

Comparison of direct and indirect constraints on anomalous tWb couplings

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Introduction

- The top quark is the heaviest of all known elementary particles and is expected to have large couplings with non-SM physics
- One of the possibilities is to study the possible **anomalous interactions** of the **Wtb** vertex

The most general, lowest-dimension, CP-conserving Lagrangian for the Wtb vertex has the following form:

$$\mathcal{L} = \frac{g}{\sqrt{2}} \bar{b} \gamma^\mu (f_V^L P_L + f_V^R P_R) t W_\mu^- - \frac{g}{\sqrt{2}} \bar{b} \frac{\sigma^{\mu\nu} \partial_\nu W_\mu^-}{M_W} (f_T^L P_L + f_T^R P_R) t + \text{h.c.},$$

where $P_{L,R} = (1 \mp \gamma_5)/2$, $\sigma_{\mu\nu} = i(\gamma_\mu \gamma_\nu - \gamma_\nu \gamma_\mu)/2$, g is the coupling constant of the weak interaction, the form factor f_V^L (f_V^R) represents the left-handed (right-handed) vector coupling, and f_T^L (f_T^R) represents the left-handed (right-handed) tensor coupling.

The SM has the following set of coupling values:

$$f_V^L = V_{tb}, f_V^R = f_T^L = f_T^R = 0.$$

Experimental Data

Experimental data on anomalous Wtb couplings can be divided into two types:

- direct constraints (e.g. from single top quark production)
- indirect constraints from B-physics

In B-physics the following three channels give the most stringent constraints:

- Data on $B_{d,s} - \bar{B}_{d,s}$ oscillations
- The branching ratio of the $\bar{B} \rightarrow X_s \gamma$ decay
- The branching ratio of the $\bar{B} \rightarrow X_s \mu^+ \mu^-$ decay

Scenarios

We considered the following scenarios when two or three couplings can have non-SM values:

$$(f_V^L; f_V^R)$$

$$(f_V^L; f_T^L)$$

$$(f_V^L; f_T^R)$$

$$(f_V^L; f_T^L; f_T^R)$$

With these scenarios we compared the results from indirect constraints obtained from B-physics and direct constraints obtained from single top quark production [1]

 [1] V. Khachatryan *et al.* [CMS Collaboration], JHEP **1702**, 028 (2017)

Theoretical framework for B-physics

In case of B-physics the constraints on anomalous couplings can be obtained through the corresponding constraints on the **non-SM additions to the Wilson coefficients**. We used the scheme proposed in [Drobnak et al.](#)

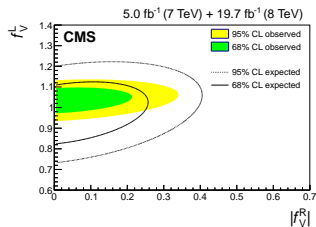
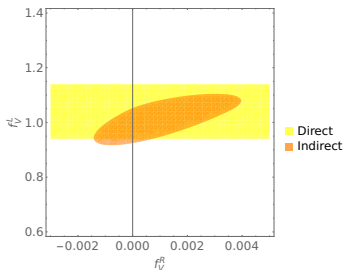
[10.1016/j.nuclphysb.2011.10.004](#), where in the Lagrangian at the scale $\mu \sim m_b$ of the B-meson decays the Wilson coefficients have the form $C_i = C_i^{SM} + \delta C_i$:

$$\begin{aligned} \mathcal{L} = \mathcal{L}_{QCD \times QED} &+ \frac{4G_F}{\sqrt{2}} \left[\sum_{i=1}^2 C_i \left(V_{ub}V_{us}^* \mathcal{O}_i^{(u)} + V_{cb}V_{cs}^* \mathcal{O}_i^{(c)} \right) \right] + \\ &+ \frac{4G_F}{\sqrt{2}} V_{tb}V_{ts}^* \sum_{i=3}^{10} C_i \mathcal{O}_i, \end{aligned}$$

where the first term consists of the kinetic terms of the light SM particles and their QCD and QED interactions

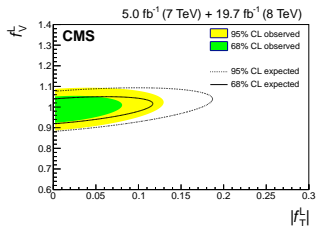
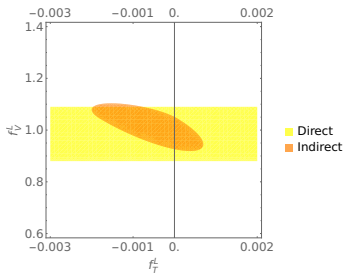
Results: constraints on anomalous tWb couplings

$$f_V^L - f_V^R$$



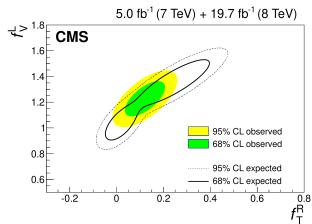
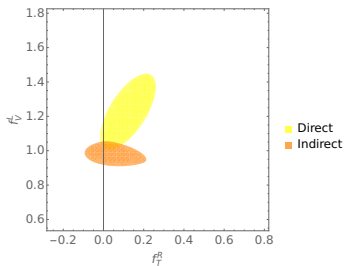
Results: constraints on anomalous tWb couplings

$$f_V^L - f_T^L$$



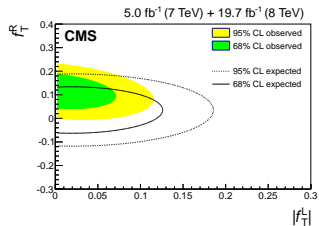
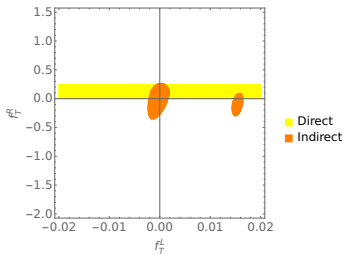
Results: constraints on anomalous tWb couplings

$$f_V^L - f_T^R$$



Results: constraints on anomalous tWb couplings

$$f_V^L - f_T^R - f_T^L$$



Conclusions

- Indirect constraints from B-physics in most scenarios are more rigid than direct constraints from single top quark production
- So far no deviations from the Standard Model have been observed