

# Positivity Bounds on Vector Boson Scattering

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Cen Zhang & **SYZ**, arXiv:1808.00010  
Bi Qi, Cen Zhang & **SYZ**, arXiv:1902.08977

# Effective field theories (EFTs)

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- EFTs are widely used in physics

BSM physics, GR, inflation, dark energy,...

- Separation of physics at different scales

write down all local operators consistent with symmetries suppressed by cut-off scale

$$\mathcal{L} = \sum_i \Lambda^4 f_i \mathcal{O}_i \left( \frac{\text{boson}}{\Lambda}, \frac{\text{fermion}}{\Lambda^{3/2}}, \frac{\partial}{\Lambda} \right)$$

# Are all EFTs allowed?

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Answer: No!

Not every set of Wilson coefficients are allowed!

$$e^{\frac{i}{\hbar}S_W[\text{light}]} = \int D[\text{heavy}] e^{\frac{i}{\hbar}S_{UV}[\text{light,heavy}]}$$

UV completion satisfies:

Lorentz invariance, unitarity,  
locality, causality, **analyticity**, ...



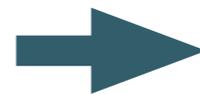
**Positivity bounds on Wilson coefficients**

# Simple example: P(X)

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$$\mathcal{L} = -\frac{1}{2}\partial_\mu\phi\partial^\mu\phi + \frac{\lambda}{\Lambda^4}(\partial_\mu\phi\partial^\mu\phi)^2 + \dots$$

Positivity bound



$$\lambda > 0$$

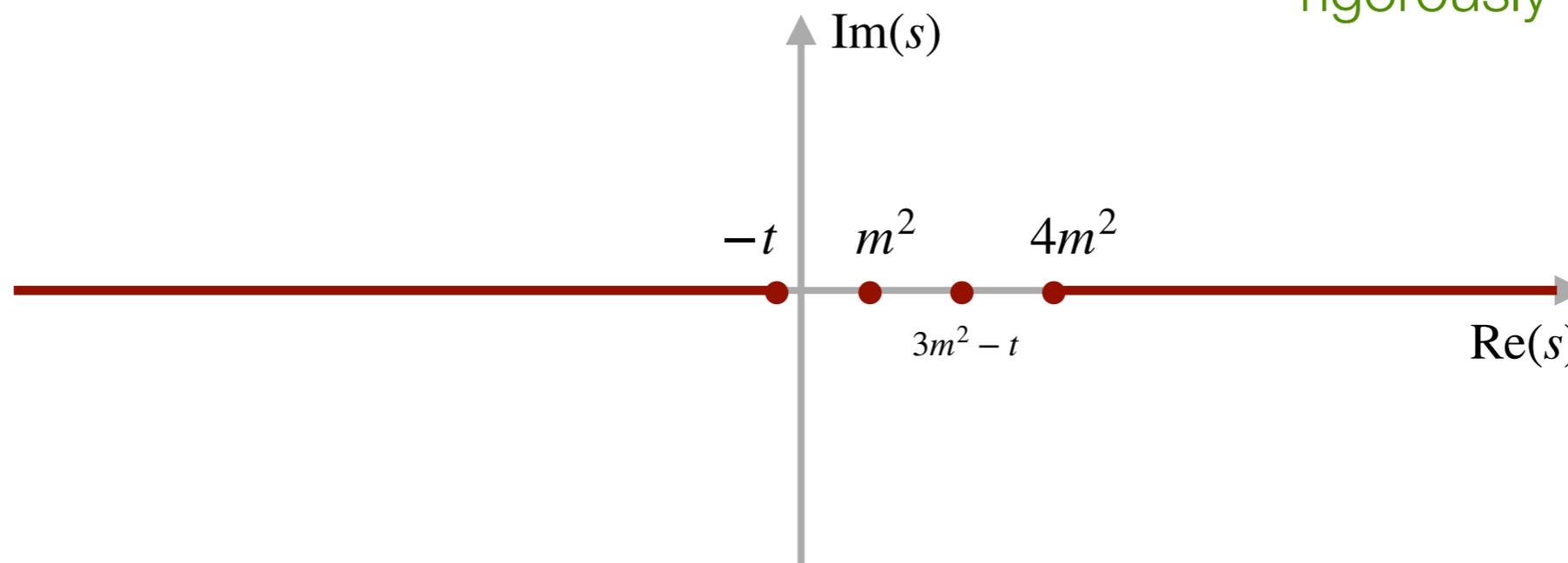
Theories with  $\lambda < 0$  **do not** have an analytic UV completion

How does it work?

# Analyticity and unitarity

Complex amplitude  $A(s, t)$

2 to 2 scattering  
rigorously proven in 60'



Optical theorem:

$$\text{Im}[A(s, 0)] = \sqrt{s(s - 4m^2)}\sigma(s) > 0$$

Froissart bound: as  $s \rightarrow \infty$ ,  $|A(s, 0)| < Cs \ln^2 s$

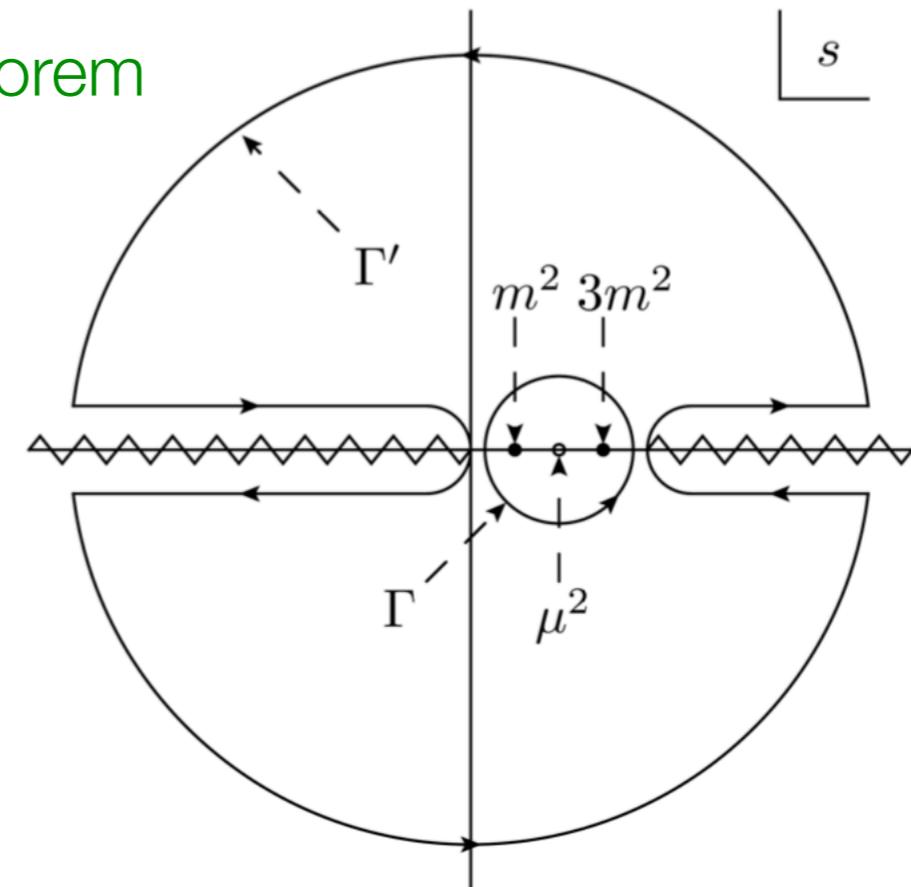
# Dispersion relation

$$f := \frac{1}{2\pi i} \oint_{\Gamma} ds \frac{A(s,0)}{(s - \mu^2)^3}$$

Cauchy's theorem

Froissart bound: as  $s \rightarrow \infty$ ,  $|A(s,0)| < Cs \ln^2 s$

$$f = \frac{1}{2\pi i} \left( \int_{-\infty}^0 + \int_{4m^2}^{+\infty} \right) ds \frac{\text{Disc } A(s,0)}{(s - \mu^2)^3}$$



Crossing and  $\text{Disc } A(s,0) = 2i \text{Im}A(s,0)$

$$f = \frac{1}{\pi} \int_{4m^2}^{\infty} ds \left[ \frac{\text{Im}A(s,0)}{(s - \mu^2)^3} + \frac{\text{Im}A^{(u)}(s,0)}{(s + \mu^2 - 4m^2)^3} \right]$$

# Positivity bound

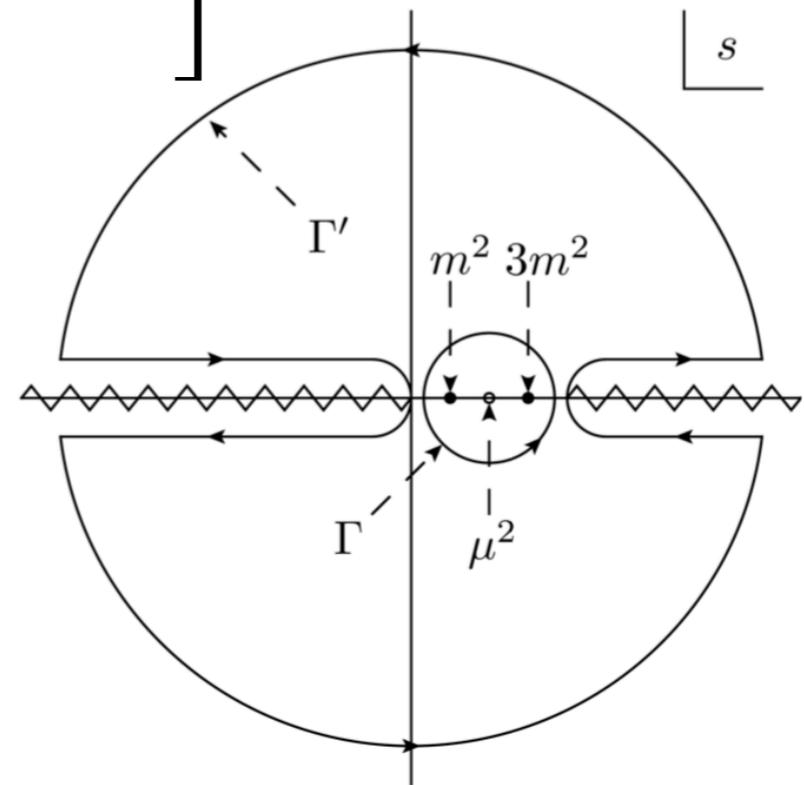
Optical theorem:  $\text{Im}[A(s, 0)] = \sqrt{s(s - 4m^2)}\sigma(s) > 0$

$$f = \frac{1}{\pi} \int_{4m^2}^{\infty} ds \left[ \frac{\sqrt{s(s - 4m^2)}}{(s - \mu^2)^3} \sigma(s) + \frac{\sqrt{s(s - 4m^2)}}{(s + \mu^2 - 4m^2)^3} \sigma^{(u)}(s) \right]$$

For  $s > 4m^2$ ,  $0 < \mu^2 < 4m^2$

$$f > 0$$

Adams, Arkani-Hamed, Dubovsky,  
Nicolis, Rattazzi, 2006



$$f = \sum_{\Gamma} \text{Res} \left[ \frac{A(s, 0)}{(s - \mu^2)^3} \right]$$

**Calculable within low energy EFT!**

# Application to the SMEFT

# Standard Model Effective Field Theory

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- Write down all possible operators with
  - SM particle contents and global symmetries
  - SM gauge group structure  $SU(3)_c \times SU(2)_L \times U(1)$

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_j \frac{f_j^{(6)} \mathcal{O}_j^{(6)}}{\Lambda^2} + \sum_i \frac{f_i^{(8)} \mathcal{O}_i^{(8)}}{\Lambda^4} + \dots$$

# Vector boson scattering (VBS)

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$$V_1 + V_2 \rightarrow V_3 + V_4, \quad V_i \in \{Z, W^+, W^-, \gamma\}$$

$$\begin{aligned}
 O_{S,0} &= [(D_\mu \Phi)^\dagger D_\nu \Phi] \times [(D^\mu \Phi)^\dagger D^\nu \Phi] \\
 O_{S,1} &= [(D_\mu \Phi)^\dagger D^\mu \Phi] \times [(D_\nu \Phi)^\dagger D^\nu \Phi] \\
 O_{S,2} &= [(D_\mu \Phi)^\dagger D_\nu \Phi] \times [(D^\nu \Phi)^\dagger D^\mu \Phi] \\
 O_{M,0} &= \text{Tr} \left[ \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \right] \times [(D_\beta \Phi)^\dagger D^\beta \Phi] \\
 O_{M,1} &= \text{Tr} \left[ \hat{W}_{\mu\nu} \hat{W}^{\nu\beta} \right] \times [(D_\beta \Phi)^\dagger D^\mu \Phi] \\
 O_{M,2} &= \left[ \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \right] \times [(D_\beta \Phi)^\dagger D^\beta \Phi] \\
 O_{M,3} &= \left[ \hat{B}_{\mu\nu} \hat{B}^{\nu\beta} \right] \times [(D_\beta \Phi)^\dagger D^\mu \Phi] \\
 O_{M,4} &= \left[ (D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} D^\mu \Phi \right] \times \hat{B}^{\beta\nu} \\
 O_{M,5} &= \left[ (D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} D^\nu \Phi \right] \times \hat{B}^{\beta\mu} \\
 O_{M,7} &= \left[ (D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} \hat{W}^{\beta\mu} D^\nu \Phi \right]
 \end{aligned}$$

$$\begin{aligned}
 O_{T,0} &= \text{Tr} \left[ \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \right] \times \text{Tr} \left[ \hat{W}_{\alpha\beta} \hat{W}^{\alpha\beta} \right] \\
 O_{T,1} &= \text{Tr} \left[ \hat{W}_{\alpha\nu} \hat{W}^{\mu\beta} \right] \times \text{Tr} \left[ \hat{W}_{\mu\beta} \hat{W}^{\alpha\nu} \right] \\
 O_{T,2} &= \text{Tr} \left[ \hat{W}_{\alpha\mu} \hat{W}^{\mu\beta} \right] \times \text{Tr} \left[ \hat{W}_{\beta\nu} \hat{W}^{\nu\alpha} \right] \\
 O_{T,5} &= \text{Tr} \left[ \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \right] \times \hat{B}_{\alpha\beta} \hat{B}^{\alpha\beta} \\
 O_{T,6} &= \text{Tr} \left[ \hat{W}_{\alpha\nu} \hat{W}^{\mu\beta} \right] \times \hat{B}_{\mu\beta} \hat{B}^{\alpha\nu} \\
 O_{T,7} &= \text{Tr} \left[ \hat{W}_{\alpha\mu} \hat{W}^{\mu\beta} \right] \times \hat{B}_{\beta\nu} \hat{B}^{\nu\alpha} \\
 O_{T,8} &= \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \times \hat{B}_{\alpha\beta} \hat{B}^{\alpha\beta} \\
 O_{T,9} &= \hat{B}_{\alpha\mu} \hat{B}^{\mu\beta} \times \hat{B}_{\beta\nu} \hat{B}^{\nu\alpha},
 \end{aligned}$$

# Leading order bounds

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If dim-6 ops alone, all positivity bounds are violated

$$\mathcal{O}\left(\frac{1}{\Lambda^4}\right) : \quad \sum_i (-C_i) \left( \sum_j D_j f_j^{(6)} \right)^2 \geq 0, \quad C_i > 0$$

Positivity bounds require the existence of higher-D ops!

Then, add dim-8 ops

$$(\text{dim-8 part}) + (\text{dim-6 part}) > 0$$

$$\mathcal{O}\left(\frac{1}{\Lambda^4}\right) :$$

$$(\text{dim-8 part}) > 0$$



weak but simpler bounds



$$\sum_i E_i f_i^{(8)} \geq 0$$

# Leading order bounds for dim-8

$$M_{S,ij} F_{S,j} > 0$$

$$M_{M,ij} F_{M,j} > 0$$

$$M_{T,ij} F_{T,j} > 0$$

+ a few nonlinear bounds

$$M_S = \begin{pmatrix} 2 & 1 & 1 \\ 1 & 1 & 1 \\ 1 & 0 & 1 \end{pmatrix}$$

$$M_M = \begin{pmatrix} 0 & -2c_W^4 & 0 & -s_W^4 & 0 & s_W^2 c_W^2 & c_W^4 \\ 0 & -2c_W^4 & 0 & -s_W^4 & 0 & -s_W^2 c_W^2 & c_W^4 \\ 0 & -2 & 0 & 0 & 0 & 0 & 1 \\ 0 & -2 & 0 & -1 & 0 & 1 & 1 \\ 0 & -2 & 0 & -1 & 0 & -1 & 1 \end{pmatrix}$$

$$F_{S,i} = (F_{S,0}, F_{S,1}, F_{S,2})^T$$

$$F_{M,i} = (F_{M,0}, F_{M,1}, F_{M,2}, F_{M,3}, F_{M,4}, F_{M,5}, F_{M,7})^T$$

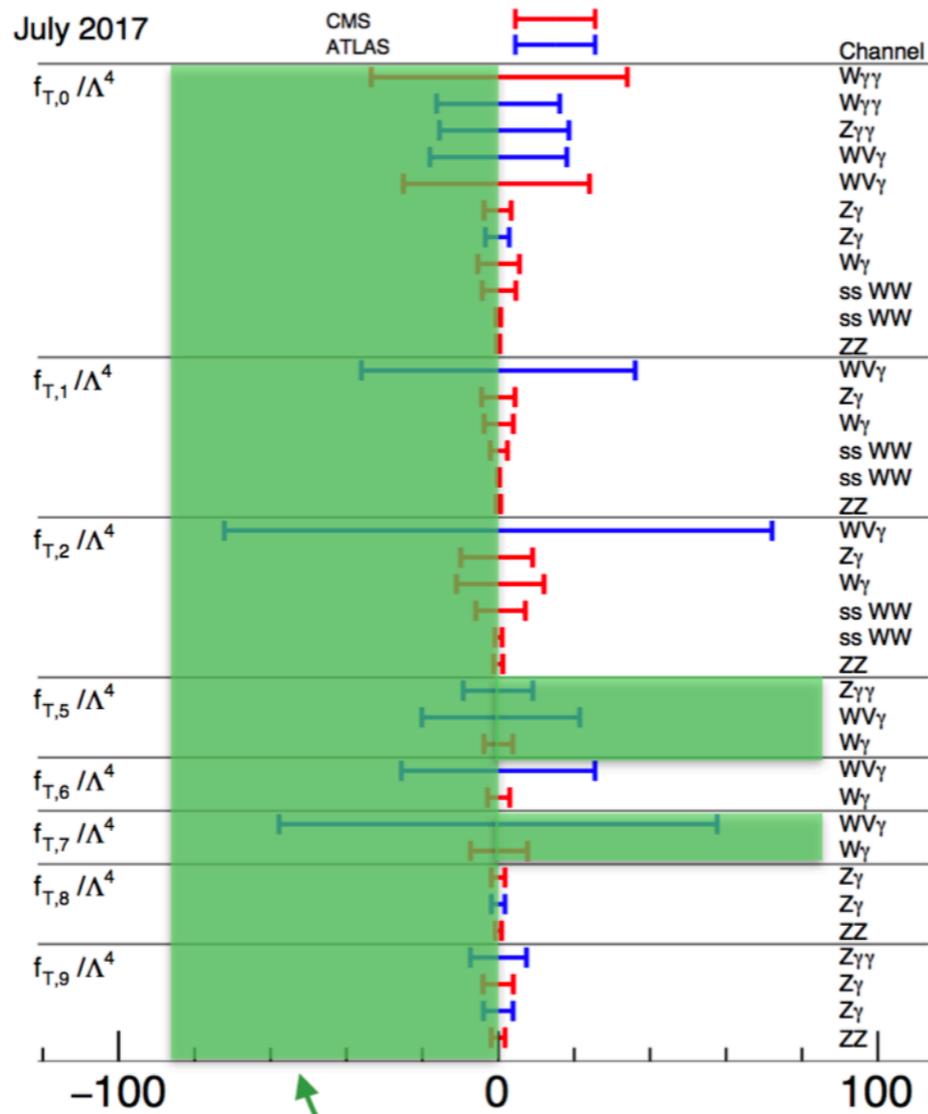
$$F_{T,i} = (F_{T,0}, F_{T,1}, F_{T,2}, F_{T,5}, F_{T,6}, F_{T,7}, F_{T,8}, F_{T,9})^T$$

$$M_T = \begin{pmatrix} 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 2 & 1 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 8 & 12 & 5 & 0 & 0 & 0 & 0 & 0 & 0 \\ 8c_W^8 & 8c_W^8 & 4c_W^8 & 4c_W^4 s_W^4 & 4c_W^4 s_W^4 & 2c_W^4 s_W^4 & 2s_W^8 & s_W^8 \\ 0 & 0 & 4c_W^8 & 0 & 0 & 2c_W^4 s_W^4 & 0 & s_W^8 \\ 0 & 0 & 4c_W^4 & 0 & 0 & s_W^4 & 0 & 0 \\ 0 & 16c_W^4 & 4c_W^4 & 0 & 4s_W^4 & s_W^4 & 0 & 0 \\ 0 & 0 & 4 & 0 & 0 & 1 & 0 & 0 \\ 0 & 16 & 4 & 0 & 4 & 1 & 0 & 0 \\ 0 & 0 & 8c_W^4 & 0 & 0 & (c_W^2 - s_W^2)^2 & 0 & 2s_W^4 \\ 32c_W^4 & 32c_W^4 & 16c_W^4 & -16c_W^2 s_W^2 & 4(c_W^2 - s_W^2)^2 & 1 - 8s_W^2 c_W^2 & 8s_W^4 & 4s_W^4 \\ 0 & 0 & 4 & 0 & 0 & 2 & 0 & 1 \\ 8 & 8 & 4 & 4 & 4 & 2 & 2 & 1 \end{pmatrix}$$

# 1D bounds on aQGCs

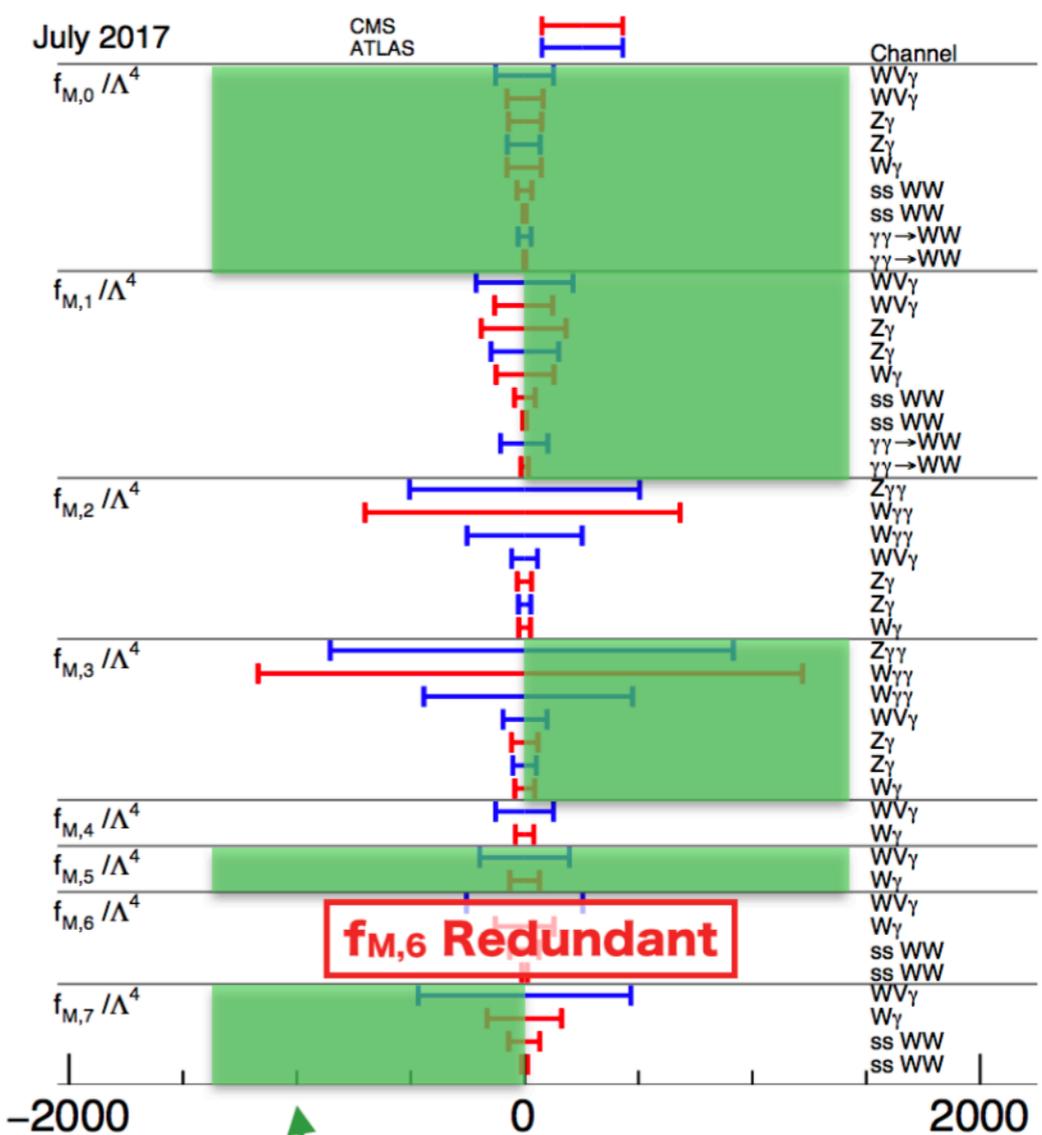
Cen Zhang & **SYZ**, arXiv:1808.00010

## Transversal coefficients, positivity



No UV completion

## Mixed coefficients, positivity

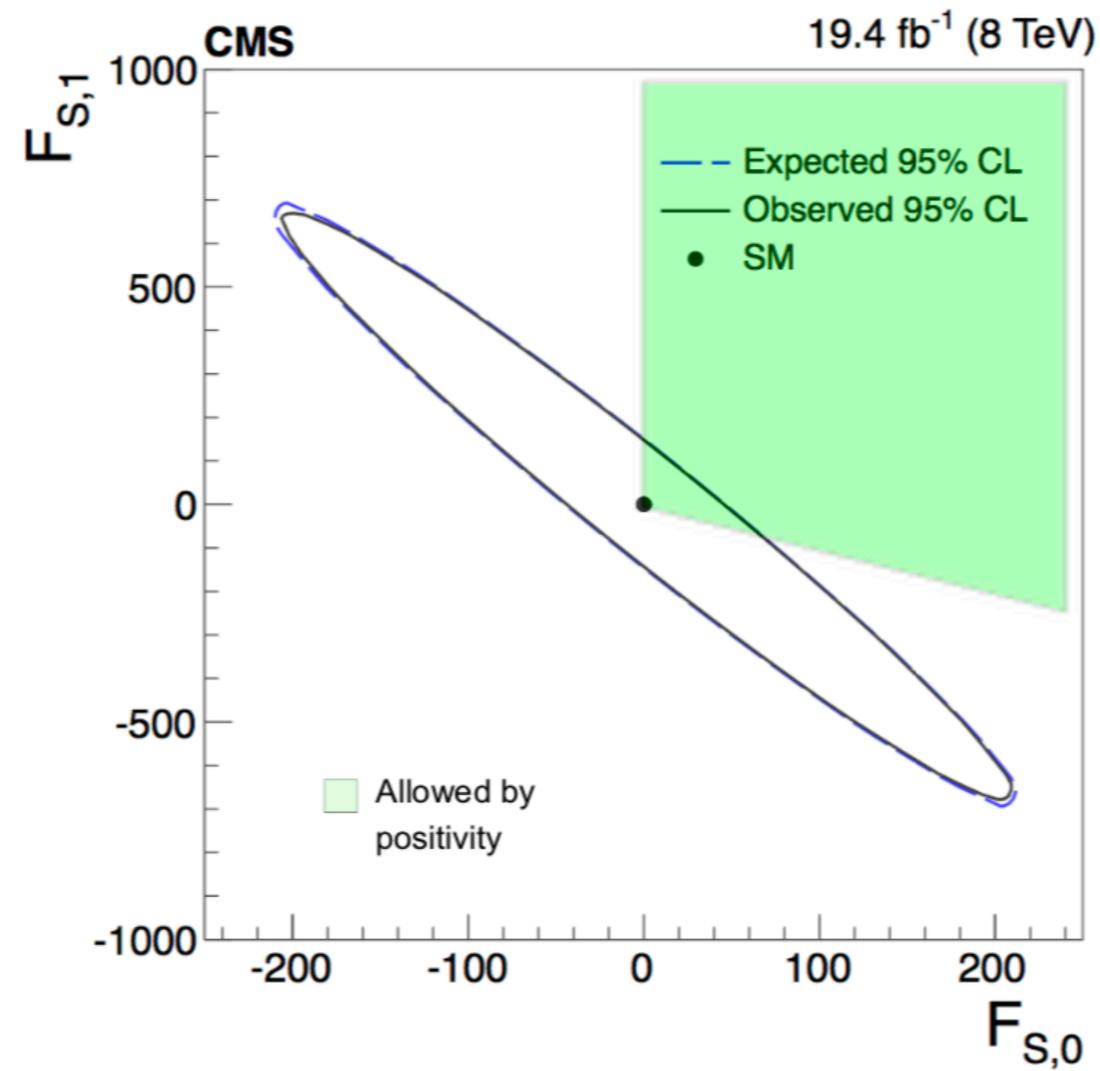


$f_{M,6}$  Redundant

No UV completion

# 2D bounds: an example

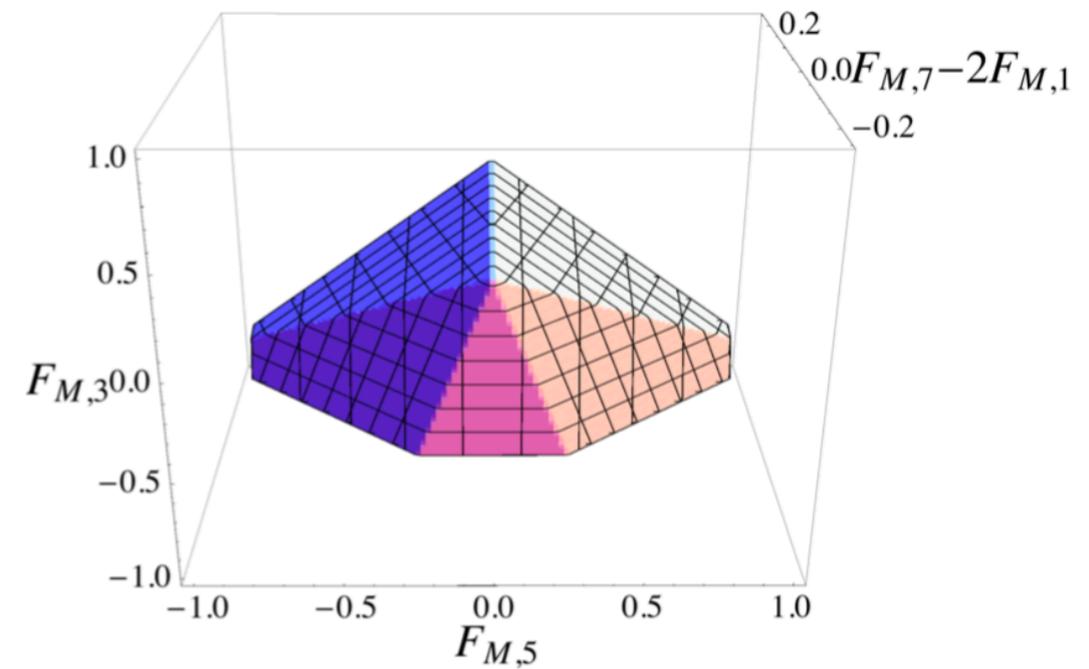
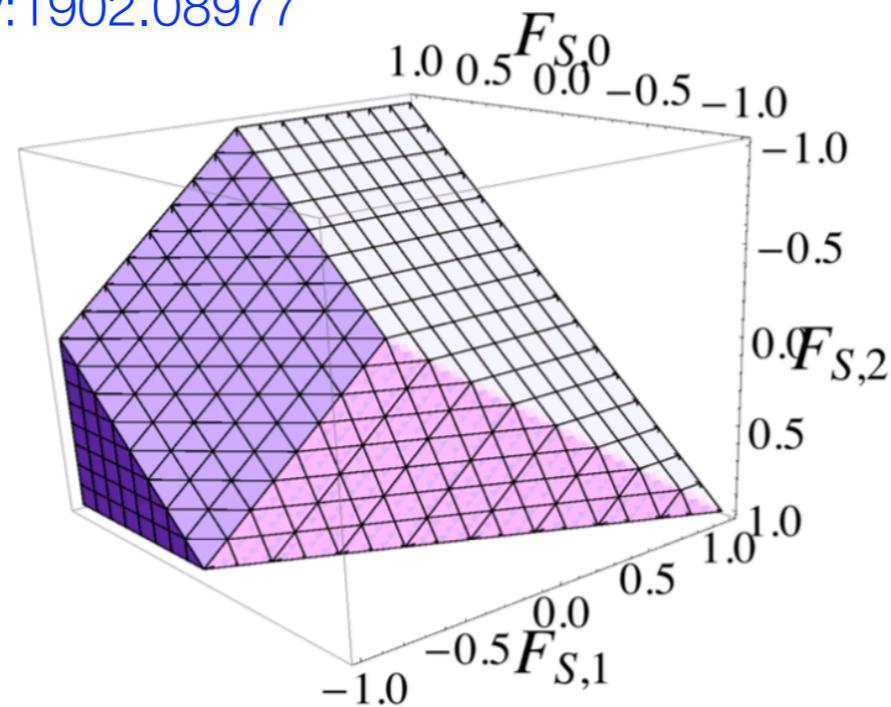
The case of  $O_{S0}$  and  $O_{S1}$



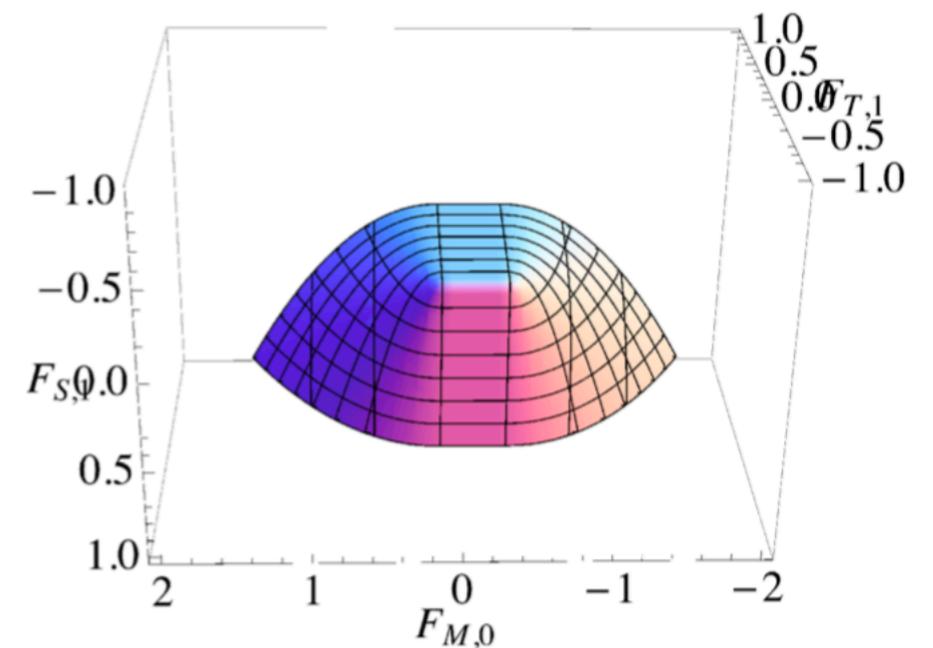
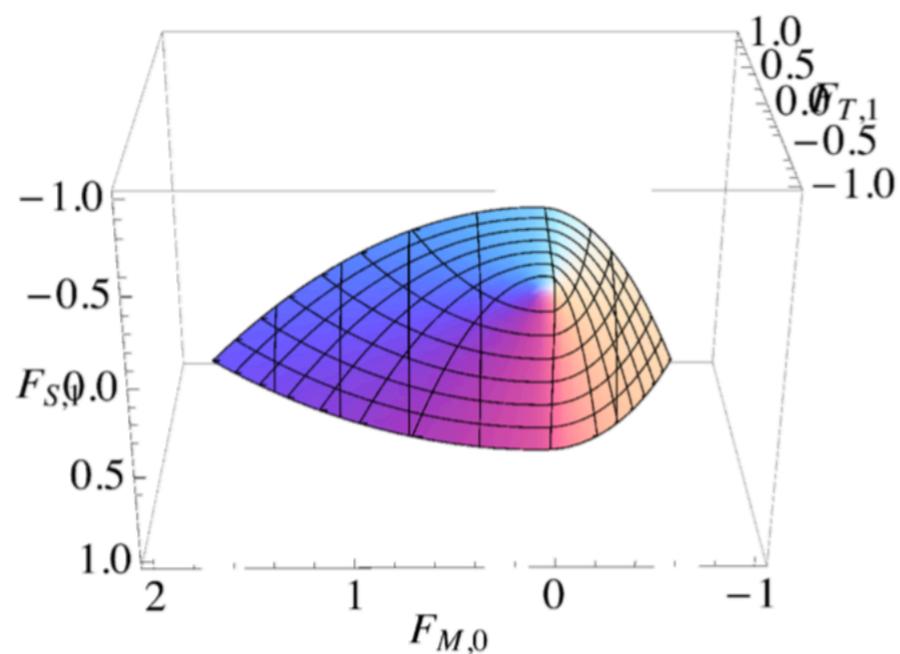
# Morphology of higher-D bounds

Bi, Zhang & **SYZ**, arXiv:1902.08977

Pyramids



Cones

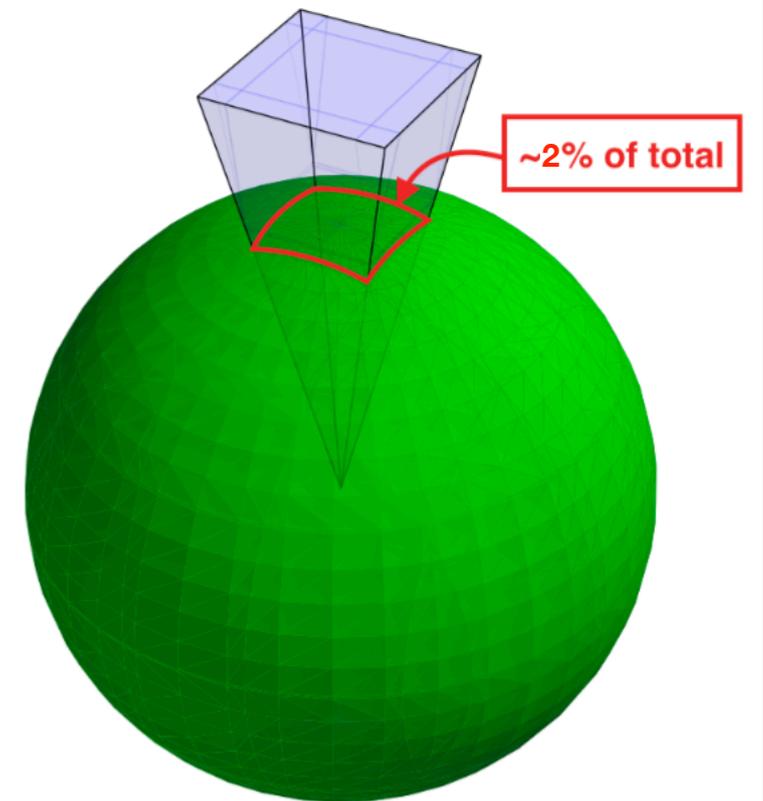


# General bounds on aQGCs

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All 18 parameters:

Randomly take points on 18D parameter space and count how many of them satisfy all the positivity bounds.



Only ~2% of the total parameter space admit an analytic UV completion!

# Summary

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Not all EFTs have a UV completion!

Positivity bounds: new constraints on coupling constants

Most of the VBS parameter space in SMEFT  
do not admit a UV completion.