J/ψ production in hadron scattering: three-pomeron contribution

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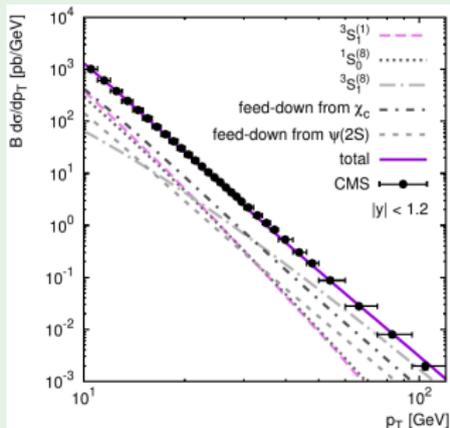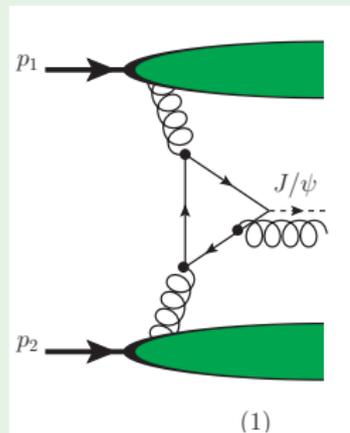FEDERICO SANTA MARIA

J/ψ in pp collisions-how is it produced?

Color Singlet Model & NRQCD corrections

(1)

(Reggeized) gluon=(cut) pomeron

Other mechanisms at large p_T ?

- Co-production ($J/\psi + \bar{Q}Q$, ...)

(PRL 101 (2008) 152001)

- Quark and gluon fragmentation
(EPJC 79 (2019) 241)

- ... ?

[Phys. Rev. D 96, no. 3, 034019 (2017)]:

(CSM alone):
- Reasonable predictions for:
  - small-p_T observables
  - p_T-integrated observables
- Wrong behaviour for $p_T \gg m_{J/\psi}$
(EPJC 79 (2019) 241)
(NRQCD):
- Description for large-p_T

Challenge for NRQCD: LDMEs not universal ?!

| | $\langle \mathcal{O}^\psi\left[^3S_1^{(1)}\right]\rangle/\text{GeV}^3$ | Color octet contributions | | |
		$\langle \mathcal{O}^\psi\left[^1S_0^{(8)}\right]\rangle/\text{GeV}^3$	$\langle \mathcal{O}^\psi\left[^3S_1^{(8)}\right]\rangle/\text{GeV}^3$	$\langle \mathcal{O}^\psi\left[^3P_0^{(8)}\right]\rangle/\text{GeV}^5$
kT-factorization				
A0	1.97	0.0	9.01×10^{-4}	0.0
JH	1.62	1.71×10^{-2}	2.83×10^{-4}	0.0
KMR	1.58	8.35×10^{-3}	2.32×10^{-4}	0.0
Collinear factorization				
[11]	1.32	3.04×10^{-2}	1.68×10^{-3}	-9.08×10^{-3}
[15]	1.16	9.7×10^{-2}	-4.6×10^{-3}	-2.14×10^{-2}

Multiplicity dependence of charmonia production

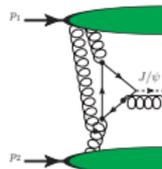- Observable: J/ψ + charged hadrons, study multiplicity $dN_{J/\psi}$ vs. dN_{ch}
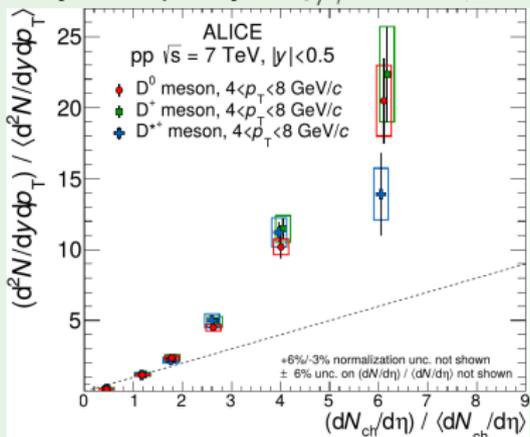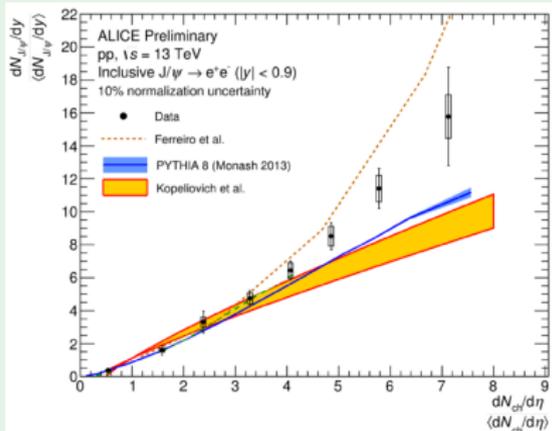
- Enhancement (deviation from linear) seen both for J/ψ and D mesons
- So far not clear if the effect exists for $\psi(2S)$, χ_c, Υ?
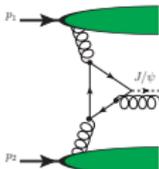
- Difficult to explain in terms of gluon-gluon fusion approach: if each reggeized gluon (cut pomeron) contributes approx. equal number $\bar{n}$ of charged hadrons, expect linear dependence
- Data hints that multipomeron mechanisms are pronounced (usually discarded as corrections).
- Production on nuclei: $\sim A^{1/3}$ enhancement

3-pomeron contributions
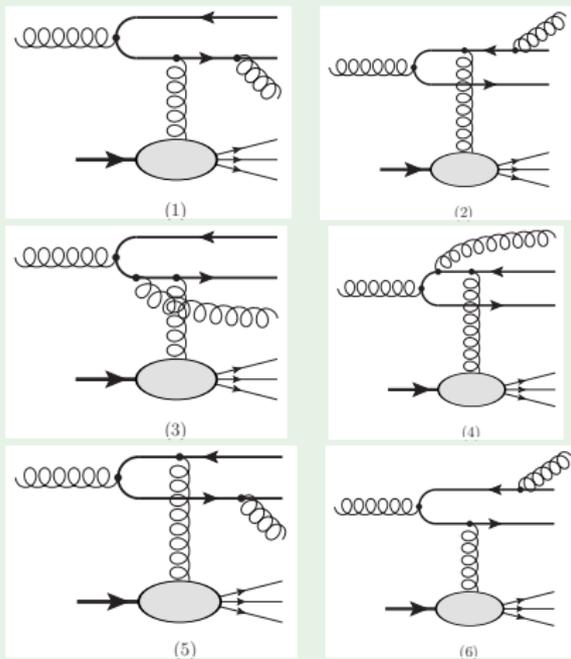

- Diagram (2) contributes additively to $d\sigma$
- Diagrams (1) and (3) might interfere:

- Diagrams (1) and (2) include Parton Distributions with odd number of gluons
- ... such distributions are sensitive if the proton is polarized [PLB 668 (2008), 216, PRD 78 (2008), 014024 , PRD 78 (2008), 114013].
- ... corresponding asymmetries are small, might neglect them [PHENIX Collaboration, PRD 82, 112008 (2010), PRD 86 (2012), 099904]

$\Rightarrow$ Additive contribution to the LO cross-section, as shown in (3).

CSM in dipole approach

LO diagrams

(1) (2) (3) (4) (5) (6)

● Same diagrams as in k_T factorization, we just express everything in terms of dipole framework.
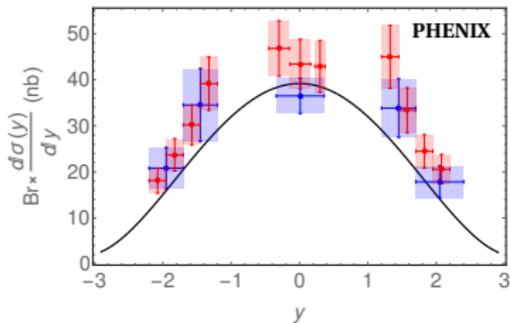
[PRC95 (2017), 065203]

Large-m_Q limit

● Dipole cross-section related to $k_\perp$ uPDF, e.g. for color singlet

$$\sigma_d(x, r_\perp) = \int \frac{d^2 k_\perp}{(2\pi)^2} \mathcal{F}(x, k_\perp) \left(1 - e^{ik_\perp \cdot r_\perp}\right).$$

● Color octet dipole cross-section could be expressed as linear combinations of color singlets

Reasonable description of pp data

# J/ψ production cross-section in the dipole approach

Meson WFs

$$\frac{d\sigma_{pp}}{dy} = \frac{9}{8}\, g\left(x_1(y)\right) \int d\alpha_G \, d\alpha_1 d^2 r_1 \, d\alpha_2 d^2 r_2 \, d^2\rho \, \overbrace{\Psi_M^*\left(\alpha_1 \, r_1\right) \Psi_M\left(\alpha_2 \, r_2\right)}^{} \times$$

$$\times \sum_{n,n'=1}^{6} \underbrace{\eta_n \eta_{n'}}_{\text{numerical factor}} \quad \text{Tr}\left[\Lambda_M \Phi_{g\to\bar{Q}Q}\left(\epsilon_n,\, \vec{r}_n^{(1)}\right) \Phi_{Q\to Q\, g}\left(\delta_n,\, \vec{\rho}_n\right)\right]$$

$$\times \quad \text{Tr}\left[\Lambda_M \Phi_{g\to\bar{Q}Q}\left(\epsilon_{n'},\, \vec{r}_{n'}^{(2)}\right) \Phi_{Q\to Q\, g}(\delta_{n'},\, \vec{\rho}_{n'})\right]^*$$

$$\times \quad \underbrace{\sigma\left(x_2,\, \vec{r}_n - \vec{r}_{n'}\right)}_{\text{dipole cross-section}}$$

- Λ_M-spin projector on meson WF
- $\Phi_{g\to\bar{Q}Q}$, $\Phi_{Q\to Q\,g}$ are evaluated perturbatively ($m_c \to \infty$ limit)
- $\vec{r}_n^{(1,2)} \approx \vec{r}_{1,2} + \delta\vec{r}_n\left(\alpha,\, \alpha_G,\, r,\, \rho\right)$, $\qquad \vec{\rho}_n \approx \vec{\rho} + \delta\vec{\rho}_n\left(\alpha,\, \alpha_G,\, r,\, \rho\right)$
- Sum over 6 diagrams in amplitude and its conjugate is implied
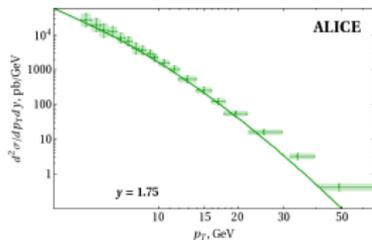
- For p_T-dependent cross-section, additional Fourier over difference of dipole impact parameter $\Delta\vec{b} = \vec{b}_1 - \vec{b}_2 \neq 0$; terms $\delta\vec{r}_n$, $\delta\vec{\rho}_n$ also depend on $\vec{b}_{1,2}$

3-gluon fusion in dipole approach
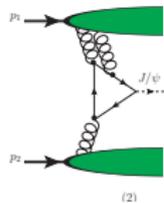
$$\frac{d\sigma\left(Y, Q^2\right)}{dy\, d^2 q_T} = 4 \int \frac{d^2 Q_T}{(2\pi)^2}\, S_h^2\left(Q_T\right)\ x_g\, G\left(x_g, M_{J/\psi}\right)$$

$$\times \int_0^1 dz \int_0^1 dz' \int \frac{d^2 r}{4\pi}\, \frac{d^2 r'}{4\pi}\, d^2 b\ e^{-i q_T \cdot b}$$

$$\times \left\langle \Psi_g\left(r, z\right)\, \Psi_{J/\psi}\left(r, z\right)\right\rangle \left\langle \Psi_g\left(r', z'\right)\, \Psi_{J/\psi}\left(r', z'\right)\right\rangle \left(N\left(y;\, b - \frac{1}{2}\left(r - r'\right)\right) \right.$$

$$\left. + N\left(y;\, b + \frac{1}{2}\left(r - r'\right)\right) - N\left(y;\, b + \frac{1}{2}\left(r + r'\right)\right) - N\left(y;\, b - \frac{1}{2}\left(r + r'\right)\right) \right)^2$$

$$+ (y \to Y - y)$$

- Note: Results are defined up to nonperturbative normalization factor $\sim \int \frac{d^2 Q_T}{(2\pi)^2} S_h^2\left(Q_T\right)$, related to Fourier image of cut pomeron Green function. Can be related to σ_{eff} in DPDF.

Explanation of uncertainty in k_T-factorization language

- For gluon uPDF use KMR parametrization, $\mu_R = \mu_F = \sqrt{p_T^2 + M_{J/\psi}^2}$
- For DPDFs $\mathcal{F}^{c_1 c_2, c_1' c_2'}(x_{1a}, k_{1a,T}, x_{1b}, k_{1b,T}, \Delta)$ can have different color structures,

$$\mathcal{F}^{aa', bb'} = \frac{1}{64}\left[{}^1F\, \delta^{aa'}\delta^{bb'} - \frac{\sqrt{8}}{3}\, {}^AF\, f^{aa'c}f^{bb'c} + \frac{3\sqrt{8}}{5}\, {}^SF\, d^{aa'c}d^{bb'c} \right.$$
$$\left. + \frac{4}{\sqrt{27}}\, {}^{27}F\, t_{27}^{aa',bb'} + \frac{2}{\sqrt{10}}\, {}^{10}F\left(t_{10}^{aa',bb'} + \left(t_{10}^{aa',bb'} \right)^* \right) \right],$$

- We *assume* that the parameterization is given by

$$\mathcal{F}^{c_1 c_2, c_1' c_2'}(x_{1a}, k_{1a,T}, x_{1b}, k_{1b,T}) \sim \mathcal{F}(x_{1a}, k_{1a,T})\, \mathcal{F}(x_{1b}, k_{1b,T}),$$

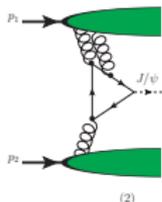- ... reproduces "pocket formula" for DPS

$$d\sigma_{pp \to h_1 h_2 X} = \frac{1}{\sigma_{\text{eff}}}\, d\sigma_{pp \to h_1 X}\, d\sigma_{pp \to h_2 X}.$$

$\Rightarrow \int \frac{d^2 Q_T}{(2\pi)^2}\, S_h^2(Q_T) \sim (\text{const})\, \sigma_{\text{eff}}^{-1}$

- <u>Caveat</u>: the "universal" value of σ_{eff} is not that universal; typical values are between 5 mb and 20 mb. [Eur.Phys.J. C77 (2017) no.2, 76]

Fixing uncertainty in 3-gluon fusion cross-section

$$\frac{d\sigma\left(Y, Q^2\right)}{dy} = \left(\int \dots\right) \langle \Psi_g\left(r, z\right) \Psi_{J/\psi}\left(r, z\right)\rangle \langle \Psi_g\left(r', z'\right) \Psi_{J/\psi}\left(r', z'\right)\rangle$$

$$\times \left(N\left(y; \frac{r - r'}{2}\right) - N\left(y; \frac{r + r'}{2}\right)\right)^2$$

- Use diffractive photoproduction ($\gamma^* p \to J/\psi\, X$) data to fix normalization:

$$\frac{d\sigma_{\text{diff}}\left(Y, Q^2\right)}{dy} = \left(\int \dots\right) \langle \Psi_{\gamma^*}\left(r, z\right) \Psi_{J/\psi}\left(r, z\right)\rangle \langle \Psi_{\gamma^*}\left(r', z'\right) \Psi_{J/\psi}\left(r', z'\right)\rangle$$

$$\times N\left(y;\, r\right) N\left(y;\, r'\right)$$

For $Q^2 \to \infty$, $m_c \to \infty$, $r \to 0$, can use pQCD expression for both $\Psi_{\gamma^*} \sim \Psi_g$.
- The dipole amplitude $N\left(y;\, r\right) \sim \left(Q_s^2 r^2\right)^{\bar{\gamma}}$ at small r, so rapidity dependence of $d\sigma$ and $d\sigma_{\text{eff}}$ is the same $\Rightarrow$ Can fix *hadro*production cross-section $d\sigma$ from diffractive *photo*production $d\sigma_{\text{diff}}$ studied at HERA:

| $d\sigma\left(y, Q^2\right)/dy$	Theoretical estimates	Experiment
$\sqrt{s} \approx 1.96$ TeV	2.1-2.6 μb	2.38 μb [67]
$\sqrt{s} \approx 7$ TeV	3.8-5.6 μb	5.8 μb [68]

- Theoretical estimates: three-pomeron fusion gives sizeable contribution
- [Eur.Phys.J. C75 (2015) no.5, 213] (in k_T-fact): qualitatively agree; smaller due to larger σ_{eff}
- In what follows will show only shapes for p_T, rapidity, multiplicity dependence

Contribution of $3g \to J/\psi$ to phenomenological observables
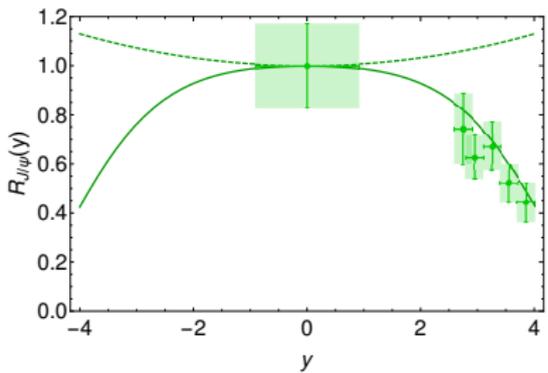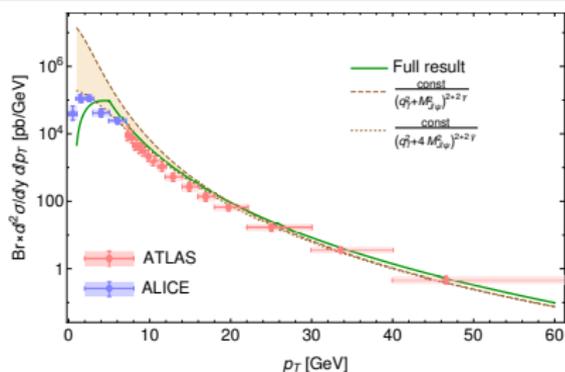
Multiplicity distributions
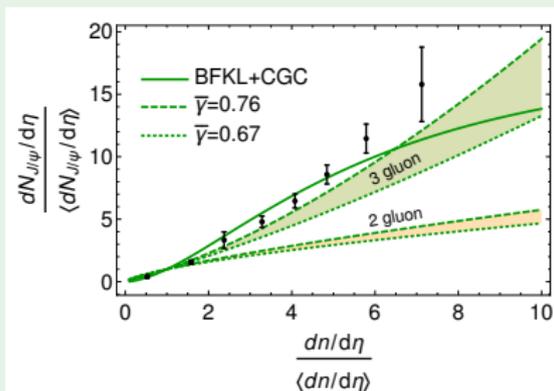
- BFKL: $\bar{\gamma} \approx 0.63$, phenomenology: $\bar{\gamma} \in [0.67, 0.76]$
- Data clearly favour 3-pomeron contribution.

- $\Rightarrow$ Overall, reasonable description of pp data might be achieved

## p_T- and rapidity distribution

- BFKL: each gluon $\sim x^{\alpha}$ (dashed)
- Experiment: add $\sim (1-x)^5$ (solid)

Contribution of the $g + gg \rightarrow g + J/\psi$ process

- Use k_T factorization, too difficult in dipole approach
  - Result expressed in terms of the gluon uPDFs and DPDFs of the proton

$$
\begin{aligned}
d\sigma \quad &= \frac{\pi \alpha_s^3 (\mu) \, |\mathcal{R}(0)|^2}{(x_{1a} x_{1b} x_2 s)^2} \frac{d^2 k_{1a,T}}{\pi} \frac{d^2 k_{1b,T}}{\pi} \frac{d^2 k_{2,T}}{\pi} \\
&\times \frac{dy d^2 p_T}{\pi} \, d\Delta^+ \frac{d^2 \Delta_T}{\pi} \, dy_g \, dx_{1a} \\
&\times \mathcal{M}_{gg+g \rightarrow g \, J/\psi}^{c_1 c_2, \, c_3} \left(k_{1a}, \, k_{1b}, \, k_2, \, p_{J/\psi} \right) \\
&\times \left(\mathcal{M}_{gg+g \rightarrow g \, J/\psi}^{c_1' c_2', \, c_3} \left(k_{1a} - \frac{\Delta}{2}, \, k_{1b} + \frac{\Delta}{2}, \, k_2, \, p_{J/\psi} \right) \right)^* \\
&\times \underbrace{\mathcal{F}^{c_1 c_2, c_1' c_2'} \left(x_{1a}, k_{1a,T}, x_{1b}, k_{1b,T}, \Delta \right)}_{\text{DPDF}} \underbrace{\mathcal{F} \left(x_2, k_{2\perp} \right)}_{\text{uPDF}}
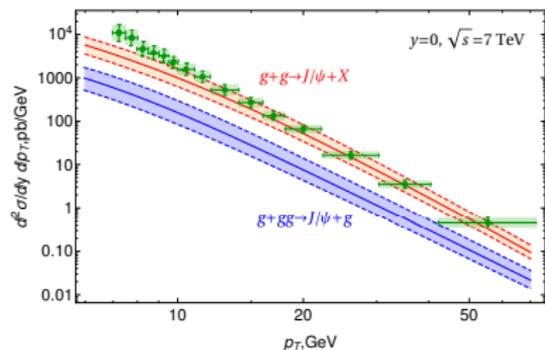\end{aligned}
$$
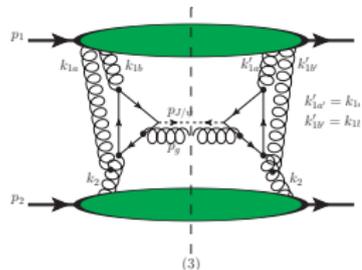
$\mathcal{M}$ = Sum over all possible permutations.

$\mathcal{M}_{gg+g \rightarrow J/\psi}^{} = \sum_{P_i} \left(\text{⟶} \quad \text{⟶} \right)$
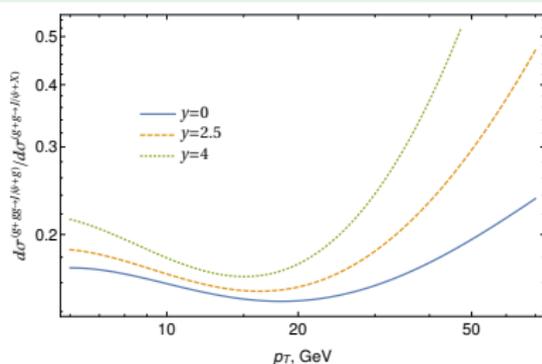
- For gluon uPDF use KMR parameterization.
- For gluon DPDF use factorized ansatz mentioned earlier

3-pomeron J/ψ production with gluon emission
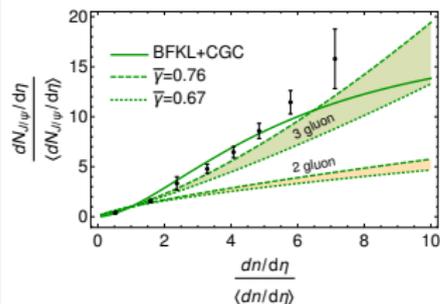
- Result for the cross-section:

- Relative w.r.t. CSM+NRQCD

- $\Rightarrow$ Similar to $g + gg \rightarrow J/\psi$, we get <u>sizeable</u> contribution from the three-pomeron fusion to J/ψ production (slightly smaller due to $\mathcal{O}\left(\alpha_s\left(m_c\right)\right)$)

Summary

- The observed multiplicity of charged particles signals that 3-pomeron process gives a sizeable contribution to J/ψ production.
- We believe that similar studies of multiplicity dependencies for $\psi(2S)$, $\Upsilon(nS)$ should shed light on the role of the suggested mechanism in their formation.

Thank You for your attention