

# Recent results of the azimuthal anisotropic flow measurements from STAR experiment

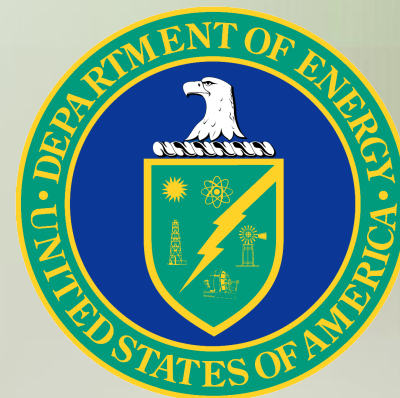


**Initial Stages 2017**

Polish Academy of Arts and Sciences  
September 18-22 2017, Cracow, Poland

Niseem Magdy  
Stony Brook University  
For the STAR Collaboration

[niseem.abdelrahman@stonybrook.edu](mailto:niseem.abdelrahman@stonybrook.edu)



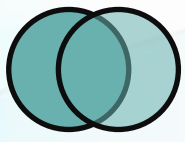
# Data Studied

- Collected data for Au+Au at  $\sqrt{s_{NN}} = 200 - 7.7 \text{ GeV}$

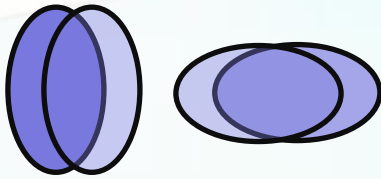
# Data Studied

➤ Collected data for Au+Au at  $\sqrt{s_{NN}} = 200 - 7.7 \text{ GeV}$

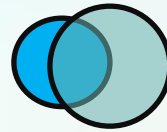
➤ Collected data for different systems at  $\sqrt{s_{NN}} \sim 200$



Au + Au



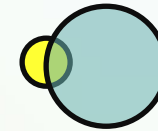
U + U



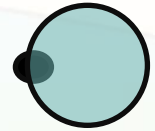
Cu + Au



Cu + Cu



d + Au

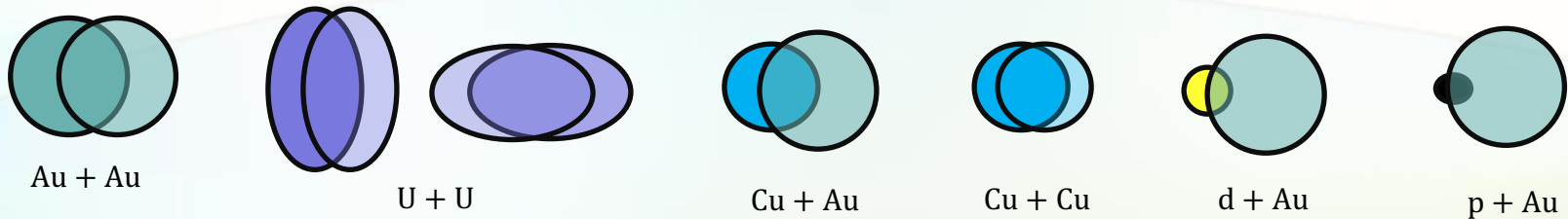


p + Au

# Data Studied

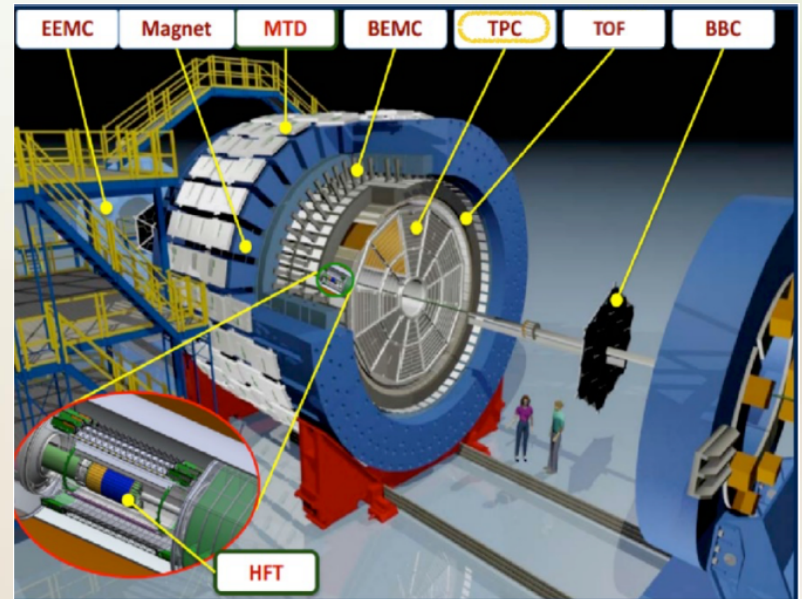
➤ Collected data for Au+Au at  $\sqrt{s_{NN}} = 200 - 7.7 \text{ GeV}$

➤ Collected data for different systems at  $\sqrt{s_{NN}} \sim 200$



## STAR Detector at RHIC

➤ TPC detector mainly get used in the current analysis





# Azimuthal anisotropy measurements

## Correlation function

Two-particle correlation function  $Cr(\Delta\varphi = \varphi_a - \varphi_b)$ ,

$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

# Azimuthal anisotropy measurements

## Correlation function

Two-particle correlation function  $Cr(\Delta\varphi = \varphi_a - \varphi_b)$ ,

$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

Flow

Non-flow

# Azimuthal anisotropy measurements

## Correlation function

Two-particle correlation function  $Cr(\Delta\varphi = \varphi_a - \varphi_b)$ ,

$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

$$n > 1$$
$$v_n^{ab} = v_n^a v_n^b + \delta_{short}$$

$$n = 1$$
$$v_1^{ab} = v_1^a v_1^b + \delta_{long}$$

Flow

Non-flow

# Azimuthal anisotropy measurements

## Correlation function

Two-particle correlation function  $Cr(\Delta\varphi = \varphi_a - \varphi_b)$ ,

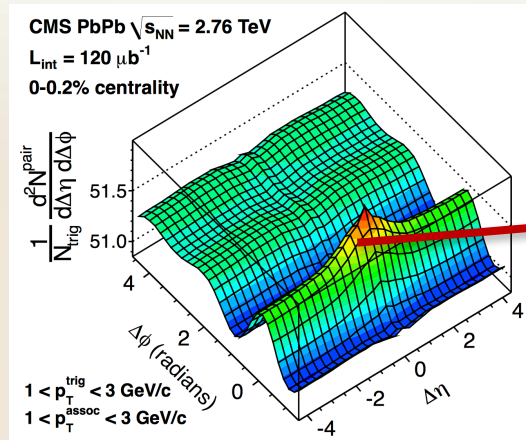
$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

$$n > 1 \\ v_n^{ab} = v_n^a v_n^b + \delta_{short}$$

$$n = 1 \\ v_1^{ab} = v_1^a v_1^b + \delta_{long}$$

Flow

Non-flow



Short – range

# Azimuthal anisotropy measurements

## Correlation function

Two-particle correlation function  $Cr(\Delta\varphi = \varphi_a - \varphi_b)$ ,

$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

$$n > 1$$

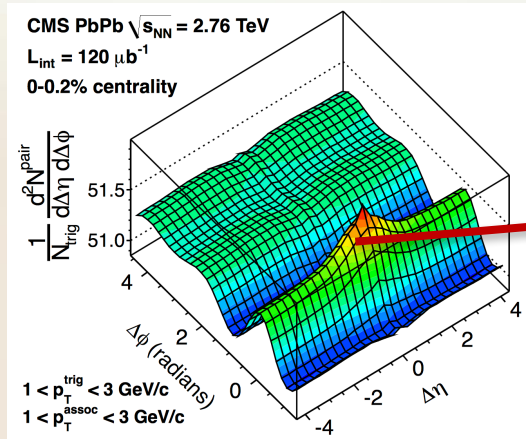
$$v_n^{ab} = v_n^a v_n^b + \delta_{short}$$

$$n = 1$$

$$v_1^{ab} = v_1^a v_1^b + \delta_{long}$$

Flow

Non-flow



Short – range

HBT

# Azimuthal anisotropy measurements

## Correlation function

Two-particle correlation function  $Cr(\Delta\varphi = \varphi_a - \varphi_b)$ ,

$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

$$n > 1$$

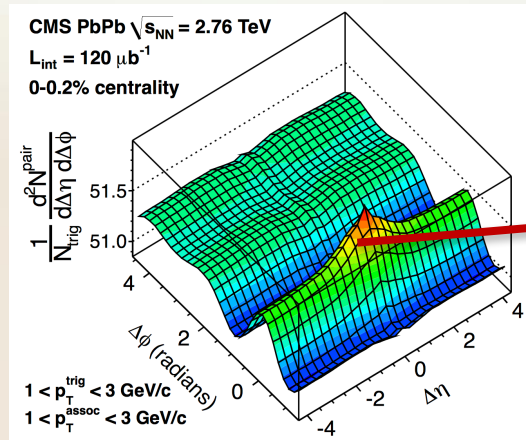
$$v_n^{ab} = v_n^a v_n^b + \delta_{short}$$

$$n = 1$$

$$v_1^{ab} = v_1^a v_1^b + \delta_{long}$$

Flow

Non-flow



Short – range

HBT

Decay

Charge

# Azimuthal anisotropy measurements

## Correlation function

Two-particle correlation function  $Cr(\Delta\varphi = \varphi_a - \varphi_b)$ ,

$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

$$n > 1$$

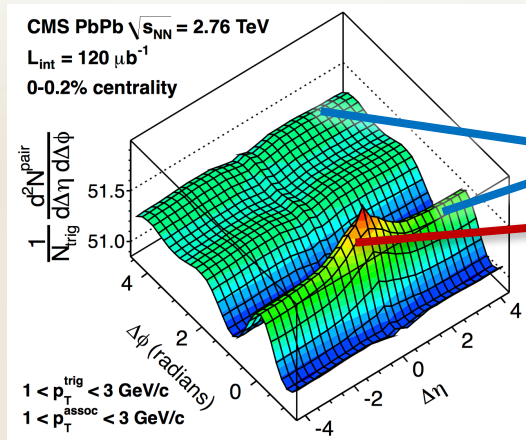
$$v_n^{ab} = v_n^a v_n^b + \delta_{short}$$

$$n = 1$$

$$v_1^{ab} = v_1^a v_1^b + \delta_{long}$$

Flow

Non-flow



Long – range

Short – range

HBT

Decay

Charge



# Azimuthal anisotropy measurements

## Correlation function

Two-particle correlation function  $Cr(\Delta\varphi = \varphi_a - \varphi_b)$ ,

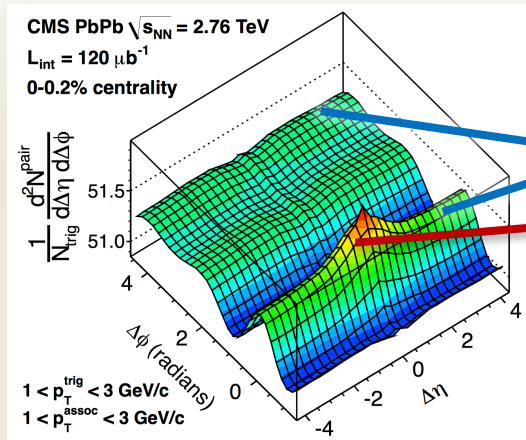
$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

$$n > 1 \\ v_n^{ab} = v_n^a v_n^b + \delta_{short}$$

$$n = 1 \\ v_1^{ab} = v_1^a v_1^b + \delta_{long}$$

Flow

Non-flow



Long – range

Short – range

Momentum Conservation

HBT

Decay

Charge

# Azimuthal anisotropy measurements

## Correlation function

Two-particle correlation function  $Cr(\Delta\varphi = \varphi_a - \varphi_b)$ ,

$$Cr(\Delta\varphi) = dN/d\Delta\varphi \text{ and } v_n^{ab} = \frac{\sum_{\Delta\varphi} Cr(\Delta\varphi) \cos(n \Delta\varphi)}{\sum_{\Delta\varphi} Cr(\Delta\varphi)}$$

$$n > 1$$

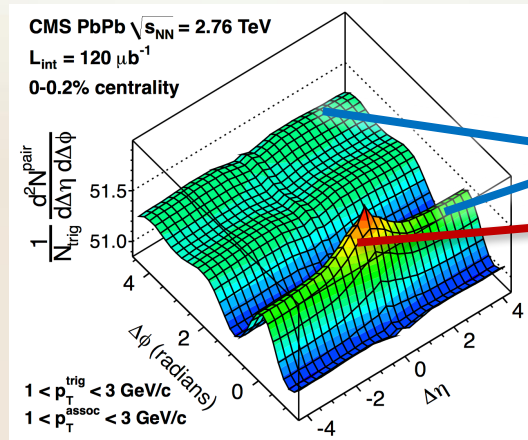
$$v_n^{ab} = v_n^a v_n^b + \delta_{short}$$

$$n = 1$$

$$v_1^{ab} = v_1^a v_1^b + \delta_{long}$$

Flow

Non-flow



Long – range

Short – range

Momentum Conservation

HBT

Di-jets

Decay

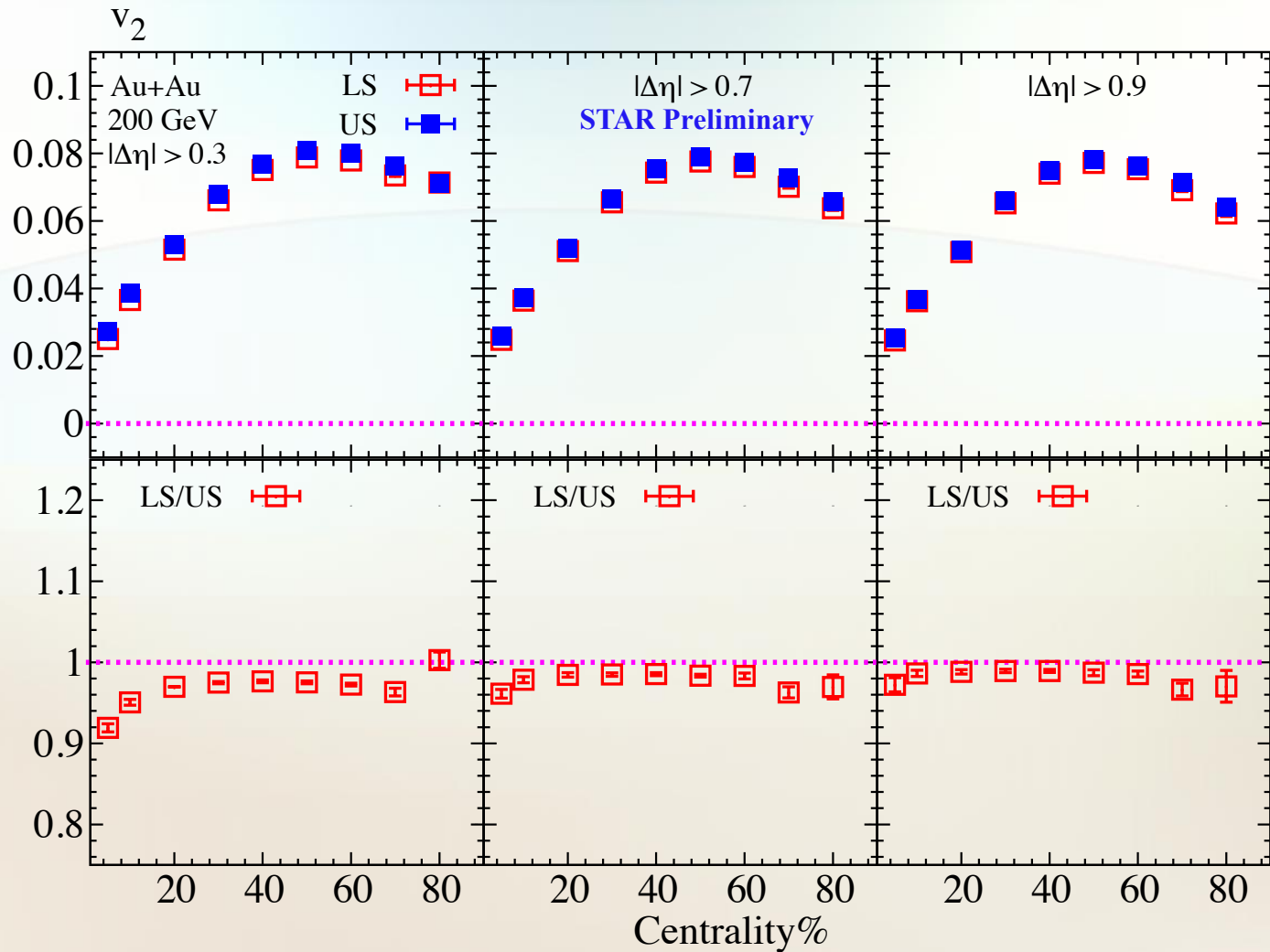
Charge

Non-flow suppression is needed

Short – range  
Non–flow

## Short-range non-flow suppression

The  $v_2$  vs centrality at  $\sqrt{s_{NN}} = 200$  using different  $\Delta\eta$  cuts



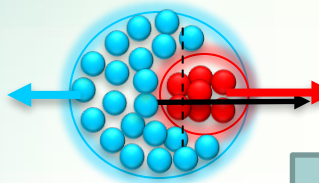
➤ Short-range non-flow effect get reduced using  $|\Delta\eta| > 0.7$  cut

Long – range

# Long-range non-flow suppression

arXiv:1203.0931  
arXiv:1203.3410  
arXiv:1208.1874  
arXiv:1208.1887  
arXiv:1211.7162

$$v_1^{ab} = v_1^a v_1^b + \delta_{long} \quad n = 1$$



1

$$v_{11}(p_T^a, p_T^b) = v_1^{even}(p_T^a) v_1^{even}(p_T^b) - C p_T^a p_T^b$$

$C \propto \langle Mult \rangle^{-1}$

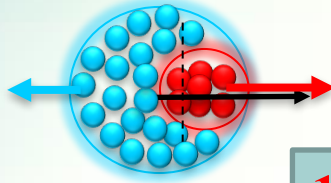
Momentum  
Conservation

Long – range

# Long-range non-flow suppression

arXiv:1203.0931  
arXiv:1203.3410  
arXiv:1208.1874  
arXiv:1208.1887  
arXiv:1211.7162

$$v_1^{ab} = v_1^a v_1^b + \delta_{long} \quad n = 1$$



1

$$v_{11}(p_T^a, p_T^b) = v_1^{even}(p_T^a) v_1^{even}(p_T^b) - C p_T^a p_T^b$$

$C \propto (\langle Mult \rangle)^{-1}$

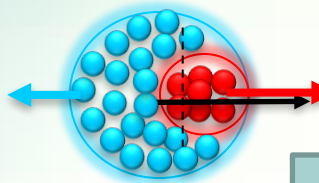
$v_{11}$  in Eq(1) represents NxM matrix which we fit with N+1 parameters

Momentum  
Conservation

# Long-range non-flow suppression

Momentum Conservation

$$v_1^{ab} = v_1^a v_1^b + \delta_{long} \quad n = 1$$

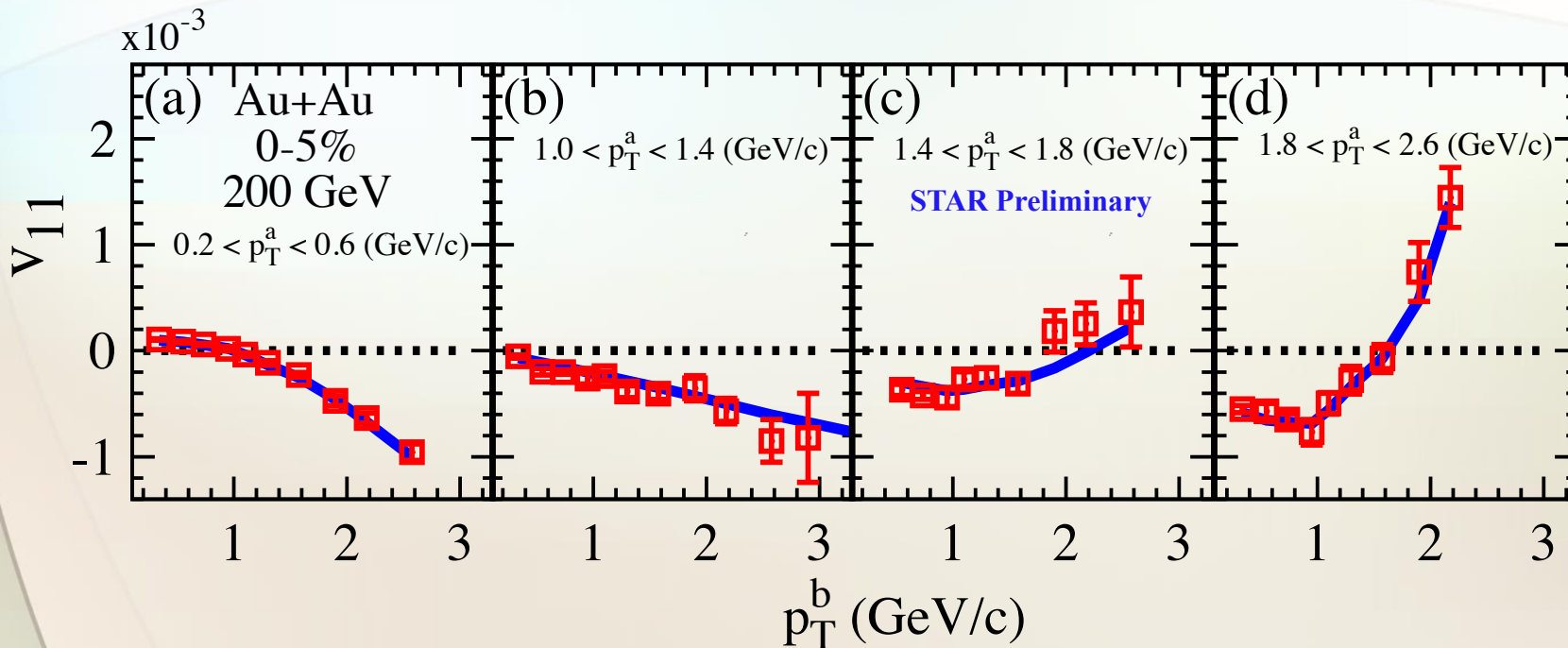


1

$$v_{11}(p_T^a, p_T^b) = v_1^{even}(p_T^a) v_1^{even}(p_T^b) - C p_T^a p_T^b$$

$$C \propto (\langle Mult \rangle)^{-1}$$

$v_{11}$  in Eq(1) represents NxM matrix which we fit with N+1 parameters

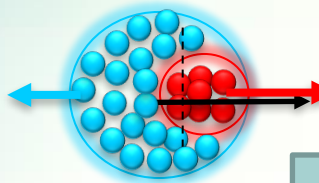


➤ Good simultaneous fit ( $\frac{\chi^2}{ndf} \sim 1.1$ ) obtained with Eq. 1

# Long-range non-flow suppression

Momentum Conservation

$$v_1^{ab} = v_1^a v_1^b + \delta_{long} \quad n = 1$$

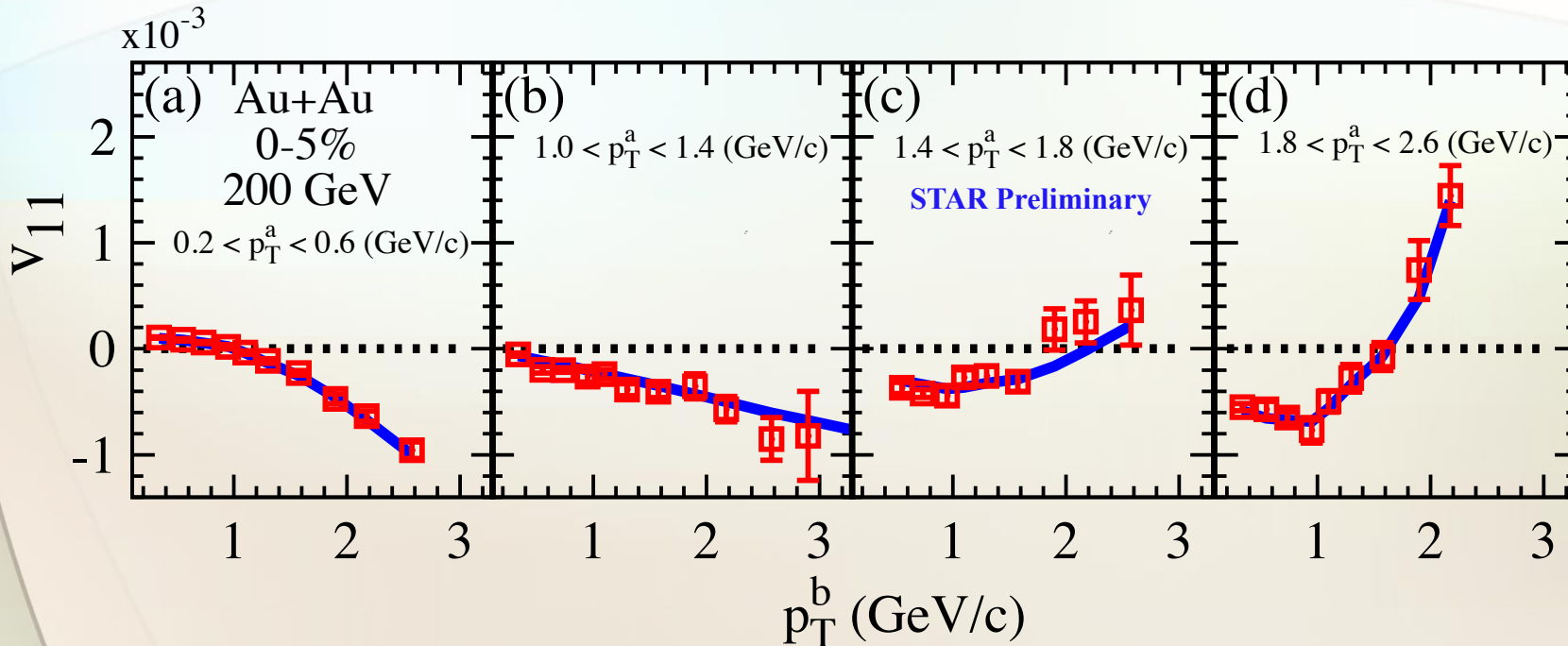


1

$$v_{11}(p_T^a, p_T^b) = v_1^{even}(p_T^a) v_1^{even}(p_T^b) - C p_T^a p_T^b$$

$$C \propto (\langle Mult \rangle)^{-1}$$

$v_{11}$  in Eq(1) represents NxM matrix which we fit with N+1 parameters



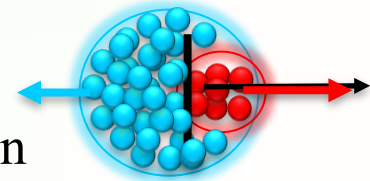
➤ Good simultaneous fit ( $\frac{\chi^2}{ndf} \sim 1.1$ ) obtained with Eq. 1

➤  $v_{11}$  characteristic behavior gives a good constraint for  $v_1^{even}(p_T)$  extraction



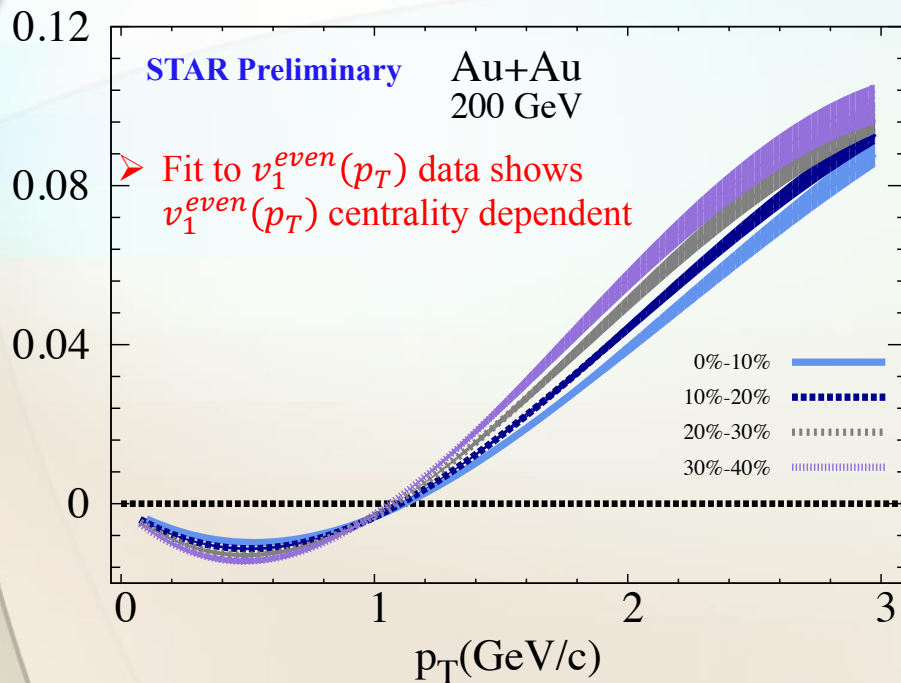
$$v_{11}(p_T^a, p_T^b) = v_1^{even}(p_T^a)v_1^{even}(p_T^b) - C p_T^a p_T^b$$

The extracted  $v_1^{even}(p_T)$  and the momentum conservation parameter  $C$  at  $\sqrt{s_{NN}} = 200$



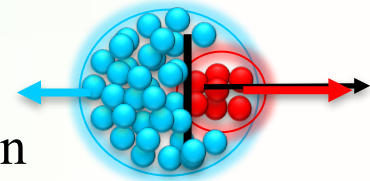
Momentum Conservation

$v_1^{even}$



➤ The characteristic behavior of  $v_1^{even}(p_T)$  shows a weak centrality dependence

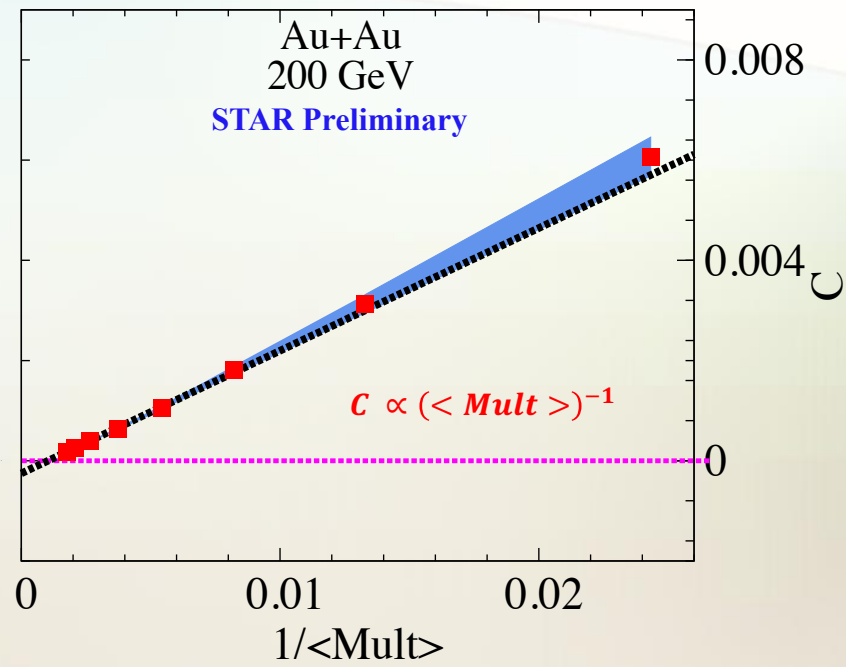
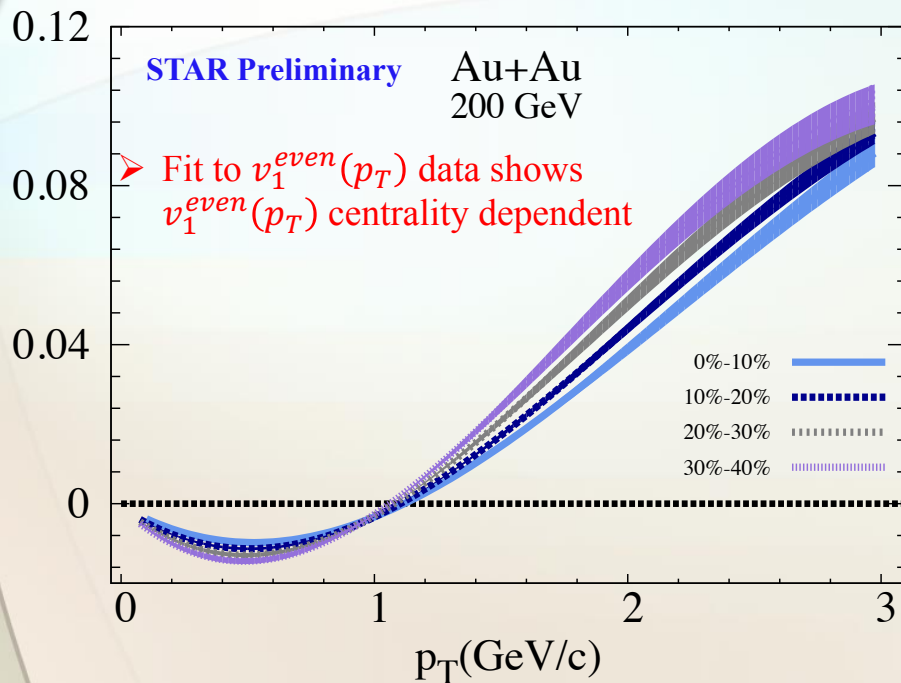
$$v_{11}(p_T^a, p_T^b) = v_1^{even}(p_T^a)v_1^{even}(p_T^b) - C p_T^a p_T^b$$



The extracted  $v_1^{even}(p_T)$  and the momentum conservation parameter  $C$  at  $\sqrt{s_{NN}} = 200$

Momentum Conservation

$v_1^{even}$



➤ The characteristic behavior of  $v_1^{even}(p_T)$  shows a weak centrality dependence

➤ The momentum conservation parameter  $C$  scales as  $\langle \text{Mult} \rangle^{-1}$

# Non-flow suppression

*Long – range*

*Short – range*

*Momentum  
Conservation*

*HBT*

*Di-jets*

*Decay*

- $\eta$  gap between particles in each pair used to suppress the short-range non-flow

# Non-flow suppression

*Long – range*

*Short – range*

*Momentum  
Conservation*

*HBT*

*Di-jets*

*Decay*

- $\eta$  gap between particles in each pair used to suppress the short-range non-flow
- Simultaneous fit used to suppress long-range non-flow associated with momentum conservation

# Non-flow suppression

*Long – range*

*Short – range*

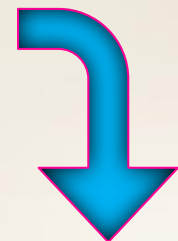
*Momentum  
Conservation*

*HBT*

*Di-jets*

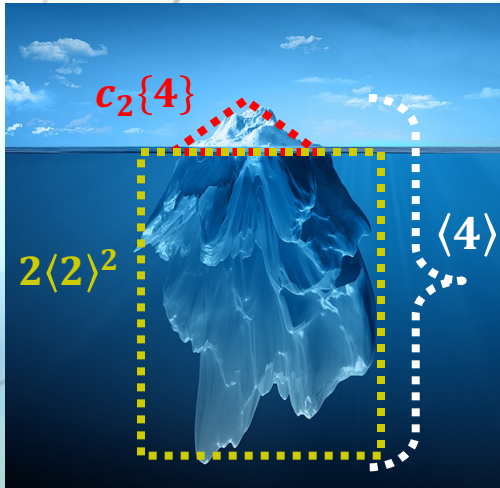
*Decay*

- $\eta$  gap between particles in each pair used to suppress the short-range non-flow
- Simultaneous fit used to suppress long-range non-flow associated with momentum conservation
- Multi-particle correlations also used to further suppress non-flow



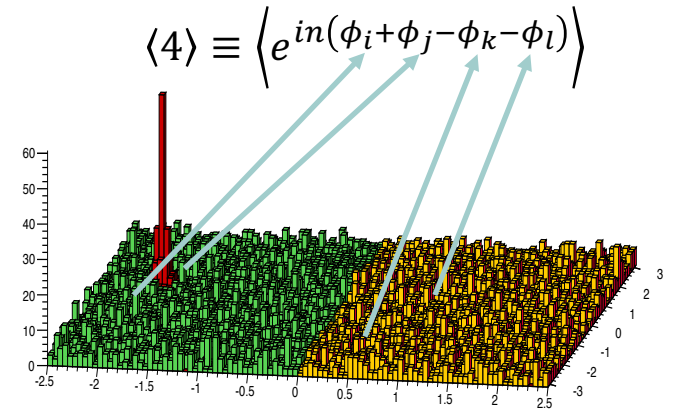
# Multi-particle correlations and the non-flow suppression

$$c_2\{4\} \equiv \langle 4 \rangle - 2\langle 2 \rangle^2$$



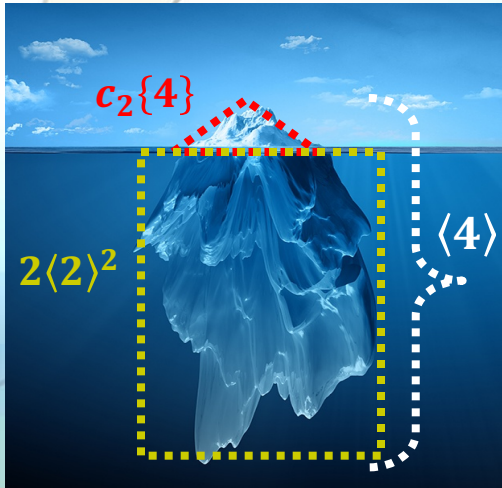
In the subevent method,  
particles are correlated  
across all subevents  
(long-range)

Short-range non-flow dominate



# Multi-particle correlations and the non-flow suppression

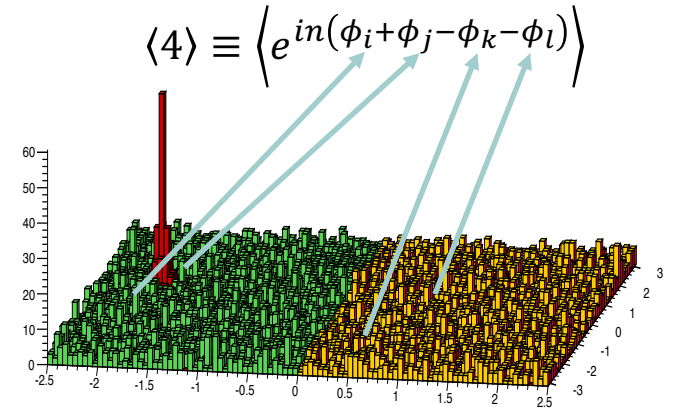
$$c_2\{4\} \equiv \langle 4 \rangle - 2\langle 2 \rangle^2$$



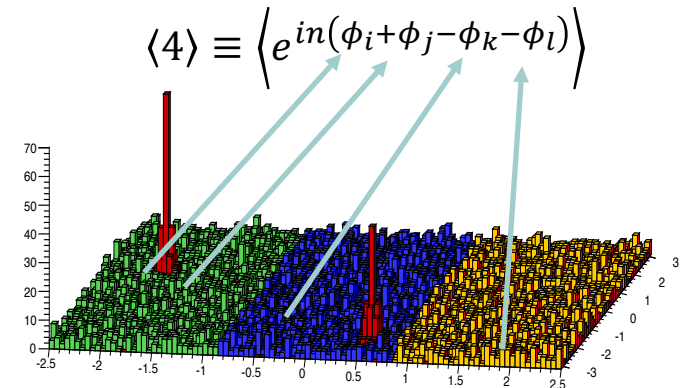
In the subevent method, particles are correlated across all subevents (long-range)

Three subevent cumulant can further suppress away-side jet contribution

Short-range non-flow dominate



Long-range non-flow dominate

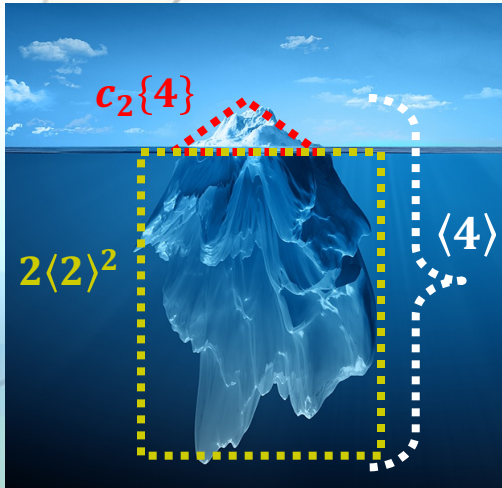


[arXiv: 1701.03830](https://arxiv.org/abs/1701.03830)



# Multi-particle correlations and the non-flow suppression

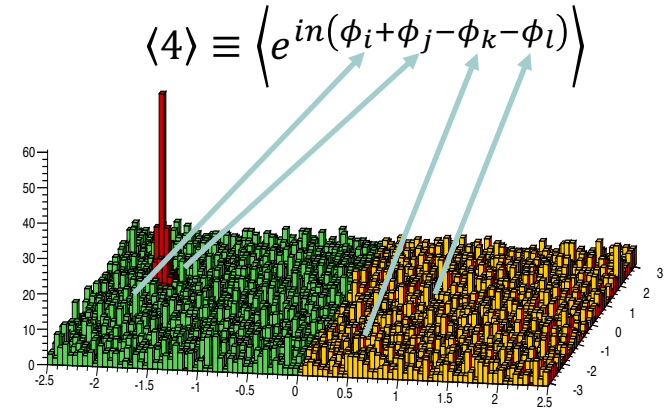
$$c_2\{4\} \equiv \langle 4 \rangle - 2\langle 2 \rangle^2$$



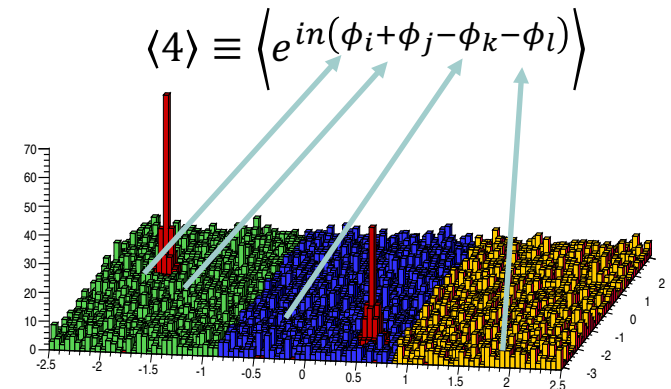
In the subevent method, particles are correlated across all subevents (long-range)

Three subevent cumulant can further suppress away-side jet contribution

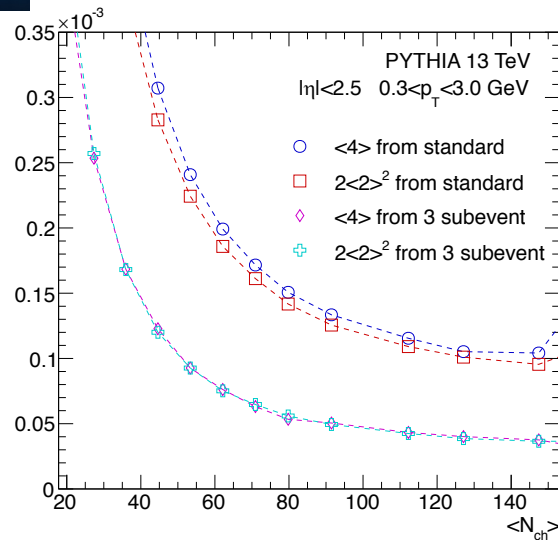
Short-range non-flow dominate



Long-range non-flow dominate



Both  $\langle 4 \rangle$  and  $2\langle 2 \rangle^2$  are much smaller with subevent



[arXiv: 1701.03830](https://arxiv.org/abs/1701.03830)

Subevent cumulant measures long-range collectivity.

# Results

## Au+Au Beam Energy Scan

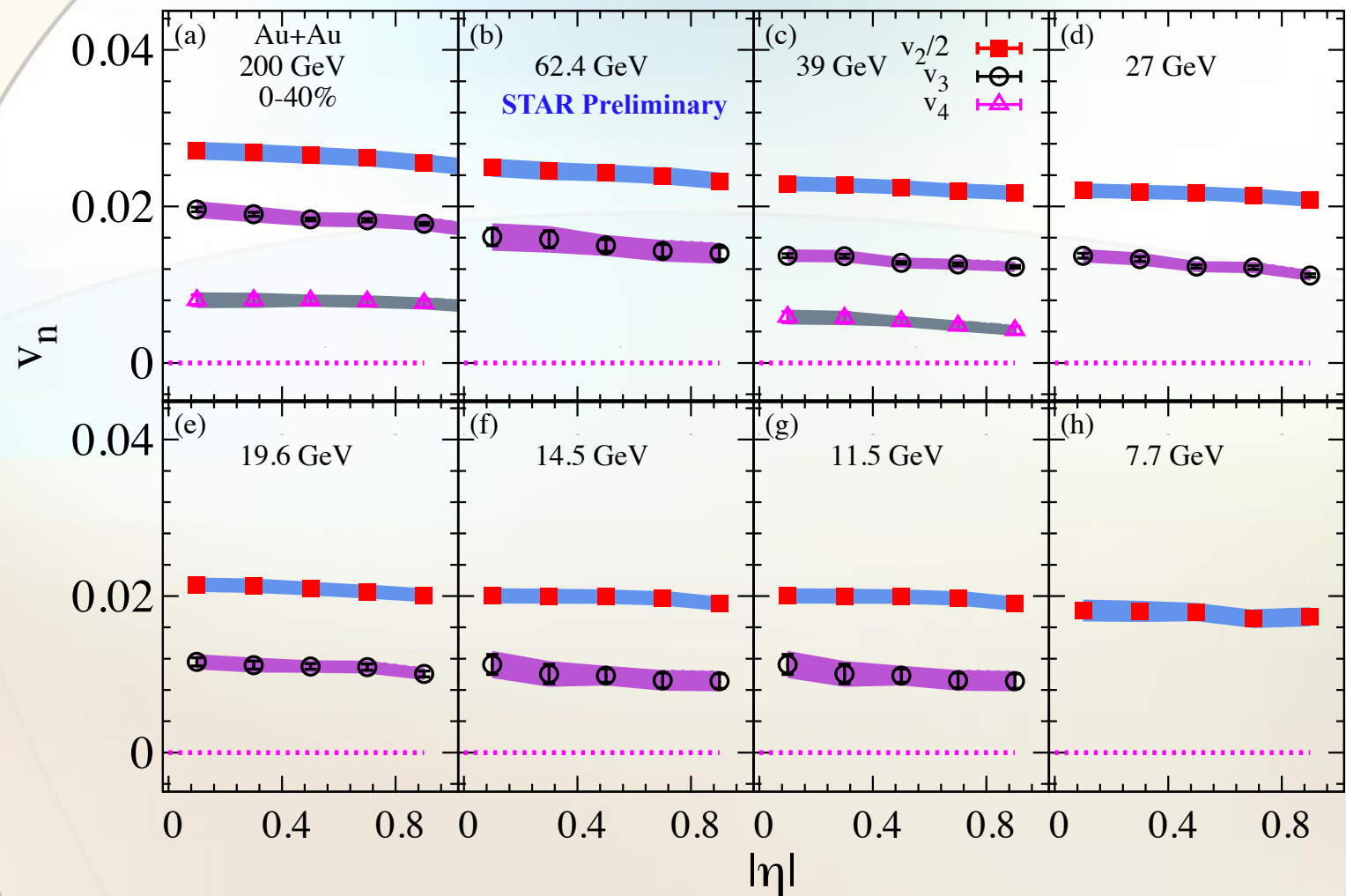
$v_n(\eta)$

$v_n(\sqrt{s_{NN}})$

# Pseudorapidity dependence of $v_{n>1}$

$|\eta| < 1$  and  $|\Delta\eta| > 0.7$   
 $0.2 < p_T < 4 \text{ GeV}/c$

The extracted  $v_{n>1}(\eta)$  at all BES energies



➤  $v_n(\eta)$  has similar trends for different beam energies.

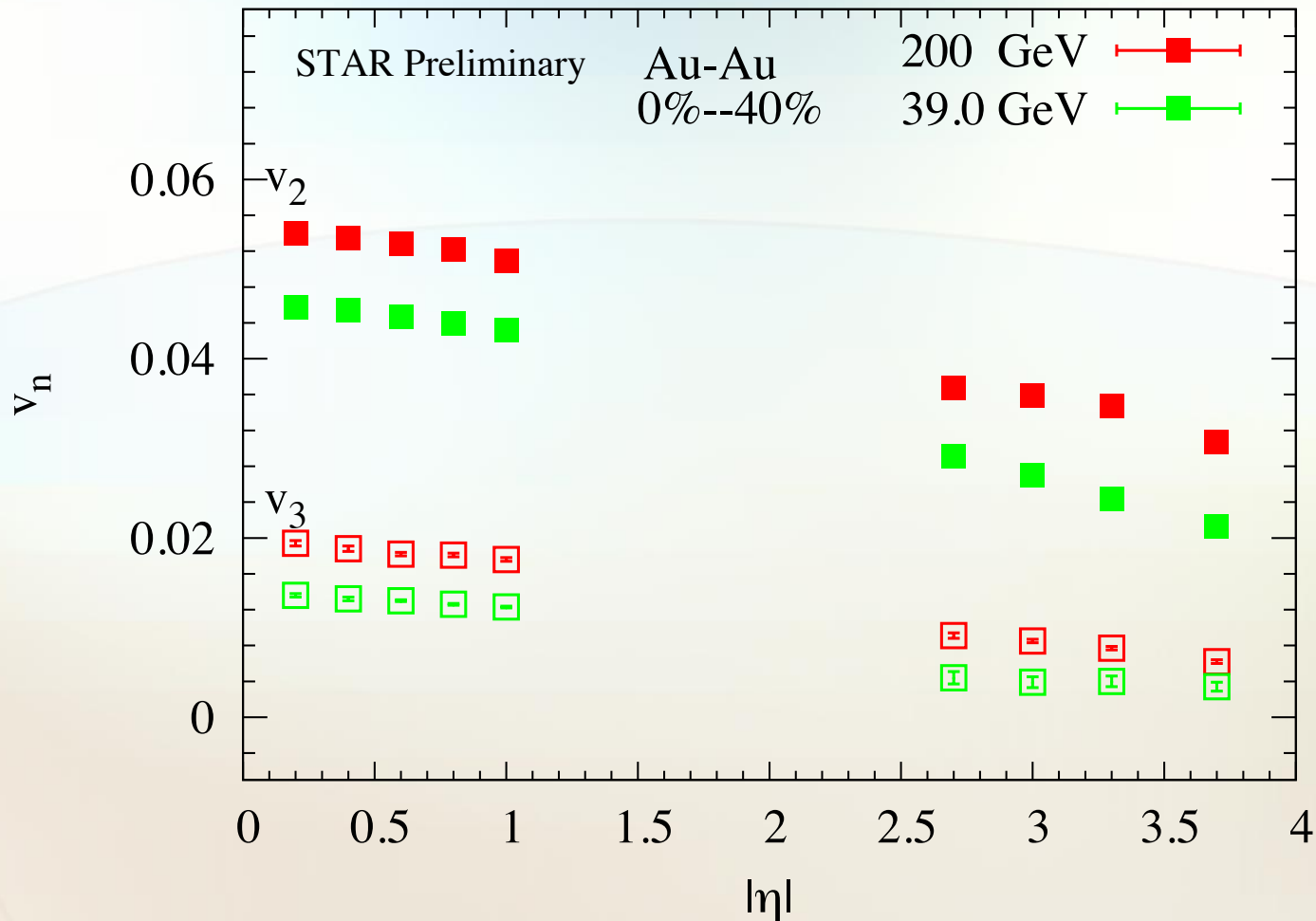
➤  $v_n(\eta)$  decreases with harmonic order  $n$ .

# Pseudorapidity dependence of $v_{n>1}$

$|\eta| < 1$  and  $|\Delta\eta| > 0.7$

$0.2 < p_T < 4 \text{ GeV}/c$

The extracted  $v_{n>1}(\eta)$

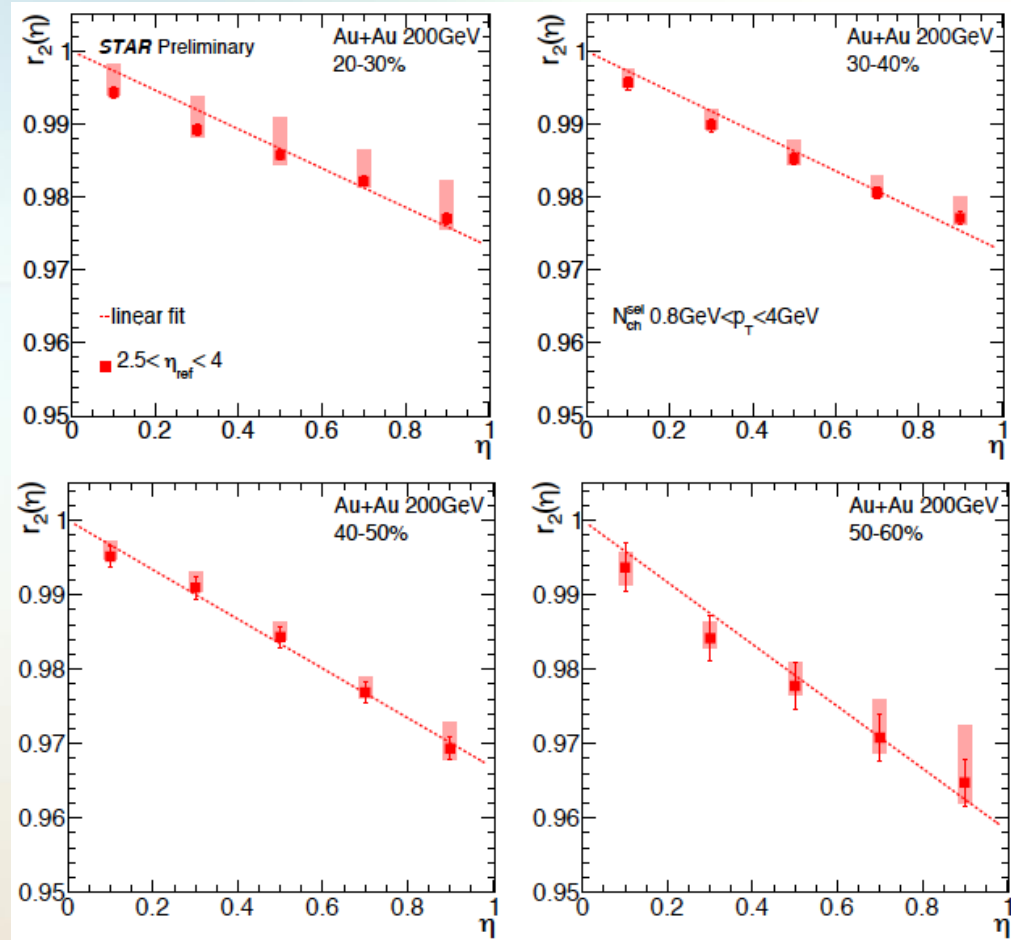


- $v_n(\eta)$  has similar trends for different beam energies.
- $v_n(\eta)$  decreases with harmonic order  $n$ .

# Longitudinal decorrelation of $v_2$ in Au+Au 200 GeV

$$r_n(\eta) = \frac{\langle v_n(-\eta)v_n^*(\eta_{\text{ref}}) \rangle}{\langle v_n(\eta)v_n^*(\eta_{\text{ref}}) \rangle}$$

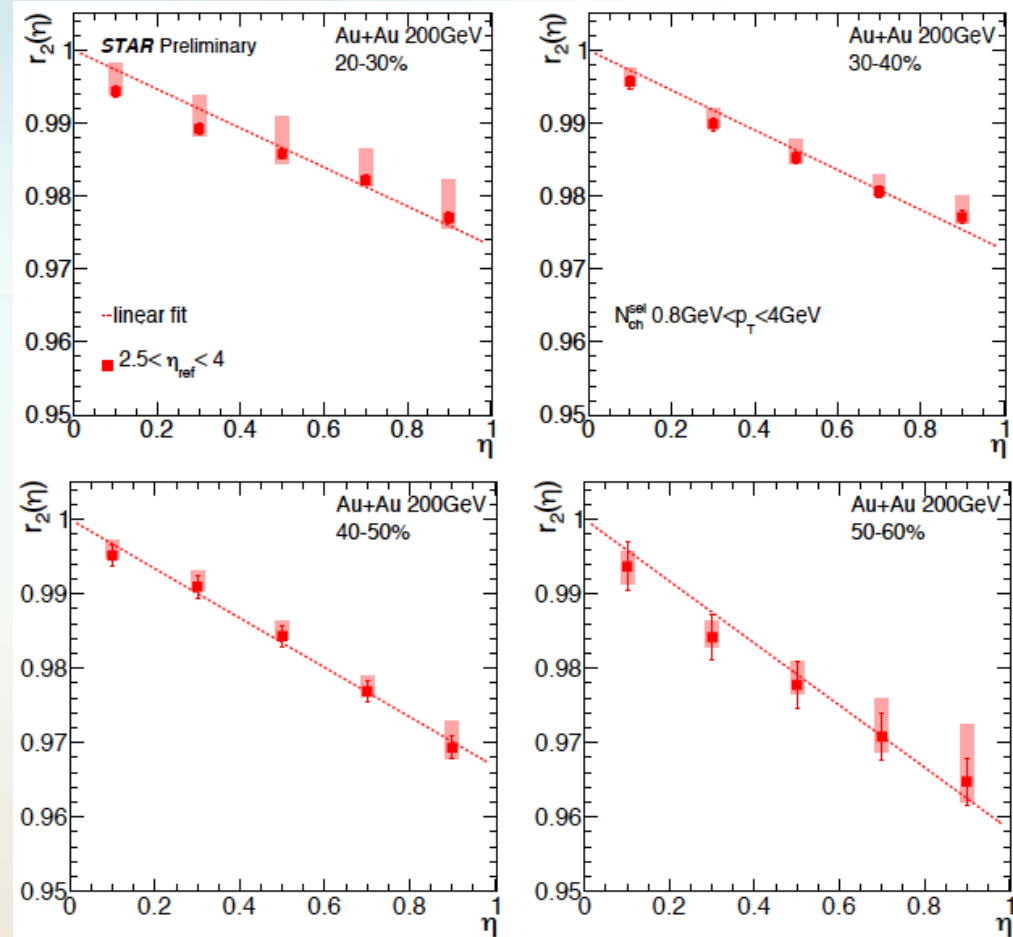
- $r_n(\eta)$  decrease linearly for different centrality studied



# Longitudinal decorrelation of $v_2$ in Au+Au 200 GeV

$$r_n(\eta) = \frac{\langle v_n(-\eta)v_n^*(\eta_{\text{ref}}) \rangle}{\langle v_n(\eta)v_n^*(\eta_{\text{ref}}) \rangle}$$

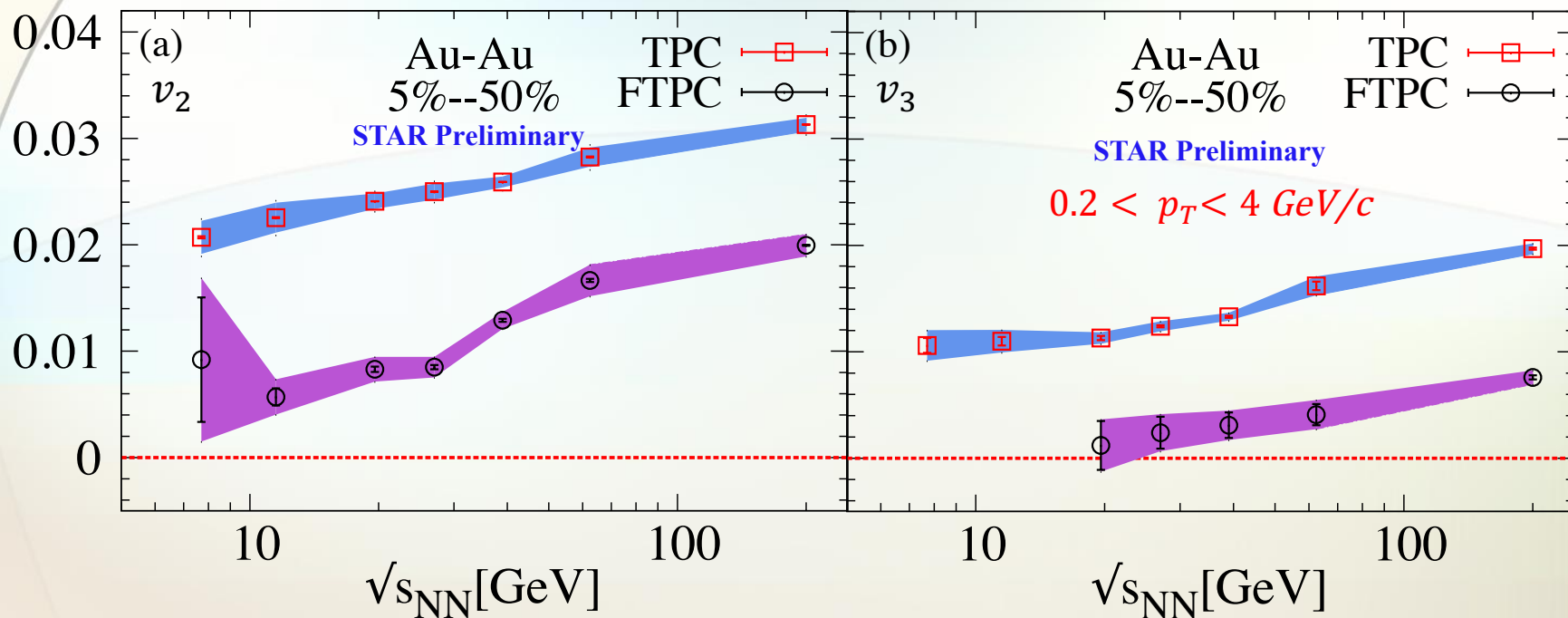
- $r_n(\eta)$  decrease linearly for different centrality studied
- The decorrelation effect gets stronger as the collision become more peripheral



- More information will be given today by Maowu Nie

# Beam-energy dependence of $v_{n>1}$

The extracted  $v_{n>1}$  vs  $\sqrt{s_{NN}}$  for TPC ( $|\eta| < 1$ ) and FTPC ( $2.5 < |\eta| < 4$ )



➤ At mid and forward rapidity;

- ✓  $v_n(\sqrt{s_{NN}})$  shows a monotonic increase with beam-energy.
- ✓  $v_n(\sqrt{s_{NN}})$  decreases with harmonic order  $n$  (viscous effects).



# Summary-1

- $\eta$  gap between particles in each pair used to suppress the short-range non-flow

# Summary-1

- $\eta$  gap between particles in each pair used to suppress the short-range non-flow
- Simultaneous fit used to suppress long-range non-flow associated with momentum conservation

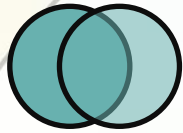
# Summary-1

- $\eta$  gap between particles in each pair used to suppress the short-range non-flow
- Simultaneous fit used to suppress long-range non-flow associated with momentum conservation
- The longitudinal decorrelation of  $v_2$  was observed for Au+Au at 200 GeV

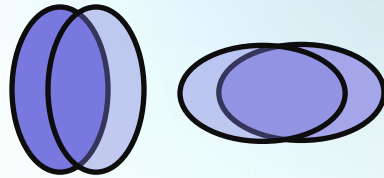
# Summary-1

- $\eta$  gap between particles in each pair used to suppress the short-range non-flow
- Simultaneous fit used to suppress long-range non-flow associated with momentum conservation
- The longitudinal decorrelation of  $v_2$  was observed for Au+Au at 200 GeV
- At mid and forward rapidity;
  - ✓  $v_n(\sqrt{s_{NN}})$  shows a monotonic increase with beam-energy.
  - ✓  $v_n(\sqrt{s_{NN}})$  decreases with harmonic order  $n$  (viscous effects).

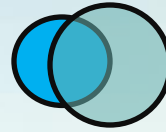
# Collectivity in small systems



Au + Au



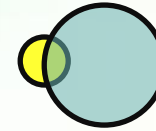
U + U



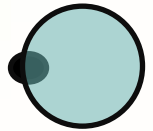
Cu + Au



Cu + Cu



d + Au



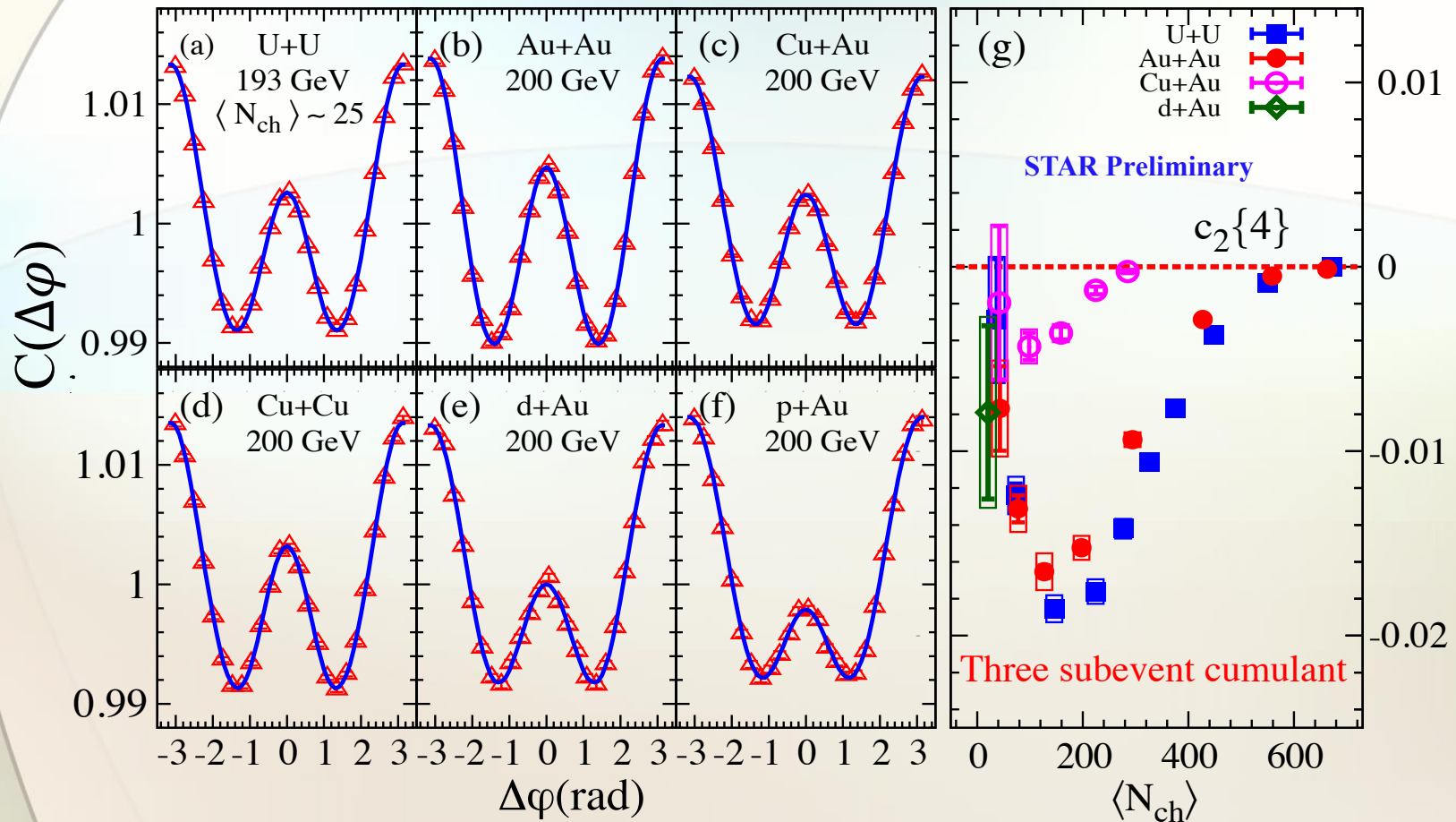
p + Au

$$v_n(p_T)$$

$$v_n\{2,4\}(N_{ch})$$

# Collectivity in small systems

- The correlation function  $C(\Delta\varphi)$  for all systems at one  $\langle N_{Ch} \rangle$  value
  - $c_2\{4\}$  vs  $\langle N_{Ch} \rangle$  for 4 systems  $\times 10^{-3}$



- $C(\Delta\varphi)$  shows similar trend for all systems
- $c_2\{4\}$  shows negative value for different systems

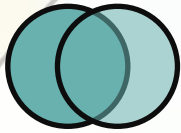
# Acoustic ansatz

PRC 84, 034908 (2011)  
P. Staig and E. Shuryak.

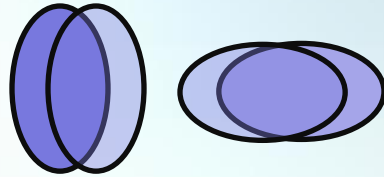
arXiv:1305.3341  
Roy A. Lacey, et al.

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

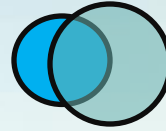
arXiv:1601.06001  
Roy A. Lacey, et al.



Au + Au



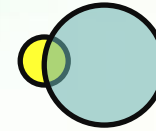
U + U



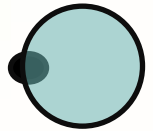
Cu + Au



Cu + Cu



d + Au



p + Au

- $v_n$  measurements for different systems are sensitive to system shape ( $\varepsilon_n$ ), dimensionless size ( $RT$ ) and transport coefficients  $\left(\frac{\eta}{s}, \frac{\zeta}{s}, \dots\right)$ .



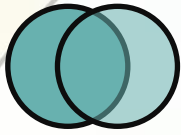
# Acoustic ansatz

PRC 84, 034908 (2011)  
P. Staig and E. Shuryak.

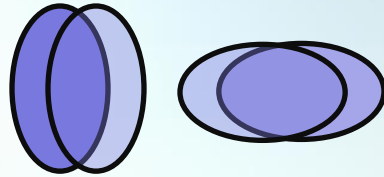
arXiv:1305.3341  
Roy A. Lacey, et al.

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

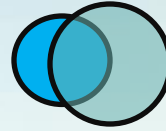
arXiv:1601.06001  
Roy A. Lacey, et al.



Au + Au



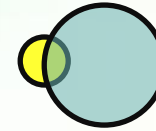
U + U



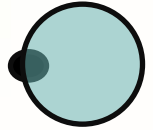
Cu + Au



Cu + Cu



d + Au



p + Au

- $v_n$  measurements for different systems are sensitive to system shape ( $\epsilon_n$ ), dimensionless size ( $RT$ ) and transport coefficients  $\left(\frac{\eta}{s}, \frac{\zeta}{s}, \dots\right)$ .

$$v_n/\epsilon_n \propto e^{-A \left(\frac{\eta}{s} \frac{n^2}{RT}\right)}$$

$A$  is a constant

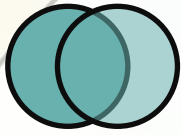
# Acoustic ansatz

PRC 84, 034908 (2011)  
P. Staig and E. Shuryak.

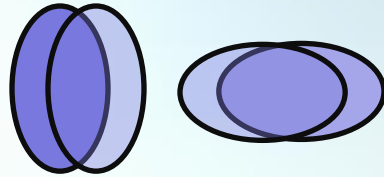
arXiv:1305.3341  
Roy A. Lacey, et al.

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

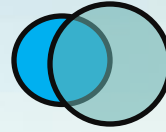
arXiv:1601.06001  
Roy A. Lacey, et al.



Au + Au



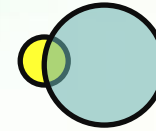
U + U



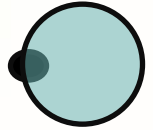
Cu + Au



Cu + Cu



d + Au



p + Au

- $v_n$  measurements for different systems are sensitive to system shape ( $\epsilon_n$ ), dimensionless size ( $RT$ ) and transport coefficients  $\left(\frac{\eta}{s}, \frac{\zeta}{s}, \dots\right)$ .

$$v_n/\epsilon_n \propto e^{-A \left(\frac{\eta}{s} \frac{n^2}{RT}\right)}$$

$A$  is a constant

$$S \sim (RT)^3 \sim \langle N_{Ch} \rangle \text{ then } RT \sim \langle N_{Ch} \rangle^{1/3}$$

$$\ln \left(\frac{v_n}{\epsilon_n}\right) \propto -A \left(\frac{\eta}{s}\right) \langle N_{Ch} \rangle^{-1/3}$$

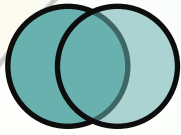
# Acoustic ansatz

PRC 84, 034908 (2011)  
P. Staig and E. Shuryak.

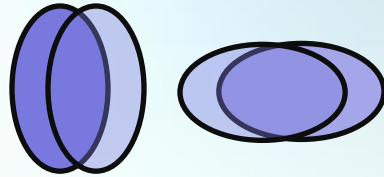
arXiv:1305.3341  
Roy A. Lacey, et al.

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

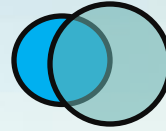
arXiv:1601.06001  
Roy A. Lacey, et al.



Au + Au



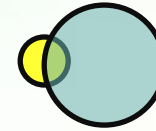
U + U



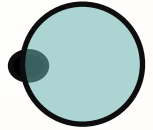
Cu + Au



Cu + Cu



d + Au



p + Au

➤  $v_n$  measurements for different systems are sensitive to system shape ( $\epsilon_n$ ), dimensionless size ( $RT$ ) and transport coefficients  $\left(\frac{\eta}{s}, \frac{\zeta}{s}, \dots\right)$ .

$$v_n/\epsilon_n \propto e^{-A \left(\frac{\eta}{s} \frac{n^2}{RT}\right)}$$

$A$  is a constant

$$S \sim (RT)^3 \sim \langle N_{Ch} \rangle \text{ then } RT \sim \langle N_{Ch} \rangle^{1/3}$$

$$\ln\left(\frac{v_n}{\epsilon_n}\right) \propto -A \left(\frac{\eta}{s}\right) \langle N_{Ch} \rangle^{-1/3}$$

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

At the same  $\frac{\eta}{s}$  and  $\langle N_{Ch} \rangle^{-1/3}$

driven by

$$v_n \longrightarrow \epsilon_n + \dots$$

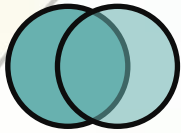
# Acoustic ansatz

PRC 84, 034908 (2011)  
P. Staig and E. Shuryak.

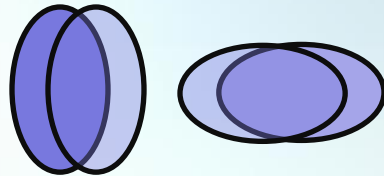
arXiv:1305.3341  
Roy A. Lacey, et al.

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

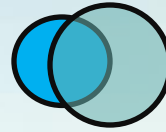
arXiv:1601.06001  
Roy A. Lacey, et al.



Au + Au



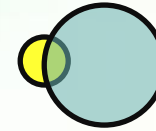
U + U



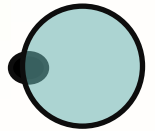
Cu + Au



Cu + Cu



d + Au



p + Au

➤  $v_n$  measurements for different systems are sensitive to system shape ( $\epsilon_n$ ), dimensionless size ( $RT$ ) and transport coefficients ( $\frac{\eta}{s}, \frac{\zeta}{s}, \dots$ ).

$$v_n / \epsilon_n \propto e^{-A \left( \frac{\eta}{s} \frac{n^2}{RT} \right)}$$

$A$  is a constant

$$S \sim (RT)^3 \sim \langle N_{Ch} \rangle \text{ then } RT \sim \langle N_{Ch} \rangle^{1/3}$$

$$\ln \left( \frac{v_n}{\epsilon_n} \right) \propto -A \left( \frac{\eta}{s} \right) \langle N_{Ch} \rangle^{-1/3}$$

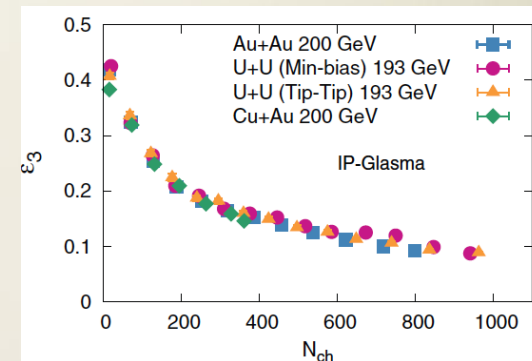
PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

At the same  $\frac{\eta}{s}$  and  $\langle N_{Ch} \rangle^{-1/3}$

$$v_n \xrightarrow{\text{driven by}} \epsilon_n + \dots$$

Odd Harmonic  $v_3$

$$\epsilon_3 \propto \frac{1}{\sqrt{N}}$$



B.Schenke, et al.  
PRC 89, 064908 (2014)

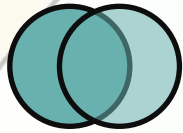
# Acoustic ansatz

PRC 84, 034908 (2011)  
P. Staig and E. Shuryak.

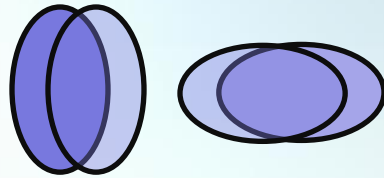
arXiv:1305.3341  
Roy A. Lacey, et al.

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

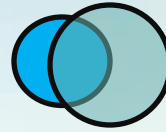
arXiv:1601.06001  
Roy A. Lacey, et al.



Au + Au



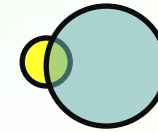
U + U



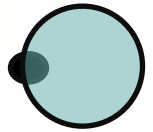
Cu + Au



Cu + Cu



d + Au



p + Au

➤  $v_n$  measurements for different systems are sensitive to system shape ( $\epsilon_n$ ), dimensionless size ( $RT$ ) and transport coefficients ( $\frac{\eta}{s}, \frac{\zeta}{s}, \dots$ ).

$$v_n / \epsilon_n \propto e^{-A \left( \frac{\eta}{s} \frac{n^2}{RT} \right)}$$

$A$  is a constant

$$S \sim (RT)^3 \sim \langle N_{Ch} \rangle \text{ then } RT \sim \langle N_{Ch} \rangle^{1/3}$$

$$\ln \left( \frac{v_n}{\epsilon_n} \right) \propto -A \left( \frac{\eta}{s} \right) \langle N_{Ch} \rangle^{-1/3}$$

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

Even Harmonic  $v_2$

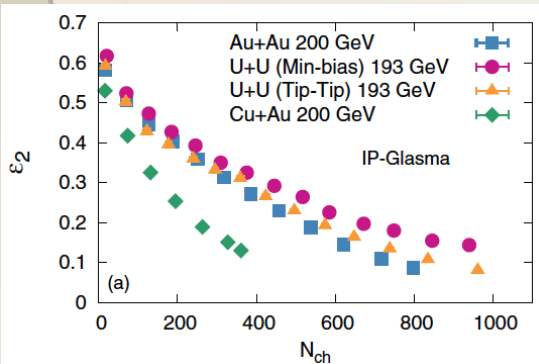
$\epsilon_2$  scaling is needed

At the same  $\frac{\eta}{s}$  and  $\langle N_{Ch} \rangle^{-1/3}$

$$v_n \xrightarrow{\text{driven by}} \epsilon_n + \dots$$

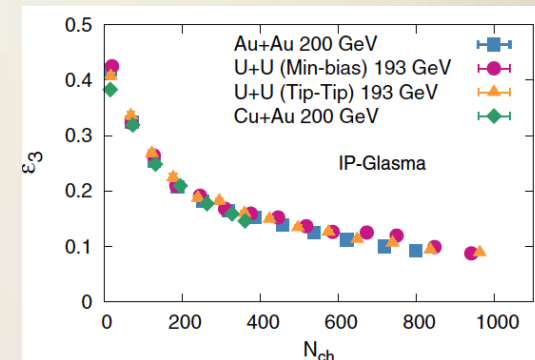
Odd Harmonic  $v_3$

$$\epsilon_3 \propto \frac{1}{\sqrt{N}}$$



B.Schenke, et al.

PRC 89, 064908 (2014)



B.Schenke, et al.

PRC 89, 064908 (2014)

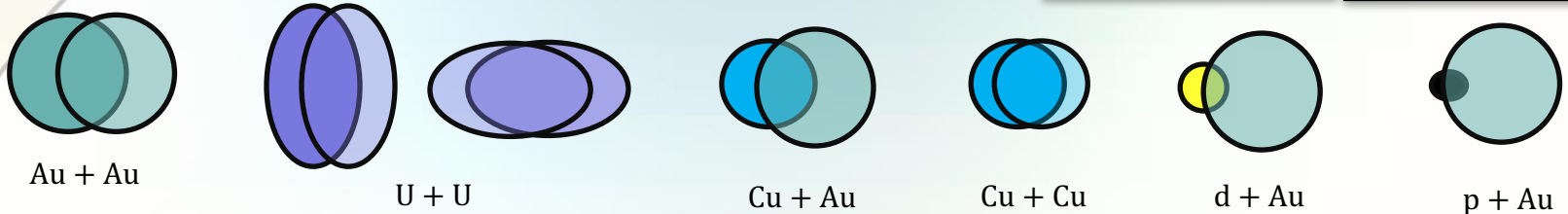
# Acoustic ansatz

PRC 84, 034908 (2011)  
P. Staig and E. Shuryak.

arXiv:1305.3341  
Roy A. Lacey, et al.

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

arXiv:1601.06001  
Roy A. Lacey, et al.



➤  $v_n$  measurements for different systems are sensitive to system shape ( $\epsilon_n$ ), dimensionless size ( $RT$ ) and transport coefficients ( $\frac{\eta}{s}, \frac{\zeta}{s}, \dots$ ).

$$v_n/\epsilon_n \propto e^{-A \left( \frac{\eta}{s} \frac{n^2}{RT} \right)}$$

$A$  is a constant

$$S \sim (RT)^3 \sim \langle N_{Ch} \rangle \text{ then } RT \sim \langle N_{Ch} \rangle^{1/3}$$

$$\ln \left( \frac{v_n}{\epsilon_n} \right) \propto -A \left( \frac{\eta}{s} \right) \langle N_{Ch} \rangle^{-1/3}$$

PRC 88, 044915 (2013)  
E. Shuryak and I. Zahed

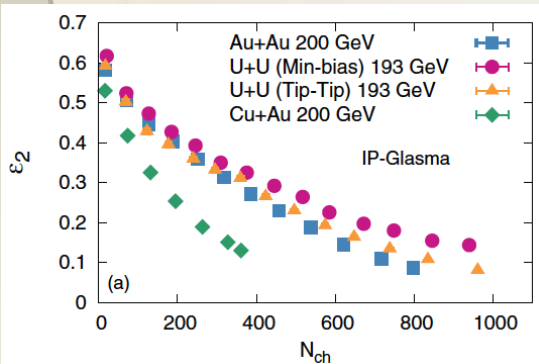
At the same  $\frac{\eta}{s}$  and  $\langle N_{Ch} \rangle^{-1/3}$   
driven by  
 $v_n \longrightarrow \epsilon_n + \dots$

Even Harmonic  $v_2$

$\epsilon_2$  scaling is needed

Odd Harmonic  $v_3$

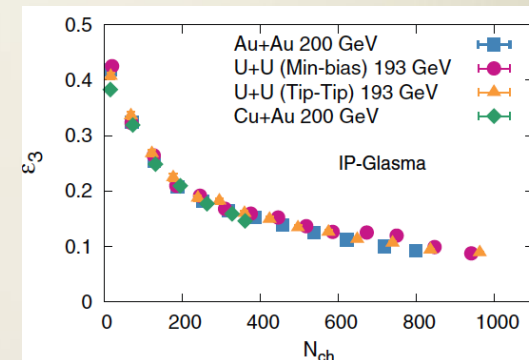
$$\epsilon_3 \propto \frac{1}{\sqrt{N}}$$



B.Schenke, et al.  
PRC 89, 064908 (2014)

**Expectations**

- $v_1^{even}$  and  $v_3$  are system independent
- $v_2$  is system dependent

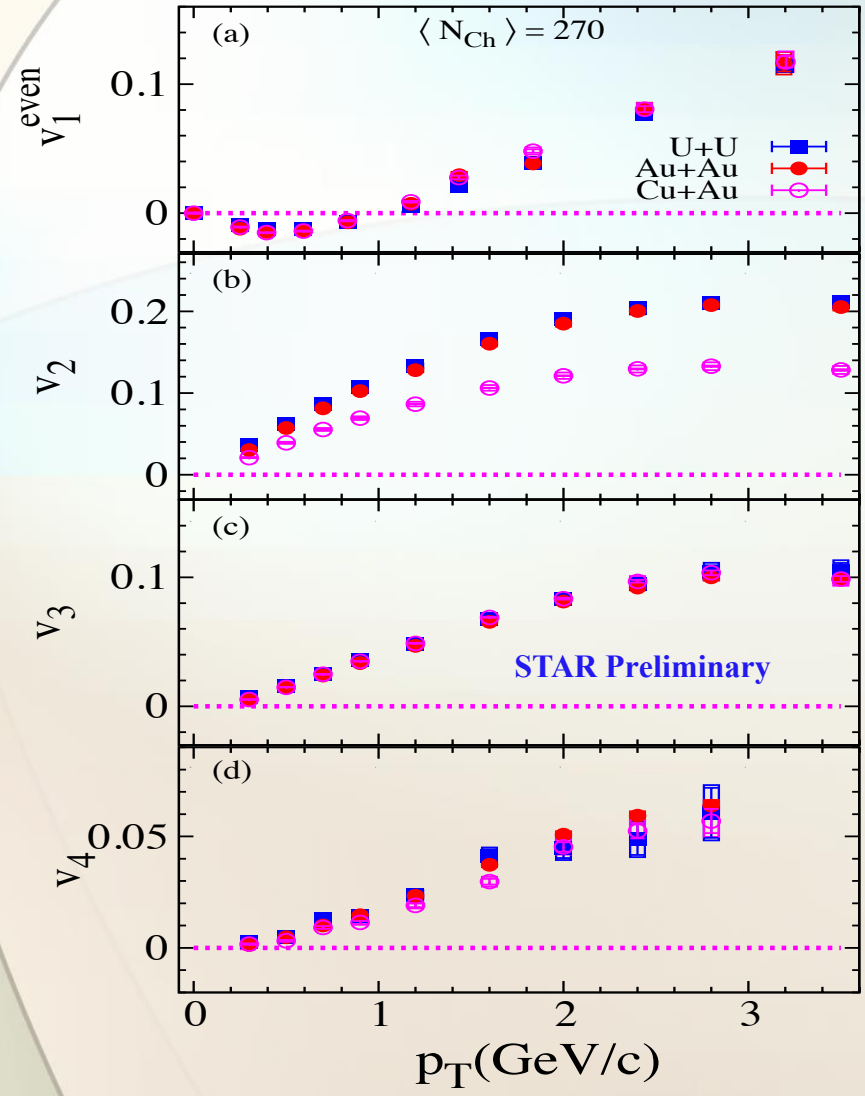


B.Schenke, et al.  
PRC 89, 064908 (2014)

# $v_n$ for large systems (A + B)

$$|\eta| < 1 \text{ and } |\Delta\eta| > 0.7$$

$v_n$  vs  $p_T$  at fixed  $\langle N_{Ch} \rangle = 270$



$$\ln\left(\frac{v_n}{\epsilon_n}\right) \propto -A (\eta/s) \langle N_{Ch} \rangle^{-1/3}$$

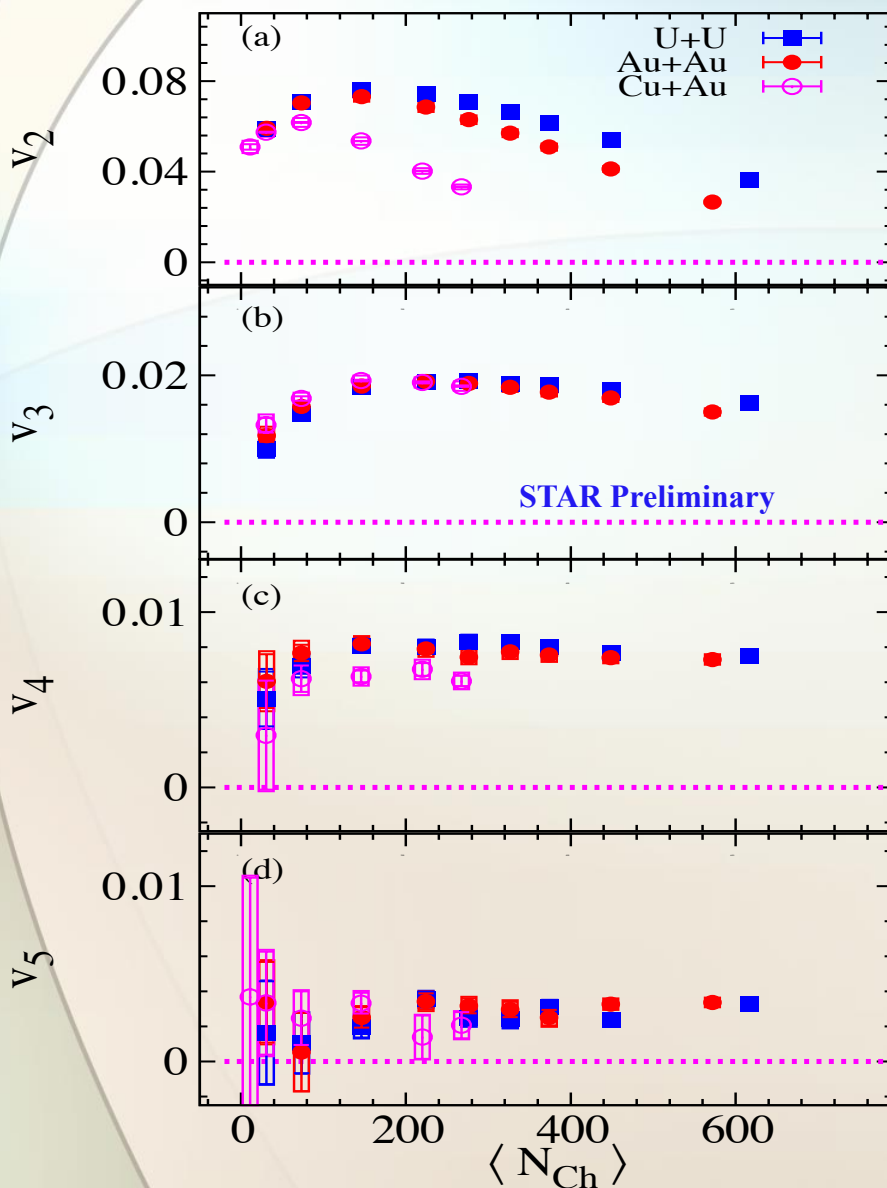
- Odd harmonics are system independent.
- Even harmonics are system dependent, with weak system dependence for the higher harmonics



# $v_n$ for large systems ( $A + B$ )

$|\eta| < 1$  and  $|\Delta\eta| > 0.7$

$v_n$  vs  $\langle N_{Ch} \rangle$  at fixed  $p_T [0.2:4 \text{ GeV}/c]$



$$\ln\left(\frac{v_n}{\varepsilon_n}\right) \propto -A (\eta/s) \langle N_{Ch} \rangle^{-1/3}$$

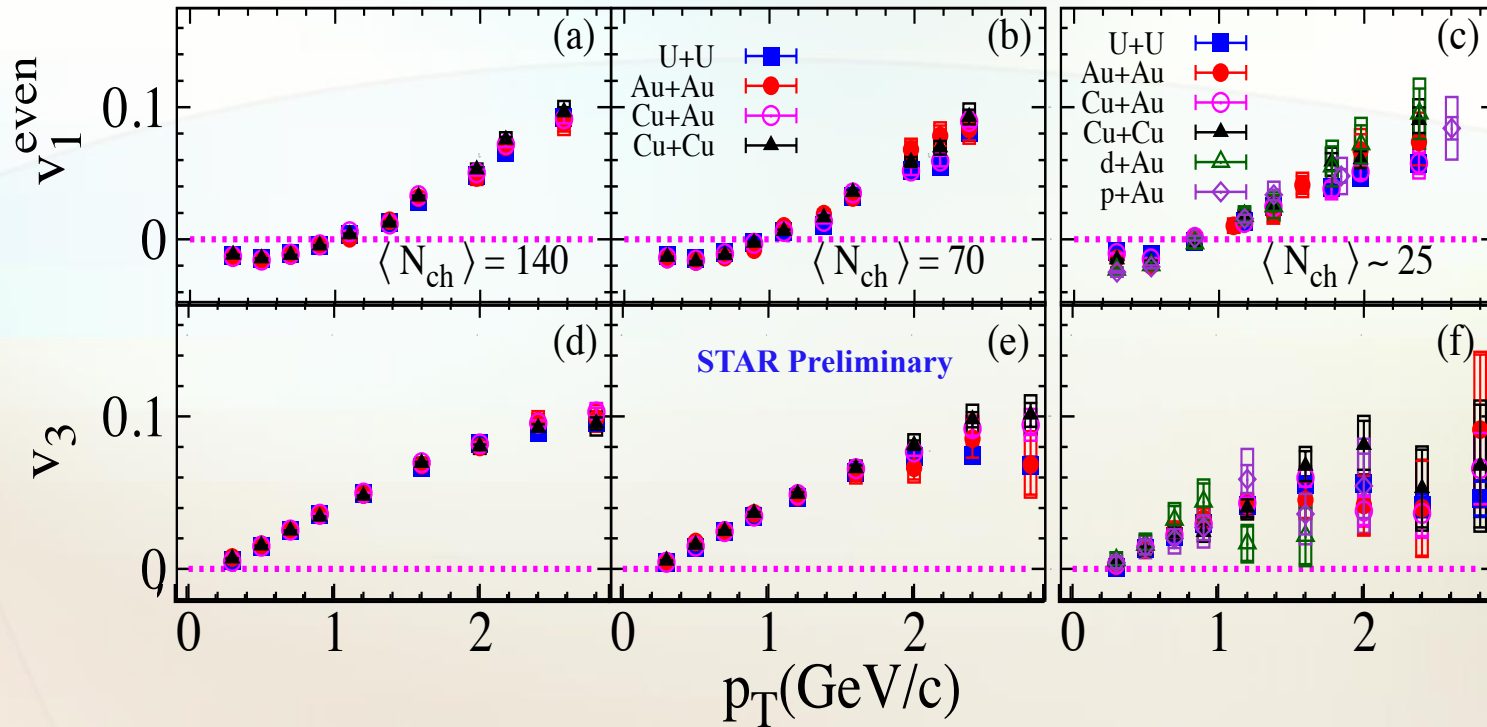
- Odd harmonics are system independent.
- Even harmonics are system dependent, with weak system dependence for the higher harmonics

# $v_n$ for different systems

$|\eta| < 1$  and  $|\Delta\eta| > 0.7$

$$\ln\left(\frac{v_n}{\epsilon_n}\right) \propto -A (\eta/s) \langle N_{Ch} \rangle^{-1/3}$$

$v_1^{even}$  and  $v_3$  vs  $p_T$  at different  $\langle N_{Ch} \rangle$  for all systems



➤  $v_1^{even}$  and  $v_3$  show similar trends and magnitudes for all systems.

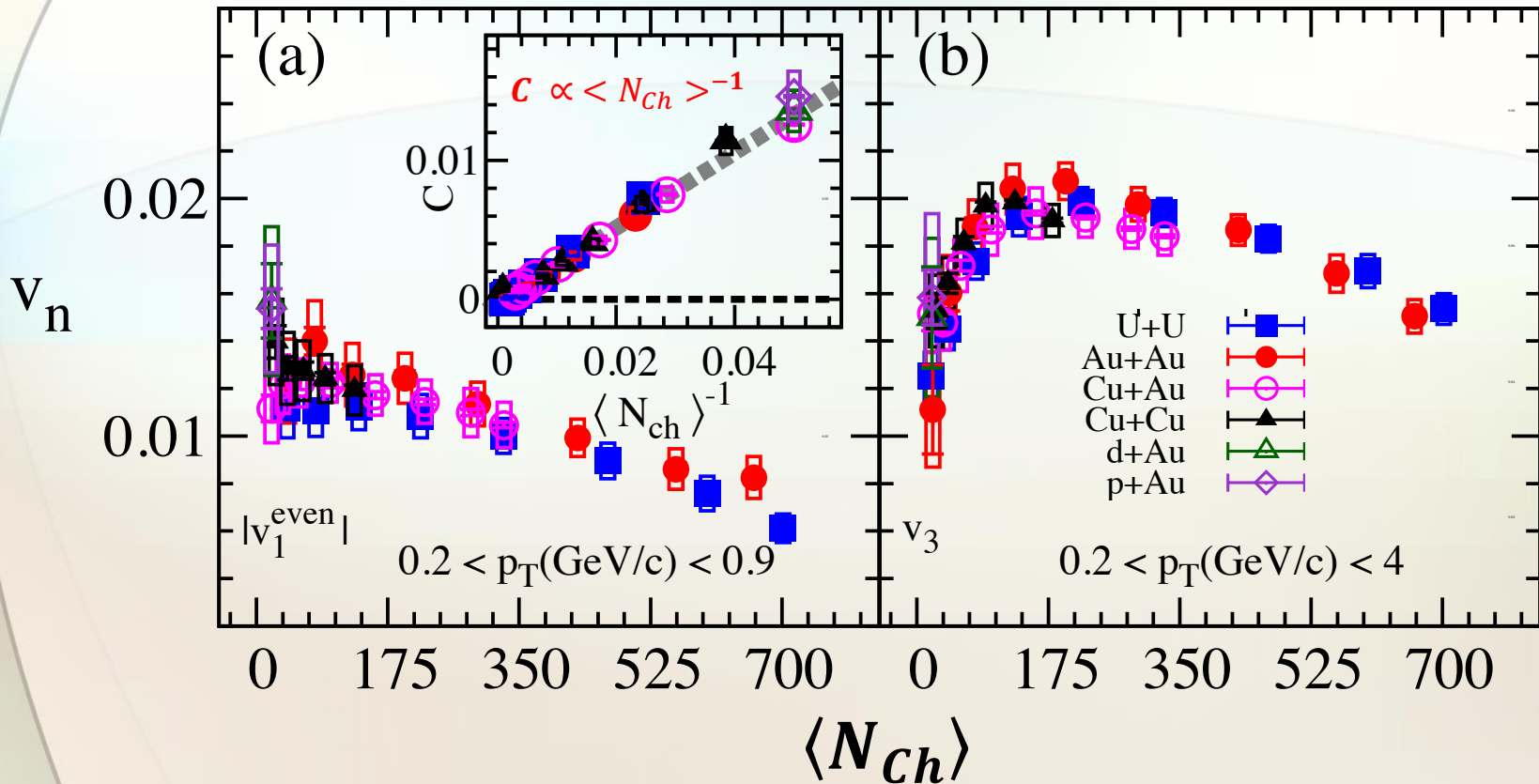
➤  $v_1^{even}$  and  $v_3$  are system independent.

# $v_n$ for different systems

$|\eta| < 1$  and  $|\Delta\eta| > 0.7$

$$\ln\left(\frac{v_n}{\epsilon_n}\right) \propto -A (\eta/s) \langle N_{Ch} \rangle^{-1/3}$$

$v_1^{even}$  and  $v_3$  vs  $\langle N_{Ch} \rangle$  for all systems



➤  $v_1^{even}$  and  $v_3$  show similar trends and magnitudes for all systems.

➤  $v_1^{even}$  and  $v_3$  are system independent.

# $v_n$ for different systems

$|\eta| < 1$  and  $|\Delta\eta| > 0.7$

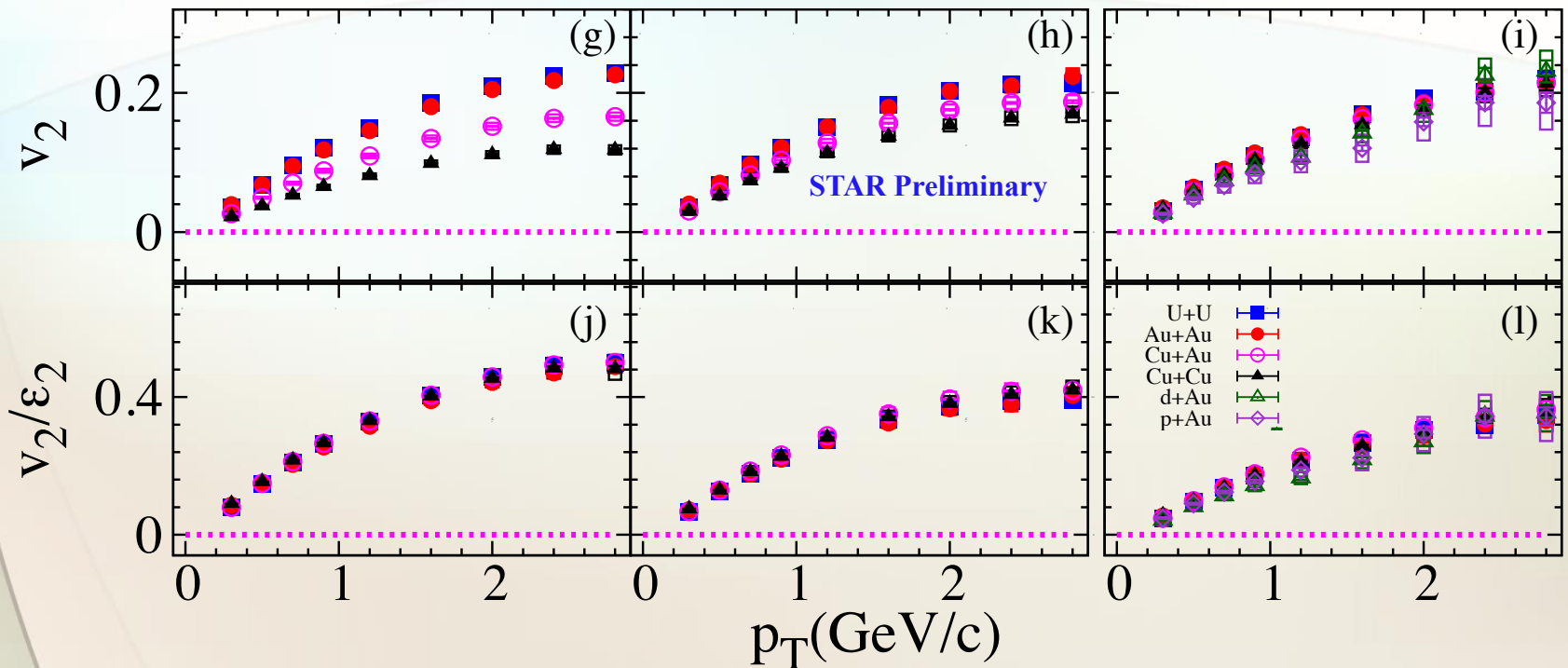
$$\ln\left(\frac{v_n}{\epsilon_n}\right) \propto -A (\eta/s) \langle N_{Ch} \rangle^{-1/3}$$

$v_2$  vs  $p_T$  at different  $\langle N_{Ch} \rangle$  for all systems

$\langle N_{Ch} \rangle = 140$

$\langle N_{Ch} \rangle = 70$

$\langle N_{Ch} \rangle \sim 25$



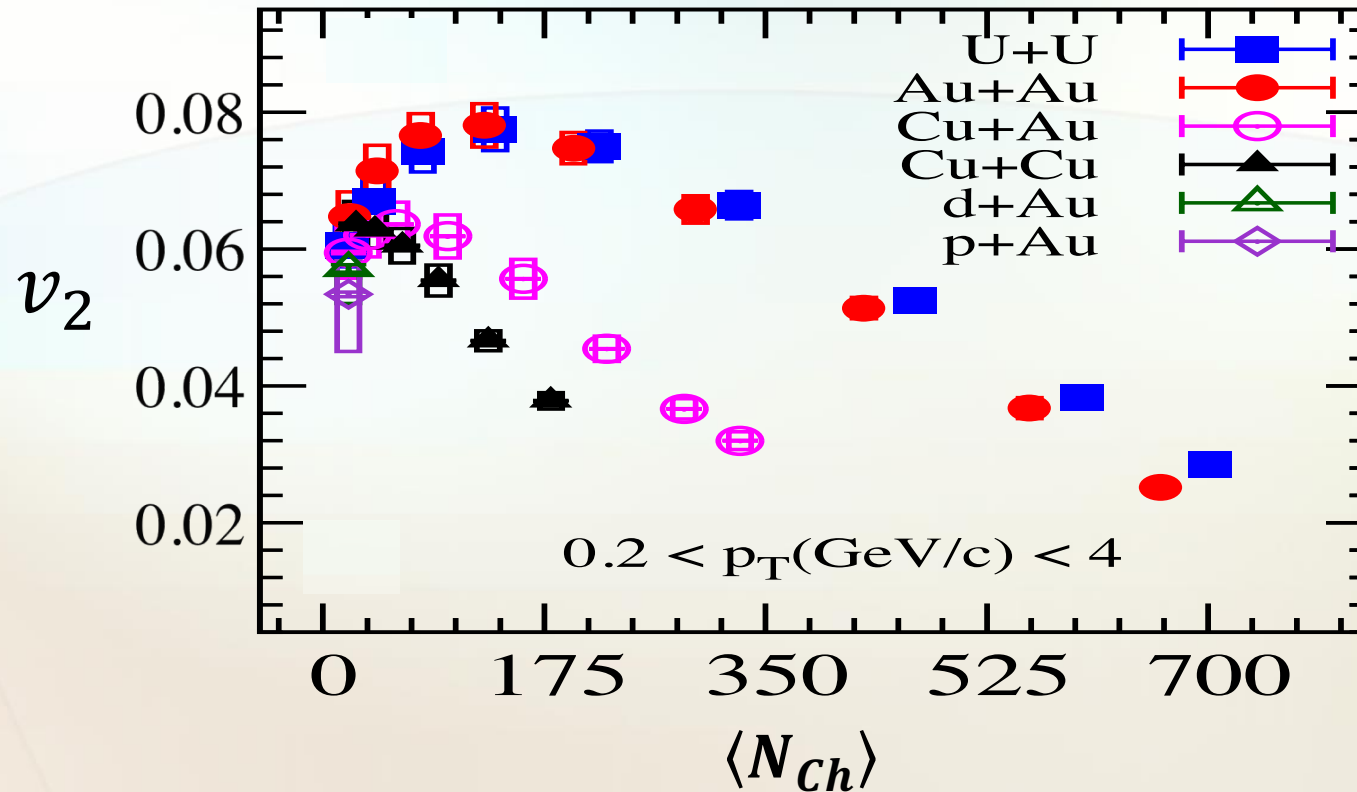
➤  $v_2$  is system dependent (shape).

➤  $\frac{v_2}{\epsilon_2}(p_T)$  for all systems scales to a single curve.

# $v_n$ for different systems

$$\ln\left(\frac{v_n}{\varepsilon_n}\right) \propto -A (\eta/s) \langle N_{Ch} \rangle^{-1/3}$$

$v_2$  vs  $\langle N_{Ch} \rangle$  for all systems



➤  $v_2$  show similar trends but different magnitudes for different systems.

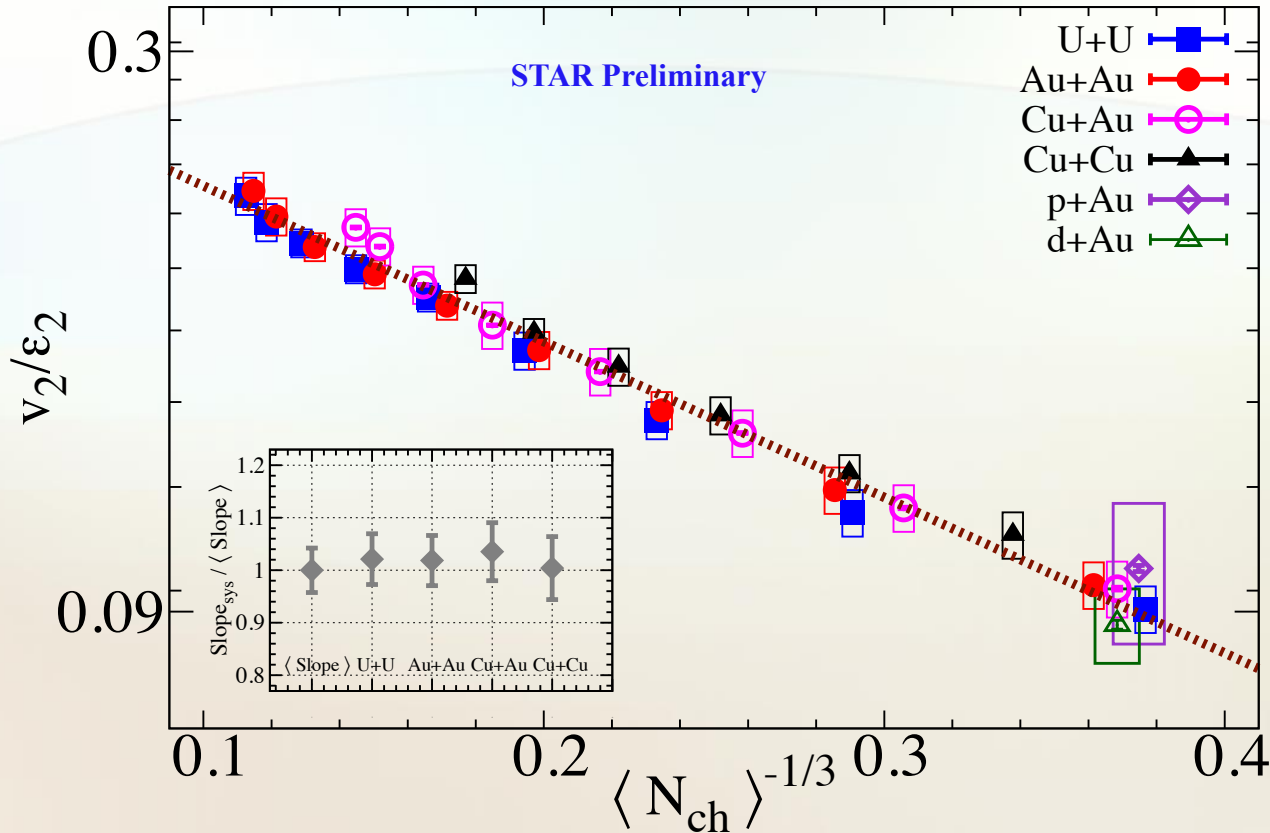
➤  $v_2$  is system dependent (shape).

# $v_n$ for different systems

$|\eta| < 1$  and  $|\Delta\eta| > 0.7$

$$\ln\left(\frac{v_n}{\epsilon_n}\right) \propto -A (\eta/s) \langle N_{ch} \rangle^{-1/3}$$

$\ln\left(\frac{v_2}{\epsilon_2}\right)$  vs  $\langle N_{ch} \rangle^{-1/3}$  for all systems



➤  $\frac{v_2}{\epsilon_2}$  for all systems scales to a single curve.

➤ Similar slopes implies similar viscous coefficient ( $A \eta/s$ ) for all systems.

# Conclusion

Comprehensive set of STAR measurements presented for  $v_n(\eta, p_T, \langle N_{Ch} \rangle)$  for several collision systems/energies.

## ➤ Non-flow suppression

- ✓  $\Delta\eta$  cut used to suppress the short range non-flow
- ✓  $c_2\{4\}$  shows negative value for all presented systems.



# Conclusion

Comprehensive set of STAR measurements presented for  $v_n(\eta, p_T, \langle N_{Ch} \rangle)$  for several collision systems/energies.

- Non-flow suppression
  - ✓  $\Delta\eta$  cut used to suppress the short range non-flow
  - ✓  $c_2\{4\}$  shows negative value for all presented systems.
- Scaling the system size;
  - ✓ The odd harmonics are system independent

# Conclusion

Comprehensive set of STAR measurements presented for  $v_n(\eta, p_T, \langle N_{Ch} \rangle)$  for several collision systems/energies.

- Non-flow suppression
  - ✓  $\Delta\eta$  cut used to suppress the short range non-flow
  - ✓  $c_2\{4\}$  shows negative value for all presented systems.
  
- Scaling the system size;
  - ✓ The odd harmonics are system independent
  - ✓  $v_2$  is system dependent

# Conclusion

Comprehensive set of STAR measurements presented for  $v_n(\eta, p_T, \langle N_{Ch} \rangle)$  for several collision systems/energies.

- Non-flow suppression
  - ✓  $\Delta\eta$  cut used to suppress the short range non-flow
  - ✓  $c_2\{4\}$  shows negative value for all presented systems.
  
- Scaling the system size;
  - ✓ The odd harmonics are system independent
  - ✓  $v_2$  is system dependent
  - ✓  $\frac{v_2}{\epsilon_2}$  for all systems scaled onto one curve to  $\sim 10\%$  in slope

# Conclusion

Comprehensive set of STAR measurements presented for  $v_n(\eta, p_T, \langle N_{Ch} \rangle)$  for several collision systems/energies.

➤ Non-flow suppression

- ✓  $\Delta\eta$  cut used to suppress the short range non-flow
- ✓  $c_2\{4\}$  shows negative value for all presented systems.

➤ Scaling the system size;

- ✓ The odd harmonics are system independent
- ✓  $v_2$  is system dependent
- ✓  $\frac{v_2}{\epsilon_2}$  for all systems scaled onto one curve to  $\sim 10\%$  in slope

At the same energy, the scaling features suggest similar viscous coefficient ( $A \frac{\eta}{s}$ ) for different systems.

**THANK YOU**