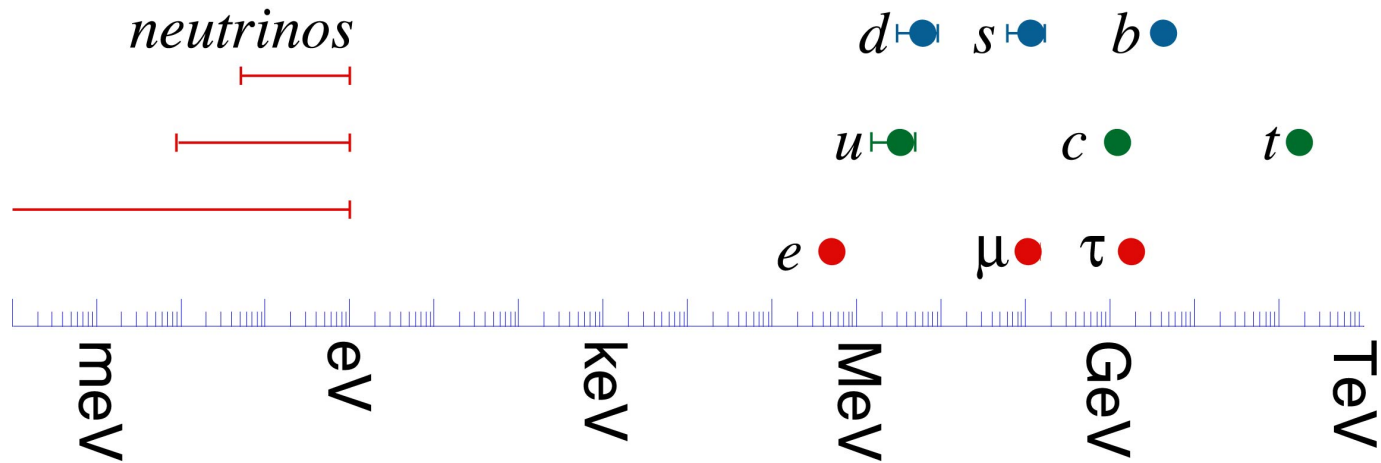


The seesaw path to leptonic CP violation

P. Hernandez
(CERN & University of Valencia)

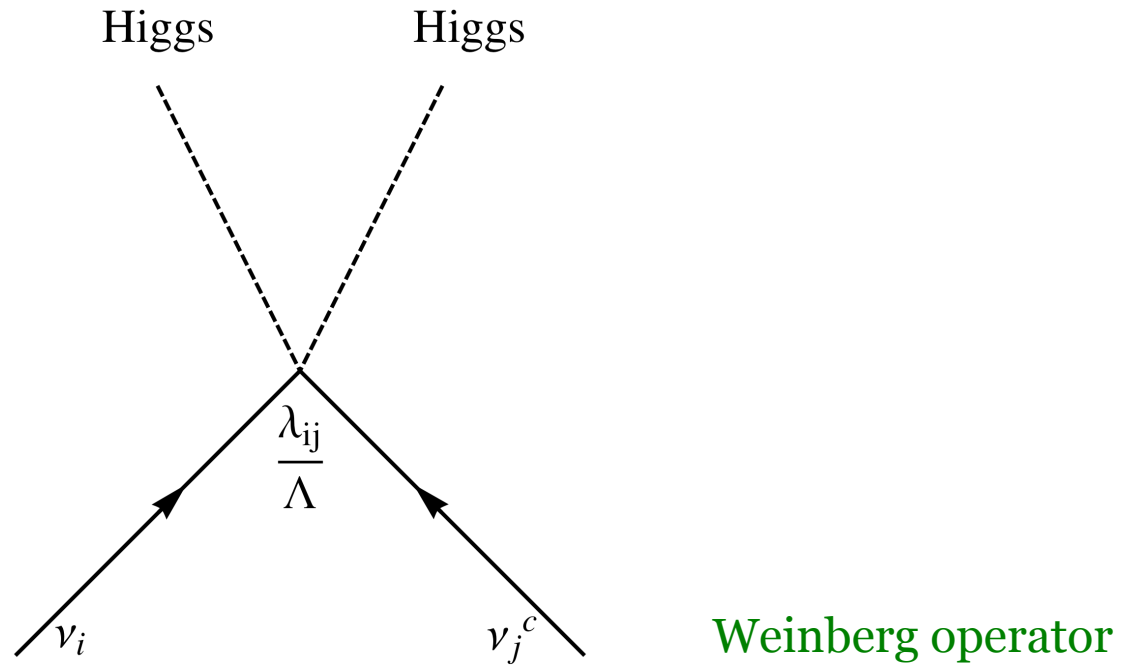
Caputo, PH, Kekic, Lopez-Pavón, Salvado arXiv:1611.05000

For the time being the only BSM signal: **neutrino masses**



Strongly suggests the existence of **a new physics scale**

Fermi-era of neutrino physics

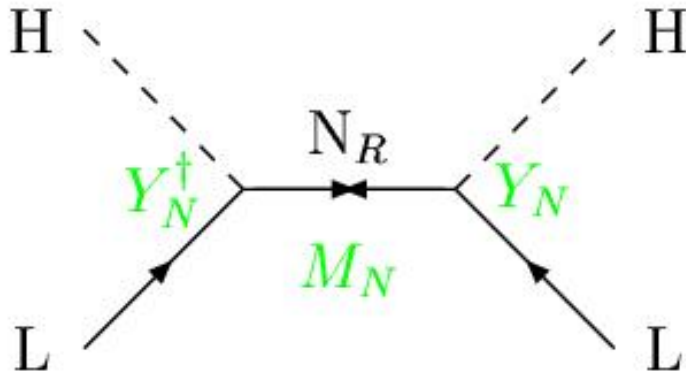


Neutrino masses = first higher dimensional operator

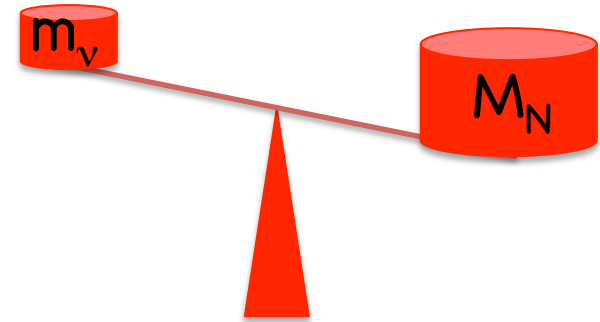
Why the seesaw ?

Minimality: SM+right-handed neutrinos

$$\mathcal{L} = \mathcal{L}_{SM} - \sum_{i=1}^{n_R} \bar{l}_L^\alpha Y^{\alpha i} \tilde{\Phi} \nu_R^i - \sum_{i,j=1}^{n_R} \frac{1}{2} \bar{\nu}_R^{ic} M_N^{ij} \nu_R^j + h.c.$$



$$n_R \geq 2$$

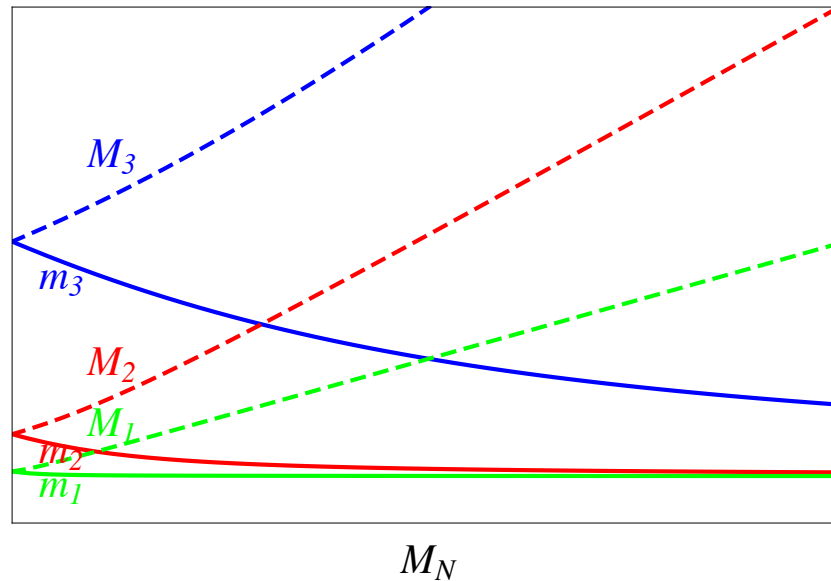


$$m_\nu = \frac{\alpha v^2}{\Lambda} \equiv Y_N^T \frac{v^2}{M_N} Y_N$$

Minkowski; Yanagida; Glashow; Gell-Mann, Ramond Slansky; Mohapatra, Senjanovic...

Type I seesaw models

$n_R = 3$: 18 free parameters (6 masses+6 angles+6 phases) out of which we have measured 2 masses and 3 angles...

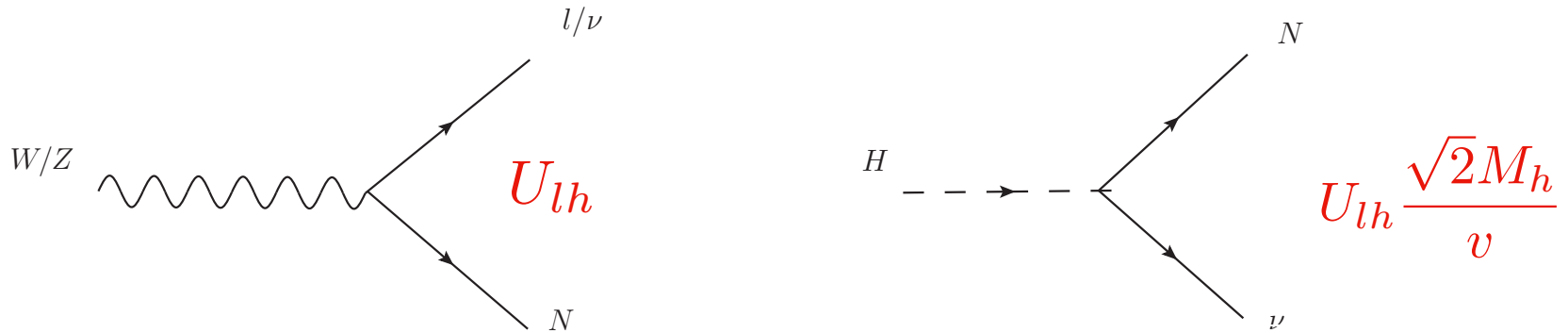


Type I seesaw models

Phenomenology (beyond neutrino masses) of these models depends on the heavy spectrum and the size of active-heavy mixing:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U_{ll} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} + U_{lh} \begin{pmatrix} N_1 \\ N_2 \\ N_3 \end{pmatrix}$$

Type I seesaw models



$$U_{lh} \simeq iU_{\text{PMNS}} \sqrt{m_l} R \frac{1}{\sqrt{M_h}} \quad \text{Casas-Ibarra}$$

R: general orthogonal complex matrix (contains all the parameters we cannot measure in neutrino experiments)

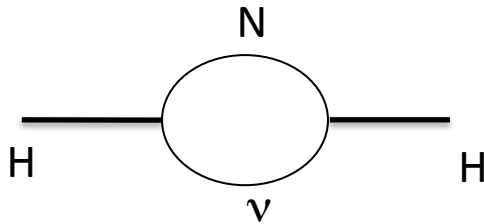
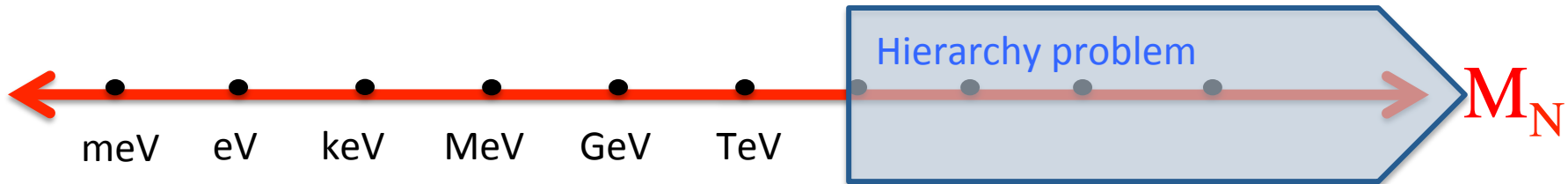
Strong correlation between active-heavy mixing and neutrino masses, but the naïve scaling ($|U_{lh}|^2 \sim m_l/M_h$) too naïve...

Why low-scale seesaw ?

$$M_N \leq v$$

To avoid fine-tuning

The new scale is stable under radiative corrections due to Lepton Number symmetry but the EW is not!



$$\delta m_H^2 = \frac{Y^\dagger Y}{4\pi^2} M_N^2 \log \frac{M_N}{\mu}$$

Vissani

$M_N \gg m_H$ not natural in the absence of SUSY

Testability

eg. leptogenesis in the low-scale minimal model $n_R=2$

PH, Kekic, López-Pavón, Racker, Salvadó 1606.06719

Bayesian posterior probabilities (using nested sampling Montecarlo Multinest)

$$\mathcal{L} = - \left(\frac{Y_B(\text{param}) - Y_B^{\text{obs}}}{\sigma_{Y_B}} \right)^2$$

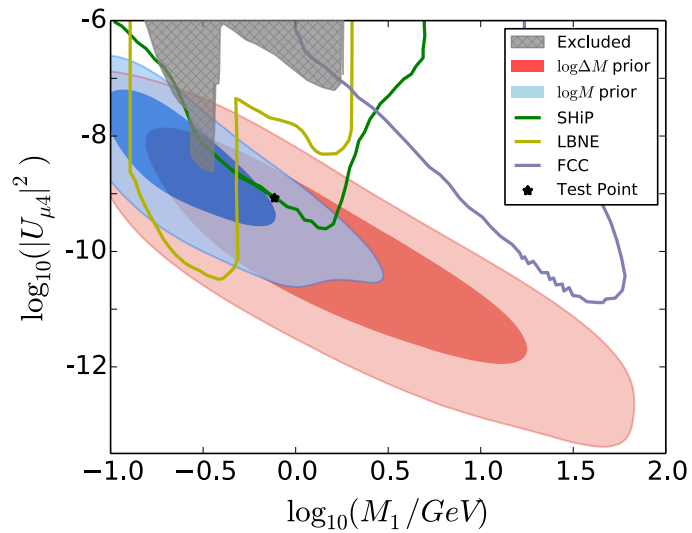
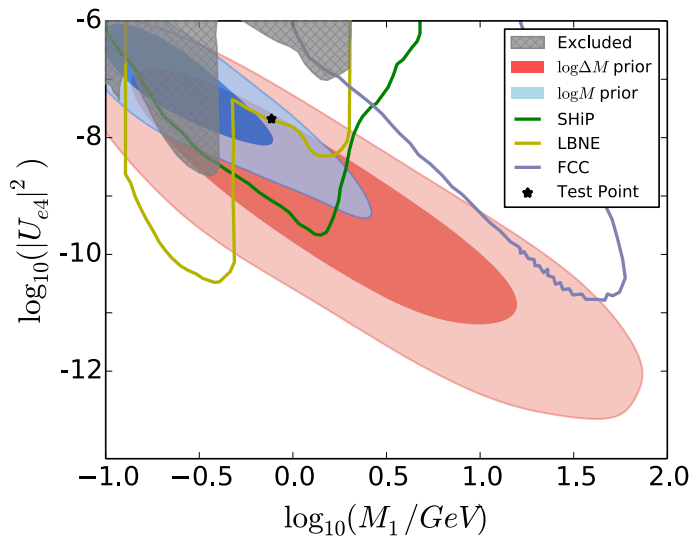
Use Casas-Ibarra parametrization: fix light neutrino masses and mixings to the best fit oscillation points (IH/NH) and vary

$$R(\theta + i\gamma); U_{PMNS}(\delta, \phi_1); M_1, M_2$$

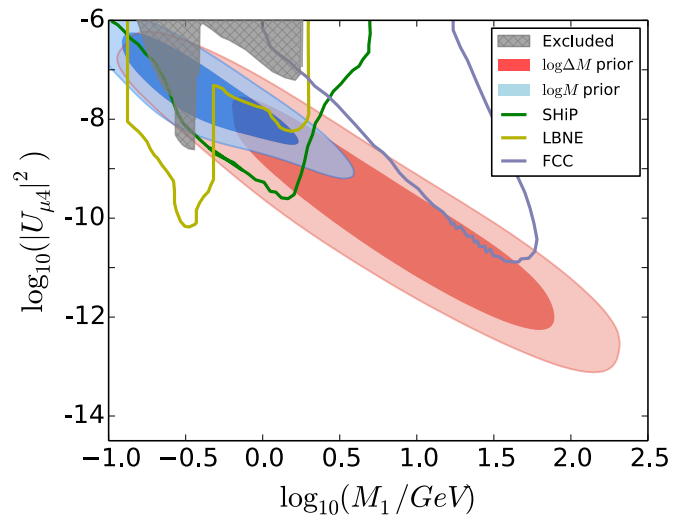
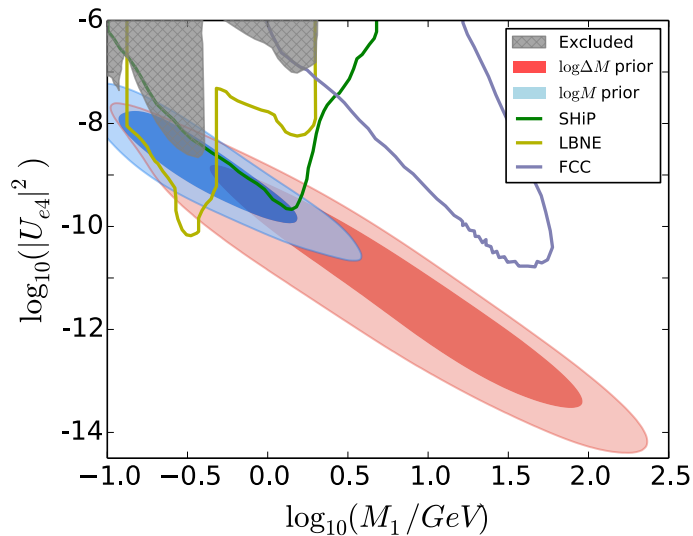
Flat priors in:

$$\theta = [0, \pi]; \delta = [0, 2\pi]; \phi_1 = [0, 2\pi]; \gamma = [-9, 9]; \\ \log_{10} M_1 \text{ and } \log_{10} M_2 / \log_{10}(M_2 - M_1)$$

Full exploration of the minimal model $n_R=2$

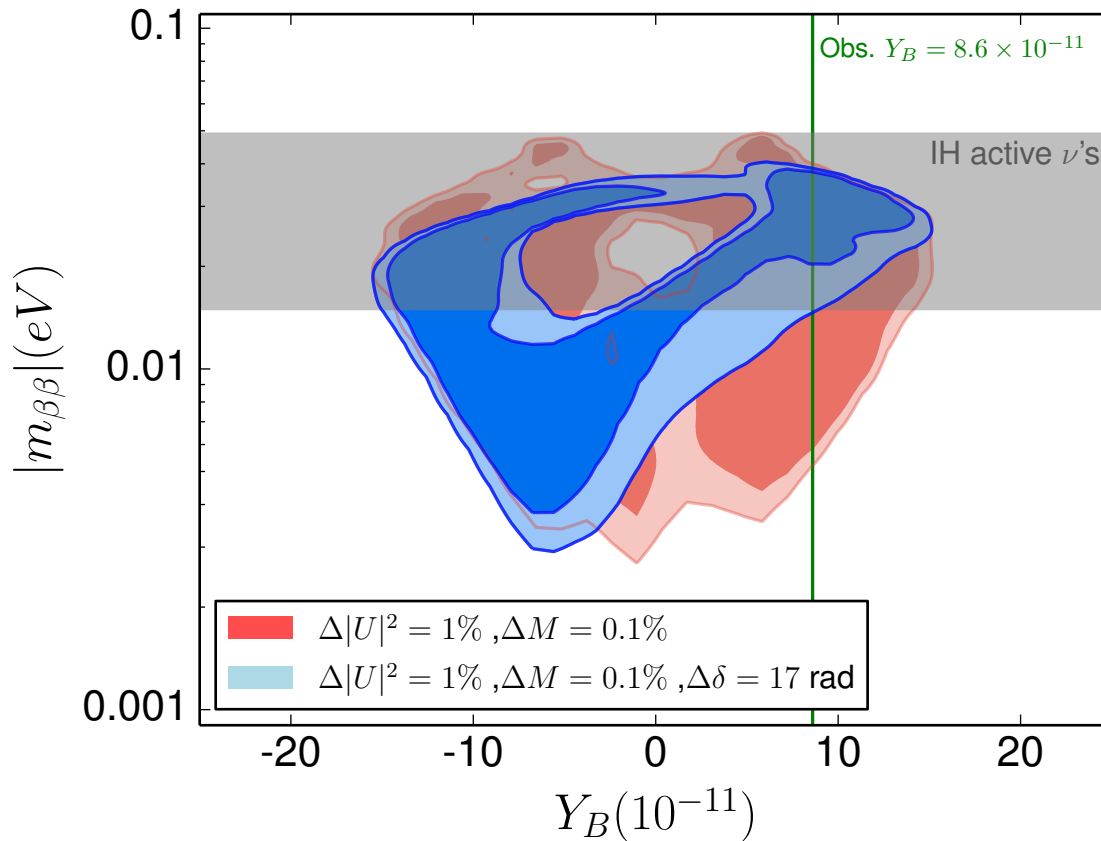


IH



NH

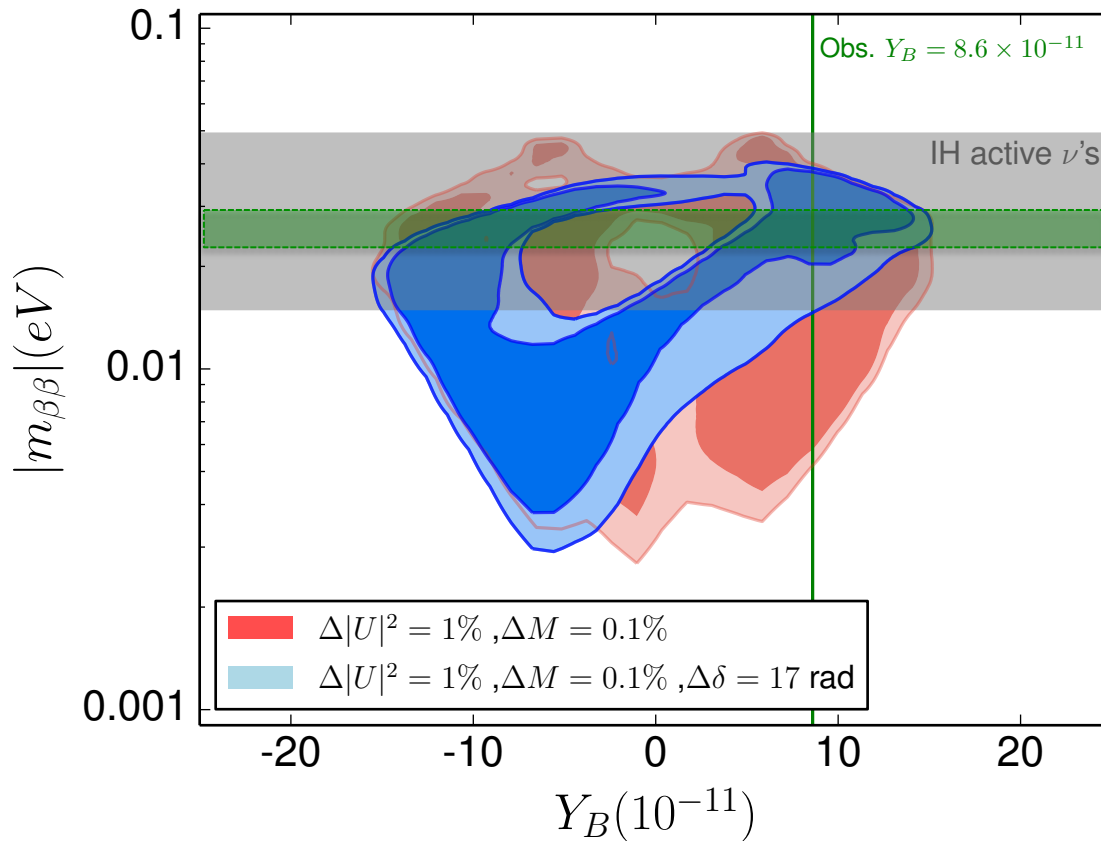
Predicting Y_B in the minimal model $n_R=2$ (IH)



PH, Kekic, López-Pavón, Racker, Salvadó 1606.06719

The GeV-miracle: the measurement of the mixing to e/μ of the sterile states, neutrinoless double-beta decay and δ in neutrino oscillations have a chance to give a prediction for Y_B

Predicting Y_B in the minimal model $n_R=2$ (IH)



PH, Kekic, López-Pavón, Racker, Salvadó 1606.06719

The GeV-miracle: the measurement of the mixing to e/μ of the sterile states, neutrinoless double-beta decay and δ in neutrino oscillations have a chance to give a prediction for Y_B

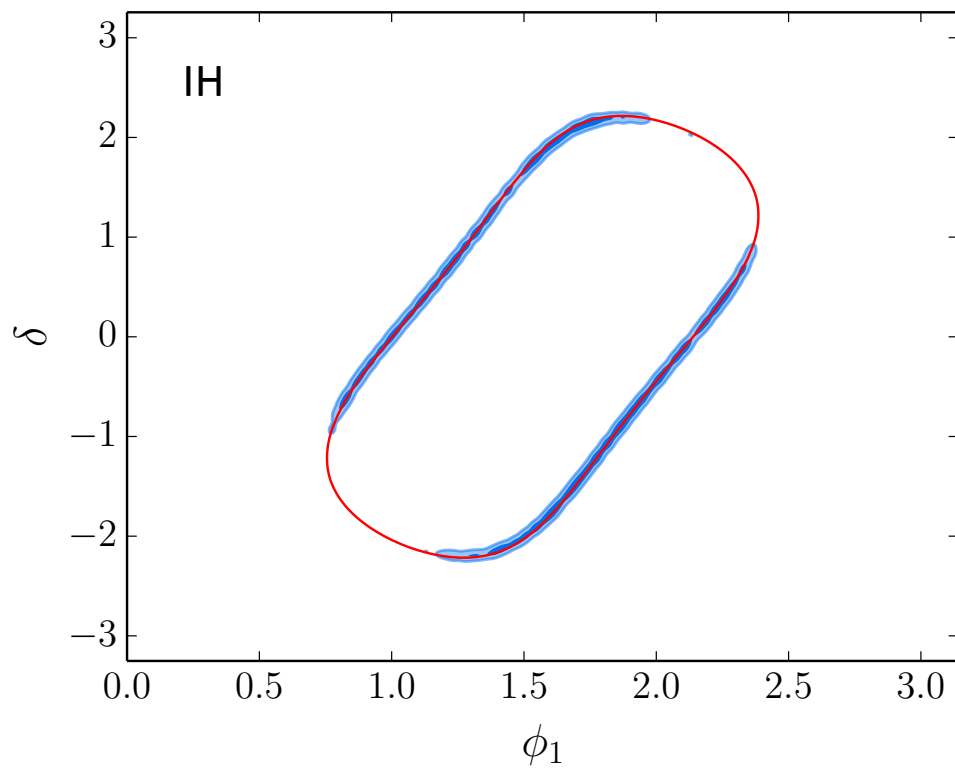
Leptonic CP violation
(whether Y_B or not)

In the regime of sensitivity of SHIP/FCC-ee , the flavour structure of the heavy mixings is strongly correlated with that of the U_{PMNS} matrix.

$$\begin{aligned}
 |U_{e4}|^2 M_1 \simeq |U_{e5}|^2 M_2 &\simeq A \left[(1 + \sin \phi_1 \sin 2\theta_{12})(1 - \theta_{13}^2) + \frac{1}{2}x^2 s_{12}(c_{12} \sin \phi_1 + s_{12}) + \mathcal{O}(\epsilon^3) \right], \\
 |U_{\mu 4}|^2 M_1 \simeq |U_{\mu 5}|^2 M_2 &\simeq A \left[\left(1 - \sin \phi_1 \sin 2\theta_{12} \left(1 + \frac{1}{4}x^2 \right) + \frac{1}{2}x^2 c_{12}^2 \right) c_{23}^2 \right. \\
 &\quad + \theta_{13}(\cos \phi_1 \sin \delta - \sin \phi_1 \cos 2\theta_{12} \cos \delta) \sin 2\theta_{23} \\
 &\quad \left. + \theta_{13}^2(1 + \sin \phi_1 \sin 2\theta_{12})s_{23}^2 + \mathcal{O}(\epsilon^3) \right],
 \end{aligned}$$

$$\epsilon \sim e^{-\gamma} \sim \theta_{13} \sim x \equiv \sqrt{\frac{\Delta m_{\text{sol}}^2}{\Delta m_{\text{atm}}^2}} \qquad A \equiv \frac{e^{2\gamma} \sqrt{\Delta m_{\text{atm}}^2}}{4},$$

For $n_{\text{R}}=2$ the ratio of e/μ mixings depends on the two phases of the U_{PMNS} matrix:
 δ, ϕ_1



From a measurement of the ratio

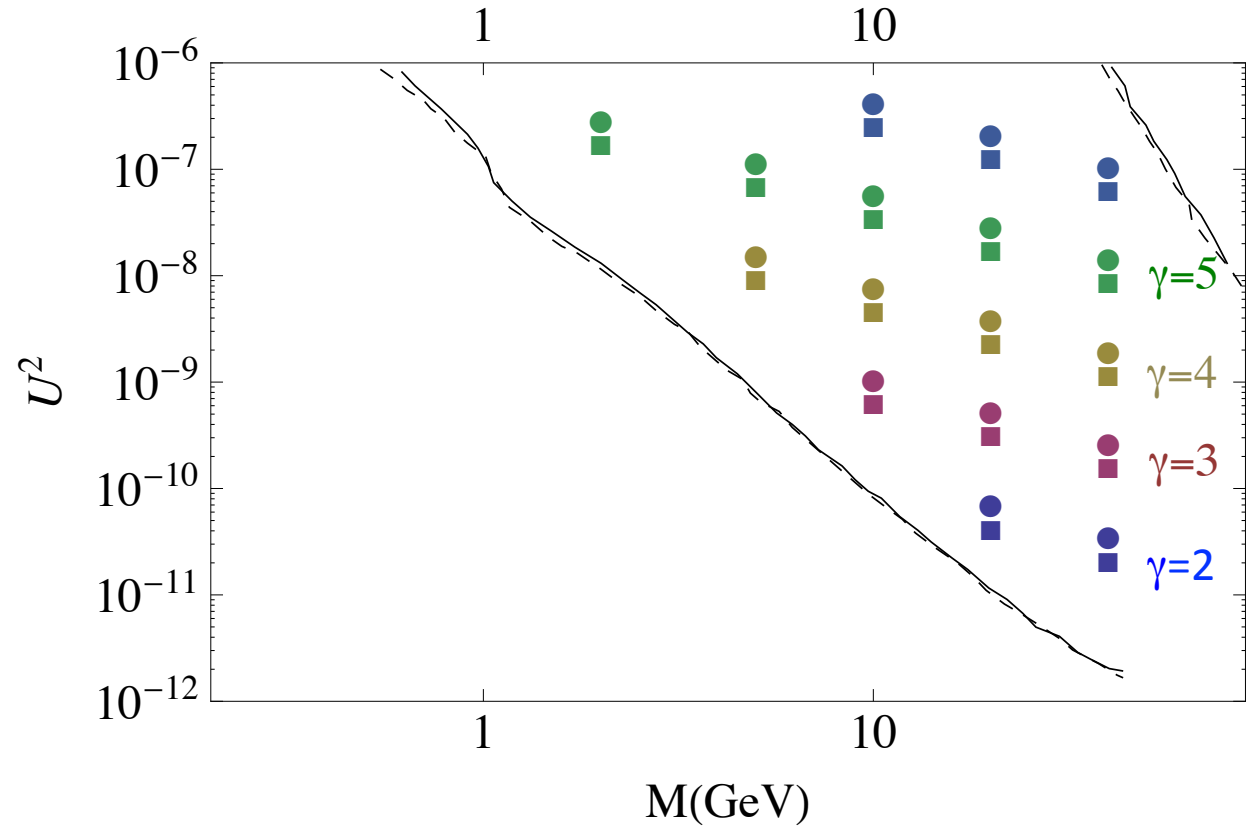
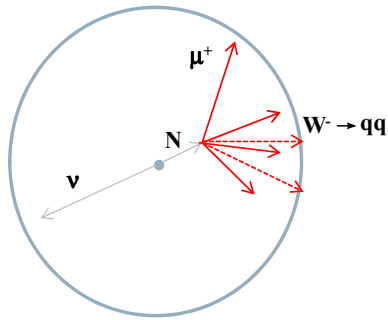
$$\frac{|U_{eh}|^2}{|U_{\mu h}|^2}$$

If SHIP/FCC-ee measures the heavy neutrinos and their mixings to e/μ :

Can we exclude a real U_{PMNS} matrix ie.
discover leptonic CP violation in mixing ?

$$(\delta, \phi_1) \neq (0/\pi, 0/\pi)$$

FCC-ee sensitivity region (10^{13} Z's, 0.1mm-5m DV search)



Blondel, Graverini, Serra, Shaposhnikov, FCC-ee study Team Collaboration, arXiv:1411.5230;
 A. Abada et al arXiv:1412.6322; S. Antusch and O. Fischer arXiv:1502.05915

-> Talk by O. Fischer

Test statistics (shown to be distributed like χ^2 (1dof))

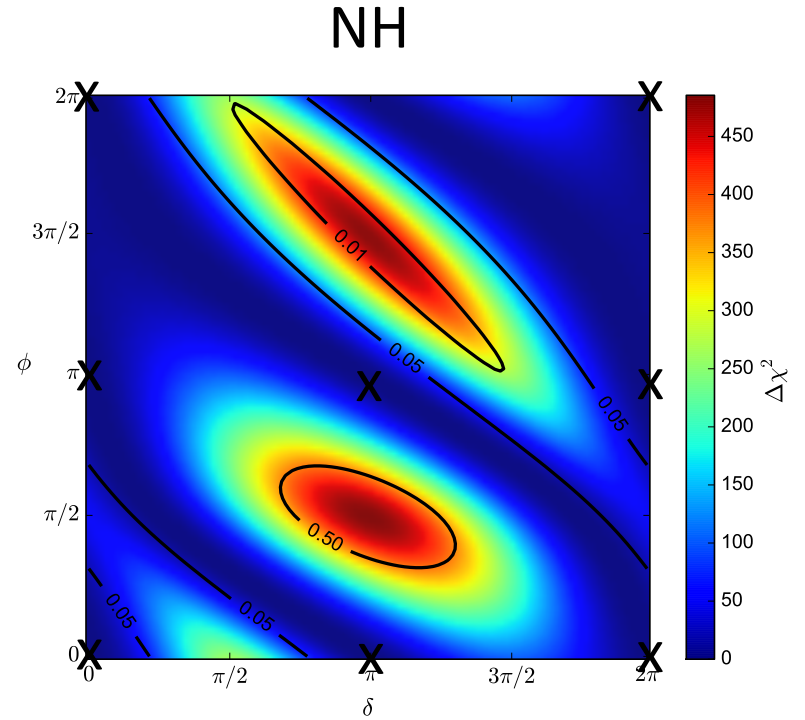
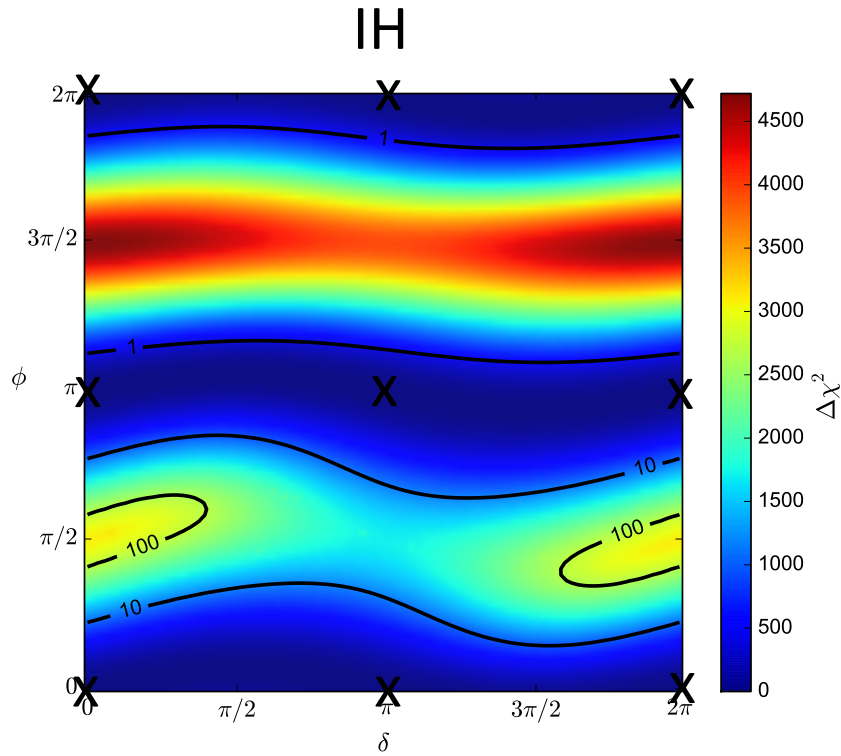
$$\Delta\chi^2 \equiv -2 \sum_{\alpha=e,\mu} N_{\alpha}^{\text{true}} - N_{\alpha}^{\text{CP}} + N_{\alpha}^{\text{true}} \log \left(\frac{N_{\alpha}^{\text{CP}}}{N_{\alpha}^{\text{true}}} \right) + \left(\frac{M_1 - M_1^{\text{min}}}{\Delta M_1} \right)^2 .$$

$N_{\alpha}^{\text{true}}(\gamma, \theta, M_1, \delta, \phi) = \#$ events for true parameters

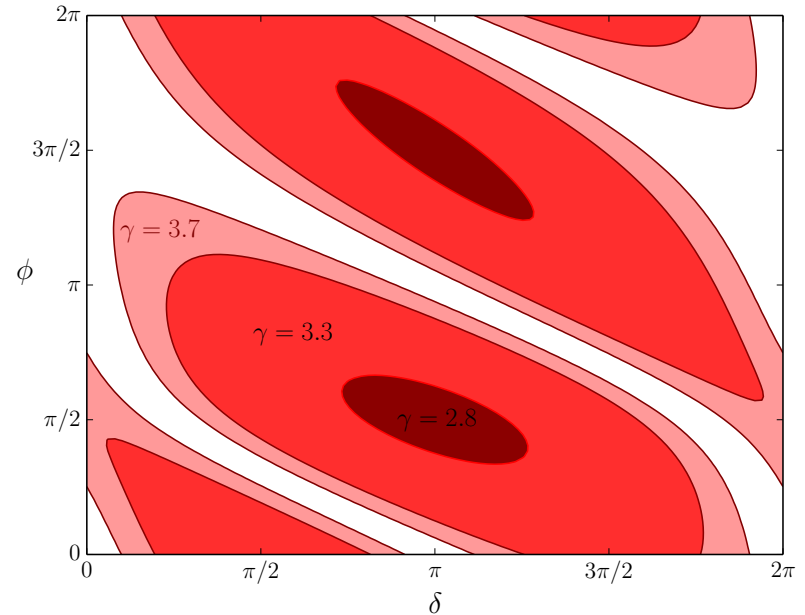
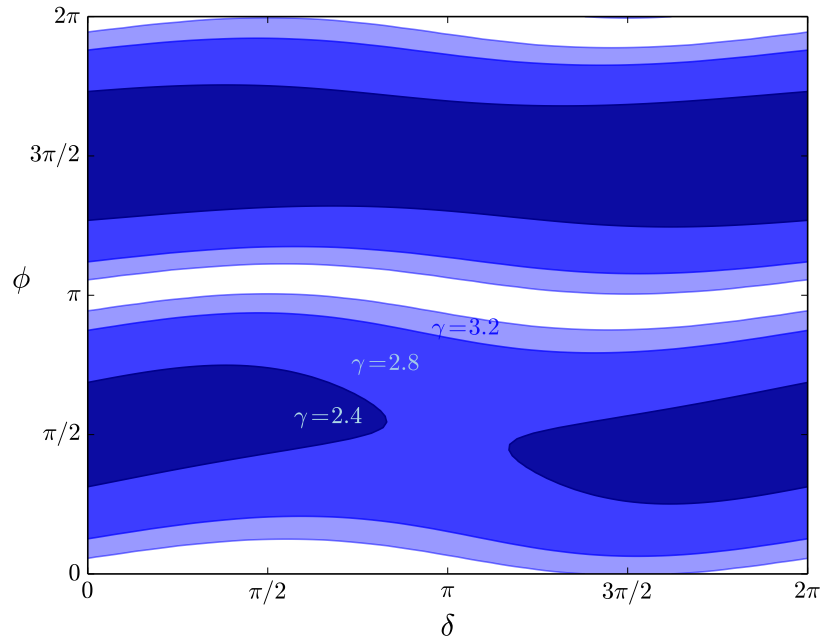
$N_{\alpha}^{\text{CP}} = \#$ events for best fit CP conserving hypothesis

CP: $(\delta, \phi) = (0/\pi, 0/\pi)$ any γ, θ, M_1

Assumed the neutrino ordering is known:



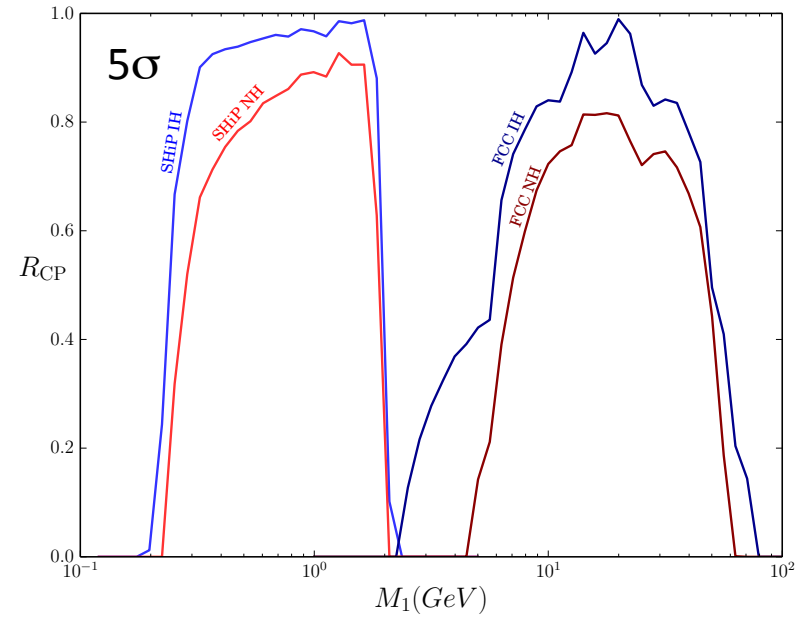
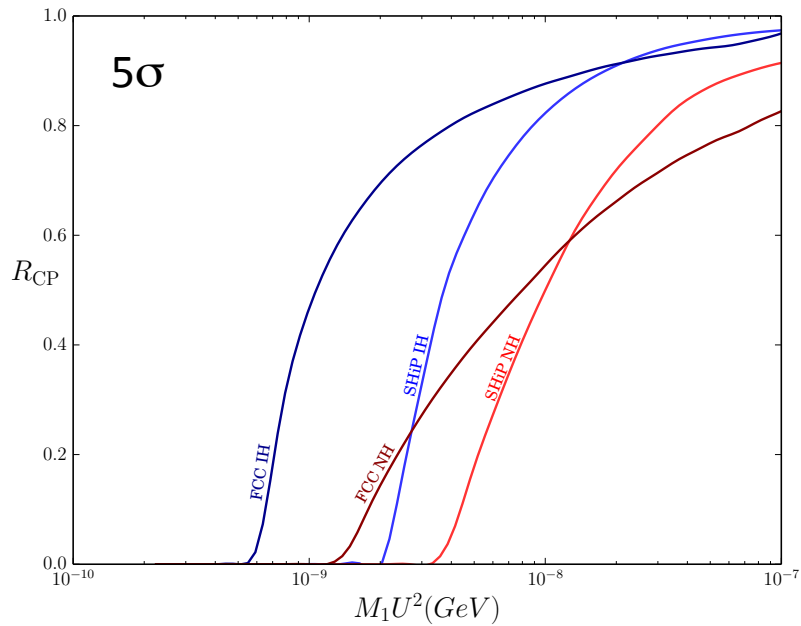
Leptonic CP violation 5σ CL discovery regions



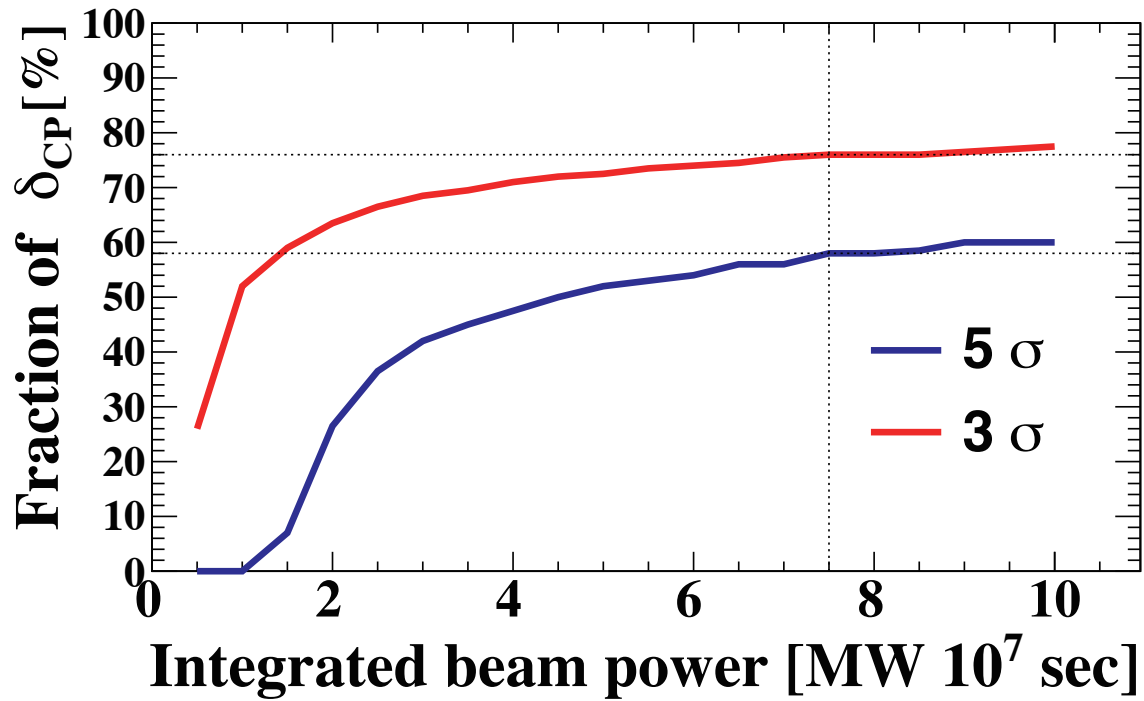
(no systematic error included)

$R_{\text{CP}=5\sigma}$ CP-fraction =
fraction of the area of the CP rectangle which is colored

Discovery potential for leptonic CP violation in mixing



CP fraction of δ from future neutrino oscillation experiments:



HyperKamiokande LOI

BACKUPS

Testability: Leptogenesis

$$\Gamma_s(T) \sim y^2 T \sim \frac{M_N m_\nu}{v^2} T \quad H(T) = \sqrt{\frac{4\pi^3 g_*(T)}{45}} \frac{T^2}{M_P}$$

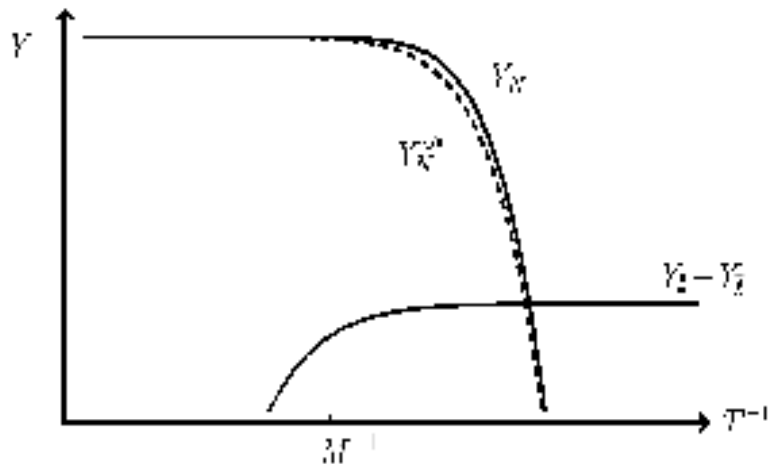
$$\frac{\Gamma_s(T_{EW})}{H(T_{EW})} \sim 5 \left(\frac{M}{1\text{GeV}} \right) \left(\frac{m_\nu}{0.05\text{eV}} \right)$$

CP asymmetries arise in **production** of sterile states that can seed the baryon asymmetry

Akhmedov, Rubakov, Smirnov

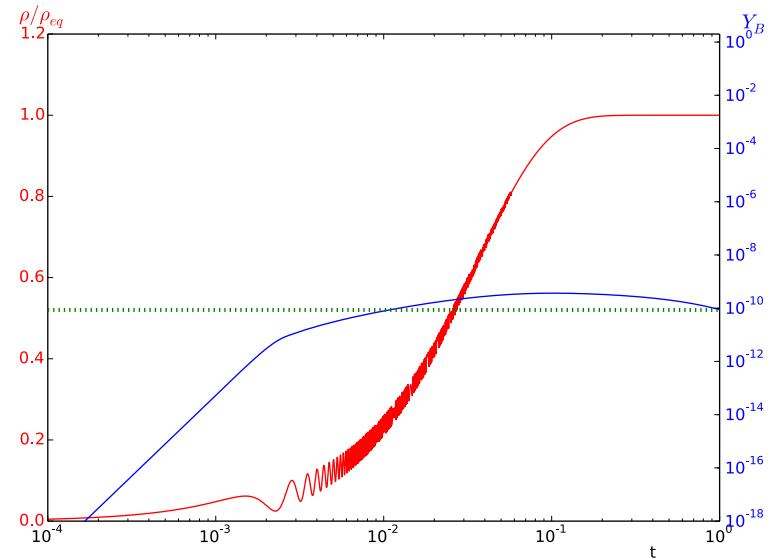
High-scale leptogenesis

$$M_N \gg v$$



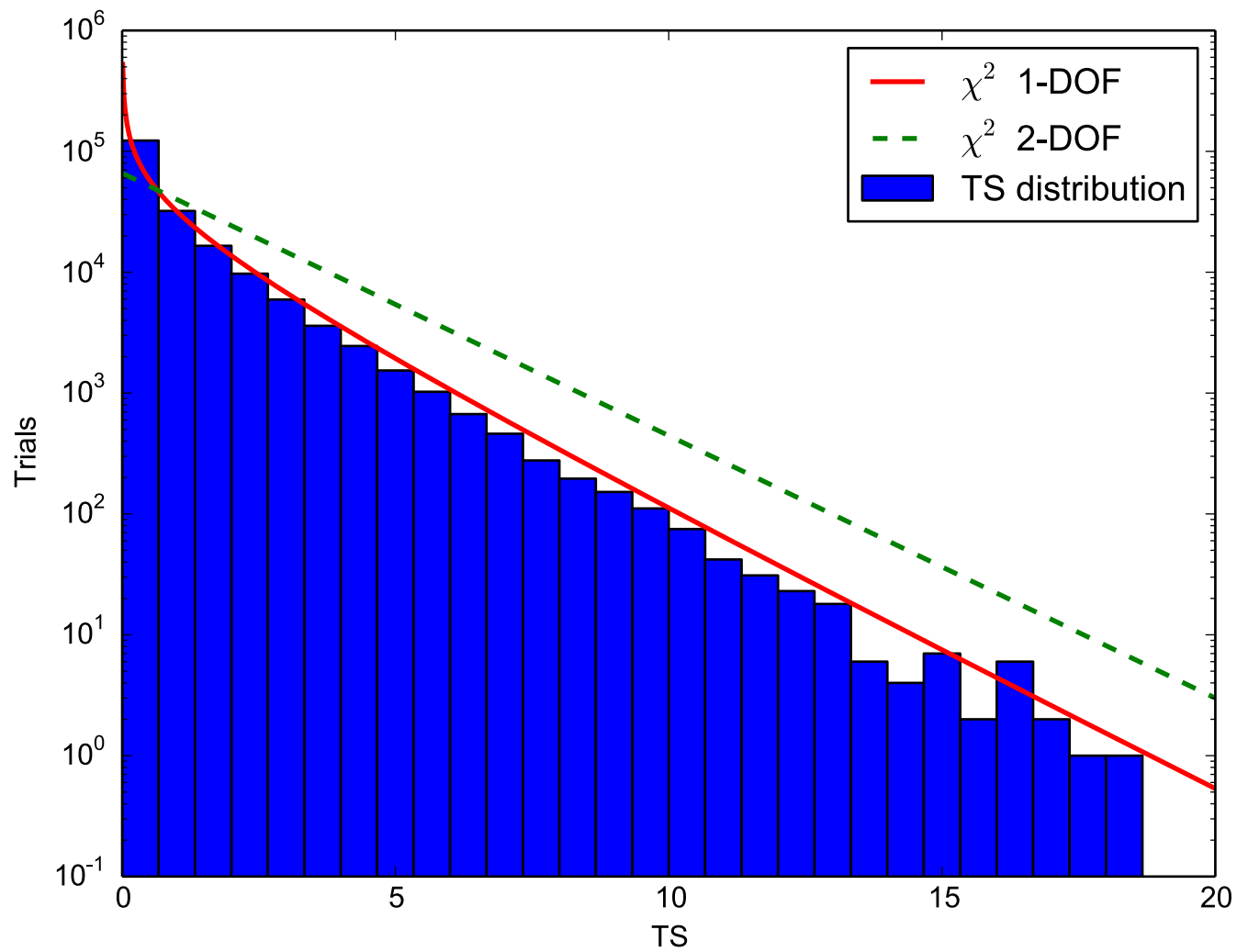
Low-scale leptogenesis

$$M_N \sim v$$



T_{EW}

Courtesy of M. Kevic



Adding information from tau mixings

