



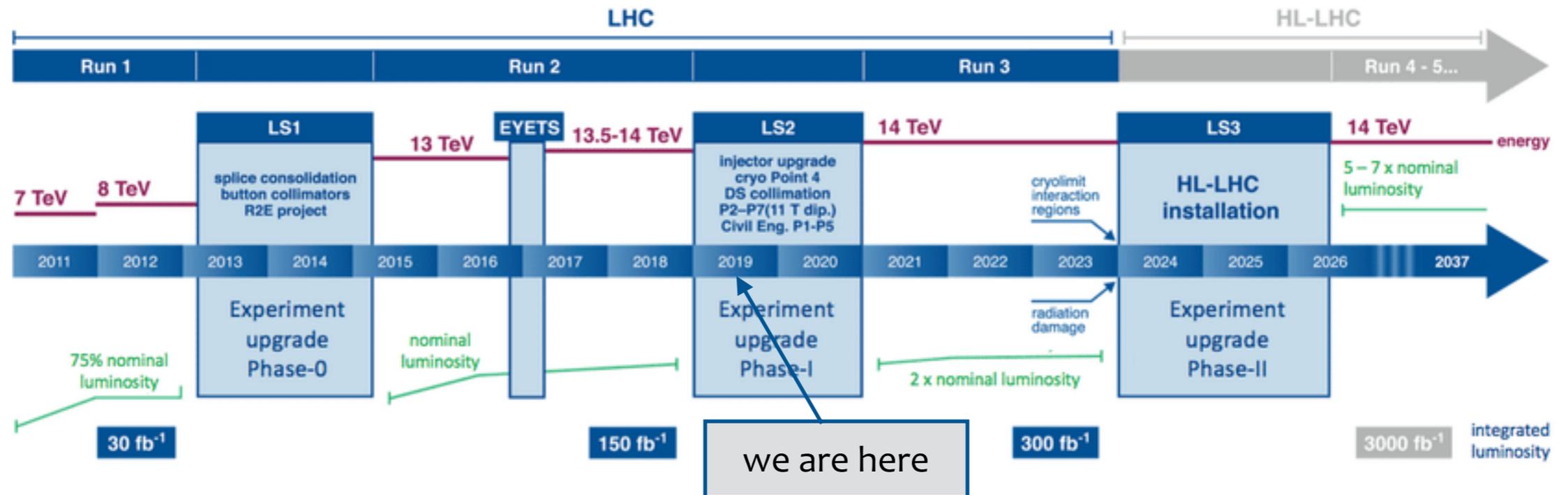
Higgs boson measurements at the High Luminosity LHC with CMS

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- After Higgs discovery, mandate of the LHC experiments is to characterize the resonance by measuring its properties precisely
- Coupling of this scalar to other particles determines if it belongs to standard model (SM) or corresponds to *New Physics*
- Status of relevant measurements performed by the LHC experiments:
 - ➔ couplings to Vector bosons: largely constrained to SM values with Run 1+Run 2 data
 - ➔ couplings to fermions: only to third gen particles measured with large uncertainty
 - ➔ self-coupling: not accessible till now
- 2026 onward LHC will operate in the high luminosity mode (HL-LHC)
 - ➔ Instantaneous luminosity $\sim 5-7.5 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
 - ➔ $\sim 300 \text{ fb}^{-1}/\text{year} \rightarrow 3000 \text{ fb}^{-1}$ by 2036
 - ➔ Access to rare processes, like di-Higgs production
 - ➔ Higgs factory: HL-LHC will produce $\sim 150\text{M}$ Higgs
 - ➔ Average pileup anticipated 150-200, harsh experimental condition



- Very busy environment, radiation damage of the detectors
- Phase II upgrade ongoing to maintain physics performance
 - ➔ Pixel and strips system granularity to be increased, acceptance increase to $|\eta| < 4$
 - ➔ ECAL and HCAL endcaps to be replaced by High Granularity Calorimeter (HGC)
 - ➔ Increase in muon acceptance; GEM detectors already being installed
 - ➔ Addition of MIP Timing Detector for pileup suppression
 - ➔ Improved L1T, HLT, DAQ...
 - ➔ Detector upgrades described in detail in dedicated TDRs
- Physics potential with Phase II improved detectors has been estimated ([YR2018](#))

- Based on results or analysis strategies for 2016 data
- Most of the time at HL-LHC statistical uncertainty is not an issue, evolution of systematic uncertainties studied with time
- Consider two different integrated lumi (L): 300 fb^{-1} and 3000 fb^{-1}
- Uncertainty scenarios following YR recommendation:
 - ➔ **Run 2 systematic uncertainties (S1)**: event yields scale with luminosity, systematics uncertainties remain unchanged w.r.t. Run 2
 - ➔ **YR2018 systematic uncertainties (S2)**: event yields scale with luminosity, theoretical uncertainties reduce by factor $1/2$, experimental uncertainties reduce by $1/\sqrt{L}$ until they reach a lower threshold
- Uncertainty determination methods:
 - ➔ Simple projection: scale Run 2 analysis to Phase II statistics, assume detector conditions are unchanged, incorporating the above uncertainty scenarios
 - ➔ Analyses using samples with HL-LHC detector simulation (Delphes): re-evaluate Run 2 analysis strategies using upgraded detector conditions and pileup effects, S2 uncertainty values

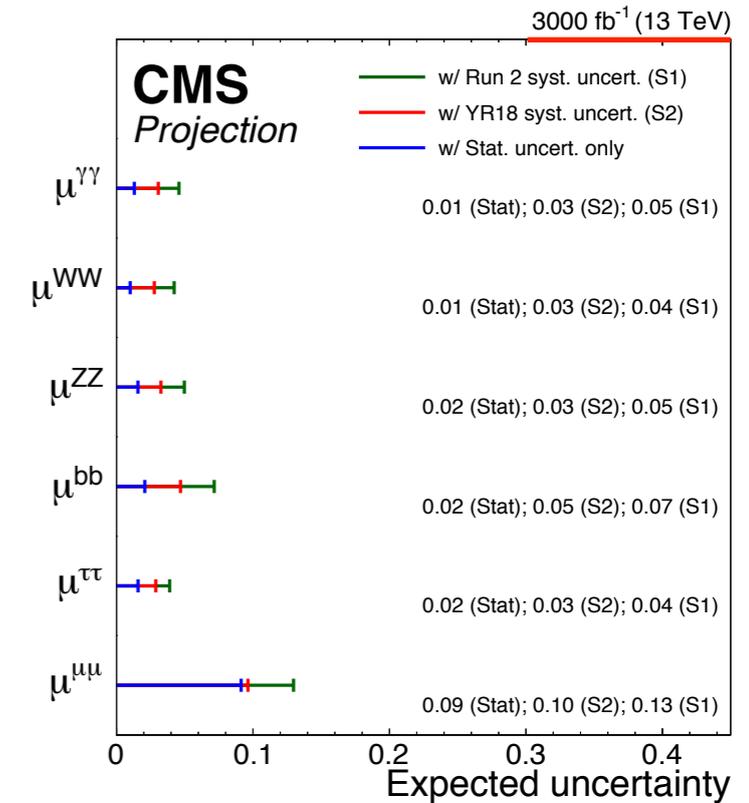
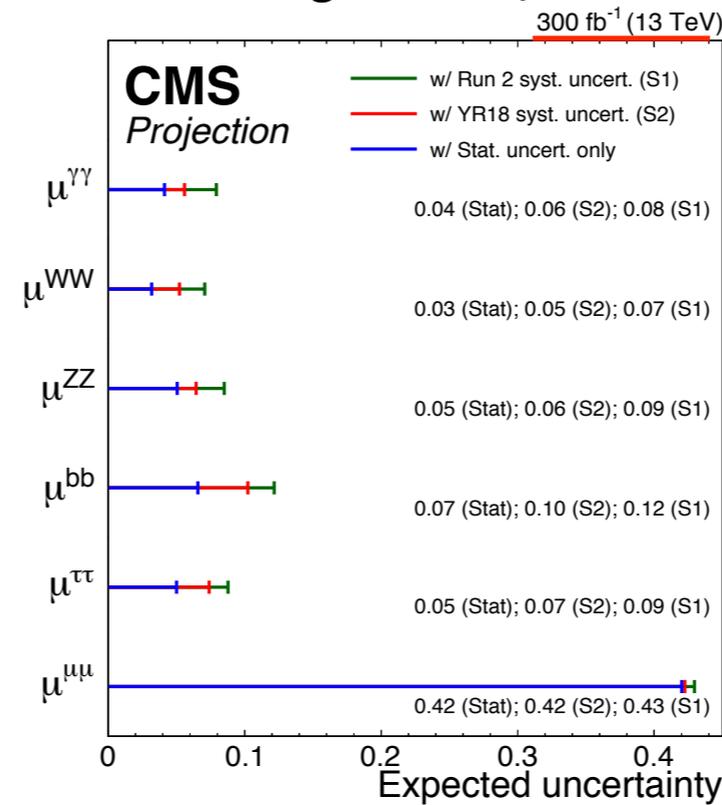
- Signal strength parameters defined to scale cross section (production) and branching ratio (decay):

$$\mu_i^f = \mu_i \times \mu^f = \frac{\sigma_i}{\sigma_i^{\text{SM}}} \times \frac{\text{BR}_f}{\text{BR}_f^{\text{SM}}}$$

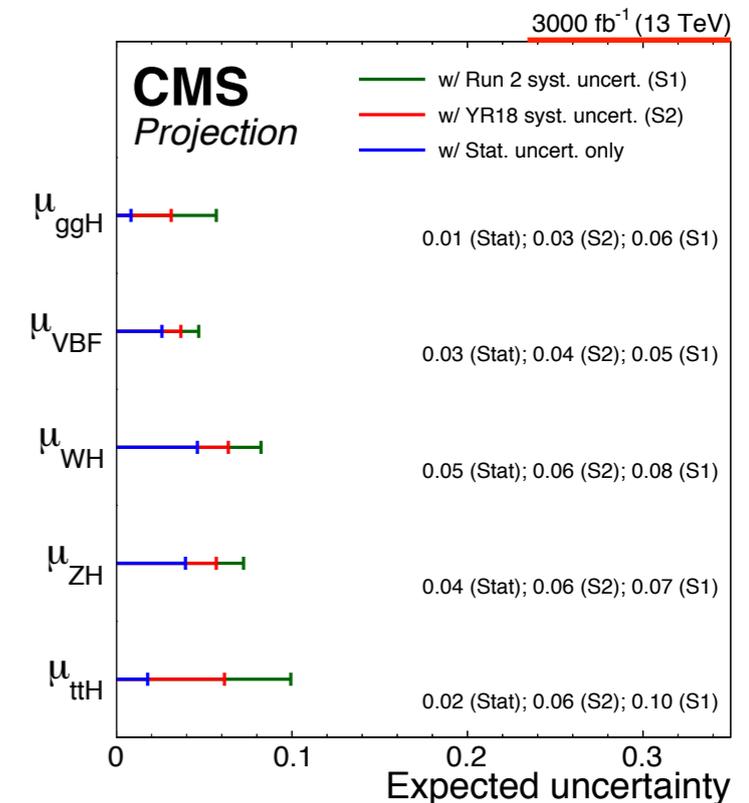
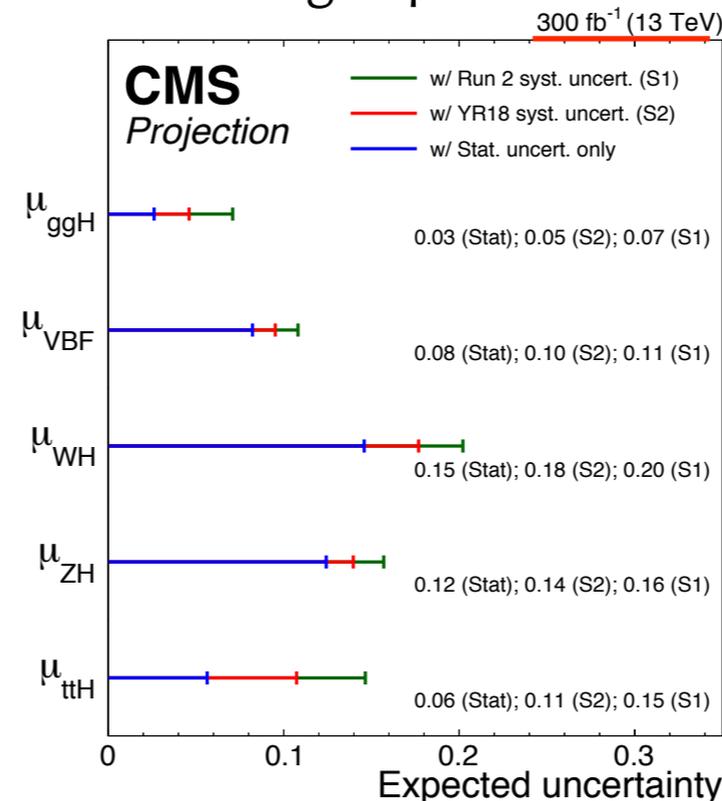
Precision reaches upto 3-10%,
main uncertainty due to signal
theory systematics

To be compared with Run 2
achievement of 10-50%

according to decay:



according to production:



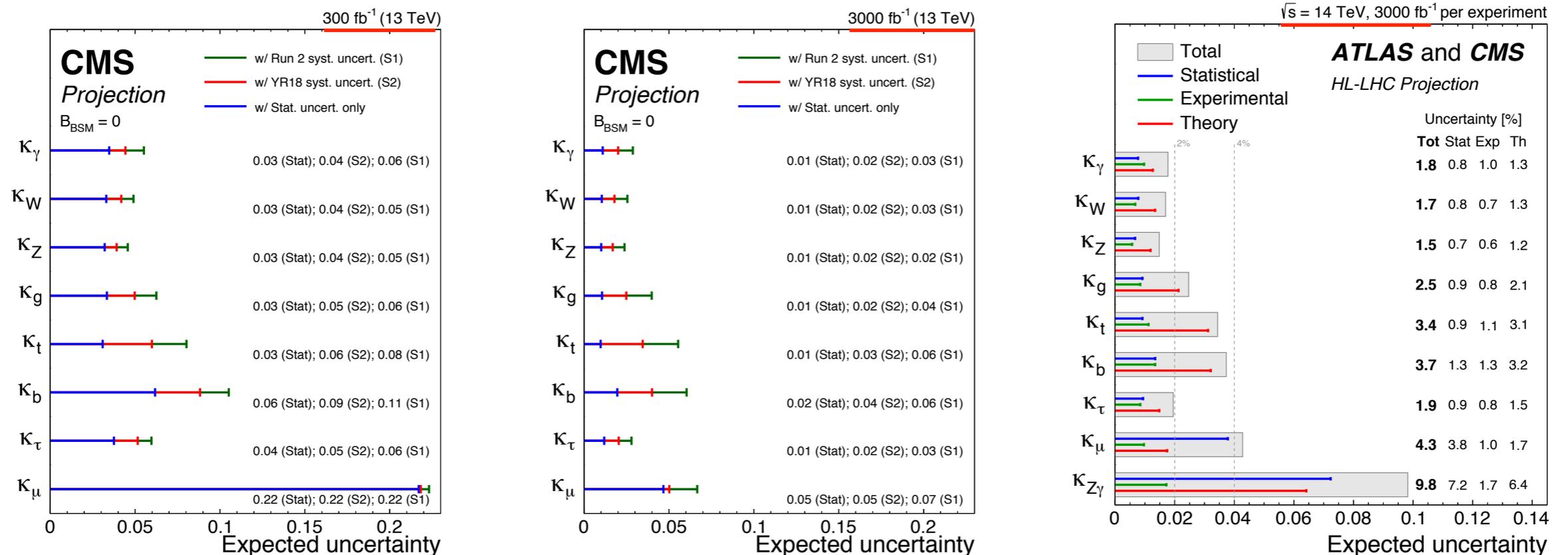
- Coupling strength modifier or kappa model: scale factors used to determine deviation from standard model (SM) couplings:

$$\kappa_i^2 = \frac{\sigma_i}{\sigma_i^{\text{SM}}} \quad \kappa_i^2 = \frac{\Gamma_i}{\Gamma_i^{\text{SM}}}$$

- Total width determined as:

$$\Gamma_H = \frac{\kappa_H^2 \cdot \Gamma_H^{\text{SM}}}{1 - B_{\text{BSM}}} \quad \kappa_H^2 = \sum_j B_{\text{SM}}^j \kappa_j^2$$

(YR2018)

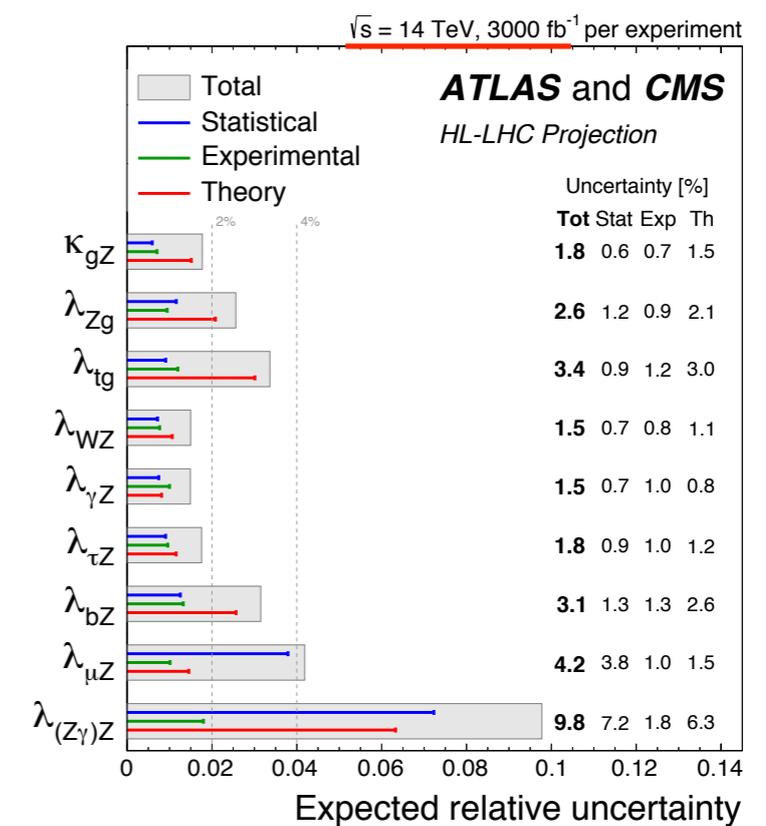
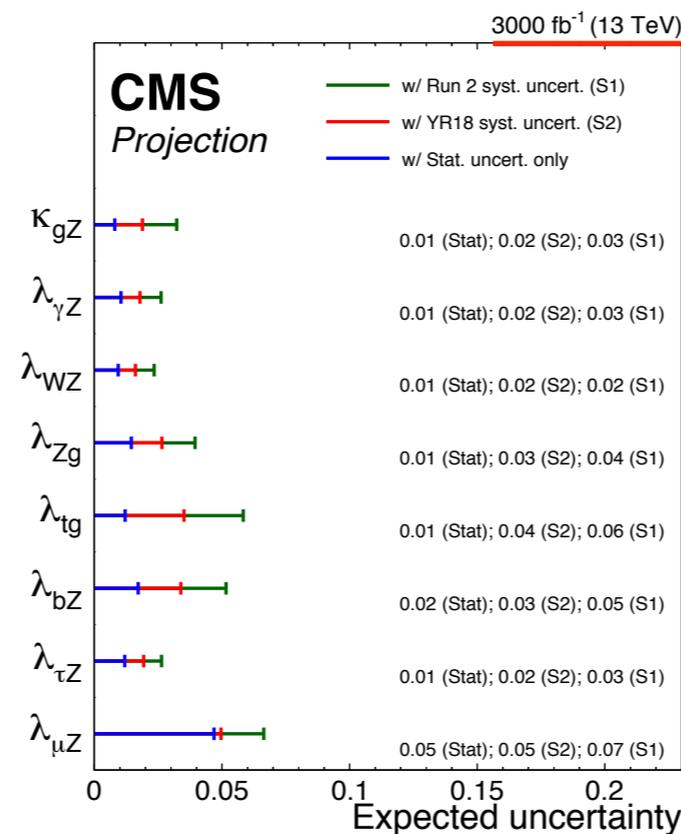
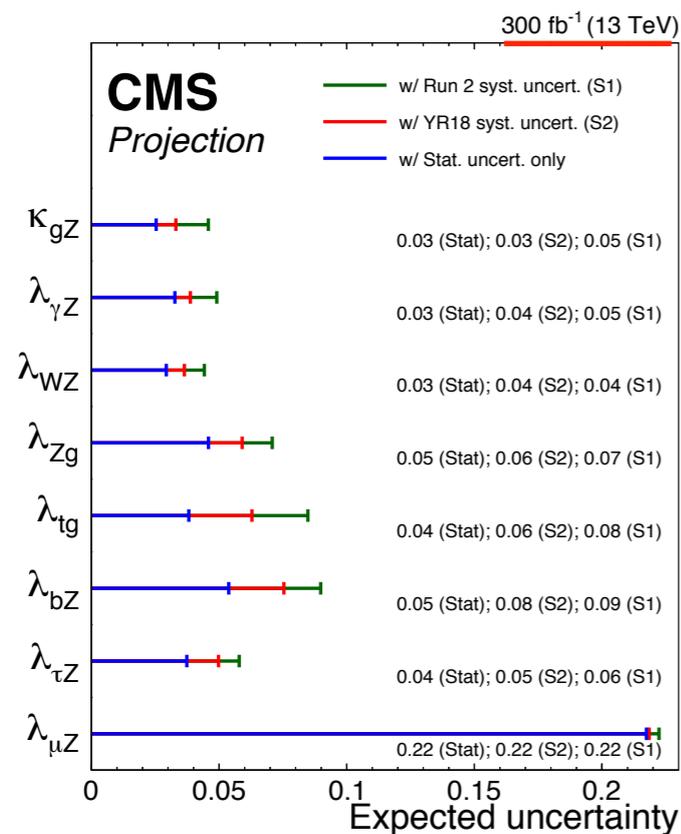


Most of the uncertainties reach the range 2-5% from 10-100% at Run 2

Uncertainty on κ_μ statistically dominated

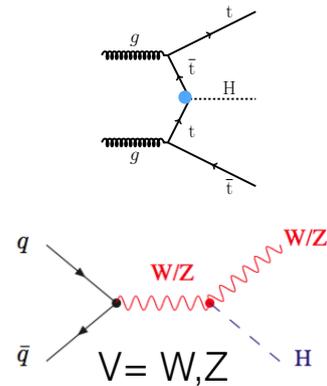
- Ratios of coupling modifiers: $\lambda_{ij} = \kappa_i/\kappa_j$, assumptions in Γ_H not needed
- Given the reference ratio: $\kappa_{gZ} = \kappa_g \cdot \kappa_Z/\kappa_H$
- Common uncertainties also cancel out in the ratio

(YR2018)



Precision improves to 2-5% from 10-100% at Run 2

- $BR(H \rightarrow b\bar{b})$ is largest $\sim 58\%$, useful to study rare Higgs production modes
- $t\bar{t}H$ provides direct probe of top-Higgs Yukawa coupling observed in Run 2
- VH production is most sensitive to study, $H \rightarrow b\bar{b}$ decay observed in Run 2

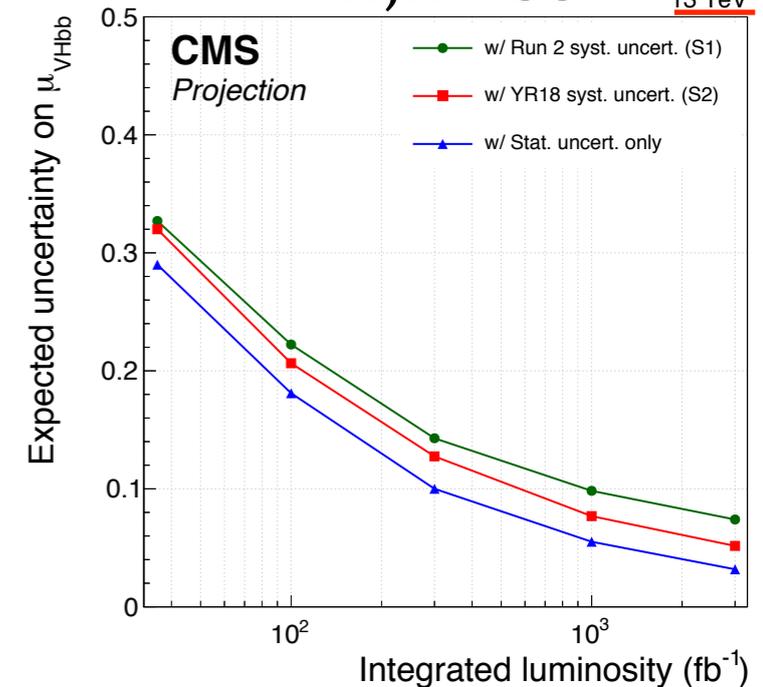
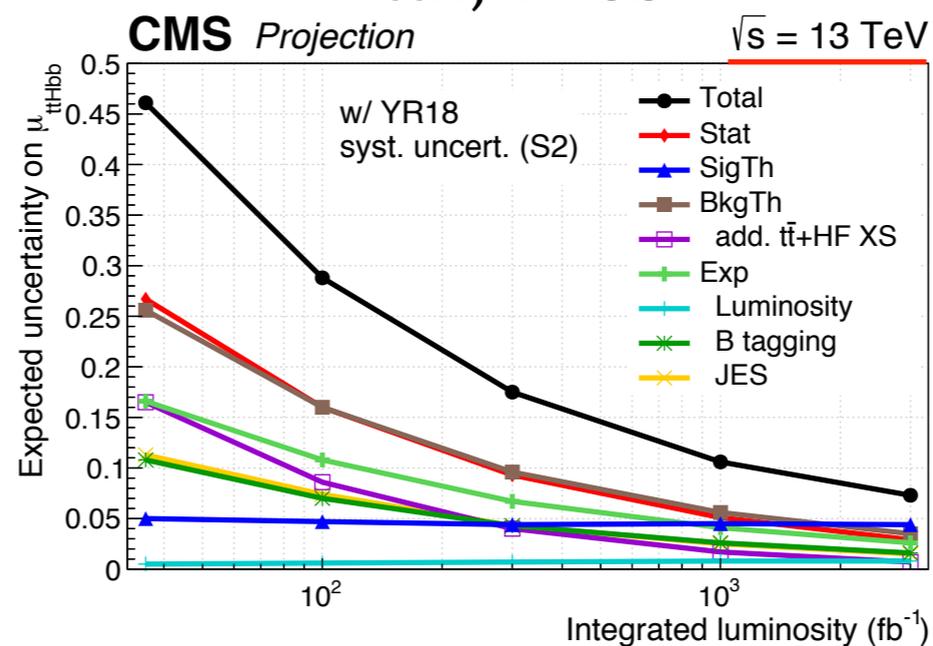
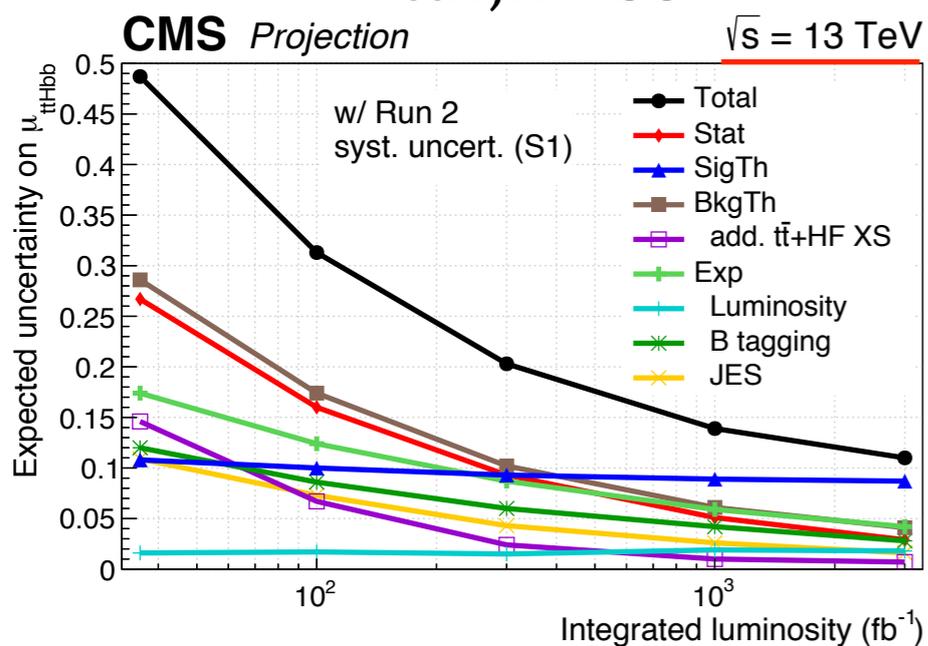


$t\bar{t}H, H \rightarrow b\bar{b}$

$t\bar{t}H, H \rightarrow b\bar{b}$

$VH, H \rightarrow b\bar{b}$

13 TeV

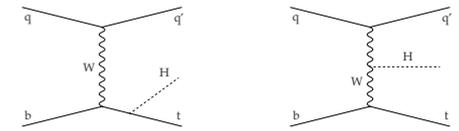


$t\bar{t}H$ analysis will gain from better background estimation from control region, e.g. $t\bar{t}+HF$ process cross section uncertainties

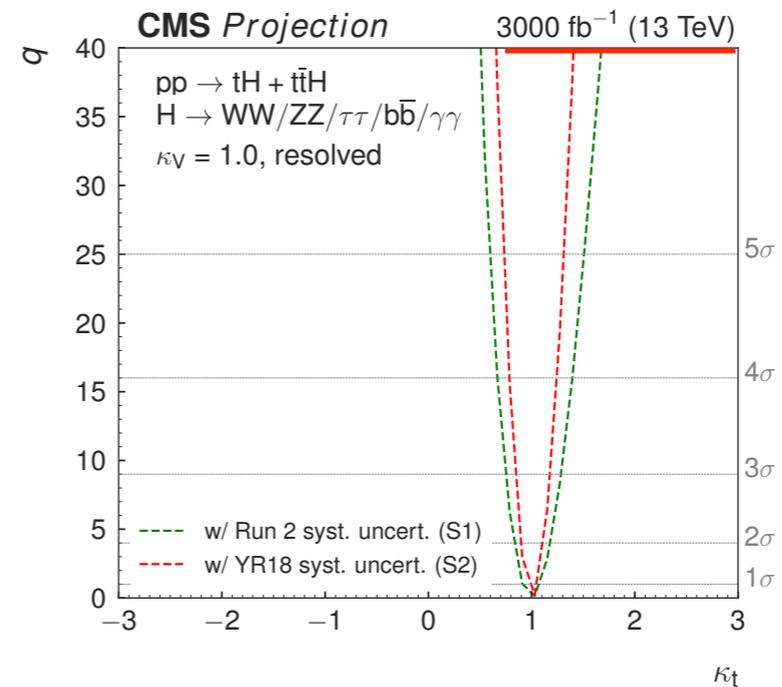
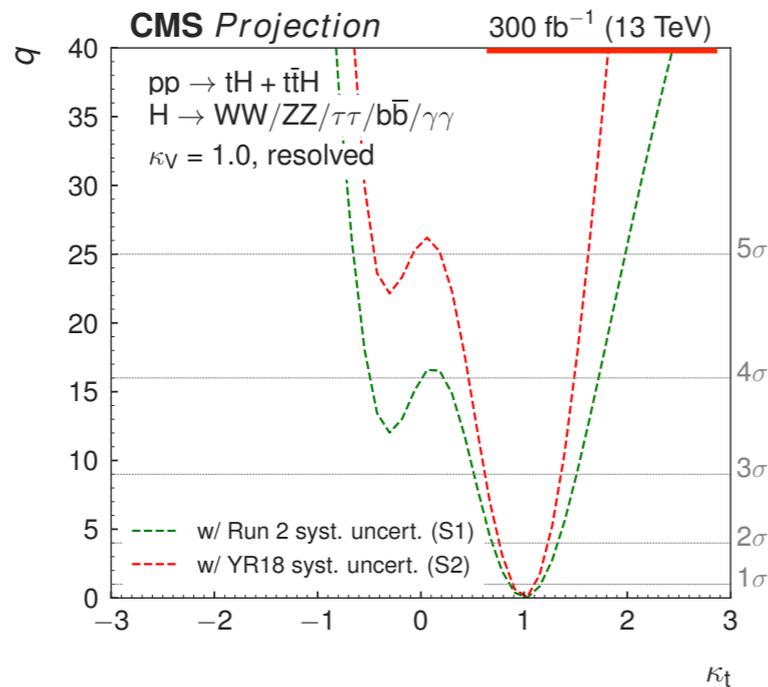
VH signal strength will improve with better knowledge of QCD scale uncertainties of signal ($ggZH$ production) and backgrounds ($V+jets$)

tH production process

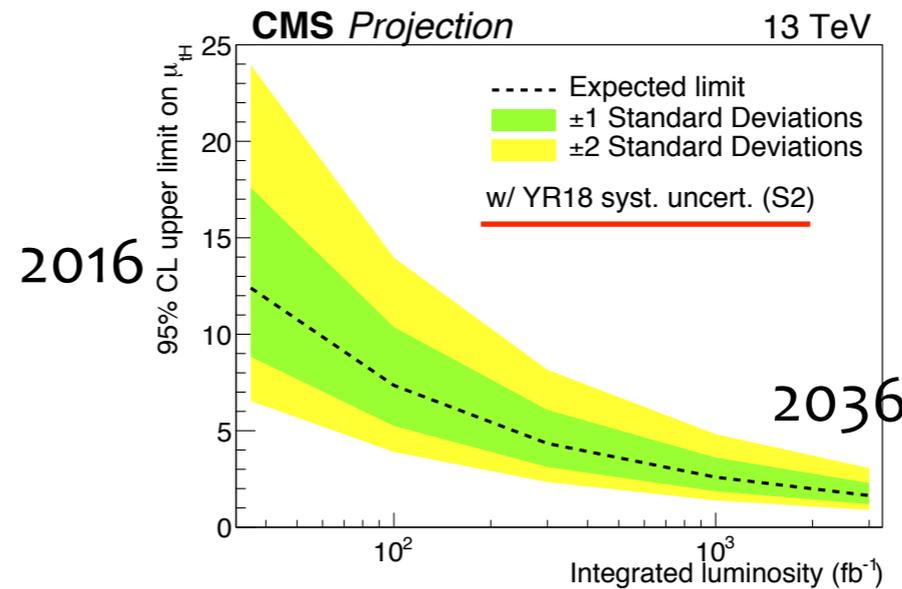
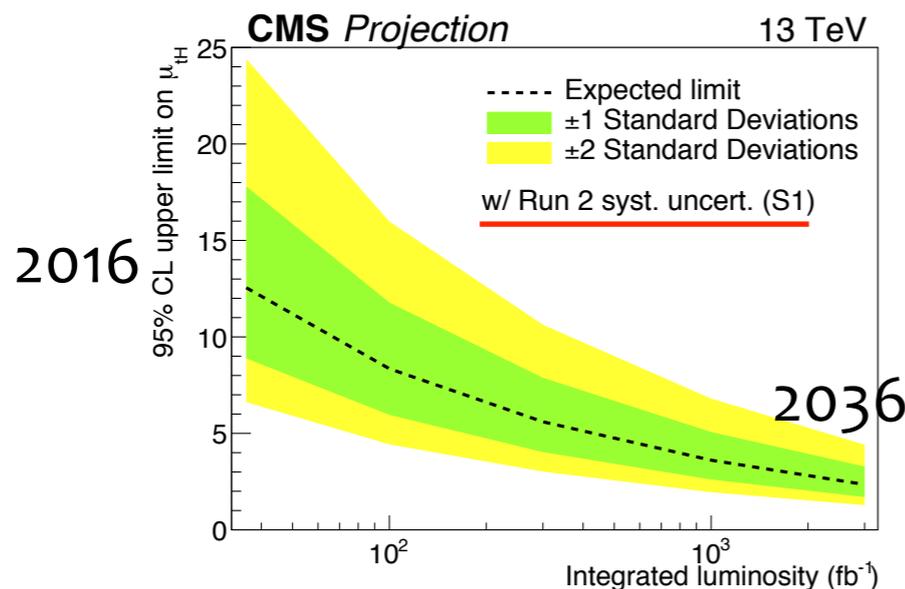
- tH (tHq+tHW) production sensitive to relative sign of κ_t and κ_V
- SM production rate very small due to interference, Run 2 sensitivity limited
- Higgs decay modes considered: $H \rightarrow bb$, multilepton, $\gamma\gamma$



$$q(\vec{\alpha}) = -2 \ln \left(\frac{L(\vec{\alpha}, \hat{\theta}_\alpha)}{L(\hat{\alpha}, \hat{\theta})} \right)$$

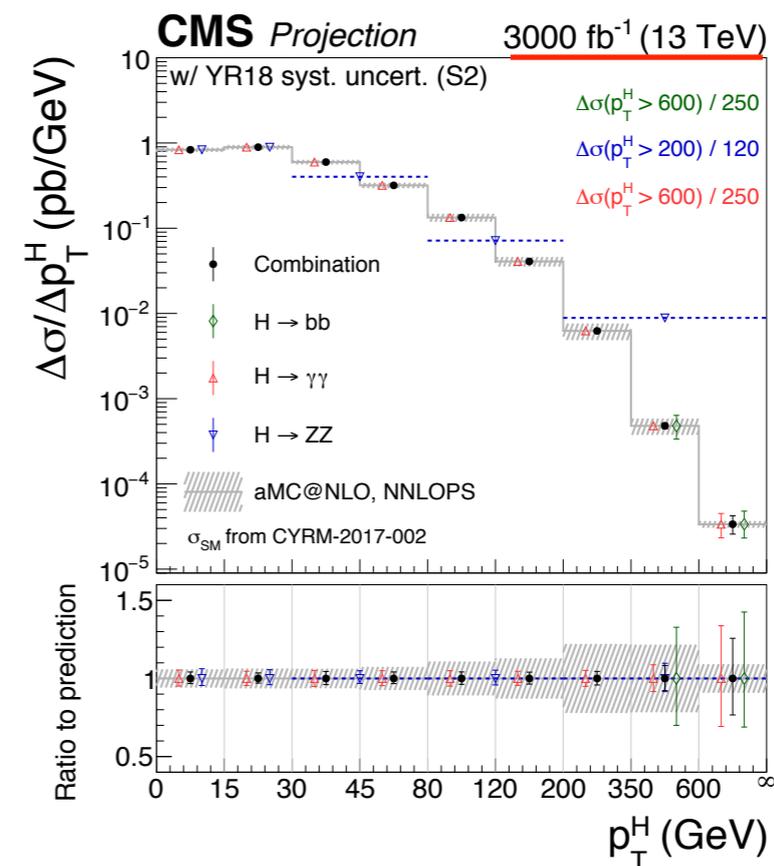
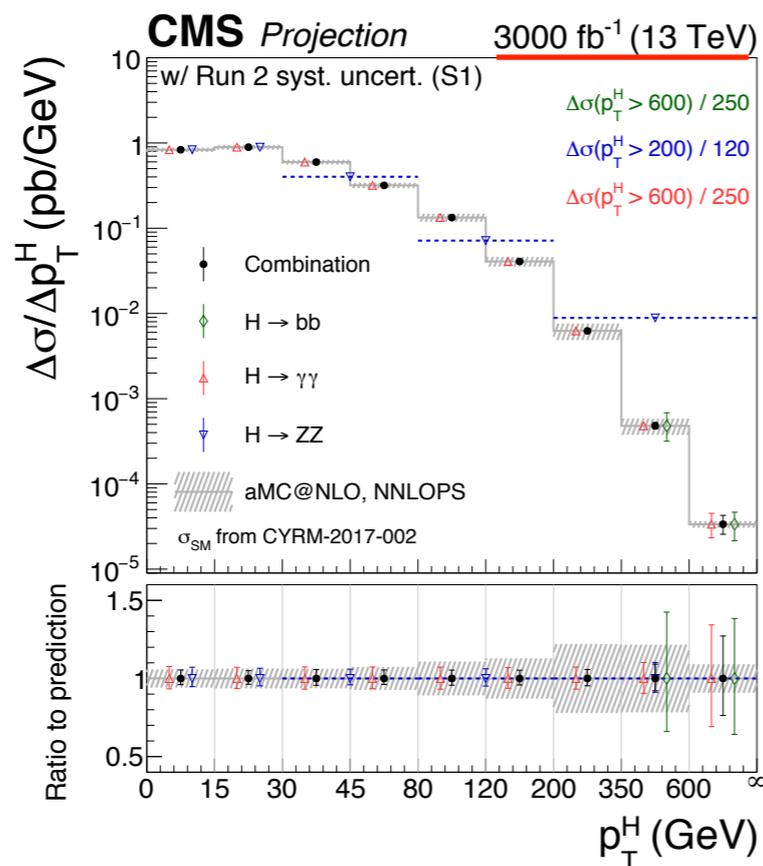


More data \rightarrow -ve κ_t values more constrained



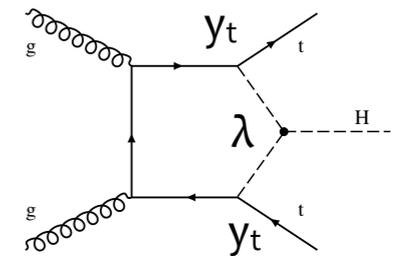
Factor of 8 improvement on μ_{tH} at HL-LHC, sensitive to SM production

- Higgs p_T is moderate in SM, but beyond SM (BSM) affects higher values
- Differential distribution only possible with large statistics
- Constrain coupling modifiers using $p_T(\text{Higgs})$ spectrum
- New physics at high energies can modify event kinematics



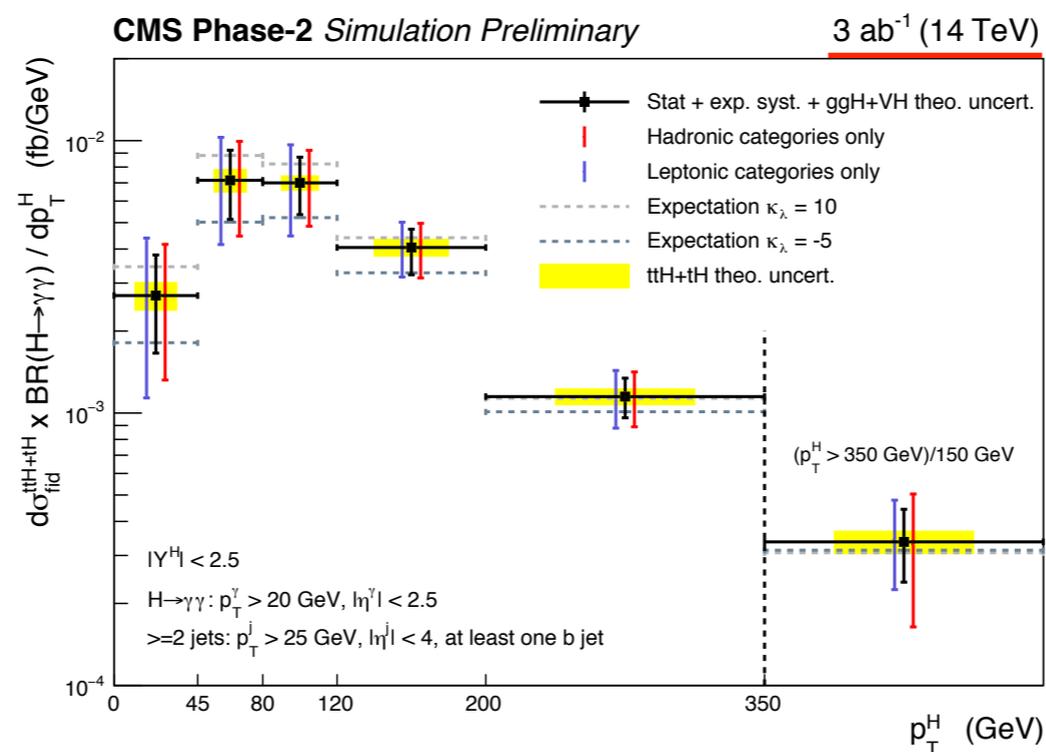
Uncertainty in high p_T bins statistically dominated → factor of 10 improvement possible with S1 alone at 3000 fb⁻¹, additional 25% at S2

- Differential cross section measurement using $p_T(\text{Higgs})$ can disentangle self-coupling (λ) effects from anomalous top-Higgs coupling
- Delphes simulated samples for Phase II geometry used



$$\lambda_{\text{SM}} = 0.13 \text{ for } m_H = 125 \text{ GeV}$$

- $t\bar{t}H$ and tH events with $H \rightarrow \gamma\gamma$ decay selected in hadronic or leptonic final states depending on top decay



Analysis results on κ_λ in Sylvie's talk

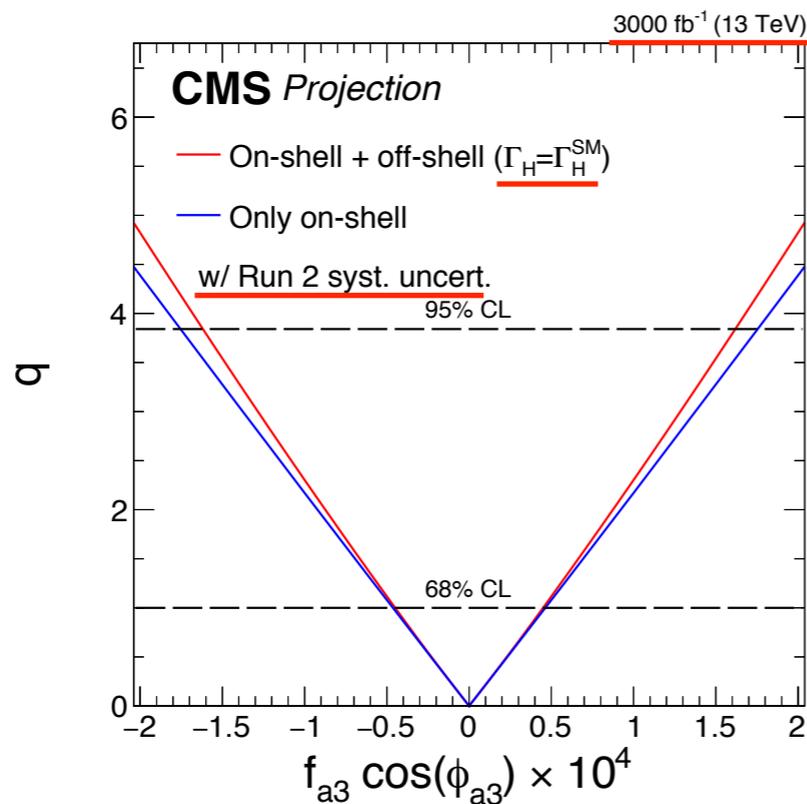
Uncertainties on differential cross section 20-40% at 3000 fb^{-1}

- Physical Higgs width (~ 4 MeV) much smaller than experimental resolution (~ 100 MeV)
- Width can be constrained comparing $H \rightarrow ZZ \rightarrow 4\ell$ on-shell and off-shell contributions
- Off-shell contributions can be enhanced due to:
 - ➔ larger width
 - ➔ anomalous HVV coupling
 - ➔ new resonances

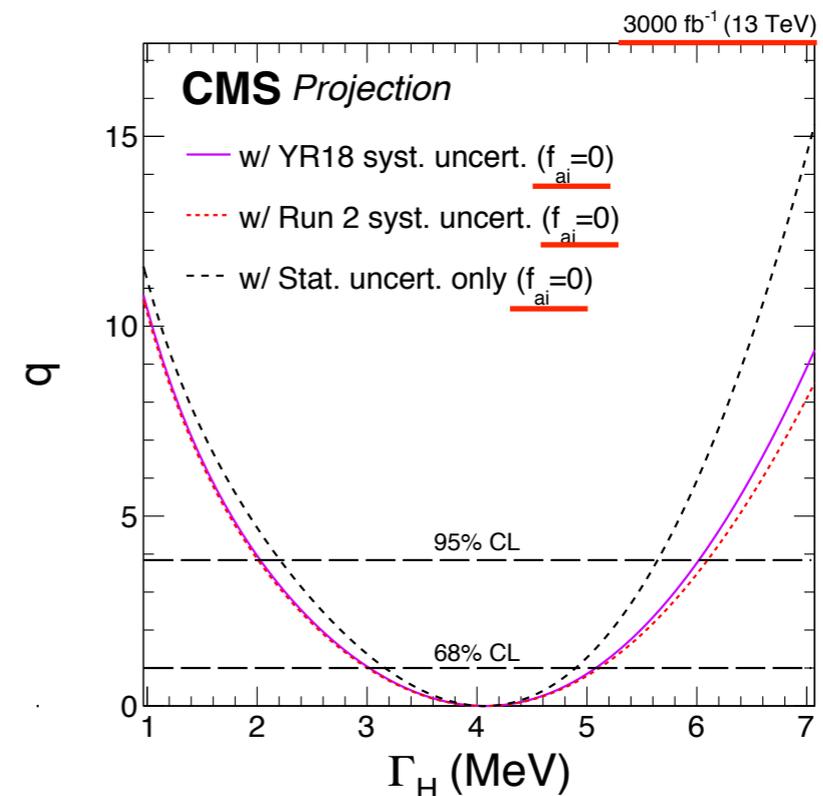
$$\sigma_{H \rightarrow VV \rightarrow 4\ell}^{\text{on-shell}} \propto \mu_{VVH}$$

$$\sigma_{H \rightarrow VV \rightarrow 4\ell}^{\text{off-shell}} \propto \mu_{VVH} \cdot \Gamma_H$$

tensor structure of HVV coupling can be probed using anomalous coupling a_3 , corresponding phase φ_3

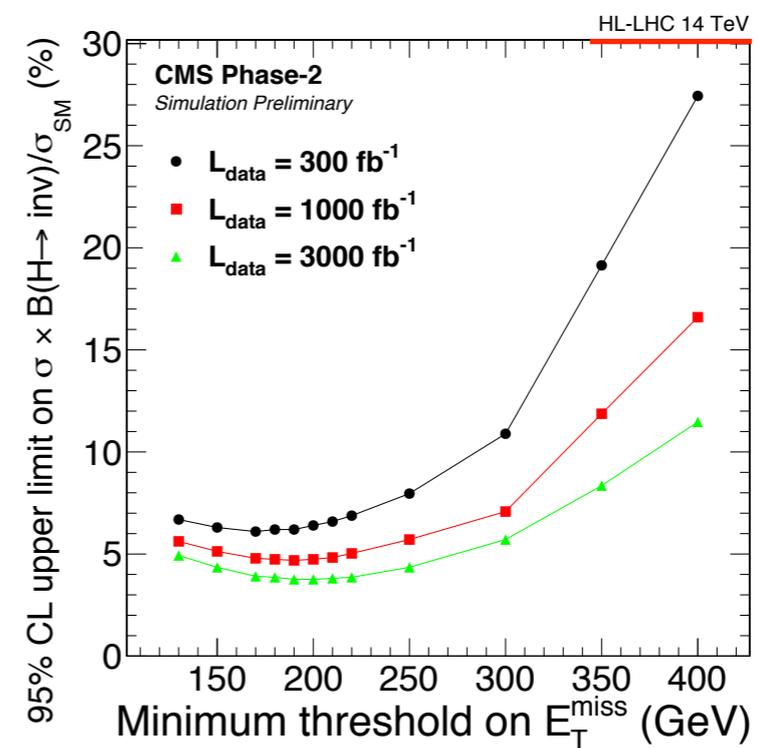
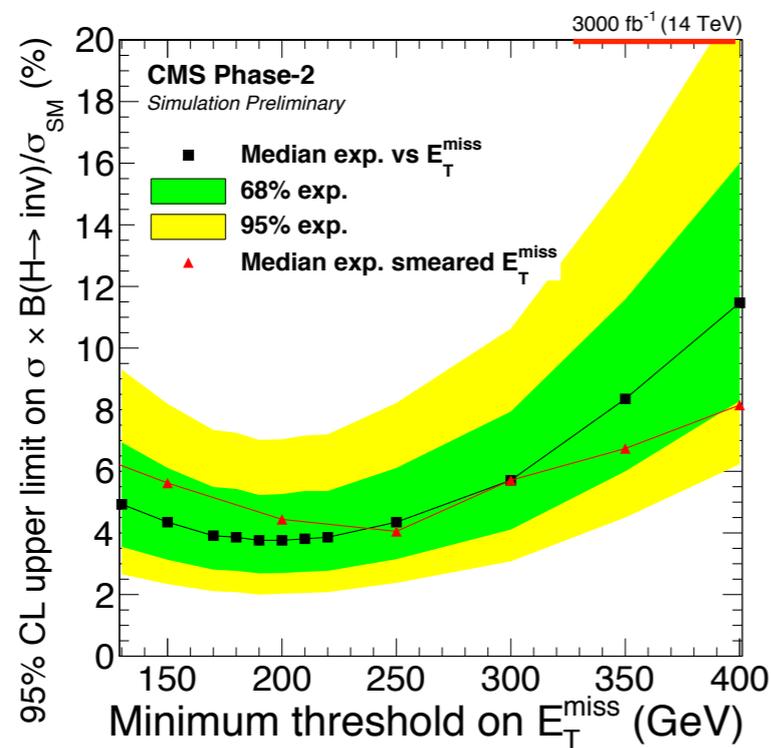


95% CL exclusion outside $(-0.00016, 0.00016)$
 at S1, Run 2 measurement $(-0.0076, 0.0050)$



95% CL exclusion outside $(2, 6)$ at S2,
Run 2 measurement $(0.08, 9.16)$

- Higgs decays to massive invisible particles is sensitive probe of DM coupling
- Limited sensitivity to Higgs to invisible decays with Run 2 data
- VBF topology has the best sensitivity
- **95% CL upper limit on $H \rightarrow \text{invisible}$ BR is 3.8%** assuming SM production at 3000 fb^{-1} , SM prediction of higgs BR to invisible $\sim 10^{-3}$
- Run 2 measurement in VBF channel yields an upper limit of 33%



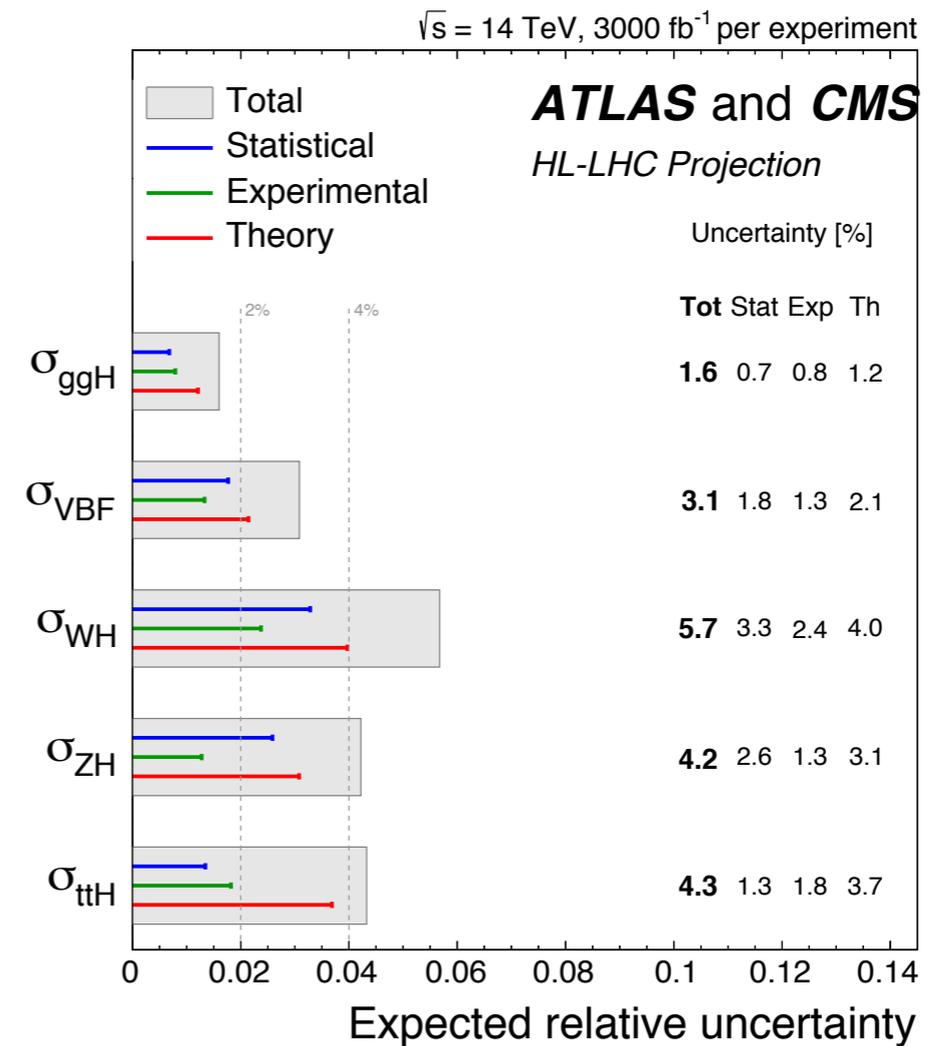
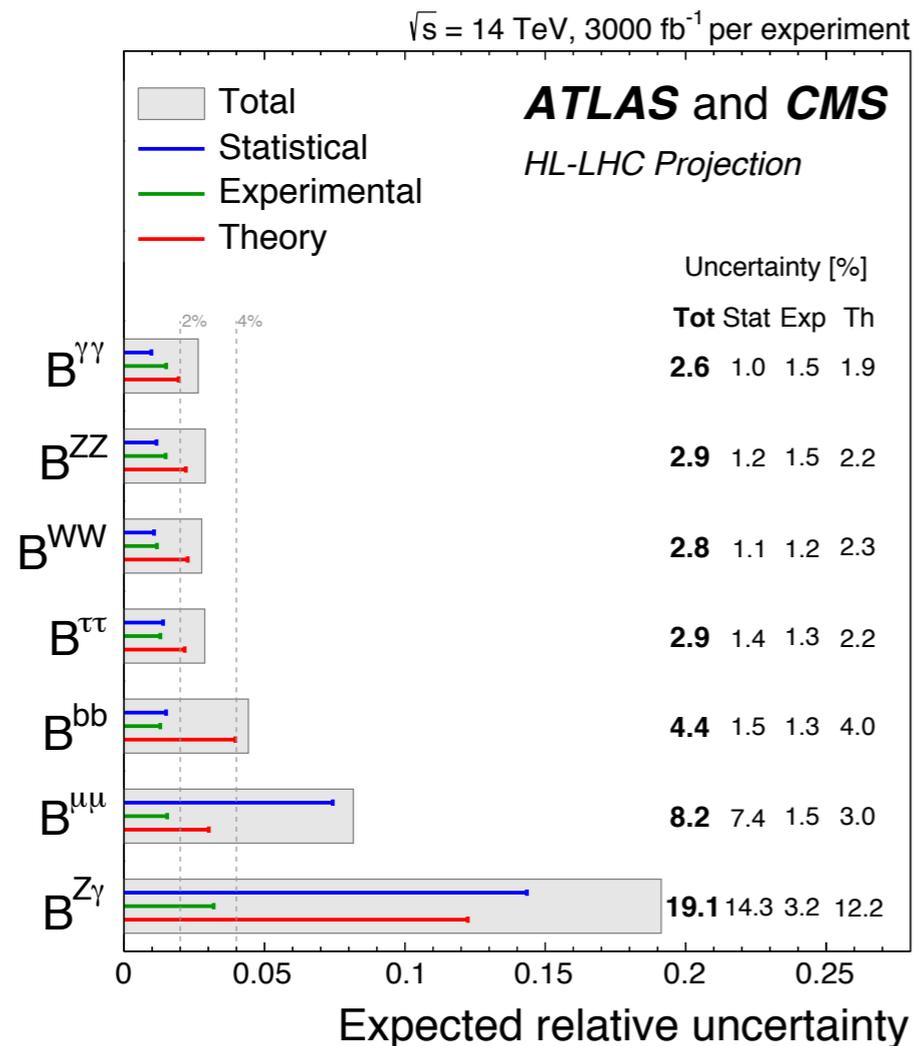
HL-LHC will put more stringent constraint on new physics in Higgs \rightarrow invisible sector

- Run 2 saw plethora of Higgs measurements → reaching precision frontier, some still statistics dominated
- Analyses are well established, fine-grained measurements will be possible at HL-LHC
- More data will provide access to rare production and decay modes, e.g., $H \rightarrow \mu\mu$, $Z\gamma$
- To reach the goal, need improvement in theoretical and experimental uncertainties
- Roadmap of LHC essentially includes building new detector to tackle pileup, radiation damage

Thank You.

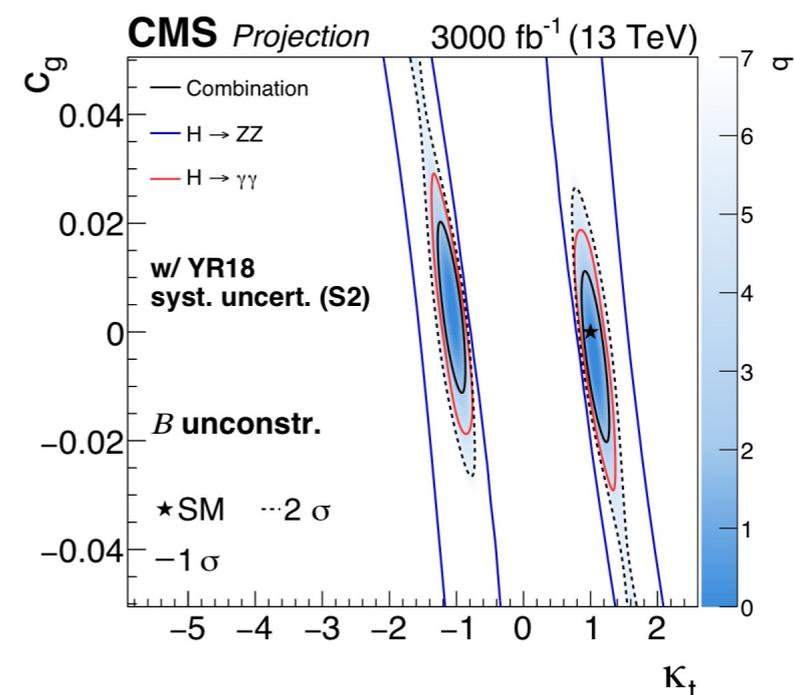
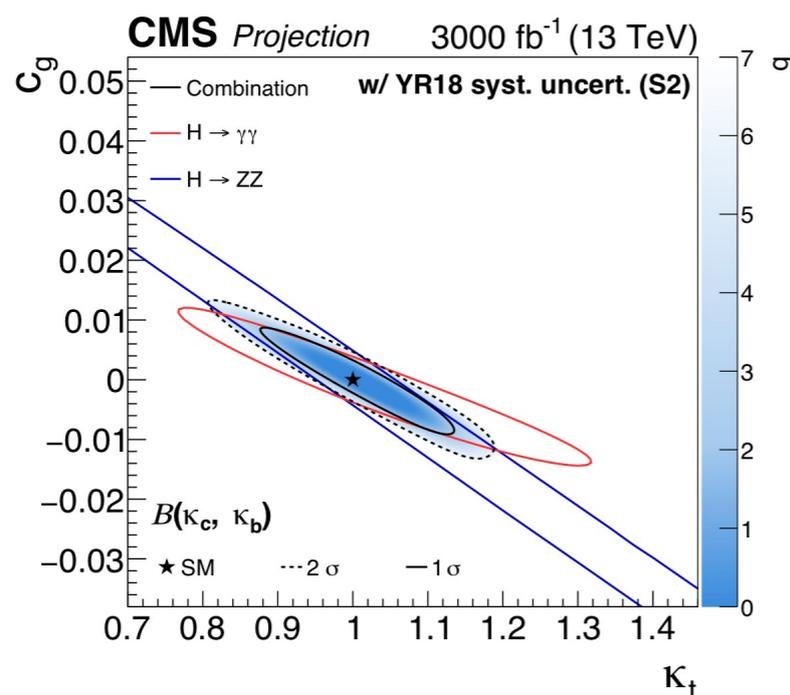
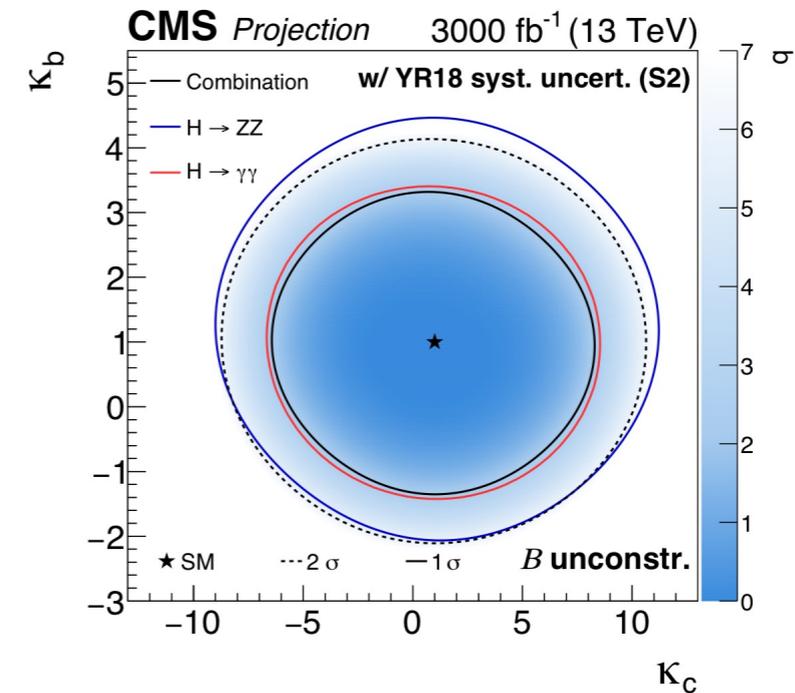
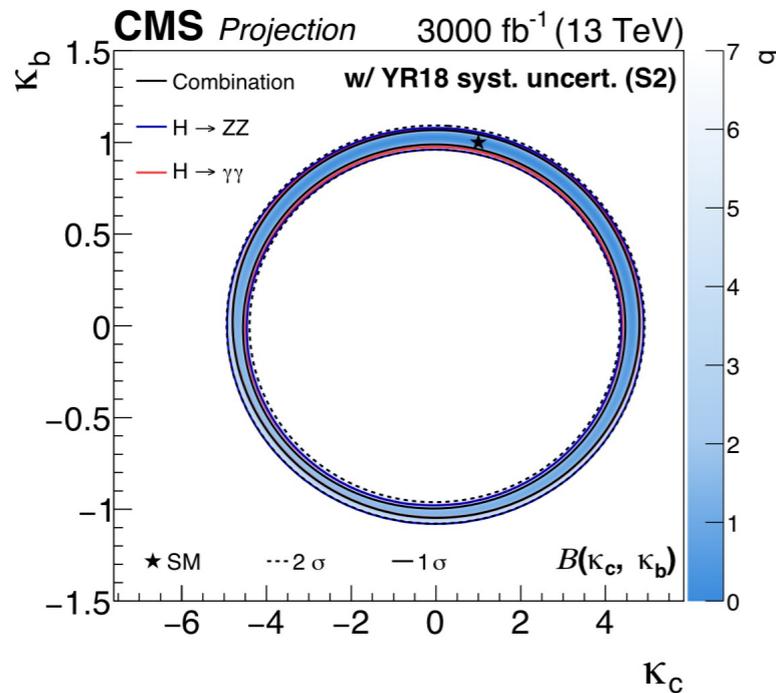
Table 1: The sources of systematic uncertainty for which limiting values are applied in scenario S2.

Source	Component	Run 2 uncertainty	Projection minimum uncertainty
b-Tagging	b-/c-jets (syst.)	Varies with p_T and η	Same as Run 2
	light mis-tag (syst.)	Varies with p_T and η	Same as Run 2
	b-/c-jets (stat.)	Varies with p_T and η	No limit
	light mis-tag (stat.)	Varies with p_T and η	No limit
Muon ID		1–2%	0.5%
Electron ID		1–2%	0.5%
Photon ID		0.5–2%	0.25–1%
Hadronic tau ID		6%	Same as Run 2
Jet energy scale	Absolute	0.5%	0.1–0.2%
	Relative	0.1–3%	0.1–0.5%
	Pileup	0–2%	Same as Run 2
	Method and sample	0.5–5%	No limit
	Jet flavour	1.5%	0.75%
	Time stability	0.2%	No limit
Jet energy res.		Varies with p_T and η	Half of Run 2
MET scale		Varies with analysis selection	Half of Run 2
Integrated lumi.		2.5%	1%

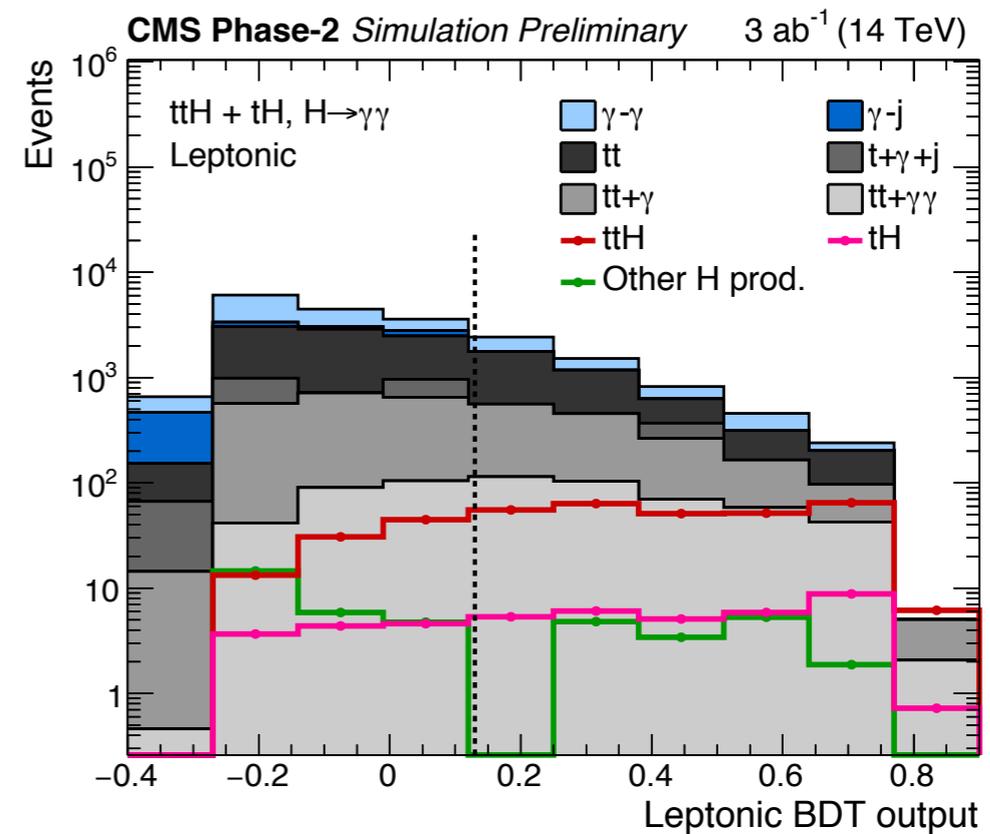
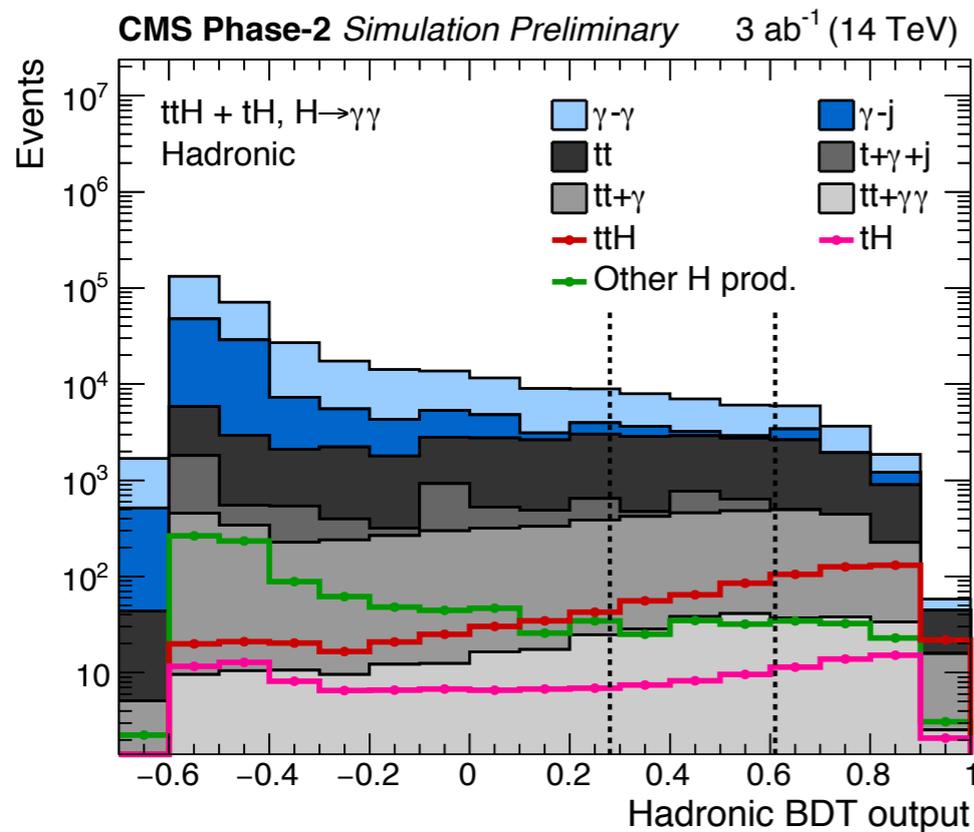


- Combined total expected $\pm 1\sigma$ uncertainties in S2 on per-decay-mode BR and per-production-mode cross section normalized to SM prediction

- Theoretical prediction for the p_T (Higgs) fit to data can be used to constrain Higgs couplings
- Use parametrisation dependent on κ_b , κ_t and κ_b , c_g



- In the $ttH+tH, H\rightarrow\gamma\gamma$ analysis, BDT classifiers are trained with variables such as event kinematics and photon quality
- Variables that may distort $p_T(\text{Higgs})$ spectrum are avoided, such as di-photon rapidity



- Difference in Delphes and FullSim distributions adjusted accordingly
- Final state contains two jets with high rapidity gap, high dijet mass (M_{jj}), small angle in transverse plane, missing transverse momentum (MET)
- Reject events with identified electrons, muons, photons, taus, b jets
- Four control regions dominated with $W \rightarrow e\nu$, $W \rightarrow \mu\nu$, $Z \rightarrow ee$, $Z \rightarrow \mu\mu$ events

