

A data-drive analysis for heavy quark diffusion coefficient

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Abstract. We apply a Bayesian model-to-data analysis on an improved Langevin framework to estimate the temperature and momentum dependence of the heavy quark diffusion coefficient in the quark-gluon plasma (QGP). The spatial diffusion coefficient is found to have a minimum around 1-3 near T_c in the zero momentum limit, and has a non-trivial momentum dependence. With the estimated diffusion coefficient, our improved Langevin model is able to simultaneously describe the D -meson R_{AA} and v_2 in three different systems at RHIC and the LHC.

1 Introduction

Over the past few decades, facilities at RHIC and the LHC have been exploring the strongly coupled QGP medium by conducting heavy-ion collision experiments. Since the QGP is not directly observable, the study of its properties relies on the comparison between experimental measurements and theoretical modeling. Open heavy flavor, as one of the hard probes of the QGP medium, is valuable yet difficult to study, since it involves measuring ‘rare process’, as well as modeling the complex evolution of heavy quarks during different stages of the collision. Nevertheless, significant progress has been made in recent years: a number of theoretical models are able to describe some of the heavy quark observables, and provide qualitative agreement on the heavy quark diffusion coefficient [1]. However, the simultaneous description of different heavy quark observables in different collision systems still poses challenges for some of the models. Moreover, the current studies of the heavy quark often rely on ‘eye-ball fit’ or ‘ χ^2 -goodness fit’ comparisons between the model calculations and the experimental measurements, leading to very limited distinguishability between the performance of different models, and possible bias on the estimate of the heavy quark diffusion coefficient.

A rigorous and systematic approach to a model-to-data comparison requires optimizing the model and simultaneously describing different experimental observables (which are often highly interdependent and have correlated uncertainties). A Markov chain Monte Carlo (MCMC) algorithm and Bayesian statistics can be used to optimize the process, and perform the comparison and constrain the parameters in a data-driven manner. A Bayesian analysis has already been applied to the study of soft medium properties with great success, for example, the determination of the temperature dependence of shear and bulk viscosities $\eta/s(T)$, $\zeta/s(T)$ [2], or the constraint of the QCD matter equation of state purely from experimental measurements [3]. In this work, we shall extend this type of analysis to the heavy quark sector.

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2 Heavy quark in-medium evolution

Our study of the full space-time evolution of heavy quarks utilizes the well-established improved Langevin framework developed by the Duke QCD group [4]: The initial entropy density and the heavy quark positions are generated by an effective parametric initial condition model TRenTo – mimicking the scaling behavior of IP-Glasma [5]. The heavy quark initial transverse momentum distribution is provided by FONLL [6]; After being produced, heavy quarks propagate in the QGP medium according to an improved Langevin equation [4], which takes into account both collisional and radiative energy loss; at the critical temperature $T_c = 154$ MeV, heavy quarks hadronize into heavy mesons via a hybrid model of recombination and fragmentation. After hadronization, the heavy mesons continuously experience hadronic re-scattering with light hadrons – all hadronic interactions are modeled by UrQMD [8].

The improved Langevin equation that describes the dynamics of heavy quarks inside the QGP medium is:

$$\frac{d\vec{p}}{dt} = -\eta_D(p)\vec{p} + \vec{\xi} + \vec{f}_g, \quad (1)$$

where $-\eta_D(p)\vec{p}$ and $\vec{\xi}$ are the drag and thermal random forces. They are responsible for the heavy quark collisional energy loss and related to the transport coefficient \hat{q} via $\eta_D(p) = \hat{q}/(4TE)$ and $\langle \xi_i(t)\xi_j(t') \rangle = 1/2\hat{q}\delta_{ij}\delta(t-t')$. The third term $\vec{f}_g = -d\vec{p}_g/dt$ is the recoil force when the heavy quark emits bremsstrahlung gluons, with \vec{p}_g being the gluon momentum. We adopt the higher twist formalism for the medium-induced gluon radiation [7] and the recoil force is dependent on \hat{q} . With sufficient large gluon energy limit, the system will equilibrate when the Einstein relationship is satisfied [4]. Under this construction, the drag force, the thermal random force and the recoil force are dependent on the heavy quark transport coefficient \hat{q} . In literature, the spatial diffusion coefficient $D_s = 4T^2/\hat{q}$ is often used to characterize the interacting strength between heavy quarks and the medium.

At high temperature or large momentum, the heavy quark transport coefficient can be calculated using the perturbative QCD approach (pQCD). However, it has been found that the pQCD calculation is not sufficient to explain the suppression experienced by heavy quarks when propagating in the medium, nor able to simultaneously describe both the heavy flavor R_{AA} and v_2 . In order to introduce a non-perturbative effect, we parameterize the spatial diffusion coefficient as a combination of linear temperature dependent component and a pQCD component:

$$D_s 2\pi T(T, \mathbf{p}; \alpha, \beta, \gamma) = \frac{1}{1 + (\gamma^2 p)^2} (D_s 2\pi T)^{\text{soft}} + \frac{(\gamma^2 p)^2}{1 + (\gamma^2 p)^2} (D_s 2\pi T)^{\text{pQCD}}. \quad (2)$$

The soft component $(D_s 2\pi T)^{\text{soft}} = \alpha * (1 + \beta * (T/T_c - 1))$ is the heavy quark diffusion coefficient in the $p = 0$ GeV/c limit. It accounts for the non-perturbative effects. The $(D_s 2\pi T)^{\text{pQCD}} = 8\pi T^3/\hat{q}^{\text{pQCD}}$ is calculated by summing over the leading-order $2 \rightarrow 2$ scattering between heavy quarks ($M_c = 1.3$ GeV) and massless light partons at fixed coupling $\alpha_s = 0.3$. The value of heavy quark diffusion coefficient is determined by three parameters $\mathbf{x} = (\alpha, \beta, \gamma)$, which can be interpreted as following: Parameter α is $D_s 2\pi T(p = 0)$ around T_c , parameter β is its slope above T_c . Parameter γ controls the ratio between the linear component and the pQCD component. A small value of γ indicates the non-perturbative effects are non-negligible towards the high momentum region, and a large value of γ indicates a quick conversion to the pQCD dominated region.

The evolution of heavy quarks depends on the surrounding matter. In this study, the QGP medium is simulated by an event-by-event (2+1)-dimensional viscous hydrodynamical model VISHNEW [9], which also includes shear and bulk viscous corrections. All parameters related to the medium evolution (such as $\eta/s, \zeta/s$) as well as the initial entropy deposition have been calibrated to experimental data of charged particles by an independent Bayesian analysis [2].

3 Statistical analysis and calibration

According to Bayes' theorem, the posterior probability distribution of the model parameters ($P(\mathbf{x}|\mathbf{y}_{\text{exp}})$) is proportional to the product of the parameters' prior distribution ($P(\mathbf{x})$), and the likelihood of observing experimental results given parameter ($L(\mathbf{y}_{\text{exp}}|\mathbf{x})$). In this case, we assume the prior parameter have a uniform distribution, which implies that we don't have any initial constraint/knowledge on the parameters. Thus the posterior distribution of the parameters is proportional to the likelihood:

$$L(\mathbf{y}_{\text{exp}}|\mathbf{x}) \propto \exp\left[-\frac{1}{2}(\mathbf{y} - \mathbf{y}_{\text{exp}})\Sigma^{-1}(\mathbf{y} - \mathbf{y}_{\text{exp}})^T\right] \quad (3)$$

where \mathbf{y}_{exp} is the experimental data (D -meson R_{AA} , v_2 at different centralities at AuAu collisions (200 GeV) and PbPb collisions (2.76 and 5.02 TeV)) [10], \mathbf{y} is the model calculations at parameter $\mathbf{x} = (\alpha, \beta, \gamma)$, Σ is the covariance matrix which includes quantified uncertainties (experimental and theoretical).

In practice, we explore the parameter space and sample the Bayesian posterior distribution by performing a MCMC random walk, which requires evaluating the model at thousands or even millions of parameter values. However, the simulation of heavy quark evolution in heavy-ion collision is computationally highly expensive. Therefore, efficient techniques such as Gaussian process emulators are applied to solve the problem – these act as a fast surrogate model of the full Langevin simulation. A detailed description of the construction of the emulators and performance of the MCMC algorithm can be found in Ref. [2].

4 Posterior results and discussion

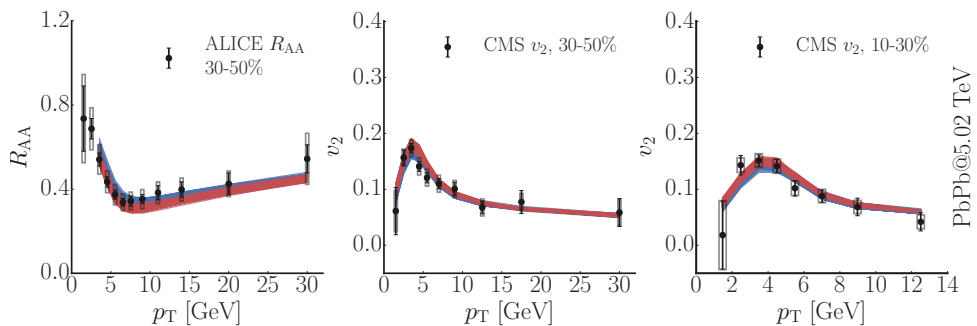


Figure 1: (Color online) 200 random parameters $\mathbf{x} = (\alpha, \beta, \gamma)$ sampled from the posterior distribution and evaluated by emulators. A selections of the output – D -meson R_{AA} and v_2 at PbPb collisions at 5.02 TeV are compared with experimental data.

To verify our analysis, 200 sets of parameters are randomly selected from the posterior distributions and evaluated by the GP emulators. Figure 1 compares a selection of observables with experimental data. Two different colors correspond to two independent Bayesian analysis: the blue lines correspond to the results obtained by calibrating on a single collision system – PbPb collision at 5.02 TeV, the red lines correspond to the results obtained by calibrating on the three collision systems at RHIC and the LHC. After the calibration, our improved Langevin approach is capable of simultaneously describing the experimental data. The spreads of the posterior R_{AA} and v_2 visualize the uncertainties kept in our analysis, which come from the experimental statistic and uncorrelated systematic uncertainties, as well as the uncertainty arising from the prediction made by the GP emulators.

The posterior distributions of parameters (α, β, γ) are plotted in Fig. 2, where the diagonal histograms are the marginal distribution of each parameter with the others integrated out. The off-

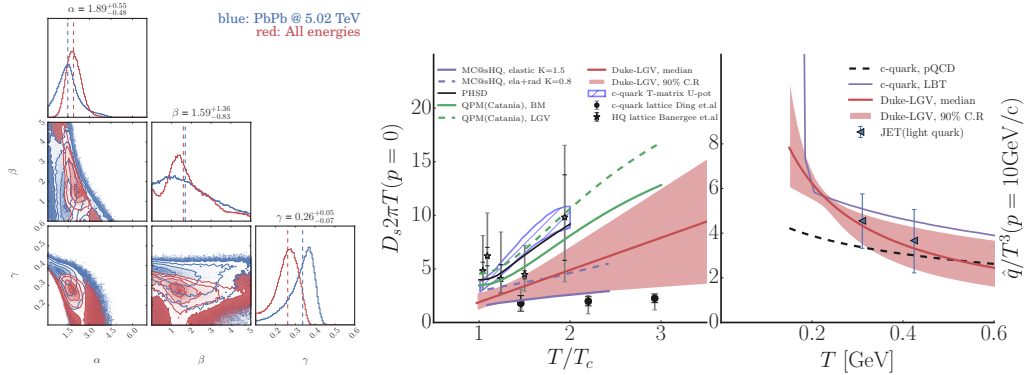


Figure 2: (Color online) **left:** Posterior distributions of parameters (α, β, γ) ; **right:** the heavy quark diffusion coefficients estimated by taking the posterior samples, compared with other models' calculations.

diagonal plots show the correlation between each pair. We observe a narrow distribution for parameters α and γ , implying those two are constraint with current experimental data and uncertainties. Although parameter β is poorly constraint, the strong negative correlation between α and β indicates some collinear interdependency and should not be ignored.

In Fig. 2 we also estimate the diffusion coefficient by taking the posterior samples from the combined analysis (with 90% credibility). Our estimate of $D_s 2\pi T(p=0)$ shows a minimum between (1-3) around T_c , and has a positive temperature dependent. Comparing our estimate with the coefficient used/calculated by a number of other models in the market [1, 11], as well as the lattice QCD calculations [12], our estimate is consistent with lattice QCD calculations within the large uncertainties current inherent in lattice QCD calculations. Although the diffusion coefficient used in different models are rather different, they all have a minimum near T_c with value range $D_s 2\pi T(p=0) \sim (1-7)$. Nevertheless, we should note that in order to make an apple-to-apple comparison among different models, it is essentially important for heavy quarks to propagate in the same QGP medium, especially for those transport models, which is not the case for the current comparison. A detailed investigation is needed, but beyond the scope of this work.

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