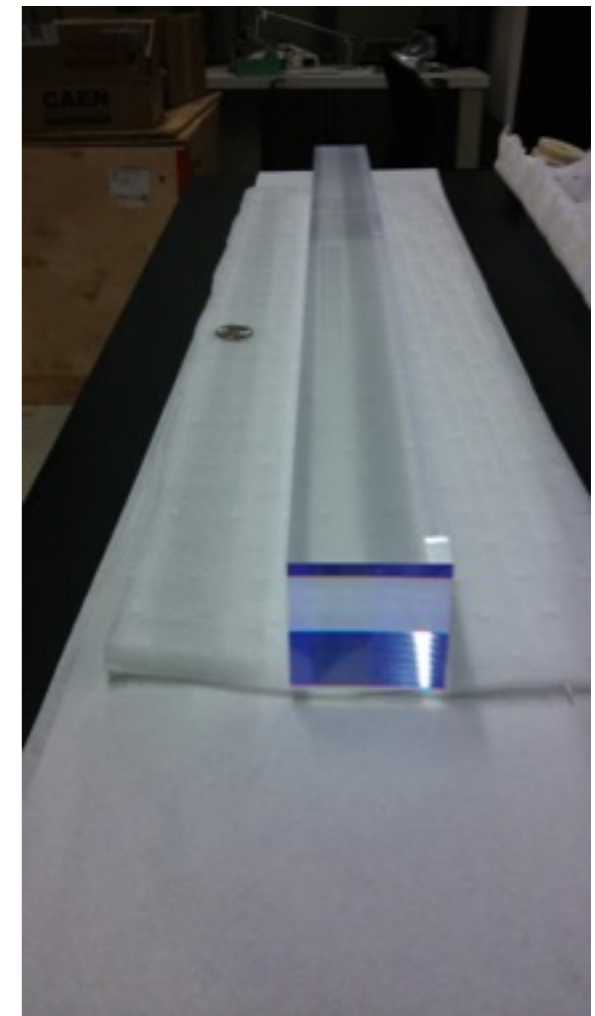




# OSU Activities

- OSU has several prototype devices of interest to MilliQan
  - Four 2"x2"x80cm bars of Saint-Gobain BC408 plastic scintillator
  - A fifth scintillator bar incorporated with a Hamamatsu R7725 PMT
    - Assembled by Saint-Gobain



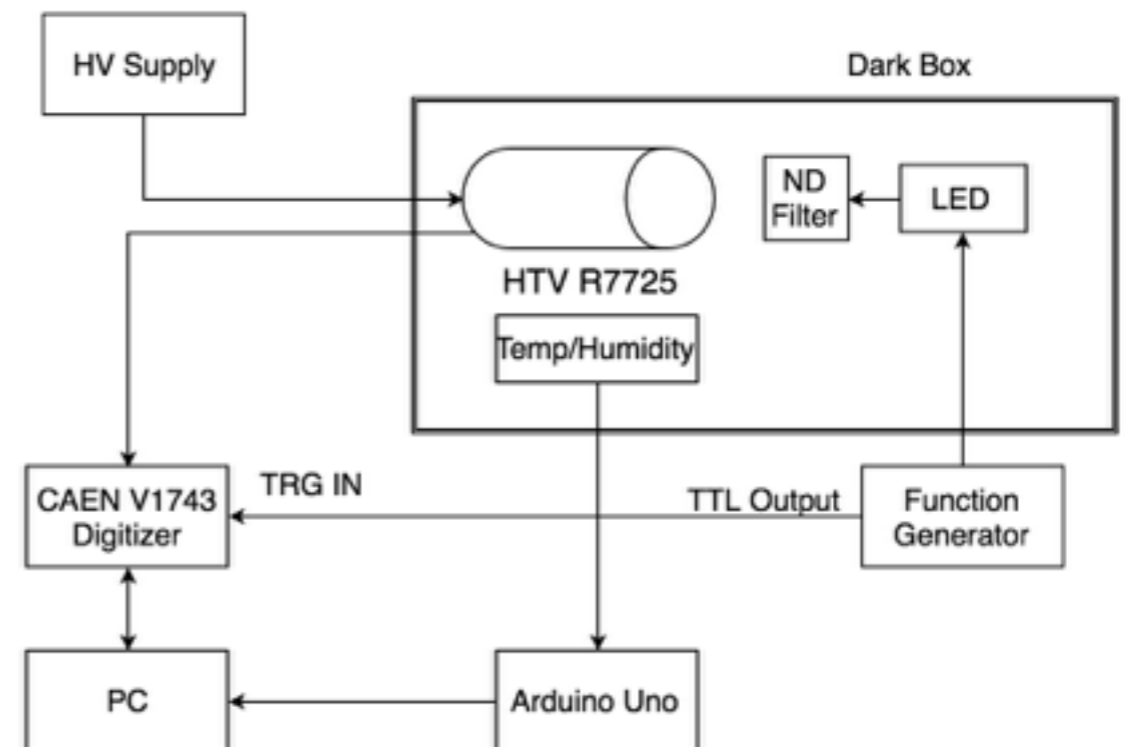
- In addition to the R7725 PMT coupled to scintillator from Saint-Gobain:
  - Stand-alone R7725 PMT from Hamamatsu
  - Hamamatsu E5859-11 voltage divider base



# OSU Facilities



- We read out each device with a CAEN V1743 digitizer
- Connected to a PC by fiber optic connection to a CAEN A2818 optical controller

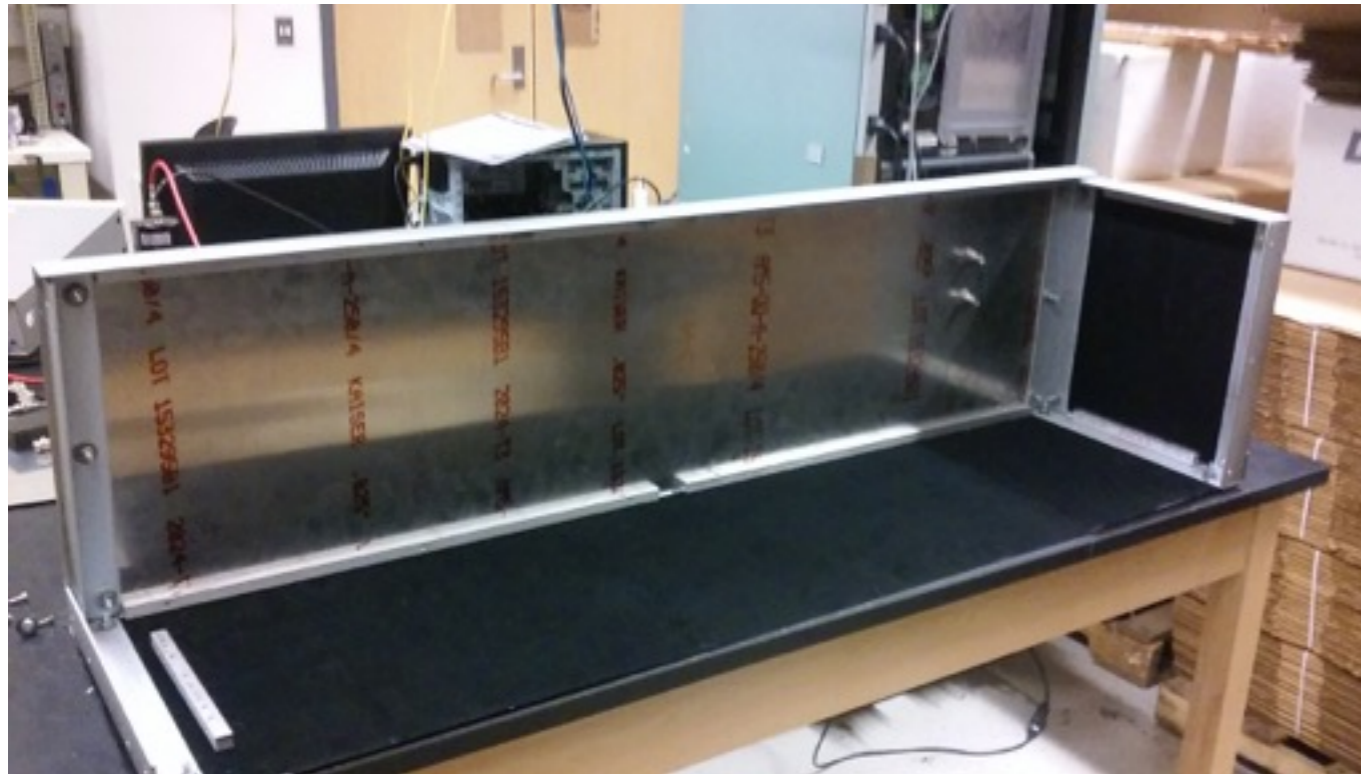


# R7725 Dark Rate Measurements

# Light Shielding



- Over several months we have improved our HTV R7725 PMT
- Working to prevent ambient light from disturbing dark rate measurement



- Initially a metal tube with small amounts of electrical tape, within a larger metal box
  - Most of the box internally is covered with “black-out” paper

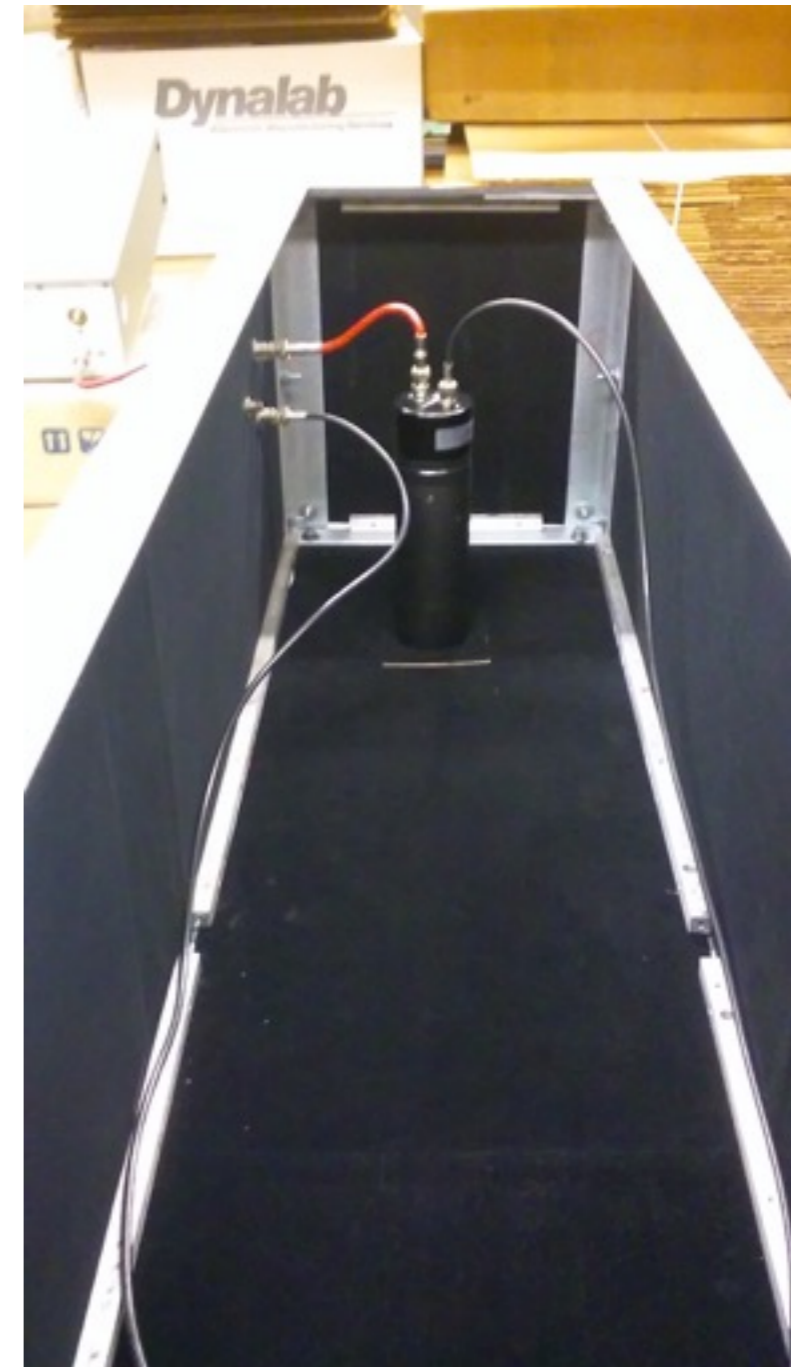
# Light Shielding



- Box designed to be large enough to accommodate PMT and scintillator



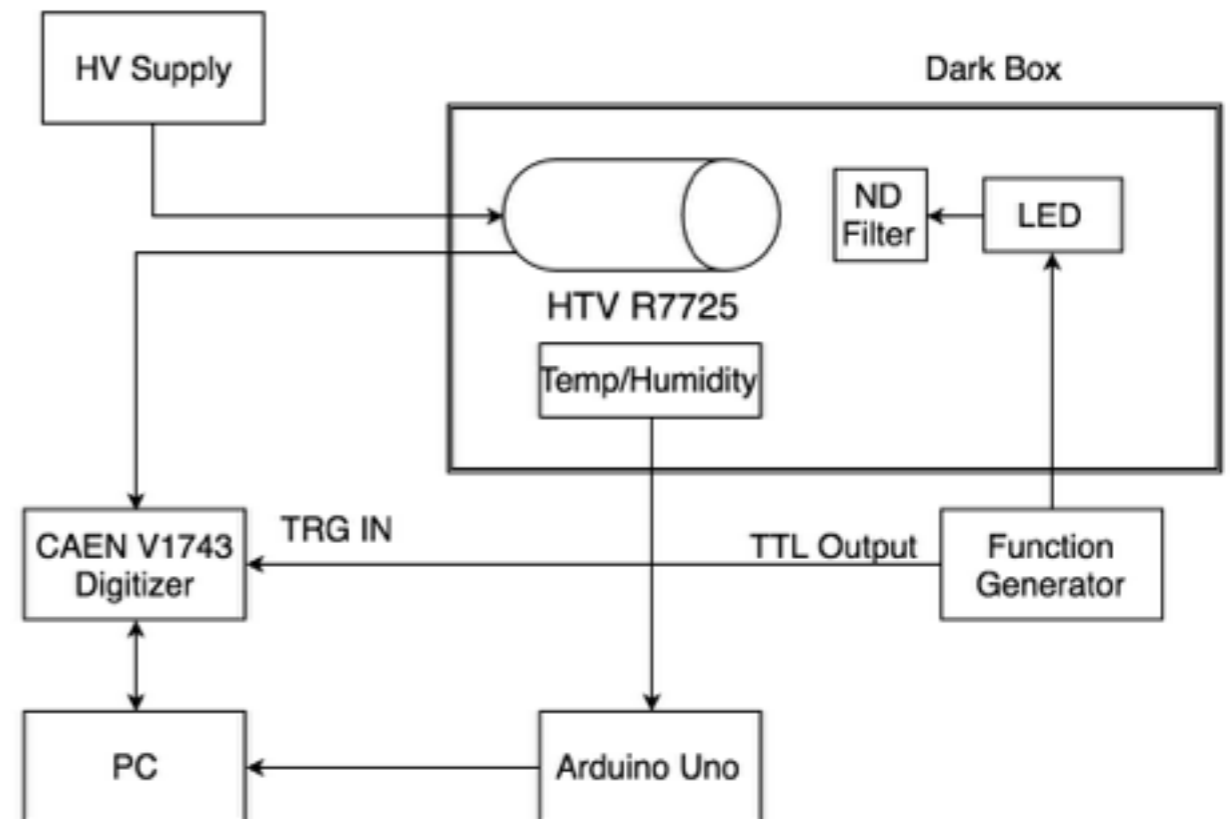
- Studio photography type “black-out” cloth



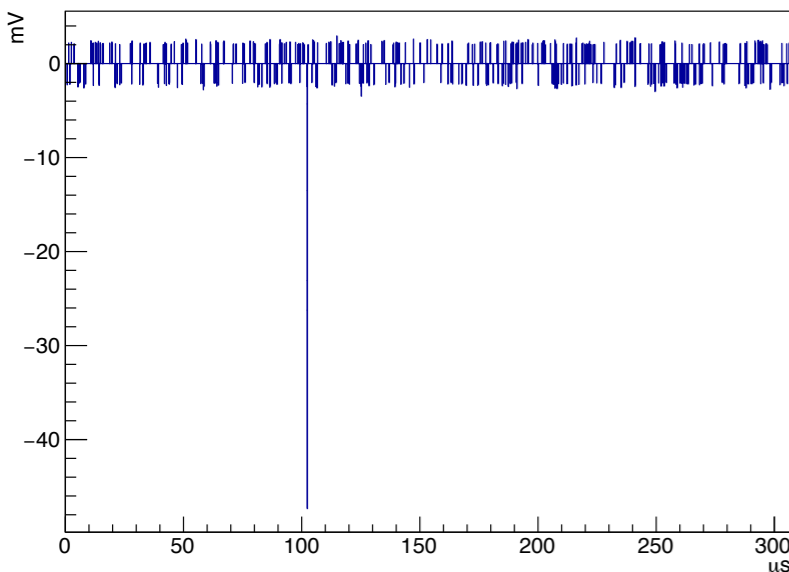
# Dark Rate Measurement



- To observe the dark rate:
  - Trigger digitizer on TTL output of function generator (1500 Hz)
  - Configure digitizer for longest acquisition possible
    - 1024 samples @ 0.4 GS/s — 2.56  $\mu$ s/event
  - Observe the amplitude of each event as the minimum sample in mV
  - Then the dark rate, above some threshold, is:



1750 V



$|\text{sample}| > 2 \text{ mV}$  for 120 events

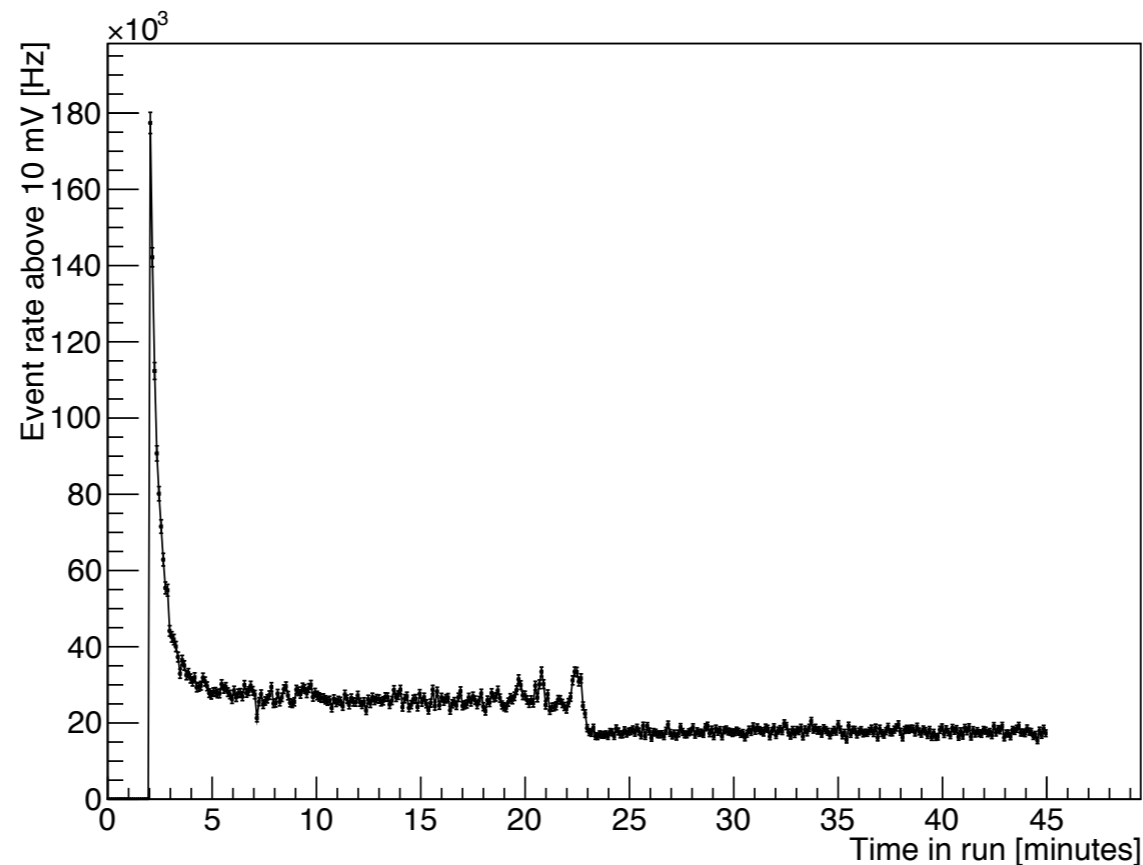
Events (amplitude  $>$  threshold) over the total acquisition time  
where total acquisition time = N events times 2.56  $\mu$ s / event



# Dark Rate Measurement



- Initial rate measurements were very high (20-30 kHz) and very unstable

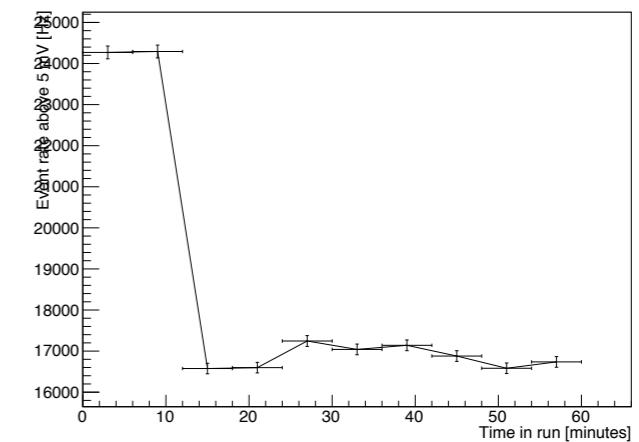
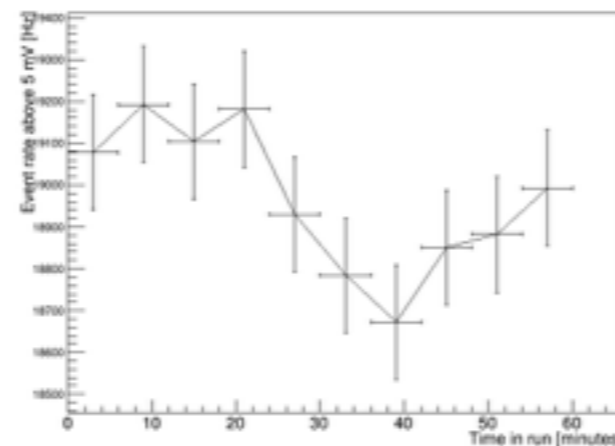
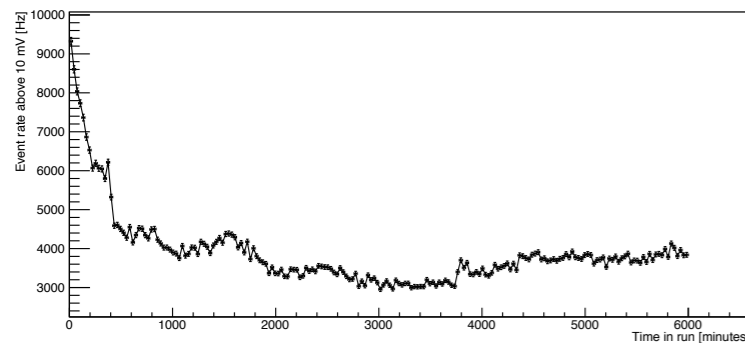
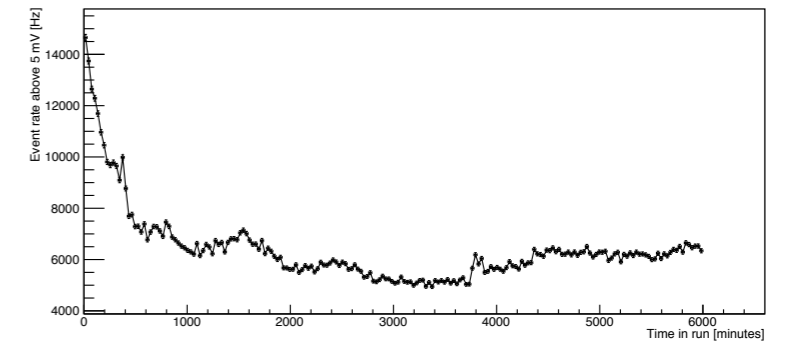
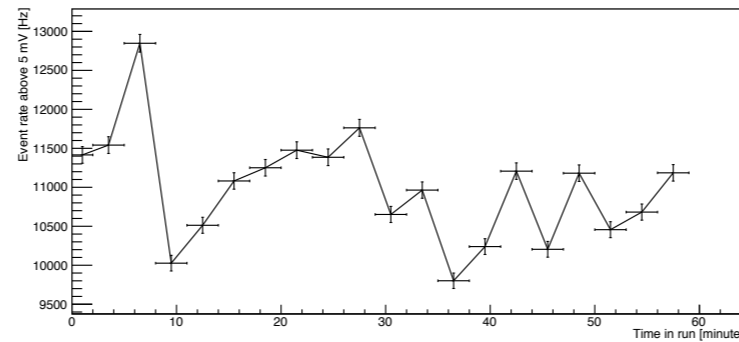
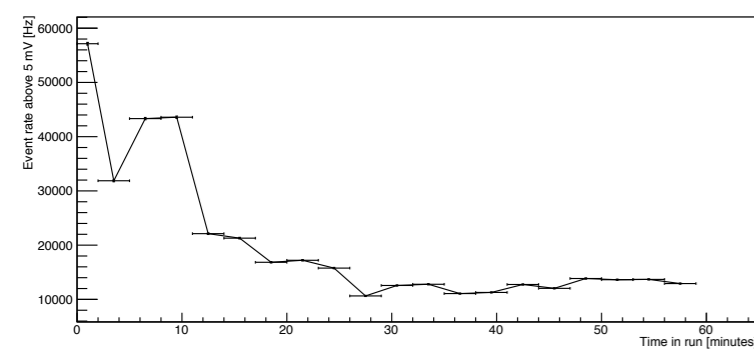


- Sharp increase at the beginning is when the tube bias voltage is turned on
- We spent considerable time examining how repeatable this measurement was, and investigating the cause of the large drops in rate

# Dark Rate Measurement



- Many repeated measurements by several different users:



- After observing different rates for different users, and even rate changes within data-taking runs, we convinced ourselves that environmental factors were not being controlled
  - In particular we suspected a very large light leak

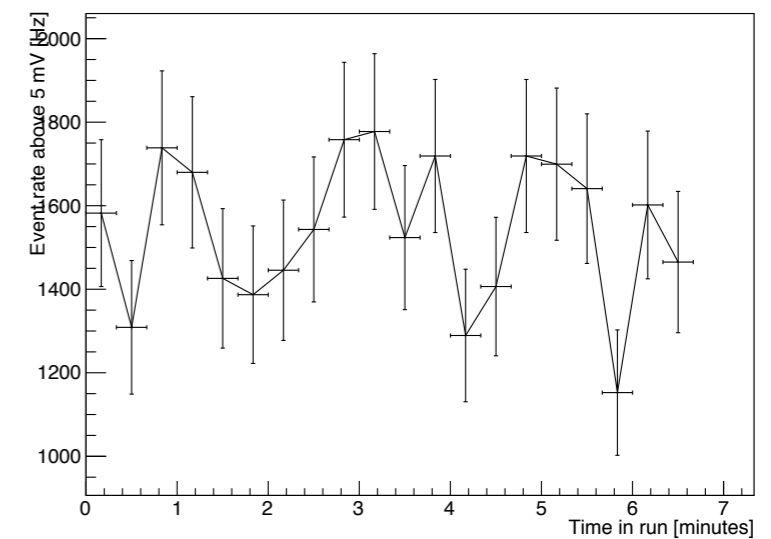
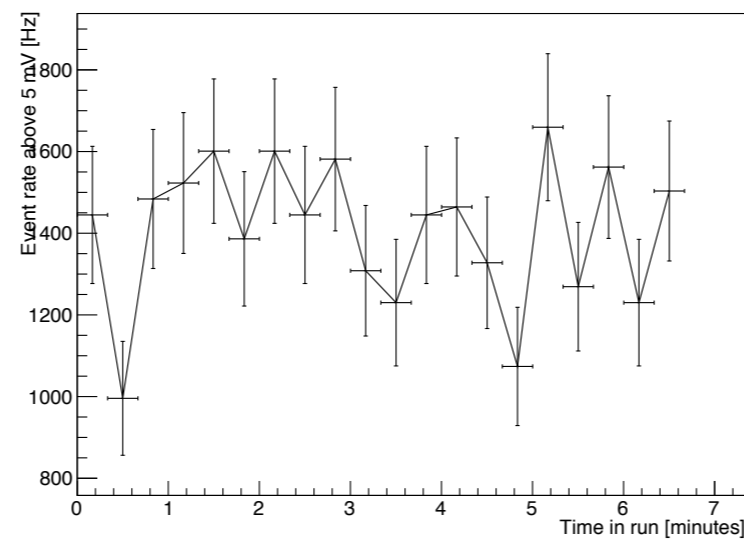
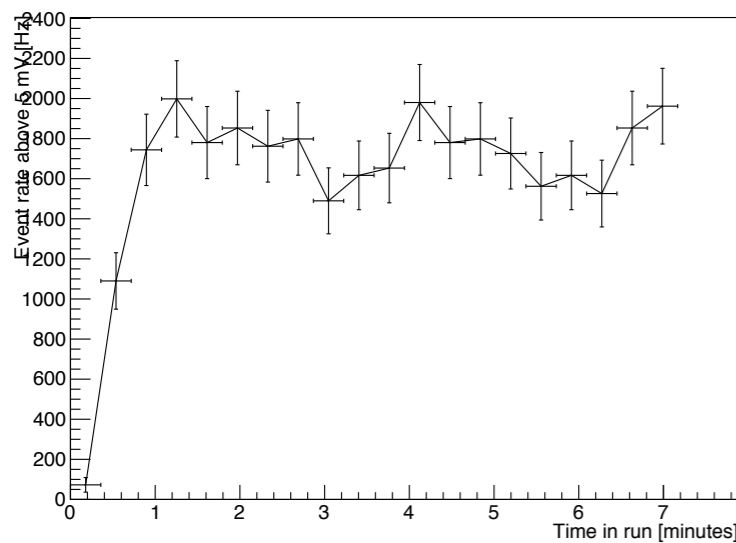
# Dark Rate Measurement



- Applied a very aggressive amount of black electrical tape
- Not visible here: cathode window itself is covered with black-out paper and directly taped



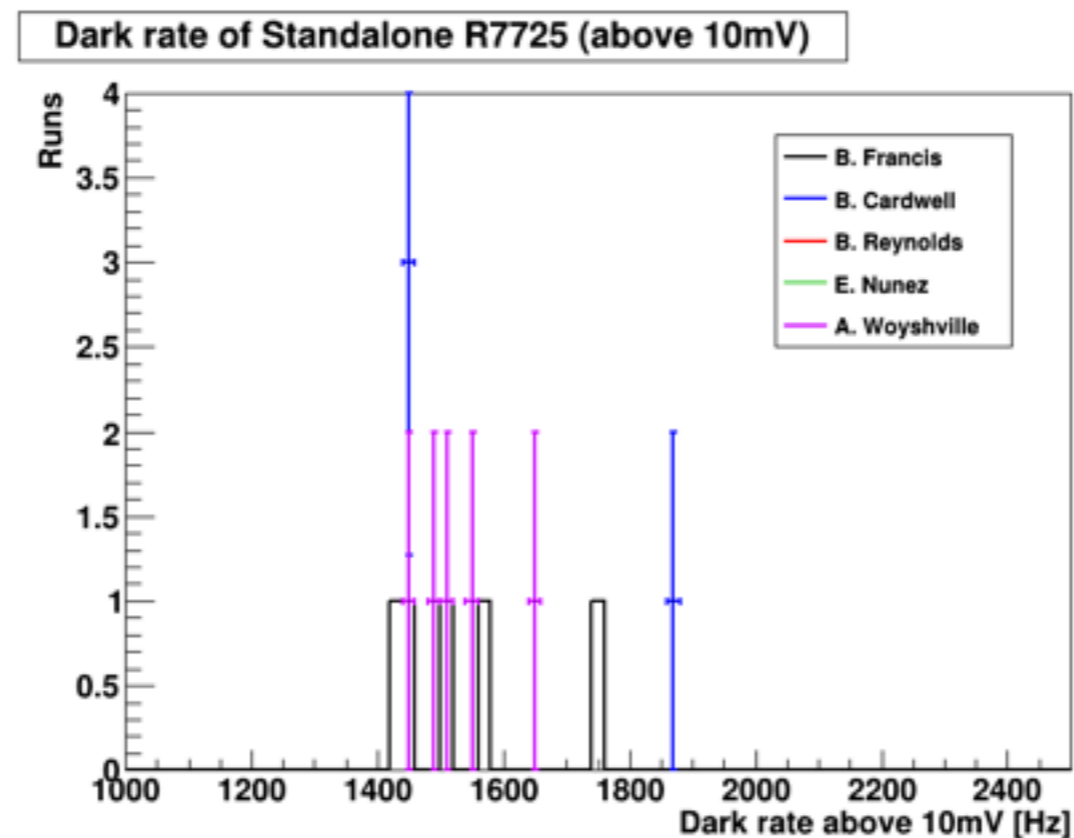
- A large decrease in rate resulted, confirming the existence of large light leaks



# Dark Rate Measurement



- Here, several separate measurements recorded over several days, each for 30 minutes of real time (6.2 seconds total acquisition live time)



# New PMT Enclosure



- Seeing that light leaks are easily missed, the next challenge was in coupling an LED and optical fiber without re-introducing large light leaks
- Constructed a new enclosure aiming to better control for environmental factors
  - Discrete neutral density filters (rather than a continuously variable one)
  - Tightly-fitting lens tube and LED mount from Thorlabs
    - No need for fiber coupling, LED directly faces PMT window
  - Temperature and humidity measurements
  - Tightly-fitting aluminum enclosure that can accept future cooling elements



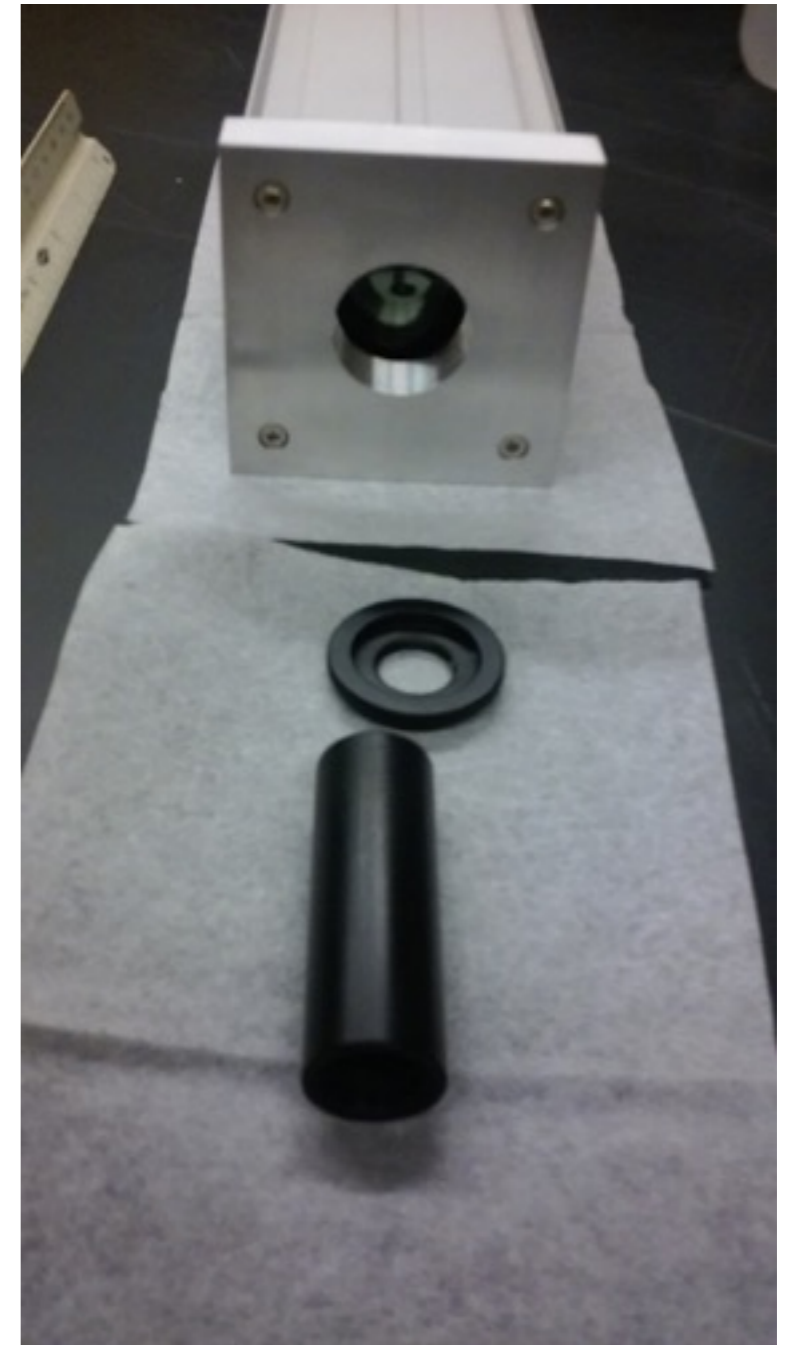
# New PMT Enclosure



- Thorlabs LED and filters:
  - LED430L — 430nm, 8 mW



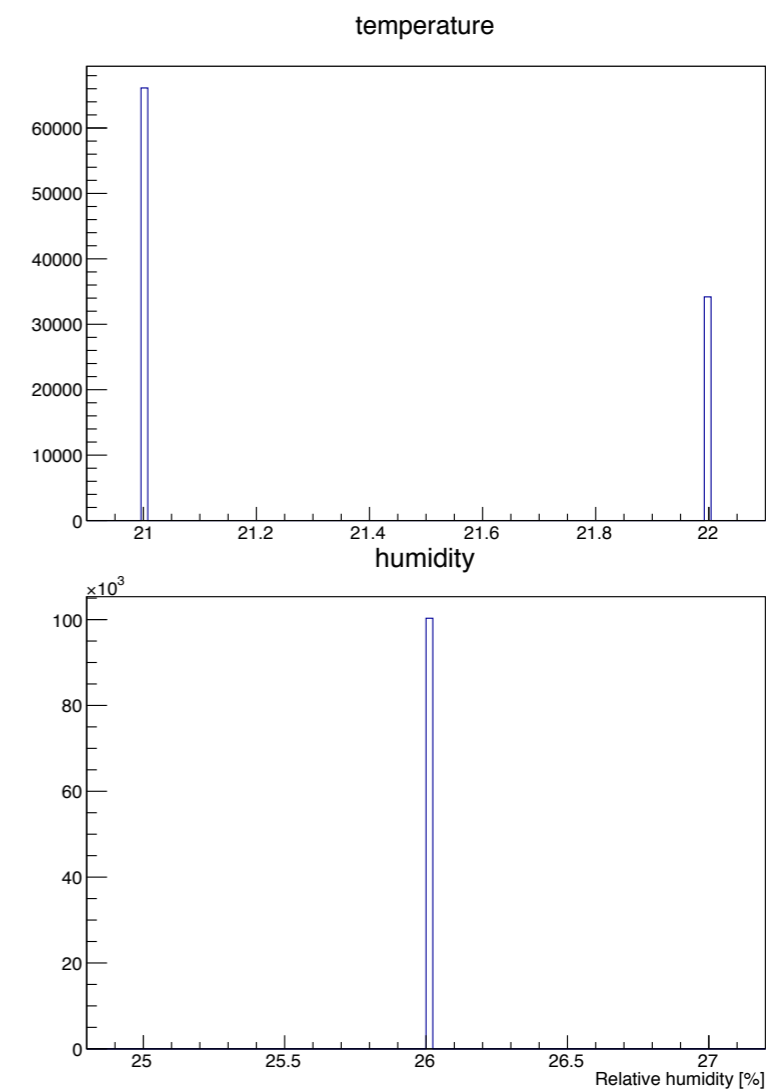
- NE500A series neutral density filters:
  - NE510A ( $T=10^{-1}$ ), NE530A ( $T=10^{-3}$ ), NE550A ( $T=10^{-5}$ )
  - Note: “OD5”  $\rightarrow$  optical density 5  $\rightarrow T=10^{-5}$
  - I will make heavy use of this notation!
    - “OD4” is a combination of OD3 and OD1



# New PMT Enclosure



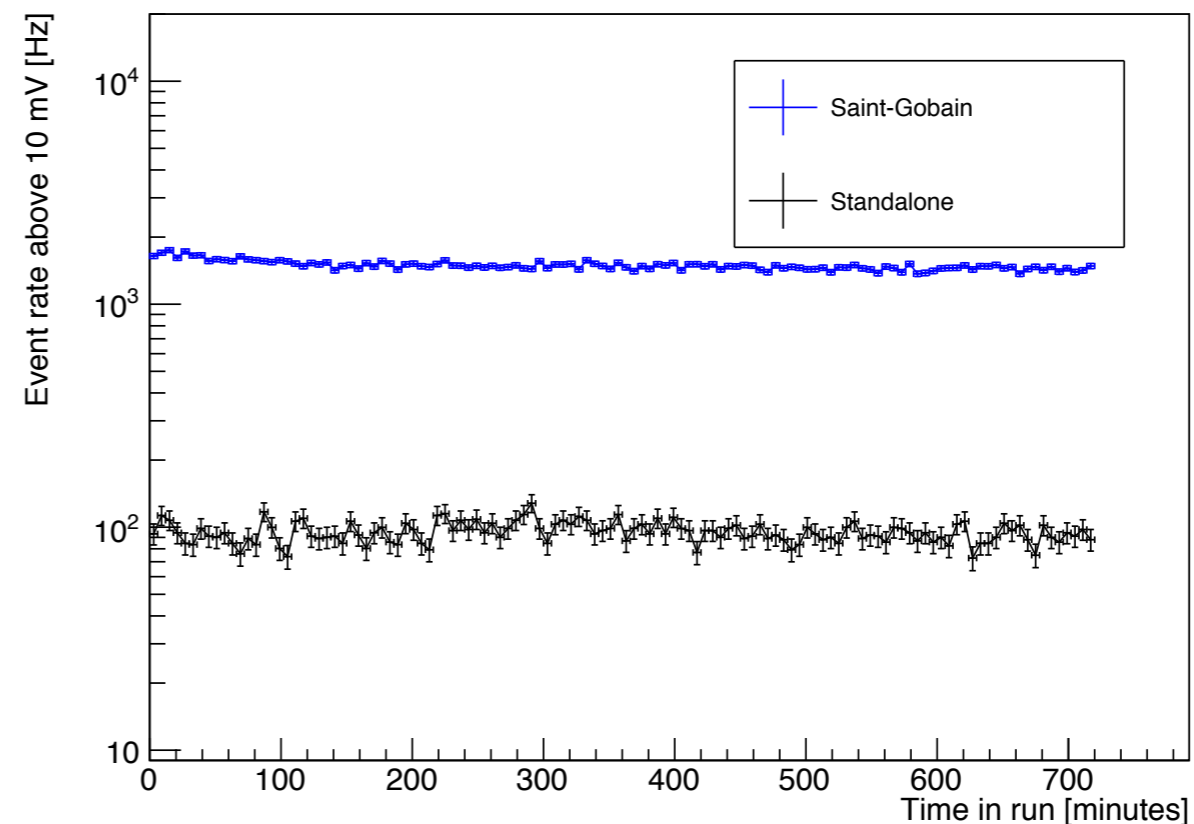
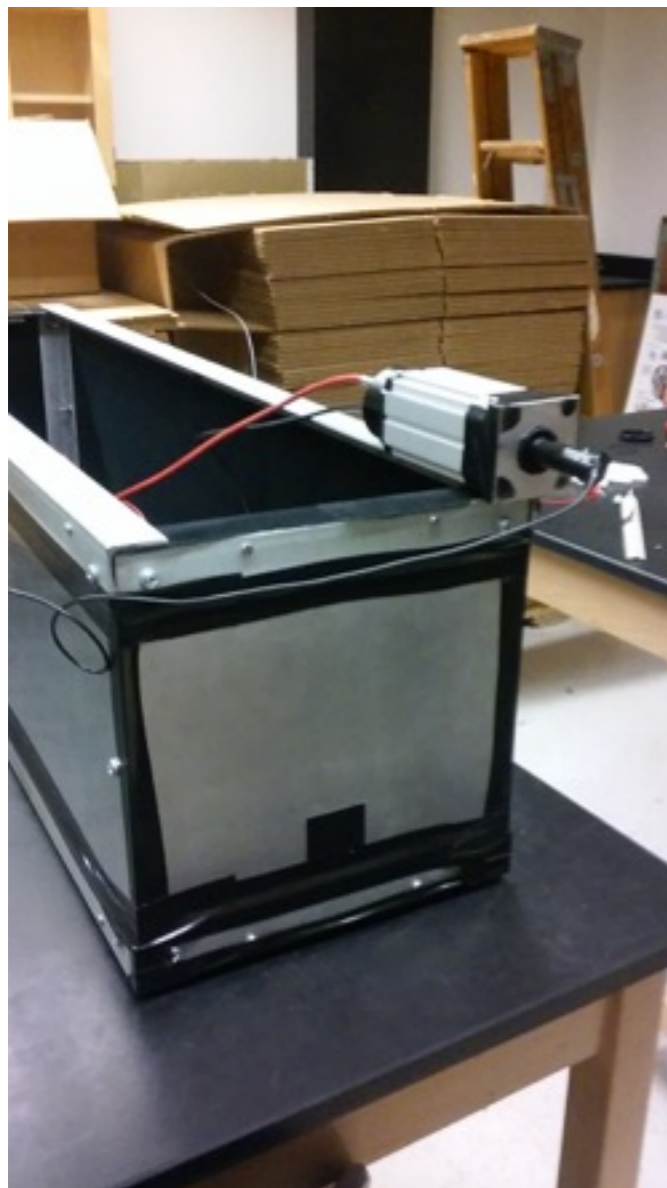
- A simple DHT11 sensor provides temperature ( $\pm 1^\circ \text{C}$ ) and relative humidity ( $\pm 1\%$ ) measurements from inside the larger box
- Read out with an Arduino Uno that is integrated into the DAQ software and data stream



~60 seconds

# New PMT Enclosure

- With this new enclosure intended to be better light-proofed, we first measured the dark rate
- The new enclosure was intentionally left outside of the larger dark box to test for light-tightness



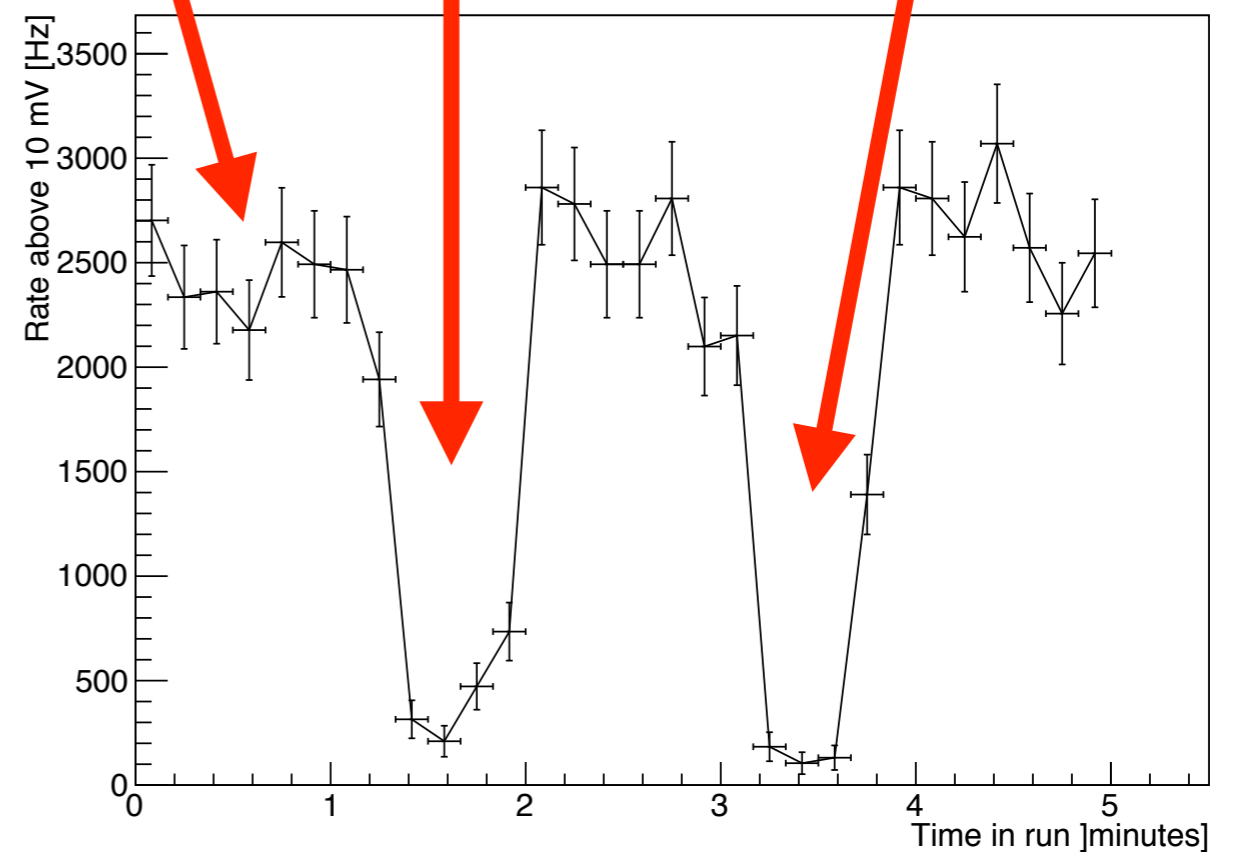
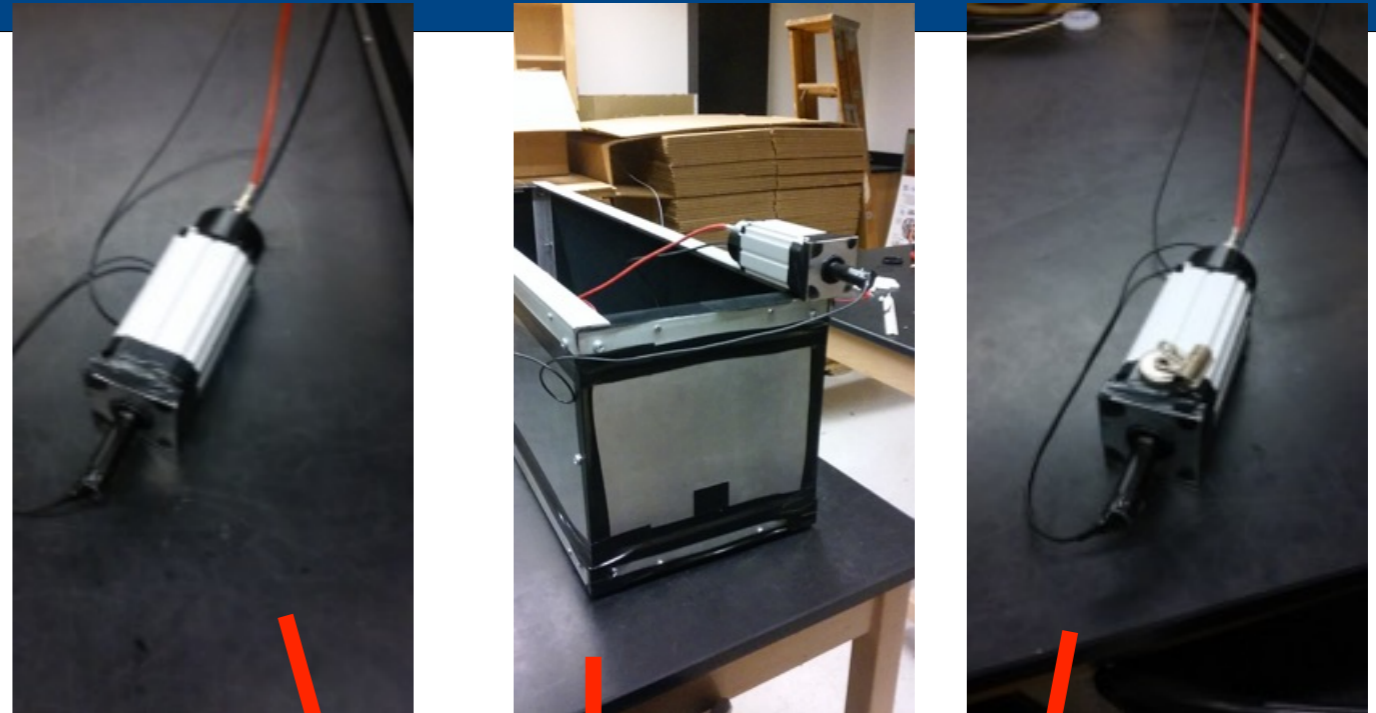
- A very low rate (100 Hz)!
- We very quickly observed that simply moving the enclosure to a different location increased this



# New PMT Enclosure



- Happened to have placed photocathode above the edge of a steel component of the dark box
- Placing a magnet near the photocathode similarly reduces this
- In both cases, applying LED light does not give observable signals
- Not surprising that a magnetic field would reduce the gain and dark rate
  - Is however surprising that this material would cause a similar effect
- Need to acquire a gauss meter and measure the field in this particular areas
- For the time being we avoid placing the enclosure on large metal objects

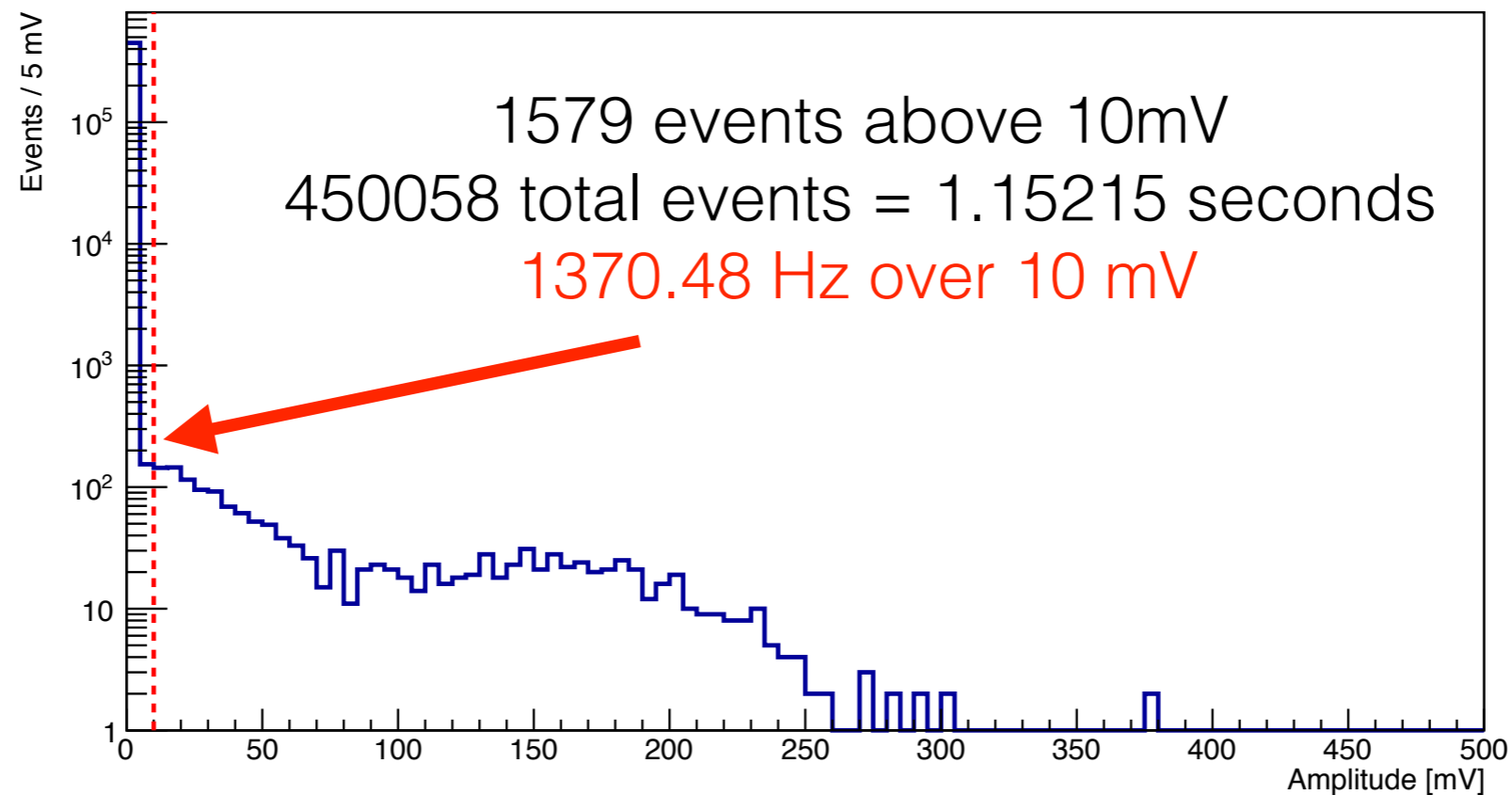


# Dark Rate Measurement



- With this new enclosure (avoiding known areas that drastically reduce the gain of the PMT), the observed dark rates are slightly lower than with the previous enclosure (very taped-up)
- An example at the recommended bias voltage, -1750 V:

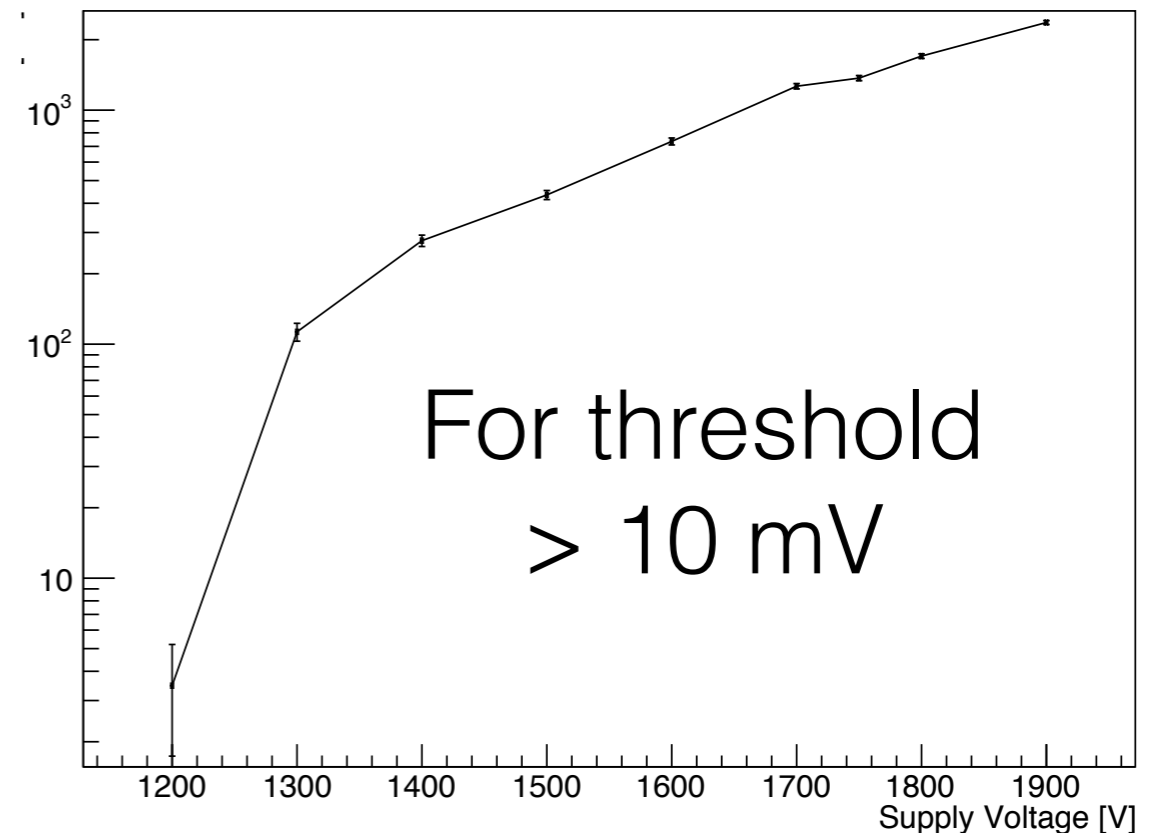
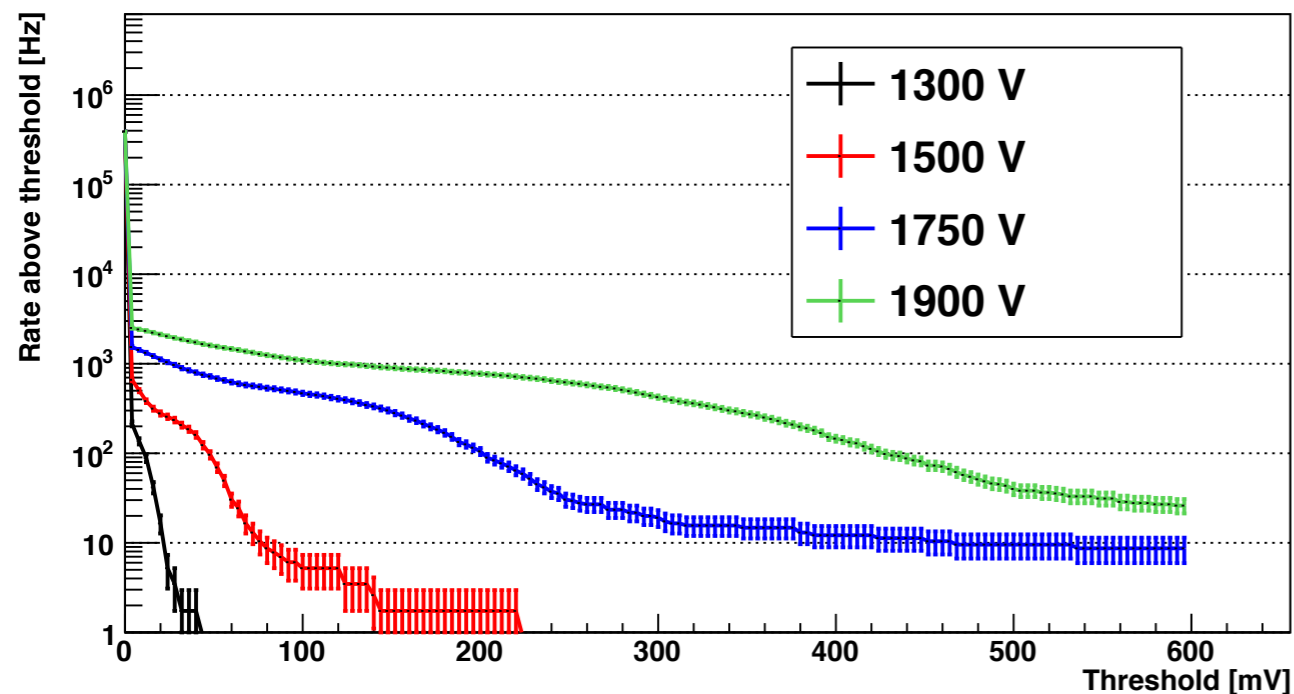
1750 V



# Dark Rate Measurement



- Performing this measurement over a range of PMT bias voltages:



- Ultimately, a rate of 1 kHz should be possible given appropriate choice in trigger threshold
- The LOI quotes an expected rate of 500 Hz per channel, which should be achievable with some cooling

# Dark Rate Measurement



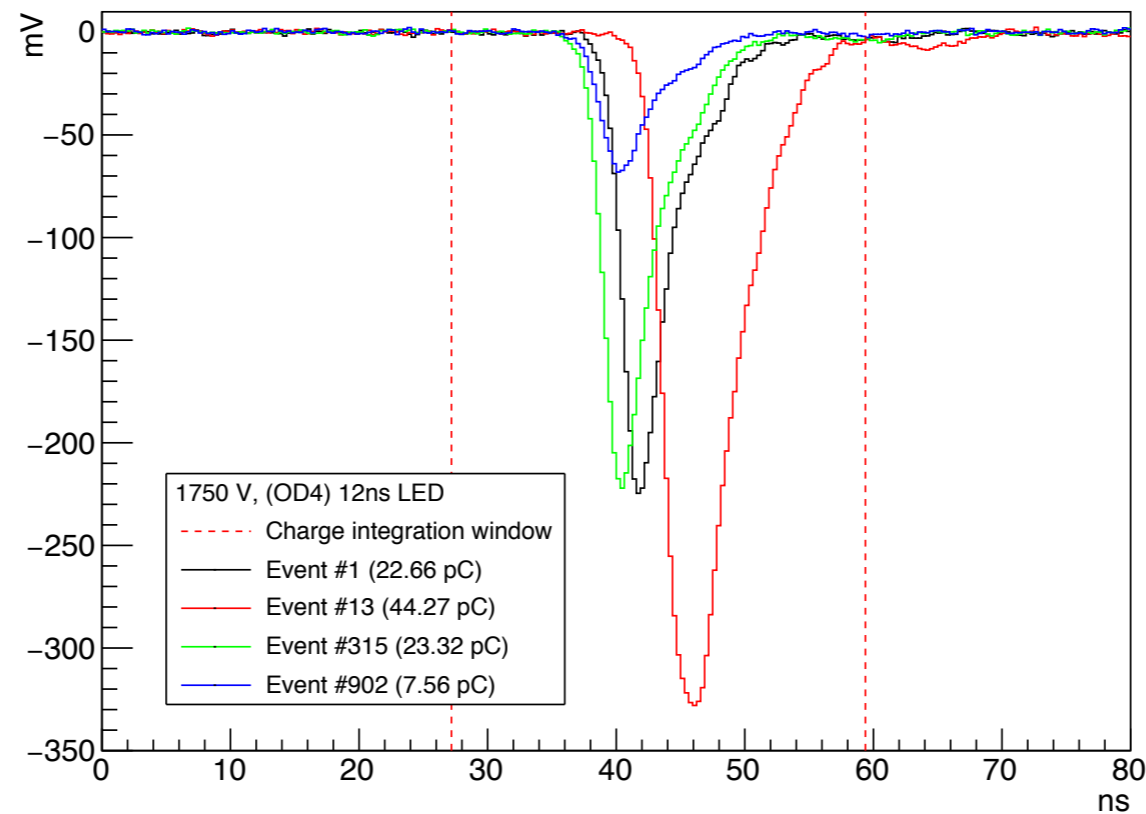
- Of course, this threshold is an arbitrary choice
- A more meaningful threshold, and thus dark rate measurement, is found relative to the single photoelectron (SPE) response

# R7725 Single Photoelectron Response

# Calculating charge with V1743



- While waveform amplitude is a directly triggerable quantity and of interest for the dark rate, the total charge is of more interest when calculating PMT gain
- To calculate the charge for a given waveform, define a charge integration window:



- And integrate as:

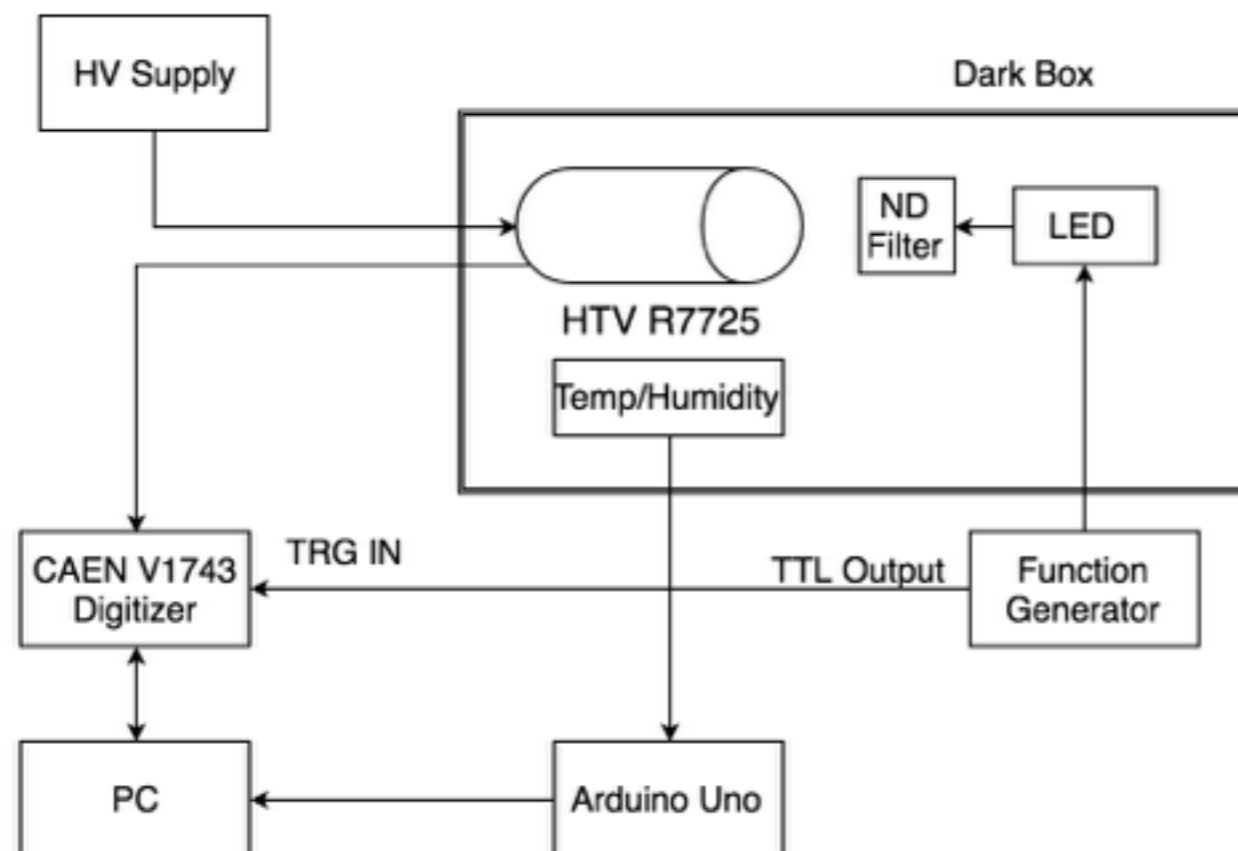
$$\text{Charge [pC]} = 10^9 \frac{\text{pC}}{\text{mC}} \cdot \sum_i \frac{V_i [\text{mV}] - \text{baseline}}{50 \Omega} \frac{1 \text{ s}}{3.2 \times 10^9}$$

- where the “baseline” is the average voltage in the first 16 samples of the waveform

# Observing the SPE



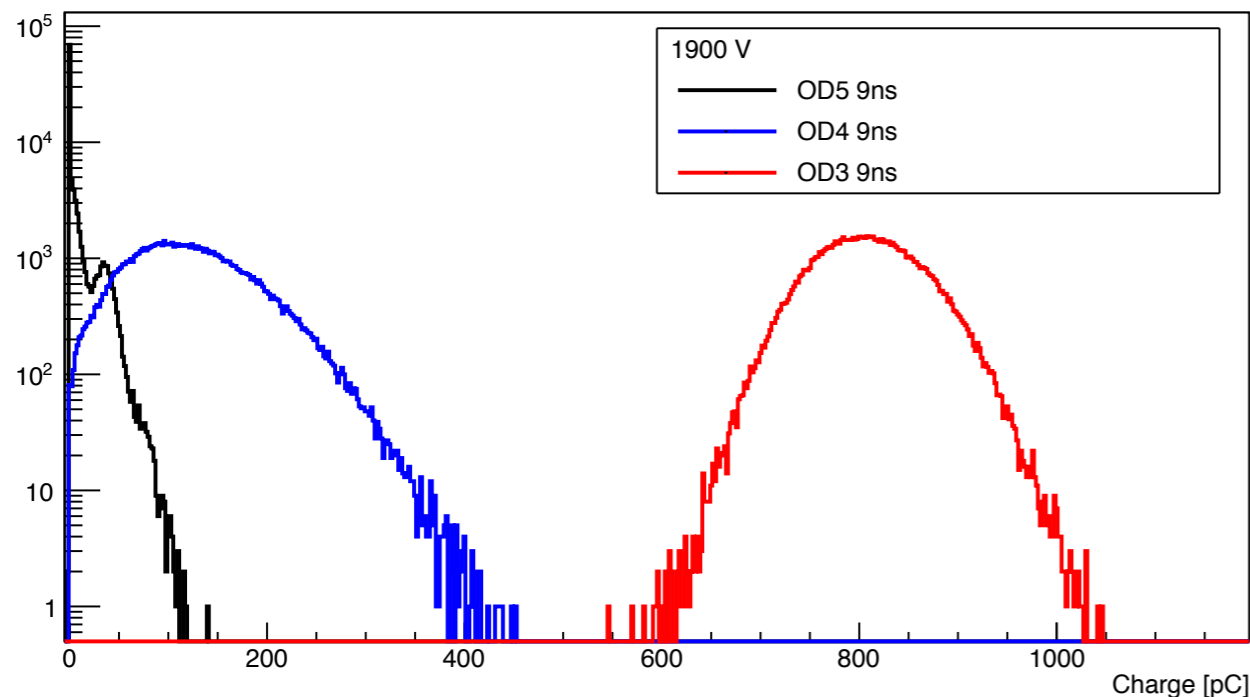
- To observe the single photoelectron (SPE) response:
  - Drive the LED with 3V “square” pulses at 1800 Hz for short periods of time
    - Function generator has rise/fall time of 2.5ns, less square at short widths
  - Configure the digitizer for its maximal sampling rate (3.2 GS/s) and shortened record length (256 samples) to handle 1800 Hz incoming trigger rate
  - Trigger on the LED driving signal (TTL output to V1743 TRG IN)



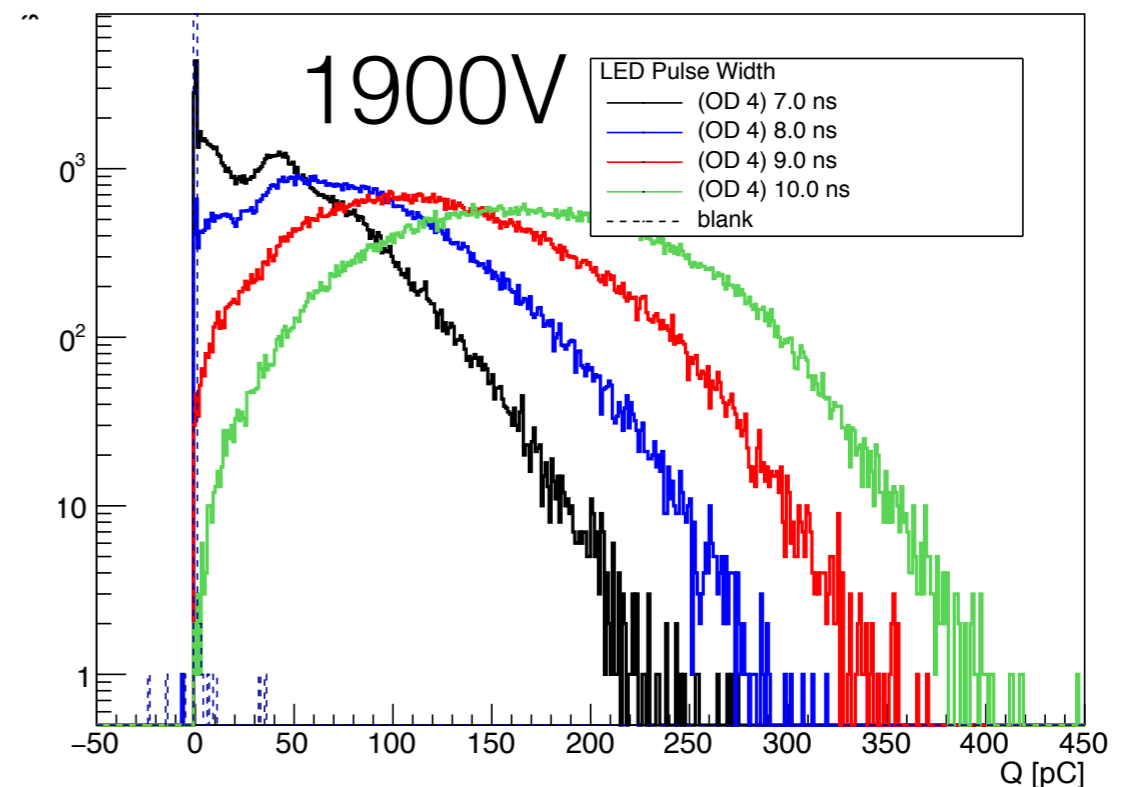
# Observing the SPE



- Several “knobs” to adjust the light intensity:
  - LED driver pulse width
  - Optical density of neutral density filters
    - Have 3 different filters to combine (OD1, OD3, OD5 — Transmission ~ 10%, 0.1%, 0.001%)
- Additionally a “knob” to adjust the PMT gain, the supply voltage
  - Range from 1200 V to 1900 V (max rating is 2000 V, recommended 1750 V)



Adjusting ND filters



Adjusting LED width



# Calculating SPE Response

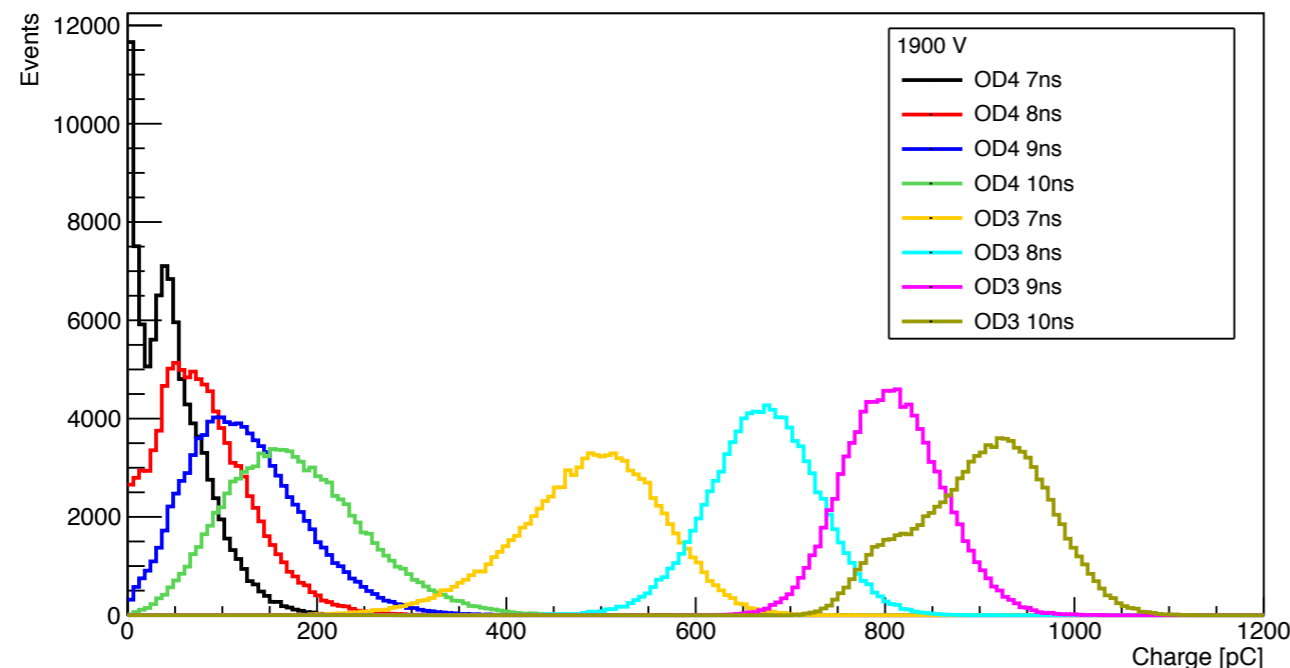


- We attempt several separate methods to determine the single photoelectron (SPE) response:
  - Method 1 — in the many-PE (gaussian) regime, scale the mean charge of distributions with the light intensity
    - Altering intensity with the LED pulse width or by changing ND filters
  - Method 2 — functional fits of the charge distribution
  - Method 3 — “Model Independent Approach” paper method
- Some definitions first:
  - $\mu, \langle NPE \rangle$  — average number of photoelectrons
  - $Q_1, \sigma_1$  — single photoelectron charge and width
  - $Q_\infty, \sigma_\infty$  — large NPE, overall distribution mean
  - c.f. Bellamy et al — *Nucl. Inst. and Meth. for Phys. Res. A 339 (1994) 468-476*

# SPE: Gaussian Regime



- Method 1:
- In the limit of large  $\langle NPE \rangle$ , the poisson distribution of NPE approaches gaussian
- Charge distribution approaches a gaussian with\*  
$$Q_\infty = \mu Q_1$$
$$\sigma_\infty = \sqrt{\mu(\sigma_1^2 + Q_1^2)} \quad \longrightarrow \quad \left(\frac{Q_\infty}{\sigma_\infty}\right)^2 = \frac{Q_1^2}{\sigma_1^2 + Q_1^2} \cdot \mu$$
- So we take several datasets at high gain (1900 V) and higher LED intensities
  - Using ND filter densities of OD4 and OD3

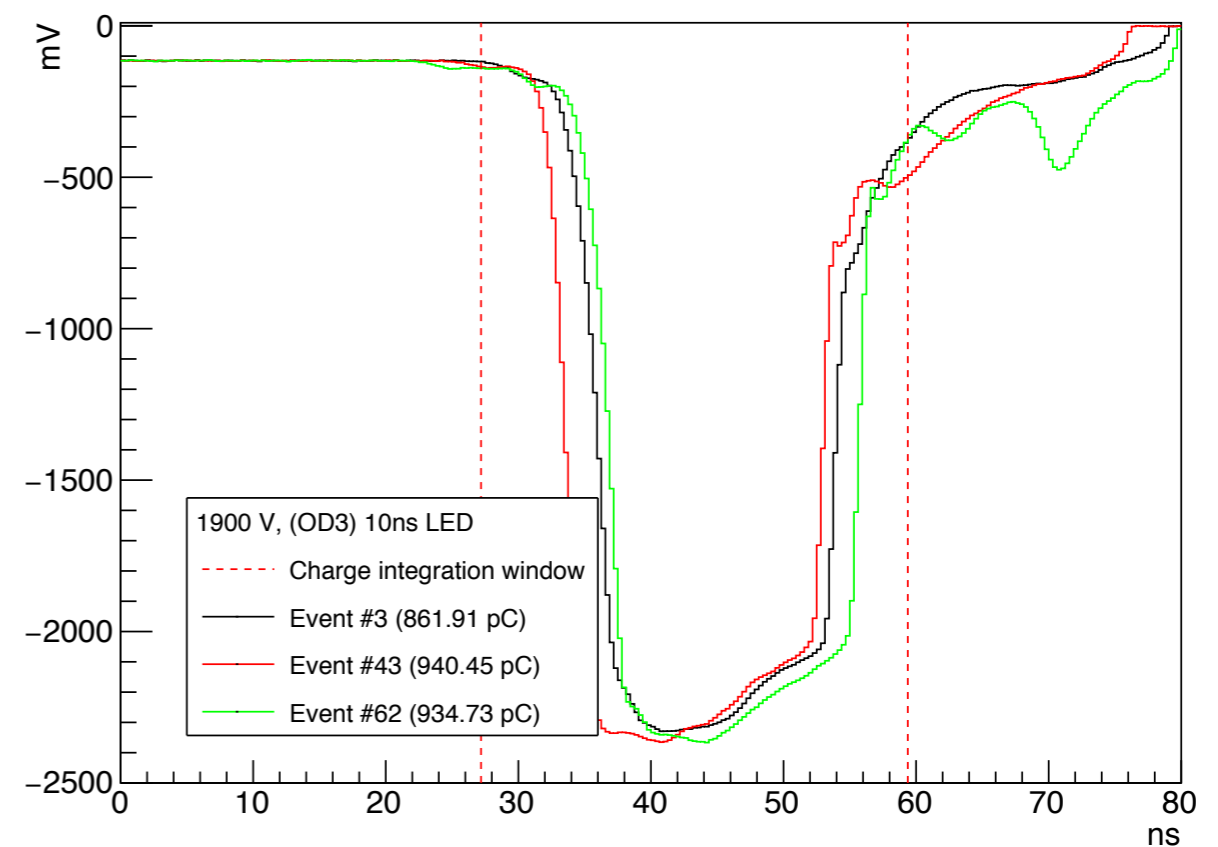
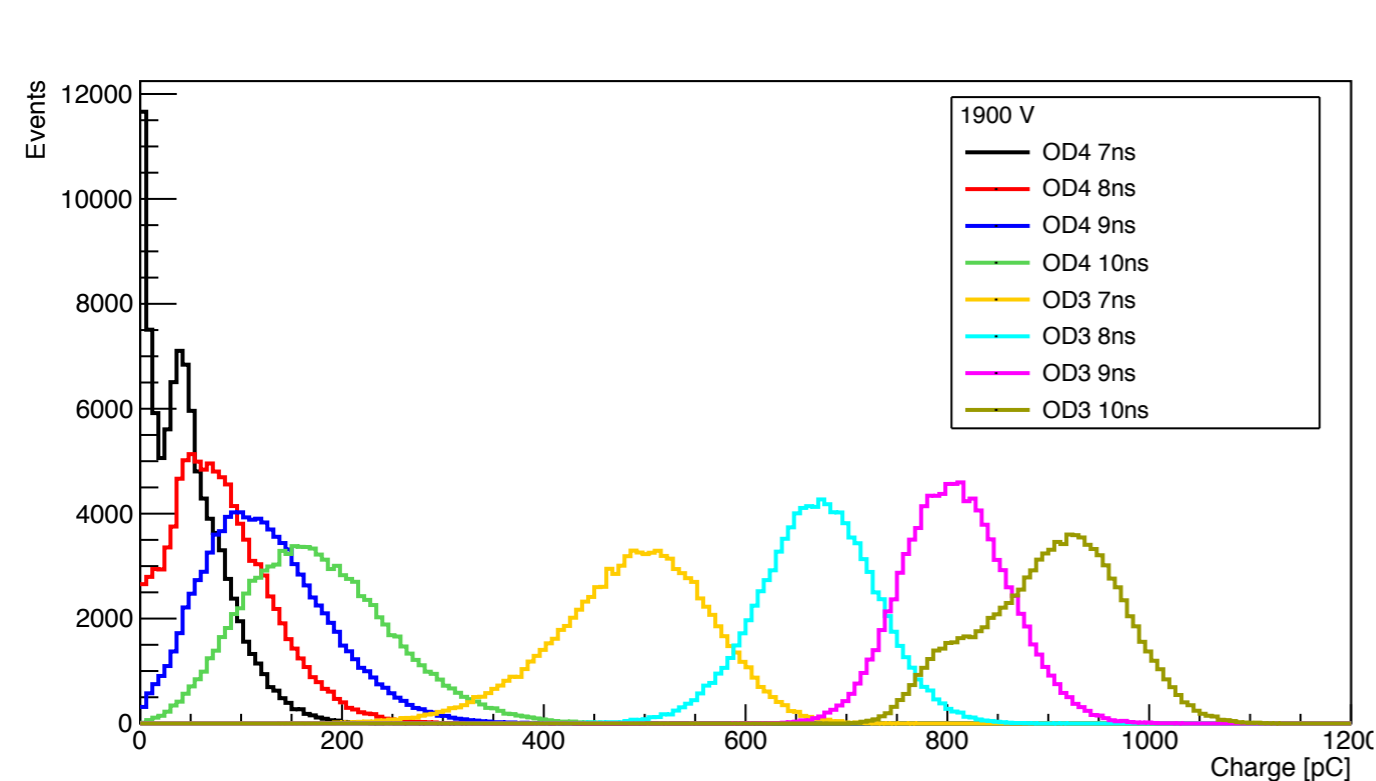


\*c.f. Bellamy et al

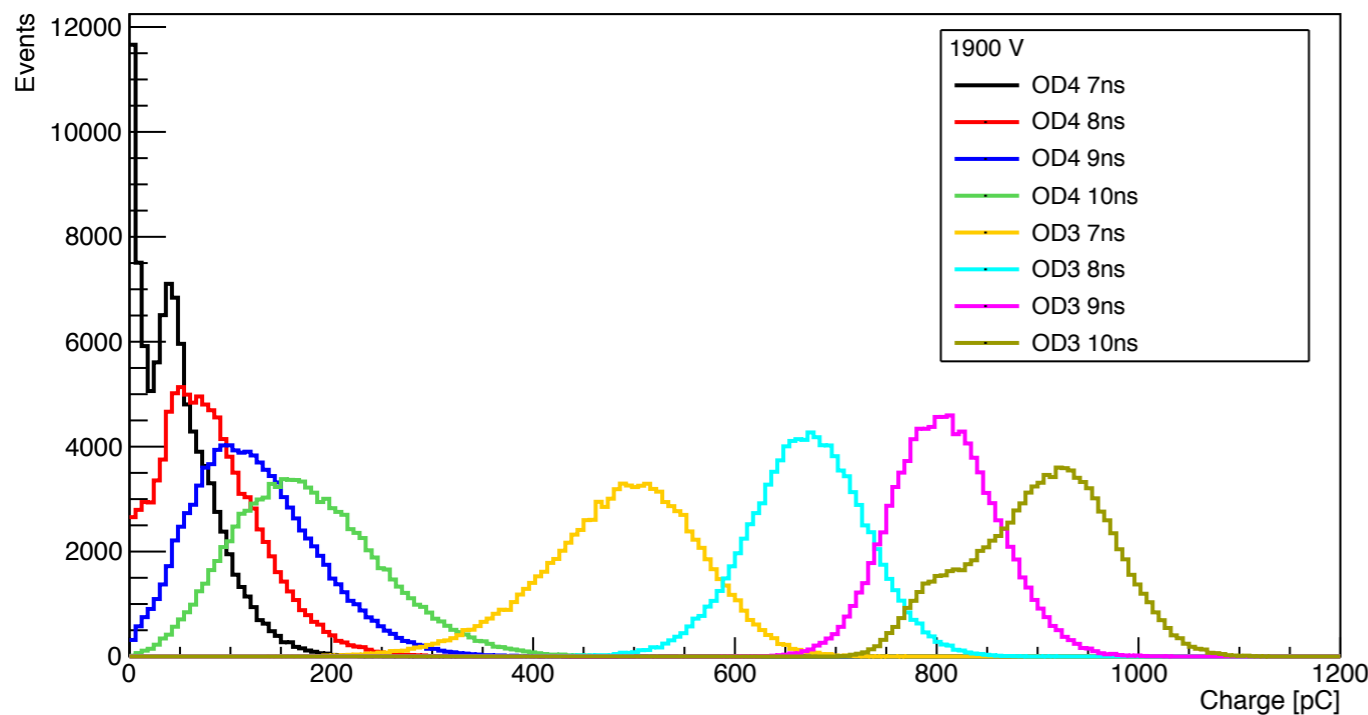
# SPE: Gaussian Regime



- In this range of LED intensities, only OD3 7-9ns and OD4 10ns appear fairly gaussian
- For the OD3 at 10ns dataset, very long and large pulses disturb the charge calculation by falling outside the defined integration window



# SPE: Gaussian Regime



	Mean (pC)	Width (pC)	(Mean/Width) <sup>2</sup>
<b>OD4 9ns</b>	124.6	61.2	4.1
<b>OD4 10ns</b>	177.5	72.7	6.0
<b>OD3 7ns</b>	487.0	76.9	40.1
<b>OD3 8ns</b>	670.7	58.1	133.4
<b>OD3 9ns</b>	808.2	52.3	239.2

For gaussians:

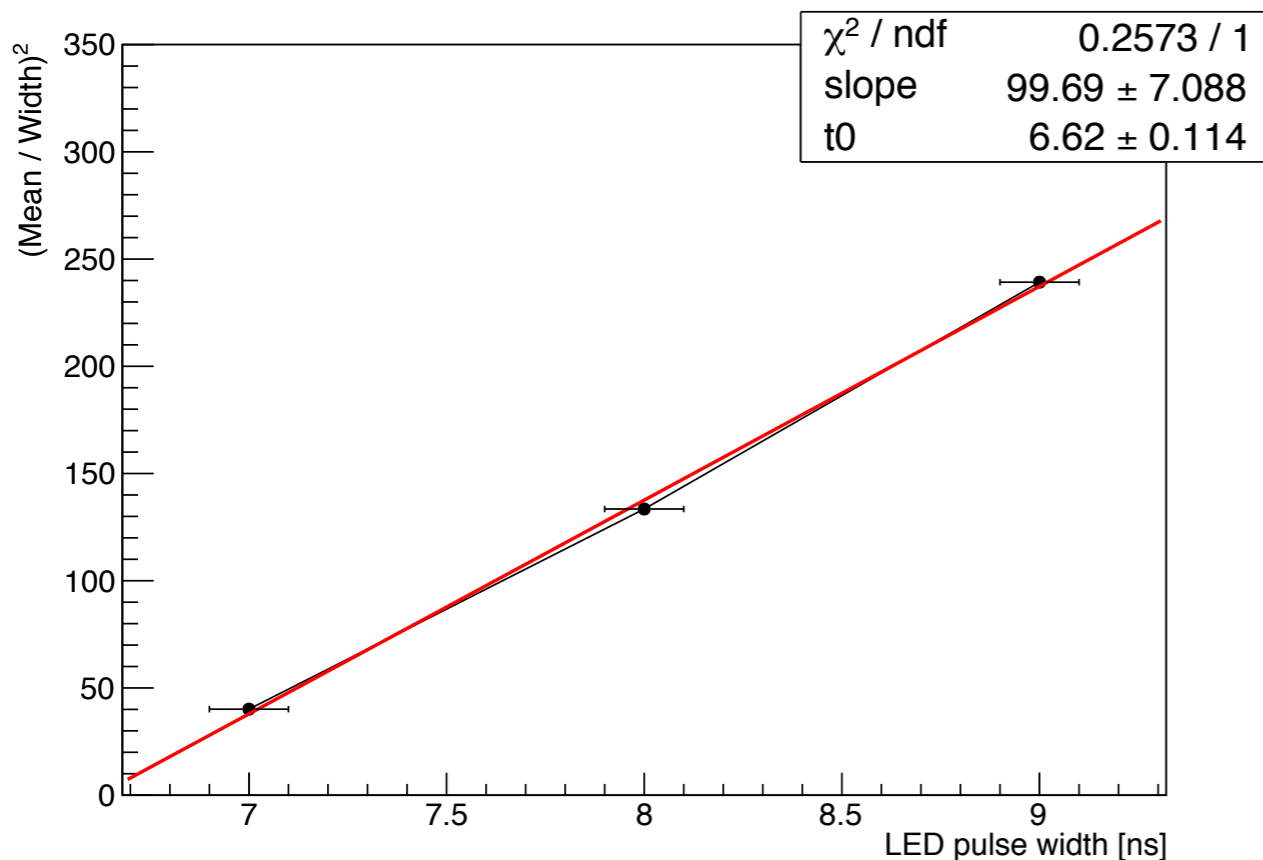
$$\left(\frac{Q_{\infty}}{\sigma_{\infty}}\right)^2 = \frac{Q_1^2}{\sigma_1^2 + Q_1^2} \cdot \mu$$

# SPE: Gaussian Regime



- Varying LED pulse width:
  - Assuming it takes time  $t_0$  for the LED to reliably begin emitting light
  - and assuming each nanosecond of LED light gives  $\alpha$  pC of charge, then expect

$$\left(\frac{Q_\infty}{\sigma_\infty}\right)^2 = \frac{Q_1^2}{\sigma_1^2 + Q_1^2} \cdot \alpha \cdot [\text{Pulse width} - t_0]$$

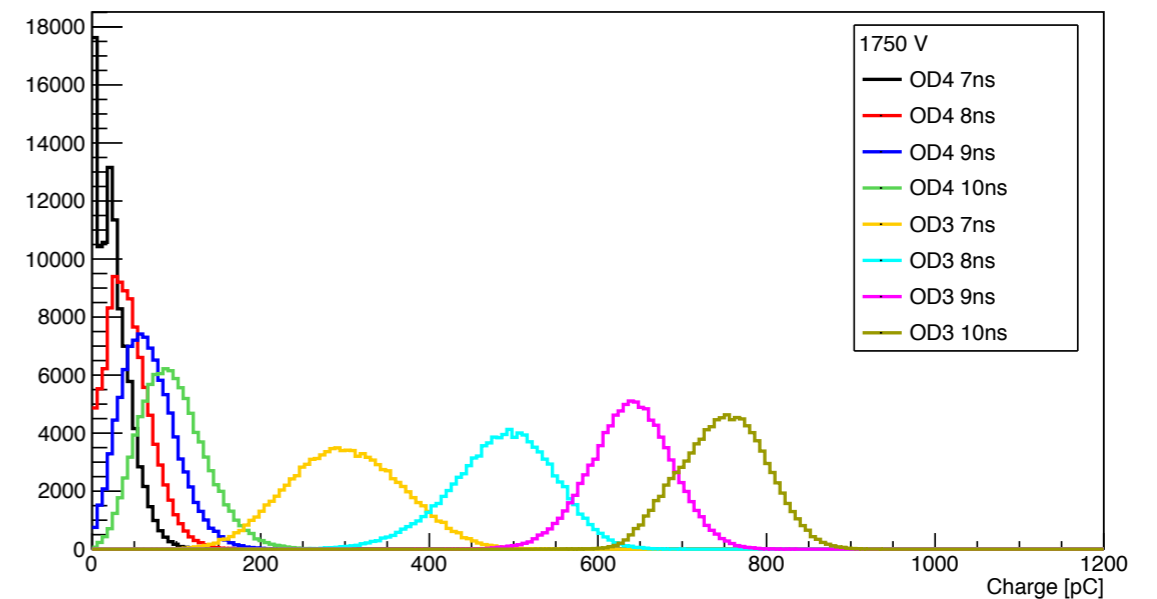
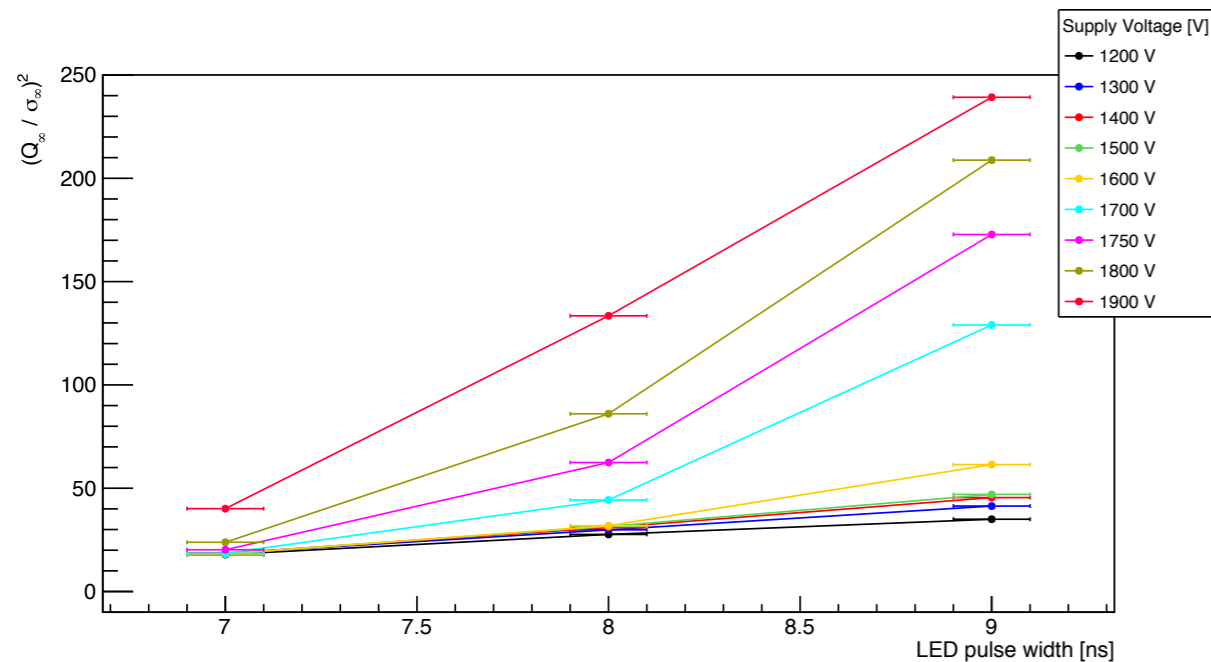


- Varying LED pulse width for OD3 at 1900V:
  - Quite linear
  - However the slope of this line is still a product of the light intensity and the SPE parameters
  - Cannot extract SPE here without calibrating the light source

# SPE: Gaussian Regime



- Lower PMT supply voltages however do not show this linear behavior:



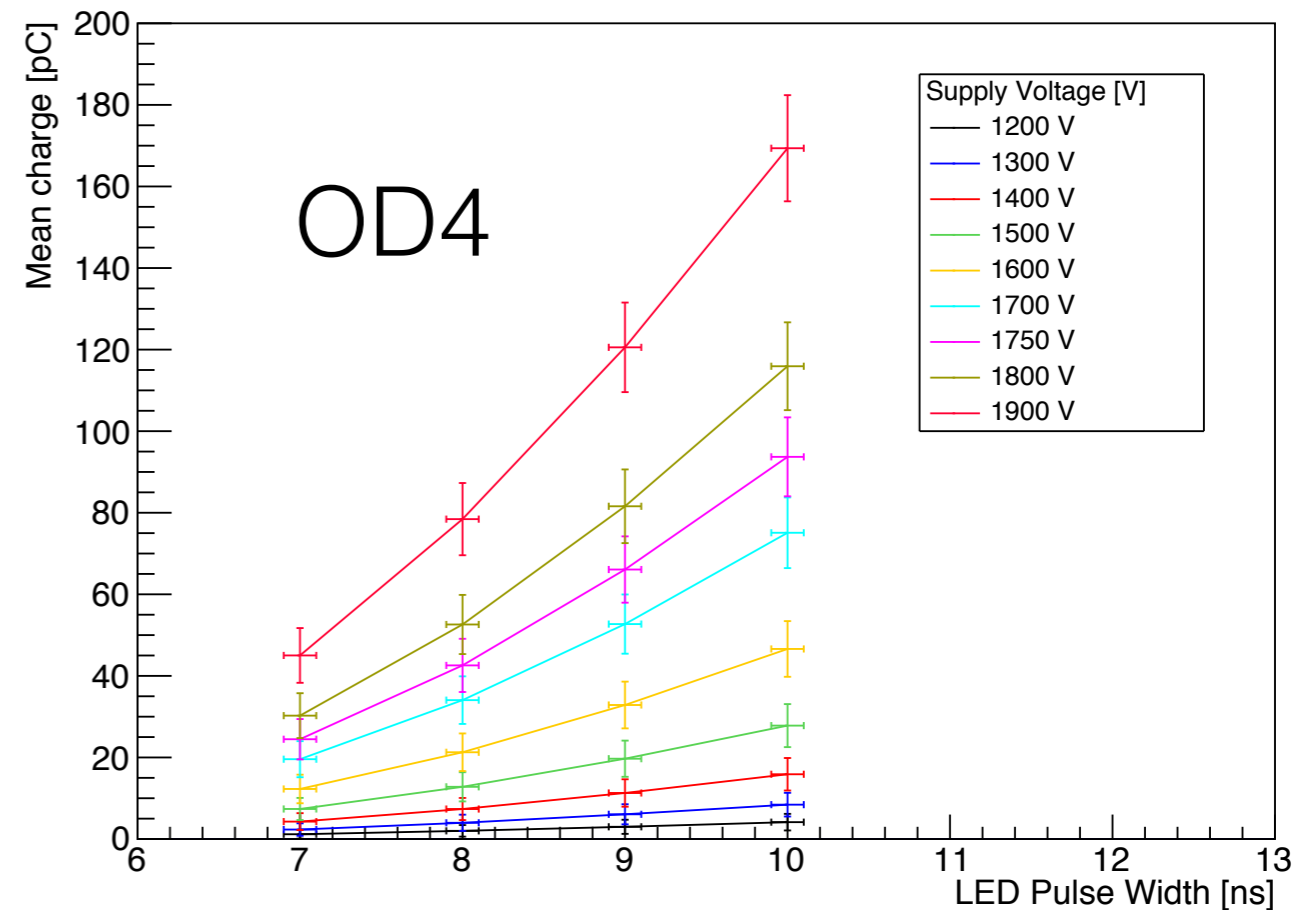
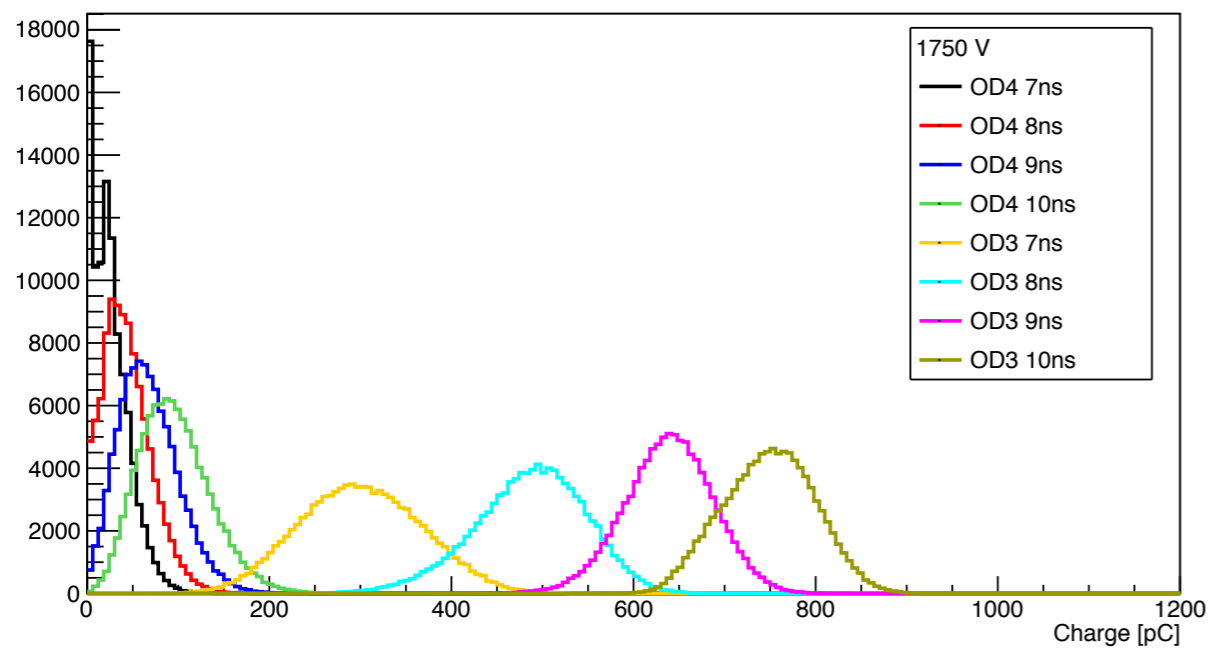
- Not clear why this is the case
- Something else that should scale with light intensity is the means of these distributions:

$$Q_\infty = \mu Q_1$$

# SPE: Gaussian Regime



- Plot the evolution of the mean charge with light intensity
- Appears to be fairly linear with LED pulse width



$$Q_{\infty} = \mu Q_1$$

# SPE: Gaussian Regime



- Fitting each of these to a line:

- Scales well at all supply voltages

$$Q_{\infty} = Q_1 \cdot \mu$$
$$= Q_1 \cdot \alpha \cdot [\text{Pulse width} - t_0]$$

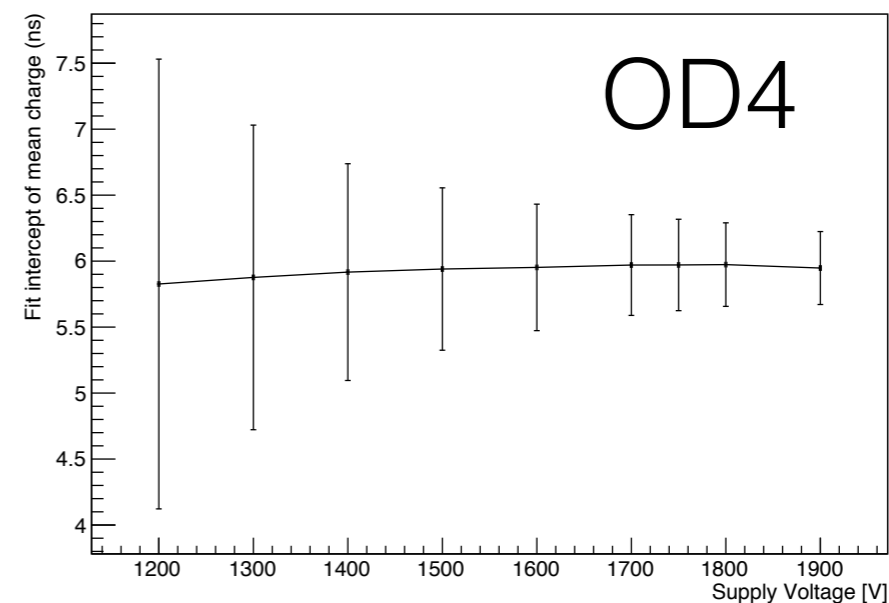
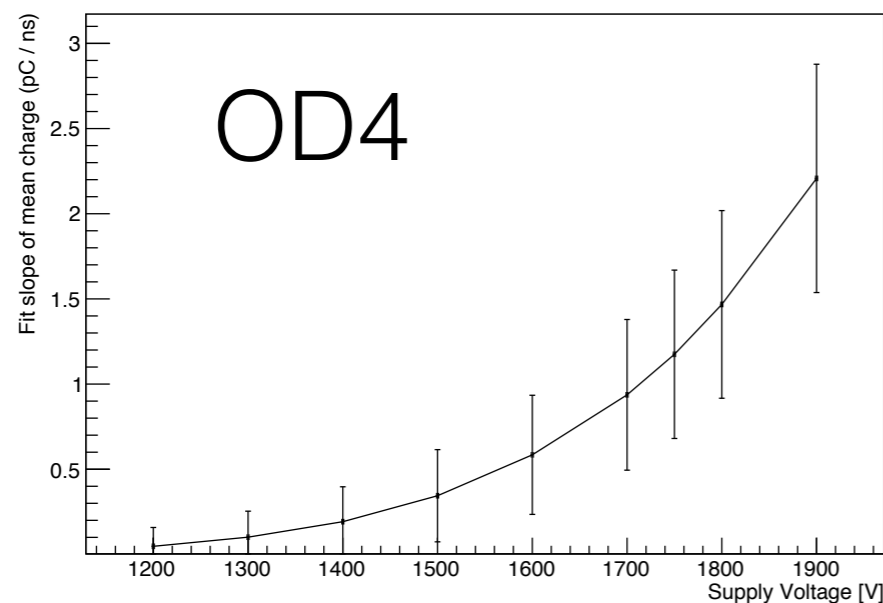
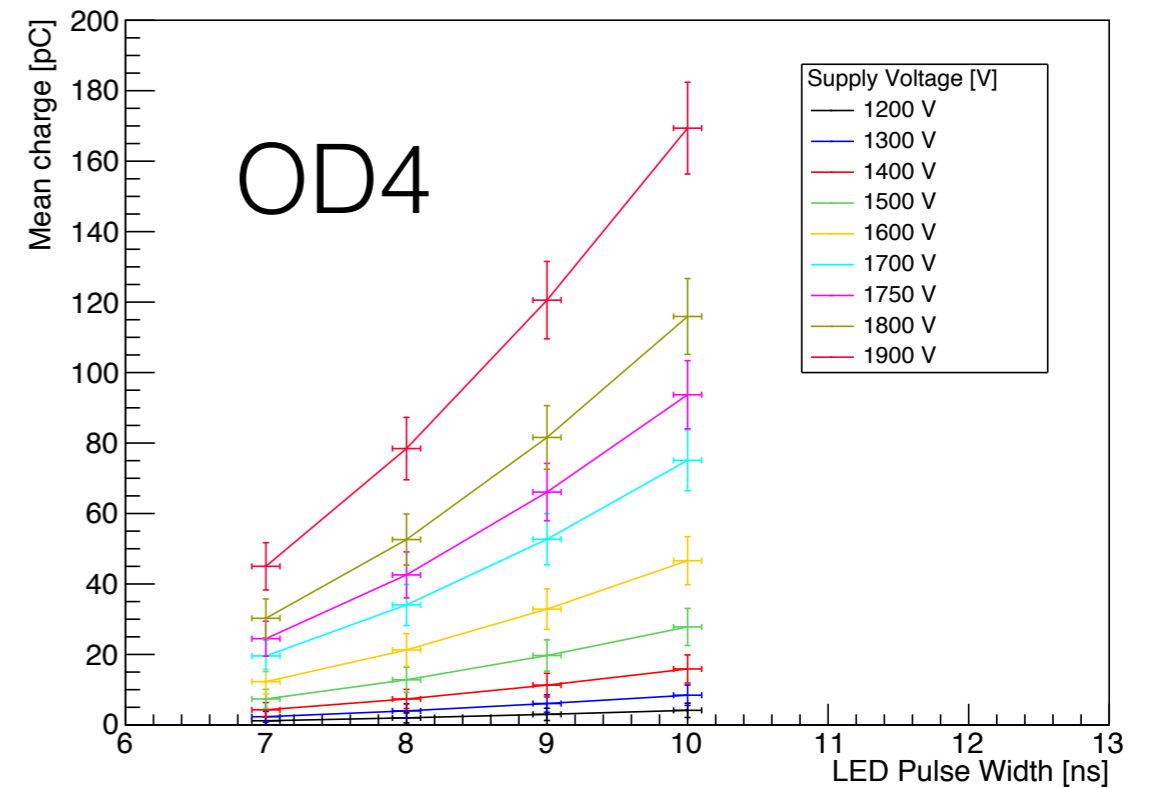
- For 1900V:

- Slope =  $2.2 \pm 0.7$  pC/ns

- Intercept =  $6.2 \pm 0.7$  ns

- Compare to intercept of previous fit:  $6.6 \pm 0.1$

- Again cannot separate SPE from light source



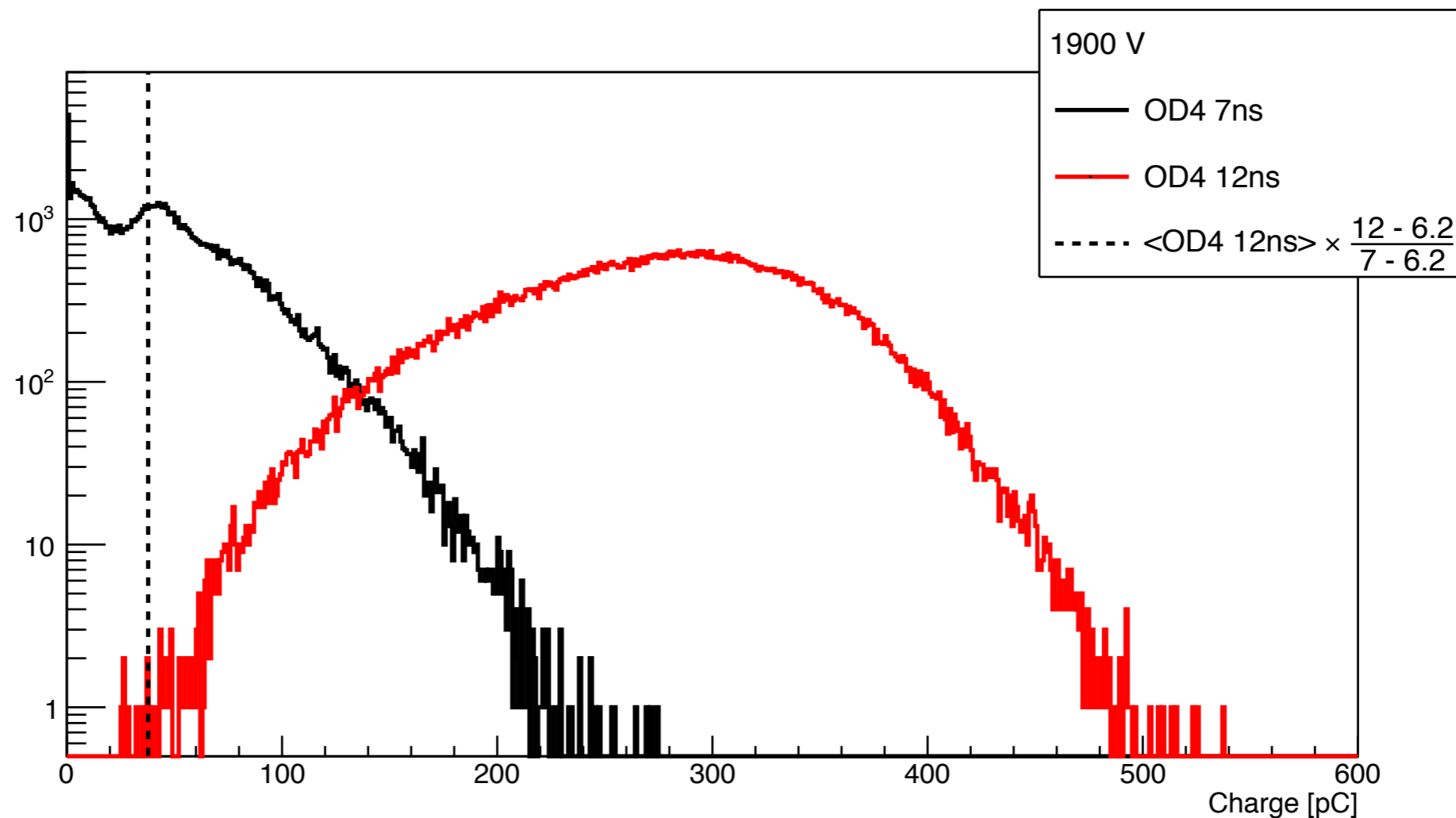


# SPE: Gaussian Regime



- Using the previous fit as a relationship between LED pulse width and expected mean charge, perhaps this can isolate the SPE
- With the same ND filters, scale the mean charge of a large-NPE distribution down to a small-NPE distribution

$$\langle Q_{12\text{ns}} \rangle \approx \frac{2.2(12-6.2)}{2.2(7-6.2)} \langle Q_{7\text{ns}} \rangle$$

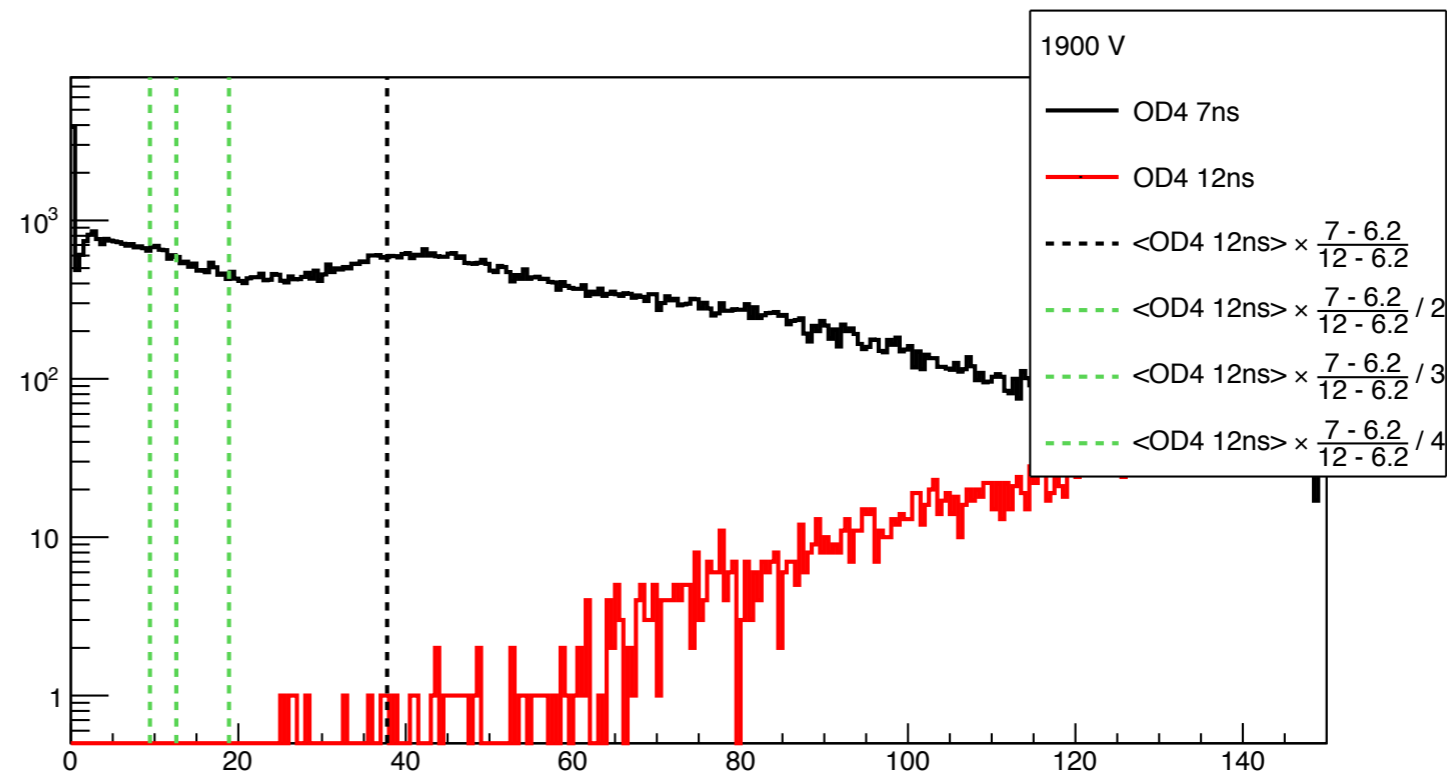


# SPE: Gaussian Regime



$$\langle Q_{12\text{ns}} \rangle \approx \frac{1}{7.3} \langle Q_{7\text{ns}} \rangle$$

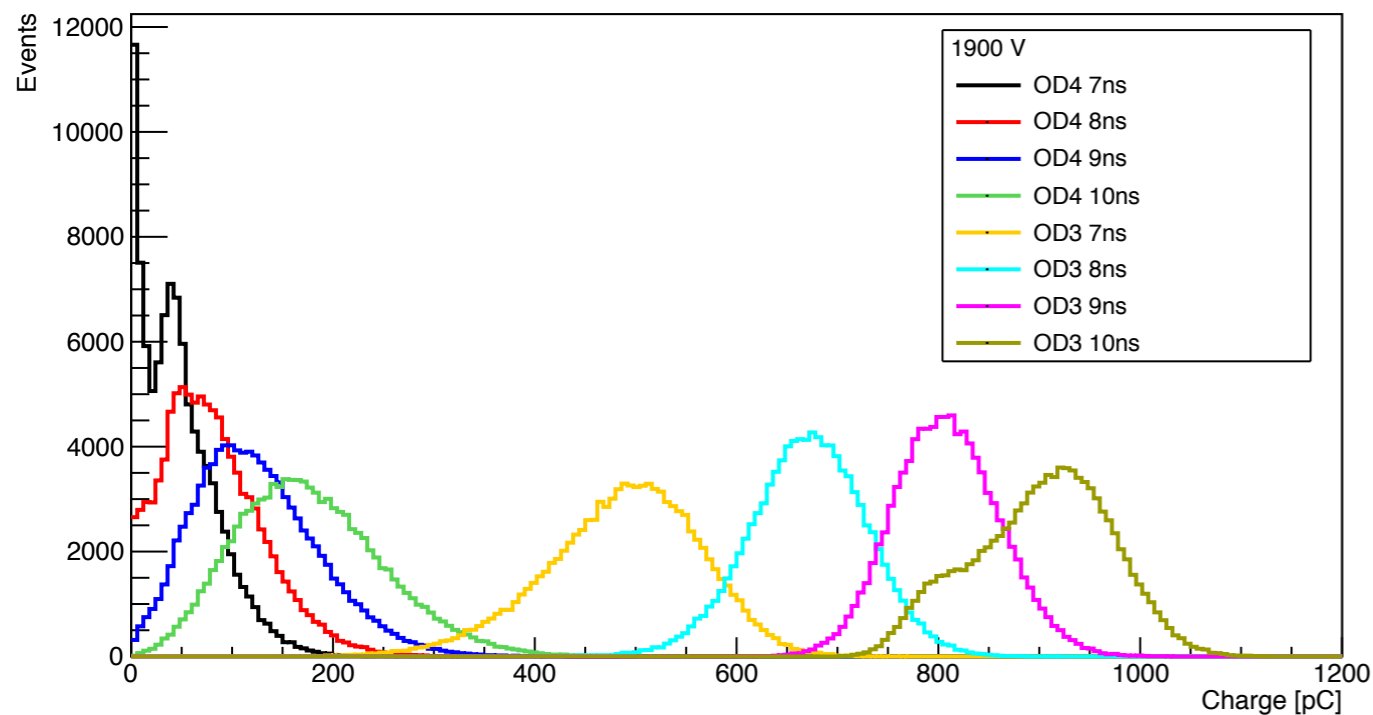
- This appears to scale down to the peak charge, but this does not guarantee you the SPE
- If the larger-NPE distribution had 30 PEs on average, for example, this method gives you a mean of  $\sim 4.1$  PE
- In fact if the larger-NPE distribution somehow has less than 7.3 NPE on average, this method gives you something less than the SPE
  - Projected mean: 37.8 pC
  - Actual mean: 45 pC



# SPE: Gaussian Regime



- One other scaling to consider, changing ND filters
- “OD4” is the combination of OD3 and OD1, so to change from OD3 to OD4 the OD1 filter is added
  - Expect a 10-fold decrease in light intensity



$$\text{Log}_{10}(0.15) = -0.82$$

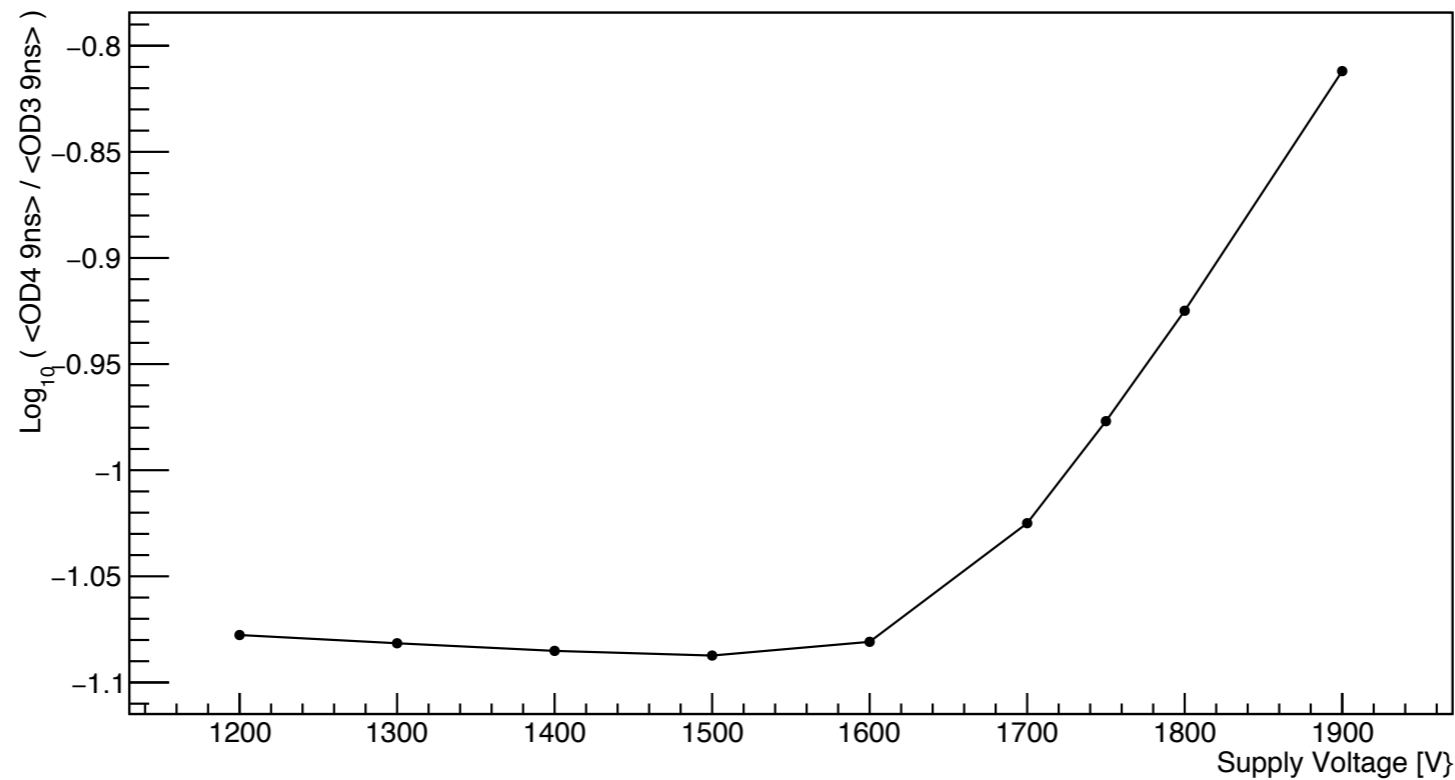
OD1 ~ OD 0.82 ?

Mean (pC)	
OD4 9ns	124.6
OD3 9ns	808.2
Ratio of OD4/ OD3	0.15

# SPE: Gaussian Regime



- Keeping the same LED pulse widths, adding an OD1 neutral density filter should decrease the mean charge by a factor of 10
  - This decrease should be the same at all PMT supply voltages

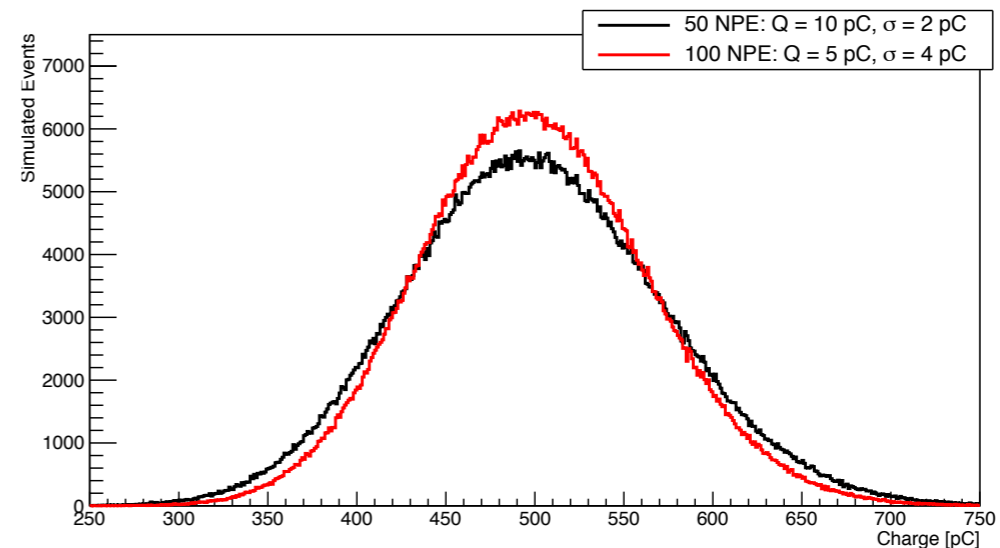
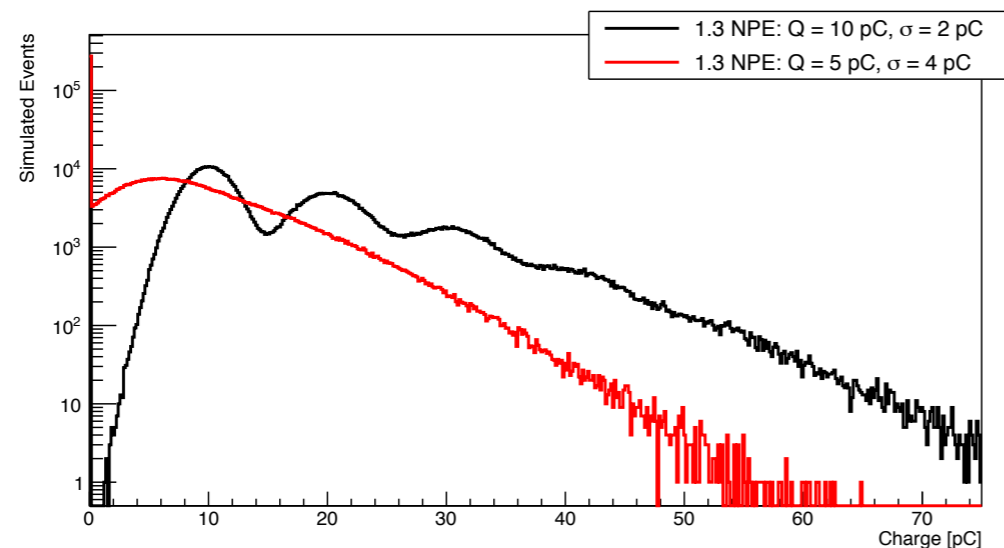


- Not clear why this is not a constant, or why it deviates from -1.0
  - For OD3 and OD4, 9ns should be well into larger NPE values so as to be fairly gaussian

# SPE: Gaussian Regime



- Difficulties aside, working in the gaussian regime does not seem to isolate the SPE response
- Shown below, two different SPE responses with different  $\langle \text{NPE} \rangle$ :
  - Simulated by RNG shooting NPE from poisson, and charge per PE from gaussian distributions

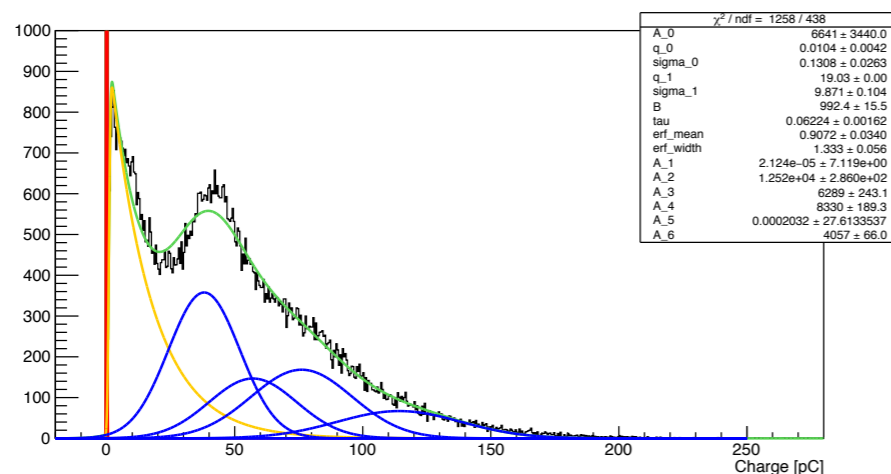


- Easily distinguished at low NPE, but possible for distributions to be very similar at large (yet different) NPE

# SPE: Functional Fit

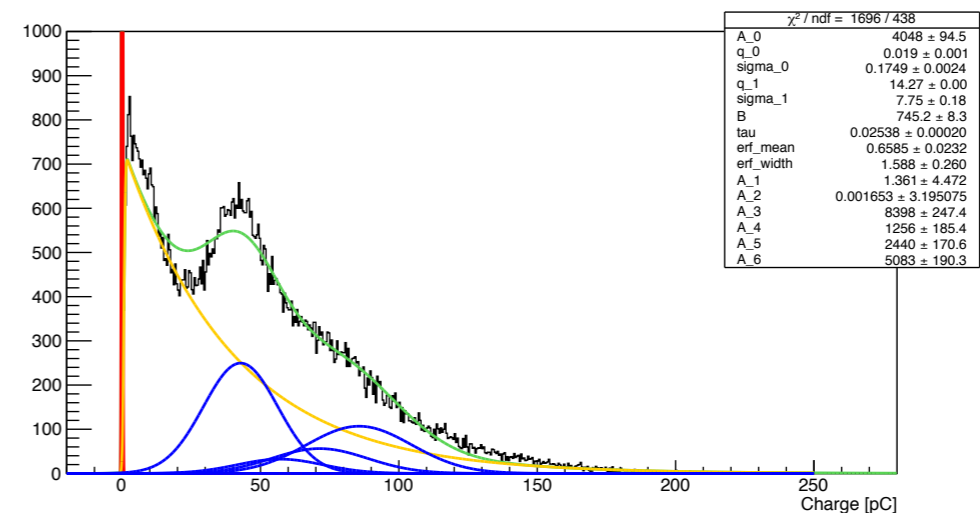


- Method 2: assume a functional form of the expected distribution and extract the SPE parameters
  - Gaussian pedestal
  - Exponential background
  - Sum of several signal gaussians:  $\text{Gaus}(x, i * Q1, \text{sqrt}(i) * \text{sigma}_1)$ 
    - Without the constraint on peak spacings and width, fits fail to converge as yet
- Even with these assumptions and constraints, fit is very dependent on initial conditions
- For OD4 7ns at 1900 V:



Initial Q1 value = 16 pC

→  
Q1 = 19 pC



Initial Q1 value = 12 pC

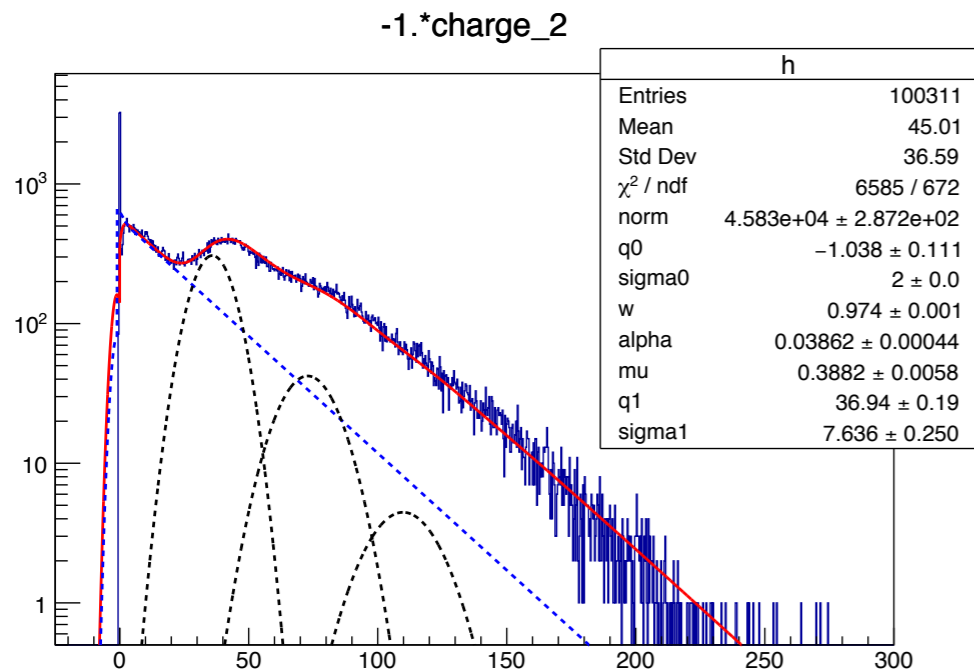
→  
Q1 = 14 pC

- Safe to say that this method is not yet working
- Consistently if you initialize the SPE peak below the peak at ~45, the fit returns lower-N peaks as having nearly zero events

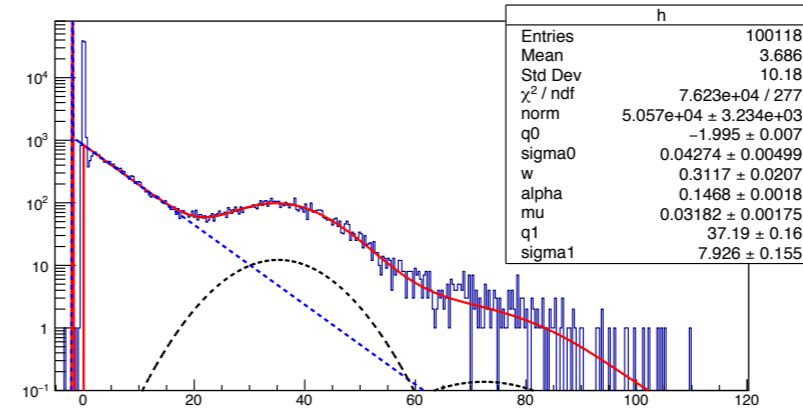
# SPE: Functional Fit



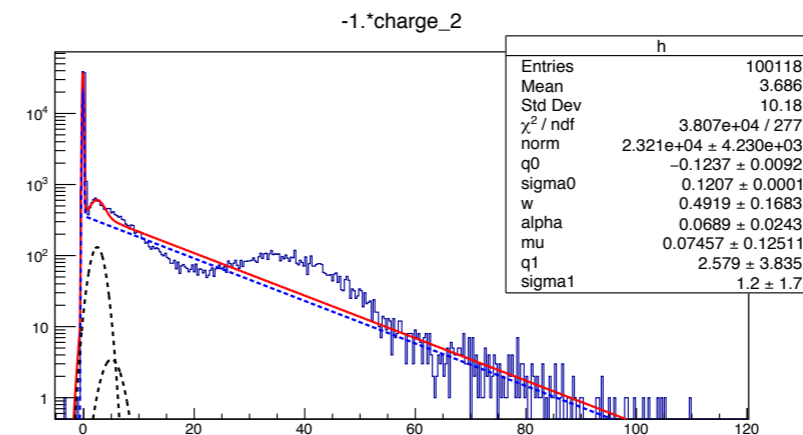
- A much more complicated functional form is described in Bellamy et al
  - Fully convolutes exponential background with each signal peak



OD4 7ns 1900V



Initial peak  
40 pC

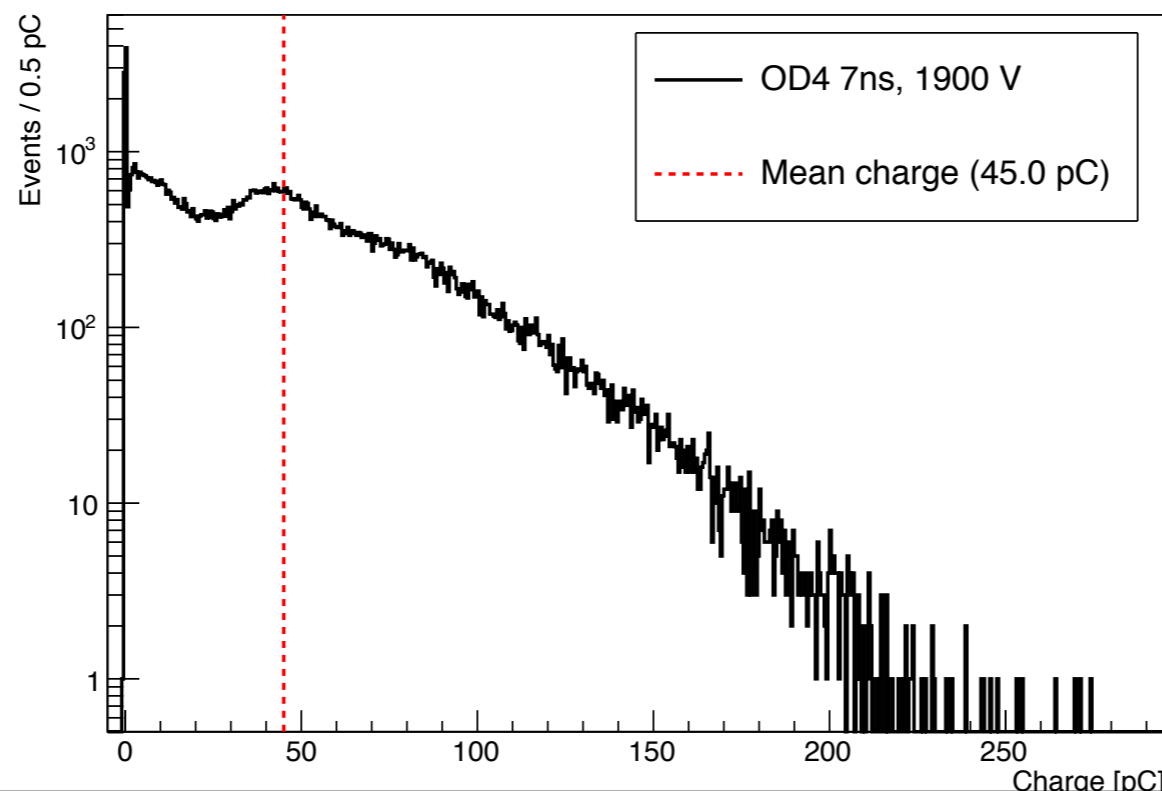


Initial peak  
20 pC

OD5 8ns 1900V

- In some instances this seems reasonable, but still is highly dependent on initial conditions
- Still involves some bias of suggesting the visible peak is the SPE peak

- Method 3: “Model Independent Approach” (MIA)
  - <https://arxiv.org/abs/1602.03150>
  - Method essentially has two ingredients:
    - 1) From a “blank”, no-LED dataset, use a low-charge cut to define ‘zero PE triggers’
    - 2) Compare the mean and variance of the “blank” dataset to that of an LED-on dataset

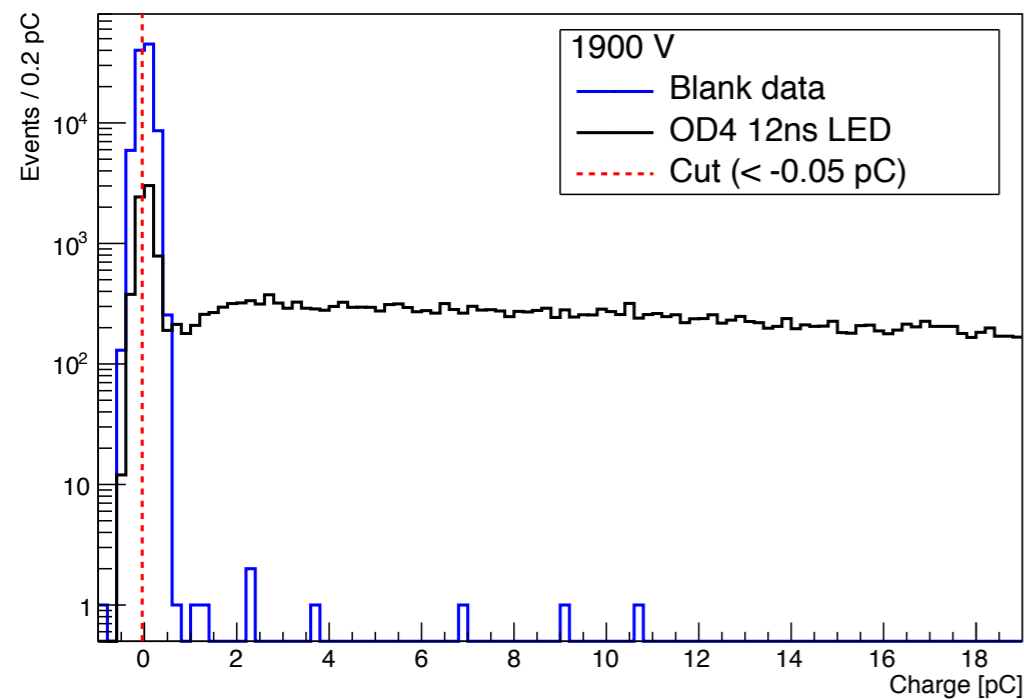




# MIA SPE Response



- To calculate (<https://arxiv.org/abs/1602.03150>):
  - Define a charge cut such that  $\epsilon \sim 1/3$  in blank data (LED is turned off)
    - This 1/3 is arbitrary and its choice a source of systematic uncertainty
  - Take the same number of blank and LED events
  - Assuming the average  $N(\text{PE})$   $\lambda$  is low and Poisson-distributed, the occupancy  $\lambda$  is directly related to the likelihood of observing zero PEs:



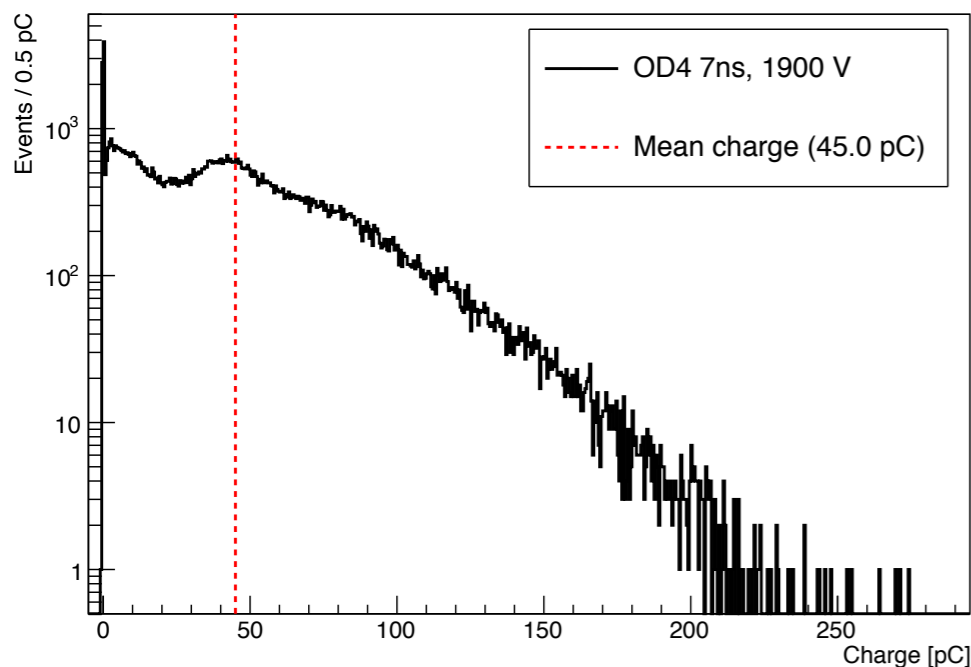
$$\begin{aligned}\lambda &= -\ln[L(0)] \\ &= -\ln\left[N_{\text{LED}}(Q < \text{cut}) \times \frac{1}{\epsilon_{\text{blank}} N_{\text{blank}}}\right] \\ &= -\ln\left[2010 \times \frac{1}{0.32975 \cdot 100,000}\right] \\ &= 2.80\end{aligned}$$

# MIA SPE Response



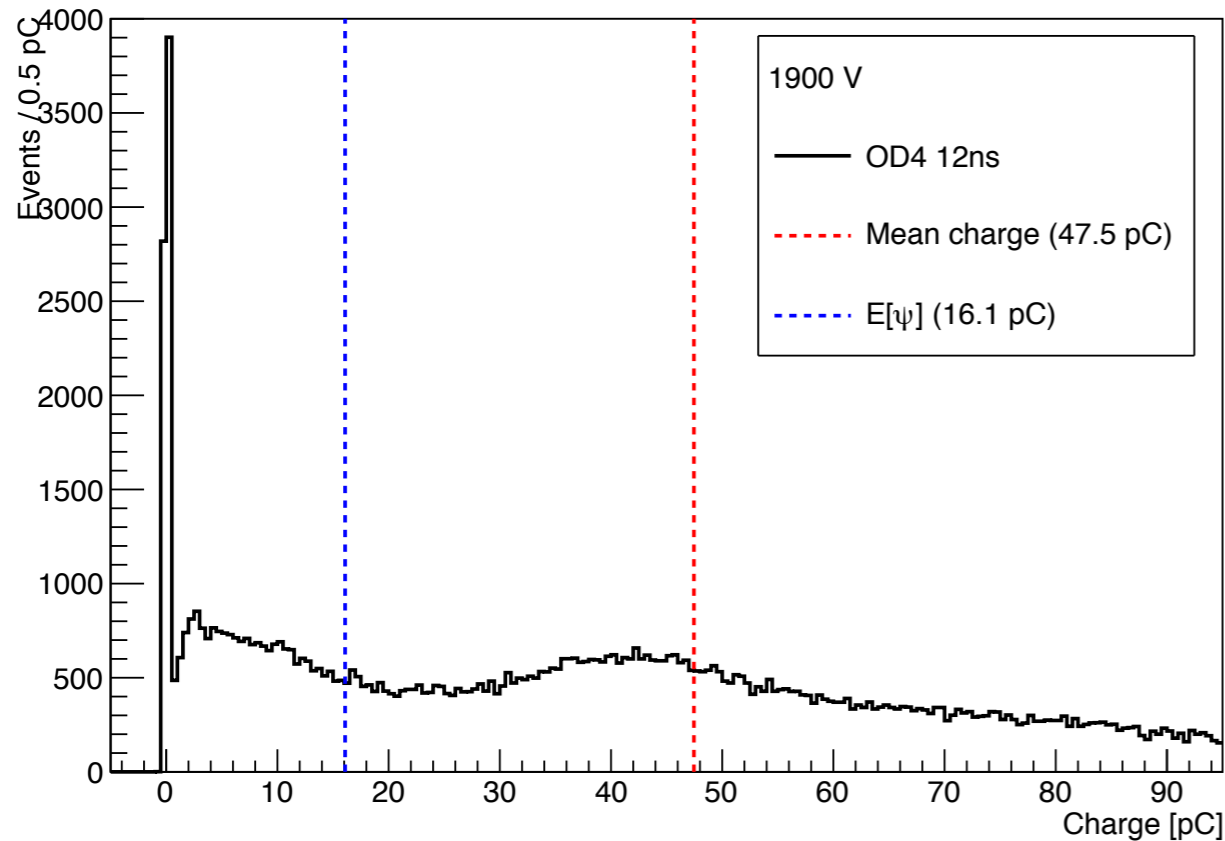
- Now with  $\lambda$ , and the assumptions that:
  - The signal and background distributions are uncorrelated and
  - Signal is a repeated convolution of the SPE:

- Then:
$$E[\psi] = \frac{E[\text{LED}] - E[\text{blank}]}{\lambda}$$
$$V[\psi] = \frac{V[\text{LED}] - V[\text{blank}] - E^2[\psi] \cdot V[\lambda]}{\lambda}$$



$$E[\psi] = \frac{45.0 - 0.013}{2.80} = 16.09 \text{ pC}$$
$$V[\psi] = \frac{1338.6 - 0.43 - 16.09^2 \cdot 2.80}{2.80} = 219.6 \text{ pC}^2$$

# MIA SPE Response

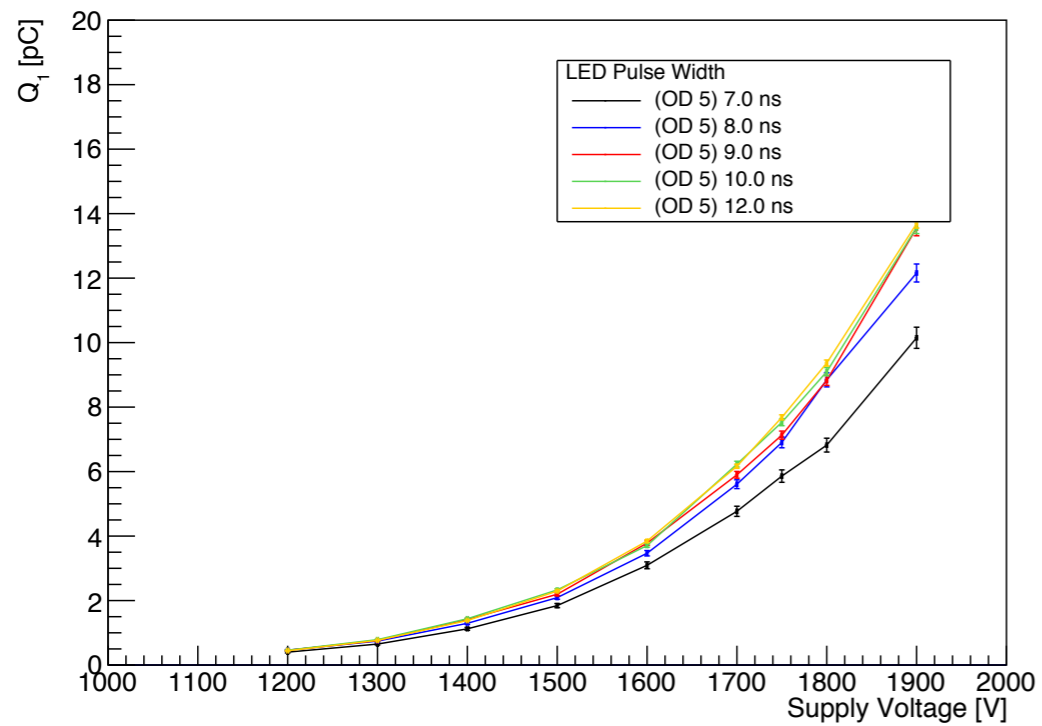


$$E[\psi] = \frac{45.0 - 0.013}{2.80} = 16.09 \text{ pC}$$

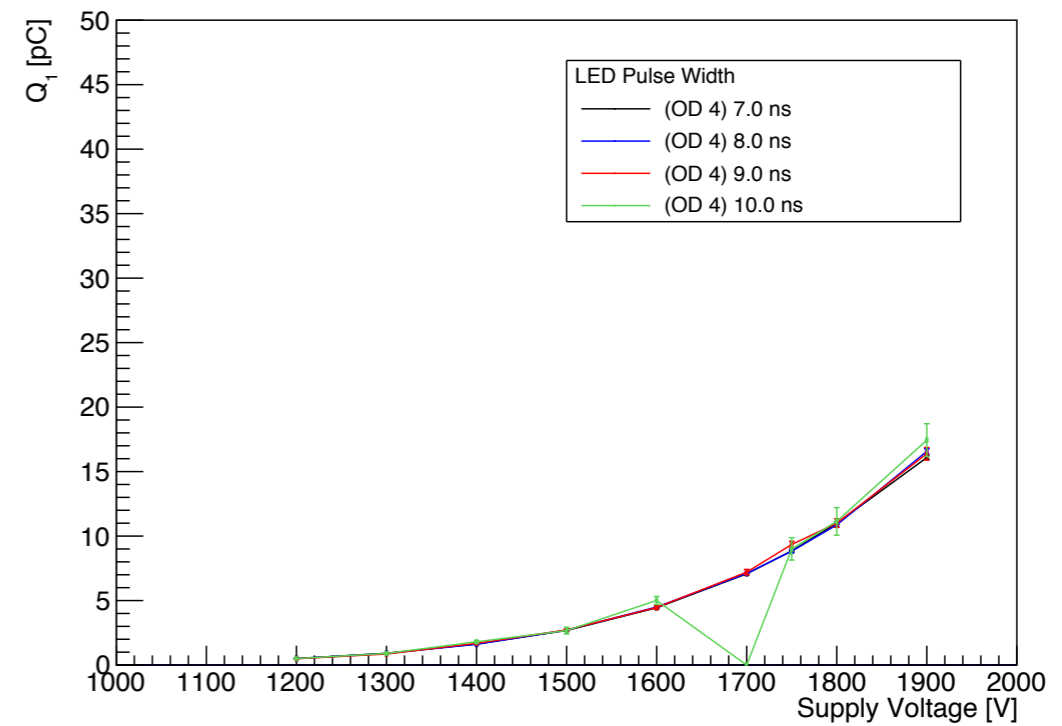
$$V[\psi] = \frac{1338.6 - 0.43 - 16.09^2 \cdot 2.80}{2.80} = 219.6 \text{ pC}^2$$

- This method consistently derives an SPE expectation well below the observable peak, with a very large width

# MIA SPE Charge

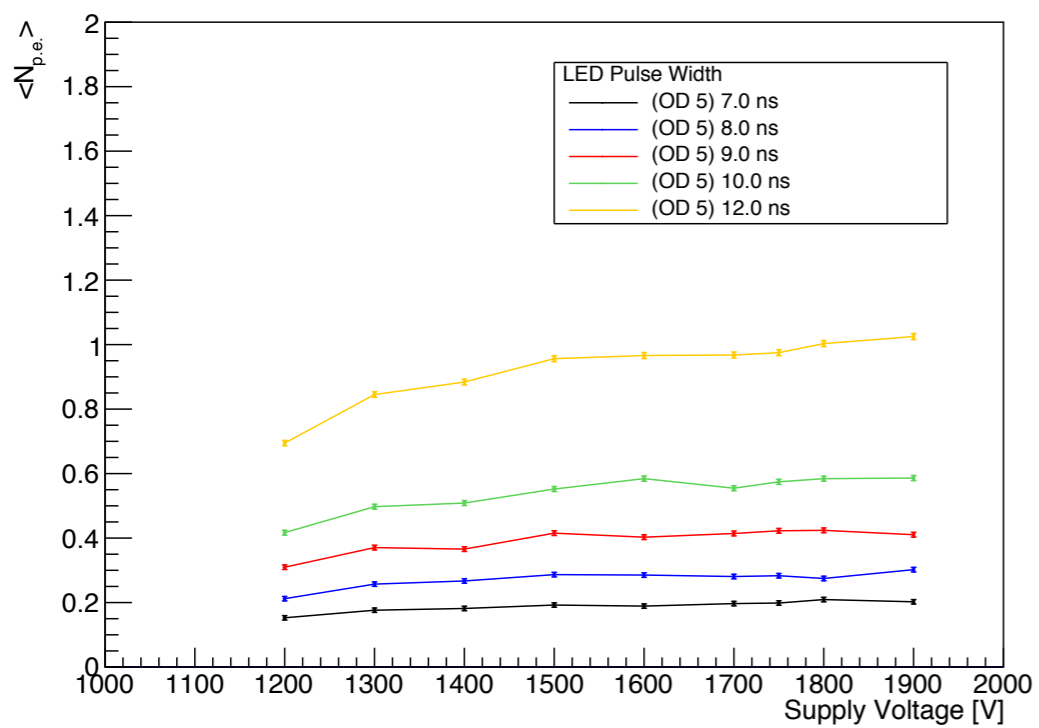


OD5

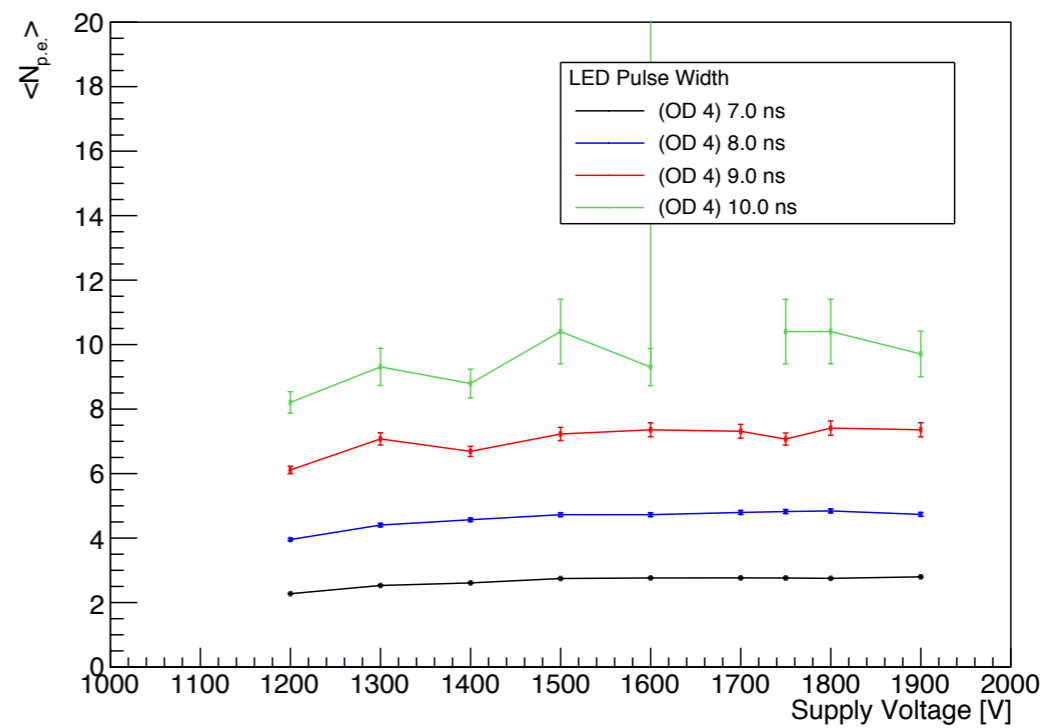


OD4

# MIA SPE NPE

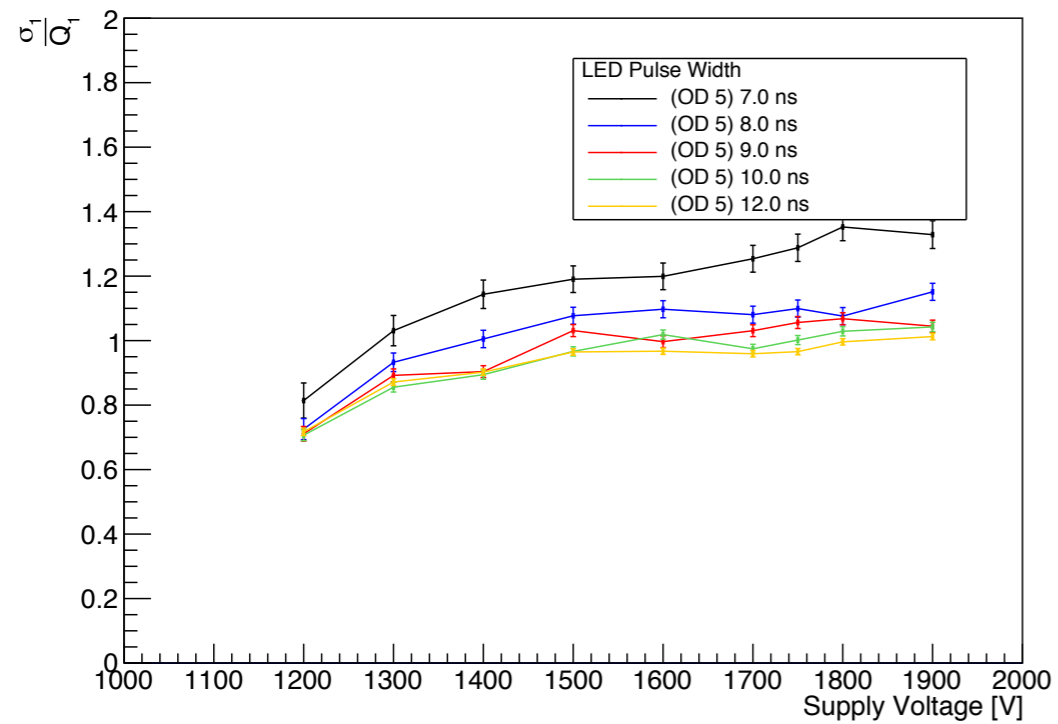


OD5

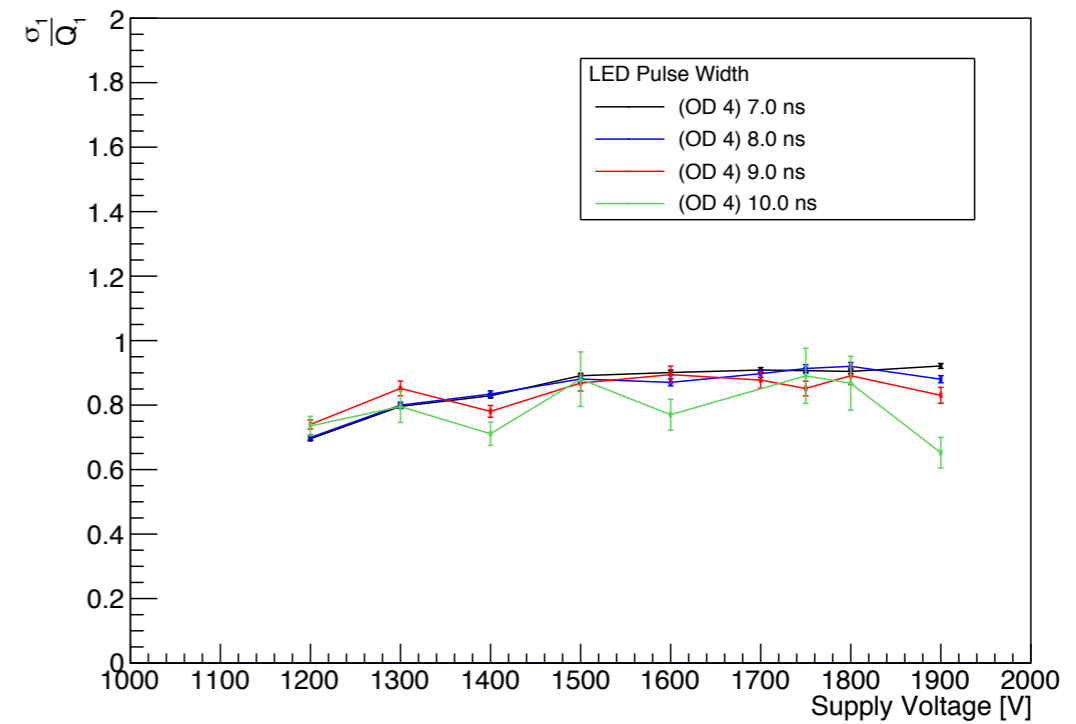


OD4

# MIA SPE Resolution (Width / Peak)



OD5

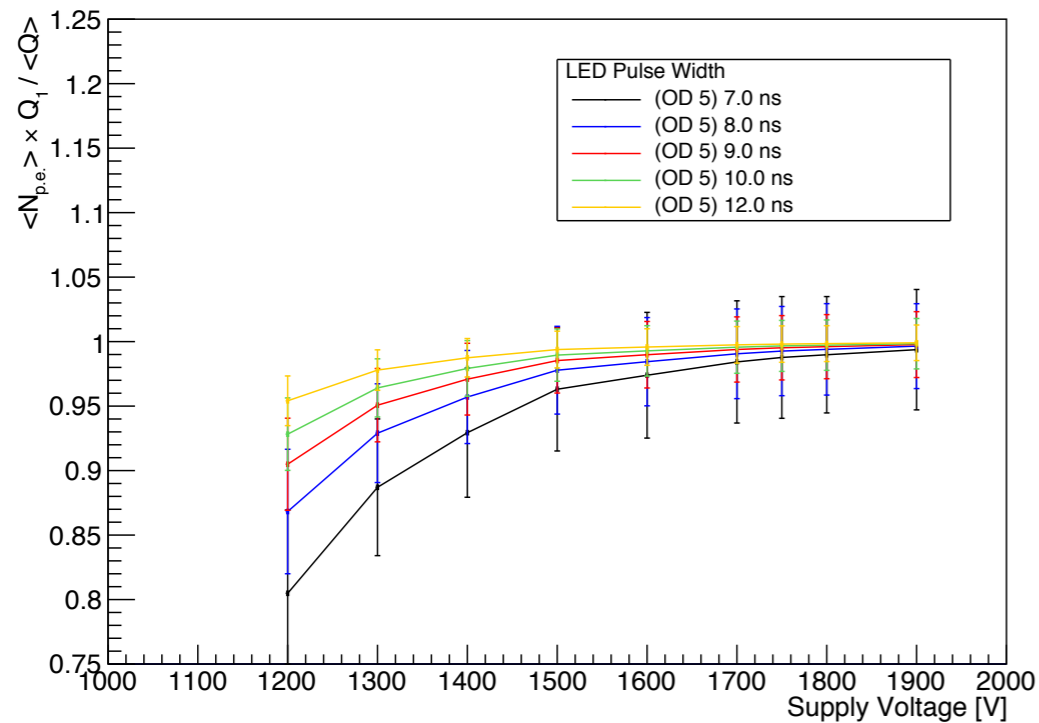


OD4

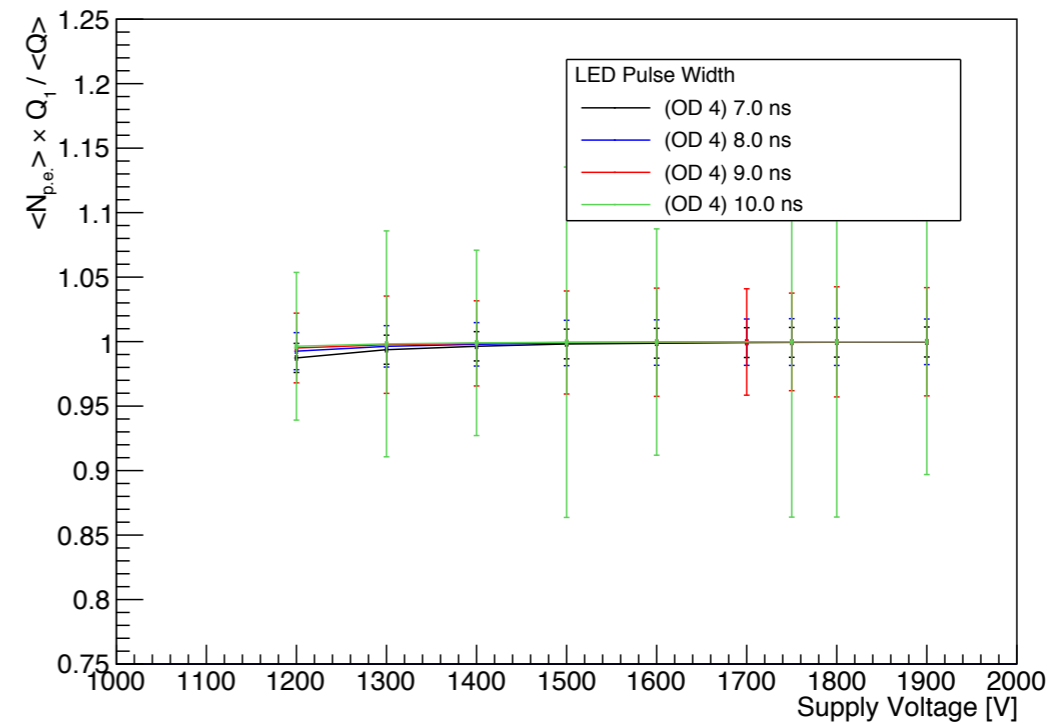
# MIA SPE Approaching Distribution Mean



- As the gain and light intensity are increased, this method does approach  $\langle N_{pe} \rangle * Q_1 \rightarrow$  mean charge of the distribution

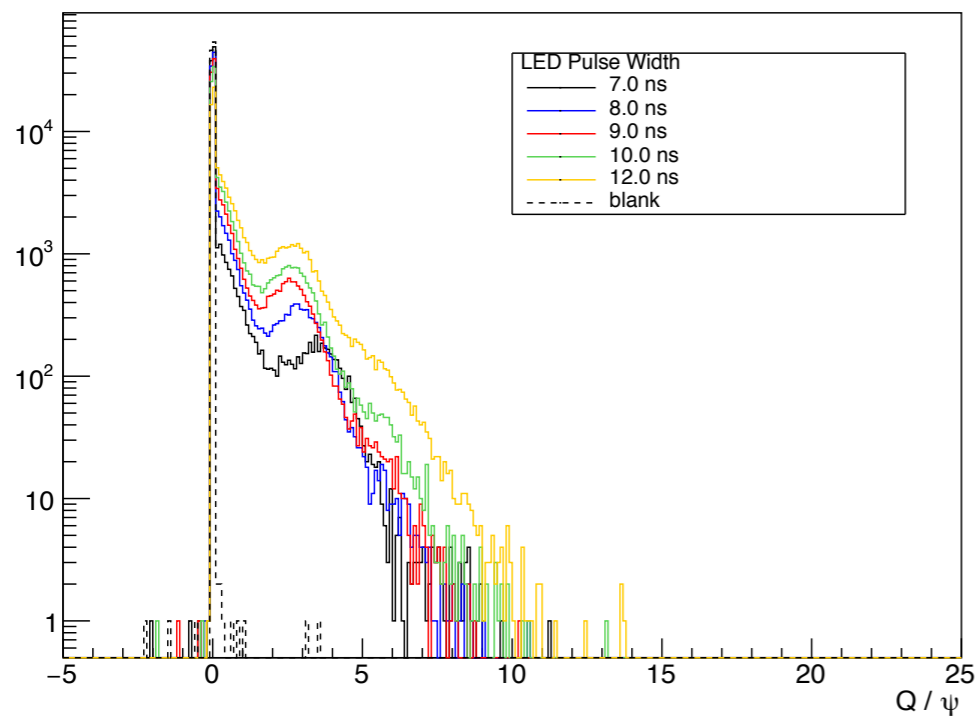


OD5

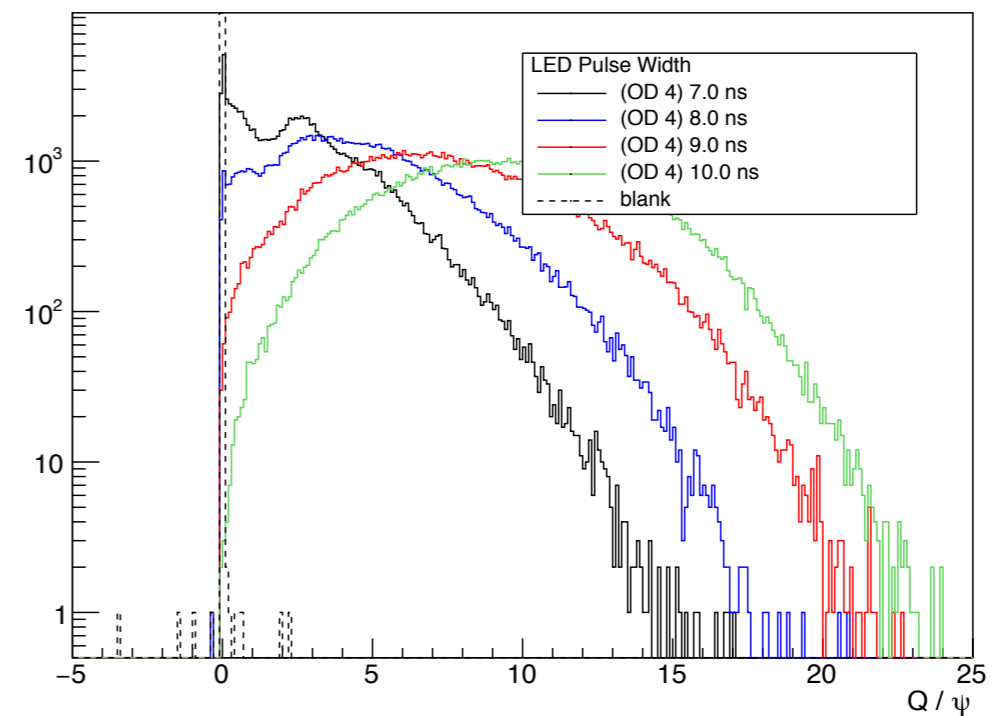


OD4

- Scaling the charge distributions by the calculated single photoelectron charge does not seem to produce noticeable photoelectron peaks
- Charge / (SPE charge) should be  $\sim$  NPE



OD5



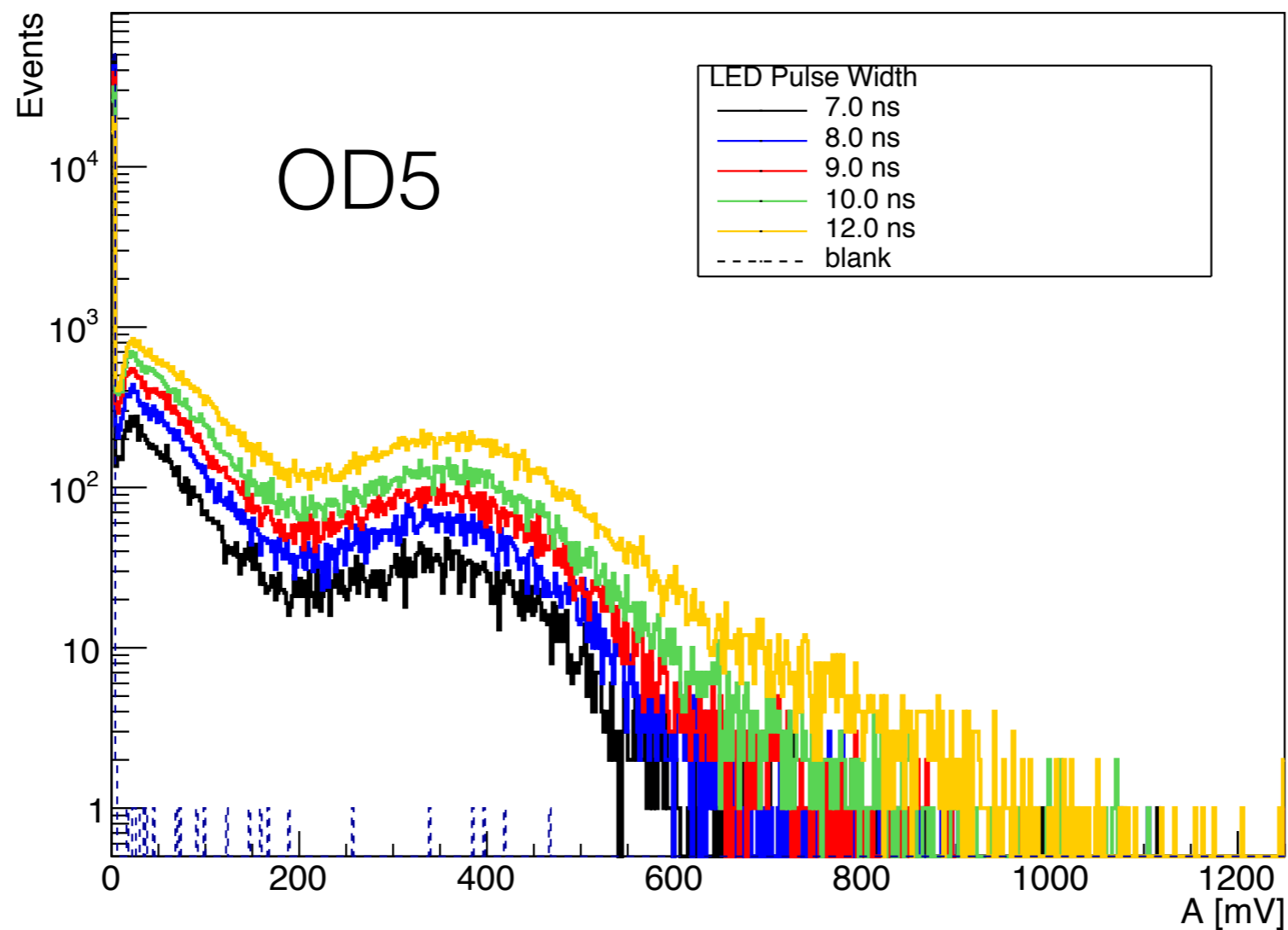
OD4



# Dark Rate from MIA SPE



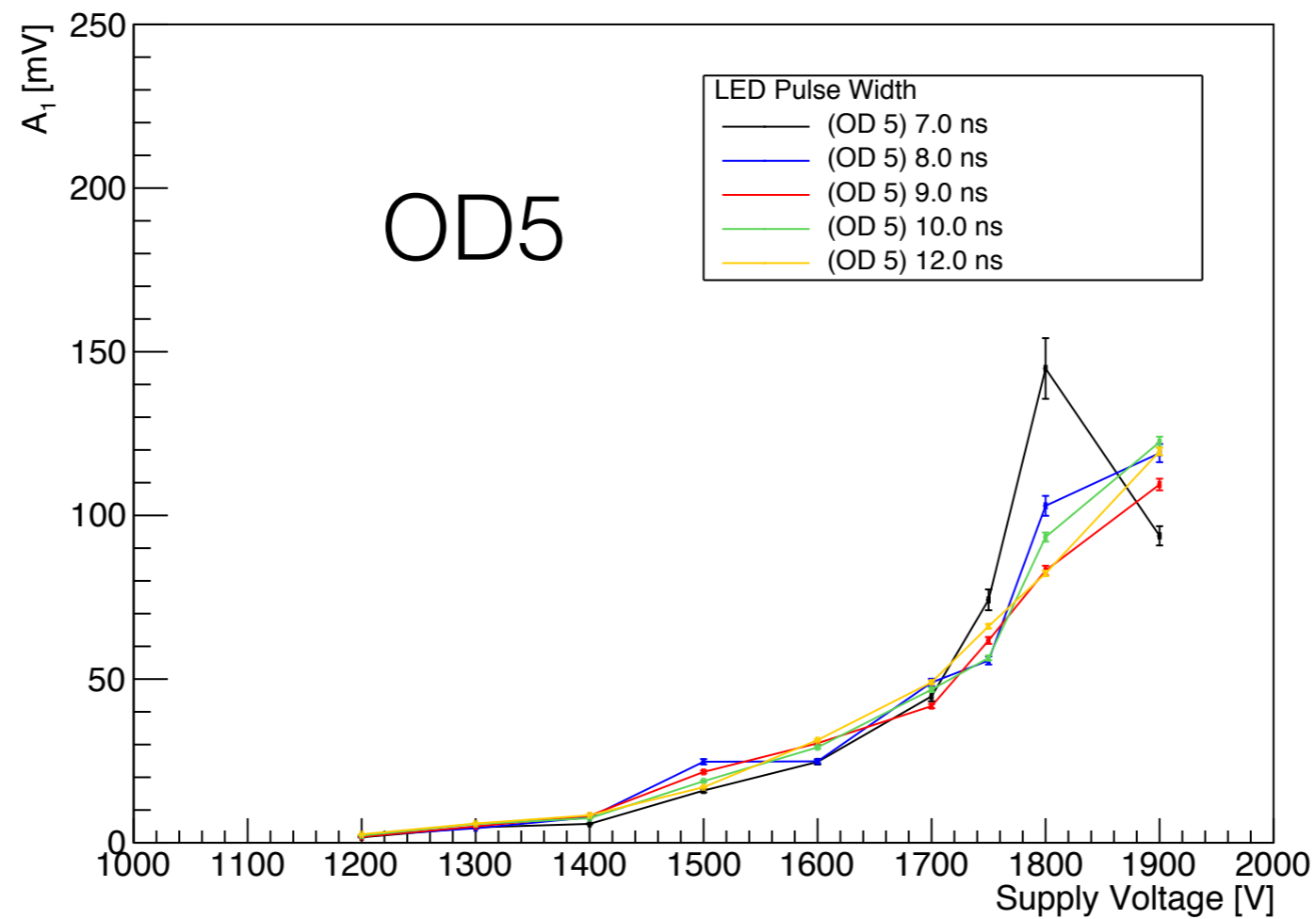
- To tie these results into the dark rate, which results from a trigger threshold in amplitude, can apply this method also to amplitude:



# Dark Rate from MIA SPE



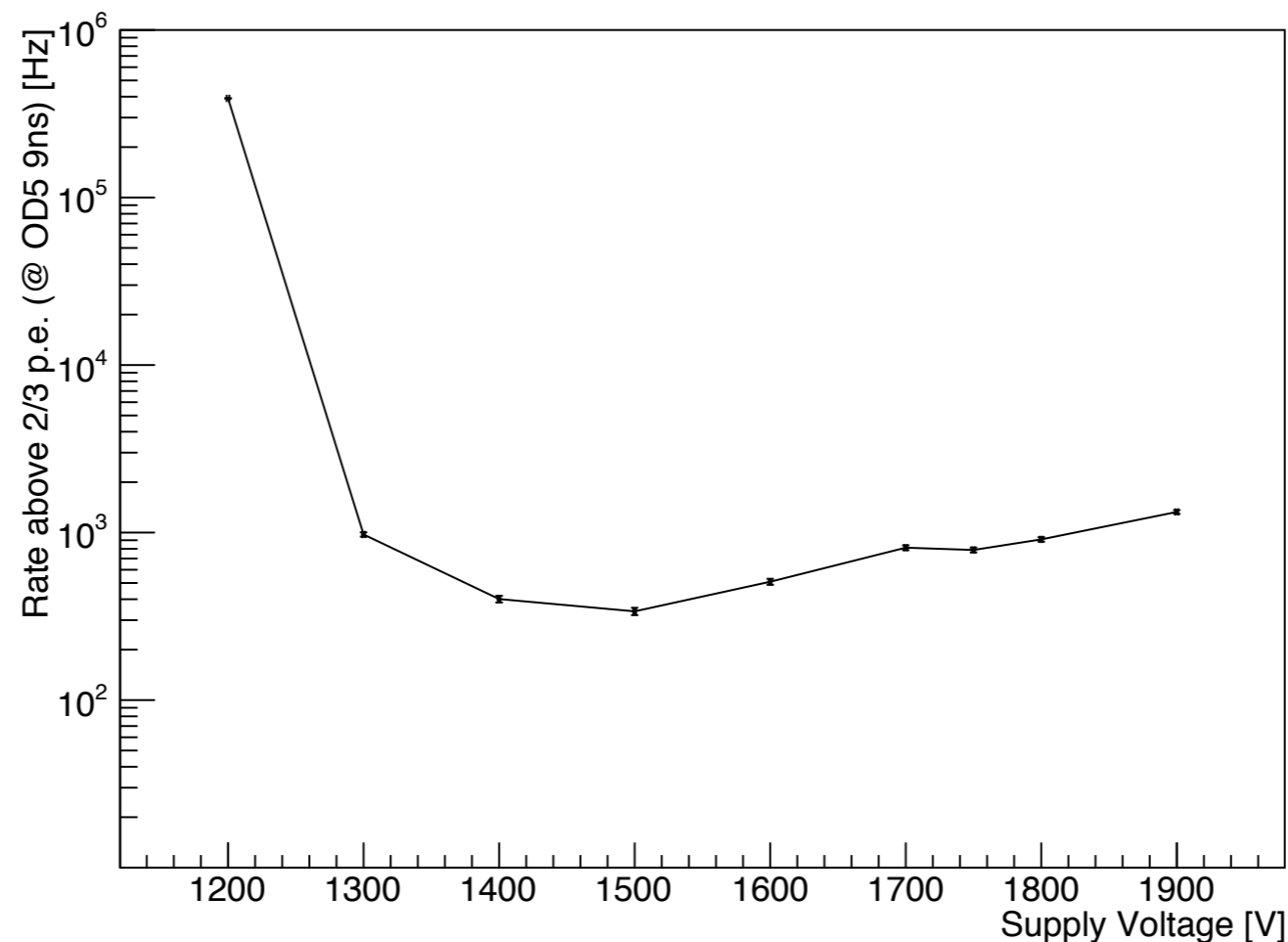
- The results of this method in amplitude less clean in their evolution with voltage, LED width



# Dark Rate from MIA SPE



- Despite this, still attempt to use these expectations to form a trigger threshold
- Using one set of results (OD5 9ns) versus voltage, express the dark rate relative to a threshold 2/3 of these values



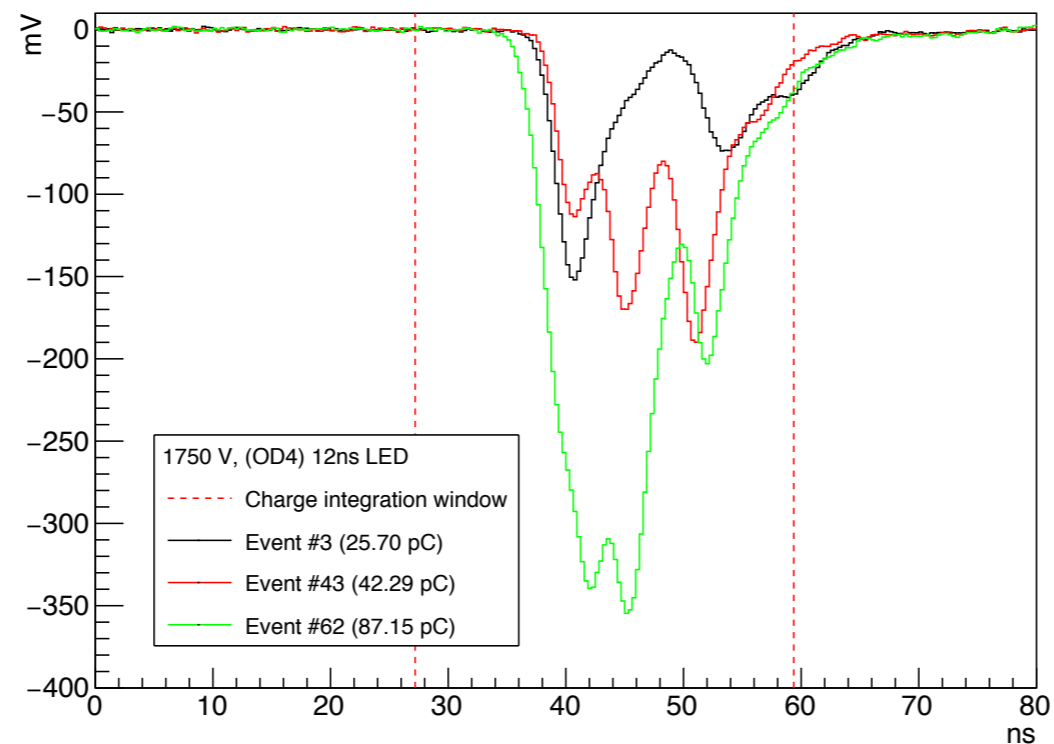
- Considering the MIA likely underestimates the distribution peaks, this is likely an overestimate of the rate at 2/3 PE

# Backup

# Calculating charge with V1743



- At higher gains and larger light intensities, you begin to observe very large-amplitude, long width signals
- For example at 1750 V, a 12ns LED pulse with OD4 gives roughly 5% of events looking as:



- And for even larger signals the dynamic range of the V1743 can become saturated
- For these types of signals the amplitude underestimates the total amount of “signal” present
- It is best to examine both amplitude and charge

# Dark Rate from Peak Amplitude



- Simply just look at the amplitude of the peak

