

# Confronting Jet Quenching with Jet Grooming: Splitting Function and Jet Mass Distribution in Heavy Ion Collisions

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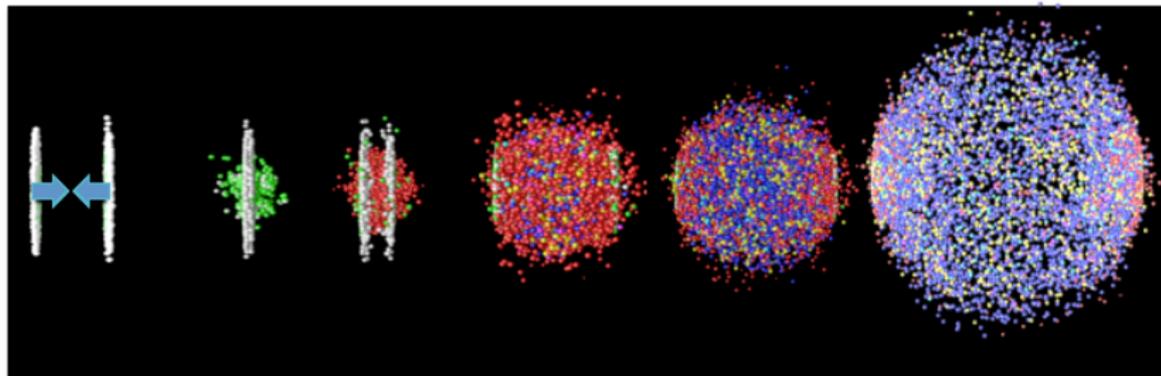
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# Outline

- ▶ Heavy ion jet physics
  - ▶ hard probes of quark-gluon plasma
  - ▶ jet quenching and jet modification
  - ▶ the necessity and utility of jet grooming
- ▶ Hard and soft jet substructure
  - ▶ splitting function and subjet distribution
  - ▶ groomed jet mass with small jet radius
- ▶ Conclusion and outlook

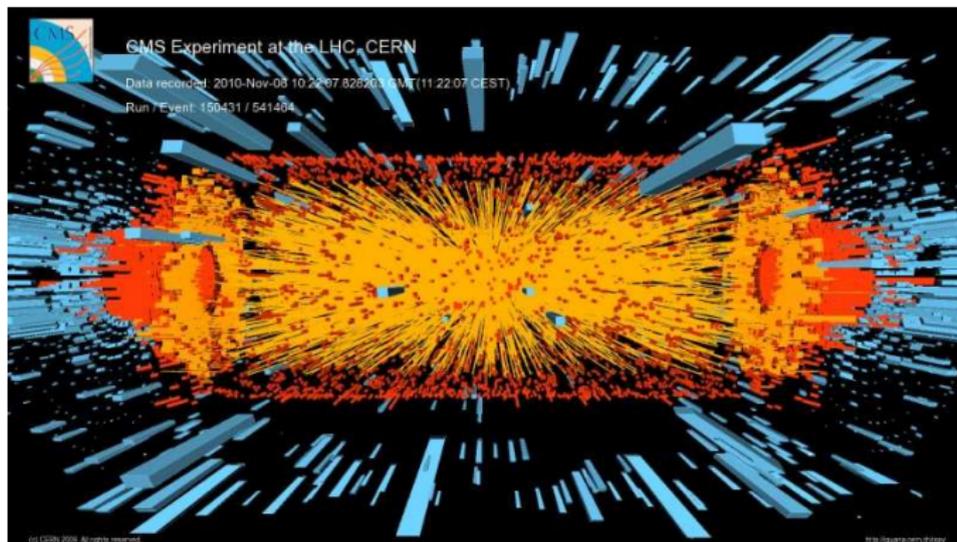
# The creation of the Quark Gluon Plasma (QGP)

- ▶ A hot and dense medium is created during the collision
  - ▶ The medium quickly thermalizes and allows a hydrodynamic description of its spacetime evolution, eventually turning into soft hadrons
  - ▶ Energetic jets are also produced abundantly in the medium



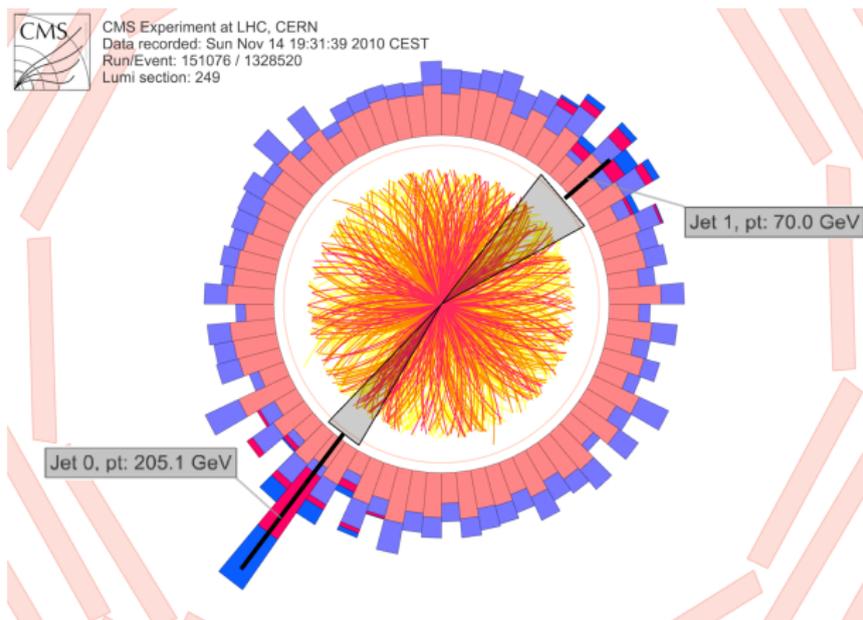
# Jet physics gets the richest in heavy ion collisions

- ▶ Thousands of particles are produced and the underlying event backgrounds are enormous



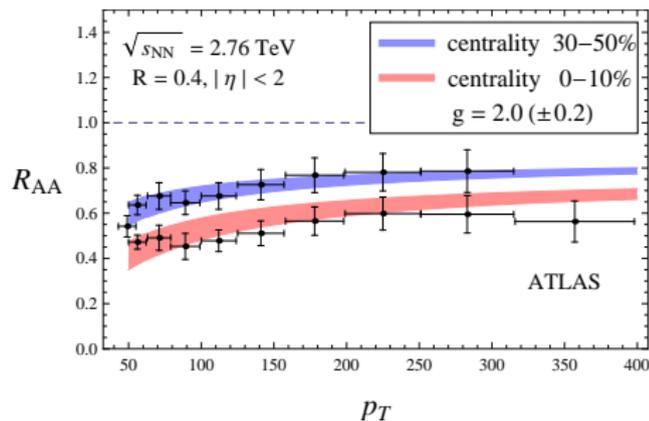
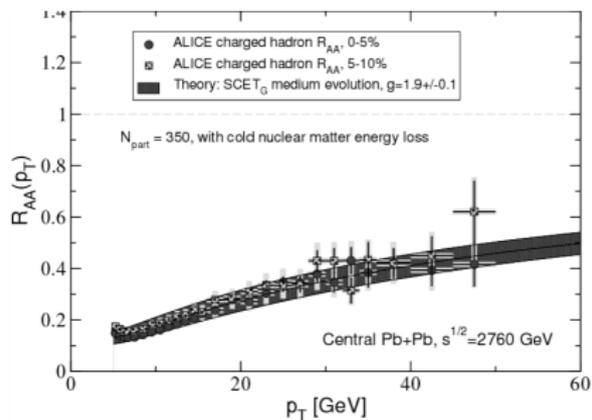
# Jets are "quenched" and modified in heavy ion collisions

- ▶ Jets are not only embedded in an enormous underlying event background but also significantly modified
- ▶ Because of the huge background, one needs to do both background subtraction and jet grooming and measure jets with small radii ( $0.2 < R < 0.4$ )
- ▶ Dramatic suppression of jets and momentum imbalance is observed

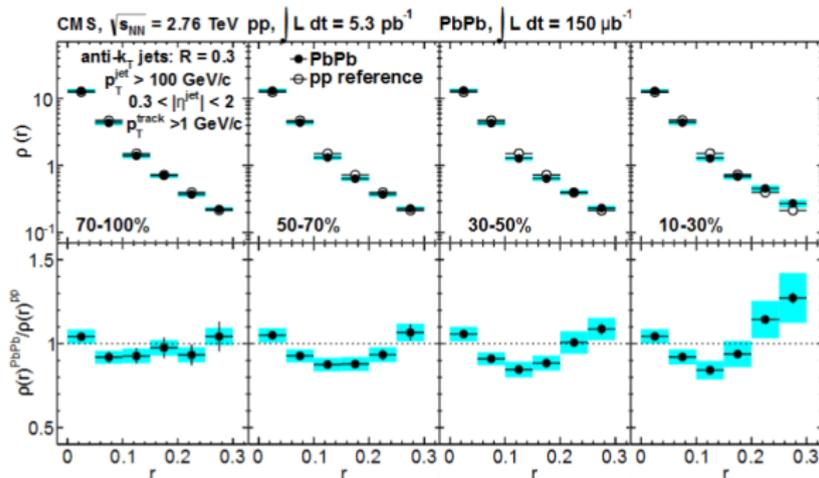
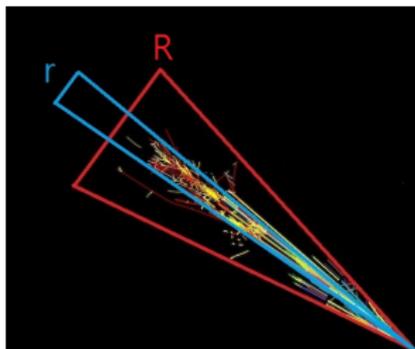


# Hadron and jet cross section suppression

- ▶  $R_{AA} < 1$  is the ratio of the cross sections in  $AA$  and  $pp$  collisions



# Jet spectroscopy of the QGP



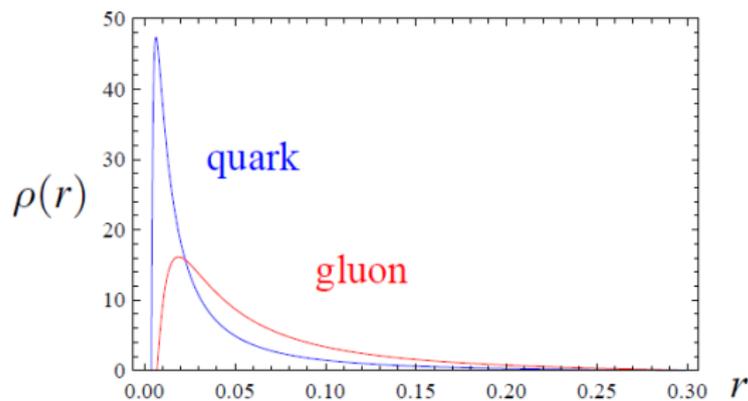
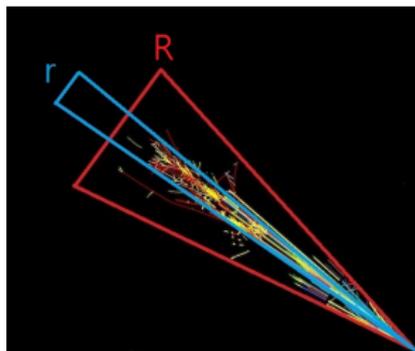
$$\Psi_J(r) = \frac{\sum_{r_i < r} E_{Ti}}{\sum_{r_i < R} E_{Ti}}$$

$$\langle \Psi \rangle = \frac{1}{N_J} \sum_J \Psi_J(r, R)$$

$$\rho(r) = \frac{d\langle \Psi \rangle}{dr}$$

- ▶ Jets have become essential tools to probe the quark-gluon plasma produced in heavy ion collisions
- ▶ One typically evaluates the observable modification by the ratio of the curves in AA and pp collisions  $\frac{\mathcal{O}^{AA}}{\mathcal{O}^{pp}}$
- ▶ With detailed understanding of jets and their structures we can relate their modifications to the medium properties: the need of precise jet substructure studies

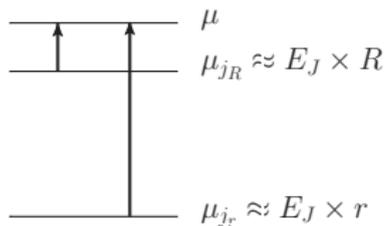
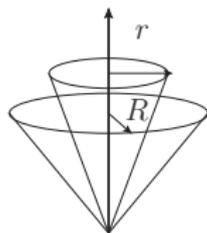
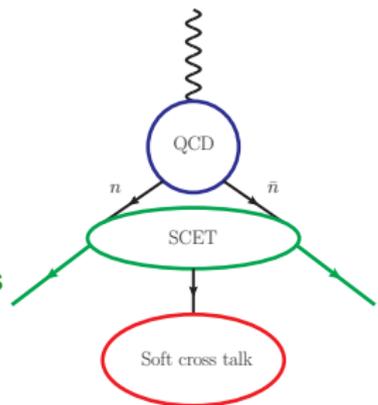
# Jet substructure calculation and resummation



- ▶ Jet shapes probe the averaged energy distribution inside a jet
- ▶ The infrared structure of QCD induces Sudakov logarithms
- ▶ Fixed order calculation breaks down at small  $r$
- ▶ Large logarithms of the form  $\alpha_s^m \log^m r/R$  ( $m \leq 2n$ ) need to be resummed
- ▶ Sensitive to the partonic origin of jets and the quark/gluon jet fraction

# Resummation using Soft-Collinear Effective Theory (SCET)

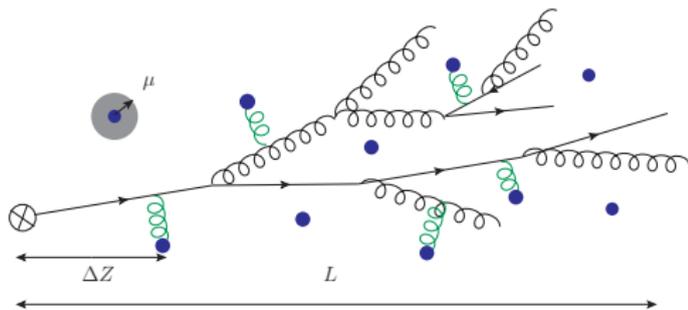
- ▶ Effective field theory techniques are most useful when there is hierarchy between characteristic energy scales
- ▶ SCET factorizes physical degrees of freedom in QCD by a systematic expansion in power counting
  - ▶ Match SCET with QCD at the hard scale by integrating out the **hard** modes
  - ▶ Integrating out the off-shell modes gives **collinear Wilson lines** which describe the collinear radiation
  - ▶ The soft sector is described by **soft Wilson lines** along the jet directions



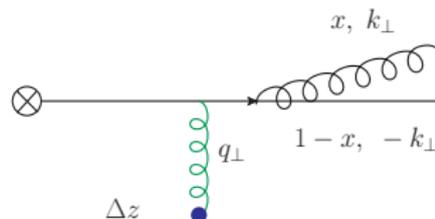
Renormalization group evolution  
between  $\mu_{j_r}$  and  $\mu_{j_R}$  resums  
 $\log \mu_{j_r} / \mu_{j_R} = \log r / R$   
(Chien et al 1405.4293)

# Multiple scattering in a medium and QCD bremsstrahlung

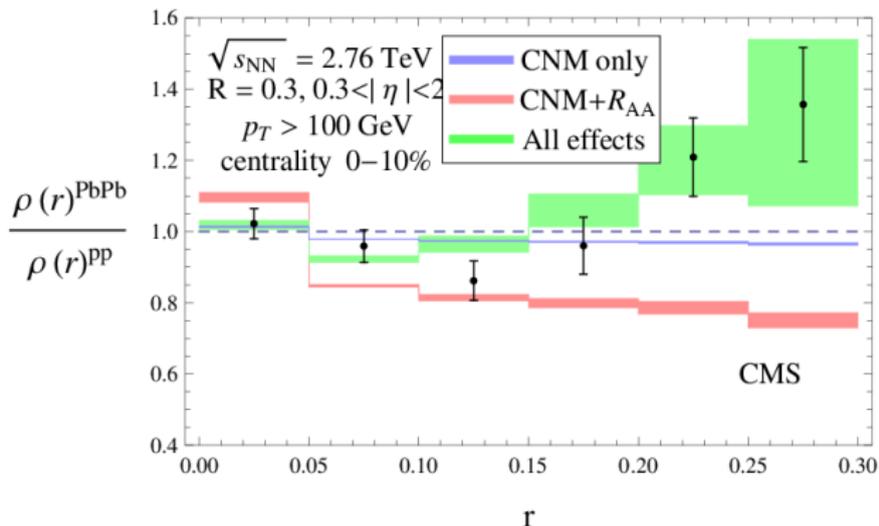
- ▶ Coherent multiple scattering and induced bremsstrahlung are the qualitatively new ingredients in the medium parton shower
- ▶ Interplay between multiple characteristic scales:
  - ▶ Debye screening scale  $\mu$
  - ▶ Parton mean free path  $\lambda$
  - ▶ Radiation formation time  $\tau$



- ▶ Jet-medium interaction using SCET with background Glauber gluon fields SCET<sub>G</sub> (Glauber-collinear: Majumder et al, Vitev et al. Glauber-soft: work in progress)
- ▶ Leading-order medium induced splitting functions  $\mathcal{P}_{i \rightarrow j l}^{med}(x, k_{\perp})$  were calculated using SCET<sub>G</sub> (Vitev et al)



# First quantitative understanding of jet shape modification



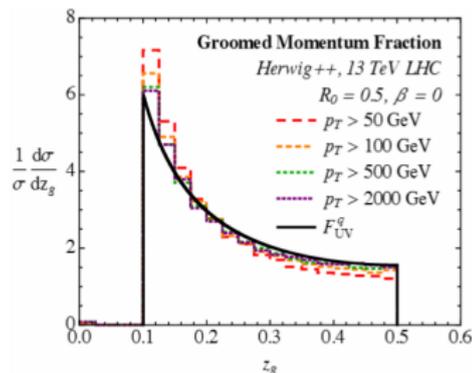
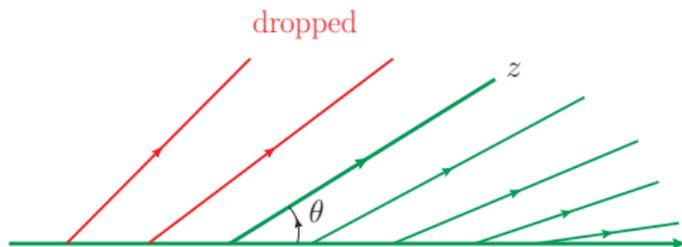
- ▶ Cold nuclear matter effect is negligible
- ▶ Jet quenching increases the quark jet fraction
- ▶ Jet-by-jet the shape is broadened
- ▶ Chien et al 1509.07257 and CMS data 1310.0878

## How do we isolate physics and distinguish jet quenching models?

- ▶ Jets are multi-scaled objects with rich information about the physics across the entire energy spectrum
- ▶ Jet observables have different sensitivities to physics at different energy scales
- ▶ Through a series of jet measurements we can map out the whole jet formation history
- ▶ Whether the model relies on the low scale physics corresponds to two rough pictures of jet quenching
  - ▶ Yes. Parton showers are not affected much until the later stages. The medium depletes the partons out of the jet
  - ▶ No. The medium effects open up more channels in the jet formation process, all the way from the hard process through hadronization
- ▶ Can we test the two pictures and the role of medium response?
  - ▶ We are able to dissect radiations and pick out the components of interest
  - ▶ The idea: come up with an observable as insensitive to low scale physics as possible
  - ▶ The tool: jet grooming

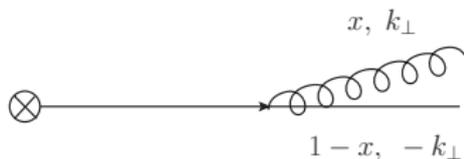
# Jet grooming is actually artificial jet quenching

- ▶ It is a controlled way to remove soft radiation
- ▶ How does a jet quenching model confront with jet grooming?
  - ▶ Do they add up or interfere?

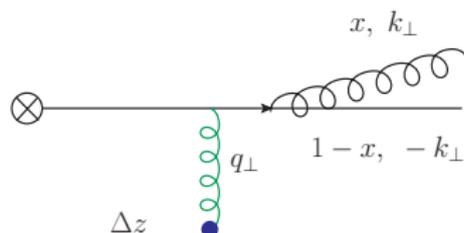
Groomed momentum fraction  $z_g$ 

- ▶ Soft Drop: a tree-based procedure to drop soft radiation (Larkoski et al 1402.2657)
  - ▶ Recluster a jet using  $C/A$  algorithm: angular ordered
  - ▶ For each branching, consider the  $p_T$  of each branch and the angle  $\theta$
  - ▶ Drop the soft branch if  $z < z_{cut} \theta^\beta$ , where  $z = \frac{\min(p_{T,1}, p_{T,2})}{p_{T,1} + p_{T,2}}$
  - ▶ CMS used  $\beta = 0$ ,  $z_{cut} = 0.1$ ,  $R = 0.4$ ,  $\Delta R_{12} > \Delta = 0.1$  and measured  $z_g$
- ▶  $z_g$ : the momentum fraction of the soft branch.  $r_g$ : the angle between the branches

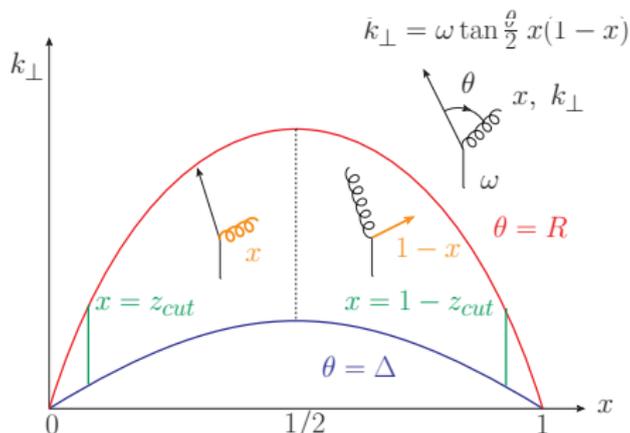
## $z_g$ and splitting functions



$$P(x, k_{\perp}) \propto \frac{1}{x k_{\perp}}$$



- ▶ In vacuum, the soft branch kinematics is closely related to the Altarelli-Parisi splitting function
- ▶ In the medium, the bremsstrahlung component modifies the soft branch kinematics

Analysis of  $z_g$  (Chien et al 1608.07283)

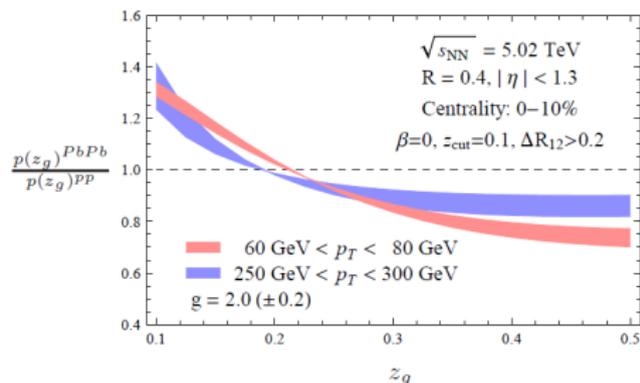
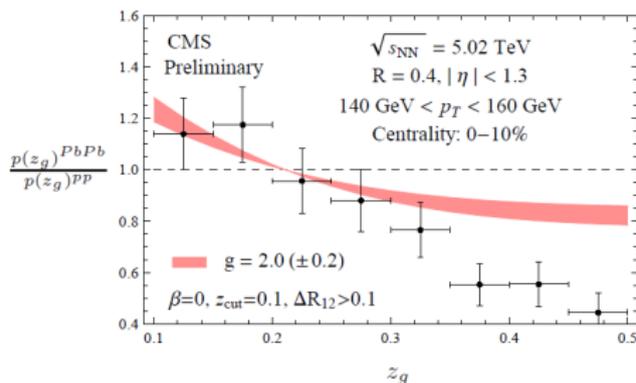
- ▶ The partonic phase space is constrained by  $R$  (jet algorithm),  $\Delta$  (jet selection) and  $z_{cut}$  (jet grooming)
- ▶ At leading order, the  $1 \rightarrow 2$  branching probability directly affects the subjet distribution

$$\mathcal{P}_{i \rightarrow jl}(x, k_{\perp}) = \mathcal{P}_{i \rightarrow jl}^{vac}(x, k_{\perp}) + \mathcal{P}_{i \rightarrow jl}^{med}(x, k_{\perp})$$

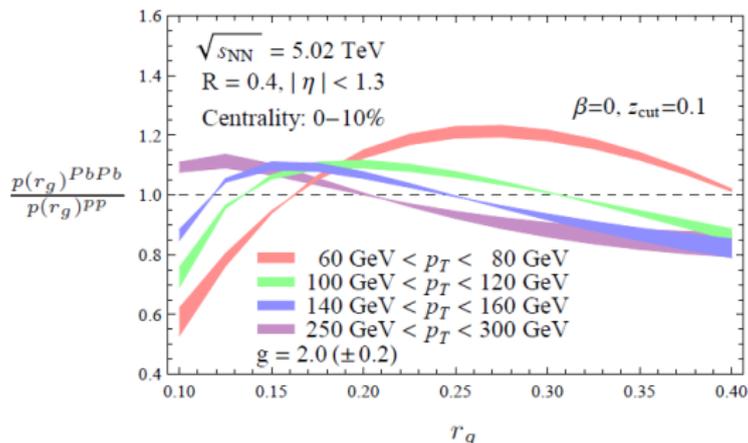
- ▶ The distributions of  $z_g$  and  $r_g$  are calculated ( $\overline{\mathcal{P}}(x) = \mathcal{P}(x) + \mathcal{P}(1-x)$ )

$$p_i(z_g) = \frac{\int_{k_{\Delta}}^{k_R} dk_{\perp} \overline{\mathcal{P}}_i(z_g, k_{\perp})}{\int_{z_{cut}}^{1/2} dx \int_{k_{\Delta}}^{k_R} dk_{\perp} \overline{\mathcal{P}}_i(x, k_{\perp})}, \quad p_i(r_g) = \frac{\int_{z_{cut}}^{1/2} dx p_T x(1-x) \overline{\mathcal{P}}_i(x, k_{\perp}(r_g, x))}{\int_{z_{cut}}^{1/2} dx \int_{k_{\Delta}}^{k_R} dk_{\perp} \overline{\mathcal{P}}_i(x, k_{\perp})}$$

# Theory calculation of $z_g$



- ▶ The medium enhances the soft branches, and the effect becomes smaller for higher  $p_T$  jets
- ▶ Qualitatively expected and quantitatively surprising (CMS is updating the data (CMS-HIN-16-006) soon. Stay tuned!)
- ▶ Cutting on the angle between branches selects a special subset of the jet sample
  - ▶ Jets with a two prong structure not typical for QCD jets
  - ▶ The scale of this subset branching is high: hard jet substructure

Theory prediction for  $r_g$ 

- ▶ The subjet angular distribution will reveal the nature of QCD bremsstrahlung
- ▶ It will be a direct probe of the medium scale
- ▶ The next step is to measure the groomed jet mass (CMS measurement coming out. Ungroomed charged jets done by ALICE)

# Power counting of modes for groomed jet mass (work in progress)

- ▶ Factorization in SCET (Larkoski et al 1603.09338 "in jets"):

- ▶ **In-jet soft** mode

$$p_s = E_J z_{cut}(1, R^2, R), \text{ with } \mu_s = E_J R z_{cut}$$

- ▶ **Collinear** mode

$$p_c = (E_J, \frac{m^2}{E_J}, m), \text{ with } \mu_j = m$$

- ▶ **Soft-collinear** mode respecting the measurement  $x\theta^2 \sim m^2/E_J^2$  and jet grooming  
 $z_{cut} \sim x(\theta/R)^{-\beta}$

$$p_{sc} = (E_J z_{cut} \left( \frac{m}{E_J R \sqrt{z_{cut}}} \right)^{\frac{2\beta}{2+\beta}}, \frac{m^2}{E_J}, m \sqrt{z_{cut}} \left( \frac{m}{E_J R \sqrt{z_{cut}}} \right)^{\frac{\beta}{2+\beta}}), \text{ with } \mu_{sc} = m \sqrt{z_{cut}} \left( \frac{m}{E_J R \sqrt{z_{cut}}} \right)^{\frac{\beta}{2+\beta}}$$

- ▶ **Hard collinear** mode from pure jet reconstruction

$$p_{jR} = E_J(1, R^2, R), \text{ with } \mu_{jR} = E_J R$$

## Groomed jet mass function

- ▶ The process-independent groomed jet mass function  $J_M^{i\cancel{f}}(m^2, \mu)$  captures all the soft-collinear radiation inside jets ( $i = q, g$ )

$$J_M^{i\cancel{f}}(m^2, \mu) = \int dp^2 dk J_i(p^2, \mu) S_i^{i\cancel{f}}(k, R, z_{cut}, \mu) \delta(m^2 - p^2 - 2EJk)$$

where  $S_i^{i\cancel{f}}(k, R, z_{cut}, \mu) = S_i^C(k, R, z_{cut}, \mu) S_i^{IN}(R, z_{cut}, \mu)$

- ▶ Medium-induced splitting functions are used to calculate the modification of  $J_M^{i\cancel{f}}(m^2, \mu)$ . At  $\mathcal{O}(\alpha_s)$ ,

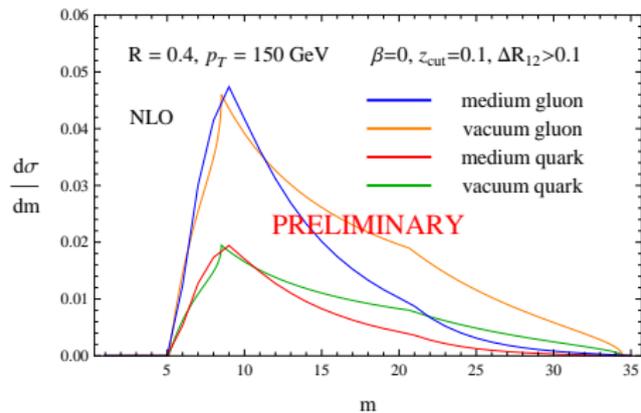
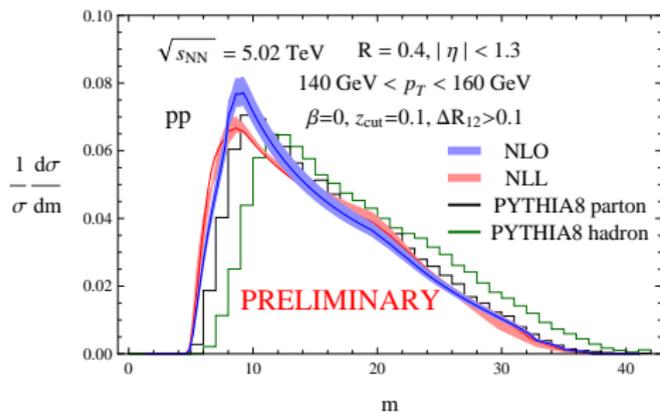
$$J_M^{i\cancel{f}}(m^2, \mu) = \sum_{j,k} \int_{PS} dx dk_{\perp} \mathcal{P}_{i \rightarrow jk}(x, k_{\perp}) \delta(m^2 - M^2(x, k_{\perp})) \Theta_{\text{alg.}} \Theta_{\cancel{f}}$$

$$M^2(x, k_{\perp}) = \frac{k_{\perp}^2}{x(1-x)}, \Theta_{\text{kT}} = \Theta(EJR x(1-x) - k_{\perp}), \Theta_{\cancel{f}} = \Theta(EJR x(1-x) \left(\frac{x}{z_{cut}}\right)^{1/\beta} - k_{\perp}).$$

- ▶ The full jet mass distribution can be calculated by weighing the groomed jet mass functions with jet cross sections

$$\frac{d\sigma}{dm^2} = \sum_{i=q,g} \int_{PS} dp_T dy \frac{d\sigma^i}{dp_T dy} P_i^{i\cancel{f}}(m^2, \mu), \text{ where } P_i^{i\cancel{f}}(m^2, \mu) = \frac{J_M^{i\cancel{f}}(m^2, \mu)}{J_{un}^i(\mu)}$$

# Preliminary results



- ▶ The  $\Delta R_{12} > 0.1$  cut cuts out the Sudakov peak
- ▶ The lower and upper limits of jet mass are essentially dictated by kinematics.  $r_g$  and jet mass are highly correlated
- ▶ The medium perturbative contributions enhance small mass region
- ▶ Soft contributions and hadronization are still under examination

# Conclusion

- ▶ What we have learned: flavor dependence of jet quenching and the role of quark/gluon jet fraction in jet substructures
- ▶ Subjet distribution provides an opportunity to test the modification of hard splitting within jets
- ▶ Groomed jet mass is resummed with small radius, and the medium enhances the small mass region (preliminary)
- ▶ Design specific grooming procedure along the physics goal of jet quenching studies
- ▶ Stay tuned for the jet mass calculation with different soft drop parameters ( $\beta = 1.5$ ,  $z_{cut} = 0.5$ )!