

New TMD approach and xFitter

Hannes Jung (DESY)

- Why TMDs are needed
 - TMDs for hadron-hadron collisions
- New developments (together with F. Hautmann, A. Lelek, V. Radescu, R. Zlebcik)
 - parton branching algorithm to solve evolution equations
 - benchmark tests
 - advantages for integrated PDFs
 - determination of TMD densities at LO and NLO with xFitter

TMDs – what is it ?

- TMDs (Transverse Momentum Dependent parton distribution)
 - at very small transverse momenta
 - typically for small q_t in DY production, or semi-inclusive DIS
 - at very small x – un-integrated PDFs
 - essentially only gluon densities (CCFM, BFKL etc)
 - new approach to cover all transverse momenta from small k_t to large k_t as well as to cover all x and all μ^2
 - parton branching method (described here)
 - for an overview of different approaches and state-of-the-art discussion see R. Angeles-Martinez et al
Transverse momentum dependent (TMD) parton distribution functions: status and prospects. Acta Phys. Polon., B46(12):2501–2534, 07 2015, arXiv 1507.05267

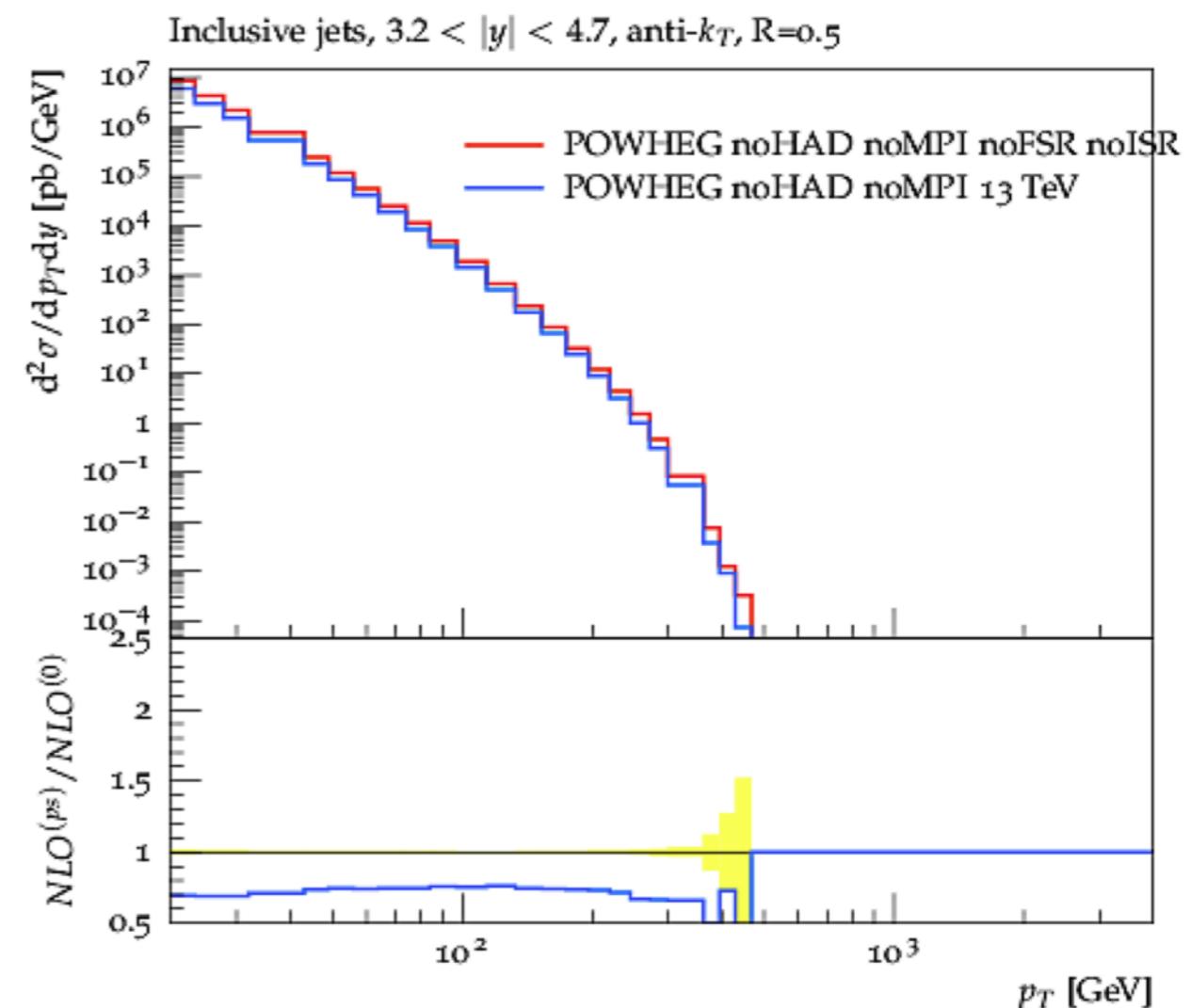
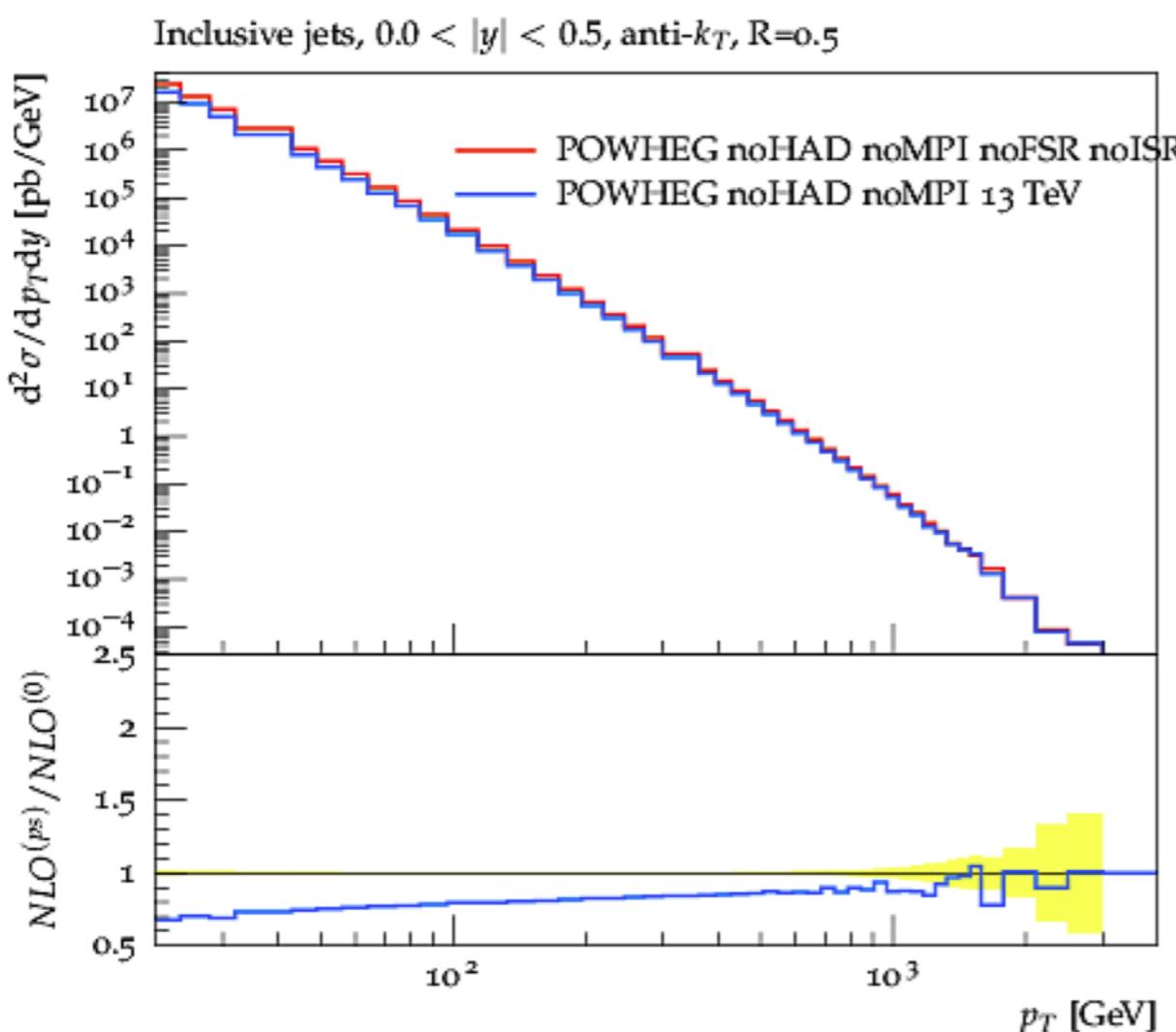
Why TMDs are needed ?

- use NLO+PS to calculate:

$$K^{PS} = \frac{N_{NLO-MC}^{(ps)}}{N_{NLO-MC}^{(0)}}$$

Approach described in: S. Dooling et al
Phys.Rev., D87:094009, 2013.

- Corrections to be applied to fixed order NLO calculations:
 - kinematic effects: TMDs !
 - radiation outside of jet-cone



TMDs – how to determine ?

- Transverse momentum effects are naturally coming from intrinsic k_t and parton showers
- TMD effects can be significant in all distributions, even for inclusive (or semi-inclusive) distributions at large p_t
- New: parton branching method
 - perform evolution using a parton branching method
 - determine integrated PDF from parton branching solution of evolution eq.
 - check consistency with standard evolution on integrated PDFs
 - at LO, NLO and NNLO
 - determine TMD:
 - since each branching is generated explicitly, energy-momentum conservation is fulfilled and transverse momentum distributions can be obtained

for similar approaches see also:

W. Placzek, K. J. Golec-Biernat, S. Jadach, M. Skrzypek. Acta Phys. Polon., B38:2357–2368, 2007.
H. Tanaka. Prog. Theor. Phys., 110:963–973, 2003.

How to obtain TMDs – the evolution equation

- work in progress, in collaboration with:
 - F. Hautmann (Oxford, Antwerp)
 - A. Lelek (DESY)
 - V. Radescu (CERN)
 - R. Zlebcik (DESY)
- preliminary results reported at REF2016 by
 - A. Lelek
- stay tuned for
 - QCD Moriond (A. Lelek)
 - DIS2017 (A. Lelek, R. Zlebcik)

DGLAP evolution – solution with parton branching method

- differential form: $\mu^2 \frac{\partial}{\partial \mu^2} f(x, \mu^2) = \int \frac{dz}{z} \frac{\alpha_s}{2\pi} P_+(z) f\left(\frac{x}{z}, \mu^2\right)$

$$\Delta_s(\mu^2) = \exp\left(- \int_{\mu_0^2}^{z_M} dz \int_{\mu_0^2}^{\mu^2} \frac{\alpha_s}{2\pi} \frac{d\mu'^2}{\mu'^2} P^{(R)}(z)\right)$$

- differential form using f/Δ_s with

$$\mu^2 \frac{\partial}{\partial \mu^2} \frac{f(x, \mu^2)}{\Delta_s(\mu^2)} = \int \frac{dz}{z} \frac{\alpha_s}{2\pi} \frac{P^{(R)}(z)}{\Delta_s(\mu^2)} f\left(\frac{x}{z}, \mu^2\right)$$

- integral form

$$f(x, \mu^2) = f(x, \mu_0^2) \Delta_s(\mu^2) + \int \frac{dz}{z} \int \frac{d\mu'^2}{\mu'^2} \cdot \frac{\Delta_s(\mu^2)}{\Delta_s(\mu'^2)} P^{(R)}(z) f\left(\frac{x}{z}, \mu'^2\right)$$



no – branching probability from μ_0^2 to μ^2

DGLAP re-sums leading logs...

$$f(x, \mu^2) = f(x, \mu_0^2) \Delta_s(\mu^2) + \int \frac{dz}{z} \int \frac{d\mu'^2}{\mu'^2} \cdot \frac{\Delta_s(\mu^2)}{\Delta_s(\mu'^2)} P^{(R)}(z) f\left(\frac{x}{z}, \mu'^2\right)$$

- solve integral equation via iteration:

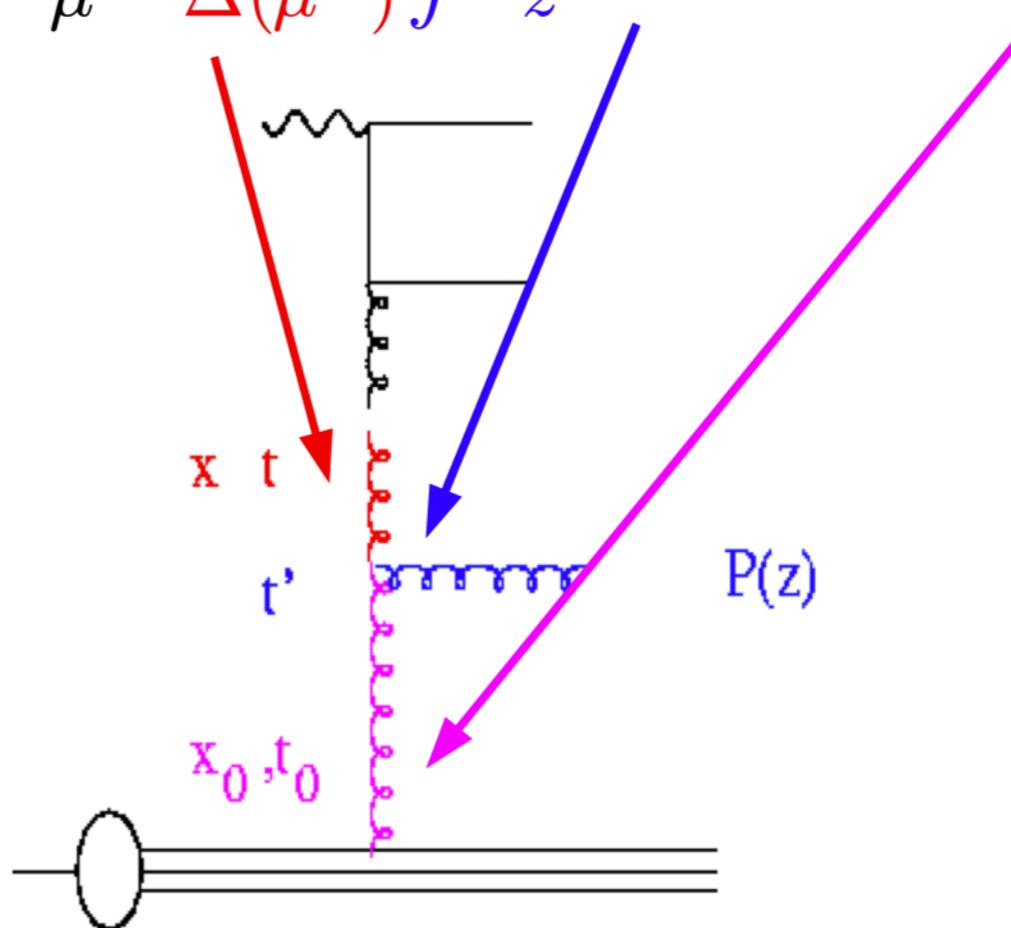
$$f_0(x, \mu^2) = f(x, \mu_0^2) \Delta(\mu^2)$$

from t' to t
w/o branching

branching at t'

from t_0 to t'
w/o branching

$$f_1(x, \mu^2) = f(x, \mu_0^2) \Delta(\mu^2) + \int_{\mu_0^2}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \frac{\Delta(\mu^2)}{\Delta(\mu'^2)} \int \frac{dz}{z} P^{(R)}(z) f(x/z, \mu_0^2) \Delta(\mu'^2)$$



Evolution equation and parton branching method

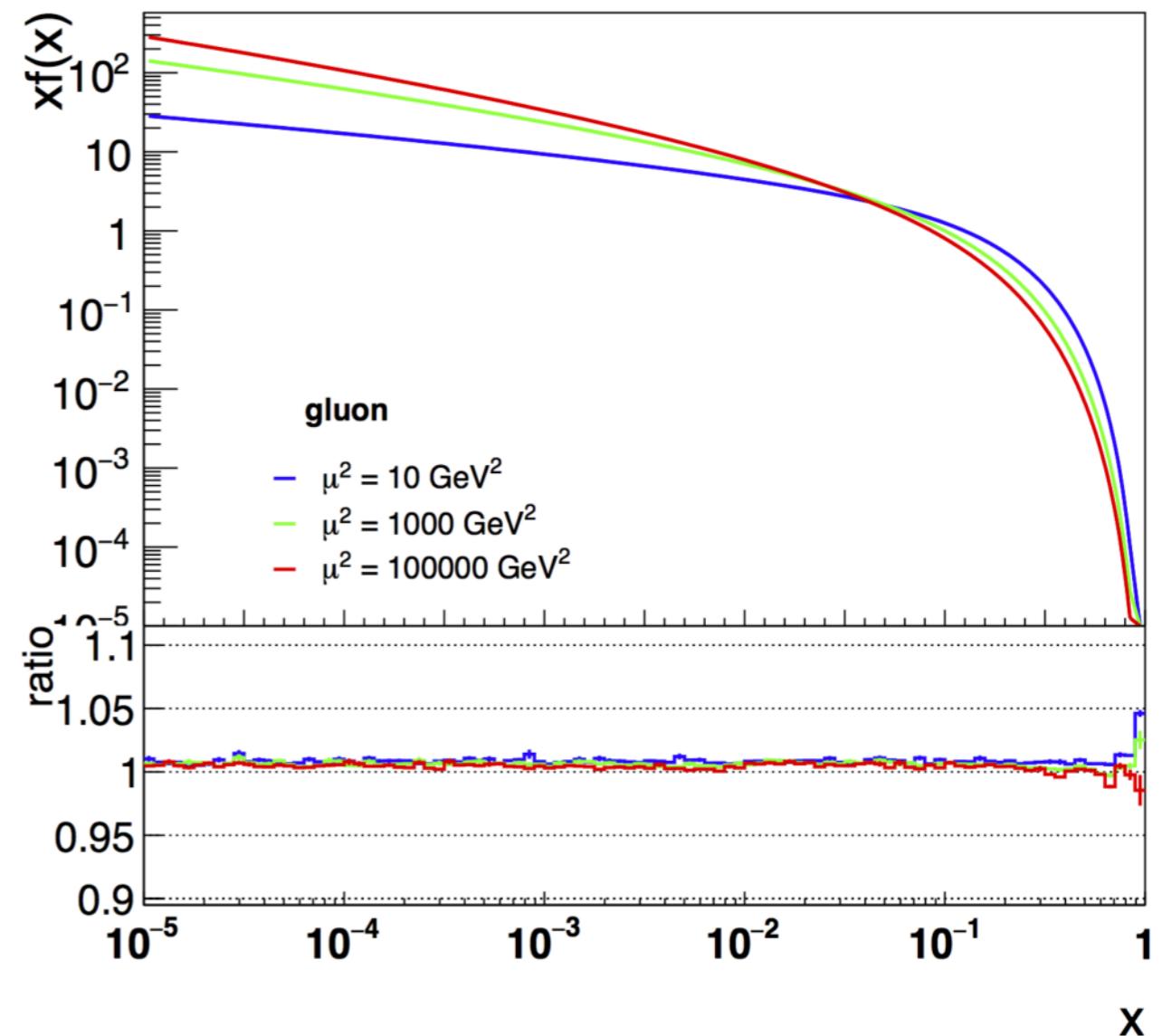
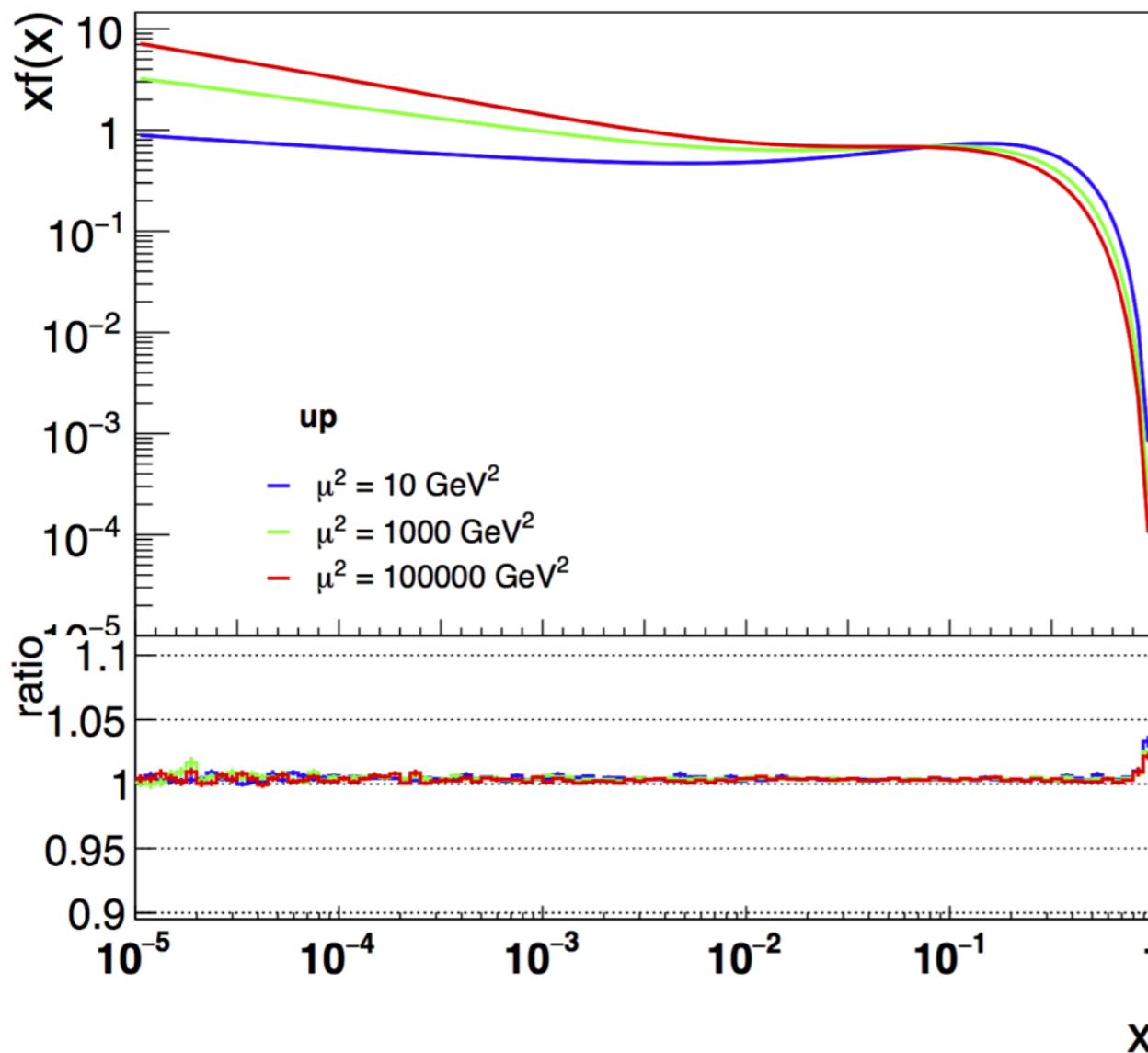
- use momentum weighted PDFs: $xf(x,t)$

$$xf_a(x, \mu^2) = \Delta_a(\mu^2) xf_a(x, \mu_0^2) + \sum_b \int_{\mu_0}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \frac{\Delta_a(\mu^2)}{\Delta_a(\mu'^2)} \int_x^{z_M} dz P_{ab}^{(R)}(\alpha_s, z) \frac{x}{z} f_b\left(\frac{x}{z}, \mu'^2\right)$$

- with $P_{ab}^{(R)}(\alpha_s(t'), z)$ real emission probability (without virtual terms)
 - z_M introduced to separate real from virtual and non-emission probability
- make use of momentum sum rule to treat virtual corrections
 - use Sudakov form factor to treat non-resolvable and virtual corrections

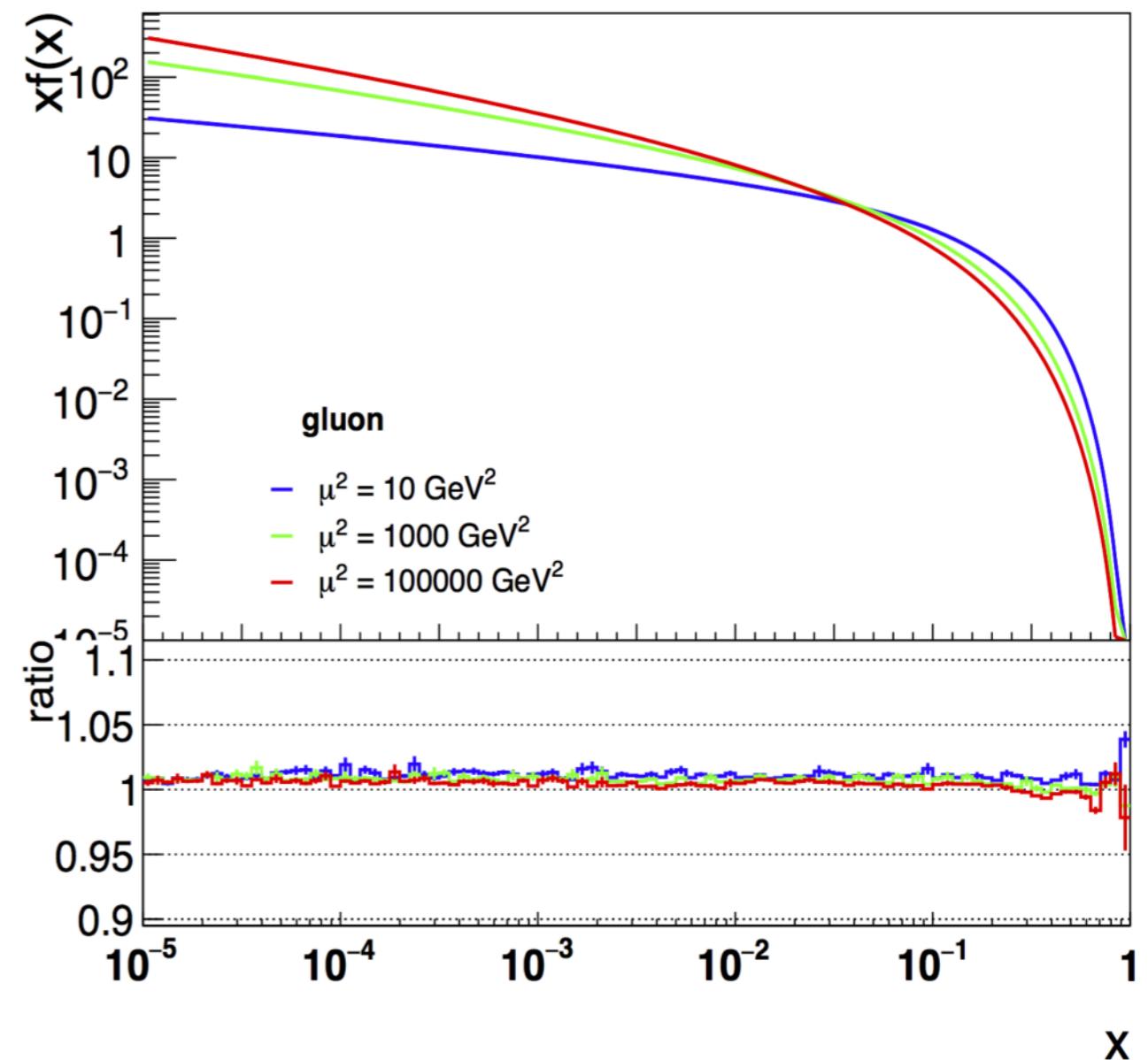
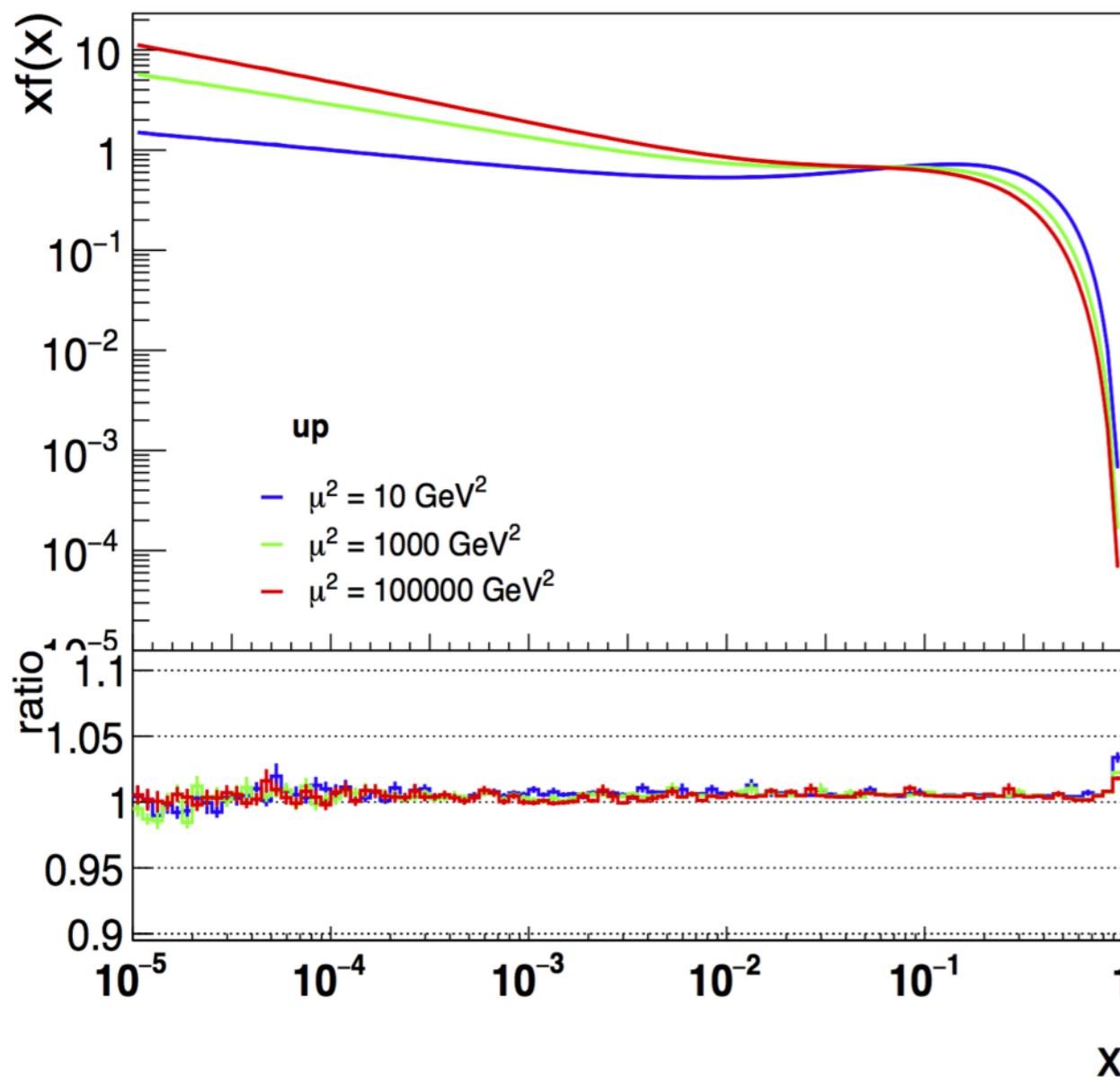
$$\Delta_a(z_M, \mu^2, \mu_0^2) = \exp\left(-\sum_b \int_{\mu_0^2}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \int_0^{z_M} dz z P_{ba}^{(R)}(\alpha_s), z\right)$$

Validation of method with QCDnum at LO



- Very good agreement with LO - QCDnum over all x and μ^2

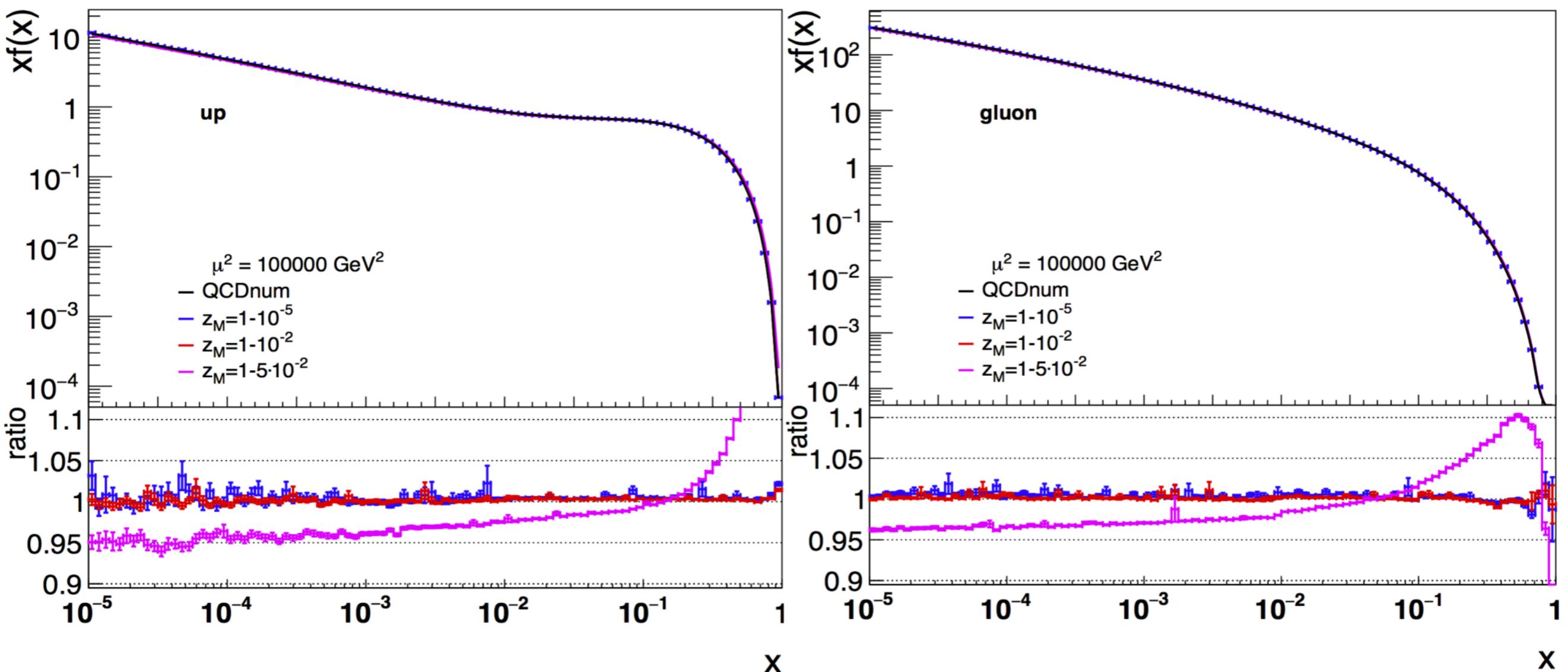
Validation of method with QCDnum at NLO



- Very good agreement with NLO - QCDnum over all x and μ^2
 - the same approach work also at NNLO !

Resolvable branching – at LO and NLO

- Investigate dependence on z_M : separate resolvable from virtual and non-resolvable branchings



- for large enough z_M : results are stable, both at LO and NLO (shown)
- Sudakov treats non-resolvable and virtual branchings to all orders !

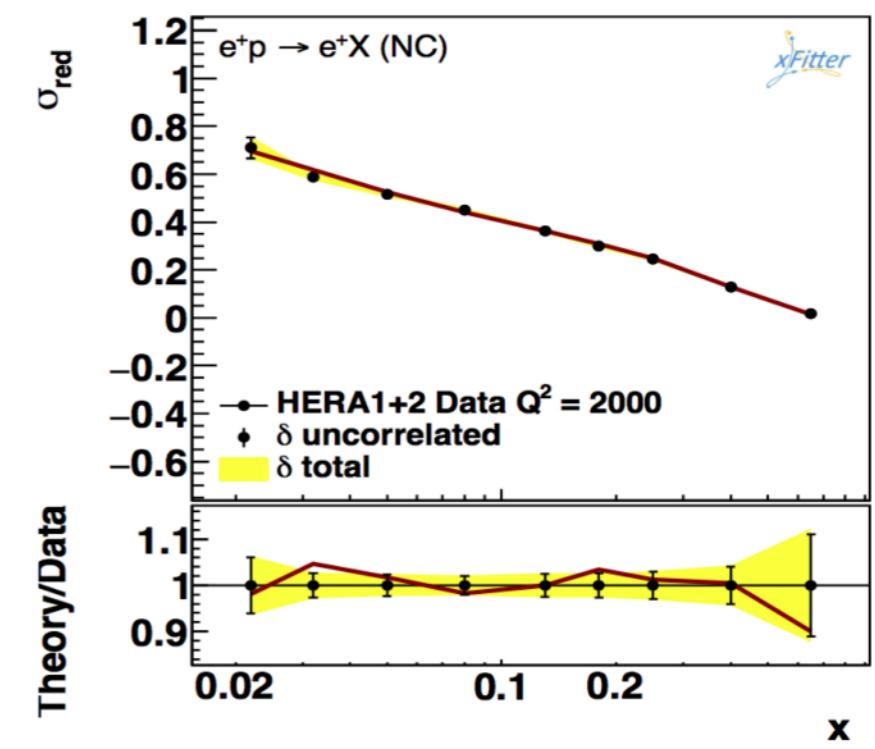
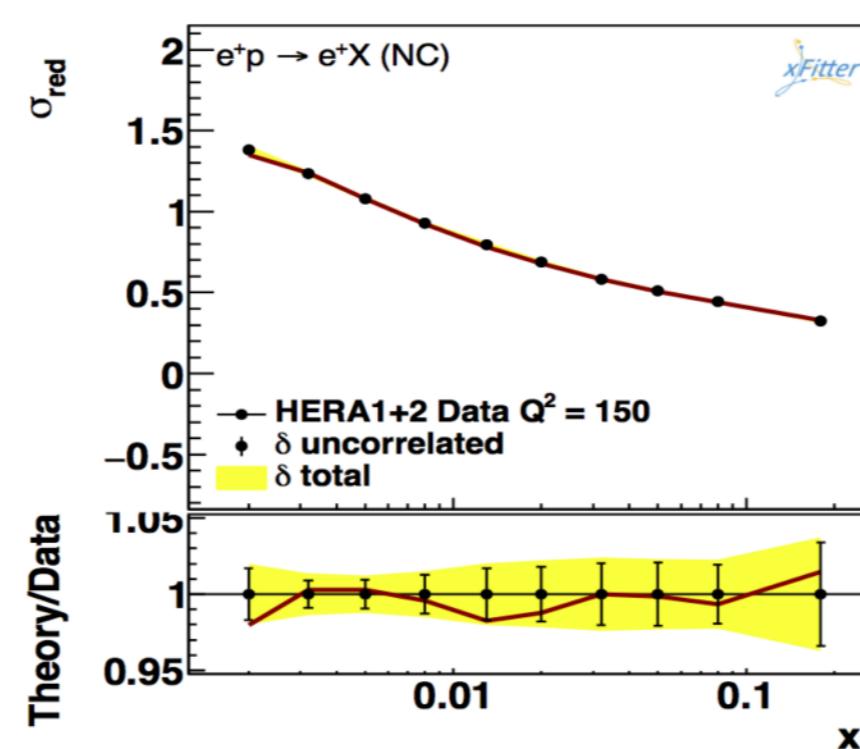
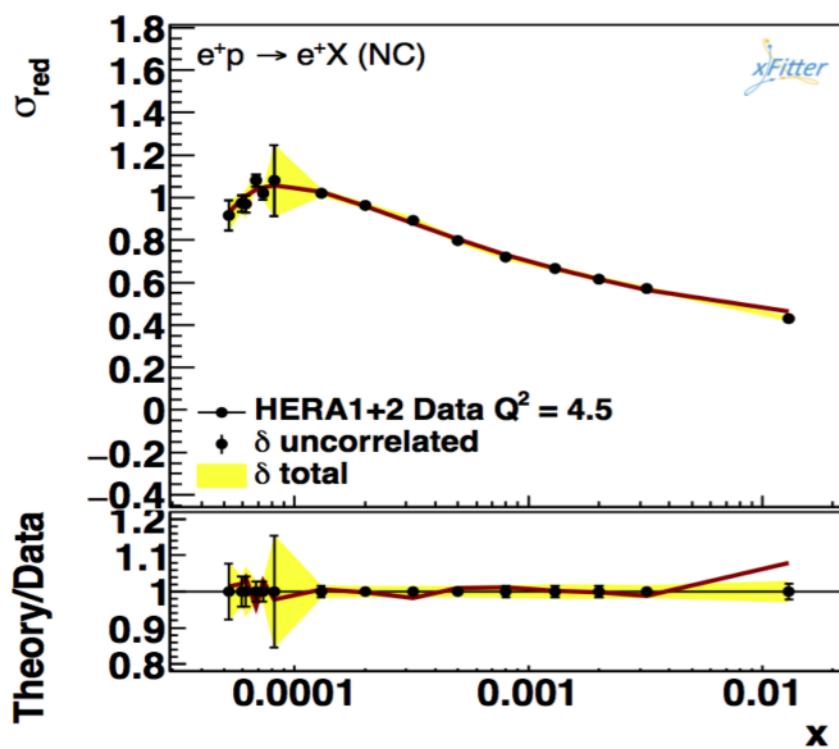
Parton branching method in xFitter

- Determine starting distribution

A. Lelek et al REF 2016

$$\begin{aligned}
 xf_a(x, \mu^2) &= x \int dx' \int dx'' \mathcal{A}_{0,b}(x') \tilde{\mathcal{A}}_a^b(x'', \mu^2) \delta(x'x'' - x) \\
 &= \int dx' \mathcal{A}_{0,b}(x') \cdot \frac{x}{x'} \tilde{\mathcal{A}}_a^b\left(\frac{x}{x'}, \mu^2\right)
 \end{aligned}$$

- fit to HERA data (using xFitter) with $Q^2 \geq 3.5 \text{ GeV}^2$ gives $\chi^2/ndf \sim 1.2$

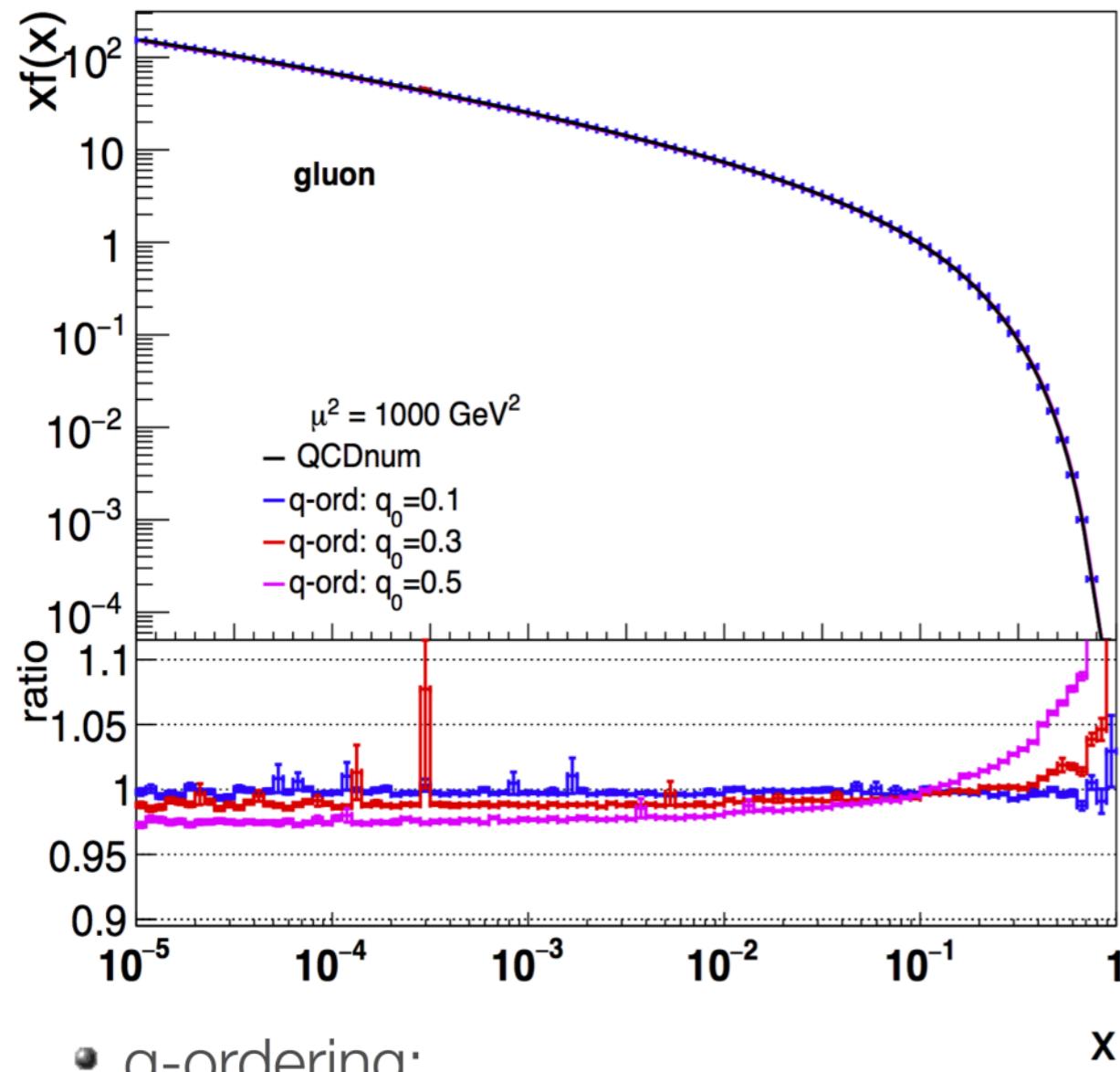


- procedure to fit initial distribution is working and producing results as expected

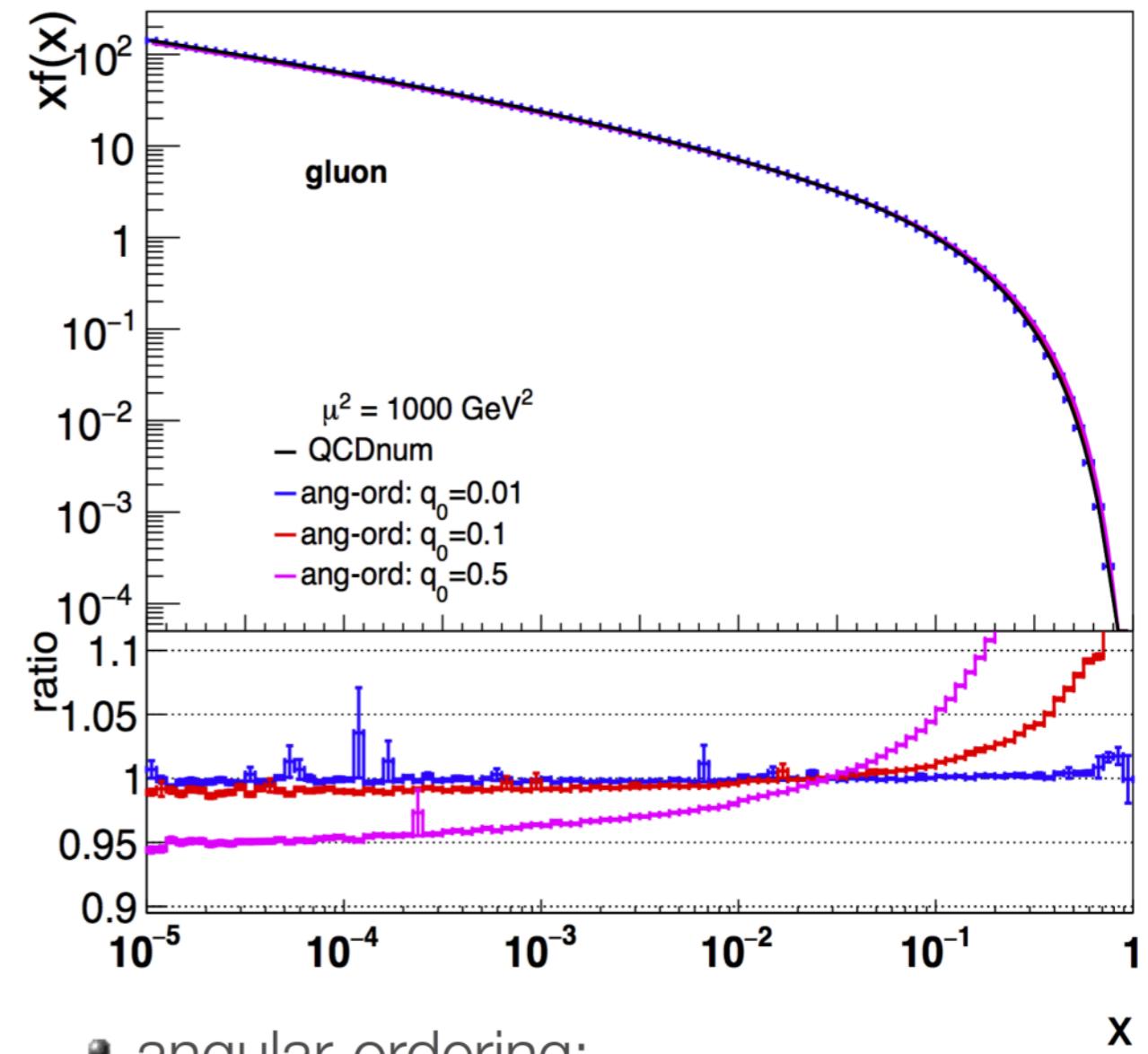
Advantages of parton branching method

- Consistency checks with QCDnum show agreement for inclusive distributions at the 0.1% level
- Advantages of parton branching method for collinear PDFs:
 - studies of different ordering conditions possible for the first time
 - resolvable branchings as defined by:
 - angular ordering with varying $z_M = 1 - q_0/Q$
 - Q^2 ordering with varying $z_M = 1 - q_0^2/Q^2$

parton branching: different ordering conditions



- q-ordering:
 - dependence on parameter for resolvable branching $z_M = 1 - q_0^2/Q^2$
- large x resummation effects in parton densities: resolvable branchings !



- angular-ordering:
 - strong dependence on parameter for resolvable branching $z_M = 1 - q_0/Q$

Advantages of parton branching method

- Consistency checks with QCDnum show agreement for inclusive distributions at the 0.1% level
- Advantages of parton branching method for collinear PDFs:
 - studies of different ordering conditions possible for the first time
 - resolvable branchings as defined by:
 - angular ordering with varying $z_M = 1 - q_0/Q$
 - Q^2 ordering with varying $z_M = 1 - q_0^2/Q^2$
 - different choices of scales in $\alpha_s(\mu)$ possible
 - any investigation which involves details of parton branching kinematics
 - further advantages – determination of TMD parton densities
 - since parton branching kinematics are known, transverse momenta of propagating partons can be calculated – determine TMD

Advantages of parton branching method

- Consistency checks with QCDnum show agreement for inclusive distributions at the 0.1% level
- Advantages of parton branching method for collinear PDFs:
 - studies of different ordering conditions possible for the first time
 - resolvable branchings as defined by:
 - angular ordering with varying $z_M = 1 - q_0/Q$
 - Q^2 ordering with varying $z_M = 1 - q_0^2/Q^2$
 - different choices of scales in $\alpha_s(\mu)$ possible
 - any investigation which involves details of parton branching kinematics
 - further advantages – determination of TMD parton densities
 - since parton branching kinematics are known, transverse momenta of propagating partons can be calculated – determine TMD

Determination of TMD distribution

$$f(x, \mu^2) = f(x, \mu_0^2) \Delta_s(\mu^2) + \int \frac{dz}{z} \int \frac{d\mu'^2}{\mu'^2} \cdot \frac{\Delta_s(\mu^2)}{\Delta_s(\mu'^2)} P^{(R)}(z) f\left(\frac{x}{z}, \mu'^2\right)$$

- solve integral equation via iteration:

$$f_0(x, \mu^2) = f(x, \mu_0^2) \Delta(\mu^2)$$

from t' to t
w/o branching

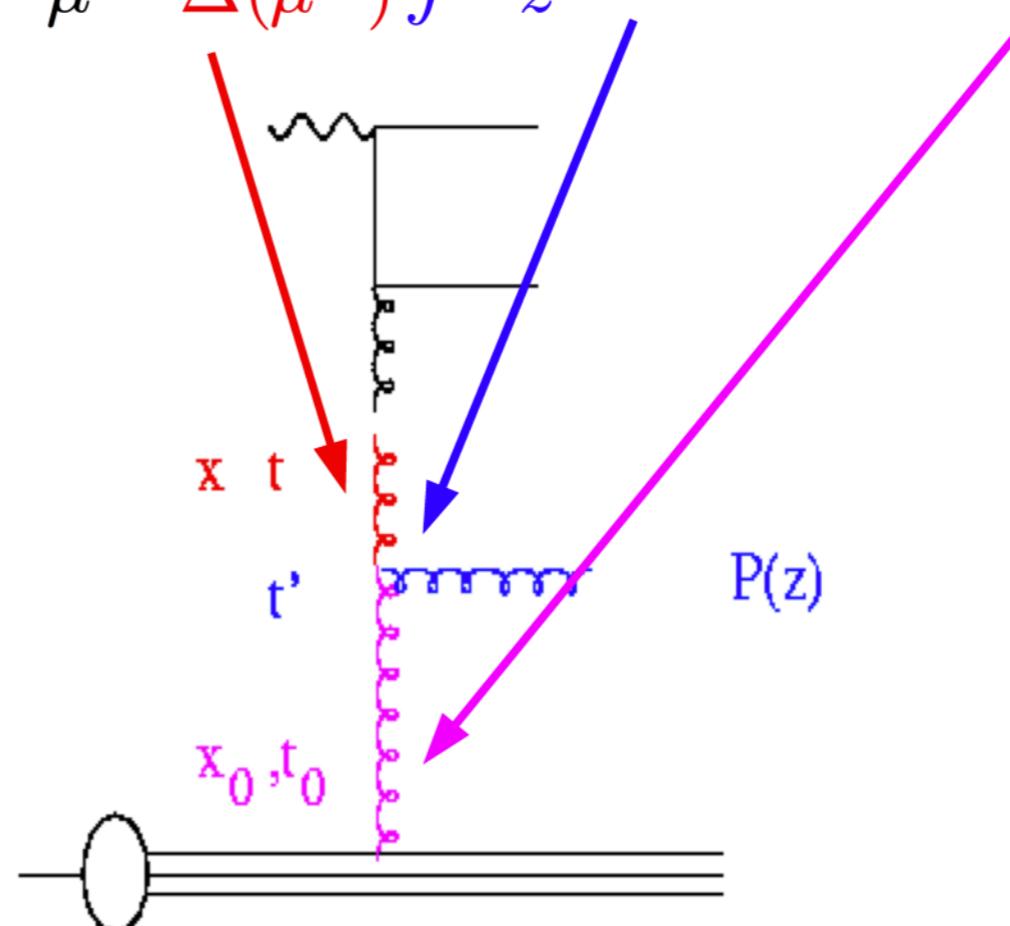
branching at t'

from t_0 to t'
w/o branching

$$f_1(x, \mu^2) = f(x, \mu_0^2) \Delta(\mu^2) + \int_{\mu_0^2}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \frac{\Delta(\mu^2)}{\Delta(\mu'^2)} \int dz P^{(R)}(z) f(x/z, \mu_0^2) \Delta(\mu'^2)$$

- in every step, kinematics are known:

- calculate k_t of propagator



Determination of TMD distribution

$$f(x, \mu^2) = f(x, \mu_0^2) \Delta_s(\mu^2) + \int \frac{dz}{z} \int \frac{d\mu'^2}{\mu'^2} \cdot \frac{\Delta_s(\mu^2)}{\Delta_s(\mu'^2)} P^{(R)}(z) f\left(\frac{x}{z}, \mu'^2\right)$$

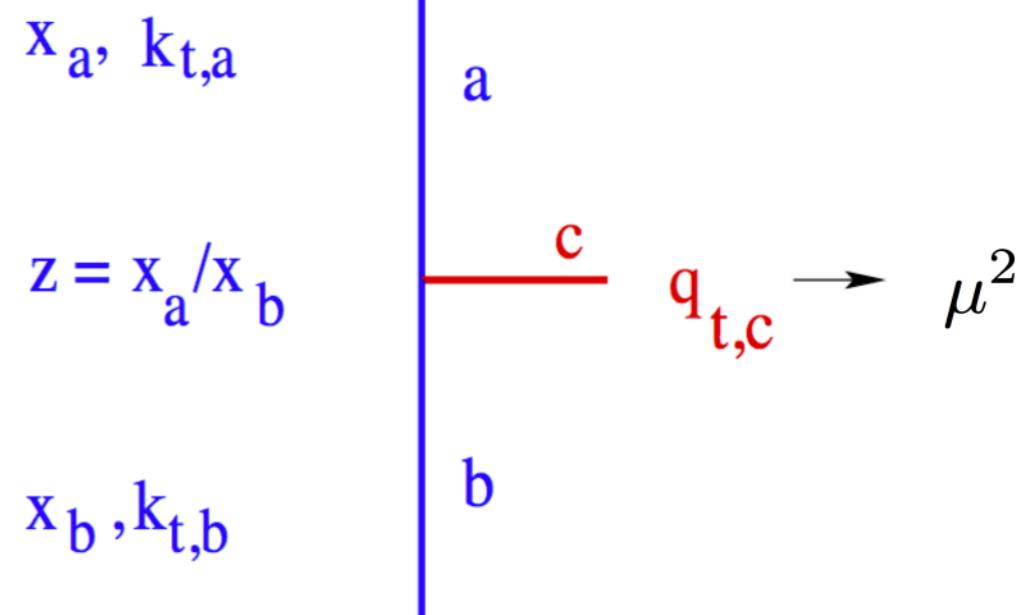
- solve integral equation via iteration:

$$f_0(x, \mu^2) = f(x, \mu_0^2) \Delta(\mu^2)$$

$$f_1(x, \mu^2) = f(x, \mu_0^2) \Delta(\mu^2) + \int_{\mu_0^2}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \frac{\Delta(\mu^2)}{\Delta(\mu'^2)} \int \frac{dz}{z} P^{(R)}(z) f(x/z, \mu_0^2) \Delta(\mu'^2)$$

- in every step, kinematics are known:

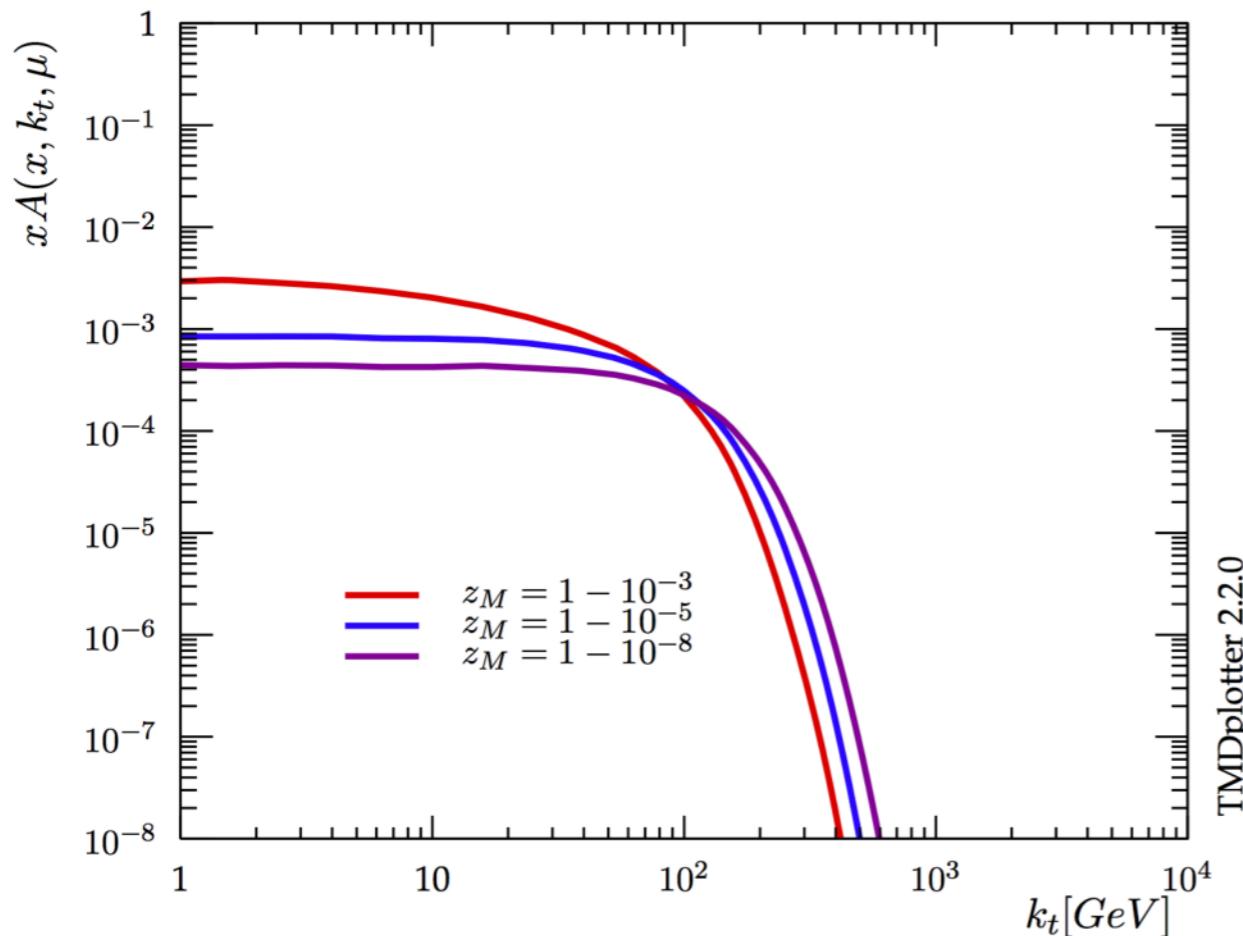
- calculate k_t of propagator
- need correspondence:
 - $q_t^2 = \mu^2$ with q_t emitted parton
OR
 - $q_t^2 = (1-z) \mu^2$, q^2 - ordering



Determination of TMD distribution

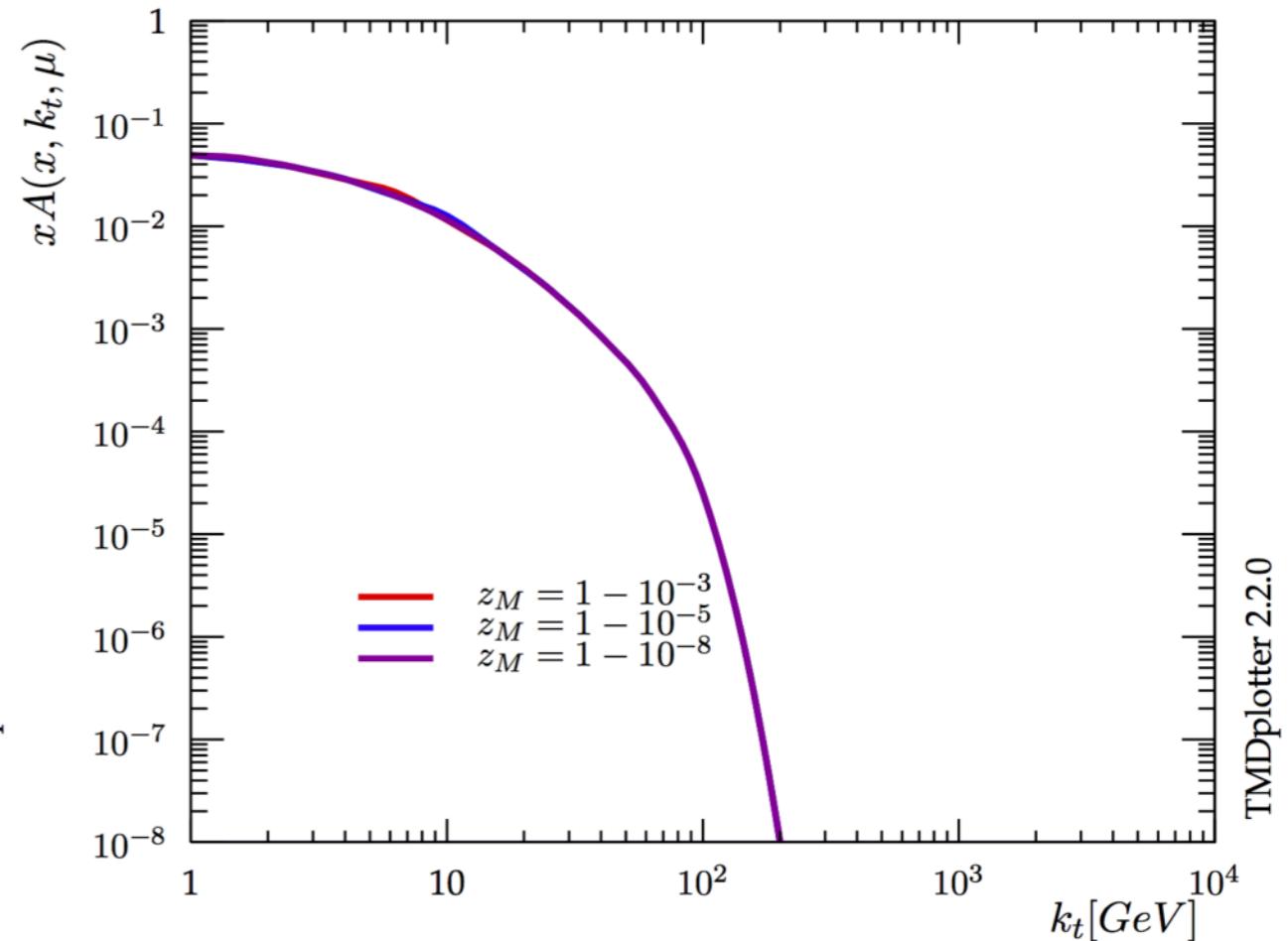
- naïve q_t - ordering
 - $q_t^2 = \mu^2$ with q_t emitted parton
 - $k_t = k'_t + q_t$

gluon, $x = 0.01, \mu = 100 \text{ GeV}$



- q^2 - ordering
 - $q_t^2 = (1-z) \mu^2$
 - $k_t = k'_t + q_t \sqrt{(1-z)}$

gluon, $x = 0.01, \mu = 100 \text{ GeV}$

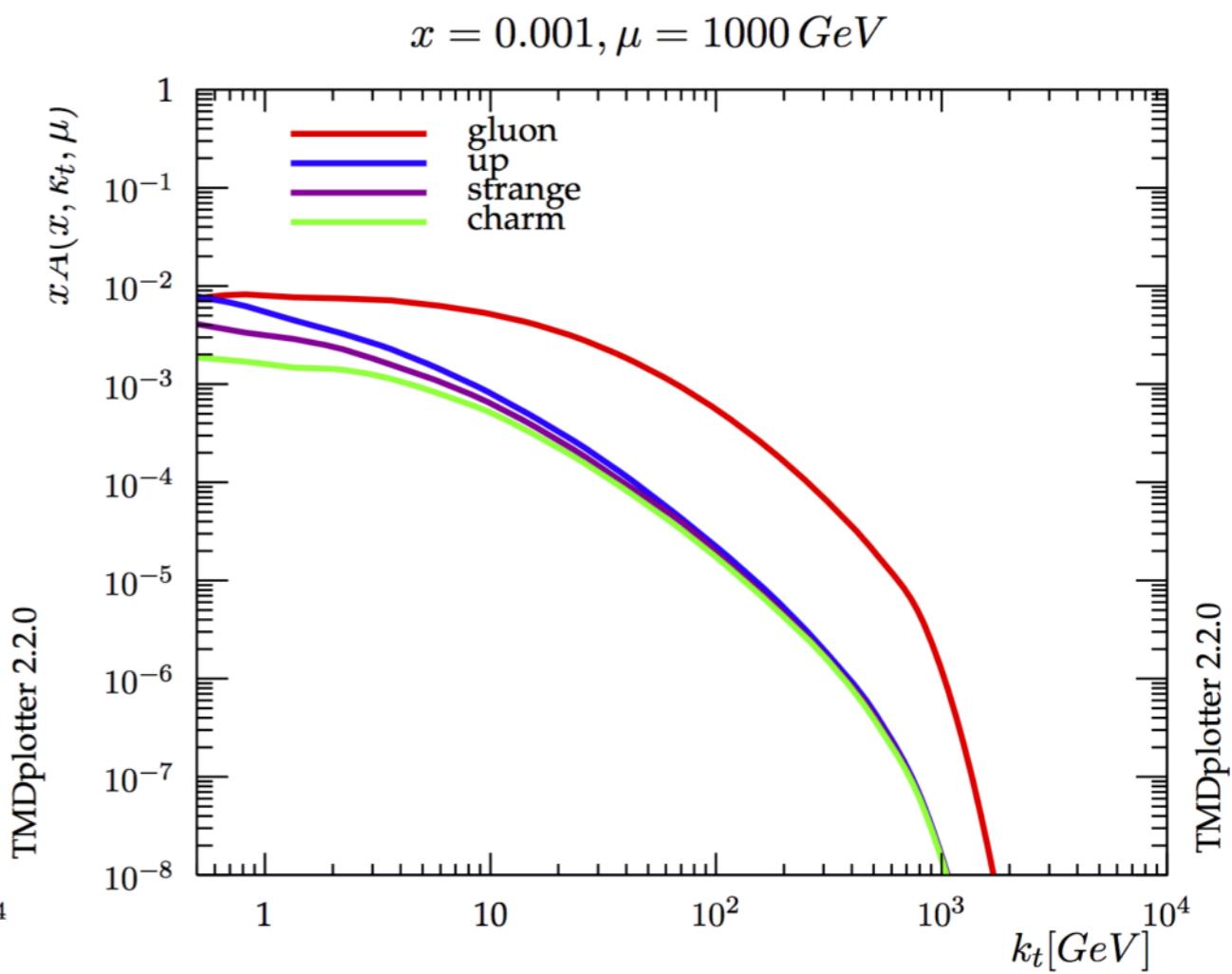
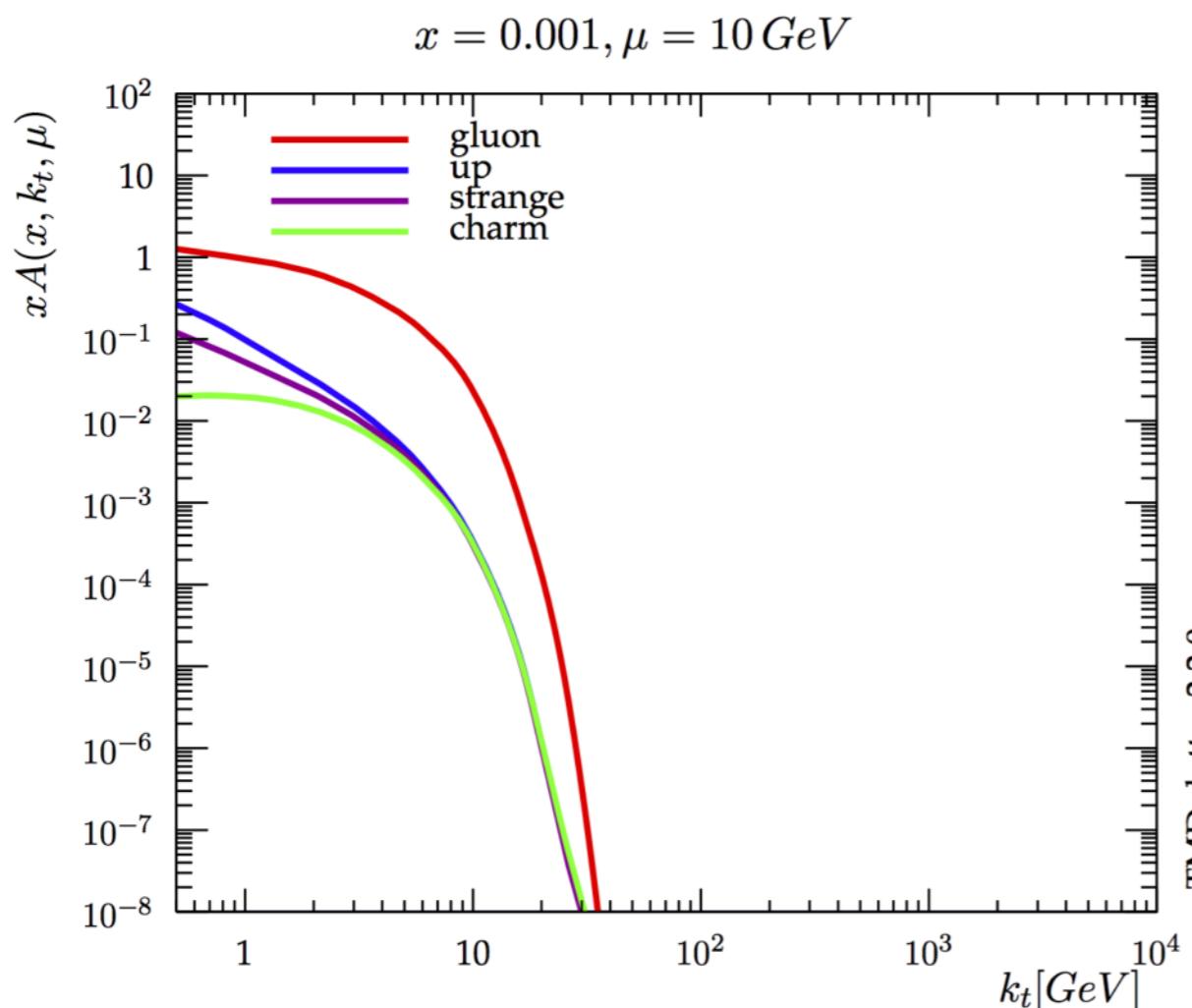


- Huge z_{max} dependence due to many soft gluons ($z_{max} \rightarrow 1$)

- Due to q^2 ordering, soft gluons are suppressed \rightarrow stable results

TMD distributions for different flavors

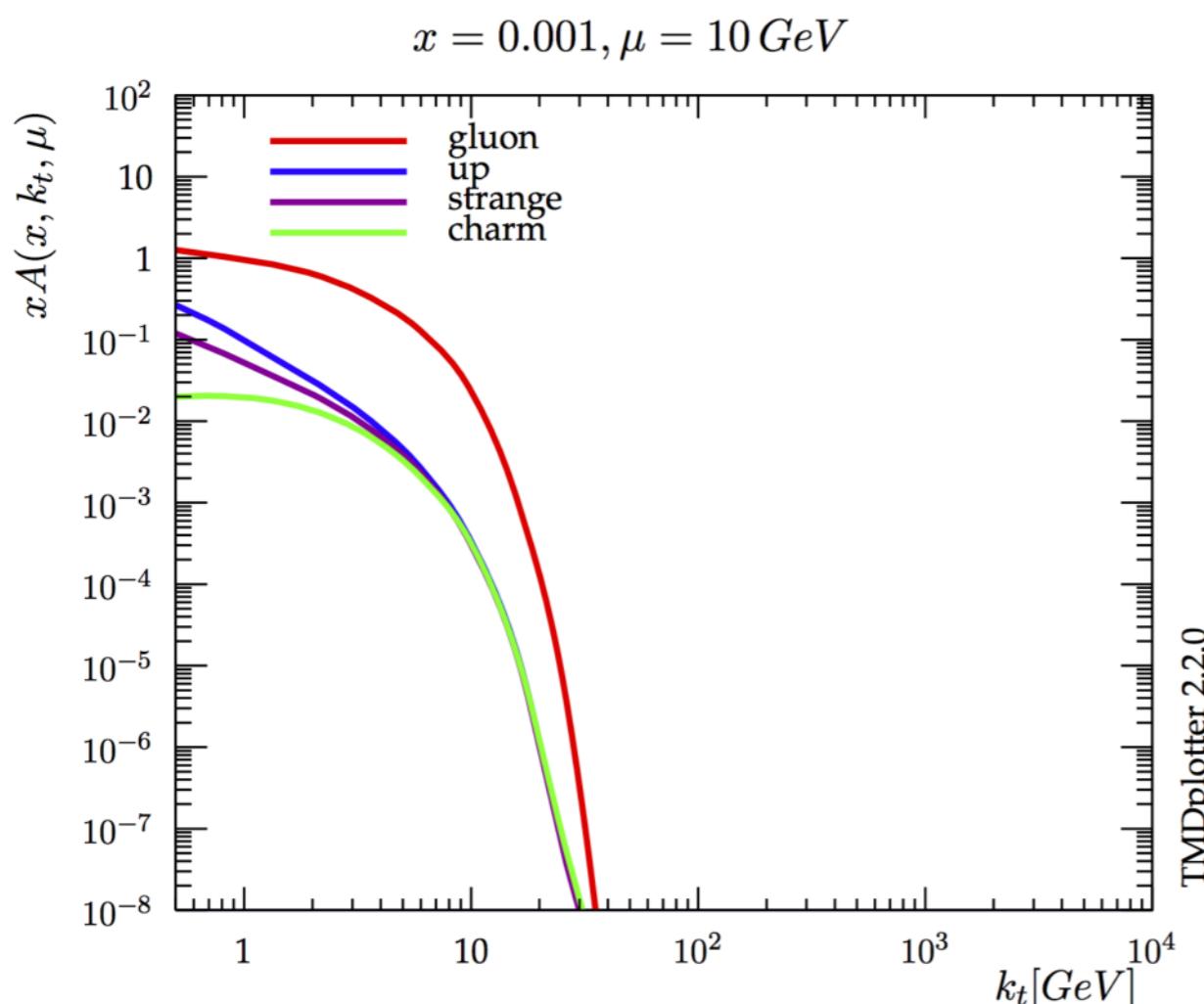
- with parton-branching method, TMD distribution for all flavors can be determined.



- at small k_t intrinsic (gauss) distribution is used \rightarrow subject to fit at small k_t
- at $k_t \geq Q_0$, k_t – distribution comes entirely from evolution,
- no free parameters, except association of evolution scale with q_t**

TMD distributions for different flavors

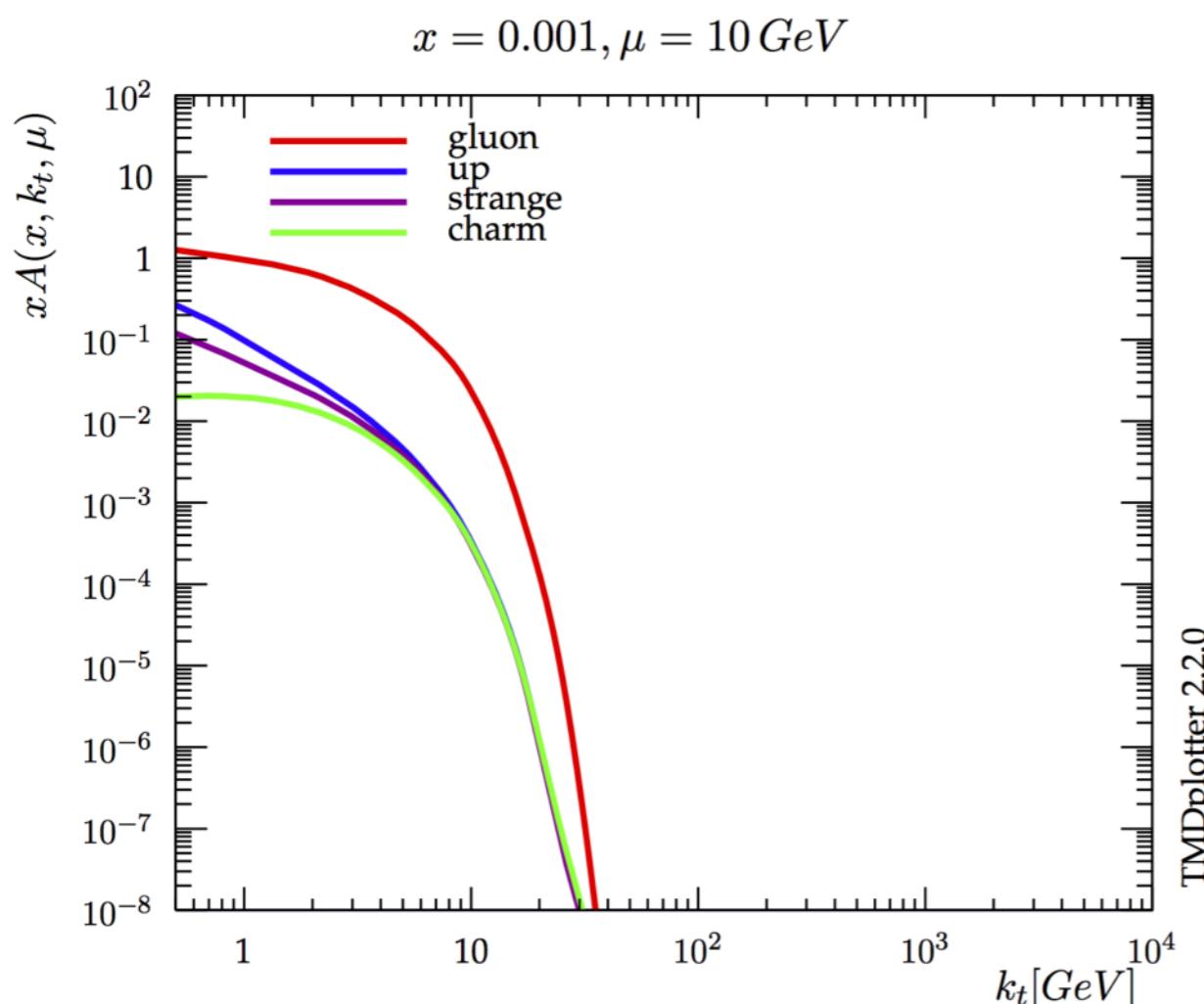
- with parton-branching method, TMD distribution for all flavors can be determined.



- small k_t : distributions are different

TMD distributions for different flavors

- with parton-branching method, TMD distribution for all flavors can be determined.



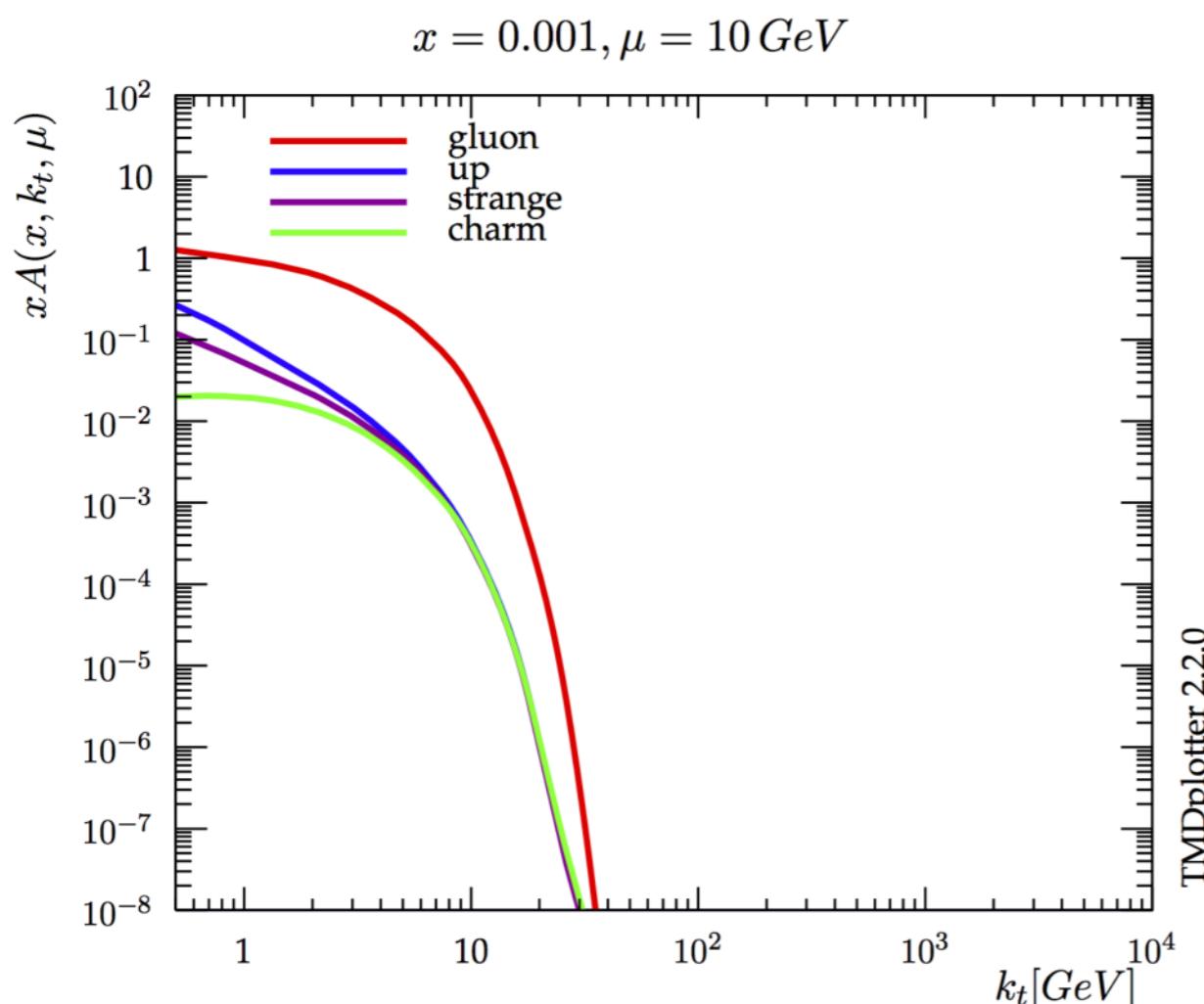
- small k_t : distributions are different

$$x f_a(x, \mu^2) = \Delta_a(\mu^2) x f_a(x, \mu_0^2) + \dots$$

→ depends on distribution at starting scale

TMD distributions for different flavors

- with parton-branching method, TMD distribution for all flavors can be determined.



- small k_t : distributions are different

$$x\mathcal{A}(x, k_\perp^2, \mu^2) = \Delta_a(\mu^2) x\mathcal{A}(x, k_\perp^2, \mu_0^2) + \dots$$

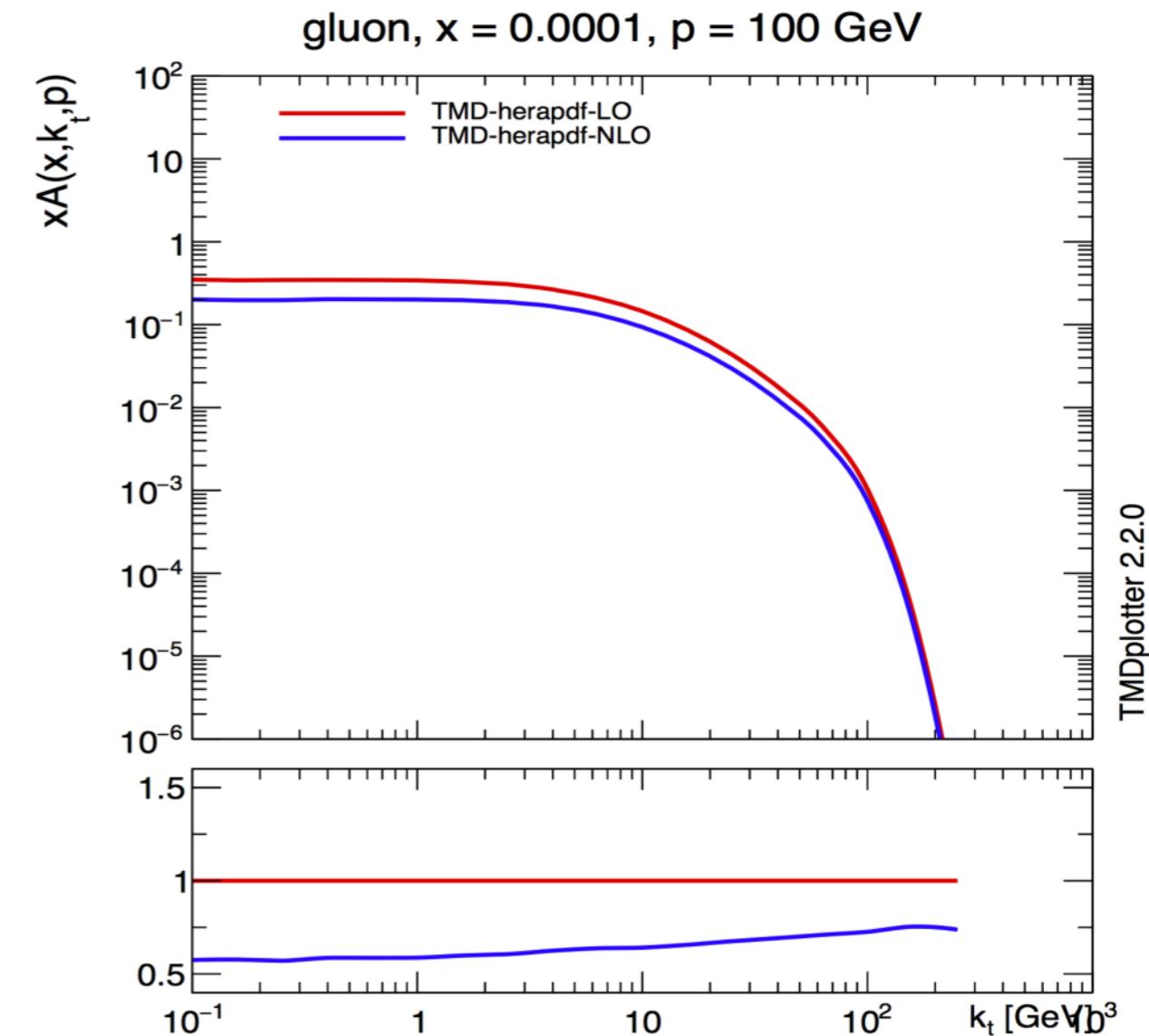
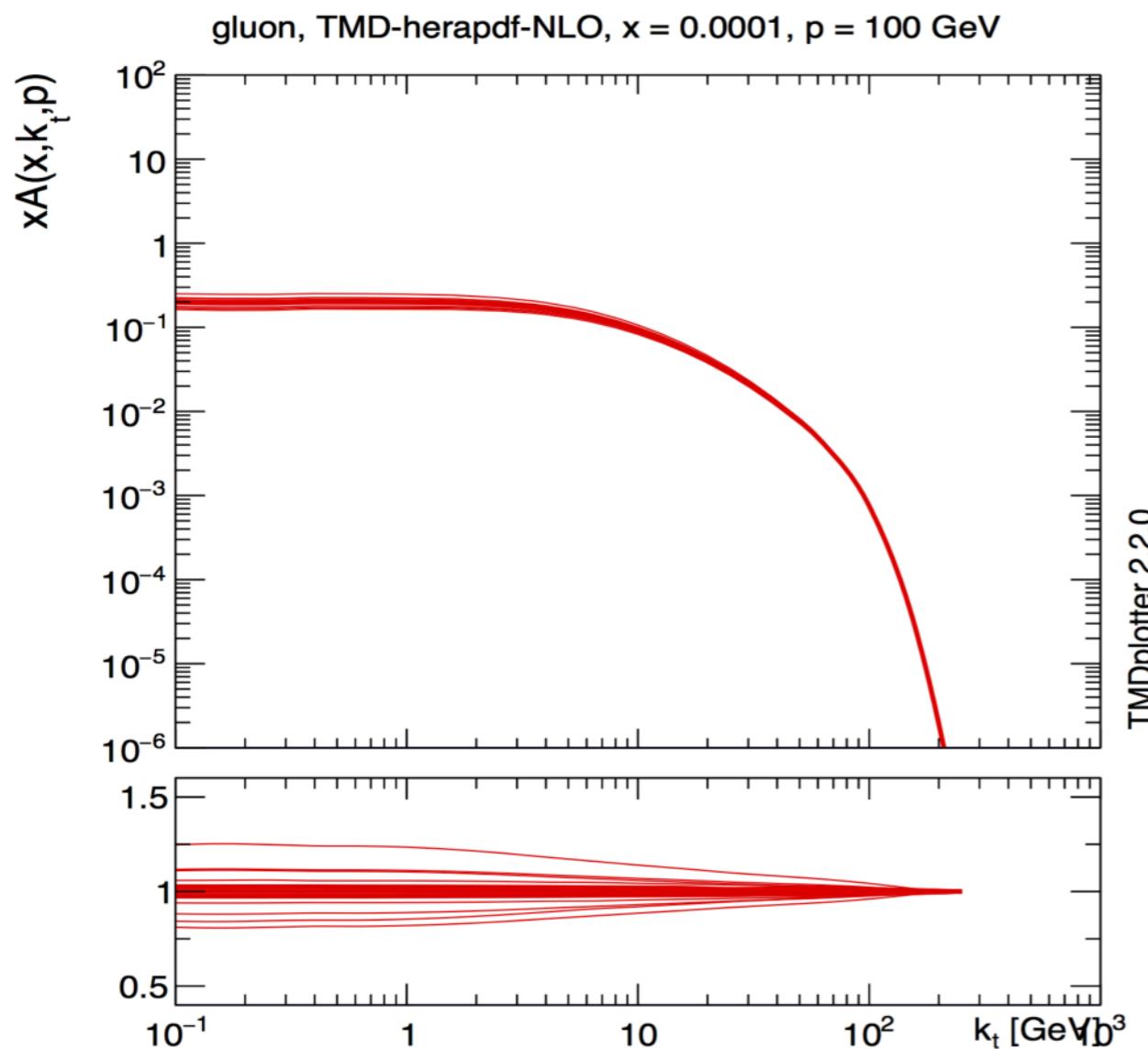
→ depends on distribution at starting scale
- large k_t : quarks have similar shape

$$x\mathcal{A}(x, k_\perp^2, \mu^2) = \dots +$$

$$\sum_b \int_{\mu_0}^{\mu^2} \frac{d\mu'^2}{\mu'^2} \int_x^{z_M} dz P_{ab}^{(R)} \frac{x}{z} \mathcal{A}\left(\frac{x}{z}, k_\perp'^2, \mu'^2\right)$$

→ distributions similar due to parton evolution
- k_t -distributions at large scales depend on initial and evolved distributions

TMD distributions from xFitter: herapdf-type



- Transverse momentum distributions including uncertainties from herapdf fit
 - only experimental uncertainties

- TMD distribution is different for NLO and LO
 - LO and NLO have different number of branchings !

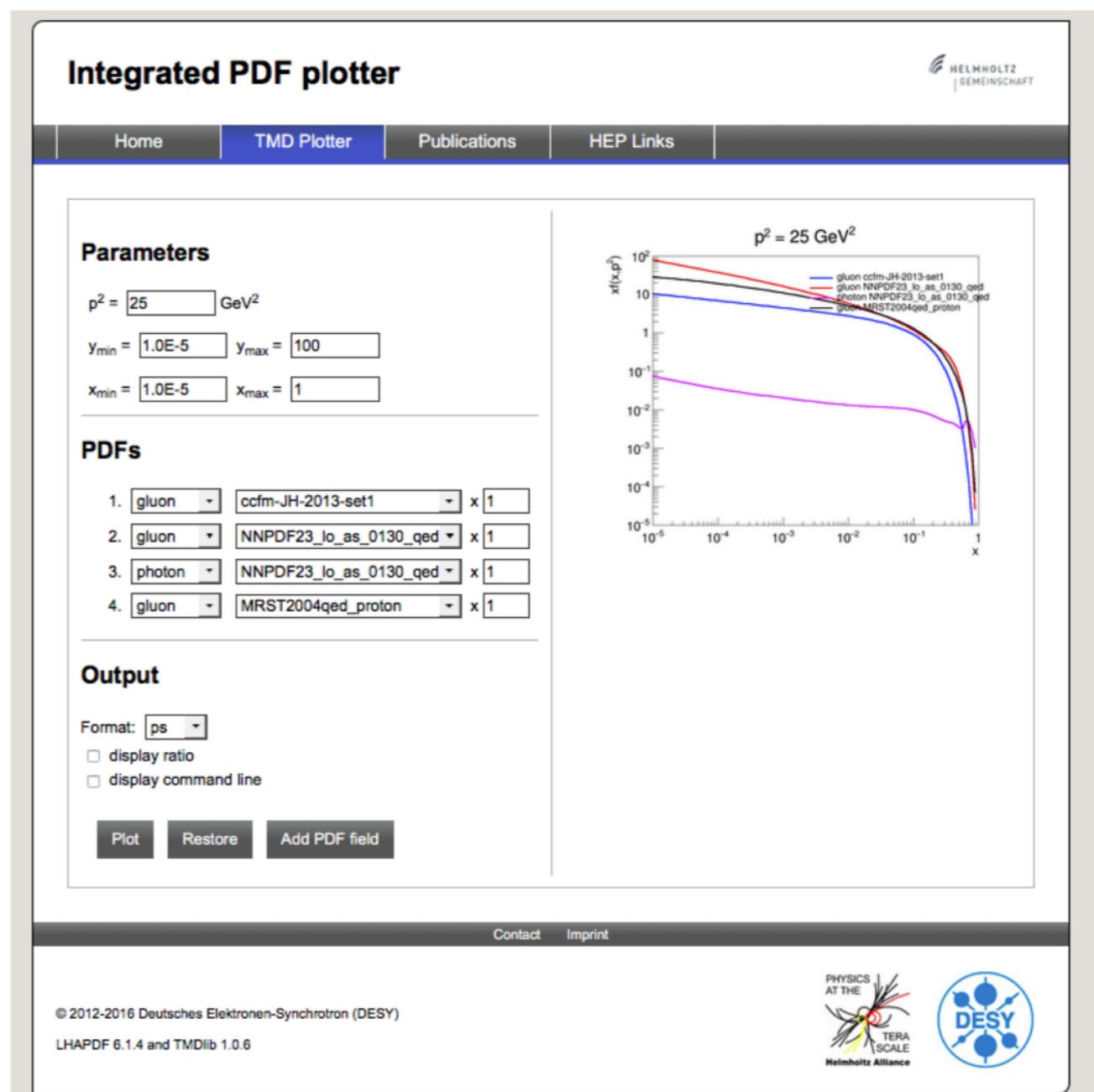
Where to find TMDs ? TMDlib and TMDplotter

- TMDlib proposed in 2014 as part of REF workshop and developed since
- combine and collect different ansaetze and approaches:

<http://tmd.hepforge.org/> and
<http://tmdplotter.desy.de>

- ➔ TMDlib: a library of parametrization of different TMDs and uPDFs (similar to LHAPdf)

TMDlib and TMDplotter: library and plotting tools for transverse-momentum-dependent parton distributions, *F. Hautmann et al.* arXiv 1408.3015, Eur. Phys. J., C 74(12):3220, 2014.



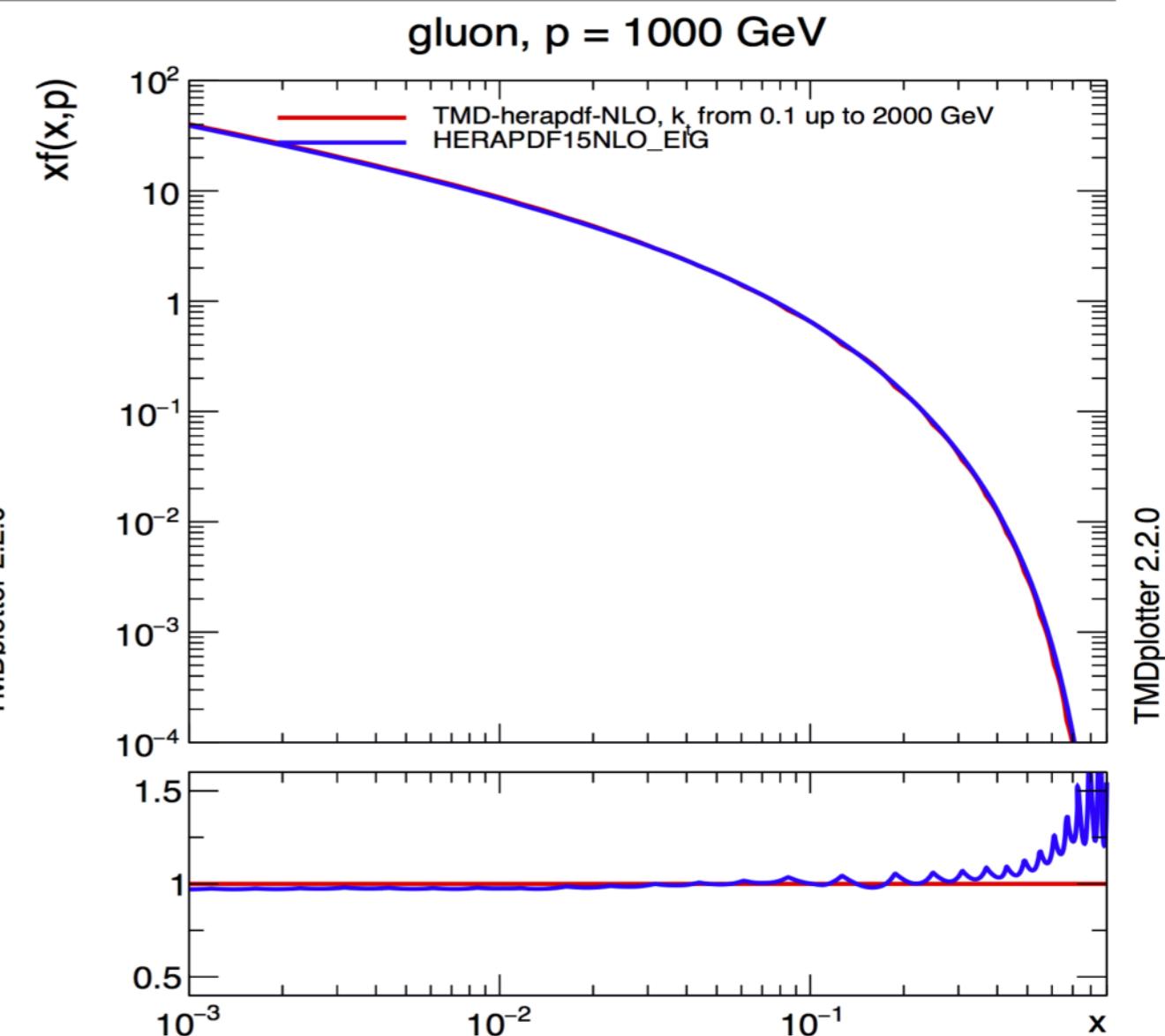
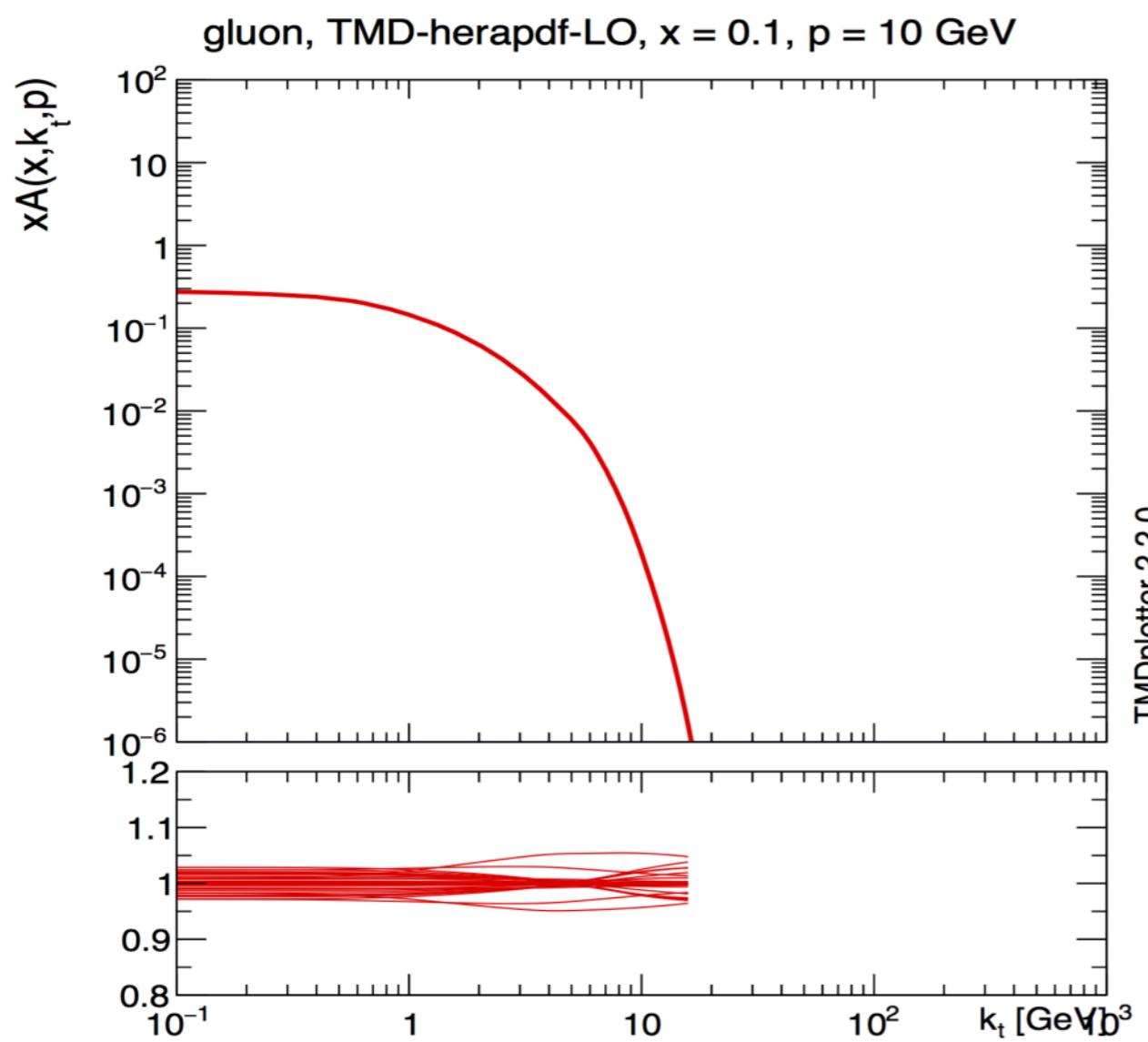
- ➔ Also integrated pdfs (including photon pdf are available via LHAPDF)
- Feedback and comments from community is needed – just use it !

Conclusion

- transverse momenta of interaction partons can be important for precision physics
 - need for TMDs
- Parton Branching method developed for solving DGLAP equation at LO, NLO and NNLO
 - consistence for collinear (integrated) PDFs shown
- method directly applicable to determine k_t distribution (as would be done in PS)
 - TMD distributions for all flavors determined at LO and NLO, without free parameters
 - TMD evolution implemented in xFitter – applicable for DIS processes
- Application in calculations for LHC processes, like DY, jets etc
 - using TMD distributions which reproduce collinear PDFs at NLO and NNLO

Appendix

TMD distributions from xFitter: herapdf-type



- Transverse momentum distributions including uncertainties from herapdf fit
 - essentially constant at fixed x and μ^2

- iTMD distribution: (integrated TMD) compared to herapdf