Introdução à Física de Partículas

Introduction to particle Physics

(3/4)

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CERN Portuguese Language Teachers Programme 2017

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1932

Modelo do átomo durante os anos 20:

Protões + electrões

Inconsistente com a "nova" MQ.

as 3 mm. at N.T.P.
These results, and others I have obtained in the course of the work, are very difficult to explain on the assumption that the radiation from beryllium is a quantum radiation, if energy and momentum are to be conserved in the collisions. The difficulties disappear, however, if it be assumed that the radiation consists of particles of mass 1 and charge 0, or neutrons. The capture of the α-particle by the Be^9 nucleus may be supposed to result in the

O prémio Nobel da Física foi atríbuido a James Chadwick em 1935;

"for the discovery of the neutron".

312 *NATURE* [February 27, 1932

Letters to the Editor

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Possible Existence of a Neutron

It has been shown by Bothe and others that beryllium when bombarded by α-particles of polonium emits a radiation of great penetrating power, which has an absorption coefficient in lead of about 0·3 (cm.)$^{-1}$. Recently Mme. Curie-Joliot and M. Joliot found, when measuring the ionisation produced by this beryllium radiation in a vessel with a thin window, that the ionisation increased when matter containing hydrogen was placed in front of the window. The effect appeared to be due to the ejection of protons with velocities up to a maximum of nearly 3×10^9 cm. per sec. They suggested that the transference of energy to the proton was by a process similar to the Compton effect, and estimated that the beryllium radiation had a quantum energy of 50×10^6 electron volts.

I have made some experiments using the valve counter to examine the properties of this radiation excited in beryllium. The valve counter consists of a small ionisation chamber connected to an amplifier, and the sudden production of ions by the entry of a particle, such as a proton or α-particle, is recorded by the deflexion of an oscillograph. These experiments have shown that the radiation ejects particles from hydrogen, helium, lithium, beryllium, carbon, air, and argon. The particles ejected from hydrogen behave, as regards range and ionising power, like protons with speeds up to about $3·2 \times 10^9$ cm. per sec. The particles from the other elements have a large ionising power, and appear to be in each case recoil atoms of the elements.

If we ascribe the ejection of the proton to a Compton recoil from a quantum of 52×10^6 electron volts, then the nitrogen recoil atom arising by a similar process should have an energy not greater than about 400,000 volts, should produce not more than about 10,000 ions, and have a range in air at N.T.P. of about 1·3 mm. Actually, some of the recoil atoms in nitrogen produce at least 30,000 ions. In collaboration with Dr. Feather, I have observed the recoil atoms in an expansion chamber, and their range, estimated visually, was sometimes as much as 3 mm. at N.T.P.

These results, and others I have obtained in the course of the work, are very difficult to explain on the assumption that the radiation from beryllium is a quantum radiation, if energy and momentum are to be conserved in the collisions. The difficulties disappear, however, if it be assumed that the radiation consists of particles of mass 1 and charge 0, or neutrons. The capture of the α-particle by the Be^9 nucleus may be supposed to result in the formation of a C^{12} nucleus and the emission of the neutron. From the energy relations of this process the velocity of the neutron emitted in the forward direction may well be about 3×10^9 cm. per sec. The collisions of this neutron with the atoms through which it passes give rise to the recoil atoms, and the observed energies of the recoil atoms are in fair agreement with this view. Moreover, I have observed that the protons ejected from hydrogen by the radiation emitted in the opposite direction to that of the exciting α-particle appear to have a much smaller range than those ejected by the forward radiation. This again receives a simple explanation on the neutron hypothesis.

If it be supposed that the radiation consists of quanta, then the capture of the α-particle by the Be^9 nucleus will form a C^{13} nucleus. The mass defect of C^{13} is known with sufficient accuracy to show that the energy of the quantum emitted in this process cannot be greater than about 14×10^6 volts. It is difficult to make such a quantum responsible for the effects observed.

It is to be expected that many of the effects of a neutron in passing through matter should resemble those of a quantum of high energy, and it is not easy to reach the final decision between the two hypotheses. Up to the present, all the evidence is in favour of the neutron, while the quantum hypothesis can only be upheld if the conservation of energy and momentum be relinquished at some point.

J. Chadwick.

Cavendish Laboratory,
Cambridge, Feb. 17.

The Oldoway Human Skeleton

A letter appeared in Nature of Oct. 24, 1931, signed by Messrs. Leakey, Hopwood, and Reck, in which, among other conclusions, it is stated that "there is no possible doubt that the human skeleton came from Bed No. 2 and not from Bed No. 4". This must be taken to mean that the skeleton is to be considered as a natural deposit in Bed No. 2, which is overlaid by the later beds Nos. 3 and 4, and that all consideration of human interment is ruled out.

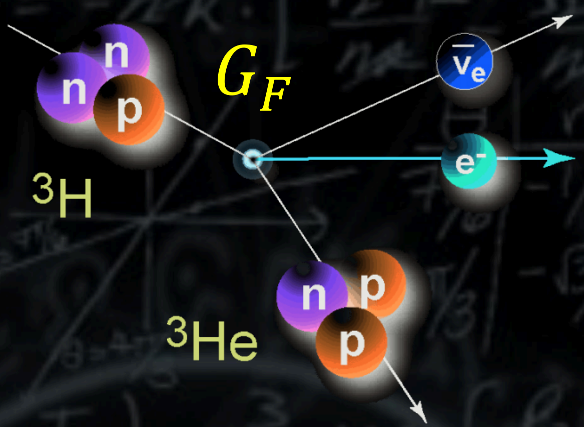
If this be true, it is a most unusual occurrence. The skeleton, which is of modern type, with filed teeth, was found completely articulated down even to the phalanges, and in a position of extraordinary contraction. Complete mammalian skeletons of any age are, as field palæontologists know, of great rarity. When they occur, their perfection can usually be explained as the result of sudden death and immediate covering by volcanic dust. Many of the more or less perfect skeletons which may be seen in museums have been rearticulated from bones found somewhat scattered as the result of death from floods, or in the neighbourhood of drying water-holes. We know of no case of a perfect articulated skeleton being found in company with such broken and scattered remains as appear to be abundant at Oldoway. Either the skeletons are all complete, as in the *Stenomylus* quarry at Sioux City, Nebraska, or are all scattered and broken in various degrees, as in ordinary bone beds. The probability, therefore, that the Oldoway skeleton represents an artificial burial is thus one that will occur to palæontologists.

The skeleton was exhumed in 1913, and published photographs show that the excavation made for its disinterment was extensive. It is, therefore, very difficult to believe that in 1931 there can be reliable evidence left at the site as to the conditions under which it was deposited. If naturally deposited in Bed No. 2, the skeleton is of the highest possible importance, because it would be of pre-Mousterian age, and would be in the company of *Pithecanthropus* and the Piltdown, Heidelberg, and Peking men, all of whose remains are fragmentary to the last degree. Of the few other human remains for which such antiquity is claimed, the Galley Hill skeleton and the Ipswich skeleton are, or apparently were, complete. The first of these was never seen *in situ* by any trained observer, and the latter has, we believe, been withdrawn by its discoverer. The other fragments, found long ago, are entirely without satisfactory evidence as to their mode of occurrence.

No. 3252, Vol. 129]

1934

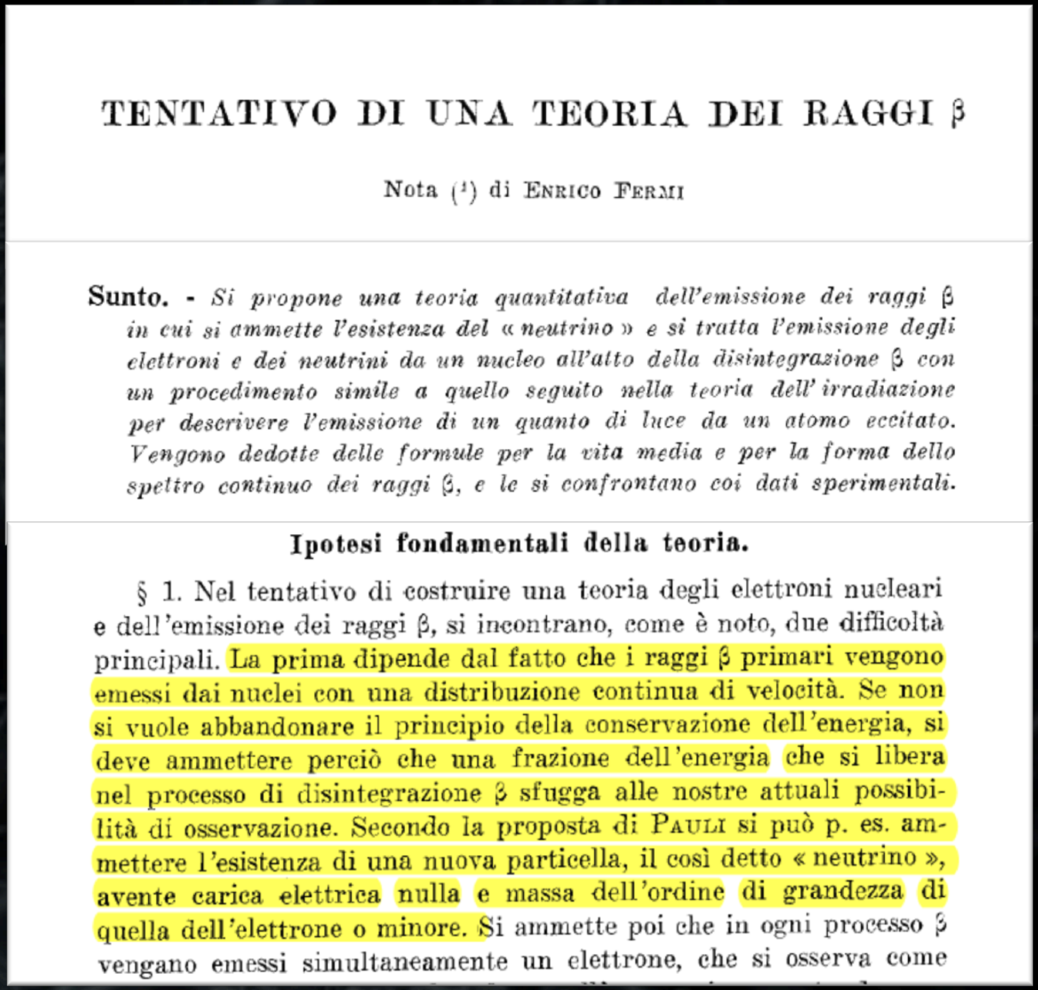
Em 1934, Fermi propõe a primeira teoria para explicar o decaimento β dos núcleos.



Esta descrição não é válida para toda a gama de energias...

Tempo de vida do muão:

$$\tau_\mu = \frac{192\pi^3}{G_F^2 m_\mu^5}$$



TENTATIVO DI UNA TEORIA DEI RAGGI β

Nota (¹) di ENRICO FERMI

Sunto. - *Si propone una teoria quantitativa dell'emissione dei raggi β in cui si ammette l'esistenza del «neutrino» e si tratta l'emissione degli elettroni e dei neutrini da un nucleo all'atto della disintegrazione β con un procedimento simile a quello seguito nella teoria dell'irradiazione per descrivere l'emissione di un quanto di luce da un atomo eccitato. Vengono dedotte delle formule per la vita media e per la forma dello spettro continuo dei raggi β, e le si confrontano coi dati sperimentali.*

Ipotesi fondamentali della teoria.

§ 1. Nel tentativo di costruire una teoria degli elettroni nucleari e dell'emissione dei raggi β, si incontrano, come è noto, due difficoltà principali. La prima dipende dal fatto che i raggi β primari vengono emessi dai nuclei con una distribuzione continua di velocità. Se non si vuole abbandonare il principio della conservazione dell'energia, si deve ammettere perciò che una frazione dell'energia che si libera nel processo di disintegrazione β sfugga alle nostre attuali possibilità di osservazione. Secondo la proposta di PAULI si può p. es. ammettere l'esistenza di una nuova particella, il così detto «neutrino», avente carica elettrica nulla e massa dell'ordine di grandezza di quella dell'elettrone o minore. Si ammette poi che in ogni processo β vengano emessi simultaneamente un elettrone, che si osserva come

Constante de Fermi:

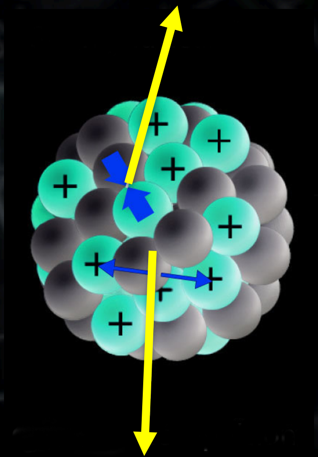
$$G_F = 1.166364\times10^5\,\mathrm{GeV}^{-2}$$

1934

O núcleo é constituido por protões e neutrões...

Mas qual a força que os mantém ligado no núcleo?

To remove this defect, it seems natural to modify the theory of Heisenberg and Fermi in the following way. The transition of a heavy particle from neutron state to proton state is not always accompanied by the emission of light particles, i. e., a neutrino and an electron, but the energy liberated by the transition is taken up sometimes by another heavy particle, which in turn will be transformed from proton state into neutron state. If the probability of occurrence of the latter process is much larger than that of the former, the interaction between



Potencial de Yukawa:

$$V(r) = -g^2 \frac{e^{-\frac{mcr}{\hbar}}}{r^2}$$

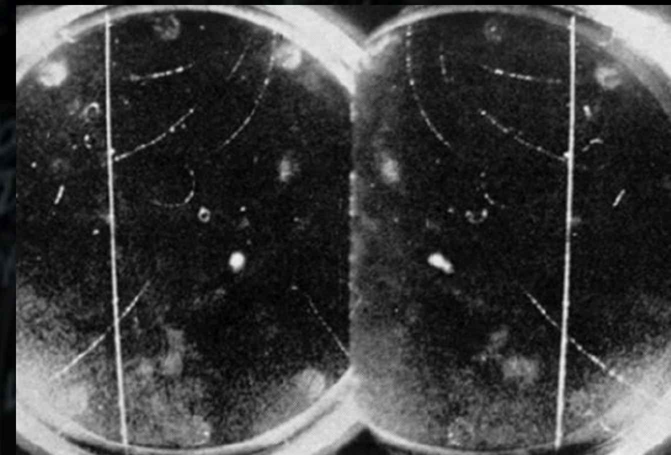
$$m \sim 100 \text{ MeV}$$

O prémio Nobel da Física foi atríbuido a Hideki Yukawa em 1949;

"for his prediction of the existence of mesons on the basis of theoretical work on nuclear forces".

1936

Em 1936 Anderson observam novos rastos numa câmara de Wilson com massa "intermédia".



Seria esta partícula de Yukawa?

Não. Esta nova partícula não tinha afinidade para se ligar ao núcleo.

Anderson chamou a esta partícula o "mesotrão" (hoje conhecida como muão μ). O pião viria a ser descoberto em 1946.

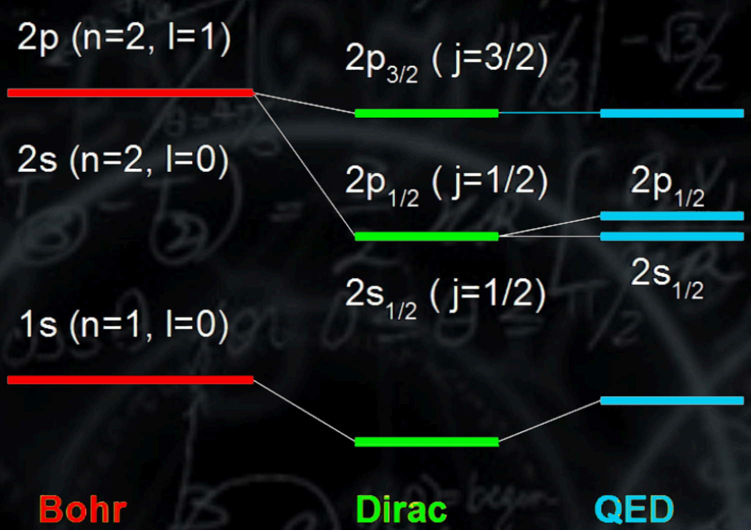
O prémio Nobel da Física foi atríbuido a Cecil Powell em 1950;

"for his development of the photographic method of studying nuclear processes and his discoveries regarding mesons made with this method".

1947

$$E_{n,j} = mc^2 \left[1 + \frac{Z^2}{(n-\varepsilon_j)^2} \right]^{\frac{1}{2}} \quad , \quad \varepsilon_j = j + \frac{1}{2} - \sqrt{\left(j+\frac{1}{2}\right)^2 - Z^2 e^4}$$

Segundo as previsões de Dirac, os estados $^2S_{1/2}$ e $^2P_{1/2}$ são degenerados



Em 1947 Willis Lamb e Robert Retherford mediram uma diferença de energia entre estes dois níveis.

O prémio Nobel da Física foi atríbuido a Willis Lamb em 1955;

"for his discoveries concerning the fine structure of the Hydrogen spectrum".

1947

1936

Bethe calculou pela primeira vez o desvio de Lamb obtendo o valor de:

1040 MHz

Conferência em Shelter Island (1947)

figure and $K = mc^2$, the logarithm has the value 7.63, and we find

$$W_{ns}' = 136 \ln[K/(E_n - E_m)]$$

$$= 1040 \text{ megacycles.} \quad (12)$$

[7] I am indebted to Dr. Stehn and Miss Steward for the numerical calculations.

Bethe foi o primeiro a determinar as correcções relevantes para o desvio de Lamb.

PHYSICAL REVIEW VOLUME 72, NUMBER 4 AUGUST 15, 1947

The Electromagnetic Shift of Energy Levels

H. A. Bethe
Cornell University, Ithaca, New York
(Received June 27, 1947)

By very beautiful experiments, Lamb and Retherford[1] have shown that the fine structure of the second quantum state of hydrogen does not agree with the prediction of the Dirac theory. The $2s$ level, which according to Dirac's theory should coincide with the $2p_{\frac{1}{2}}$ level, is actually higher than the latter by an amount of about 0.033 cm^{-1} or 1000 megacycles. This discrepancy had long been suspected from spectroscopic measurements.[2,3] However, so far no satisfactory theoretical explanation has been given. Kemble and Present, and Pasternack[4] have shown that the shift of the $2s$ level cannot be explained by a nuclear interaction of reasonable magnitude, and Uehling[5] has investigated the effect of the "polarization of the vacuum" in the Dirac hole theory, and has found that this effect also is much too small and has, in addition, the wrong sign.

Schwinger and Weisskopf, and Oppenheimer have suggested that a possible explanation might be the shift of energy levels by the interaction of the electron with the radiation field. This shift comes out infinite in all existing theories, and has therefore always been ignored. However, it is possible to identify the most strongly (linearly) divergent term in the level shift with an electromagnetic *mass* effect which must exist for a bound as well as for a free electron. This effect should

[1] Phys. Rev. **72**, 241 (1947).
[2] W. V. Houston, Phys. Rev. **51**, 446 (1937).
[3] R. C. Williams, Phys. Rev. **54**, 558 (1938).
[4] E. C. Kemble and R. D. Present, Phys. Rev. **44**, 1031 (1932); S. Pasternack, Phys. Rev. **54**, 1113 (1938).
[5] E. A. Uehling, Phys. Rev. **48**, 55 (1935).

Mais problemas...

Apesar de resolver algumas questões em aberto, a MQR ainda não era uma teoria satisfatória.

$$E_c \sim mc^2 \rightarrow p \sim mc \rightarrow \lambda \sim \frac{h}{mc} = \lambda_c$$

λ_c - Comprimento de onda de Compton

PRINCIPIO DA INCERTEZA: $\Delta p \geq \frac{\hbar}{\Delta x} = mc$

Conclusão: Se tentarmos localizar uma partícula de massa m numa região do espaço de dimensões menores que λ_c, então as flutuações na energia são suficientes para criar um par partícula-antíparticula.

UMA TEORIA COMPLETAMENTE RELATIVISTA NÃO SE PODE BASEAR NA IDEIA DE QUE UM SISTEMA PODE SER DESCRITO PELA FUNÇÃO DE ONDA A 1 PARTÍCULA.

TEORIA QUÂNTICA DE CAMPOS (TQC)

MECÂNICA QUÂNTICA: Posição e momento são tratados como operadores:

$$\hat{x}\,\psi = x\,\psi \quad , \quad \hat{p}\,\psi = -i\hbar\frac{\partial\psi}{\partial x}$$

E O TEMPO t ?

Numa teoria relativista espaço deveriam estar em pé de igualdade. Mas em MQ o tempo é um parâmetro.

$$\hat{T}\,\psi = t\,\psi$$

EM TQC, A POSIÇÃO E O MOMENTO SÃO DESPROMOVIDOS A PARÂMETROS QUE SÃO ARGUMENTOS DE UM CAMPO

$\hat{\phi}\,(x,t)$ SEGUNDA QUANTIFICAÇÃO

TEORIA QUÂNTICA DE CAMPOS (TQC)

$$\hat{\phi}\,(x,t) = f(\hat{a},\hat{a}^{+})$$

$\hat{a}, \hat{a}^{+}$ - operadores de criação e aniquilação

Estados em TQC: $|\vec{p}\rangle = \hat{a}^{+}(\vec{p})\;|\;0\;\rangle$

Estados com várias partículas:

Estados com vária partículas: $|\vec{p}_1,\vec{p}_2\rangle = \hat{a}^{+}(\vec{p}_1)\;\hat{a}^{+}(\vec{p}_2)|\;0\;\rangle$

Aniquilação: $\hat{a}(\vec{p}_1)|\vec{p}_1,\vec{p}_2\rangle = |\vec{p}_2\rangle$

O operador campo actua em estados de tal modo que se podem criar ou aniquilar partículas!

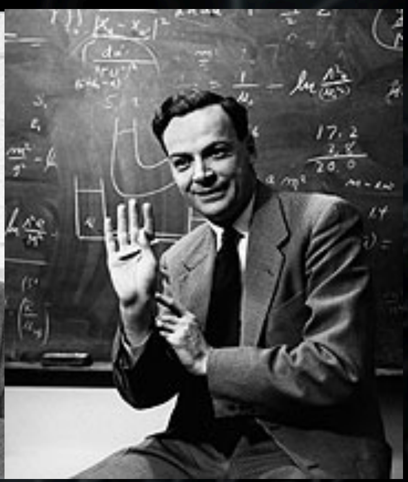
TQC DOS FOTÕES, ELECTRÕES, POSITRÕES E SUAS INTERAÇÕES

ElectroDinâmica Quântica - EDQ (QED)

Schwinger

Tomonaga

Dyson



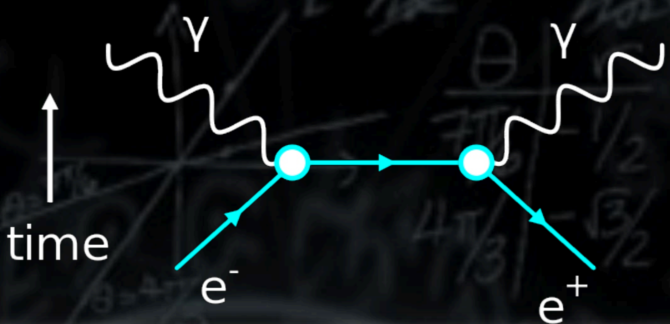
Feynman

The Nobel Prize in Physics 1965 was awarded jointly to Sin-Itiro Tomonaga, Julian Schwinger and Richard P. Feynman "for their fundamental work in quantum electrodynamics, with deep-ploughing consequences for the physics of elementary particles".

1948

A furgoneta Feynman

Em 1948 Feynman desenvolveu um método "gráfico" que permite calcular processos que envolvem interacções entre particulas elementares.



Para o diagram acima: $\mathcal{M} = \varepsilon^*_\mu(p_1,\lambda_1)\varepsilon^*_\nu(p_2,\lambda_2)\bar{v}(e^+)(ie\gamma^\nu)\frac{i}{\not{q}-m}(ie\gamma^\mu)u(e^-)$

A TÉCNICA DOS DIAGRAMAS DE FEYNMAN É USADA PARA CALCULAR AS "PROBABILIDADES" DE OCORRÊNCIA DE PROCESSOS ENVOLVENDO PARTÍCULAS ELEMENTARES.

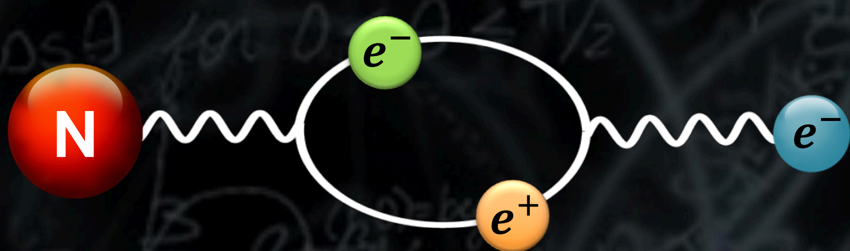
vácuo
(latim *vacuus, -a, -um*)
adj.
1. Que não contém nada; que não se acha ocupado por coisa alguma. = OCO, VAZIO
2. [Jurídico, Jurisprudência] Que se possui mas que não se desfruta ou não se goza.
s. m.
3. Espaço circunscrito que não contém ar ou que se supõe vazio.
4. O espaço entre os corpos celestes que se supõe vazio.

EM FÍSICA QUÂNTICA O VÁCUO ESTÁ LONGE DE SER ESPAÇO VAZIO...

PIH Energia-tempo

$$\Delta E \, \Delta t \geq \frac{\hbar}{2}$$

2ª Quantização: O campo E.M. pode ser interpretado como um conjunto infinito de osciladores harmónicos cujo estado fundamental tem energia não nula.



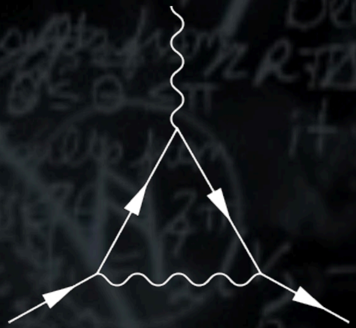
Polarização do vácuo

À semelhança da polarização de um dieléctrico.

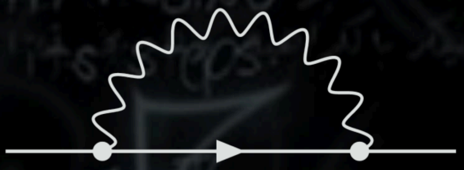
O DESVIO DE LAMB

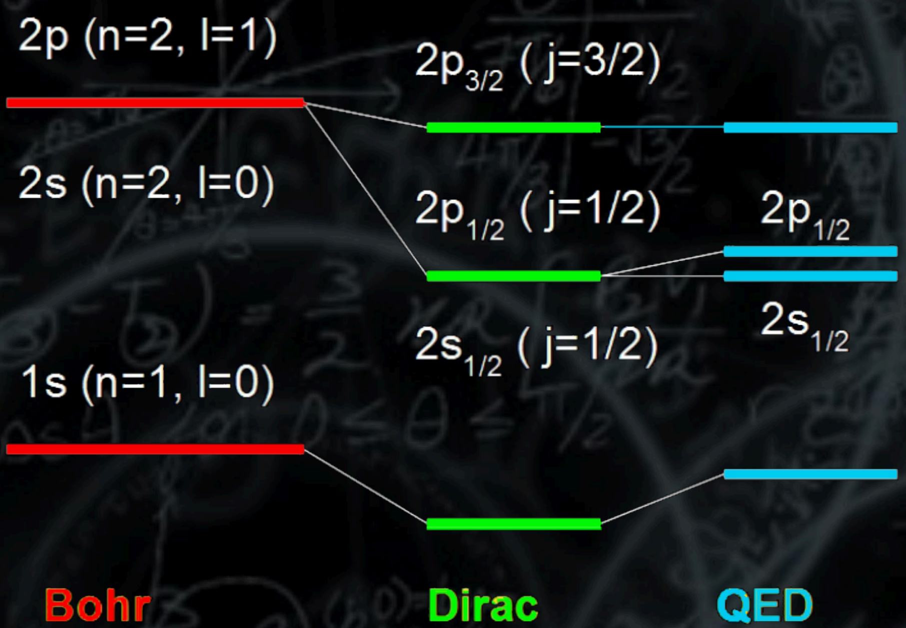
Polarização do vácuo

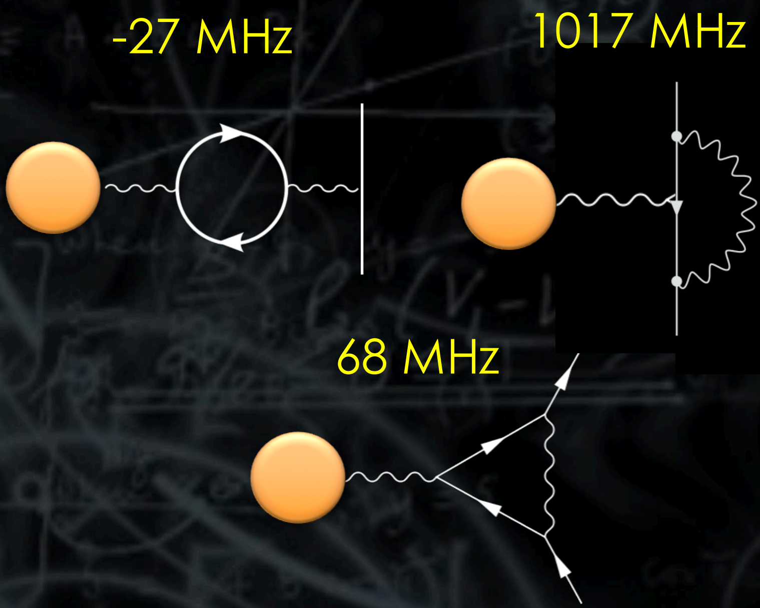
Correcção ao vértice



Correcção à energia própria







Teoria: 1058 MHz , Exp: 1057.9

O EFEITO CASIMIR

Segunda Quantização: O campo electromagnético consiste num conjunto de osciladores harmónicos.

CASIMIR (1948): $$\frac{F(a)}{A} = \frac{\pi^2}{240}\frac{\hbar c}{a^4}$$

Estre as placas condutoras só alguns modos dão permitidos.
EFEITO: Força entre as placas condutoras.

VOLUME 78, NUMBER 1 PHYSICAL REVIEW LETTERS 6 JANUARY 1997

Demonstration of the Casimir Force in the 0.6 to 6 μm Range

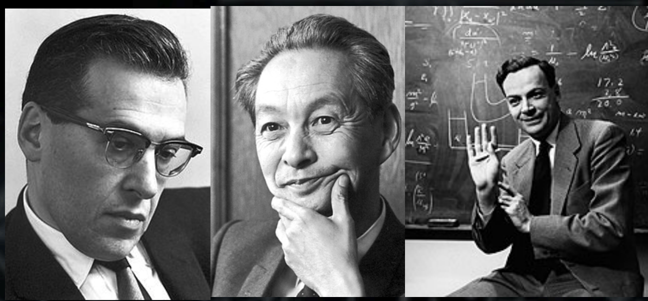
S. K. Lamoreaux*

Physics Department, University of Washington, Box 35160, Seattle, Washington 98195-1560
(Received 28 August 1996)

The vacuum stress between closely spaced conducting surfaces, due to the modification of the zero-point fluctuations of the electromagnetic field, has been conclusively demonstrated. The measurement employed an electromechanical system based on a torsion pendulum. Agreement with theory at the level of 5% is obtained. [S0031-9007(96)02025-X]

PACS numbers: 12.20.Fv, 07.07.Mp

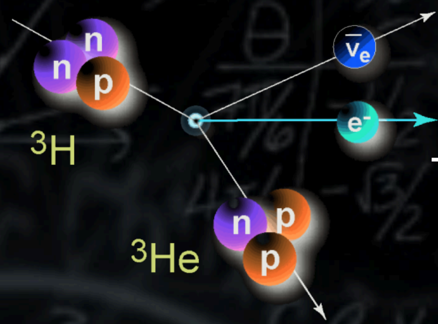
SITUAÇÃO NO FINAL DOS ANOS 40



ELECTRODINÂMICA QUÂNTICA

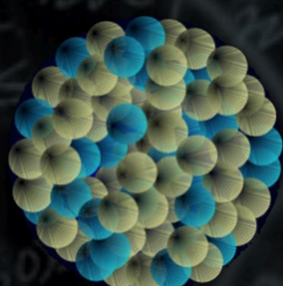
Teoria quântica dos electrões, positrões, fotões e da interacção electromagnética.

(e^-, e^+, γ)



FORÇA FRACA

Teoria do decaimento radioactivo descrita pela interacção de Fermi.



FORÇA FORTE

Força responsável pela coesão do núcleo descrita pelo potencial de Yukawa.

(n, p, π)


SIMETRIA

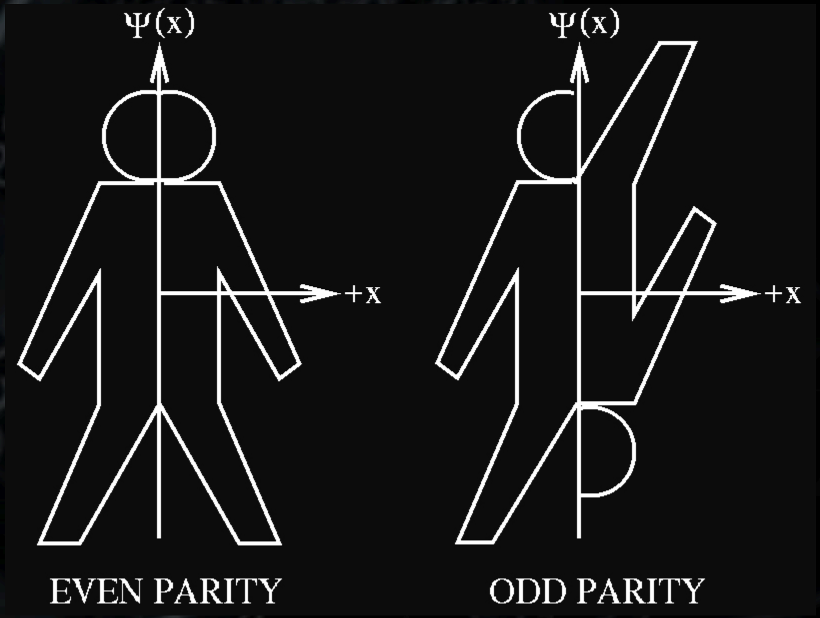
PARIDADE

Transformação de paridade P:

$$\mathrm{P}: \vec{r} \rightarrow -\vec{r}$$

Uma dupla transformação de paridade corresponde a não fazer nada…

$$\mathrm{P}^2\psi(\vec{r}) = \psi(\vec{r}) \implies \mathrm{P}\,\psi(\vec{r}) = \pm\,\psi(\vec{r})$$



Vector: $\mathrm{P}(\vec{p}) = \mathrm{P}\left(\frac{d\vec{r}}{dt}\right) = -\vec{p}$

Pseudo-vector: $\vec{L} = \vec{r}\times\vec{p}\ ,\ \vec{S}$

A interacção electromagnética, a força forte e a gravidade são invariantes debaixo de paridade.

1954

Paridade intrínseca: As partículas elementares (e não só) têm uma paridade intrínseca.

$$P_T = P\,P'$$

PHYSICAL REVIEW VOLUME 95, NUMBER 6 SEPTEMBER 15, 1954

Absorption of Negative Pions in Deuterium: Parity of the Pion*

W. Chinowsky and J. Steinberger
Columbia University, New York, New York

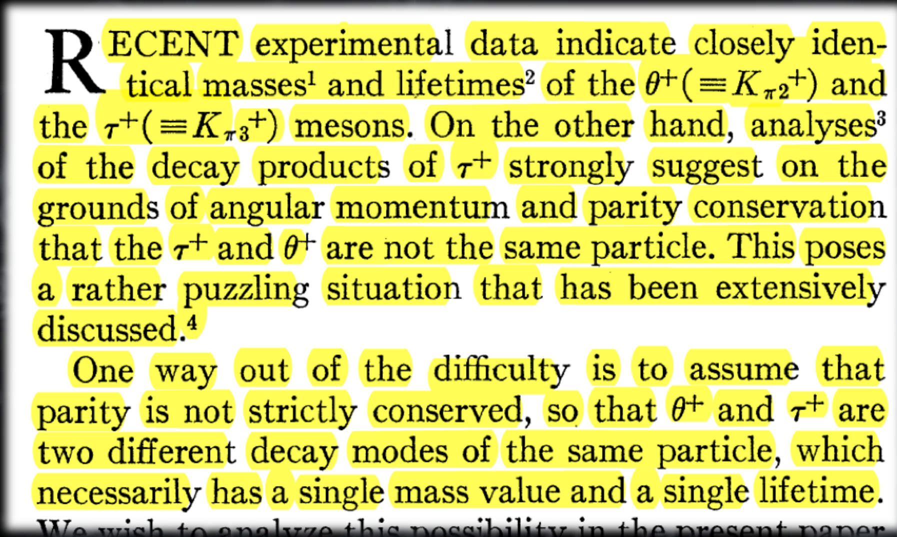
(Received June 8, 1954)

The reaction $\pi^- + d \rightarrow 2n$ has been observed by detecting the two neutrons in coincidence with slow negative mesons incident on a liquid deuterium target. The observed angular correlation of the two neutrons confirms the identification of the process. The process is therefore not forbidden, and this fact may be used to establish the odd relative parity of the pion and the nucleon.

Dalitz (1954): Puzzle $\theta - \tau$ Duas partículas (mesões θ e τ) com a mesma massa decaíam para estados de paridade diferente.

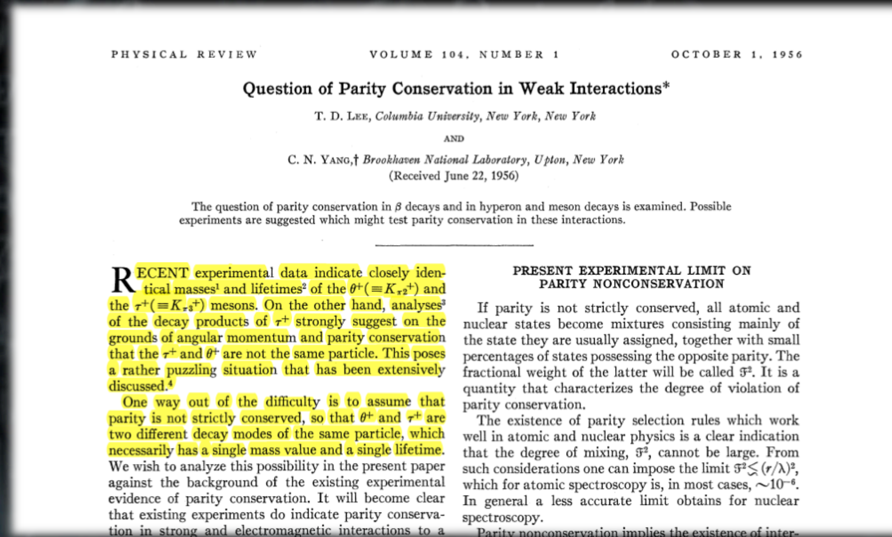
$$\tau^+ \rightarrow \pi^+\pi^-\pi^+ : P_T = (-1)^3 = -1 \quad \theta^+ \rightarrow \pi^+\pi^0 : P_T = (-1)^2 = 1$$

1956



RECENT experimental data indicate closely identical masses[1] and lifetimes[2] of the $\theta^+(\equiv K_{\pi 2}{}^+)$ and the $\tau^+(\equiv K_{\pi 3}{}^+)$ mesons. On the other hand, analyses[3] of the decay products of τ^+ strongly suggest on the grounds of angular momentum and parity conservation that the τ^+ and θ^+ are not the same particle. This poses a rather puzzling situation that has been extensively discussed.[4]

One way out of the difficulty is to assume that parity is not strictly conserved, so that θ^+ and τ^+ are two different decay modes of the same particle, which necessarily has a single mass value and a single lifetime.



PHYSICAL REVIEW VOLUME 104, NUMBER 1 OCTOBER 1, 1956

Question of Parity Conservation in Weak Interactions*

T. D. LEE, *Columbia University, New York, New York*

AND

C. N. YANG,† *Brookhaven National Laboratory, Upton, New York*

(Received June 22, 1956)

The question of parity conservation in β decays and in hyperon and meson decays is examined. Possible experiments are suggested which might test parity conservation in these interactions.

RECENT experimental data indicate closely identical masses[1] and lifetimes[2] of the $\theta^+(\equiv K_{\pi 2}{}^+)$ and the $\tau^+(\equiv K_{\pi 3}{}^+)$ mesons. On the other hand, analyses[3] of the decay products of τ^+ strongly suggest on the grounds of angular momentum and parity conservation that the τ^+ and θ^+ are not the same particle. This poses a rather puzzling situation that has been extensively discussed.[4]

One way out of the difficulty is to assume that parity is not strictly conserved, so that θ^+ and τ^+ are two different decay modes of the same particle, which necessarily has a single mass value and a single lifetime. We wish to analyze this possibility in the present paper against the background of the existing experimental evidence of parity conservation. It will become clear that existing experiments do indicate parity conservation in strong and electromagnetic interactions to a

PRESENT EXPERIMENTAL LIMIT ON PARITY NONCONSERVATION

If parity is not strictly conserved, all atomic and nuclear states become mixtures consisting mainly of the state they are usually assigned, together with small percentages of states possessing the opposite parity. The fractional weight of the latter will be called $\mathfrak{F}^2$. It is a quantity that characterizes the degree of violation of parity conservation.

The existence of parity selection rules which work well in atomic and nuclear physics is a clear indication that the degree of mixing, $\mathfrak{F}^2$, cannot be large. From such considerations one can impose the limit $\mathfrak{F}^2 \lesssim (r/\lambda)^2$, which for atomic spectroscopy is, in most cases, $\sim 10^{-6}$. In general a less accurate limit obtains for nuclear spectroscopy.

Parity nonconservation implies the existence of inter-

O prémio Nobel da Física foi atríbuido a C. N. Yang e T.D. Lee em 1957;

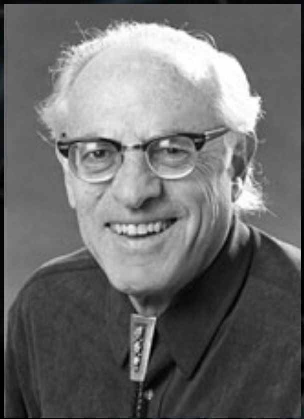
"for their penetrating investigation of the so-called parity laws which has led to important discoveries regarding the elementary particles"".

1956

Em 1956 Cowan e Reines detectaram o neutrino do electrão usando como fonte os neutrinos provenientes de um reactor nuclear.

$$\bar{\nu}_e + p \rightarrow n + e^+$$

present work was done (*3*). This work confirms the results obtained at Hanford and so verifies the neutrino hypothesis suggested by Pauli (*4*) and incorporated in a quantitative theory of beta decay by Fermi (*5*).



O prémio Nobel da Física foi atríbuido a F. Reines em 1995;

"for the detection of the neutrino"

20 July 1956, Volume 124, Number 3212 SCIENCE

Detection of the Free Neutrino: a Confirmation

C. L. Cowan, Jr., F. Reines, F. B. Harrison, H. W. Kruse, A. D. McGuire

A tentative identification of the free neutrino was made in an experiment performed at Hanford (*1*) in 1953. In that work the reaction

$$\nu_- + p^+ \rightarrow \beta^+ + n^0 \quad (1)$$

was employed wherein the intense neutrino flux from fission-fragment decay in a large reactor was incident on a detector containing many target protons in a hydrogenous liquid scintillator. The reaction products were detected as a delayed pulse pair; the first pulse being due to the slowing down and annihilation of the positron and the second to capture of the moderated neutron in cadmium dissolved in the scintillator. To identify the observed signal as neutrino-induced, the energies of the two pulses, their time-delay spectrum, the dependence of the signal rate on reactor power, and its magnitude as compared with the predicted rate were used. The calculated effectiveness of the shielding employed, together with neutron measurements made with emulsions external to the shield, seemed to rule out reactor neutrons and gamma radiation as the cause of the signal. Although a high background was experienced due to both the reactor and to cosmic radiation, it was felt that an identification of the free neutrino had probably been made.

Design of the Experiment

To carry this work to a more definitive conclusion, a second experiment was designed (*2*), and the equipment was taken to the Savannah River Plant of the U.S. Atomic Energy Commission, where the present work was done (*3*). This work confirms the results obtained at Hanford and so verifies the neutrino hypothesis suggested by Pauli (*4*) and incorporated in a quantitative theory of beta decay by Fermi (*5*).

In this experiment, a detailed check of each term of Eq. 1 was made using a detector consisting of a multiple-layer (club-sandwich) arrangement of scintillation counters and target tanks. This arrangement permits the observation of prompt spatial coincidences characteristic of positron annihilation radiation and of the multiple gamma ray burst due to neutron capture in cadmium as well as the delayed coincidences described in the first paragraph.

The three "bread" layers of the sandwich are scintillation detectors consisting of rectangular steel tanks containing a purified triethylbenzene solution of terphenyl and POPOP (*6*) in a chamber 2 feet thick, 6 feet 3 inches long, and 4 feet 6 inches wide. The tops and bottoms of these chambers are thin to low-energy gamma radiation. The tank interiors are painted white, and the solutions in the chambers are viewed by 110 5-inch Dumont photomultiplier tubes connected in parallel in each tank. The energy resolution of the detectors for gamma rays of 0.5 Mev is about 15 percent half-width at half-height.

The two "meat" layers of the sandwich serve as targets and consist of polyethylene boxes 3 inches thick and 6 feet 3 inches by 4 feet 6 inches on edge containing a water solution of cadmium chloride. This provides two essentially independent "triad" detectors, the central scintillation detector being common to both triads. The detector was completely enclosed by a paraffin and lead shield and was located in an underground room of the reactor building which provides excellent shielding from both the reactor neutrons and gamma rays and from cosmic rays.

The signals from a bank of preamplifiers connected to the scintillation tanks were transmitted via coaxial lines to an electronic analyzing system in a trailer van parked outside the reactor building. Two independent sets of equipment were used to analyze and record the operation of the two triad detectors. Linear amplifiers fed the signals to pulse-height selection gates and coincidence circuits. When the required pulse amplitudes and coincidences (prompt and delayed) were satisfied, the sweeps of two triple-beam oscilloscopes were triggered, and the pulses from the complete event were recorded photographically. The three beams of both oscilloscopes recorded signals from their respective scintillation tanks independently. The oscilloscopes were thus operated in parallel but with different gains in order to cover the requisite pulse-amplitude range. All amplifier pulses were stored in long low-distortion delay lines awaiting electronic decision prior to this acceptance.

Manual analysis of the photographic record of an event then yielded the energy deposited in each tank of a triad by both the first and second pulses and the time-delay between the pulses. Using this system, various conditions could be placed on the pulses of the pair comprising an acceptable event. For example, acceptance of events with short time delays (over ranges up to 17 microseconds, depending on the cadmium concentration used) resulted in optimum signal-to-background ratios, while analysis of those events with longer time delays yielded relevant accidental background rates. Spectral analyses of pulses comprising events with short time delays were also made and compared with those with long delays.

This method of analysis was also employed to require various types of energy deposition in the two tanks of a triad. For instance, the second pulse of an event

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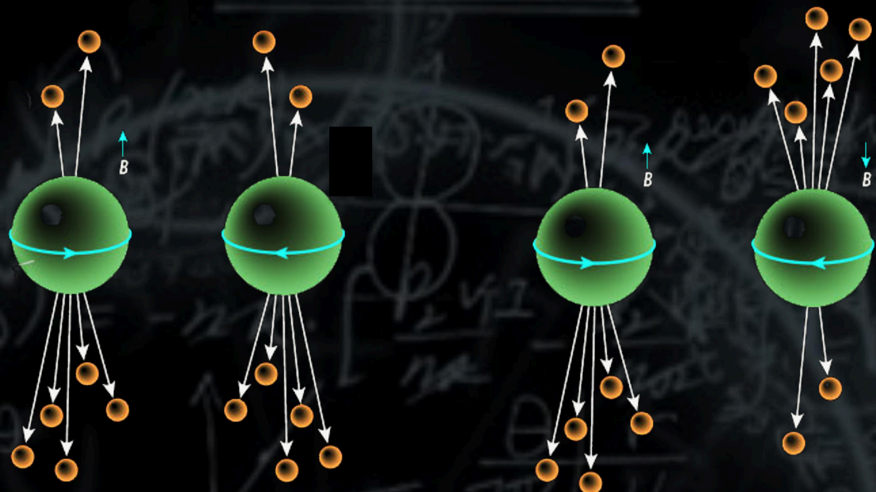
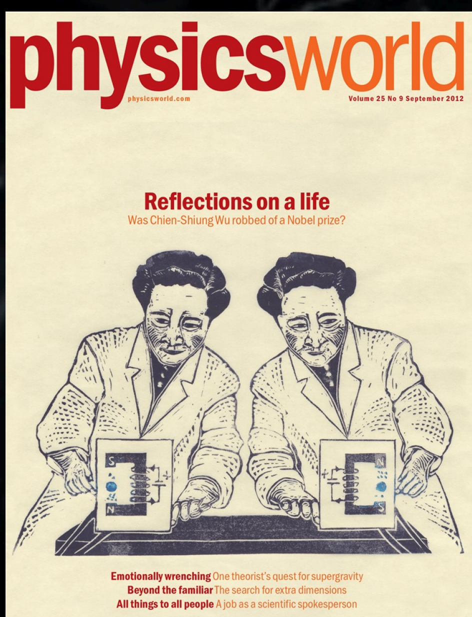
A experiência de Wu: um Nobel não atribuido?

1957



Em 1957 Wu obtém a prova experimental de que a paridade não é conservada pela interacção fraca.



Experimental Test of Parity Conservation in Beta Decay*

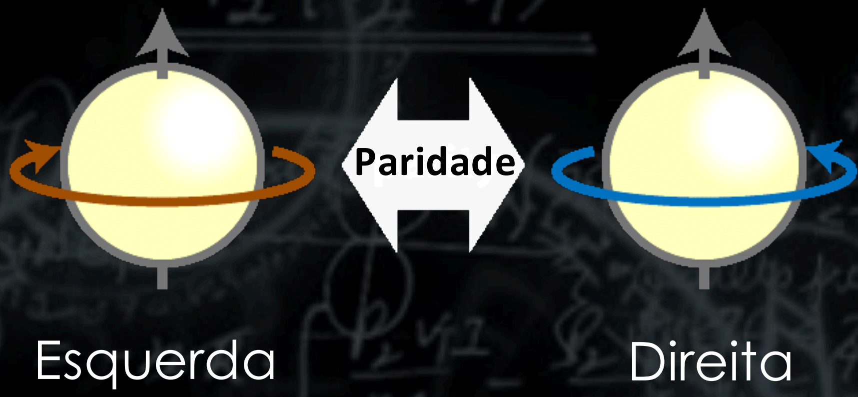
C. S. Wu, *Columbia University, New York, New York*

AND

E. Ambler, R. W. Hayward, D. D. Hoppes, and R. P. Hudson, *National Bureau of Standards, Washington, D. C.*

(Received January 15, 1957)

In a recent paper[1] on the question of parity in weak interactions, Lee and Yang critically surveyed the experimental information concerning this question and reached the conclusion that there is no existing evidence either to support or to refute parity conservation in weak interactions. They proposed a number of experiments on beta decays and hyperon and meson decays which would provide the necessary evidence for parity conservation or nonconservation. In beta decay, one could measure the angular distribution of the electrons coming from beta decays of polarized nuclei. If an asymmetry in the distribution between θ and $180°-\theta$ (where θ is the angle between the orientation of the parent nuclei and the momentum of the electrons) is observed, it provides unequivocal proof that parity is not conserved in beta decay. This asymmetry effect has been observed in the case of oriented Co^{60}.

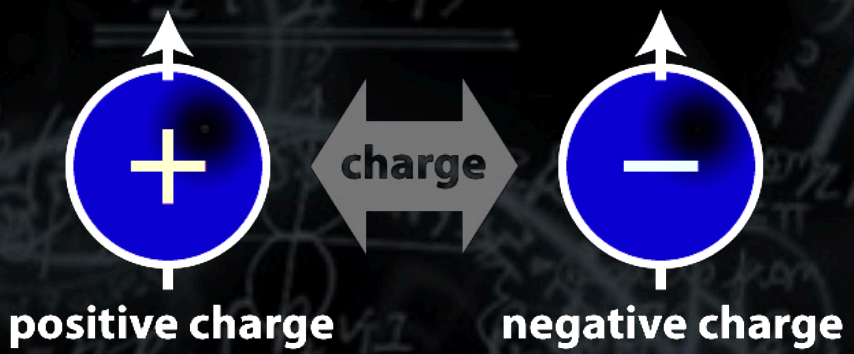


HELICIDADE: $h = \frac{\vec{s}\cdot\vec{p}}{|\vec{p}|}$

h=1 : Direita

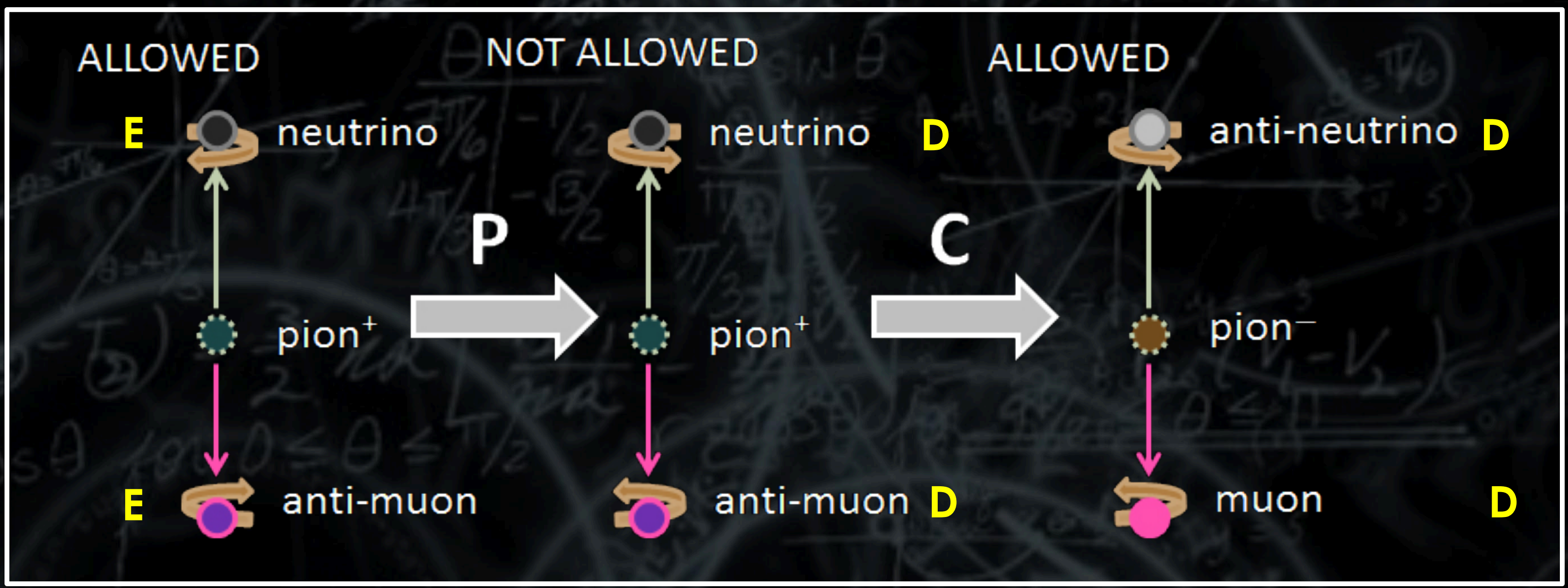
h=-1 : Esquerda

SIMETRIAS C (conjugação de carga) E CP

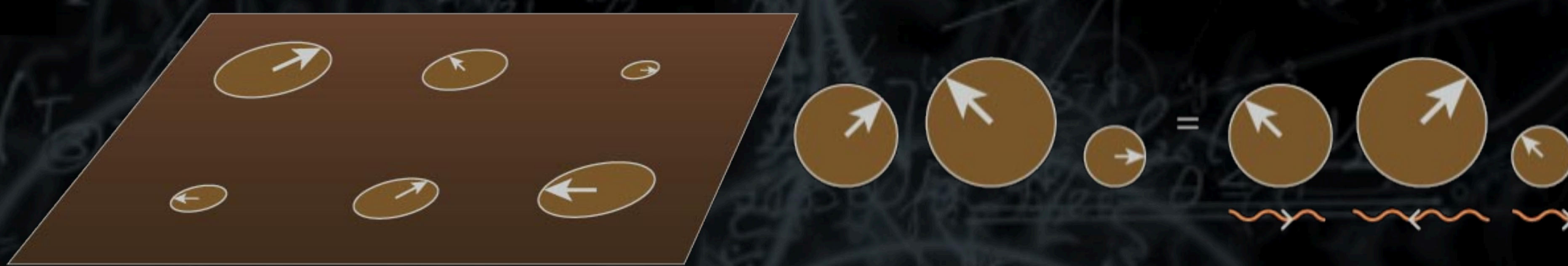


A conjugação de carga transforma uma partícula na sua antipartícula

$$C|\,p\,\rangle = |\,\bar{p}\,\rangle$$



TALVEZ CP SEJA CONSERVADA...



Os bosões de gauge (e respectivas forças) surgem como consequência de impôr invariância debaixo de simetrias locais

O MODELO PADRÃO: A RECEITA

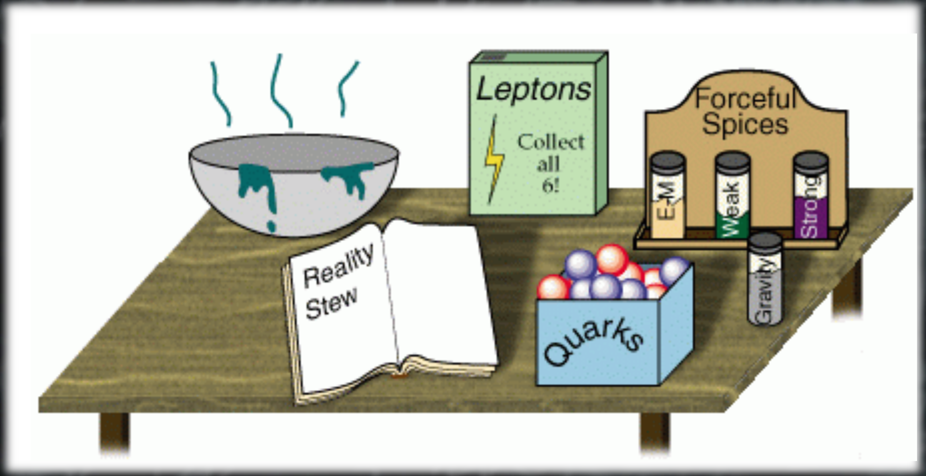
O GRUPO DE SIMETRIA DO MODELO PADRÃO É:

$$SU(2)_L \times U(1)_Y$$

Glashow

Weinberg

Salam



1) Distribua as partículas elementares pelas "representações" do grupo de simetria.

2) Escreva todas as interacções que são invariantes debaixo do grupo de simetria local.

RESULTADO: Teoria que descreve a interacção dos quarks, leptões e bosões de gauge.

mas… Todas as partículas têm massa nula !

Até agora a chávena do CERN é assim: