REASSESSING THE EXCLUSIVE DETERMINATION OF  $V_{cb}$ 

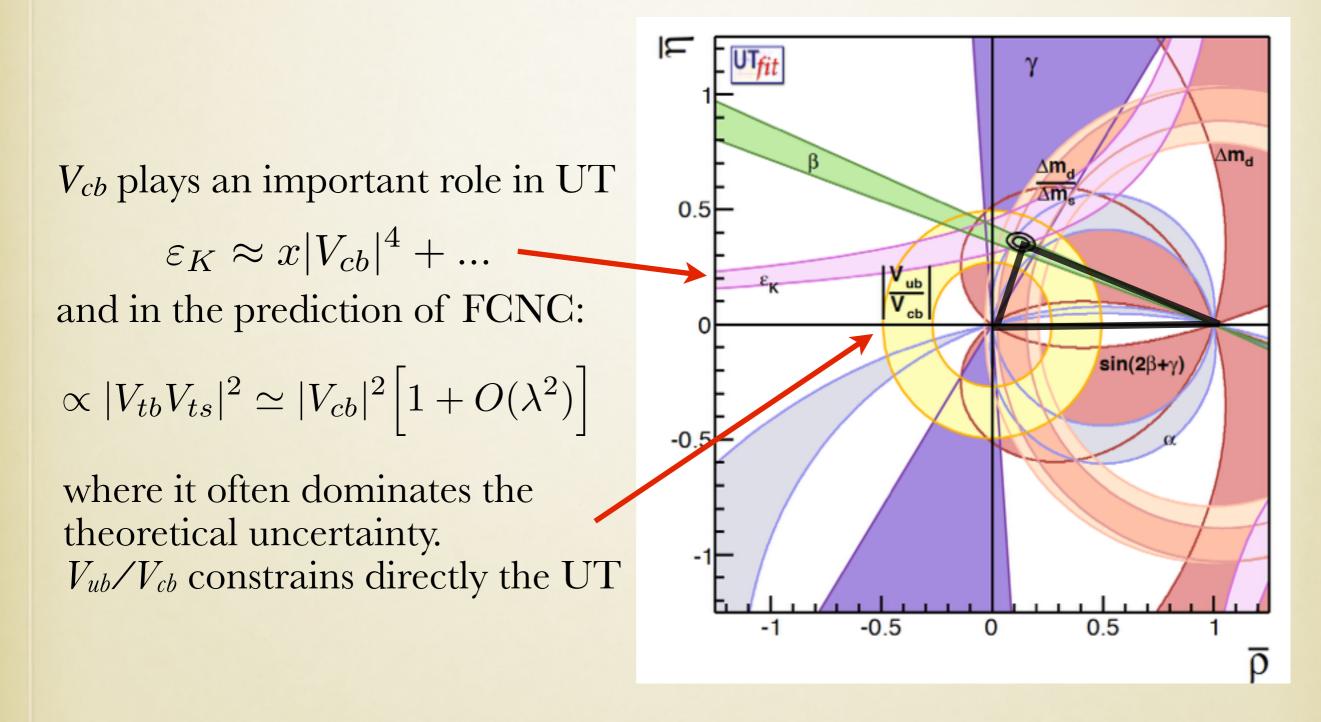
1703.06124 with D.Bigi and S. Schacht

PAOLO GAMBINO UNIVERSITÀ DI TORINO & INFN

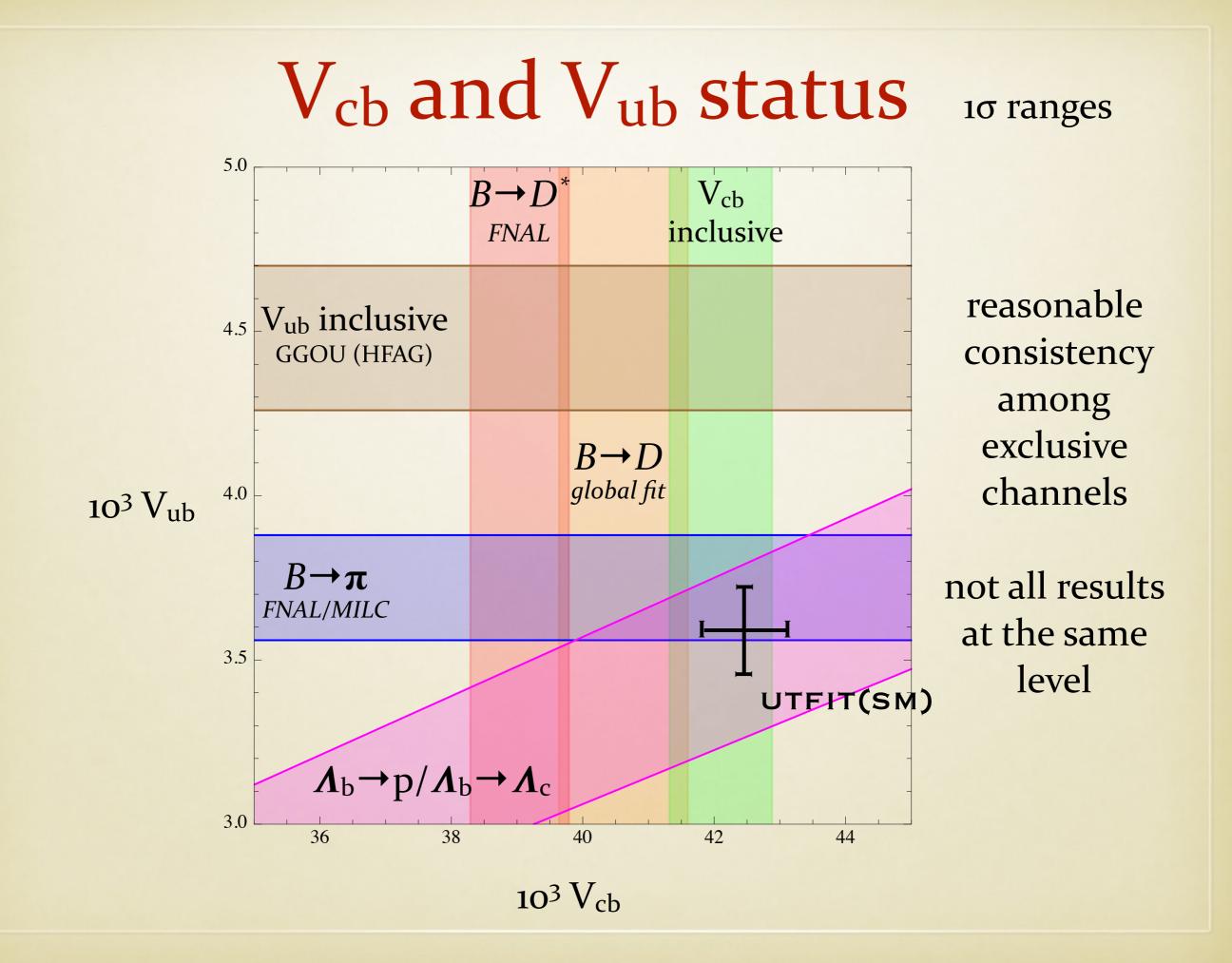


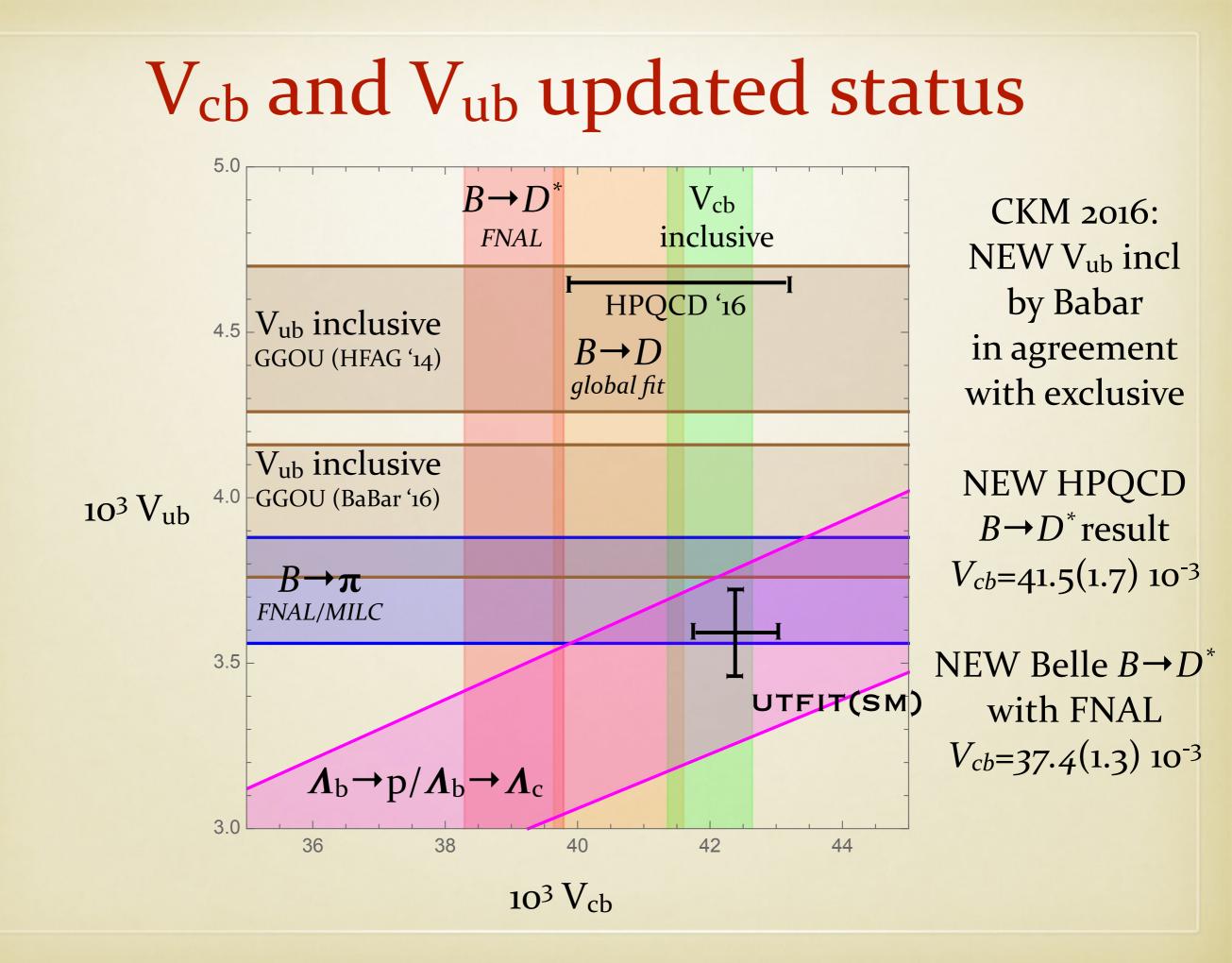
PORTOROZ 2017, 18 APRIL 2017

# **IMPORTANCE OF** $|V_{xb}|$



Since several years, exclusive decays prefer smaller  $|V_{ub}|$  and  $|V_{cb}|$ 





## SEMITAUONIC ANOMALY

$$R(D^{(*)}) = \frac{\mathcal{B}(B \to D^{(*)}\tau\nu)}{\mathcal{B}(B \to D^{(*)}\mu\nu)}$$

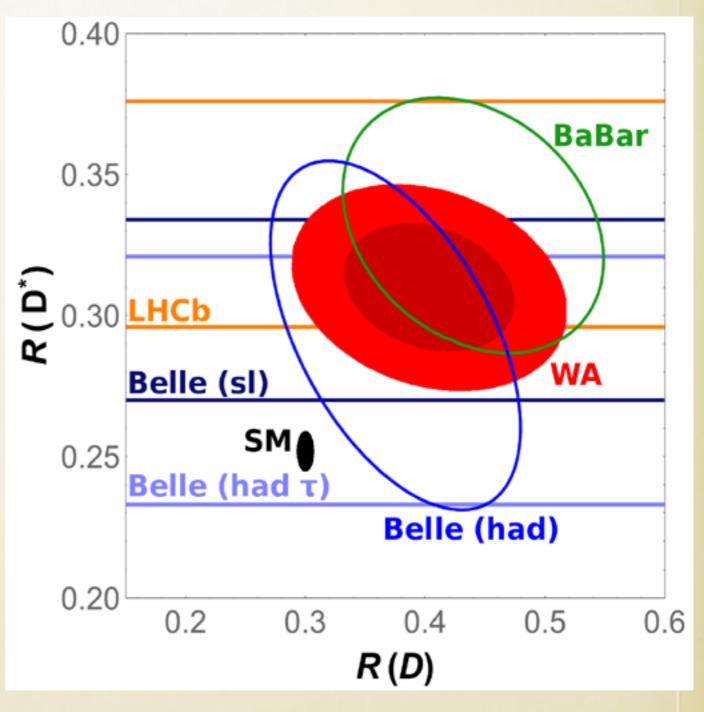
Combined discrepancy with SM  $4.0\sigma$ 

about 30% effect on tree-level process!

Lepton flavour universality violation: new scalars, leptoquarks, W'... possible connection with lepton flavour violation in  $b \rightarrow sll$ 

Inconsistent with LEP inclusive measurement

SM predictions?



Celis et al., 1612.07757

# **EXCLUSIVE** $B \rightarrow D^* \ell v$

<u>At zero recoil</u>, w=1, where rate vanishes, the ff is

$$\mathcal{F}(1) = \eta_A \left[ 1 + O\left(\frac{1}{m_c^2}\right) + \dots \right]$$

Thanks to measurement of slopes and shape parameters, exp error only ~1.3% when extrapolation to zero recoil with CLN parameterization

The ff F(1) has been computed in Lattice QCD. Only one unquenched Lattice calculation is published:

Bailey et al 1403.0635 (FNAL/MILC)

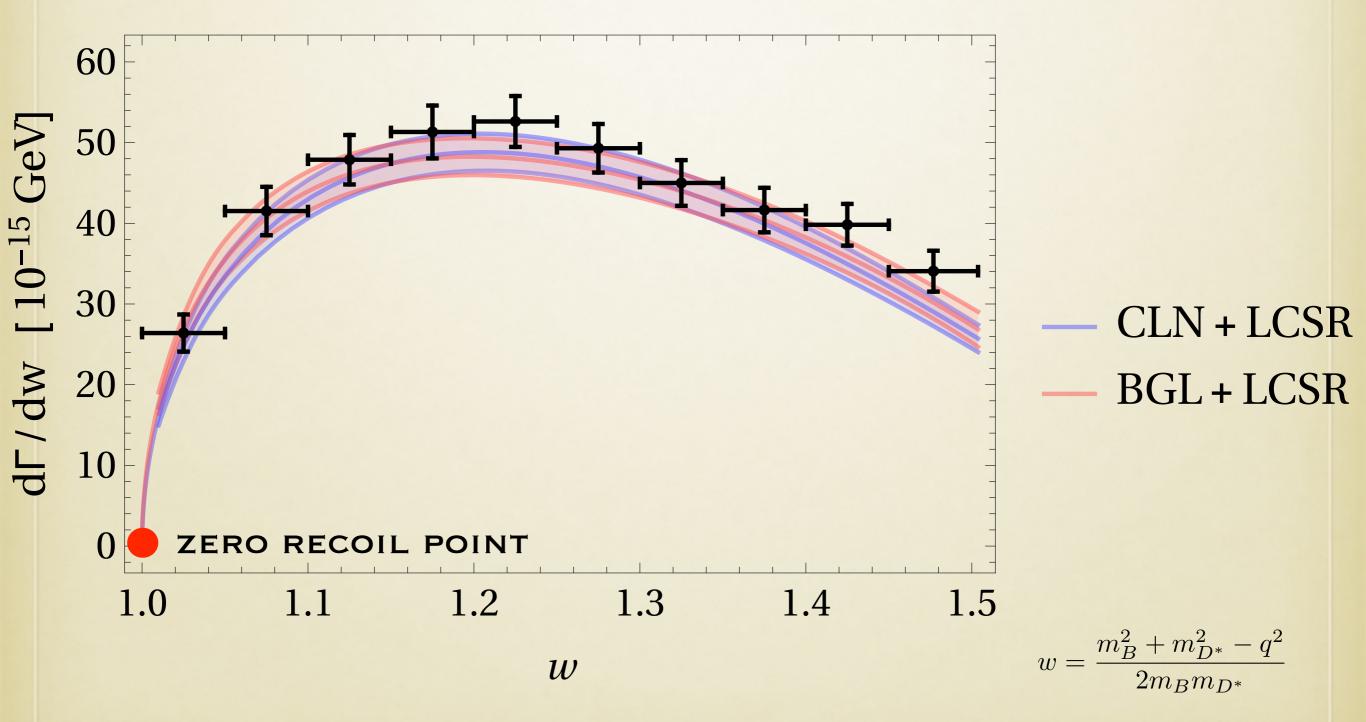
 $F(I) = 0.906(I3) \implies |V_{cb}| = 38.71(75) 10^{-3}$ 

HFAG 2016

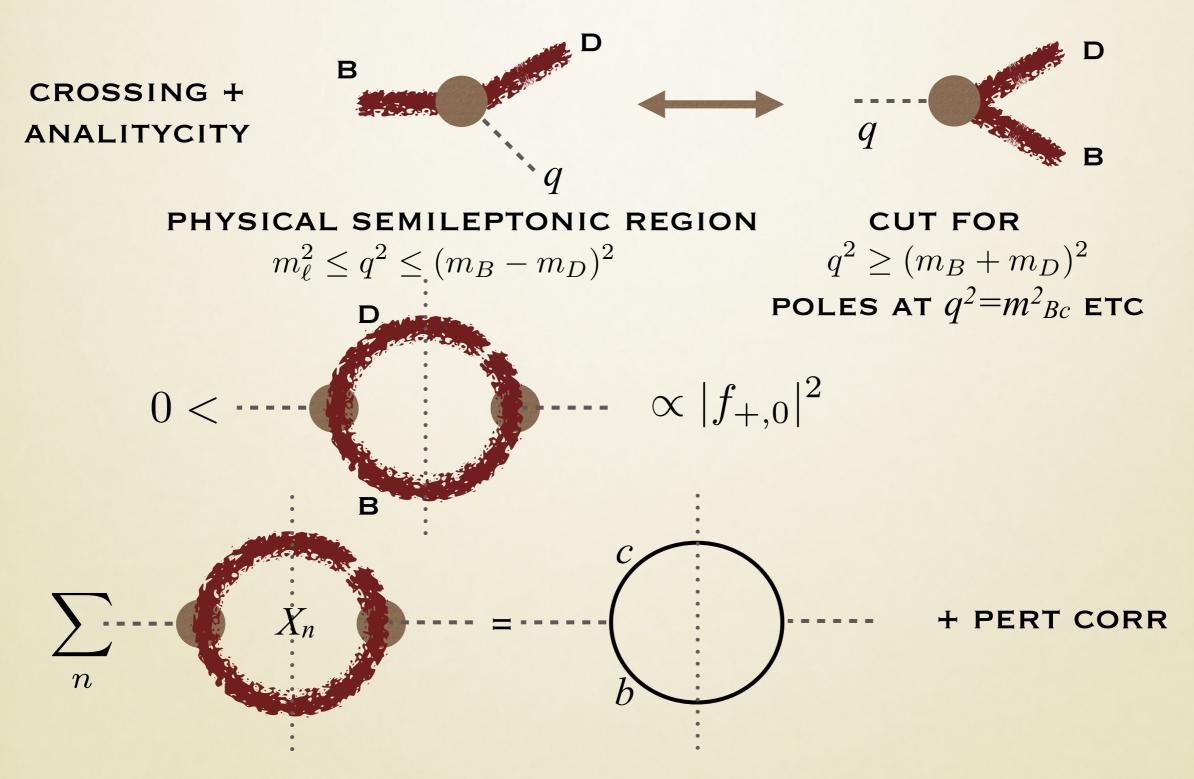
1.9% error

~ 3.3 (3.1)  $\sigma$  or ~ 8% from inclusive determination 0.04200(65) PG, Healey, Turczyk 2016 NEW HPQCD F(1)=0.862(35) preliminary, CKM 2016 NB Heavy Quark Sum Rules estimate F(1)=0.86(2)PG, Mannel, Uraltsev 2012

#### New preliminary Belle analysis 1702.01521 for the first time *w* and angular deconvoluted distributions independent of parameterization. All previous analyses are CLN based.



#### **UNITARITY CONSTRAINTS**



USING QUARK-HADRON DUALITY. DISPERSION RELATIONS→ GLOBAL QHD

### UNITARITY CONSTRAINTS

$$\begin{pmatrix} -g^{\mu\nu} + \frac{q^{\mu}q^{\nu}}{q^2} \end{pmatrix} \Pi^T(q^2) + \frac{q^{\mu}q^{\nu}}{q^2} \Pi^L(q^2) \equiv i \int d^4x \, e^{iqx} \langle 0|TJ^{\mu}(x)J^{\dagger\nu}(0)|0\rangle$$
$$\chi^L(q^2) = \frac{\partial \Pi^L}{\partial q^2}, \qquad \chi^T(q^2) = \frac{1}{2} \frac{\partial^2 \Pi^T}{\partial (q^2)^2}$$

SATISFY UNSUBTRACTED DISP REL, PERT CALCULATION FOR  $q^2=0$  Boyd, Grinstein, Lebed 1995

$$\chi_V^T(0) = [5.883 + 0.552_{\alpha_s} + 0.050_{\alpha_s^2}] \ 10^{-4} \,\text{GeV}^{-2} = 6.486(48) \ 10^{-4} \,\text{GeV}^{-2} = \chi_V^T(0) = [5.456 + 0.782_{\alpha_s} - 0.034_{\alpha_s^2}] \ 10^{-3} = 6.204(81) \ 10^{-3}$$

USING UP-TO-DATE QUARK MASSES AND 3LOOP CALCULATION Grigo et al 2012

$$\tilde{\chi}^T(0) = \chi^T(0) - \sum_{n=1,2} \frac{f_n^2(B_c^*)}{M_n^4(B_c^*)}$$

#### BOUND STATE CONTRIBUTIONS

Type	Mass $(GeV)$	Decay constants (GeV)
1-	6.329(3)	0.422(13)
1-	6.920(20)	0.300(30)
1-	7.020	
1-	7.280	
0+	6.716	
0+	7.121	

## UNITARITY CONSTRAINTS

$$z = \frac{\sqrt{1+w} - \sqrt{2}}{\sqrt{1+w} + \sqrt{2}} \qquad w = \frac{m_B^2 + m_{D^*}^2 - q^2}{2m_B m_{D^*}} \qquad 0 < z < 0.056$$

$$f_i(z) = \frac{\sqrt{\chi_i}}{P_i(z)\phi_i(z)} \sum_{n=0}^{\infty} a_n^i z^n \qquad \text{(BGL)}$$

$$f_i(z) = \frac{\sqrt{\chi_i}}{P_i(z)\phi_i(z)} \sum_{n=0}^{\infty} a_n^i z^n \qquad \text{(BGL)}$$

$$\text{TRUNCATED}$$

$$\text{AT ORDER N}$$

$$\sum_{n=0}^{\infty} (a_n^i)^2 < 1 \qquad \text{WEAK UNITARITY} \\ \text{constraints}$$
For massless leptons only 3 form factors f,g,F, contribute to B \rightarrow D'lv
$$\sum_{i=0}^{\infty} (a_n^g)^2 < 1, \qquad \sum_{i=0}^{\infty} [(a_n^f)^2 + (a_n^{\mathcal{F}_1})^2] < 1$$

$$\text{vector current} \qquad \text{axial vector current}$$

## STRONG UNITARITY CONSTRAINTS

Using information about the other channels the constraints become tighter In the heavy quark limit all  $B^{(*)} \rightarrow D^{(*)}$  form factors either vanish or are prop to the Isgur-Wise function  $H \propto$ 

$$\sum_{i=1}^{M} \sum_{n=0}^{\infty} b_{in}^2 \le 1$$

CAPRINI LELLOUCH NEUBERT CLN 1998

$$h_{A1}(z) = h_{A1}(1) \left[ 1 - 8\rho^2 z + (53\rho^2 - 15)z^2 - (231\rho^2 - 91)z^3 \right]$$
  

$$R_1(w) = R_1(1) - 0.12w_1 + 0.05w_1^2$$
  

$$R_2(w) = R_2(1) + 0.11w_1 - 0.06w_1^2$$
  

$$w_1 = w - 1$$

CLN exploit NLO HQET relations between form factors to reduce to only 2 parameters... *up to less than 2% (never included in any exp analysis)* moreover  $1/m^2$ ,  $\alpha_s^2$  and  $\alpha_s/m$  corrections can be sizable For ex at zero recoil

$$\frac{F_{D^*}(z=0)}{f_+(z=0)} = 0.966 \neq 0.860(14) \qquad \frac{f_+(0)}{f_0(0)} = 0.775 \neq 0.753(3)$$

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CLN parameterization has intrinsic uncertainties that can no longer be neglected. Recent update of HQET NLO relations by Bernlochner et al. 1703.05330

# THE FITS

#### CLN

#### BGL-2

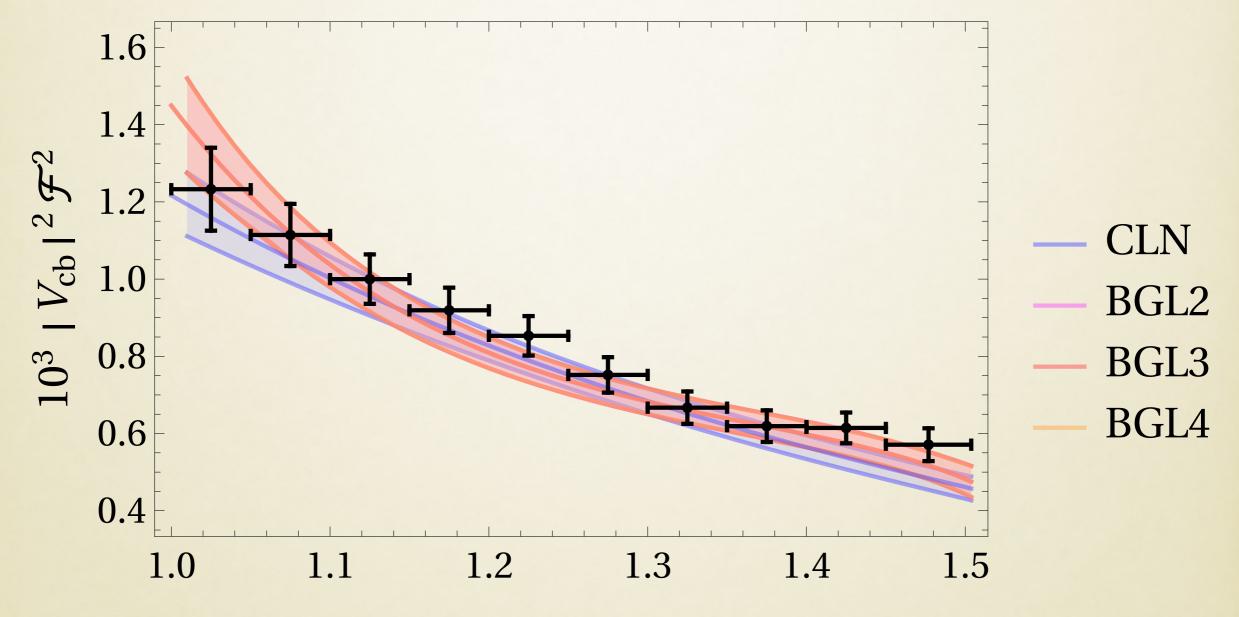
CLN Fit:	Data + lattice	Data + lattice + LCSR		Data + lattice	Data + lattice + LCSR
$\chi^2/dof$	34.3/36	34.8/39	$\chi^2/dof$	27.9/32	31.4/35
$ V_{cb} $	0.0382(15)	0.0382(14)	$ V_{cb} $	$0.0417 \left( {}^{+20}_{-21}  ight)$	$0.0404 \begin{pmatrix} +16 \\ -17 \end{pmatrix}$
$\rho_{D^*}^2$	$1.17 \begin{pmatrix} +15 \\ -16 \end{pmatrix}$	1.16(14)	$a_0^f$	0.01223(18)	0.01224(18)
$R_1(1)$	$1.391 \begin{pmatrix} +92 \\ -88 \end{pmatrix}$	1.372(36)	$a_1^f$	$-0.054\left(^{+58}_{-43} ight)$	$-0.052\left(^{+27}_{-15} ight)$
$R_{2}(1)$	$0.913 \begin{pmatrix} +73 \\ -80 \end{pmatrix}$	$0.916 \left( {}^{+65}_{-70} \right)$	$a_2^f$	$0.2 \begin{pmatrix} +7 \\ -12 \end{pmatrix}$	$1.0 \begin{pmatrix} +0 \\ -5 \end{pmatrix}$
$h_{A_1}(1)$	0.906 (13)	0.906(13)	$a_1^{\mathcal{F}_1}$	$-0.0100\left(^{+61}_{-56} ight)$	$-0.0070\left(^{+54}_{-52} ight)$
			$a_2^{\mathcal{F}_1}$	0.12(10)	$0.089 \left( ^{+96}_{-100}  ight)$
reproduces Belle's deconvoluted			$a_0^g$	$0.012 \left( {}^{+11}_{-8} \right)$	$0.0289 \begin{pmatrix} +57 \\ -37 \end{pmatrix}$
results. Best CLN			$a_1^g$	$0.7 \begin{pmatrix} +3 \\ -4 \end{pmatrix}$	$0.08 \left( {}^{+8}_{-22} \right)$
analysis V <sub>cb</sub> =0.0374(13)			$a_2^g$	$0.8 \left( {}^{+2}_{-17} \right)$	$-1.0\left(^{+20}_{-0} ight)$

see also Grinstein & Kobach, 1703.08170

#### 9% and 6% (with LCSR) difference in V<sub>cb</sub>

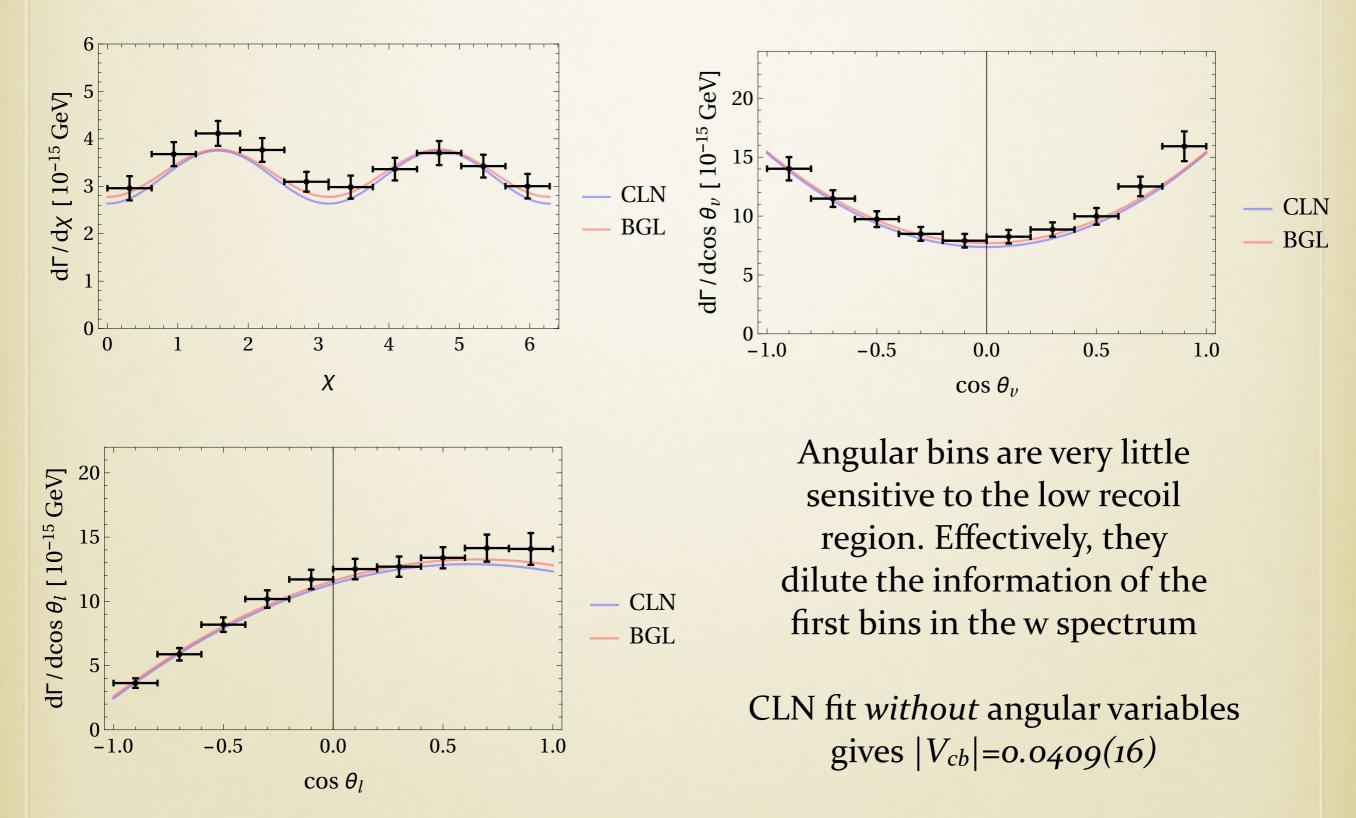
LCSR: Light Cone Sum Rule results from Faller et al, 0809.0222

 $h_{A_1}(w_{max}) = 0.65(18),$  $R_1(w_{max}) = 1.32(4), \quad R_2(w_{max}) = 0.91(17)$ 

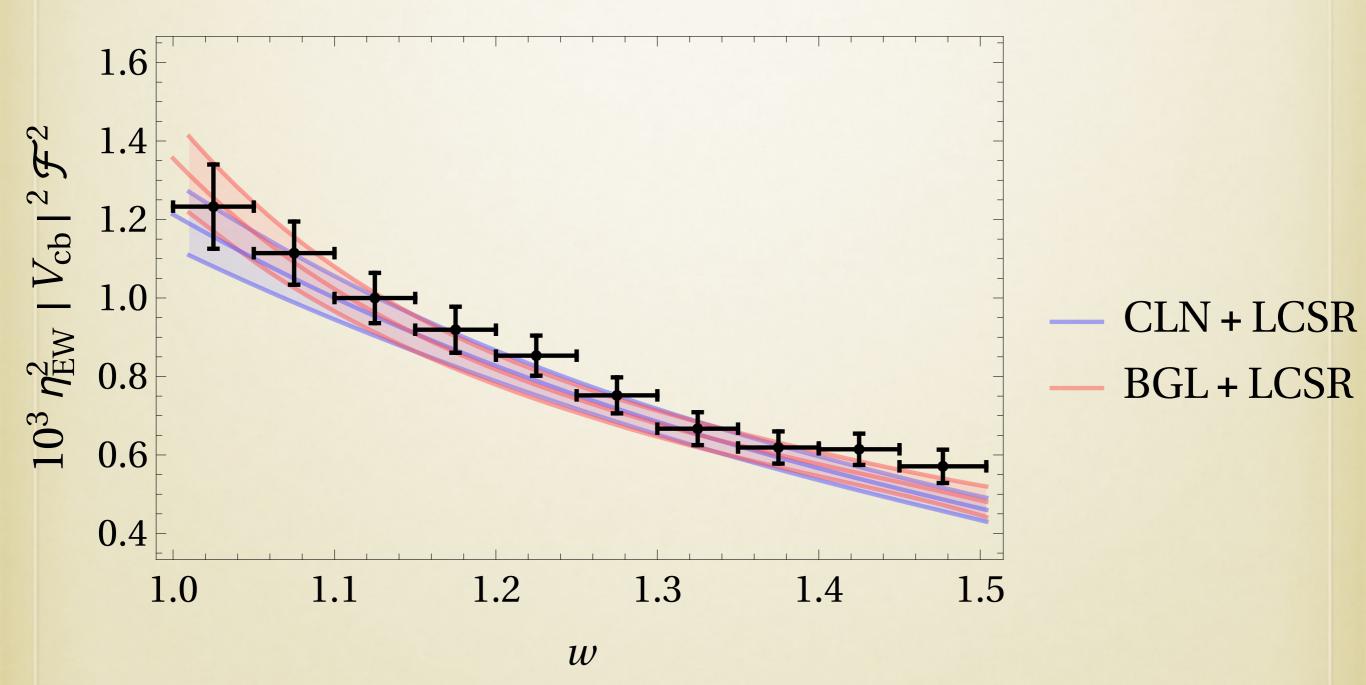


W

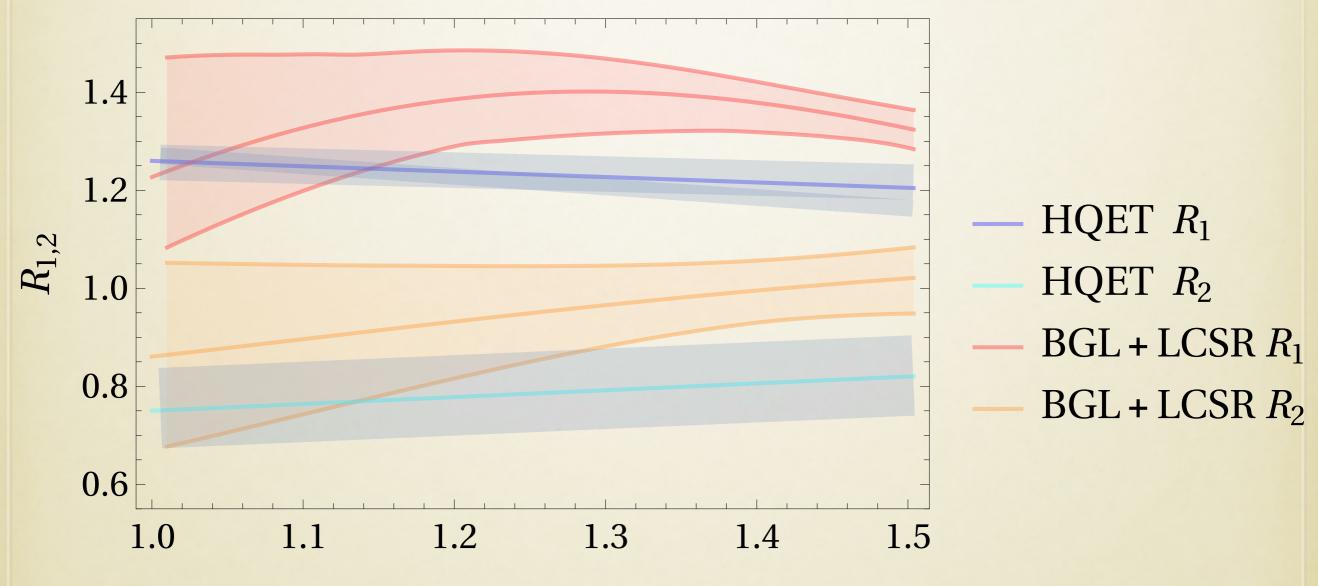
#### **ANGULAR DEPENDENCE**



### WITH LCSR CONSTRAINTS



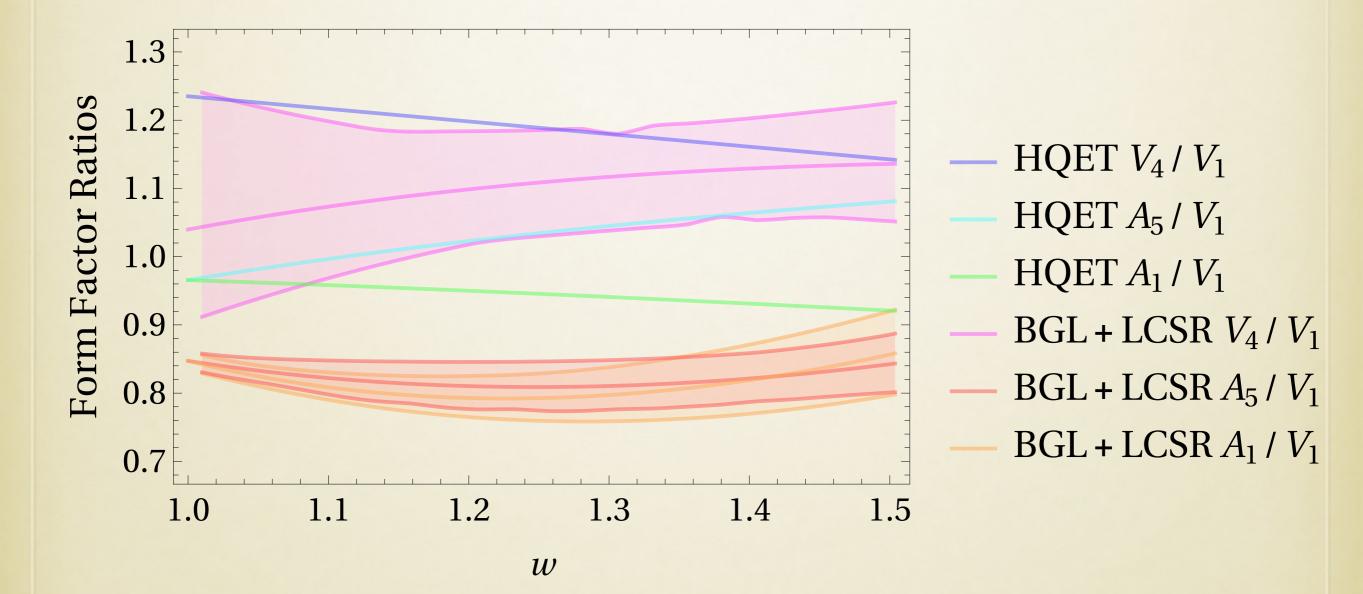
### **COMPARING WITH HQET**



W

HQET NLO predictions from Bernlochner et al. Uncertainties ONLY due to QCD sum rules

#### **COMPARING WITH HQET**



HQET NLO predictions from Bernlochner et al (uncertainties not shown) Bottom line: BGL fit compatible with HQET within uncertainties

## FUTURE SCENARIO

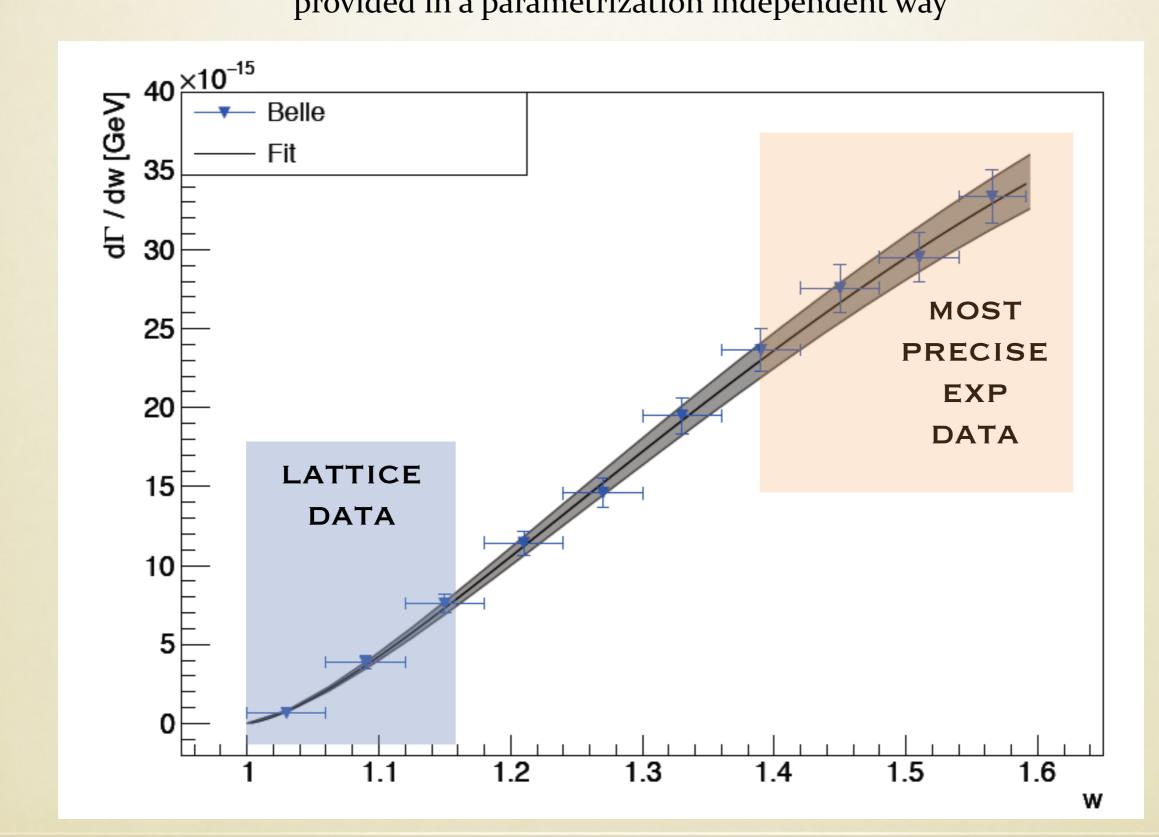
Future lattice fits	$\chi^2/dof$	$ V_{cb} $
	-	0.0407(12)
		0.0406(12)
$\operatorname{BGL}$	28.2/33	0.0409(15)
BGL+LCSR	31.4/36	0.0404(13)

assuming Lattice QCD will provide an estimate of the slope with 5% accuracy

$$\frac{\partial \mathcal{F}}{\partial w}|_{w=1} = -1.44 \pm 0.07$$

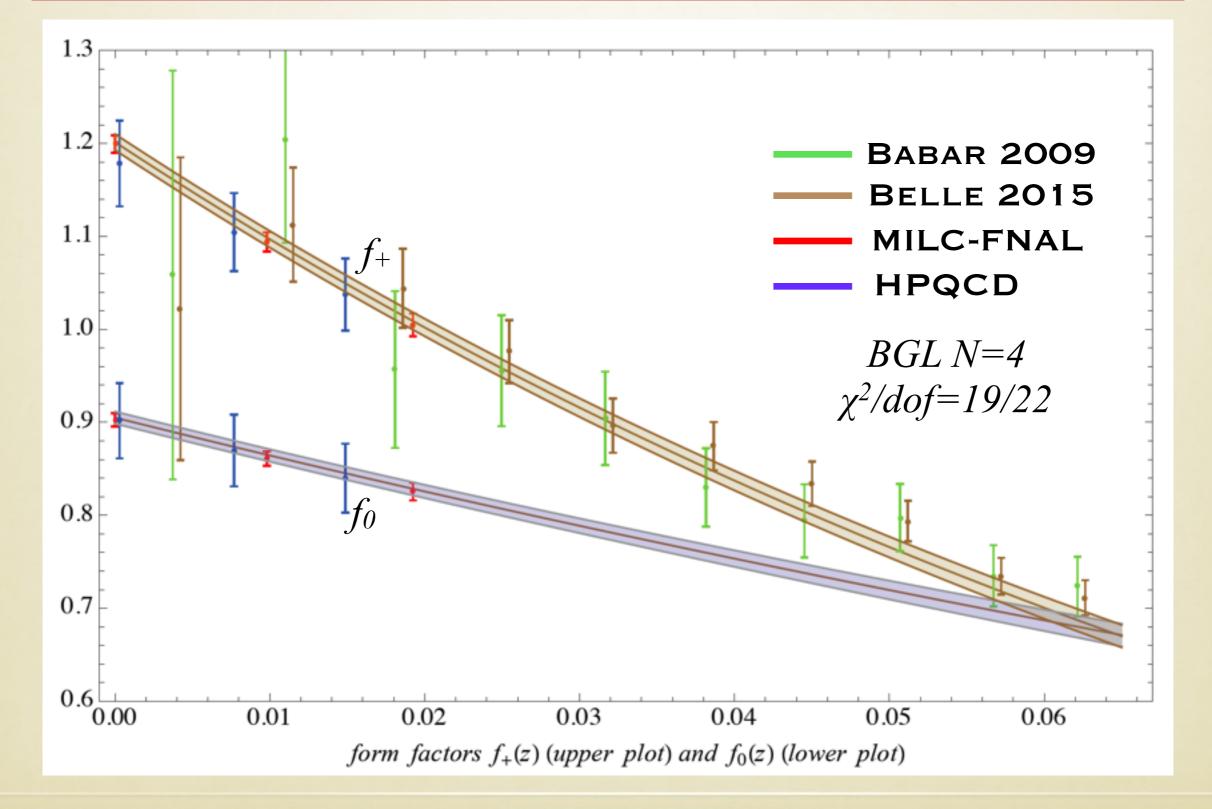
## BELLE SPECTRUM 1510.03657

provided in a parametrization independent way



Global fit to  $B \rightarrow Dlv$ 

#### D.Bigi, PG arXiv:1606.08030



## RESULTS

exp data	lattice data	N,par	$10^3 \times  V_{cb} $	$\chi^2/dof$	R(D)
all	all	2,BGL	40.62(98)	22.1/26	0.302(3)
all	all	3,BGL	40.47(97)	18.2/24	0.299(3)
all	all	4, BGL	40.49(97)	19.0/22	0.299(3)
Belle	all	3,BGL	40.92(1.12)	11.6/14	0.300(3)
BaBar	all	3,BGL	40.11(1.55)	12.6/14	0.301(4)
all	FNAL	3,BGL	40.17(1.05)	10.4/18	0.293(4)
all	HPQCD	3,BGL	$40.51^{+1.82}_{-1.71}$	10.1/18	0.299(7)
all	all	CLN	40.85(95)	77.1/29	0.305(3)
all	$f_+$ only	CLN	40.33(99)	20.0/23	0.305(3)
all	all	2, BCL	40.49(98)	18.2/26	0.299(3)
all	all	3,BCL	40.48(96)	18.2/24	0.299(3)
all	all	4, BCL	40.48(97)	17.9/22	0.299(3)



- Fitting the latest Belle B → D<sup>\*</sup>lv data with BGL and CLN leads to quite different values of V<sub>cb</sub>. While this may be related to this specific set of data, it certainly calls for a reanalysis of previous Babar and Belle data with a more flexible parameterization.
- HQET input is still useful, but it carries non-negligible uncertainty which can no longer be neglected.
- Future lattice determinations of the zero recoil slope of the form factors will improve significantly the situation, but it will be impossible to combine them with the present HFAG averages based on CLN. This is already the case for the B → D channel.
- We did not resolve the V<sub>cb</sub> puzzle, just pointed out a neglected systematics