

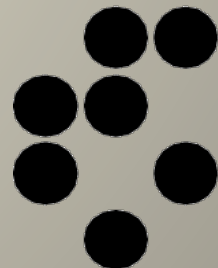
# Understanding pulse shapes at (very) high fluences

**M.Mikuž, G.Kramberger, V.Cindro, I.Mandić,  
M.Zavrtanik**

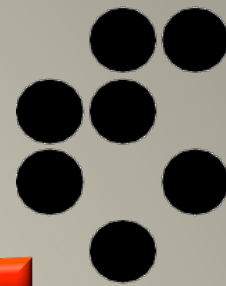
**University of Ljubljana & Jožef Stefan Institute**

**2<sup>nd</sup> TCT Workshop**

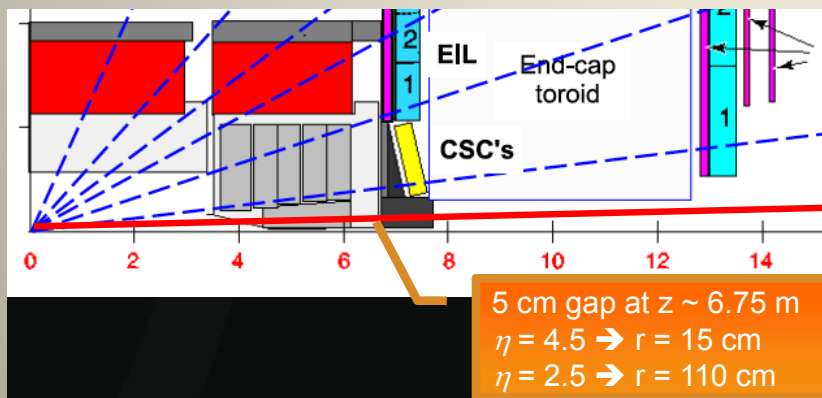
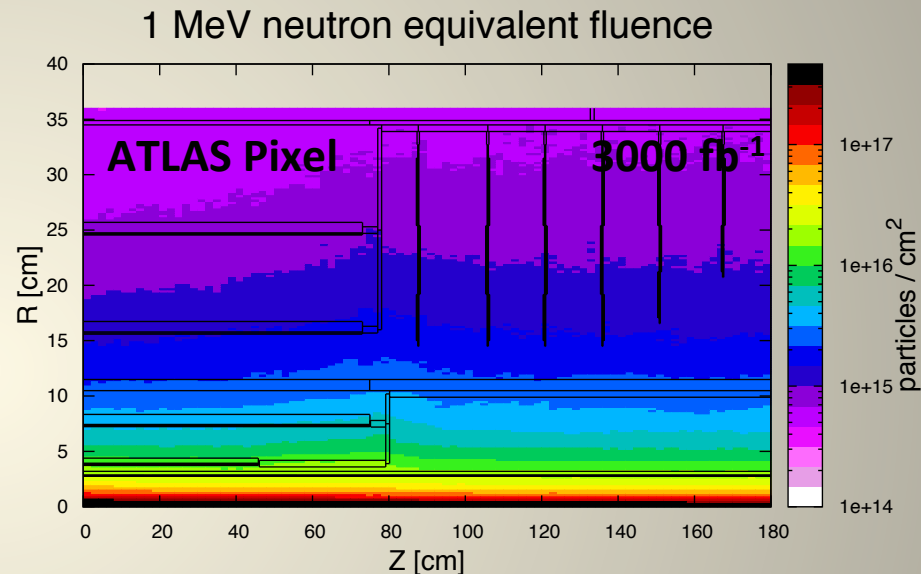
**Ljubljana, October 17<sup>th</sup>, 2016**



# Why the $10^{17}$ Ballpark ?

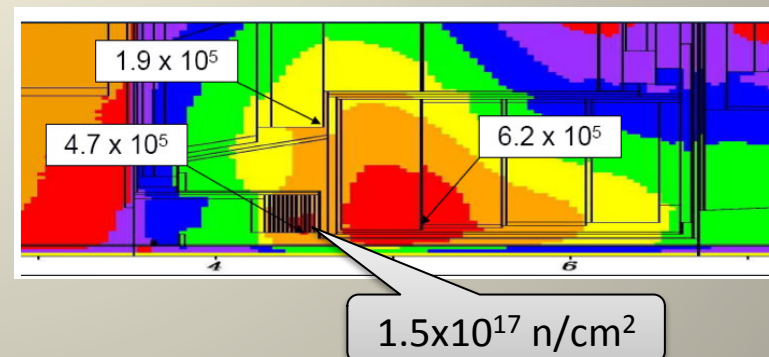


- Run1 at LHC finished, 2 under way
  - LHC trackers designed for  $730 \text{ fb}^{-1}$  of 14 TeV pp collisions,  $\sim 35 \text{ fb}^{-1}$  up to now
  - Will probably get  $\sim 1/2$  of planned
- HL-LHC in advanced planning
  - $3000 \text{ fb}^{-1}$  i.e.  $\sim 10 \times \text{LHC}$ 
    - $\sim 10^{15} n_{\text{eq}}/\text{cm}^2$  for strips (neutrons&pions)
    - $\sim 10^{16} n_{\text{eq}}/\text{cm}^2$  for pixels (pions)
    - $n \times 10^{16} n_{\text{eq}}/\text{cm}^2$  for vFW pixels ( $\pi$  &  $n$ )
    - $\sim 10^{17} n_{\text{eq}}/\text{cm}^2$  for FCAL (neutrons)
- Can (tracking) sensors survive in these extreme environments ?



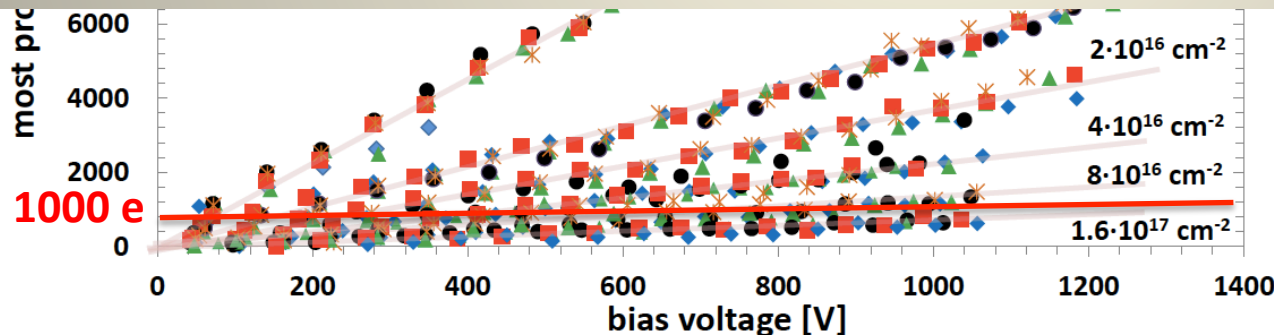
ATLAS FCAL

$3000 \text{ fb}^{-1}$



# Expectations for $10^{17} n_{eq}/\text{cm}^2$

- Linear extrapolation from low fluence data
  - Current:  $I_{leak} = 4 \text{ A/cm}^2 @ 20^\circ\text{C}$ 
    - 2 mA for 300  $\mu\text{m}$  thick 1  $\text{cm}^2$  detector @  $-20^\circ\text{C}$
  - Depletion:  $N_{eff} \approx 1.5 \times 10^{15} \text{ cm}^{-3}$ 
    - $FDV \approx 100 \text{ kV}$
  - Trapping  $\tau_{eff} \approx 1/40 \text{ ns} = 25 \text{ ps}$ 
    - $Q \approx Q_0/d v_{sat} \tau_{eff} \approx 80 \text{ e}/\mu\text{m} \cdot 200 \mu\text{m}/\text{ns} \cdot 1/40 \text{ ns} = 400 \text{ e}$  in very high electric field ( $\gg 1 \text{ V}/\mu\text{m}$ )
- Observed signal not at all compatible with expectations



From:

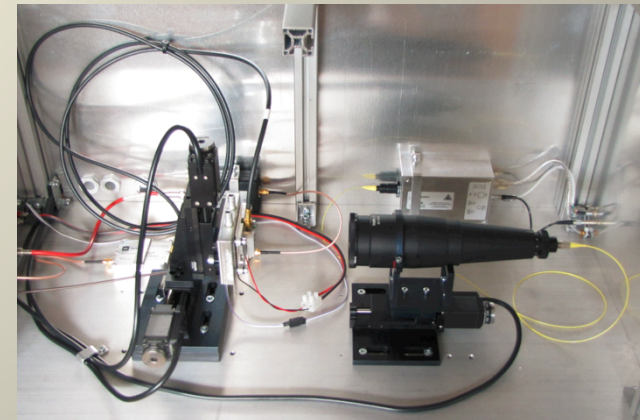
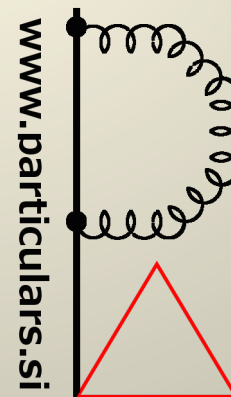
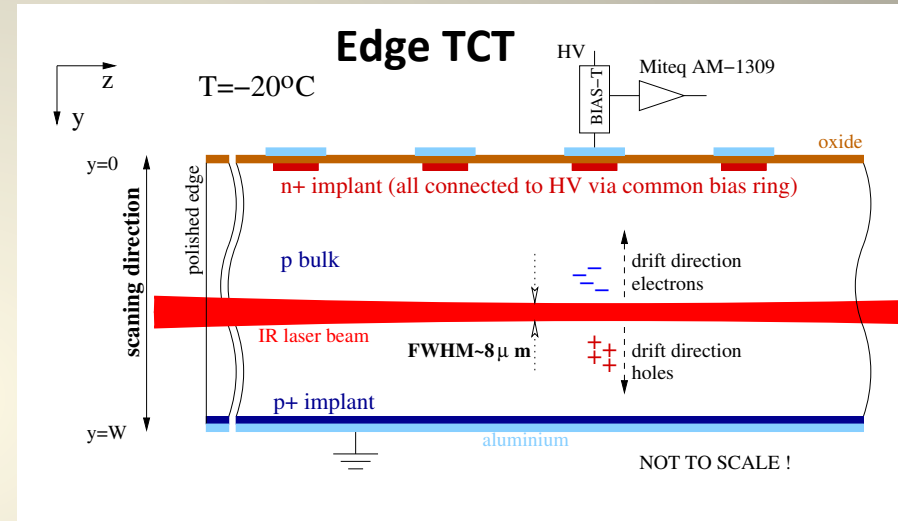
**G. Kramberger et al.,  
JINST 8 P08004 (2013).**



# Edge TCT

- Edge-TCT

- Generate charges by edge-on IR laser perpendicular to strips, detector edge polished
- Focus laser under the strip to be measured, move detector to scan
- Measure induced signal with fast amplifier with sub-ns rise-time (Transient Current Technique)
- Laser beam width  $8\text{ }\mu\text{m}$  FWHM under the chosen strip, fast (40 ps) and powerful laser
  - Caveat – injecting charge under all strips effectively results in constant weighting (albeit not electric !) field



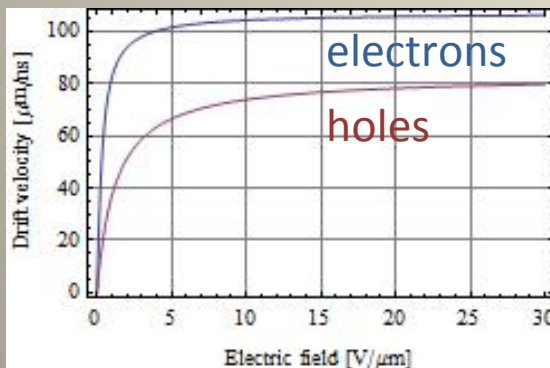


# Measured signal

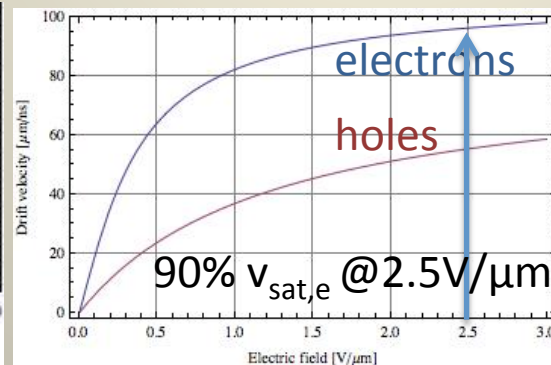
- Induced current signal given by Ramo's theorem
  - Transfer function of electronics convolutes induced current signal
  - Carriers might reach detector end
  - Velocity depends on  $E$ 
    - $v$  saturates for  $E \gg 1\text{V}/\mu\text{m}$

$$I(t=0) = q \cdot \vec{v} \cdot \vec{E}_w =$$

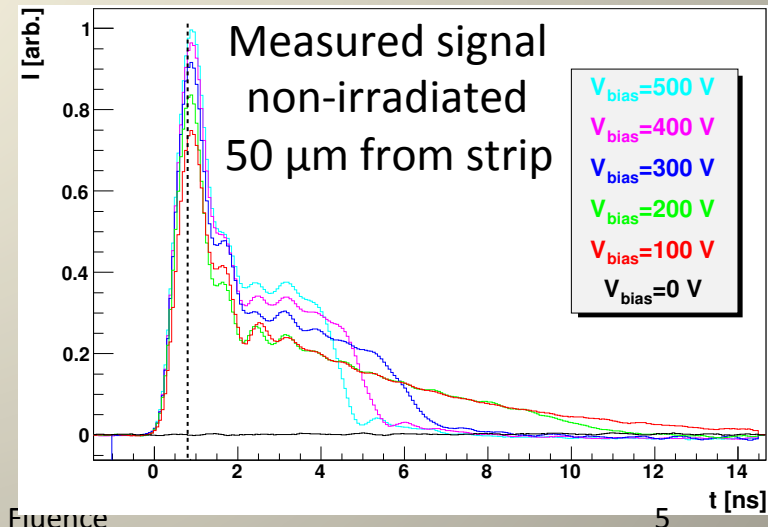
$$= N_{e-h} e_0 \cdot (v_e e^{-t/\tau_e} + v_h e^{-t/\tau_h}) / a$$



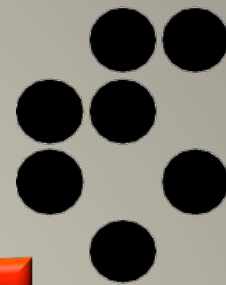
TCT, Ljubljana, Oct 17, 2016



Marko Mikuž: Pulse Shapes @ High Fluence



# Mobility Considerations FW bias

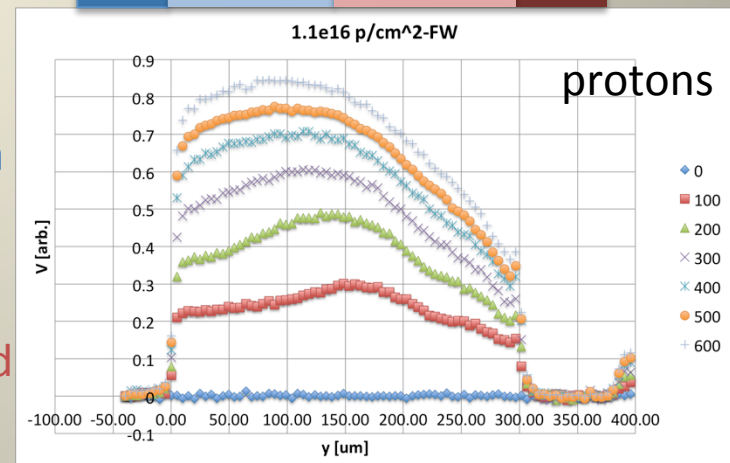
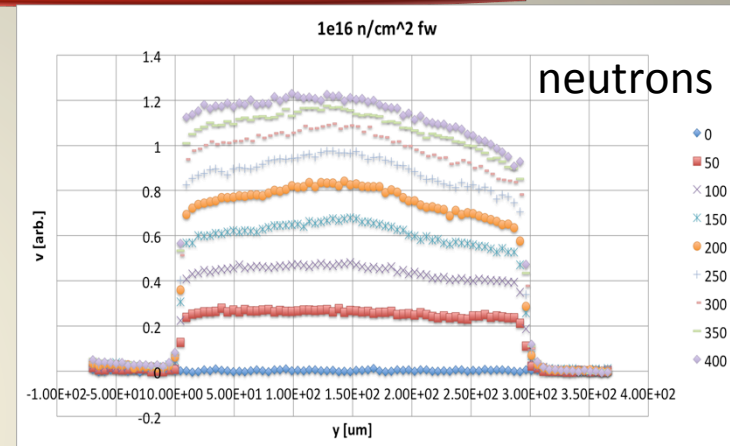


- For forward bias can extract  $v(E)$  up to a scale factor
- Observe less saturation than predicted
- Model with

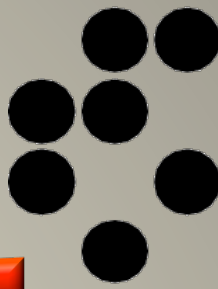
$$v_{sum}(E) = \frac{\mu_{0,e}E}{1 + \frac{\mu_{0,e}E}{v_{e,sat}}} + \frac{\mu_{0,h}E}{1 + \frac{\mu_{0,h}E}{v_{h,sat}}}$$

- keep saturation velocities at nominal values @-20°C ( $v_{e,sat} = 107 \mu\text{m/ns}$ ;  $v_{h,sat} = 83 \mu\text{m/ns}$ )
- float (common) zero field mobility degradation
- fit  $v(E)$  for  $\phi_n \geq 5 \times 10^{15}$  and  $\phi_p \geq 3 \times 10^{15}$

n.b. FW profiles less uniform for lower fluences and for protons, but departures from average field still small

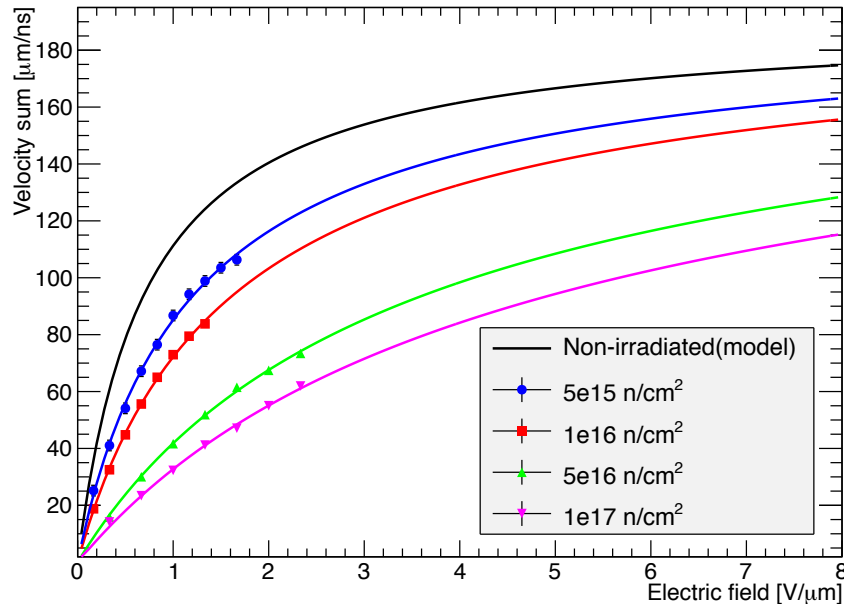


# Mobility Fits

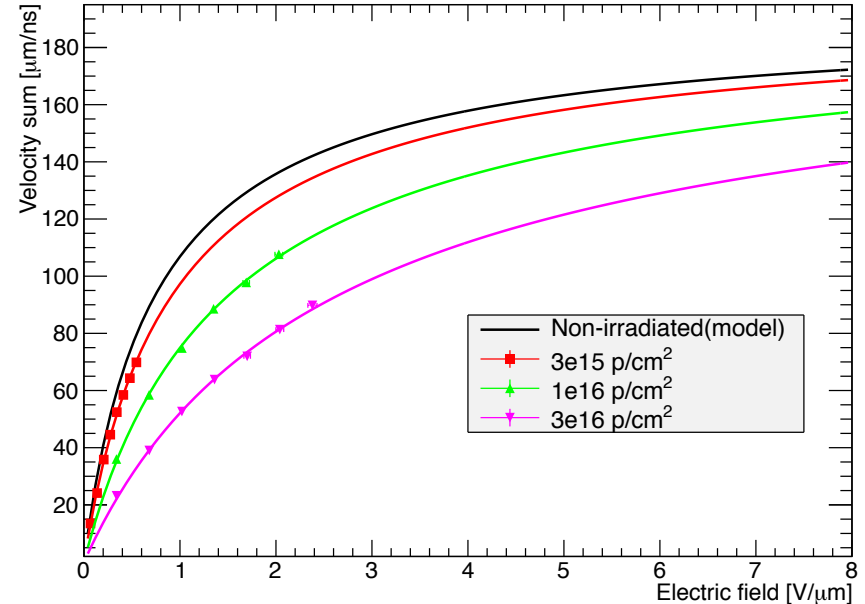


- Data follow the model perfectly
  - $\mu_0$  degradation the only free parameter, scale fixed by  $v_{sum,sat}$
  - although  $E$  range limited,  $v_{sum,max}$  still  $> 1/3$  of  $v_{sum,sat}$

Mobility neutrons

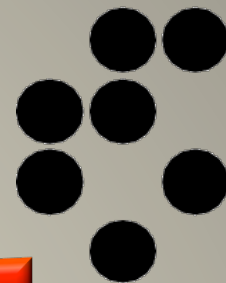


Mobility protons





# Mobility Results



- Fit to  $v_e + v_h$  with common mobility degradation factor
  - factor of **2** at  $10^{16} n_{eq}/cm^2$
  - factor of **6** at  $10^{17} n_{eq}/cm^2$
  - need **2x/6x** higher  $E$  to saturate  $v$  !

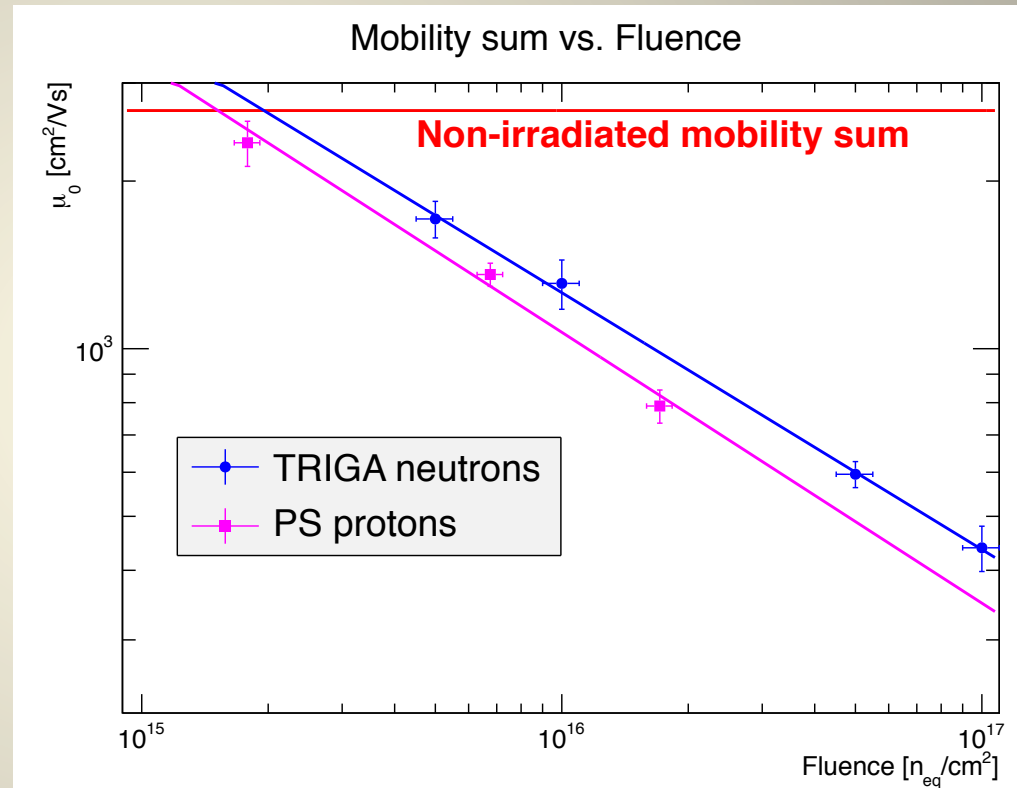
$\Phi n$	$\mu_{0,sum}$	$\Phi p$	$\mu_{0,sum}$
$[10^{15} n_{eq}/cm^2]$	$[cm^2/Vs]$	$[10^{15} n_{eq}/cm^2]$	$[cm^2/Vs]$
non-irr (model)	2680		
5	$1661 \pm 134$	1.8	$2165 \pm 212$
10	$1238 \pm 131$	6.8	$1319 \pm 67$
50	$555 \pm 32$	17	$750 \pm 54$
100	$407 \pm 40$	<b>T=-20°C</b>	

# Mobility Analysis

- Fit mobility dependence on fluence with a power law

$$\mu_{0,sum}(\Phi) = C\Phi^a$$

- Fits perfectly with  $a \approx -1/2$  indicating a single scattering process in this fluence range
  - ~same  $a$  for neutrons and protons
- Below  $\sim 10^{15} n_{eq}/cm^2$  the process gets obscured by acoustic phonon scattering
- At same equivalent fluence, mobility decrease  $\sim 20\%$  worse for protons
  - NIEL violation
- Is  $a \approx -1/2$  accidental?

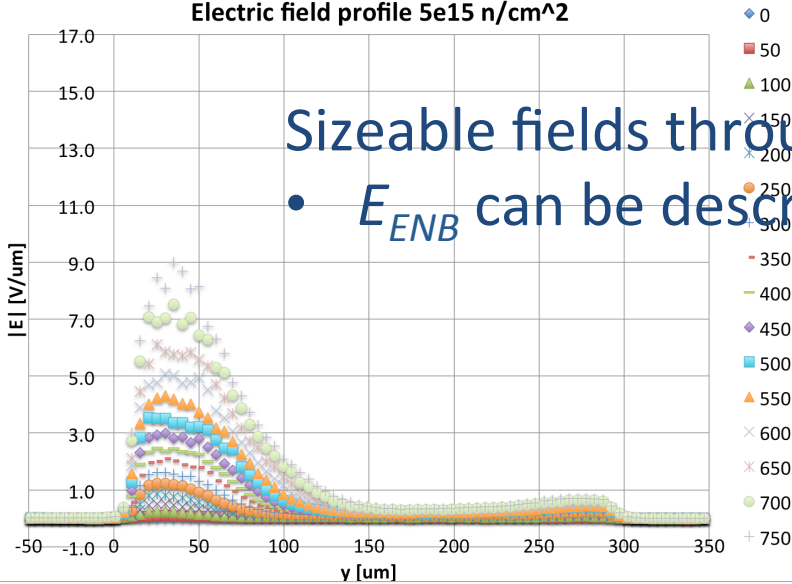


Irradiation particle	$a$	$\sigma_a$
Reactor neutrons	-0.46	0.04
PS protons	-0.49	0.05

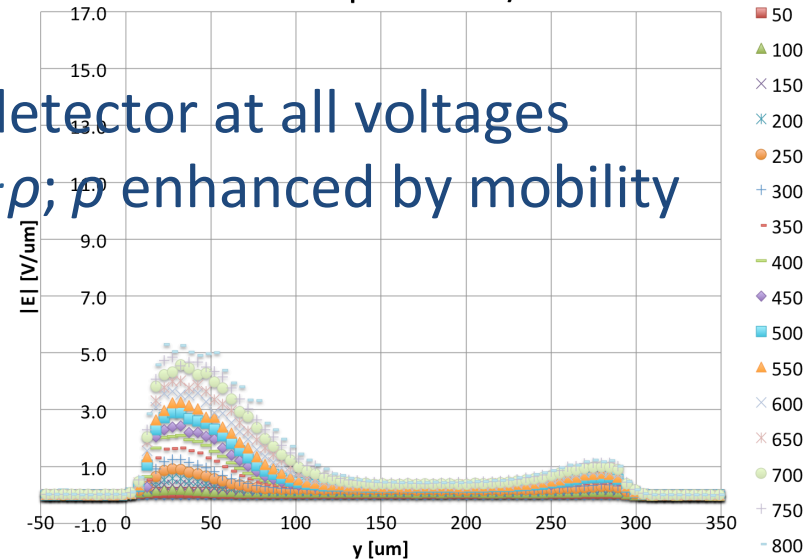
# Field Profiles Neutrons

- Sizeable fields through whole detector at all voltages
- $E_{ENB}$  can be described as  $j_{leak}/p$ ;  $p$  enhanced by mobility

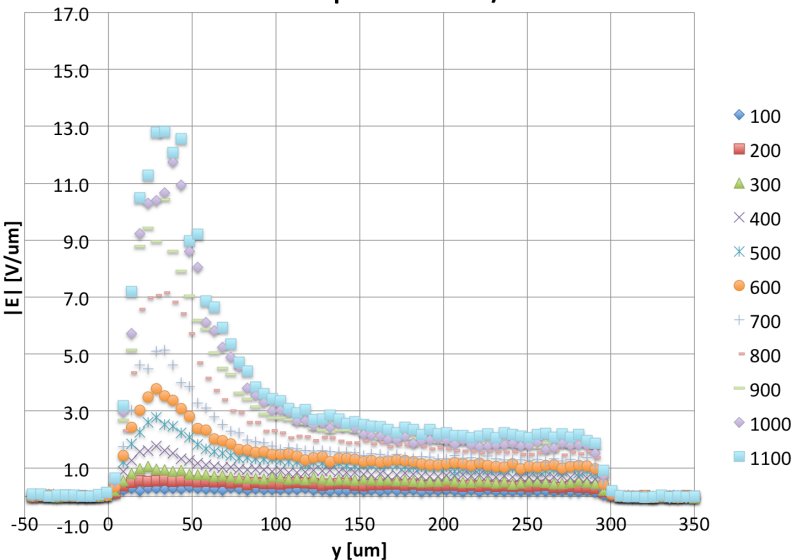
Electric field profile 5e15 n/cm^2



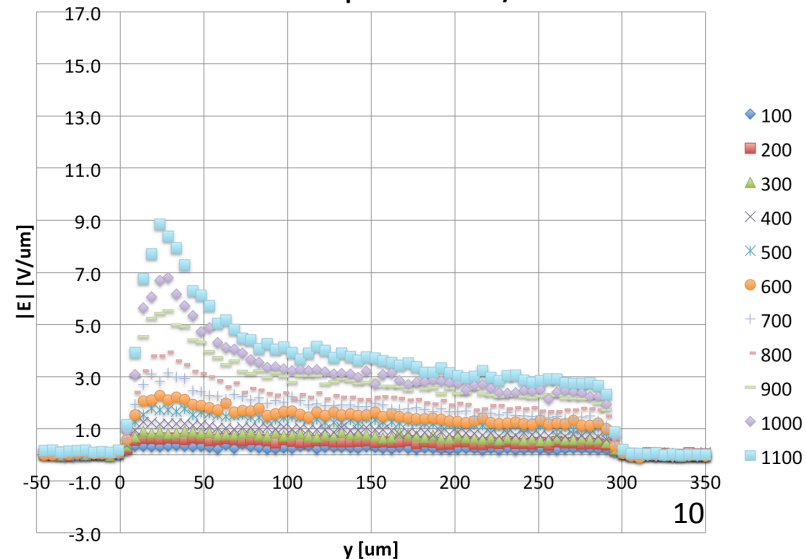
Electric field profile 1e16 n/cm^2



Electric field profile 5e16 n/cm^2

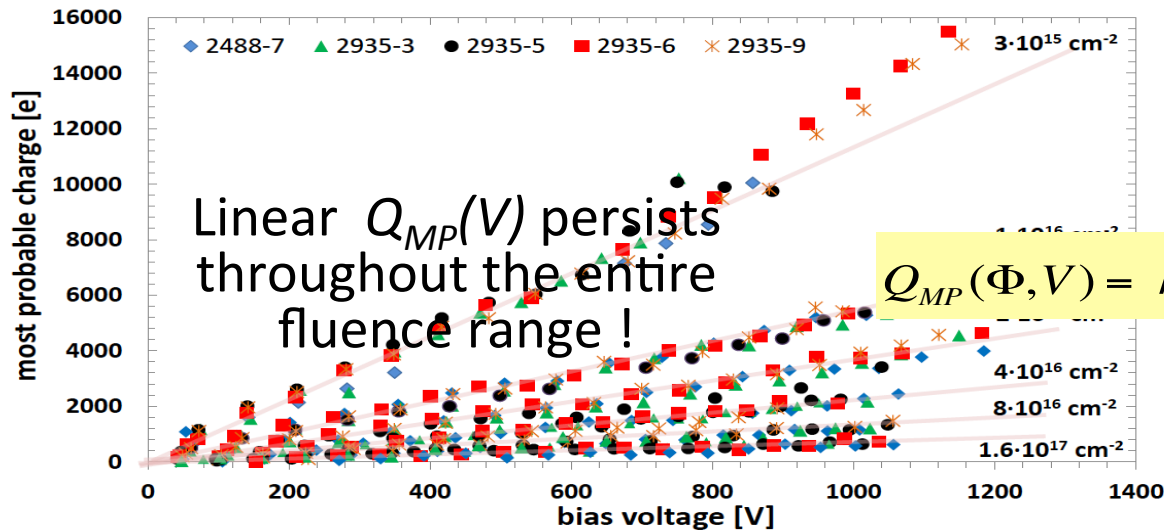


Electric field profile 1e17 n/cm^2





# Magic revisited

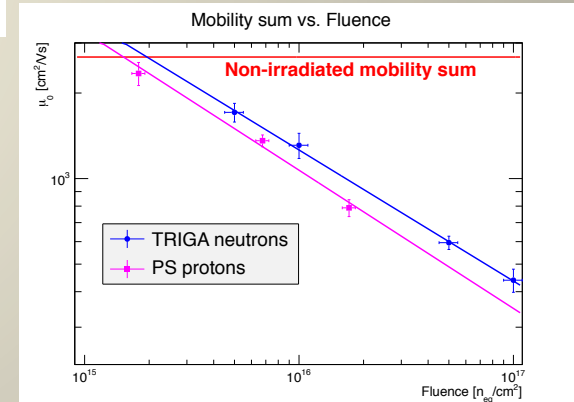


$$Q_{MP}(\Phi, V) = k \cdot (\Phi / 10^{15} \text{ n}_{eq} / \text{cm}^2)^b \cdot V$$

$$k = 26.4 \text{ e}_0 / \text{V}$$

$$b = -0.683$$

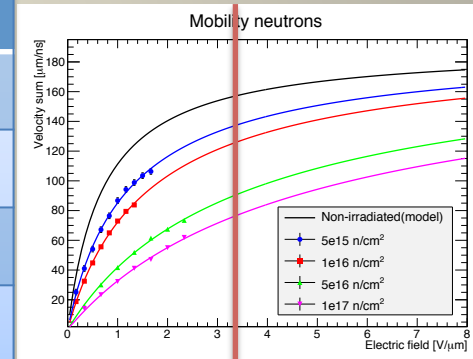
- $Q = k \cdot V$  most natural when linear  $v(E)$ 
  - not to  $E \sim 3 \text{ V}/\mu\text{m}$ , especially at low  $\Phi$
  - far from saturation, too
- Fluence dependence as  $\Phi^{-2/3}$ 
  - but mobility already decreases as  $\Phi^{-1/2}$
- Small margin left for trapping increase, certainly not linear



# Trapping Considerations

- Assumption: take  $v_{sum}$  at average  $E = 3.3 \text{ V}/\mu\text{m}$

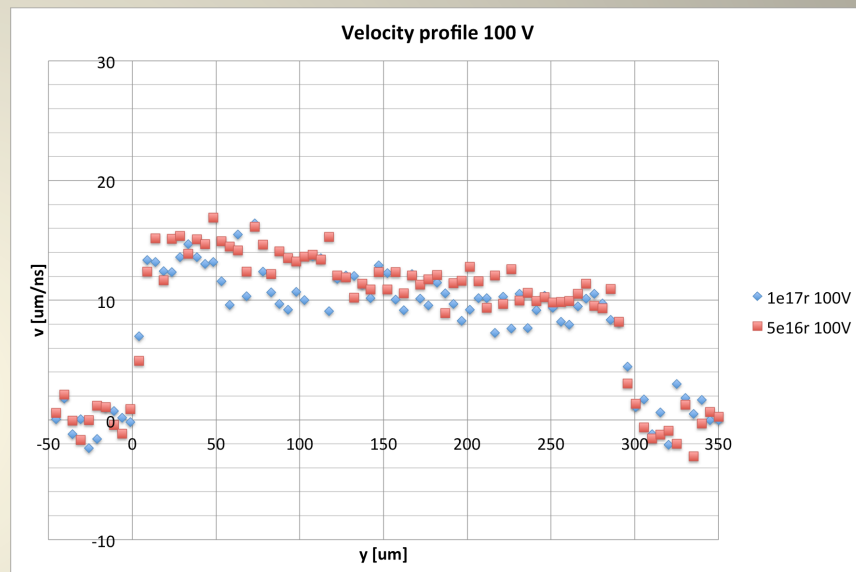
$\Phi [1\text{e}15]$	5	10	50	100
$v_{sum}(3.3 \text{ V}/\mu\text{m})$	137	126	90	77
$CCD_{1000 \text{ V}} [\mu\text{m}]$	110	70	23	14
$\tau \approx CCD/v [\text{ps}]$	800	560	260	180
$\tau_{ext} [\text{ps}]$	400	200	40	20



- Implies factor of 6-9 less trapping at highest fluences
  - weak dependence on fluence as anticipated by “-1/6” power law
  - not good when large  $E$  variations (damped by  $v(E)$ )
  - not good when  $CCD \approx$  thickness (less signal at same  $\tau$ )
  - not good when CM (more signal at same  $\tau$ )

# Another try

- Focus on cases with small and linear  $v(E) \rightarrow v(E) = v$  — —
  - 100 V at  $5 \times 10^{16}$  and  $10^{17}$  look promising – flat field
  - also the integral of  $E(x)$  yields 63/100 and 76/100 V
- Can assume linear  $v(E)$  in whole detector
  - assume same ratio as for low fluences
  - less trapping compared to linear extrapolation by factors of 3.2 and 5.4
- Might overestimate trapping as field stil peaks close to strip

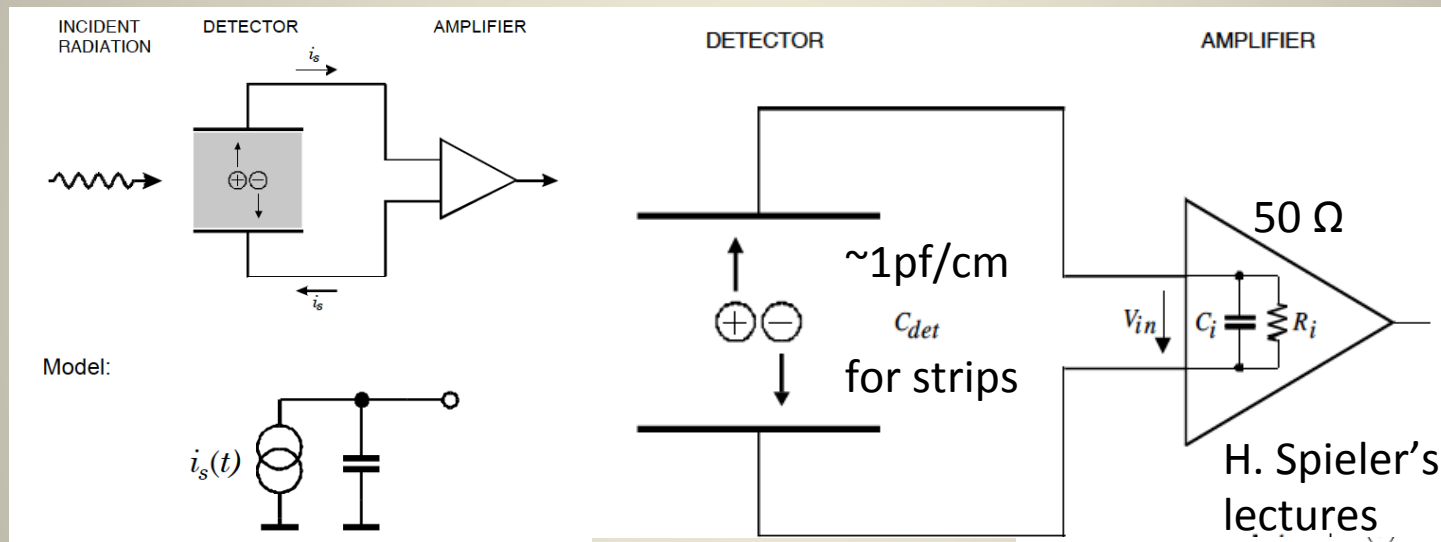


$\Phi$	$\tau_e[\text{ps}]$	$\tau_h[\text{ps}]$
5e16	145	110
1e17	122	93

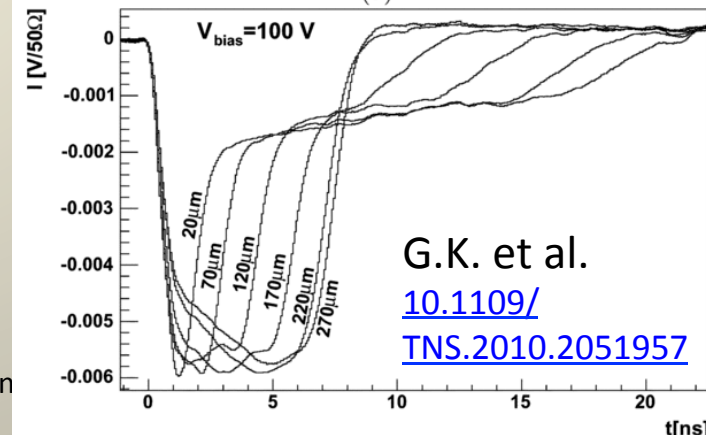


# Textbook Waveforms

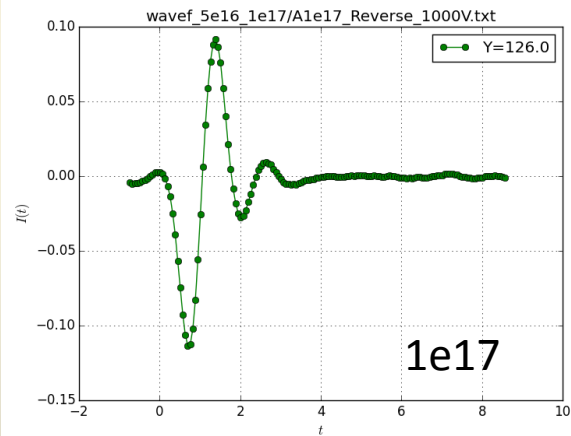
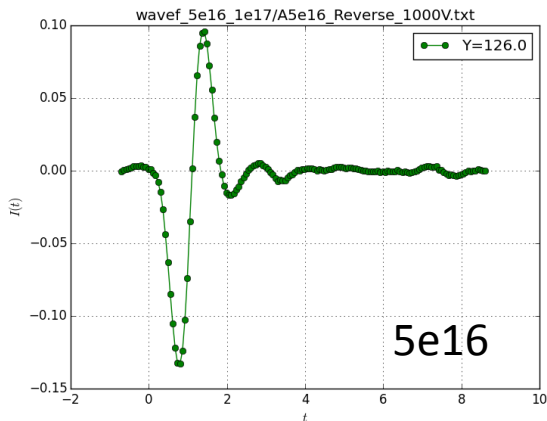
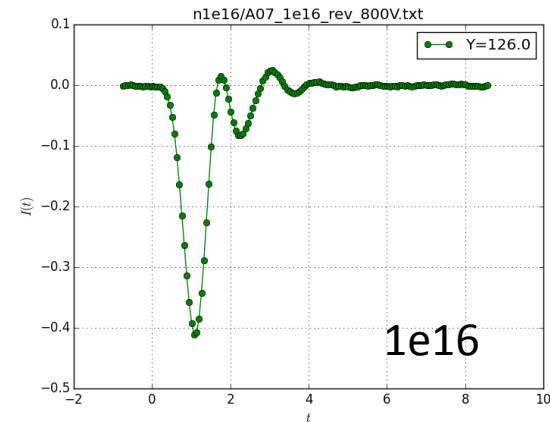
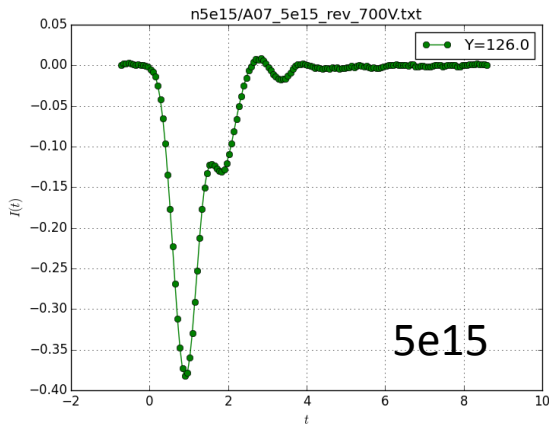
- You think of your system like this



- $H(\omega) = 1/(1+i\omega RC)$ ;  $RC \sim 0.1\text{ ns}$ ;  $\nu_{cutoff} \sim 1.5\text{ GHz}$ 
  - $\omega \gg 1/RC \rightarrow$  integrator
  - $\omega \ll 1/RC \rightarrow$  current amplifier ✓



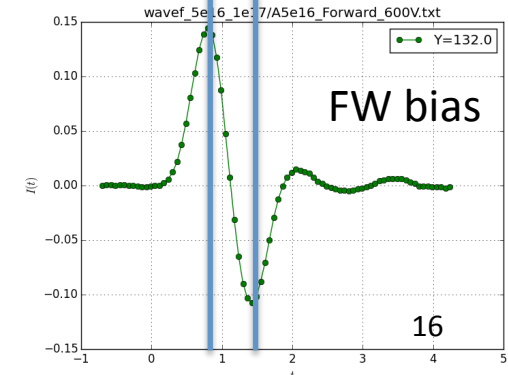
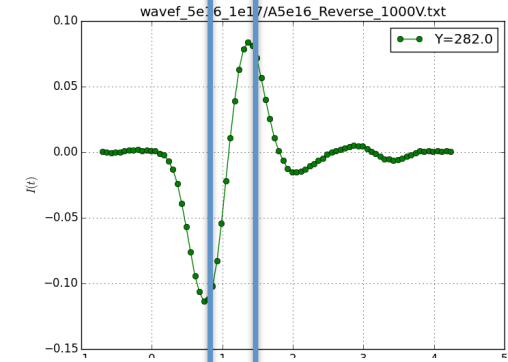
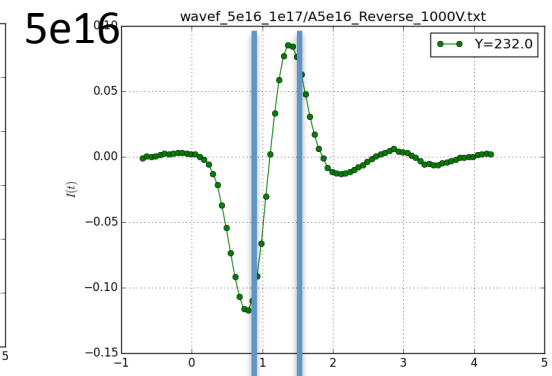
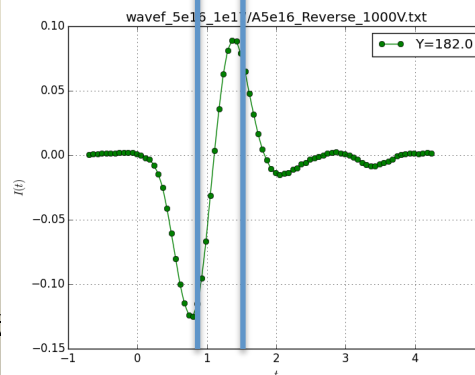
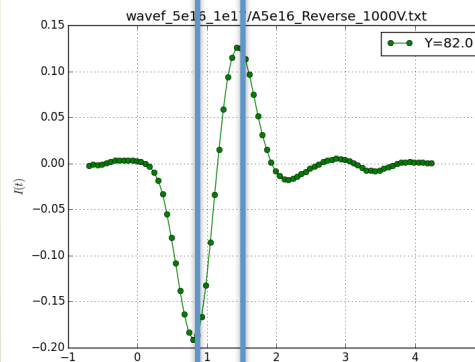
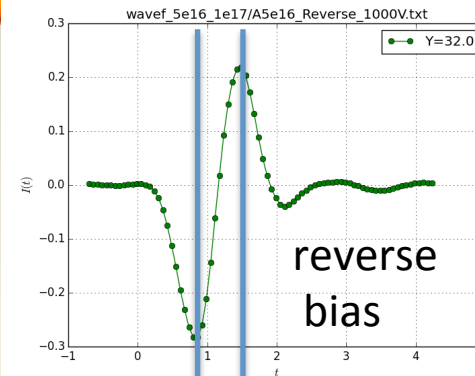
# TCT Waveforms @ High Fluence



- Nothing like textbook behaviour !
  - damped oscillations, influenced by trapping time ?

# WF – position dependence ?

- Waveforms plotted every 50  $\mu\text{m}$  in detector depth for reverse bias at 1000 V
- Forward bias in middle of detector added at 600 V
- Very little, if any, wf dependence on position observed
- Trapping not position (even not bias) dependent !?





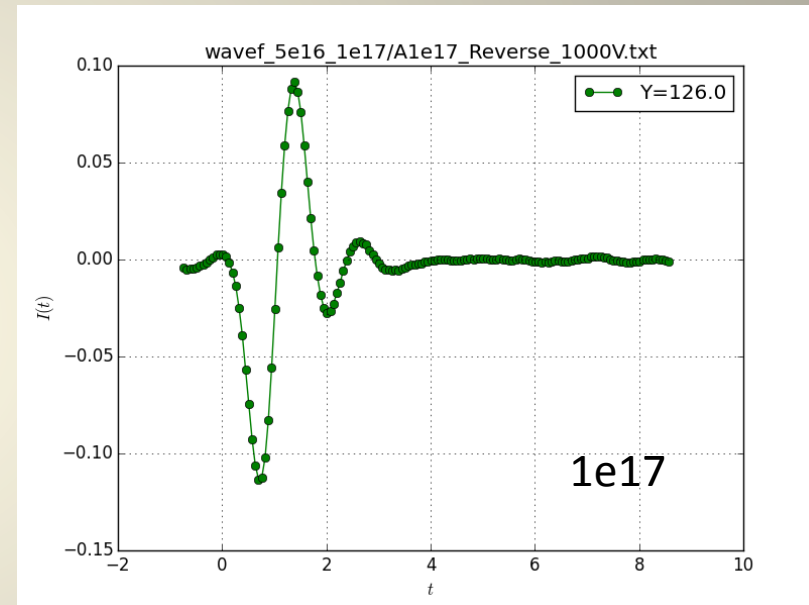
# Transfer function

- Waveforms at  $y=100\text{ }\mu\text{m}$ , 1000 V,  $1e10^{17}$

- $E \approx 3\text{ V}/\mu\text{m}$ , CCD implies signal within  $\sim 15\text{ }\mu\text{m}$  or 0.2 ns
  - the rest you see is the transfer function of the system

- Observed features

- damped oscillations with  $t_0 \sim 1.5\text{ ns}$ ,  $\tau \sim 1\text{ ns}$
- definitely under-critical
  - $\omega = 2\pi/t_0 = 4\text{ ns}^{-1} \sim 4\beta$ ;  $\omega \sim \omega_0$
- simple  $RLC$  hard to imagine with  $R=50\text{ }\Omega$  and  $C \sim 1\text{ pF}$ 
  - $L \sim 30\text{-}60\text{ nH}$  from  $\omega_0^2 = 1/LC$ ;  $\beta = R/2L$



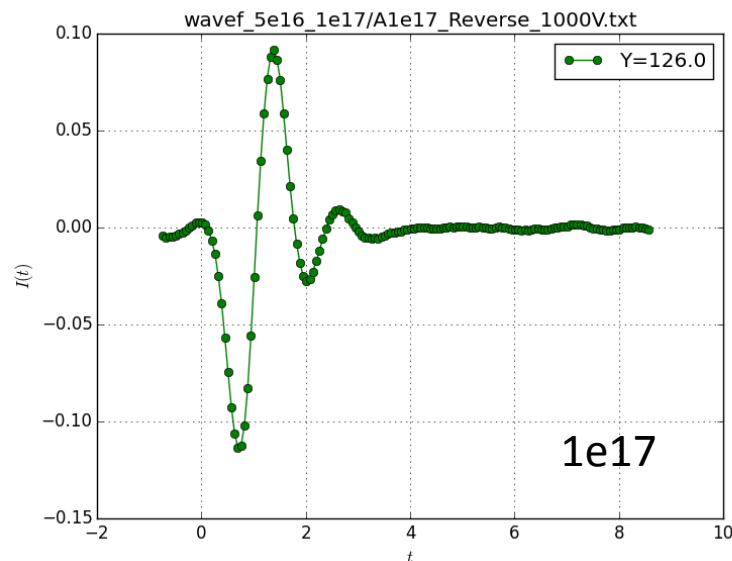
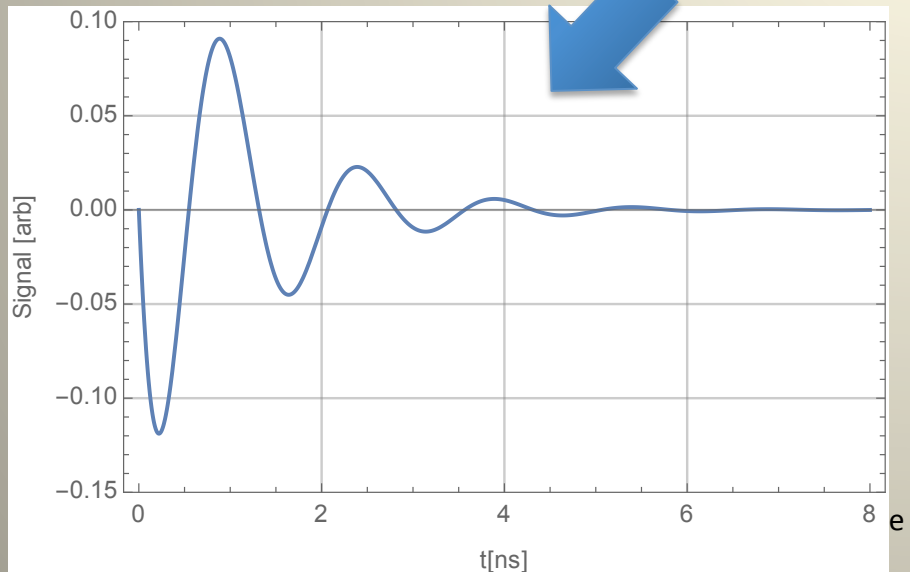
# Transfer function

- Construct  $H(t)$  with observed features
  - $t_0 = 1.5 \text{ ns}$ ,  $\tau = 1.1 \text{ ns}$
- Signal
  - $\tau = 0.2 \text{ ns}$

$$H(t) = e^{-\beta t} \cos \omega t$$

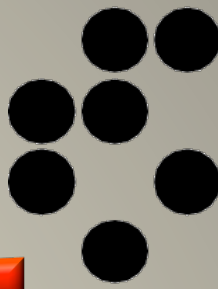
$$I(t) = \int_0^t I_0 e^{-t'/\tau} H(t-t') dt'$$

$$I(t) = A\{(\beta - 1/\tau)[e^{-t/\tau} - e^{-\beta t} \cos \omega t] + \omega e^{-\beta t} \sin \omega t\}$$





# At 5e16

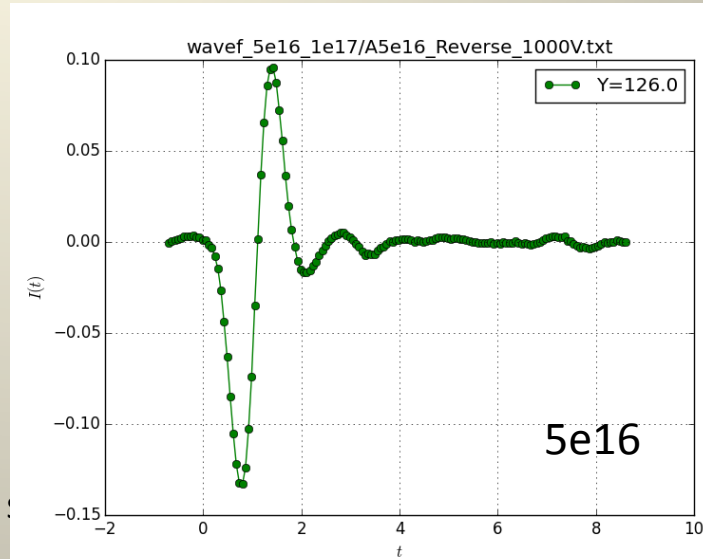
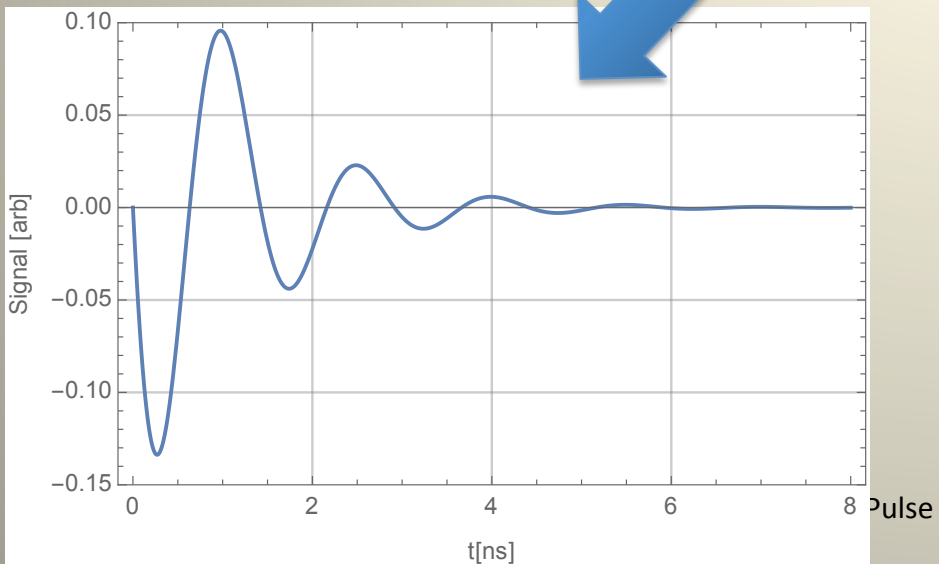


- Same  $H(t)$ 
  - $t_0 = 1.5$  ns,  $\tau = 1.1$  ns
- Signal
  - $\tau = 0.4$  ns

$$H(t) = e^{-\beta t} \cos \omega t$$

$$I(t) = \int_0^t I_0 e^{-t'/\tau} H(t-t') dt'$$

$$I(t) = A\{(\beta - 1/\tau)[e^{-t/\tau} - e^{-\beta t} \cos \omega t] + \omega e^{-\beta t} \sin \omega t\}$$

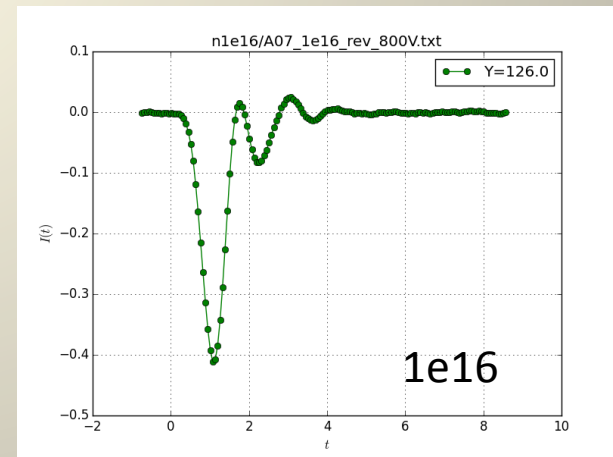
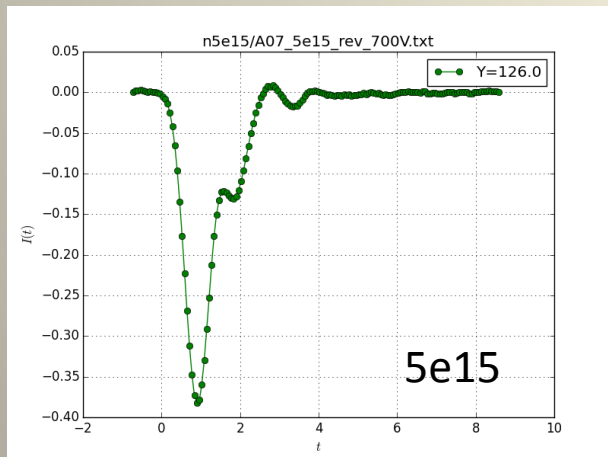
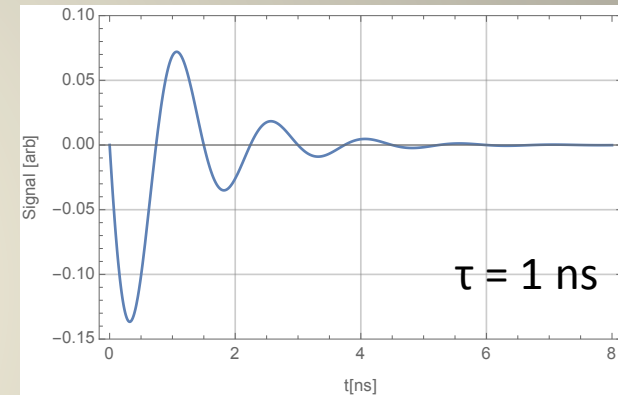


# At 5e15 and 1e16

- No way to reproduce signal, although basic features of  $H(t)$  still there

$$I(t) = A\{(\beta - 1/\tau)[e^{-t/\tau} - e^{-\beta t} \cos \omega t] + \omega e^{-\beta t} \sin \omega t\}$$

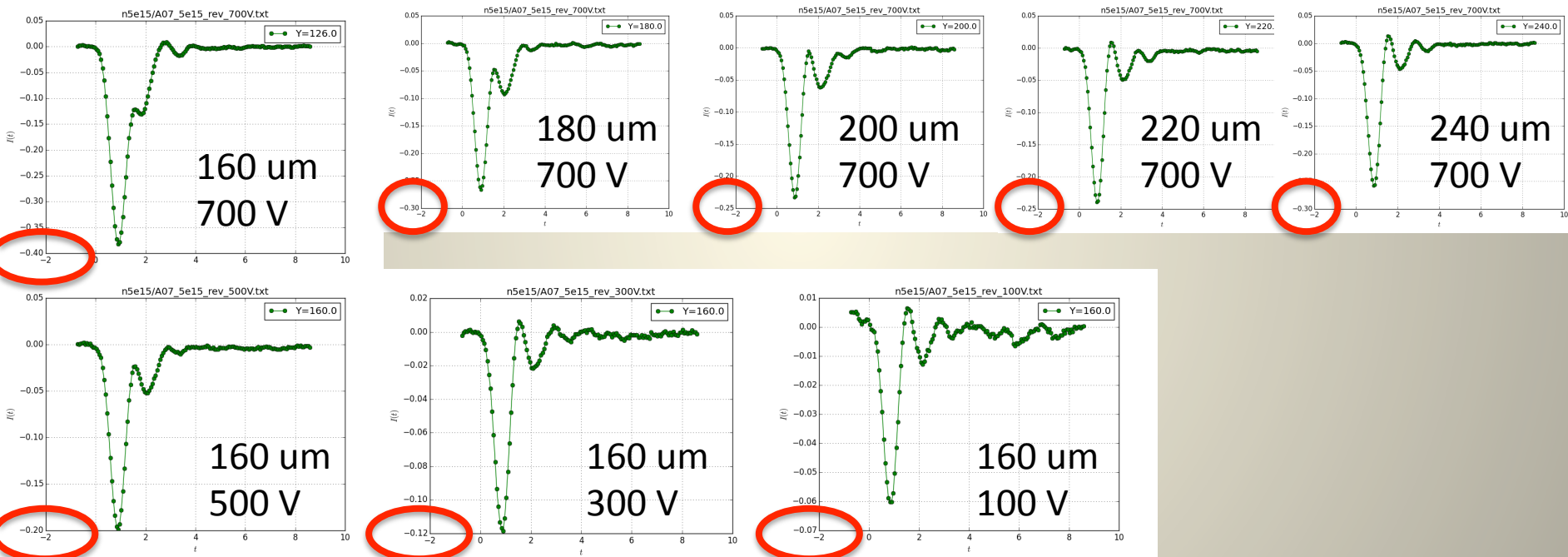
- border ? CM ?





# 5e15: position, voltage

- border (CM@strip) ? – further away
- CM ? – lower V



- qualitatively makes sense
  - still no hope fitting with same  $H(t)$  and  $\tau$  ?
- not the same run as 5e16, 1e17 -> did  $H(t)$  change ?

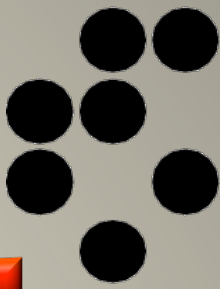


# Summary



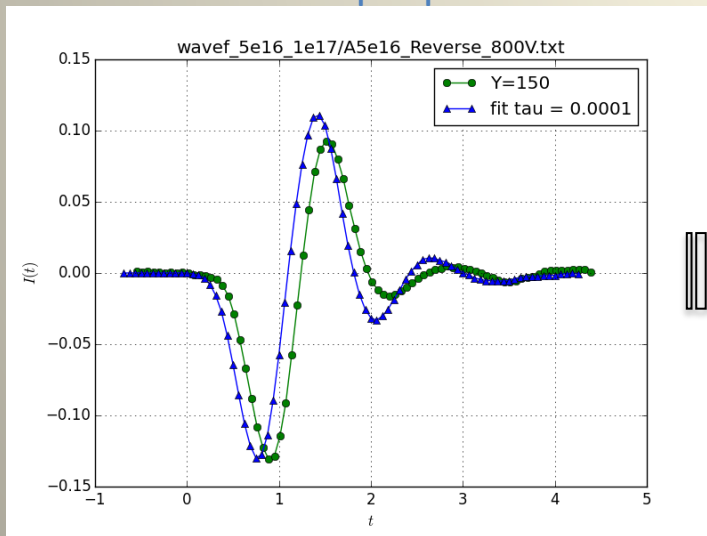
- TCT waveforms studied for Si detectors irradiated
  - with neutrons from  $5 \times 10^{15}$  to  $10^{17} n_{eq}/cm^2$
- Prior knowledge on el. field, mobility and trapping taken into account
  - expected trapping times  $O(100 \text{ ps}) @ 10^{17}$
- Transfer function constructed from observed characteristics @  $10^{17}$ 
  - no reasonable equivalent *RLC* found (so far ?)
- Trapping convoluted with  $H(t)$  reasonably matches *WF* at  $10^{17}$  and  $5 \times 10^{16}$
- No match for  $5 \times 10^{15}$  and  $10^{16}$ 
  - position/voltage dependent *WF* (CM ?)
  - different  $H(t)$  ?

# Backup Slides

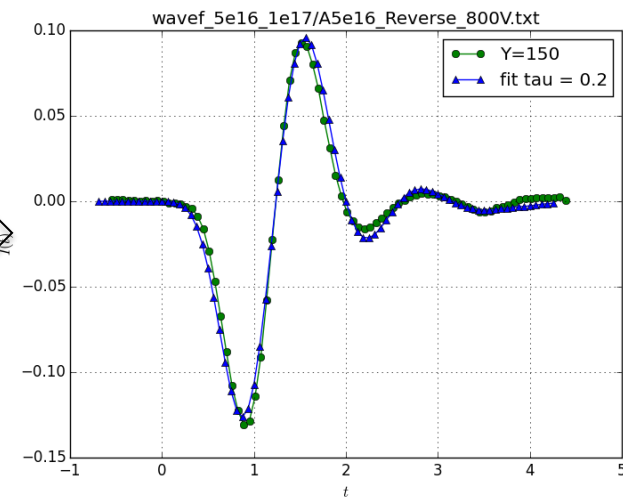


# Waveforms –direct method

- Waveforms at  $y=100\text{ }\mu\text{m}$ , 800 V,  $5 \times 10^{16}$  and  $10^{17}$ 
  - $E \approx 3\text{ V}/\mu\text{m}$ , CCD implies signal within  $\sim 15\text{ }\mu\text{m}$  or  $\sim 0.2\text{ ns}$ 
    - the rest you see is the transfer function of the system
- Still distinct signals from the two fluences
  - treat  $10^{17}$  waveform as transfer function of the system
    - convolute with  $e^{-t/\tau}$  to match  $5 \times 10^{16}$  response
    - $\tau = 0.2\text{ ns}$  provides a good match
- In fact, measure  $\Delta\tau$ , as “transfer” already convoluted with  $e^{-t/\tau(1e17)}$  !
  - Should do proper Fourier analysis... but looks consistent

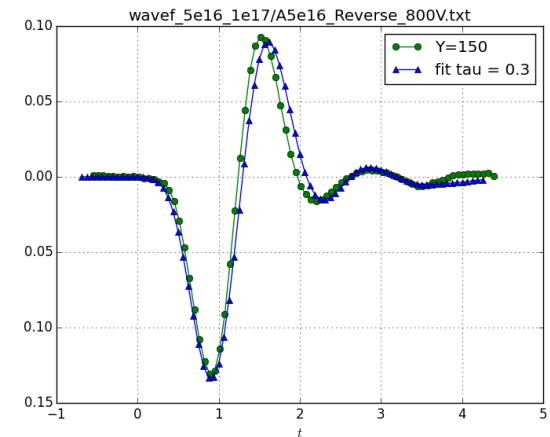
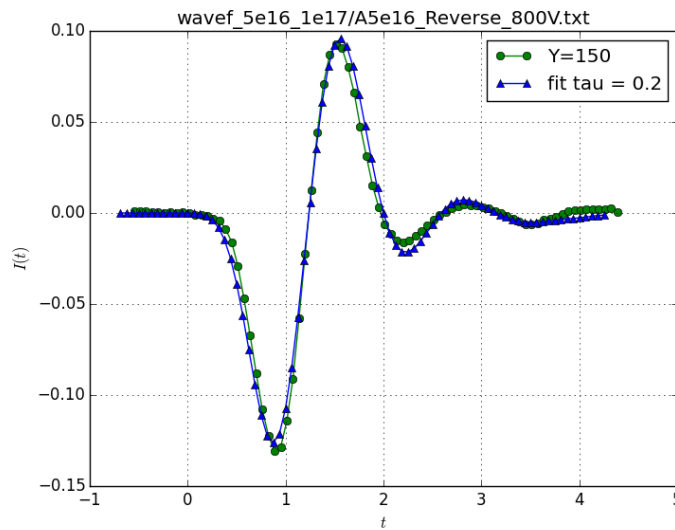
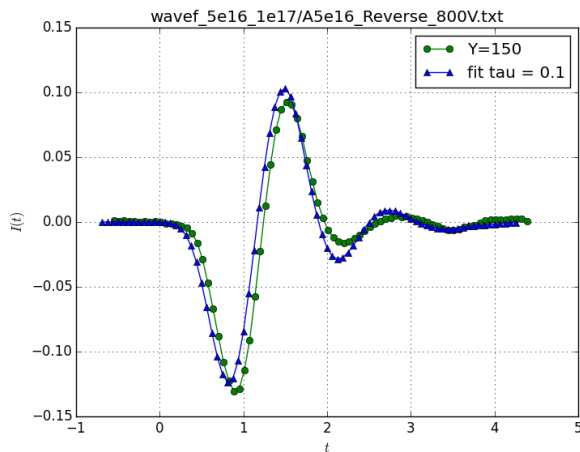
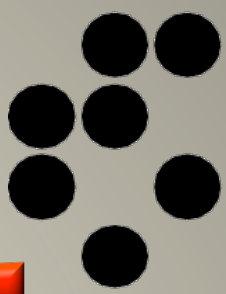


$\tau = 0.2\text{ ns}$





# Waveforms: How sensitive ?



- $\Delta\tau = 0.2$  ns certainly best fit, 0.1 too narrow, 0.3 too broad
- precision  $\sim 50$  ps