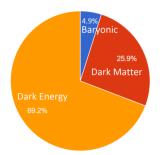
Unravelling Cosmic Acceleration with Gravitational Waves and Large-Scale Structure

Lucas Lombriser

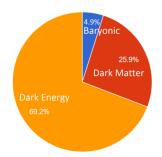
Département de Physique Théorique, Université de Genève

EPS-HEP2019, Ghent 12 Jul 2019



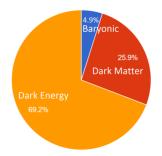


Cosmic acceleration at late times



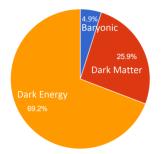
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- Best candidate:

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = M_{\rm Pl}^{-2}T_{\mu\nu}$$



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 → old and new aspects



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 → old and new aspects
- $\Lambda = 0 \& DE \text{ or MG}$?



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Horndeski scalar-tensor action (simplest MG)

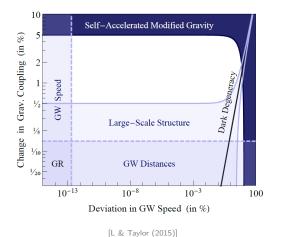
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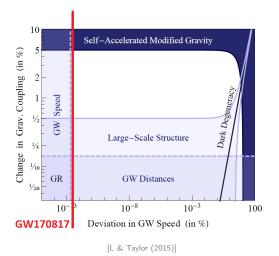
$$\begin{split} S &= \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} \left\{ G_2(\varphi,X) - G_3(\varphi,X) \Box \varphi \right. \\ &+ \left. G_4(\varphi,X) R + \frac{\partial G_4}{\partial X} \left[(\Box \varphi)^2 - (\nabla_\mu \nabla_\nu \varphi)^2 \right] \right. \\ &+ \left. G_5(\varphi,X) G_{\mu\nu} \nabla^\mu \nabla^\nu \varphi \right. \\ &- \left. \frac{1}{6} \frac{\partial G_5}{\partial X} \left[(\Box \varphi)^3 - 3 \Box \varphi (\nabla_\mu \nabla_\nu \varphi)^2 + 2 (\nabla_\mu \nabla_\nu \varphi)^3 \right] \right. \\ &+ \mathcal{L}_m(g_{\mu\nu},\psi_i) \left. \right\}, \end{split}$$

where
$$X \equiv -\frac{1}{2}(\partial_{\mu}\varphi)^2$$

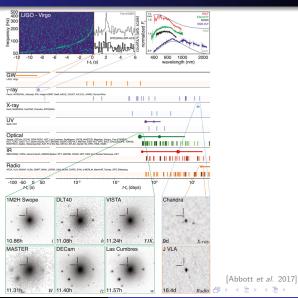
[Horndeski 1974; Deffayet, Gao, Steer, Zahariade 2011; Kobayashi, Yamaguchi, Yokoyama 2011]

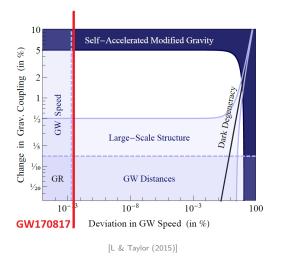


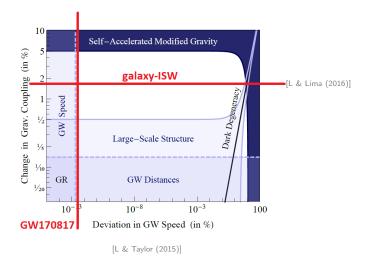


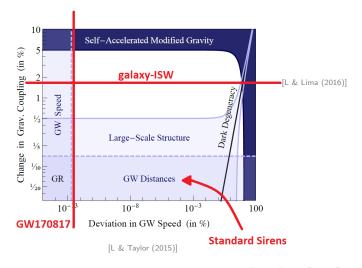


GW170817









Outlook – Standard Sirens



A gravitational-wave standard siren measurement of the Hubble constant

The LIGO Scientific Collaboration and The Virgo Collaboration, The 1M2H Collaboration, The Dark Energy Camera GW-EM Collaboration and the DES Collaboration, The DLT40 Collaboration, The Las Cumbres Observatory Collaboration, The VINROUGE Collaboration & The MASTER Collaboration

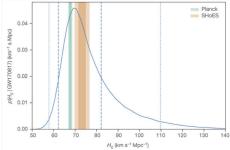
Affiliations | Contributions | Corresponding authors

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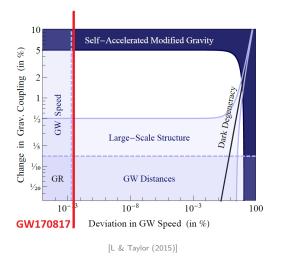


On 17 August 2017, the Advanced LIGO¹ and Virgo² detectors observed the gravitationalwave event GW170817—a strong signal from the merger of a binary neutron-star system³.

Figure 1: GW170817 measurement of Ho.



The marginalized posterior density for H_0 , $P(H_0 \mid \text{CWIT0B17})$, is shown by the blue curve. Constraints at 1 σ (darker shading) and 2 σ (lighter shading) from Planck ²⁰ and SHoES²¹ are shown in green and orange, respectively. The marginal constraints of the shading shading the shading sha



Horndeski action:

$$S = \frac{1}{2\kappa^{2}} \int d^{4}x \sqrt{-g} \left\{ G_{2}(\varphi, X) - G_{3}(\varphi, X) \Box \varphi + G_{4}(\varphi, X)R + \frac{\partial G_{4}}{\partial X} \left[(\Box \varphi)^{2} - (\nabla_{\mu} \nabla_{\nu} \varphi)^{2} \right] + G_{5}(\varphi, X) G_{\mu\nu} \nabla^{\mu} \nabla^{\nu} \varphi - \frac{1}{6} \frac{\partial G_{5}}{\partial X} \left[(\Box \varphi)^{3} - 3 \Box \varphi (\nabla_{\mu} \nabla_{\nu} \varphi)^{2} + 2 (\nabla_{\mu} \nabla_{\nu} \varphi)^{3} \right] + \mathcal{L}_{m}(g_{\mu\nu}, \psi_{i}) \right\}$$

where
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also see: [Ezquiaga & Zumalacárregui (2017); Creminelli & Vernizzi (2017); Sakstein & Jain (2017); Baker, Bellini, Ferreira, Lagos, Noller, Sawicki (2017); . . .]

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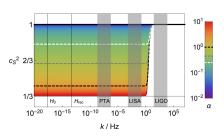
$$+ \mathcal{L}_{m}(g_{\mu\nu}, \psi_{i}) \right\} + \text{beyond?} \rightarrow \text{GWs decay to } \delta \phi$$
[Creminelli et al. (2018)]

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Linear shielding for GWs



[de Rham & Melville (2018)]

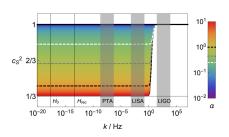
Scale-dependent propagation of GWs:

$$\begin{split} h_{ij}^{\prime\prime} + \left(3 + \frac{H^{\prime}}{H} + \frac{^{(0)}\alpha_{\mathrm{M}} + ^{(1)}\alpha_{\mathrm{M}}w(a)k_{H}^{2}}{1 + w(a)k_{H}^{2}}\right)h_{ij}^{\prime\prime} \\ + \left(1 + \frac{^{(0)}\alpha_{\mathrm{T}} + ^{(1)}\alpha_{\mathrm{T}}w(a)k_{H}^{2}}{1 + w(a)k_{H}^{2}}\right)k_{H}^{2}h_{ij} = 0 \end{split}$$

with GW170817
$$\Rightarrow$$
 $^{(1)}\alpha_{\mathrm{T}} \simeq 0$ as $k_{H} \sim 10^{19}$

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Self-acceleration driven by

$$\left|\frac{\Omega'}{\Omega}\right| = \left|^{(0)}\alpha_{\rm M} + \frac{^{(0)}\alpha_{\rm T}'}{1 + ^{(0)}\alpha_{\rm T}}\right| \gtrsim \mathcal{O}(1)$$

(cf. linear shielding for GWs) [L & Taylor (2014)]



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- Back to Λ? Cosmological constant problem?



Einstein-Hilbert action

$$S = M_{
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Field and constraint equations

$$G_{\mu\nu} + \Lambda g_{\mu\nu} = M_{\rm Pl}^{-2} T_{\mu\nu} \,, \qquad \int dV_4 \left(\frac{R}{2} - \Lambda\right) = 0$$

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angle$$

[L (2019)] (also see, e.g., vacuum energy sequestering, unimodular gravity, Barrow & Shaw (2011), 4-form field strength of a 3-form gauge field, δ -function from scalar-vector term, multiverse type II, higher-dimensional scalar-tensor theory, supergravity, string theory . . .)



Old cosmological constant problem

• Combined field equations: $\int dV_4 (R-2\Lambda) = 0$, $M_{\rm Pl}^2 \Lambda = \frac{1}{2} \frac{\int dV_4 T}{\int dV_4} \equiv \frac{1}{2} \langle T \rangle$

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• Example 3: expansion of Λ_{vac} in M_{Pl}^2 cancelled by expansion of classical Λ

• Value of the cosmological constant, $M_{\rm Pl}^2 \Lambda = \frac{1}{2} \frac{\int dV_4 T}{\int dV_4} \equiv \frac{1}{2} \langle T \rangle$ assume spatially perfectly homogeneous matter-only universe with instantaneous collapse at $t_{\rm end}$

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- Universe is inhomogeneous on small scales (backreaction)
 (final conditions clearer than initial conditions: ultimate critical matter cells)

$$\mathcal{S}_{\mathrm{m}} = \sum_{i} \int_{\mathcal{U}_{i}} \mathsf{d}V_{4} \mathcal{L}_{\mathrm{m},i} + \int_{\mathcal{M} \setminus \bigcup_{i} \mathcal{U}_{i}} \mathsf{d}V_{4} \mathcal{L}_{\varnothing}$$

Matching Λ between \mathcal{U}_i and $\mathcal{M}\setminus\bigcup_i\mathcal{U}_i$: $\mathcal{L}_\varnothing=M_{\mathrm{Pl}}^2(\Lambda/n+M_{\mathrm{Pl}}^{2n}\bar{\Lambda}_\varnothing)$



Spherical collapse equation from energy-momentum conservation in matter cell

$$y'' + \left(2 + \frac{H'}{H}\right)y' + \frac{1}{2}\Omega_{\rm m}(a)(y^{-3} - 1)y = 0$$

where $y \equiv (
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Self-consistent production of cosmological constant

$$rac{rac{1}{2} \mathcal{M}_{
m Pl}^{-2} \langle \mathcal{T}
angle}{\Lambda_{
m obs}} = rac{\Omega_{
m m}}{4 (1-\Omega_{
m m})} rac{t_{
m max}}{\int_0^{t_{
m turn}} dt \, a^3 y^3} > rac{t_{
m max}}{2 t_{
m turn}} = 1$$

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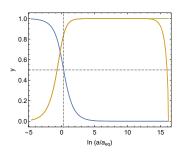
Self-consistent production of cosmological constant

$$rac{rac{1}{2} \mathcal{M}_{ ext{Pl}}^{-2} \langle T
angle}{m{\Lambda}_{ ext{obs}}} = rac{\Omega_{ ext{m}}}{4(1-\Omega_{ ext{m}})} rac{t_{ ext{max}}}{\int_0^{t_{ ext{turn}}} dt \ a^3 y^3} > rac{t_{ ext{max}}}{2t_{ ext{turn}}} = 1$$

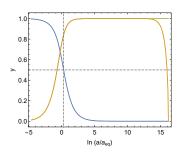
• Standard field equations (w/ non-gravitating vacuum)

$$R_{\mu\nu} - rac{1}{2}Rg_{\mu\nu} + \Lambda_{
m obs} = M_{
m Pl}^{-2}T_{\mu\nu}$$

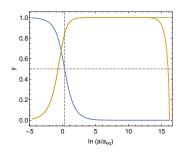




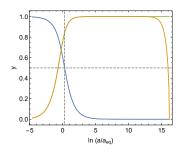
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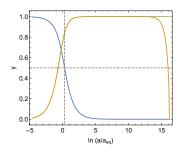


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$$\Rightarrow \Omega_{\Lambda} = 0.704$$

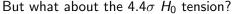
• 1σ & 3σ agreement w/ DES & Planck

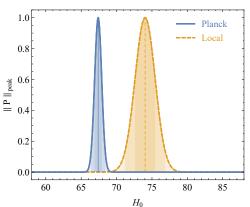
• Challenges to cosmic acceleration from DE or MG.

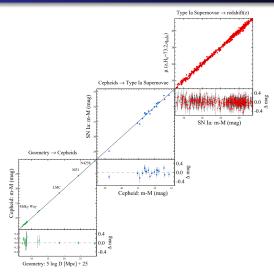
- Challenges to cosmic acceleration from DE or MG.
- Variation of Einstein-Hilbert action wrt Planck mass yields a topological constraint equation that prevents vacuum energy from gravitating (old Λ problem).

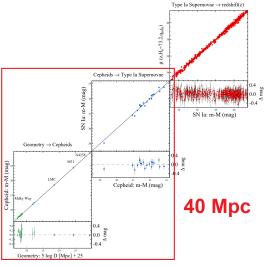
- Challenges to cosmic acceleration from DE or MG.
- Variation of Einstein-Hilbert action wrt Planck mass yields a topological constraint equation that prevents vacuum energy from gravitating (old Λ problem).
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- Variation of Einstein-Hilbert action wrt Planck mass yields a topological constraint equation that prevents vacuum energy from gravitating (old Λ problem).
- Evaluation of the constraint equation for matter patches that form ultimate critical cells yields backreaction effect that predicts $\Omega_{\Lambda}=0.704$, giving rise to cosmic late-time acceleration & coincident current $\Omega_{\rm m}\sim\Omega_{\Lambda}$ (new Λ problem).
- No problem with the cosmological constant.

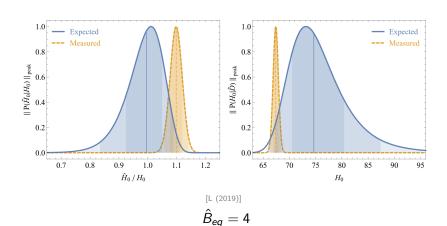








Lucas Lombriser



Thank you!