

Neutrinos: the quest for a new physics scale

27-31 March 2017, CERN

Neutrino Secret Interaction (νSI)



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Mainz University



Outline

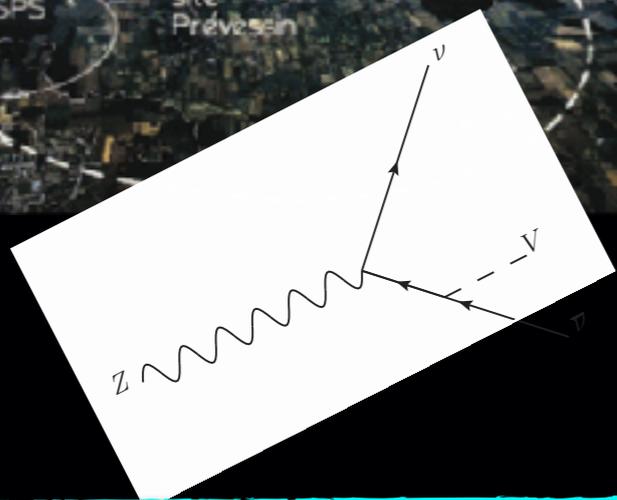
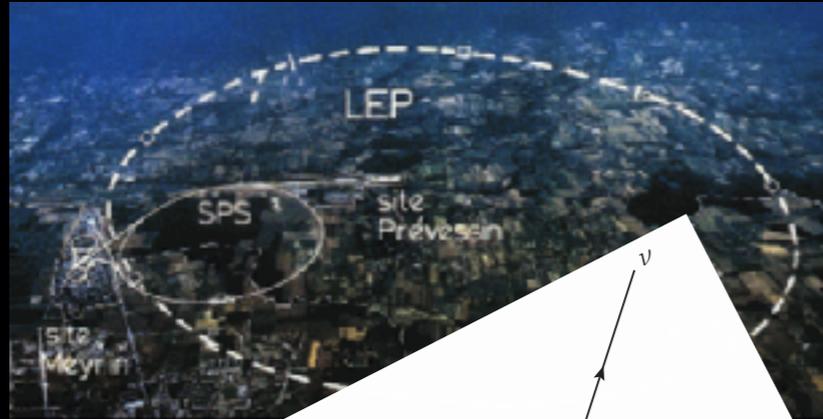
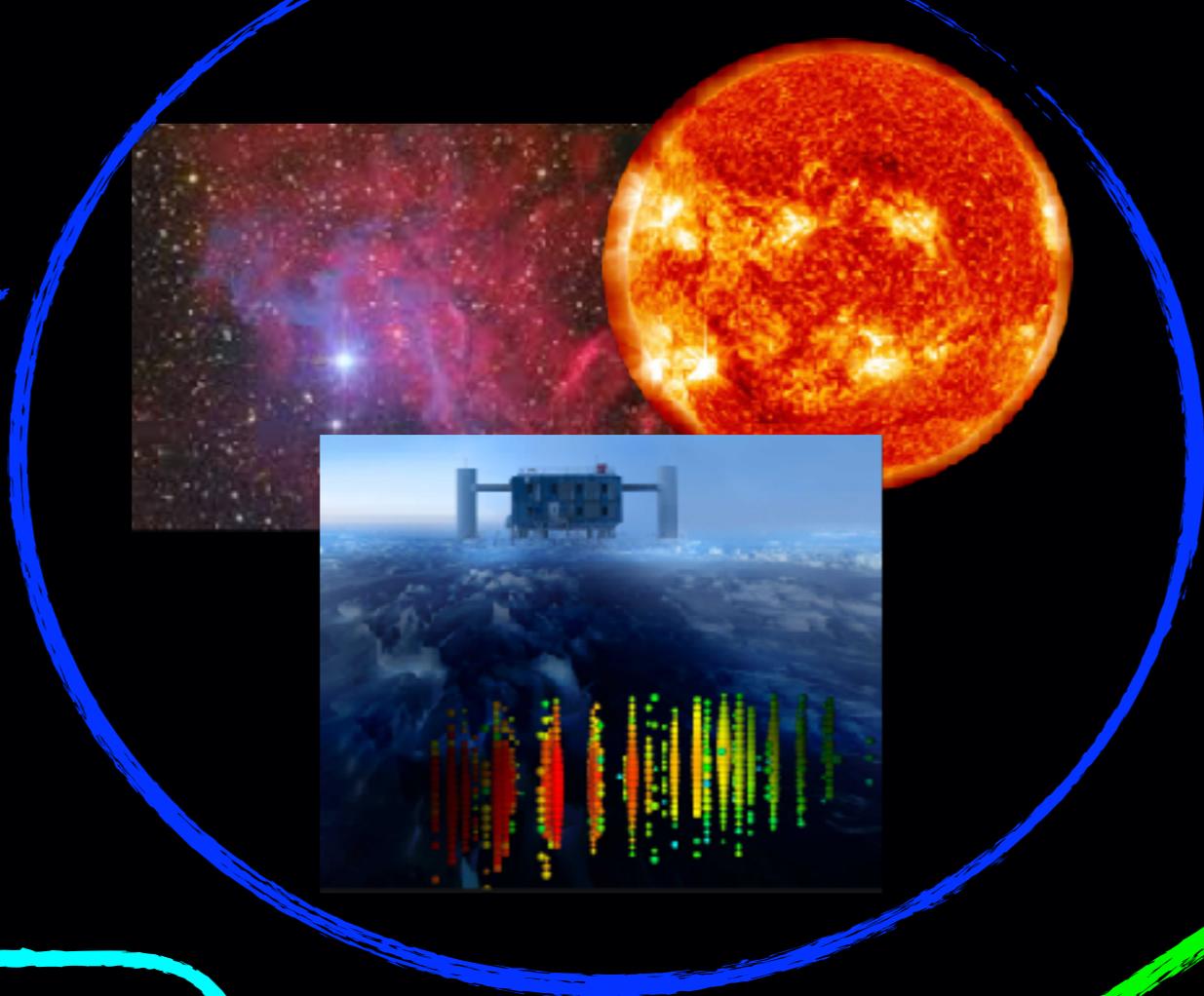
The term “secret neutrino interactions” (ν SI) indicates new physics that couples ν to ν

Several models have been studied (vector, scalar, pseudo-scalar boson) for a large range of the new mediator mass.

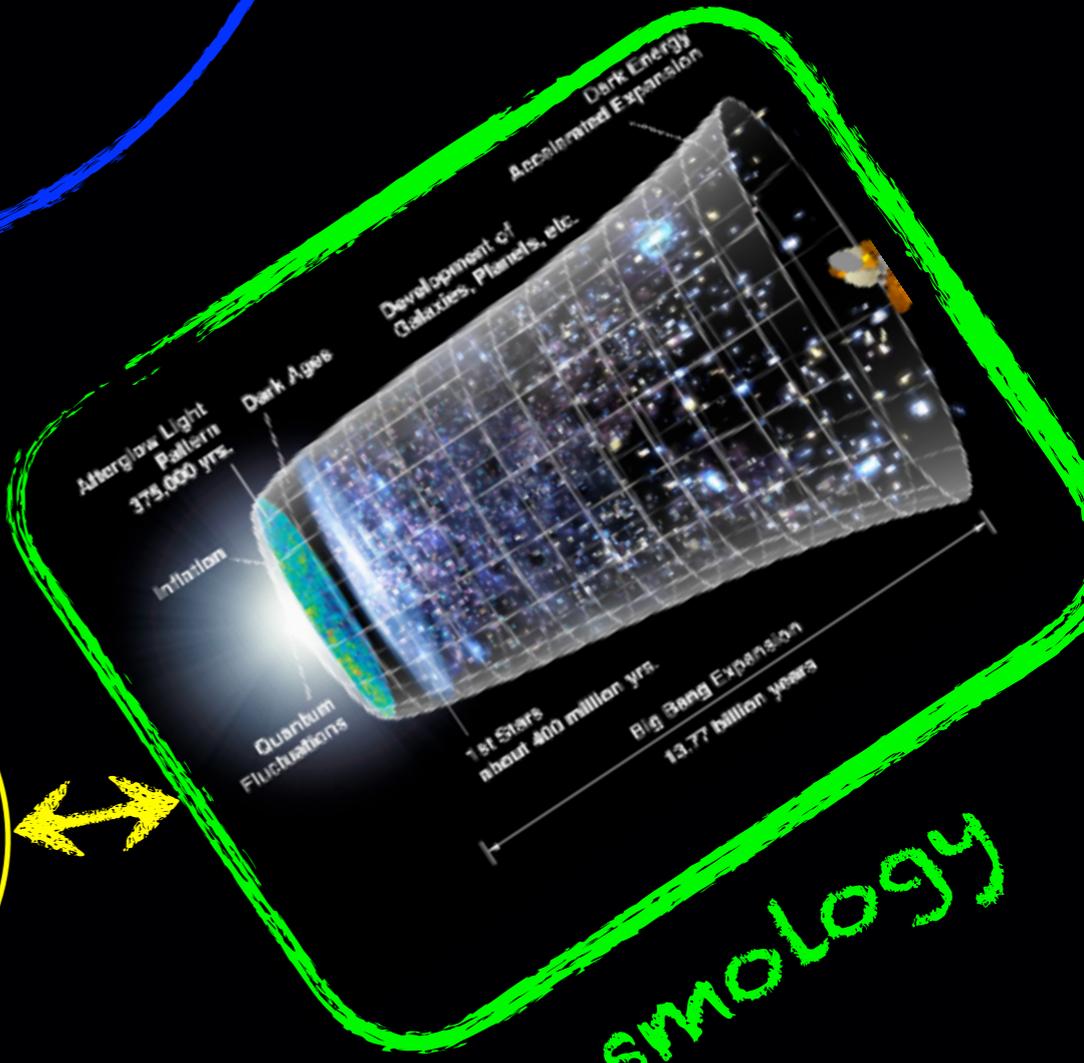
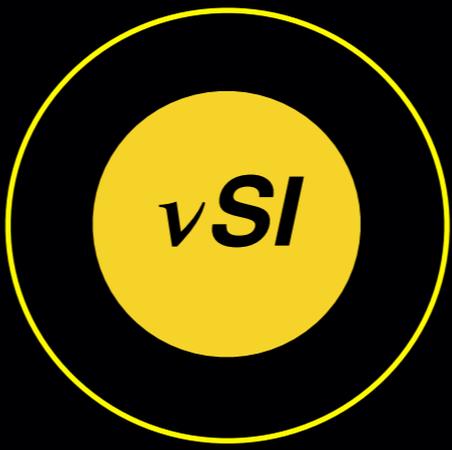
Concise overview on:

- ν SI constraints in the active sector
- ν SI constraints in the sterile sector
- Dark Matter connection

Astrophysical sources



Laboratory

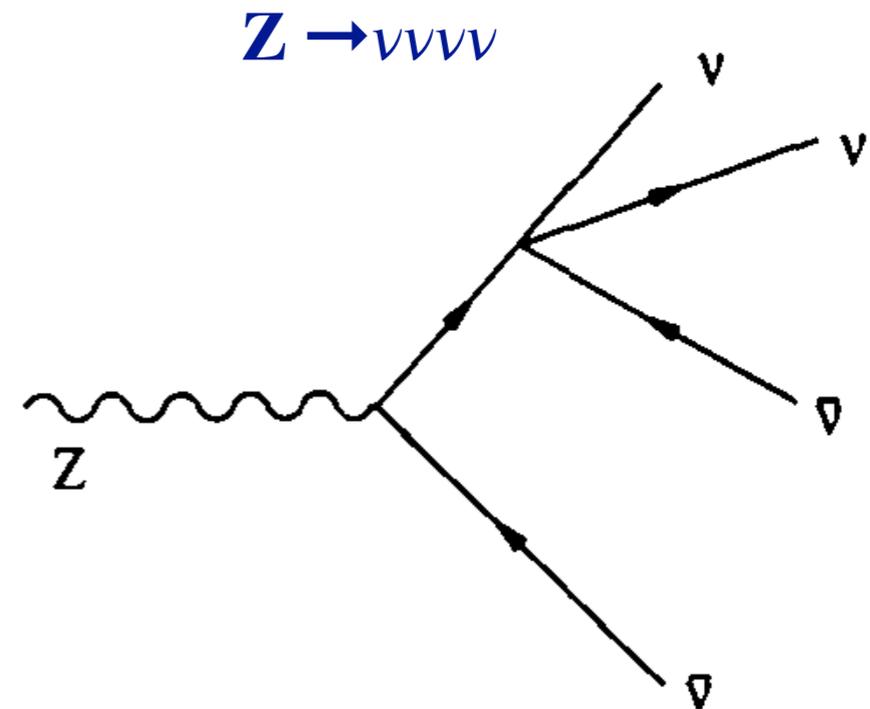
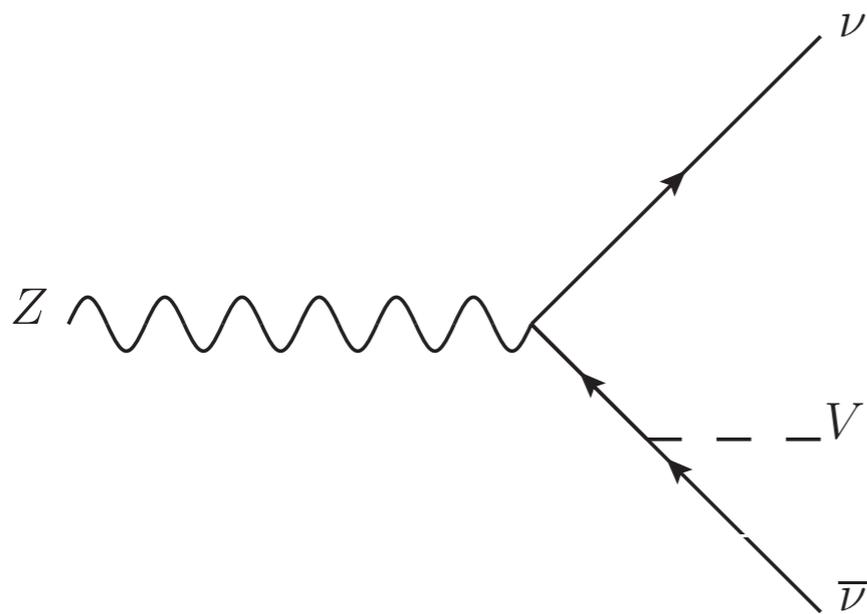


Cosmology

Constraints from laboratory

A light boson V coupling to ν can be constrained by:

**The 3- body decay of the Z increases
the total decay width of the Z**



The decays remain consistent with existing measurements

Laha, Dasgupta and Beacom, 2014

Bilenky & Santamaria, 1992

Constraints from astro-sources

SN

PHYSICAL REVIEW D PARTICLES AND FIELDS

THIRD SERIES, VOLUME 36, NUMBER 10

15 NOVEMBER 1987

Supernova 1987A and the secret interactions of neutrinos

Edward W. Kolb

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and NASA/Fermilab Astrophysics Center, Fermi National Accelerator Laboratory, P.O. Box 500, Batavia, Illinois 60510
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(Received 13 July 1987)

By using SN1987A as a “source” of neutrinos with energy ~ 10 MeV we place limits on the couplings of neutrinos with cosmic background particles. Specifically, we find that the Majoron–electron–neutrino coupling must be less than about 10^{-3} ; if neutrinos couple to a massless vector particle, its dimensionless coupling must be less than about 10^{-3} ; and if neutrinos couple with strength g to a massive boson of mass M , then g/M must be less than 12 MeV^{-1} .

Long list of publication, especially in connection to Majoron models...

... a limited selection:

Kolb and Turner 1987

Fuller, Mayle, Wilson (1988)

Choi and Santamaria, 1990

Kachelrieß, Tomas, Walle 2003

Yasaman Farzan, 2003...

Effects on SN by secret interactions:

- *absorption*: can suppress the spectrum of the observed neutrino flux from a supernova explosion.
- *Energy loss*: reduces the duration of the ν burst
- *Deleptonization*: reduced lepton number of SN core due to Majoron emission

IceCube

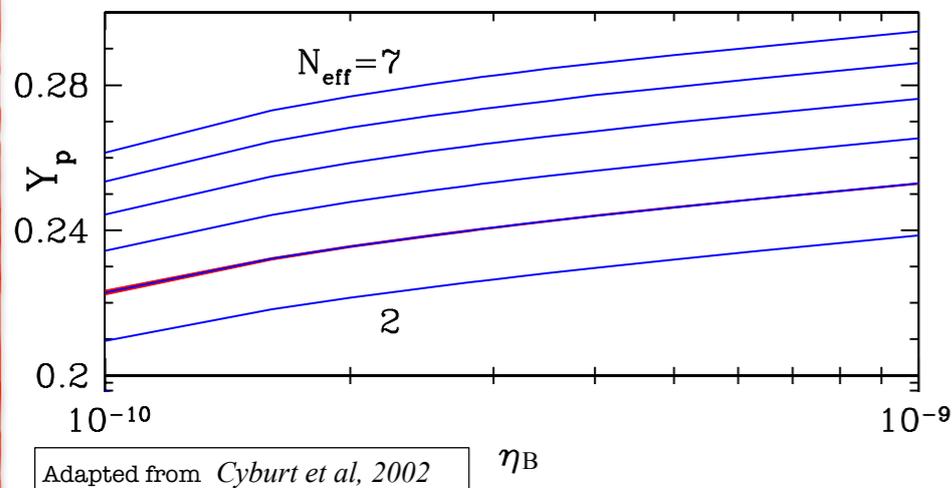
IceCube neutrinos have a much longer baseline through the CvB and much higher energies

Constraints from Cosmology

BBN

sensitive to the radiation density of the Universe, parametrised by N_{eff} and to the spectrum of ν_e (especially in the contest of ν flavour evolution).

Both effects can alter the n/p ratio.



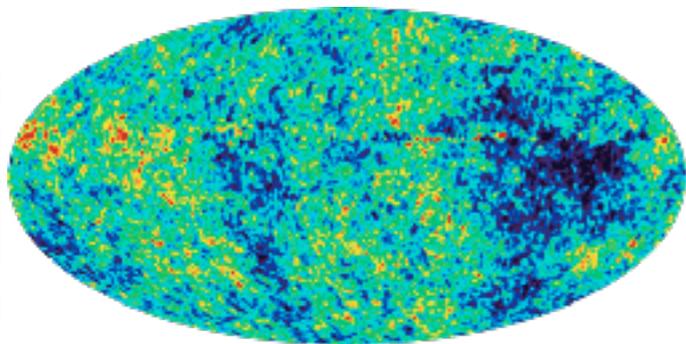
If at the BBN time new particles exist, (i.e new mediators, majorons...), if sufficiently light, they could contribute to N_{eff} , increasing the primordial abundance of light elements.

BBN does not prefer a $N_{\text{eff}} > 3$ ($\Delta N_{\text{eff}} < 1$)

Mangano and Serpico 2011, R. Cooke et al, 2013

CMB

Collisional processes induced by the new interaction would affect the evolution of perturbations in the cosmological neutrino fluid.



Consequent observational signatures on the CMB spectrum

ν SI in the active sector

Lab constraints for Vector Boson

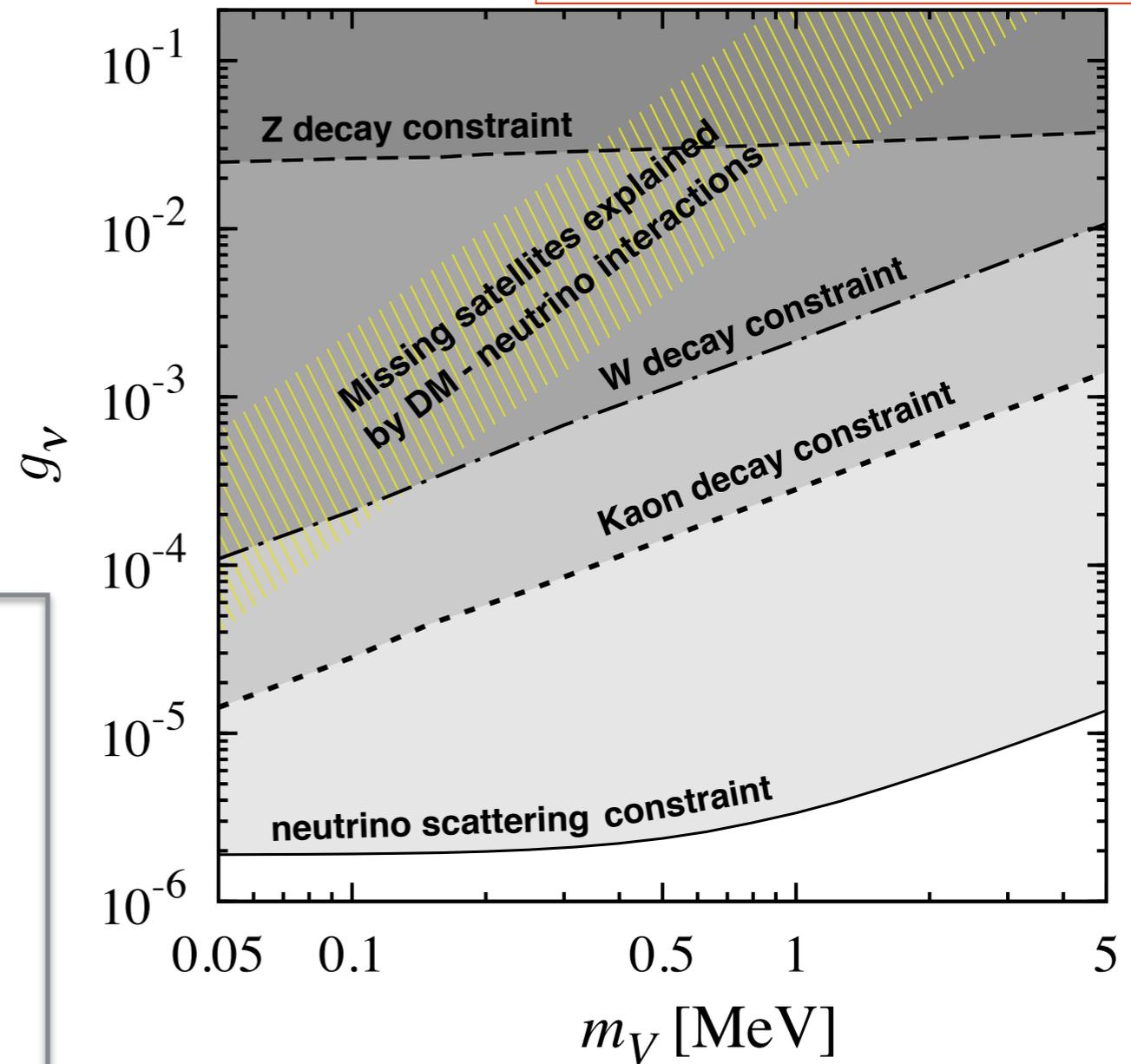
Light vector gauge boson V , with mass \mathcal{O} MeV and couples only to SM neutrinos (and charged leptons through their V-A current): $g_\nu \bar{\nu} \not{V} \nu$

$$\Gamma(Z \rightarrow \nu \bar{\nu} V) \leq 1.28 \times 0.0023 \text{ GeV} \Rightarrow g_\nu \lesssim 0.03$$

Summary of constraints on new interactions of neutrinos with light vector gauge bosons at $m_V = 1$ MeV

Process	Interaction	Constraint
$Z \rightarrow \nu \bar{\nu} V$	$g_\nu \bar{\nu} \not{V} \nu$	$g_\nu \lesssim 3 \times 10^{-2}$
$W \rightarrow \tau^- \bar{\nu}_\tau V$	$g_\nu (\bar{\nu} \not{V} \nu + \bar{\ell} \not{V} \ell)$	$g_\nu \lesssim 2 \times 10^{-3}$
$K^- \rightarrow \mu^- \bar{\nu}_\mu V$	$g_\nu (\bar{\nu} \not{V} \nu + \bar{\ell} \not{V} \ell)$	$g_\nu \lesssim 3 \times 10^{-4}$
$\nu e \rightarrow \nu e$	$g_\nu (\bar{\nu} \not{V} \nu + \bar{\ell} \not{V} \ell)$	$g_\nu \lesssim 3 \times 10^{-6}$

Laha, Dasgupta and Beacom, 2014



BBN & CMB constraints for Vector Boson

The constraints on g_V and m_V from **BBN** are very restrictive.

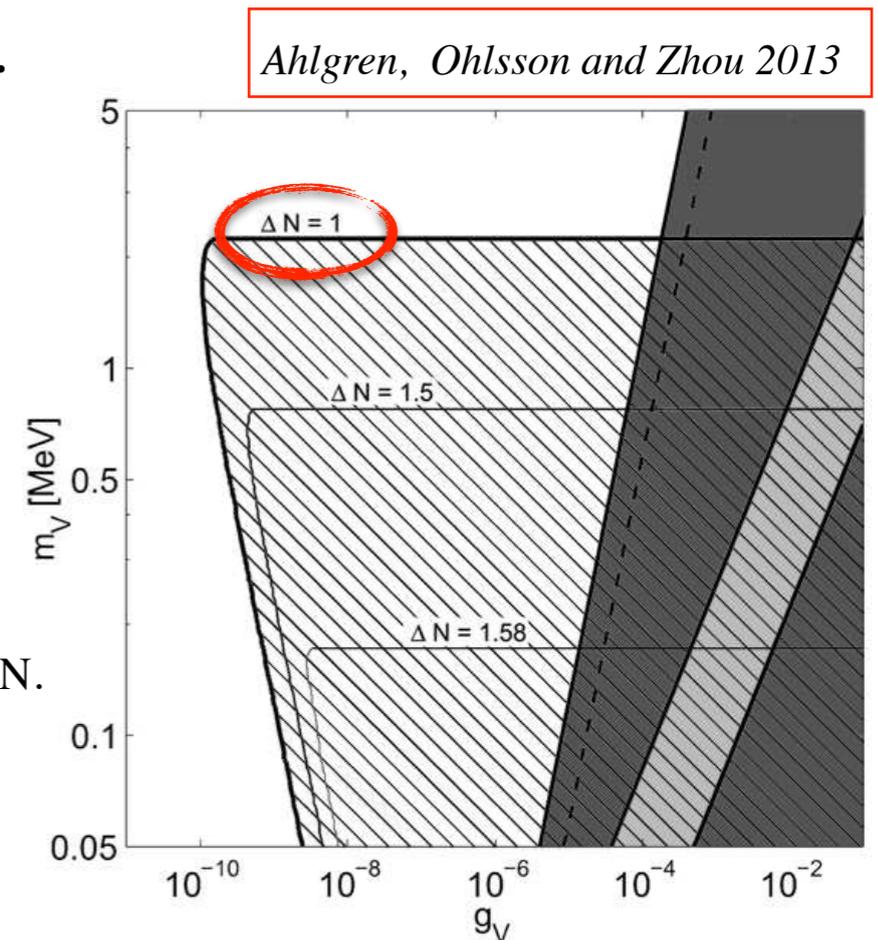
In the early Universe, V can be thermalized via

- the inverse decay $\nu + \nu \rightarrow V$
- pair annihilation $\nu + \nu \rightarrow V + V$

and contribute to the energy density

Constraints on g_V and m_V from K decays (gray, solid line), W decays (gray, dashed line), and BBN.

The hashed region bounded by the thick solid curve is excluded by $\Delta N_V < 1$ at 95% C.L.



CMB

4-fermion point-like interaction model: $G_X = g^2 / m_X^2$

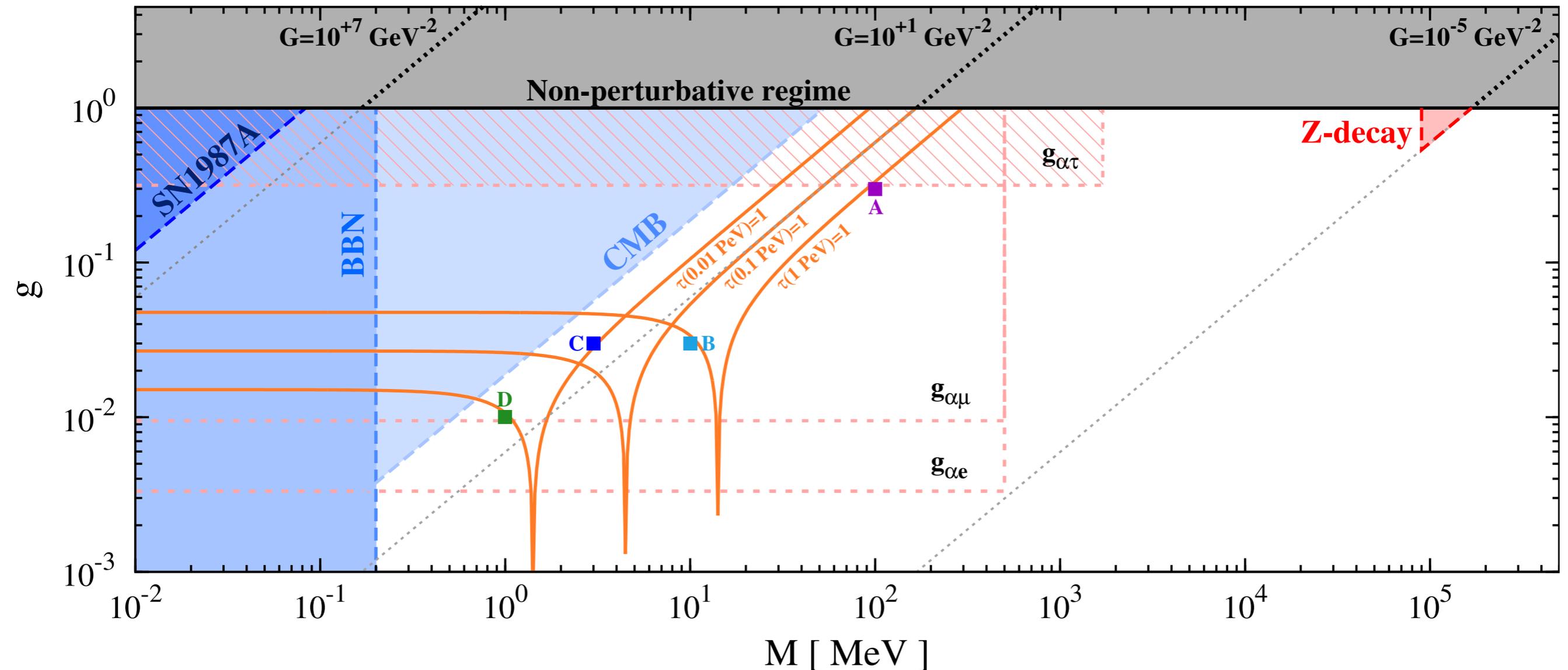
Using Planck data (combined with other data): $G_X \leq 2.5 \times 10^7 G_F$

Archidiacono & Hannestad 2014

Constraints for Scalar Boson

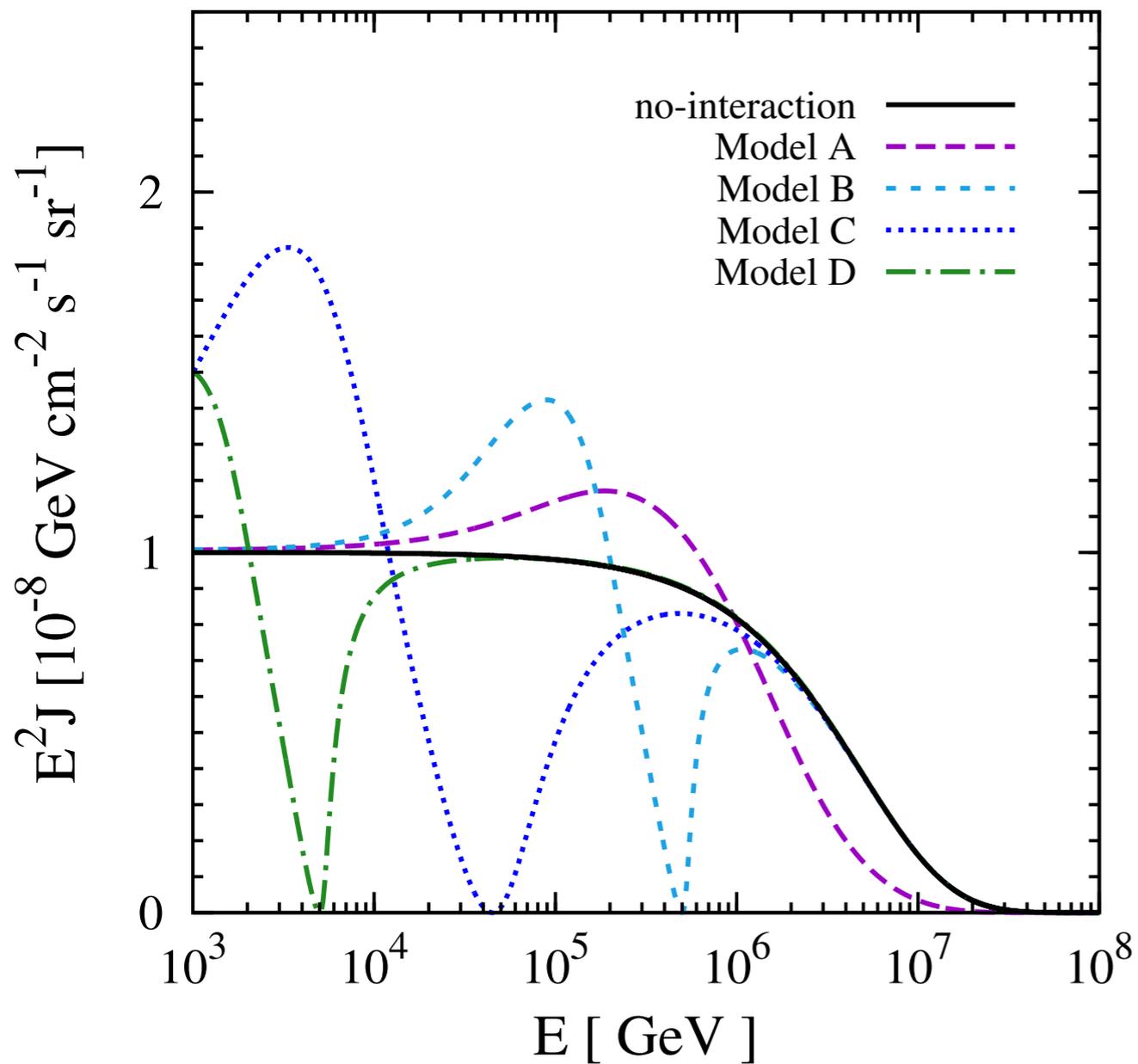
Scalar interaction term: $\mathcal{L}_{int} \sim g\phi\bar{\nu}\nu$

Ng and Beacom, 2014



Blue shaded regions are excluded by astrophysical and cosmological considerations based on SN 1987A [K&T], BBN [A,O,Z], and the CMB [A&H]. Pink dashed lines indicate flavor-dependent limits based on laboratory measurements of meson and lepton decays. Shaded pink: only the weakest limit for $\nu\tau$. Red shaded region is excluded based on measurement of Z-boson decay [B&S]. Grey shaded region indicates the non-perturbative regime. Orange lines are contours of unit optical depth for different initial neutrino energies, indicating the **approximate boundary of the parameter space above which IceCube is sensitive to ν SI**.

"scalar" ν SI impact on IceCube spectra



$$\mathcal{L}_0(E) \propto E^{-\gamma} e^{-E/E_{\text{cut}}}$$

$$\gamma = 2 \text{ and } E_{\text{cut}} = 10^7 \text{ GeV}$$

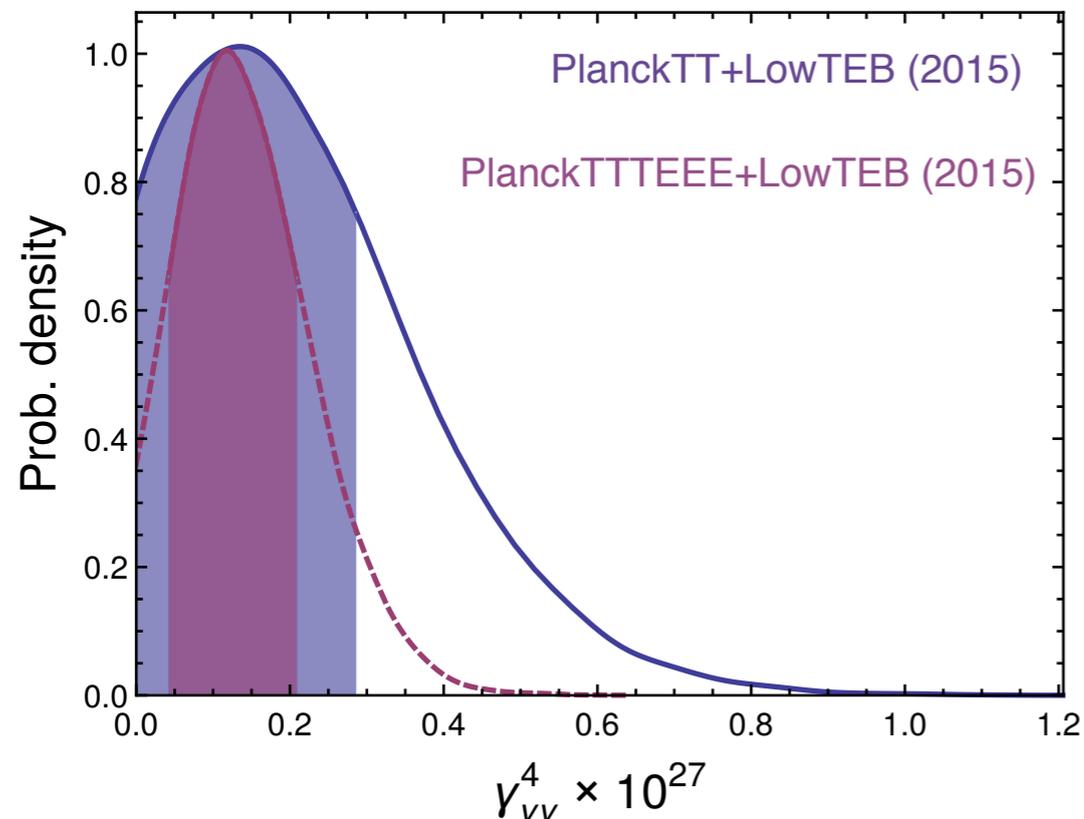
	A	B	C	D
g	0.3	0.03	0.03	0.01
M [MeV]	100	10	3	1
τ (1 PeV)	~ 0.7	~ 0.6	~ 0.2	~ 0.002
E_{res} [PeV]	50	0.5	0.045	0.005

CMB constraints for light (pseudo) scalar Boson

$$\mathcal{L} = h_{ij} \bar{\nu}_i \nu_j \phi + g_{ij} \bar{\nu}_i \gamma_5 \nu_j \phi + h.c.$$

Scalar and pseudoscalar neutrino couplings can be relevant for CMB, since collisional processes induced by the new interaction would affect the evolution of perturbations in the cosmological neutrino fluid.

For binary process mediated by a massless boson, $\sigma_{\text{bin}} \sim g^4/s \longrightarrow \Gamma_{\text{bin}} = \langle \sigma_{\text{bin}} v \rangle n_{\text{eq}} \propto g^4 T$



Constraint in the Majoron case: $g \lesssim 7.4 \times 10^{-7}$

$$g^4 = \gamma_{\nu\nu}^4 / (1.8 \times 10^{-3})$$

Forastieri, Lattanzi and Natoli 2015

ν SI in the (light)
sterile sector

eV sterile neutrinos

The investigation on Light Sterile Neutrinos has been stimulated by the presence of anomalous results from neutrino oscillation experiments

Interpretation: *1* (or more) *sterile neutrino* with $\Delta m^2 \sim O(\text{eV}^2)$ and $\theta_s \sim O(\theta_{13})$

Giunti et al 2013; J. Kopp et al, 2013

Giunti's Talk

ν_s can be produced in the *EU* by the mixing with the active species in presence of collisions

Are eV ν_s compatible with cosmology?

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Are eV ν_s compatible with cosmology? **NO** Too "many" and too heavy

Possible solutions:

- *suppress the ν_s abundance*

Feasible with lepton asymmetry, low-reheating scenario, secret interactions

Di Bari et al, 2001, Cirelli et al., 2004, Gelmini, Palomarez-Ruiz, Pascoli, 2004, Chu & Cirelli, 2006, Mirizzi, N.S., Miele, Serpico 2012, Hannestad, Tamborra and Tram 2012, Saviano et al., 2013,

- *modification of cosmological model*

Gariazzo, Giunti, Laveder, 2015

SI for sterile neutrinos

Different authors have assumed the Standard Model (SM) is augmented by one extra species of light ($\sim eV$) neutrinos ν_s , which do not couple to the SM gauge bosons but experiment a new force

Hannestad et al., 2013., Dasgupta and Kopp 2013, Bringmann et al., 2014

Such a new interaction can have profound effects on active-sterile neutrino conversion in the early Universe, **since sterile ν feel a new potential that can suppresses active-sterile mixing** (*through an effective ν_a - ν_s mixing reduced by a large matter term*)

Caveat: they also generate MSW resonance and strong collisional production, increasing their abundance, with non trivial consequences on the cosmological observables



→ ν SI constraints from cosmological probes

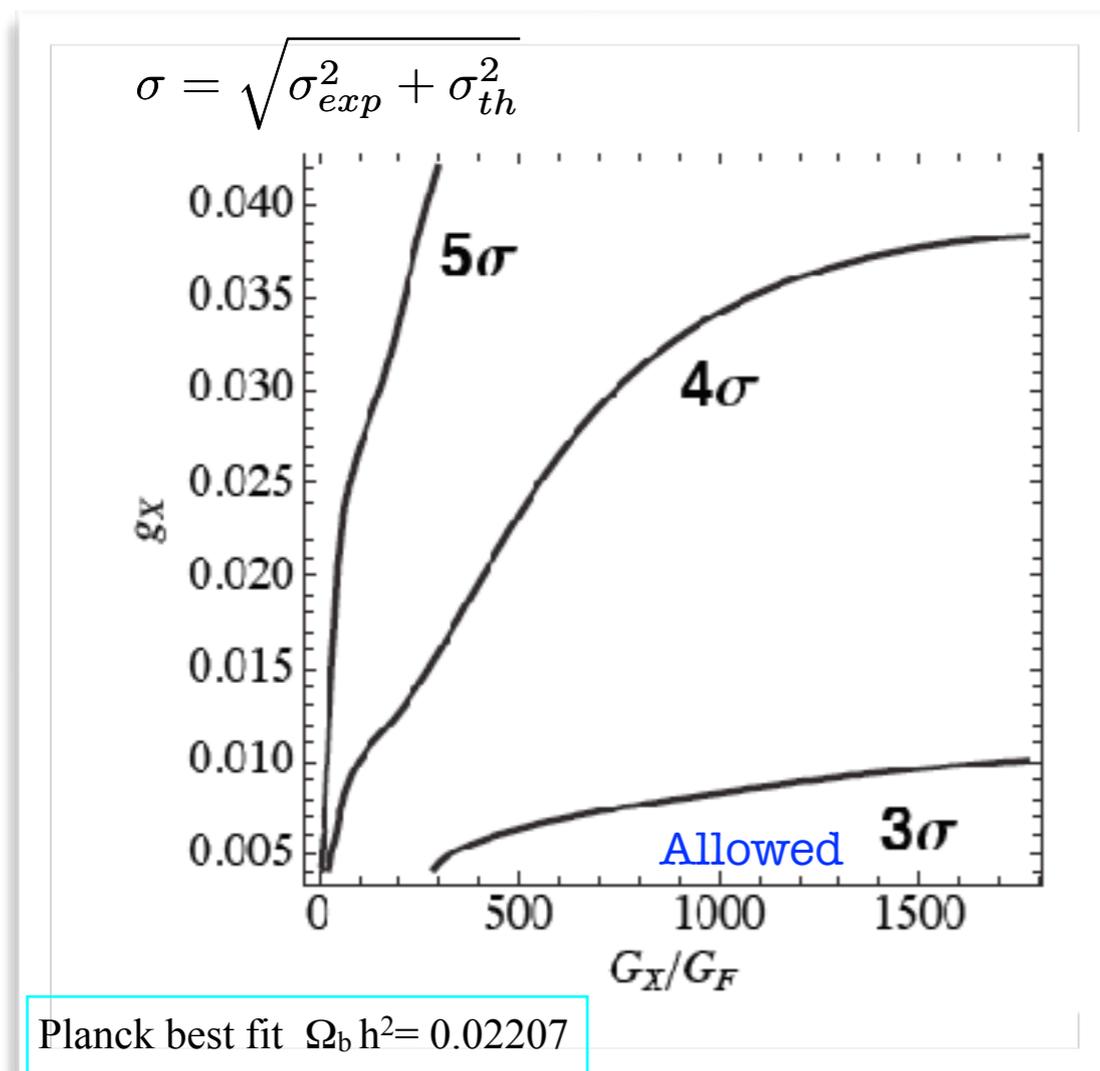
If the new mediator interaction X also couples to Dark Matter possible attenuation of some of the small scale structure problems (“missing satellites” problem...)

BBN constraints for sterile ν SI (vector)

Particular model: **4-fermion point-like interaction** $G_X = \frac{\sqrt{2}}{8} \frac{g_X^2}{M_X^2}$ for $T \ll M_X$

After the ν oscillation in the range of g_X and G_X relevant for BBN, we have both :

$\Delta N_{\text{eff}} > 0$ and distortions of the active ν_e spectra



Deuterium yield

Experimental reference value:

$${}^2\text{H}/\text{H} = (2.53 \pm 0.04) \times 10^{-5}$$

R. Cooke et al, 2013

σ_{exp}

Translating in a bound for the mediator mass:

mass permitted: $M_X \leq 40 \text{ MeV}$

Saviano, Pisanti, Mangano, Mirizzi 2014

Mass constraints for sterile ν SI (vector)

If we want to constraint lower $M_X \leftrightarrow$ very large $G_X (> 10^5 G_F)$

Very strong secret collisional term leads to a quick flavor equilibrium

Stodolsky, 1987

$$\begin{aligned} (\rho_{ee}, \rho_{\mu\mu}, \rho_{\tau\tau}, \rho_{ss})_{\text{initial}} &\longrightarrow (\rho_{ee}, \rho_{\mu\mu}, \rho_{\tau\tau}, \rho_{ss})_{\text{final}} \\ (1, 1, 1, 0) &\quad (3/4, 3/4, 3/4, 3/4) \end{aligned}$$



lower value in the 2σ range from anomalies gives $m_s^{\text{eff}} \sim 0.8 \text{ eV}$

in tension with the CMB and LSS bounds on sterile mass ($< 0.5 \text{ eV}$)

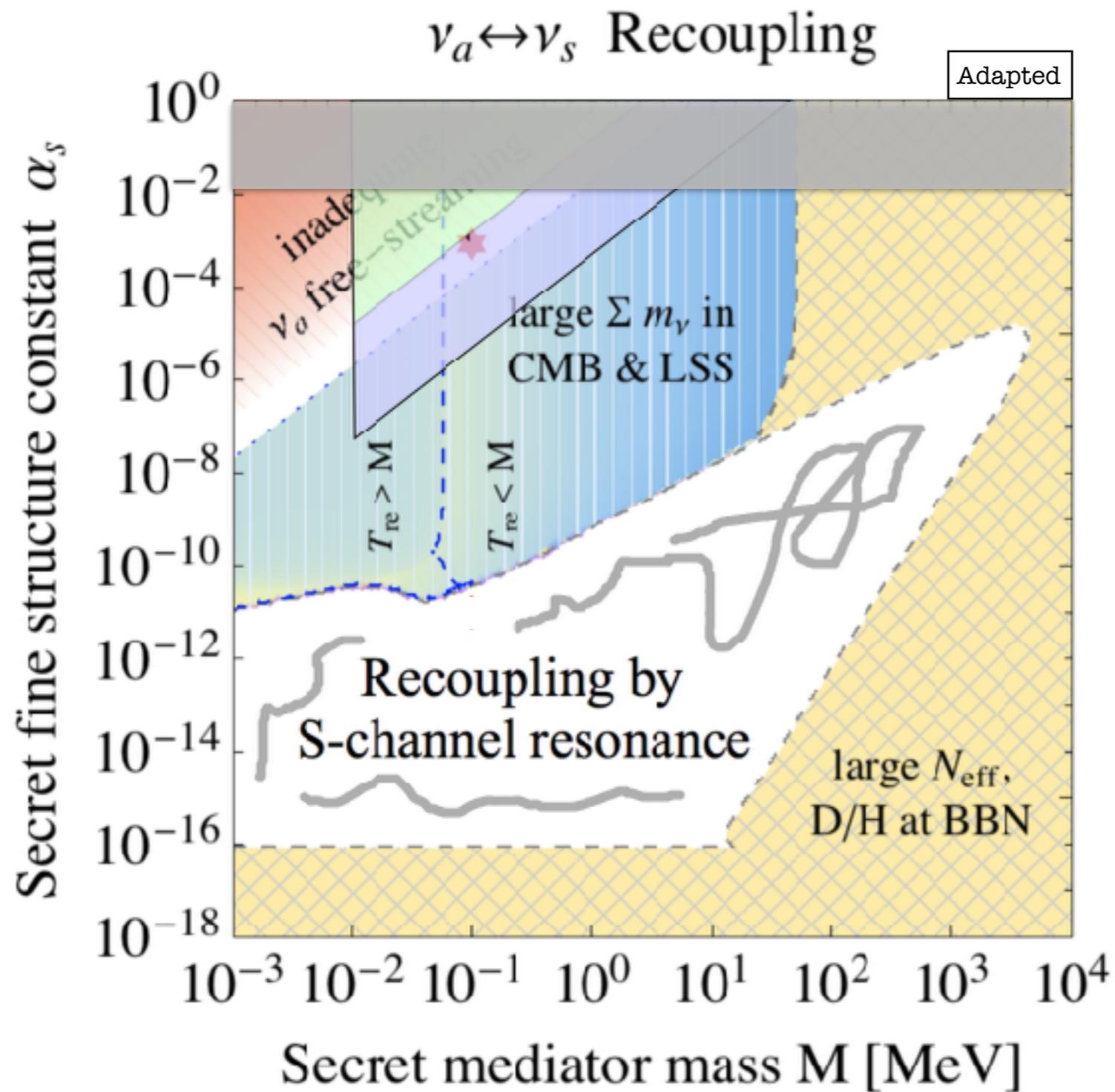
Mirizzi, Mangano, Pisanti Saviano, 2014

Secret interaction scenario: disfavored $M_X > 0.1 \text{ MeV}$ ($\Leftrightarrow \sim 10^9 G_F$)

For $M_X \leq 0.1 \text{ MeV}$ ($\geq 10^{10} G_F$) $\rightarrow \nu_s$ could be still coupled at CMB and LSS epoch \rightarrow no free-streaming...

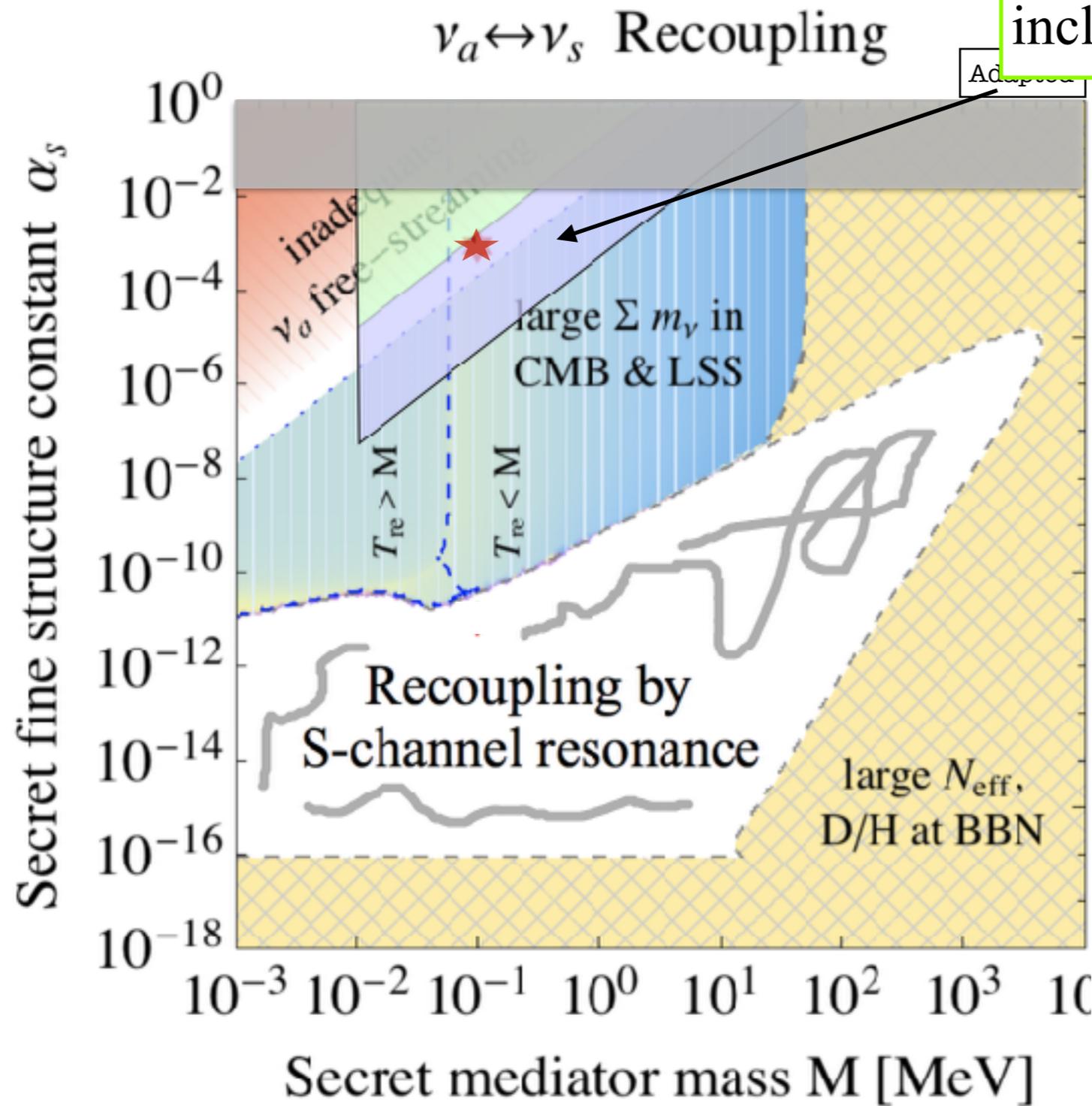
.... an appropriated analysis should be performed

Summary Plot

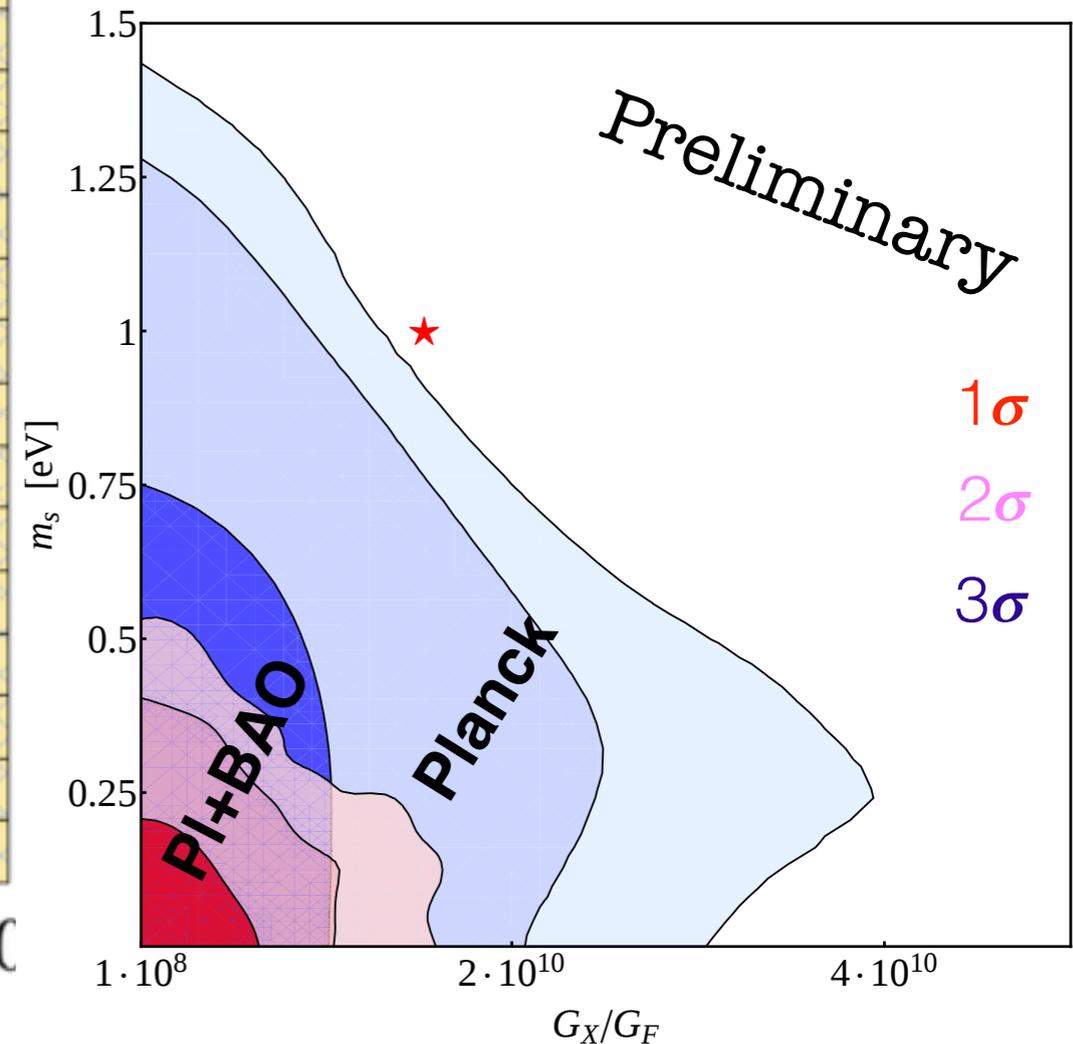


Summary Plot

Preliminary results of CMB analysis including no-free-streaming sterile ν



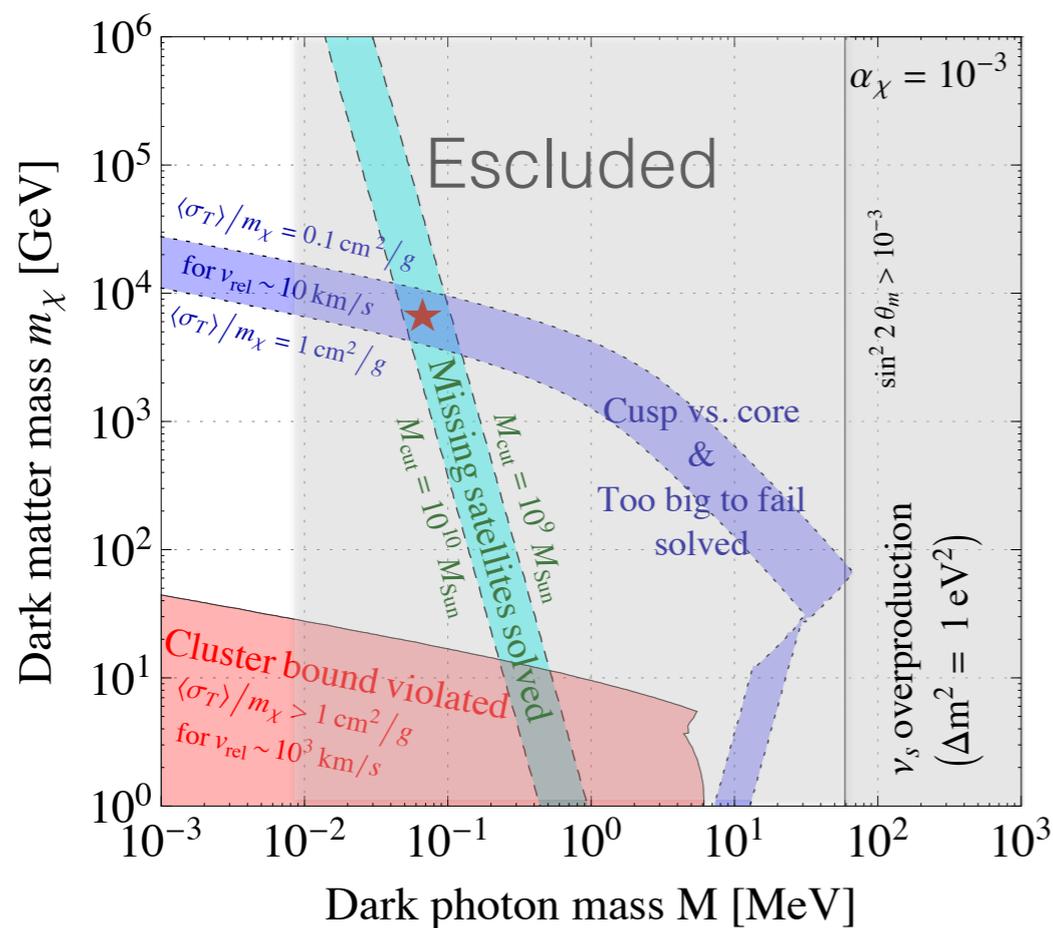
Forastieri, Lattanzi, Mangano, Mirizzi, Natoli, Saviano, 2017
to appear soon



Connection with the DM

If a new force exists, it is plausible that not only (sterile) neutrinos, but also DM particles couple to it

"*neutrinophilic DM*" $\mathcal{L}_{\text{int}} \supset -g_\chi \bar{\chi} \not{V} \chi - g_\nu \bar{\nu} \not{V} \nu$

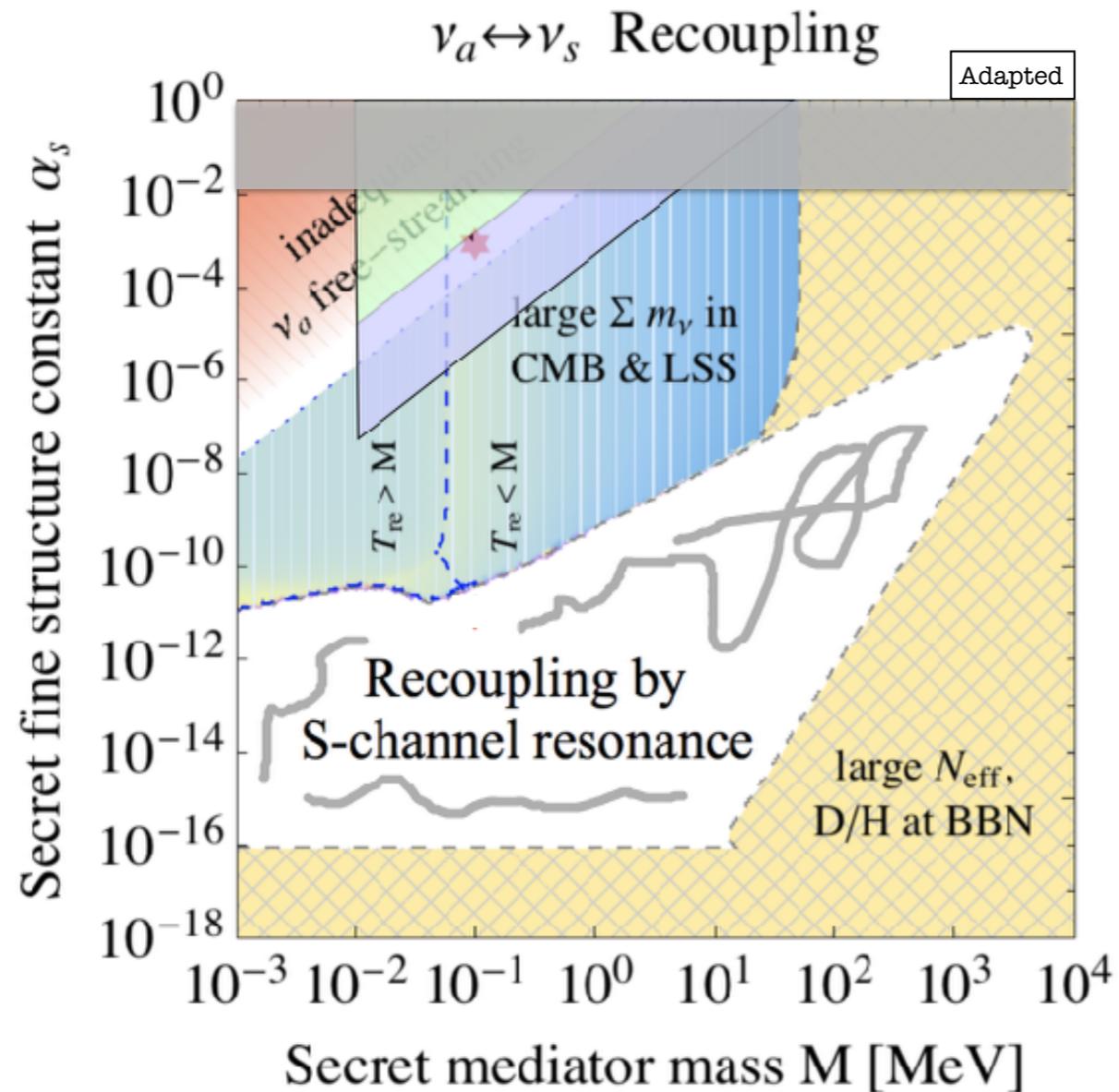


This scenario may solve all of the small-scale structure issues mentioned above. Indeed, the efficient scattering of DM would lead to late kinetic decoupling, delaying the formation of the smallest protohalos.

(Barions?)

van den Aarssen, Bringmann and Pfrommer 2012, Dasgupta and Kopp 2013, Bringmann, Hasenkamp, Kersten 2014, Cherry, Friedland, Shoemaker 2014, Archidiacono et al. 2015....

Summary Plot



Chu, Dasgupta and Kopp 2015

Similar plot obtained by *Cherry, Friedland, Shoemaker 2016* using present data and future sensitivity of IceCube

Possible hint (very dependent on the set of data used): in pseudoscalar model, $10^{-6} \lesssim g_s \lesssim 10^{-5}$ would reconcile eV sterile ν , H_0 , ν SI. Also link to the DM small scale problem.

Archidiacono et al. 2015

Open questions

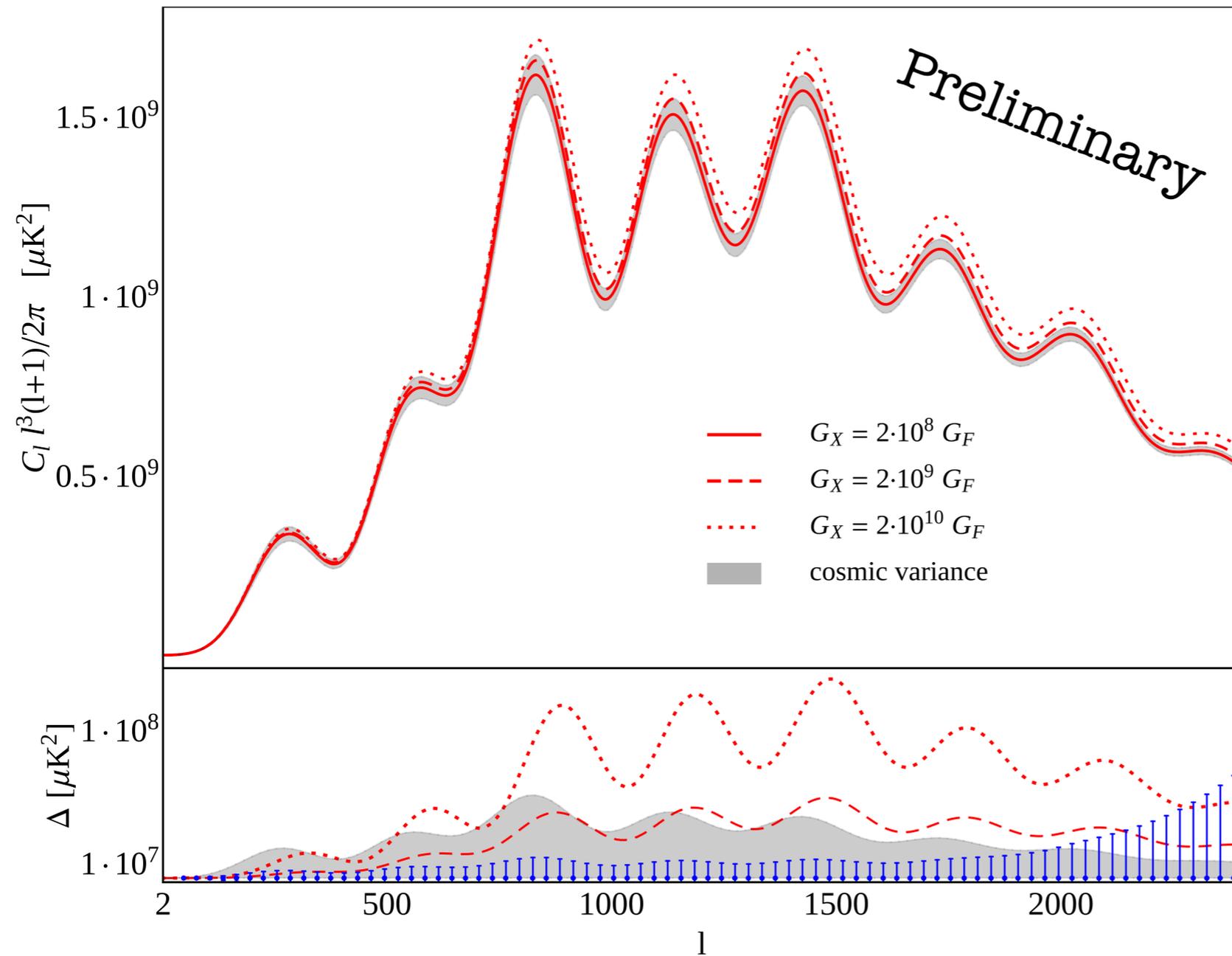
What happens considering sterile ν and ν SI in the SN?

Secret interactions can be productive for other sterile ν mass scales?

How many extra exotic ingredients we are willing to introduce in our theory? (sterile ν , dark matter, secret interactions...)

Thank you

Backup slides



*Forastieri, Lattanzi, Mangano, Mirizzi,
Natoli, Saviano, 2017
to appear soon*

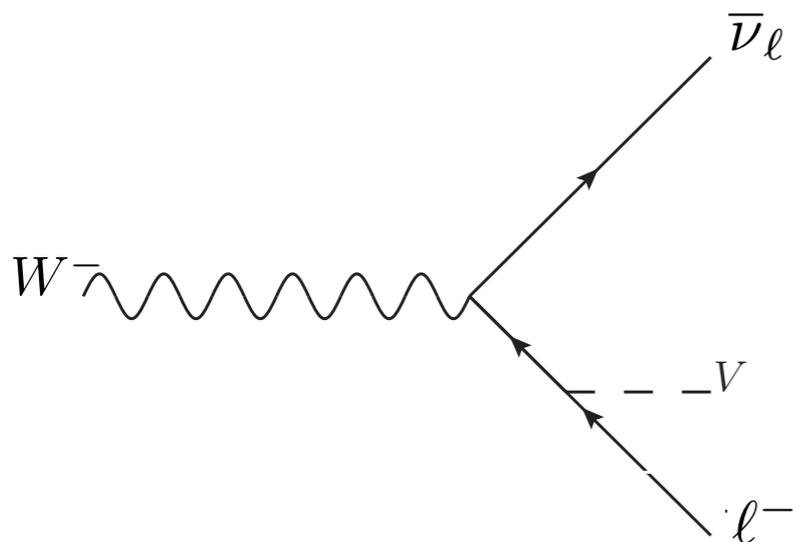
As long as the collision rate is much larger than the expansion rate, interacting neutrinos behave as perfect fluid. This means that shear and higher moments are exponentially suppressed. The net effect is that, at scales that are within the horizon during the interacting regime, density and pressure perturbations are enhanced with respect to the non-interacting case.

This enhancement propagates to the photon fluid, and thus to CMB anisotropies, through the metric perturbations.

Constraints for Vector Boson

Light vector gauge boson V , with mass \mathcal{O} MeV and couples only to SM neutrinos and charged leptons through their V-A current:

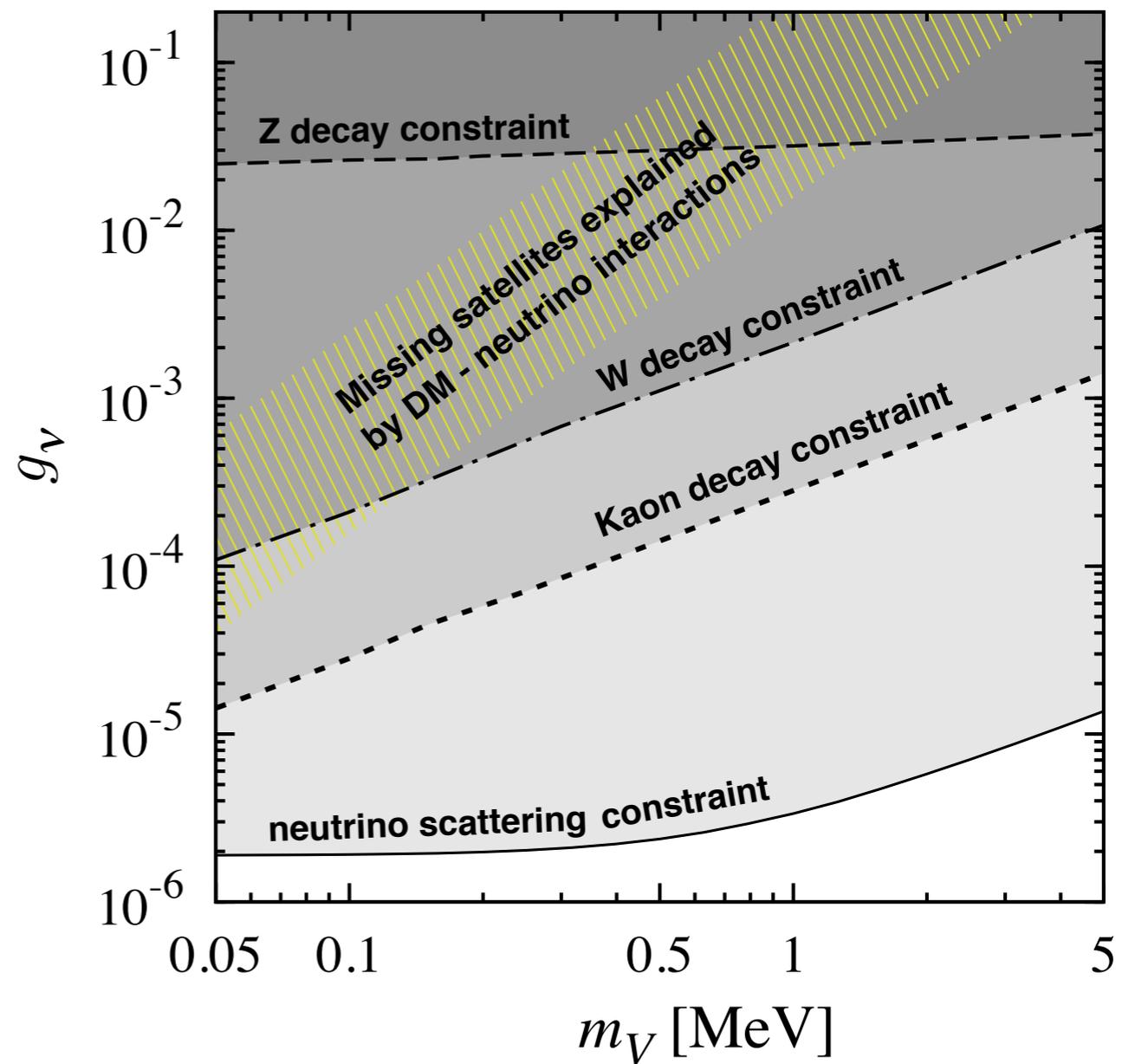
$$-g_\nu (\bar{\ell} \not{V} P_L \ell - \bar{\nu} \not{V} P_L \nu)$$



The 3-body decay of the W leads to additional events with missing energy, increasing the total decay width of the W

The decay remains consistent with existing measurements

$$\Gamma(W^- \rightarrow \ell^- \bar{\nu}_\ell V) \leq 1.28 \times 0.042 \text{ GeV}$$

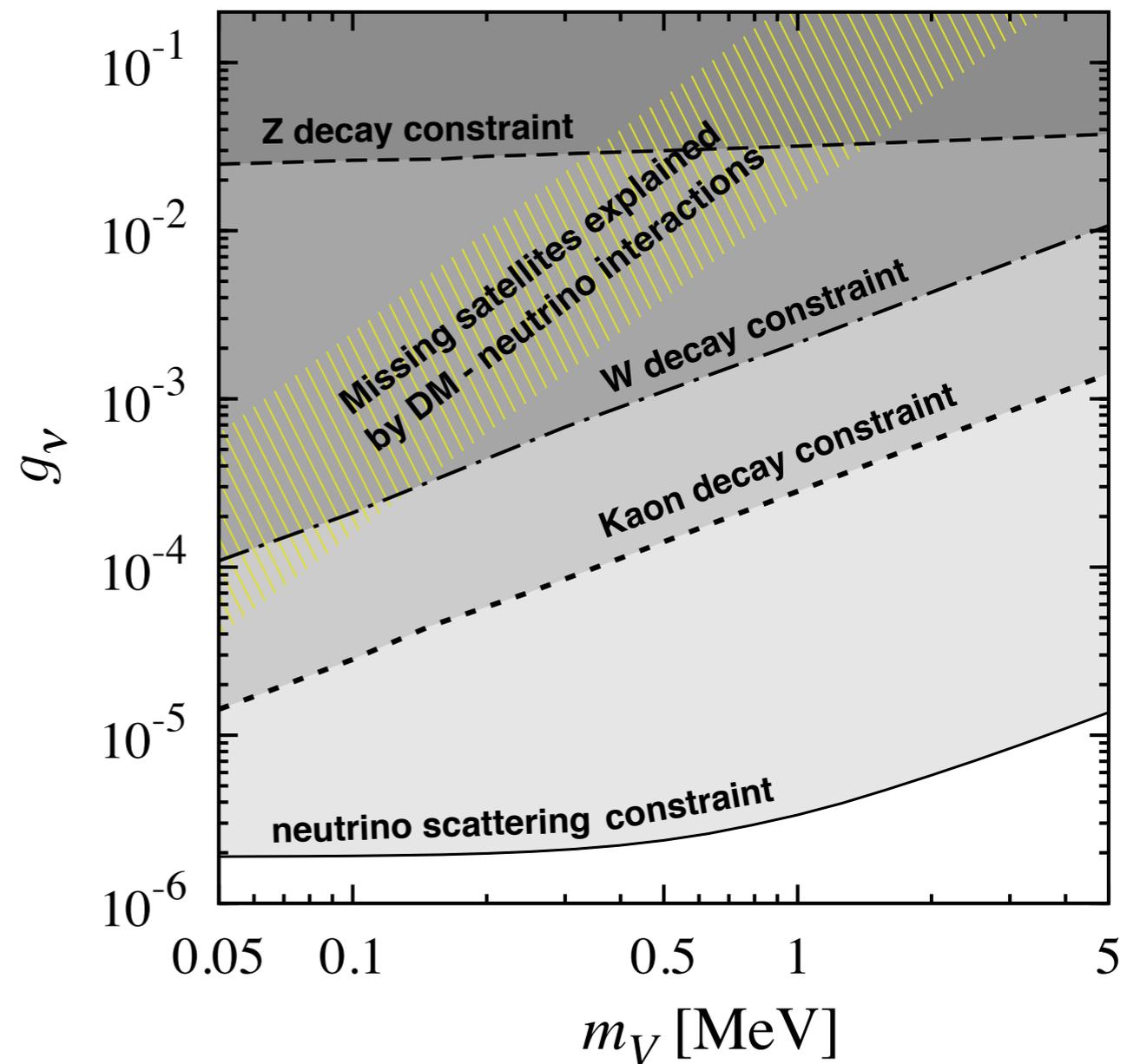
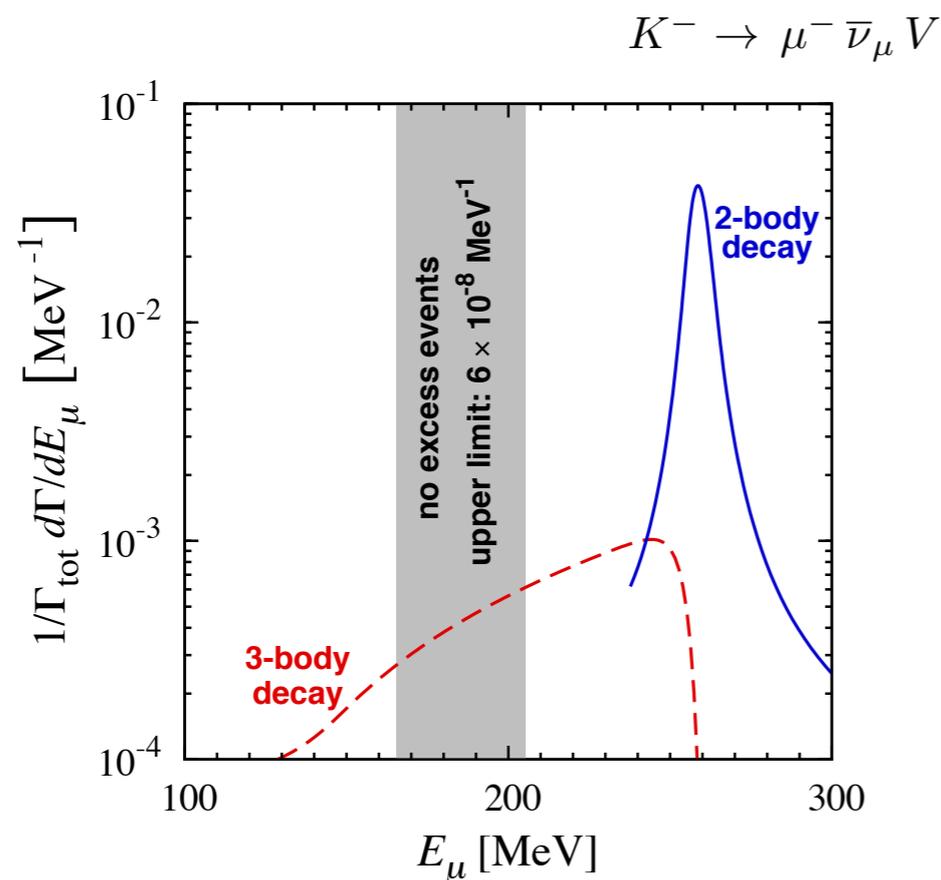


Constraints for Vector Boson

Light vector gauge boson V , with mass \mathcal{O} MeV and couples only to SM neutrinos and charged leptons through their V-A current:

$$-g_\nu (\bar{\ell} \not{V} P_L \ell - \bar{\nu} \not{V} P_L \nu)$$

Distortion of the charged lepton spectrum due to excess missing energy in kaon decays



$$\Gamma(K^- \rightarrow \mu^- + \text{inv.}) / \Gamma(K^- \rightarrow \mu^- \bar{\nu}_\mu) \leq 3.5 \times 10^{-6}$$

Constraints for Vector Boson

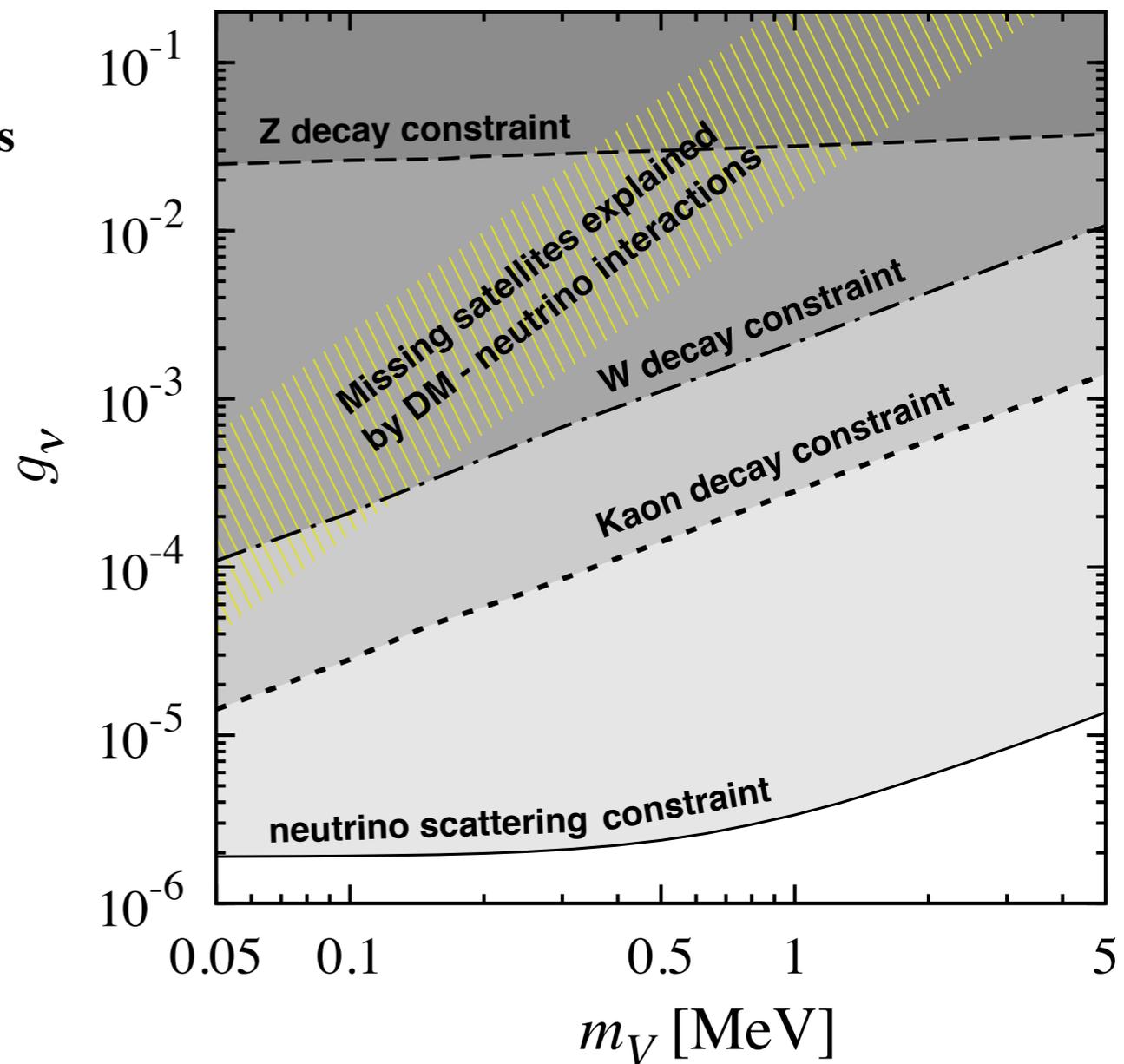
Light vector gauge boson V , with mass O MeV and couples only to SM neutrinos and charged leptons through their V-A current:

$$-g_\nu (\bar{\ell} \not{V} P_L \ell - \bar{\nu} \not{V} P_L \nu)$$

**neutrino-electron scattering at very low neutrino energies
(e.g., as in solar neutrino detection)**

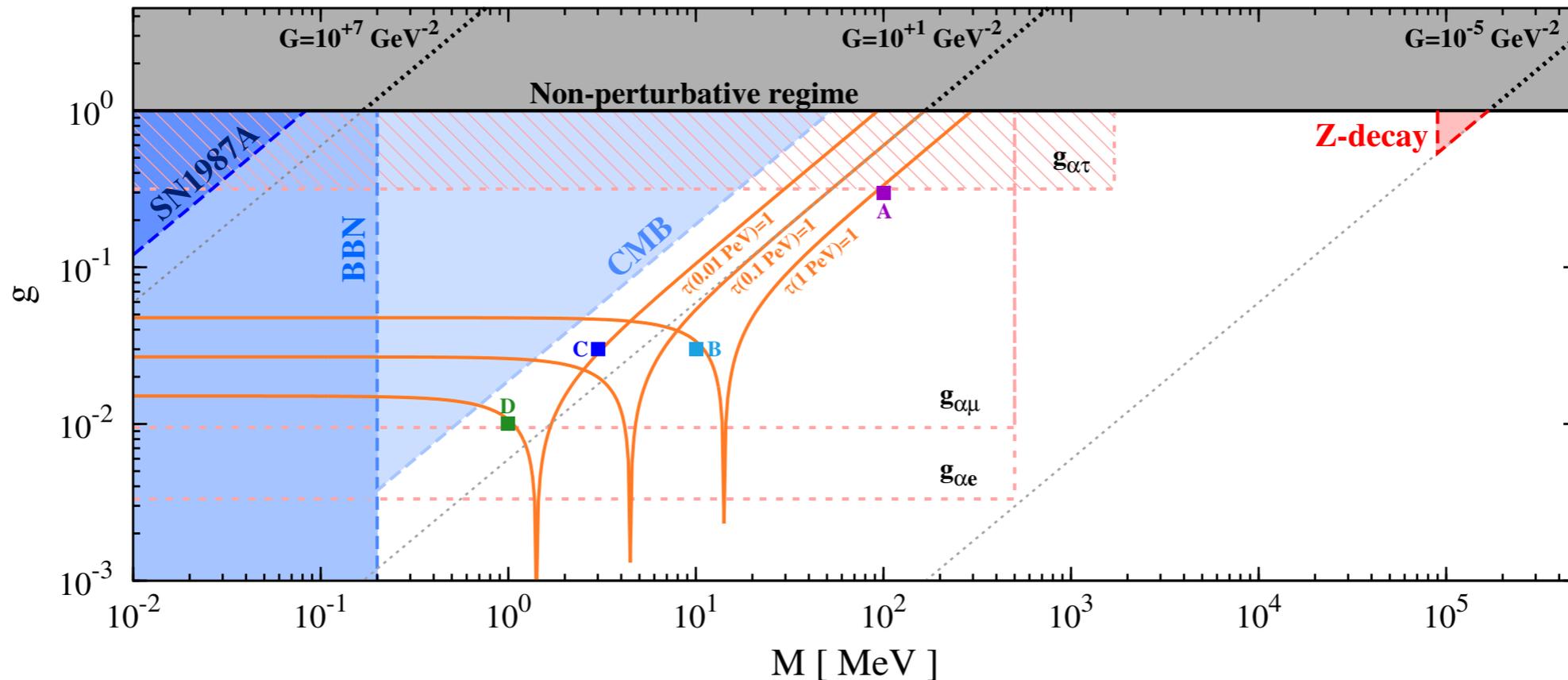
Solar neutrinos, produced as ν_e , change to ν_μ or ν_τ

The presence of the new vector boson would alter the charged current (CC) interaction between solar ν_e neutrinos and target electrons in the detector. It would also alter the ν_μ or ν_τ interaction with electrons via the weak neutral current (NC) interaction



Constraints for Scalar Boson

Scalar interaction term: $\mathcal{L}_{int} \sim g\phi\bar{\nu}\nu$



ν SI: elastic scattering of neutrinos in the s channel \rightarrow

$$\sigma(E) = \frac{g^4}{4\pi} \frac{s}{(s - M^2)^2 + M^2\Gamma^2}$$

$$s = 2Em_\nu$$

$$\Gamma = g^2 M / 4\pi$$

3 kinematic regimes:

- **mediator mass $< \sqrt{s}$** (i.e. Majoron with $m=0$), $\sigma \simeq g^4 / (4\pi s)$
- **mediator mass $\sim \sqrt{s}$** (resonance), $s = M^2$, $\sigma = 4\pi / M^2$
- **mediator mass $> \sqrt{s}$** (effective theory), $\sigma \simeq g^4 s / (4\pi M^4)$

	A	B	C	D
g	0.3	0.03	0.03	0.01
M [MeV]	100	10	3	1
τ (1 PeV)	~ 0.7	~ 0.6	~ 0.2	~ 0.002
E_{res} [PeV]	50	0.5	0.045	0.005

$$\tau(E|g, M) = c \int_0^1 dz \frac{n_t(0)}{H(z)} (1+z)^2 \sigma(E|g, M)$$

Secret interactions for sterile ν_s

4-fermion point-like interaction:
 new secret self-interactions among
 sterile ν mediated by a massive
 gauge boson X : $M_X \ll M_W$



*Suppress the thermalization of
 sterile neutrinos*
 (Effective ν_a - ν_s mixing reduced by a large
 matter term)

$$\nu_s - \nu_s \text{ interaction strength } G_X = \frac{\sqrt{2}}{8} \frac{g_X^2}{M_X^2} \quad \text{for } T \ll M_X$$

Evolution equation:

$$i \frac{d\rho}{dt} = [\Omega, \rho] + C[\rho]$$

$$\Omega = \Omega_{\text{vac}} + \Omega_{\text{mat}} + \Omega_{\nu-\nu} + \Omega_{\nu_s-\nu_s}^{\text{secr}}$$

$\swarrow \quad \searrow \quad \quad \quad \searrow$
 $\propto G_F \quad \quad \quad \propto G_X$

$$C = C_{\text{SM}} + C_{\text{secr}}$$

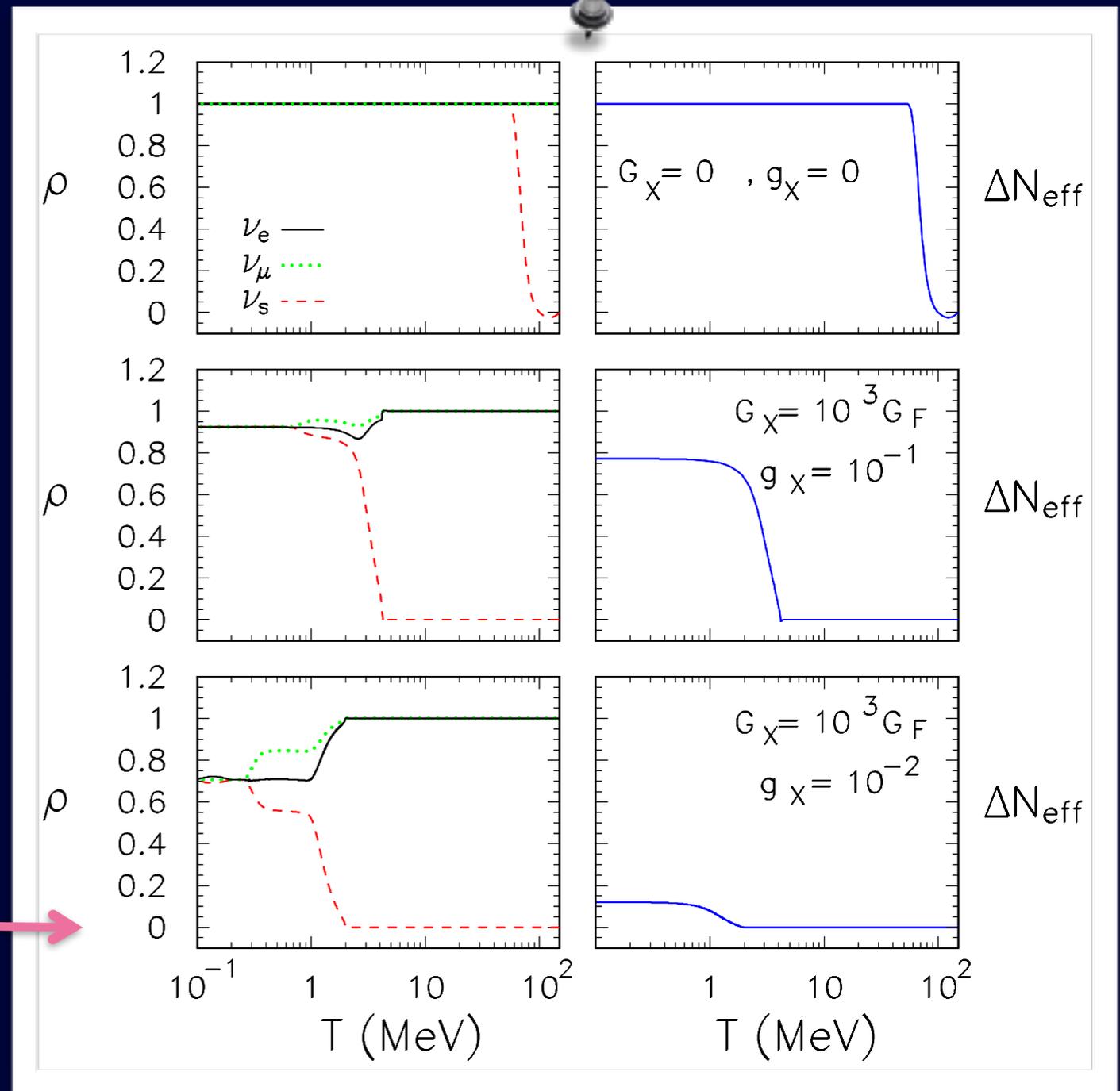
$\downarrow \quad \quad \quad \searrow$
 $\propto G_F^2 \quad \quad \quad \propto G_X^2$

Sterile production by secret interactions

Secret interactions: resonances around 1 MeV,
sterile ν suppressed.

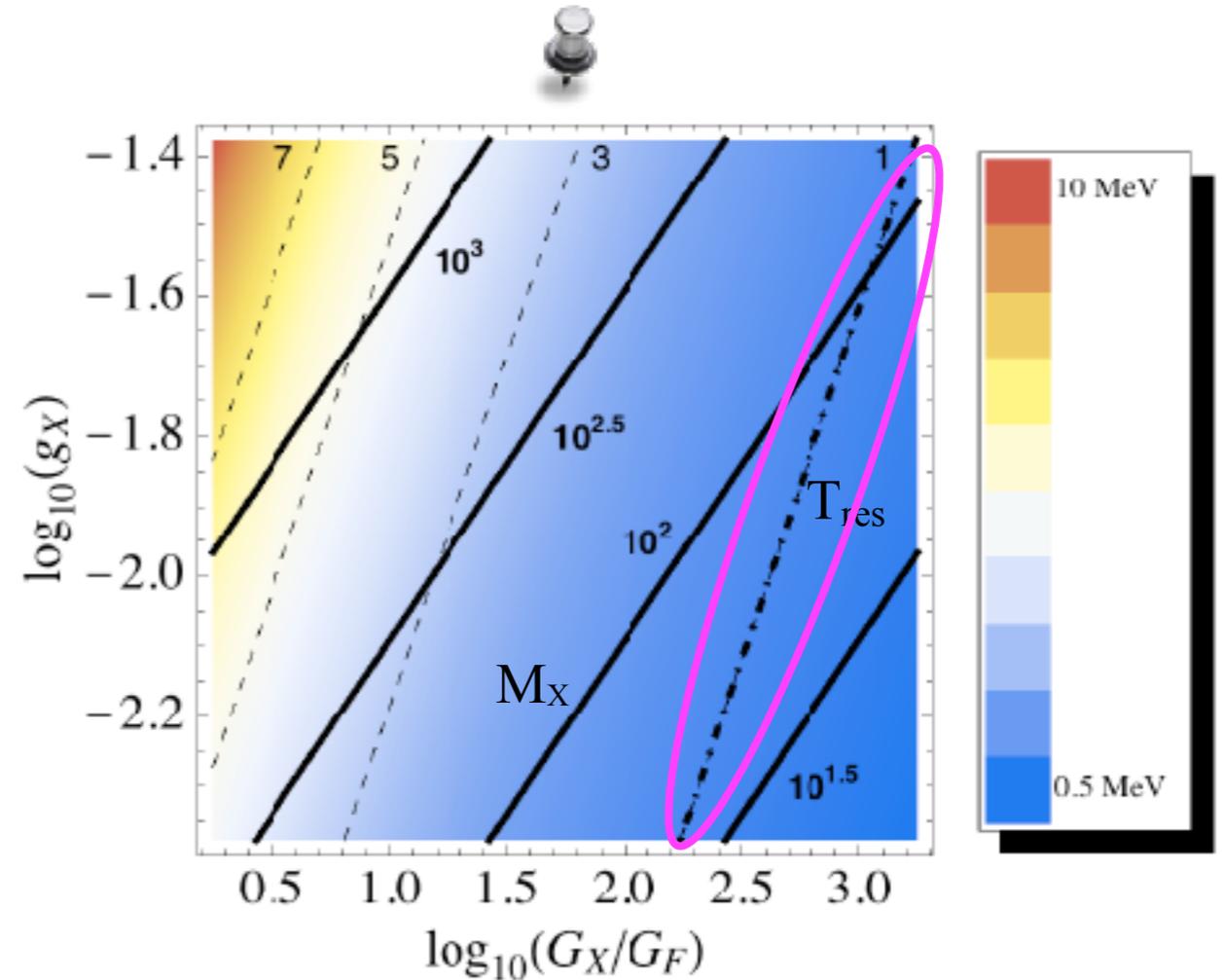
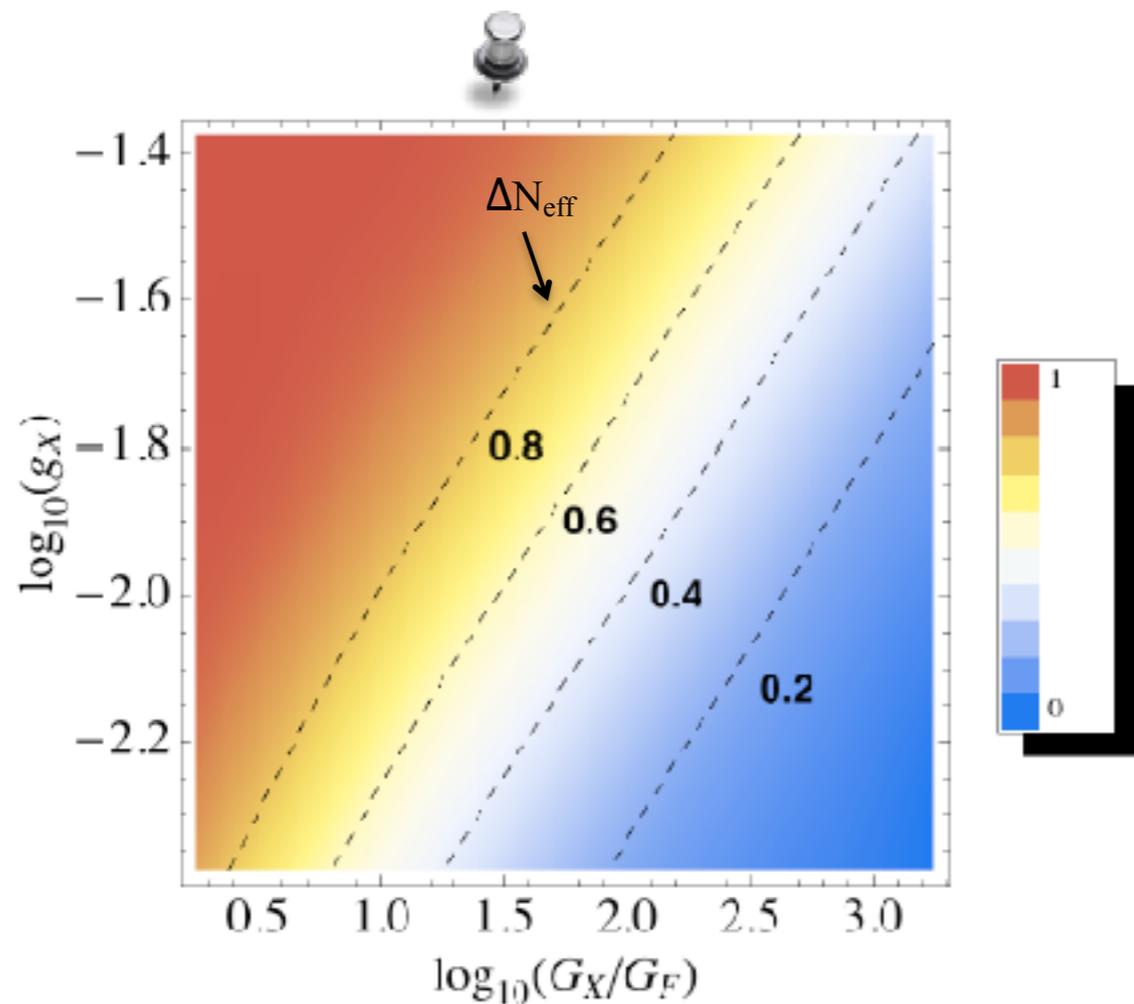
Note that also ν_e and ν_μ are depleted:
crucial for N_{eff} but also for BBN

$$\rho_{ee} = 0.7, \quad N_{\text{eff}} = 0.18$$



Secret interactions and BBN

Asymptotic values of ΔN_{eff} versus G_X and g_X



Resonance temperature in the plane (G_X, g_X)

Dashed curves: constant T_{res} contours

Solid curves: constant M_X contours

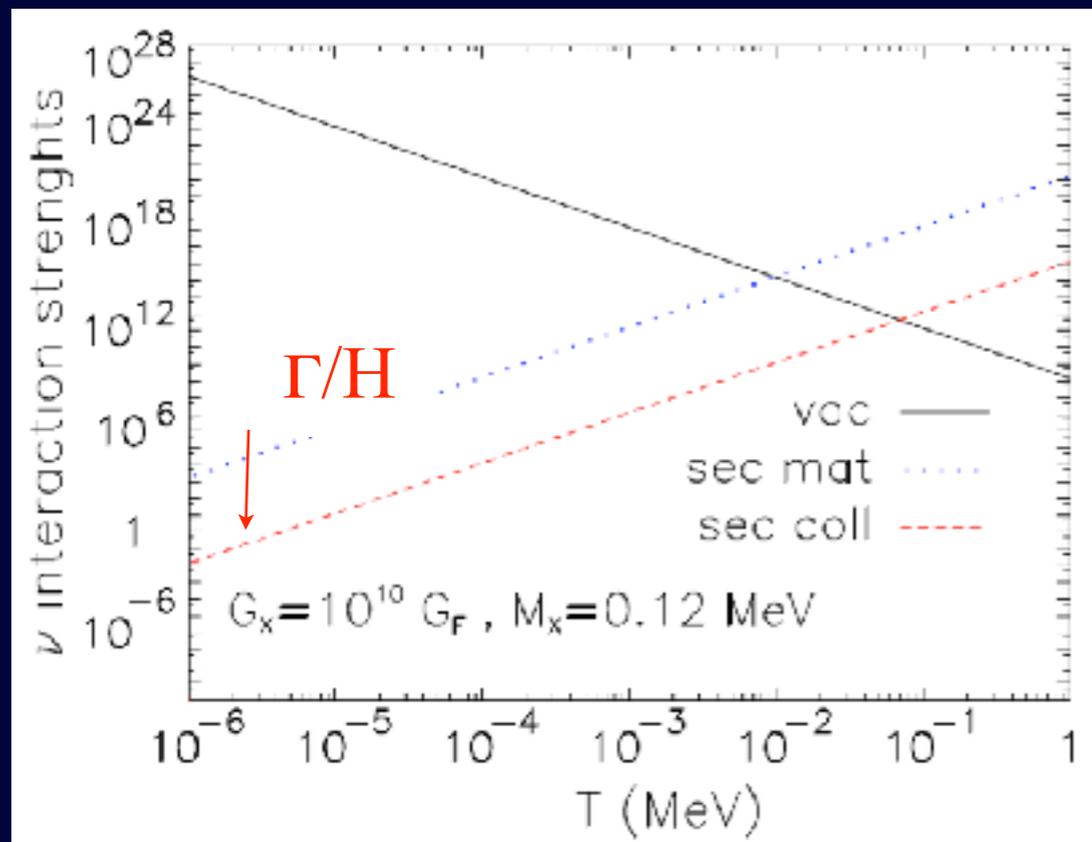
Cosmological constraints on secret interactions

Summarising:

✓ Very large $M_X \rightarrow$ thermalization of $\nu_s \leftrightarrow$ secret interactions do not have effect

✓ $400 \text{ MeV} < M_X < 40 \text{ MeV} \rightarrow$ severely constrained by BBN bounds

✓ $40 \text{ MeV} < M_X < 0.1 \text{ MeV} \rightarrow$ severely constrained by sterile mass bounds



For $M_X < 0.1 \text{ MeV} \rightarrow \nu_s$ could be still coupled at CMB and LSS epoch \rightarrow no free-streaming

Present cosmological mass bound obtained considering free-streaming ν



An appropriated analysis should be performed

A surprising feature on N_{eff}

- After the production, ν_s have a “grey-body” spectrum ($\rho_{\text{ss}} = 3/4$)....

... but the collisions and oscillations are still active pushing all neutrinos to a common FD distribution

Constraint: $n_{\nu \text{ TOT}}$ must be constant



T_ν is reduced by a factor $(3/4)^{1/3}$,
leading to an effect on the radiation density

$$\epsilon_{\nu, \text{in}} = 3 \times \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} \epsilon_\gamma$$

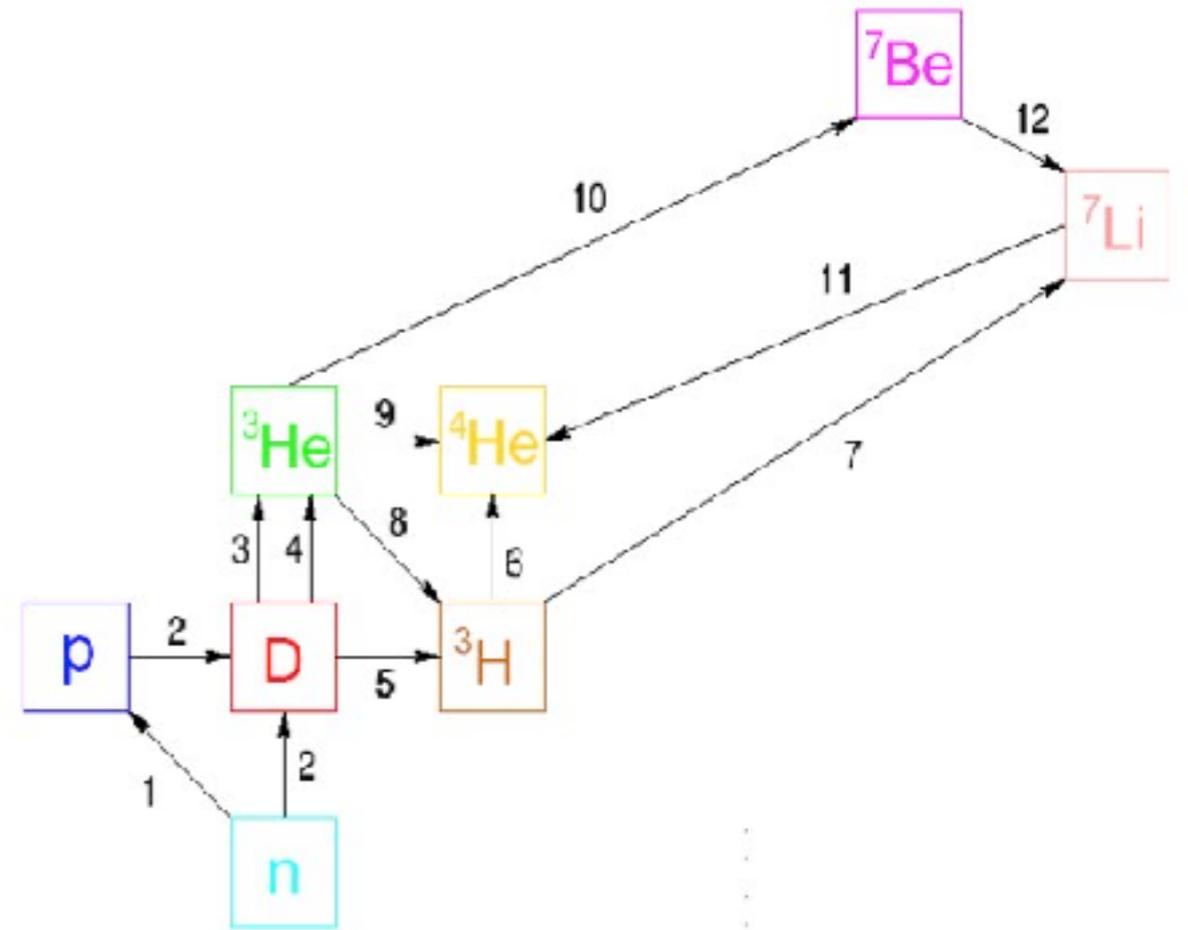
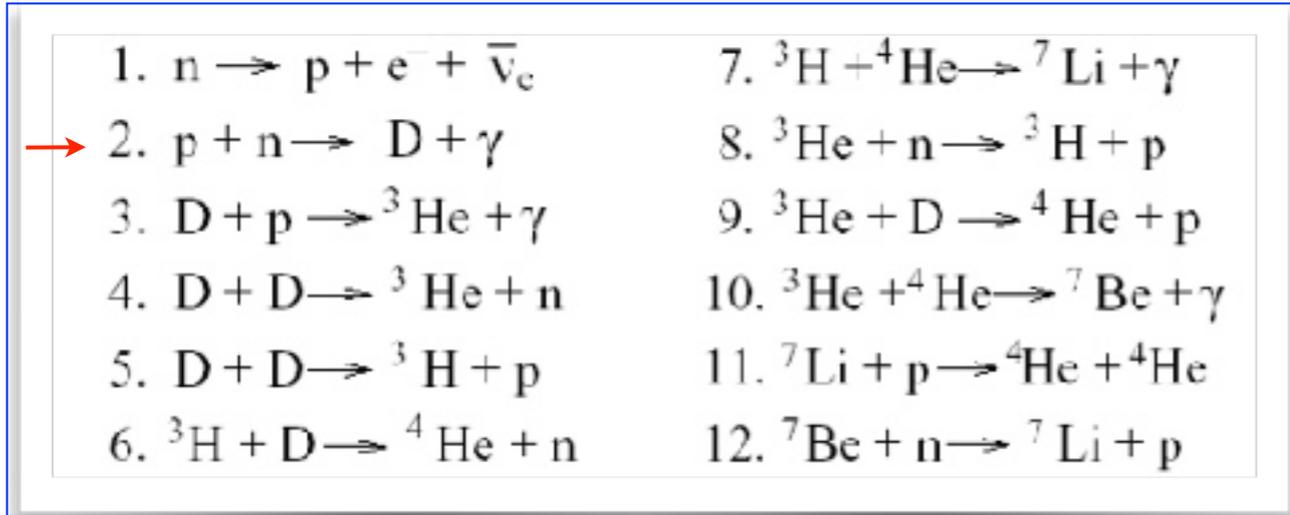


$$\epsilon_{\nu, \text{fin}} = 4 \times \left(\frac{3}{4} \right)^{4/3} \times \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} \epsilon_\gamma$$

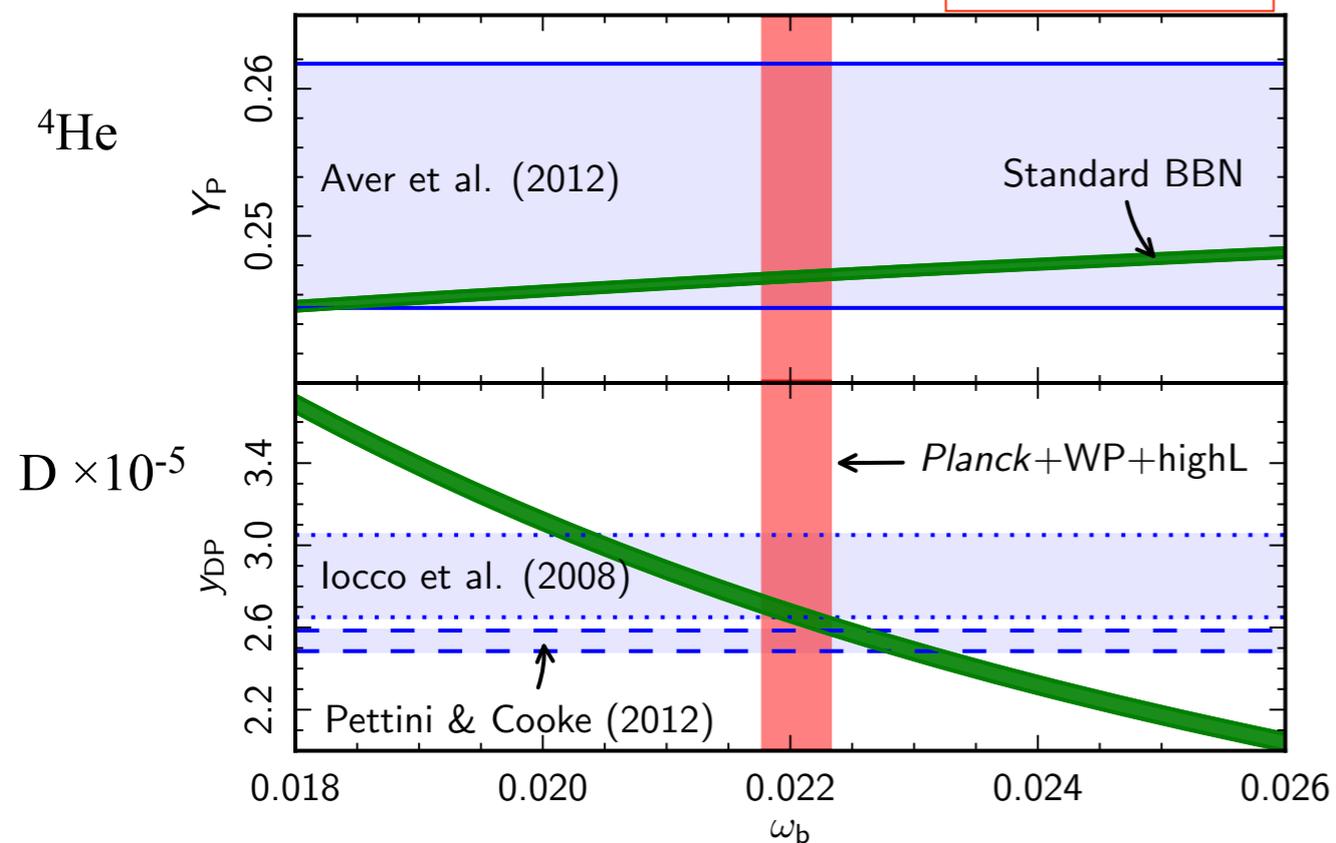
Big Bang Nucleosynthesis

* 0.1-0.01 MeV

Formation of light nuclei starting from D



Planck XVI, 2013



Prediction for ${}^4\text{He}$ and D in a **standard** BBN obtained by Planck collaboration using **PARthENoPE**

Blue regions: primordial yields from measurements performed in different astrophysical environments

$$\omega_b = 0.02207 \pm 0.00027$$