## Tau Leptons: A tool for studying SM and BSM physics at CMS

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on behalf of the CMS collaboration
CERN LPCC EP-LHC Seminar Series
Geneva, 7/2/2017


## Why taus?

- t's are precious for SM Higgs measurements:
- Theoretical reasons:
- Yukawa couplings are proportional to the mass of the interacting particle
- $\boldsymbol{\tau}$ are massive $\rightarrow$ sizeable BR
- Experimental reasons:
- relatively clean experimental signature, b-quarks couple more strongly to Higgs but are more difficult to identify
$\rightarrow$ measure fermion couplings, and investigate possible deviations
- T's are precious for BSM searches:
- BSM physics manifests as anomalous couplings, e.g. Lepton Flavour Violation
- a zoo of new particles predicted to decay in channels with t's
- often t channels are enhanced, e.g. MSSM
- T's in final states of multiple decays involving $\mathrm{H}(125)$
(some) $\mathbf{\tau}$ analyses at a glance




## Study of the 125 GeV SM Higgs boson

- Higgs coupling measurement
- triple Higgs self coupling $\lambda_{\text {hhh }}$ (for the future)
- measurement of Higgs CP (for the future)


## BSM searches

- t analyses lead MSSM H sensitivity
- Lepton Flavour Violation
in Higgs sector accessible only with T's
- wealth of BSM resonances decaying in final states with t's
- heavy resonance and boosted taus



## Why t's are challenging? - 1

- overwhelming production of jets at hadronic colliders
- need to keep the jet $\rightarrow$ T mis-identification probability as low as possible
- $\rightarrow$ isolation is the main tool to reduce such contamination. It estimates of the activity around the tau: high for jets, low for taus
- electrons and muons too can be erroneously identified as $\tau$ !
- need to cope with diverse sources of contaminations



## Why t's are challenging? - 2

- t decays always involve neutrinos
- not detectable $\rightarrow \mathrm{ET}^{\text {miss }} \rightarrow$ missing kinematic information
- impossible to directly reconstruct di-t invariant masses
$\rightarrow$ need tools to estimate di-t mass from visible products and $E_{T}{ }^{m i s s}$
- $\tau$ triggers need to efficiently select taus in such a busy environment
- they need to fit into tight constraints of
- time spent analysing each event
- number of events that can be saved per unit time
- avoid cutting phase space useful interesting physics
$\rightarrow$ often $\tau$ triggers are multi-object and analysis specific: $\mathscr{\ell} \tau, d i-\tau, \tau+E_{\top}{ }^{m i s s}$


## Outline

- tau object in CMS
- reconstruction and identification
- performance on data
- tau triggers
- tau in CMS analyses:
- Higgs: SM and BSM
- Exotic searches
$\rightarrow \tau_{\mu} \tau_{n}$ VBF candidate

$$
\begin{aligned}
m_{\pi} & =102.0 \mathrm{GeV} \\
m_{i j} & =1.4 \mathrm{TeV} \\
\Delta n_{i j} & =5.8
\end{aligned}
$$

# Tau object at CMS <br> https://cds.cern.ch/record/2196972 

## Tau reconstruction

## The $\tau$ lepton

- $\tau$ is the only lepton heavy enough to decay into hadrons $\mathrm{m}_{\tau}=1.777 \mathrm{GeV}$, lifetime 2.91E-13 s, ct = $90 \mu \mathrm{~m}$
- out of the many possible hadronic tau decays (PDG), we group them into three families or decay modes:
- 1-prong, 1-prong + $n \pi^{0}$, 3-prong
- $h$ can be either $\pi$ or $K$, but predominantly $\pi$ and $m_{\pi}{ }^{ \pm}$hypothesis is always assumed
- $\pi^{0}$ are reconstructed from e and $\gamma$ deposits in rectangular regions in ECal, called strips

| Decay mode | Meson resonance | $\mathcal{B}[\%]$ |
| :--- | :---: | :---: |
| $\tau^{-} \rightarrow \mathrm{e}^{-} \bar{v}_{\mathrm{e}} v_{\tau}$ |  | 17.8 |
| $\tau^{-} \rightarrow \mu^{-} \nu_{\mu} v_{\tau}$ |  | 17.4 |
| $\tau^{-} \rightarrow \mathrm{h}^{-} v_{\tau}$ | 11.5 |  |
| $\tau^{-} \rightarrow \mathrm{h}^{-} \pi^{0} v_{\tau}$ | $\rho(770)$ | 26.0 |
| $\tau^{-} \rightarrow \mathrm{h}^{-} \pi^{0} \pi^{0} v_{\tau}$ | $\mathrm{a}_{1}(1260)$ | 10.8 |
| $\tau^{-} \rightarrow \mathrm{h}^{-} \mathrm{h}^{+} \mathrm{h}^{-} v_{\tau}$ | $\mathrm{a}_{1}(1260)$ | 9.8 |
| $\tau^{-} \rightarrow \mathrm{h}^{-} \mathrm{h}^{+} \mathrm{h}^{-} \pi^{0} v_{\tau}$ | 4.8 |  |
| Other modes with hadrons |  | 1.8 |
| All modes containing hadrons |  | 64.8 |

## Intermission - Particle Flow - 1


make use of all CMS sub detector: redundancy and inter-calibration reconstruct all stable particles in the event: $\mu, \mathrm{e}, \gamma$, neutral/charged hadrons

## Intermission - Particle Flow - 2



PF particles are used to build high-level objects in a consistent way global event description

## $t_{h}$ reconstruction in CMS



## Hadron Plus Strip algorithm (HPS)

- run on Particle Flow inputs:
jets and their charged and neutral constituents
- can identify each $\tau_{n}$ decay mode
- exploits the $\rho(770)$ and $\mathrm{a}_{1}(1220)$ intermediate resonances through mass window requirements

|  | Generated |  |  |
| :--- | :---: | :---: | :---: |
| Reconstructed | $\tau^{-} \rightarrow \mathrm{h}^{-} v_{\tau}$ | $\tau^{-} \rightarrow \mathrm{h}^{-} \geq 1 \pi^{0} v_{\tau}$ | $\tau^{-} \rightarrow \mathrm{h}^{-} \mathrm{h}^{+} \mathrm{h}^{-} v_{\tau}$ |
| $\tau^{-} \rightarrow \mathrm{h}^{-} v_{\tau}$ | $\mathbf{0 . 8 9}$ | 0.16 | 0.01 |
| $\tau^{-} \rightarrow \mathrm{h}^{-} \geq 1 \pi^{0} v_{\tau}$ | 0.11 | $\mathbf{0 . 8 3}$ | 0.02 |
| $\tau^{-} \rightarrow \mathrm{h}^{-} \mathrm{h}^{+} \mathrm{h}^{-} v_{\tau}$ | 0.00 | 0.01 | $\mathbf{0 . 9 7}$ |

## $t_{h}$ reconstruction in CMS



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## Dynamic strip reconstruction

$$
\begin{array}{r}
\boldsymbol{\tau} \rightarrow \boldsymbol{\pi}^{ \pm}+\boldsymbol{\pi}^{0}\left(\boldsymbol{\pi}^{0}\right)+\mathbf{v}_{\boldsymbol{\tau}} \\
57 \% \text { of all had } \\
\text { decaying t's }
\end{array}
$$



## Dynamic strip reconstruction



## Boosted Taus

dedicated $\tau$ reconstruction for the high- $p_{T}$ regime

- e.g. heavy $X \rightarrow h h \rightarrow b b \pi$, where h's are highly boosted and their decay product overlap

Semileptonic


Fully hadronic


- unified approach for both $\boldsymbol{\ell t}$ and $\tau \tau$
- start from a fat jet (cone R=0.8)
- identify subjets ( $\mathrm{p}_{\mathrm{T}}>10 \mathrm{GeV}$ )
- $\ell$ are considered as subjets too
- run standard $\mathbf{t}$ reconstruction on subjets
- isolation computed within the subjet radius to avoid overlaps


## Boosted taus - Performance




- the boosted tau reconstruction significantly improves signal acceptance for high $\mathrm{p}^{\mathrm{H}}>500 \mathrm{GeV}$, especially for $\mathrm{t}_{\mathrm{h}} \mathrm{t}_{\mathrm{h}}$


## di-t mass reconstruction: SVFit algorithm

- SVFit: Maximum likelihood estimator of the di- $\tau$ system mass
- Estimated event-by-event using four-momenta of visible decay products, Ex ${ }^{\text {miss }}, \mathrm{Ey}^{\text {miss }}$, and expected $\mathrm{E}^{\text {miss }}$


$$
\begin{array}{ll}
\text { Phasespace } & \text { Expected } \square_{T} \\
\text { of } \tau \text {-decays } & \text { Resolution }
\end{array}
$$ resolution, $E_{T}$ miss is assumed to be coming only from taus



$Z / H(125)$ separation largely improved, $m_{H}$ resolution $\sim 15 \%$ essential tool for $\mathrm{H} \rightarrow \boldsymbol{\pi}$ analyses

## Tau identification

## Sources of misidentified t's

- quark/gluon initiated jets
- cut- and MVA-based isolation discriminators
- electrons can be misidentified as 1-prong or 1-prong $+(n) \pi^{0} \tau_{h}$
- both electrons and $\pi \pm$ are associated to a track and calorimetric deposits
- they can emit bremsstrahlung and the emerging $\gamma$ (possibly converting back to $\mathrm{e}^{+} \mathrm{e}^{-}$) could be identified as $\pi^{0}$
- multivariate discriminant
- muons can be misidentified as 1-prong taus
- veto discriminants based on the presence of segments in the outer muon detectors
- efficiency $>95 \%$ up to the TeV scale, $\mu \rightarrow \mathrm{T}$ rate $<10^{-4}$



## Cut-based isolation

$$
\begin{gathered}
I_{\tau}=\sum p_{T}^{\text {charged }}\left(d_{Z}<0.2 \mathrm{~cm}\right)+\max \left(0, \sum p_{T}^{\gamma}-\Delta \beta \sum p_{T}^{\text {charged }}\left(d_{Z}>0.2 \mathrm{~cm}\right)\right) \\
p_{T}^{\text {strip, outer }}=\sum p_{T}^{e / \gamma}\left(\Delta R>R_{\mathrm{sig}}\right)<0.10 \cdot p_{T}^{\tau}
\end{gathered}
$$


charged isolation

- tracks compatible with T's vertex
- pile-up robust


## neutral isolation

- pile-up corrected $\gamma$ 's
- neutral pile-up subtraction proportional to charged pile-up through the empirical $\Delta \beta$ factor
strip specific requirement
- on $\Sigma p_{T}$ of the strips far from the signal cone
- $\mathrm{R}_{\text {sig }}$ is defined as $0.05<3.0 / \mathrm{pT}^{\top}<0.1$


## MVA-based isolation

## Boosted Decision Tree discriminator

- training includes all observables used in cut-based isolation plus:
- t lifetime variables impact parameter (transverse and 3D for 3-prong) and its significance
- shape variables weighed $\Delta R, \Delta \phi$ and $\Delta \eta$ between the $e / \gamma$ in strip and the $T_{n}$ direction
- e/ $\boldsymbol{\gamma}$ multiplicities in signal and isolation cones
- training done on a mix of genuine taus from DY, H, Z' and W' and fake taus from QCD and W+Jets processes





## MVA-based isolation expected performance - 1


fakes are reduced by $\sim 2 x$ at equal efficiency

## MVA-based isolation expected performance - 2


efficiency

fake rate

## Anti-electron discriminator



Electrons can easily mimic 1-prong t's

If they emit bremsstrahlung can also be misidentified as 1-prong + ( $n$ ) $\pi^{0} \mathrm{~T}^{\prime} \mathrm{s}$
anti-electron
Boosted Decision Tree discriminator

- based on shape variables, HCal/ECal deposits, bremsstrahlung quantities and e/v multiplicities
- training done on genuine taus from $Z /$ $\gamma^{\star} \rightarrow T T$ and fake taus from $Z / \gamma^{\star} \rightarrow e e$
- medium WP: eff 80\%, FR 3E-3




## Tau identification performance on data L = $2.3 \mathrm{fb}^{-1} @ 13 \mathrm{TeV}, 2015$

## Tau measurement techniques using data

- $\mathbf{Z} \rightarrow \boldsymbol{\pi T}$ process is the standard candle for $\boldsymbol{\tau}$ measurements in data
- different techniques are used
- $Z \rightarrow T_{\mu} \tau_{n}$ Tag\&Probe main method, workhorse
- $\mathrm{Z} \rightarrow \mu \mu / \mathrm{Z} \rightarrow \mathrm{TT} \quad$ orthogonal method, different systematics
- $\mathrm{W}^{\star} \rightarrow$ TV $\quad$ to cover high p t taus phase space



## TaulD efficiency measurement via T\&P





## T\&P on $Z \rightarrow \mu \tau_{h}$ events

- $T_{n}$ isolation passing and failing probes
- two complementary observables:
- visible $\tau_{\mu} \tau_{h}$ mass
- track multiplicity in $T_{n}$
- SFs compatible with 1. within 6\% uncertainty


## TaulD efficiency measurement via $\mathbf{Z} \rightarrow \tau_{\mu} \tau_{\boldsymbol{h}} / \mathbf{Z} \rightarrow \mu \mu$



- similar selections to $\mathbf{Z} \mu \mu$ and $Z_{\tau \tau}$ to improve Phase Space overlap. Cancellation of common systematic uncertainties: complementary to T\&P
- resulting SFs compatible with those obtained via T\&P


## TaulD efficiency measurement at high $\mathrm{p}_{\boldsymbol{T}}$ via $\mathbf{W}^{*} \rightarrow$ Tv



- select highly virtual $\mathbf{W}$ boson ( $\mathrm{m}_{\mathrm{T}} \mathbf{>} \mathbf{2 0 0} \mathbf{G e V}$ ) to enrich the sample in high $\mathrm{p}_{\mathrm{T}}$ taus
- similarly to the previous slide, use both $\mathbf{W}^{*} \rightarrow \mu \mathrm{v}$ and $\mathbf{W}^{*} \rightarrow$ Tv
- resulting SFs $\sim 0.95$ with $15 \%$ uncertainty


# Tau helicity <br> https://cds.cern.ch/record/2216986 

## т helicity measurement

- the $\boldsymbol{\tau}$ polarisation is an interesting observable
- $\sin ^{2} \theta_{w} \propto A_{T}=\left(N^{+}-N^{-}\right) /\left(N^{+}+N^{-}\right)$where $N^{ \pm}$is the number
of $\tau$ leptons with helicity $\pm 1$ coming from the $Z \rightarrow \pi$ process
- Higgs CP properties can be measured through measuring the helicity of $\tau$ 's from the $\mathrm{H} \rightarrow \pi$ process
- tau spin can can be accessed through different polarisation-sensitive observables:

- angular distributions of the $\boldsymbol{\tau}$ decay products in $\boldsymbol{\tau}^{ \pm} \rightarrow \mathbf{a}_{1}{ }^{ \pm} \mathbf{V}_{\boldsymbol{\tau}} \rightarrow \boldsymbol{\pi}^{ \pm} \boldsymbol{\pi}^{ \pm} \boldsymbol{\pi}^{\mp} \mathbf{V}_{\boldsymbol{\tau}}$
- energy asymmetry $E\left(\pi^{ \pm}\right)-E\left(\pi^{0}\right) / E\left(\pi^{ \pm}\right)+E\left(\pi^{0}\right)$ in $\tau^{ \pm} \rightarrow \rho^{ \pm} v_{\tau} \rightarrow \pi^{ \pm} \pi^{0} v_{\tau}$


## $\tau$ helicity measurement - $\mathbf{\tau}^{ \pm} \rightarrow \mathbf{a}_{1}{ }^{ \pm} \mathbf{v}_{\boldsymbol{\tau}} \rightarrow \boldsymbol{\pi}^{ \pm} \boldsymbol{\pi}^{ \pm} \boldsymbol{\pi}^{\mp} \mathbf{v}_{\boldsymbol{\tau}}$



$$
\omega_{a_{1}}=\frac{\left|M_{+}\left(\theta^{*}, \beta, \gamma\right)\right|^{2}-\left|M_{-}\left(\theta^{*}, \beta, \gamma\right)\right|^{2}}{\left|M_{+}\left(\theta^{*}, \beta, \gamma\right)\right|^{2}+\left|M_{-}\left(\theta^{*}, \beta, \gamma\right)\right|^{2}}
$$

optimal 1D observable $\propto A_{T}$

## $\tau$ helicity measurement $-\tau^{ \pm} \rightarrow \rho^{ \pm} \mathbf{v}_{\boldsymbol{\tau}} \rightarrow \boldsymbol{\pi}^{ \pm} \boldsymbol{\pi}^{0} \mathbf{v}_{\boldsymbol{\tau}}$




$$
\cos \psi^{*}=\frac{m_{\rho}}{\sqrt{m_{\rho}^{2}-4 m_{\pi}^{2}}} \frac{E_{\pi^{*}}-E_{\pi^{0}}}{\vec{P}_{\pi}+\vec{P}_{\pi^{0}}}
$$

- $\mathrm{E}_{\mathrm{ch}}-\mathrm{E}_{\mathrm{n}}$ asymmetry $\alpha \cos \boldsymbol{\psi}^{*}$ and polarisation-sensitive
- detector effects (i.e. lower efficiency for soft $\pi^{0}$ ) only marginally smear the distribution shown here for the generator level


## т helicity measurement - validation on data

$$
\mathbf{\tau}^{ \pm} \rightarrow \mathbf{a}_{1}{ }^{ \pm} \mathbf{v}_{\boldsymbol{\tau}} \rightarrow \boldsymbol{\pi}^{ \pm} \boldsymbol{\pi}^{ \pm} \boldsymbol{\pi}^{\mp} \mathbf{v}_{\boldsymbol{\tau}}
$$



$$
\mathbf{T}^{ \pm} \rightarrow \boldsymbol{\rho}^{ \pm} \mathbf{V}_{\mathbf{T}} \rightarrow \boldsymbol{\pi}^{ \pm} \boldsymbol{\pi}^{\mathbf{0}} \mathbf{v}_{\mathbf{T}}
$$



- in both measurements the measured polarisation values are well compatible with MC prediction
- indication of robustness of PF-based $\tau$ reconstruction algorithm


## Tau triggers

## $\tau$ triggers

- the CMS trigger system comprises two distinct subsystems:
- Level-1: hardware based, fast but coarse, $40 \mathrm{MHz} \rightarrow 100 \mathrm{kHz}$
- High Level Trigger: software based, sophisticated but slower, $100 \mathrm{kHz} \rightarrow 1 \mathrm{kHz}$
- the goal is to maximise the signal efficiency at minimum cost in terms of rate and CPU time
- taus are reconstructed at both levels and a variety of triggers are used by the analyses, often as multi object triggers

Tau triggers in 2016

| Channel | typical trigger selection | used by |
| :---: | :---: | :---: |
| $\mu \mathrm{t}$ | iso $\mu 19 \mathrm{GeV}$, iso t 20 GeV | SM \& MSSM H $\rightarrow \tau \tau$, hh $\rightarrow$ bbtr, $\mathrm{Z}^{\prime}$ |
| et | iso e 22 GeV , iso ¢ 29 GeV | SM \& MSSM H $\rightarrow \tau \tau$, hh $\rightarrow$ bbtt, $\mathrm{Z}^{\prime}$ |
| $\tau$ | double iso ¢ 35 GeV | SM \& MSSM $H \rightarrow \tau \tau$, hh $\rightarrow$ bbt , $Z^{\prime}$ |
| $\boldsymbol{\tau}+\mathrm{E}_{T}^{\text {miss }}$ | iso t 50 GeV , ETmiss 100 GeV | $H^{ \pm} \rightarrow \tau \nu, W^{\prime}$ |
| T | iso $\tau 140 \mathrm{GeV}$ | $H^{ \pm} \rightarrow \tau v, W^{\prime}$ |

## Upgraded CMS Level 1 trigger

- in 2016 the upgrade of CMS L1 trigger system has been completed
- the calorimetric system can now read out at much finer granularity
- L1 taus are purely calorimetric objects and vastly profit from the upgraded system



## t's at Level-1

## $\tau$ identification

- build clusters around local maxima
- template shapes
- contiguous clusters can be merged
- gather all the $\pi^{0 / \pm}$ from t decay


## $\tau$ isolation

- compute energy in a 6x9 пф region


## pile-up subtraction

- estimate average PU by counting \# of trigger towers with $\mathrm{E}_{\mathrm{T}}>0 \mathrm{GeV}$






## T's at HLT

- based on ParticleFlow@HLT (simplified tracking)
- streamlined reconstruction:
- $\pi^{ \pm}$and $\pi^{0} \rightarrow \gamma \gamma$ are built as in offline, but decay modes not explicitly enforced
- main goal is to minimise CPU consumption preserving efficiency
- cut-based isolation



## Tau in CMS analyses

## CMS analyses using taus

- $\tau_{h}$ are used in several realms, both in SM and BSM scenarios
- couplings and search of new resonances
- I will discuss only a subset of representative analyses
- $\mathrm{SM} H \rightarrow \mathrm{TT}$ : couplings and properties
- MSSM $\phi \rightarrow$ тт @13 TeV
- $\mathrm{H}^{ \pm} \rightarrow \mathrm{TV}$ @13 TeV
- $X \rightarrow h_{125} h_{125} \rightarrow T T$ bb @13 TeV
- $\operatorname{LFV~H} \rightarrow \mu \mathrm{t}_{(e, h)}$
- all CMS public results are collected here
- many of the analyses presented in the following are being updated for coming Moriond17... stay tuned!


## SM H $\boldsymbol{T} \boldsymbol{T \tau}$

## $H \rightarrow \tau$ analysis motivations

## Boson with mass 125 GeV discovered by ATLAS \& CMS

is this the SM Higgs boson? $\Rightarrow$ measure the couplings



- $\mathrm{H} \rightarrow \mathrm{Tt}$ main probe to test Higgs Yukawa coupling to fermions
- $\tau$ is heavy: sizeable $B R(H \rightarrow \pi)=6.3 \%$ at $m_{H}=125 \mathrm{GeV}$
- cleaner experimental signature than $\mathrm{H} \rightarrow \mathrm{bb}$


## SM H $\rightarrow$ tr analysis

- complex analysis:
- (6 di-т final states) $\times(\mathrm{VBF}, \mathrm{VH}, \mathrm{ggF}) \times$ (categories)






- relies on reconstructing $\mathrm{m}_{\boldsymbol{\pi}}$ using SVfit
- Run2 analysis will be published soon:
- Run1 sensitivity surpassed only with the full 2016 statistics
- $8 \rightarrow 13 \mathrm{TeV}$ made BSM searches more interesting with the first data


## $H \rightarrow \pi \tau$ as a probe to fermion couplings


$\mathbf{H} \boldsymbol{\rightarrow} \boldsymbol{\tau \tau}$ very sensitive
channel to $\mathbf{k}_{\mathbf{f}}$ $k_{f}$ measured with $30 \%$ precision

sensitivity to kv from VBF and VH

## $H \rightarrow \tau \tau$ in the big Higgs picture

J. High Energy Phys. 08 (2016) 045




- ATLAS + CMS: $5 \sigma \mathbf{H} \rightarrow \pi \tau$ observation!
- provides strong constraints on:
- fermion coupling modifier $\boldsymbol{\kappa}$ F
serves both as measurement and as search through deviations
- VBF H cross section


## MSSM A/H and $\mathbf{H}^{ \pm}$

## $H / A \rightarrow \tau$ and $\mathrm{H}^{ \pm} \rightarrow \tau v$ as a probe for MSSM

the Minimal Supersymmetric Standard Model is the simplest extension to the SM including SUSY partners of the SM particles

- 5 Higgs bosons
- 2 charged $\mathbf{H}^{ \pm}$
- 3 neutral (collectively labelled as $\phi$ )
- 1 light $\mathbf{h}$ (SM-like)
- 2 heavy A/H (CP-odd/even)
- two parameters describe the model at the tree level
- $\mathbf{M A}_{\mathbf{A}}$ mass of the heavy CP-odd Higgs
- $\boldsymbol{\operatorname { t a n }} \boldsymbol{\beta}$ related to ratio of the couplings to up/down-type fermions
for moderate-to-large values of $\tan \beta$, the couplings to $\tau$ 's are greatly enhanced


## MSSM H/A $\rightarrow$ ti - strategy

- four final states considered $\boldsymbol{\mu \tau}, \mathbf{e t}, \boldsymbol{\tau \tau}$ and $\mathbf{e} \boldsymbol{\mu}$
- two categories based on the presence of b-jets to address the two dominant production modes

- fit to the total transverse mass distribution $\mathrm{m}^{\text {tot }}$
$m_{\mathrm{T}}^{\text {tot }}=\sqrt{m_{\mathrm{T}}\left(E_{\mathrm{T}}^{\text {miss }}, \tau_{1}^{\text {vis }}\right)^{2}+m_{\mathrm{T}}\left(E_{\mathrm{T}}^{\text {miss }}, \tau_{2}^{\text {vis }}\right)^{2}+m_{\mathrm{T}}\left(\tau_{1}^{\text {vis }}, \tau_{2}^{\text {vis }}\right)^{2}}$


## MSSM H/A $\rightarrow$ TT - model independent results




- $\tau_{h} \tau_{h}$ channel is the most sensitive especially for $m_{A}>200 \mathrm{GeV}$ every improvement on $\tau$ reconstruction/ID impacts here directly


## MSSM H/A $\rightarrow \boldsymbol{T \tau}$ - model dependent results




- exclusion contour in the MSSM vs SM hypothesis test
- already surpassed Run-1 performance at high mass thanks to $8 \rightarrow 13 \mathrm{TeV}$ and both analysis and $\tau_{n}$ improvements


## $\mathbf{H}^{ \pm} \rightarrow$ Tv - strategy

- two different scenarios:

| heavy charged Higgs: | $\mathbf{m}_{\mathbf{H}}>\mathbf{m}_{\mathbf{t}}-\mathbf{m}_{\mathbf{b}}$ |
| :--- | :--- |
| light charged Higgs: | $\mathbf{m}_{\mathbf{H} \pm}<\mathbf{m}_{\mathbf{t}}-\mathbf{m}_{\mathbf{b}}$ |

sensitivity dominated
by $\mathrm{H}^{ \pm} \rightarrow \mathrm{tb}$
sensitivity dominated by $\mathrm{H}^{ \pm} \rightarrow \mathrm{TV}$

heavy $\mathrm{H}^{ \pm}$


- signature: $1 \mathrm{~T}_{\mathrm{h}}, \geq 3$ jets, $\geq 1 \mathrm{~b}$-jets, $\mathrm{E}_{\mathrm{T}}^{\text {miss }}$ and $\mathrm{O} \ell$
- different kinematic selections for the two scenarios
- topological selections: $\Delta \phi\left(\mathrm{T}_{\mathrm{h}}, \mathrm{E}_{T}{ }^{\text {miss }}\right)$ vs $\Delta \phi\left(j e t, \mathrm{E}_{T}{ }^{\text {miss }}\right)$
- signal extraction through a fit to $\mathrm{m}_{\mathbf{T}}$ distribution



## $H^{ \pm} \rightarrow$ tv - results




- light $\mathbf{H}^{ \pm}$almost ruled out in MSSM $\mathbf{m}_{\mathbf{h}}{ }^{\text {mod+ }}$ (and most of other) scenario(s)


## Summary of MSSM analyses - 8 TeV


$\mathrm{H} \rightarrow \boldsymbol{\tau}$ dominates the exclusion in the large $\tan \beta$ region
$\mathrm{H}^{ \pm} \rightarrow$ Tv rules out the low mass region
$t$ analyses driving force of MSSM H bosons searches

## Lepton Flavour Violation

## Search for Lepton Flavour Violating $\mathbf{H} \boldsymbol{\rightarrow} \boldsymbol{\mu} \boldsymbol{\tau}_{(\mathrm{e}, \mathrm{h})}$

- LFV Higgs decays are prohibited in the SM, but can arise in a vast number of BSM models
- 2HDM, SUSY, composite Higgs, Randall-Sudrum, ...
- why taus? taus are the heaviest leptons and couple favourably to the Higgs
- analysis strategy:
- similar to SM H $\boldsymbol{H} \boldsymbol{T r}$, but $\boldsymbol{\mu}$ has higher $\mathrm{p}_{\mathrm{T}}$ being prompt
- events are sorted in categories based on the number of jets
- signal is extracted through a ML fit to the collinear mass distribution
- Mcoll is bult assuming that neutrino(s) arising from the t decay are highly boosted and their direction can be approximated to that of the visible products of the $\tau$


## Search for Lepton Flavour Violating $\mathbf{H} \boldsymbol{\rightarrow} \boldsymbol{\mu} \boldsymbol{\tau}_{(\mathrm{e}, \mathrm{h})}$



- observed $B(H \rightarrow \mu \tau)<1.2 \%$ : most stringent limit on LFV Yukawa coupling
- 13 TeV analysis does not confirm the $\mathbf{2 . 4 \sigma}$ excess observed in the 8 TeV analysis, but the Run2 analysis not as sensitive as Run1 yet


## Resonant (and non-resonant) $\mathrm{X}(\mathrm{h}) \rightarrow \mathrm{hh} \rightarrow \mathrm{bb} \boldsymbol{T} \boldsymbol{T}$

## Motivations for $\mathbf{X} / \mathrm{h} \rightarrow \mathrm{hh} \rightarrow$ bbrt

- non resonant:
- predicted SM process, access to Higgs trilinear self coupling $\lambda_{\text {hhh }}$ (Higgs potential) (beyond reach in Run2, need HL-LHC)
- BSM contributions can modify the coupling: $\kappa_{\lambda}=\lambda_{\text {hhh }} / \lambda_{\text {hhh }}{ }^{\mathrm{SM}}$
- resonant:
- in several models, heavy particles can decay in hh, e.g.: Radion, Graviton, MSSM H
- the bbtr final state is chosen because it has good BR 7.3\% and sufficiently clean experimental signature
- $\mu \mathrm{T}$, et, TT channels considered



## $\mathrm{X} / \mathrm{h} \rightarrow \mathrm{hh} \rightarrow \mathrm{bb} \pi$ - strategy

- strategy largely common for the two analyses
- events with a good di-т pair and a good b-jet pair
- $m_{\pi T}$ and $m_{b b}$ mass windows around 125 GeV
- BDT discriminator to reduce ttbar (angular variables, NR only)
- categorisation:
- 2 b-jet resolved


HH (SM) - bb $\mu \tau_{h}$ channel


- 1 b-jet resolved + 1 jet
- boosted b-jets (relevant for $m x>600 \mathrm{GeV}$ )
- signal extraction through a fit to the 4-body $\mathrm{m}_{\text {тtbb }}$ invariant mass
- 4-body mass computed using a kinematic fit imposing $m_{b b}=m_{\pi}=125 \mathrm{GeV}$


## Non-resonant $\mathbf{h} \rightarrow \mathbf{h h} \rightarrow \mathbf{b b} \pi$ - results

## https://cds.cern.ch/record/2204934

$12.9 \mathrm{fb}^{-1}(13 \mathrm{TeV})$


- exclusion limit on $\sigma \times$ BR as a function of the coupling modifier $\mathbf{k}_{\boldsymbol{\lambda}}$


## Resonant $\mathrm{X} \rightarrow \mathrm{hh} \rightarrow \mathrm{bb} t \mathrm{r}$ - results

https://cds.cern.ch/record/2204936
$12.9 \mathrm{fb}^{-1}(13 \mathrm{TeV})$


- exclusion limit on $\sigma \times B R$ as a function of the mass of the resonance


## High mass $X \rightarrow h h \rightarrow$ bbtr with boosted taus

- extend the search to masses up to 2.5 TeV
- the two h125 from the heavy resonance are boosted and their decay products very close to each other $\rightarrow$ boosted b's and t's



## Exotic searches

## W' $^{\prime} \boldsymbol{T}$ Tv and Z' $\rightarrow$ Tr searches

- heavy gauge bosons are foreseen in several BSM models, e.g. Sequential Standard Model (SSM)
- similar signature and properties to SM W \& Z


- model dependent limits exclude $\mathrm{Mw}^{\prime}<3 \mathrm{TeV}, \mathrm{Mz}$ < 2.1 TeV (SSM), Mz' < 1.7 TeV (TAT)


## Heavy neutrinos and $3^{\text {rd }}$ generation leptoquark

- neutrino oscillations $\rightarrow \mathrm{m}_{\mathrm{v}} \boldsymbol{>} \boldsymbol{0} \boldsymbol{\rightarrow}$ seesaw mechanism $\rightarrow$ right-handed neutrinos
- can be accommodated in Left Right Symmetric Extensions (LRSE) of the SM, which predict the existence of heavy gauge bosons $\mathbf{W}^{ \pm} \mathbf{R}$ and $Z^{\prime}$
- typical process in $\mathbf{W}^{ \pm} \mathbf{R} \rightarrow \mathbf{\tau} \mathbf{N}_{\boldsymbol{\tau}} \rightarrow \mathbf{\tau} \mathbf{T} \mathbf{v}_{\mathbf{\tau}}$, where $N_{\boldsymbol{T}}$ is the heavy neutrino
- Lepto-Quark (LQ) models also foresee signatures with two t's

- under model dependent assumptions, limits are set on $m_{W R}, m_{L Q}$ and $m_{N_{\tau}}$


## Outlook to CMS upgrades impact on taus

## CMS detector upgrades towards HL-LHC

- the name of the game: higher luminosity $\longleftrightarrow$ higher PU (up to 200!)
- 2017 phase-1: $4^{\text {th }}$ pixel barrel layer and $3^{\text {rd }}$ endcap disk
- better track/vertex resolution
- improvement in tau reconstruction and performance, e.g. lifetime variables
- HL-LHC phase-2:
- $4 x$ tracker granularity
- High Granularity Calorimeter
- tracking at L1 trigger
- more L1 bandwidth
- muon system up to $|n| \sim 3$



## Projections to HL-LHC, 300-3000 fb-1



- precision measurement of the SM Higgs couplings
- at HL-LHC will reach sensitivity to $\mathbf{S M} \mathbf{h} \rightarrow \mathbf{h h}$
- push the exclusion limits in the MSSM scenario


## Conclusions

# $\tau_{h}$ is a fundamental tool for a broad physics programme 

look for new results at Moriond17 with full Run2 stats

## $\tau_{h}$ reconstruction vastly improved in Run2

established algorithms, triggers, performance

CMS upgrades will ensure this programme to continue

even in extremely harsh PU conditions

## Backup

## The Compact Muon Solenoid detector

CMS DETECTOR


## Particle signatures in CMS



## Anti-electron discriminator expected performance


efficiency

fake rate

## jet $\rightarrow \tau_{h}$ rate measurement




- data/MC comparison shown form in $\mathbf{W} \rightarrow \mu \mathrm{v}+$ Jets enriched regions
- mis-ID probability strongly depends on parton flavour and on jet- $\tau_{h}$ charge mismatch
- measurements in other regions are performed too
- the data/MC disagreement is understood to be due ultimately to the MC hadronisation tune


## Tau energy scale measurement



- $\boldsymbol{m}_{\text {vis }}(\boldsymbol{\mu} \boldsymbol{\tau})$ and $\tau_{\boldsymbol{h}}$ mass (not shown here) distributions are sensitive to $\tau_{E S}$
- produce $Z \rightarrow T_{\mu} T_{n}$ templates in the $\mathbf{- 6 \% + 6 \%}$ energy scale range
- maximum likelihood fit to data with tes (per decay mode) as POI
- results range from $-1.5 \%$ to $+1.5 \%$ depending on the decay mode


## t charge misidentification measurement






- selection of $Z \tau_{\mu} T_{n}$ enriched $O S$ events and the corresponding $S S$ sideband
- simultaneous ML fit to extract the charge misID probability
- upper limit PCF $=0.22 \%$


## Embedded sample



- hybrid data + MC events
- Z kinematics, Jets, $\mathrm{E}_{\mathrm{T}}{ }^{\text {miss }}$ underlying event from data
- better modelling and small (if not absent) uncertainties
- only tau decay is left to the simulation


## $\mathbf{H}^{ \pm \rightarrow T v}$ - final distributions




## $\mathbf{H}^{ \pm} \rightarrow \mathrm{Tv}$ - model independent results




- Low Mass: model is the $\mathrm{BR} \quad \mathcal{B}\left(\mathrm{t} \rightarrow \mathrm{bH} \mathrm{H}^{ \pm}\right) \times \mathcal{B}\left(\mathrm{H}^{ \pm} \rightarrow \tau \nu\right)$
- High mass: $\sigma\left(\mathrm{pp} \rightarrow \mathrm{H}^{ \pm} \mathrm{W}^{\mp} \mathrm{b} \overline{\mathrm{b}}\right) \times \mathcal{B}\left(\mathrm{H}^{ \pm} \rightarrow \tau v\right)$


## $\mathbf{H}^{ \pm} \rightarrow$ Tv - results in MSSM $\mathbf{m}^{\text {mod }+}$ scenario



- light $\mathbf{H}^{ \pm}$almost ruled out in this (and most of other) scenario(s)


## Non-resonant $\mathbf{h} \rightarrow \mathbf{h h}$ - theory motivation

- in the SM this arises from the interference of these two processes

- where $\lambda_{\text {hhh }}$ indicates the trilinear Higgs coupling which is expected to be too small to be measured at LHC in Run2
- BSM contributions can be present and are represented by the coupling modifier $\mathbf{k}_{\boldsymbol{\lambda}}=\boldsymbol{\lambda}_{\text {hhh }} / \lambda_{\text {hhh }}{ }^{\text {SM }}$ that affects both the cross section and the kinematics of the hh process




## Resonant MSSM H $\rightarrow$ hh $\rightarrow$ bbtr - 8 TeV



- the results of the 8 TeV analysis were interpreted in the MSSM (together with $\mathrm{A} \rightarrow \mathrm{Zh} \rightarrow \ell \ell \mathrm{T})$, low $\tan \beta$ and 2 HDM type-II models, further reducing the non excluded parameter space
- update at 13 TeV foreseen for Moriond17

