Portoroz, 19.IV. 2017

Extensions of the IDM



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In coll. with I. Ginzburg, K. Kanishev, D.Sokołowska, B. Świeżewska, G. Gil, P.Chankowski, N. Darvishi, A. Ilnicka, T. Robens, L. Diaz-Cruz, C. Bonilla

Higgs particle at LHC -summer 2016ATLAS+CMS Run 1arXiv:1606.02266v1 [hep-ex]

SM-like scenario observed

- Mass 125.09 \pm 0.24 GeV ZZ \rightarrow 4 I, $\gamma \gamma$
- Total width < 23 MeV (95%CL); SM ~4 MeV</p>
- Signal strengths $\mu = R = \sigma \times Br/(\sigma \times Br)|_{SM}$; SM =1 global 1.09 ± 0.11/0.10 $\gamma\gamma$ 1.14 ± 0.19/0.18 $\rightarrow R_{\gamma\gamma}$
- Invisible decay BR = 0.00^{+0.16} (< 0.32 at 95% CL)
 Spin/CP J^{CP} 0 +



LHC 2016



LHC 2016







DM: Z₂ 2HDM potential for 2HDM Branco, Rebelo ,85 (CP conserved) Potential V = $\frac{1}{2}\lambda_{1}(\Phi_{1}^{\dagger}\Phi_{1})^{2} + \frac{1}{2}\lambda_{2}(\Phi_{2}^{\dagger}\Phi_{2})^{2} - \frac{1}{2}m^{2}_{11}(\Phi_{1}^{\dagger}\Phi_{1}) - \frac{1}{2}m^{2}_{22}(\Phi^{\dagger}\Phi_{2})$ $+ \lambda_{3}(\Phi_{1}^{\dagger}\Phi_{1})(\Phi_{2}^{\dagger}\Phi_{2}) + \lambda_{4}(\Phi_{1}^{\dagger}\Phi_{2})(\Phi_{2}^{\dagger}\Phi_{1}) + \frac{1}{2}[\lambda_{5}(\Phi_{1}^{\dagger}\Phi_{2})^{2} + h.c]$ $\lambda_{345} = \lambda_3 + \lambda_4 + \lambda_5$ Z_2 symmetry transf.: $\Phi_1 \rightarrow \Phi_1 \quad \Phi_2 \rightarrow - \Phi_2$ Yukawa interaction **Model I** – one doublet Φ_1 couples to all fermions

Vacuum state ? various possible M. Krawczyk, Portoroz 2017 **positivity (stability) constraints** $\lambda_1 > 0, \quad \lambda_2 > 0, \quad R+1 > 0, \quad R_3+1 > 0$

$$\lambda_{345} = \lambda_3 + \lambda_4 + \lambda_5, \quad R = \lambda_{345} / \sqrt{\lambda_1 \lambda_2}, \quad R_3 = \lambda_3 / \sqrt{\lambda_1 \lambda_2},$$

Extrema \rightarrow **Vacua** $\langle \phi_S \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_S \end{pmatrix}, \quad \langle \phi_D \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} u \\ v_D \end{pmatrix}$

Symmetry
EWs:
$$v_D = 0$$
, $v_S = 0$, $\mathcal{E}_{EWs} = 0$;
 I_1 , $v_D = 0$, $v_S^2 = v^2 = \frac{m_{11}^2}{\lambda_1}$, $\mathcal{E}_{I_1} = -\frac{m_{11}^4}{8\lambda_1}$
Inerverting
 I_2 : $v_S = 0$, $v_D^2 = v^2 = \frac{m_{22}^2}{\lambda_2}$, $\mathcal{E}_{I_2} = -\frac{m_{22}^4}{8\lambda_2}$
 $w_S^2 = \frac{m_{11}^2\lambda_2 - \lambda_{345}m_{22}^2}{\lambda_1\lambda_2 - \lambda_{345}^2}$, $v_D^2 = \frac{m_{22}^2\lambda_1 - \lambda_{345}m_{11}^2}{\lambda_1\lambda_2 - \lambda_{345}^2}$;
M:
 $\mathcal{E}_M = -\frac{m_{11}^4\lambda_2 - 2\lambda_{345}m_{11}^2m_{22}^2 + m_{22}^4\lambda_1}{8(\lambda_1\lambda_2 - \lambda_{345}^2)}$.
 $\mathcal{E}_{I_1} - \mathcal{E}_M = \frac{(m_{11}^2\lambda_{345} - m_{22}^2\lambda_1)^2}{8\lambda_1^2\lambda_2(1 - R^2)}$
 $v_S^2 = \frac{m_{11}^2\lambda_2 - \lambda_3m_{22}^2}{\lambda_1\lambda_2 - \lambda_3^2}$, $v_D = 0$, $u^2 = \frac{m_{22}^2\lambda_1 - \lambda_3m_{11}^2}{\lambda_1\lambda_2 - \lambda_3^2}$,
 \mathcal{CB} :
 $\mathcal{E}_{CB} = -\frac{m_{11}^4\lambda_2 - 2\lambda_3m_{11}^2m_{22}^2 + m_{22}^4\lambda_1}{8(\lambda_1\lambda_2 - \lambda_3^2)}$.
 $U=0$
 $u \neq 0$
 $u \neq 0$
 $\mathcal{L} = 0$

Inert Doublet Model

 Φ_{s} as in SM (BEH)



Higgs boson h (SM-like)

 $\Phi_{\mathbf{D}}$ – no vev

$$\Phi_{\rm D} = \begin{pmatrix} {\rm H}^+ \\ {\rm H} + {\rm i} \ {\rm A} \end{pmatrix}$$
 (no Higgses!)

Ma,...'78

4 scalars H+,H-,H, A no interaction with fermions

D symmetry
$$\Phi_{s} \rightarrow \Phi_{s} \quad \Phi_{D} \rightarrow \Phi_{D}$$
 exact
 \blacktriangleright D parity
 \lnot only Φ_{D} has odd D-parity
 \lnot only Φ_{D} has odd D-parity
 \lnot the lightest scalar stable - DM candidate (H)
 \lnot (Φ_{D} dark doublet with dark scalars)

IDM: An Archetype for Dark Matter, Lopez Honorez,...Tytgat..07 LHC phenomenology (Barbieri., Ma.. 2006,...)

Ma'2006, .Barbieri 2006, Dolle, Su, **Testing IDM** Gorczyca(Świeżewska), MSc T2011,... Theoretical constraints: Posch 2011, Arhrib..2012, Chang, Stal. vacuum stability, pert.unitarity $\frac{m_{11}^2}{\sqrt{\lambda_1}} \ge \frac{m_{22}^2}{\sqrt{\lambda_2}}$ *condition for Inert vacuum* Swiezewska Detailed study of the SM-like h $M_{h}^{2} = m_{11}^{2} = \lambda_{1} v^{2} = (125 \text{ GeV})^{2}$ Study of dark scalars D = (H, A, H+, H-) - in pairs! $M_{H+}^2 = -\frac{m_{22}^2}{2} + \frac{\lambda_3}{2}v^2 \quad M_A^2 = -\frac{m_{22}^2}{2} + \frac{\lambda_3 + \lambda_4 - \lambda_5}{2}v^2$ m₂₂² arbitrary ! (decoupling...) $M_{H}^{2} = -\frac{m_{22}^{2}}{2} + \frac{\lambda_{3} + \lambda_{4} + \lambda_{5}}{2}v^{2}$ H – dark matter ($\lambda_5 < 0$) D couple to V = W/Z (eg. AZH, H⁻W⁺H), not DVV! Quartic selfcouplings D⁴ proportional to λ_2 Couplings with Higgs: hHH ~ λ_{345} h H+H- ~ λ_3 10 M. Krawczyk, Portoroz 2017

LHC – Higgs H_{125} data \rightarrow h (IDM) Direct couplings to W/Z and fermions - as in SM

- Loop coupling hgg as in SM
- Loop coupling hyy, h $Z\gamma$ extra H⁺ (λ_3) contribution
- Total width extra contributions $h \rightarrow HH, AA, H+H-$
- Invisible decay h \rightarrow HH ($\sim\lambda_{345}$)



$R_{\gamma\gamma}$ as a function of mass H, H +



Invisible h decay ->coupling hHH

- *h* → *HH* invisible decay (*H* is stable)
- augmented total width of the Higgs boson, $\Gamma(h \rightarrow HH) \sim \lambda_{345}^2$







WMAP window for very light H (DM)

using MicrOmegas



M. Krawczyk, Portoroz 2017

here $\lambda_{345} \sim 0.5$ in contradiction to LHC₅!

Relic density for DMD. Sokołowskawith mass > 64 GeV $M_{A,H^{\pm}} = M_H + \delta_{A,\pm}$



For 64 GeV distribution still symmetric, above 76 GeV asymmetry due to annihilation to gauge bosons

M. Krawczyk, Portoroz 2017

Two scales: M_h/2 and M_W

Using PLANCK data

[Planck update: D. Sokołowska, P. Swaczyna, 2014]

$h \rightarrow HH$ open



- light DM $(M_H < 10 \text{ GeV})$ \Rightarrow excluded
- intermediate DM 1 (50 GeV $< M_H < M_H/2$) $\Rightarrow M_H > 53$ GeV
- intermediate DM 2 $(M_h/2 < M_H \lesssim 82 \,\text{GeV})$ $\Rightarrow R_{\gamma\gamma} < 1$

• heavy DM $(M_H > 500 \text{ GeV})$ $\Rightarrow R_{\gamma\gamma} \approx 1$

Full scan for IDM A. Ilnicka, T. Robens, MK Phys.Rev. D93 (2016) Theor. constraints – stability of the potential (positivity), pert.unitarity, condition for the Inert vacuum +LEP constraints STU (from 2014) h total width W/Z total width **Higgssignal/Higgs bounds** Lifetime of H+ (< 10^{-7} s to decay inside detector) Relic density Planck $\Omega < 0.1241$ (95% CL) and "exact" Direct detection LUX (2015) \rightarrow scan over M_H up to 1 TeV Benchmarks other analyses Stahl.., Blinov ... Cline ...Arhrib, ..Belayev...,Poulose, ...Banerjee 18 M. Krawczyk, Portoroz 2017

Low mass H (DM)

1505.04734,1508.01671



Limit on mass of DM: M_H > 45 GeV !

T² corrections:evolution of the Universe
→ rays from EWs phase to Inert phase
one, two or three stages of Universe
(2nd order PT, one 1st order)

Ginzburg, Kanishev,MK, Sokołowska PRD 2010

$$R = \frac{\lambda_{345}}{\sqrt{\lambda_1 \lambda_2}}.$$



beyond T² corrections: strong 1st order PT

G. Gil MsThesis'2011, G.Gil, P. Chankowski, MK 1207.0084 [hep-ph] PLB 2012

We applied one-loop effective potential at T=0 (Coleman-Wienberg term) and temperature dependent effective potential at T≠0 (with sum of ring diagrams)

$$V_T^{(1L)}(v_1, v_2) = V_{\text{eff}}^{(1L)}(v_1, v_2) + \Delta^{(1L)} V_{T \neq 0}(v_1, v_2)$$

Results for v(T_{EW})/T_{EW}>1 Mh=125 GeV, MH=65 GeV, λ2=0.2

strong 1st order phase transition



IDMS Bonilla, Diaz-Cruz, Darvishi, Sokołowska, MK – J.Phys. G43 (2016)

 IDM + extra neutral complex singlet χ with a complex vev

 → towards CP violation and baryogenesis
 SM-like doublet - singlet interaction → mixing in the neutral scalar sector
 3 neutral Higgses: h1 (SM-like), h2, h3

Small change in h₁ couplings to SM particles
 Dark doublet as before → H is a good DM candidate, modifications due to h2 and h3

Fields and potential of the IDMS

$$\Phi_{\mathsf{S}} = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}} (v + \phi_1 + i\phi_6) \end{pmatrix}, \quad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_4 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_5 + i\phi_5) \end{pmatrix}, \quad \Phi_{\mathsf{D}} = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}} (\phi_5 + i\phi_5)$$

$$Z_2$$
: $\Phi_1 \to \Phi_1$, $\Phi_2 \to -\Phi_2$, SM fields \to SM fields, $\chi \to \chi$.

$$V = -\frac{1}{2} \left[\frac{m_{11}^2 \Phi_1^{\dagger} \Phi_1 + m_{22}^2 \Phi_2^{\dagger} \Phi_2}{m_{11}^2 \Phi_1 \Phi_2} \right] + \frac{1}{2} \left[\lambda_1 \left(\Phi_1^{\dagger} \Phi_1 \right)^2 + \lambda_2 \left(\Phi_2^{\dagger} \Phi_2 \right)^2 \right]$$

$$+ \lambda_3 \left(\Phi_1^{\dagger} \Phi_1 \right) \left(\Phi_2^{\dagger} \Phi_2 \right) + \lambda_4 \left(\Phi_1^{\dagger} \Phi_2 \right) \left(\Phi_2^{\dagger} \Phi_1 \right) + \frac{\lambda_5}{2} \left[\left(\Phi_1^{\dagger} \Phi_2 \right)^2 + \left(\Phi_2^{\dagger} \Phi_1 \right)^2 \right]$$

$$- \frac{m_3^2}{2} \chi^* \chi + \lambda_{s1} (\chi^* \chi)^2 \cdot \left(- \frac{m_4^2}{2} (\chi^{*2} + \chi^2) + \kappa_2 (\chi^3 + \chi^{*3}) + \kappa_3 [\chi(\chi^* \chi) + \chi^*(\chi^* \chi)] \right]$$

$$+ \Lambda_1 (\Phi_1^{\dagger} \Phi_1) (\chi^* \chi)$$
with coffly broken $L(1)$ $U(1)$: $\Phi_1 \to \Phi_2 \to \Phi_3 \chi \to \delta^{i\alpha} \chi$

M. Krawczyk, Portoroz 2017 M. Krawczyk, Portoroz 2017 Droken U(1) $U(1): \Phi_1 \rightarrow \Phi_1, \Phi_2 \rightarrow \Phi_2, \chi \rightarrow e^{i\alpha}\chi.$

Remarks

1

$$Z_2$$
: $\Phi_1 \to \Phi_1$, $\Phi_2 \to -\Phi_2$, SM fields \to SM fields, $\chi \to \chi$,

respected by vacuum -> no domain problem

The general singlet part of the potential is equal to:

$$V_{S} = -\frac{m_{3}^{2}}{2}\chi^{*}\chi - \frac{m_{4}^{2}}{2}(\chi^{*2} + \chi^{2}) + \lambda_{s1}(\chi^{*}\chi)^{2} + \lambda_{s2}(\chi^{*}\chi)(\chi^{*2} + \chi^{2}) + \lambda_{s3}(\chi^{4} + \chi^{*4}) + \kappa_{1}(\chi + \chi^{*}) + \kappa_{2}(\chi^{3} + \chi^{*3}) + \kappa_{3}(\chi(\chi^{*}\chi) + \chi^{*}(\chi^{*}\chi)).$$

The doublet-singlet interaction terms are:

$$V_{DS} = \Lambda_1(\Phi_1^{\dagger}\Phi_1)(\chi^*\chi) + \Lambda_2(\Phi_2^{\dagger}\Phi_2)(\chi^*\chi) + \Lambda_3(\Phi_1^{\dagger}\Phi_1)(\chi^{*2} + \chi^2) + \Lambda_4(\Phi_2^{\dagger}\Phi_2)(\chi^{*2} + \chi^2) + \kappa_4(\Phi_1^{\dagger}\Phi_1)(\chi + \chi^*) + \kappa_5(\Phi_2^{\dagger}\Phi_2)(\chi + \chi^*).$$

To simply model we use U(1) But with non-zero vev for singlet \rightarrow massless Nambu-Goldstone boson. So we softly break it... In order to have DM ~ IDM we neglect terms with dark doublet

Higgs sector – $h_1(125 \text{ GeV})$, h_2 , h_3

$$\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}$$

$$R = R_1 R_2 R_3 = \begin{pmatrix} c_1 c_2 & c_3 s_1 - c_1 s_2 s_3 & c_1 c_3 s_2 + s_1 s_3 \\ -c_2 s_1 & c_1 c_3 + s_1 s_2 s_3 & -c_3 s_1 s_2 + c_1 s_3 \\ -s_2 & -c_2 s_3 & c_2 c_3 \end{pmatrix}$$

 $h_1 = c_1 c_2 \phi_1 + (c_3 s_1 - c_1 s_2 s_3) \phi_2 + (c_1 c_3 s_2 + s_1 s_3) \phi_3,$

$$R_{11} = R_{11}^{-1} = c_1 c_2 \sim 1$$

M h₁ ~125 GeV, w=300-1000 GeV

$$\begin{split} M_{h_3} > M_{h_2} > 150 \text{ GeV}. \\ \kappa_{2,3} = w\rho_{2,3}, \\ -1 < \Lambda_1 < 1, \quad 0 < \lambda_{s1} < 1, \quad -1 < \rho_{2,3} < 1, \quad 0 < \xi < 2\pi. \end{split}$$

DMS - $h_1(125 \text{ GeV}_{\Gamma(h_1 \to XX)} = R_{11}^2 \Gamma(\phi_{SM} \to XX)$



Relic density - interference and second light Higgs



(a) A1-A3





(c) IDM



IDMS – heavy DM





A1-A4 -> similar results

SM+complex singlet

Branco .. Espinosa

Darvishi, Sokolowska,MK 1512.06437 (APP B47 2016); Darvishi, MK 1603.00598 SM SU(2) doublet + complex singlet with non-zero complex vev (in agreement with LHC)

- Important cubic terms
- Possibility of spontaneous CP violation
- Strong 1st order phase transition

Darvishi, JHEP 2016

Baryogenesis with vector-like quarks (iso-doublet) Darvishi, JHEP 2016; McDonald 1996

Fields and potential of the SMCS

$$\Phi_{s} = \Phi_{1} = \begin{pmatrix} \phi_{1}^{+} \\ \frac{1}{\sqrt{2}} (v + \phi_{1} + i\phi_{6}) \end{pmatrix},$$

$$\chi = \frac{1}{\sqrt{2}} (we^{i\xi} + \phi_{2} + i\phi_{3}).$$

$$W_{1} = w \cos w^{2} = w \sin w^{2} = w \sin w^{2} = w \sin w^{2}$$
Symmetry transformation $\chi \rightarrow \chi^{*}$

$$V = -\frac{1}{2} \begin{bmatrix} m_{11}^{2} \Phi_{1}^{\dagger} \Phi_{1} \end{bmatrix} + \frac{1}{2} \begin{bmatrix} \lambda_{1} (\Phi_{1}^{\dagger} \Phi_{1})^{2} \end{bmatrix}$$

$$W_{1} = \frac{1}{2} \begin{bmatrix} m_{11}^{2} \Phi_{1}^{\dagger} \Phi_{1} \end{bmatrix} + \frac{1}{2} \begin{bmatrix} \lambda_{1} (\Phi_{1}^{\dagger} \Phi_{1})^{2} \end{bmatrix}$$

$$W_{1} = \frac{1}{2} \begin{bmatrix} m_{11}^{2} \Phi_{1}^{\dagger} \Phi_{1} \end{bmatrix} + \frac{1}{2} \begin{bmatrix} \lambda_{1} (\Phi_{1}^{\dagger} \Phi_{1})^{2} \end{bmatrix}$$

$$W_{1} = \frac{1}{2} \begin{bmatrix} m_{11}^{2} \Phi_{1}^{\dagger} \Phi_{1} \end{bmatrix} + \frac{1}{2} \begin{bmatrix} \lambda_{1} (\Phi_{1}^{\dagger} \Phi_{1})^{2} \end{bmatrix}$$

$$W_{1} = \frac{1}{2} \begin{bmatrix} m_{11}^{2} \Phi_{1}^{\dagger} \Phi_{1} \end{bmatrix} + \frac{1}{2} \begin{bmatrix} \lambda_{1} (\Phi_{1}^{\dagger} \Phi_{1})^{2} \end{bmatrix} + \frac{$$

Vacuum: v, $w_1 = \cos\xi$, $w_2 = \sin\xi \neq 0$ Spont. CP violation in region

$$-4m_4^2\cos\xi + 3R_2(1+2\cos 2\xi) + R_3 \cdot = 0$$





Scan – results similar to IDMS but here low $\langle \chi \rangle \sim w$ possible

 $M_{h_1} \in [124.00, 127.00] \text{ GeV}, \ M_{h_3} \gtrsim M_{h_2} > 150 \text{ GeV}$ $0.2 < \lambda_1 < 0.3$

 $-1 < \Lambda < 1, \ 0 < \lambda_s < 1, \ -1 < \rho_{2,3} < 1, \ 0 < \xi < \pi,$

 $-90000 \text{ GeV}^2 < \mu_1^2, \mu_2^2, m_{11}^2 < 90000 \text{ GeV}^2.$ $\mu_1^2 = m_s^2 + 2m_4^2, \quad \mu_2^2 = m_s^2 - 2m_4^2.$

 $\rho_{2,3} = \kappa_{2,3}/w$



singlet self coupling λ_s , to be greater than 0.2, a the doublet-singlet coupling $|\Lambda|$, to be below 0.2.

w not too large

mass 125 GeV 33

Strong 1st order PT

T2 – corrections



Benchmarks

Benchmark	α_1	α_2	α_3	M_{h_1}	M_{h_2}	M_{h_3}	S	Т	J_{1}/v^{6}
A1	-0.047	-0.053	1.294	124.64	652.375	759.984	-0.072	-0.094	-2.2×10^{-4}
A2	-0.048	0.084	0.084	124.26	512.511	712.407	-0.001	-0.039	7.2×10^{-4}
A3	0.078	0.297	0.364	124.27	582.895	650.531	0.003	-0.046	4.5×10^{-4}
A4	0.006	-0.276	0.188	125.86	466.439	568.059	-0.013	-0.169	-9.5×10^{-4}
A5	0.062	-0.436	0.808	125.21	303.545	582.496	0.002	-0.409	5.0×10^{-6}
A6	-0.210	0.358	0.056	124.92	181.032	188.82	0.003	-0.010	-4.0×10^{-5}
A7	-0.205	0.403	0.057	125.01	175.45	178.52	0.002	-0.020	-3.5×10^{-5}

Table I. Benchmark points A1 - A7, masses are given in GeV.

Benchmark	$R^{h_1}_{\gamma\gamma}$	$R^{h_2}_{\gamma\gamma}$	$R^{h_3}_{\gamma\gamma}$	$\Gamma^{h_1}_{tot}$	$\Gamma_{tot}^{h_2}$	$\Gamma_{tot}^{h_3}$
A1	0.98	0.0021	0.0028	0.0042	0.304	0.781
A2	0.98	0.0021	0.0070	0.0042	0.145	1.31
A3	0.98	0.0055	0.085	0.0042	0.566	12.24
A4	0.92	3.3×10^{-5}	0.074	0.0043	0.001	7.08
A5	0.81	0.0029	0.17	0.0043	0.002	17.51
A6	0.82	0.19	0.11	0.0043	0.119	0.163
A7	0.81	0.18	0.15	0.0043	0.871	0.083

Baryogenesis with heavy iso - doublet vector - like quarks

Neda Dravishi, JHEP 2016 (1608.02820)

$$\mathcal{L}_Y(V_q, \chi) = \lambda_V \chi \overline{Q}_L V_R + M \overline{V}_L V_R + h.c,$$

$$\Delta \mathcal{L}_k = -\frac{\lambda_V^2 w^2}{M^2} \dot{\xi} (\overline{Q'}_L \gamma^0 Q'_L - \overline{V'}_L \gamma^0 V'_L).$$

$$\frac{n_B}{s} = \frac{225 K \alpha_W^4}{4\pi^2 g^*} \frac{\lambda_V^2 w^2}{M^2} \delta\xi,$$



Summary

- Doublets and singlets extensions of SM rich phenomenology
- Higgs and Dark Matter in IDM and IDMS
 in agreement with data
- Various stages of the Universe ?
- Strong first order phase transition → baryogenesis with vector - quarks