

UV completion of the radiative neutrino mass models

work in progress with K. Kumerički and I. Picek

Timon MEDE

University of Zagreb

CROATIA

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GOALS :

- BOTTOM-UP approach to address m_ν (and DM) problems (in the SM gauge group context)
- LHC rationale: focusing on mechanisms whose required states can get tested @ LHC
- possible UV completions of such models: embedding into GUT scenarios

Open problems of Standard Model

- neutrino masses (neutrino oscillations: existence, lightness)
- dark matter
- hierarchy problem - Higgs naturalness problem
- ...

Motivates search for BSM physics: requires introducing new “heavy” particles (no gauge anomaly: e.g. vector-like matter)

Open problems of Standard Model

“accidental” SM matter (used representations seemingly arbitrary)

SM Periodic Table

Three Generations
of Matter (Fermions) spin $\frac{1}{2}$

	I	II	III	
mass →	2.4 MeV	1.27 GeV	173.2 GeV	0
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0
name →	u up	c charm	t top	g gluon
	Left Right	Left Right	Left Right	
	d down	s strange	b bottom	0
Quarks	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	0
	Left Right	Left Right	Left Right	γ photon
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	91.2 GeV
	0	0	0	0
	e electron	μ muon	τ tau	Z weak force
Leptons	-1	-1	-1	126 GeV
	Left Right	Left Right	Left Right	0
	W[±] weak force			H Higgs boson
	80.4 GeV			spin 0
	± 1			

Bosons (Forces) spin 1

Top-down BSM approach: GUTs

PROs :

- unified framework \rightarrow less parameters, connection between different quantities
- electric charge quantization
- prediction of new phenomena (proton decay, ...)
- “explanation” for the heavy spectrum

CONs :

- simplest models ruled out (e.g. no unification, massless neutrinos in Georgi-Glashow minimal non-SUSY $SU(5)$)
- infinitely many possibilities (GUT gauge group, representations, scalar potential): no clear selection rule
- phenomenology (neutrino masses, DM, proton decay, light fermion mass relations ...)

Bottom-up BSM approach: neutrino mass models

Focusing on established deficiencies of SM: construction of mechanisms which try to explain experimental findings

Seesaw mechanism

Weinberg dim 5 effective operator $LLHH$ after integrating out BSM particles (M)

$$\Delta\mathcal{L} = \frac{\lambda}{M} LLHH + \mathcal{O}\left(\frac{1}{M^2}\right) \quad \rightarrow \quad m_\nu \sim \lambda \frac{v^2}{M} \sim 0.1 \text{ eV}$$

@ tree-level: only 3 possible realizations (with single BSM particle) :

- type I:** $3 \times N_R^f \sim (1, 1, 0)$
 [Minkowski '77; Yanagida '79; Glashow '79; Gell-Mann, Ramond, Slansky '79; Mohapatra, Senjanović '79]
- type II:** $1 \times \Delta^s \sim (1, 3, -1) \equiv \begin{pmatrix} \Delta^- & -\sqrt{2}\Delta^0 \\ \sqrt{2}\Delta^{--} & -\Delta^- \end{pmatrix}$
 [Magg, Wetterich '80; Schechter, Valle '80; Lazarides et al. '80; Mohapatra, Senjanović '80; Gelmini, Roncadelli '80]
- type III:** $3 \times \Sigma^f \sim (1, 3, 0) \equiv \begin{pmatrix} \frac{1}{\sqrt{2}}\Sigma^0 & \Sigma^+ \\ \Sigma^- & -\frac{1}{\sqrt{2}}\Sigma^0 \end{pmatrix}$
 [Foot, Lew, He, Joshi '89]

Seesaw mechanism

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If $\lambda \sim \mathcal{O}(1)$ and $v = \langle H \rangle = 246 \text{ GeV}$, scale of new physics (seesaw scale)
 $M \sim 10^{14} \text{ GeV} \rightarrow$ out of reach of LHC.

Either lowering seesaw scale by going to higher dimensional operators:

$$m_\nu \sim v \left(\frac{v}{M}\right)^{\dim-4}$$

- dim 7: $(LLHH)(H^\dagger H)$
- dim 9: $(LLHH)(H^\dagger H)(H^\dagger H) \rightarrow M \sim \text{TeV}$
- ...

or ...

Zee model

... or consider radiative neutrino mass models (mass suppression by loop factor)

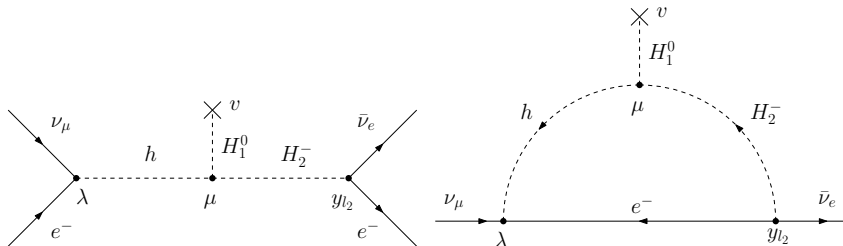
Zee '80: realization of radiative neutrino mass by 2HD + charged Higgs singlet

$$\text{SM} + H_2 \sim (1, 2, +\frac{1}{2}) \subset 5_{H_2} + h^+ \sim (1, 1, +1) \subset 10_H$$

$$V_{Zee} \sim \lambda LLh^+ + y_l \bar{L} H_a e + \mu H_1 H_2 h^- + h.c.$$

(where λ is antisymmetric in flavour space) breaks global $B - L$

$\Delta L = 2$ processes (e.g. $\nu_\mu e^- \rightarrow \bar{\nu}_e e^-$) break L_e, L_μ, L_τ and contribute to Majorana neutrino mass (oscillation of neutrino into antineutrino; Majorana neutrino is its own antineutrino):



Zee model

... or consider radiative neutrino mass models (mass suppression by loop factor)

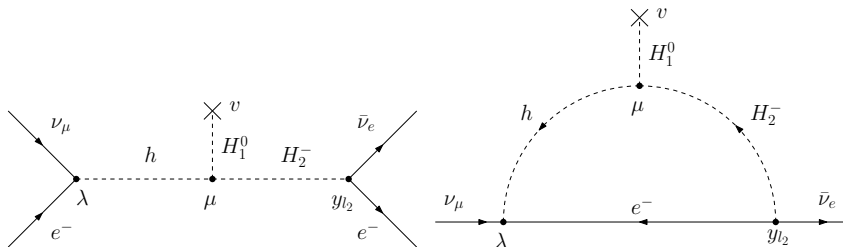
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(where λ is antisymmetric in flavour space) breaks global $B - L$

$$m_\nu \sim m_e \lambda y_{l_2} \mu v m_h^{-2} \log \frac{m_h^2}{m_{H_2}^2}$$



Synthesis: UV complete neutrino mass models

Potentially realistic extensions of Georgi-Glashow model :

▷▷▷ Georgi, Glashow '74:

minimal non-SUSY $SU(5)$: $24_H \oplus 5_H \oplus 3 \times (10_F \oplus \bar{5}_F)$ ✗

- no unification, massless neutrinos

Synthesis: UV complete neutrino mass models

Potentially realistic extensions of Georgi-Glashow model :

▷▷▷ $SU(5) + 3 \times 1_F + \text{type I seesaw}$ ✗

- no unification

Synthesis: UV complete neutrino mass models

Potentially realistic extensions of Georgi-Glashow model :

▷▷▷ Doršner, Fileviez Perez et al. '05:

$SU(5) + 15_H \supseteq (1, 3, +1) \oplus (3, 2, +\frac{1}{6}) \oplus (6, 1, -\frac{2}{3}) + \text{type II seesaw} \checkmark$

- unification - light leptoquarks
- non-renormalizable model - higher dimensional operators needed to correct $SU(5)$ fermion mass relations

Synthesis: UV complete neutrino mass models

Potentially realistic extensions of Georgi-Glashow model :

- ▷▷▷ Bajc, Nemevšek, Senjanović '07: $SU(5) + 24_F \supset (1, 3, 0)_F \oplus (1, 1, 0)_F \oplus \dots$
 + combination of type I and type III seesaw + 1 massless neutrino ✓
- unification - $m_{(1,3,0)_{H,F}} \lesssim \text{TeV}$
 - non-renormalizable model - higher dimensional operators needed to correct the $SU(5)$ fermion mass relations & to split fields in 24_F
 - fine-tuned but predictive: proton decay & unification constraints

Synthesis: UV complete neutrino mass models

Potentially realistic extensions of Georgi-Glashow model :

▷▷▷ Fileviez Perez, Murgui '16: $SU(5) + 10_H + 45_H \supseteq (1, 2, +\frac{1}{2}) \oplus (3, 1, -\frac{1}{3}) \oplus (3, 3, -\frac{1}{3}) \oplus (\bar{3}, 1, +\frac{4}{3}) \oplus (\bar{3}, 2, -\frac{7}{6}) \oplus (\bar{6}, 1, -\frac{1}{3}) \oplus (8, 2, +\frac{1}{2})$ ✓

- unification & proton decay (main contribution from gauge boson mediation) - $m(8, 2, +\frac{1}{2}) \sim 1 - 10 \text{ TeV}$
- 45_H - correct mass relations between charged leptons & down-type quarks @ renormalizable level

$$M_d = Y_1^T \frac{v_5}{\sqrt{2}} - \frac{1}{3} Y_2^T \frac{v_{45}}{\sqrt{2}}$$

$$M_e = Y_1 \frac{v_5}{\sqrt{2}} + Y_2 \frac{v_{45}}{\sqrt{2}}$$

+ contains 2nd Higgs doublet - generates neutrino masses @ quantum level through Zee mechanism

$$V_{Zee} \supseteq \lambda \bar{5}_F \bar{5}_F 10_H + \bar{5}_F 10_F (Y_1^* 5_H^* - \frac{1}{6} Y_2^* 45_H^*) - \frac{1}{6} \mu 5_H 45_H 10_H^* + h.c.$$

Brdar, Picek, Radovčić '13: 1-loop radiative neutrino mass model - variant of Zee model: triplet instead of 2nd Higgs doublet

$$\Delta \sim (1, 3, 0) = \begin{pmatrix} \frac{1}{\sqrt{2}}\Delta^0 & \Delta^+ \\ \Delta^- & -\frac{1}{\sqrt{2}}\Delta^0 \end{pmatrix}$$

$$h^+ \sim (1, 1, +1)$$

and 3 generations of vector-like lepton doublets: $E_{L,R} \sim (1, 2, -\frac{1}{2})$

$$\mathcal{L} = M\bar{E}_L E_R + \tilde{M}\bar{L}_L E_R + y\bar{E}_L H I_R + y_1\bar{L}_L^c E_L h^+ +$$

$$+ y_2\bar{L}_L \Delta E_R + y_3\bar{E}_L \Delta E_R + y_4\bar{L}_L^c L_L h^+ + h.c.$$

$$V(H, \Delta, h^+) = -\mu_H^2 H^\dagger H + \lambda_1 (H^\dagger H)^2 + \mu_h^2 h^- h^+ + \lambda_2 (h^- h^+)^2$$

$$+ \mu_\Delta^2 \text{Tr}[\Delta^2] + \lambda_3 (\text{Tr}[\Delta^2])^2 + \lambda_4 H^\dagger H h^- h^+ + \lambda_5 H^\dagger H \text{Tr}[\Delta^2]$$

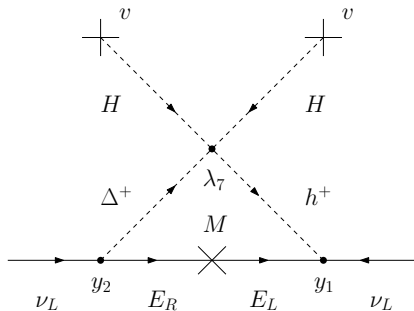
$$+ \lambda_6 h^- h^+ \text{Tr}[\Delta^2] + (\lambda_7 H^\dagger \Delta \tilde{H} h^+ + h.c.) + \mu H^\dagger \Delta H$$

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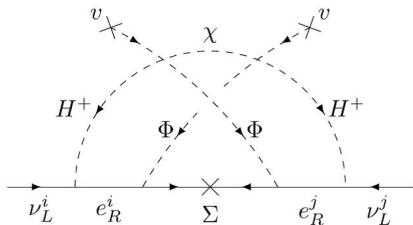
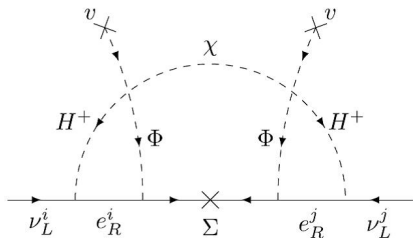
$$m_\nu \propto \frac{1}{4\pi^2} y_1 y_2 \lambda_7 M v^2 \times (\text{loop factor})$$

Čuljak, Kumerički, Picek '15: 3-loop radiative neutrino mass model - Yukawas of order $\mathcal{O}(1)$: 2HD + exotic matter

$$\phi^s \sim (1, 5, -1)$$

$$\chi^s \sim (1, 7, 0)$$

$$\Sigma^f \sim (1, 5, 0)$$



multiplets	$\subset SU(5)$	$\subset SU(5)'_{Z'}$	$\subset SO(10)$
m_ν @ 1-loop			
$\Delta = (1, 3, 0)$ $h^+ = (1, 1, +1)$	24 10	$24_0, 15_1, \overline{15}_{-1}, 175_1$ $1_5, 24_5, \overline{10}_4, 10_6,$ $50_3, \overline{50}_7, 75_5$	45, 54, 144, 210, 560 16, 45, 120, 126, 144, 210, $560^{(3)}$, $\overline{560}$
$E_{L,R} = (1, 2, -\frac{1}{2})$	$\overline{5}, 45, 70$	$5_{-2}, \overline{5}_{-3}, 40_{-4}, 40_{-1},$ $45_{-2}, \overline{45}_{-3}, 70_{-2}, \overline{70}_{-3}$	10, 16, $120^{(2)}$, 126, 126, $144^{(2)}$, $\overline{144}$, 210, $210'$, $320^{(3)}$, $560^{(3)}$, $\overline{560}$
m_ν @ 3-loops			
$H_{1,2} = (1, 2, +\frac{1}{2})$	5, 45, 70	$\overline{5}_2, 5_3, \overline{40}_4, 40_1,$ $\overline{45}_2, 45_3, \overline{70}_2, 70_3$	10, $\overline{16}$, $120^{(2)}$, 126, $\overline{126}$, 144, $\overline{144}^{(2)}$, 210, $210'$, $320^{(3)}$, 560 , $\overline{560}^{(3)}$
$\phi = (1, 5, -1)$	> 75	$70'_{-7}, \overline{70}'_{-3}$	> 560
$\chi = (1, 7, 0)$	> 75	> 75	> 560
$\Sigma = (1, 5, 0)$	> 75	$70'_{-2}, \overline{70}'_2$	> 560

SM singlets $(1, 1, 0)$ with potentially non-zero VEVs $\in 24, 75$ ($SU(5)$); $1_0, 10_1, \overline{10}_{-1}, 24_0, 50_{-2}, \overline{50}_2, 75_0$ ($SU(5)'_{Z'}$); 16, $45^{(2)}$, 54, 126, 144, $210^{(3)}$, $560^{(3)}$ ($SO(10)$)

$SU(5)$ embedding of the 1-loop (BPR) model

Particle content:

	SM multiplets	$\subset SU(5)$
scalar	$H = (1, 2, +\frac{1}{2})$	$5^a = (1, 2, +\frac{1}{2}) \oplus (3, 1, -\frac{1}{3})$; or 45, 70
	$\Delta = (1, 3, 0)$ $h^+ = (1, 1, +1)$	$24_b^a = (1, 3, 0) \oplus (8, 1, 0) \oplus (1, 1, 0) \oplus (3, 2, -\frac{5}{6}) \oplus (3, 2, +\frac{5}{6})$ $10^{ab} = (1, 1, +1) \oplus (\bar{3}, 1, -\frac{2}{3}) \oplus (3, 2, +\frac{1}{6})$
fermion	$3 \times Q = (3, 2, +\frac{1}{6})$ $3 \times u^c = (\bar{3}, 1, -\frac{2}{3})$ $3 \times e^c = (1, 1, +1)$ $3 \times L = (1, 2, -\frac{1}{2})$ $3 \times d^c = (\bar{3}, 1, +\frac{1}{3})$	$3 \times 10^{ab} = (3, 2, +\frac{1}{6}) \oplus (\bar{3}, 1, -\frac{2}{3}) \oplus (1, 1, +1)$ $3 \times \bar{5}_a = (1, 2, -\frac{1}{2}) \oplus (\bar{3}, 1, +\frac{1}{3})$
	$3 \times E_R = (1, 2, -\frac{1}{2})$ $3 \times E_L = (1, 2, -\frac{1}{2})$	$3 \times \bar{5}_a = (1, 2, -\frac{1}{2}) \oplus (\bar{3}, 1, +\frac{1}{3})$; or 45, 70 $3 \times \bar{5}_a = (1, 2, -\frac{1}{2}) \oplus (\bar{3}, 1, +\frac{1}{3})$; or $\overline{45}, \overline{70}$
	gauge	$G_\mu = (8, 1, 0)$ $W_\mu = (1, 3, 0)$ $B_\mu = (1, 1, 0)$

Realistic GUT scenarios

Potential problems of $SU(5)$:

- unification of gauge couplings (non-SUSY!!!)
- proton decay bounds: $\tau_p(p^+ \rightarrow \pi^0 e^+) \rightarrow M_{\text{GUT}} \gtrsim 10^{15.6} \text{ GeV}$
(if unification scale M_{GUT} too low, X, \bar{X} gauge bosons of masses $M_g \sim g_{\text{GUT}} v$ which mediate proton decay are too light + colour charged particles need to be sufficiently heavy)
- realistic mass relations between down-type quarks & charged leptons (@ renormalizable level 45_H is needed)
- doublet-triplet splitting (not only in 5_H , but also in 45_H or whenever some multiplets are required to be light):
fine-tuning, missing partner mechanism, ... ?

RGEs (@ 1-loop): B test

$$\frac{d\alpha_i^{-1}(m)}{d \ln \frac{m}{\mu_r}} = -\frac{b_i}{2\pi} \rightarrow \alpha_i^{-1}(m) = \alpha_i^{-1}(m_Z) - \frac{1}{2\pi} \underbrace{\left(b_i^{(SM)} + \sum_{m_k < m} \Delta b_i^{(k)} \overbrace{\frac{\ln \frac{m}{m_k}}{\ln \frac{m}{m_Z}} }^{r_k} \right)}_{B_i \text{ for } m=M_{\text{GUT}}} \ln \frac{m}{m_Z}$$

$$\frac{B_{23}}{B_{12}} \equiv \frac{B_2 - B_3}{B_1 - B_2} = \frac{\alpha_2^{-1}(m_Z) - \alpha_3^{-1}(m_Z)}{\alpha_1^{-1}(m_Z) - \alpha_2^{-1}(m_Z)} = \frac{5}{8} \frac{\sin^2 \theta_w - \frac{\alpha_{EM}(m_Z)}{\alpha_3(m_Z)}}{\frac{3}{8} - \sin^2 \theta_w} = 0.718$$

$$= 0.528 + (\text{BSM})$$

$$M_{\text{GUT}} = m_Z \exp \left(\frac{2\pi(\alpha_1^{-1}(m_Z) - \alpha_2^{-1}(m_Z))}{B_{12}} \right) = m_Z \exp \left(\frac{184.87}{B_1 - B_2} \right)$$

$$= 10^{13} \text{ GeV} + (\text{BSM})$$

$$\alpha_{\text{GUT}}^{-1} = -\frac{\frac{3}{5}B_3}{B_1 + \frac{3}{5}B_2 - \frac{8}{5}B_3} \alpha_{EM}^{-1}(m_Z) + \frac{B_1 + \frac{3}{5}B_2}{B_1 + \frac{3}{5}B_2 - \frac{8}{5}B_3} \alpha_3^{-1}(m_Z)$$

$$= 41.48 + (\text{BSM})$$

BSM contributions to running

SCALARS:

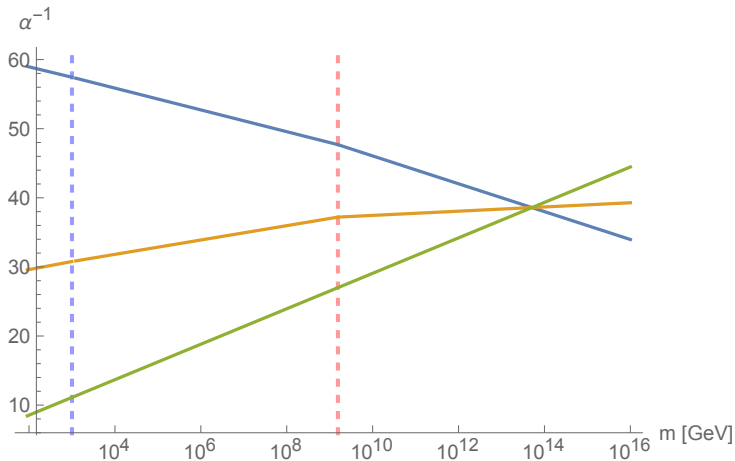
k	5_H		24_H					10_H		
	H	T_5	Δ	O_{24}	S_{24}	LQ_{24}	\overline{LQ}_{24}	h^+	\overline{T}_{10}	LQ_{10}
ΔB_{23}	$+\frac{1}{6}$	$-\frac{1}{6}r_k$	$+\frac{1}{3}r_k$	$-\frac{1}{2}r_k$	0	$+\frac{1}{12}r_k$	$+\frac{1}{12}r_k$	0	$-\frac{1}{6}r_k$	$+\frac{1}{6}r_k$
ΔB_{12}	$-\frac{1}{15}$	$+\frac{1}{15}r_k$	$-\frac{1}{3}r_k$	0	0	$+\frac{1}{6}r_k$	$+\frac{1}{6}r_k$	$+\frac{1}{5}r_k$	$+\frac{4}{15}r_k$	$-\frac{7}{15}r_k$

FERMIONS:

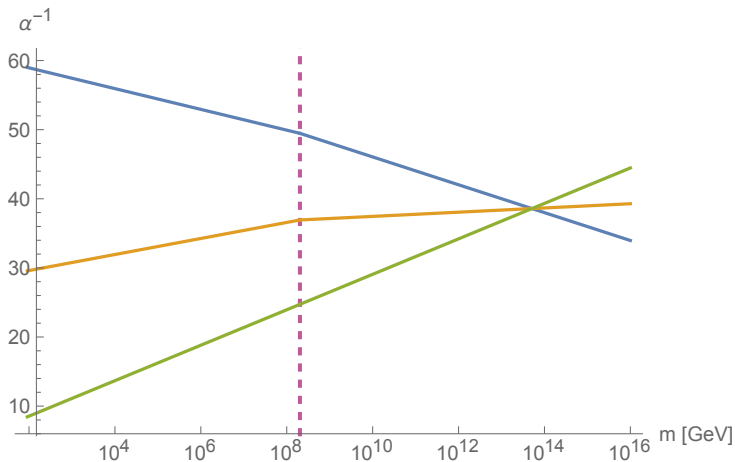
k	5_F	
	$E_{L,R}$	\overline{T}_5
ΔB_{23}	$+\frac{1}{3}r_k \times n_{\text{gen}}$	$-\frac{1}{3}r_k \times n_{\text{gen}}$
ΔB_{12}	$-\frac{2}{15}r_k \times n_{\text{gen}}$	$+\frac{2}{15}r_k \times n_{\text{gen}}$

GAUGE BOSONS:

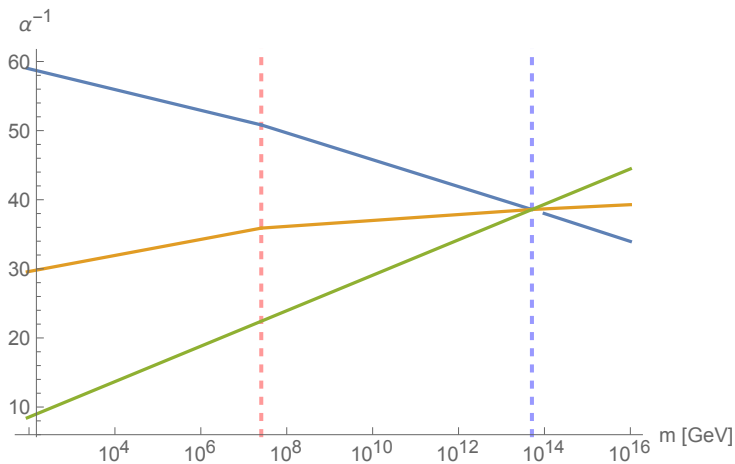
k	24_g	
	LQ_{24}	\overline{LQ}_{24}
ΔB_{23}	$-\frac{11}{6}r_k$	$-\frac{11}{6}r_k$
ΔB_{12}	$-\frac{11}{3}r_k$	$-\frac{11}{3}r_k$



$$m_{\Delta} = m_{h^+} = 10^3 \text{ GeV}, \quad m_{E_{L,R}} = 1.55 \times 10^9 \text{ GeV}, \quad M_{\text{GUT}} = 5.11 \times 10^{13} \text{ GeV}$$



$$m_{\Delta} = m_{h^+} = m_{E_{L,R}} = 2.03 \times 10^8 \text{ GeV}, \quad M_{\text{GUT}} = 5.11 \times 10^{13} \text{ GeV}$$



$$m_{E_{L,R}} = 2.55 \times 10^7 \text{ GeV}, \quad m_{\Delta} = m_{h^+} = M_{\text{GUT}} = 5.11 \times 10^{13} \text{ GeV}$$