

Model Selection for Pion Photoproduction

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Analysis of photoproduction reactions

- **HBChPT, RBChPT**, low-energy neutral pion photoproduction, $E_\gamma \leq 170$ MeV. Bernard, Kaiser, Meißner, Scherer, Tiator
- Polynomial parametrizations which incorporate unitarity in the S -wave, $E_\gamma \sim 185$ MeV. Hornidge, C. Fernández-Ramírez, Bernstein
- ChPT calculations including isospin breaking. Varu, Hanhart, Hoferichter, Nogga, Kubis, Nogga
- **RBChPT** with $\Delta(1232)$, $E_\gamma \sim 200$ MeV. Hiller Blin, Ledwig, Vicente Vacas
- Effective field theories, **K -matrix** parametrizations, **Regge** parametrizations, higher energies. Anisovich, Drechsel, Kamalov, Tiator, Haidenbauer, Krewald, Meißner
- Phenomenological parametrizations, basic principles of S matrix, unitarity, correct threshold behavior, Fermi-Watson Theorem, **SAID** approach. Workman, Paris, Briscoe, Strakovsky

Shrinkable methods: Ridge and Lasso

Linear model:

$$Y = \beta_0 + \sum_{j=1}^p \beta_j X_j + \epsilon \quad (1)$$

$$\beta_j = \frac{\partial Y}{\partial X_j}; \quad \text{Penalty } \lambda \sum_{j=1}^p f(\beta_j) \quad (2)$$

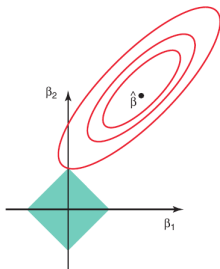
n observations, (x_{ij}, y_i) . Find $\hat{\beta}_i^\lambda$ which minimize,

$$\sum_{i=1}^n \left(y_i - \beta_0 - \sum_{j=1}^p \beta_j x_{ij} \right)^2 + \lambda \sum_{j=1}^p \beta_j^2 \quad \text{RIDGE}$$

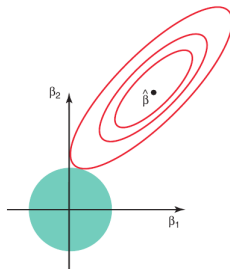
$$\sum_{i=1}^n \left(y_i - \beta_0 - \sum_{j=1}^p \beta_j x_{ij} \right)^2 + \lambda \sum_{j=1}^p |\beta_j| \quad \text{LASSO}$$

Shrinkable methods: Ridge and Lasso

$$\text{minimize}_{\beta} \left\{ \sum_{i=1}^n \left(y_i - \beta_0 - \sum_{j=1}^p \beta_j x_{ij} \right)^2 \right\} \text{ subject to } \sum_{j=1}^p f(\beta_j) \leq s$$



LASSO



RIDGE

Bayesian Interpretation

Linear model:

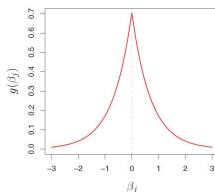
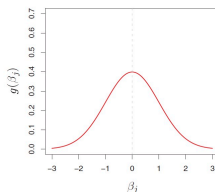
$$Y = \beta_0 + \sum_{j=1}^p \beta_j X_j + \epsilon \quad (3)$$

Assume

$$\epsilon \sim N(\mu, \sigma), \quad p(\beta) = \prod_{j=1}^p g(\beta_j), \quad \text{with } \beta = (\beta_0, \beta_1, \dots, \beta_p)^T$$
$$p(\beta|X, Y) \propto f(Y|X, \beta)p(\beta|X) = f(Y|X, \beta)p(\beta) \quad (4)$$

It follows

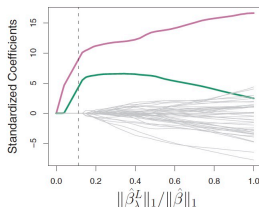
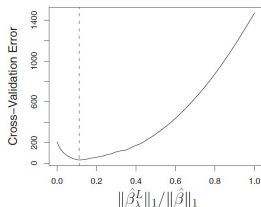
Ridge $g(\beta_j) \sim$ gaussian **LASSO** $g(\beta_j) \sim$ double exponential



Cross validation

How to fix λ ?

1. Discretize λ
2. Split the data set into $k \sim 5$ (or 10) parts
3. In every iteration, set 4 parts to be the **Training set**, 1 part the **Validation set**.
4. Minimize χ^2_T
5. Evaluate Cross Validation Error (CV), $CV = \frac{1}{k} \sum_{i=1}^k (y_i - \hat{y}_i)$ or χ^2_V for every i part and average it for every λ .



Information Criteria

Nested models $\{M_i\}$ with i parameters, $M_1 \prec M_2 \prec \dots \prec M_k$

Example

$$M_1 : y = \beta_0 + \beta_1 x_1 + \epsilon$$

$$M_2 : y = \beta_0 + \beta_1 x_1 + \beta_2 x_2 + \epsilon$$

$M_1 \prec M_2$. How to *evaluate* the models given a set of data?

Training error, $\sum_{i=1}^n (y - \hat{y})^2$, χ^2_T ? *but..* size of test error?.

n , number of observations; k , number of parameters;

$$AIC = -2 \max \log(L(\hat{\theta}|data)) + 2k = \chi^2 + 2k$$

$$AIC_c = AIC + \frac{2k(k+1)}{n-k-1}$$

$$BIC = -2 \max \log(L(\hat{\theta}|data)) + 2\log(n) = \chi^2 + k \log(n)$$

All of them penalize overfitting

Scan a large set of models and compare them.

Advantages

- Ability to recover the true model (least squares needs a large amount of data)
- Minimal predicted errors
- Good variance-bias trade-off
- Reduces the correlations between the parameters
- Select “simple” models

Parametrization for $\gamma p \rightarrow \pi^0 p$

For \mathcal{M} the electric $E_{L_{\pm}}$, and magnetic $M_{L_{\pm}}$ multipoles,

$$\text{Re, Im } \mathcal{M}_{L_{\pm}} = \frac{q_{\pi^0}^l}{m_{\pi^+}^{l+1}} \sum_{i=0}^{i_{\max}} \frac{a_i}{10^{-i}} \left(\frac{\omega_{\pi^0} - m_{\pi^0}}{m_{\pi^+}} \right)^i \quad (5)$$

q_{π^0} c. m. momentum, ω_{π^0} energy, a_i fitting parameters

$$E_{1+} = \frac{1}{6}(P_1 + P_2); M_{1+} = \frac{1}{6}(P_1 - P_2 + 2P_3); M_{1-} = \frac{1}{3}(P_3 + P_2 - P_1) \quad (6)$$

Partial waves Includes P and D waves, and

$$\Delta E_{0+} = i \frac{q_{\pi^+}}{m_{\pi^+}^2} \sum_{i=0}^2 \frac{a_i}{10^{-i}} \left(\frac{q_{\pi^+}}{m_{\pi^+}} \right)^{2i} \quad (7)$$

supplements the S wave multipole to take into account the $\pi^+ n$ threshold cusp. **Total 46 parameters.**

Parametrization for $\gamma p \rightarrow \pi^0 p$

Differential cross section and asymmetry

$$\begin{aligned}\frac{d\sigma}{d\Omega}(\mathbf{s}, \theta) &= \frac{q_{\pi^0}}{k_\gamma} W_T(\mathbf{s}, \theta) \\ \Sigma(\mathbf{s}, \theta) &\equiv \frac{\sigma_\perp - \sigma_\parallel}{\sigma_\perp + \sigma_\parallel} = -\frac{W_S(\mathbf{s}, \theta)}{W_T(\mathbf{s}, \theta)} \sin^2 \theta\end{aligned}\quad (8)$$

Electromagnetic responses, W_T and W_S :

$$\begin{aligned}W_T &= T_0(\mathbf{s}) + T_1(\mathbf{s})\mathcal{P}_1(\theta) + T_2(\mathbf{s})\mathcal{P}_2(\theta) + \dots \\ W_S &= S_0(\mathbf{s}) + S_1(\mathbf{s})\mathcal{P}_1(\theta) + \dots\end{aligned}\quad (9)$$

C. Fernández Ramírez and A. M. Bernstein, PLB724, PRC80; Legendre Polynomials $\mathcal{P}_j(\theta)$, and

$$T_n(\mathbf{s}) = \sum_{ij} \text{Re} \{ \mathcal{M}_i^*(\mathbf{s}) T_n^{ij} \mathcal{M}_j(\mathbf{s}) \} \quad S_n(\mathbf{s}) = \sum_{ij} \text{Re} \{ \mathcal{M}_i^*(\mathbf{s}) S_n^{ij} \mathcal{M}_j(\mathbf{s}) \}$$

being $\mathcal{M}_j(\mathbf{s}) = E_{0+}, E_{1+}, E_{2+}, E_{2-}, M_{1+}, M_{1-}, M_{2+}, M_{2-}$

LASSO

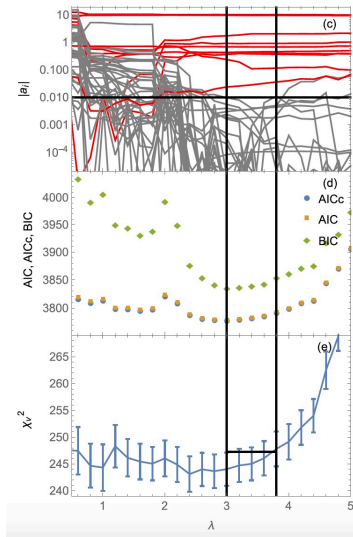
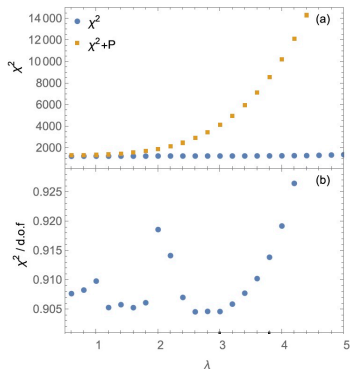
$$\chi_T^2(\lambda) = \chi^2(\lambda) + \lambda^4 \sum_{i=1}^{i_{\max}} |a_i| \quad (10)$$

Test: Lasso in a benchmark model

B_0 model synthetic data
9 parameters, no D waves

D. Hornidge, PRL111 (2013)

Analysis with the 46 parameter model



Test: Lasso in a benchmark model

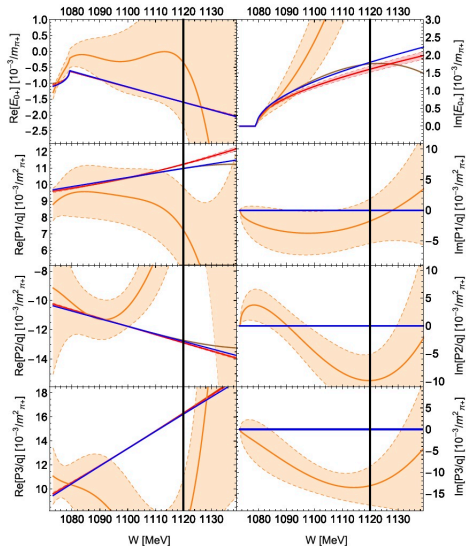
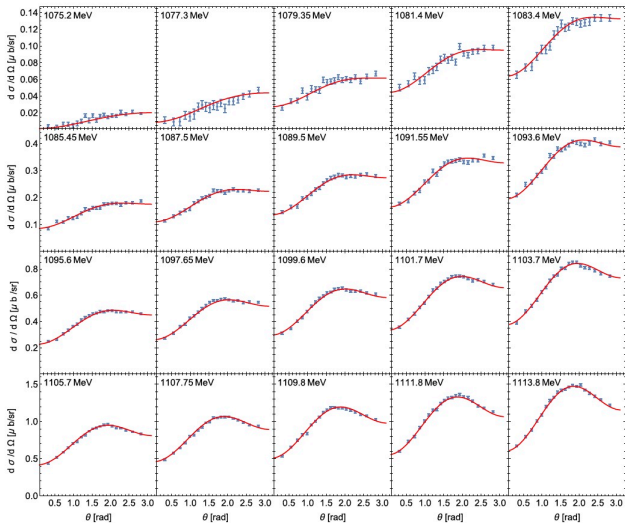


Figure. S and P partial waves. Red curves and bands indicates the Lasso solution. Blue lines stand for the benchmark solution. Orange curve and bands are the solution for $\lambda = 0$. The solution with $\lambda = 3$ is indicated with brown color.

Lasso in a B_0 model: Differential cross section



Lasso in a \mathcal{B}_0 model: Differential cross section

Result:

The analysis with LASSO, cross validation and information criteria, recovers the true model with only one additional parameter, reducing the number of parameters from 46 to 10.

Compare two models

- Model 1, k parameters
- Model 2, $m+k$ parameters

If the true values of the m extra parameters vanish,

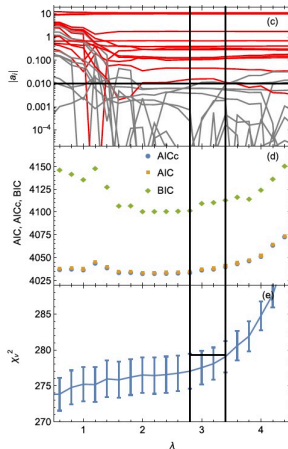
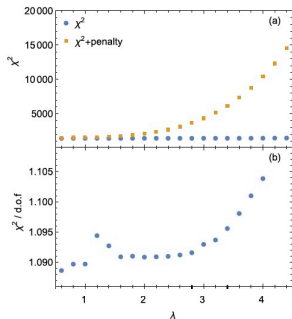
$$y = \frac{(\chi_1^2 - \chi_2^2)/k}{\chi_2^2/(n - m - k)} \quad (11)$$

is $F(k, n - m - k)$ distributed. A value of y beyond a chosen CL limit indicates that the more complex model 2 is significantly better than 1. **We find $y = 1.64$, below the 90% CL interval ending at $y = 2.63$, the overfit is not significantly better than the simplest fit.**

Real data for the $\gamma p \rightarrow \pi^0 p$ reaction

MAMI $d\sigma/d\Omega$, Σ , $d\sigma_T/d\Omega$.

D. Hornidge et al., PRL111,062004 (2013), S. Schumann et al., PLB750, 252 (2015)



$\gamma p \rightarrow \pi^0 p$: partial waves, beam asymmetry

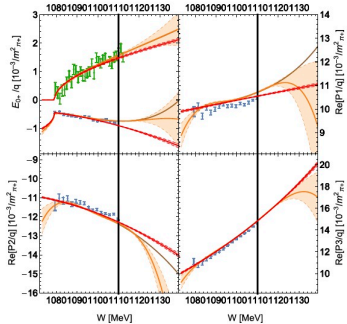


Figure. S and P partial waves. Red curve with bands indicates the Lasso solution with bootstrap.

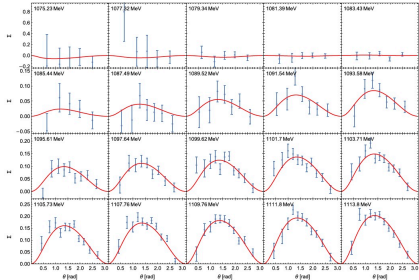
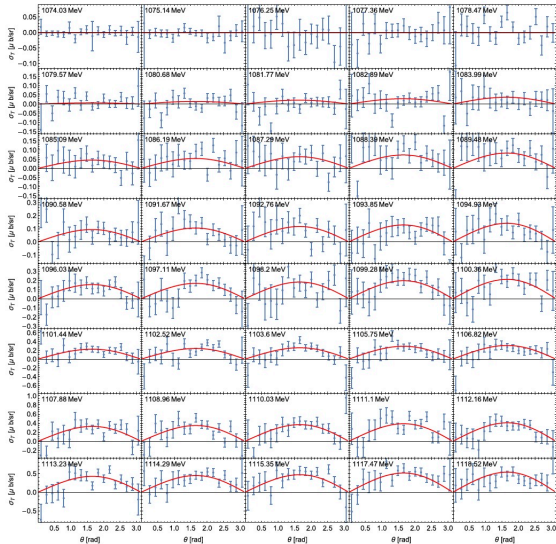


Figure. Beam Asymmetry

$\gamma p \rightarrow \pi^0 p$: Polarized differential cross section $\sigma_T = d\sigma/d\Omega T$



β_0 , Cusp parameter at threshold

$$\beta_0 = (2.41 \pm 0.05) \times 10^{-3} m_{\pi^+}^{-1}$$

Agrees with the value from, S. Schumann et al., PLB750,252,
 $\beta_0 = (2.2 \pm 0.2[\text{stat.}] \pm 0.6[\text{syst.}]) \times 10^{-3} m_{\pi^+}^{-1}$.
(A2 Collaboration, 2015)

Conclusions

- LASSO, in combination with cross validation and criteria from information theory provides a tool to scan large classes of models, selecting the *simplest* model with minimum number of parameters and prediction error
- It has a wide range of applicability, being a promising tool for the analysis of data of the excited baryon and meson spectra in future experiments