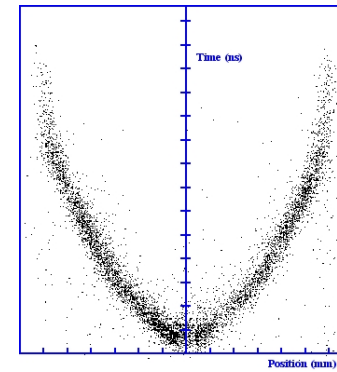
A drift chamber is a particle tracking detector that measure the drift time of ionization electrons in a gas to calculate the spatial position of ionizing (charged) particle. Similar to MWPC, but with a better accuracy.

Measure of the position of the particle by mesuring the drift time of the electrons

Need : Precise knoweledge of drift velocities
Precise timing (trigger)



$$X = Vd \Delta t$$



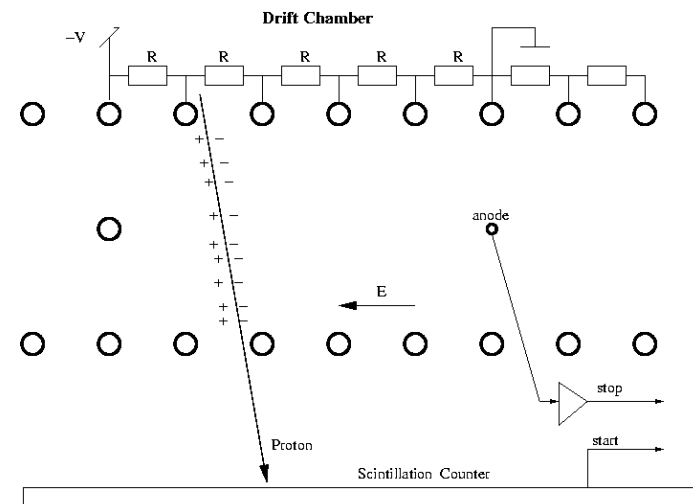
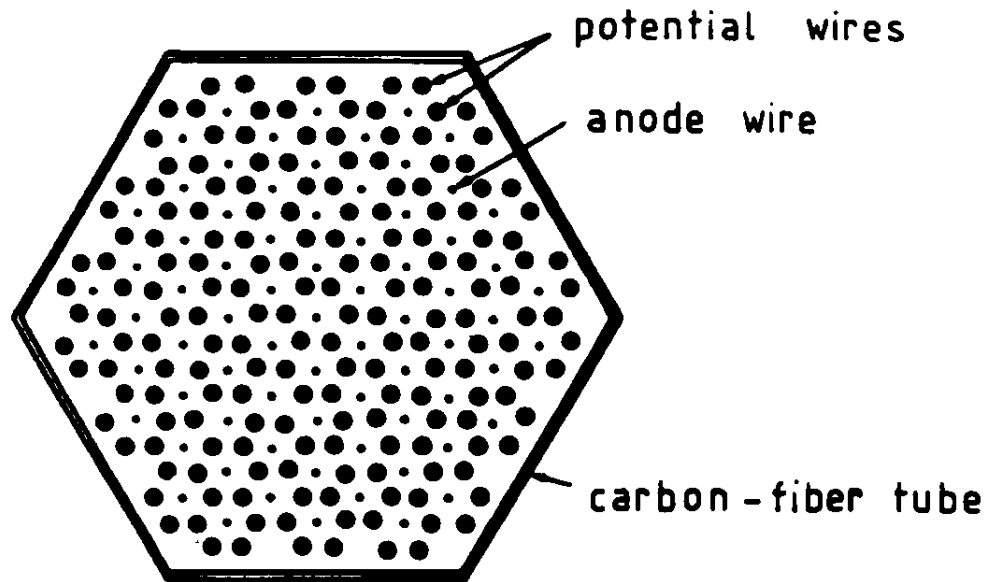
SPACE-TIME CORRELATION
(RIGHT-LEFT AMBIGUITY)

Possible resolution down to 50 μm

Gaseous detectors : Drift Chambers

Main limitation : non-uniformity of the field (non uniformity of the drift time)

Solution : defining a cell (basic unit of field)



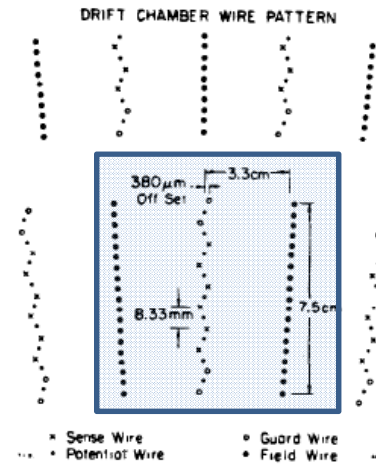
$$\text{Drift Distance} = \text{Drift Velocity} * \text{Time}$$

PHENIX (RHIC) Drift Chamber

Gaseous detectors : Drift Chambers



MarkII big Drift Chamber



1 cell = 7.5 cm x 6.6 cm

12 layers

Length : 2.3 m

Radius : 1.6 m

5732 sense wires (anodes)

31104 potential wires

Gas mixture : 89% Ar, 10% CO₂ ,1% methane

Gain : $2 \cdot 10^4$

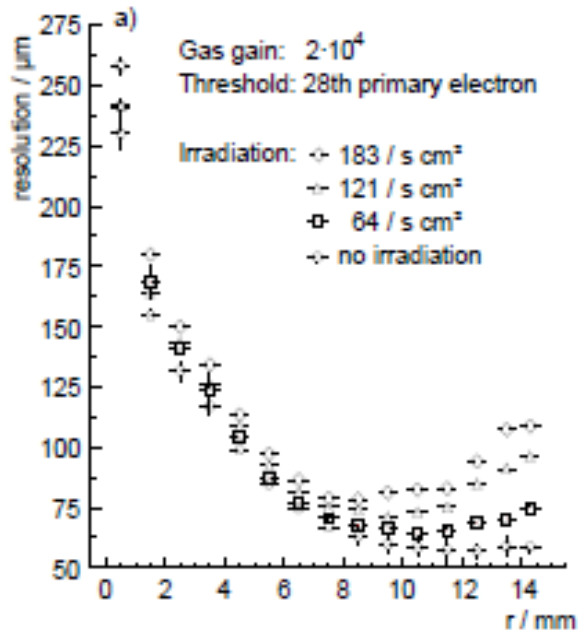
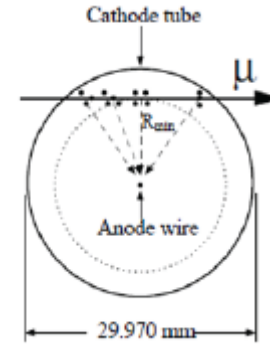
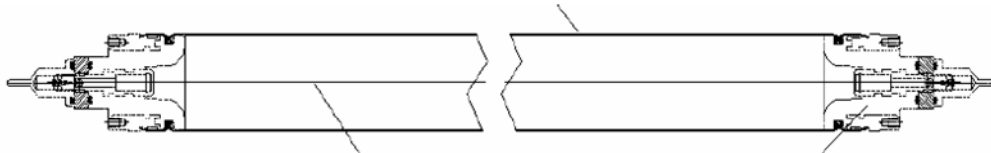
Drift field : 900 V/m

Spatial resolution : $\approx 150 \mu\text{m}$

Gaseous detectors : Drift Chambers

ATLAS DRIFTS TUBES (MUON SPECTROMETER)

370 000 tubes
Surface 5500 m²



Parameter	Design value
Tube material	Al
Outer tube diameter	29.970 mm
Tube wall thickness	0.4 mm
Wire material	gold-plated W/Re (97/3)
Wire diameter	50 μm
Gas mixture	Ar/CO ₂ /H ₂ O (93/7/≤ 1000 ppm)
Gas pressure	3 bar (absolute)
Gas gain	2×10^4
Wire potential	3080 V
Maximum drift time	~ 700 ns
Average resolution per tube	~ 80 μm

Max counting rate : 20 Hz / m

Gas leak < 10⁻⁸ Bar.l / sec.

Wire tension tolerance 17g

Wire position tolerance : 25 μm

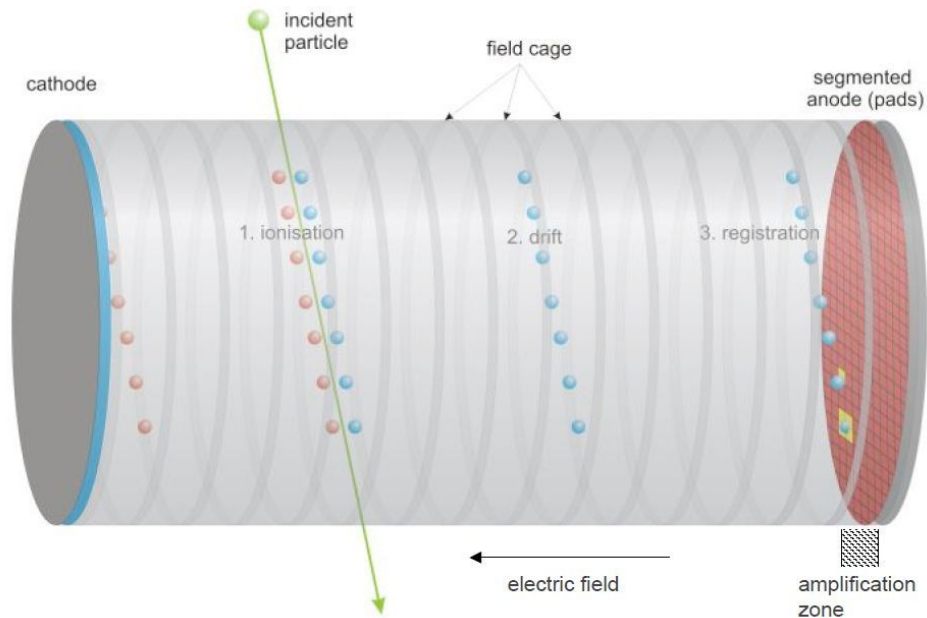
Gaseous detectors : Drift Chambers



Gaseous detectors : Time Projection Chambers

TPC : « The best of the best evolution »

Combination of a Drift Chamber and a MWPC (Pad chamber)



Position (X and Y measurements)

Drift time (Z measurement)

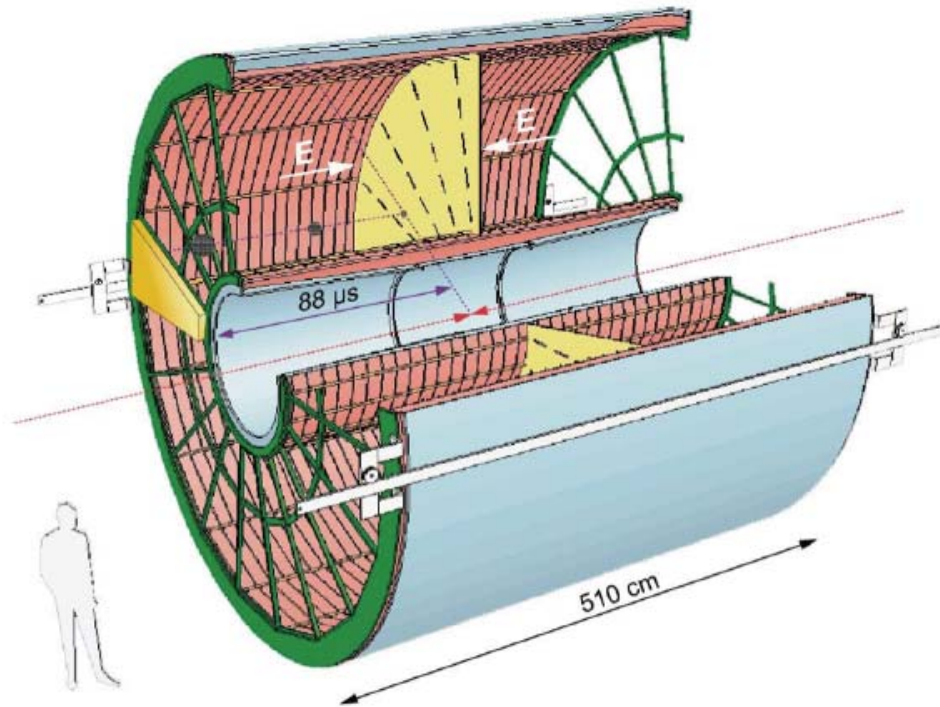
Huge (and empty) drift volume

Multi channels (read out pads)

Drift volume full of ions

→ séparation for the avalanche region

Gaseous detectors : Time Projection Chambers – The ALICE TPC at LHC



Length: 5 meter
 Radius: 2.5 meter
 Gas volume: 88 m³

Total drift time: 92 μs
 High voltage: 100 kV

End-cap detectors: 32 m²
 Readout pads: 557568

159 samples radially
 1000 samples in time

Gas: Ne/CO₂/N₂ (90-10-5)
 Low diffusion (cold gas)

Gain: > 10⁴

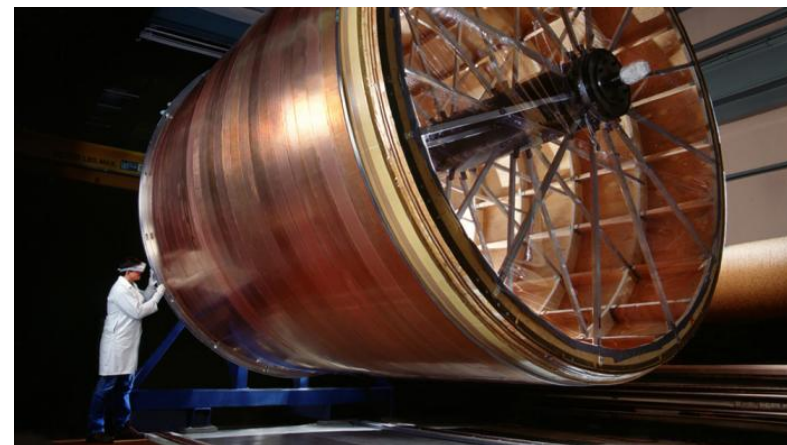
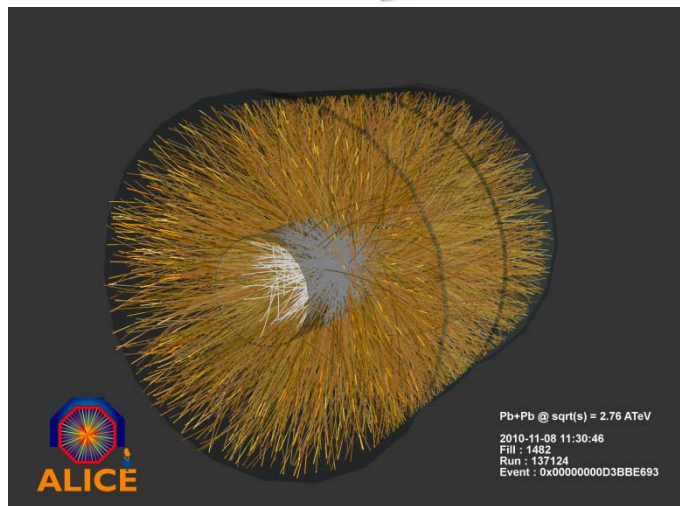
Diffusion: $\sigma_t = 250 \mu\text{m}$
 Resolution: $\sigma \approx 0.2 \text{ mm}$

$\sigma_p/p \sim 1\% p$; $\epsilon \sim 97\%$
 $\sigma_{dE/dx}/(dE/dx) \sim 6\%$

Magnetic field: 0.5 T

Pad size: 5x7.5 mm² (inner)
 6x15 mm² (outer)

Temperature control: 0.1 K



Gaseous detectors : Time Projection Chambers – The ALICE TPC at LHC

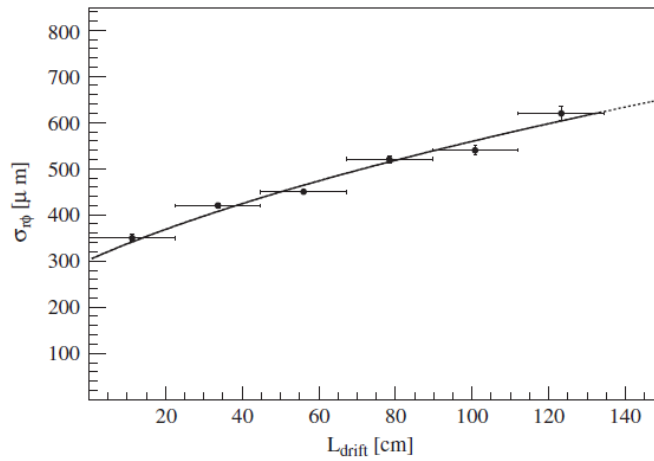


Fig. 6. Space point resolution in pad direction (momentum plane) as a function of drift length. Solid line shows the fit and dashed line the extrapolation.

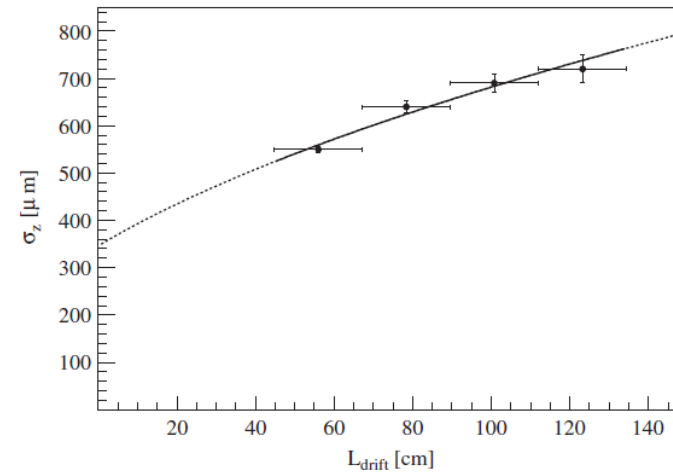
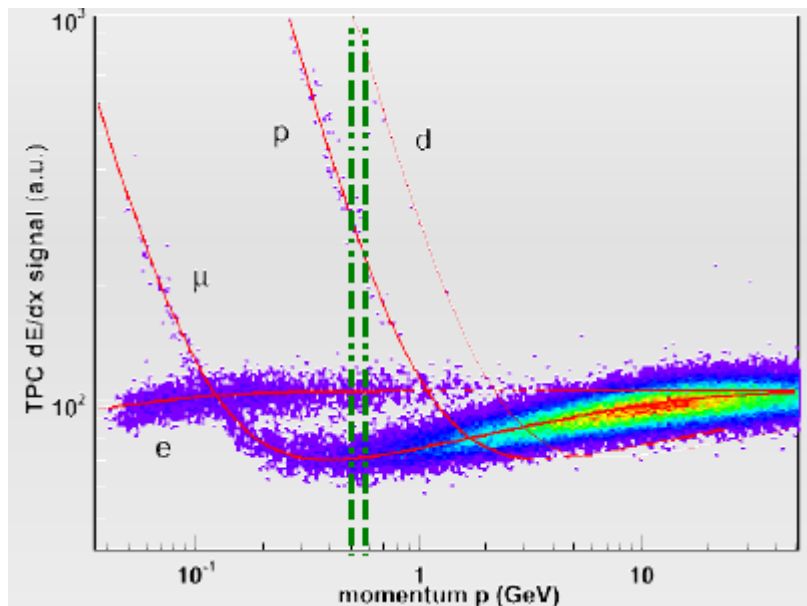


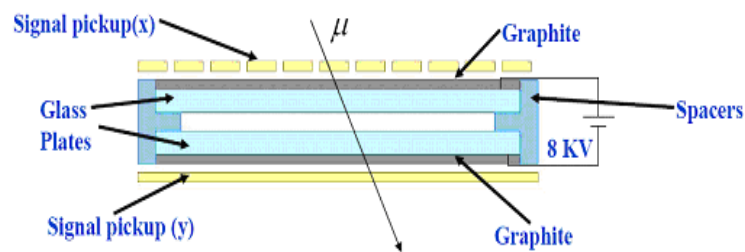
Fig. 7. Space point resolution in drift direction as a function of drift length. Solid line shows the fit and dashed line the extrapolation.



	Experiment	Monte Carlo
$\sigma_{r\phi}$ (μm)	800 ± 80	900
σ_z (μm)	900 ± 100	900

Gaseous detectors : Resistive Plate Chamber

RPC : Thin (2 mm) drift volume sandwiched between
Two highly resistive ($2.5 \cdot 10^{10} \Omega \cdot \text{cm}$) plates
Simple – inexpensive – fast (used as trigger)

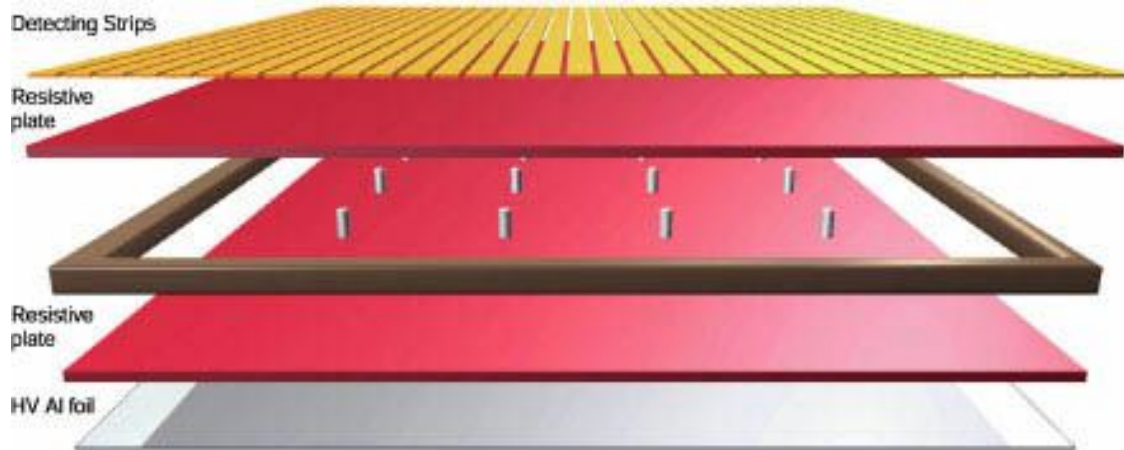


Thin drift volume and High field

Works usually in streamer mode
Large signal, but slow (100 nsec)

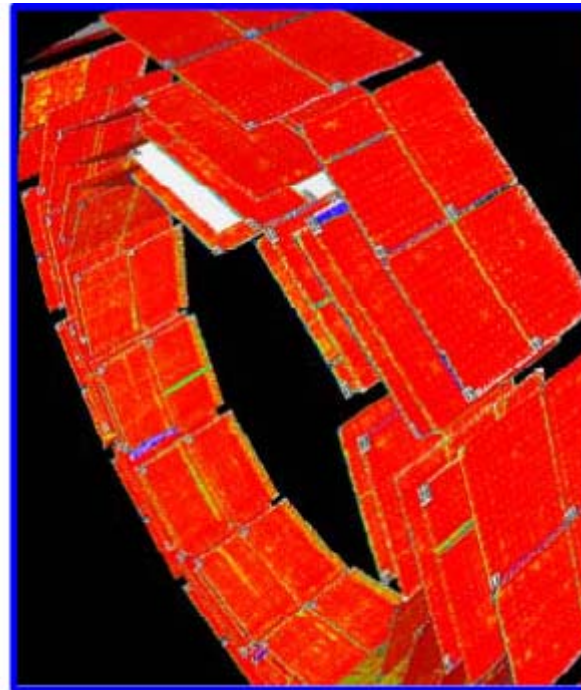
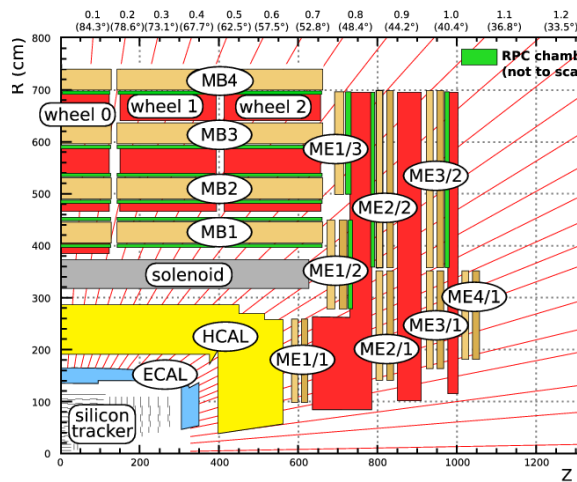
At LHC : in avalanche mode
Lower signal, but fast (1-10 nsec)

Gaseous detectors : Resistive Plate Chamber for CMS - LHC



Resistive plates : $1 \cdot 10^{12} \Omega \cdot \text{cm}$
 Gas mixture :
 $\text{C}_2\text{H}_2\text{F}_4$ (92.5 %) – C_4H_{10} (4.5%) – SF_6 (0.3%)
 HV : 8.5 – 9.7 kV

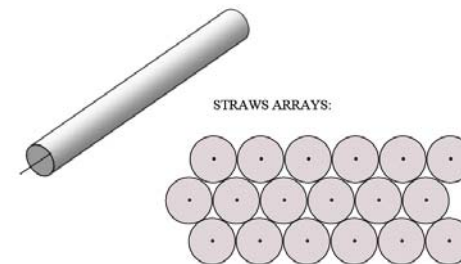
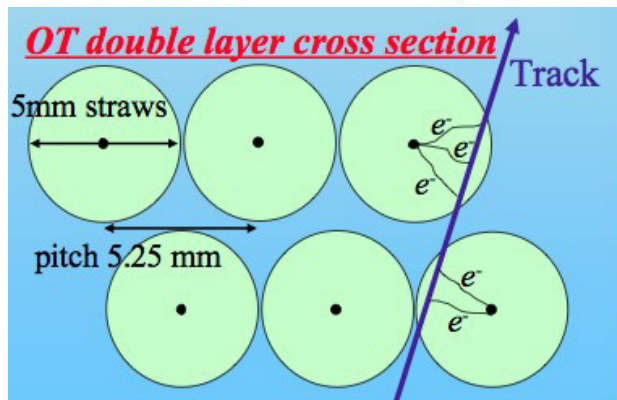
RPC efficiency > 97%
 Rate capability > 1 KHz/cm²
 Operation efficiency plateau > 400V
 Time resolution < 3ns
 Cluster size < 3
 Dead time should be few nano seconds



Gaseous detectors : Straw Chambers

SWPC : Single Wire Proportional Chamber
Single proportional counter in an array

- Cheap and simple to build
- Capable of withstanding very high fluxes



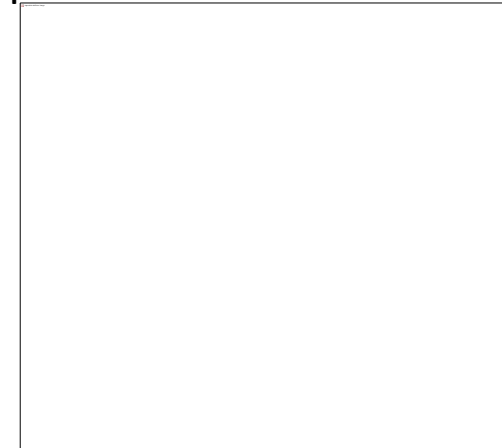
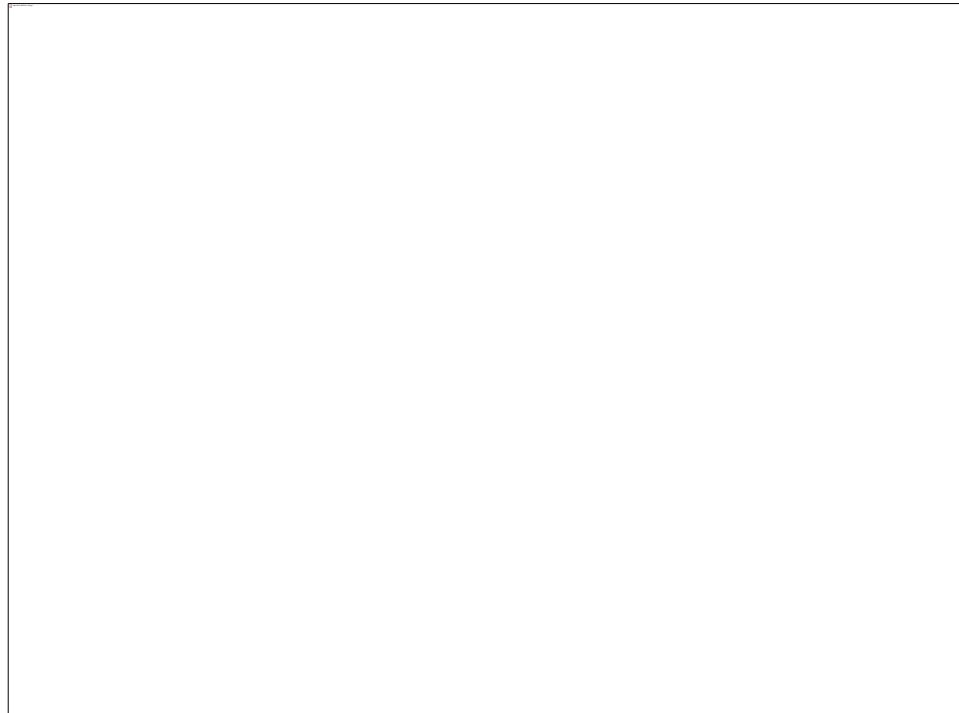
Straw Chambers : the ATLAS read-out system for the Transition Radiation Tracker



Gaseous detectors : MicroStrips Gas Detector

MSGC : following an idea of Oed (1989) :

A MWPC where the wires are replaced by strips deposited on an insulating substrate (glass)



High field : gain $\approx 10^4$

Spatial resolution : $\sigma \approx 20 \mu\text{m}$

Dead time $\approx 10^{-5}$ sec.

(Short distance for the ions)

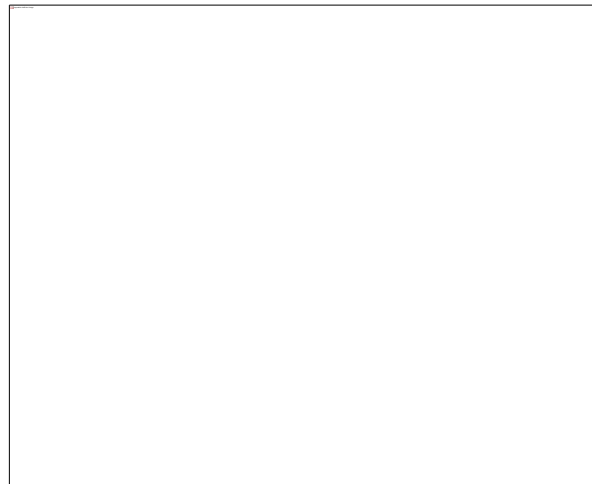
Gas : Ar – DME / Ne - DME

Gaseous detectors : MicroStrips Gas Detector



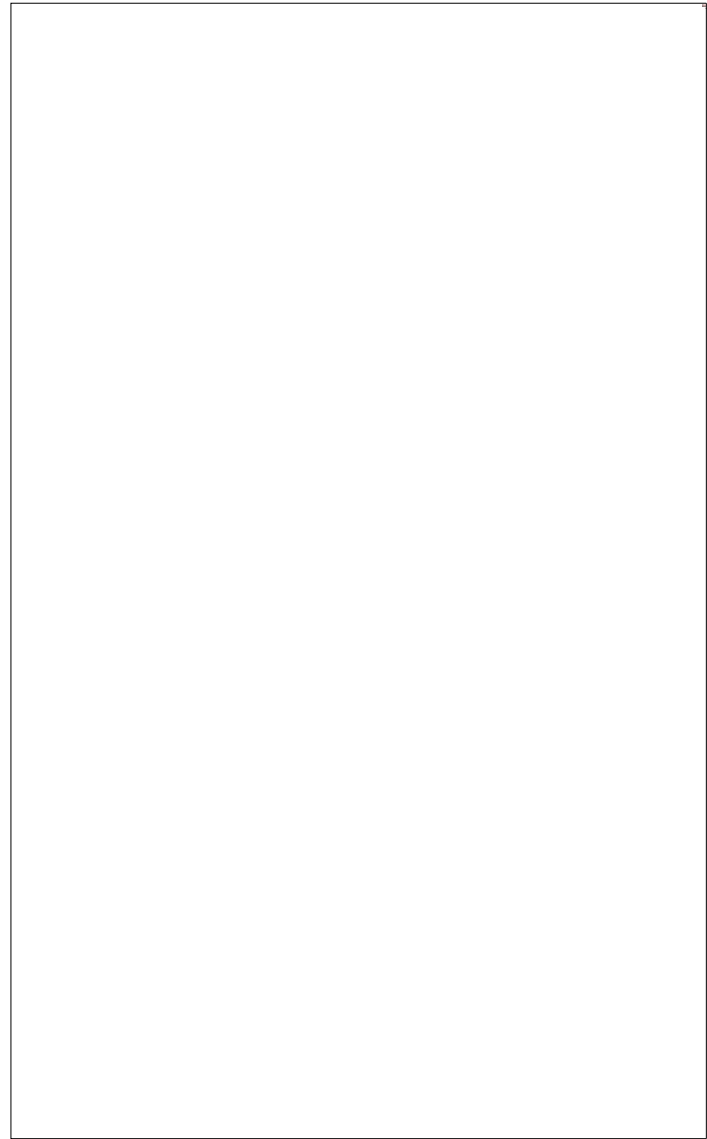
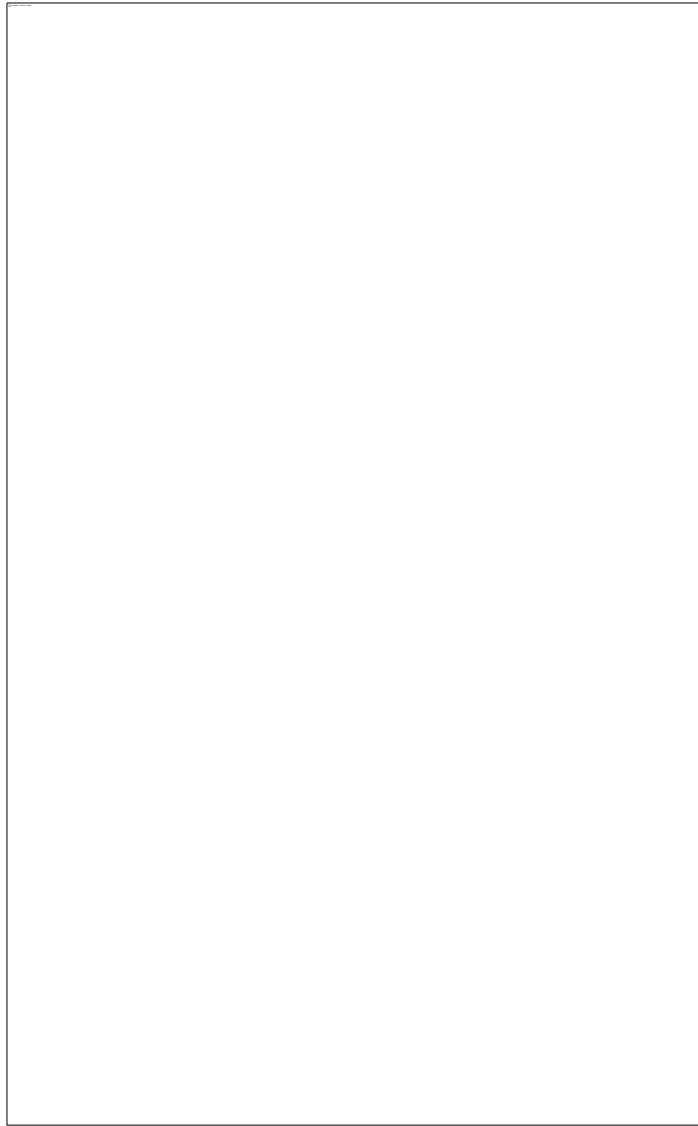
MWPC
anode spacing : 1-3 mm

MSGC
anode spacing 200 μm



Rate capability comparison
MSGC seems to be well adapted
to high fluxes of particles (LHC)

Gaseous detectors : MicroStrips Gas Detector

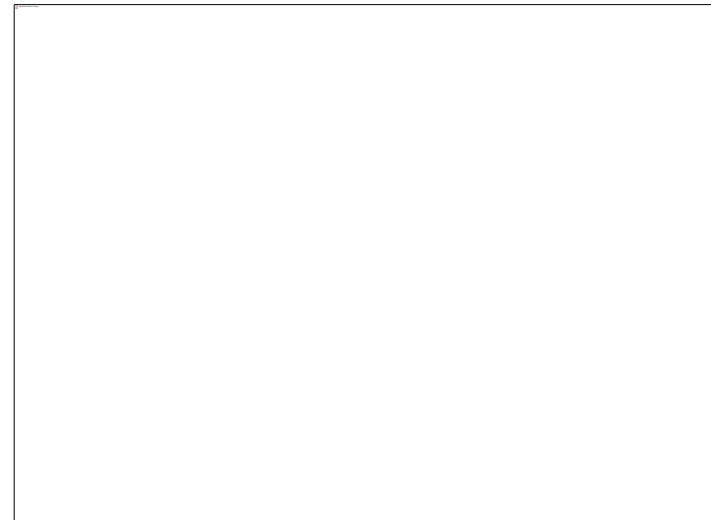
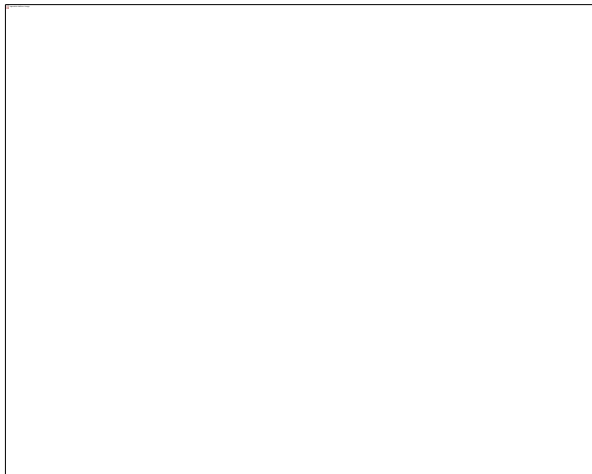


Gaseous detectors : MicroStrips Gas Detector

Surface charging : Bulk or surface resistivity of the support material is modified by irradiation (flux)
Choice of support (special glass or doping)

Ageing : Polymerization due to construction material
(DME is a solvent)
Choice of non-ovable material in DME

Discharges: Possible with higher flux or low energetic particles
Certain with dust (short between anode and cathode (50 μm))

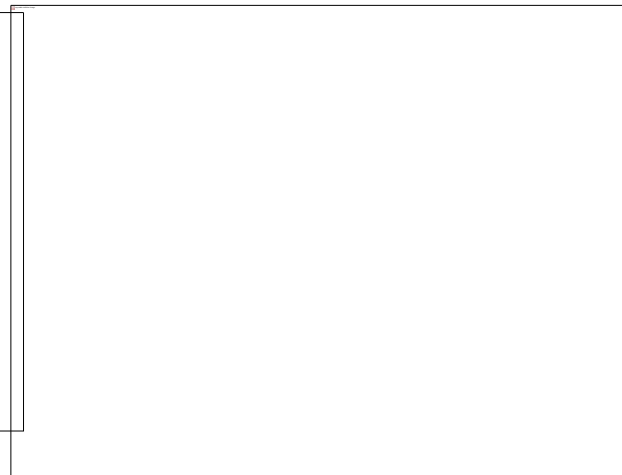
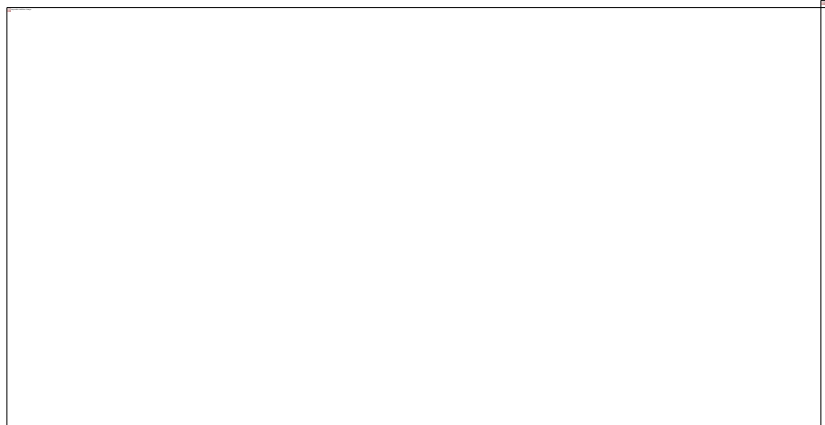
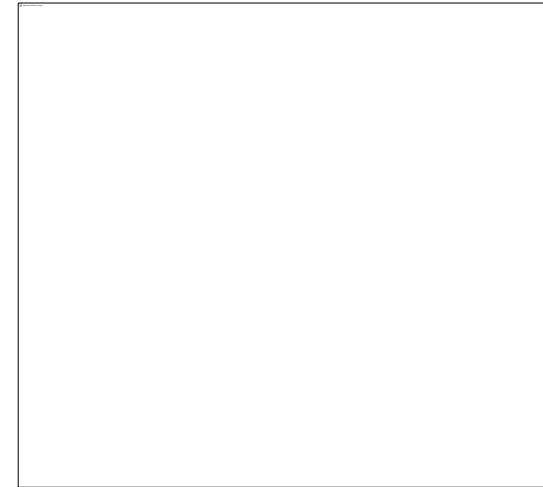
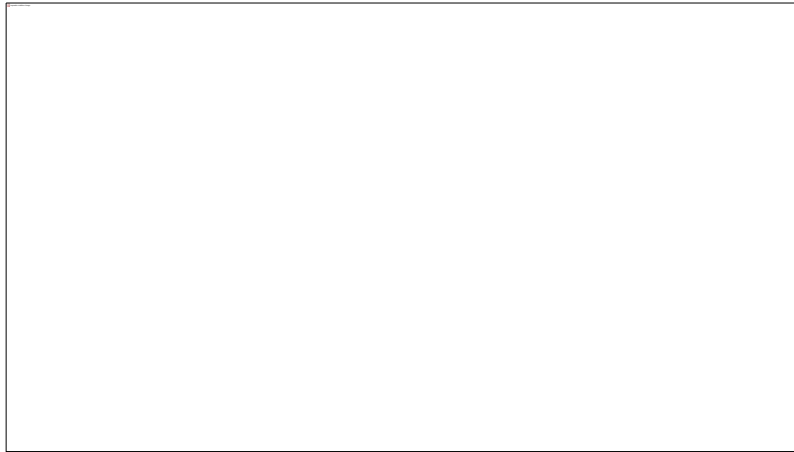


Gaseous detectors : MicroStrips Gas Detector



Gaseous detectors : Gas Electron Multiplier

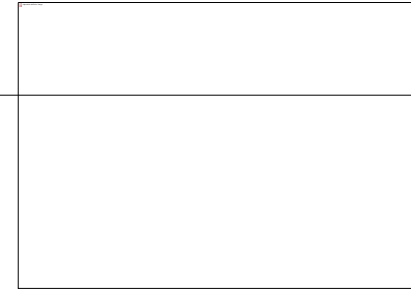
GEM (on a MSCG) : préamplification at $\approx 100 \mu\text{m}$ above the substrate
kapton foil (copper coated) with amplification holes



Discharge probability
using single, double
and triple GEMs

Gaseous detectors : MSGC + GEMs

Compass

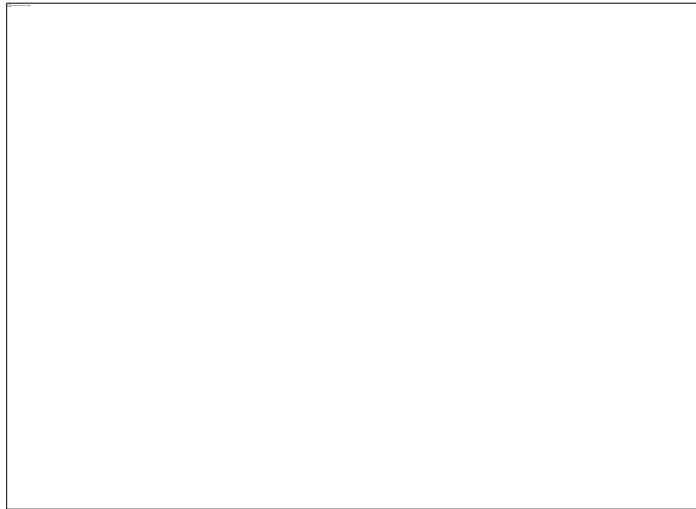


Gaseous detectors : Micromegas

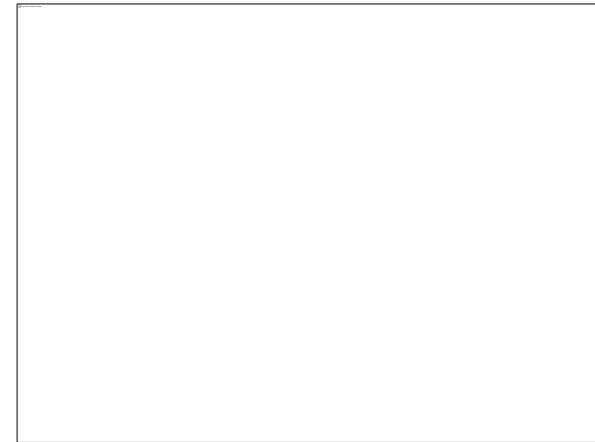
Micromegas : (G. Charpak and Y.Giomataris – 1992) is similar to a MSGC+GEM and a drift chamber. The cathode is a mesh at 100mm from the anodes (strip deposited on a substrate)



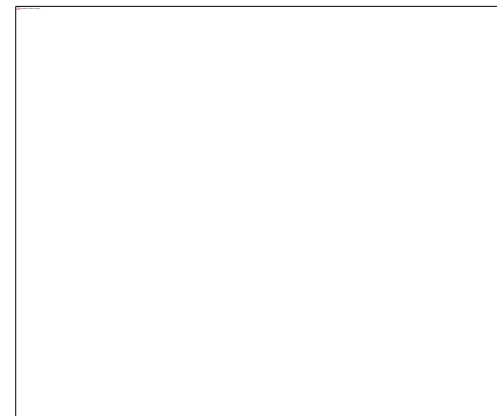
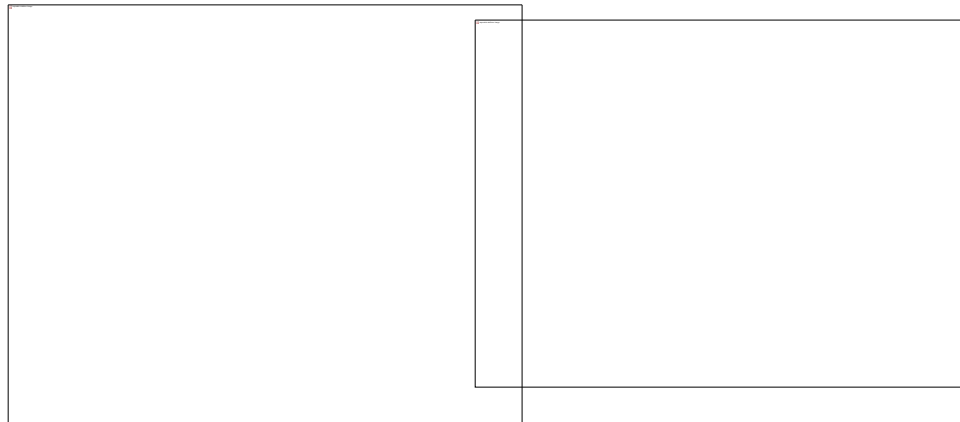
Gaseous detectors : Micromegas



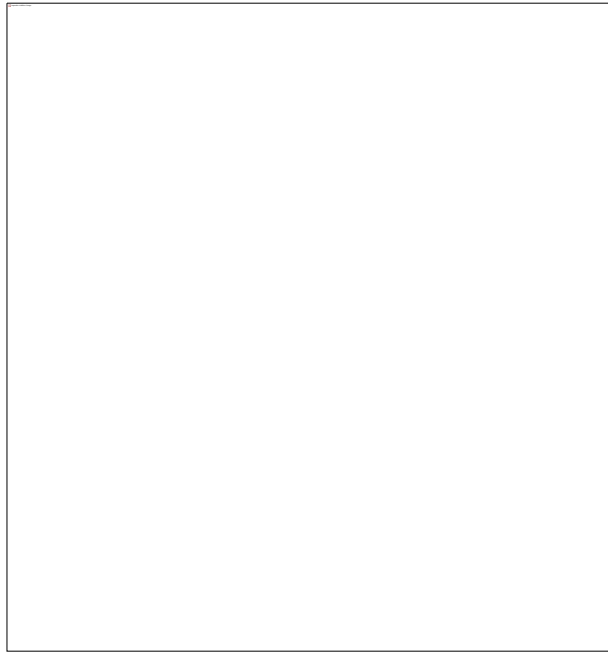
Limitation : Mesh spacers (loss of acceptance)



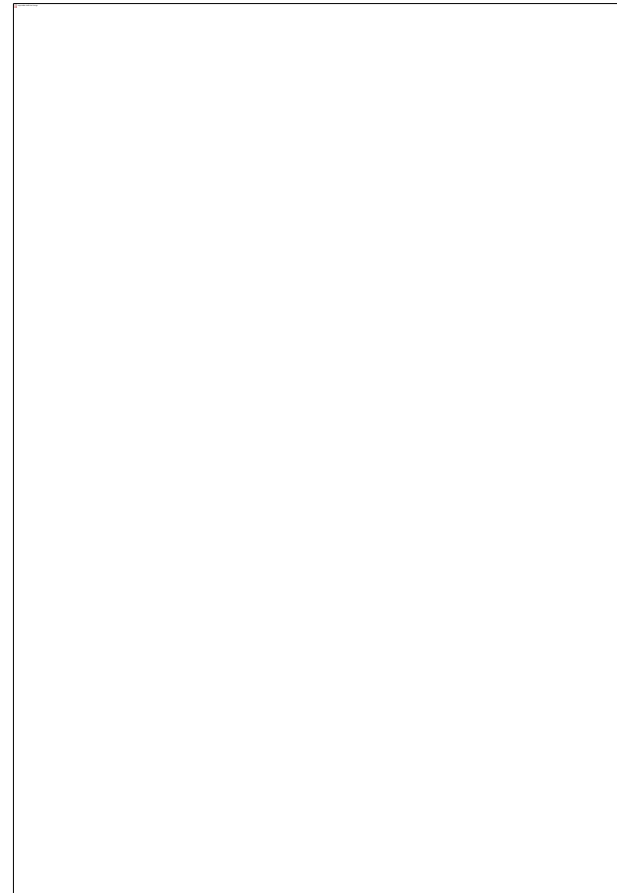
Prototype (1997)



Gaseous detectors : Micromegas



Compass
Set of 12 plates 40x40 cm²



T2K / TPC
Set of 12 plates 40x40 cm²

Gaseous detectors : Aging problems

Example : The GE1/1 project for CMS :

Adding CF_4 to the « classic Ar/CO_2 mixture to increase the time response (5 nsec)

Effect : dissociation of CF_4 leads to HF (hydrofluoric acid) which etch the copper...

Example :
Etching of the GEM
Holes (GE1/1 project)

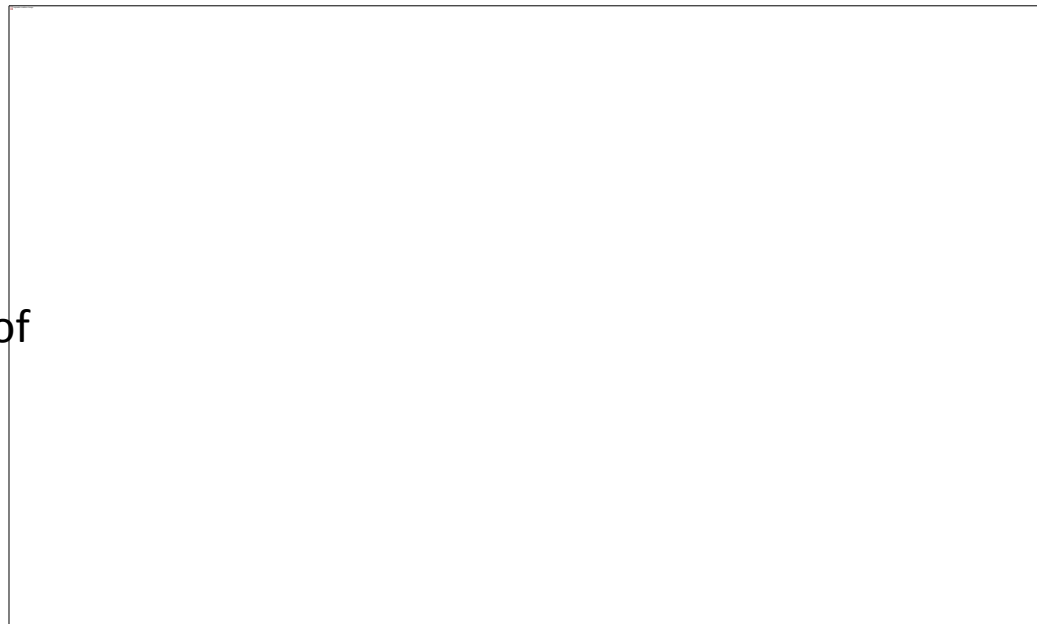


Gaseous detectors : « New » Aging problems

In future machine (HL-LHC, FCC...) at very high Luminosity, the particle flux will degrade the performances of the detectors.

One has to test the irradiation effects !
But simulation is « impossible »....

Example :
Gain versus irradiation
(corresponds to 10 years of
Operation at CMS



Gaseous detectors : Conclusions

Conclusion :

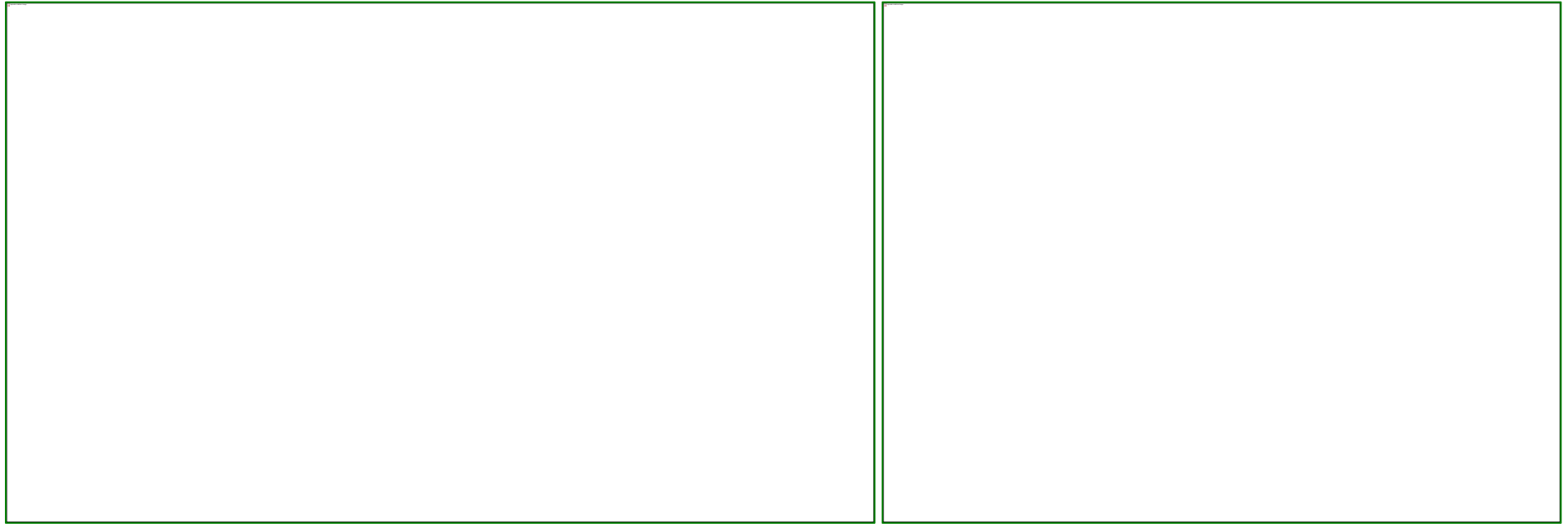
- Evolution of gaseous detectors : MGD (Micro Gap Detectors)
 - GEM
 - Micromégas
- Main problem for these detectors :
 - Cleanliness
 - Long term operation

EASY TO DESIGN

COMPLICATED TO BUILD

DIFFICULT TO OPERATE ON A LONG TERM BASIS

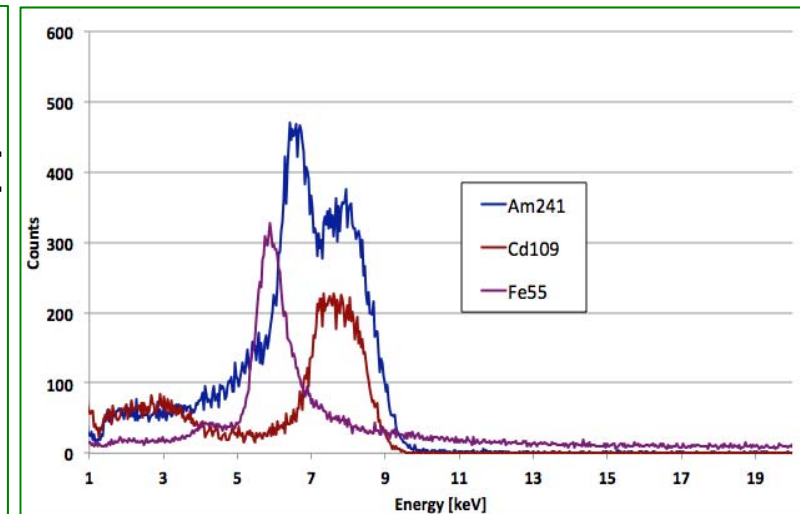
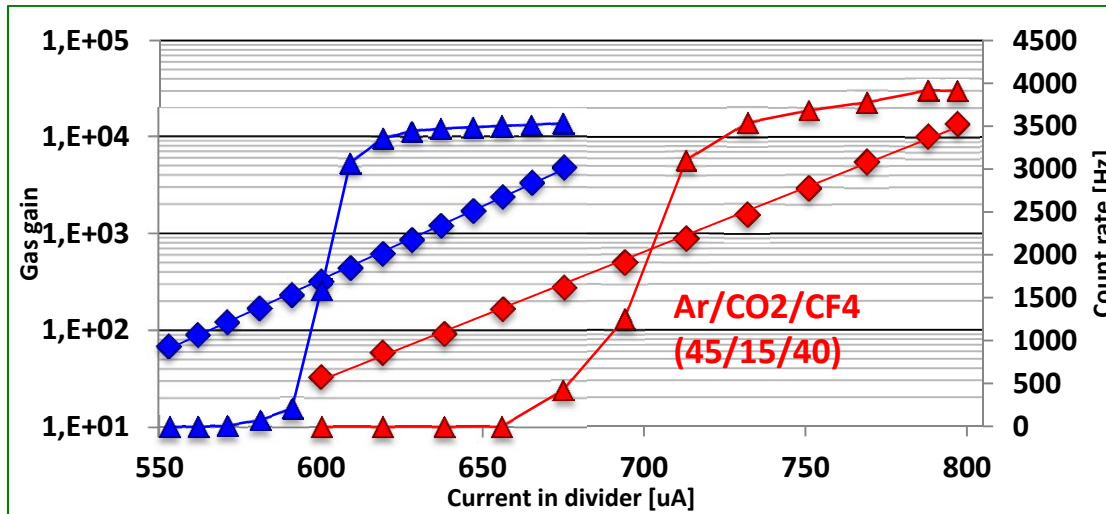
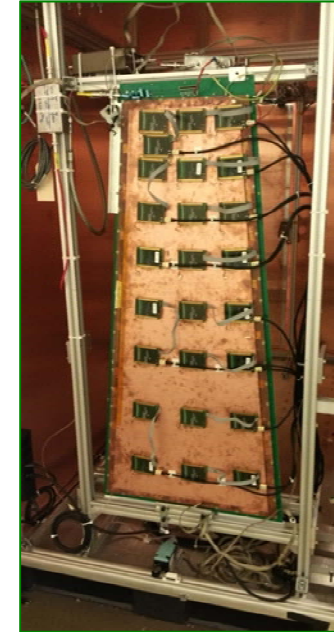
R&D on Big GEM Chambers : a nice nightmare !



R&D on Big GEM Chambers : 1. DEFINING THE PARAMETERS

Extensive study on GEM detectors :

- Basic operation with Xray sources
- Calibration tests with different :
 - gas mixtures
 - GEM geometries
 - HV power systems
 - sizes of detectors
- Defining the best configuration



R&D on Big GEM Chambers : 2. VERIFYING THE PARAMETERS

Advanced measurements and characterization:

- General understanding on GEM technology
- Comparison with past measurements
- Charging up effects (short-term stability)

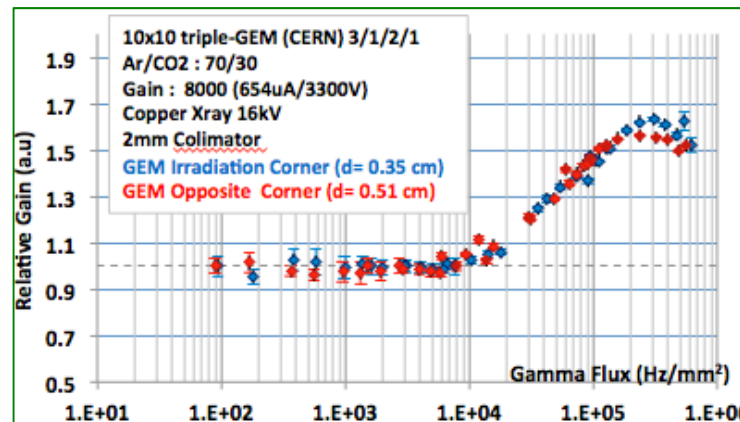
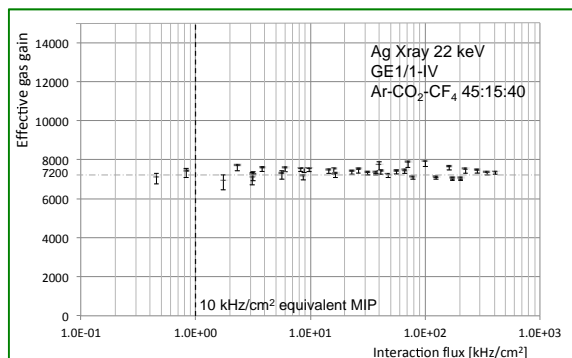
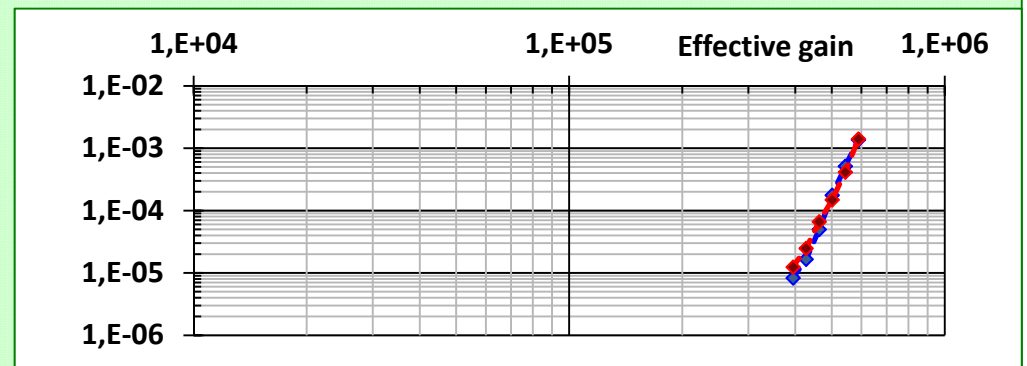
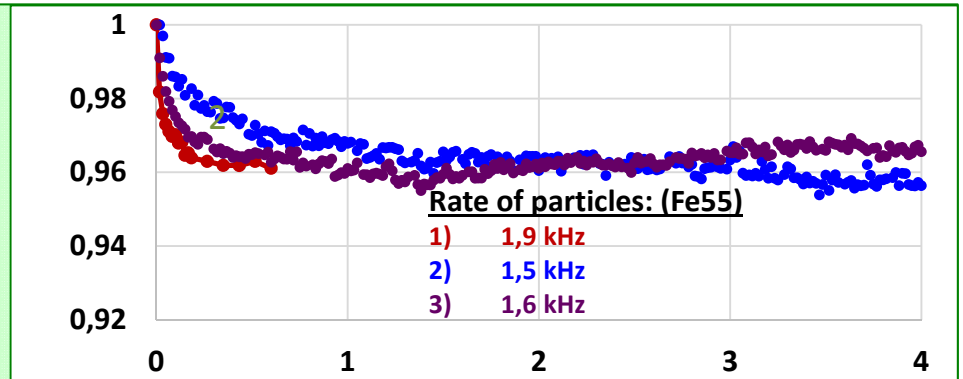
5-10 % after 1 hour

- Discharge probability

$< 10^{-12}$ at a gain of 10^4

- Rate capability

No gain loss up to 1 MHz/mm^2



R&D on Big GEM Chambers : 3. BEAM TESTS

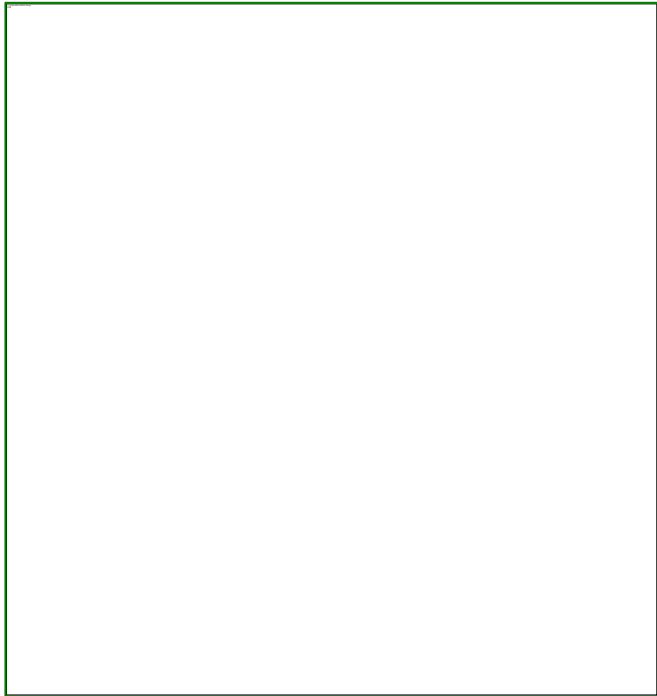
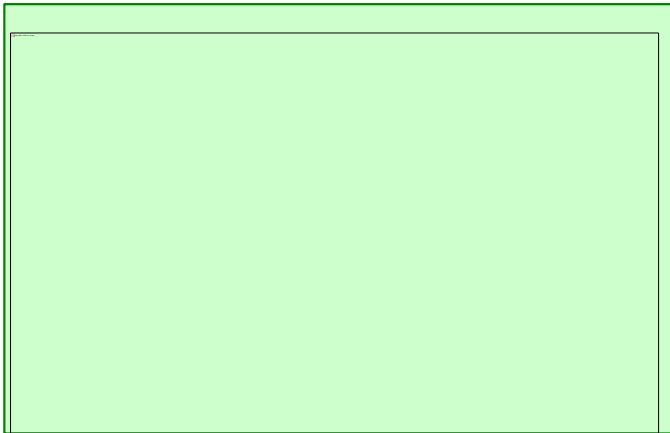
Detectors performances :

- Intense beam of **charged particles**
- All generations tested in different config. And B field.
- **Characterization** of the beam and comparison with the detector response
- Information about :
 - **Efficiency**
 - **Space resolution**
 - **Time resolution**
- Characterization of new electronics and DAQ systems

$\sigma_x < 300 \mu\text{m}$ (digital readout)

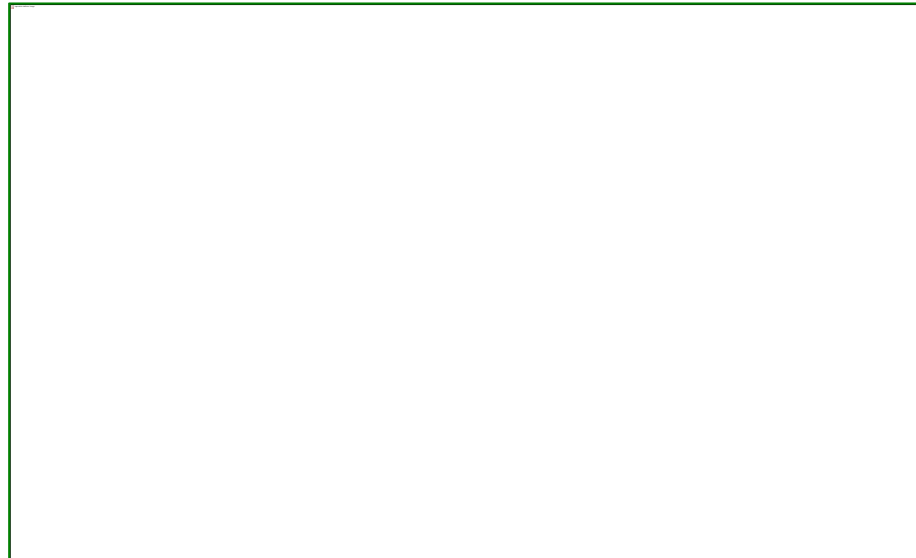
Detector efficiencies above 98%

R&D on Big GEM Chambers : 4. LONG TERM TESTS...



Aging Test at CERN GIF

- ^{137}Cs source 566 GBq
- Gamma emission 662 keV



R&D on Big GEM Chambers : 5. AGING TESTS

Outgassing Study :

- select “clean” materials to prevent self-contamination and increase longevity
- 9 materials already tested / 8 approved