A experiência

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IF INVESTIGADOR





Sumário

- Revisão da matéria dada
 - As partículas e as forças fundamentais
 - O mecanismo de Higgs
- A experiência ATLAS
 - O LHC e ATLAS
 - A descoberta do Higgs
- Depois da descoberta

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O Modelo Padrão da física de partículas Ou O que pensamos saber do funcionamento do Universo...



Há cerca de 100 anos... A pré-história...





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O átomo

99.999 999 999 9% do volume do átomo é vazio!

núcleo x100





A idade heroica

Observação de raios cósmicos, Plaza de Mulas Argentina, 1950







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E as forças fundamentais?





As forças fundamentais











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Finalmente, o Modelo Padrão da Física de Partículas!



Teste de Avaliação Contínua: O que está errado nestas imagens?





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Agora chegamos aos problems...



Isto não devia funcionar!...



Л



Isto não devia funcionar!...



Л





O Mecanismo de Higgs em bonecos

Robert Brout (1928 – 2011)

Peter Higgs (b. 1929)

François Englert (b. 1932)













A Física Experimental de Partículas de hoje

A história começa há mais de 30 anos...

what is the origin of mass?

what kind of unification may exist beyond the standard model?

what is the origin of flavour?

is there a deeper reason for gauge symmetry?

We have simply too many a priori plausible hypotheses concerning the nature of symmetry breaking in the standard model. Experimentation in the TeV range at the constituent level is bound to provide most essential clues, and the present successes of the pp collider are a very strong encouragement to go to higher energies and to higher luminosities in hadron-hadron collisions.

PROCEEDINGS OF THE ECFA-CERN WORKSHOP

held at Lausanne and Geneva, 21-27 March 1984 LINAC2 \rightarrow **50 MeV** Booster \rightarrow **1.4 GeV** Proton Synchroton (PS) \rightarrow **25 GeV** Super Proton Synchrotron (SPS) \rightarrow **450 GeV** Large Hadron Collider (LHC) \rightarrow **7 TeV**



O que é um protão?



O protão no modelo dos quarks



O protão como o conhecemos hoje



F. Wilczek [Nobel Prize 2004]

Como se distribui a energia dentro dos protões

Partões: quarks e gluões constituintes dos hadrões



Em detalhe: funções de distribuição dos partões



A energias maiores, como as do LHC, a contribuição dos gluões e dos quarks de "mar" aumenta – o LHC colide quarks e gluões!

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Feixes de protões



F








- Tudo contido num feixe de ≈16µm
- Runs típicos duram cerca de 8 horas
- Intensidade diminui devido a perdas
- Depois voltamos a injectar novos feixes

- Energia do feixe:
- 2802 bunches de 1.15x10¹¹ protões
- 7TeV / protão (2015) = 7x10¹² x 1.602x10⁻¹⁹ J
- Dá 362 MJ por feixe...
- Igual à energia cinética de um porta-aviões de 20,000t a viajar a 11,7 nós (21.7 km/h)



TeV





Muon Spectrometer: $|\eta| < 2.7$ Air-core toroids and gas-based muon chambers $\sigma/p_T = 2\%$ @ 50GeV to 10% @ 1TeV (ID+MS)

EM calorimeter: $|\eta| < 3.2$ Pb-LAr Accordion $\sigma/E = 10\%/\sqrt{E \oplus 0.7\%}$

> Hadronic calorimeter: $|\eta| < 1.7$ Fe/scintillator $1.3 < |\eta| < 4.9$ Cu/W-Lar $\sigma/E_{iet} = 50\%/\sqrt{E \oplus 3\%}$

L = 44 m, Ø ≈ 25 m
7000 tonnes
≈10⁸ electronic channels
3-level trigger reducing
40 MHz collision rate to
200 Hz of events to tape

Inner Tracker: $|\eta| < 2.5$, B=2T Si pixels/strips and Trans. Rad. Det. $\sigma/p_T = 0.05\% p_T$ (GeV) $\oplus 1\%$





O calorímetro de telhas (TileCal)





Sistema de ímanes de ATLAS





- Campos magnéticos muito intensos (2T, 4T max) gerados por ímanes supercondutores
- Curvam as partículas carregadas através da força de Lorenz:

$$\mathbf{F} = \mathbf{v} \times \mathbf{B}$$

Map contours: BMOD

Detetores de muões

Thin-gap chambers (TGC)

Cathode strip chambers (CSC)

Resistive-plate chambers (RPC)

End-cap toroid

Monitored drift tubes (MDT) Escola de Professores CERN - 8/9/2017

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Barrel toroid

Extrato de filme: "ATLAS Episode II: the particles strike back" https://cds.cern.ch/record/1096390?ln=ka

Etapas do projecto Tilecal (1993-2009)



1993-1995 R&D-protótipos



1996-2002:construção

1999-2002 Instrumentação



1999-2004: Electrónica



2002-2004: calibrações



2004-2006 Instalação



2007-2009 certificação (raios cósmicos)



2009--→: aquisição/análise dados LHC

Um longo percurso para conseguir as excelentes características do Tilecal no detetor ATLAS Com participação portuguesa: LIP, FCUL, UMinho

Jactos hadrónicos





Trigger: sistema de seleção em tempo real

- 25 ns entre pacotes
 - (i.e. \approx 7.5m à velocidade de *c*)
 - 40 milhões de cruzamentos de feixes por segundo
 - Cada colisão daria ≈1.5Mb
 - => 60Tb por segundo
- Impossível guardar todos os dados
 - E desnecessário!
 - A maioria das colisões é sem interesse
- O sistema de trigger guarda apenas ≈10-15 colisões por cada milhão
- Mas tem que decidir em 2,5µs!!





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10 de Setembro 2008: primeiros feixes!



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Filme em: https://cds.cern.ch/record/2020780?ln=en

A Colaboração ATLAS

3000 cientistas(1000 estudantes)33 países177 universidadese laboratórios

ATLAS em Português



Em Portugal: ~20 investigadores, ~10 – 15 estudantes LIP – Coimbra, Minho, Lisboa FCUL, FCTUC, UM, CEFITEC/UNL, INESC, CFMC Programa de treino de engenheiros no CERN No Brasil:

Cerca de 30 colegas de várias instituições UFRJ-COPPE, UFJF, USP, UFSJ (UFBA, CEFET-RJ, UFF)

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Standard Model Total Production Cross Section Measurements Sta

Status: August 2016



Breve história da descoberta do Higgs

A PHENOMENOLOGICAL PROFILE OF THE HIGGS BOSON

John ELLIS, Mary K. GAILLARD * and D.V. NANOPOULOS ** CERN, Geneva





We should perhaps finish with an apology and a caution. We apologize to experimentalists for having no idea what is the mass of the Higgs boson, unlike the case with charm [3,4] and for not being sure of its couplings to other particles, except that they are probably all very small. For these reasons we do not want to encourage big experimental searches for the Higgs boson, but we do feel that people performing experiments vulnerable to the Higgs boson should know how it may turn up.

Por volta de Julho de 2010, depois de décadas de buscas...



At the LHC



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Combinação estatística



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 $m_{\rm T}$ [GeV]

(a)

Nem sempre se acerta à primeira...



Conferência europeia de Física de Partículas (EPS-HEP) 2011 [6]



"The Brazil Plot"



O Gráfico dos 6 mil milhões de francos suíços



Probabilidade de 1 em 3.5 milhões de o sinal vir de uma flutuação estatística do ruído de fundo



A Descoberta do bosão de Higgs



A Descoberta do bosão de Higgs... ...premiada com o Prémio Nobel 2013:

Peter Higgs, Inglês, nascido em 1929, Univ. Edimburgo



François Englert, Belga, nascido em 1932, U. Libre de Bruxelles



"for the **theoretical discovery** of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the **discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider**"



What if?...









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Notícias frescas


"Há mais coisas no céu e na terra, Horácio, do que sonha a nossa filosofia."





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KEEP CALM AND CHECK BACKUP SLIDES

Estrutura e princípio de funcionamento do calorímetro Tilecal Estrutura periódica, ferro e cintiladores de plástico





• Ok, so we reject background and take only signal events Maybe not so simple:

- Bunch spacing is 25ns: not much time to decide! (25ns x C = 7.5m)
- Put event fragments in memory pipeline to buy time for Level 1 decision
- Pileup of minimum-bias events means longer reconstruction time and higher occupancy
- Not only pileup from same bunch crossing! ATLAS sub-detector response varies from a few ns to about 700 ns (= 28 bunch crossings!)
- Try to rely mostly on high-pT particles







Very good correlation between RPC (trigger chambers) and MDT (precision chambers) hits

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ATLAS

Real name: Johnny Rammond Occupation: Head of the O.P.T.C.- Optic

City Planning Committee Identity: Known to his best friend Toby Whey

Legal status: Johnny Rammond is a United States citizen, Atlas is considered "above the law"

Place of birth: Heffield, Nebraska Marital status: Single

Known relatives: Sturgill Rammond (father, deceased), Dorothy Rammond (mother, deceased), Toby Whey (best friend)

Base of operations: The Optic City Planning Center, as well as his secret laboratory in a hidden chamber of his house **First appearance:** All Time Comics:

Crime Destroyer #1

Origin: Johnny Rammond came to the big city from the Midwest with a couple bucks in his pocket and stars in his eyes. He soon found the reality of Optic City a harsher one than he could have dreamed. After months of hardship and difficulty, he was reduced to sleeping in the sewers of the city with the other underground denizens of the underground communities. One night he came upon a mugging and in trying to stop it, he was assaulted and badly hurt. He managed to crawl back into the sewers. He was washed into a drain which took him into a deeper level of the under city than any human had ever been, one which was constructed by ancient aliens many aeons before human beings arrived on the continent. There, mysterious, cosmic machines whirred to life and, sensing Rammond's innate good nature and well-meaning spirit, imbued him with one of the most powerful weapons on the planet, his multi-channel, anti-matter-powered communicator / signal ring. Rammond's second life as Optic City's mightiest costumed hero began that day. His only weakness is when uncertainty and fear take hold in his mind

Height: 6'1" Weight: 215 lbs. Eyes: Blue Hait: Blonde Powers: Atlas's anti-matter ring gives him multiple powers, including super strength, super speed, the ability to fly, anti-matter rays which shoot out of his eyeballs, the ability to absorb tremendous trauma, the ability to communicate with whoever wears his other ring from anywhere in the universe, etc.

Weapons: Anti-matter power ring

Voltemos um pouco atrás...

Atenção! Os próximos slides contêm equações potencialmente chocantes.

Leonhard Euler(1707-1783)

Emmy Noether (1882 – 1935

Joseph-Louis Lagrange (1736–1813)

Lagrangianos e mecânica clássica

The equations of motion of a system are derived from a scalar Lagrangian function of generalized coordinates and velocities (time derivatives)

$$L(q, \dot{q}) = T - V$$

and from the Euler-Lagrange's equations:

$$\frac{\partial L}{\partial q_j} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} = 0$$

Example

Particle in a conservative potential V. The Lagrangian

$$L = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - V(x, y, z)$$

has derivatives (e.g. for x)

$$\frac{\partial L}{\partial x} = -\frac{\partial V}{\partial x}, \frac{\partial L}{\partial \dot{x}} = m\dot{x}, \frac{d}{dt}(\frac{\partial L}{\partial \dot{x}}) = m\ddot{x}$$

and Euler-Lagrange's equations

$$\frac{\partial L}{\partial q_j} - \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_j} = 0$$

finally give us Newton's familiar 2nd law!

$$m\ddot{x} = -\frac{\partial V}{\partial x}, m\ddot{y} = -\frac{\partial V}{\partial y}, m\ddot{z} = -\frac{\partial V}{\partial z} \Leftrightarrow m\vec{a} = \vec{F}$$

Symmetries and conservation laws

Noether's theorem:

If a system has a continuous symmetry property, then there are corresponding quantities whose values are conserved in time.

Simplest case: Coordinates not explicitly appearing in the Lagrangian

⇒ Lagrangian invariant over a continuous transformation of the coordinates

Example: mass **m** orbiting in the field of a fixed mass **M**
$$L(r, \phi, \dot{r}, \dot{\phi}) = T - V = \frac{1}{2}m\dot{r}^2 + \frac{1}{2}mr^2\dot{\phi}^2 + \frac{GMm}{r}$$

Since the lagrangian doesn't depend explicitly on $\boldsymbol{\phi}$ (symmetry with respect to rotations in space), the Euler-Lagrange equation gives $\frac{d}{dt}\left(\frac{\partial \mathcal{F}_{I}}{\partial \dot{\phi}}\right) \stackrel{\text{space}}{=} 0 \Leftrightarrow \frac{\partial \mathcal{F}_{I}}{\partial \dot{\phi}} \stackrel{\text{space}}{=} mr^{2} \phi \stackrel{\text{space}}{=} J$

Where the **angular momentum J** is a constant of motion!

Agora nos campos quânticos...

Richard Feynman (1918 - 1988)

Schrödinger's cat (?-?)

Erwin Schrödinger (1887 - 1961)



Partículas são ondas num "mar" (campo quântico)

Generalized coordinates are now **fields** (dislocation of each spring)

$$q_i \to \phi_i(x^\mu)$$

In a relativistic theory we must treat space and time coordinates on an equal footing, so the derivatives in the classical equations are now

 $\frac{d}{dt}, \nabla \to \partial_{\mu} = \left(\frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}\right)$ In place of a Lagrangian we have a **Lagrangian density** (we call it Lagrangian anyway, just to be confusing) $L(q_i, \frac{dq_i}{dt}) \to \mathcal{L}(\phi_i, \partial_{\mu}\phi_i) \text{ with: } L = \int \mathcal{L}d^3x$ The new Euler-Lagrange equation now becomes

$$\partial_{\mu} \left(\frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \phi_i)} \right) - \frac{\partial \mathcal{L}}{\partial \phi_i} = 0$$

Gauge invariance

Take the Dirac Lagrangian for a field ψ representing a spin- $\frac{1}{2}$ particle, for example an electron:

$$\mathcal{L} = i\hbar\bar{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\bar{\psi}\psi$$

It is invariant under a global U(1) phase transformation like:

$$\psi(x) \to \psi'(x) = e^{iq\chi}\psi(x)$$

Where $\mathbf{\chi}$ is a constant

$$\mathcal{L}' = e^{-iq\chi} e^{iq\chi} (i\hbar\bar{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\bar{\psi}\psi) = \mathcal{L}$$

Local gauge invariance and interactions

If $\mathbf{\chi} = \mathbf{\chi}(x)$ then we get extra terms in the Lagrangian:

$$\mathcal{L}' = i e^{-iq\chi} \bar{\psi} \gamma^{\mu} [e^{iq\chi} \partial_{\mu} \psi + iq(\partial_{\mu} \chi) e^{iq\chi} \psi] - m e^{-iq\chi} e^{iq\chi} \bar{\psi} \psi$$
$$= \mathcal{L}' - q \bar{\psi} \gamma^{\mu} (\partial_{\mu} \chi) \psi$$

But we can now make the Lagrangian invariant by adding an *interaction term* with a new **gauge** field A_{μ} which transforms as:

$$A_{\mu} \to A'_{\mu} = A_{\mu} - \partial_{\mu}\chi$$

We get:

$$\mathcal{L} = i\bar{\psi}\gamma^{\mu}\partial_{\mu}\psi - m\bar{\psi}\psi - q\bar{\psi}\gamma^{\mu}A_{\mu}\psi$$

A few things to note:

- 1. Gauge theories are renormalizable, i.e. calculable without infinities popping up everywhere (Nobel prize of t'Hooft and Veltman)
- 2. The new gauge field A_{μ} is the photon in QED
- 3. The mass of the fermion is the coefficient of the term on $\psi\overline{\psi}$
- 4. There is no term in $A_{\mu}A^{\mu}$ (the photon has zero mass) \rightarrow this is the beginning of the Higgs story...

$$\psi(x) \to \psi'(x) = e^{iq\chi}\psi(x)$$

Esfera original



Transformação global



Transformação local



χ = constante

 $\mathbf{\chi}=\mathbf{\chi}\;(\mathbf{x})$

Mass of elementary particles and gauge bosons

 $\mathcal{L}_{QED} = \bar{\psi}(i\gamma^{\mu}\partial_{\mu} - m_e)\psi - e\bar{\psi}\gamma^{\mu}\psi A_{\mu} - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}m_{\gamma}A_{\mu}A^{\mu}$

To keep the Lagrangian gauge invariant (against a U(1) local phase transformation) the photon field transforms as:

$$A_{\mu} \to A'_{\mu} = A_{\mu} - \partial_{\mu}\chi$$

But the A^{μ} mass term breaks the invariance of the Lagrangian:

$$\frac{1}{2}m_{\gamma}A_{\mu}A^{\mu} \to \frac{1}{2}m_{\gamma}(A_{\mu} - \partial_{\mu})(A^{\mu} - \partial^{\mu}\chi) \neq \frac{1}{2}m_{\gamma}A_{\mu}A^{\mu}$$

For the SU(2)_L gauge symmetry transformations of the **weak** interaction the fermion mass term $m_e \overline{\Psi} \psi$ also breaks invariance!

Bottom line: the SM (without the Higgs mechanism) results in wrong calculations and breaks down for massive particles

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O Mecanismo de Higgs

Robert Brout (1928 – 2011)

Peter Higgs (b. 1929)

François Englert (b. 1932) • Introduce a SU(2) doublet of spin-0 complex fields

 $\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1 + i\phi_2 \\ \phi_3 + i\phi4 \end{pmatrix}$

 $\mathcal{L} = (\partial_{\mu}\phi)^{\dagger}(\partial^{\mu}\phi) - V(\phi)$

 $V(\phi) = \mu^2 \phi^{\dagger} \phi + \lambda (\phi^{\dagger} \phi)^2$

- With a potential
- For $\lambda > 0$, $\mu^2 < 0$ the potential has a minimum at the origin
- For $\lambda > 0$, $\mu^2 < 0$ the potential has an infinite number of minima at:

$$|\phi| = \frac{v}{\sqrt{2}} = \sqrt{-\frac{\mu^2}{2\lambda}}$$
 .

The choice of vacuum (lowest energy state of the field) breaks the symmetry of the Lagrangian Ricardo Gonçalo

Electroweak symmetry breaking

- In the Standard Model with no Higgs mechanism, interactions are symmetric and particles do not have mass
- Electroweak symmetry is broken:
 - Photon does not have mass
 - W, Z have a large mass
- Higgs mechanism: mass of W and Z results from the Higgs mechanism
- Masses of fermions come from a direct interaction with the Higgs field



see J.Varela's lecture ⁹⁵

EWK Symmetry Breaking in Pictures



We have at this point a massive scalar field with vacuum expectation value *v* and mass

$$m_h = \sqrt{2\lambda}v$$

4 gauge fields: $W^{(1)}$, $W^{(2)}$, $W^{(3)}$, and $B^{(1)}$ which transform to give the massive W⁺, W⁻ and Z, and massless A (the photon)

$$\begin{split} m_{W^{(1)}} &= m_{W^{(2)}} = m_W = \frac{1}{2} g_W v \\ m_A &= 0 \qquad \Leftrightarrow v = 246 \text{GeV} \\ m_Z &= \frac{1}{2} v \sqrt{g_W^2 + g^2} \end{split}$$

with g, g_w the couplings of electromagnetic and weak forces

Defining the Weinberg angle as

we also get the relation between the masses of W and Z

$$\frac{g}{g_W} = \tan \theta_W$$

$$\frac{m_W}{m_Z} = \cos \theta_W$$

Fermions get their masses from interaction terms with the Higgs field (Yukawa coupling)

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Finally! What we think we know:

- Higgs mass (was) the only unknown parameter
- We can give mass to W[±] and Z while keeping the photon massless
- Relation between masses of W and Z
- Higgs couples to W and Z with strengths proportional to their masses
- Higgs couples to all fermions with a strength proportional to their mass



 $m_h = \sqrt{2\lambda v}$

 $\frac{m_W}{m_Z}$ $=\cos\theta_W$

 $g_f = \sqrt{2 \frac{m_f}{m_f}}$