

Constraining the gluon nuclear content with heavy-flavour production at the LHC

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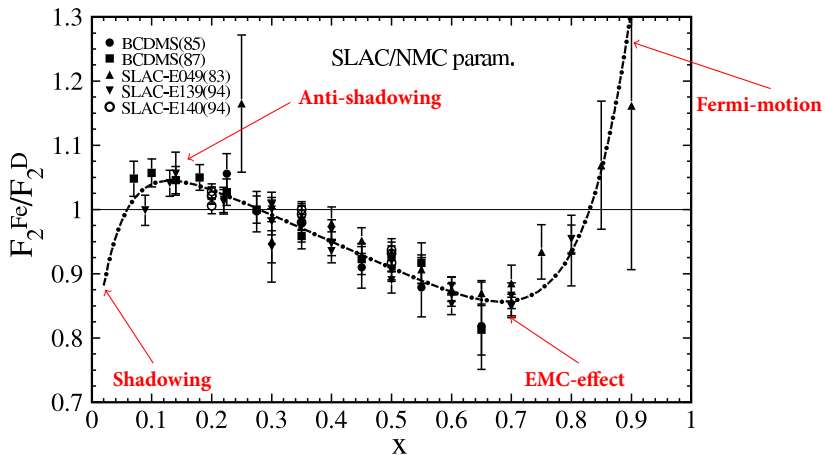
in collaboration with A. Kusina, I. Schienbein and H.S. Shao

based on PRL 121 (2018) 052004 & EPJC 77 (2017) 1 & more to appear

Introduction

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- Cross-sections in nuclear collisions are modified



- We translate these modifications into universal quantities: **nuclear PDFs (nPDFs)**

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Nuclear PDFs (nPDFs) allow one to get information on:

- the **structure of the nuclei** in terms of quarks and gluons;
- the **initial state** of heavy-ion collisions at the LHC and RHIC
esp. to use **perturbative probes** of the QGP to study its properties.

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fit to experimental data. Only the evolution is perturbative

→ Can heavy-flavour data help us better constrain the nPDFs ?

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Differences with the free-proton PDFs

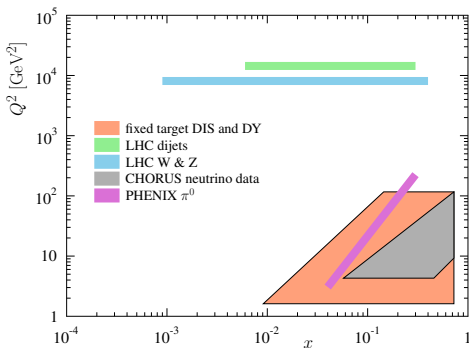
- Theoretical status of Factorisation
- Parametrisation – more parameters to model the A -dependence
- Different data sets – much less data:

- Less data \rightarrow less constraining power \rightarrow **more assumptions** (fixing) about fitting parameters

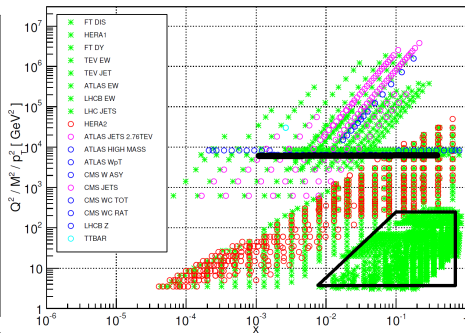
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EPPS16 dataset

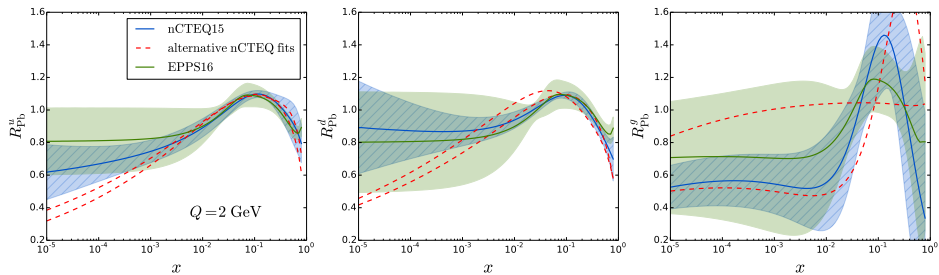


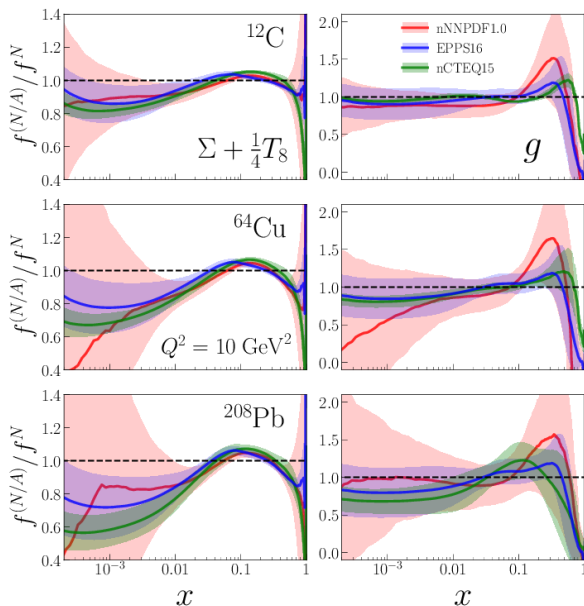
NNPDF3.0 NLO dataset



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SET		EPS09 JHEP 0904 (2009) 065	DSSZ PRD85 (2012) 074028	nCTEQ15 PRD93 (2016) 085037	KA15 PRD93 (2016) 014036	EPPS16 EPJC C77 (2017)163	nNNPDF1.0 1904.00018
data	eDIS	✓	✓	✓	✓	✓	✓
	DY	✓	✓	✓	✓	✓	✗
	π^0	✓	✓	✓	✗	✓	✗
	vDIS	✗	✓	✗	✗	✓	✗
	pPb	✗	✗	✗	✗	✓	✗
# data		929	1579	740	1479	1811	451
order		NLO	NLO	NLO	NNLO	NLO	NNLO
proton PDF		CTEQ6.1	MSTW2008	~CTEQ6.1	JR09	CT14NLO	NNPDF3.1
mass scheme		ZM-VFNS	GM-VFNS	GM-VFNS	ZM-VFNS	GM-VFNS	FONLL-B
comments		$\Delta\chi^2=50$, ratios, huge shadowing-antishadowing	$\Delta\chi^2=30$, ratios, medium-modified FFs for π^0	$\Delta\chi^2=35$, PDFs, valence flavour sep., not enough sensitivity	PDFs, deuteron data included	$\Delta\chi^2=52$, flavour sep., ratios, LHC pPb data	NNPDF methodology, isoscalarity assumed





Heavy-flavour LHC data

Available heavy-flavour p Pb LHC data*

	D^0	J/ψ	$B \rightarrow J/\psi$	$Y(1S)$
μ_0	$\sqrt{4M_{D^0}^2 + P_{T,D^0}^2}$	$\sqrt{M_{J/\psi}^2 + P_{T,J/\psi}^2}$	$\sqrt{4M_B^2 + \left(\frac{M_B}{M_{J/\psi}} P_{T,J/\psi}\right)^2}$	$\sqrt{M_{Y(1S)}^2 + P_{T,Y(1S)}^2}$
$p+p$ data	LHCb [1]	LHCb [2,3]	LHCb [2,3]	ALICE [4], ATLAS [5], CMS [6], LHCb [7,8]
R_{pPb} data	ALICE [9], LHCb [15]	ALICE [10,11], LHCb [16,12]	LHCb [12]	ALICE [13], ATLAS [14], LHCb [17]

- [1] LHCb, R. Aaij et al., JHEP 06, 147 (2017), 1610.02230.
- [2] LHCb, R. Aaij et al., Eur. Phys. J. C71, 1645 (2011), 1103.0423.
- [3] LHCb, R. Aaij et al., JHEP 06, 064 (2013), 1304.6977.
- [4] ALICE, B. B. Abelev et al., Eur. Phys. J. C74, 2974 (2014), 1403.3648.
- [5] ATLAS, G. Aad et al., Phys. Rev. D87, 052004 (2013), 1211.7255.
- [6] CMS, S. Chatrchyan et al., Phys. Lett. B727, 101 (2013), 1303.5900.
- [7] LHCb, R. Aaij et al., Eur. Phys. J. C72, 2025 (2012), 1202.6579.
- [8] LHCb, R. Aaij et al., JHEP 11, 103 (2015), 1509.02372.
- [9] ALICE, B. B. Abelev et al., Phys. Rev. Lett. 113, 232301 (2014), 1405.3452.
- [10] ALICE, J. Adam et al., JHEP 06, 055 (2015), 1503.07179.
- [11] ALICE, B. B. Abelev et al., JHEP 02, 073 (2014), 1308.6726.
- [12] LHCb, R. Aaij et al., (2017), 1706.07122.
- [13] ALICE, B. B. Abelev et al., Phys. Lett. B740, 105 (2015), 1410.2234.
- [14] The ATLAS collaboration, (2015), ATLAS-CONF-2015-050.
- [15] LHCb, R. Aaij et al., JHEP 1710 (2017) 090, 1707.02750.
- [16] LHCb, R. Aaij et al., JHEP 02, 072 (2014), 1308.6729.
- [17] LHCb, R. Aaij et al., JHEP 07, 094 (2014), 1405.5152.

* : At the time we performed our study (\rightarrow we can add more in the future; [see the talks in this session](#))

Expected nuclear effects on heavy quark(onium) production in pA collisions

- Nuclear modification of **PDFs**: initial-state effect
- **Energy loss** (w.r.t. pp collisions): initial-state or final-state effect
- **Break up** of the quarkonium in the **nuclear matter**: final-state effect
- **Break up** by **comoving particles**: final-state effect
- **Colour filtering** of intrinsic QQ pairs: initial-state effect
- ...

We assume that the leading twist factorisation is valid – **ONLY** modifications of PDFs are present → “shadowing-only” hypothesis.

→ we de facto do not consider excited quarkonium data.

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Automating the computation of nuclear PDF effects on HF production

An automated code to evaluate the impact of nuclear PDF I

JPL, H.S. Shao Eur.Phys.J. C77 (2017) 1

- **Partonic scattering cross section fit from pp data** with a Crystal Ball function parametrising $|\mathcal{A}_{gg \rightarrow \mathcal{H}X}|^2$

C.H. Kom, A. Kulesza, W.J. Stirling PRL 107 (2011) 082002

- A way to evade the quarkonium-production-mechanism controversy ?

for a recent review: JPL, 1903.09185, submitted to Phys. Rept.

To some extent, I would say "yes".

- Applied to J/ψ , Υ , D and B : it can be extended to all the probes produced in $2 \rightarrow 2$ partonic processes with a single partonic contribution

- The key point to compute nPDF effect is to have a **partonic** cross section

- Validated (for D^0) by the state-of-the-art pQCD code FONLL

- Any nPDF set available in LHAPDF5 or 6 can be used

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(gg or $q\bar{q}$, ...)

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which make sense if nPDFs are the only nuclear effect
- Conversely, one can test this hypothesis by comparing our curves to data
Global agreement $\stackrel{?}{\Rightarrow}$ only nPDFs matter
- One can go further in the data comparison with reweighting (see later)
and then HF-data inclusion in nPDF fits
- Last but not least: the automation of the evaluation allows one to study different nPDF sets AND the scale uncertainties: better control of the theory uncertainties

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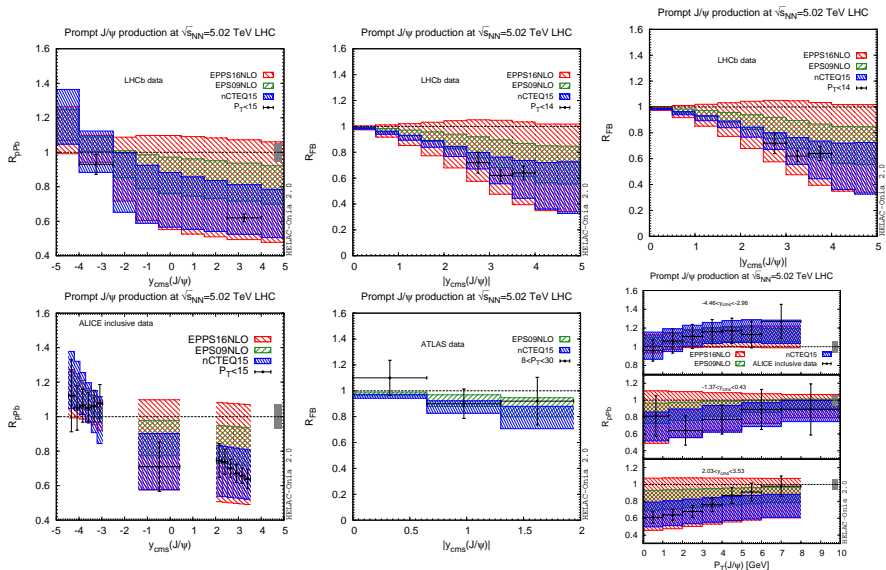
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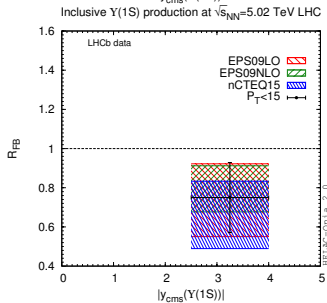
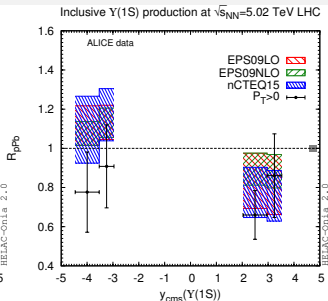
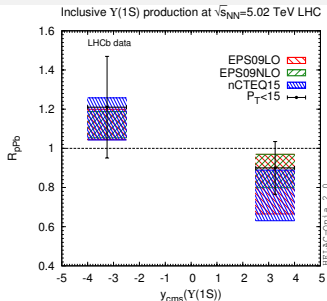
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Results for pA collisions using nCTEQ15 & EPPS16 out-of-the-box

Some J/ψ comparisons



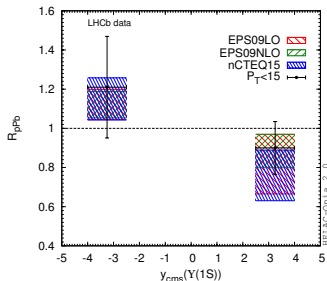
More results: $\Upsilon(1S)$ and ... η_c



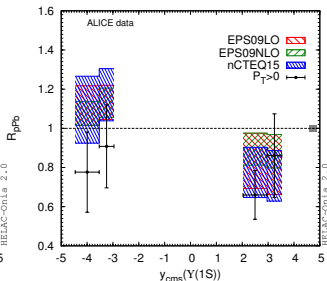
Our 8 TeV predictions
in J.L. Albacete, *et al.*
NPA 972 (2018) 18

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Inclusive $\Upsilon(1S)$ production at $\sqrt{s_{NN}}=5.02$ TeV LHC

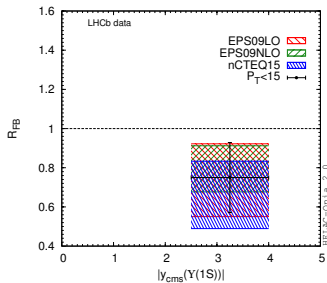


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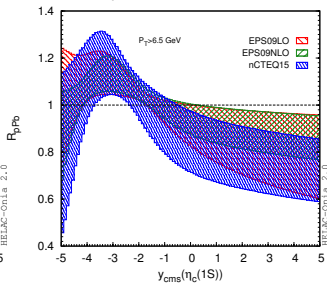


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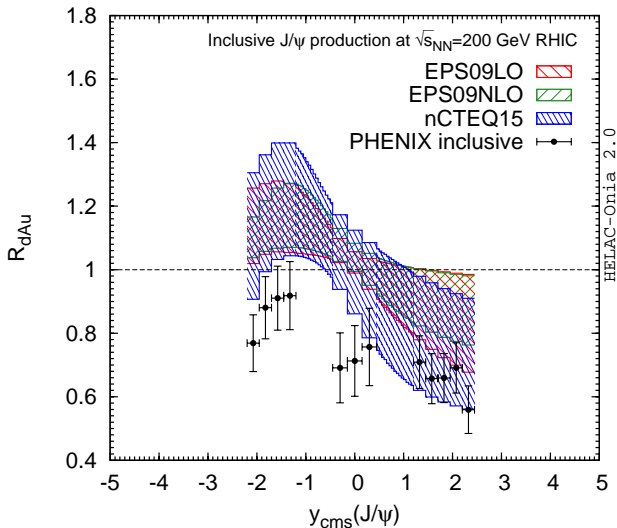
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Prompt $\eta_c(1S)$ production at $\sqrt{s_{NN}}=5.02$ TeV LHC



J/ψ 's at RHIC: no more tension with forward J/ψ data !



- : Backward data likely affected by break-up or comover interactions;
- : Question mark on the $y \sim 0$ data: no antishadowing ?

Impact on nPDFs: reweighting analysis¹

¹From now on, all nPDF uncertainties are 68%CL

Reweighting: the principle behind

1. Convert Hessian error PDFs into replicas

$$f_k = f_0 + \sum_i^N \frac{f_i^{(+)} - f_i^{(-)}}{2} R_{ki},$$

2. Calculate weights for each replica

$$w_k = \frac{e^{-\frac{1}{2}\chi_k^2/T}}{\frac{1}{N_{\text{rep}}} \sum_i^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_i^2/T}}, \quad \chi_k^2 = \sum_j^{N_{\text{data}}} \frac{(D_j - T_j^k)^2}{\sigma_j^2}$$

3. Calculate observables with new (reweighted) PDFs

$$\langle \mathcal{O} \rangle_{\text{new}} = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} w_k \mathcal{O}(f_k),$$
$$\delta \langle \mathcal{O} \rangle_{\text{new}} = \sqrt{\frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} w_k (\mathcal{O}(f_k) - \langle \mathcal{O} \rangle)^2}.$$

- * N is the # of eigensets, N_{rep} is the # of constructed replicas
- * f_0 is the "central-value" of the nPDF vector (i.e. of functions of x) in N_{flavour} dimension
- * $f_i^{(\pm)}$ ($i \in [1 : N]$) is the "upper/lower value" function of a given eigenset i
- * R_{ki} is a number randomly chosen for each set of (k, i) (thus fixed for all N_{flavour}) according to a standard Normal distribution
- * f_k is the constructed vector
- * T_j^k is the predicted theory for data j and replica k
- * T is the tolerance factor (for 68% CL: 13 for nCTEQ15; 19 for EPPS16)

- Global data (or theory) uncertainties can be dealt with adjusting T_j^k
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- Any other observables can also be redrawn ($p_A d\sigma$, R_{pA} , R_{FB} , ...)

Reweighting: the principle behind

1. Convert Hessian error PDFs into replicas

$$f_k = f_0 + \sum_i^N \frac{f_i^{(+)} - f_i^{(-)}}{2} R_{ki},$$

2. Calculate weights for each replica

$$w_k = \frac{e^{-\frac{1}{2}\chi_k^2/T}}{\frac{1}{N_{\text{rep}}} \sum_i^{N_{\text{rep}}} e^{-\frac{1}{2}\chi_i^2/T}}, \quad \chi_k^2 = \sum_j^{N_{\text{data}}} \frac{(D_j - T_j^k)^2}{\sigma_j^2}$$

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$$\langle \mathcal{O} \rangle_{\text{new}} = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} w_k \mathcal{O}(f_k),$$
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- * N is the # of eigensets, N_{rep} is the # of constructed replicas
- * f_0 is the "central-value" of the nPDF vector (i.e. of functions of x) in N_{flavour} dimension
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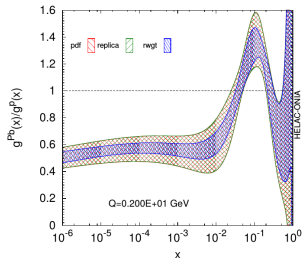
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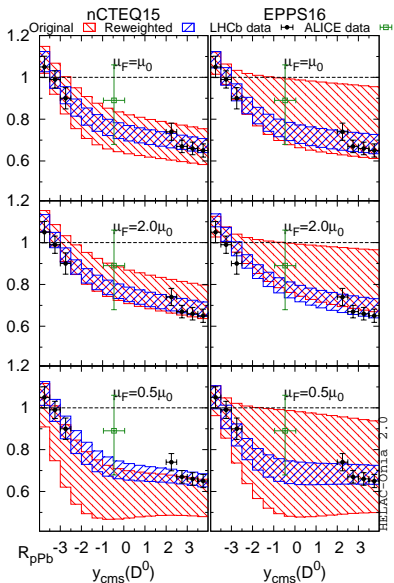
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Reweighting with D^0 data

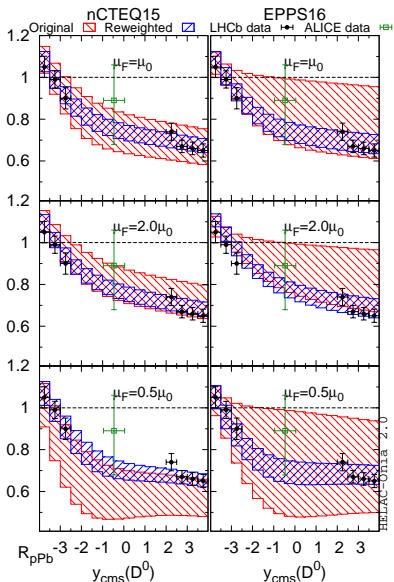


LHCb [JHEP 1710 (2017) 090, 1707.02750]

ALICE [PRL113, 232301 (2014), 1405.3452]

- Initial description of data is good for both nCTEQ15 and EPPS16.
- Substantial reduction of uncertainty especially for EPPS16.

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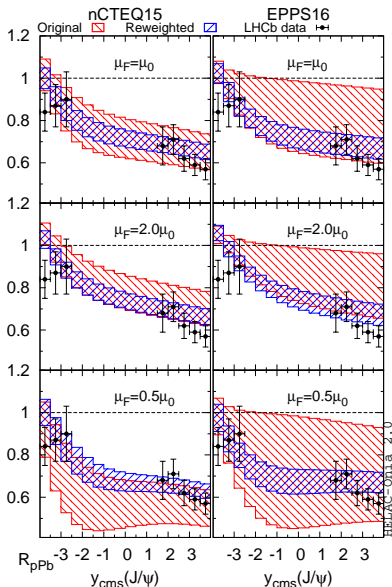


LHCb [JHEP 1710 (2017) 090, 1707.02750]

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- If we include factorisation scale uncertainty errors increase and it can become the dominant uncertainty.

Reweighting with J/ψ data

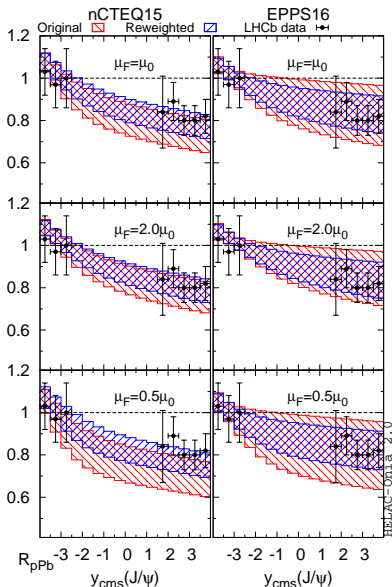


LHCb [JHEP 02, 072 (2014), 1308.6729; PLB 774 (2017) 159, 1706.07122]

ALICE [JHEP 06, 055 (2015), 1503.07179; JHEP 02, 073 (2014), 1308.6726]

- Again we observe a good agreement with the data; the **scale uncertainty** becomes important.

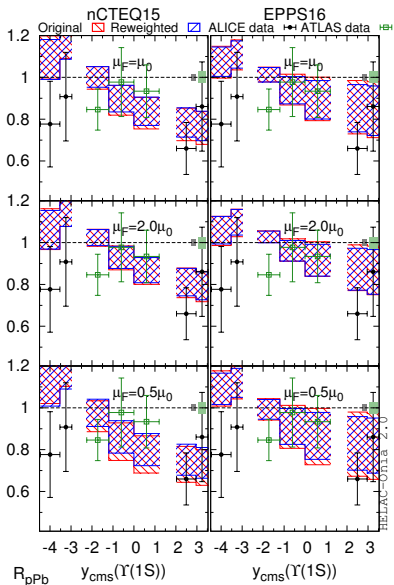
Reweighting with $B \rightarrow J/\psi$ data



LHCb [PLB 774 (2017) 159, 1706.07122]

- Scale uncertainty is reduced compared to the D^0 and J/ψ case.
- Data are not yet precise enough to give substantial constraints on nPDFs (but if the precision rises there is big potential).

Reweighting with $\Upsilon(1S)$ data

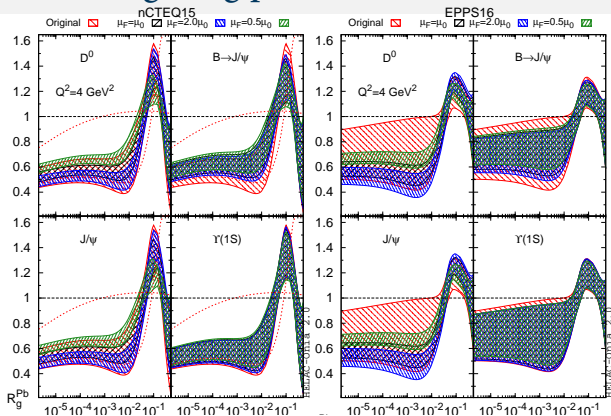


ALICE [PLB 740, 105 (2015), 1410.2234]

ATLAS [ATLAS-CONF-2015-050 (updated in: 1709.03089)]

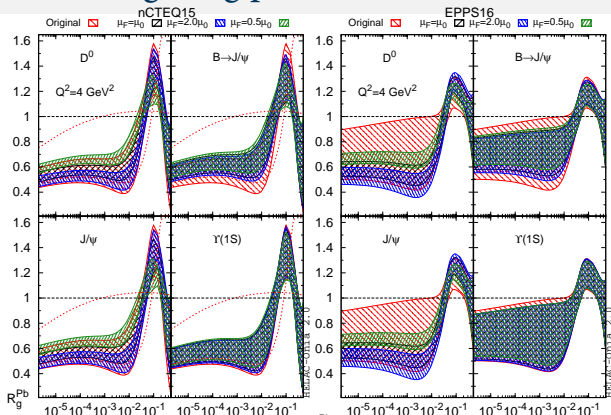
- With the current precision we do not get any additional constraints on the nPDFs.

Results of the reweighting process



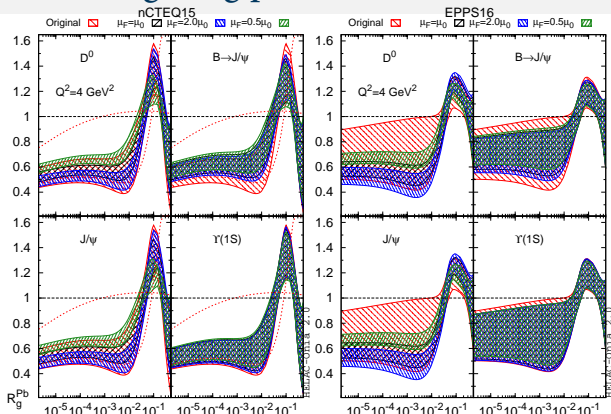
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- Confirmation of the existence of a gluon anti-shadowing: $R_g(0.05 \leq x \leq 0.1) > 1$

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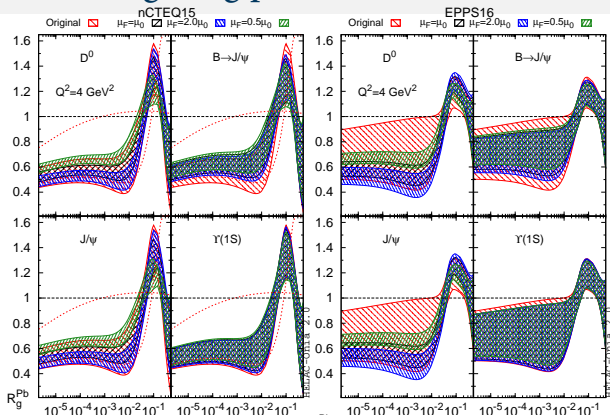
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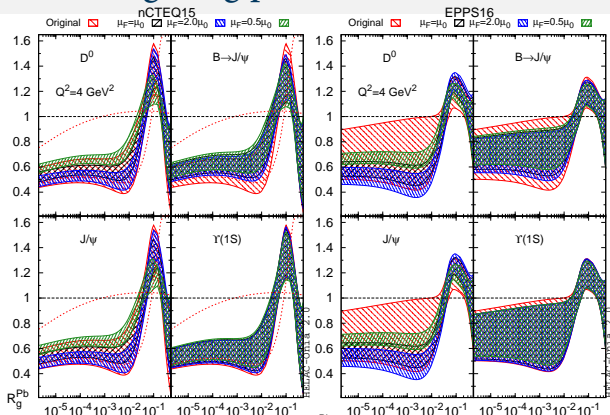
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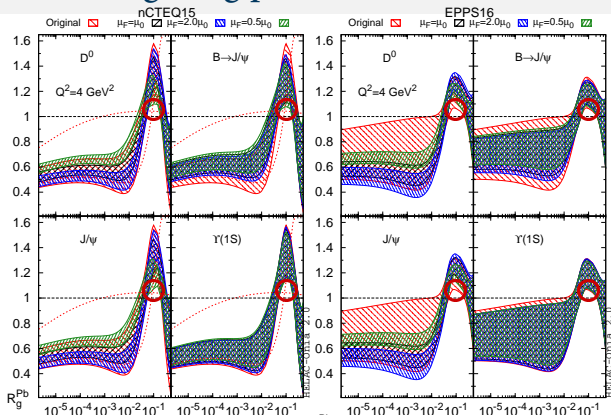
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Consistency with other data

We checked the consistency of the reweighted (nCTEQ15) nPDFs with other data sets entering global analyses:

- DIS data (the most precise set NMC Sn/C [NPB 481 (1996) 23]).
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This is very non-trivial and further confirms the “shadowing-only” hypothesis of the leading-twist factorisation is valid within the current data precision!

Table: χ^2/N_{data} values for nCTEQ15 nPDFs before and after reweighting using B , D , J/ψ data for different scale choices.

	nCTEQ15	after reweighting								
		$B \mu_0$	B up	B down	$D \mu_0$	D up	D down	$J/\psi \mu_0$	J/ψ up	J/ψ down
W/Z LHC (102)	1.43	0.10	0.15	0.26	0.39	0.11	0.14	0.66	0.25	0.35
NMC $F_2^{\text{Sn}}/F_2^{\text{C}}$ (III)	0.38	0.71	0.64	0.77	0.39	0.56	0.64	0.60	0.26	0.86
NMC $F_2^{\text{Pb}}/F_2^{\text{C}}$ (14)	0.37	0.51	0.53	0.56	0.30	0.63	0.58	0.44	0.56	0.39
PHENIX $J/\psi R_{dAu}$	1.89							1.93	0.43	0.35

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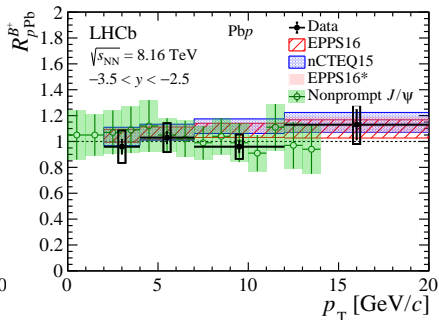
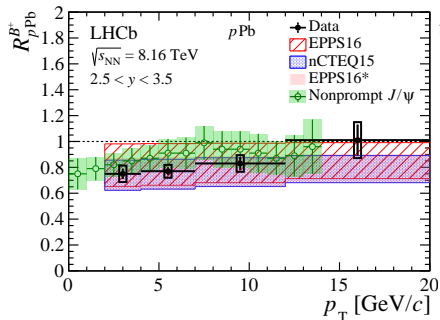
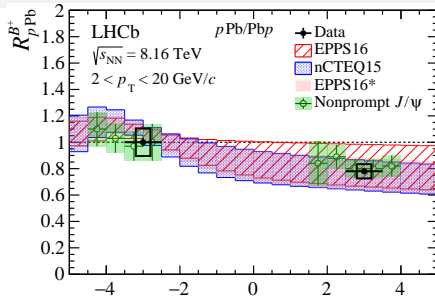
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Predictions using the reweighted nPDFs

Improved agreement with the reweighted nPDFs for B^+

LHCb, PRD 99 (2019) 052011



Conclusions

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→ need for new constraints at $x \lesssim 10^{-3}$
- We have proposed a **quick and robust method** to evaluate effects of heavy-quark data on nPDFs
– complementary to full (but time consuming) pQCD computations
→ we used it to analyse D^0 , J/ψ , $B \rightarrow J/\psi$ and $Y(1S)$ data.
- We proposed a **“shadowing-only” hypothesis** based on the leading twist factorisation
→ suggested by the good agreement with the LHC data
- Under this hypothesis, we made, using a Bayesian reweighting, the first direct **observation of gluon shadowing**. Additionally, our analysis corroborates the existence of **gluon antishadowing**.
- Studying the coherence with **AA UPC J/ψ** data is in our to do list BUT
(i) the nuclear GPDs are unknown (likely not directly connected to nPDFs) &
(ii) the scale uncertainties is likely huge (cf. NLO $\gamma p \rightarrow J/\psi p$ analyses)
- We will release very soon the LHAPDF grids of the reweighted nPDFs presented here.
- Overall, the heavy-quark data can improve our knowledge of the nuclear gluon distribution and it should be included in future nPDF fits.
- Our work is part of a clear **renewal our interest for HF data** for (n)PDF studies
[EPJC 75 (2015) 396, JHEP 1704 (2017) 044, PRL 118 (2017) 072001, 1906.06971, 1906.02512, 1907.01400, ...].
- and will contribute to the elaboration of NLOACCESS within the EU network STRONG-2020.

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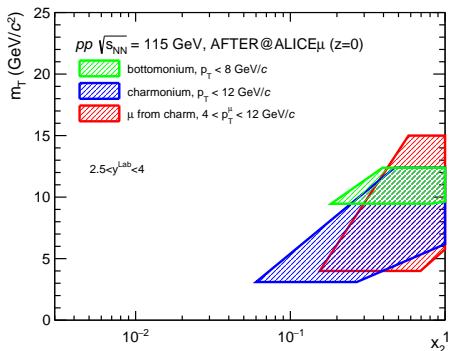
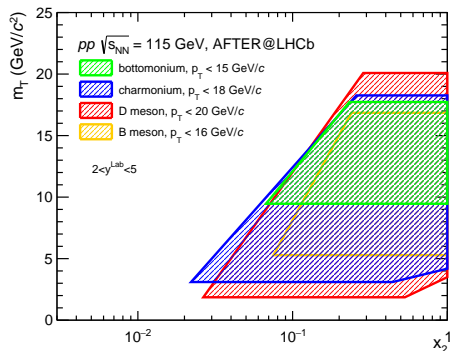
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BACKUP SLIDES

Application to the LHC in the fixed-target mode (AFTER@LHC)

HF kinematical coverage of LHCb and ALICE in the fixed-target mode [\[arXiv:1807.00603\]](https://arxiv.org/abs/1807.00603)

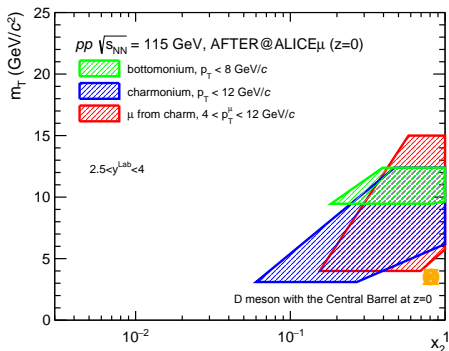
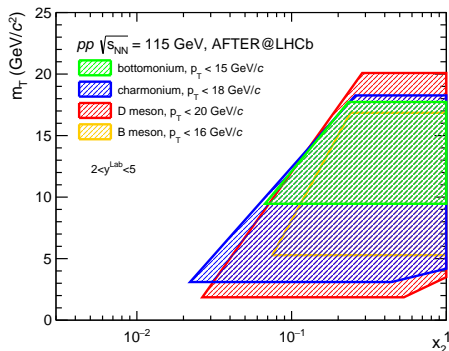


ALICE could cover $\eta \sim 1 - 2$ for quarkonium into dileptons with one muon in the muon arm and another in the central barrel

[done for UPCs in the collider mode]

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- EMC gluon effect totally unknown
- This is the realm of the FT LHC experiments
- First projections are extremely promising
[NB: initial nPDF uncertainties for $x > 0.1$ are underestimated;
→ simply no data exist there]
- This assumes that we control other potential CNM effects: global study required !
- Similar studies for the *proton* PDFs are yet to be done along the lines of the studies carried out for low- x gluon at the LHC
PROSA Coll. Eur.Phys.J. C75 (2015) 396; R. Gauld, J. Rojo PRL 118 (2017) 072001
- In the *pp* case, contrary to nPDF studies bearing on nuclear modification factors, one needs ways to reduce the systematical uncertainties

Reward: unique constraints on gluon PDFs at high x and low scales

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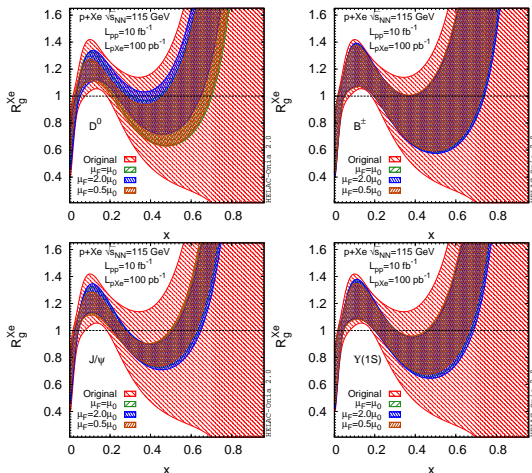
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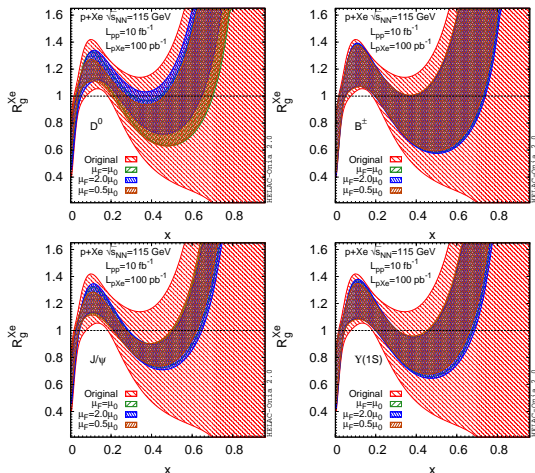
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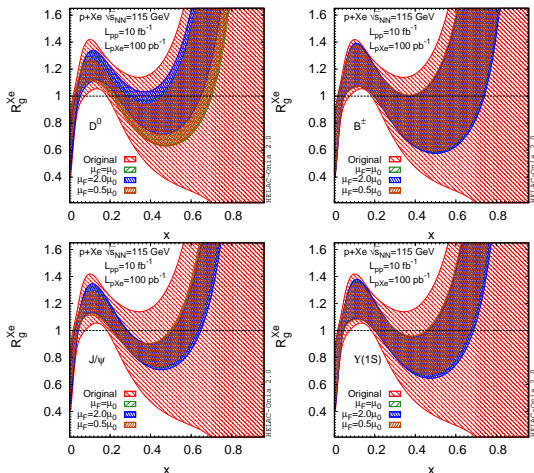
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A quick comparison between nCTEQ15 and EPPS16

• Parametrisation

- PDF of nucleus (A - mass, Z - charge)

$$f_i^{(A,Z)}(x, Q) = \frac{Z}{A} f_i^{p/A}(x, Q) + \frac{A-Z}{A} f_i^{n/A}(x, Q)$$

- bound proton PDFs

nCTEQ15 [arXiv:1509.00792]

$$x f_i^{p/A}(x, Q_0) = c_0 x^{c_1} (1-x)^{c_2} e^{c_3 x} (1 + e^{c_4 x})^{c_5}$$

$$c_k \rightarrow c_k(A) \equiv c_{k,0} + c_{k,1} (1 - A^{-c_{k,2}})$$

EPPS16 [arXiv:1612.05741]

$$f_i^{p/A}(x, Q) = R_i^A(x, Q) f_i^p(x, Q),$$

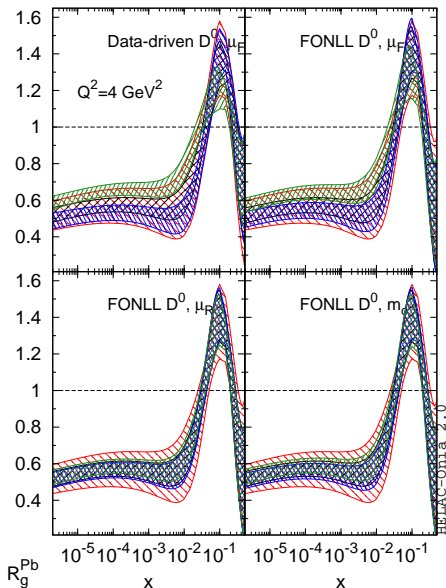
$$R_i^A(x, Q_0) = \begin{cases} a_0 + a_1(x - x_a)^2 & x \leq x_a \\ b_0 + b_1 x^\alpha + b_2 x^{2\alpha} + b_3 x^{3\alpha} & x_a \leq x \leq x_e \\ c_0 + (c_1 - c_2 x)(1-x)^{-\beta} & x_e \leq x \leq 1 \end{cases}$$

$$d_i \rightarrow d_i(A) = d_i(A_{\text{ref}}) \left(\frac{A}{A_{\text{ref}}} \right)^{\gamma_i [d_i(A_{\text{ref}}) - 1]},$$

with $d_i = a_i, b_i, \dots$ and $A_{\text{ref}} = 12$

Data-driven vs. FONLL

nCTEQ15



nCTEQ15

