

Do we understand the Initial Stages?

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BNL



IS 2017 Krakow, September 18-22, 2017

No...

...because the physics of initial stages in hadron-hadron collisions is nonperturbative for most quantities, and nonperturbative QCD is complicated...



PHYSICAL REVIEW

VOLUME 96, NUMBER 1

OCTOBER 1, 1954

Conservation of Isotopic Spin and Isotopic Gauge Invariance*

C. N. YANG † AND R. L. MILLS

Brookhaven National Laboratory, Upton, New York

(Received June 28, 1954)

... but very interesting, because it teaches us about *strongly correlated real-time dynamics of a non-Abelian gauge theory*

Caveat: This talk is NOT a summary but my limited personal perspective on some recent developments



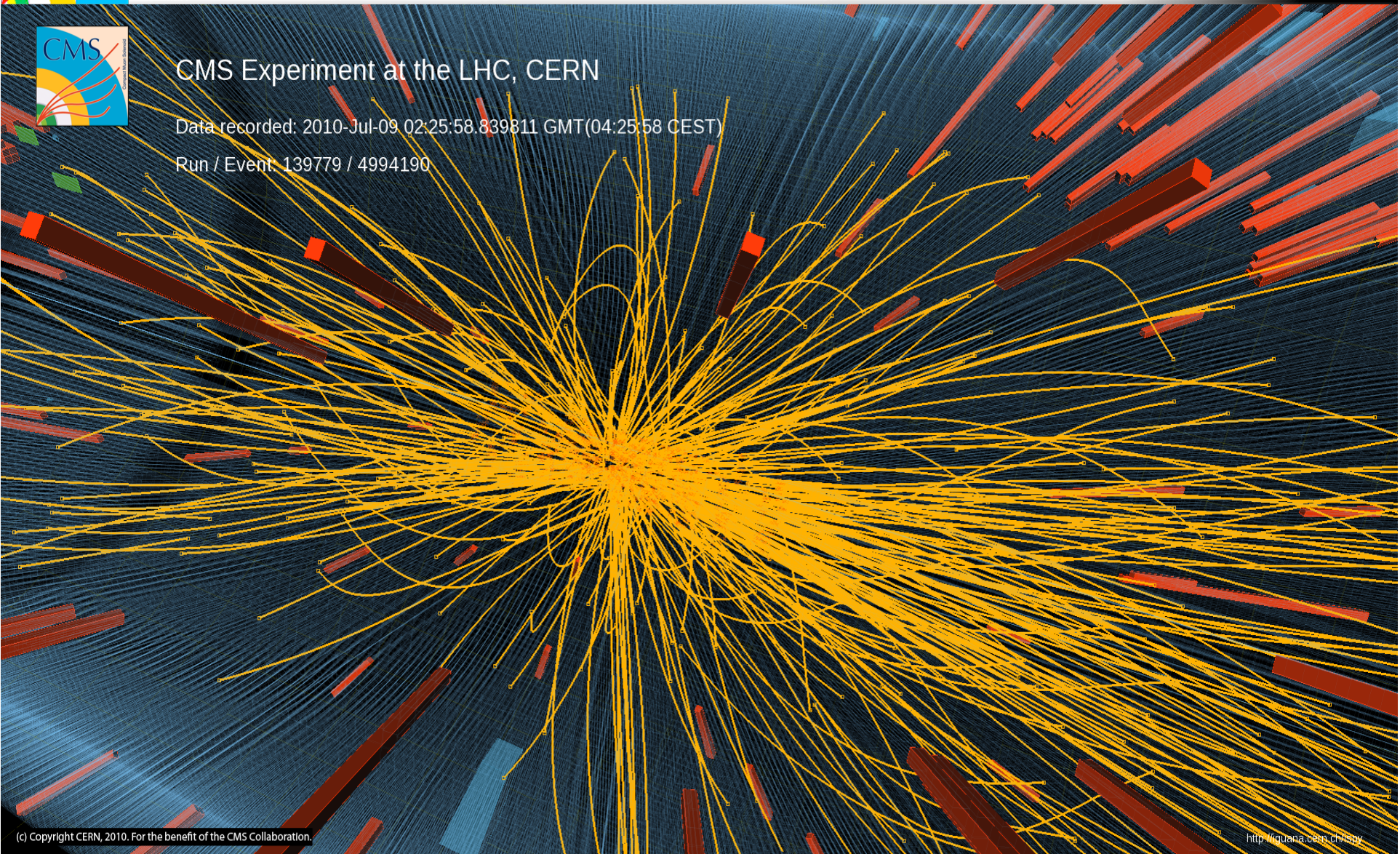
High Multiplicity pp collisions



CMS Experiment at the LHC, CERN

Data recorded: 2010-Jul-09 02:25:58.839811 GMT(04:25:58 CEST)

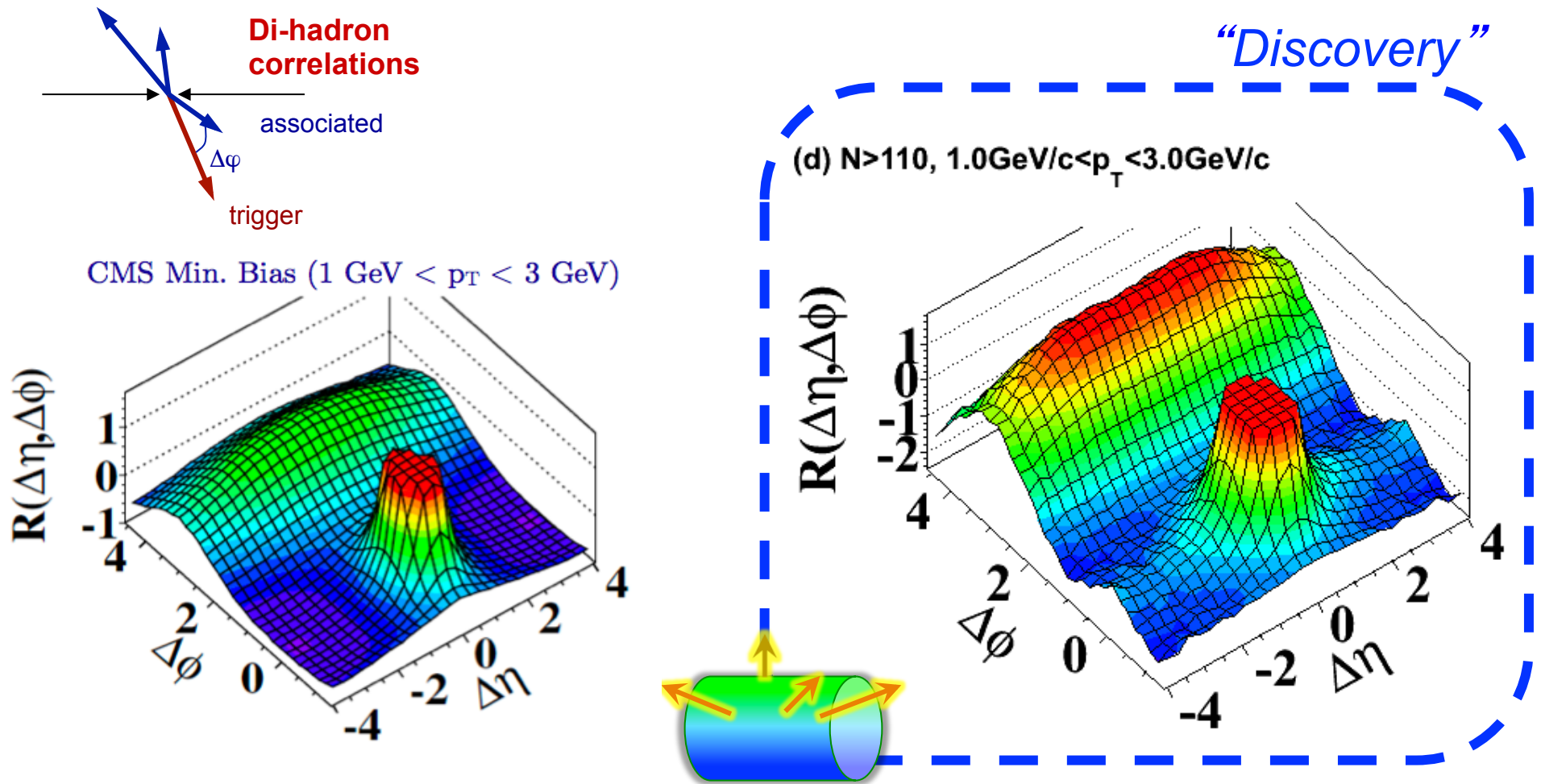
Run / Event: 139779 / 4994190



(c) Copyright CERN, 2010. For the benefit of the CMS Collaboration.

<http://gqana.cern.ch/Spy>

Two particle correlations: CMS results

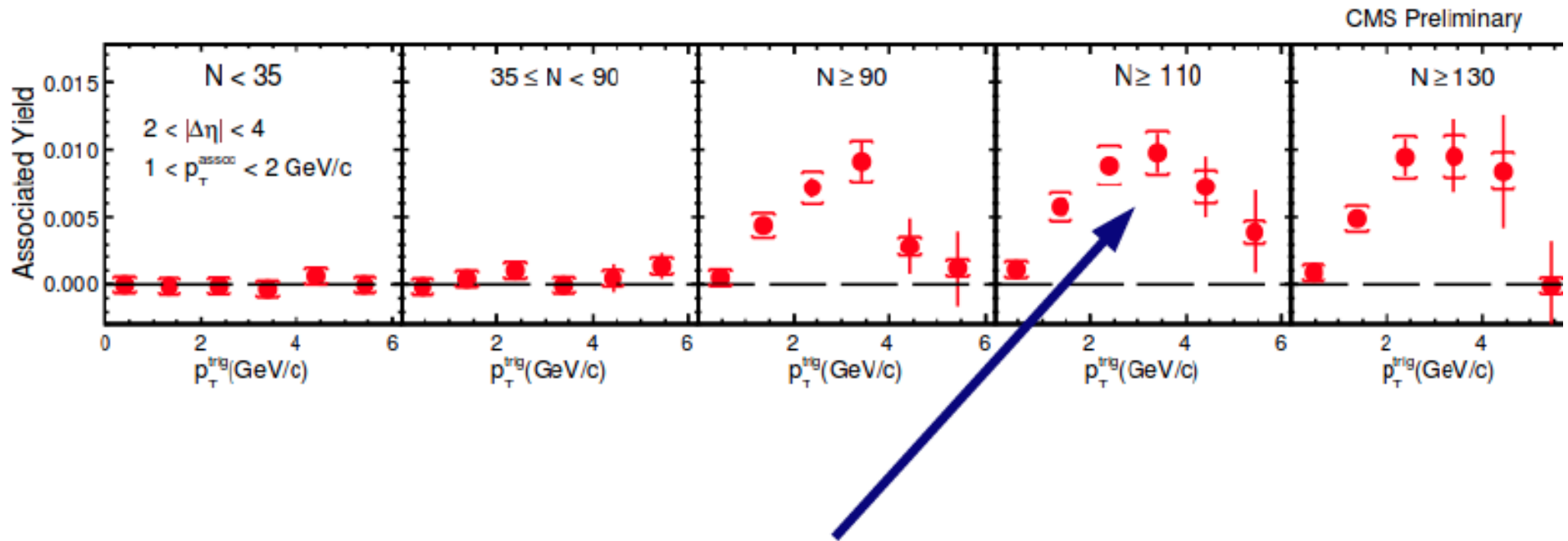


Observation of Long-Range Near-Side Angular Correlations in Proton-Proton Collisions at the LHC [CMS Collaboration \(Vardan Khachatryan \(Yerevan Phys. Inst.\) et al.\)](#). JHEP 1009 (2010) 091

[Cited by 597 records](#)

5th most cited CMS physics paper to date!

My initial interest (in 2010) was piqued by the sub-structure of the ridge shown:

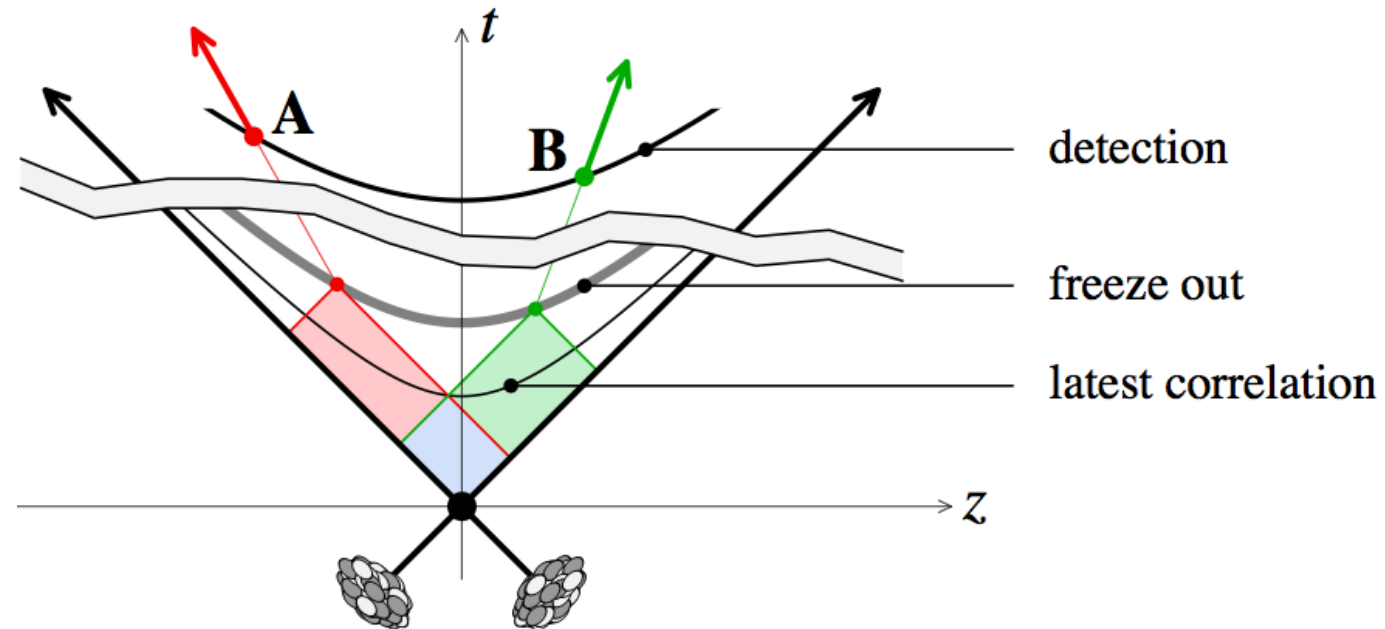


Evidence of a semi-hard scale in the data?

If this scale is $\approx Q_s$, since $\alpha_s(Q_s^2) \ll 1$
could examine this nonperturbative strongly correlated
phenomenon in weak coupling

will return to this point later...

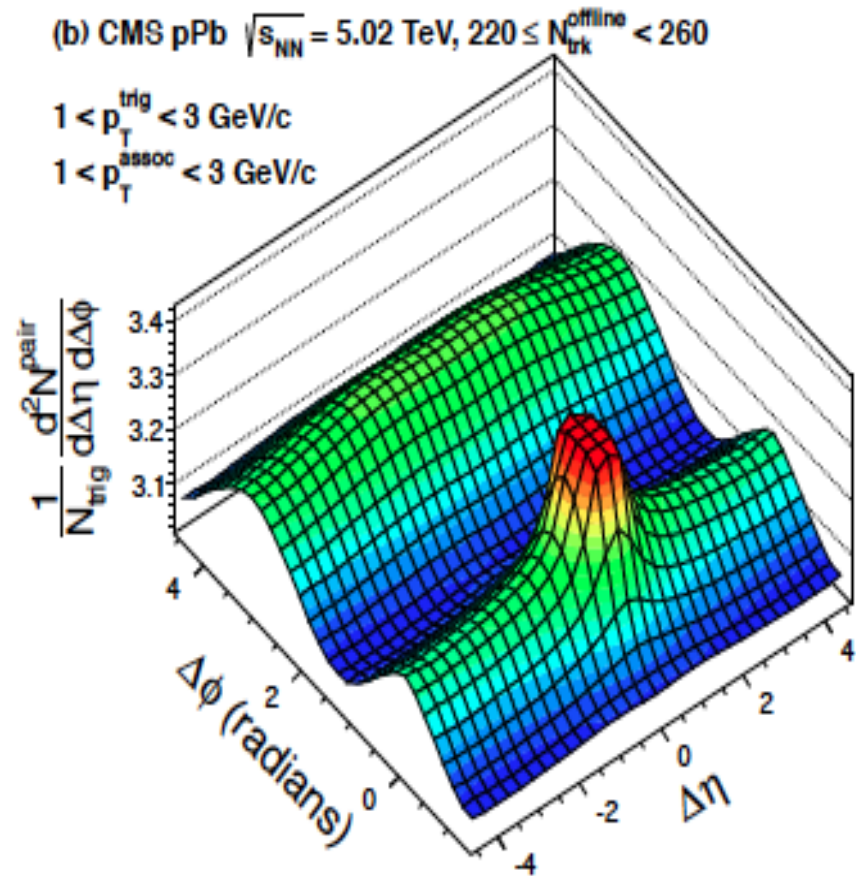
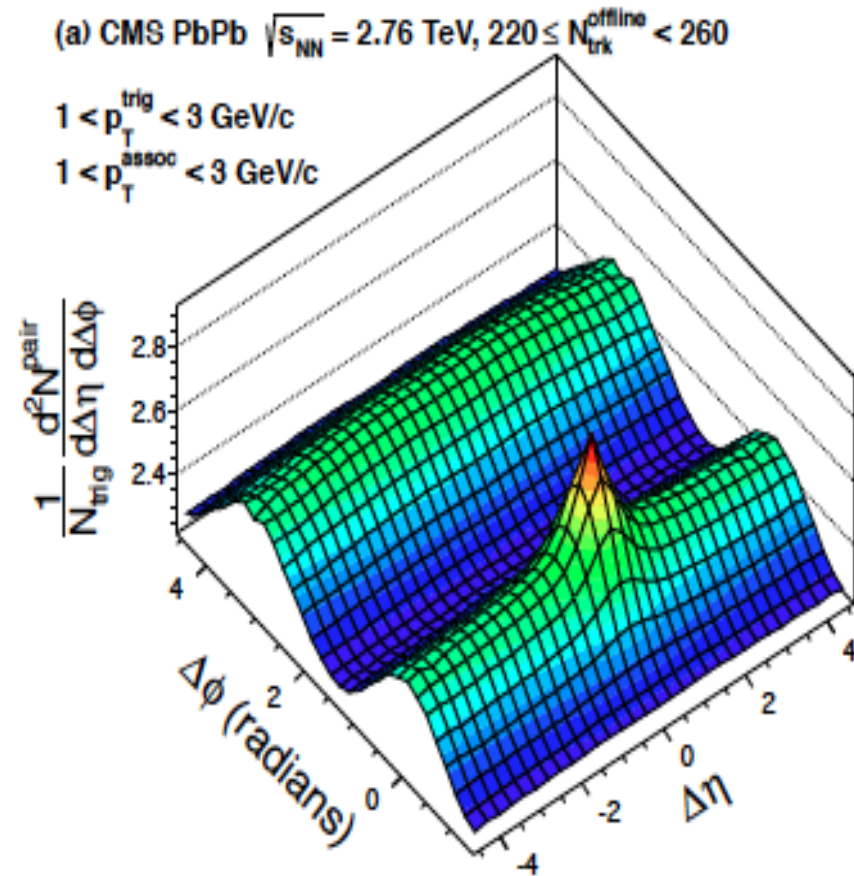
Long range rapidity correlations are a chronometer



$$\tau \leq \tau_{\text{frz-out}} \exp \left(-\frac{1}{2} \underbrace{|y_A - y_B|} \right)$$

Long range correlations sensitive to very early time
(fractions of a femtometer $\sim 10^{-24}$ seconds) dynamics in collisions

Another surprise: central p+A looks like peripheral A+A



p+A ridge seen --much larger than p+p at same multiplicity
and nearly as large as that in peripheral Pb+Pb collisions

First indication that hydrodynamics may play a role in such events

Several early papers by Bozek and Broniowski



See Inside

Particles That Flock: Strange Synchronization Behavior at the Large Hadron Collider

Scientific American, February (2011)

Scientists at the Large Hadron Collider are trying to solve a puzzle of their own making: why particles sometimes fly in sync

The high-energy collisions of protons in the LHC may be uncovering “a new deep internal structure of the initial protons,” says Frank Wilczek of the Massachusetts Institute of Technology, winner of a Nobel Prize

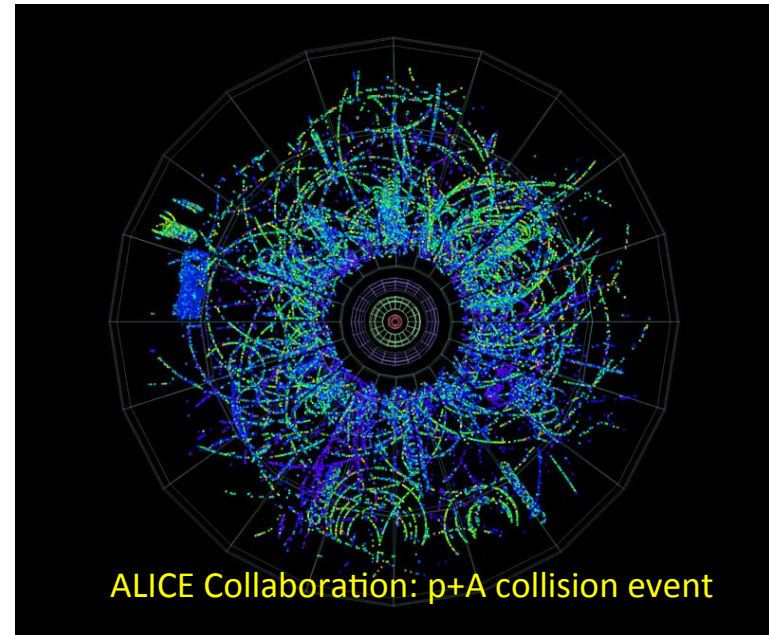
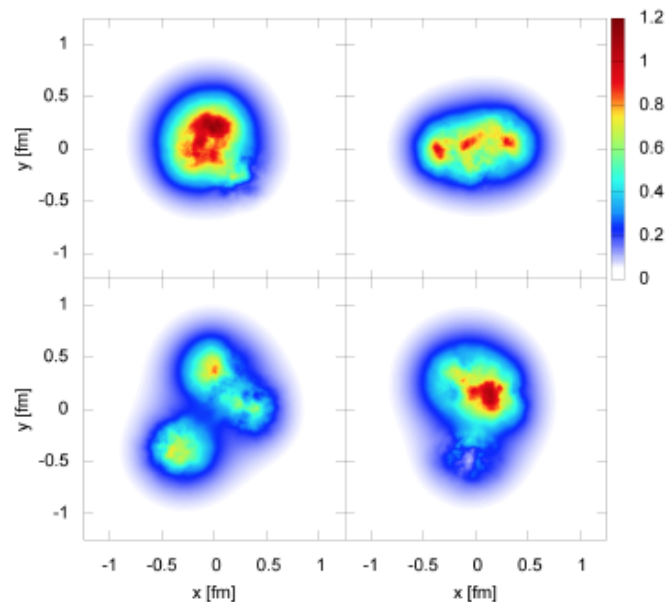
“At these higher energies [of the LHC], one is taking a snapshot of the proton with higher spatial and time resolution than ever before”



Deep internal structure of the proton...

A piece of initial state physics that may matter for the final state

Eccentric protons



Bjorken, Brodsky, Goldhaber, PLB726 (2013) 344

HERA data on incoherent diffractive vector meson production favor this

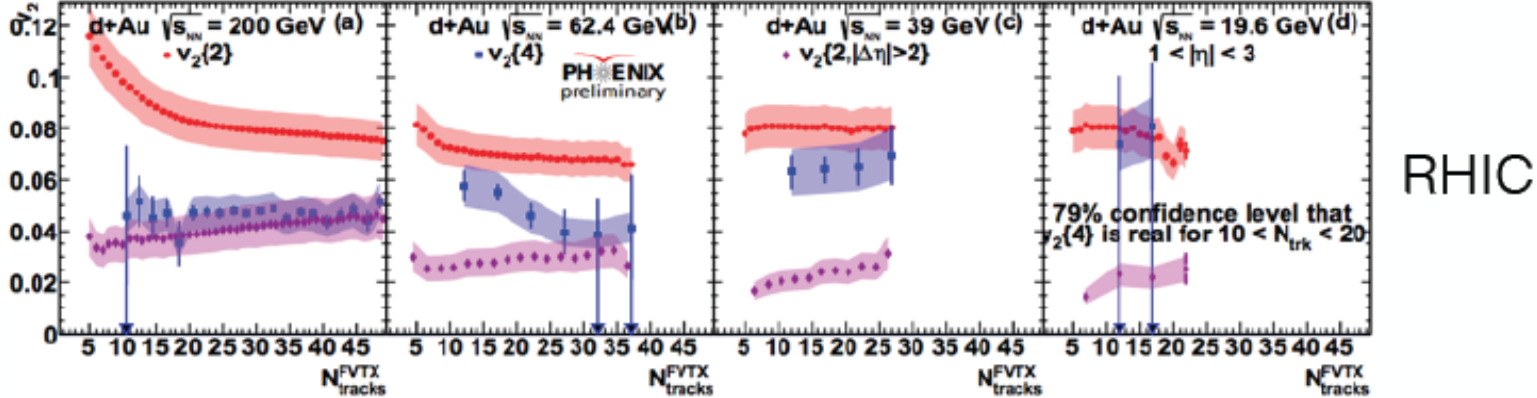
Talks by Schenke and Mantysaari

Timeline of our (mis?) understanding...

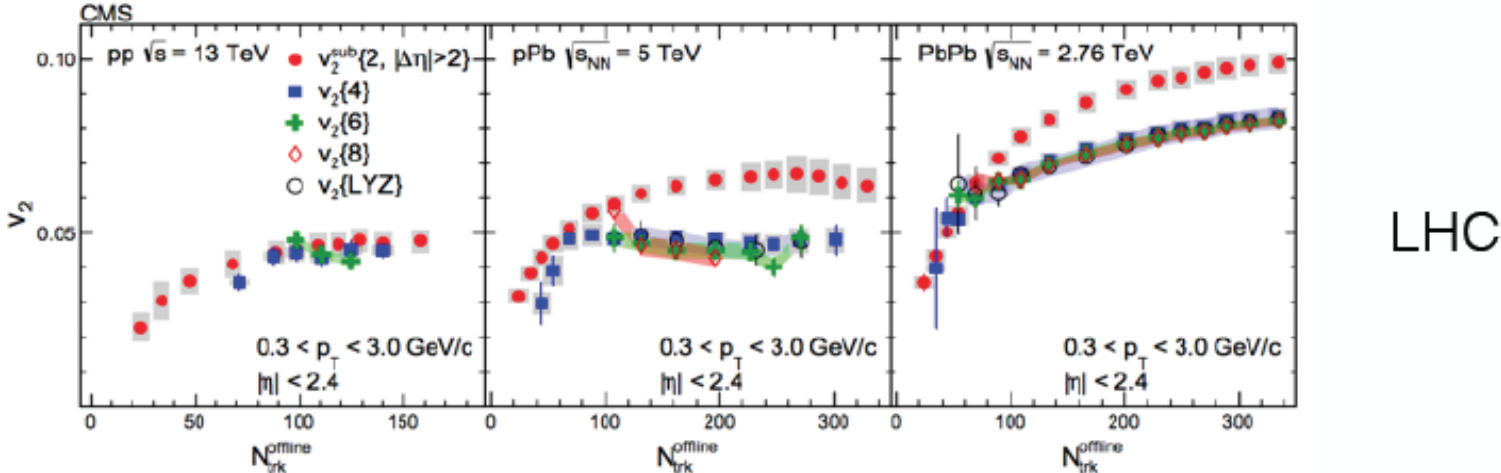
- 2010—Ridge in high multiplicity p+p (LHC)! Probably CGC!
- Early 2010s—QGP in p+Pb!
- Early 2010s—QGP in d+Au!
- Mid 2010s and now-ish—QGP in high multiplicity p+p? QGP in mid-multiplicity p+p?? QGP in d+Au even at low energies???

From Ron Belmont's talk

Collectivity in the sense of $v_2\{2\} \geq v_2\{4\} \approx v_2\{6\} \approx v_2\{8\}$ is now ubiquitous
 ...widely believed when first seen to be “proof” of final state response to
 initial state geometry



R. Belmont PHENIX arXiv:1704.04570

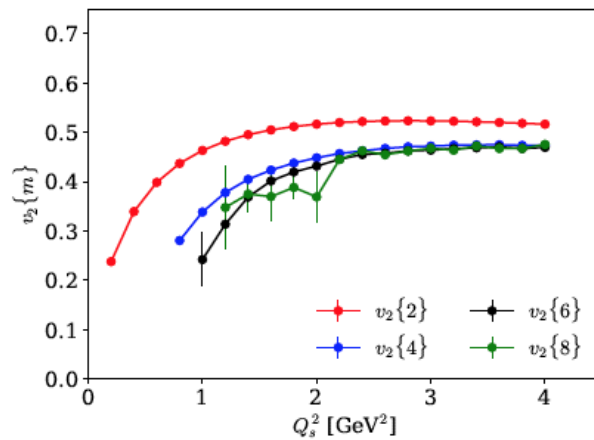


CMS PRL 115 (2015) 012301

Our understanding of cumulant measurements now more nuanced...

- Multi-particle long-range correlation, i.e. collectivity, could produce $c_2\{4\}$, $c_2\{6\}$.. with any sign: it is the property of the shape of $p(v_2)$
 - Also possible in final state, most models has $p(\epsilon)$ that give $\epsilon_n\{2k\}$ with “correct” sign. But the initial geometry in small systems is not settled
 - sign-change is possible if $c_2\{4\} = \langle v_2^4 \rangle - 2\langle v_2^2 \rangle^2 \propto \langle \epsilon^4 \rangle - 2\langle \epsilon^2 \rangle^2 > 0$ **From talk by Jia**
 - or $p(\epsilon)$ may be engineered such that the signs of $\epsilon_n\{2k\}$ flip

Especially relevant to understanding whether positive $c_2\{4\}$ seen in p+A at RHIC is consistent with flow



Convergence of higher moments also seen in a very simple initial state “color domain” model



**Talk by Mace
See also, talk by Kovner**

**$v_2\{4\}$ also seen from MPI model of initial state
QCD interference contributions**

Talk by Wiedemann

Our understanding of initial/final state correlations is also more nuanced...

COLLECTIVITY: POSSIBLE ORIGIN

Final state

- Driven by initial state **geometric** correlations
- Develops gradually during (hydro) evolution
- Requires large multiplicities to facilitate final state interaction
- Requires non-trivial initial state geometry (proton shape fluctuations)

~~VS~~

and

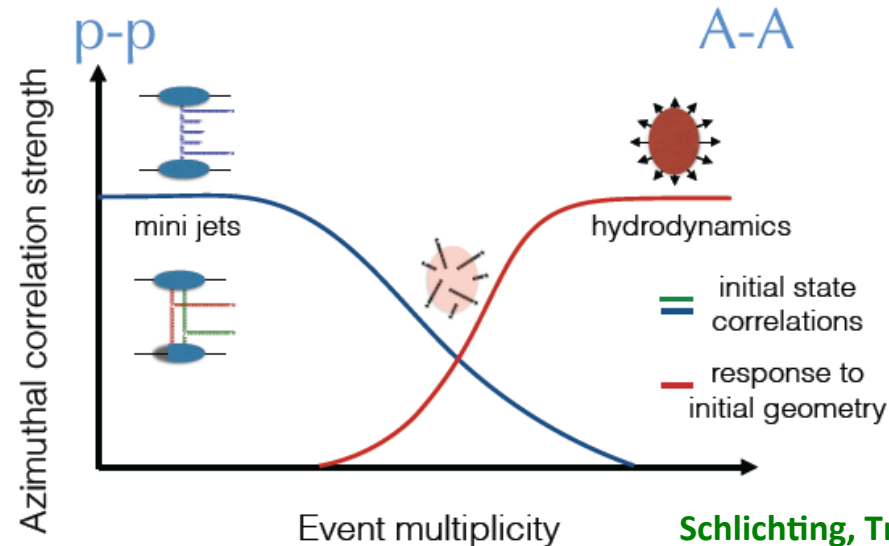
Initial state

- Driven by initial **momentum** correlations
- Pre-exist before collisions or develops very soon after
- High-multiplicities are not required, but allowed
- Momentum correlations are present (suppressed by $1/N_c^2$) for “round” p

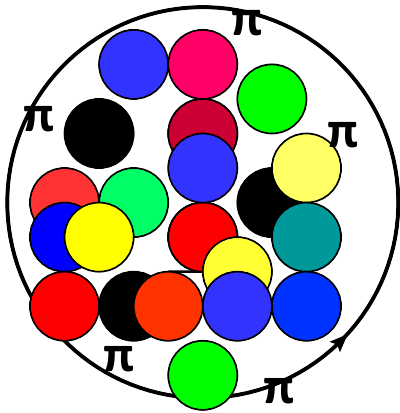
Talk by Skokov

Stated more bluntly...

Saturated wavefunctions in hadron-hadron collisions will generate both non-flow and flow contributions



Their relative contributions depend on Q_s , p_T , multiplicity and system size



Saturated gluons have the maximal occupancy possible in QCD $\sim 1/\alpha_s$

-- no quasi-particle description possible at early times for modes with $p_T < Q_s$

The spacetime evolution of these modes described by solving the equations of classical QCD: the Yang-Mills equations
QCD coupling drops out for these modes

Talks by Boguslavski, Fries, Gelis, Mueller², Schenke

Modes with $p_T \gg Q_s$:

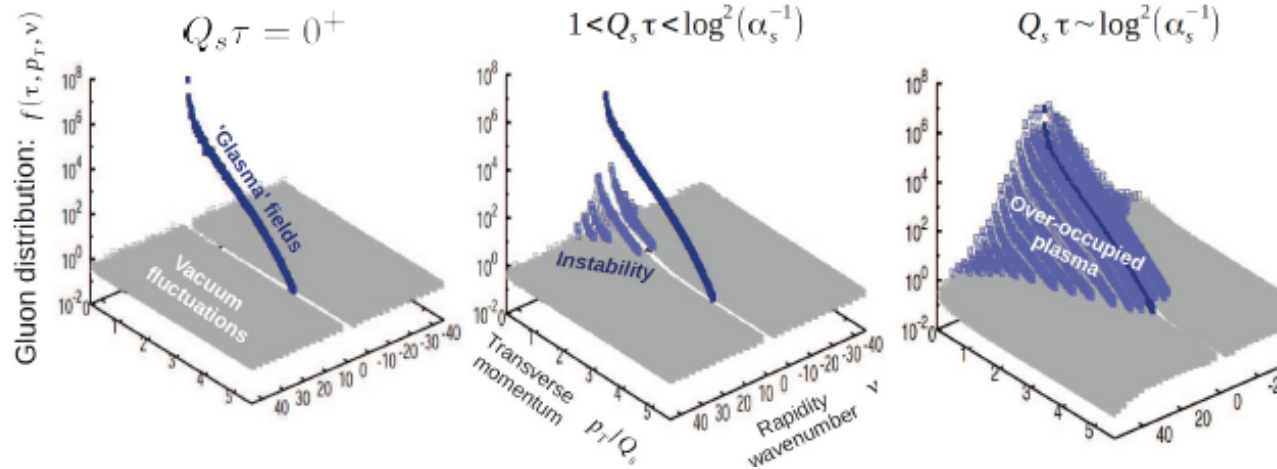
Have quasi-particle description, and match

to pQCD ... matching can be done order by order in pQCD – jet physics

Flow component at $\tau \approx 1/Q_s$ described by $p_T < Q_s$

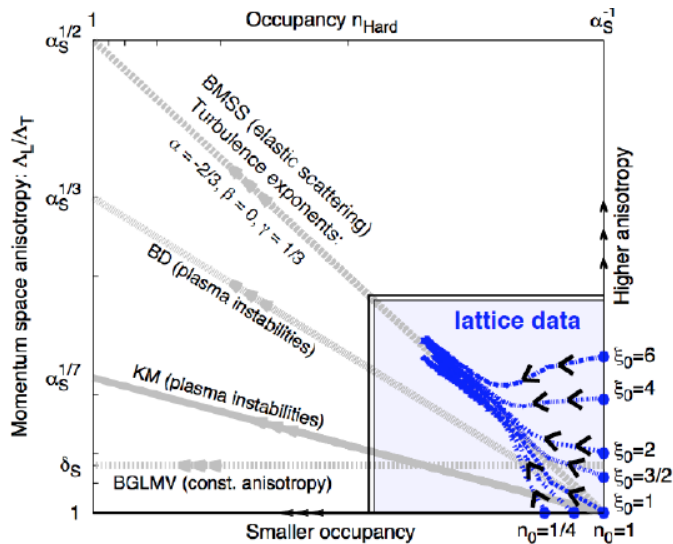
Non flow component by $p_T > Q_s$

Subsequent spacetime evolution of “flow component” for $\tau \gg 1/Q_s$



Plasma instabilities cause the system to become overoccupied in p_T and p_z very quickly... and system flows to an attractor, which is that of the bottom-up kinetic scenario

Talk by Boguslavski

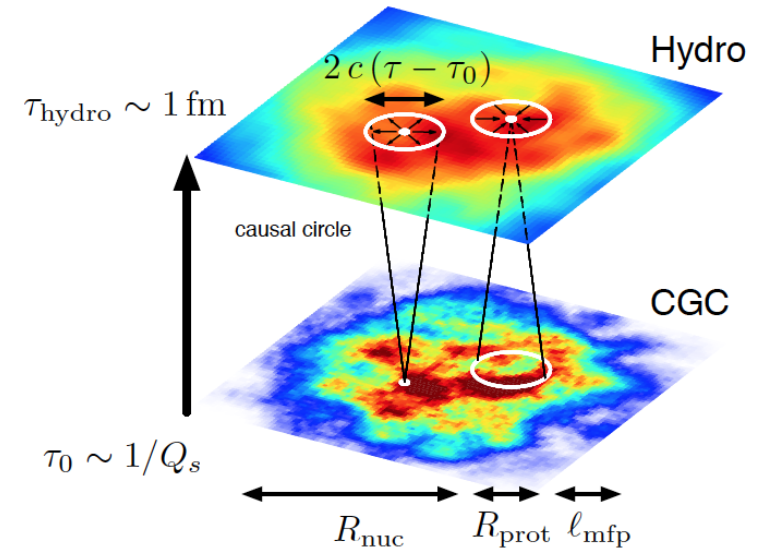


Many open questions and caveats regarding the role of so-called “quantum 1/2” contributions

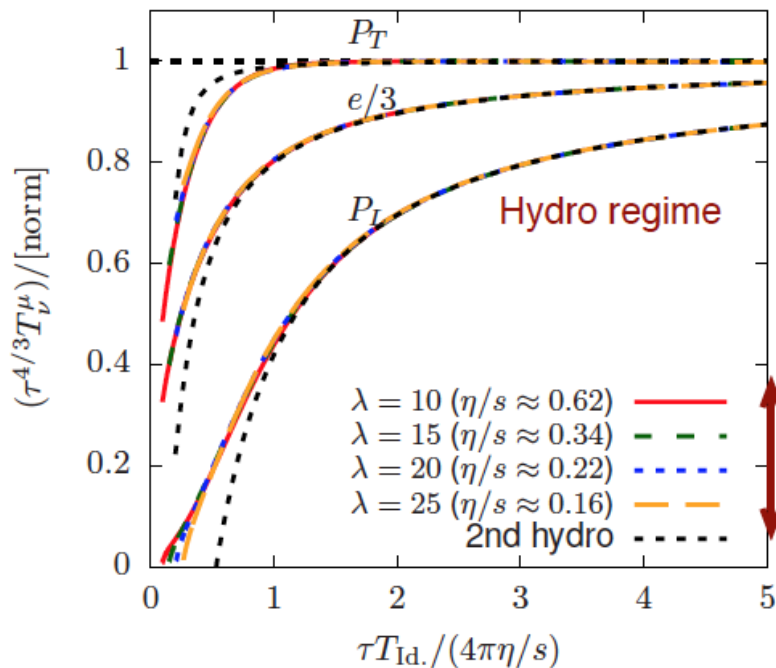
Discussed in talks by Gelis and Wu

Recent developments in bottom-up thermalization...

Lumpy “hot spot” transverse structure in IP-Glasma/EKRT survives until matching with hydro



When does the background stress tensor approach second order hydrodynamics?



Different values of coupling give different η/s

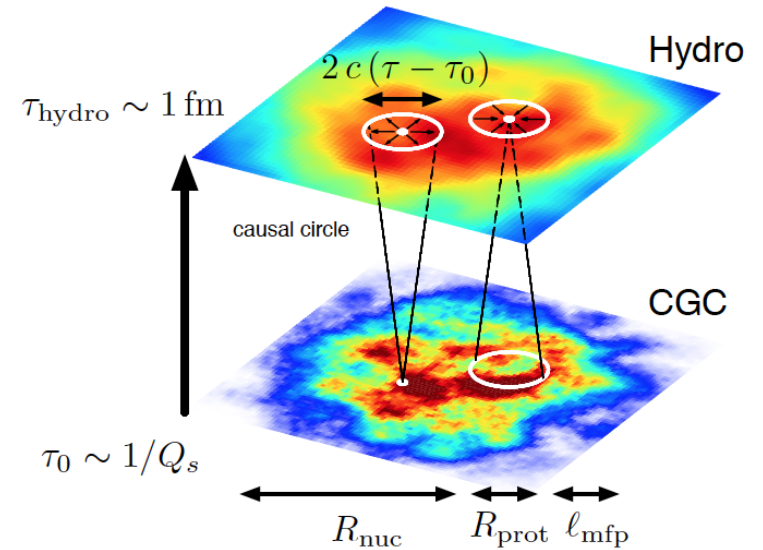
In terms of η/s , all couplings thermalize at same scaled time
Keegan, Kurkela, Romatschke, Schee, Zhu

Gives a basis for interpolating from weak coupling results to stronger coupling

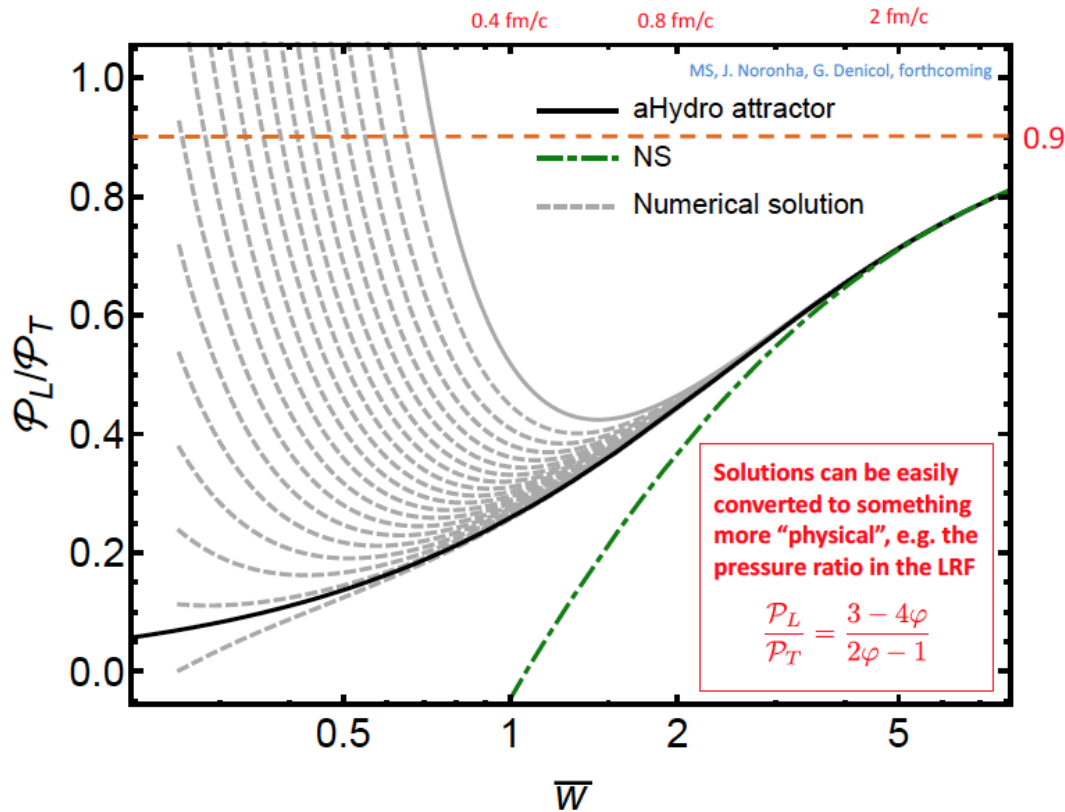
Talks by Teaney and Mazeliauskas

Recent developments in bottom-up thermalization...

Lumpy “hot spot” transverse structure in IP-Glasma/EKRT survives until matching with hydro



The attractor concept – 0+1d



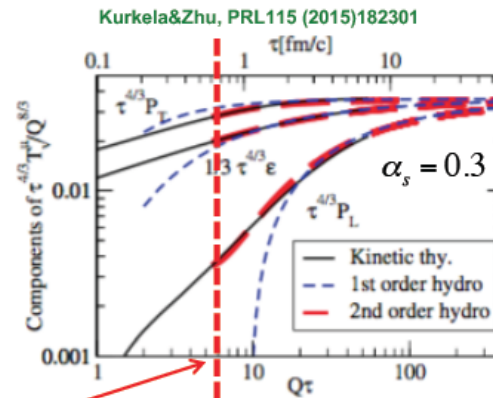
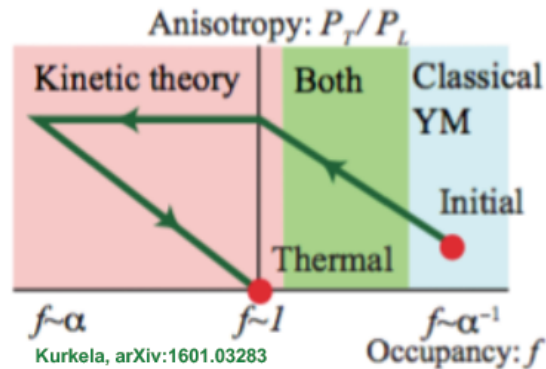
Talk by Strickland

Can one realize the “aHydro Attractor” in the bottom-up framework?

What about “non-flow” mini-jets and jets at the scale $p_T \geq Q_s$?

Fluid dynamics \Leftrightarrow final state interactions \Leftrightarrow Jet quenching

- Bottom-up thermalization formalizes relation between fluid dynamics and jet quenching R.Baier, A.H. Mueller, D. Schiff, D.T. Son, 2001



Slide from Wiedemann's talk

Partonic distributions $f(p)$ governed by Boltzmann equation.

We know:

- $f(p)$ hydrodynamizes on sub-fermi time scale
- “Hydro” & “jet quenching” arise from the same collision kernels

Berges, Eppelbaum, Kurkela, Moore, Schlichting, Venugopalan, ...

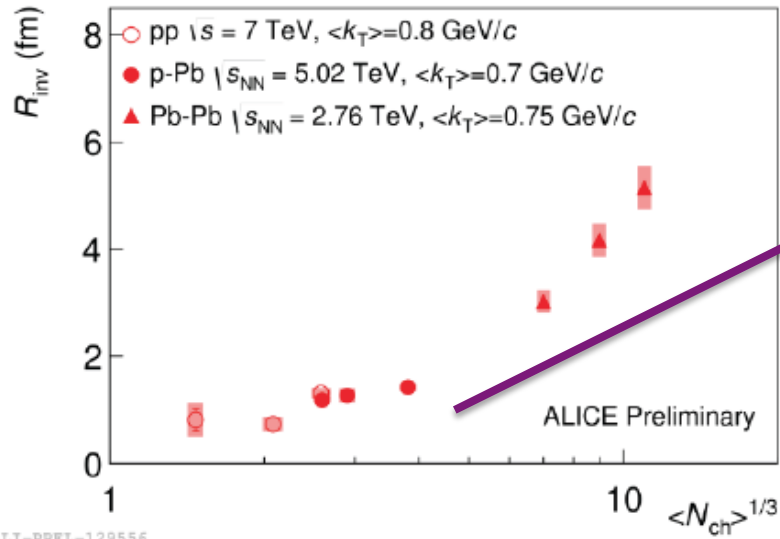
$$\partial_t f(p,t) = -C_{2 \leftrightarrow 2}[f] - C_{1 \leftrightarrow 2}[f]$$

2->2 collision kernel

LPM splitting kernel

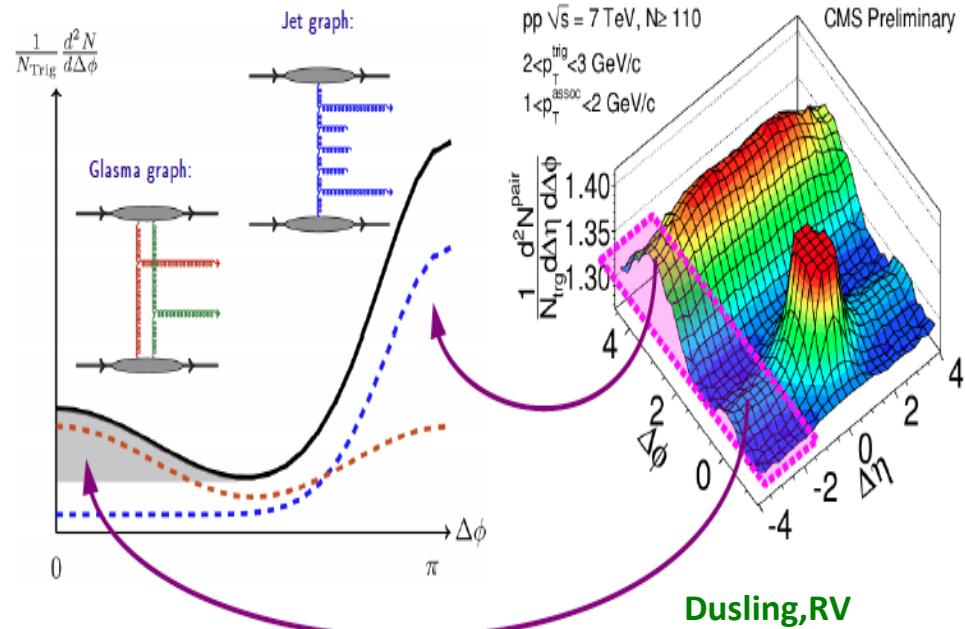
Jet quenching and fluid dynamics = two manifestations of the same physics

What about this flow+non-flow framework for the smaller systems?



Talk by Zhou

HBT radii indicate that lifetimes of the smaller systems is quite short ...little time for rescattering

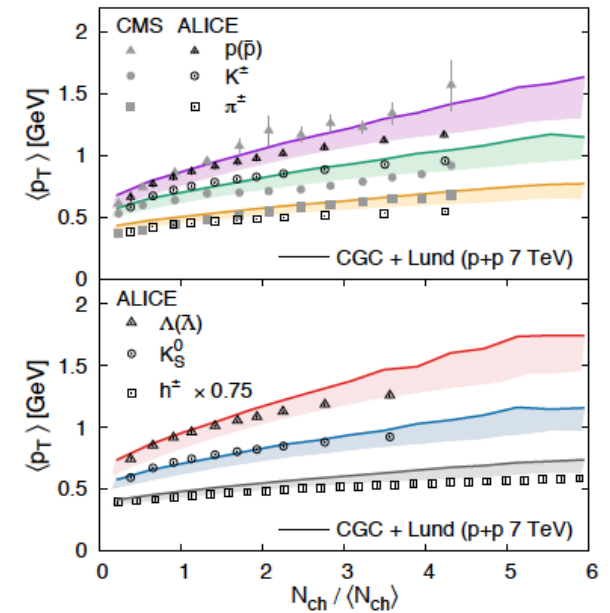
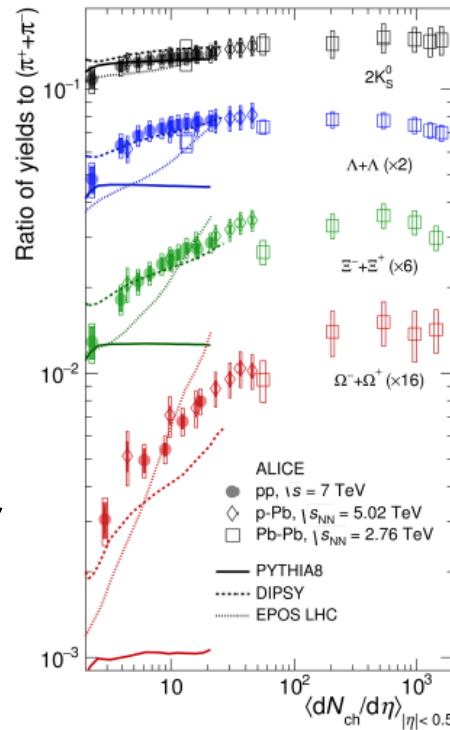
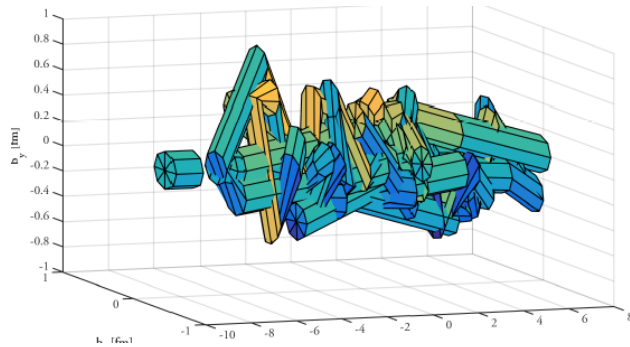
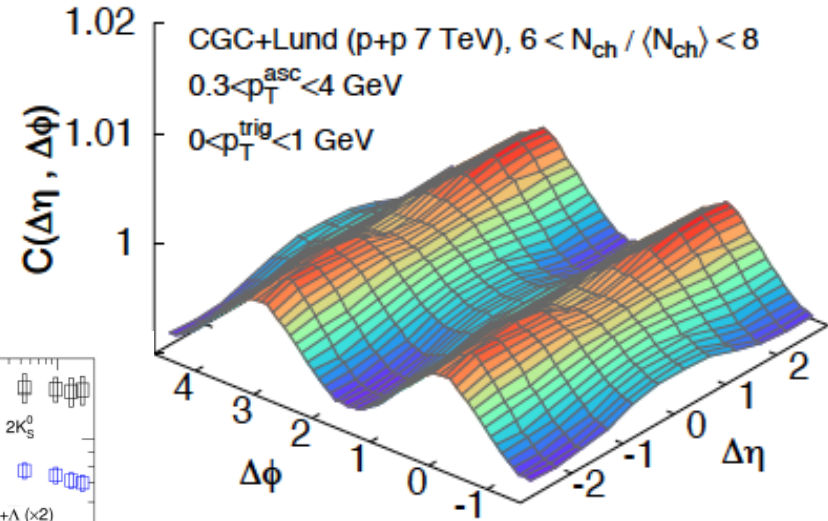


Non-flow $p_T \geq Q_s$ piece of the initial state generate the Glasma graphs – QCD interference “Bose Enhancement” contributions...little jet quenching

Talks by Altinouk and Kovner

What about this flow+non-flow framework for the smaller systems?

The $p_T < Q_s$ modes flow – described by the Yang-Mills equations – freeze-out before “quasi-particle” rescattering can occur... generate v_2, v_3 and mass ordering...

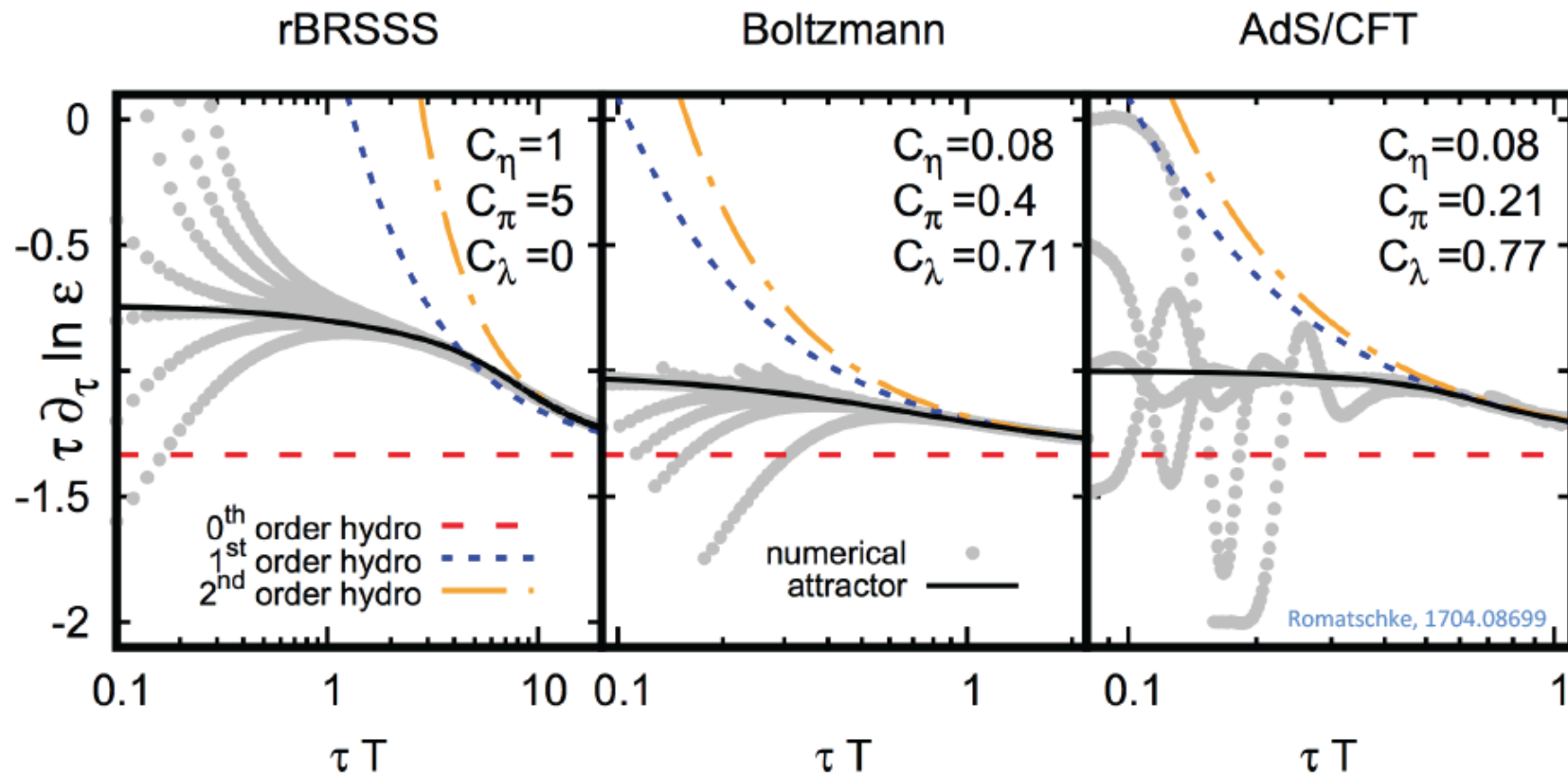


“String shoving” and rope hadronization in DIPSY

Talk by Bierlich

Can one understand this conformal strong field Yang-Mills flow in small systems within the “hydrodynamization” paradigm?

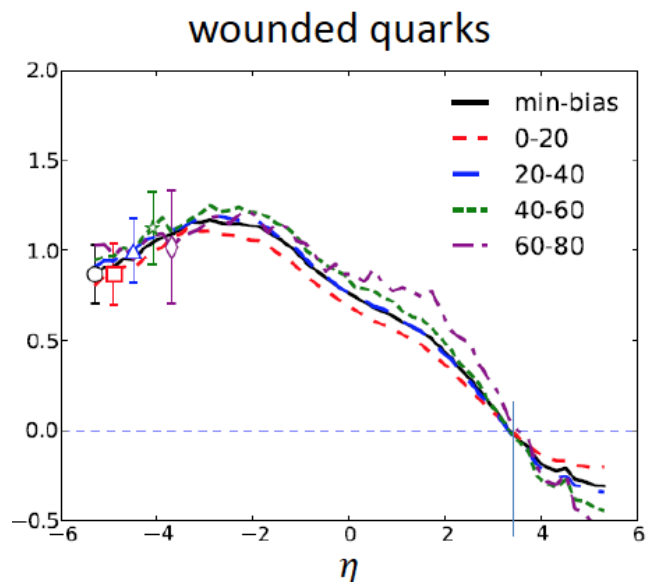
Talks by Romatschke and Spalinski



Simulations of early time dynamics employing holographic ideas

Talks by Attems, Lublinsky, van der Schee

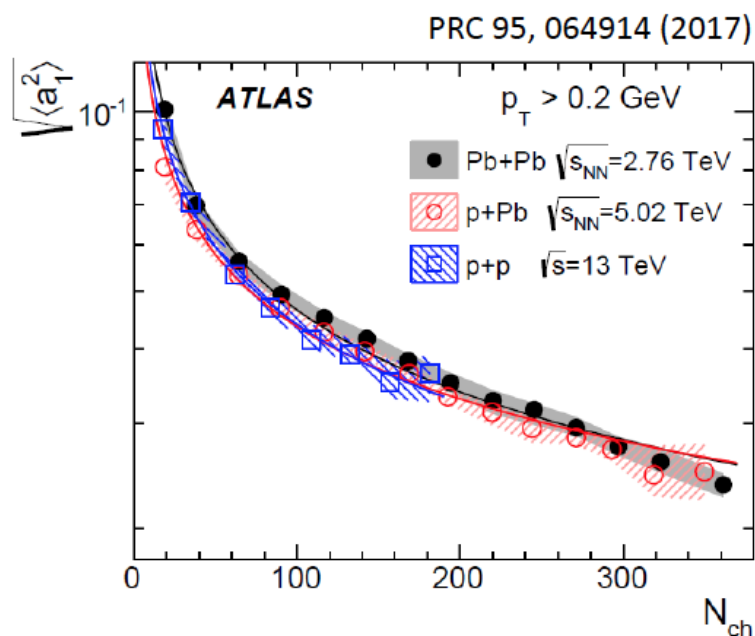
Rapidity correlations and puzzling data on the underlying stringy structure...



Wounded quark model – Can multiparticle production in this framework –developed by the Krakow school -- be related to the CGC?

$$\frac{C_2(y_1, y_2)}{\langle \rho(y_1) \rangle \langle \rho(y_2) \rangle} \sim \langle a_0^2 \rangle + \langle a_1^2 \rangle \frac{y_1 y_2}{Y^2} + \dots$$

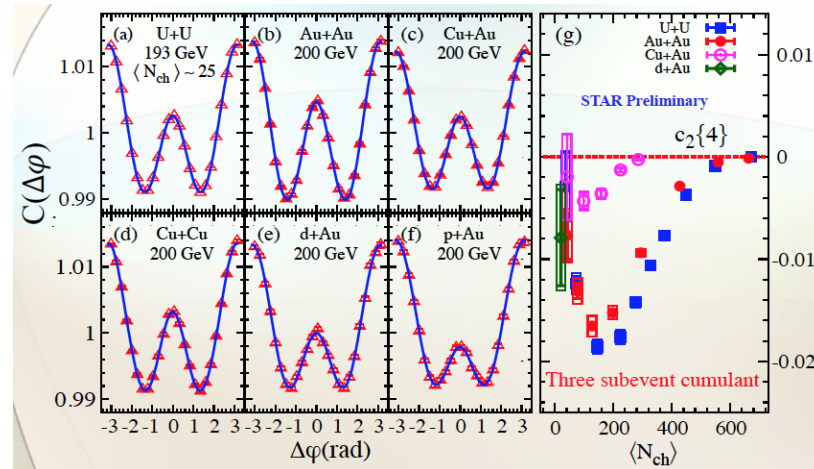
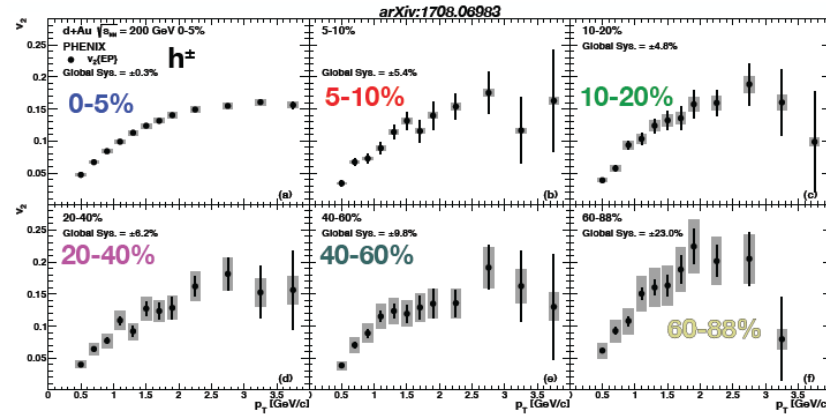
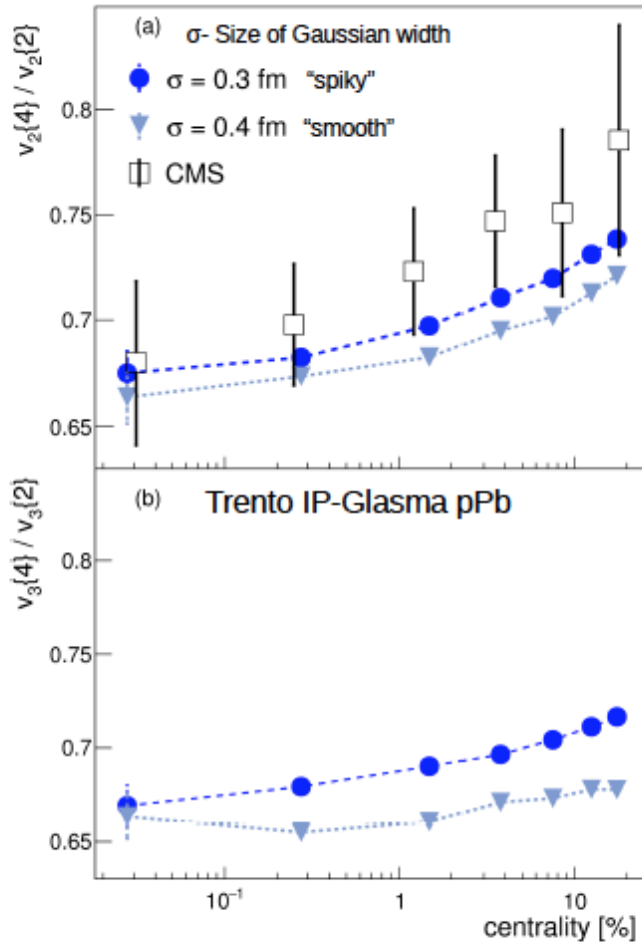
Talk by Bzdak



surprising scaling

$$\sqrt{\langle a_1^2 \rangle} \sim \frac{1}{N_{ch}^{0.5}}$$

Going ahead... important benchmarks in the CGC flow/hydro debate

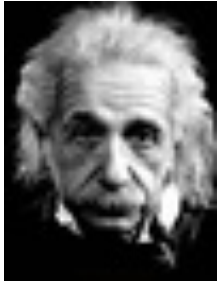


Giacalone, JNH, Ollitrault Phys.Rev. C95 (2017) no.1, 014913

Event engineering & d/3He +Au very important to understand in both frameworks

Theory Talks by Noronha-Hostler, Niemi, Soto-Antoso, Tribedy, ...
 Experimental talks by McGlinchey, Hill, Nagle, Magdy, Nie, Lacey

Apologies to anyone whose talks I could not cover...but this is not a summary!



“A theory is something nobody believes, except the person who made it. An experiment is something everybody believes, except the person who made it.”

Significant advances in both theory and experiment. Keeping this wise remonstrance in mind, we may uncover more remarkable things!

Significant advances in both theory and experiment. Keeping this wise remonstrance in mind, we may uncover more remarkable things!



I am sure I speak for all participants in offering warm thanks to the organizers for a very interesting conference and for their outstanding hospitality in this beautiful city!

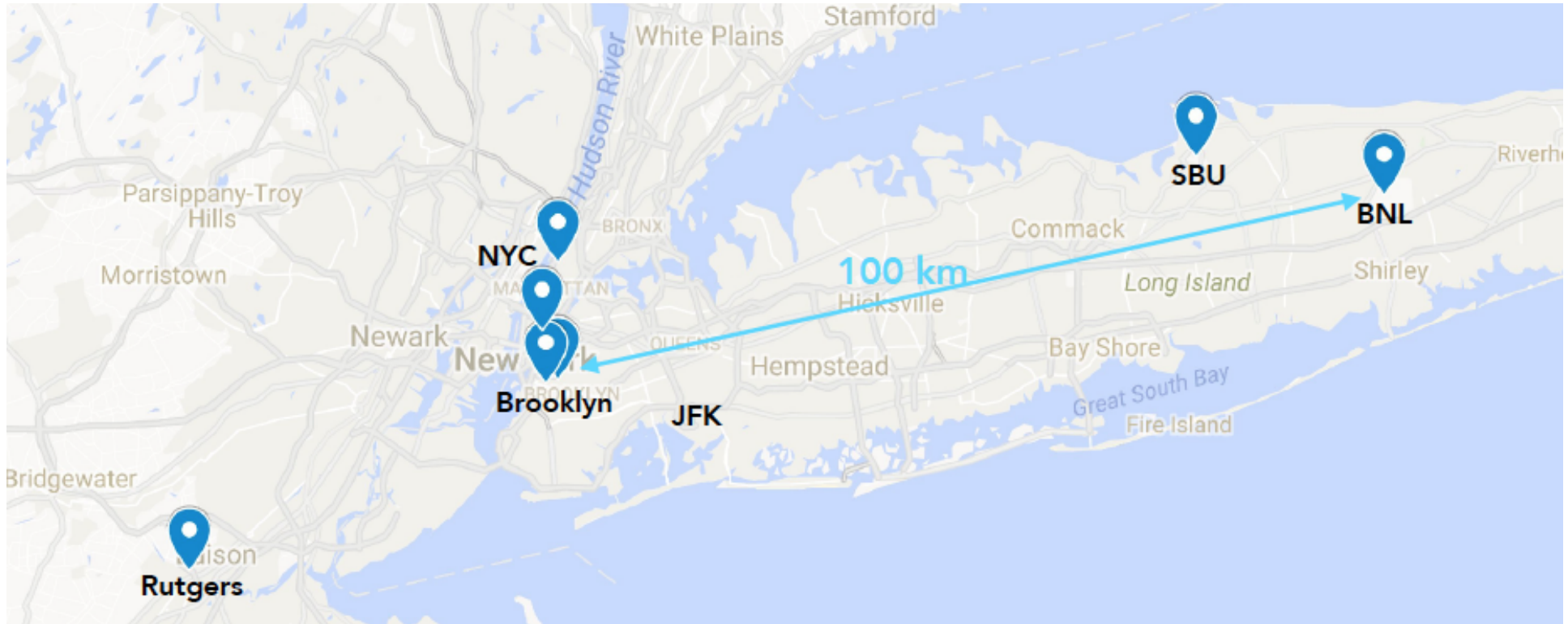
Dziękuję Ci bardzo!!

Announcing... IS2019NY



**Late May or early June 2019,
in the New York City area**

Several locations being considered: *Stony Brook* and *Brooklyn* on Long Island & *Manhattan* (Columbia/CUNY) in NYC and *New Brunswick* (Rutgers) in NJ



Interesting possibilities at each location...we expect a final site selection & fixing of the date soon...

Local organizing committee:

Peter Steinberg & Raju Venugopalan (BNL, co-Chairs)

- | | |
|--|--------------------------|
| • Bjoern Schenke | BNL, theory & experiment |
| • Dave Morrison | |
| • Lijuan Ruan | |
| • Thomas Ullrich | |
| • Agnes Mocsy | Pratt |
| • Adrian Dumitru | Baruch |
| • Stefan Bathe | |
| • Brian Cole | Columbia |
| • Derek Teaney | |
| • Tom Hemmick | Stony Brook |
| • Jiangyong Jia | |
| • Sevil Salur | Rutgers |
| • Jaki Noronha Hostler | |
| • Helen Caines | Yale |
| • 2 Temple faculty as JLab representatives | |

10 EXPERIMENT
6 THEORY

All the major experiments
are represented on the LOC

Springtime in New York:



We look forward to welcoming you all to the next edition of Initial Stages!