

# Sequential Strangeness Freeze-out

R. Bellwied (University of Houston)

Thanks to: Paolo Alba, Livio Bianchi, Szabolcs Borsanyi,  
Zoltan Fodor, Anders, Knospe, Valentina Mantovani-Sarti,  
Jackie Noronha-Hostler, Paolo Parotto, Attila Pasztor,  
Claudia Ratti

17<sup>th</sup> International Conference on  
**Strangeness in  
Quark Matter**



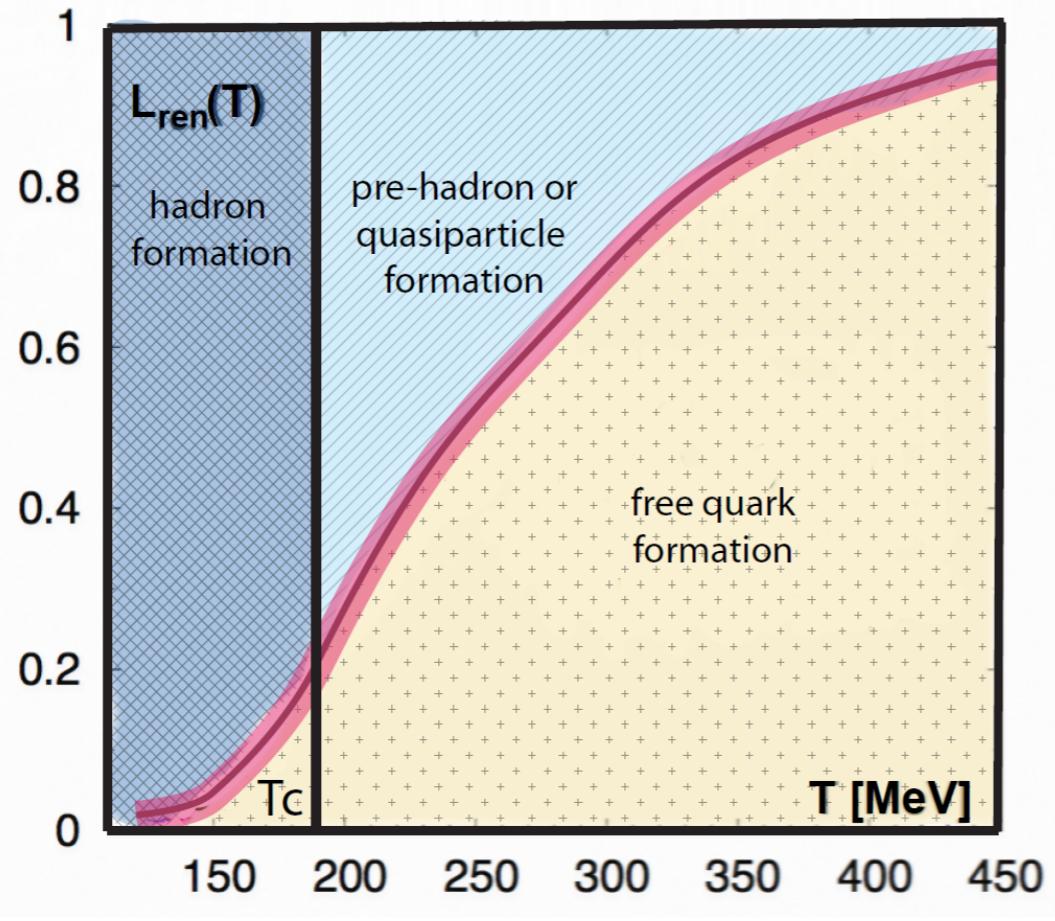
Universiteit Utrecht

10-15 July 2017  
Utrecht, the Netherlands

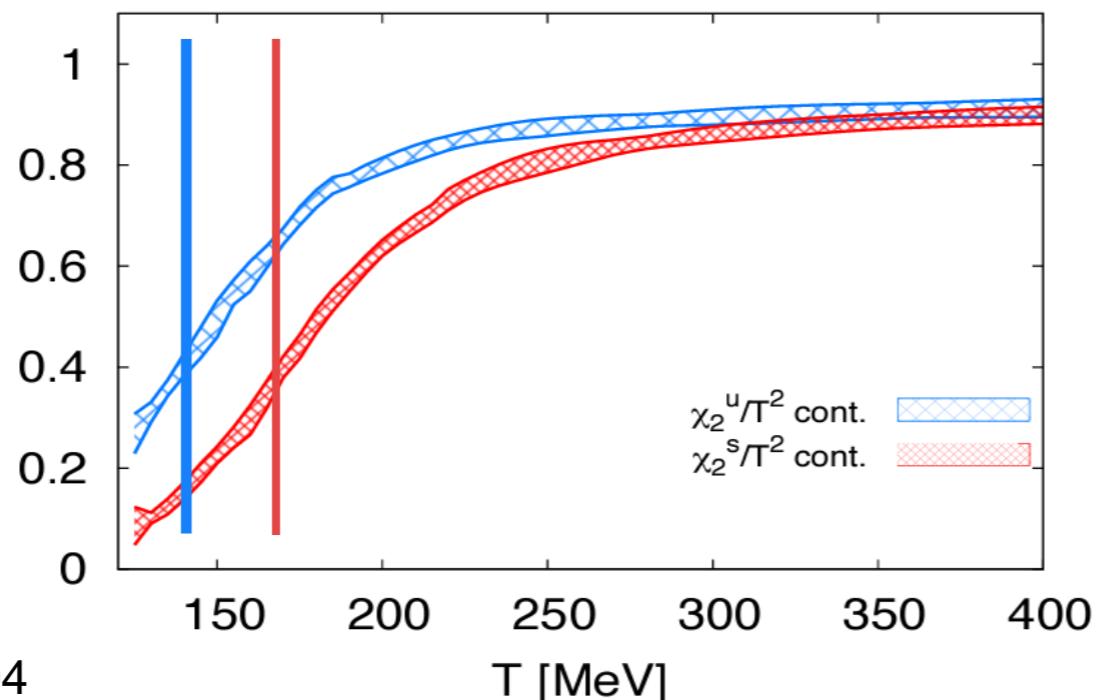
Utrecht 2017

# Lattice order parameters in the QCD cross-over

e.g. a re-interpretation of the Polyakov Loop calculation in lattice QCD



RB, C. Markert,  
PLB691 (2010) 208



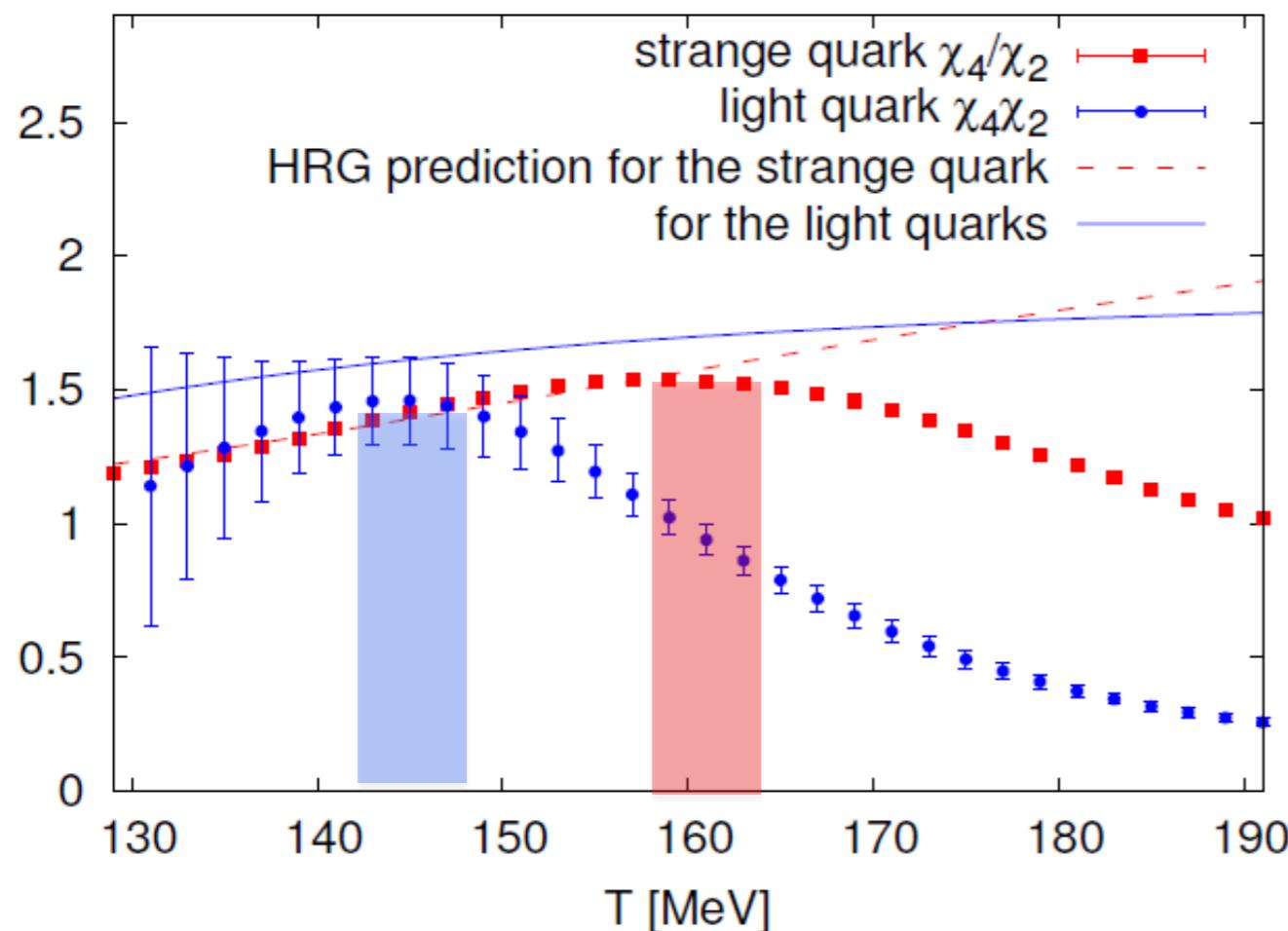
C. Ratti et al.,  
PRD 85 (2012) 014004

- In a regime where we have a smooth crossover and where quark masses (even for the s-quark) could play a role why would there be a single freeze-out surface ?
- We can calculate thermodynamic quantities for a static equilibrated system at a fixed temperature
- But are pseudo-critical temperatures extracted from flavor dependent susceptibilities as relevant for hadronization properties as chirality ?

# Direct determination of freeze-out parameters from first principles (lattice QCD)

$$\kappa_B \sigma_B^2 \equiv \frac{\chi_{4,\mu}^B}{\chi_{2,\mu}^B} = \frac{\chi_4^B(T)}{\chi_2^B(T)} \left[ \frac{1 + \frac{1}{2} \frac{\chi_6^B(T)}{\chi_4^B(T)} (\mu_B/T)^2 + \dots}{1 + \frac{1}{2} \frac{\chi_4^B(T)}{\chi_2^B(T)} (\mu_B/T)^2 + \dots} \right]$$

Susceptibility ratios are a model independent measure of the chemical freeze-out temperature near  $\mu=0$ . (Karsch, arXiv:1202.4173)



R. Bellwied & WB Collab., PRL (2013), arXiv:1305.6297

**Indication of sequential hadronization**

Either based on the peak position in the lattice QCD calculation or on the point of deviation from the hadron resonance gas (HRG)

Needs experimental verification

# Lattice QCD vs. HRG

The lattice calculations favor a flavor separation.

They use the proper pion and quark masses and they are continuum extrapolated.

Nothing can be changed or improved, except for the possibility that

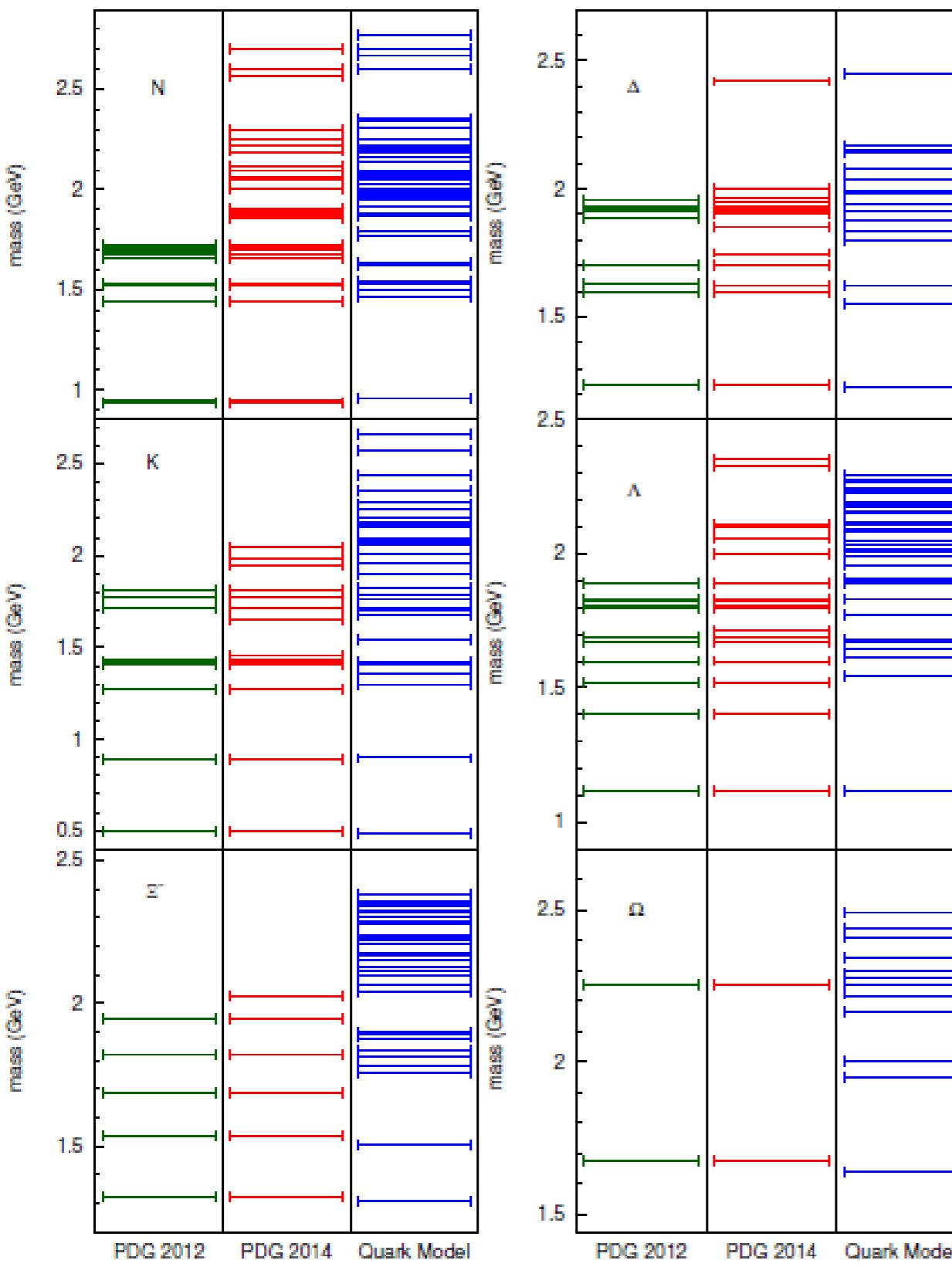
- a.) the action is not perfect
- b.) the system is not in equilibrium.

On the other hand the HRG comparison is subject to a lot of discussion in the recent past.

Two avenues have been pursued:

- a.) a non-interacting HRG can only describe ‘reality’ if the resonance spectrum is complete.
- b.) an interacting hadron gas is a better proxy as long as the interaction is properly modeled.

# The incomplete HRG input spectrum

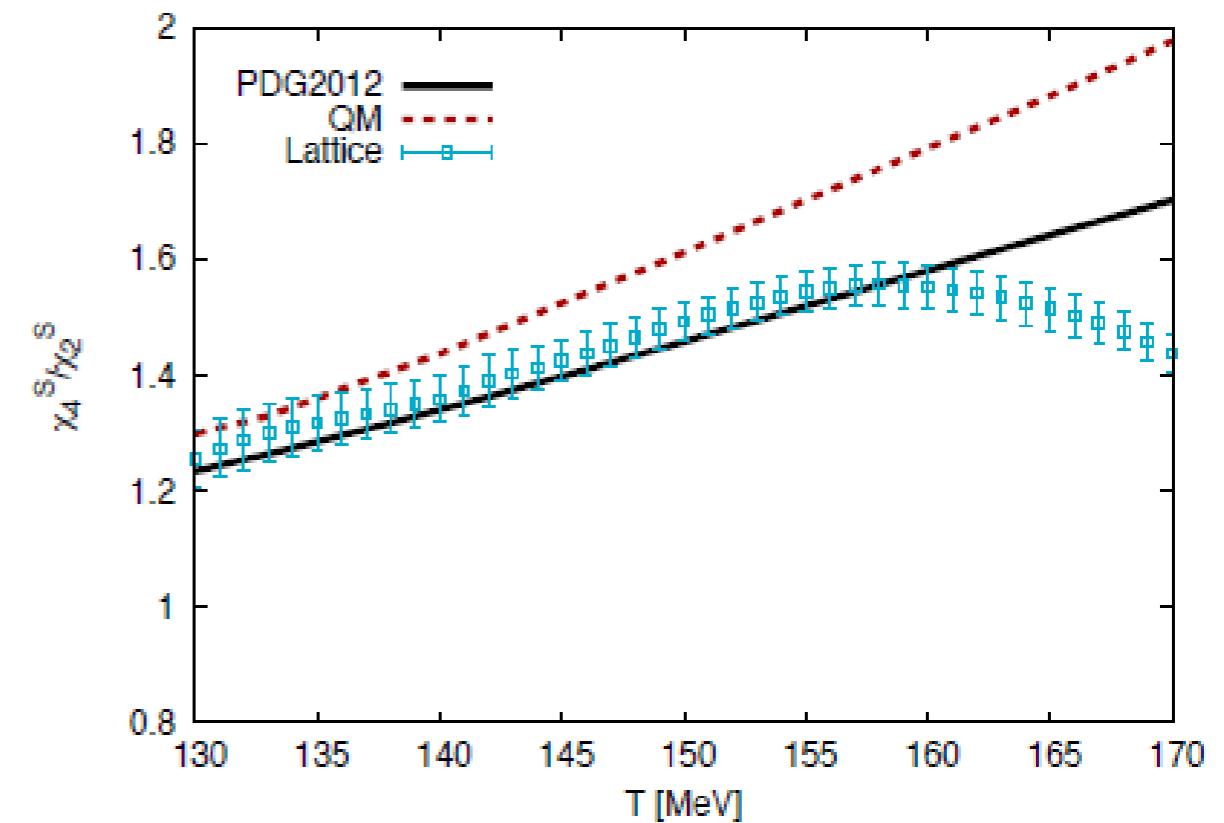
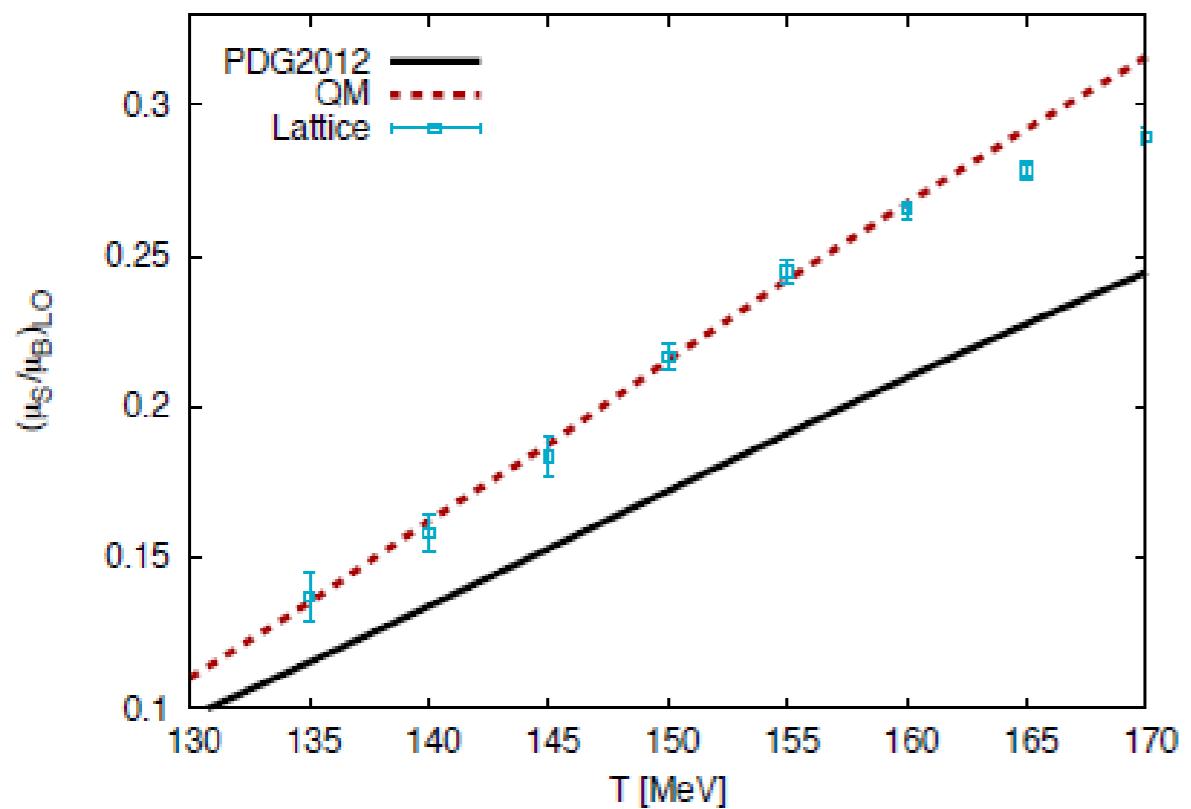


First suggested as a possible  
'solution' to the different flavor  
surfaces by Bazavov et al. (PRL  
(2014), arXiv:1404.6511)

A more detailed study by P.Alba et  
al., arXiv:1702.01113)

See talk by C. Ratti on  
Wednesday

# The problem: It impacts the comparison to different lattice parameters in different ways



Furthermore: there are many different Quark Models with different numbers of ‘extra states’ depending on the quark interaction

Best compromise seems to be PDG2016+ (incl. all 1-star resonances), see P.Alba et al., arXiv:1702.01113

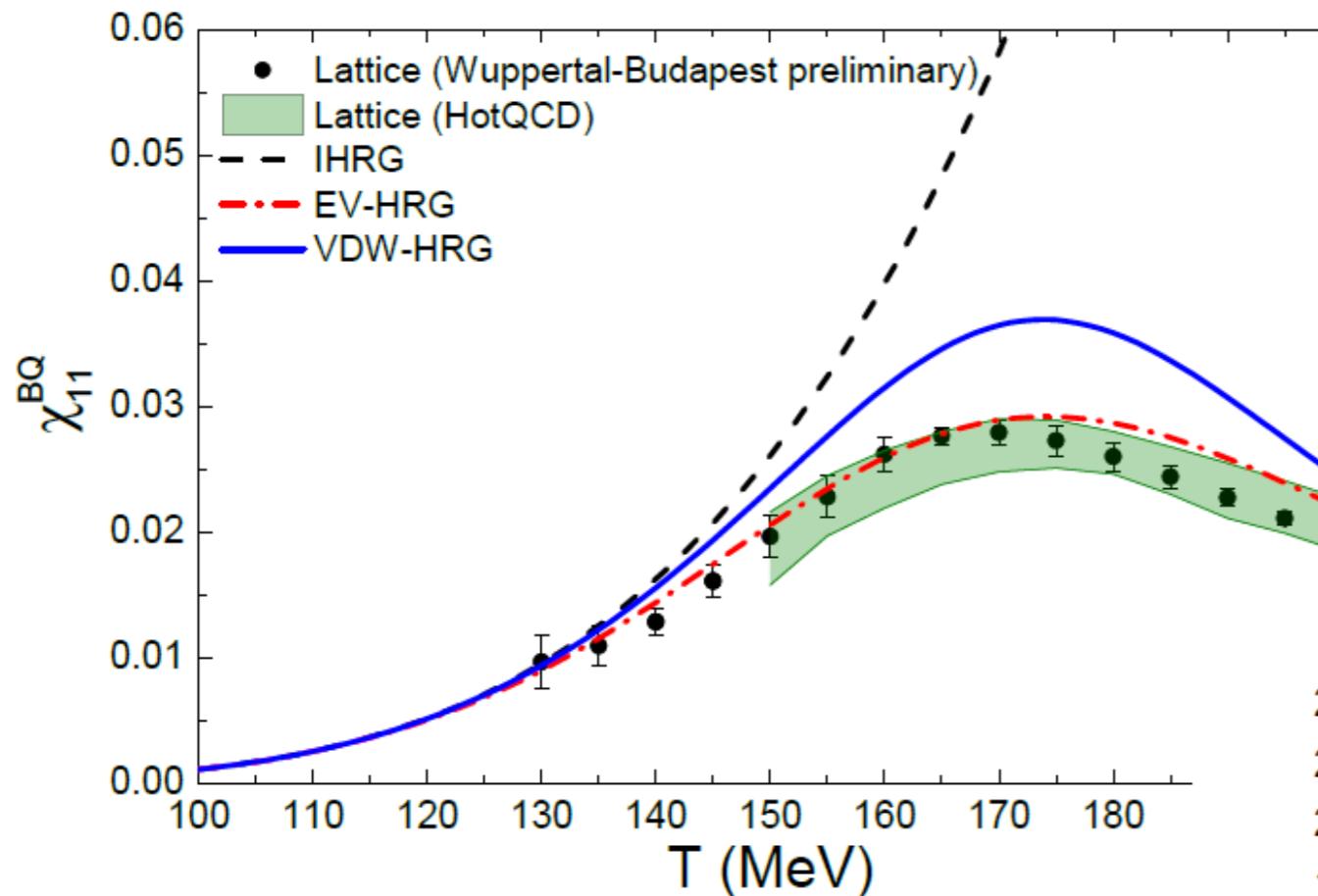
# Conclusions for the experimental program

We need more experimental verification of higher mass baryonic resonances, especially in the  $S = 1$  sector (higher level  $\Lambda$  and  $\Sigma$  resonances)

This is an interesting challenge for RHIC and LHC experiments, but also for the future JLab program

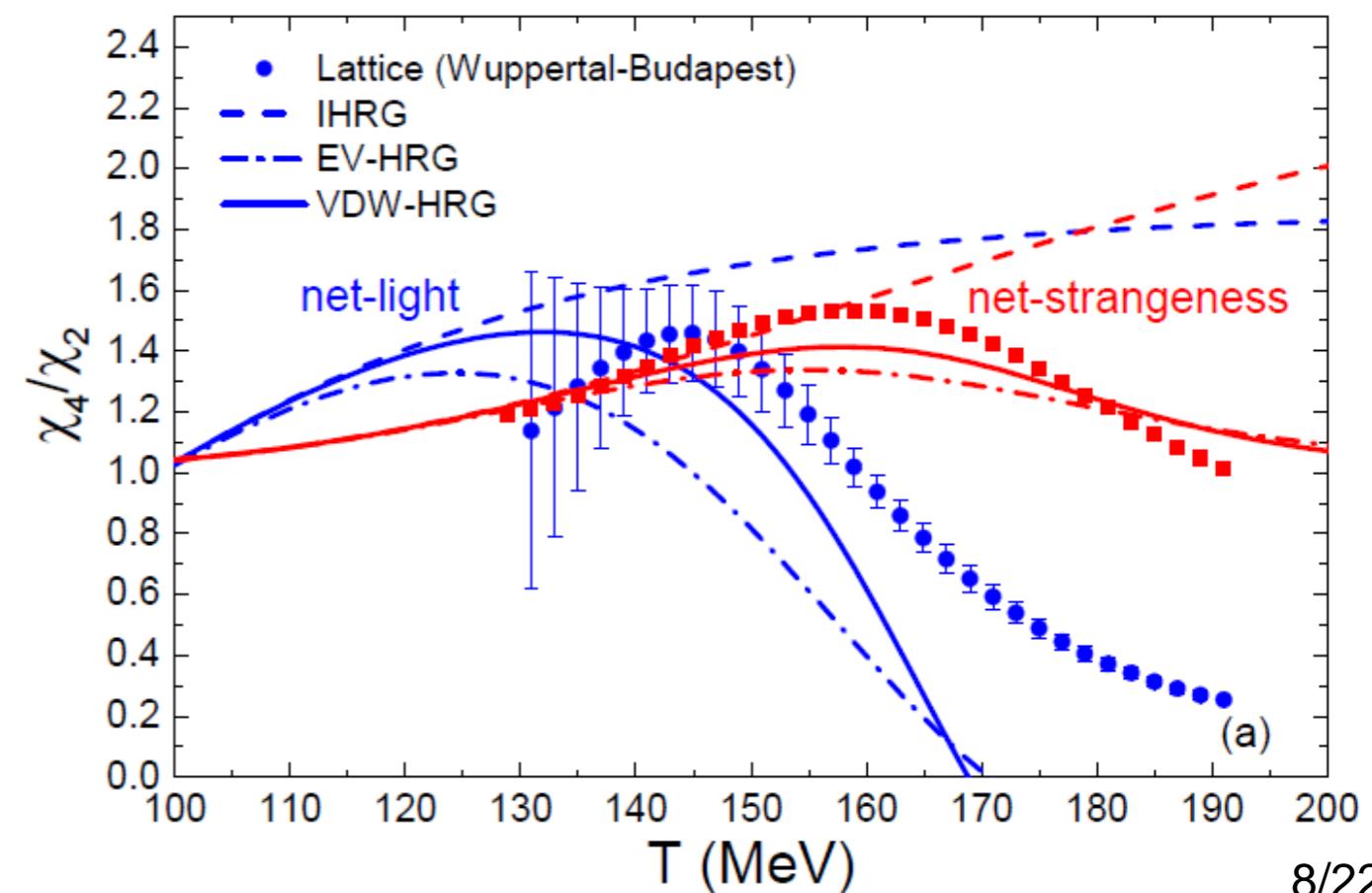
(new JLab experiment KLF was just proposed for this purpose, arXiv:1701.07346)

An alternative: the interacting hadron gas (excluded volume or van der Waals interaction), Vovchenko et al., PRL 2016, arXiv:1609.03975

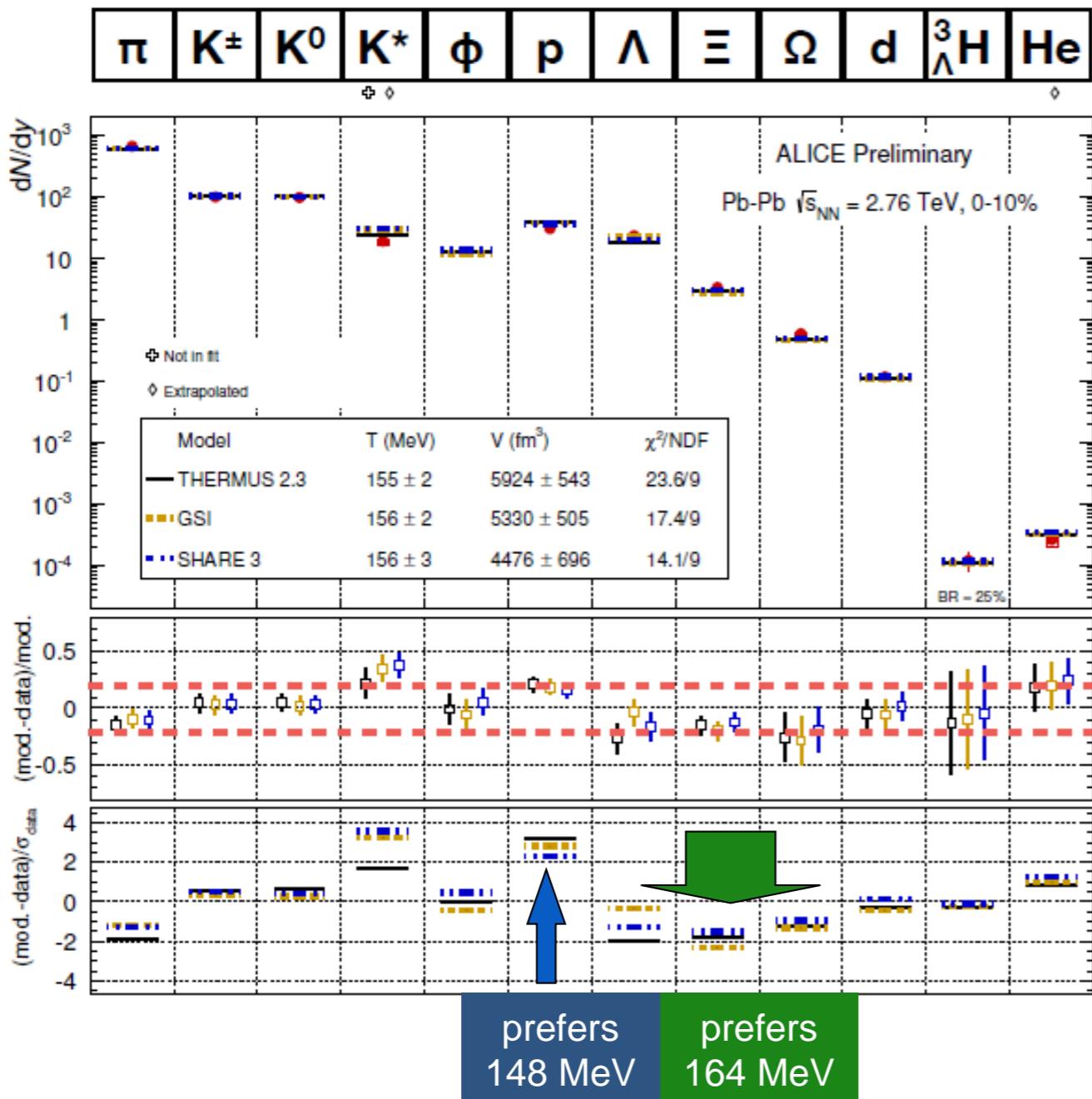


Seems to get some of the features right (peak structures in lattice susceptibility ratios)

- Problems:
- excluded volume calculation debatable
    - pushes ‘agreement’ with lattice way beyond the pseudo-critical temperature



# Experimental evidence: HRG (PDG 2010) model comparison based on yields



Data: ALICE, QM 2014, arXiv:1408.6403  
This looks like a good fit, but it is not

$\chi^2/\text{NDF}$  improves from 2 to 1 when pions and protons are excluded.

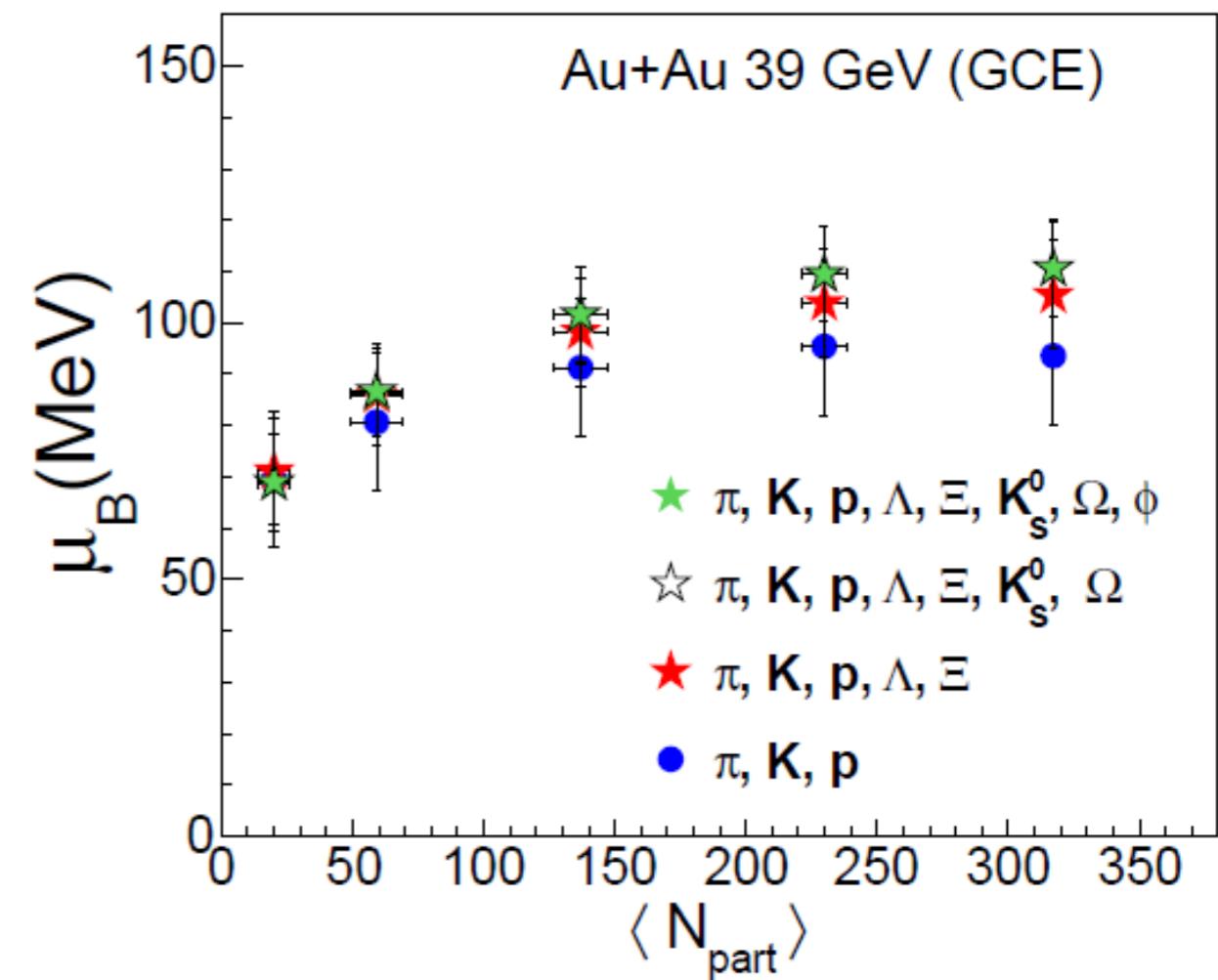
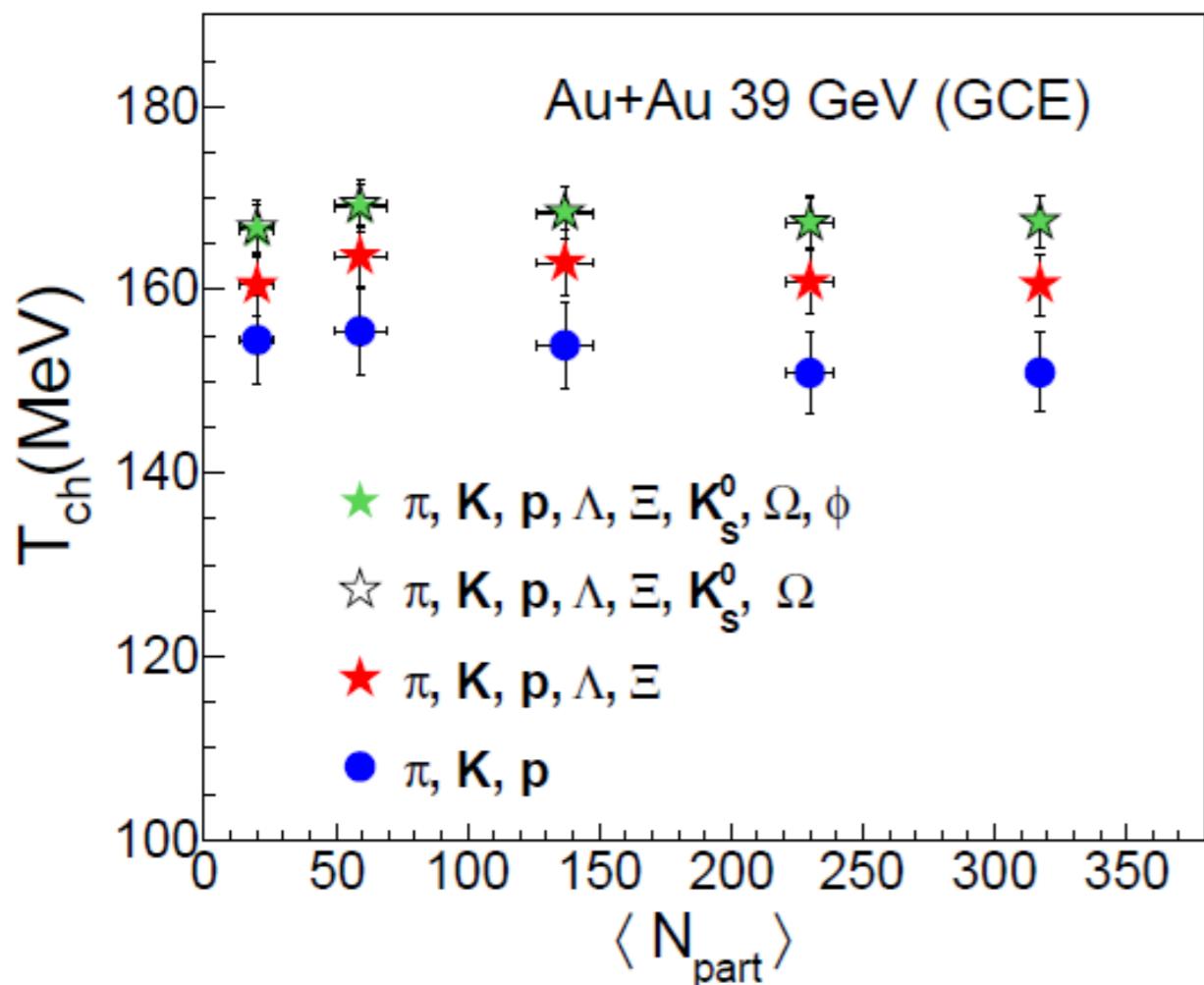
Fit to pions and protons alone yield a temperature of 148 MeV. Strange baryons yield 164 MeV

Several alternate explanations:

- Inclusion of Hagedorn states
  - Non-equilibrium fits
  - Baryon annihilation
- Different  $T_{ch}$  for light and strange

# Experimental evidence from varying the input particles into the chemical fit

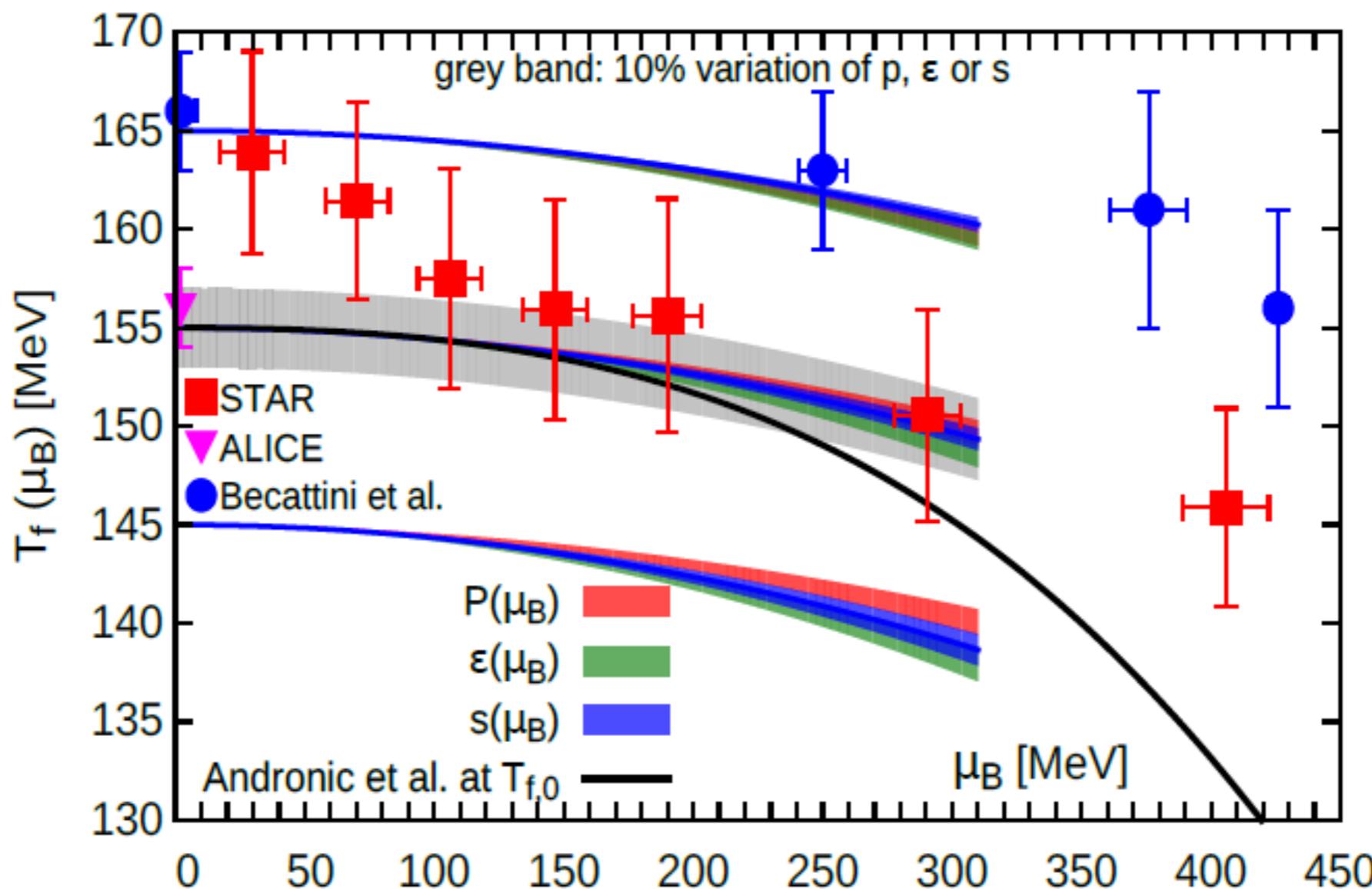
Latest example: Beam Energy Scan data from STAR  
(arXiv:1701.07065)



This is a long known fact in SHM, always argued as ‘the more states the better’, but all additional states (to  $\pi, k, p$ ) are strange states

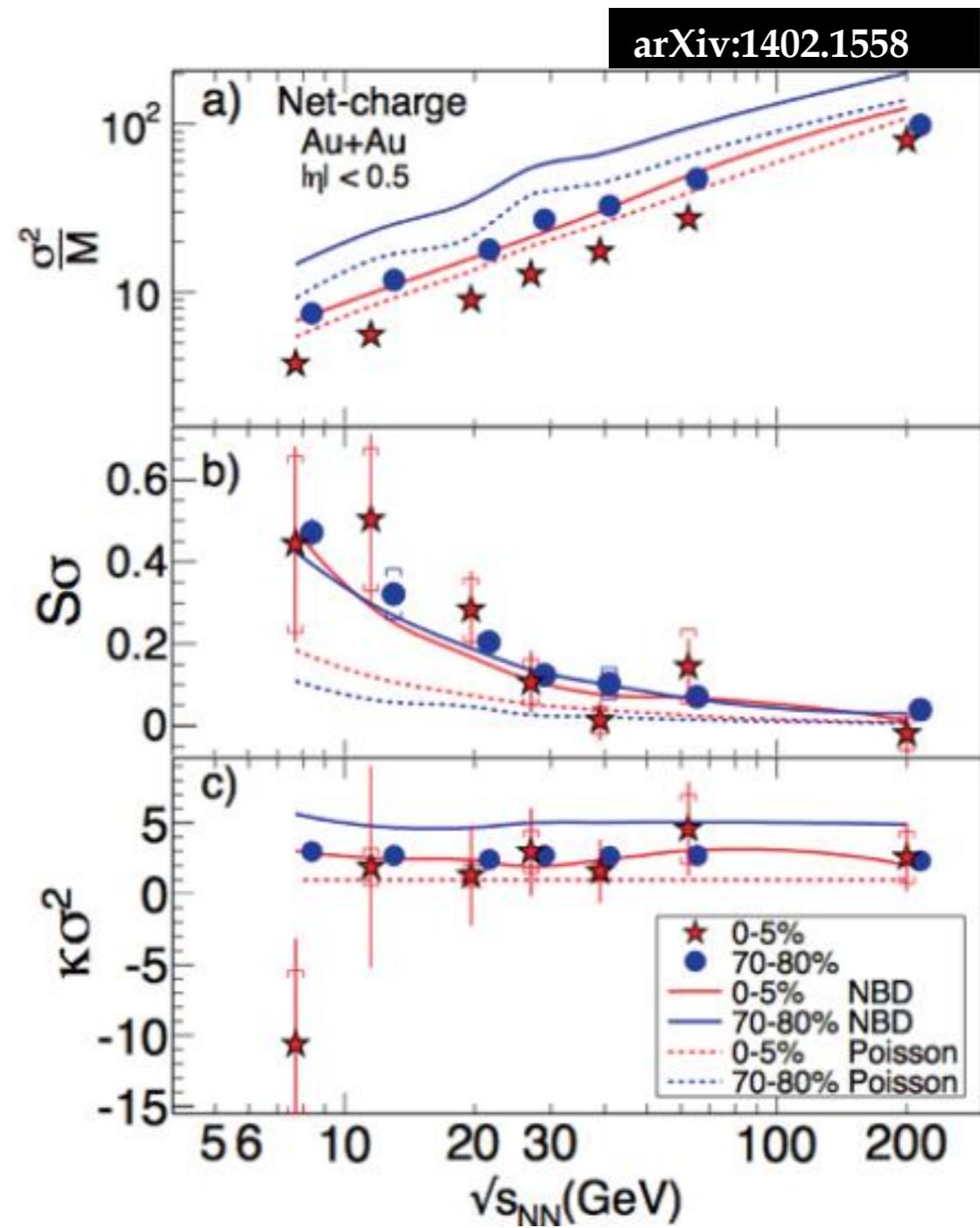
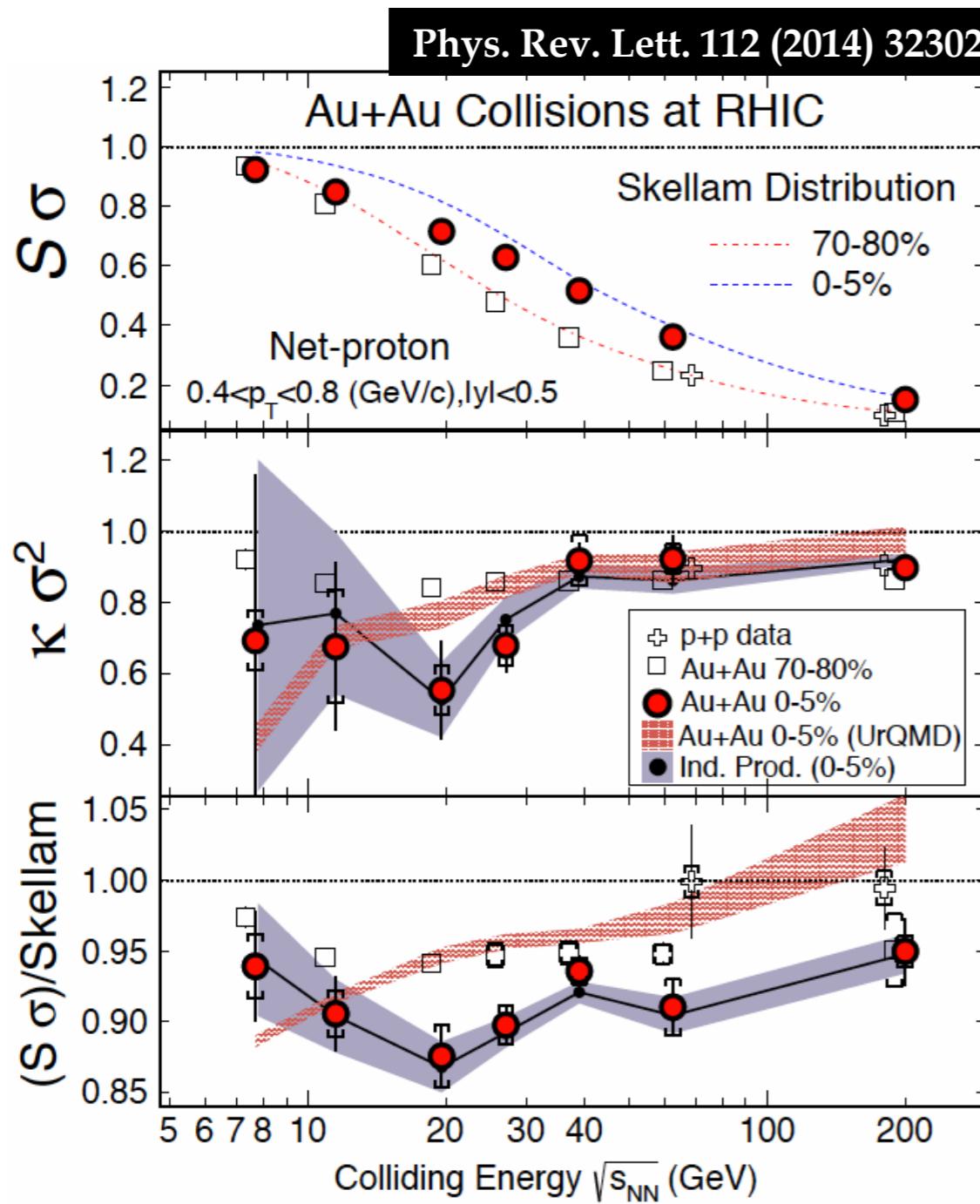
# There is an upper limit.....

One important fact to remember: The chemical freeze-out surface cannot be above the pseudo-critical temperature from the lattice (154+9 MeV)  
(both moved down over the years....)



F.Karsch, arXiv:1611.01973, SQM 2016  
spread in  $T$  given by uncertainty in chiral pseudo-critical  $T$ ,  
but also by difference between  $T$  from  $\chi_u$  and  $\chi_s$

# Higher moment ratios for net-charge and net-proton distributions (STAR 2014)



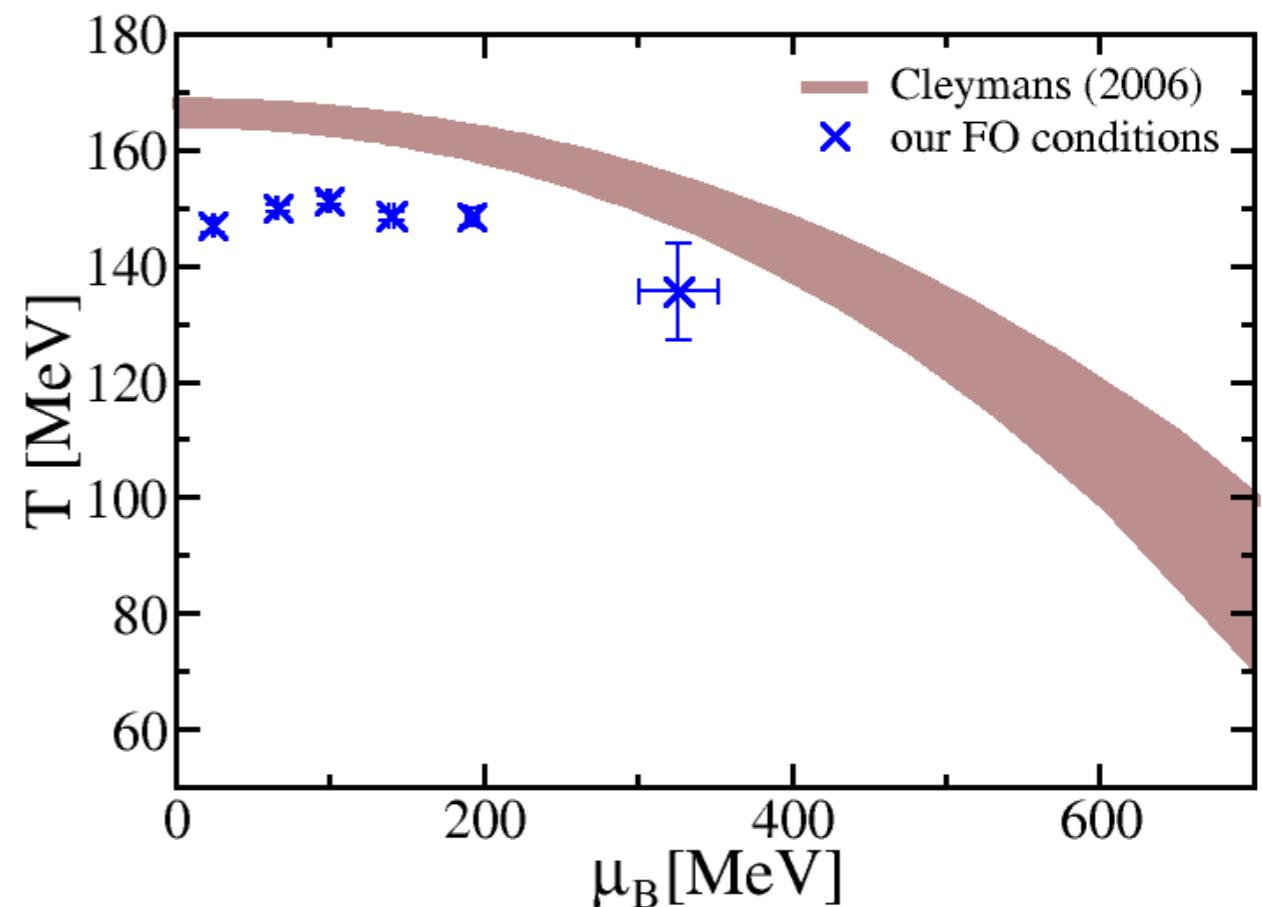
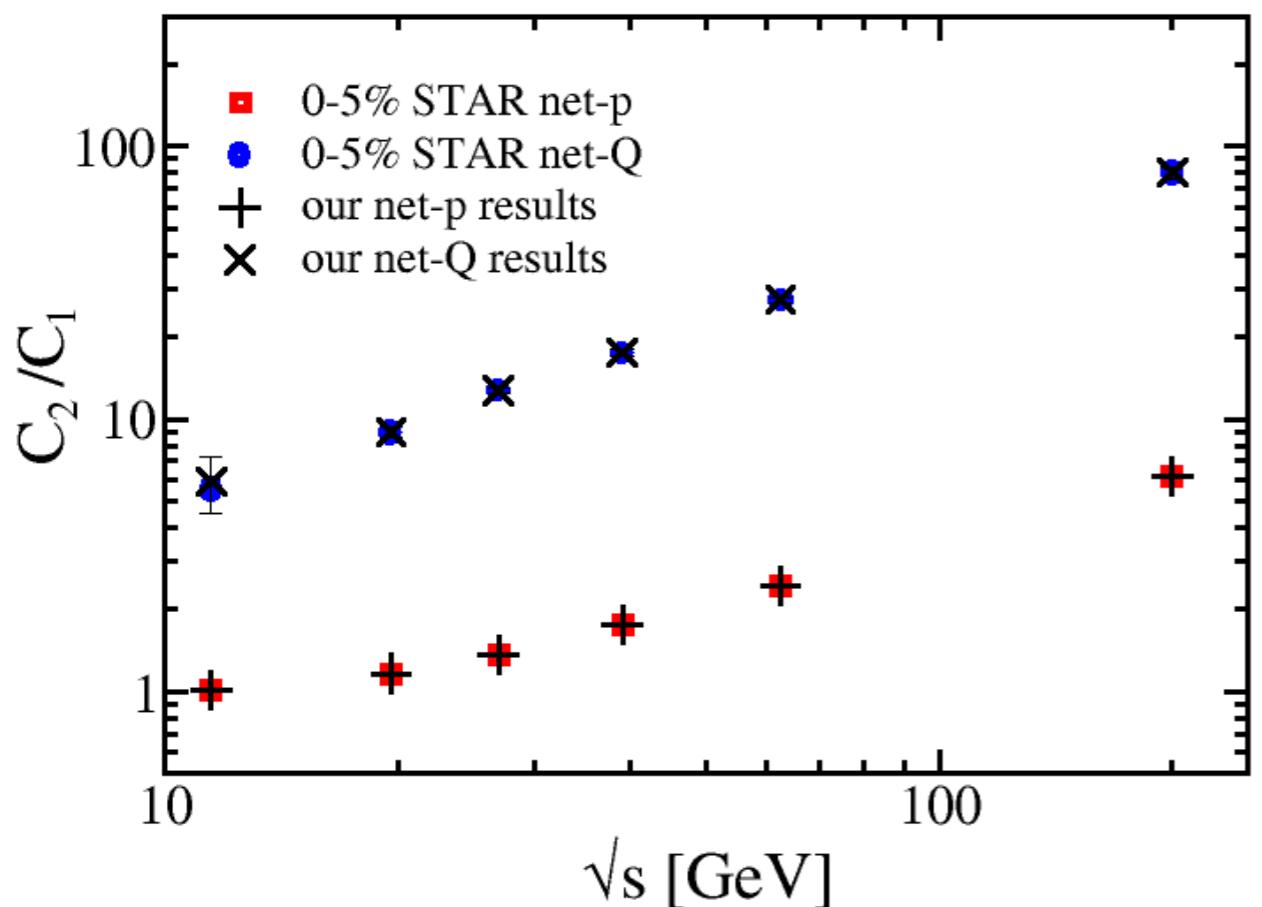
Fluctuations are more sensitive to chemical freeze-out as simple yields. They can be directly compared to susceptibilities on the lattice (P.Alba et al., PRC, (arXiv:1504.03262))

# HRG (PDG 2008) analysis of STAR results (charge & proton)

Alba, Bellwied, Bluhm, Mantovani, Nahrgang, Ratti, PLB (2014),  
arXiv:1403.4903

HRG in partial chemical equilibrium: resonance decays for resonances up to  $2 \text{ GeV}/c^2$  and weak decays taken into account), experimental cuts applied.

Use the lowest moments with the smallest errors and least ‘criticality’, i.e.  $\sigma^2/M$

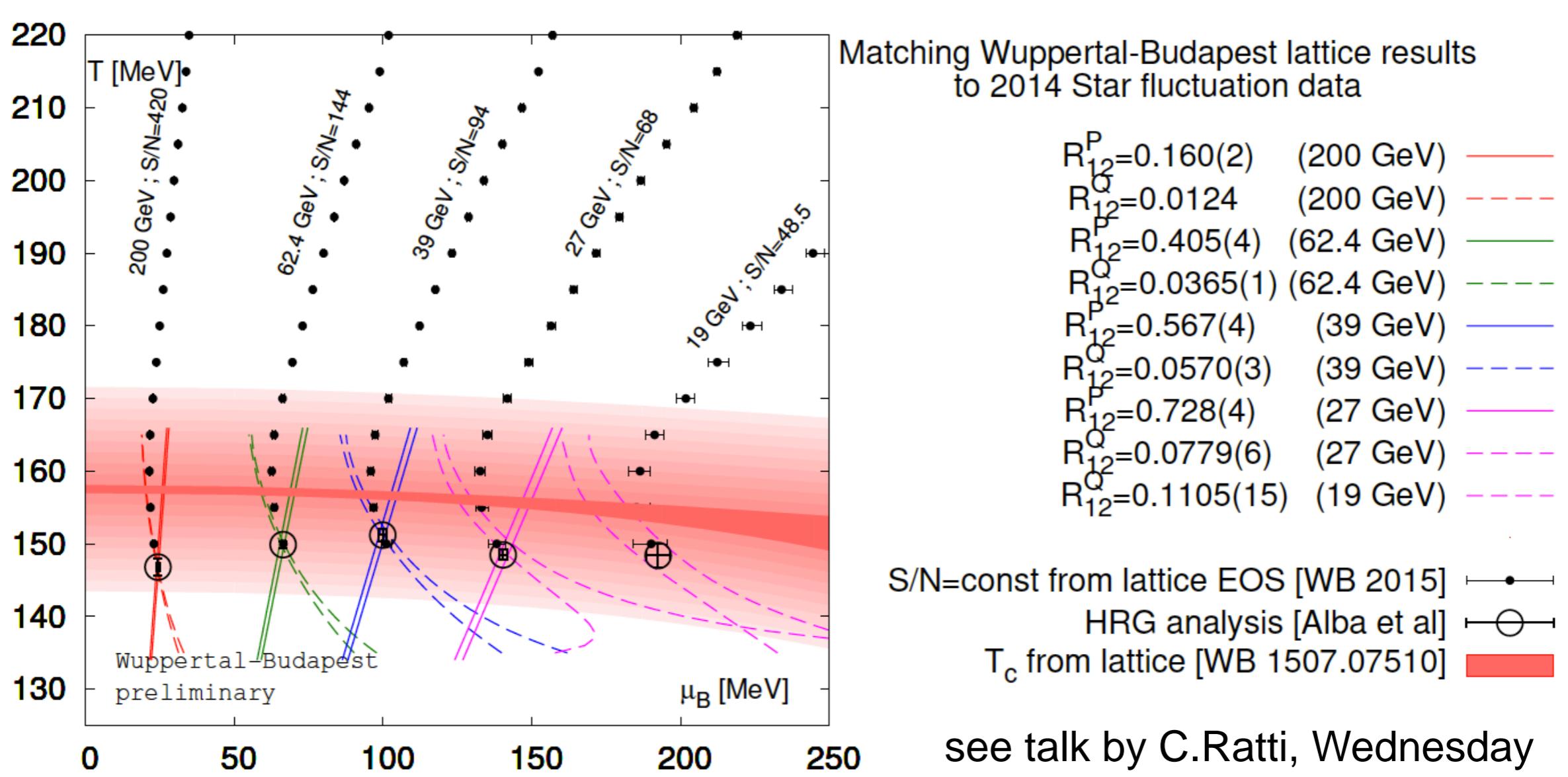


Result: intriguing ‘lower’ freeze-out temperature (compared to SHM yield fits)  
with very small error bars (due to good determination of  $c_2/c_1$ )

# The light quark freeze-out surface is well defined

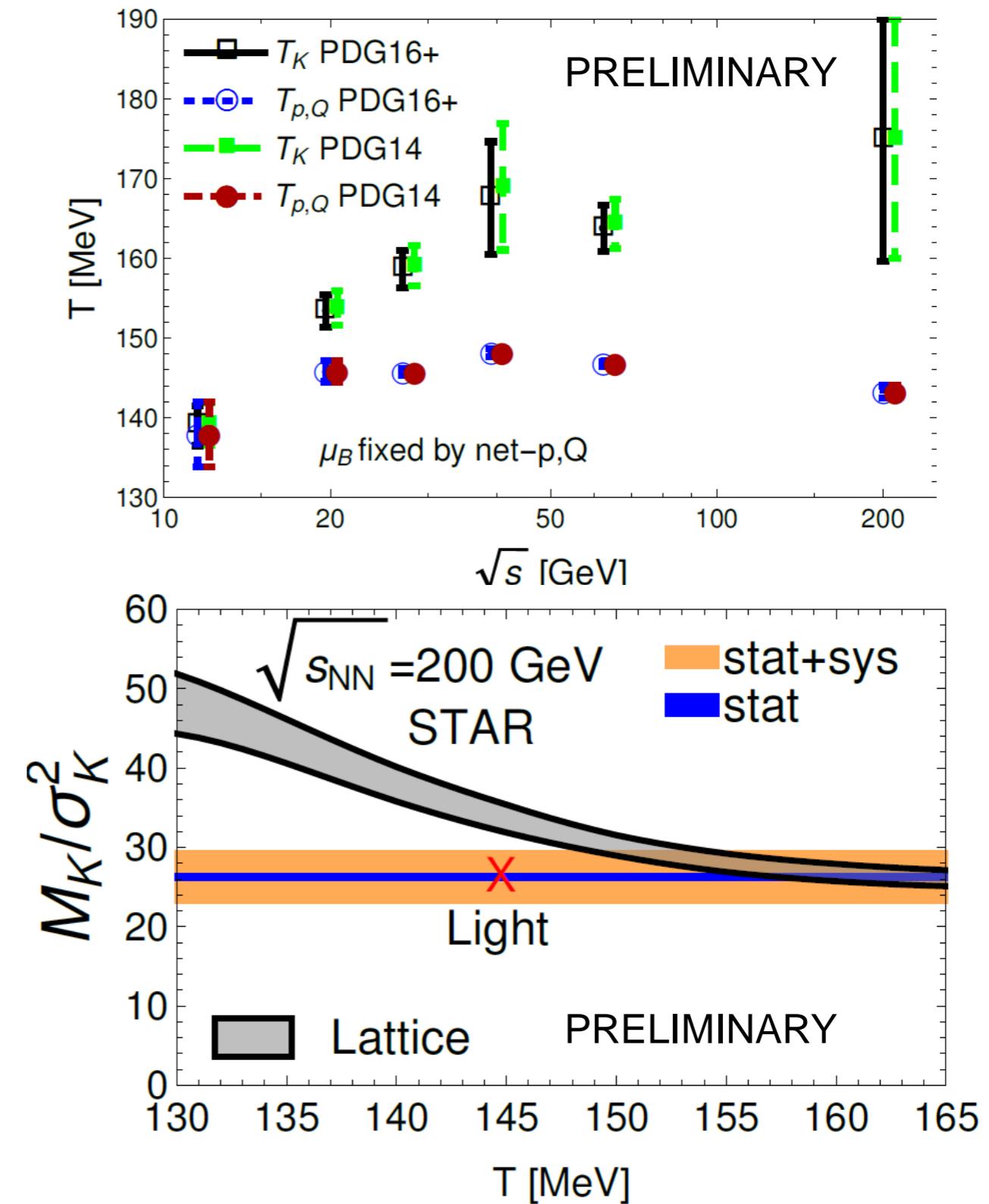
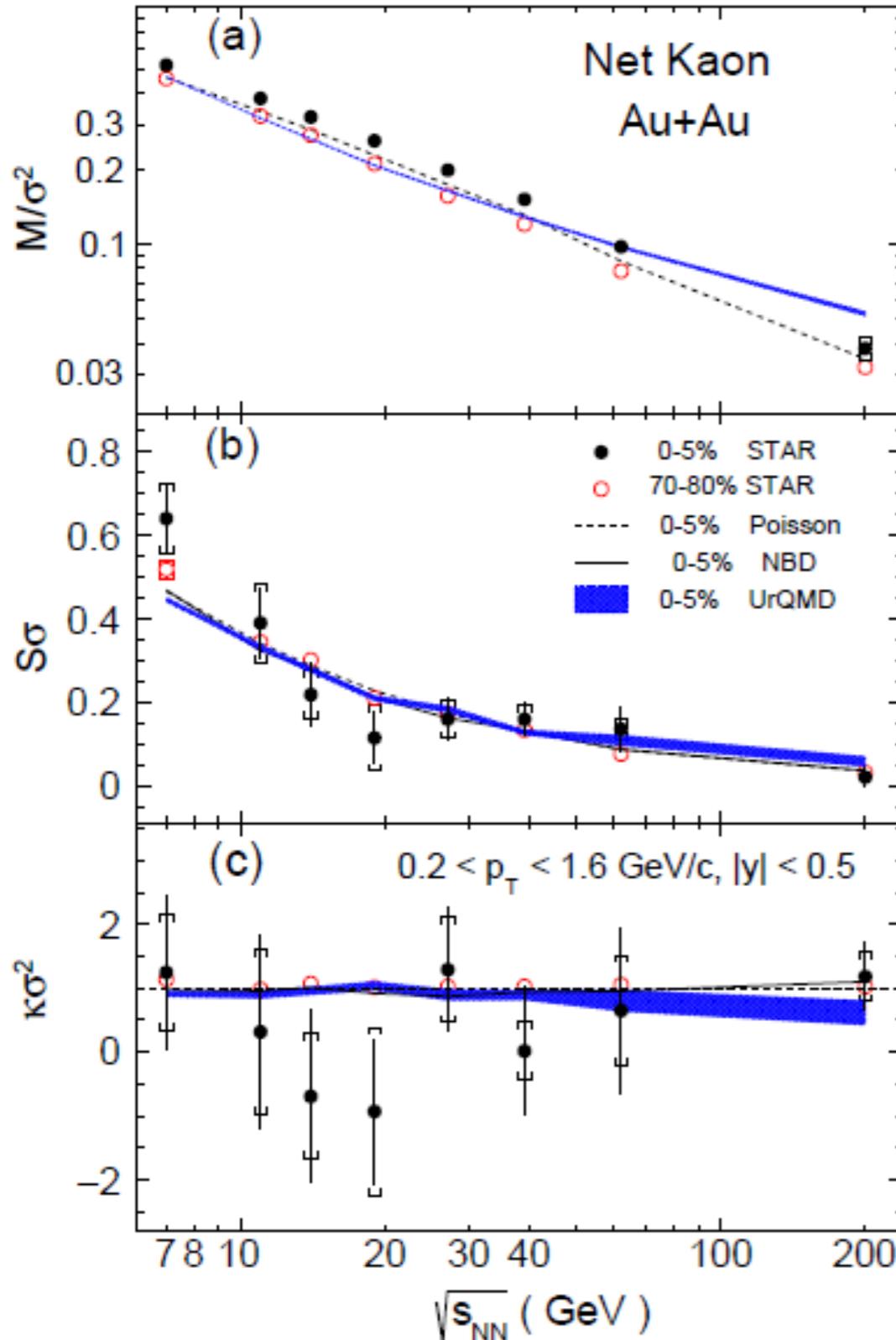
Lattice: improved actions, realistic masses, continuum extrapolations –  
DONE

HRG: implementations improved, better PDG, more states – DONE



The spread in the lattice range is effectively covered by the light vs strange quark susceptibilities

# Fit $\sigma^2/M$ for net-kaons in the same fashion than for net-proton and net-charge (Data: STAR, arXiv:1611.07132)

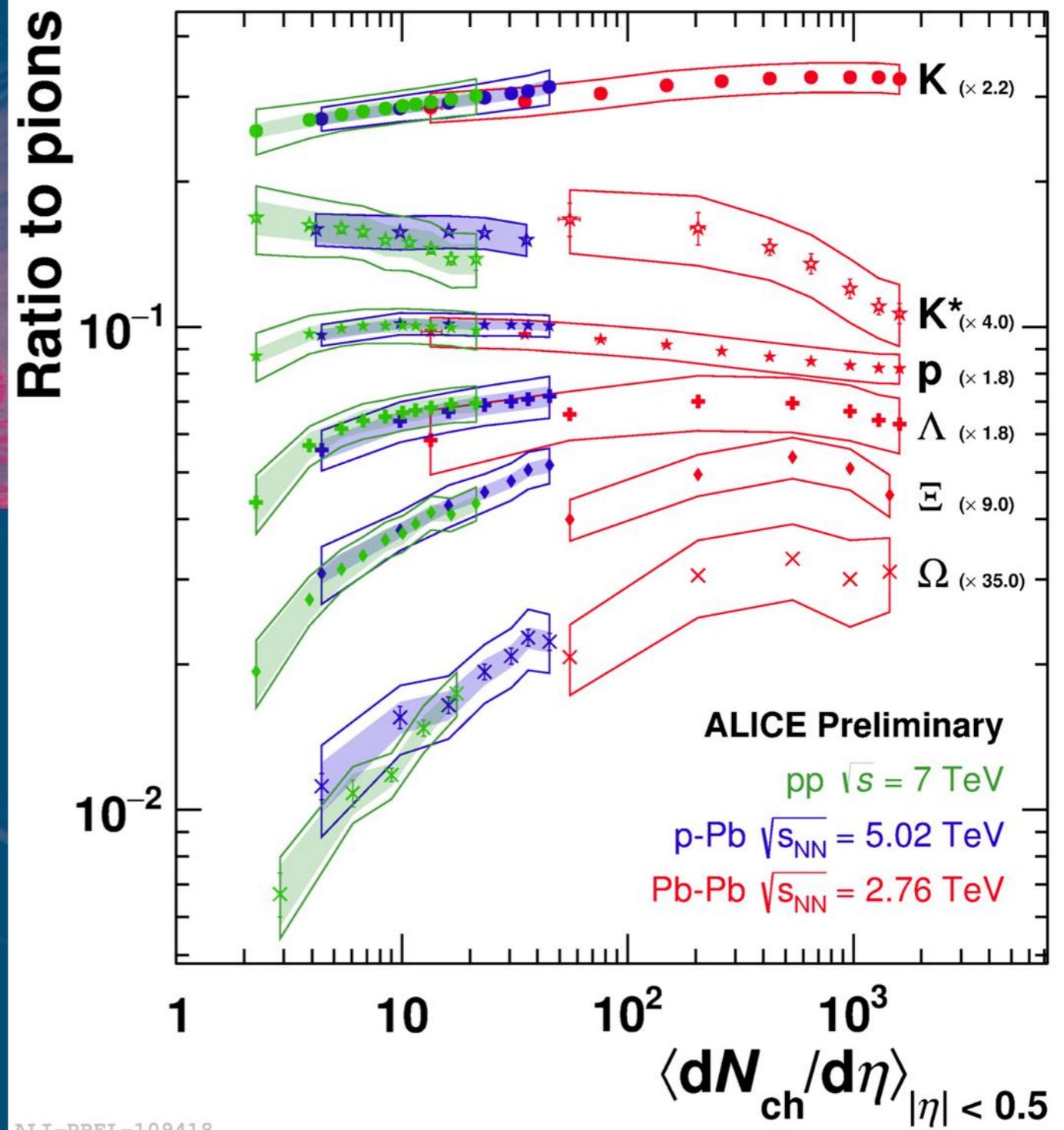
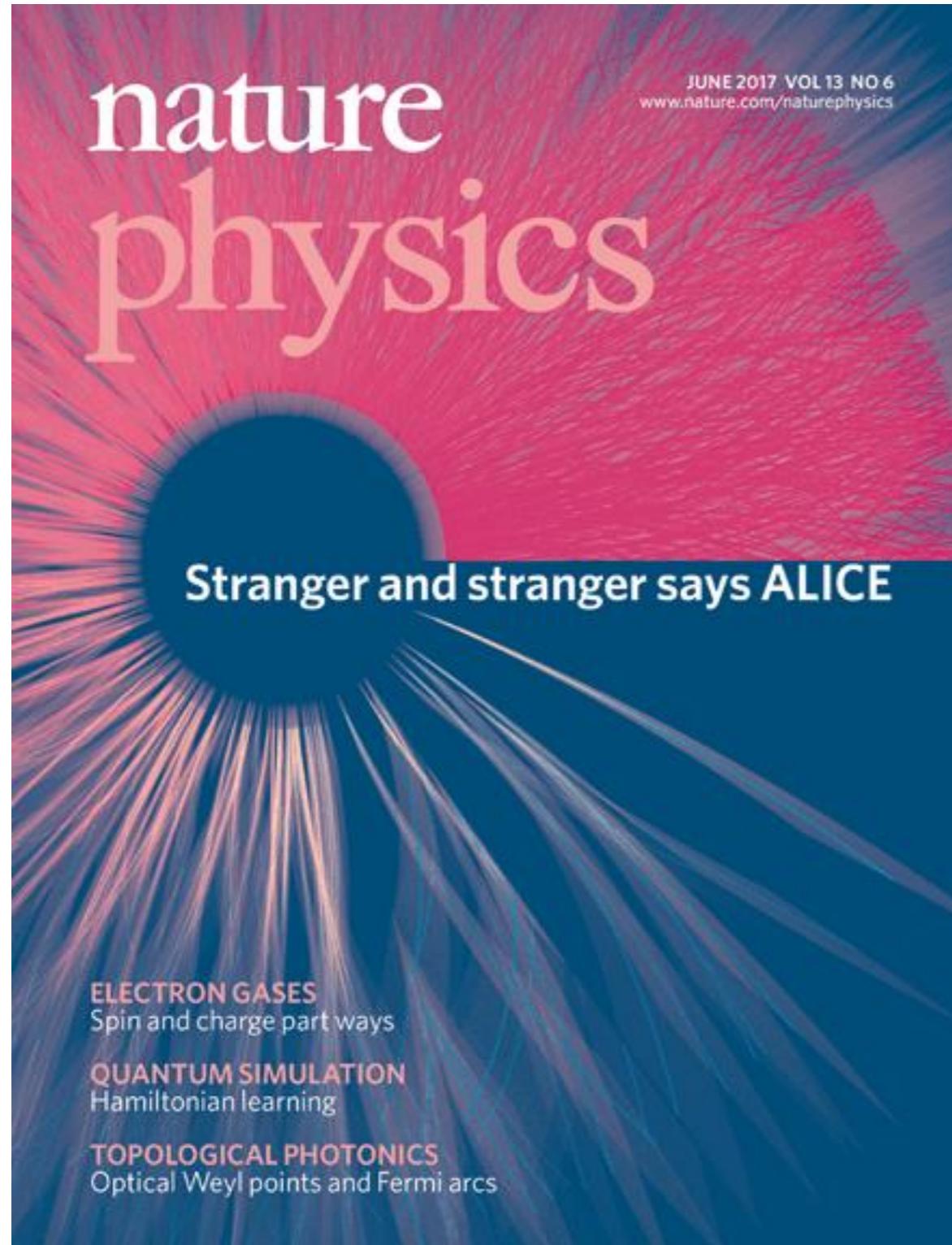


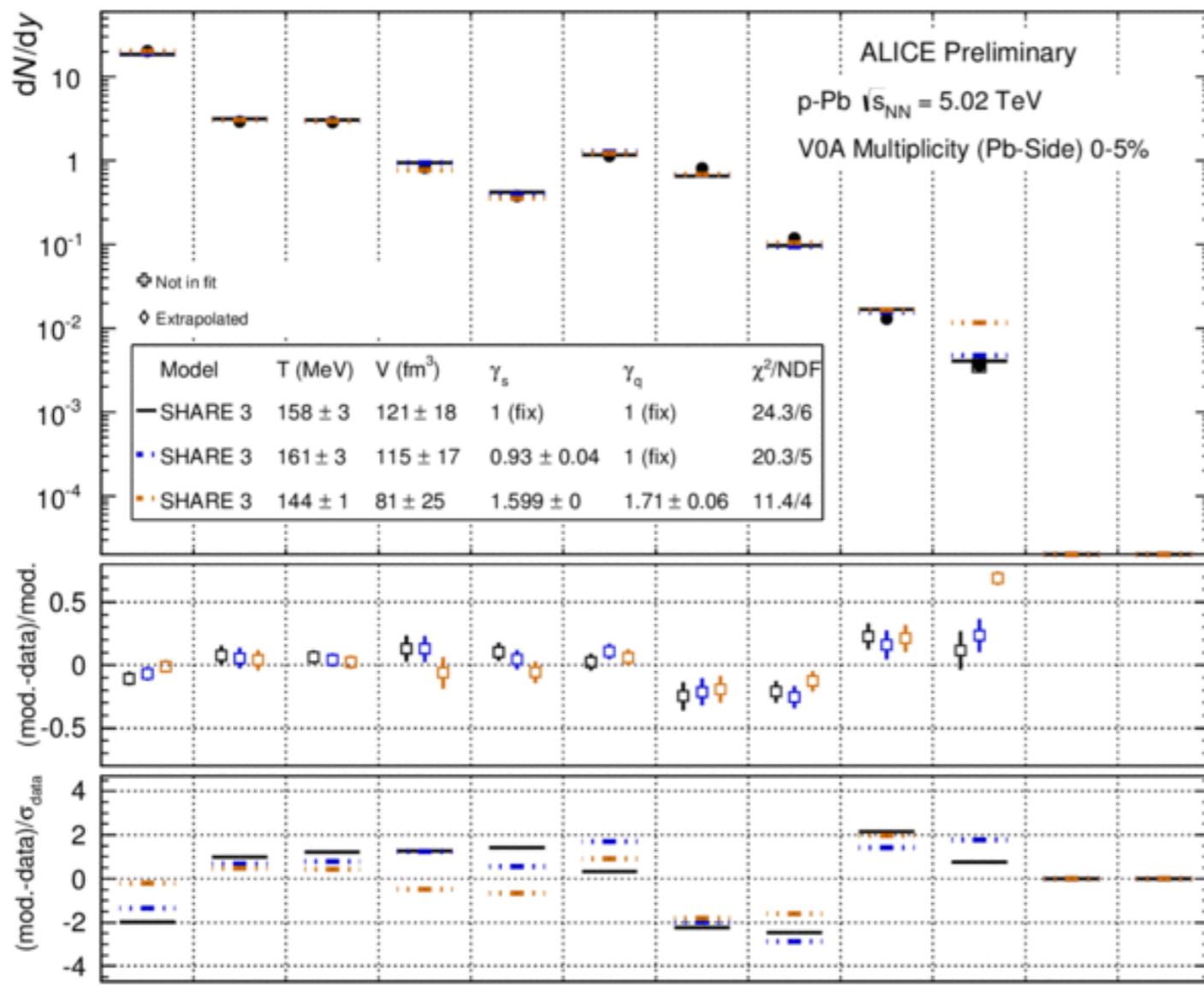
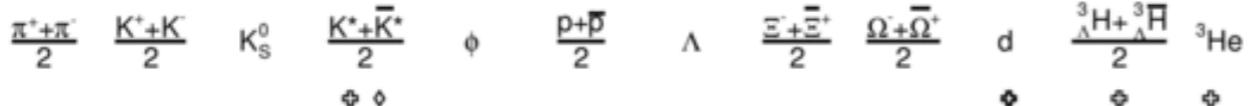
So let's assume there is a separate  
freeze-out hypersurface for strangeness –  
do we care ?

Strange matter creation ?:

- 1.) strangeness enhancement vs.  
suppression
- 2.) strange resonance formation
- 3.) exotica

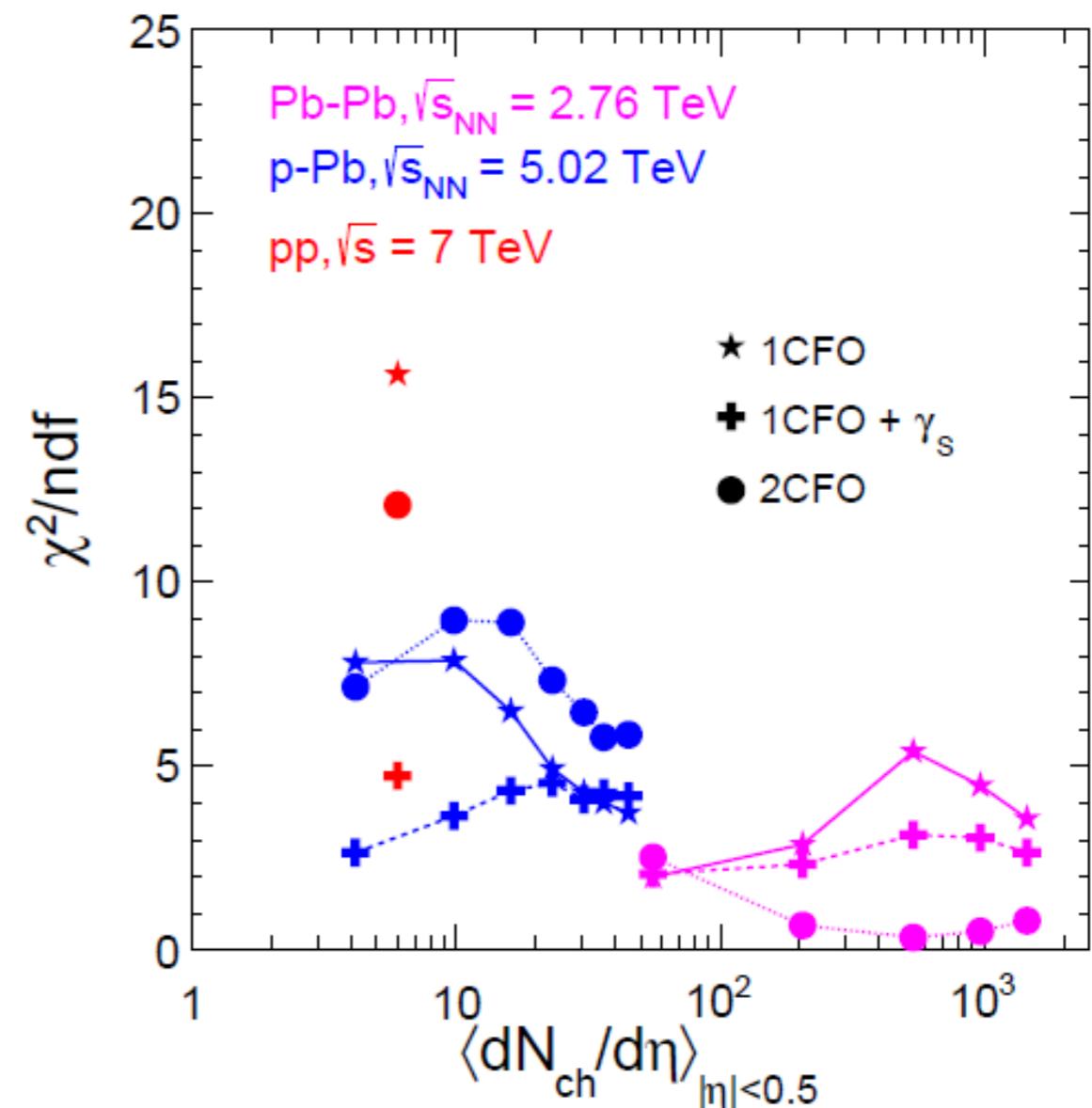
# Stranger and stranger from small to large systems (ALICE, arXiv:1606.07424), see talk by L.Bianchi





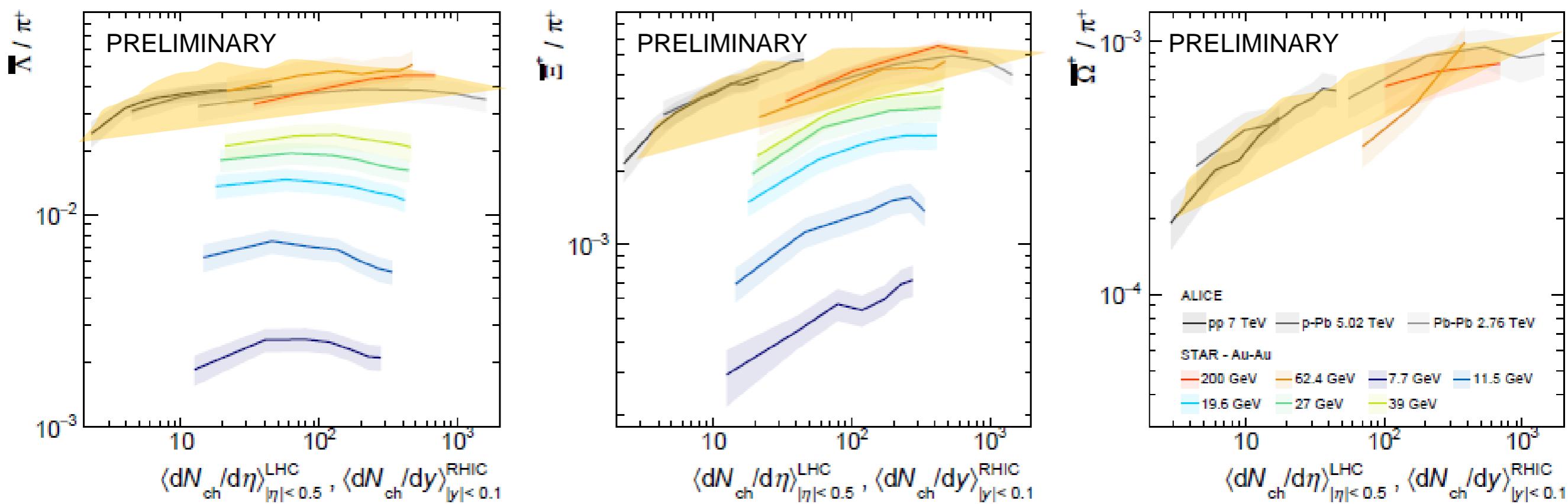
- HRG fit works rather well and yield similar T to PbPb fit (dependent on  $\gamma_s$ )
- Less evidence for flavor dependent freeze-out (Chatterjee, Dash, Mohanty, arXiv:1608.00643)

## Small systems chemistry



# Is it time to re-evaluate strangeness suppression/enhancement ?

Canonical suppression reduces as a function of energy and as a function of system size (Tounsi, Redlich (2001)). Is suppression over at LHC energies ?  
Do we only see enhancement ? Can we distinguish ?



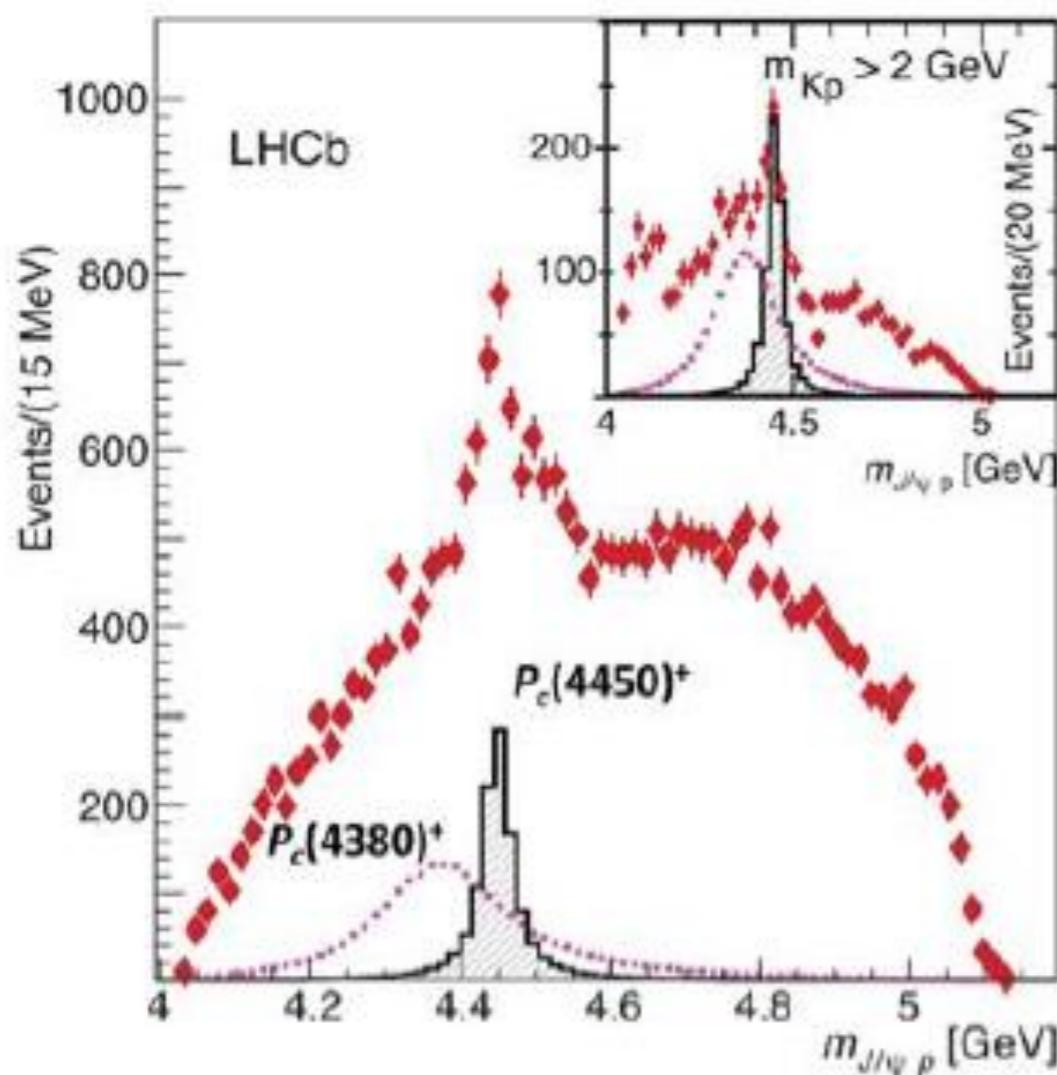
Above 39 GeV the curves seem to fall together, no more energy dependence. The volume dependence is still there ( $\gamma_s$  dependence ?). A higher T freeze-out surface in PbPb will lead to actual strangeness enhancement

# Exotica: Penta- and Tetra-quarks from LHCb

Penta-quark in 2015,  $9\sigma$  evidence by 2016

In the charm sector:  $J/\psi p$  resonance

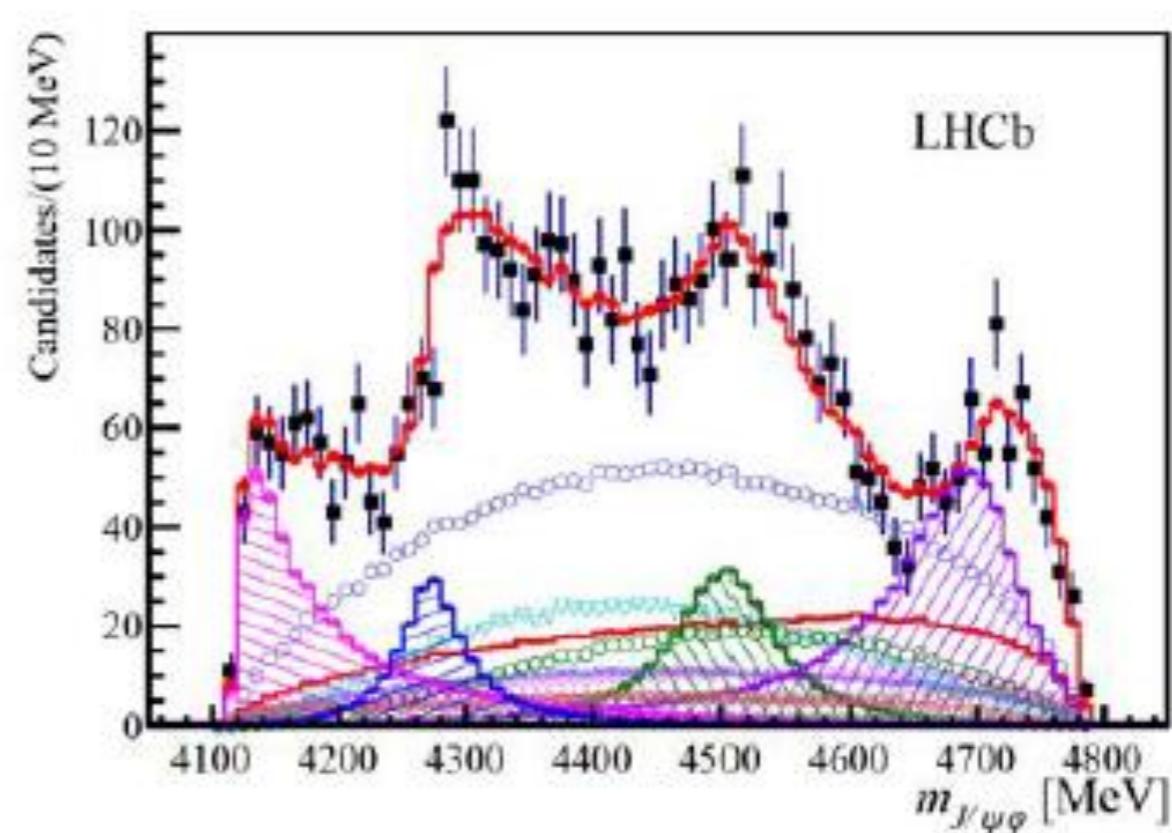
In  $\Lambda_b$  decays to  $J/\psi p K^-$



Tetra-quarks in 2016

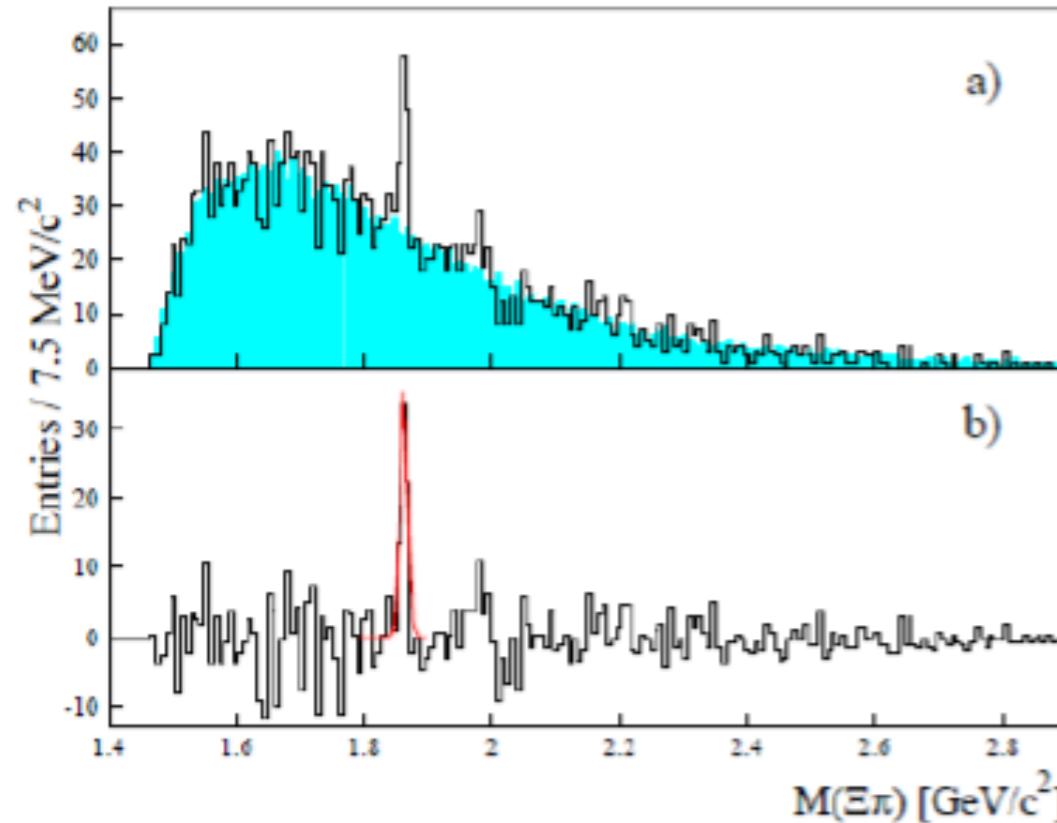
In the charm sector:  $J/\psi \phi$  resonance

In  $B^+$  decays to  $J/\psi \phi K^+$

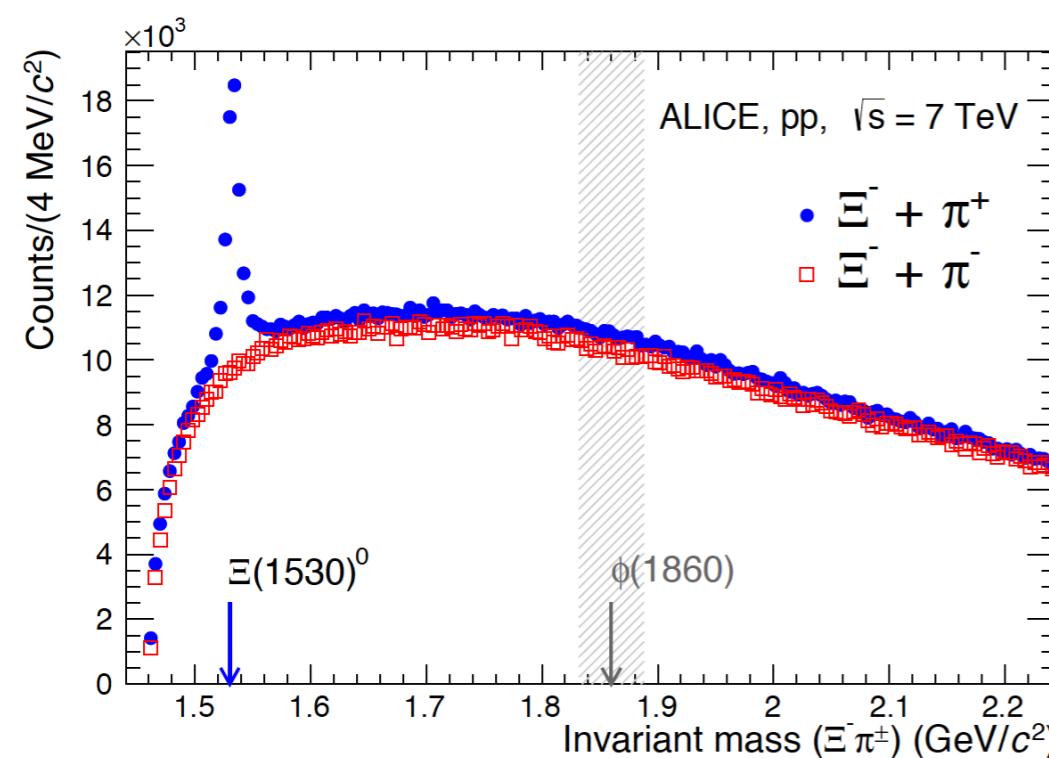


Why nothing in the strange sector ?

# Exotica in strange sector ?



Famous pentaquark candidate from NA49  
(2008) in  $\Xi\pi$  channels ( $\phi(1860)$ ) (dsdsubar)  
Never retracted, never confirmed



No evidence for H-dibaryon or  $\phi(1860)$  in ALICE data.

Maybe we are looking in the wrong channels. In the charm sector all tetra- and penta-quarks seem to require closed charm components.

Keep looking !!

# Conclusions / Outlook

- High precision (continuum limit) lattice QCD susceptibility ratios indicate *flavor separation in the crossover from the partonic to the hadronic matter*.
- There are hints, when comparing to hadron resonance gas and PNJL calculations, that this could lead to a short phase during the crossover in which strange particle formation is dominant.
- If the abundance of strange quarks is sufficiently high (LHC) this could lead to *enhancements in the strange hadron yields (evidence from ALICE)* and it could lead to *strangeness clustering* (exotic states: *dibaryons, strangelets*) or *higher mass strange Hagedorn states* (as predicted by Quark Models).
- Dynamic quantities that evolve during the deconfined phase will be affected as long as the hadronization temperature plays a significant role, i.e. quark phase is shortened for heavier flavors, which could explain flavor effects in  $R_{AA}$  if energy loss builds up near  $T_c$ .
- Ongoing project (JNH, Ratti, Parotto): The phases can be linked in a hydrodynamic calculation by using a mixed EOS from lattice and HRG with varying flavor-dependent switching temperatures.