

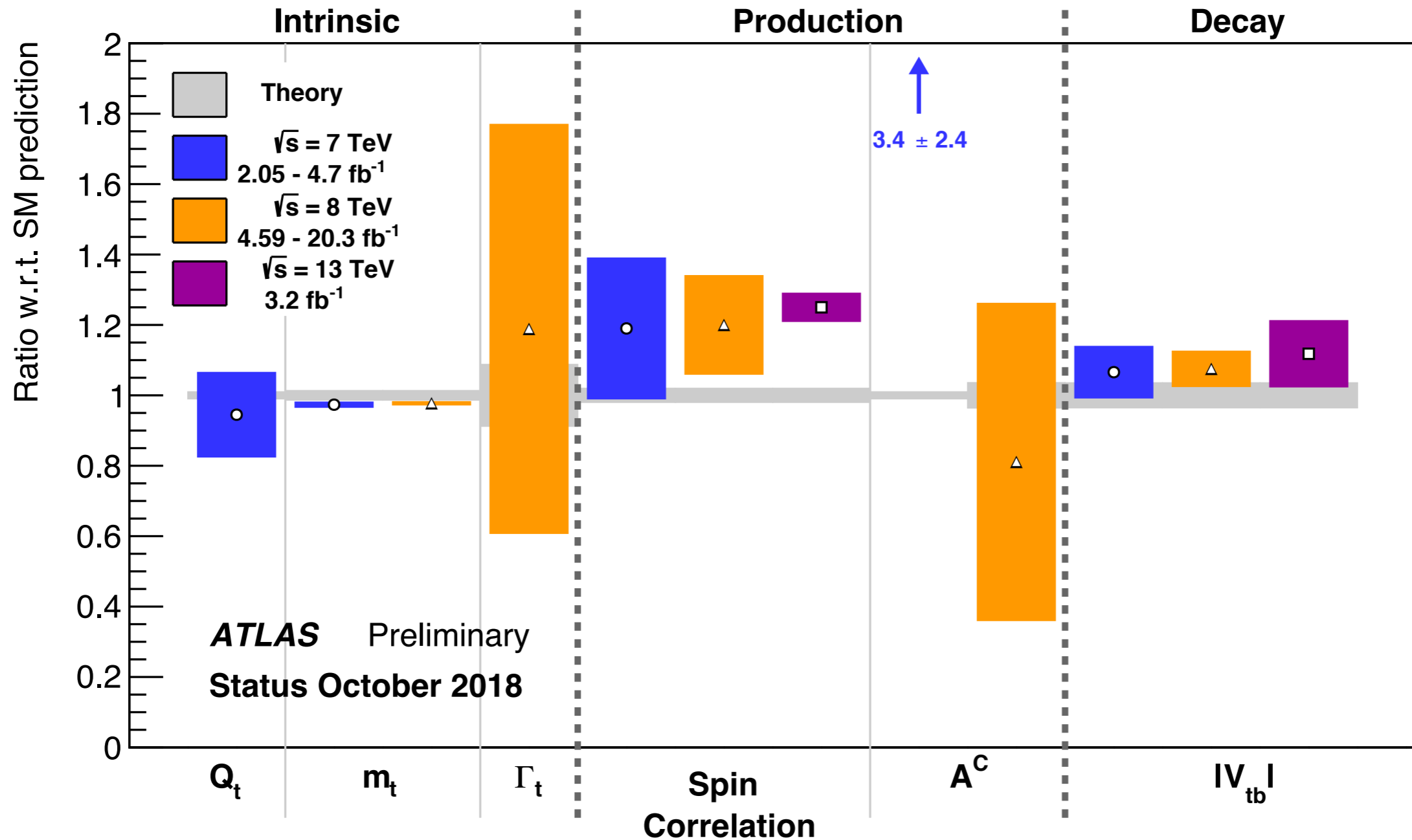
Top quark pair property measurements using the ATLAS detector at the LHC

EPS 2019, Ghent

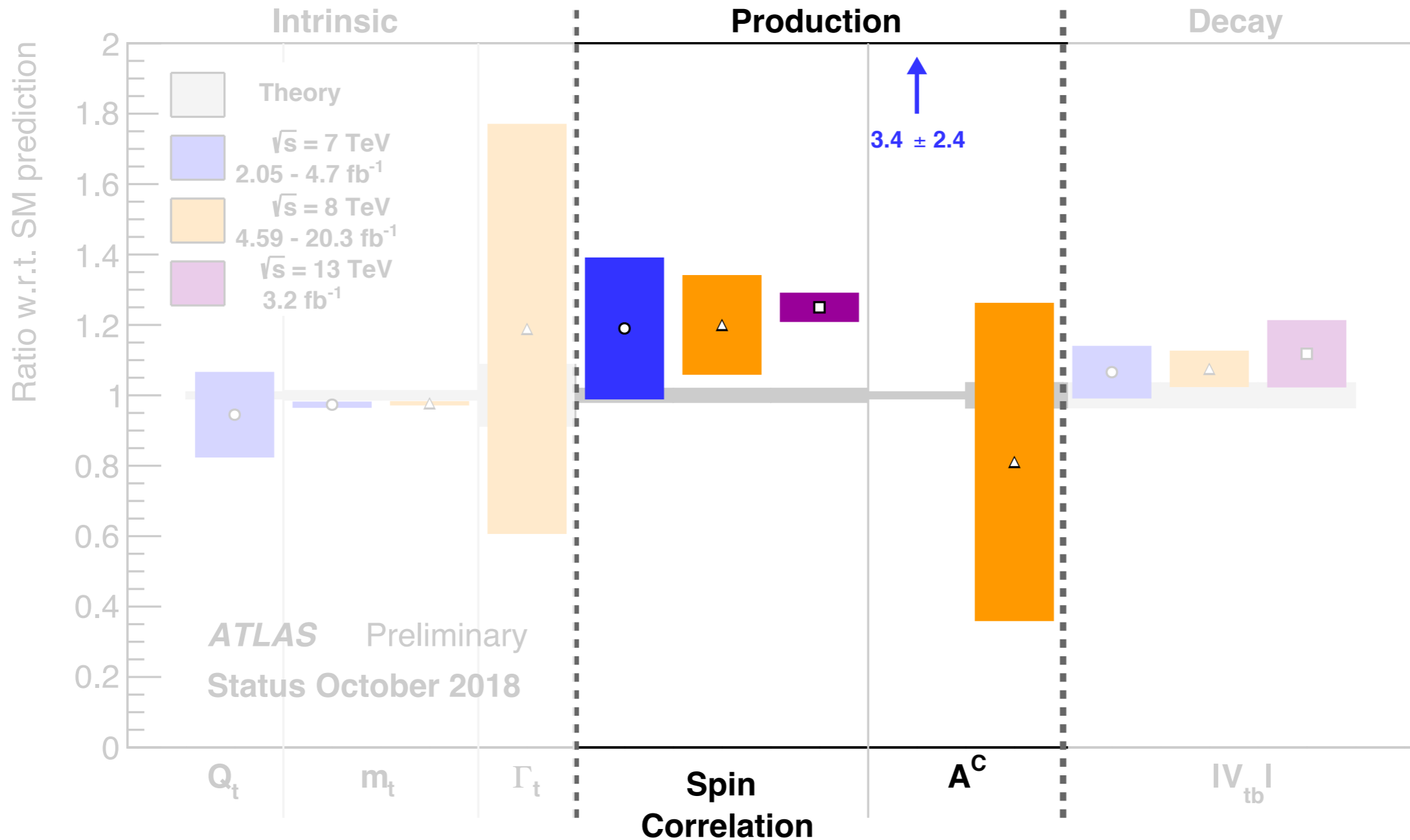
Jay Howarth, on behalf of the ATLAS collaboration

Production time	Decay time	Hadronisation time	Spin-Decor. time
$\frac{1}{m_t}$	$\frac{1}{\Gamma_t}$	$\frac{1}{\Lambda}$	$\frac{m_t}{\Lambda}$
$\sim 10^{-27} \text{ s}$	$\sim 10^{-25} \text{ s}$	$\sim 10^{-23} \text{ s}$	$\sim 10^{-22} \text{ s}$

- Top quark is the heaviest particle in the SM, which leads to unique features:
 - ➔ Decay timescale is orders of magnitude shorter than the hadronisation or spin de-correlation timescale → **top acts like bare quark.**
 - ➔ Top yukawa term ~ 1 → **possibly plays a special role in EWSB?**
 - ➔ Very clear signal with little background and high production cross-section at the LHC → **possible to make precision measurements.**



- Large suite of measurements from ATLAS probing the top quark's properties

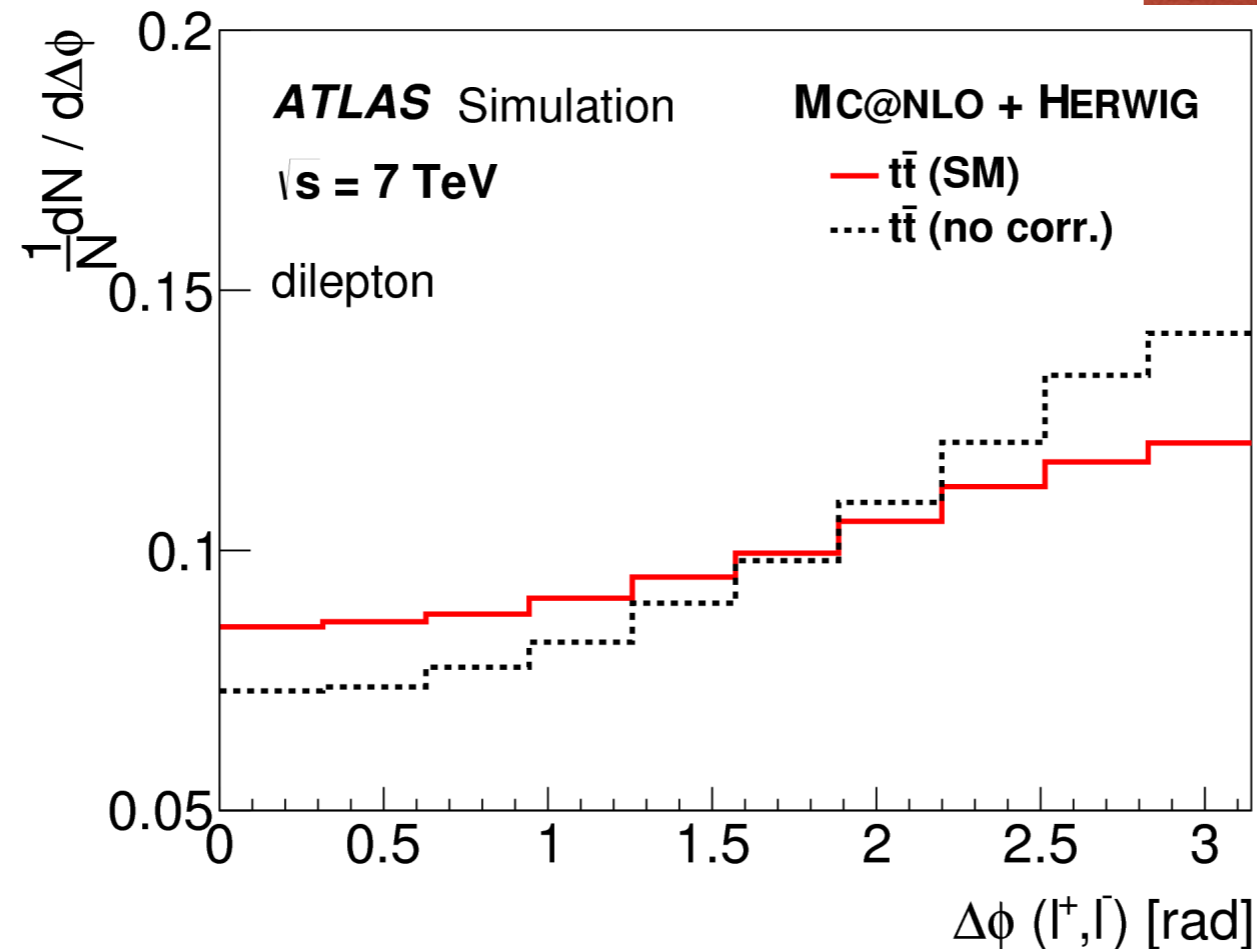


- Today, I'll focus on production properties of $t\bar{t}$ pairs.

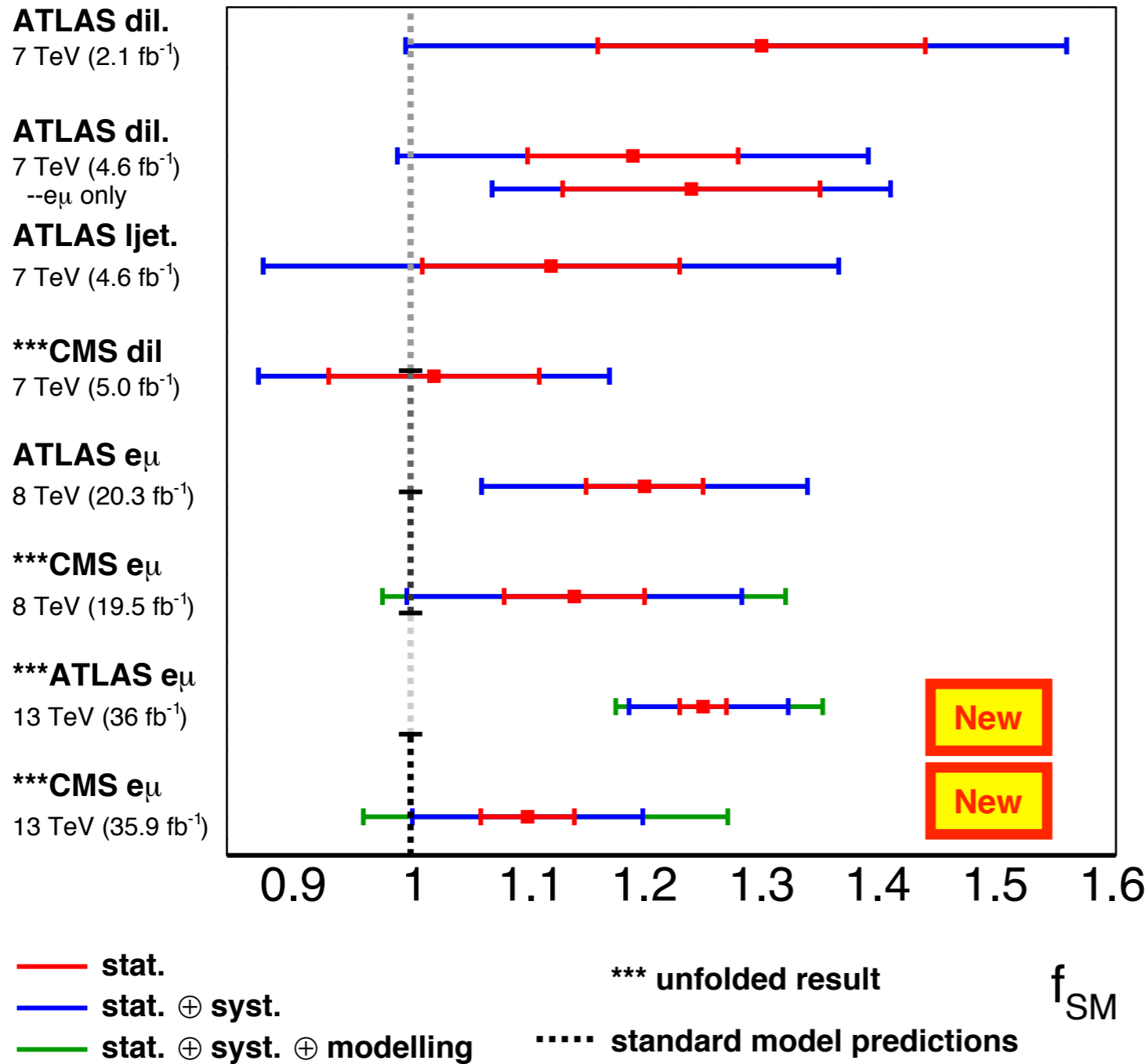
$$C = A\alpha_1\alpha_2 = \frac{N(\uparrow\uparrow) + N(\downarrow\downarrow) - N(\uparrow\downarrow) - N(\downarrow\uparrow)}{N(\uparrow\uparrow) + N(\downarrow\downarrow) + N(\uparrow\downarrow) + N(\downarrow\uparrow)}$$

- Tops produced via QCD are not intrinsically polarised, but spins between top pairs are correlated.
- Tops decay before they can hadronise:
 - ➔ Spin information is transferred directly to decay particles.
 - ➔ Leptons carry the full spin information ($\alpha_\ell \sim 1$)
 - ➔ Dilepton $t\bar{t}$ decays are the best choice for accessing spin information.
- Can be measured directly using complex observables involving the angles between the tops and their decay products:
 - ➔ Incurs large top reconstruction uncertainties in dilepton channel due to ν 's.

Phys. Rev. D90 (2014) 112016

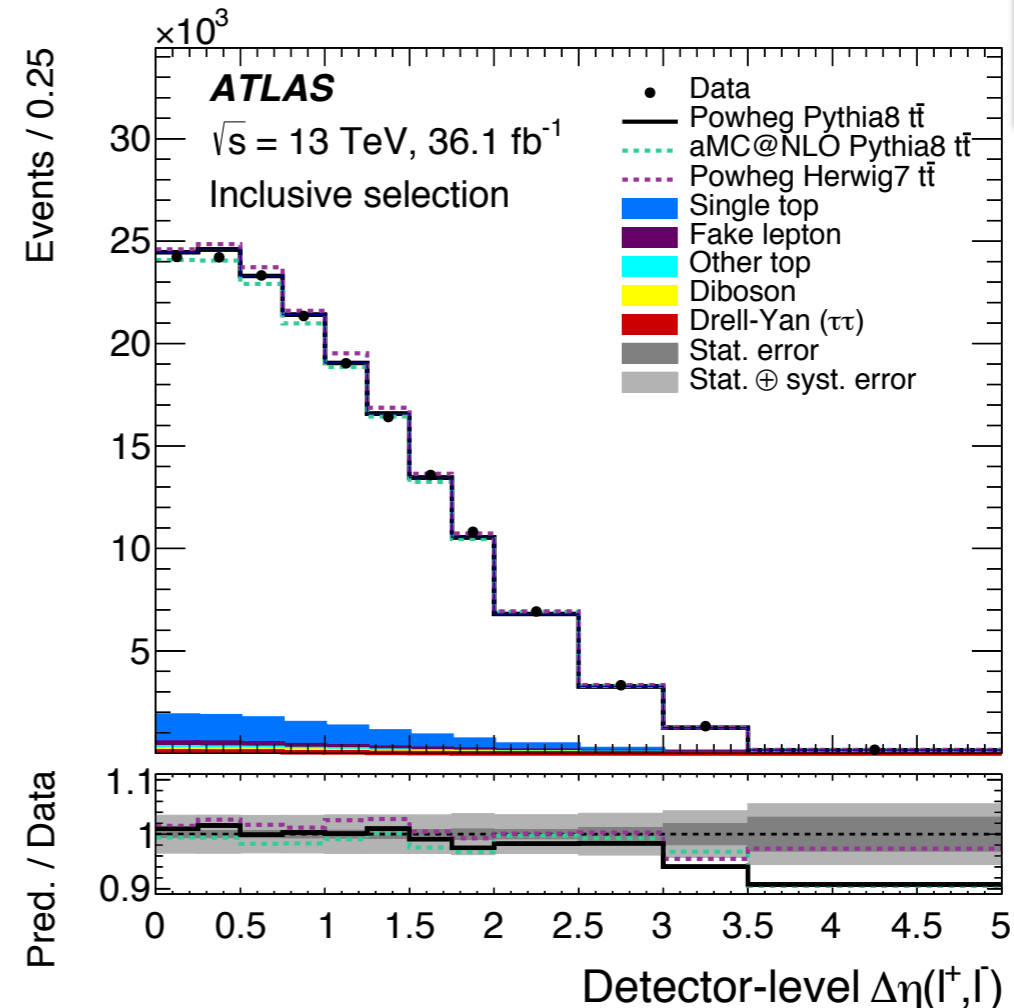
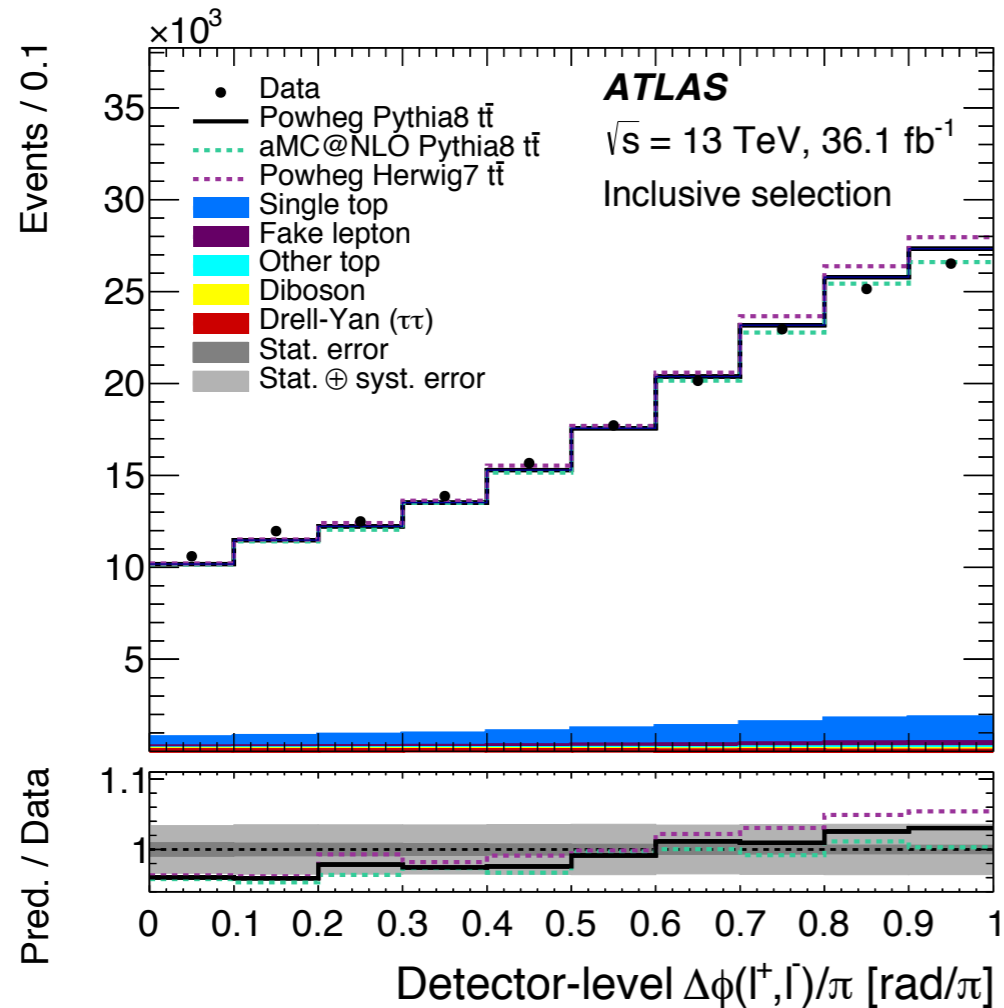


- Fortunately, there is a lab-frame observable that does not require any top reconstruction that is sensitive to spin correlation:
 - ➔ The difference in the azimuthal angle between the leptons from the top decay
 - ➔ Usually, spin is extracted using a template fit using a SM spin and NoSpin hypothesis.

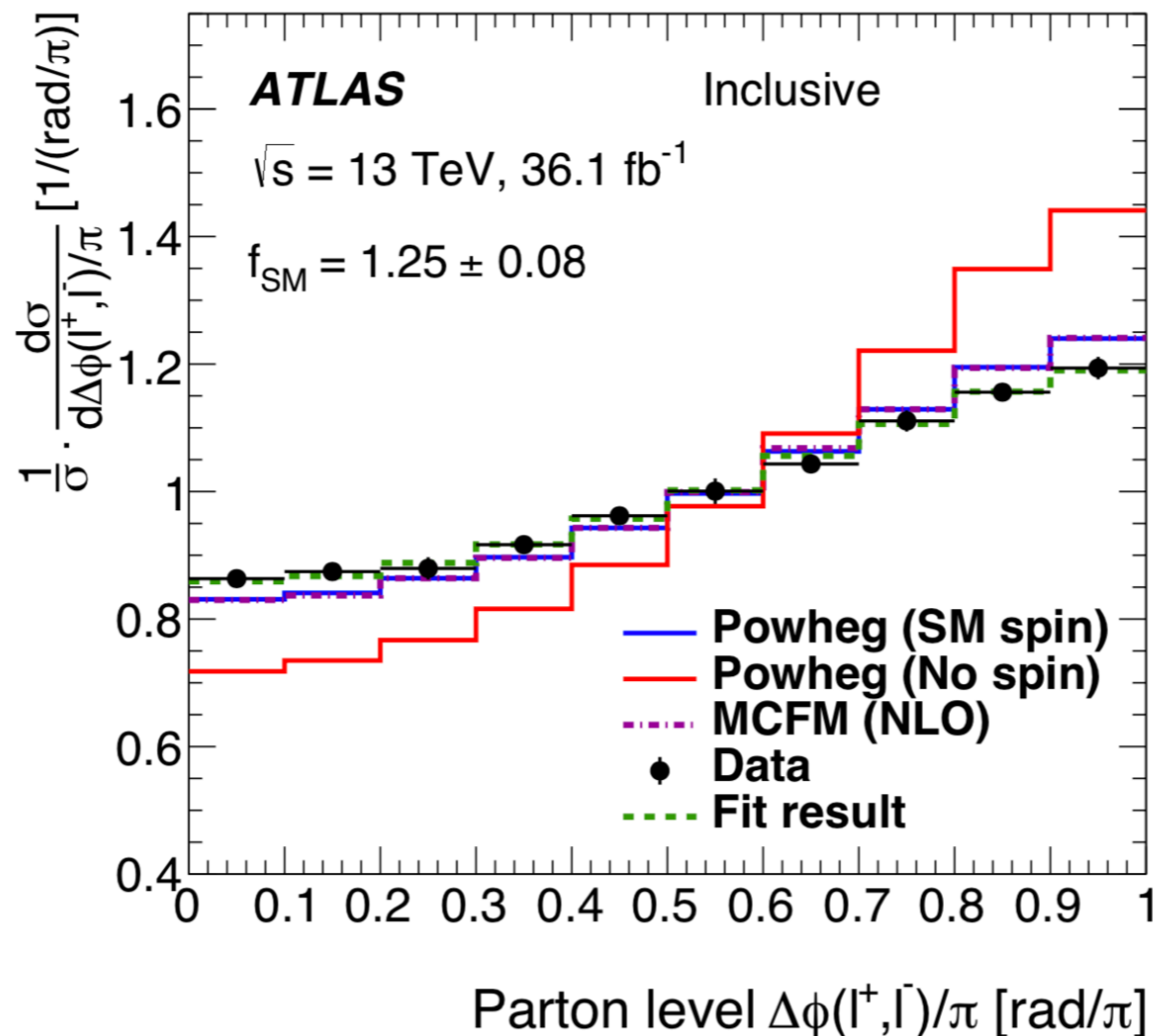


- This property has been measured many times by ATLAS and CMS, at each collision energy.
- f_{SM} is “fraction of SM-like spin correlation”:
 - ➔ $f_{SM} = 1$ is SM-like
 - ➔ $f_{SM} = 0$ is uncorrelated
- Both ATLAS and CMS consistently measure stronger than SM spin correlations using the $|\Delta\phi_{\ell\ell}|$ observable.

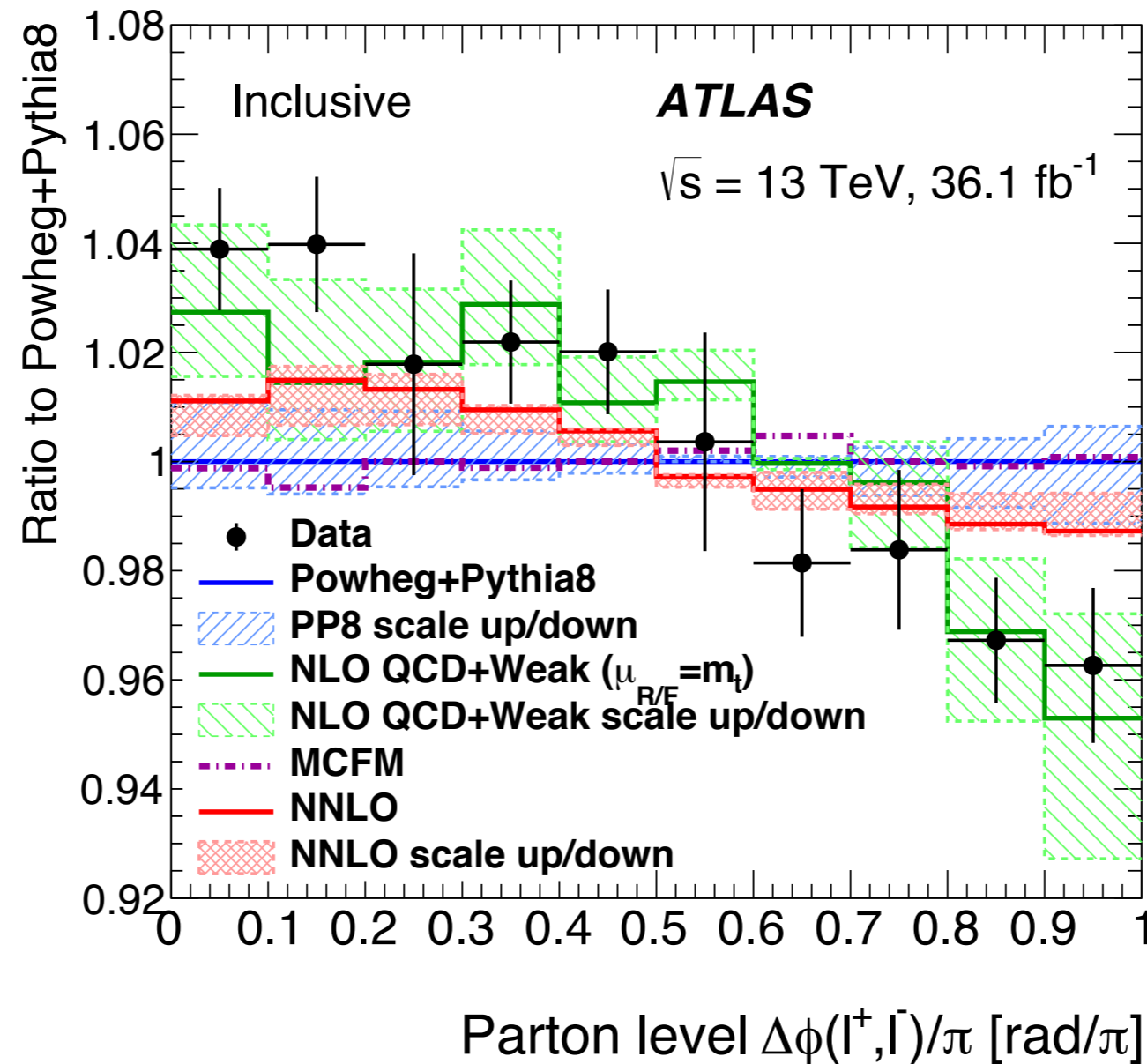
Submitted
to EPJC



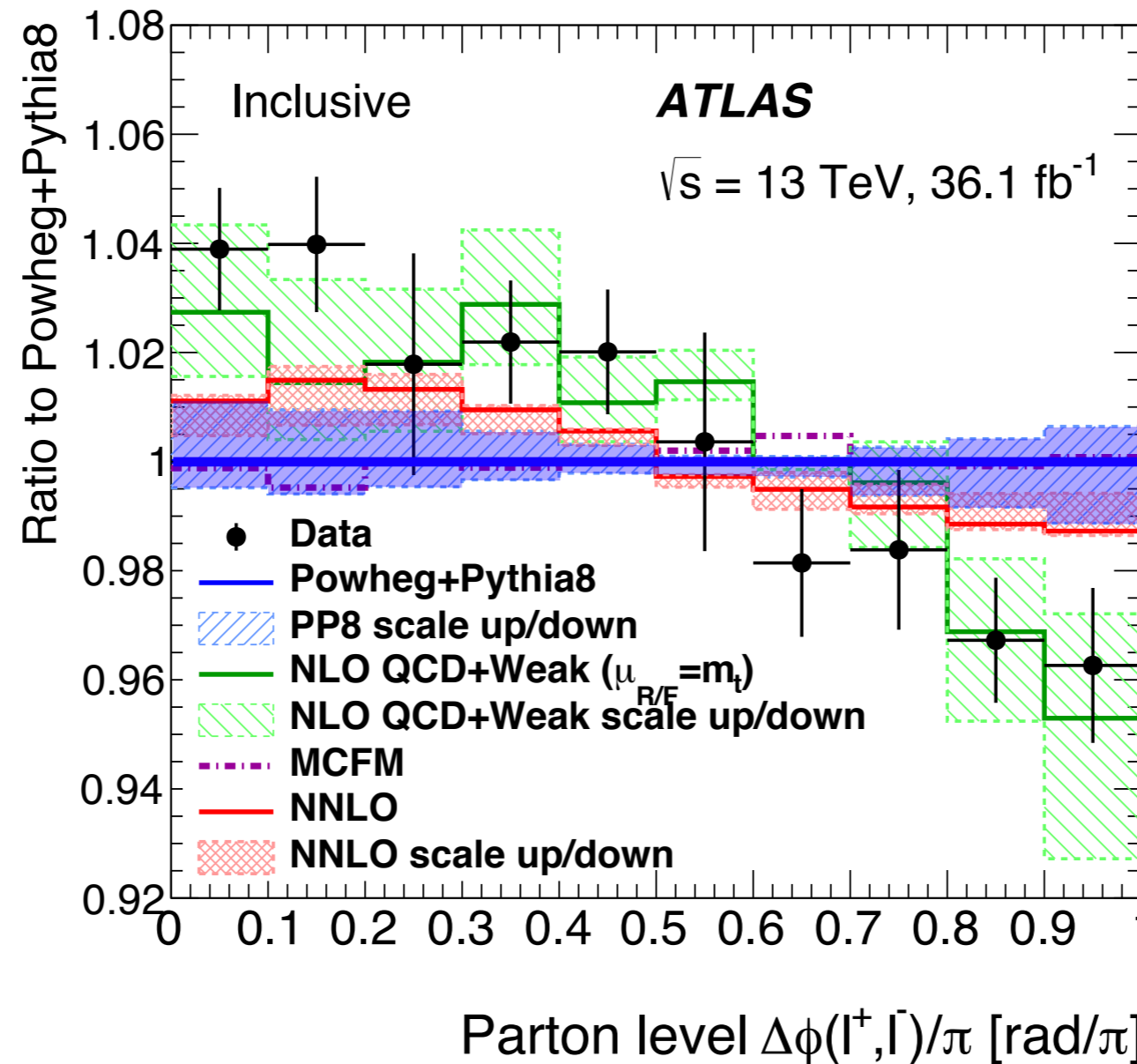
- Measured the $|\Delta\phi_{\ell\ell}|$ using 36 fb^{-1} of 13 TeV Run2 data and a $e\mu + 2b$ selection.
➡ Also measured differentially vs. $m(t\bar{t})$
- Also measured the $|\Delta\eta_{\ell\ell}|$ observable, which is sensitive to SUSY production.
- All results unfolded to **fiducial particle level** and **full phase-space parton level** using Bayesian Iterative Unfolding.



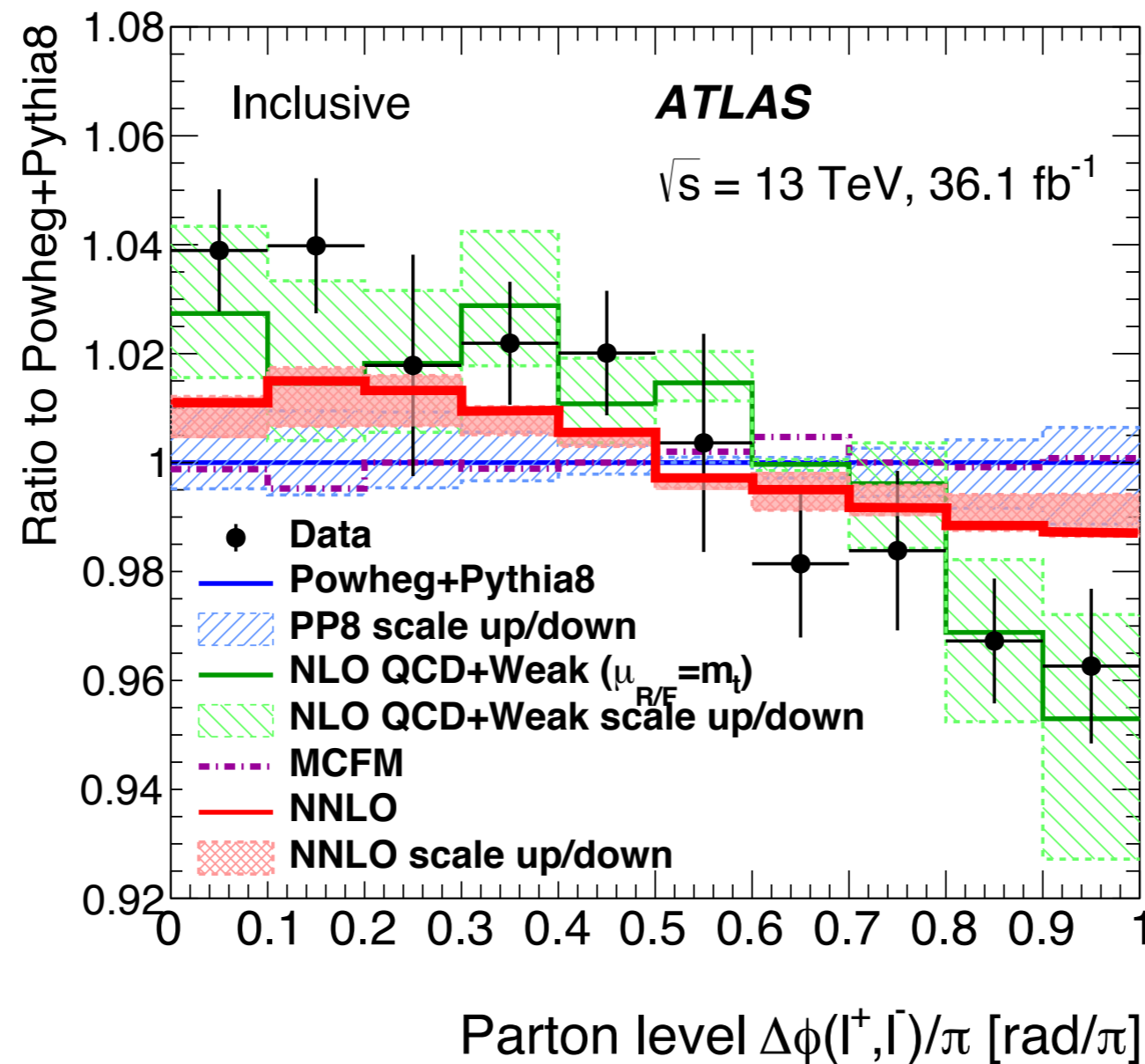
- As with previous results, the $|\Delta\phi_{\ell\ell}|$ shows a stronger slope than the data, and we measure $f_{\text{SM}} = 1.25 \pm 0.08$, relative to NLO predictions.
➔ This is more than 3σ discrepant from what we expect from the SM at NLO.



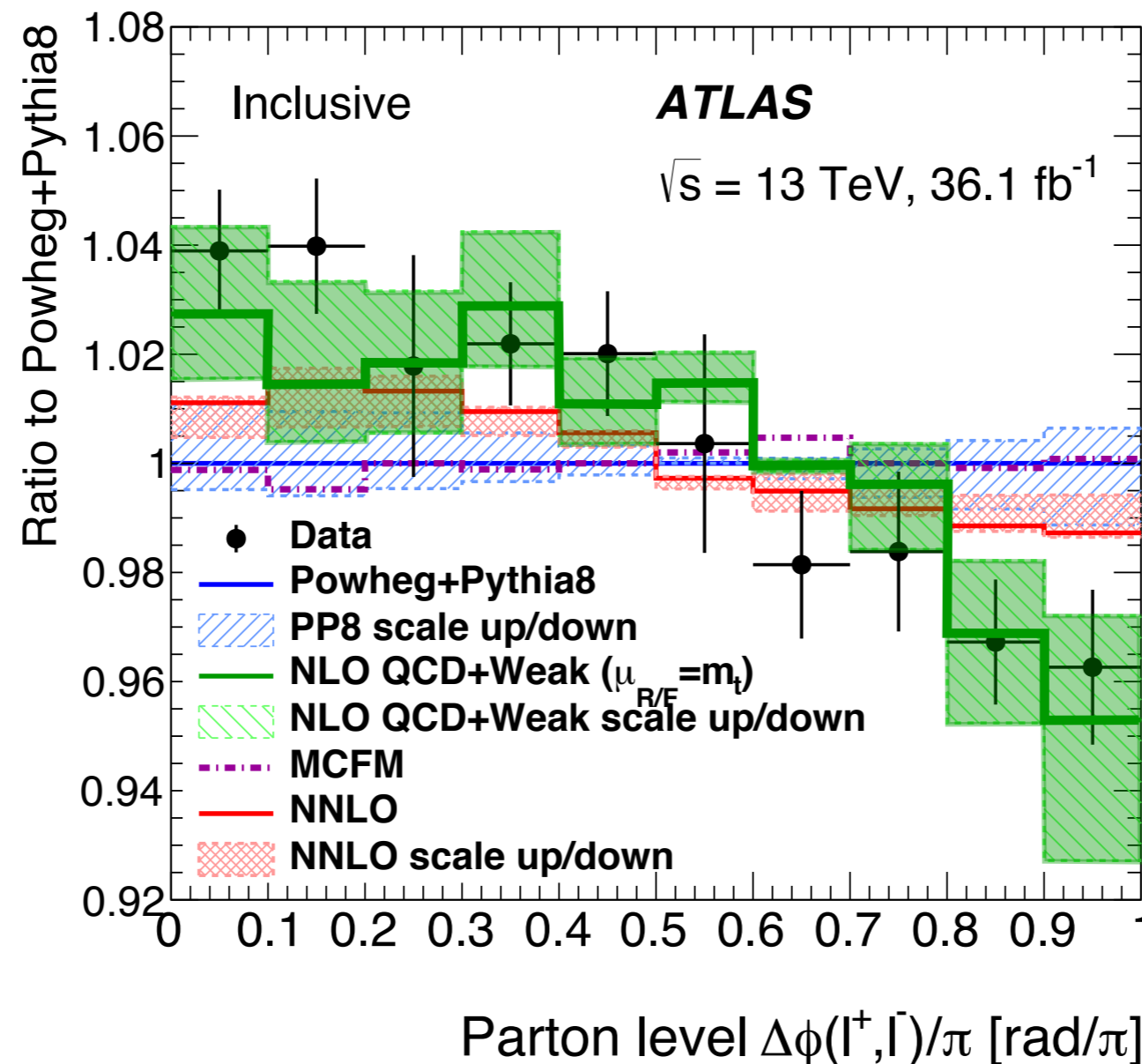
- Lots of discussions in the theory community after the CONF note for this result.
- ATLAS focused on understanding the assumptions involved in the template hypotheses.



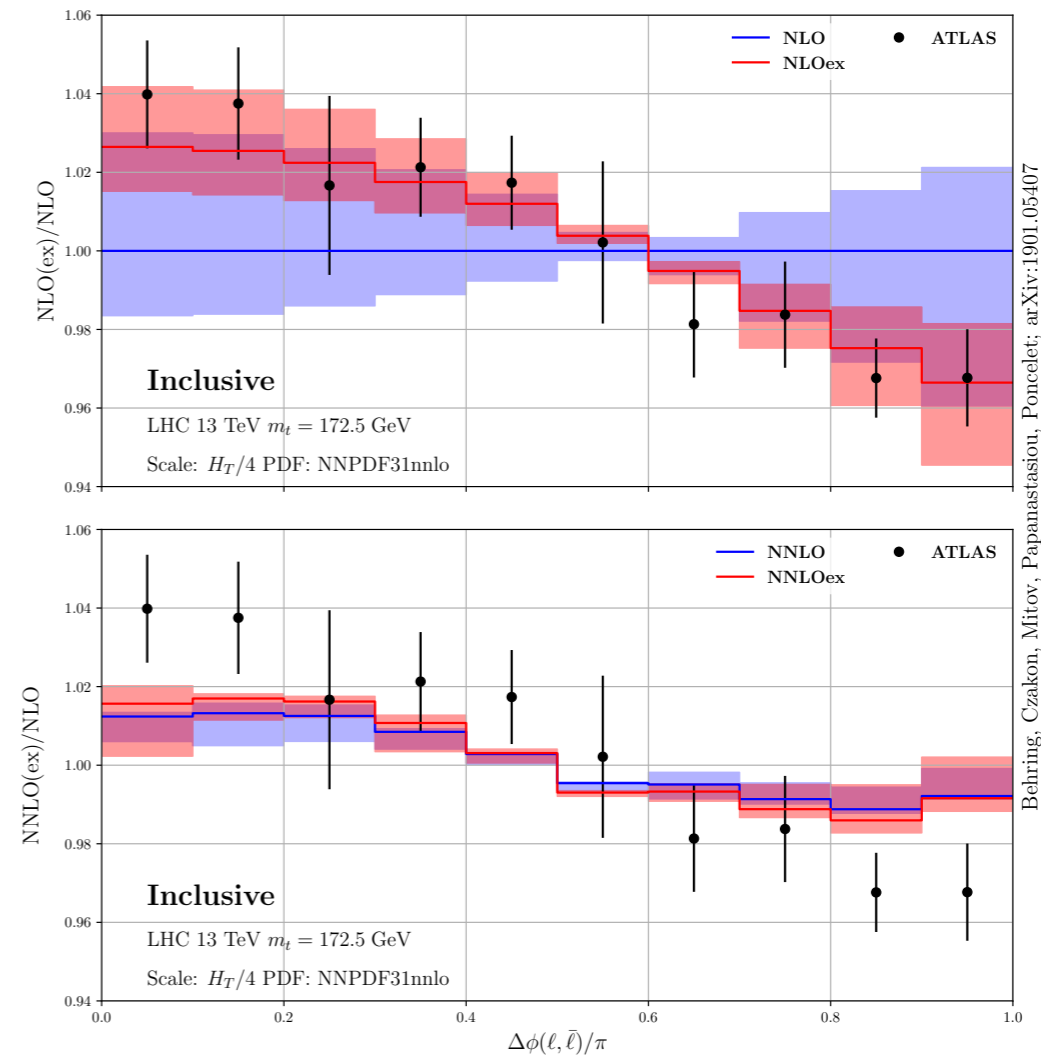
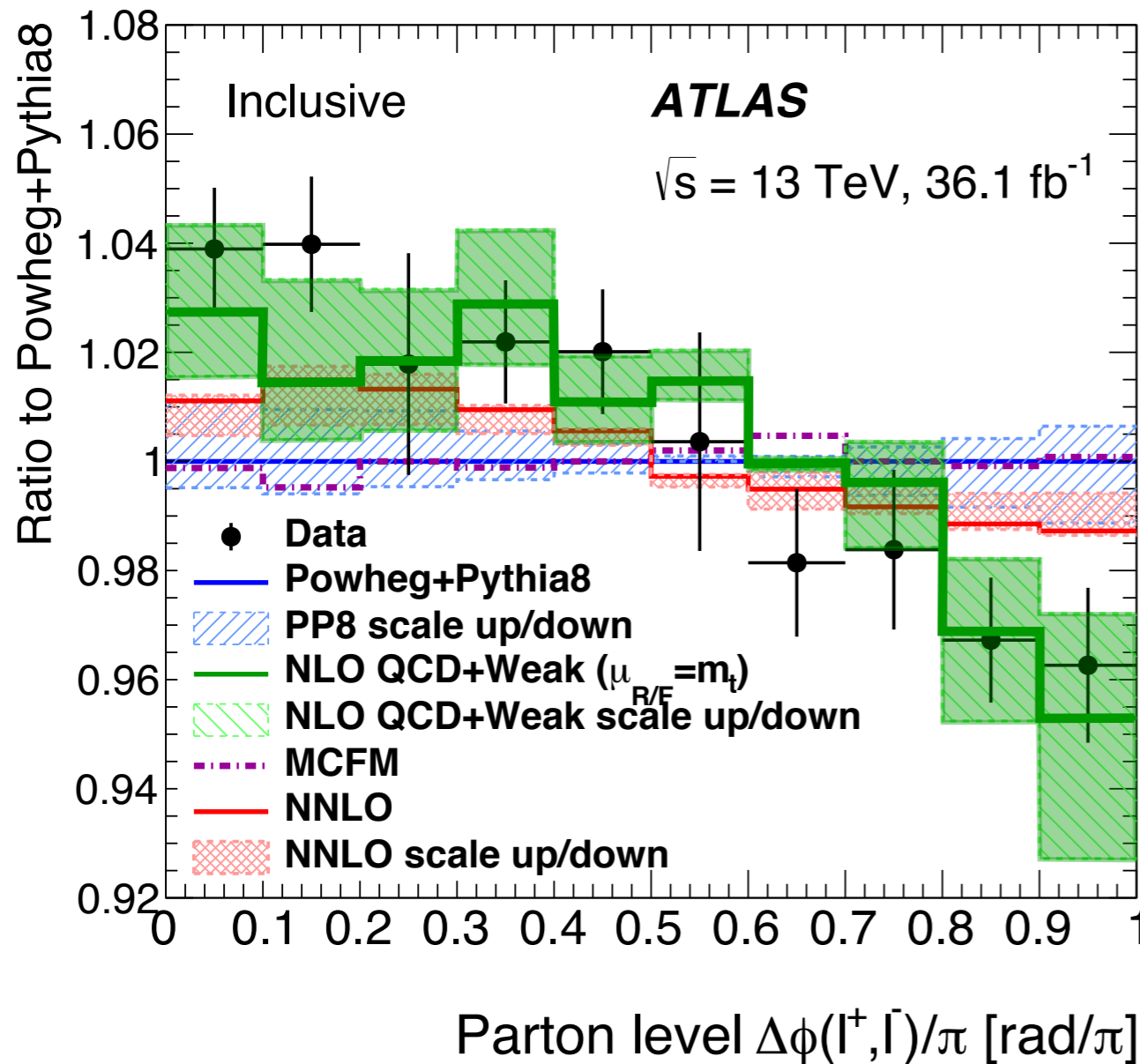
- **NLO + Parton** shower MC consistent with fixed-order calculations from **MCFM**.
- We also tested assumptions such as LO vs. NLO top decays and the use of the narrow-width approximation, none were significant shifts.



- ATLAS investigated alternative predictions, including bespoke state-of-the-art **NNLO-QCD** predictions (Brun et. al.).
- These are closer to the data, but not all the way (significance would be 2.2 sigma).

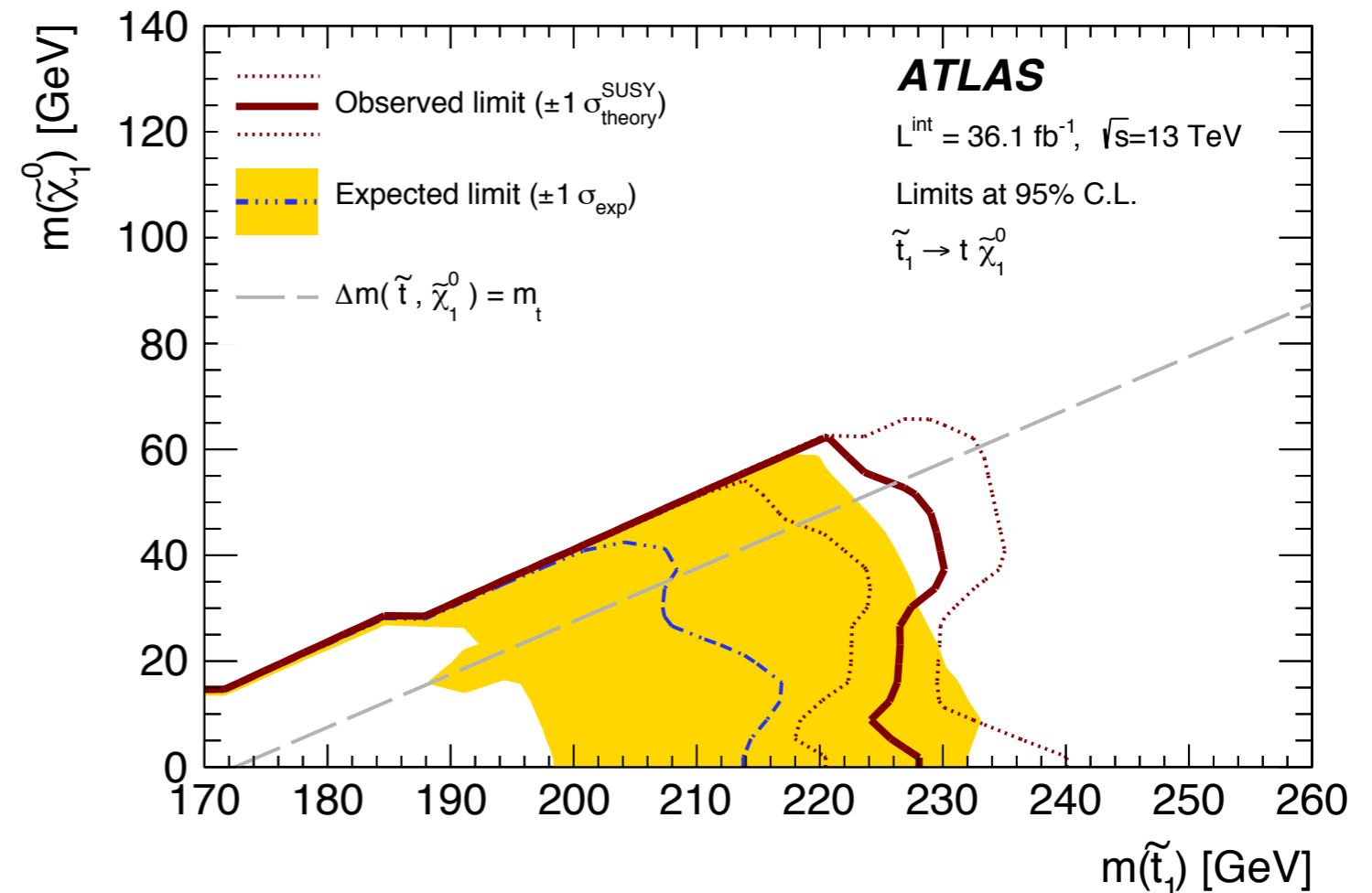
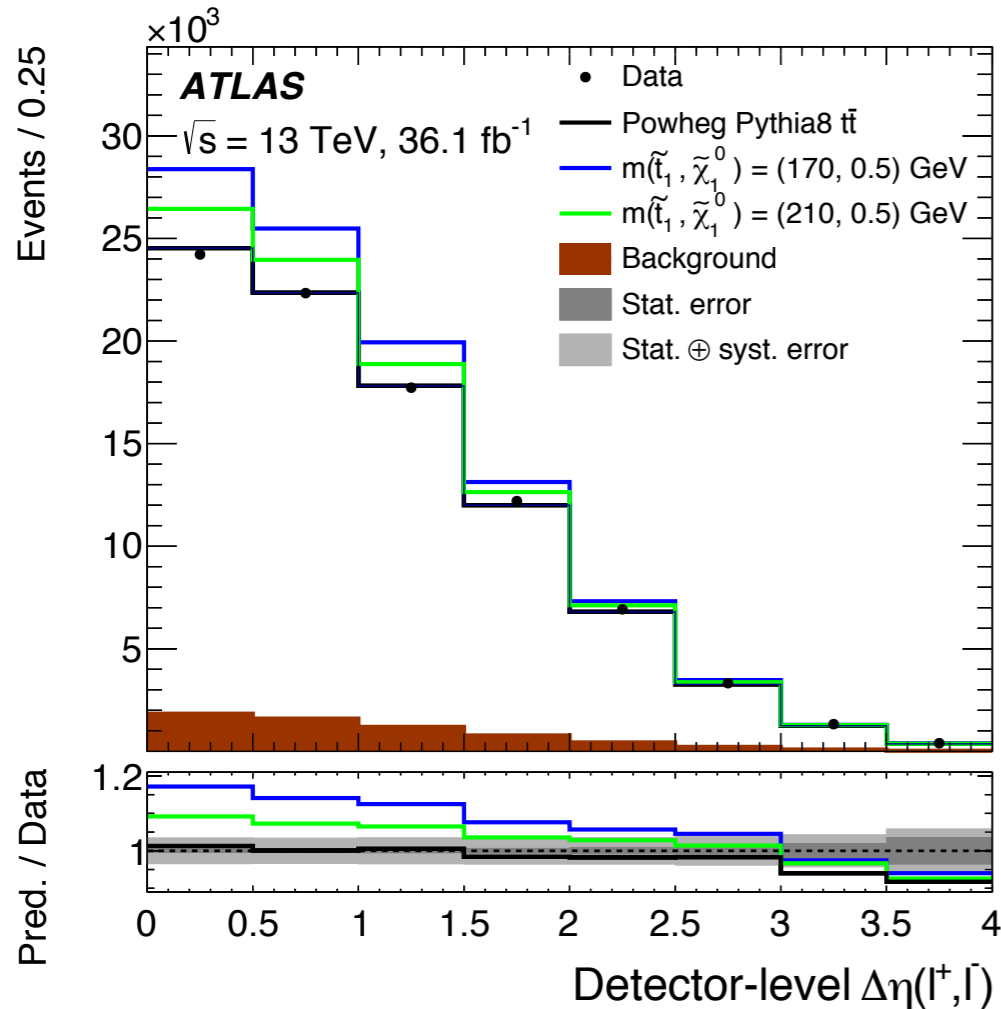


- A **NLO-QCD + Weak Ratio-Expanded** prediction was also available.
- Agrees with the data but with large scale uncertainties.
- Further studies imply that the Weak corrections are not driving the agreement but rather the ratio expansion method. However, at NNLO, this may not be true...



Paper: [:1901.05407](https://arxiv.org/abs/1901.05407)
Extension: [Web Page](#)

- It appears that when the ratio is expanded at NNLO, the agreement disappears.
- Discrepancy remains unresolved, but is stimulating state-of-the-art predictions and development in the theory community.



- Use both the $|\Delta\phi_{\ell\ell}|$ and $|\Delta\eta_{\ell\ell}|$ to set limits on SUSY stop production.
- Exclude Stops with a mass below $\sim 220 \text{ GeV}$ for all kinematically-allowed neutralino masses:
 - ➡ Limit is driven by $|\Delta\eta_{\ell\ell}|$ but additional modelling uncertainties are included to account for the Data/Prediction disagreement in $|\Delta\phi_{\ell\ell}|$.

May 2018

$\tilde{t}_1\tilde{t}_1$ production, $\tilde{t}_1 \rightarrow b f \tilde{\chi}_1^0$ / $\tilde{t}_1 \rightarrow c \tilde{\chi}_1^0$ / $\tilde{t}_1 \rightarrow W b \tilde{\chi}_1^0$ / $\tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$

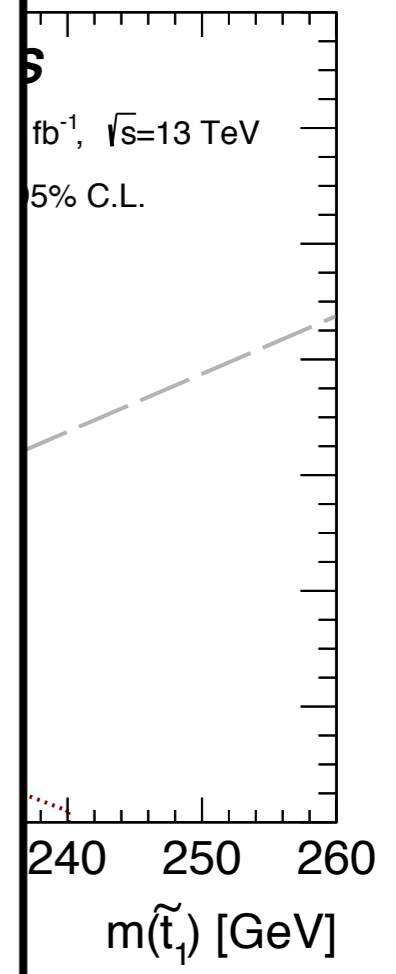
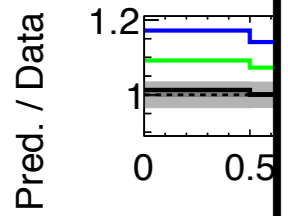
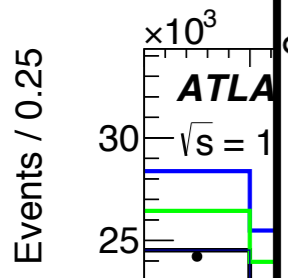
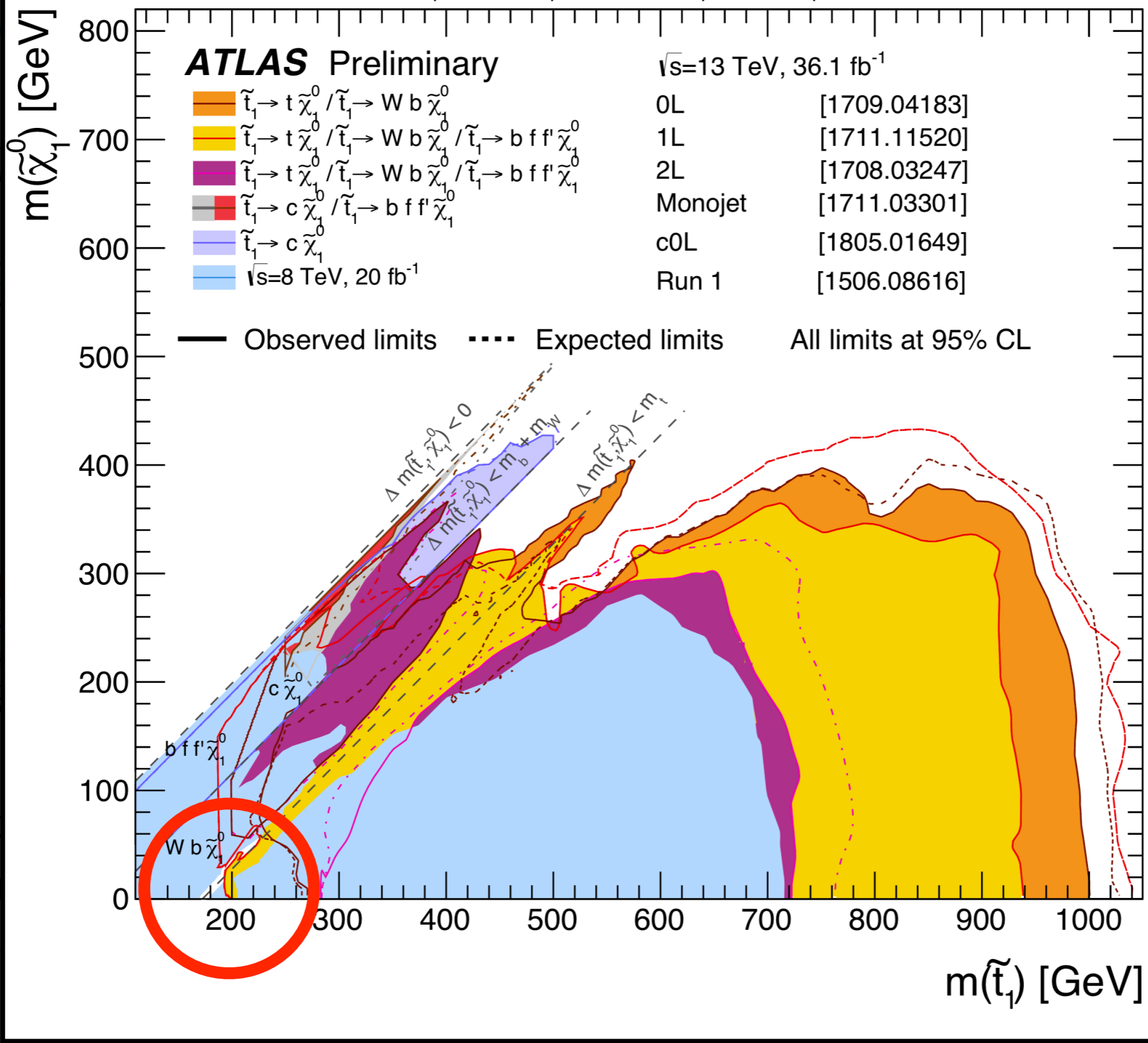
ATLAS Preliminary

$\sqrt{s}=13$ TeV, 36.1 fb $^{-1}$

- $\tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$ / $\tilde{t}_1 \rightarrow W b \tilde{\chi}_1^0$
- $\tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$ / $\tilde{t}_1 \rightarrow W b \tilde{\chi}_1^0$ / $\tilde{t}_1 \rightarrow b f \tilde{\chi}_1^0$
- $\tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$ / $\tilde{t}_1 \rightarrow W b \tilde{\chi}_1^0$ / $\tilde{t}_1 \rightarrow b f \tilde{\chi}_1^0$
- $\tilde{t}_1 \rightarrow c \tilde{\chi}_1^0$ / $\tilde{t}_1 \rightarrow b f \tilde{\chi}_1^0$
- $\tilde{t}_1 \rightarrow c \tilde{\chi}_1^0$
- $\sqrt{s}=8$ TeV, 20 fb $^{-1}$

0L	[1709.04183]
1L	[1711.11520]
2L	[1708.03247]
Monojet	[1711.03301]
c0L	[1805.01649]
Run 1	[1506.08616]

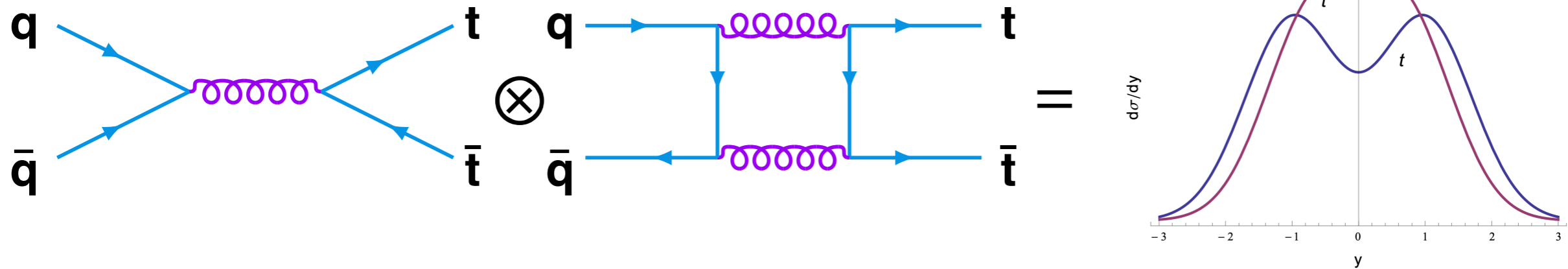
— Observed limits - - - - Expected limits All limits at 95% CL



- Use b
- Exclu
- neutr
- ➡ L

to account for the Data/Prediction disagreement in $|\Delta\phi_{\ell\ell}|$.

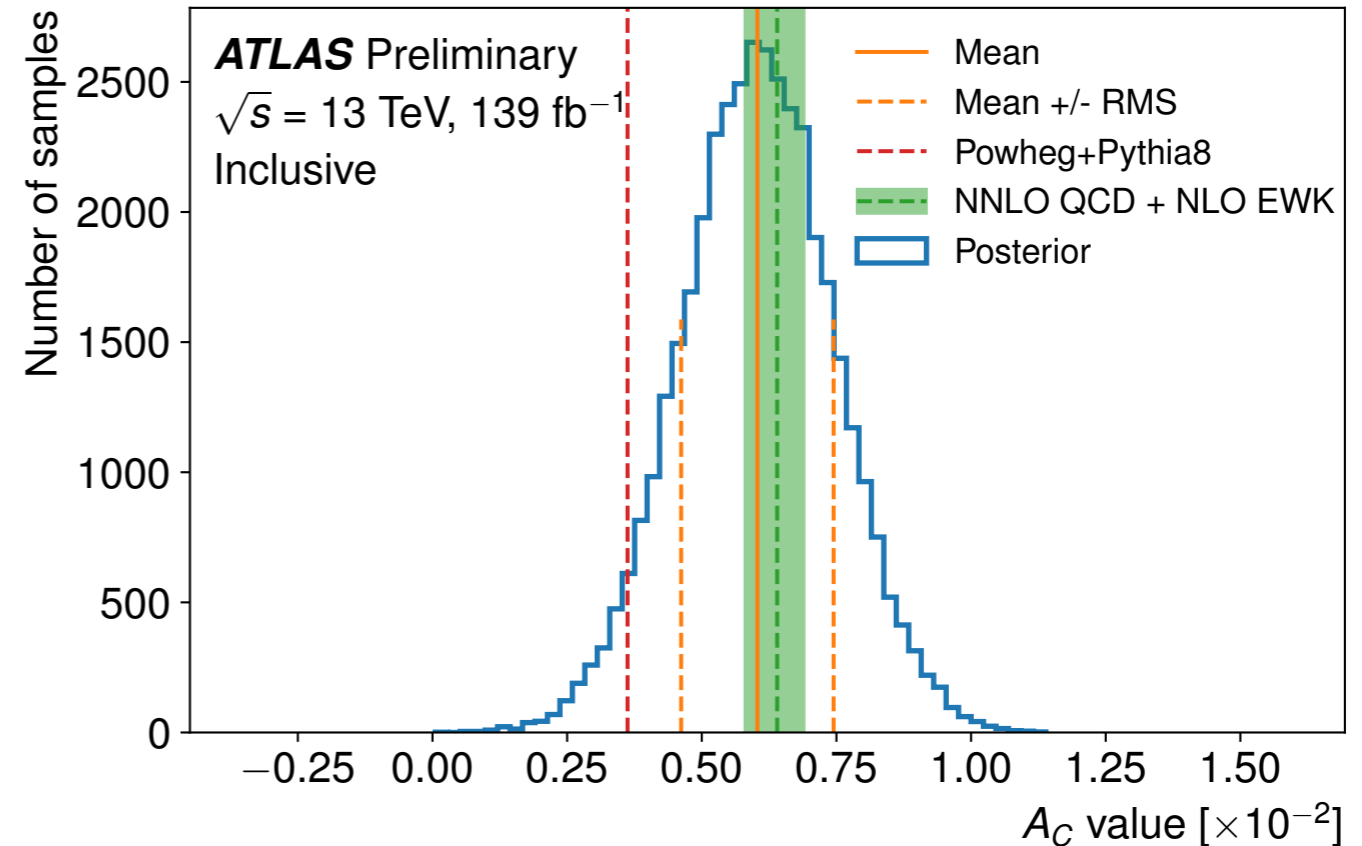
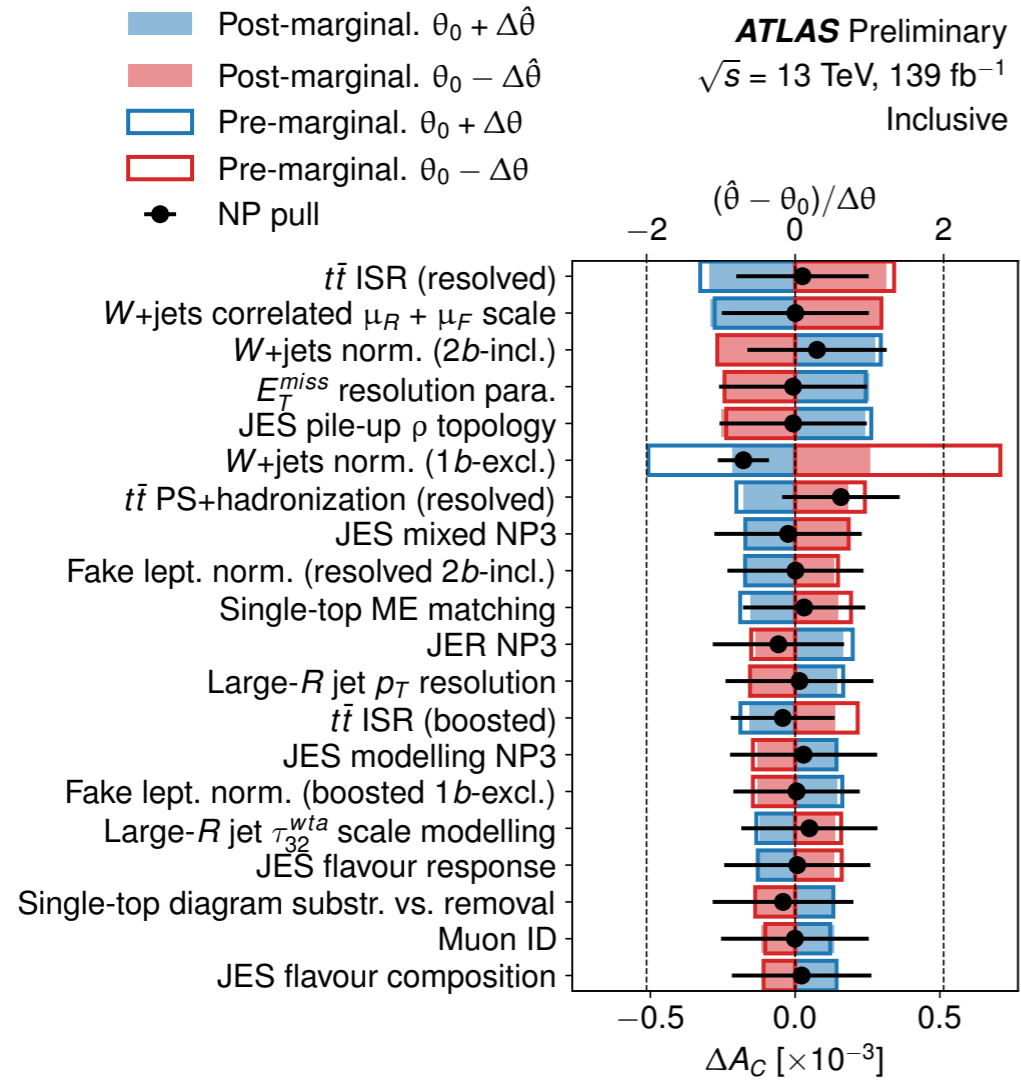
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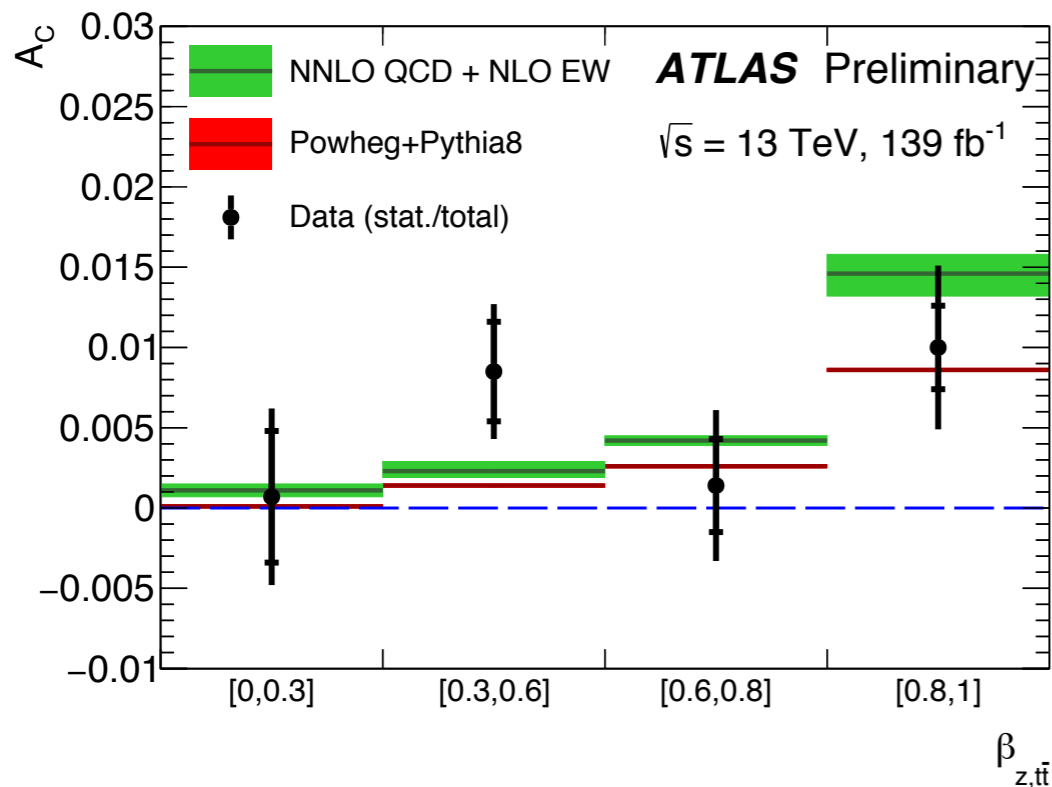
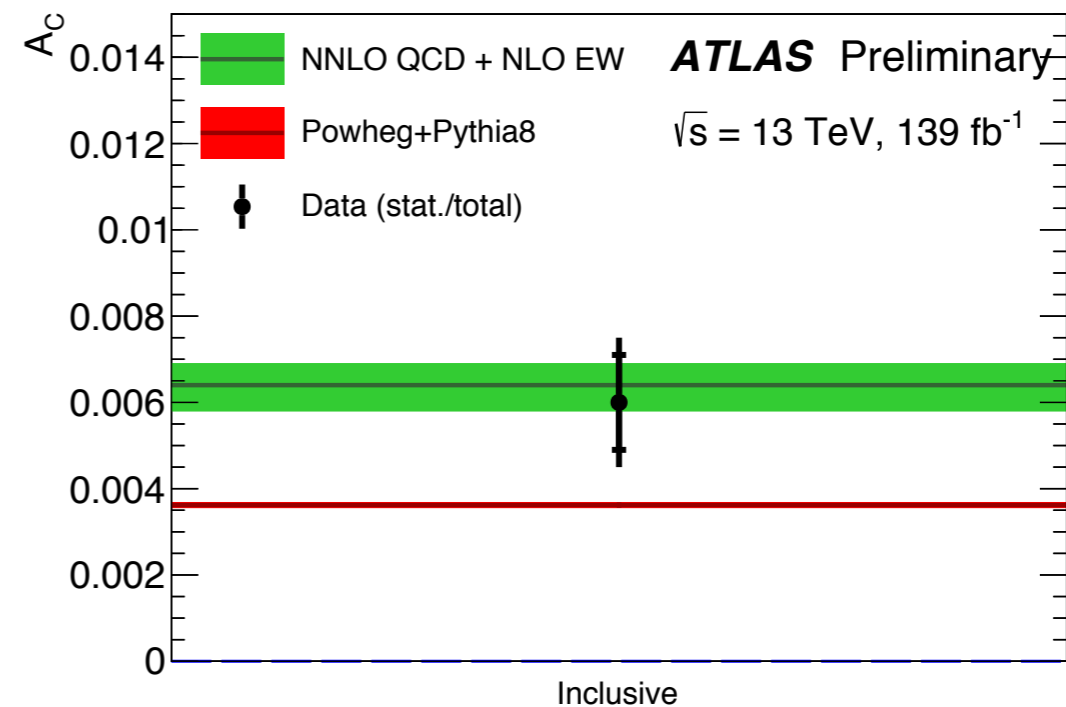
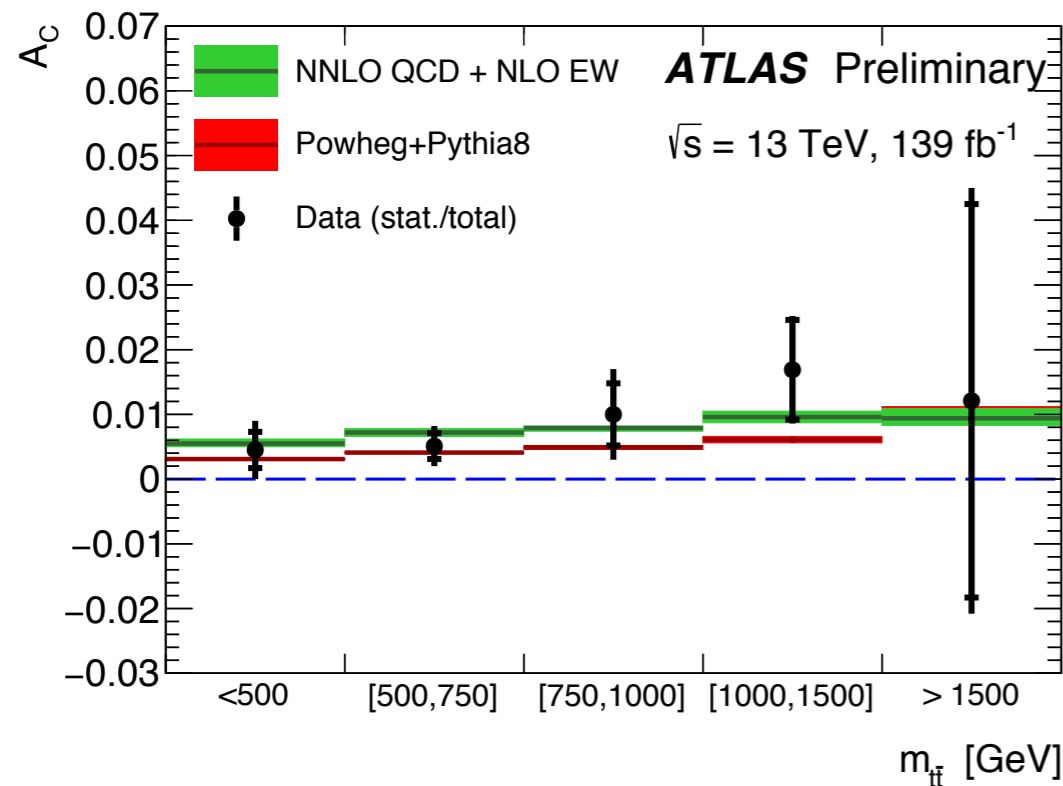
- Interference between Born and Box diagrams, and to a lesser extent ISR and FSR, induce an asymmetry in the direction of top and anti-top quarks:
 - ➡ Top quarks produced more forward than anti-top quarks.
 - ➡ (mention Tevatron ppbar asym.)
- This “*charge asymmetry*” only exists in higher-order qqbar production, gg-fusion is symmetric to all orders:
 - ➡ Extremely challenging to measure at the LHC (qqbar ~10% of production fraction at 13 TeV).

$$A_C^{t\bar{t}} = \frac{N(\Delta |y| > 0) - N(\Delta |y| < 0)}{N(\Delta |y| > 0) + N(\Delta |y| < 0)} \quad \Delta |y| = |y(t)| - |y(\bar{t})|$$

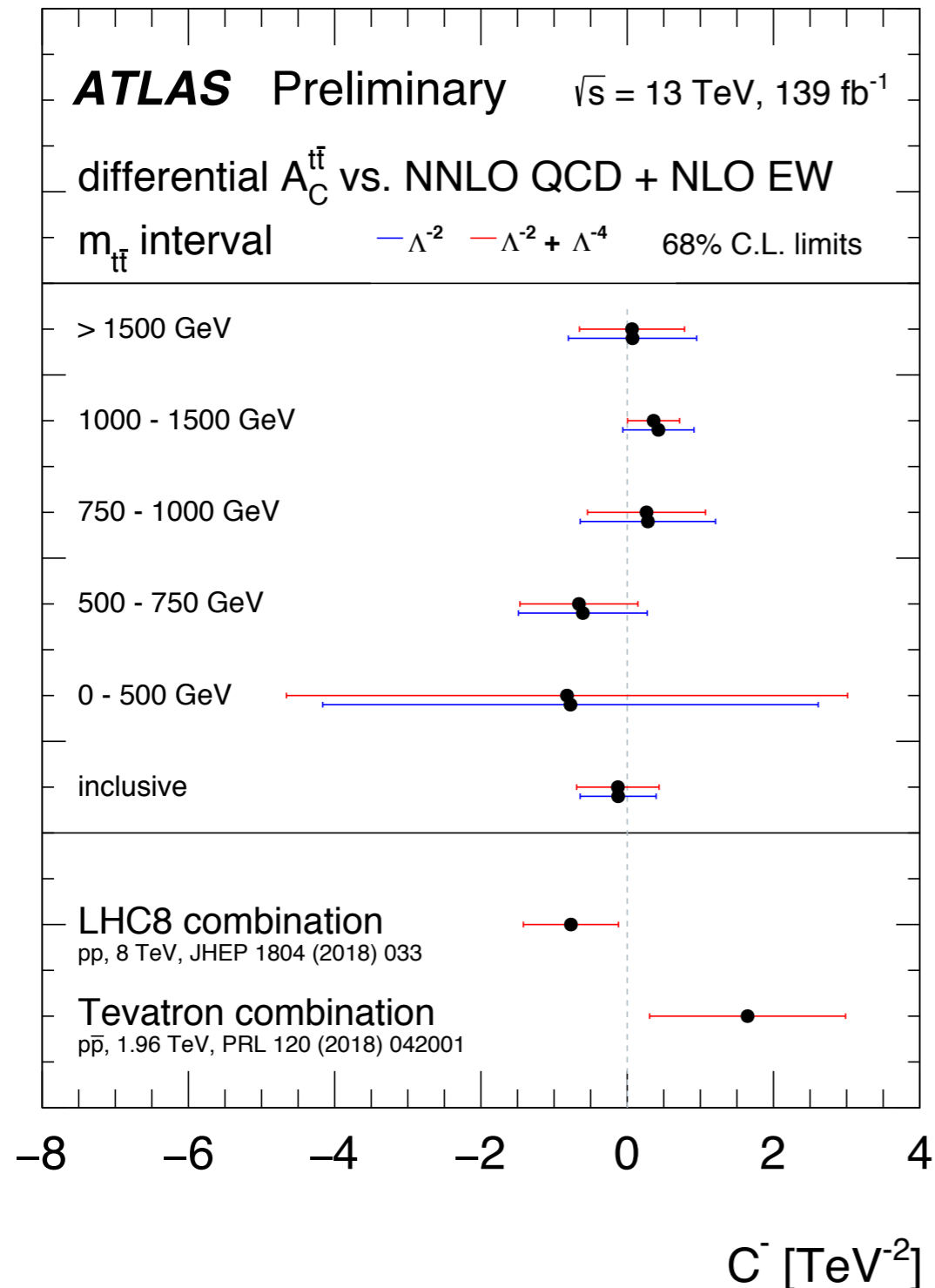
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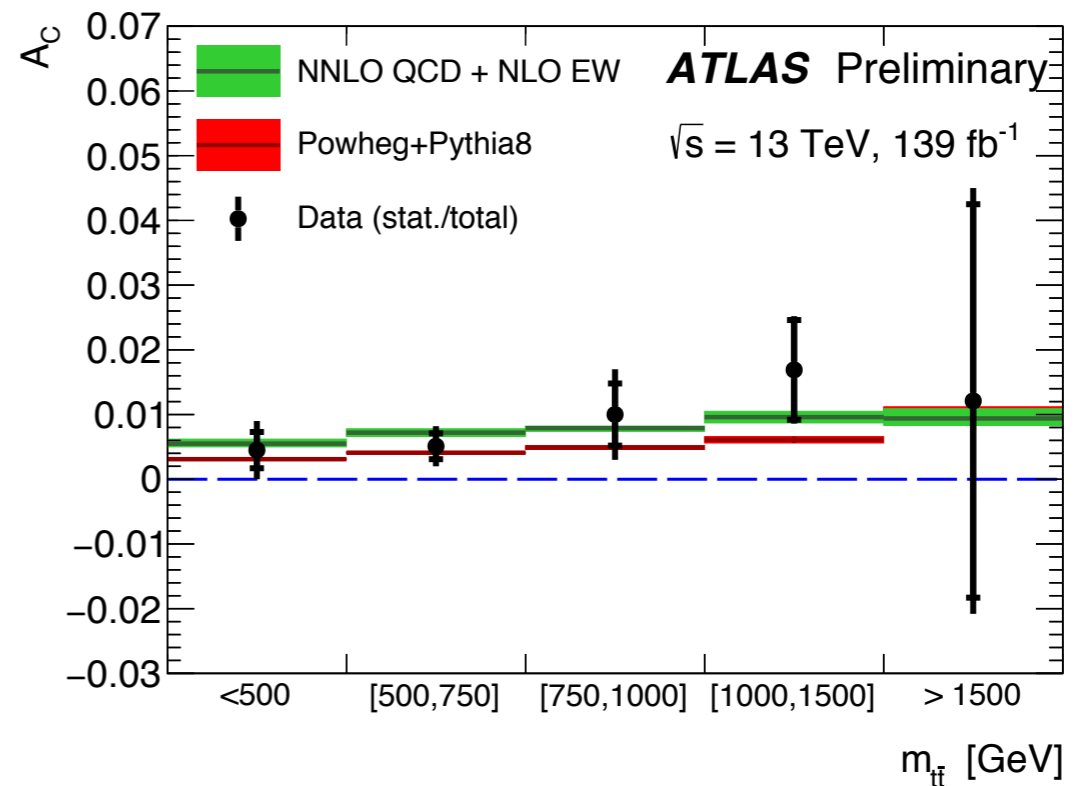
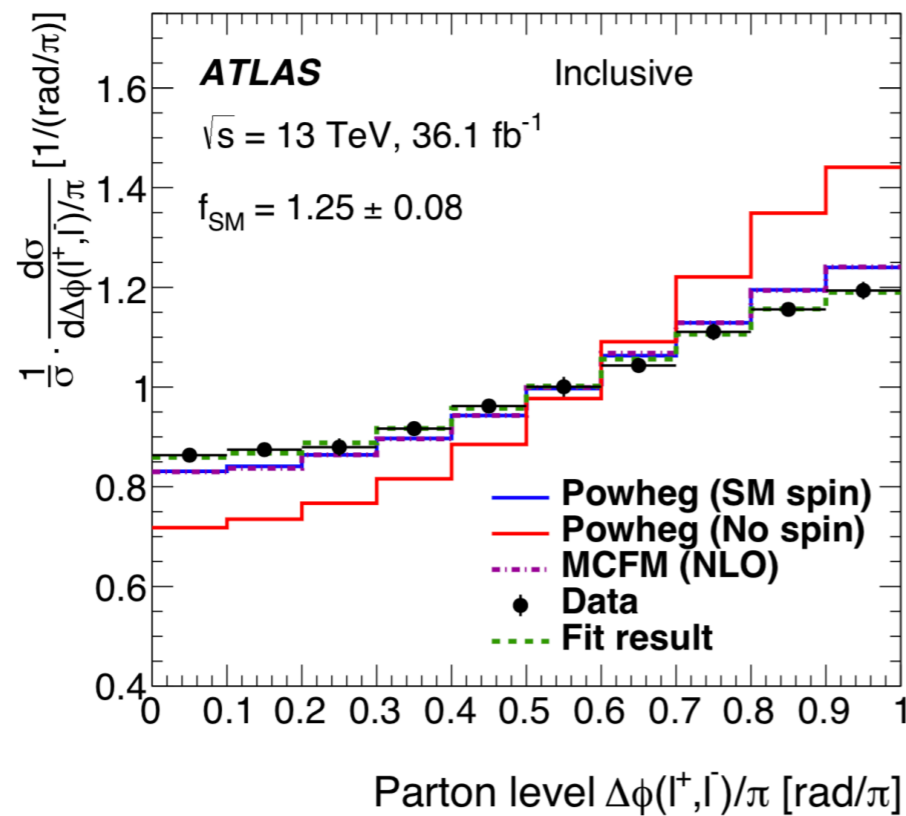
- Charge asymmetry is extracted from 139 fb^{-1} of 13 TeV data, using a resolved ($p_T(t) < 400 \text{ GeV}$) and boosted ($p_T(t) > 400$) single lepton (e/μ) selection.
- The $|\Delta y|$ distribution is unfolded using a likelihood-based technique called “*fully bayesian unfolding*” [ref].
- ➡ Systematic uncertainties are profiled as nuisance parameters.



- Charge Asymmetry measured inclusively to be 0.6% +/- 0.15%, in agreement with the NNLO QCD + NLO EW predictions and 4σ from 0.
➔ **First evidence for charge asymmetry in pp collisions.**
- Also measured as a function of $m(t\bar{t})$ and the boost of the $t\bar{t}$ system.

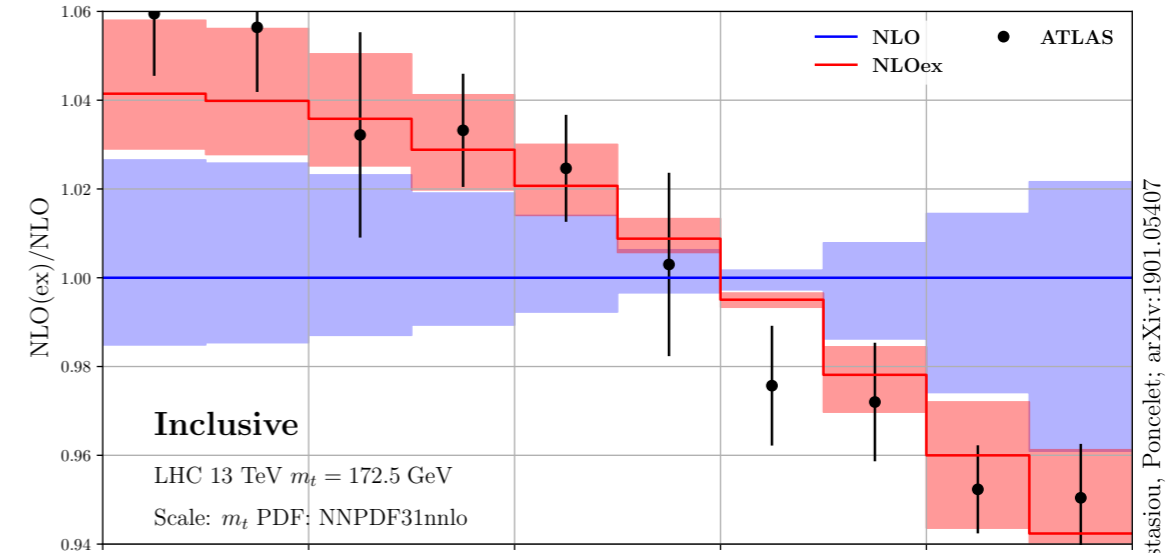
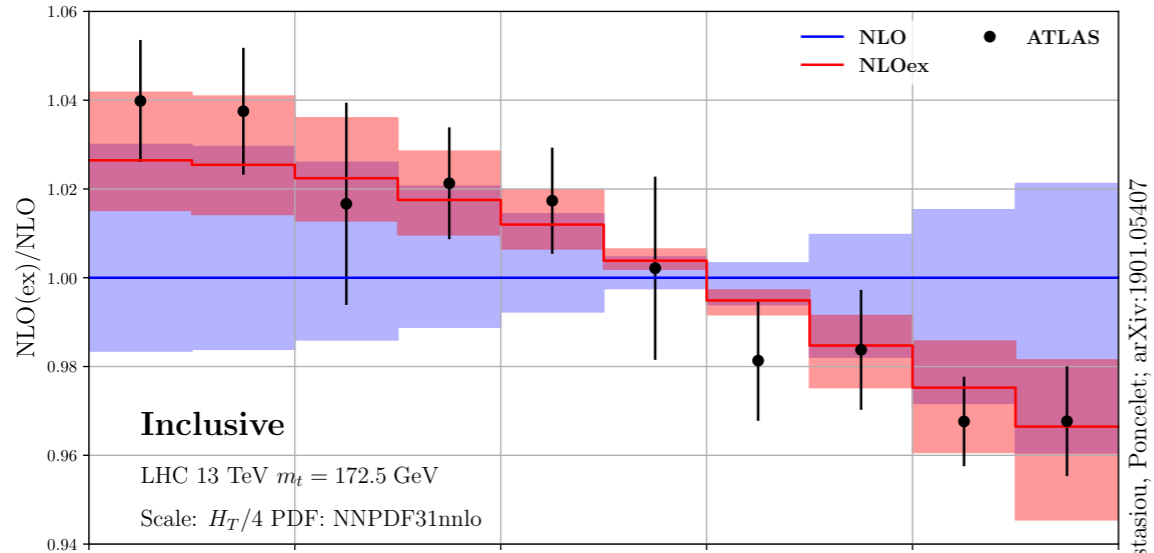


- These measurements can also be used to set limits on EFT operators:
- C^- = four fermion operator assuming flavour conservation and equal up-down type couplings (simple axion model).
- Inclusive and differential results are surpassing those set via ATLAS+CMS combination in Run1.
 - Not a large dependence on quadratic terms:
 - ➡ dimension 6 approach is stable and appropriate.



- ATLAS has a large suite of top properties measurements that are probing the SM's heaviest particle to ever-greater precision.
- We are now encountering significant tensions between data and predictions:
 - ➡ Perhaps this is highlighting the limitations in our understanding of $t\bar{t}$ production and decay.
 - ➡ Perhaps it's something more interesting...
- We are now also able to see subtle higher-order effects.

Backup



Behring, Czakon, Mitov, Papanastasiou, Poncellet; arXiv:1901.05407

Behring, Czakon, Mitov, Papanastasiou, Poncellet; arXiv:1901.05407

