Top Quark Mass
-
What It Is and How to Interpret Reconstruction Measurement ?

André H. Hoang

University of Vienna

$\int dk \, \Pi$ Doktoratskolleg Particles and Interactions

universität wien

FWF Der Wissenschaftsfonds.

JetMET 100fb-1 Workshop, Helsinki, May 10-12, 2017

Outline

- Introduction
- Monte Carlo generators and the top quark mass
- Theory issues for the LHC
- Calibration of the Monte Carlo top mass parameter: $e^+e^- \rightarrow t\,\bar{t}$ (2-jettiness)

  Butenschön, Dehnadi, Mateu, Preisser, Stewart,AH; PRL 117 (2016) 153

- Relation of $M_t^{\text{Pythia 8.2}}$ and m_t^{pole}

- Factorization for $pp \rightarrow t\,\bar{t}$ with and w/o jet grooming
- Studies for LHC top mass measurements with SoftDrop

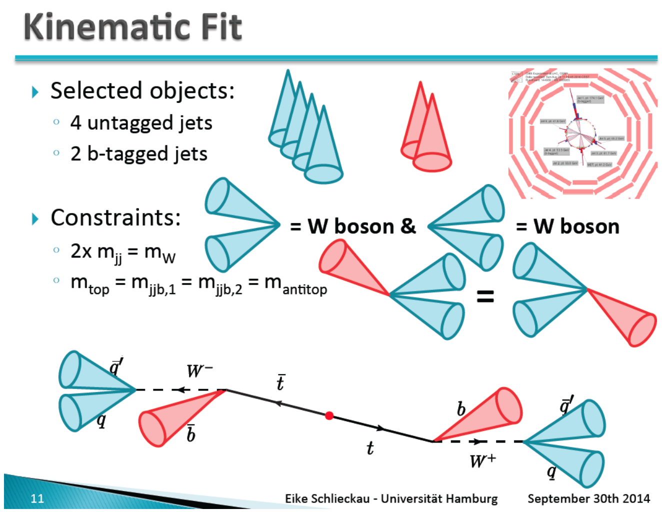
  Mantry, Pathak, Stewart, AH; to appear soon

- Summary, future plans

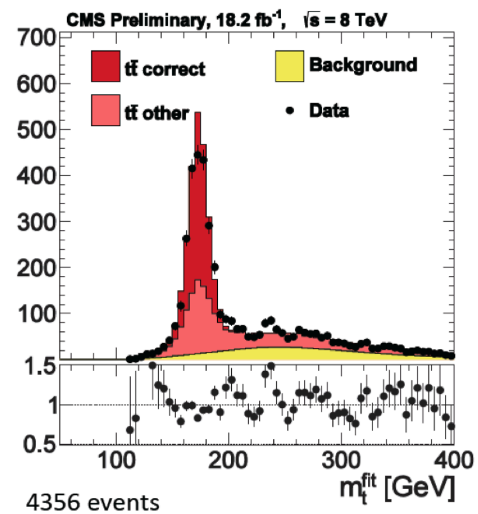
Main Top Mass Measurements Methods

LHC+Tevatron: Direct Reconstruction

kinematic mass determination

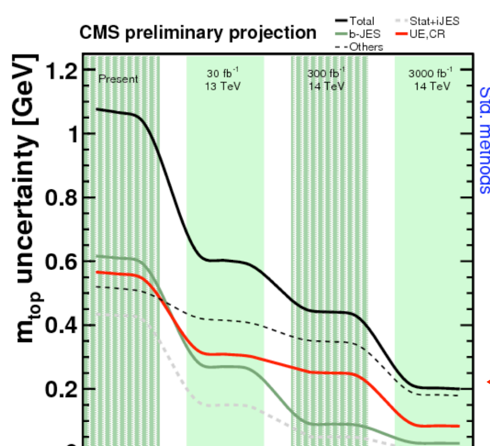


CMS-PAS-TOP-14-002



4356 events

Determination of the best-fit value of the Monte-Carlo top quark mass parameter



⊕ High top mass sensitivity

⊖ Precision of MC ?

⊖ Meaning of m_t^{MC} ?

Δm_t ~ 0.5 GeV

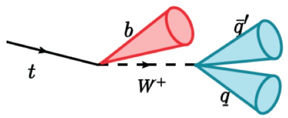
Δm_t ~ 200 MeV (projection)

$$m_t^{MC} = 174.34 \pm 0.64 \quad \text{(Tevatron final, 2014)}$$
$$m_t^{MC} = 172.44 \pm 0.49 \quad \text{(CMS Run-1 final, 2015)}$$
$$m_t^{MC} = 172.84 \pm 0.70 \quad \text{(ATLAS Run-1 final, 2016)}$$

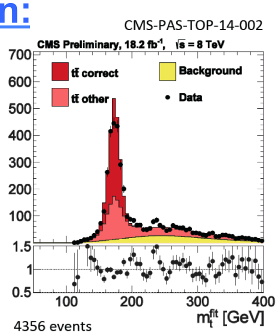
Top Mass Measurements Methods

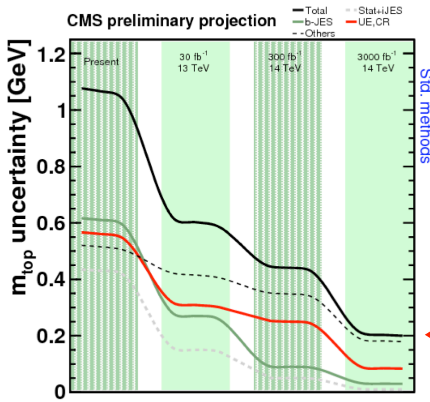
LHC+Tevatron

Direct Reconstruction:



kinematic mass determination





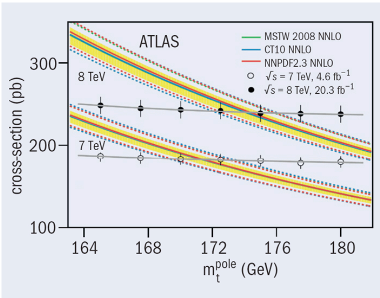
- ⊕ High top mass sensitivity
- ⊖ Precision of MC ?
- ⊖ Meaning of m_t^{MC} ?

$\Delta m_t \sim 0.5$ GeV

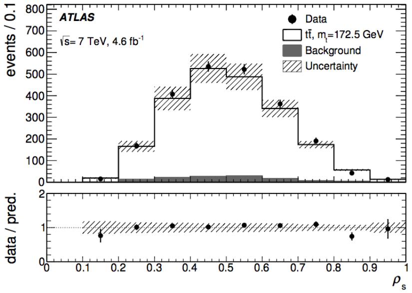
$\Delta m_t \sim 200$ MeV (projection)

Indirect Mass Fit:

global mass dependence



Total cross section



tt+jet invariant mass

- ⊕ pQCD calculations dominate
- ⊕ Control of mass scheme
- ⊖ Lower top mass sensitivity
- ⊖ High sensitivity to norm errors

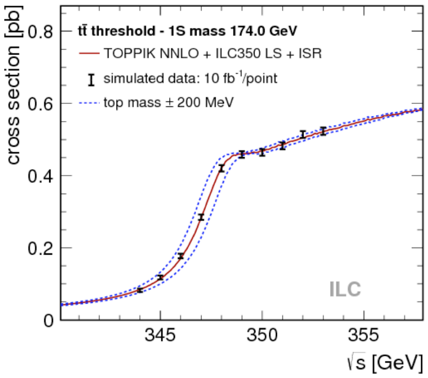
$\Delta m_t \sim$ 1-2 GeV

Future Linear Collider:

Top Pair Threshold:

kinematic mass determination

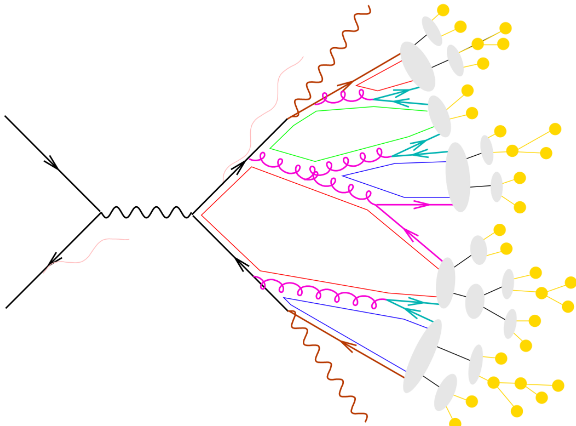
perturbative toponium



- ⊕ High top mass sensitivity
- ⊕ pQCD calculations dominate
- ⊕ Control of mass scheme

$\Delta m_t \sim 100$ MeV

Monte-Carlo Event Generators



1) Matrix elements (LO/NLO)
2) Parton shower (LL)
3) Hadronization model

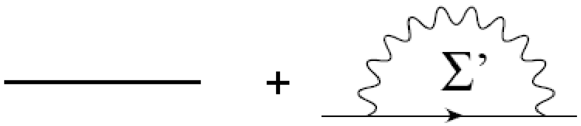
- Full simulation of all processes (all experimental aspects accessible)
- QCD-inspired: partly first principles QCD ⇔ partly model
- Description power of data better than intrinsic theory accuracy.
- Top quark in parton shower: treated like a real particle ($m_t^{MC} \approx m_t^{pole}$ +?).
- Top quark in matrix elements: $m_t^{MC} = m_t^{pole}$

BUT: parton showers sum (real & virtual !) perturbative corrections only above the shower cut and not pickup any corrections from below.

Uncertainty (a): But how precise is modelling? ⟶ Part of exp. Analyses

Unvertainty (b): What is the meaning of MC QCD parameters? ⟶ Calibration & Theory

Top Quark Mass



$$= p - m^0 - \Sigma(p, m^0, \mu)$$

$$\Sigma(m^0, m^0, \mu) = m^0 \left[\frac{\alpha_s}{\pi\epsilon} + \ldots \right] + \Sigma^{\rm fin}(m^0, m^0, \mu)$$

$\overline{\rm MS}$ scheme: $$m^0 = \overline{m}(\mu) \left[1 - \frac{\alpha_s}{\pi\epsilon} + \ldots \right]$$

Like running "strong coupling"

→ $\overline{m}(\mu)$ is pure UV-object without IR-sensitivity

→ Useful scheme for $\mu > m$

→ Far away from a kinematic mass of the quark

Pole scheme: $$m^0 = m^{\rm pole} \left[1 - \frac{\alpha_s}{\pi\epsilon} + \ldots \right] - \Sigma^{\rm fin}(m^{\rm pole}, m^{\rm pole}, \mu)$$

→ Absorbes all self energy corrections into the mass parameter

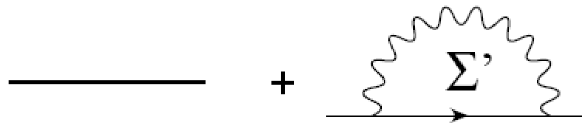
→ Close to the notion of the quark rest mass (kinematic mass)

→ Renormalon problem: infrared-sensitive contributions from < 1 GeV that cancel between self-energy and all other diagrams cannot cancel.

→ Has perturbative instabilities due to sensitivity to momenta < 1 GeV (Λ_{QCD})

Should not be used if uncertainties are below 1 GeV !

Heavy Quark Mass



$$= p - m^0 - \Sigma(p, m^0, \mu)$$

$$\Sigma(m^0, m^0, \mu) = m^0 \left[\frac{\alpha_s}{\pi\epsilon} + \ldots \right] + \Sigma^{\rm fin}(m^0, m^0, \mu)$$

$\overline{\rm MS}$ scheme: $$m^0 = \overline{m}(\mu) \left[1 - \frac{\alpha_s}{\pi\epsilon} + \ldots \right]$$

Pole scheme: $$m^0 = m^{\rm pole} \left[1 - \frac{\alpha_s}{\pi\epsilon} + \ldots \right] - \Sigma^{\rm fin}(m^{\rm pole}, m^{\rm pole}, \mu)$$

Jain, Lepenik, Mateu, Preisser, Scimemi, Stewart, AHH arXiv:1704.01580

Jain, AH, Scimemi, Stewart (2008)

MSR scheme: $$m^{\rm MSR}(R) = m^{\rm pole} - \Sigma^{\rm fin}(R, R, \mu)$$

→ Like pole mass, but self-energy correction from <R are not absorbed into mass

→ Interpolates between MSbar and pole mass scheme

$$m_t^{\rm MSR}(R = 0) = m^{\rm pole}$$

$$m_t^{\rm MSR}(R = \overline{m}(\overline{m})) = \overline{m}(\overline{m})$$

→ More stable in perturbation theory.

→ $m_t^{\rm MSR}(R = 1\,{\rm GeV})$ close to the notion of a kinematic mass, but without renormalon problem.

MC Top Quark Mass (for reconstruction)

Stewart, AHH, 2008 AHH, 2014

$$m_t^{\rm MC} = m_t^{\rm MSR}(R = 1\ {\rm GeV}) + \Delta_{t,{\rm MC}}(R = 1\ {\rm GeV})$$

$$\Delta_{t,{\rm MC}}(1\ {\rm GeV}) \sim \mathcal{O}(1\ {\rm GeV})$$

- small size of $\Delta_{t,MC}$
- Renormalon-free
- little parametric dependence on other parameters

MSR Mass Definition

Jain, Lepenik, Mateu, Preisser, Scimemi, Stewart, AHH
arXiv:1704.01580

MS Scheme: $(\mu > \overline{m}(\overline{m}))$

$$\overline{m}(\overline{m}) - m^{\rm pole} = -\overline{m}(\overline{m})\,[0.42441\,\alpha_s(\overline{m}) + 0.8345\,\alpha_s^2(\overline{m}) + 2.368\,\alpha_s^3(\overline{m}) + \ldots]$$

MSR Scheme: $(R < \overline{m}(\overline{m}))$

$$m_{\rm MSR}(R) - m^{\rm pole} = -R\,[0.42441\,\alpha_s(R) + 0.8345\,\alpha_s^2(R) + 2.368\,\alpha_s^3(R) + \ldots]$$

$$m_{\rm MSR}(m_{\rm MSR}) = \overline{m}(\overline{m})$$

⇒ $m_{\rm MSR}(R)$ Short-distance mass that smoothly interpolates all R scales

≃ "pole mass subtraction for momentum scales larger than R"

- Precision in relation to any other short-distance mass: $\lesssim$ 20 MeV @ $O(\alpha_S^4)$

"Ultimate" Precision of the Pole Mass

Answer: Δm_t^{pole} = 70 MeV

Beneke, Marquard, Nason, Steinhauser
arXiv:1605.03609

$$m_t^{\rm pole} - m_t^{\rm nMSR}(163\,{\rm GeV}) = 7.505 + 1.581 + 0.481 + 0.193$$
$$+ 0.111 + 0.079 + 0.066 + 0.064 + 0.071 + \ldots {\rm GeV}$$

region of constant correction

Size of minimal term: 64 MeV

BUT:

Lepenik, Preisser, AHH, to appear

$$m_t^{\rm nMSR}(2\,{\rm GeV}) - m_t^{\rm nMSR}(163\,{\rm GeV}) = 9.838 + 0.623 + 0.072 - 0.026$$
$$\pm 0.025\ \ {\rm GeV}$$
$$m_t^{\rm pole} - m_t^{\rm nMSR}(2\,{\rm GeV}) = 0.246 + 0.139 + 0.113 + 0.122$$
$$+ 0.187 + 0.353 + 0.791 + 2.053 + 6.053 + \ldots {\rm GeV}$$

Size of minimal term: 122 MeV

- Approach-independent ambiguity estimate
- Bottom and charm mass effects: heavy quark symmetry (Δm_t^{pole} = Δm_c^{pole}) (increases IR sensitivity)

Calibration of the MC Top Mass

Method:

✓ 1) Strongly mass-sensitive hadron level observable (as closely as possible related to reconstructed invariant mass distribution !)

✓ 2) Accurate hadron level QCD predictions at ≧ NLL/NLO with full control over the quark mass scheme dependence.

✓ 3) QCD masses as function of m_t^{MC} from fits of observable.

4) Cross check observable independence / universality

$$m_t^{\rm MC} = m_t^{\rm MSR}(R = 1\ {\rm GeV}) + \Delta_{t,{\rm MC}}(R = 1\ {\rm GeV})$$

$$\Delta_{t,{\rm MC}}(1\ {\rm GeV}) = \bar{\Delta} + \delta\Delta_{\rm MC} + \delta\Delta_{\rm pQCD} + \delta\Delta_{\rm param}$$

Experimental systematics

Monte Carlo dependence:

- different tunings
- parton showers
- color reconnection
- Intrinsic error, …

QCD errors:

- perturbative error
- scale uncertainties
- electroweak effects

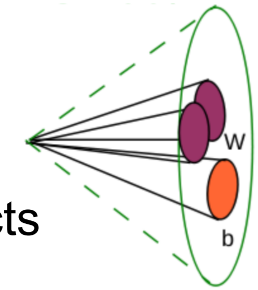
Parametric errors:

- strong coupling α_s
- Non-perturbative parameters

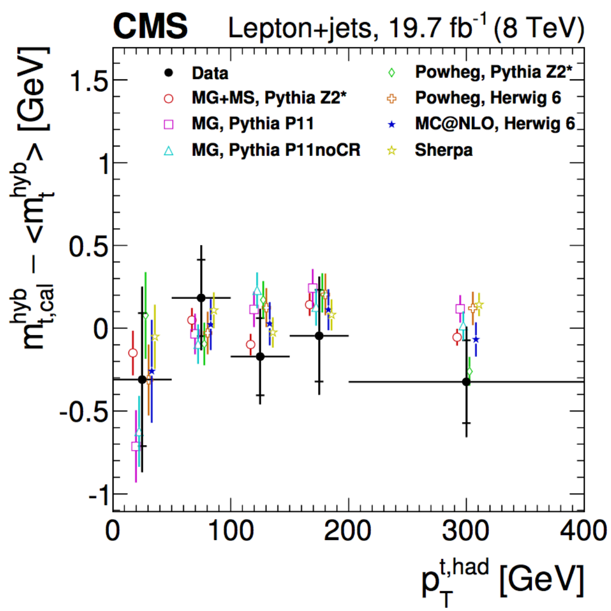
Treated in our analysis

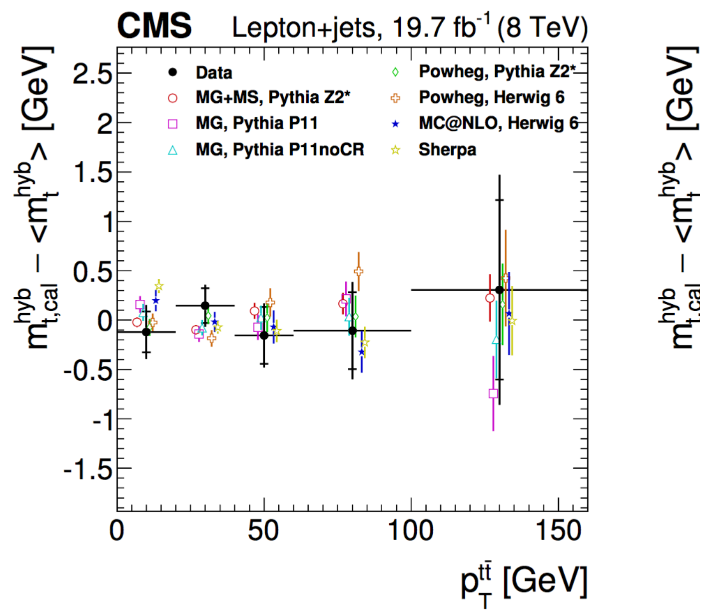
Boosted Top Quarks

First simplification: $Q = 2p_T \ll m_t$



- Enables us to be inclusive w.r. to the hard-collinear decay products

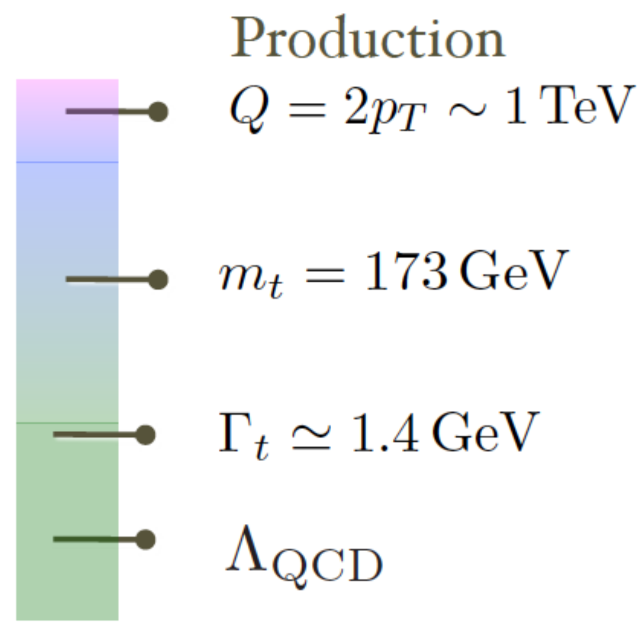




- Top mass from reconstruction of boosted tops consistent with low p_T results.
- More precise studies possible with more statistics from Run-2.
- Meaning of m_t^{MC} for boosted tops and slow top quarks consistent.

Theory Issues for $pp \to t\,\bar{t}\,X$

- jet observable
- suitable top mass for jets
- initial state radiation
- final state radiation
- underlying event
- color reconnection
- beam remnant
- parton distributions
- sum large logs $Q \gg m_t \gg \Gamma_t$



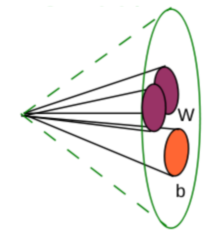
Theory Issues for $pp \to t\bar{t}X$

- jet observable ★★
- suitable top mass for jets ★
- initial state radiation
- final state radiation ★
- underlying event
- color reconnection (★) ⟵ Only final-final state color reconnection
- beam remnant
- parton distributions
- sum large logs $Q \gg m_t \gg \Gamma_t$ ★

First
$e^+e^- \to t\bar{t}X$
and the issues ★

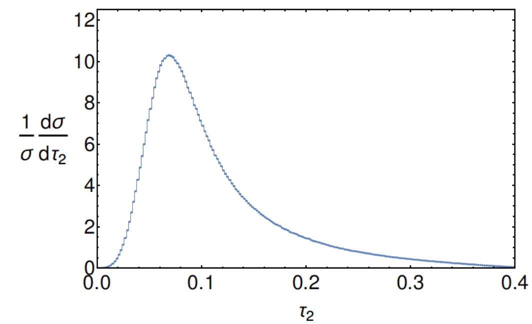
Thrust Distribution

Observable: 2-jettiness in e+e- for $Q = 2p_T \gg m_t$ (boosted tops)



$$\tau \;=\; 1 - \max_{\vec{n}} \frac{\sum_i |\vec{n}\cdot\vec{p}_i|}{Q}$$

$$\tau_2 \overset{\text{peak}}{\underset{\rightarrow}{\approx}} \frac{M_1^2 + M_2^2}{Q^2}$$

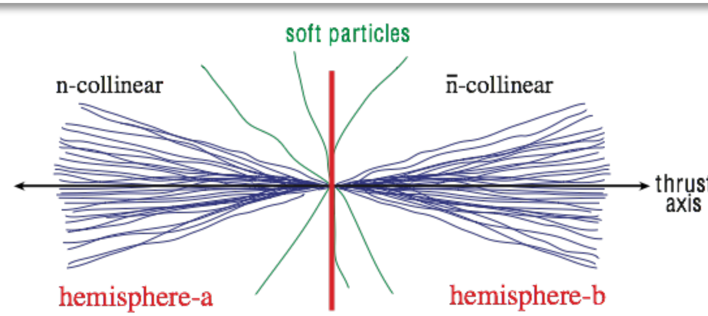


Invariant mass distribution in the resonance region of wide hemisphere jets !

$$\frac{d\sigma}{d\tau} = Q^2 \sigma_0 H_0(Q,\mu) \int d\ell \; J_0(Q\ell,\mu)\, S_0\left(Q\tau - \ell, \mu\right)$$

Excellent mass sensitivity:

$$\tau_2^{\text{peak}} = 1 - \sqrt{1 - \frac{4m_t^2}{Q^2}}$$ (tree level)



universität wien

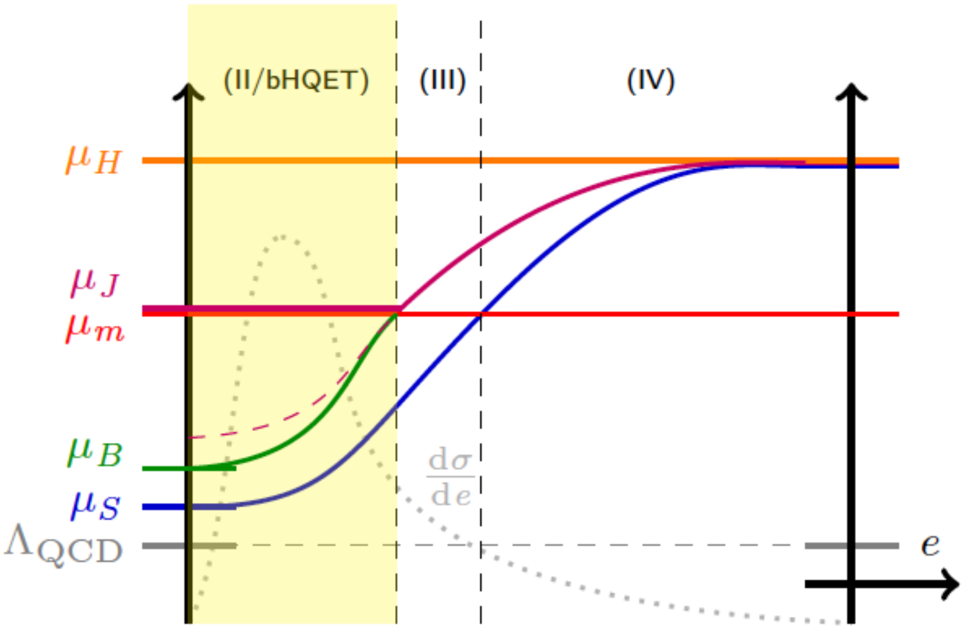
Factorization: EFT Treatment

- Boosted top jets

[Fleming, Hoang, Mantry, Stewart 2007]

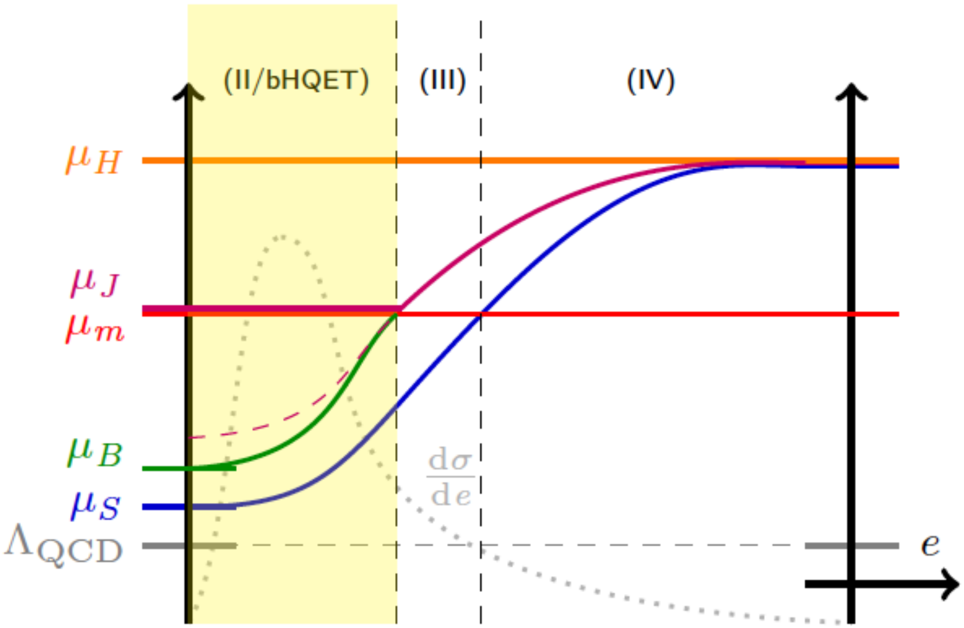
$$n_f = n_\ell + 1$$

$$\frac{d\sigma^{\mathrm{bHQET}}}{d\tau} = Q\, H(Q, m, \mu_H)\, U_H^{(n_f)}(Q, \mu_H, \mu_m)\, H_m^{(n_f)}(Q, \mu_m)\, U_m^{(n_l)}(Q, m, \mu_m, \mu_B)$$
$$\times \int ds\, d\ell\, B_e^{(n_l)}(s, m, \mu_B)\, U_S^{(n_l)}(\ell, \mu_B, \mu_S)\, S_e^{(n_l)}\Big(Q(\tau - \tau_{\min}) - \frac{s}{Q} - \ell, \mu_S\Big)$$



Factorization: EFT Treatment

- Developments:
  - VFNS for final state jets (with massive quarks)
    [Gritschacher, Hoang, Jemos, Mateu, Pietrulewicz '13 '14]
    [Butenschön, Dehnadi, Hoang, Mateu '16 (to appear)]



- Non-perturbative power-corrections are included via a shape function
  [Korchemsky, Sterman 1999]
  [Hoang, Stewart 2007]
  [Ligeti, Stewart, Tackmann 2008]

$$\frac{d\sigma}{d\tau} = \frac{d\sigma^{\rm part}}{d\tau} \otimes F_{\rm mod}(\Omega_1, \Omega_2, \ldots)$$

- Gap-scheme
- MSR mass & R-evolution
  [Hoang, Jain, Scimemi, Stewart 2010]

Jain, Lepenik, Mateu, Preisser, Scimemi, Stewart, AHH :1704.01580

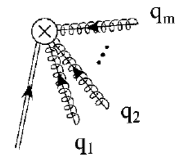
➡ NNLL + NLO + non-singular
+ hadronization
+ renormalon-subtraction
+ top quark decay

- Good convergence
- Reduction of scale variation (NLL vs. NNLL)

Why the Observed Pole is not at the Pole Mass

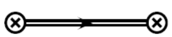
Jet function: $$B_+(2v_+\cdot k) = \frac{-1}{8\pi N_c m}\,\mathrm{Disc}\int d^4x\, e^{ik\cdot x}\langle 0|\mathrm{T}\{\bar h_{v_+}(0)W_n(0)W_n^\dagger(x)h_{v_+}(x)\}|0\rangle$$

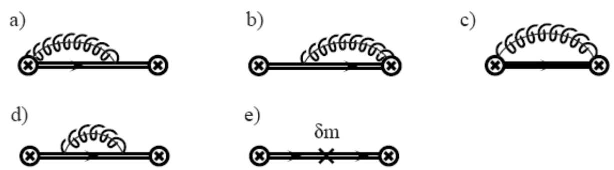
- perturbative, any mass scheme
- depends on m_t, Γ_t
- Breit-Wigner at tree level
- Gauge-invariant off-shell top quark dynamics



$$W = \sum_{m=0}^{\infty}\sum_{\text{perms}}\frac{(-g)^m}{m!}\frac{\bar n\cdot A^{a_1}_{n,q_1}\cdots\bar n\cdot A^{a_m}_{n,q_m}}{\bar n\cdot q_1\,\bar n\cdot(q_1+q_2)\cdots\bar n\cdot(\sum_{i=1}^m q_i)}T^{a_m}\cdots T^{a_1}$$

Singular functions encode information about where the physical pole is located





$$\hat s = \frac{M^2 - m_t^2}{m_t}$$

$$B_\pm^{\Gamma=0}(\hat s,\mu,\delta m) = \delta(s) + \frac{\alpha_s(\mu)C_F}{\pi}\left\{\frac{2}{m\mu}\left[\frac{\theta(z)\ln(z)}{z}\right]_+ - \frac{1}{m\mu}\left[\frac{\theta(z)}{z}\right]_+ + \delta(s)\left[1-\frac{\pi^2}{8}\right]\right\} - \frac{2\,\delta m}{m}\delta'(\hat s)$$

$$\mathcal{B}_\pm(\hat s,0,\mu,\delta m) = -\frac{1}{\pi m}\frac{1}{\hat s+i0}\left\{1+\frac{\alpha_s C_F}{4\pi}\left[4\ln^2\left(\frac{\mu}{-\hat s-i0}\right)+4\ln\left(\frac{\mu}{-\hat s-i0}\right)+4+\frac{5\pi^2}{6}\right]\right\} - \frac{1}{\pi m}\frac{2\delta m}{(\hat s+i0)^2}$$

Fleming, AHH, Mantry, Stewart 2007

Why the Observed Pole is not at the Pole Mass

Is the pole mass determining the top single particle pole?

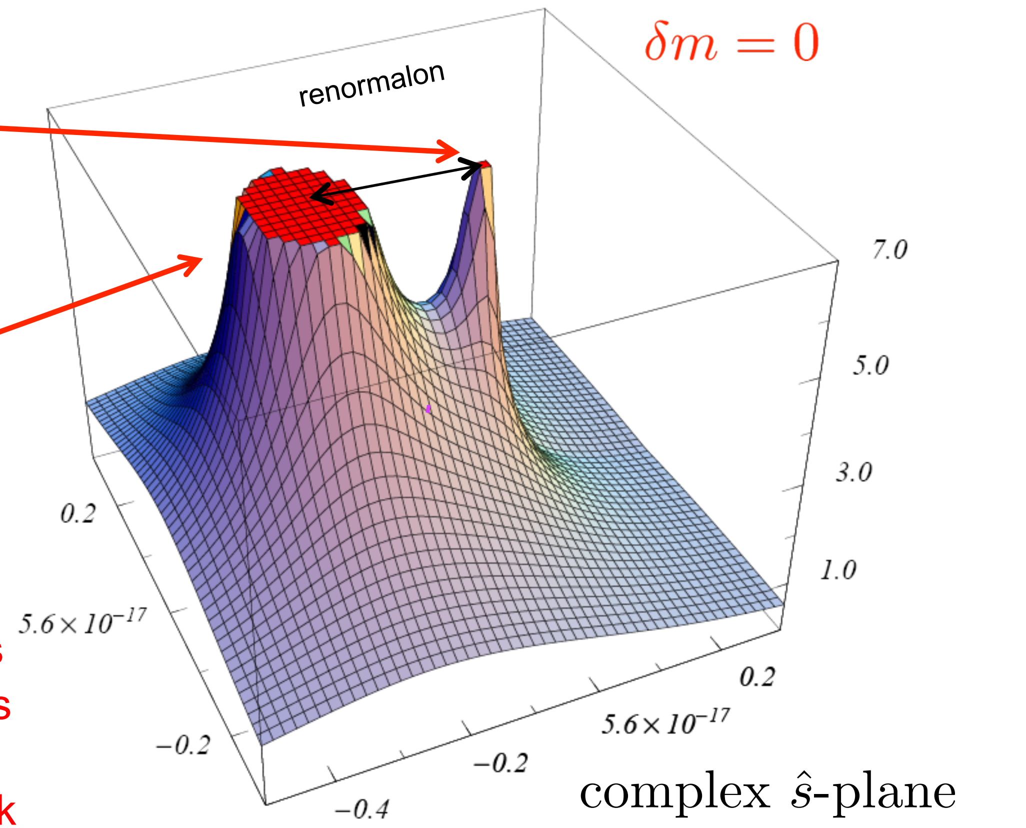
NO !

$$\hat{s} = \frac{M_t^2 - m_t^2}{m_t}$$

pole mass peak

Invisible for $\Gamma_t > 0.5$ GeV

observable peak



→ pole mass and observable peak separated by non-converging series

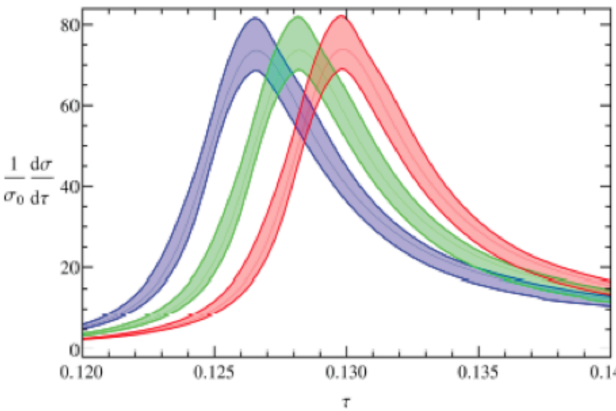
→ pole mass peak (residue) decreases with order

→ MSR mass close to observable peak

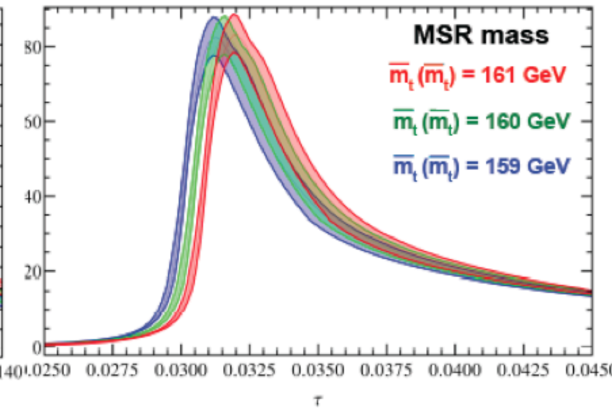
2-Jettiness for Top Production (QCD)

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\tau_2} = f(\underbrace{m_t^{\mathrm{MSR}}(R)}_{\text{any scheme possible}}, \alpha_s(M_Z), \underbrace{\Omega_1, \Omega_2, \ldots}_{\text{Non-perturbative}}, \underbrace{\mu_h, \mu_j, \mu_s, \mu_m}_{\text{renorm. scales}}, R, \underbrace{\Gamma_t}_{\text{finite lifetime}})$$

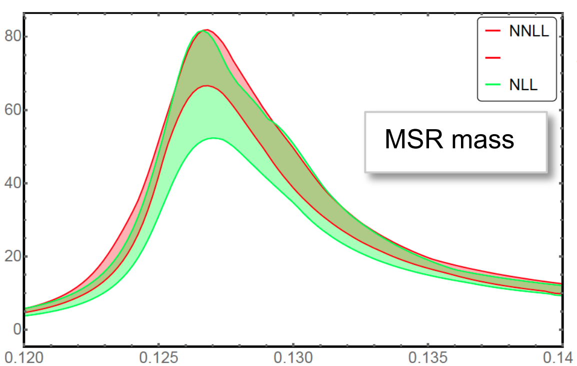
Q=700 GeV (p_T= 350 GeV)



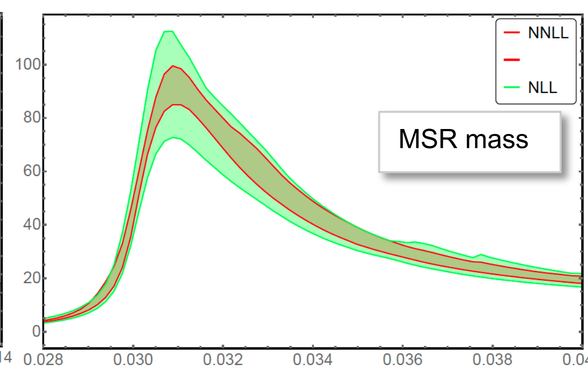
Q=1400 GeV (p_T = 700 GeV)



Q=700 GeV



Q=1400 GeV



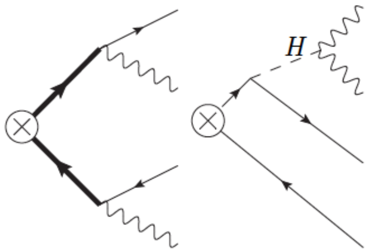
- Higher mass sensitivity for lower Q (p_T)
- Finite lifetime effects included
- Dependence on non-perturbative parameters
- Convergence: $\Omega_{1,2,\ldots}$

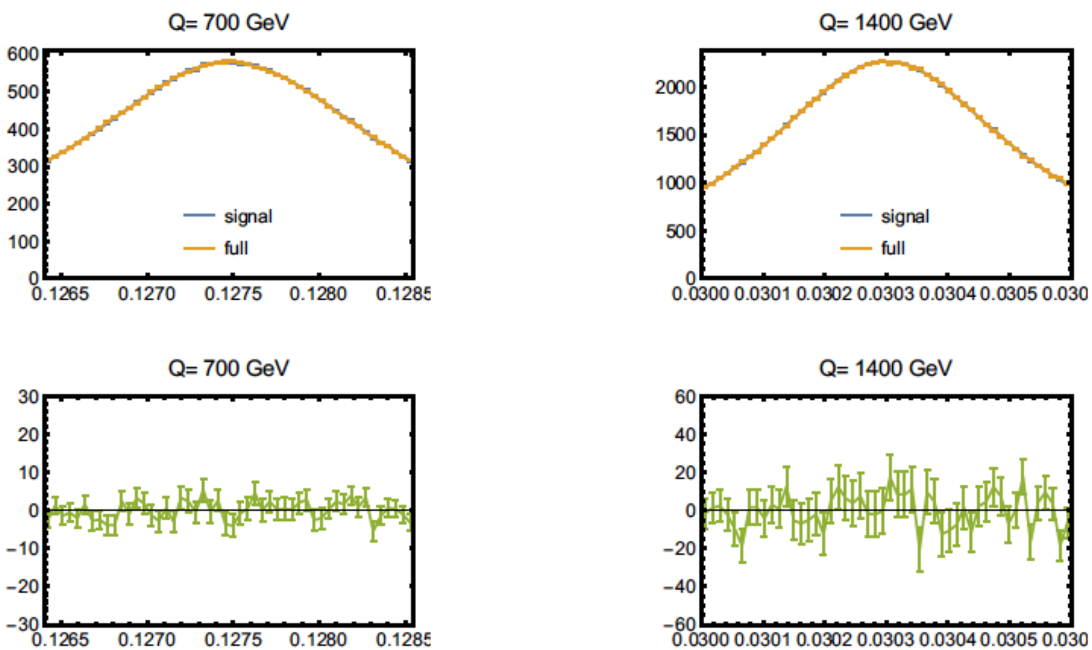
- Good convergence
- Reduction of scale uncertainty (NLL to NNLL)
- Control over whole distribution

Signal ttbar vs full ee→WWbb

MadGraph 5 study:

- Non-resonant contributions are irrelevant for τ_2 distribution
  - PYTHIA (or similar MCs) will give a good description of the production process at LO
  - hemisphere invariant mass $\sim$ top invariant mass (no pollution from background)

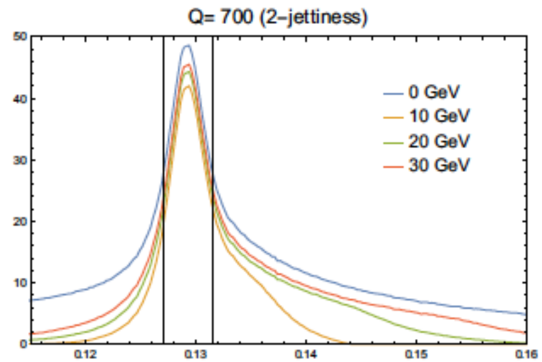


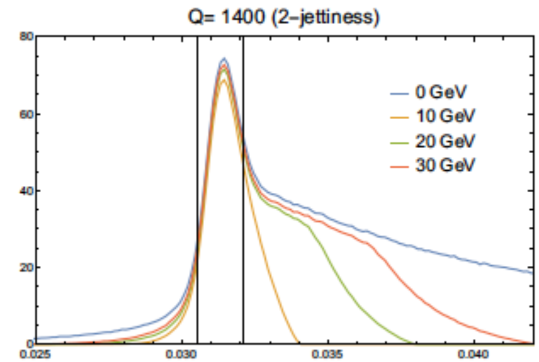


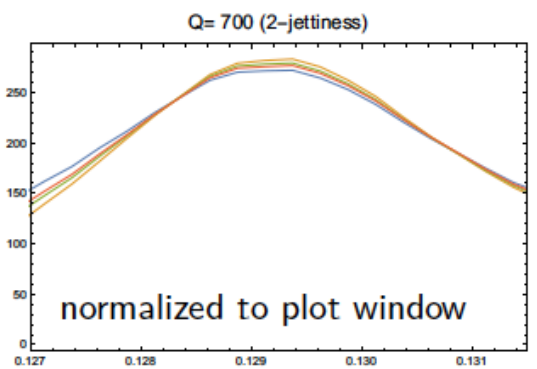
Pythia Study: Hemisphere Mass Cuts

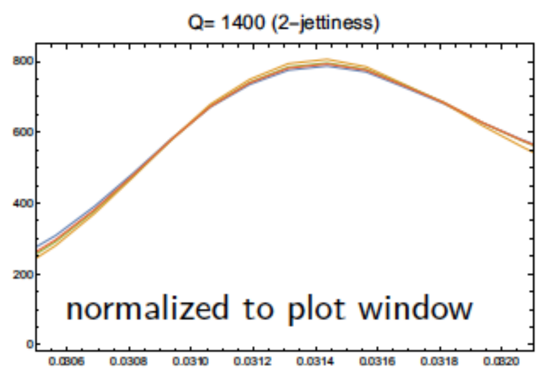
- In our theory description we treat the top decay as inclusive w.r.t. hemisphere
  - violated by decay products which cross to the other hemisphere
  - no differential impact in resonance region (irrelevant when normalized to signal region)

Cuts on hemisphere invariant mass above and below:
$M_i^{\mathrm{cut}} = m_t^{\mathrm{MC}} \pm \Delta^{\mathrm{cut}}$









Fit Procedure Details

- $\frac{d\sigma}{d\tau} = f(m_t^{\mathrm{MSR}}, \alpha_s(m_Z), \Omega_1, \Omega_2, ..., \mu_H, \mu_J, \mu_S, \mu_M, R, \Gamma_t)$

  any scheme   non-perturbative   renorm. scales   finite lifetime

- Generating PYTHIA Samples: (PYTHIA 8.205)

  at different energies: $Q = 600, 700, 800, ..., 1400$ GeV

  - masses: $m_t^{\mathrm{MC}} = 170, 171, 172, 173, 174, 175$ GeV
  - width: $\Gamma_t = 1.4$ GeV
  - Statistics: 10^7 events for each set of parameters

  - Tune 7 (Monash)

- Feed MC data into **Fitting Procedure:** all ingredients are there

  Fit parameters: m_t^{MSR}, $\alpha_s(m_Z)$, $\Omega_1, \Omega_2, \ldots$

  - Take $\alpha_S(M_Z)$ as input from world average. (Sensitivity to strong coupling very weak.)

  - standard fit based on χ^2 minimization
  - analysis with 500 sets of profiles (τ_2 dependent renorm. scales) for the each MC sample
  - different Q-sets: 7 sets with energies between 600 - 1400 GeV

    different n-sets: 3 choices for fitranges - (xx/yy)% of maximum peak height

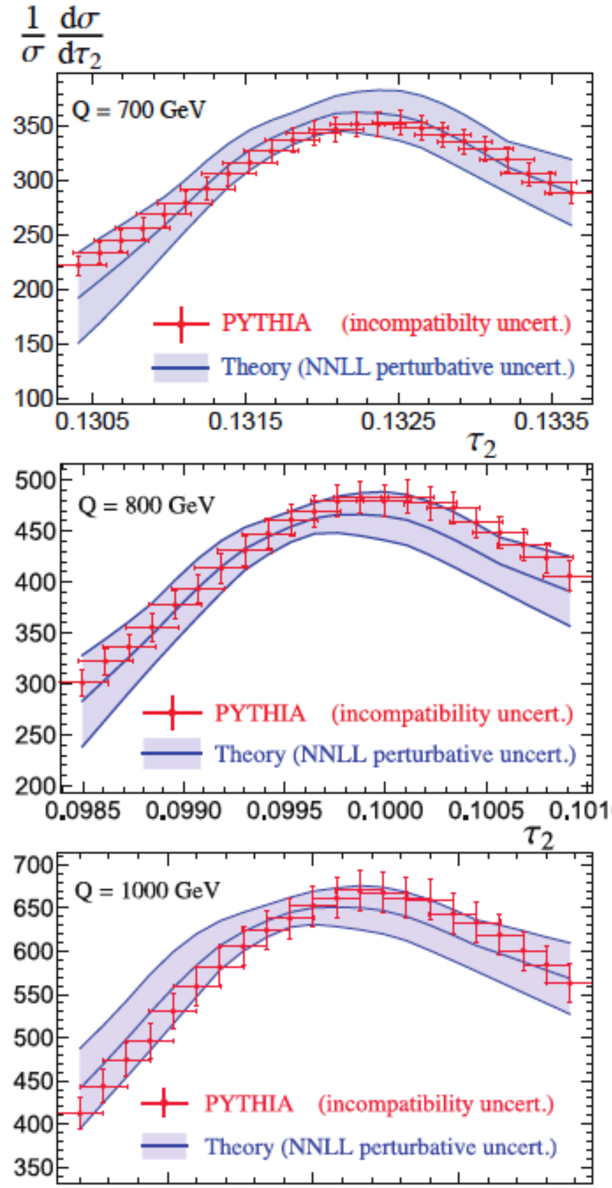
    } 21 fit setups

Fit Result: Pythia 8.205 vs. Theory

Γ_t=1.4 GeV, tune 7,
m_t^{MC} = 173 GeV

Ω_1 = 0.44 GeV,
m_t^{MSR}(1GeV) = 172.81 GeV

- Good agreement of PYTHIA with NNLL/NLO theory predictions
- Perturbative uncertainties of theory predictions based on scale uncertainties (profiles)
- MC uncertainties:
  - Vertical: rescaled statistical error (PDF rescaling method) → independent on statistics
  - Horizontal: fit coverage from 21 fit setups (incompatiblity uncertainty)



Convergence & Stability: MSR vs. Pole Mass

500 profiles; $\alpha_s = .118$; $\Gamma_t = 1.4$ GeV; tune 7;
$Q = 700, 1000, 1400$ GeV; peak$(60/80)\%$

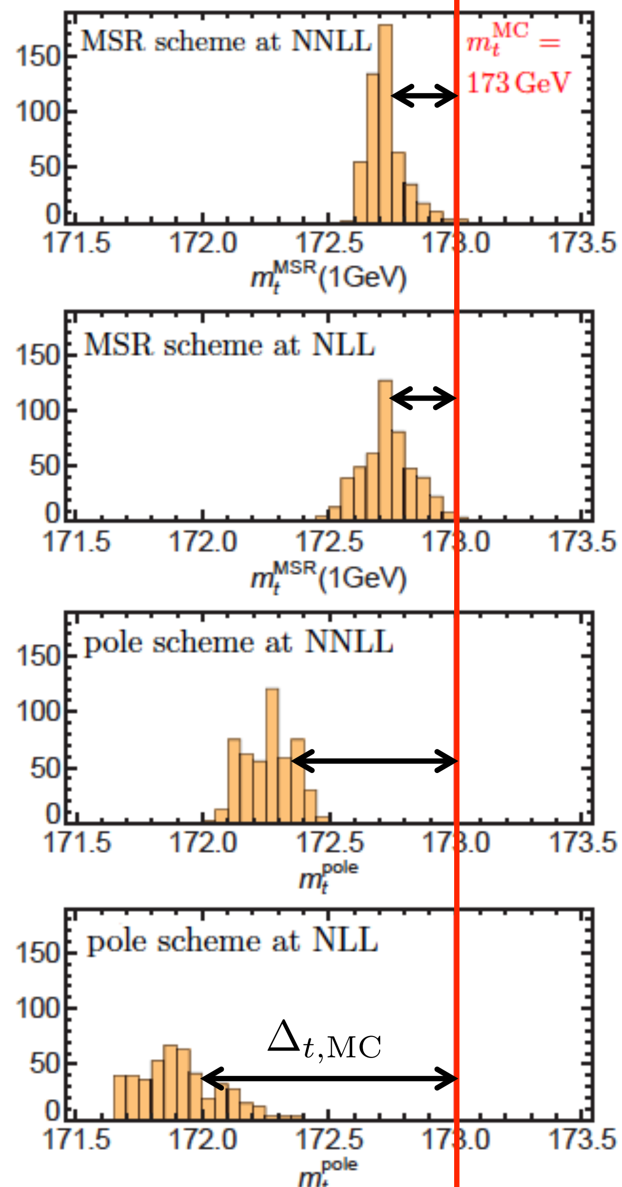
Input: $m_t^{\mathrm{MC}} = 173$ GeV

fit to find $m_t^{\mathrm{MSR}}(1\mathrm{GeV})$ or m_t^{pole}

- Good convergence & stability for MSR mass
- Mass m_t^{MSR}(1GeV) mass definition closest to the MC top mass m_t^{MC}.

- Pole mass shows worse convergence.
- Pole mass not compatible with MC mass within errors
- 1100/700 MeV difference at NLL/NNLL
- $m_t^{pole} \neq M_t^{Pythia\ 8.2}$

Similar analyses from the 20 other Q-set and n-range setups.

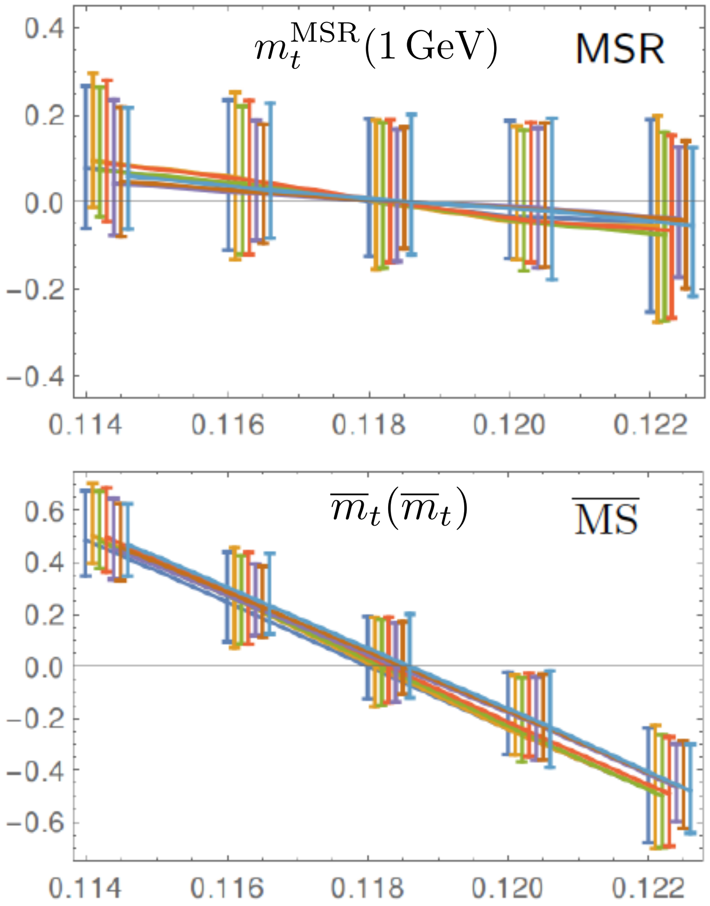


MSR/MS Parametric Dependence on α_s

500 profiles; $\Gamma_t = 1.4$,-1 GeV;tune 7;
diff. Q-sets; peak$(60/80)\%$

$m_t^{\texttt{PYTHIA}} = 173$ GeV

- α_s dependence:
  $m^{\text{scheme}}[\alpha_s] - m^{\text{scheme}}[.118]$
- small dependence of MSR mass on α_s
  ~ 50 MeV error ($\delta\alpha_s = .002$)
- large sensitivity of $\overline{\text{MS}}$ mass on α_s
- not an error:
  calculated from MSR

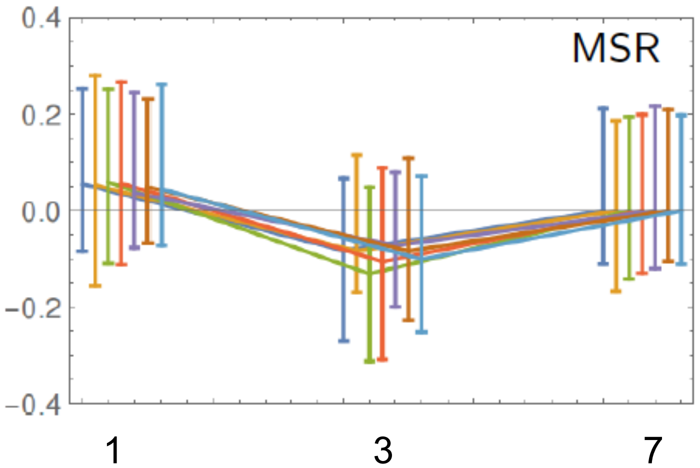


MSR Mass Tune Dependence

500 profiles; $\Gamma_t = 1.4$,-1 GeV;tune $1, 3, 7$;
diff. Q-sets; peak$(60/80)\%$

$$m_t^{\mathtt{PYTHIA}} = 173 \text{ GeV}$$

- tune dependence:
  $m^{\rm MSR}[\rm tune] - m^{\rm MSR}[7]$



- clear sensitivity to tune
- $m^{\rm MC}$ will depend on tune
- tune dependence is not a calibration uncertainty:

  (different tune $\Rightarrow$ different MC $\Rightarrow m_t^{\rm MC}$)

Final Result for m_t^{MSR}(1 GeV)

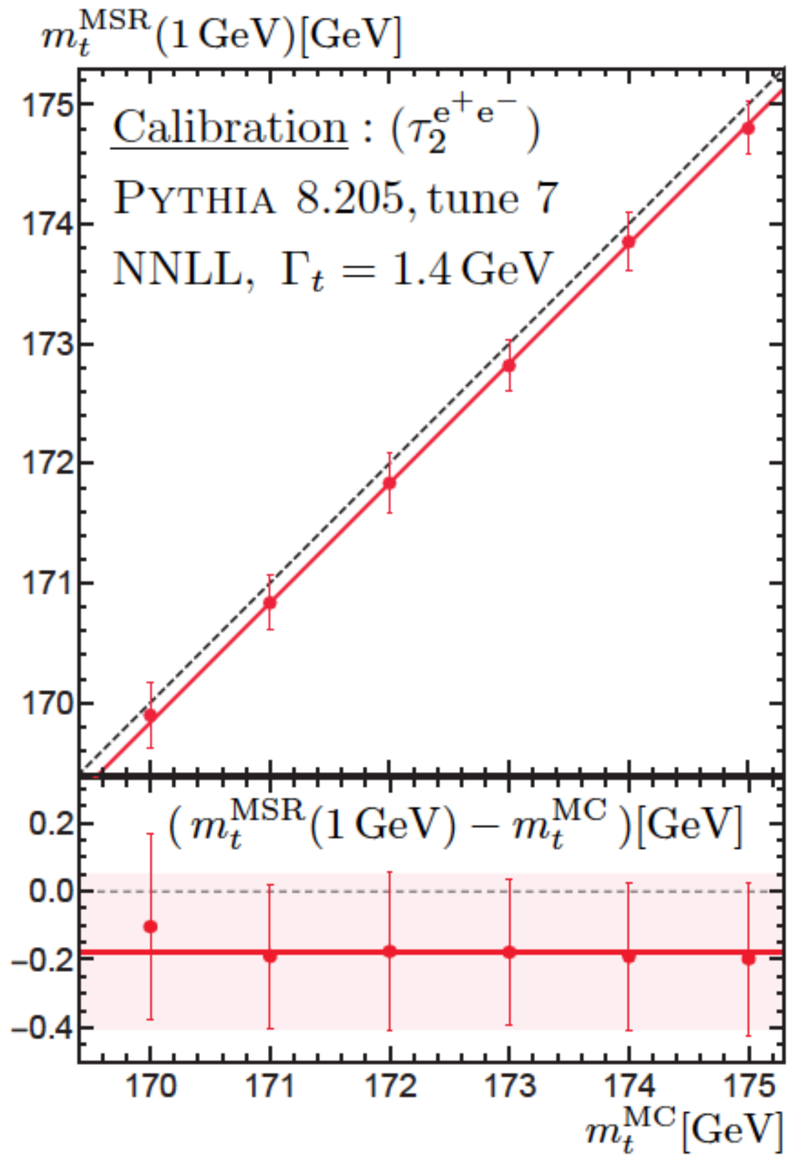
- All investigated MC top mass values show consistent picture
- MC top quark mass is indeed closely related to MSR mass

  within uncertainties:
  $m_t^{MC} \simeq m_t^{MSR}(1\text{GeV})$

$m_t^{MC} = 173\,\text{GeV} \quad (\tau_2^{e^+e^-})$

mass	order	central	perturb.	incompatibility	total
$m_{t,1\,\text{GeV}}^{MSR}$	NLL	172.80	0.26	0.14	0.29
$m_{t,1\,\text{GeV}}^{MSR}$	N^2LL	172.82	0.19	0.11	0.22
m_t^{pole}	NLL	172.10	0.34	0.16	0.38
m_t^{pole}	N^2LL	172.43	0.18	0.22	0.28

↓ Spread of results from 21 fit setups



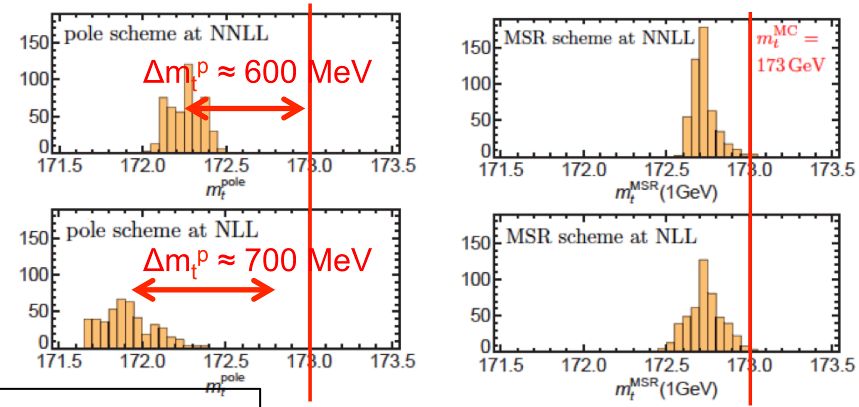
$\Omega_1^{PY} = 0.41 \pm 0.07 \pm 0.02\,\text{GeV}$ at NLL
$\Omega_1^{PY} = 0.42 \pm 0.07 \pm 0.03\,\text{GeV}$ at N^2LL

Pole Mass Determination

1) Pole mass implemented in code:

$m_t^{MC} = 173$ GeV ($\tau_2^{e^+e^-}$)

Mass	Order	Central	Perturb.	Incompatibility	Total
$m_{t,1\,\mathrm{GeV}}^{\mathrm{MSR}}$	NLL	172.80	0.26	0.14	0.29
$m_{t,1\,\mathrm{GeV}}^{\mathrm{MSR}}$	NNLL	172.82	0.19	0.11	0.22
m_t^{pole}	NLL	172.10	0.34	0.16	0.38
m_t^{pole}	NNLL	172.43	0.18	0.22	0.28



$$(m_t^{\mathrm{pole}})_{\mathrm{NLL/NNLL}} < m_t^{\mathrm{MSR}}(1\,\mathrm{GeV}) < M_t^{\mathrm{MC}}$$

2) Pole mass determined from MSR mass:

$$\alpha_s(M_Z) = 0.118 \qquad n_f = 5$$

$$m_t^{\mathrm{pole}} - m_t^{\mathrm{MSR}}(1\,\mathrm{GeV}) = \underbrace{0.173}_{\mathcal{O}(\alpha_s)} + \underbrace{0.138}_{\mathcal{O}(\alpha_s^2)} + \underbrace{0.159}_{\mathcal{O}(\alpha_s^3)} + \underbrace{0.23}_{\mathcal{O}(\alpha_s^4)}\ \mathrm{GeV}$$ ← calculated

$$+\ 0.53 + 1.43 + 4.54 + 16.6\ \mathrm{GeV}$$
$$+\ 68.6 + 317.7 + 1629 + 9158\ \mathrm{GeV}$$ ← extrapolated

$$m_t^{\mathrm{pole}} > m_t^{\mathrm{MSR}}(1\,\mathrm{GeV})$$

- Calibration in terms of the pole mass involves large higher-order perturbative corrections
- Additional uncertainty on pole mass: $(m_t^{\mathrm{pole}})_{\mathrm{NLL}}$ = 172.45 ± 0.52 GeV,
  (added quadratically) $(m_t^{\mathrm{pole}})_{\mathrm{NNLL}}$ = 172.72 ± 0.41 GeV

Result for the Top Msbar Mass

1) Approach: $m_t^{MC} \sim m_t^{MSR}(1\text{ GeV})$

No detailed analysis!

No electroweak corrections!

$$m_t^{\rm pMSR}(1\,{\rm GeV}) \;=\; 172.82 \pm 0.022,\ {\rm GeV}$$

$$m_t^{\rm pMSR}(1\,{\rm GeV}) - m_t^{\rm pMSR}(163.018\,{\rm GeV}) \;=$$

$$= 8.913 + 0.906 + 0.052 - 0.070 \pm 0.035\ {\rm GeV}$$

$$= 9.802 \pm 0.035\ {\rm GeV}$$

$$\overline{m}_t(\overline{m}_t) = 163.020 \pm 0.230\ {\rm GeV}$$

→ Can be improved by next order

2) Approach: $m_t^{MC} \sim m_t^{Pole}$

$$m_t^{\rm pole} = 172.72 \pm 0.410\ {\rm GeV}$$

$$m_t^{\rm pole} - m_t^{\rm nMSR}(163\,{\rm GeV}) \;=\; 7.505 + 1.581 + 0.481 + 0.193$$

$$+\, 0.111 + 0.079 + 0.066 + 0.064 + 0.071 + \ldots {\rm GeV}$$

$$\overline{m}_t(\overline{m}_t)_{\rm 2loop} = 163.634 \pm 0.890\ {\rm GeV}$$

$$\overline{m}_t(\overline{m}_t)_{\rm 3loop} = 163.153 \pm 0.475\ {\rm GeV}$$

$$\overline{m}_t(\overline{m}_t)_{\rm 4loop} = 162.960 \pm 0.430\ {\rm GeV} \qquad \overline{m}_t(\overline{m}_t)_{\rm 8loop} = 162.640 \pm 0.430\ {\rm GeV}$$

Theory Issues for $pp \to t\bar{t}X$

- jet observable ★★
- suitable top mass for jets ★
- initial state radiation
- final state radiation ★
- underlying event
- color reconnection ★
- beam remnant
- parton distributions
- sum large logs $Q \gg m_t \gg \Gamma_t$ ★

Can apply this to current measurements if we trust Pythia extrapolation for remaining items

Theory Issues for $pp \to t\,\bar{t}\,X$

- jet observable ★ ★ Jet Mass in Jet of radius R
- suitable top mass for jets ★
- initial state radiation ★
- final state radiation ★

Better: factorization for pp

- underlying event ← Note: no star here
- color reconnection ★
- beam remnant ★ Jet veto
- parton distributions ★ multiple channels
- sum large logs $Q \gg m_t \gg \Gamma_t$ ★

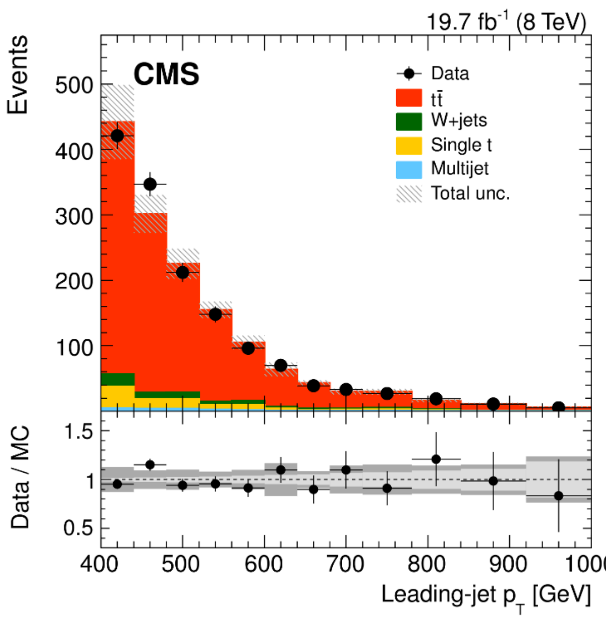
Jet Mass of Boosted Top Quarks

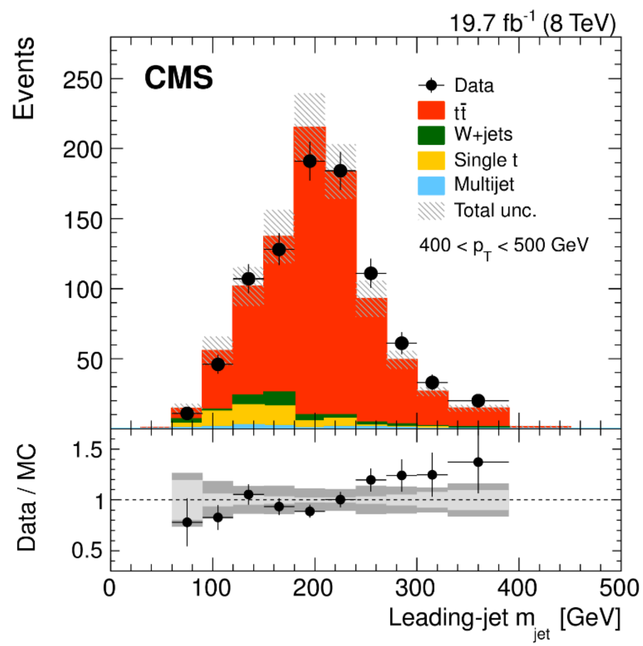
http://arxiv.org/abs/1703.06330

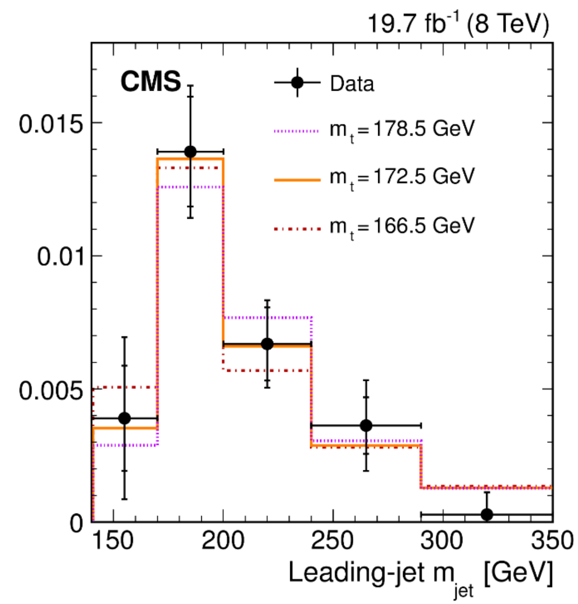
Top mass from boosted jet mass

Cambridge-Aachen jet with distance parameter $R = 1.2$, and $p_T > 400$ GeV.

$m_t = 170.8 \pm 9.0$ GeV







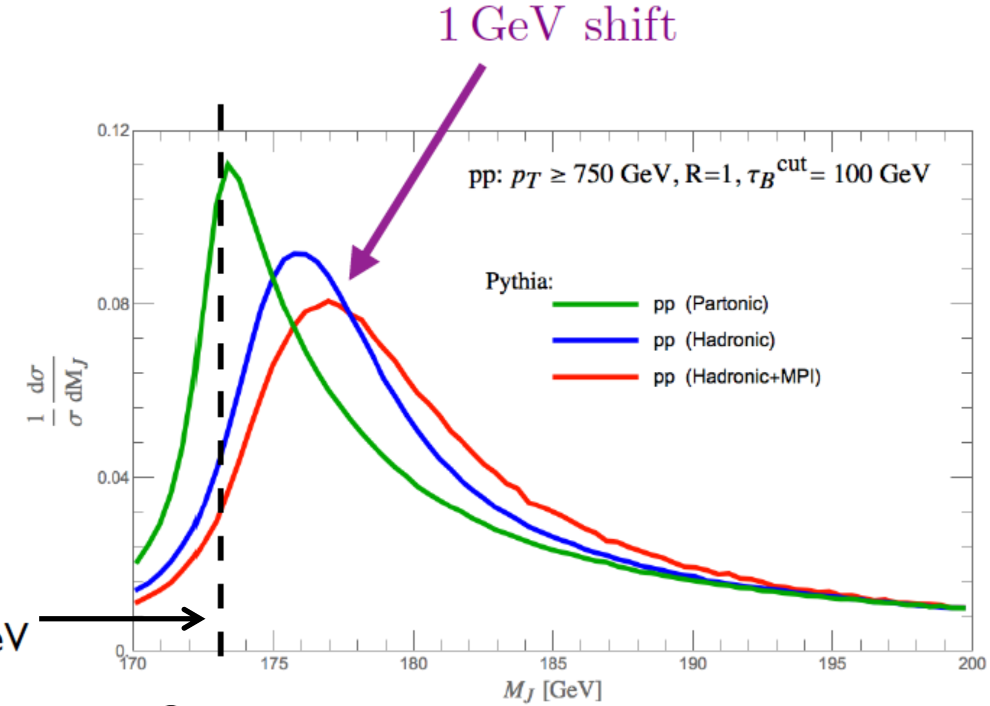
Theory Issues for $pp \to t\bar{t}X$

Extension to pp straightforward: (e.g. N-jettiness & X-Cone jets)

$$\frac{d^2\sigma}{dM_{J1}^2 dM_{J2}^2 d\mathcal{T}^{cut}} = \mathrm{tr}\left[\hat{H}_{Qm}\hat{S}(\mathcal{T}^{cut}, R, \ldots)\otimes F\right]\otimes J_B \otimes J_B \otimes \mathcal{I}\mathcal{I} \otimes f f$$

Same jet functions as e^+e^-

Issue is that UE / MPI is significant:



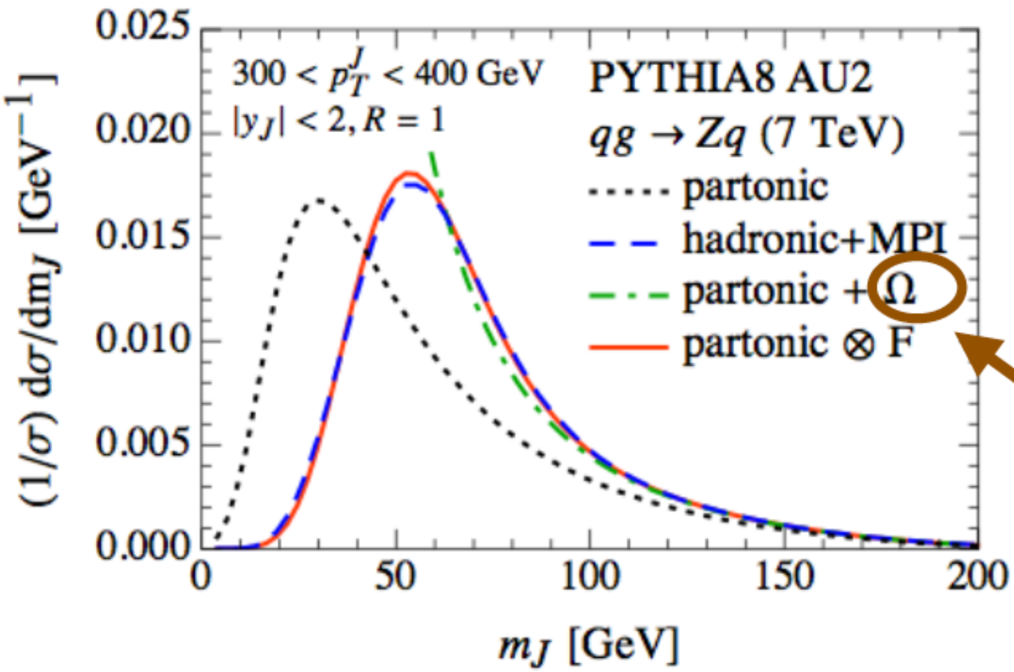
Theory Issues for $pp \to t\,\bar{t}\,X$

Extension to pp straightforward: (e.g. N-jettiness & X-Cone jets)

$$\frac{d^2\sigma}{dM_{J1}^2\, dM_{J2}^2\, d\mathcal{T}^{\rm cut}} = {\rm tr}\big[\hat{H}_{Qm}\hat{S}(\mathcal{T}^{\rm cut}, R, \ldots)\otimes F\big]\otimes J_B \otimes J_B \otimes \mathcal{I}\mathcal{I} \otimes f\,f$$

Same jet functions as e⁺e⁻

Issue is that UE / MPI is significant:



BUT control of Underlying Event is model dependent.

Same model used for Hadronization can describe UE by (primarily) tuning one parameter Ω.

$$\Omega = \int \mathrm{d}k\, k\, F(k)$$

Stewart, Tackmann, Waalewijn, 2015

Grooming with SoftDrop

- Grooms soft radiation from the jet

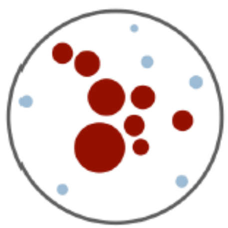
Larkowski, Marzani, Soyez, Thaler, 2014

$$\frac{\min(p_{Ti}, p_{Tj})}{p_{Ti} + p_{Tj}} > z_{\text{cut}} \left(\frac{\Delta R_{ij}}{R_0} \right)^{\beta}$$

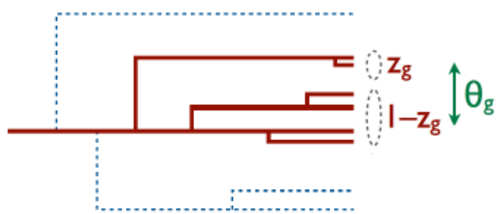
$$z > z_{\text{cut}} \; \theta^{\beta}$$

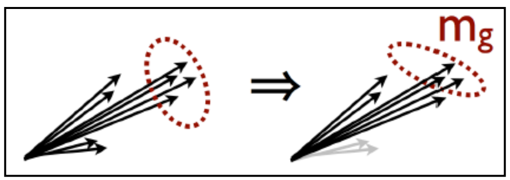
two grooming parameters

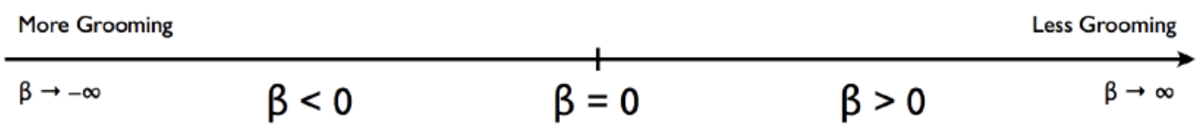
Groomed Jet



Groomed Clustering Tree



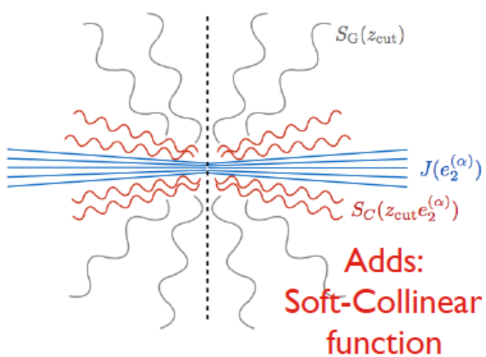




- Allows for factorization calculations

Frye, Larkowski, Schwartz, Yan, 2016

Mode separation: additional soft-collinear modes



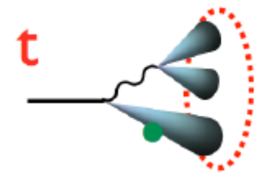
Theory Set Up with SoftDrop

AH, Mantry, Pathak, Stewart; to appear

$$p_T \gg m_t \gg \Gamma_t > \Lambda_{QCD}$$

- Boosted Tops $\quad p_T \gg m_t \quad$ retain top decay products

- Fat Jets $\quad R \gg \dfrac{m_t}{p_T}$

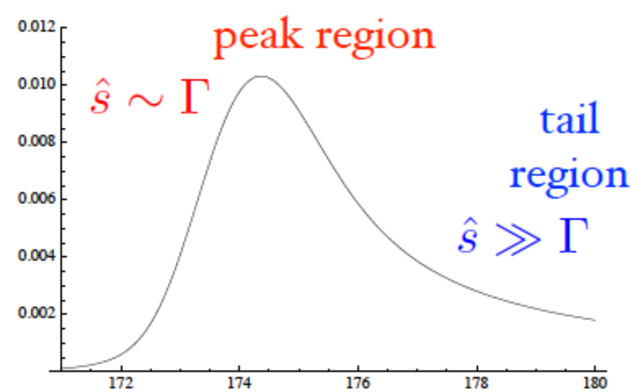


- Sensitivity $\quad \hat{s} \sim \Gamma_t \quad$ for measurement of jet-mass m_J

$$\hat{s} = \frac{m_J^2 - m_t^2}{m_t}$$

- Grooming $\quad z_{\text{cut}}, \beta$

- Jet Veto $\quad \mathcal{T}^{\text{cut}}$ or p_T^{cut}

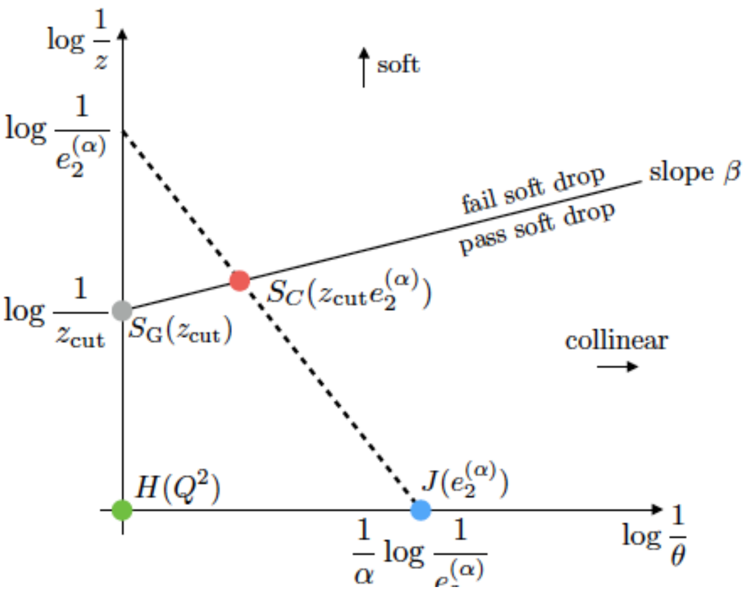


(Perturbative and Nonperturbative effects give $\Gamma > \Gamma_t$)

Theory Set Up with SoftDrop

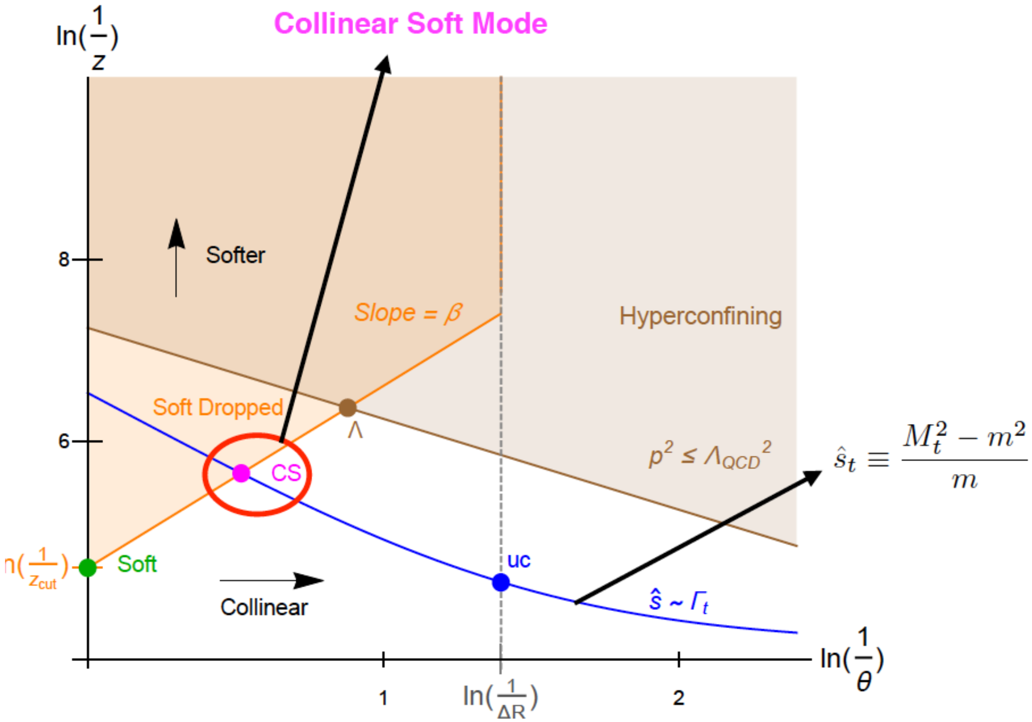
Modes:

massless quarks:

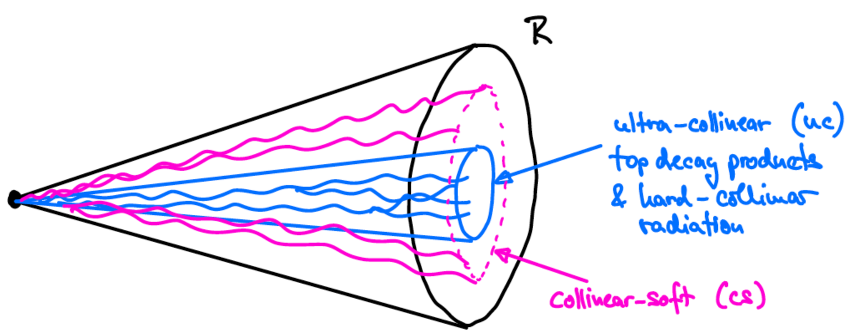


Frye, Larkowski, Schwartz, Yan, 2016

top quarks:



Mantry, Pathak, Stewart, AHH; to appear



Theory Set Up with SoftDrop

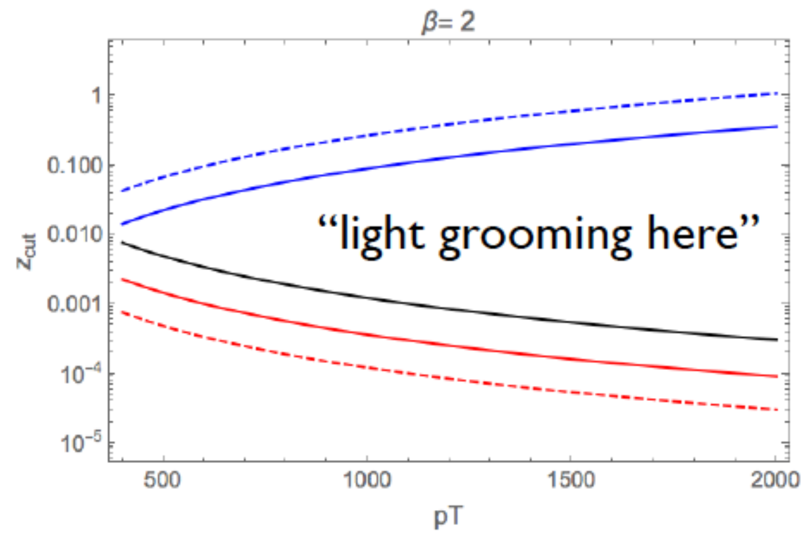
AH, Mantry, Pathak, Stewart; to appear

Can only apply a "light soft drop" for tops:

$$\frac{\Gamma_t}{m}\left(\frac{Q}{2m}\right)^{\beta} \gg z_{\rm cut} \gg \frac{2m\Gamma_t}{Q^2}$$

Ensure soft drop does not touch J_B

Ensure soft drop removes global soft radiation from measurement



Factorization with Soft Drop on one jet:

$$\frac{d^2\sigma}{dM_J^2 d\mathcal{T}^{\rm cut}} = {\rm tr}\big[\hat{H}_{Qm}\hat{S}(\mathcal{T}^{\rm cut}, Qz_{\rm cut}, \beta, \ldots)\otimes F\big]\otimes J_B \otimes \mathcal{I}\mathcal{I} \otimes f\,f$$

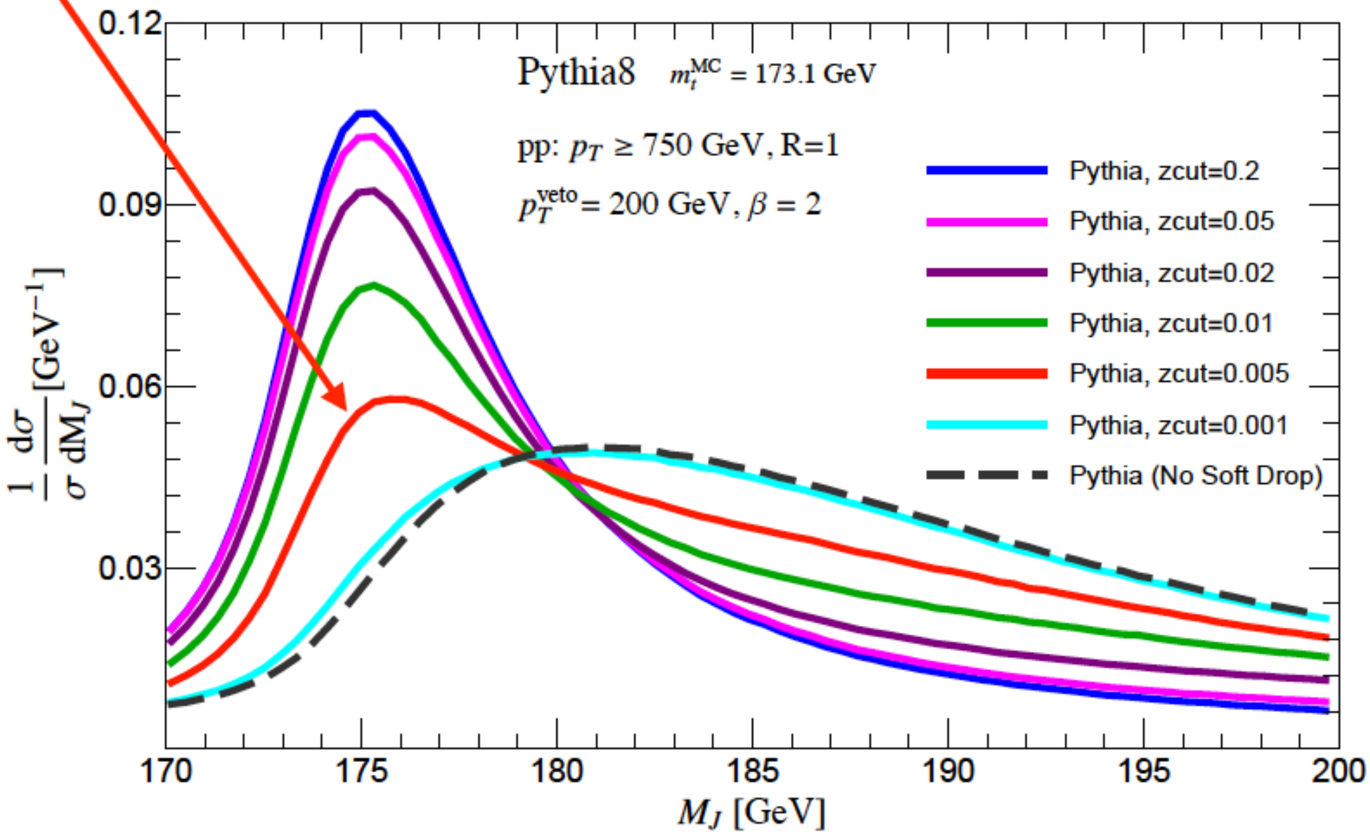
$$\times\left\{\int d\ell dk\, J_B\Big(\hat{s}_t - \frac{Q\ell}{m}, \Gamma_t, \delta m\Big) S_C\Big[\quad \ell, k, m, Q, z_{\rm cut}, \beta \quad\Big] F_C(k)\right\}$$

Preliminary Studies of SoftDrop Effects

z_{cut} dependence

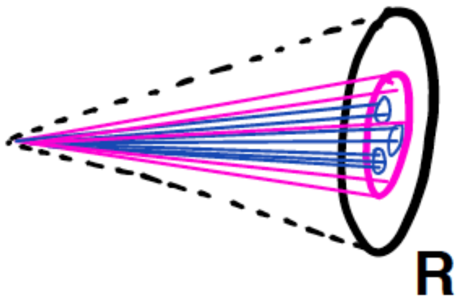
predict transition for "light Soft Drop" ✓

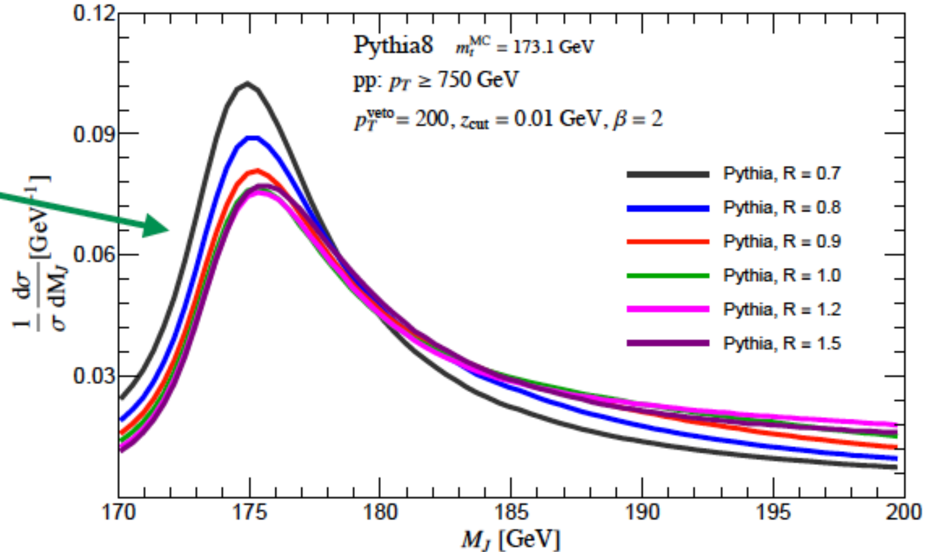
most contamination is removed



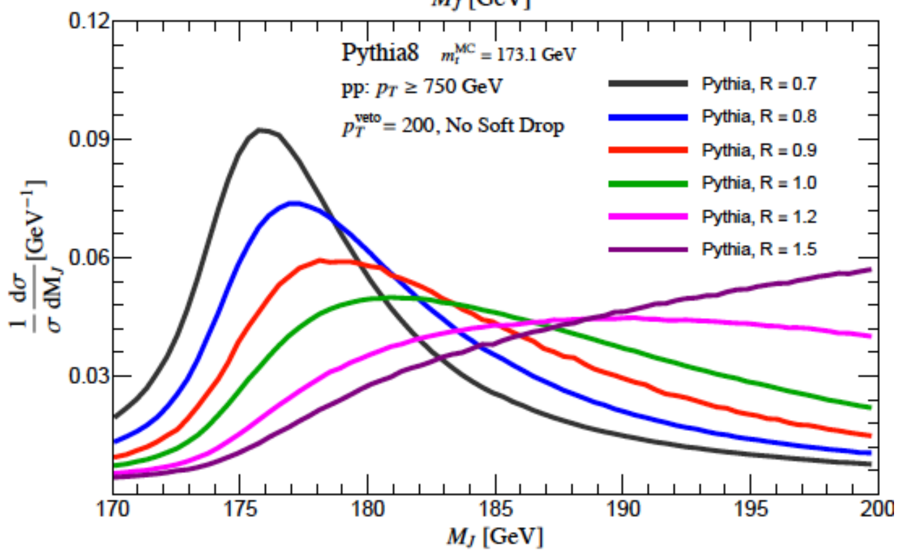
Preliminary Studies of SoftDrop Effects

predict: independent of Jet Radius ✓



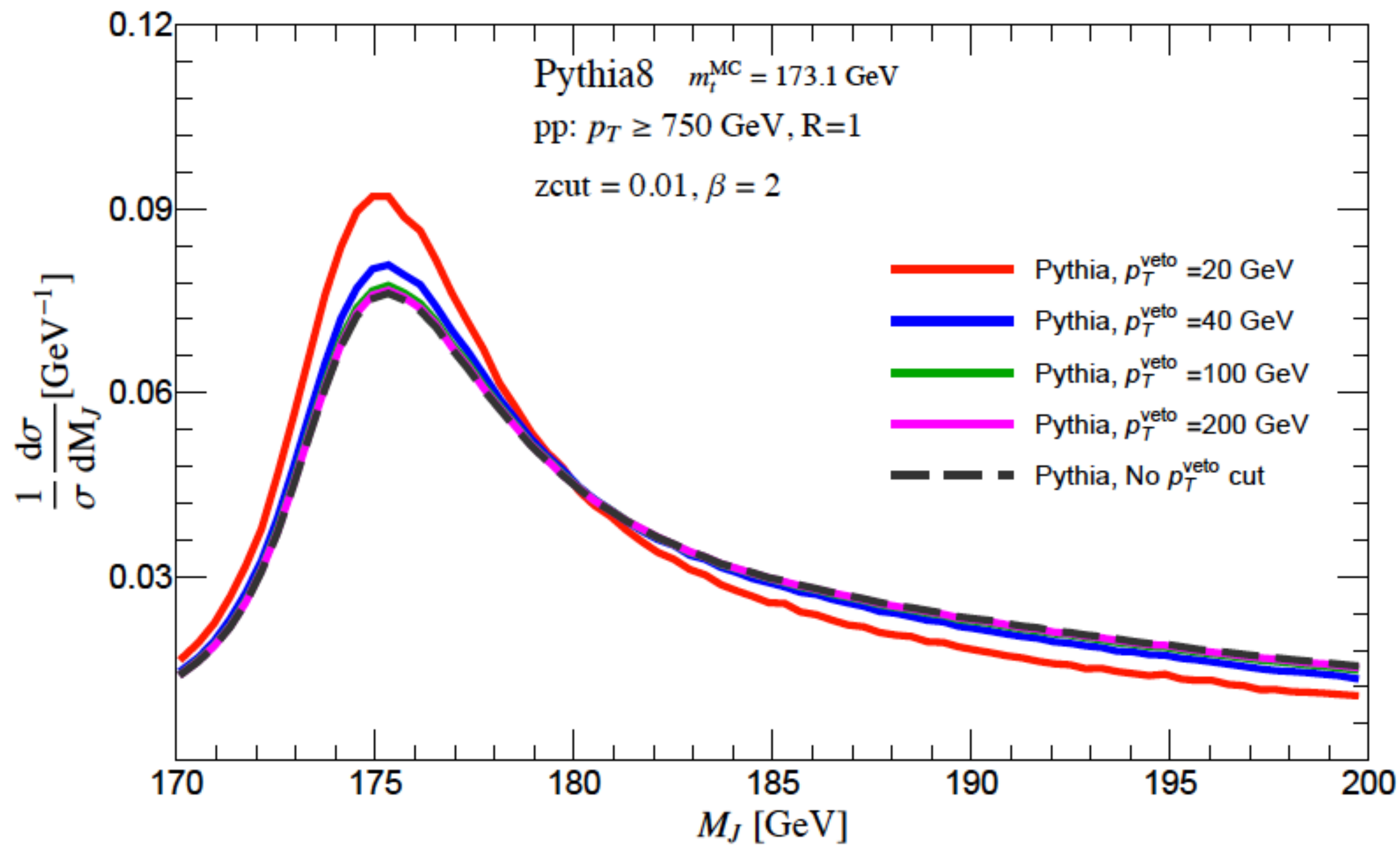


Without Soft Drop (huge):



Preliminary Studies of SoftDrop Effects

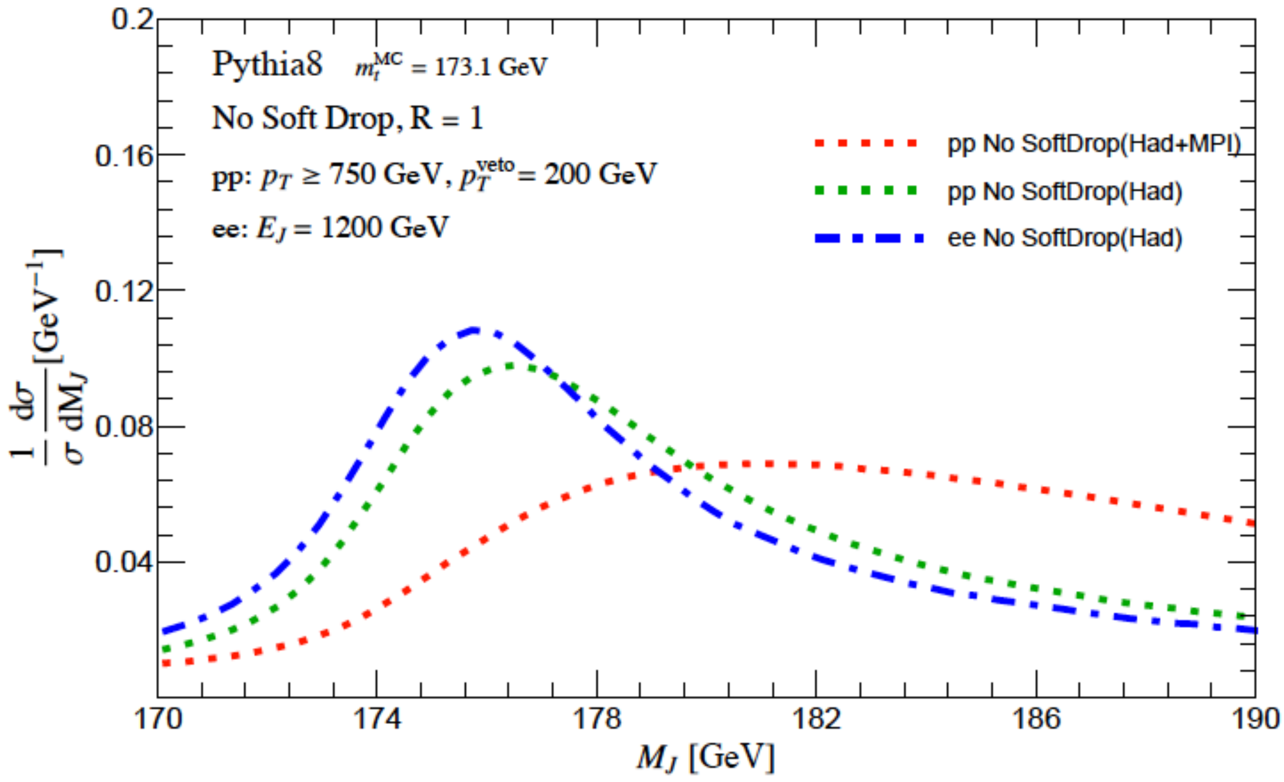
Predict independent of cutoff on radiation outside the jet ("jet veto"): ✓



Preliminary Studies of SoftDrop Effects

Soft Drop prediction: Same Result for e^+e^- and pp collisions

Without Soft Drop (differ):

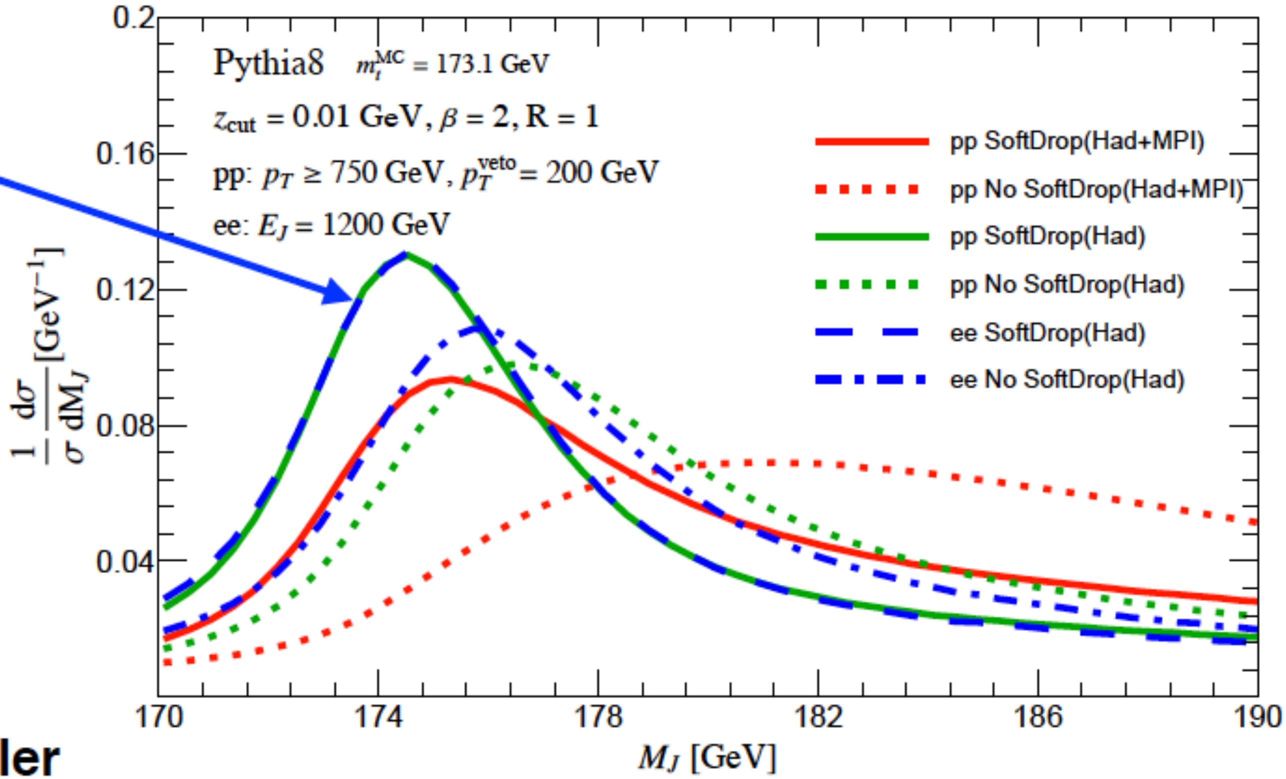


Preliminary Studies of SoftDrop Effects

Soft Drop prediction: Same Result for e^+e^- and pp collisions

With Soft Drop:

Great!



much smaller contamination

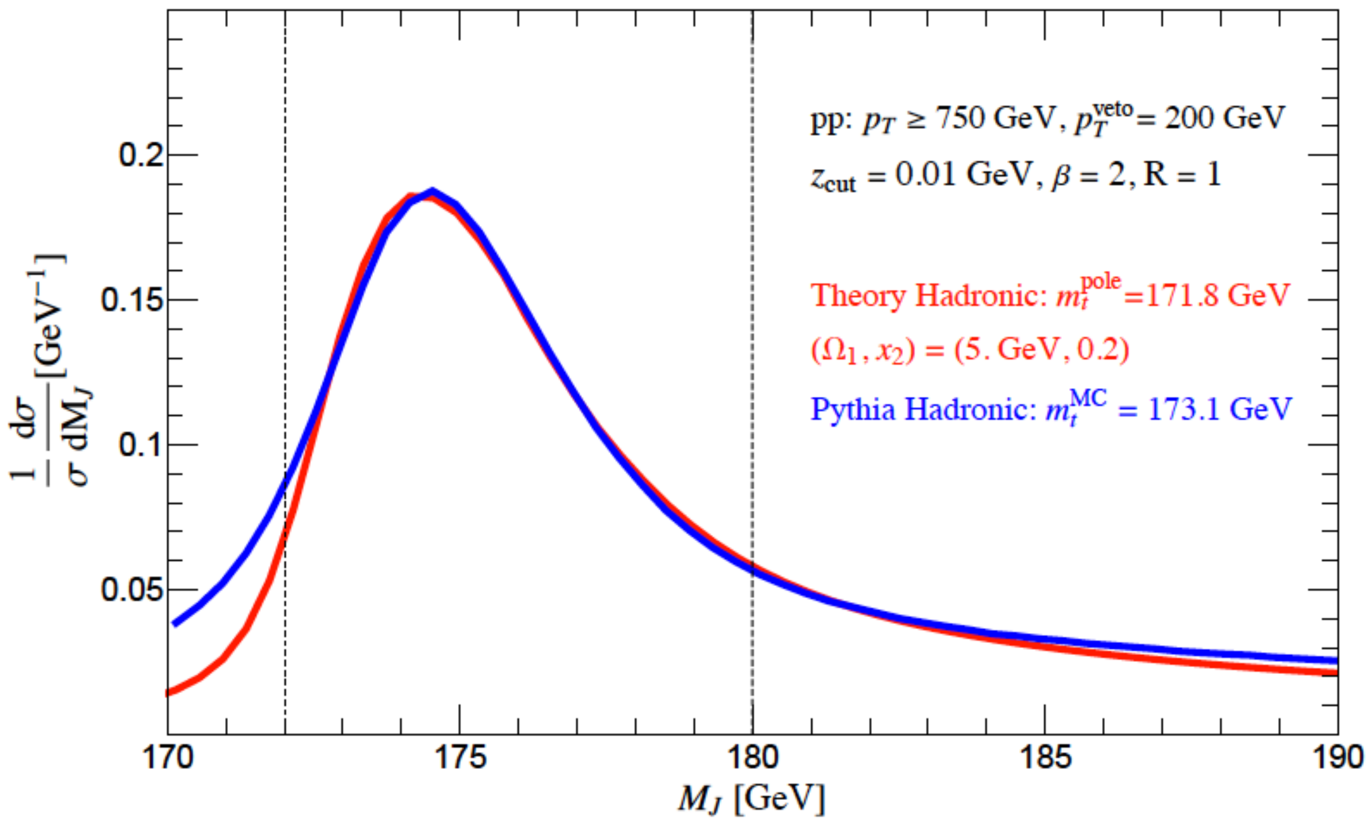
Preliminary Studies of SoftDrop Effects

Pythia Simulation vs. Theory (with Soft Drop)

without contamination:

$$m_t^{\text{pole}} = 171.8 \text{ GeV}$$

$$m_t^{\text{MC}} = 173.1 \text{ GeV}$$



universität wien

Preliminary Studies of SoftDrop Effects

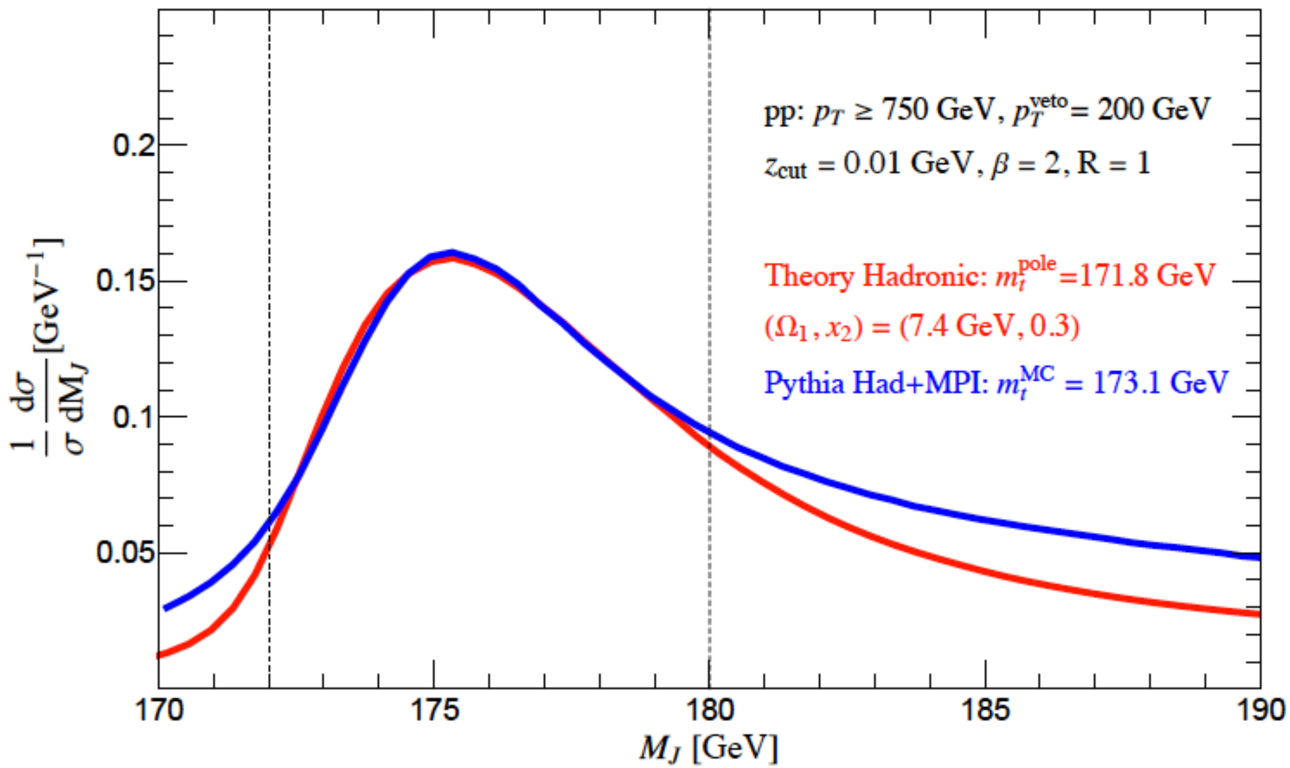
Pythia Simulation vs. Theory (with Soft Drop)

with contamination:

$m_t^{\text{pole}} = 171.8\ \text{GeV}$

$m_t^{\text{MC}} = 173.1\ \text{GeV}$

Same!



Dominant change is expected: Ω_1 (hadronization)

Summary

- First systematic MC top quark mass calibration based on e^+e^- 2-jettiness (large p_T): related to observables dominating the reconstruction method

▸ $m_t^{Pythia8.2}$ = 173 GeV	▸ $m_t^{MSR}(1\text{GeV})$ = 172.82 ± 0.22 GeV ▸ $(m_t^{pole})_{NLO}$ = 172.71 ± 0.41 GeV

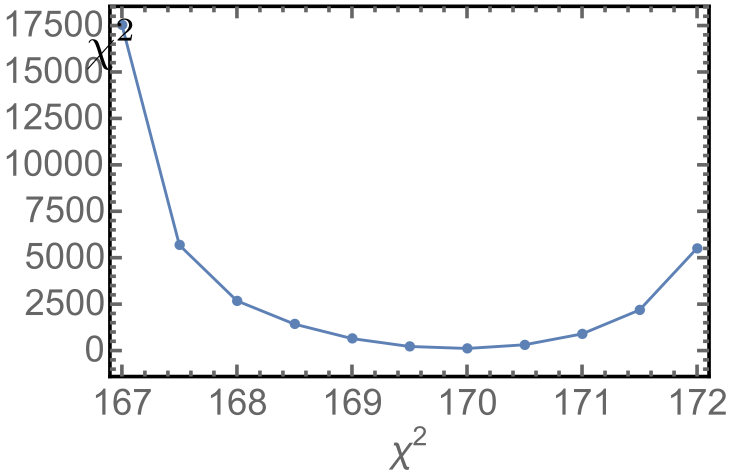
- NNLL+NLO QCD calculations based on an extension of the SCET approach concerning massive quark effects (all large logs incl. Ln(m)'s summed systematically) describing boosted top quarks.

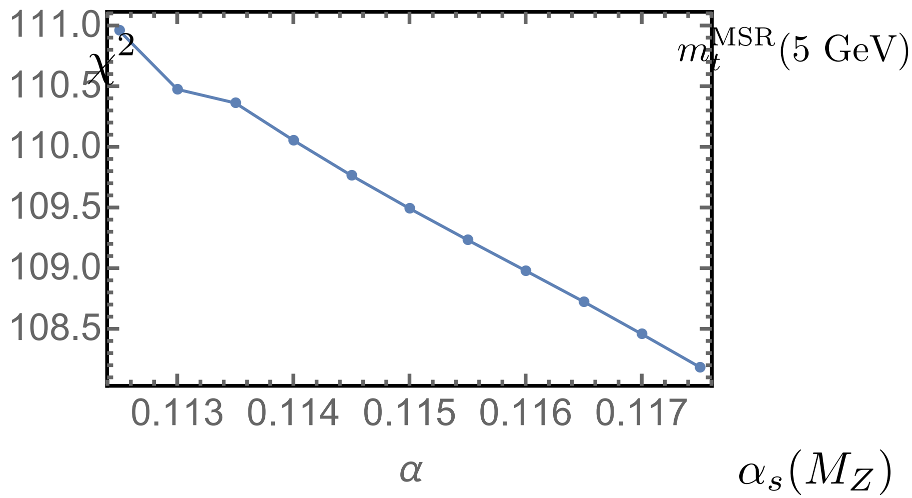
Future: consolidation & extension to pp collisions & MC studies

- Extension to pp collisions looks very promising with SoftDrop grooming to suppress MPI effects (boosted top quarks essential as well).
- Provides new ways to test and improve MC event generators.
- Plans:
  - Public code for calibration (CALIPER)
  - Other e^+e^- eventshapes (C-parameter, HJM)
  - NNNLL for e^+e^-
  - pp with SoftDrop (at NNLL)
  - Electroweak corrections
- Theory of the MC top quark mass: parton shower, hadronization model, NLO matching

Backup Slides

Peak Fits Parameter Sensitivity





Default renormalization scales; Γ_t=1.4 GeV, tune 7, $\Omega_{1,\mathrm{smear}}$=2.5 GeV, m_t^{Pythia}=171 GeV, Q={700, 1000, 1400} GeV, peak fit (60/80)%

→ $\chi^2_{\mathrm{min}} \sim O(100)$

- Very strong sensitivity to m_t
- Low sensitivity to strong coupling

- Take PDF strong coupling as input: $\alpha_S(M_Z)$ = 0.1181(13) (error irrelevant for m_t^{MSR}, m_t^{pole})

- χ^2_{min} and $\delta m_t^{\mathrm{stat}}$ do not have any physical meaning
- PDF rescaling method: $(\chi^2_{\mathrm{min}})^{\mathrm{rescale}} = 1$ can be used to define an incompatibility uncertainty

Top Mass Reconstruction Error Budget

Lepton+jets channel	m_t fit type 2D δm_t^{2D} (GeV)	2D δJSF	1D δm_t^{1D} (GeV)	hybrid δm_t^{hyb} (GeV)
Experimental uncertainties				
Method calibration	0.04	0.001	0.04	0.04
Jet energy corrections				
– JEC: Intercalibration	<0.01	<0.001	+0.02	+0.01
– JEC: In situ calibration	−0.01	+0.003	+0.24	+0.12
– JEC: Uncorrelated non-pileup	+0.09	−0.004	−0.26	−0.10
– JEC: Uncorrelated pileup	+0.06	−0.002	−0.11	−0.04
Lepton energy scale	+0.01	<0.001	+0.01	+0.01
E_T^{miss} scale	+0.04	<0.001	+0.03	+0.04
Jet energy resolution	−0.11	+0.002	+0.05	−0.03
b tagging	+0.06	< 0.001	+0.04	+0.06
Pileup	−0.12	+0.002	+0.05	−0.04
Backgrounds	+0.05	< 0.001	+0.01	+0.03
Modeling of hadronization				
JEC: Flavor-dependent				
– light quarks (u d s)	+0.11	−0.002	−0.02	+0.05
– charm	+0.03	<0.001	−0.01	+0.01
– bottom	−0.32	<0.001	−0.31	−0.32
– gluon	−0.22	+0.003	+0.05	−0.08
b jet modeling				
– b fragmentation	+0.06	−0.001	−0.06	<0.01
– Semileptonic b hadron decays	−0.16	<0.001	−0.15	−0.16
Modeling of perturbative QCD				
PDF	0.09	0.001	0.06	0.04
Ren. and fact. scales	+0.17 ± 0.08	−0.004 ± 0.001	−0.24 ± 0.06	−0.09 ± 0.07
ME-PS matching threshold	+0.11 ± 0.09	−0.002 ± 0.001	−0.07 ± 0.06	+0.03 ± 0.07
ME generator	−0.07 ± 0.11	−0.001 ± 0.001	−0.16 ± 0.07	−0.12 ± 0.08
Top quark p_T	+0.16	−0.003	−0.11	+0.02
Modeling of soft QCD				
Underlying event	+0.15 ± 0.15	−0.002 ± 0.001	+0.07 ± 0.09	+0.08 ± 0.11
Color reconnection modeling	+0.11 ± 0.13	−0.002 ± 0.001	−0.09 ± 0.08	+0.01 ± 0.09
Total systematic	0.59	0.007	0.62	0.48
Statistical	0.20	0.002	0.12	0.16
Total	0.62	0.007	0.63	0.51

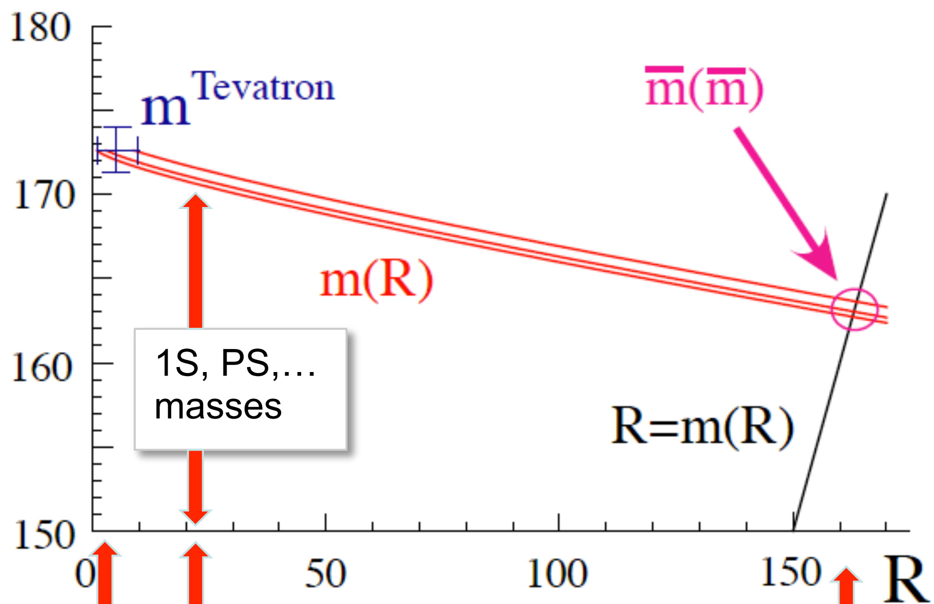
← NLO ME corrections

$m_t^{MC} = 172.44 \pm 0.49$
(CMS Run-1 final, 2015)
arXiv:1509.04044

MSR Mass Definition

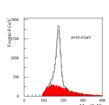
AH, Stewart: arXive:0808.0222

$$m_t^{\rm MC} = m_t^{\rm MSR}(3^{+6}_{-2}\ {\rm GeV}) = m_t^{\rm MSR}(3\ {\rm GeV})^{+0.6}_{-0.3}$$



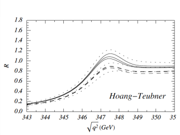
Good choice for R:

Of order of the typical scale of the observable used to measure the top mass.



Peak of invariant mass distribution, endpoints

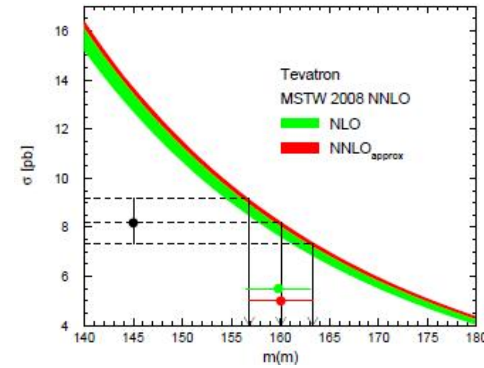
Top-antitop threshold at the ILC



Total cross section, e.w.precsion obs., Unification, MSbar mass

Masses Loop-Theorists Like to use

Langenfeld, Moch, Uwer



Total cross section (LHC/Tev):

$$m_t^{\mathrm{MSR}}(R = m_t) = \overline{m}_t(\overline{m}_t)$$

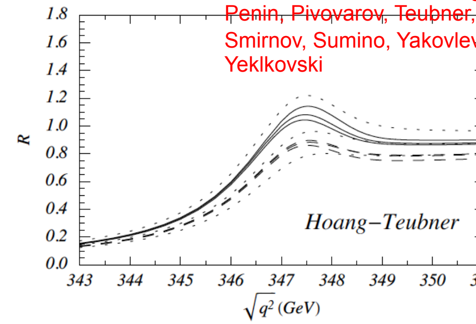
$$M_t \;=\; M_t^{(O)} \;+\; M_t(0)\alpha_s \;+\; \ldots$$

- more inclusive
- sensitive to top production mechanism (pdf, hard scale)
- indirect top mass sensitivity
- large scale radiative corrections

Mass schemes related to different computational methods

Relations computable in perturbation theory

Beneke, AH, Melnikov, Nagano, Penin, Pivovarov, Teubner, Signer, Smirnov, Sumino, Yakovlev, Yeklkovski



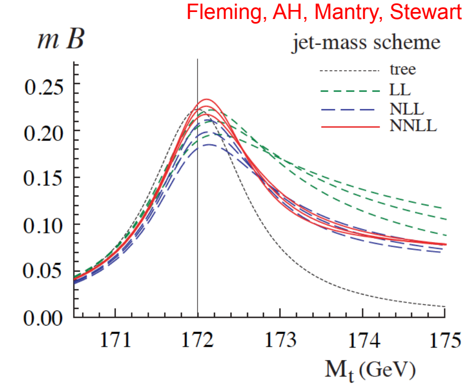
Threshold cross section (ILC):

$$m_t^{\mathrm{MSR}}(R \sim 20~\mathrm{GeV})\,,\;\; m_t^{\mathrm{1S}}\,,\;\; m_t^{\mathrm{PS}}(R)$$

$$M_t \;=\; M_t^{(O)} \;+\; \langle p_{\mathrm{Bohr}}\rangle\alpha_s \;+\; \ldots$$

$$\langle p_{\mathrm{Bohr}}\rangle \;=\; 20\,\mathrm{GeV}$$

Fleming, AH, Mantry, Stewart



Inv. mass reconstruction (ILC/LHC):

$$m_t^{\mathrm{MSR}}(R \sim \Gamma_t)\,,\;\; m_t^{\mathrm{jet}}(R)$$

$$M_t \;=\; M_t^{(O)} \;+\; \Gamma_t\alpha_s \;+\; \ldots$$

$$\Gamma_t \;=\; 1.3\,\mathrm{GeV}$$

- more exclusive
- sensitive to top final state interactions (low scale)
- direct top mass sensitivity
- small scale radiative corrections