## CP violation in multi-body charmless b-hadron decays at LHCb

### Adam Morris, on behalf of the LHCb collaboration

Aix Marseille Univ, CNRS/IN2P3, CPPM

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Adam Morris (CPPM)

Multi-body charmless CPV

EPS-HEP 2019, Ghent

### Charmless *b*-hadron decays

- Tree  $b \rightarrow u$ 
  - Cabibbo suppression  $(V_{ub})$

- Penguin  $b \rightarrow d$  or  $b \rightarrow s$ 
  - Loop-level suppression
  - Sensitive to new particles in the loop
- Similar magnitude of tree & penguin contributions
  - Relative weak phase: interference  $\rightarrow$  *CP* violation
- Rich resonant structure in multi-body decays
  - Strong-phase differences in interference between resonances  $\rightarrow$  enhanced CP violation





### Charmless three-body B decays



- Rich resonant structure warrants amplitude analysis to measure *CP* violation in different regions of the phase space
- $B^{\pm} \rightarrow \pi^{\pm}\pi^{+}\pi^{-}$  covered in Jeremy Dalseno's talk
  - LHCb-PAPER-2019-017, LHCb-PAPER-2019-018

This talk:

- Amplitude analysis of  $B^0_s 
  ightarrow K^0_S K^\pm \pi^\mp$ 
  - JHEP 06 (2019) 114
- Amplitude analysis of  $B^\pm\!
  ightarrow\pi^\pm K^+K^-$ 
  - arXiv:1905.09244

### Charmless four-body *b*-baryon decays



- CP violation in baryons not yet observed
- Potential for large CP-violating effects in multi-body charmless b-baryon decays

This talk:

- Measurements of  $\Delta \mathcal{A}^{CP}$  in charmless four-body  $\Lambda_b^0$  and  $\Xi_b$  decays
  - arXiv:1903.06792

 $B_s^0 
ightarrow K_S^0 K^{\pm} \pi^{\mp}$ 

### Background



- First observed by LHCb in  $1 \text{ fb}^{-1}$  from 2011 (JHEP 10 (2013) 143)
- Branching fraction improved in  $3 \, \text{fb}^{-1}$  from 2011+12 (JHEP 11 (2017) 027)
- Specific intermediate states studied
  - $B_s^0 \to K^{*\pm}K^{\mp}$  (New J. Phys. 16 (2014) 123001)  $B_s^0 \to K^{*0}K_s^0$  (JHEP 01 (2016) 012)
- Potential for future time-dependent CP violation measurement with larger datasets



 $B_s^0 \to K_S^0 K^{\pm} \pi^{\mp}$  JHEP 06 (2019) 114

### Introduction



- First amplitude analysis of  $B^0_s o K^0_S K^\pm \pi^\mp$ 
  - Untagged and time-integrated
  - Simultaneous amplitude fit of two final states
  - Novel approach

• 431 
$$\overleftrightarrow{B_s^0} \to K_S^0 K^+ \pi^- + 489 \, \overleftrightarrow{B_s^0} \to K_S^0 K^- \pi^+$$

- Run 1 dataset:  $3 \text{ fb}^{-1}$  from 2011+12
- Published as JHEP 06 (2019) 114

### Amplitude model

- Both  $B_s^0$  and  $\overline{B}_s^0$  can decay to each final state, although not necessarily with the same amplitude  $A_f \neq \overline{A}_f$
- Untagged analysis means  $B_s^0$  and  $\overline{B}_s^0$  cannot be distinguished
- Fit for effective amplitude that is a combination of  $A_f$  and  $\overline{A}_f$
- $K^+K_S^0$  resonances e.g.  $a_2(1320)^+$  considered but not seen in fit
- $K\pi$  P-wave and D-wave modelled with Breit–Wigners
- $K\pi$  S-wave modelled with the LASS lineshape
  - Combines  $K_0^*(1430)$  and non-resonant  $K\pi$
  - Possible to disentangle later when calculating  $\mathcal{B}(B^0_s o K^*_0(1430)K)$







 $K^0_S K^+ \pi^-$ 



### Fit results

	$K^0_S K^-$	$^{+}\pi^{-}$	$\kappa^0_{s}\kappa^-\pi^+$				
F	Resonance	Frac. (%)	Resonance	Frac. (%)			
ŀ	<*(892) <sup>_</sup>	$15.6\pm1.5$	K*(892) <sup>+</sup>	$13.4\pm2.0$			
ŀ	$\zeta_0^*(1430)^-$	$30.2 \pm 2.6$	$K_0^*(1430)^+$	$28.5\pm3.6$			
ŀ	$(\sqrt[8]{2})^{*}(1430)^{-}$	$2.9\pm1.3$	$K_{2}^{*}(1430)^{+}$	$5.8 \pm 1.9$			
ŀ	<* (892) <sup>0</sup>	$13.2\pm2.4$	$\overline{K^{*}}(892)^{0}$	$19.2\pm2.3$			
ŀ	$\zeta_0^*(1430)^0$	$33.9 \pm 2.9$	$\overline{K}_{0}^{*}(1430)^{0}$	$27.0 \pm 4.1$			
ŀ	$\zeta_{2}^{*}(1430)^{0}$	$5.9\pm4.0$	$\overline{K}_{2}^{*}(1430)^{0}$	$7.7\pm2.8$			
NB:	uncertainties a	re statistical only	-				

• Fit fractions for each resonance and its conjugate are consistent, hence no significant *CP* violation observed

Sources of systematics:

- Mismodelling in mass fit
- Efficiency and background models
- Fit bias
- Fixed parameters
- Amplitude model





### Branching fractions



Flavour-averaged fit fractions converted to branching fractions for the quasi-two-body modes First observations of  $B_s^0 \rightarrow K_0^*(1430)K$  modes

$$\begin{split} \mathcal{B}(B^0_s \to K^*(892)^{\pm} K^{\mp}) &= (18.6 \pm 1.2 \pm 0.8 \pm 4.0 \pm 2.0) \times 10^{-6} \\ \mathcal{B}(B^0_s \to K^*_0(1430)^{\pm} K^{\mp}) &= (31.3 \pm 2.3 \pm 0.7 \pm 25.1 \pm 3.3) \times 10^{-6} \\ \mathcal{B}(B^0_s \to K^*_2(1430)^{\pm} K^{\mp}) &= (10.3 \pm 2.5 \pm 1.1 \pm 16.3 \pm 1.1) \times 10^{-6} \\ \mathcal{B}(B^0_s \to \widetilde{K}^*(892)^0 \widetilde{K}^{0}) &= (19.8 \pm 2.8 \pm 1.2 \pm 4.4 \pm 2.1) \times 10^{-6} \\ \mathcal{B}(B^0_s \to \widetilde{K}^*_0(1430)^0 \widetilde{K}^{0}) &= (33.0 \pm 2.5 \pm 0.9 \pm 9.1 \pm 3.5) \times 10^{-6} \\ \mathcal{B}(B^0_s \to \widetilde{K}^*_2(1430)^0 \widetilde{K}^{0}) &= (16.8 \pm 4.5 \pm 1.7 \pm 21.2 \pm 1.8) \times 10^{-6} \end{split}$$

Uncertainties:  $\pm$  stat  $\pm$  syst  $\pm$  model  $\pm$  norm

"norm" refers to uncertainty on  ${\cal B}(B^0_s o K^0 K^\pm \pi^\mp)$ 



Branching fractions of non-resonant modes:

$$\begin{split} \mathcal{B}(B^0_s \to (\overleftarrow{K}^{!0}\pi^{\pm})_{\rm NR}K^{\mp}) &= (11.4 \pm 0.8 \pm 0.2 \pm 9.2 \pm 1.2 \pm 0.5) \times 10^{-6} \\ \mathcal{B}(B^0_s \to (K^{\mp}\pi^{\pm})_{\rm NR}\overleftarrow{K}^{!0}) &= (12.1 \pm 0.9 \pm 0.3 \pm 3.3 \pm 1.3 \pm 0.5) \times 10^{-6} \end{split}$$

Uncertainties:  $\pm$  stat  $\pm$  syst  $\pm$  model  $\pm$  norm  $\pm$  eff. range

Fifth uncertainty related to proportion of the  $(K\pi)_0^*$  component due to the effective range part of the LASS lineshape.

## $B^{\pm} \rightarrow \pi^{\pm} K^{+} K^{-}$

### Background

### Previously studied by LHCb (Phys. Rev. D 90 (2014) 112004)

- Binned model-independent analysis
- Total  $\mathcal{A}^{CP} = -0.123 \pm 0.017 \pm 0.012 \pm 0.007$
- Regions of phase space with much larger  $\mathcal{A}^{CP}$







### Introduction



- First amplitude analysis of  $B^\pm\!
  ightarrow\pi^\pm K^+K^-$ 
  - $\pi^{\pm} K^{\mp}$  resonances:  $K^{*}(892)^{0}$ ,  $K_{0}^{*}(1430)^{0}$
  - Single-pole form factor to describe non-resonant  $\pi^{\pm}K^{\mp}$
  - $K^+K^-$  resonances:  $\phi(1020), f_2(1270), \rho(1450)^0$
  - Dedicated  $\pi\pi \leftrightarrow KK$  rescattering amplitude
- Run 1 dataset:  $3 \text{ fb}^{-1}$  from 2011+12
- Candidates in signal region: 2052  $B^+$ , 1566  $B^-$
- Submitted to PRL



### Non-resonant single-pole form factor



Proposed by Alvarenga Nogueira et al. (Phys. Rev. D 92 (2015) 054010)

$$\mathcal{A}_{\mathsf{source}} = \left(1 + rac{s}{\Lambda^2}
ight)^{-1}$$

• 
$$s=m_{\pi^\pm K^\mp}^2$$
  
•  $\Lambda=1\,{
m GeV}/c^2$  sets the scale for the energy dependence

### $\pi\pi \leftrightarrow \textit{KK}$ rescattering amplitude

Based on Pelaez and Yndurain (Phys. Rev. D 71 (2005) 074016)

$$\mathcal{A}_{ ext{rescattering}} = \left(1 + rac{s}{\Lambda^2}
ight)^{-1} \sqrt{1 - 
u^2} e^{2i\delta}$$

Inelasticity,  $\nu$ :

$$u = 1 - \left(\epsilon_1 \frac{k_2}{\sqrt{s}} + \epsilon_2 \frac{k_2^2}{s}\right) \frac{M'^2 - s}{s}$$

Phase shift  $\delta$ :

$$\cot \delta = C_0 rac{(s - M_s^2)(M_f^2 - s)}{M_f^2 \sqrt{s}} rac{|k_2|}{k_2^2}$$

•  $s = m_{K^+K^-}^2$ •  $k_2 = \frac{1}{2}\sqrt{2 - 4m_K}$ •  $m_K = 0.495 \,\text{GeV}/c^2$ 

• 
$$M' = 1.5 \,\mathrm{GeV}/c^2$$

• 
$$M_s = 0.92 \, {
m GeV}/c^2$$

• 
$$M_f = 1.32 \, {
m GeV}/c^2$$

• 
$$\epsilon_1 = 2.4$$

• 
$$C_0 = 1.3$$



### $B^{\pm} \rightarrow \pi^{\pm} K^{+} K^{-}$ arXiv:1905.09244

### The $\pi^\pm {\sf K}^\mp$ spectrum





• Single-pole non-resonant is dominant contribution ( $\sim$  32%)

•  $\rho(1450)^0 - f_2(1270)$  destructive interference at high  $m_{\pi^{\pm}K^{\mp}}^2$ 

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### $B^{\pm} \rightarrow \pi^{\pm} K^{+} K^{-}$ arXiv:1905.09244

### The $K^+K^-$ spectrum

 $B^+$ 

 $B^{-}$ 



•  $ho(1450)^0 \sim 30\%$  contribution

- Unexpectedly large for  $K^+K^-$
- Further analysis with more data needed
- $\pi\pi\leftrightarrow$  *KK*  $\sim$  16% contribution
  - Large CP asymmetry



### Results



Sources of systematics:

- Mismodelling in mass fit
- Efficiency and background models
- Fit bias
- Fixed parameters

•  $\pi\pi \leftrightarrow KK$  rescattering: largest ever *CP* asymmetry for a single amplitude to date

- No significant CP asymmetry observed in the other components
- $\phi(1020)$  contribution not significant



# Four-body $\Lambda_b^0$ and $\Xi_b^0$ decays



Previous LHCb results on charmless four-body *b*-baryon decays:

- Branching fractions (JHEP 02 (2018) 098)
- Triple-product asymmetries: (Nature Phys. 13 (2017) 391-396, JHEP 08 (2018) 039)
- 3.3 $\sigma$  evidence for *CP* violation in  $\Lambda_b^0 \to p\pi^-\pi^+\pi^-$  from triple-products

•  $\mathcal{A}^{CP}$  using  $\Lambda_b^0/\Xi_b^0$  and  $\overline{\Lambda}_b^0/\overline{\Xi}_b^0$  yields obtained from fitting m(phh'h'')

$$\mathcal{A}^{CP} \equiv \frac{\Gamma(X_b^0 \to f) - \Gamma(\overline{X}_b^0 \to \overline{f})}{\Gamma(X_b^0 \to f) + \Gamma(\overline{X}_b^0 \to \overline{f})}$$

- Complementary to triple-products
- Run 1 dataset:  $3 \text{ fb}^{-1}$  from 2011+12
- Submitted to EPJC

- Six decay modes studied:

  - $\Lambda_b^0 \rightarrow p\pi^-\pi^+\pi^ \Lambda_b^0 \rightarrow pK^-\pi^+\pi^ \Lambda_b^0 \rightarrow pK^-K^+\pi^ \Lambda_b^0 \rightarrow pK^-K^+K^-$

• 
$$\Xi_b^0 \rightarrow pK^-\pi^+\pi^-$$

• 
$$\Xi_b^0 \rightarrow p K^- \pi^+ K^-$$

- For the three most abundant decays, also study specific regions of phase space
  - Low two-body mass
  - Specific intermediate resonances
- $\Lambda_{h}^{0} \rightarrow \Lambda_{c}^{+} \pi^{-}$  and  $\Xi_{h}^{0} \rightarrow \Xi_{c}^{+} \pi^{-}$  control channels to cancel production and detection asymmetries

• 
$$\Delta \mathcal{A}^{CP} = \mathcal{A}^{CP}_{charmless} - \mathcal{A}^{CP}_{charm}$$



### Phase space regions





### Mass fits

- Simultaneous maximum likelihood fit to *b*-hadron candidates under each phh'h'' hypothesis Data split by:
  - Proton charge
  - Year of data-taking
  - Hardware trigger condition
- Fit model has components for:
  - Signal
  - Cross-feed  $(\pi K \text{ mis-ID})$
  - 4-body B-meson decays  $(p \pi \text{ and } p K \text{ mis-ID})$
  - 5-body b-hadron decays
  - Combinatorial background







### Sample fit projections



*інср* 

- Total of 18  $\Delta A^{CP}$  measurements
- No indication of significant CPV
- Statistical uncertainty dominates
- $\sim 5 \times$  larger yields in Run 2 data



## Summary and conclusions

### Summary



- $B^0_s 
  ightarrow K^0_S K^\pm \pi^\mp$ 
  - No evidence of CP violation
  - Updated quasi-two-body branching fractions
  - First observation of  $B_s^0 
    ightarrow K_0^*(1430) K$  modes

 $B^\pm\!\to\pi^\pm K^+K^-$ 

- $\mathcal{A}^{CP} = (-66.4 \pm 4.2) \,\%$  in  $\pi\pi \leftrightarrow \textit{KK}$  rescattering term
- Largest CP asymmetry in a single amplitude

Four-body *b*-baryon decays

- 18  $\Delta A^{CP}$  measurements
- No evidence of CP violation





- Multi-body charmless b-hadron decays are an important area for studying CP violation
- LHCb Run 2 data and upgrade will provide improved results

## Backup slides

### Track types at LHCb





- Long tracks pass through all tracking stations
- Downstream tracks pass through the TT and T
  - A and  $K_S^0$  can decay outside VELO

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### Multi-body charmless CPV

Run 1 performance paper: LHCb-DP-2014-002

 $B_s^0 
ightarrow K_S^0 K^{\pm} \pi^{\mp}$ 





- Simultaneous fit to  $m(K_S^0 K^{\pm} \pi^{\mp})$  in 24 data categories:
  - 3 data-taking periods (change in trigger efficiency during 2012)
  - 4 final states (including  $\pi \leftrightarrow K$  mis-ID)
  - 2  $K_S^0$  reconstruction categories: decay inside ("long") or outside ("downstream") VELO
- Signal and cross-feed: sum of two Crystal Ball functions
- Combinatorial background: exponential function

### $B_s^0$ mass fit



Final	$K_S^0$	Sample	$B_s^0$ signal		Combinato	rial	Cross-feed		
state	category		Full range $2.5 \sigma$		Full range	$2.5\sigma$	Full range	$2.5\sigma$	
		2011	$73.6\pm10.6$	72.1	$108.3\pm15.1$	22.1	$8.9\pm2.8$	1.7	
	downstream	2012a	$48.2\pm8.6$	45.7	$70.1 \pm 12.1$	14.3	$7.3\pm3.8$	1.1	
$\kappa^0 \kappa^+ \pi^-$		2012b	$135.3\pm13.6$	130.0	$87.4 \pm 13.8$	17.9	$17.0\pm5.6$	3.1	
κ <sub>s</sub> κ <sup>+</sup> π		2011	$\bar{76.2 \pm 9.8}$	74.6	$\bar{44.1 \pm 9.8}$	8.4	$\overline{8.2\pm1.7}$	1.8	
	long	2012a	$38.5\pm~7.7$	36.8	$58.8 \pm 11.2$	11.2	$7.8\pm1.8$	0.9	
		2012b	$73.5\pm10.6$	71.9	$71.7 \pm 13.1$	13.6	$15.9\pm2.5$	1.7	
	total			431.1		87.5		10.3	
		2011	$72.8\pm10.3$	71.4	$78.9 \pm 12.7$	16.1	$8.2\pm2.4$	1.3	
	downstream	2012a	$68.8\pm~9.6$	65.2	$46.2\pm9.9$	9.5	$7.0\pm3.4$	1.2	
$\kappa^0 \kappa^- \pi^+$		2012b	$165.1\pm15.2$	158.6	$104.1\pm15.0$	21.3	$17.3\pm5.8$	2.9	
n <sub>š</sub> n π'		2011	$\bar{77.3}\pm \bar{9.8}$	75.7	$\bar{39.0 \pm 10.2}$	7.4	$9.6\pm1.7$	1.4	
	long	2012a	$40.3\pm8.1$	38.5	$58.9 \pm 11.9$	11.2	$8.6\pm1.8$	0.7	
		2012b	$81.7\pm10.4$	80.0	$50.1 \pm 12.3$	9.5	$15.0\pm2.5$	1.4	
	total			489.4		75.0		8.9	

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### **Systematics**



				Fit fraction (%) uncertainties							
Resonance	Yields	Bkg.	Eff.	Fit bias	Add. res.	Fixed par.	Alt. model	Method	Total		
$K^{*}(892)^{-}$	0.2	0.2	0.5	0.2	-	0.7	5.4	3.1	6.3		
$K_0^*(1430)^-$	0.1	0.2	0.6	0.3	0.1	2.1	22.0	2.9	22.3		
$K_2^*(1430)^-$	0.1	0.1	0.3	0.6	0.1	1.8	2.2	0.2	2.9		
$K^{*}(892)^{0}$	0.2	0.2	0.4	0.9	-	0.3	7.0	2.0	7.4		
$K_0^*(1430)^0$	0.2	0.3	0.9	0.4	0.1	4.4	3.3	1.3	5.7		
$K_2^*(1430)^0$	0.1	0.3	0.7	1.3	0.2	4.4	3.6	1.0	6.0		
$K^{*}(892)^{+}$	0.4	0.1	0.6	0.5	0.1	0.7	1.1	0.7	1.8		
$K_0^*(1430)^+$	0.5	0.4	0.7	0.8	0.2	6.4	13.0	4.5	15.2		
$K_{2}^{*}(1430)^{+}$	0.1	0.2	0.4	0.2	0.1	4.1	4.5	3.2	6.9		
$\overline{K}^{*}(892)^{0}$	0.4	0.3	0.4	0.2	0.2	0.5	3.0	7.9	8.5		
$\overline{K}_{0}^{*}(1430)^{0}$	0.4	0.4	0.6	0.8	0.7	0.9	3.9	5.4	6.8		
$\overline{K}_{2}^{*}(1430)^{0}$	0.1	0.2	0.4	0.8	0.1	1.0	5.5	2.7	6.3		

## $B^{\pm} \rightarrow \pi^{\pm} K^{+} K^{-}$

### $B^{\pm} \rightarrow \pi^{\pm} K^{+} K^{-}$

### *B*-factory results on $B^{\pm} \rightarrow \pi^{\pm} K^{+} K^{-}$

LHCb  $\mathcal{A}^{CP}$  results (reminder):

- Total  $\mathcal{A}^{CP}$ : -0.123 ± 0.017 ± 0.012 ± 0.007 (Phys. Rev. D 90 (2014) 112004)
- $\pi\pi \leftrightarrow \mathcal{KK} \ \mathcal{A}^{CP}$ : -66.4  $\pm$  3.8  $\pm$  1.9 (arXiv:1905.09244)
- Belle  $\mathcal{A}^{CP}$  results:
  - Total A<sup>CP</sup>: −0.170 ± 0.073 ± 0.017 (Phys. Rev. D 96 (2017) 031101)
- Branching fractions:
  - Belle:  $(5.38 \pm 0.40 \pm 0.35) \times 10^{-6}$ (Phys. Rev. D 96 (2017) 031101)
  - BaBar: (5.0  $\pm$  0.5  $\pm$  0.5)  $\times$  10^{-6}

(Phys. Rev. Lett. 99 (2007) 221801)







### Projections





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### Projections





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Contribution	Fit Fraction(%)	$\mathcal{A}^{CP}(\%)$	Magnitude $(B^+/B^-)$	Phase[°] $(B^+/B^-)$
$K^{*}(892)^{0}$	$7.5\pm0.6\pm0.5$	$+12.3 \pm 8.7 \pm 4.5$	$0.94 \pm 0.04 \pm 0.02$	0 (fixed)
			$1.06 \pm 0.04 \pm 0.02$	0 (fixed)
$K_0^*(1430)^0$	$4.5\pm0.7\pm1.2$	$+10.4 \pm 14.9 \pm 8.8$	$0.74 \pm 0.09 \pm 0.09$	$-176\pm10\pm16$
			$0.82 \pm 0.09 \pm 0.10$	$136\pm11\pm21$
Single pole	$32.3\pm1.5\pm4.1$	$-10.7 \pm 5.3 \pm 3.5$	$2.19 \pm 0.13 \pm 0.17$	$-138\pm7\pm5$
			$1.97 \pm 0.12 \pm 0.20$	$166\pm6\pm5$
$ ho(1450)^{0}$	$30.7\pm1.2\pm0.9$	$-10.9 \pm ~4.4 \pm ~2.4$	$2.14 \pm 0.11 \pm 0.07$	$-175\pm10\pm15$
			$1.92 \pm 0.10 \pm 0.07$	$140\pm13\pm20$
$f_2(1270)$	$7.5\pm0.8\pm0.7$	$+26.7 \pm 10.2 \pm 4.8$	$0.86 \pm 0.09 \pm 0.07$	$-106\pm11\pm10$
			$1.13 \pm 0.08 \pm 0.05$	$-128\pm11\pm14$
Rescattering	$16.4\pm0.8\pm1.0$	$-66.4 \pm \ 3.8 \pm \ 1.9$	$1.91 \pm 0.09 \pm 0.06$	$-56\pm12\pm18$
			$0.86 \pm 0.07 \pm 0.04$	$-81\pm14\pm15$
$\phi$ (1020)	$0.3\pm0.1\pm0.1$	$+9.8 \pm 43.6 \pm 26.6$	$0.20 \pm 0.07 \pm 0.02$	$-52\pm23\pm32$
			$0.22 \pm 0.06 \pm 0.04$	$107\pm33\pm41$



# Four-body $\Lambda_b^0$ and $\Xi_b^0$ decays

Triple-product asymmetries in  $\Lambda_b^0 \rightarrow p \pi^- \pi^+ \pi^-$ 

$$\begin{split} C_{\widehat{T}} &= \vec{p}_{p} \cdot (\vec{p}_{h_{1}^{-}} \times \vec{p}_{h_{2}^{+}}) \text{ for } \Lambda_{b}^{0} \\ \overline{C}_{\widehat{T}} &= \vec{p}_{\overline{p}} \cdot (\vec{p}_{h_{1}^{+}} \times \vec{p}_{h_{2}^{-}}) \text{ for } \overline{\Lambda_{b}^{0}} \end{split}$$

$$A_{\widehat{T}}(C_{\widehat{T}}) = \frac{N(C_{\widehat{T}} > 0) - N(C_{\widehat{T}} < 0)}{N(C_{\widehat{T}} > 0) + N(C_{\widehat{T}} < 0)}$$
$$\overline{A}_{\widehat{T}}(\overline{C}_{\widehat{T}}) = \frac{\overline{N}(-\overline{C}_{\widehat{T}} > 0) - \overline{N}(-\overline{C}_{\widehat{T}} < 0)}{\overline{N}(-\overline{C}_{\widehat{T}} > 0) + \overline{N}(-\overline{C}_{\widehat{T}} < 0)}$$

The *P*- and *CP*-violating observables are defined as

$$\begin{aligned} \mathbf{a}_{P}^{\widehat{T}\text{-odd}} &= \frac{1}{2} \left( A_{\widehat{T}} + \overline{A}_{\widehat{T}} \right) \\ \mathbf{a}_{CP}^{\widehat{T}\text{-odd}} &= \frac{1}{2} \left( A_{\widehat{T}} - \overline{A}_{\widehat{T}} \right) \end{aligned}$$





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Multi-body charmless CPV



The difference of *CP*-asymmetries measured for the charmless modes and for the control channels results in  $\Delta A^{CP}$  measurements. For each observable, the choice of the control channel is aiming at cancelling at first order production and detection asymmetries.

Charmless mode	Control channel
$\Lambda^0_b  o p \pi^- \pi^+ \pi^-$	$\Lambda^0_b  ightarrow (\Lambda^+_c  ightarrow  ho\pi^-\pi^+)\pi^-$
$\Lambda^0_b  o  ho {\cal K}^- \pi^+ \pi^-$	$arLambda_b^0  o (arLambda_c^+  o  ho {\cal K}^- \pi^+) \pi^-$
$\Lambda^0_b  o  ho {\cal K}^- {\cal K}^+ \pi^-$	$\Lambda^0_b  ightarrow (\Lambda^+_c  ightarrow  ho \pi^- \pi^+) \pi^-$
$\Lambda^0_b  ightarrow  ho K^- K^+ K^-$	$arLambda_b^0  o (arLambda_c^+  o  ho {\cal K}^- \pi^+) \pi^-$
$arepsilon_b^0  o  ho {\cal K}^- \pi^+ \pi^-$	$arepsilon_b^0  o (arepsilon_c^+  o  ho {\cal K}^- \pi^+) \pi^-$
$arepsilon_b^0  o  ho {\cal K}^- \pi^+ {\cal K}^-$	$arepsilon_b^0  ightarrow (arepsilon_c^+  ightarrow p {\it K}^- \pi^+) \pi^-$



Decay mode	Invariant-mass requirements (in $\mathrm{MeV}/c^2$ )
$\Lambda_b^0  ightarrow p \pi^- \pi^+ \pi^-$ low mass	$m(p\pi^-) < 2000$ and $m(\pi^+\pi^-) < 1640$
$arLambda_b^0  o {\it pa}_1(1260)^-$	419 $< m(\pi^+\pi^-\pi^+) < 1500$
$arLambda_b^0  o {\sf N}(1520)^0  ho^0$	$1078 < m(p\pi^-) < 1800$ and $m(\pi^+\pi^-) < 1100$
$\Lambda^0_b  ightarrow \Delta(1232)^{++} \pi^- \pi^-$	$1078 < m(p\pi^+) < 1432$



Decay mode	Invariant-mass requirements (in $\mathrm{MeV}\!/c^2)$
$\Lambda_b^0  ightarrow p K^- \pi^+ \pi^-$ low mass	$m( ho K^-) < 2000$ and $m(\pi^+\pi^-) < 1640$
$arLambda_b^0  o {\sf N}(1520)^0 {\sf K}^{*0}$	$1078 < m(p\pi^-) < 1800$ and $750 < m(\pi^+ {\cal K}^-) < 1100$
$\Lambda^0_b  ightarrow \Lambda(1520) ho^0$	$1460 < m( ho K^-) < 1580$ and $m(\pi^+\pi^-) < 1100$
$\Lambda^0_b  ightarrow \Delta$ (1232) <sup>++</sup> $K^-\pi^-$	$1078 < m(p\pi^+) < 1432$
$arLambda_b^0  o  ho {\cal K}_1(1410)^-$	$1200 < m({K^-}{\pi^+}{\pi^-}) < 1600$



Decay mode	Invariant-mass requirements (in $\mathrm{MeV}/c^2$ )
$\Lambda^0_b  o p K^- K^+ K^-$ low mass	$m( ho K^-) <$ 2000 and $m(K^+K^-) <$ 1675
$arLambda_b^{0}  ightarrow arLambda(1520) \phi$	$1460 < m( ho K^-) < 1600$ and $1005 < m(K^+K^-) < 1040$
$arLambda_{b}^{0} ightarrow( ho \mathcal{K}^{-})_{ ext{high-mass}}\phi$	$m( ho K^-) > 1600$ and $1005 < m(K^+K^-) < 1040$



*LHCb* ГНСр



*LHCb* ГНСр



*LHCb* ГНСр













Adam Morris (CPPM)





Adam Morris (CPPM)



*інср* 

### Systematics



Tracking detection efficiency

• Quantified separately for kaons  $(\sigma_K)$  and protons  $(\sigma_p)$ 

Trigger efficiency ( $\sigma_{\rm L0}$ )

- Difference in hardware-level (L0) trigger efficiency between oppositely-charged hadrons
- Production asymmetry ( $\sigma_{A_P}$ )
  - Difference in decay kinematics of signal and control channels  $\rightarrow$  incomplete cancellation
  - Estimated from measurement of  $\Lambda_b^0$  production asymmetry as a function of  $p_{\rm T}$  and  $\eta$  (Phys. Lett. B 774 (2017) 139)

PID calibration ( $\sigma_{\mathrm{PID}}$ )

• Finite size of calibration samples





Decay mode	А	bsolute	Total (%)			
	$\sigma_K$	$\sigma_{p}$	$\sigma_{ m L0}$	$\sigma_{ m PID}$	$\sigma_{A_P}$	
$\Lambda^0_b  ightarrow p \pi^- \pi^+ \pi^-$		0.20	0.06	0.42	0.28	0.54
$\Lambda^{0}_{b}  ightarrow p \pi^{-} \pi^{+} \pi^{-}$ low mass		0.16	0.06	0.36	0.28	0.49
$arLambda_b^0  o {\it pa}_1(1260)^-$		0.20	0.09	0.48	0.28	0.60
$arLambda_b^0  o {\sf N}(1520)^0  ho^0$	—	0.12	0.05	0.23	0.28	0.39
$\Lambda_b^0  ightarrow \Delta$ (1232) <sup>++</sup> $\pi^-\pi^-$		0.18	0.05	0.47	0.28	0.59



Decay mode	Absolute uncertainties (%)					Total (%)
	$\sigma_K$	$\sigma_{p}$	$\sigma_{ m L0}$	$\sigma_{ m PID}$	$\sigma_{A_P}$	
$\Lambda^0_b  ightarrow  ho K^- \pi^+ \pi^-$	0.17	0.20	0.06	0.41	0.24	0.55
$\Lambda^0_b  ightarrow p K^- \pi^+ \pi^-$ low mass	0.17	0.17	0.05	0.34	0.24	0.48
$arLambda_b^0  ightarrow arLambda(1520) ho^0$	0.12	0.12	0.04	0.36	0.24	0.49
$\Lambda^0_b  ightarrow \Delta$ (1232) <sup>++</sup> $K^-\pi^-$	0.22	0.19	0.05	0.48	0.24	0.61
$arLambda_b^0  o {\sf N}(1520)^0 {\sf K}^{*0}$	0.16	0.14	0.04	0.32	0.24	0.45
$arLambda_b^0  o  ho {\cal K}_1(1410)^-$	0.16	0.14	0.11	0.58	0.24	0.74





Decay mode	At	osolute	Total (%)			
	$\sigma_K$	$\sigma_{p}$	$\sigma_{ m L0}$	$\sigma_{ m PID}$	$\sigma_{\mathcal{A}_{\mathcal{P}}}$	
$\Lambda^0_b  ightarrow p K^- K^+ \pi^-$	—	0.21	0.06	0.40	0.55	0.72
$\Lambda^0_b  o p K^- K^+ K^-$	0.15	0.20	0.07	0.41	0.33	0.59
$\Lambda^0_b  ightarrow p K^- K^+ K^-$ low mass	0.16	0.17	0.05	0.37	0.33	0.55
$\Lambda^{0}_{b}  ightarrow \Lambda(1520)\phi$	0.11	0.10	0.05	0.30	0.33	0.34
$\Lambda^{0}_{b}  ightarrow (pK^{-})_{high-mass} \phi$	0.15	0.14	0.06	0.58	0.33	0.64
$arepsilon_b^0  o  ho {\cal K}^- \pi^+ \pi^-$	0.17	0.20	0.05	0.42	0.24	0.55
$arepsilon_b^0  o  ho K^- \pi^+ K^-$	0.15	0.20	0.05	0.41	0.55	0.73

### Results

$$\begin{split} & \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to p \pi^- \pi^+ \pi^-) = (+1.1 \pm 2.5 \pm 0.6) \,\% \\ & \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to p \mathcal{K}^- \pi^+ \pi^-) = (+3.2 \pm 1.1 \pm 0.6) \,\% \\ & \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to p \mathcal{K}^- \mathcal{K}^+ \pi^-) = (-6.9 \pm 4.9 \pm 0.8) \,\% \\ & \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to p \mathcal{K}^- \mathcal{K}^+ \mathcal{K}^-) = (+0.2 \pm 1.8 \pm 0.6) \,\% \\ & \Delta \mathcal{A}^{CP}(\Xi_b^0 \to p \mathcal{K}^- \pi^+ \pi^-) = (-17 \pm 11 \pm 1) \,\% \\ & \Delta \mathcal{A}^{CP}(\Xi_b^0 \to p \mathcal{K}^- \pi^+ \mathcal{K}^-) = (-6.8 \pm 8.0 \pm 0.8) \,\% \end{split}$$



### Results



$$\begin{array}{l} \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to p\pi^-\pi^+\pi^-) = (\pm 1.1 \pm 2.5 \pm 0.6) \,\% \\ \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to p\pi^-\pi^+\pi^-)_{\text{low mass}} = (\pm 3.7 \pm 4.1 \pm 0.5) \,\% \\ \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to pa_1(1260)^-) = (-1.5 \pm 4.2 \pm 0.6) \,\% \\ \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to N(1520)^0 \rho^0) = (\pm 2.0 \pm 4.9 \pm 0.4) \,\% \\ \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to \Delta(1232)^{++}\pi^-\pi^-) = (\pm 0.1 \pm 3.2 \pm 0.6) \,\% \\ \end{array}$$

 $\Delta A$ 

$$\begin{array}{c} \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to p \mathcal{K}^- \pi^+ \pi^-) = (+3.2 \pm 1.1 \pm 0.6) \,\% \\ \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to p \mathcal{K}^- \pi^+ \pi^-)_{\text{low mass}} = (+3.5 \pm 1.5 \pm 0.5) \,\% \\ \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to N(1520)^0 \mathcal{K}^{*0}) = (+5.5 \pm 2.5 \pm 0.5) \,\% \\ \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to \Lambda(1520) \rho^0) = (+0.6 \pm 6.0 \pm 0.5) \,\% \\ \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to \Delta(1232)^{++} \mathcal{K}^- \pi^-) = (+4.4 \pm 2.6 \pm 0.6) \,\% \\ \Delta \mathcal{A}^{CP}(\Lambda_b^0 \to p \mathcal{K}_1(1410)^-) = (+4.7 \pm 3.5 \pm 0.8) \,\% \end{array}$$



### Results



$$\Delta \mathcal{A}^{CP}(\Lambda_{b}^{0} \to pK^{-}K^{+}K^{-}) = (+0.2 \pm 1.8 \pm 0.6)\%$$

$$\Delta \mathcal{A}^{CP}(\Lambda_{b}^{0} \to pK^{-}K^{+}K^{-})_{\text{low mass}} = (+2.7 \pm 2.3 \pm 0.6)\%$$

$$\Delta \mathcal{A}^{CP}(\Lambda_{b}^{0} \to \Lambda(1520)\phi) = (+4.3 \pm 5.6 \pm 0.4)\%$$

$$\Delta \mathcal{A}^{CP}(\Lambda_{b}^{0} \to (pK^{-})_{\text{high-mass}}\phi) = (-0.7 \pm 3.3 \pm 0.7)\%$$

$$-60 -50 -40 -30 -20 -10 0 10$$

$$\Delta \mathcal{A}^{CP}[\%]$$