

Right-Handed neutrino searches at FCC-(ee)

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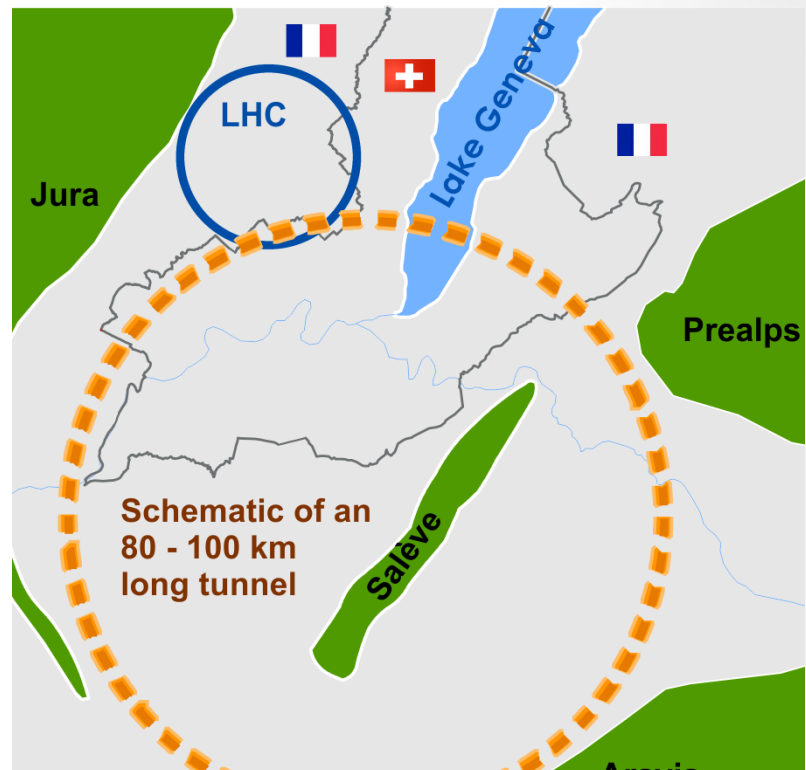
EPS 2019, Genft

On behalf of the FCC-ee physics group

Thanks to Alain Blondel and Oliver Fisher for the input material

The FCC

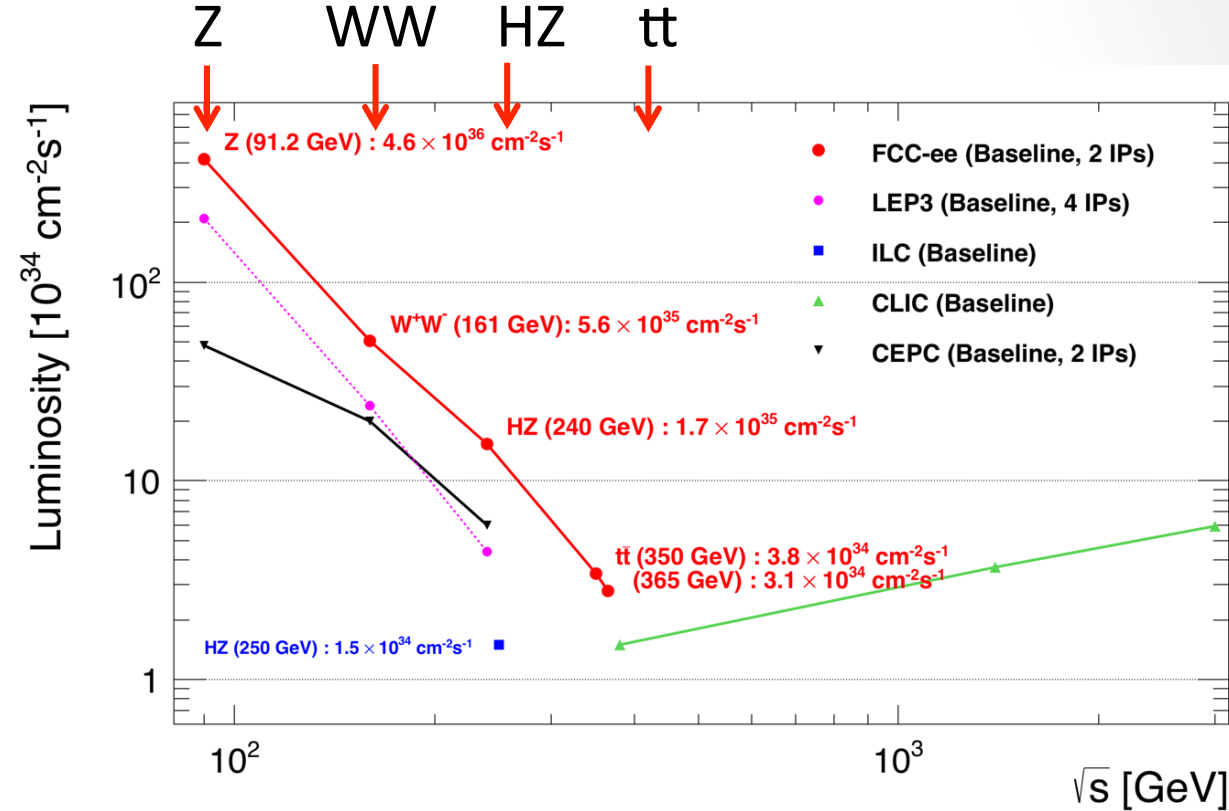
- International collaboration to Study Colliders fitting in a new ~ 100 km infrastructure, fitting in the Geneva area
- Ultimate goal:
 ≥ 100 TeV pp-collider (FCC-hh)
- defining infrastructure requirements
- Few possible first steps:
 - e^+e^- collider (FCC-ee)
High Lumi, $E_{\text{CM}} = 90\text{-}500$ GeV
 - HE-LHC $16\text{T} \Rightarrow 27$ TeV
in LEP/LHC tunnel
 - Low energy FCC $< 50\text{TeV}$ (after ESPPU)
 - Possible addition:
 - p-e (FCC-he) option
- This is the center of discussions for the European Strategy Update



The way by FCC-ee is the fastest and cheapest way to 100 TeV, also produces the most physics. Preferred scenario presented in the CDR.

<https://cerncourier.com/cern-thinks-bigger/>

It's also a good start for a $\mu\mu\text{C}$!



Z peak	E_{cm} : 91 GeV	5×10^{12}	$e^+e^- \rightarrow Z$	LEP x 10^5
WW threshold	E_{cm} : 161 GeV	10^8	$e^+e^- \rightarrow WW$	LEP x $2 \cdot 10^3$
ZH threshold	E_{cm} : 240 GeV	10^6	$e^+e^- \rightarrow ZH$	Never done
tt threshold	E_{cm} : 350 GeV	10^6	$e^+e^- \rightarrow tt$	Never done

FCC-ee running scenario

From FCC CDR Volume 2

Table 2.1: Run plan for FCC-ee in its baseline configuration with two experiments. The number of WW events is given for the entirety of the FCC-ee running at and above the WW threshold.

Phase	Run duration (years)	Center-of-mass Energies (GeV)	Integrated Luminosity (ab^{-1})	Event Statistics
FCC-ee-Z	4	88-95	150	3×10^{12} visible Z decays 10^8 WW events 10^6 ZH events 10^6 $t\bar{t}$ events
FCC-ee-W	2	158-162	12	
FCC-ee-H	3	240	5	
FCC-ee-tt	5	345-365	1.5	

FCC-ee discovery potential

Today we do not know how nature will surprise us. A few things that FCC-ee could do

- Explore

- ~20-50 (stat 400...) fold improved precision on many EW quantities
 - $\times 5-7$ in mass m_Z , m_W , m_{top} , $\sin^2\theta_W^{eff}$, R_b , $\alpha_{QED}(m_Z)$, $\alpha_s(m_Z m_W m_\tau)$, top quark couplings
- Model-independent Higgs width and couplings measurements at percent-permil level
- 10-100 TeV energy scale (and beyond) with Precision Measurements (through EFT)
- $\sim 3\sigma$ of effect of Higgs self-coupling from Vertex corrections (also maybe directly with FCC-ee 500GeV)
- Only machine with possible investigation of $H\mu\mu$ coupling at $\sqrt{s} = m_H$

- Discover

- Violation of flavour conservation or universality and unitarity of PMNS @ 10^{-5}
- FCNC ($Z \rightarrow \mu\tau$, $e\tau$) in $5 \cdot 10^{12}$ Z decays and τ BR in $2 \cdot 10^{11}$ $Z \rightarrow \tau\tau$ (also FCNC in top decays with 10^6 tops)
- Flavour physics with 10^{12} bb events ($B \rightarrow s \tau\tau$ etc..)
- Dark matter as «invisible decay» of H or Z (or in LHC loopholes)

- Direct Discovery

- Very weakly coupled particle in 5-100 GeV energy scale such as: Right-Handed neutrinos, Dark Photons etc...

- Not only a «Higgs Factory», «Z factory» and «top» are important for 'discovery potential' (also QCD)

Electroweak eigenstates

$\begin{pmatrix} e \\ \nu_e \end{pmatrix}_L$	$\begin{pmatrix} \mu \\ \nu_\mu \end{pmatrix}_L$	$\begin{pmatrix} \tau \\ \nu_\tau \end{pmatrix}_L$	$(e)_R$	$(\mu)_R$	$(\tau)_R$	Q= -1
I=1/2			$(\nu_e)_R$	$(\nu_\mu)_R$	$(\nu_\tau)_R$	Q= 0
			I=0			

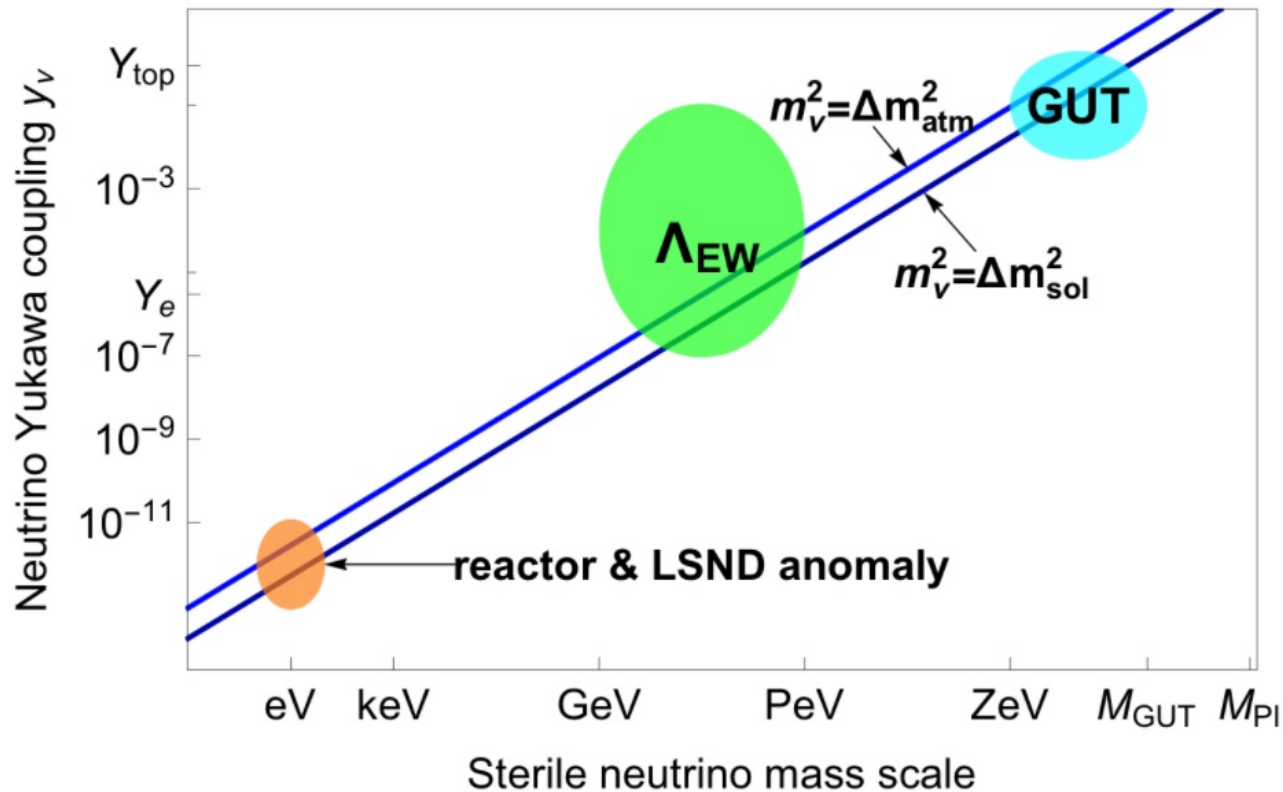
- We measure neutrino parameters, but:
 - No right-handed neutrinos in the SM
 - No mass matrix, no mixing of the neutrino flavour states
- ⇒ Neutrino oscillations are evidence for physics beyond the SM.

- Right handed neutrinos are singlets,
 - No weak interaction
 - No EM interaction
 - No strong interaction
- Can't produce them, Can't detect them
 - So why bother? (also called Sterile)

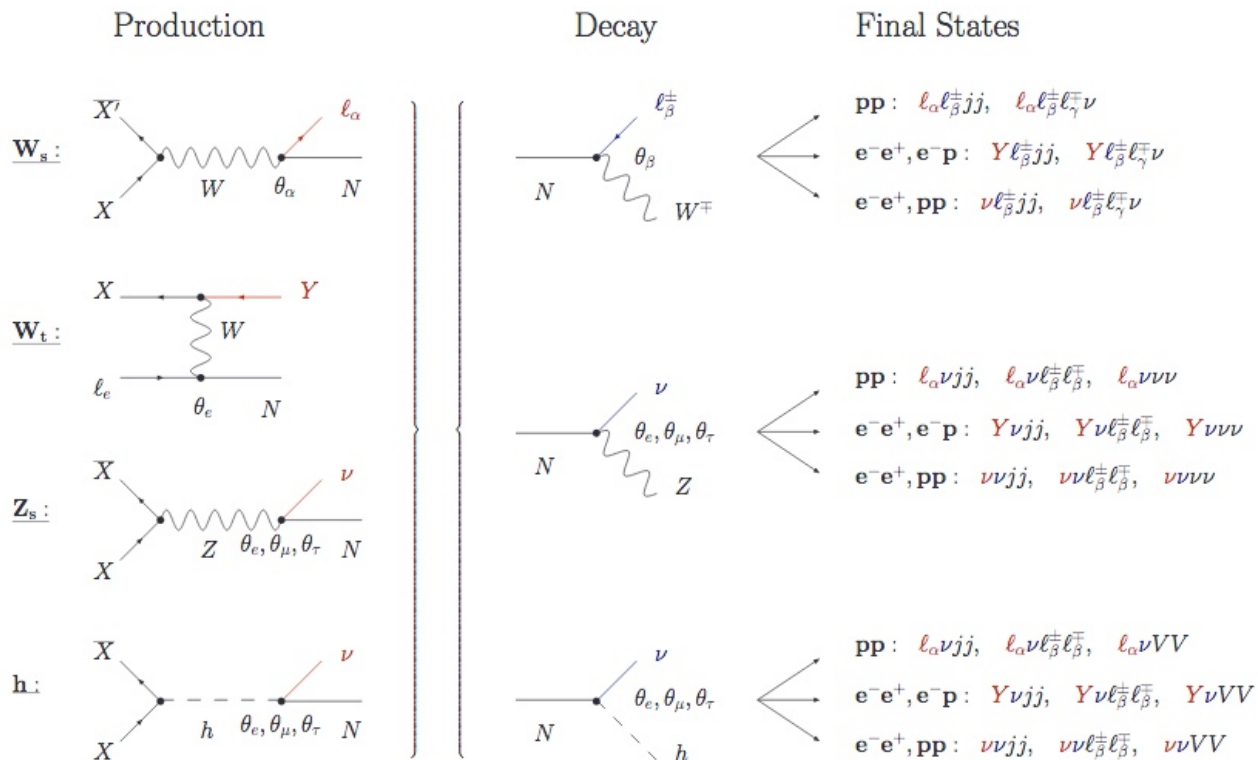
The Seesaw mechanism with RH neutrinos

- Economic extension by adding a number of Fermionic singlets
 - “Right-handed” or “sterile” neutrinos
- Two mass-differences
 - At least two sterile neutrinos
 - New mass scale, a priori unrelated to the known ones
- Many constraints from experiments on all energy scales
- May be connected to e.g. Dark Matter and Baryogenesis

The Big Picture



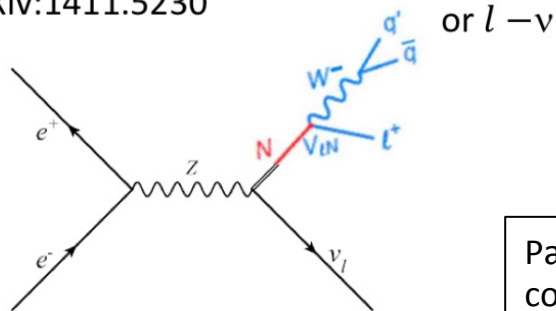
Searches at FCC



Displaced vertex searches at FCC-ee

- Test minimal type I seesaw hypothesis
- Together with ΔM also tests the compatibility with leptogenesis
- Long life time \rightarrow detached vertex for $\sim < M_Z$
- Backgrounds: four fermions
 - $e^+e^- \rightarrow W^+W^-, e^+e^- \rightarrow Z^*(\nu\nu)(Z/\gamma)^*$

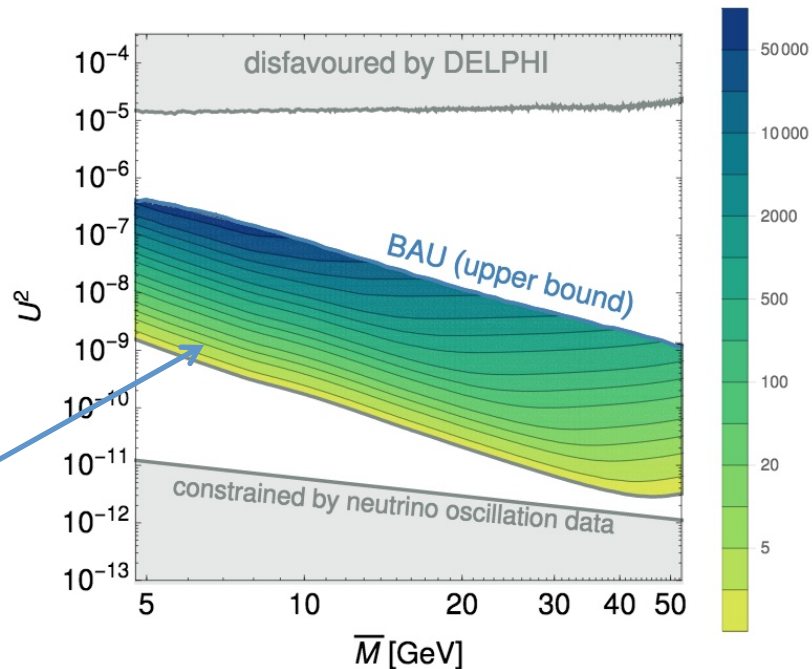
arXiv:1411.5230



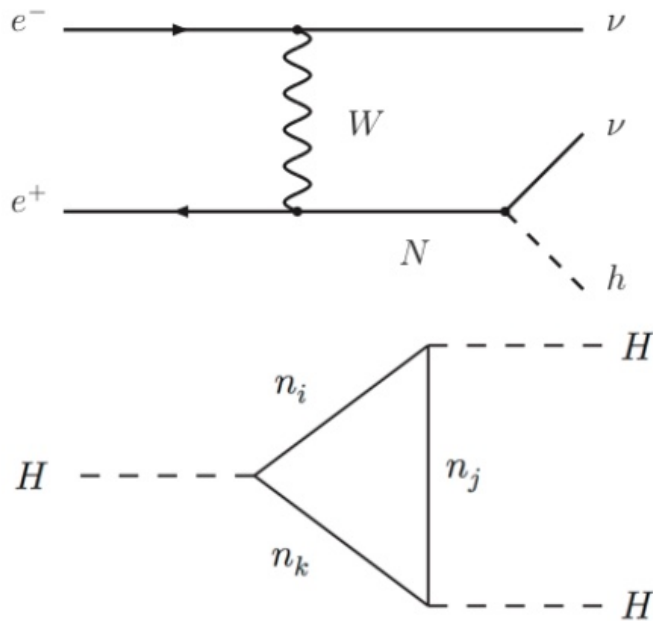
Parameter points consistent with baryogenesis

Antusch et al. JHEP 1809 (2018) 124

NO, FCC-ee at $\sqrt{s} = 90$ GeV



Indirect searches in Higgs properties



- Additional mono-Higgs production mechanism
- New Higgs decay channels:
 - Modification of Higgs branching ratios
 - New exotic decay channels: $h \rightarrow \nu N$, $N \rightarrow \text{SM}$
 - New invisible decay channels
- N contribution to the triple Higgs coupling

Outlook for FCC-hh

- Z factory like FCC-ee offers a clean method for detection of Heavy Right-Handed neutrinos
- W bosons are less abundant at the lepton colliders
 - At the 100 TeV FCC-hh W is the dominant particle: Expect 10^{13} real W's
 - There is a lot of pile-up/backgrounds/lifetime/trigger issues which need to be investigated
- But.... in the regime of long lived HNLs the simultaneous presence of
 - the initial lepton from W decays
 - the detached vertex with kinematically constrained decay
 - -> Would allow for a significant background reduction
- Could also served as a characterization both in flavour and charge of the produced neutrino
 - information of the flavour sensitive mixing angles
 - test of the fermion violating nature of the intermediate (Majorana) particle

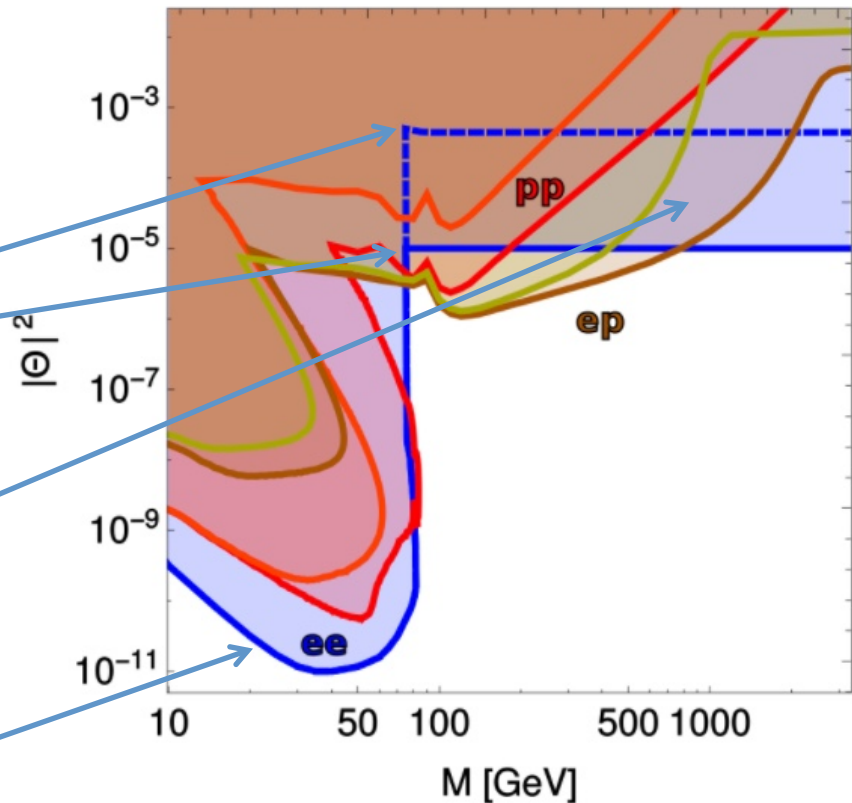
Overview of sensitivities

- At one-sigma confidence level
- ep and pp at parton level

EWPO sensitivity up to very high mass scales

Good sensitivity reach from FCC-hh and FCC-eh

Best sensitivity from displaced vertex searches at FCC-ee



Synergy and complementarity

- FCC-ee

- Highest sensitivity in the low mass regime ($M < m_W$)
 - test model predictions: seesaw, leptogenesis
- SM precision tests have high sensitivity; mass independent
 - Test heavy neutrinos up to $\sim 60\text{TeV}$
 - Not sensitive to the model details

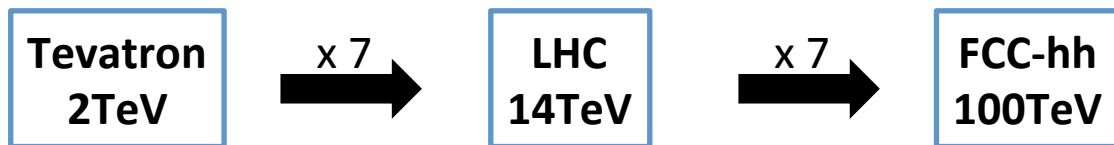
- FCC-hh and he

- Direct test of lepton-flavor and number violation
 - Number of heavy neutrino generations and their masses
- Indirect test via measurement of Higgs potential
- Sensitive to high mass regime

Conclusion

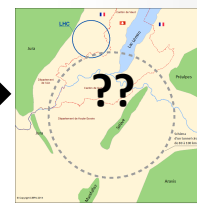
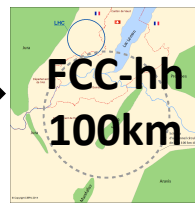
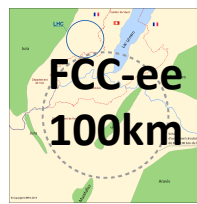
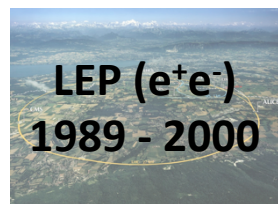
- The FCC design study is establishing the feasibility or the path to feasibility of an ambitious set of colliders after LEP/LHC, at the cutting edge of knowledge and technology.
- Both FCC-ee and FCC-hh have outstanding physics cases
 - each in their own right
 - the sequential implementation of FCC-ee, FCC-hh, would maximise the physics reach
- FCC has unique prospects of testing model predictions.
- Attractive scenarios of staging and implementation (budget!) cover more than 50 years of exploratory physics, taking full advantage of the synergies and complementarities
- Neutrino mass physics should be a benchmark for future collider studies!

A 100km circular collider as next the step



27km tunnel

The next step: 100km tunnel

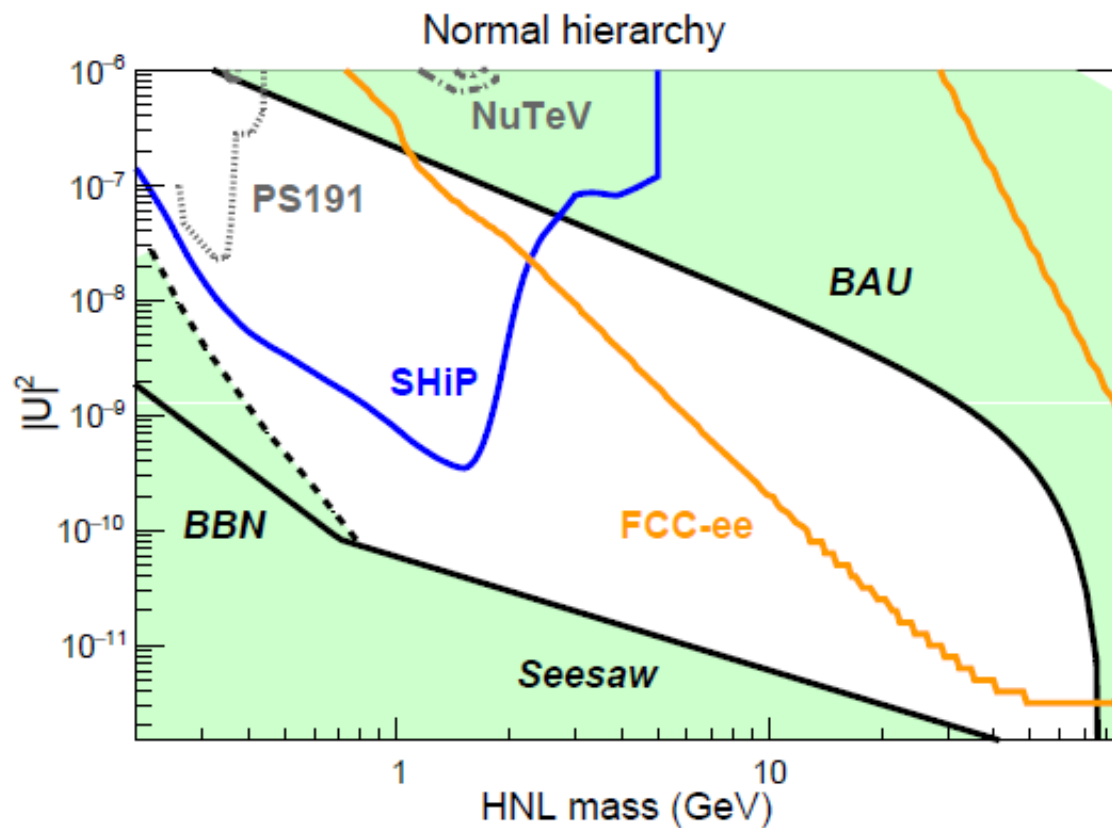


The FCC design study is establishing the feasibility of an ambitious set of colliders after LEP/LHC, at the cutting edge of knowledge and technology

Both FCC-ee and FCC-hh have outstanding physics cases

We are preparing to move to the next step, as soon as possible (EPPSU)

Bonus



(a) Decay length $500 \mu\text{m}$ to 2 m

With $5 \cdot 10^{12} Z$

Manifestation of Right-Handed neutrinos

One see saw family

$$\theta \approx (m_D / M)$$

$$m_\nu \approx m_D^2 / M$$

$$m_N \approx M$$

$$|U|^2 \propto \theta^2 \approx m_\nu / m_N$$

$$\nu = \nu_L \cos\theta - N^c_R \sin\theta$$

$$N = N_R \cos\theta + \nu_L^c \sin\theta$$

What is produced in W,Z decays is:

$$\nu_L = \nu \cos\theta + N \sin\theta$$

ν = light mass eigenstate

N = heavy mass eigenstate

$\neq \nu_L$ active neutrino which couples to weak inter

$\neq N_R$ which does not

- mixing with active neutrinos leads to various observable consequences
 - if very light (eV) , possible effect on neutrino oscillations
 - if in keV region (dark matter), monochromatic photons from galaxies with $E = m_N/2$
- possibly measurable effects at High Energy
 - If N is heavy it will decay in the detector (not invisible)
 - PMNS matrix unitarity violation and deficit in Z «invisible» width
 - Higgs, Z, W visible exotic decays $H \rightarrow \nu_i N_i$ and $Z \rightarrow \nu_i N_i$, $W \rightarrow l_i N_i$
 - also in K, charm and b decays via $W^* \rightarrow l_i \pm N$, $N \rightarrow l_j \pm$ with any of six sign and lepton flavour combination
 - violation of unitarity and lepton universality in Z, W or τ decays
- Couplings are very small (m_ν / m_N) (but *who knows?*) and generally seem out of reach at high energy colliders.

(indirect) Effect of RH ν on EW precision obs.

- The relationship $|\mathbf{U}|^2 \propto \theta^2 \approx \mathbf{m}_\nu / m_N$ is valid for one family see-saw
- For two or three families the mixing can be larger
- *Shaposhnikov, Antush and Fisher*, have shown that a slight # in Majorana mass can generate larger mixing between the left- and right-handed neutrinos
- « $\mathbf{vL} = \mathbf{v} \cos\theta + \mathbf{N} \sin\theta$ » $\rightarrow (\cos\theta)^2$ becomes parametrized as $1 + \varepsilon_{\alpha\beta}$ ($\varepsilon_{\alpha\alpha}$ is negative) the coupling to light ‘normal’ neutrinos is typically reduced.
- In the $G_F, M_Z \propto_{\text{QED}}$ scheme, G_F (extracted from $\mu \rightarrow e \nu_e \nu_\mu$) and g should be increased.
- This leads to correlated variations of all predictions upon e or μ neutrino mixing.
- Only the ‘number of neutrinos’ (R_{inv} and $\sigma_{\text{had}}^{\text{peak}}$) and the tau specific CC observables (tau decays) are sensitive to the tau-neutrino mixing.

Prediction in MUV	Prediction in the SM	Experiment
$[R_\ell]_{\text{SM}} (1 - 0.15(\varepsilon_{ee} + \varepsilon_{\mu\mu}))$	20.744(11)	20.767(25)
$[R_b]_{\text{SM}} (1 + 0.03(\varepsilon_{ee} + \varepsilon_{\mu\mu}))$	0.21577(4)	0.21629(66)
$[R_c]_{\text{SM}} (1 - 0.06(\varepsilon_{ee} + \varepsilon_{\mu\mu}))$	0.17226(6)	0.1721(30)
$[\sigma_{\text{had}}^0]_{\text{SM}} (1 - 0.25(\varepsilon_{ee} + \varepsilon_{\mu\mu}) - 0.27\varepsilon_\tau)$	41.470(15) nb	41.541(37) nb
$[R_{\text{inv}}]_{\text{SM}} (1 + 0.75(\varepsilon_{ee} + \varepsilon_{\mu\mu}) + 0.67\varepsilon_\tau)$	5.9723(10)	5.942(16)
$[M_W]_{\text{SM}} (1 - 0.11(\varepsilon_{ee} + \varepsilon_{\mu\mu}))$	80.359(11) GeV	80.385(15) GeV
$[\Gamma_{\text{lept}}]_{\text{SM}} (1 - 0.59(\varepsilon_{ee} + \varepsilon_{\mu\mu}))$	83.966(12) MeV	83.984(86) MeV
$[(s_{W,\text{eff}}^{\ell,\text{lep}})^2]_{\text{SM}} (1 + 0.71(\varepsilon_{ee} + \varepsilon_{\mu\mu}))$	0.23150(1)	0.23113(21)
$[(s_{W,\text{eff}}^{\ell,\text{had}})^2]_{\text{SM}} (1 + 0.71(\varepsilon_{ee} + \varepsilon_{\mu\mu}))$	0.23150(1)	0.23222(27)

From arXiv:1407.6607

Table 1: Experimental results and SM predictions for the EWPO, and the modification in the MUV scheme, to first order in the parameters $\varepsilon_{\alpha\beta}$. The theoretical predictions and experimental values are taken from Ref. [16]. The values of $(s_{W,\text{eff}}^{\ell,\text{lep}})^2$ and $(s_{W,\text{eff}}^{\ell,\text{had}})^2$ are taken from Ref. [17].