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Prague, Czech Republic, 26-30 June 2017



*Investigating Saturation Effects in  
Ultra - Peripheral Heavy Ion Collisions*

*Victor P. Goncalves*

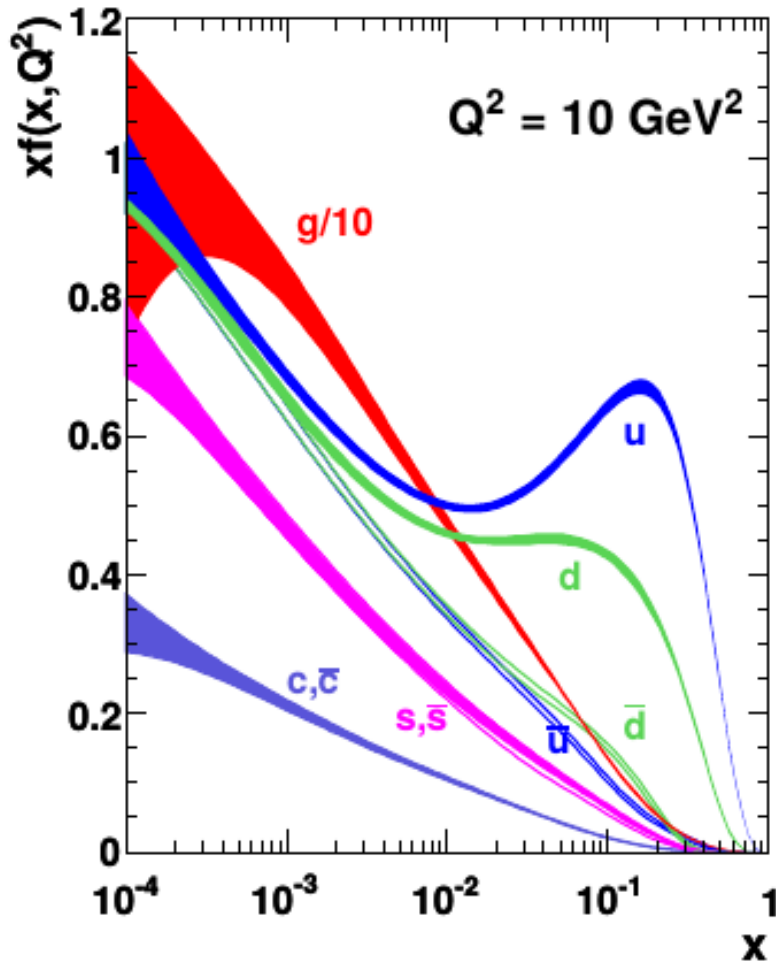
*High and Medium Energy Group - UFPel - Brazil*

**Prague  
30 June 2017**

# Outline

- ✓ Hadronic structure and QCD dynamics at high energies
- ✓ Photon - induced interactions in hadronic collisions
- ✓ Color dipole formalism for the vector meson photoproduction
- ✓ Impact of saturation effects
- ✓ Comparison with the LHC data

# Hadronic structure at high energies



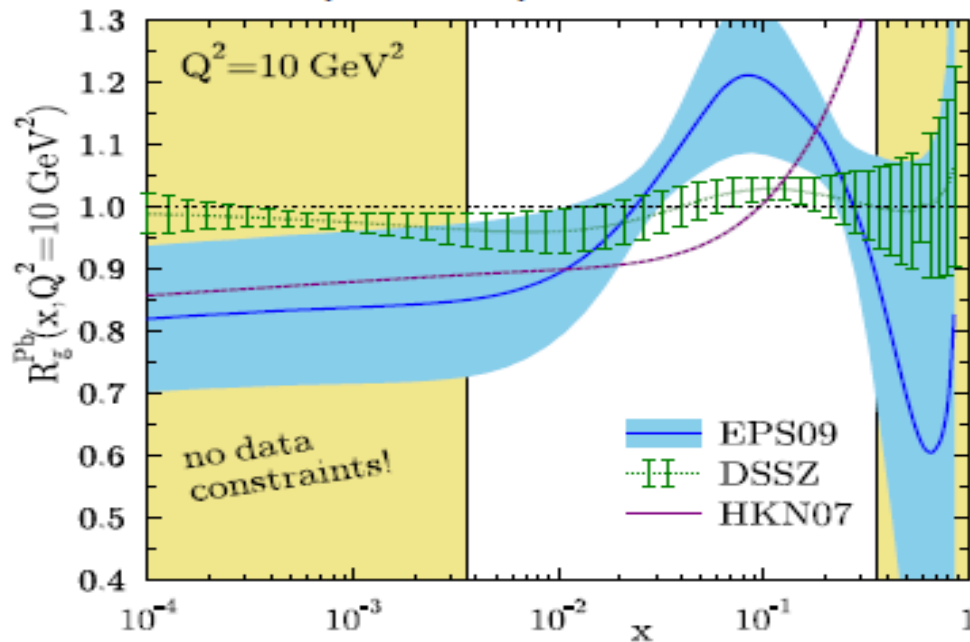
- ✓ Proton structure at high energies (small values of  $x$ ) is dominated by gluons;
- ✓ Linear QCD Evolution equations (DGLAP/BFKL) predict a power growth of gluon distribution at small  $-x$ ;
- ✓ Large uncertainty on the behaviour at small  $-x$ ;
- ✓ The current data included in the global analysis does not constrain the gluon distribution at high energies.

# Hadronic structure at high energies

$$R_g \equiv \frac{xg_A(x, Q^2)}{A \cdot xg_p(x, Q^2)}$$

- No nuclear effects  $\Rightarrow R_g = 1$ .

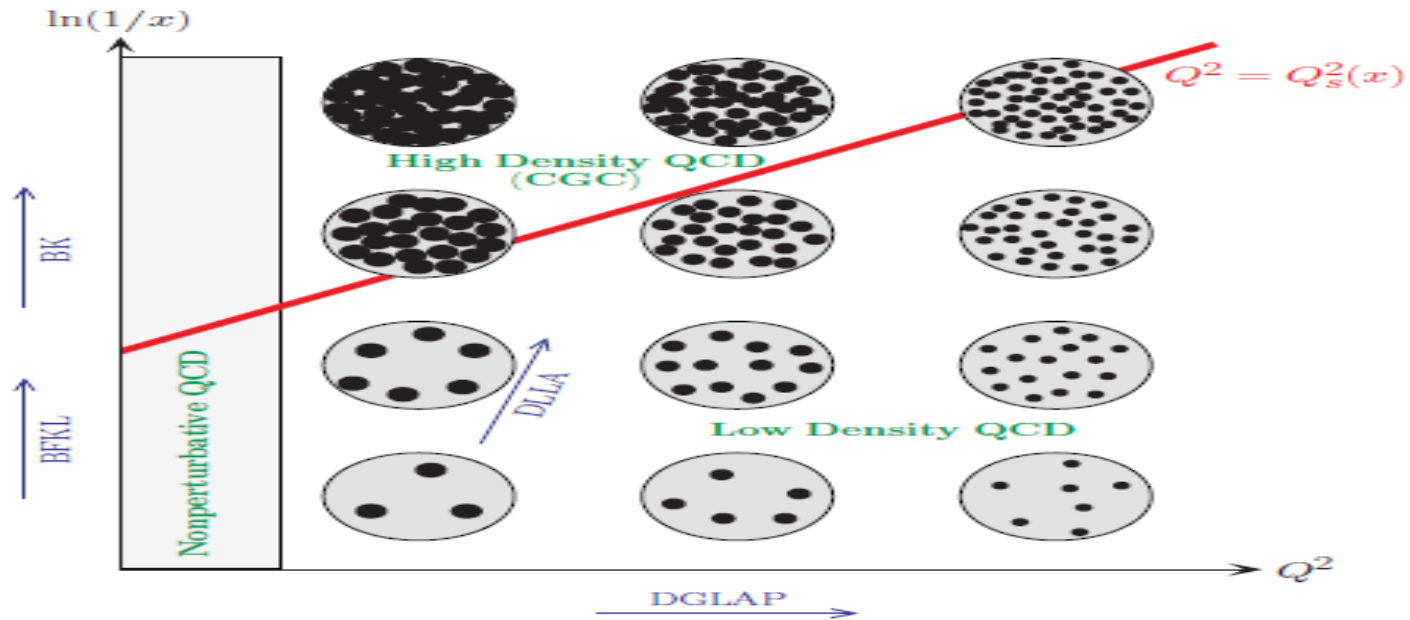
$$R = f_{j/A}(x, Q^2) / [A f_{j/N}(x, Q^2)]$$



Eskola, Puukkunen, arXiv:1401.2345

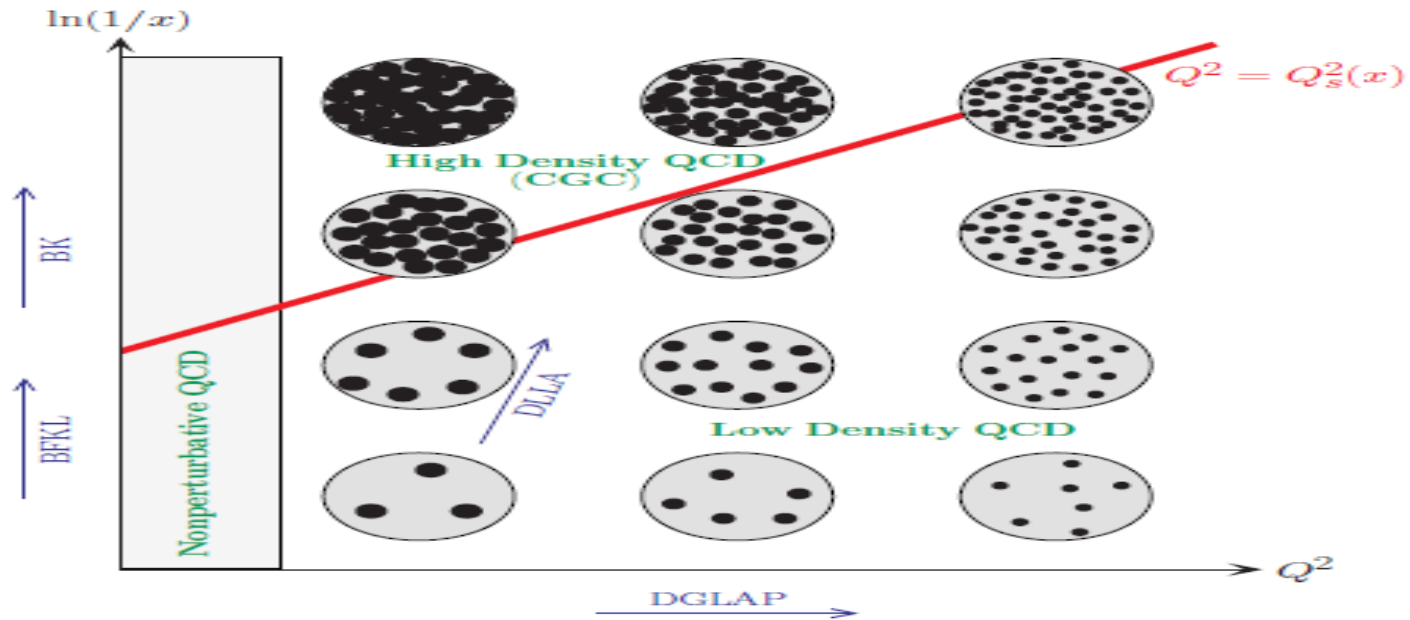
- The current electron - ion experimental data does not constrain the small  $-x$  behaviour;
- Large theoretical uncertainty present in the kinematical range probed by LHC.

# QCD dynamics at high energies



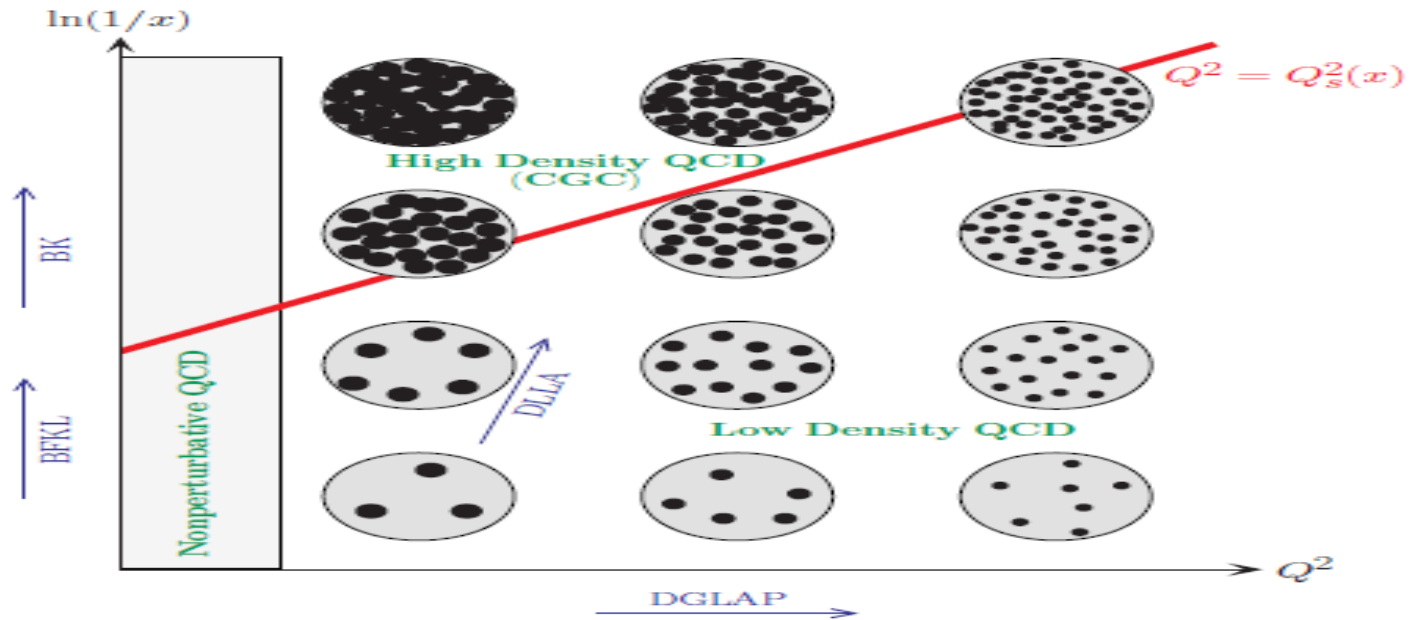
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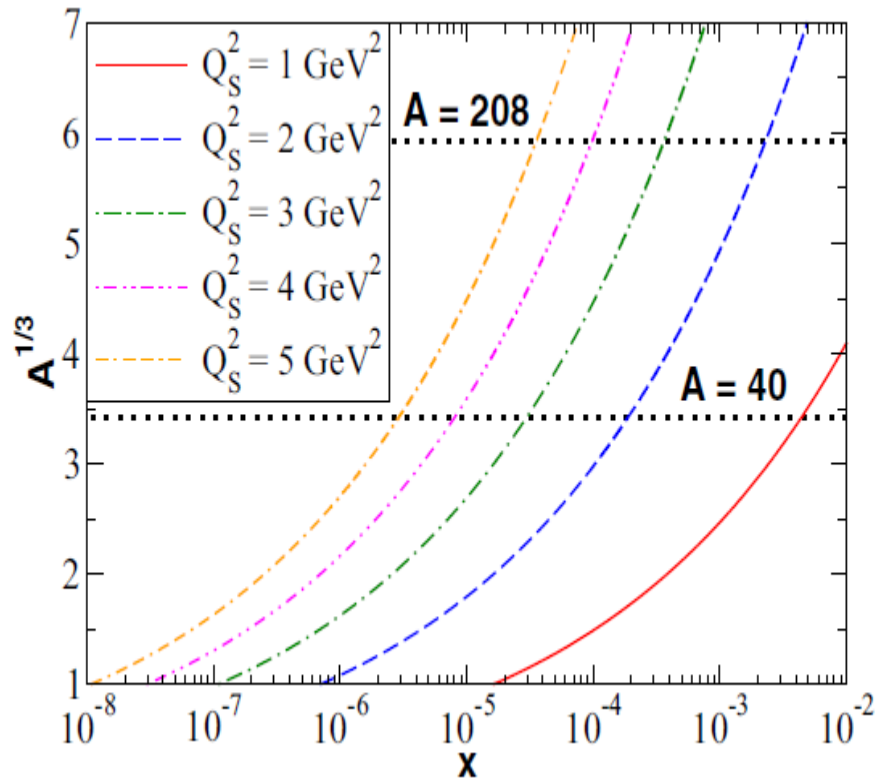
# QCD dynamics at high energies



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- ✓ Saturation scale  $Q_s$  [proportional to  $(1/x)^{\lambda/2} A^\beta$ ] defines the onset of nonlinear QCD dynamics (Gluon saturation effects).



# QCD dynamics at high energies



- ✓ Saturation scale  $Q_s$  [proportional to  $(1/x)^{\lambda/2} A^\beta$ ] defines the onset of nonlinear QCD dynamics (Gluon saturation effects);
- ✓ Nuclei are an efficient amplifier of the gluon saturation effects.



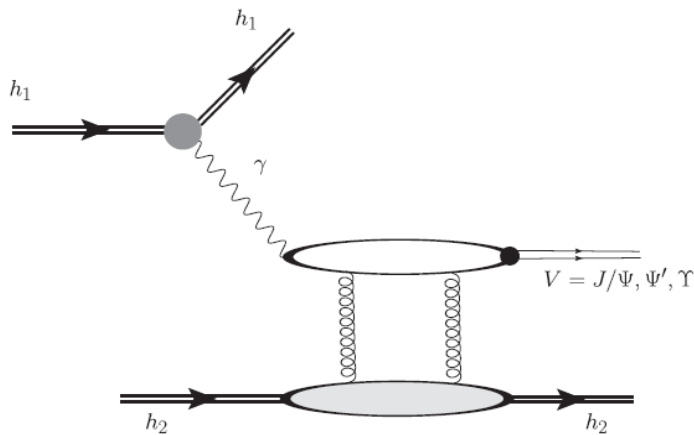
# How to probe the QCD dynamics at high energies ?

- ✓ In the last decades, several observables have been proposed and experimentally analyzed;
- ✓ **Best scenario**: Observables measured in ep and eA collisions.  
(FUTURE!)
- ✓ **Alternative**: Photon - induced interactions in hadronic collisions.

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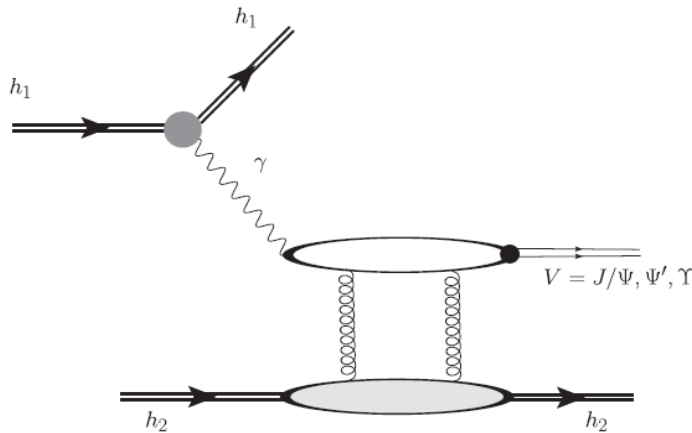
$$\frac{d\sigma [h_1 + h_2 \rightarrow h_1 \otimes V \otimes h_2]}{dY} = \left[ \omega \frac{dN}{d\omega} |_{h_1} \sigma_{\gamma h_2 \rightarrow V \otimes h_2}(\omega) \right]_{\omega_L} + \left[ \omega \frac{dN}{d\omega} |_{h_2} \sigma_{\gamma h_1 \rightarrow V \otimes h_1}(\omega) \right]_{\omega_R}$$



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At leading order in LL(1/x) approx.:

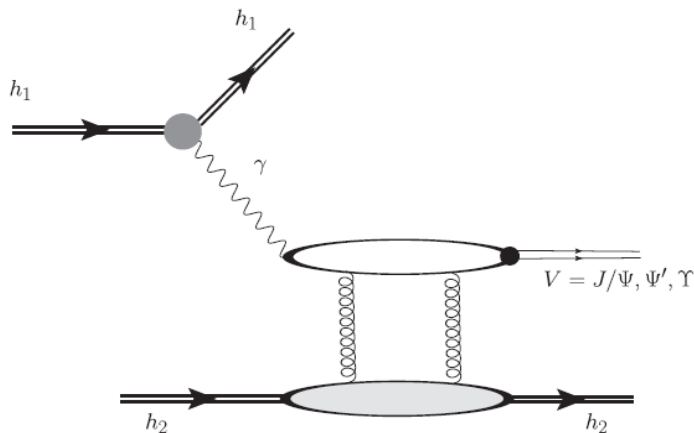
$$\frac{d\sigma^{\gamma h \rightarrow V h}}{dt} \Big|_{t=0} = \mathcal{N} \frac{\pi^3 \Gamma_{e+e-} M_V^3}{48 \alpha_{em}} \left[ \frac{\alpha_s(\bar{Q}^2)}{\bar{Q}^4} x g_h(x, \bar{Q}^2) \right]^2$$

Cross section is proportional to the **square** of the hadron gluon distribution at  $x = 4\bar{Q}^2/W^2$

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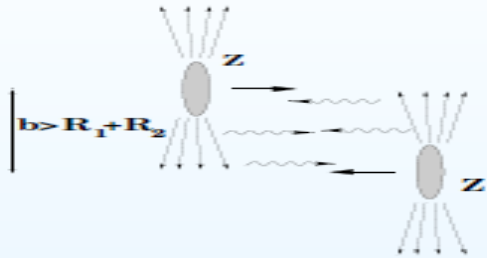
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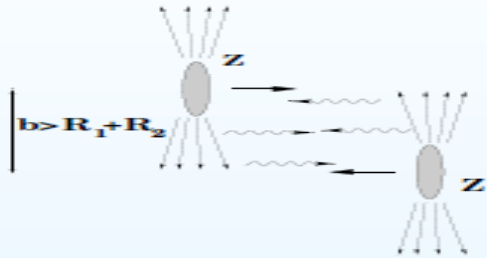
Center of mass energies

1.  $\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$
2.  $\gamma\gamma$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_1(\omega) \otimes n_2(\omega) \otimes \sigma^{\gamma\gamma \rightarrow X}(W_{\gamma\gamma})$

LHC	$pp$	$W_{\gamma p} \lesssim 8390 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 4504 \text{ GeV}$
LHC	$pPb(Ar)$	$W_{\gamma A} \lesssim 1500 (2130) \text{ GeV}$	$W_{\gamma\gamma} \lesssim 260 (480) \text{ GeV}$
LHC	$PbPb$	$W_{\gamma A} \lesssim 950 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 160 \text{ GeV}$
HERA	$ep$	$W_{\gamma p} \lesssim 200 \text{ GeV}$	-

Photoproduction in  $pp$  collisions at LHC probes photon - hadron center - of - mass energies one order of magnitude larger than HERA.

# LHC = Photon collider



Center of mass energies

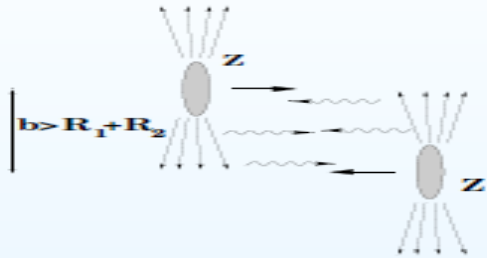
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Photoproduction in pA and AA collisions at LHC probes a unexplored regime of photon - nucleus center of mass energies.



# LHC = Photon collider



Center of mass energies

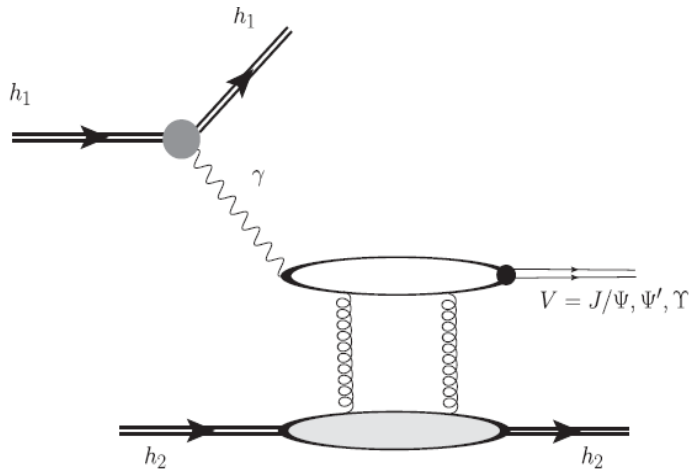
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Photon - induced interactions at LHC allows to study the high energy regime of QCD (Small - x Physics).

# Diffraction vector meson photoproduction in UPHIC: Color Dipole Formalism

$$\frac{d\sigma [h_1 + h_2 \rightarrow h_1 \otimes V \otimes h_2]}{d^2b dy} = [\omega N_{h_1}(\omega, b) \sigma_{\gamma h_2 \rightarrow V \otimes h_2}(\omega)]_{\omega_L} + [\omega N_{h_2}(\omega, b) \sigma_{\gamma h_1 \rightarrow V \otimes h_1}(\omega)]_{\omega_R}$$

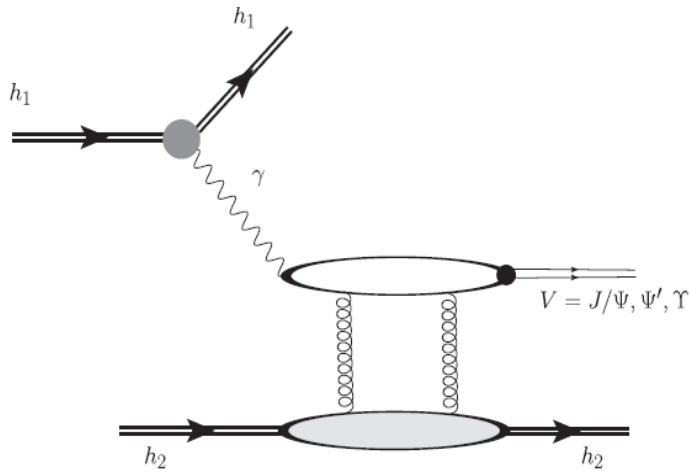


$$\sigma(\gamma h \rightarrow V h) = \int_{-\infty}^0 \frac{d\sigma}{dt} dt = \frac{1}{16\pi} \int_{-\infty}^0 |\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta)|^2 dt$$

$$\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta) = i \int dz d^2r d^2b_h e^{-i[\mathbf{b}_h - (1-z)\mathbf{r}] \cdot \Delta} (\Psi^{V*} \Psi)_T 2\mathcal{N}_h(x, \mathbf{r}, \mathbf{b}_h)$$

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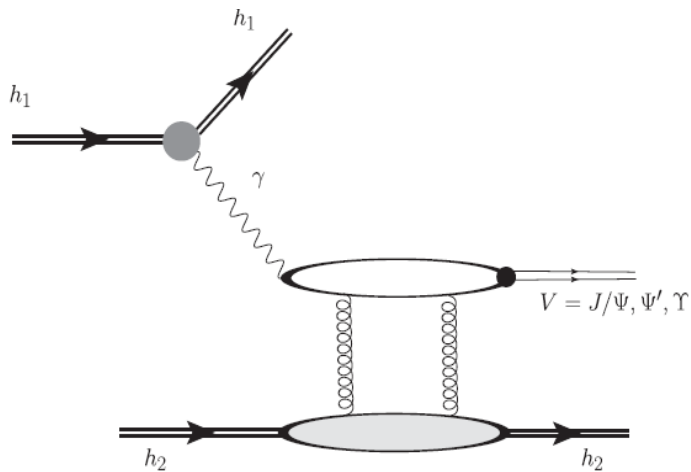
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Overlap function for Vector Mesons

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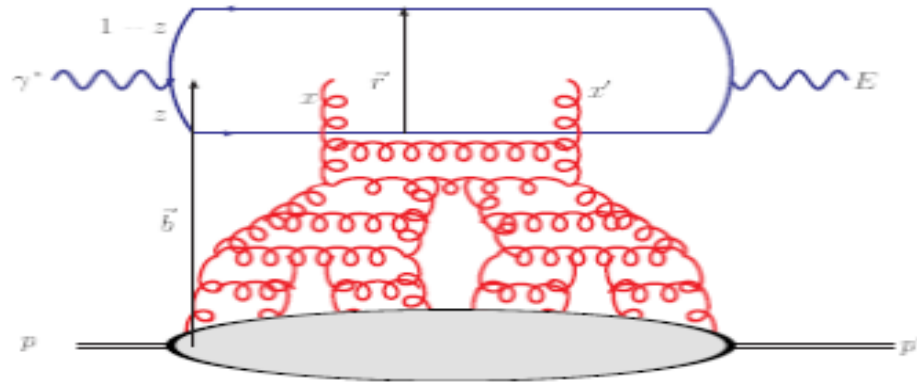
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Forward dipole - hadron scattering amplitude: Determined by the QCD dynamics

# Diffractive vector meson photoproduction in UPHIC: Color Glass Condensate Formalism

Dipole - proton scattering amplitude:



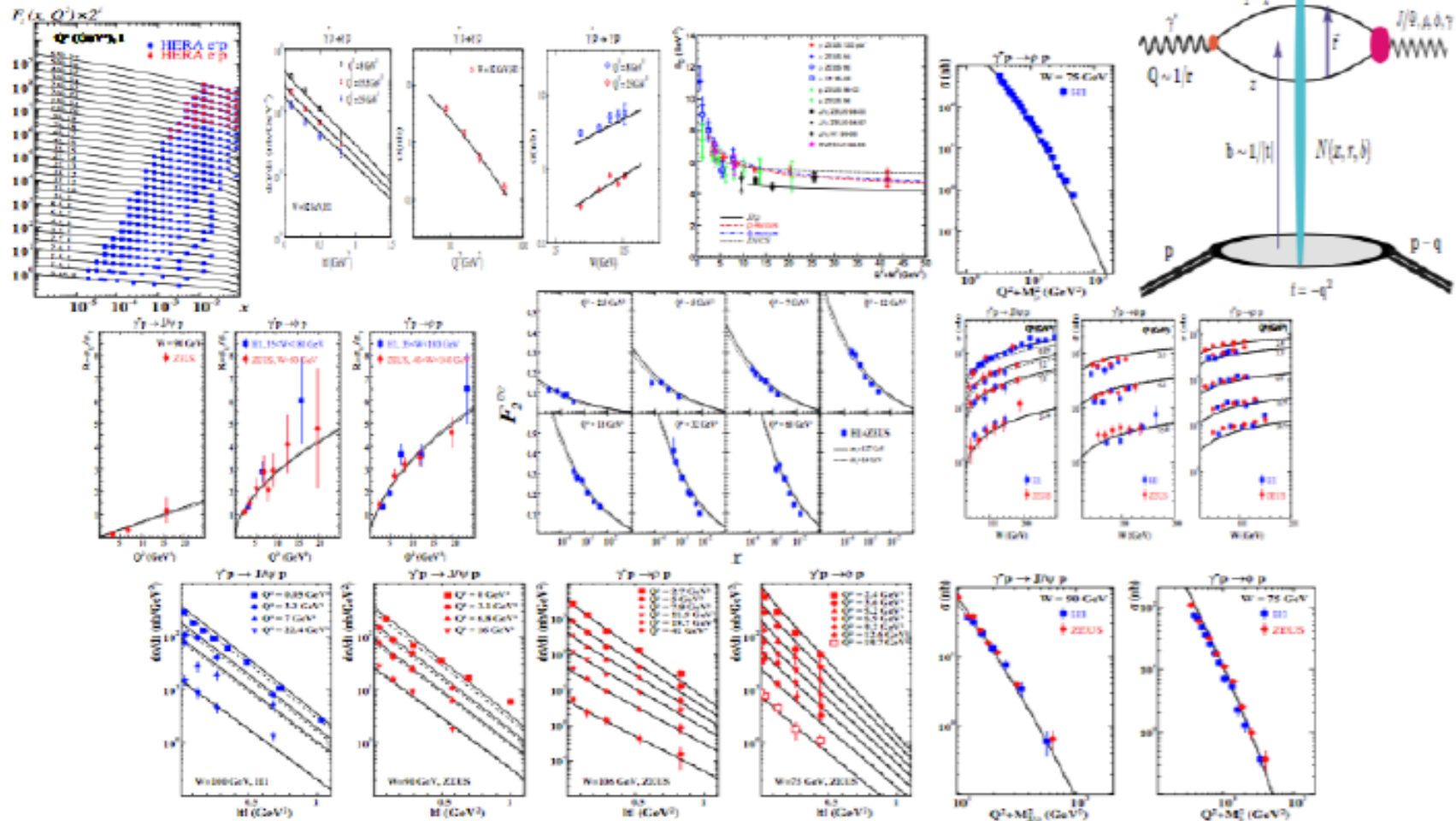
- bCGC : 
$$\mathcal{N}^P(\hat{x}, \mathbf{r}, \mathbf{b}) = \begin{cases} \mathcal{N}_0 \left( \frac{rQ_s(b)}{2} \right)^{2(\gamma_s + \frac{\ln(2/rQ_s(b))}{\kappa\lambda Y})} & rQ_s(b) \leq 2 \\ 1 - e^{-A \ln^2(BrQ_s(b))} & rQ_s(b) > 2 \end{cases}$$

- Proposed originally by Kowalski, Motyka and Watt (06)
- Parameters of the model updated considering the high precision combined HERA data (Rezaeian, Schmidt, 13)

# A unified description of combined inclusive HERA data & diffractive data in CGC

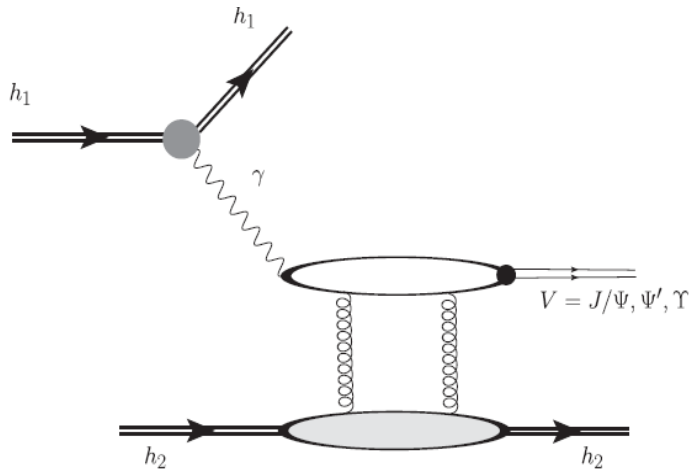
Rezaeian, Siddikov, Van de Klundert, Venugopalan, arXiv:1212.2974; Rezaeian, Schmidt, arXiv:1307.0825

The dipole scattering amplitude is the main ingredient with 3 or 4 free parameters fixed via a fit to the reduced cross-section.



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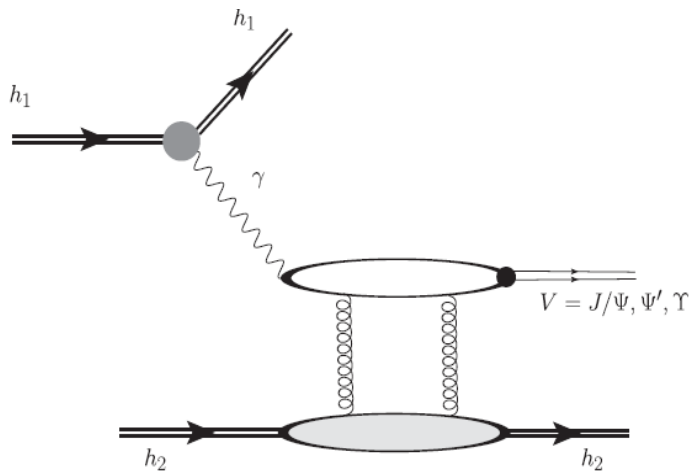
- Nucleus:  $\mathcal{N}_A(x, \mathbf{r}, \mathbf{b}_A) = 1 - \exp \left[ -\frac{1}{2} \sigma_{dp}(x, \mathbf{r}^2) T_A(\mathbf{b}_A) \right]$   $\rightarrow$  Sums all multiple elastic rescatterings of the dipole.

$$\sigma_{dp}(x, \mathbf{r}^2) = 2 \int d^2b_p \mathcal{N}_p(x, \mathbf{r}, \mathbf{b}_p)$$



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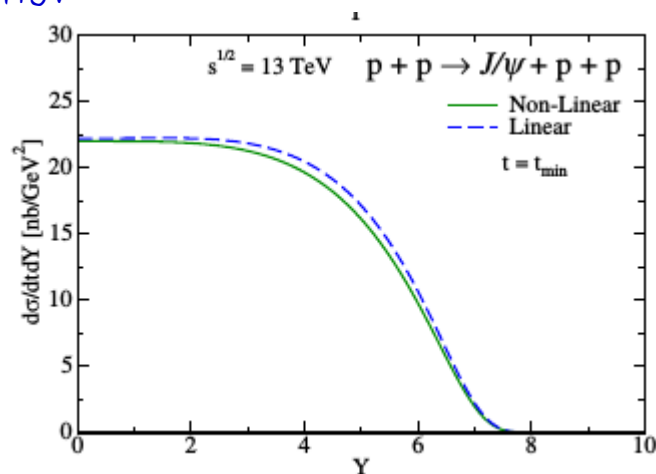
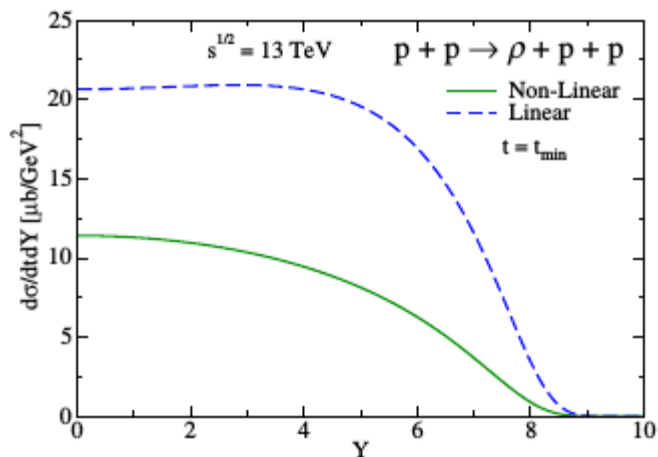
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In the dipole picture, all free parameters have been constrained by HERA data. Predictions for UPHIC are parameter free!

# Diffraction vector meson photoproduction in UPHIC: Impact of the gluon saturation effects

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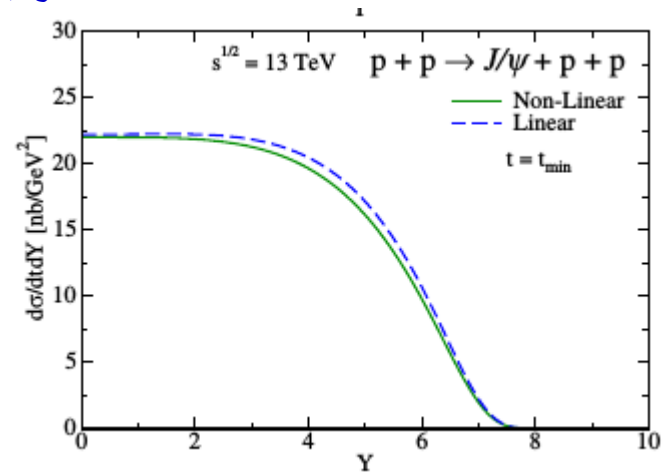
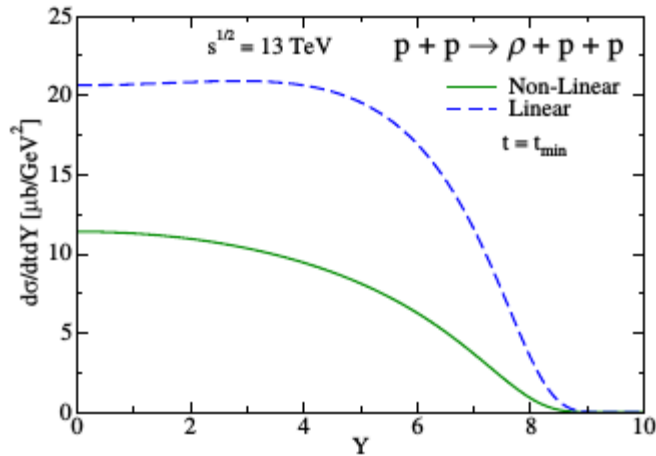
pp Collisions:



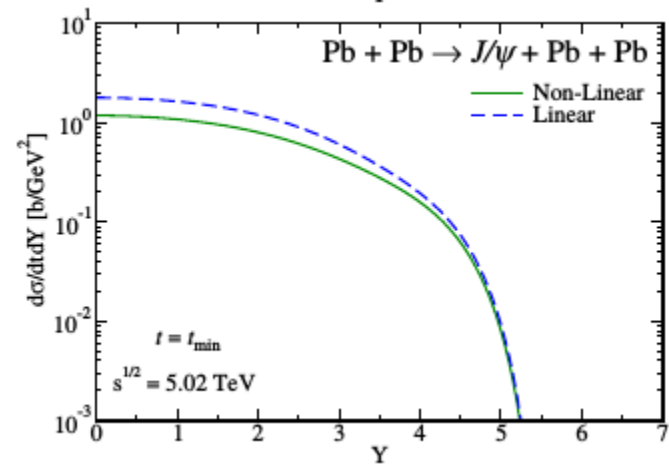
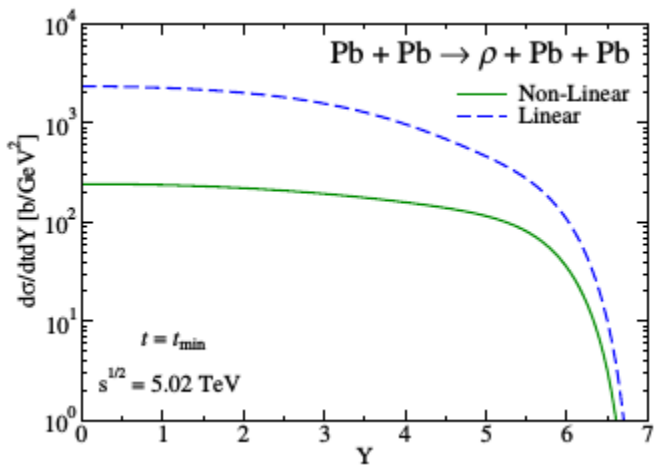
$$t_{\min} = -m_N^2 M_V^4 / W^4$$

# Diffractive vector meson photoproduction in UPHIC: Impact of the gluon saturation effects

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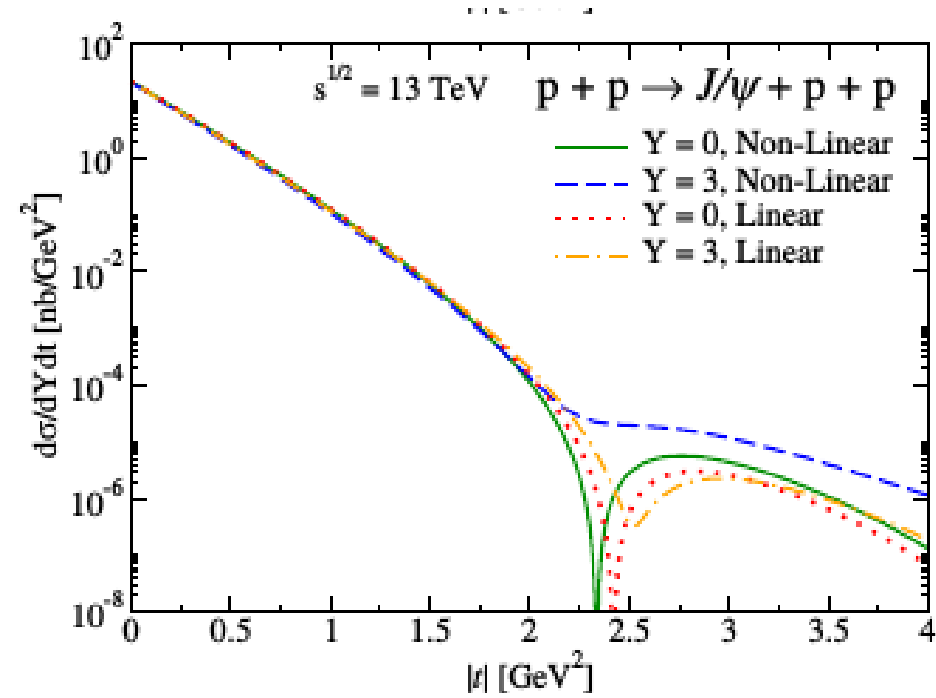
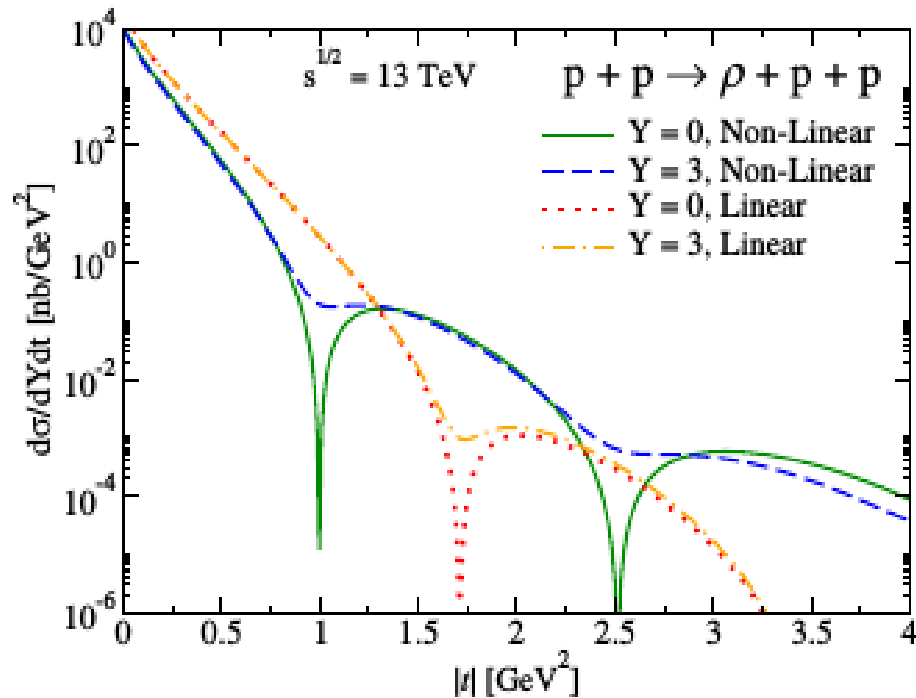
PbPb Collisions:



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pp Collisions:



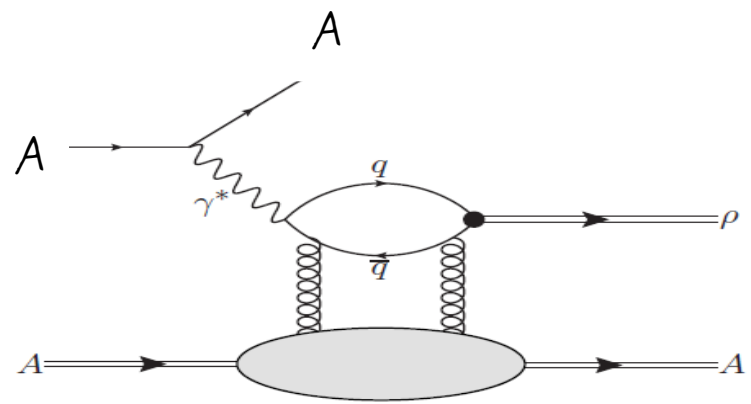
Linear model:

$$\mathcal{N}^p(x, r, b_p) = \mathcal{N}_0 \left( \frac{r Q_s(b_p)}{2} \right)^{2 \left( \gamma_s + \frac{\ln(2/r Q_s(b_p))}{\kappa \lambda Y} \right)}$$

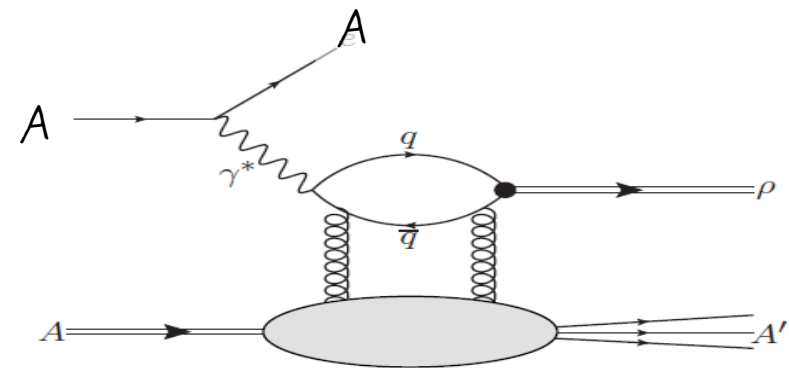
# Diffractive vector meson photoproduction in UPHIC: Impact of the gluon saturation effects

PbPb Collisions:

Coherent production:

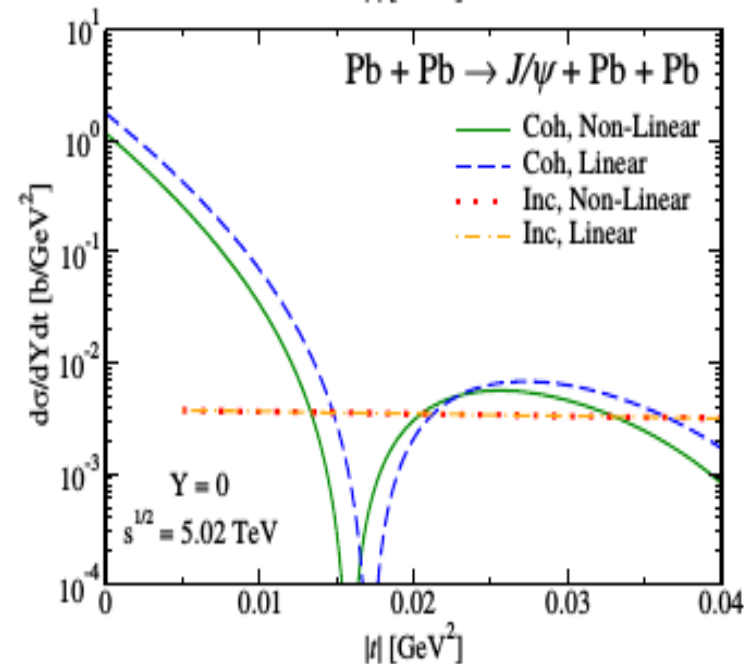
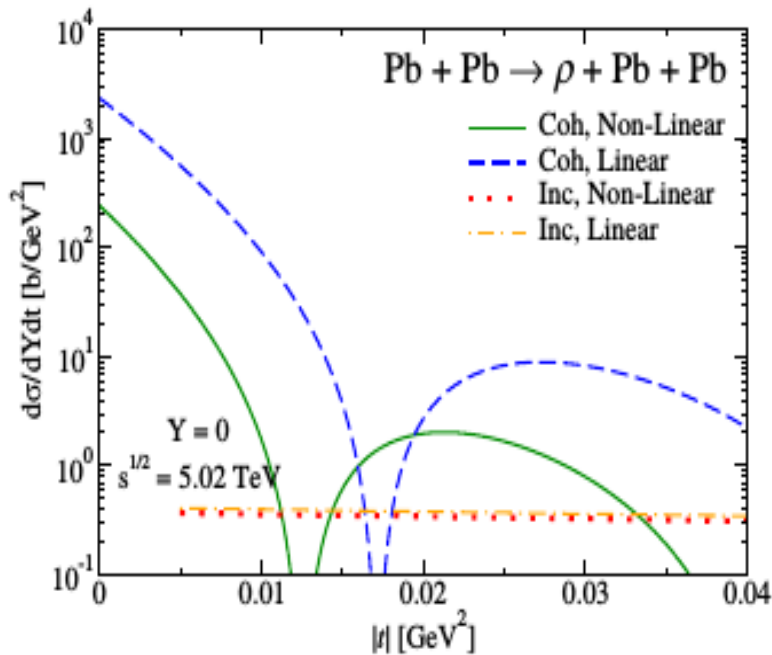


Incoherent production:



# Diffractive vector meson photoproduction in UPHIC: Impact of the gluon saturation effects

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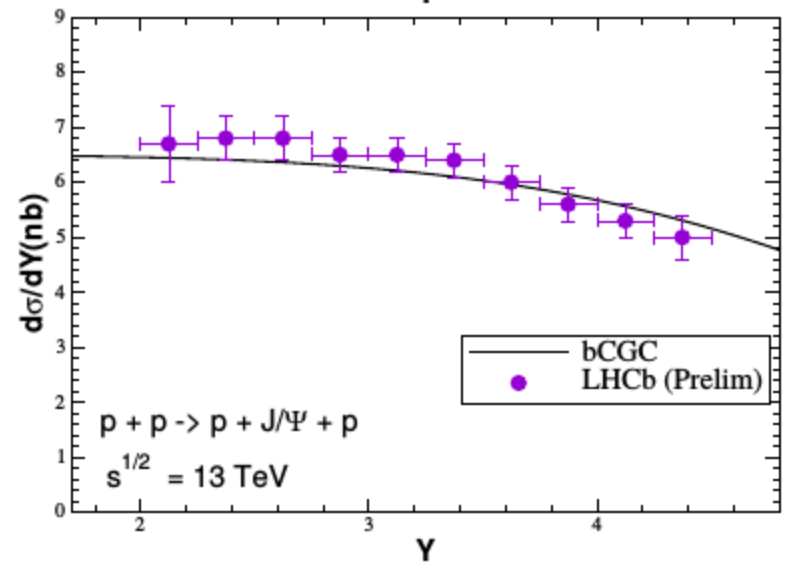
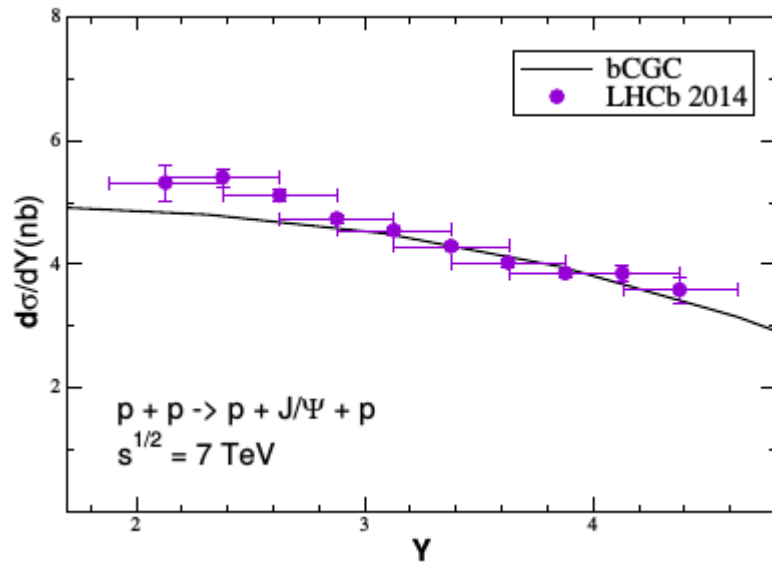
Linear model: 
$$\mathcal{N}^A(x, r, \mathbf{b}_A) = \frac{1}{2} \sigma_{dp}(x, r) A T_A(\mathbf{b}_A)$$



Comparison with the LHC data

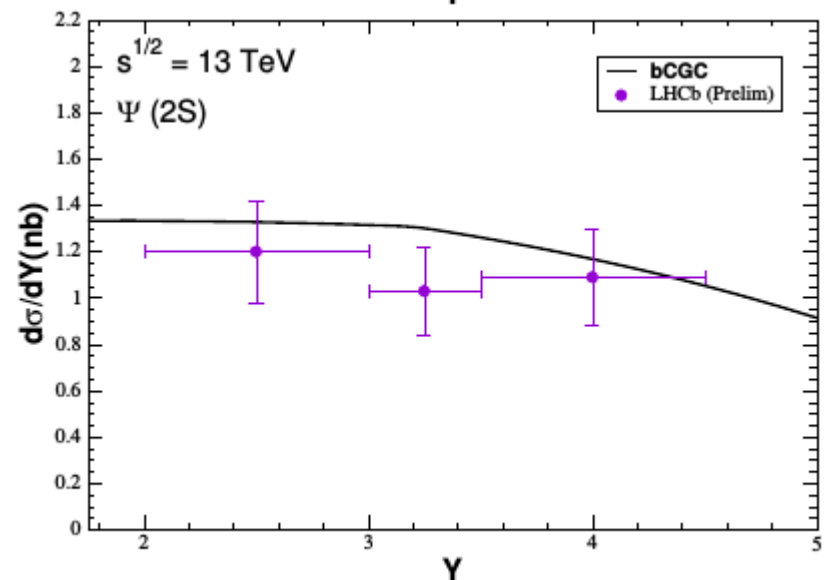
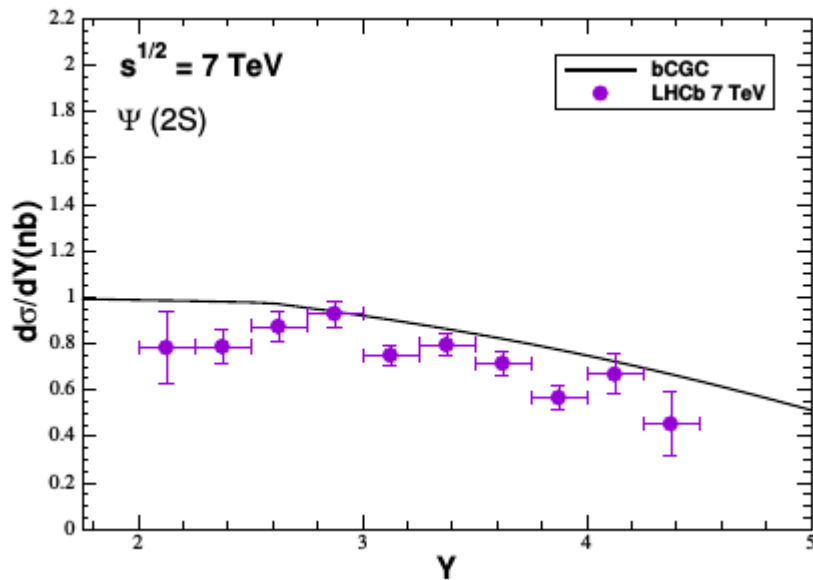
# Comparison with the LHC data

Diffractive J/Psi photoproduction in pp collisions:



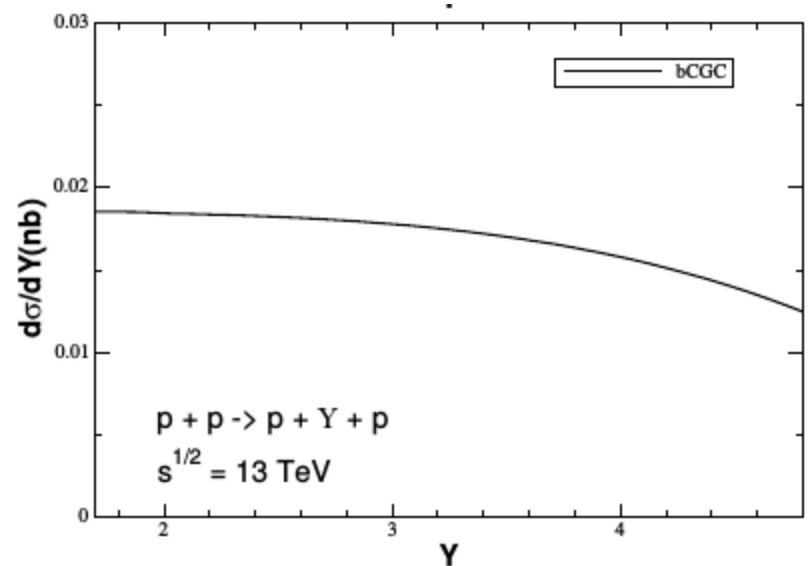
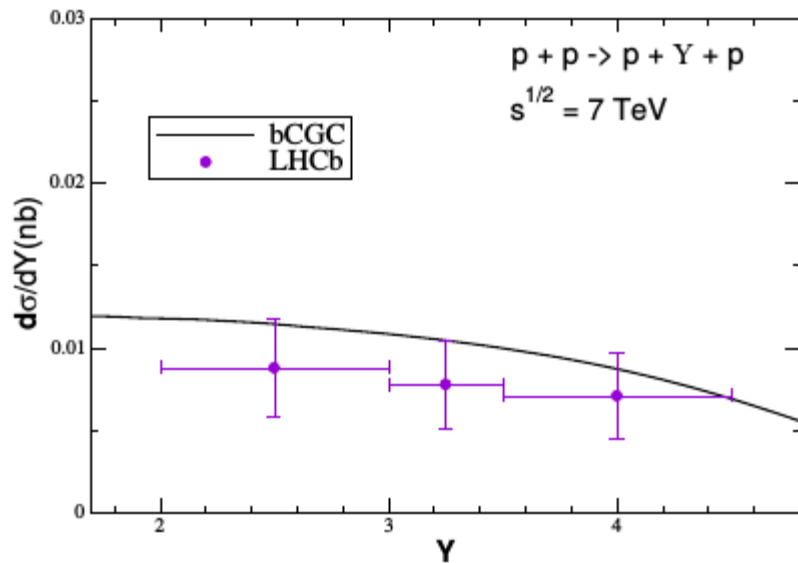
# Comparison with the LHC data

Diffractive  $\Psi(2S)$  photoproduction in pp collisions:



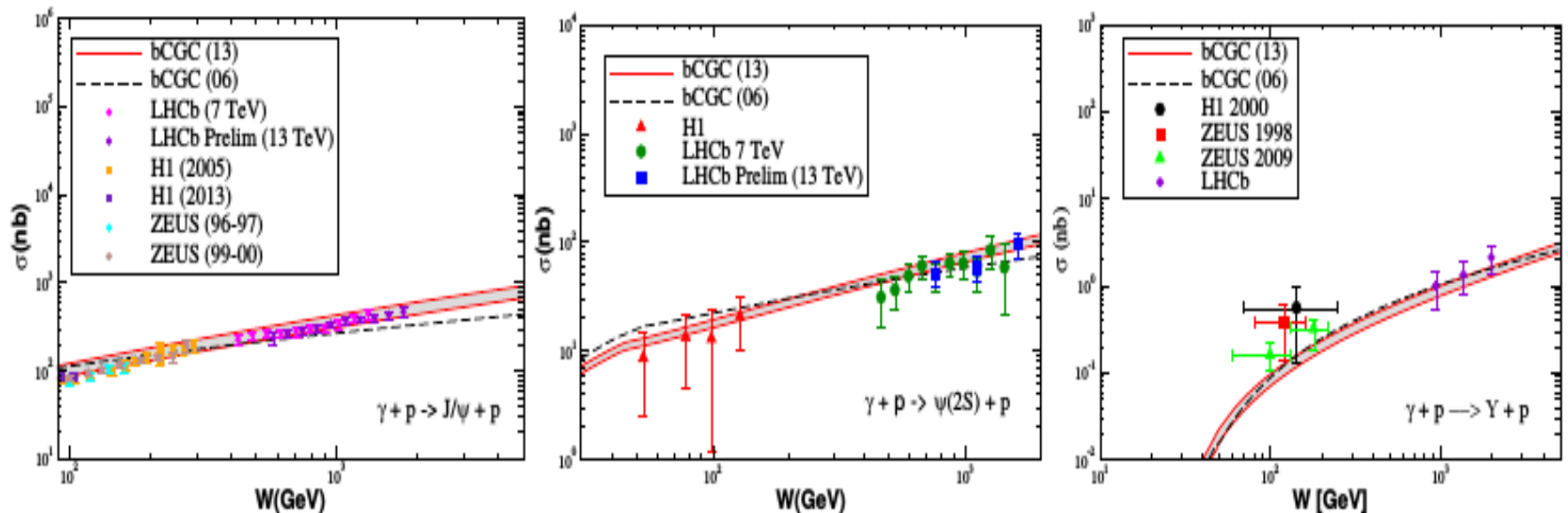
# Comparison with the LHC data

Diffractive Upsilon photoproduction in pp collisions:

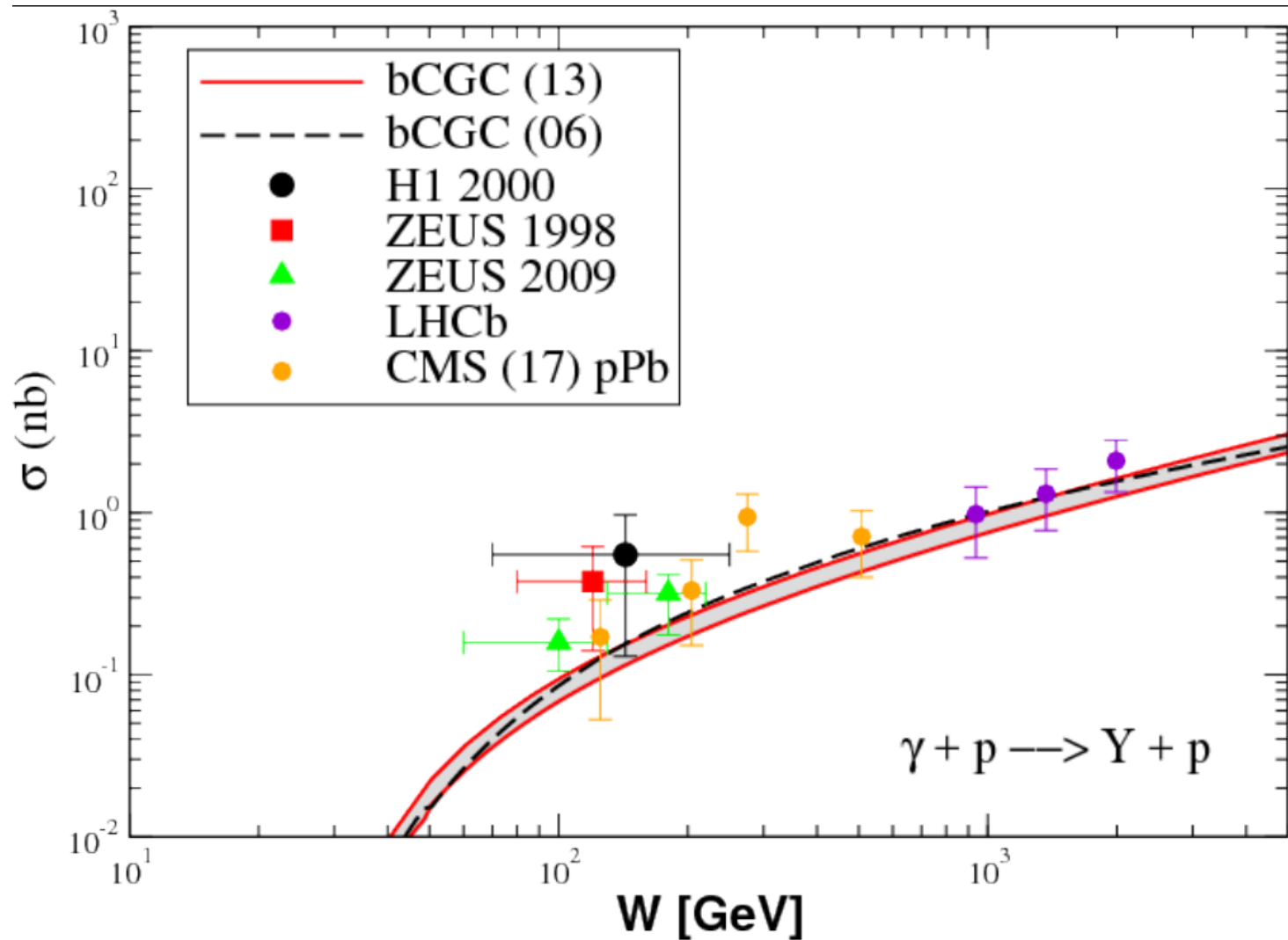


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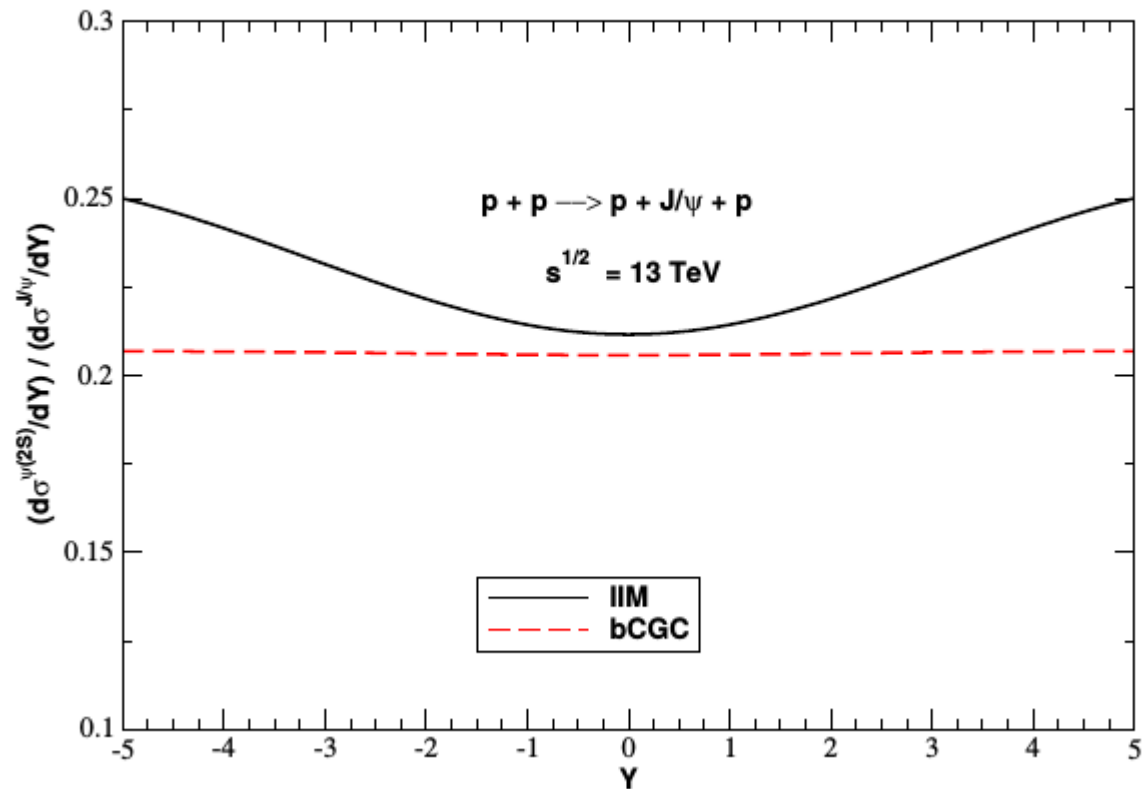
Energy dependence of the photon - proton cross section:



# Comparison with preliminary CMS pPb data



# Ratio between charmonium cross sections



# Summary



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- ✓ The diffractive vector meson photoproduction in photon - induced interactions at the LHC is an important probe of the QCD dynamics at high energies.
- ✓ The Run I data can be successfully described by the color dipole formalism taking into account the nonlinear effects in the QCD dynamics.
- ✓ The Run II data can be used to constrain the description of the dipole - hadron scattering amplitude and the vector meson wave function
- ✓ Complementary studies can be performed by analysis of the double vector meson production and the vector meson production associated to a leading neutron.

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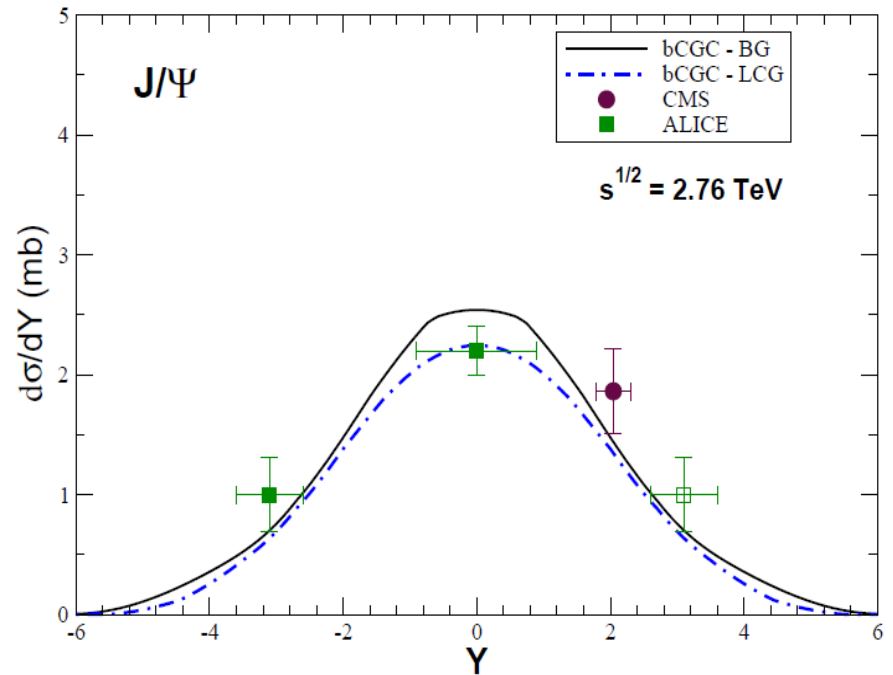
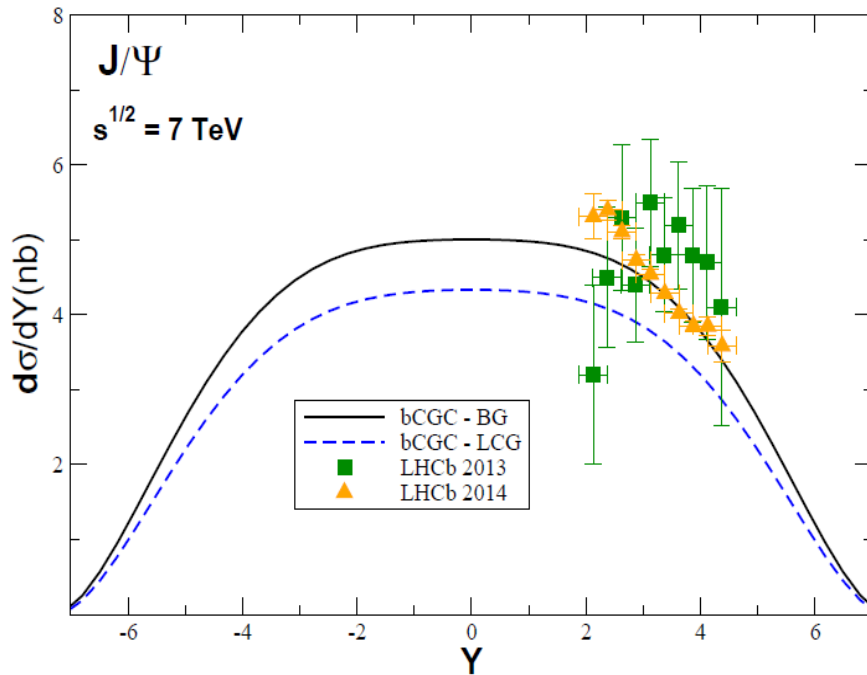
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Thank you for your attention!

# Extras

# Comparison with the Run I data

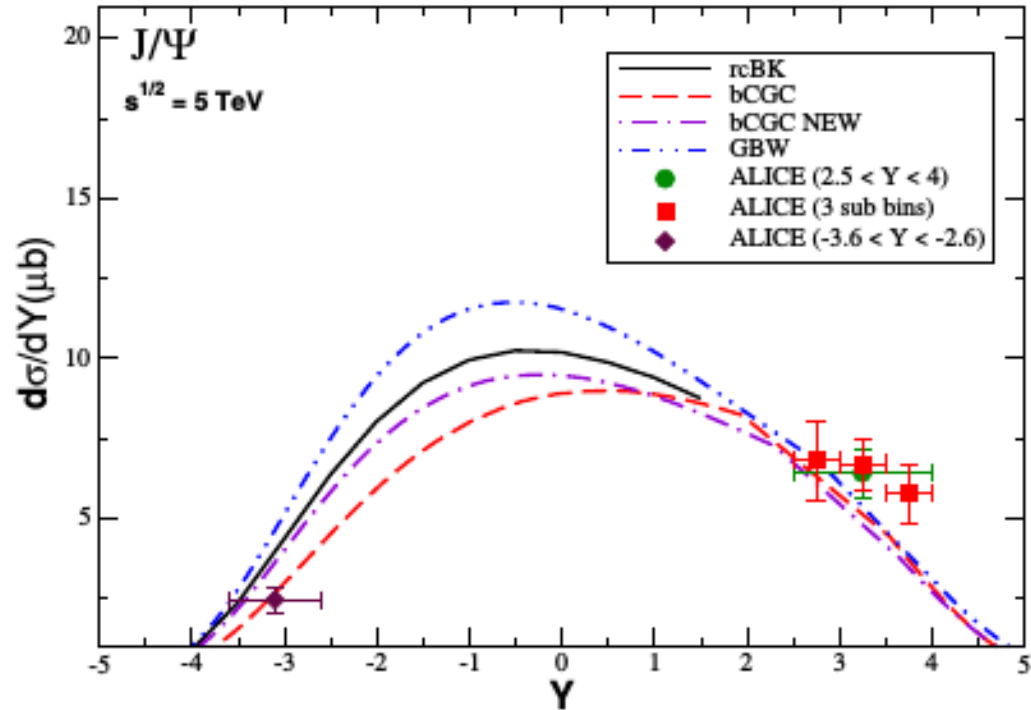
● Diffractive  $J/\Psi$  photoproduction in hadronic collisions <sup>a</sup>



(<sup>a</sup>) VPG, Moreira, Navarra, PRC90, 015203 (2014)

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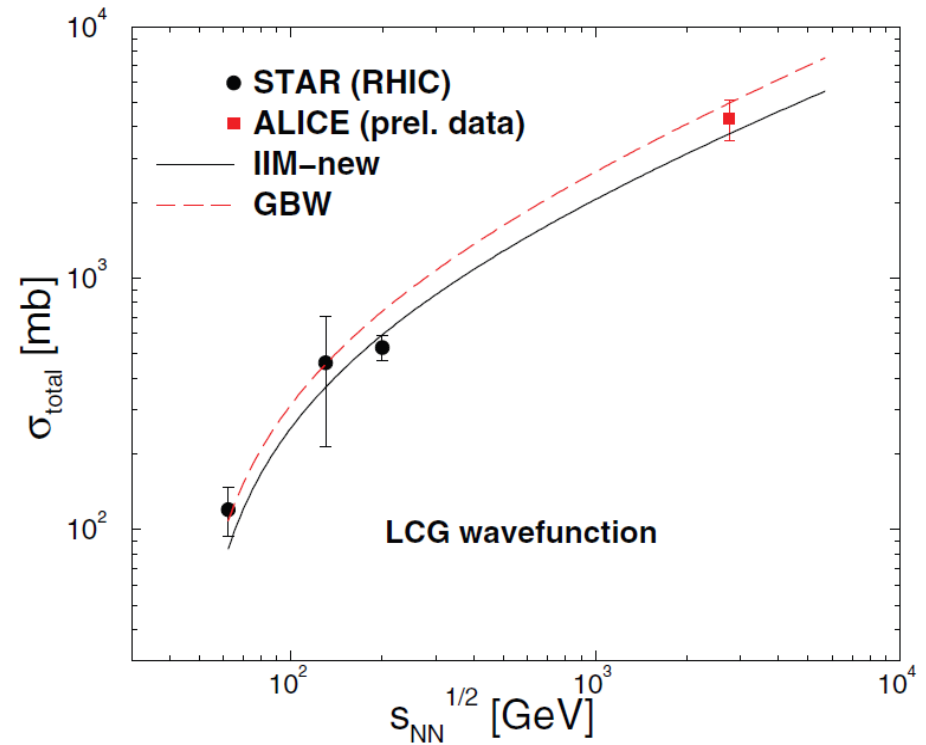
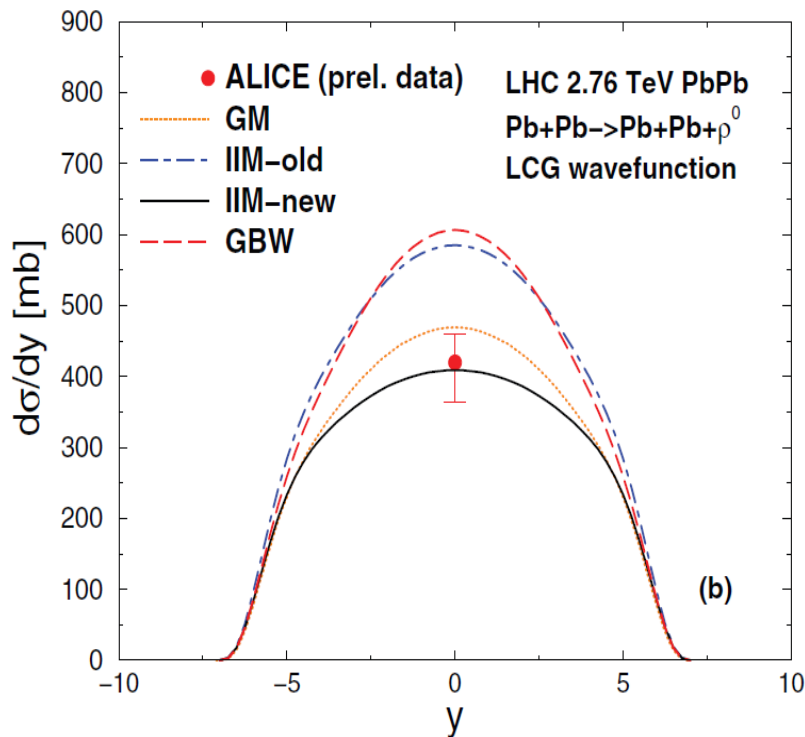
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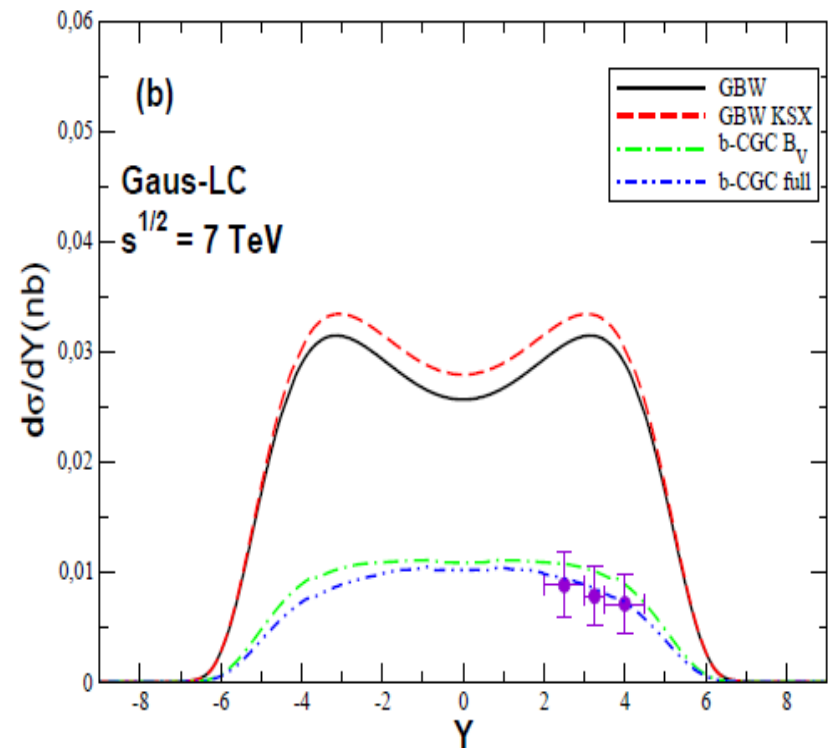
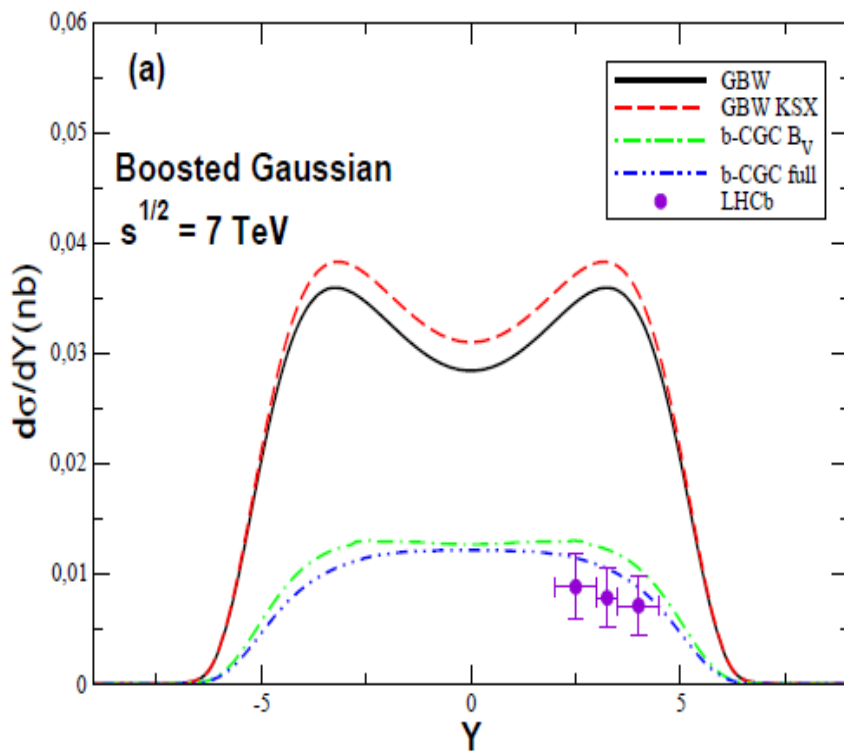
## ● Diffractive $\rho$ photoproduction in hadronic collisions <sup>c</sup>



(<sup>c</sup>) VPG, Machado, EPJC 40, 519 (2005); PRC80, 054901 (2009); PRC84, 011902 (2011); Machado, dos Santos, PRC91, 025203 (2015)

# Comparison with the Run I data

● Diffractive  $\Upsilon$  photoproduction in hadronic collisions <sup>b</sup>



<sup>b</sup>VPG, Moreira, Navarra, PLB 472, 172 (2015)