# Bounding milli-magnetic charged particles

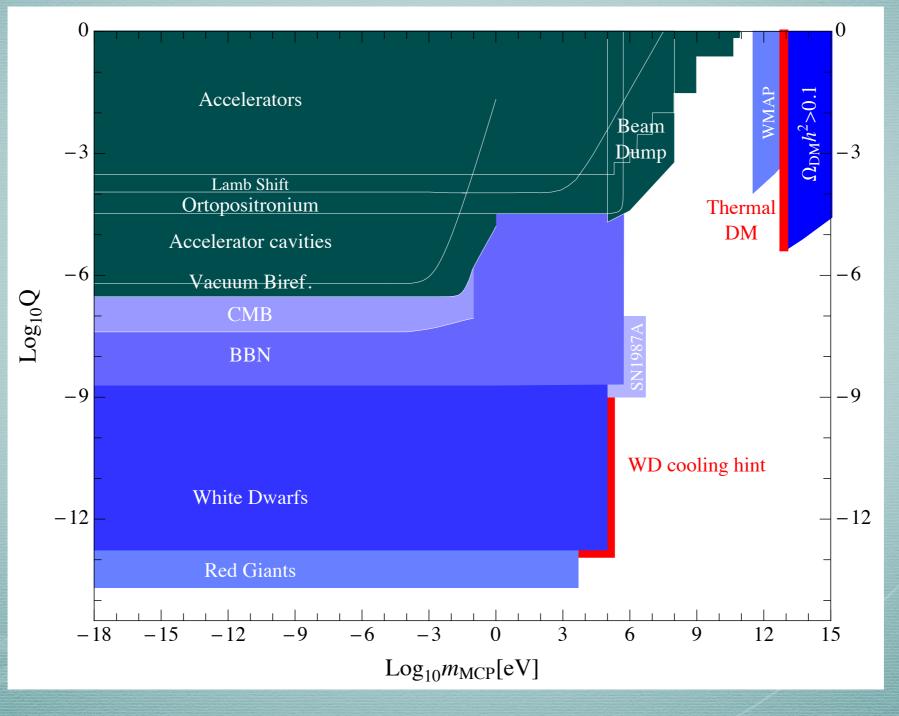
#### Anson Hook Stanford

1705.01107 : A.H. & Junwu Huang

# New light particles

- Old particle / new interactions
  - Fifth force
- New particle / old interactions
  - Gravity
  - Electricity and magnetism
- So what about new particles that carry charge and mass?

# Milli-charged particles

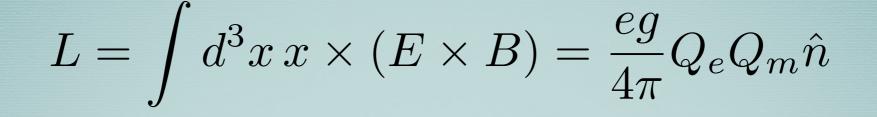


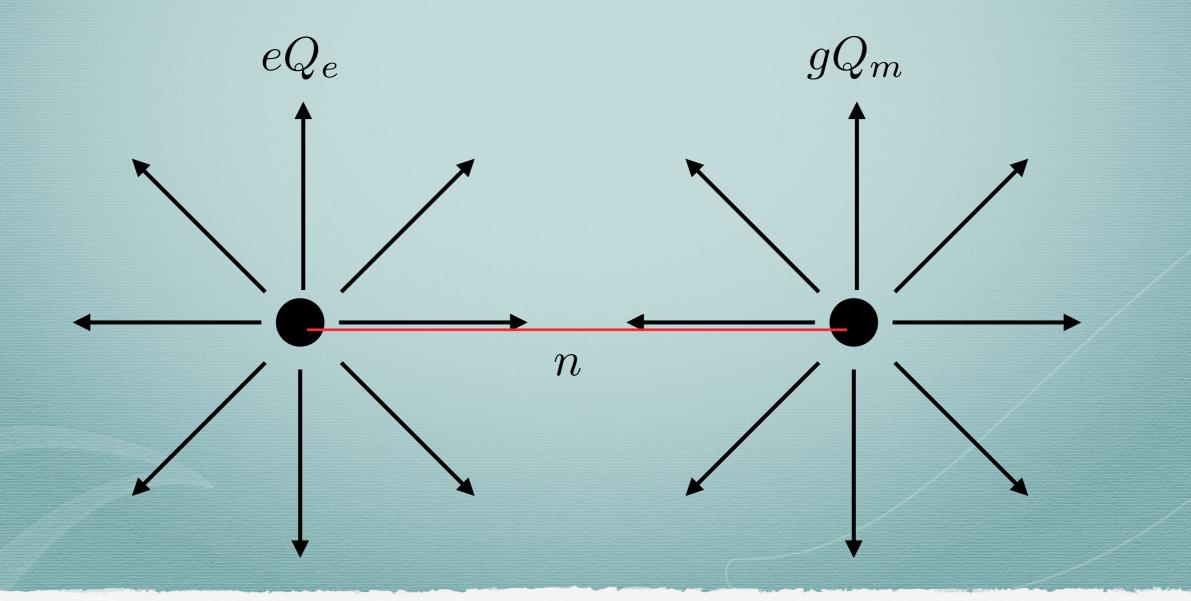
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# Milli-charged particles

- Two types of charges under the photon
- Electric and magnetic (hence E+M)
- What about milli-magnetic charged particles?
  - Quantization of angular momentum

- Electric charge can be quantized in units of electron charge but quarks have fractional charge
- Will evade constraints in a similar manner
  - Physical string
  - All finite energy objects have quantized charges





$$L = \int d^3x \, x \times (E \times B) = \frac{eg}{4\pi} Q_e Q_m \hat{n}$$

$$g = \frac{2\pi}{e}$$
  $L = \frac{\mathbb{Z}}{2}$   $Q_e Q_m = \mathbb{Z}$ 

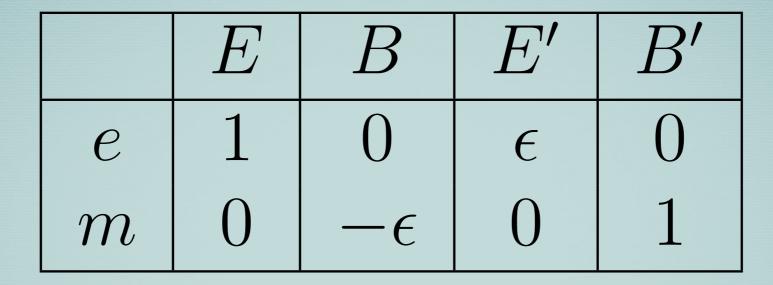
- Given that electrons exist, there exists a minimum magnetic charge
- Most work directed towards this mimimim charge

$$L = \sum_{i} \int d^3x \, x \times (E_i \times B_i) = \sum_{i} \frac{e_i g_i}{4\pi} Q_e^i Q_m^i \hat{n}$$

$$\sum_{i} Q_e^i Q_m^i = \mathbb{Z}$$

- Generalization to multiple U(1)
- Can have charges which are subminimal!
- Rest of angular momentum carried by other U(1) field

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X				
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 Angular momentum in U(1) canceled by angular momentum in U(1)'

#### Plausible?

- Milli-magnetic charged particles are possible but are they plausible?
- Claim Just as plausible as any other scenario with dark photons

Kinetic Mixing  $\mathcal{L} = -\frac{1}{A}F^2 - \frac{1}{A}F_D^2 + 2\epsilon F_D F + \frac{1}{2}m_D^2 A_D^2$ 

- Dark photon with kinetic mixing
- Naturally gives epsilon charged matter

Kinetic Mixing

 Maxwell's equations are more natural when dealing with monopoles

$$\partial_{\mu}F^{\mu\nu} - \epsilon\partial_{\mu}F^{\mu\nu}_{D} = eJ^{\nu}$$
$$\partial_{\mu}F^{\mu\nu}_{D} - \epsilon\partial_{\mu}F^{\mu\nu} = m_{D}^{2}A^{\nu}_{D} + e_{D}J^{\nu}_{D}$$
$$\partial_{\mu}\tilde{F}^{\mu\nu} = 0$$
$$\partial_{\mu}\tilde{F}^{\mu\nu}_{D} = g_{D}K^{\nu}_{D}$$

#### $A \to A + \epsilon A_D$

- Photon is not massive
- Use field redefinition to remove kinetic mixing while keeping photons massless

#### $A \to A + \epsilon A_D$

 $\partial_{\mu}F^{\mu\nu} - \epsilon\partial_{\mu}F^{\mu\nu}_{D} = eJ^{\nu}$  $\partial_{\mu}F^{\mu\nu}_{D} - \epsilon\partial_{\mu}F^{\mu\nu} = m_{D}^{2}A^{\nu}_{D} + e_{D}J^{\nu}_{D}$  $\partial_{\mu}\tilde{F}^{\mu\nu} = 0$  $\partial_{\mu}\tilde{F}^{\mu\nu}_{D} = g_{D}K^{\nu}_{D}$ 

- Our electron epsilon charged under their U(1)
- Their monopoles epsilon charged under our U(1)!

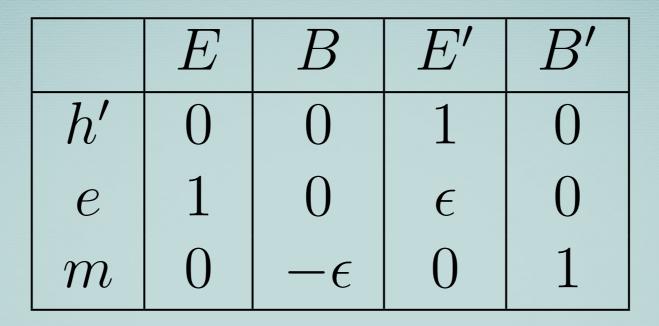
 $\partial_{\mu}F^{\mu\nu} = eJ^{\nu}$  $\partial_{\mu}F^{\mu\nu}_{D} = m_{D}A^{\nu}_{D} + e_{D}J^{\nu}_{D} + \epsilon eJ^{\nu}$  $\partial_{\mu}\tilde{F}^{\mu\nu} = -\epsilon g_{D}K^{\nu}_{D}$  $\partial_{\mu}\tilde{F}^{\mu\nu}_{D} = g_{D}K^{\nu}_{D}$ 

	E	B	E'	B'
e	1	0	$\epsilon$	0
m	0	$-\epsilon$	0	1

Exactly the charge assignments mentioned earlier

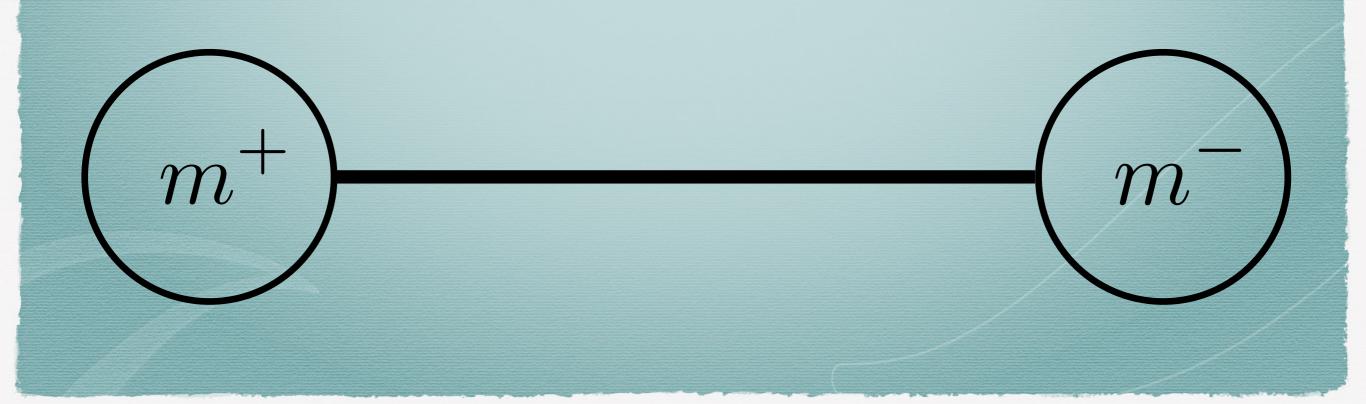
 $\mathcal{L} = -\frac{1}{4}F^2 - \frac{1}{4}F_D^2 + 2\epsilon F_D F + \frac{1}{2}m_D^2 A_D^2$ 

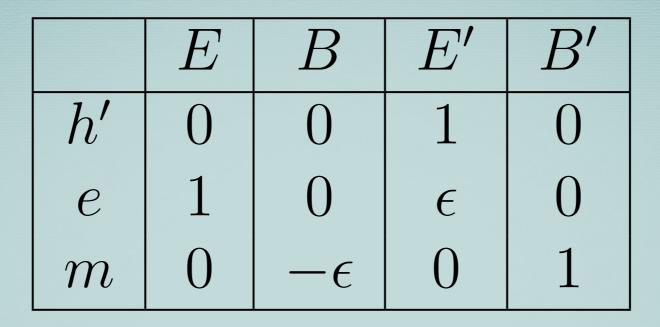
 What is the physical picture of this Lagrangian?



- Milli-magnetic charge
- Dark U(1) Higgsed so monopoles are confined
  - Like QCD, strings connecting monopole with anti-monopole

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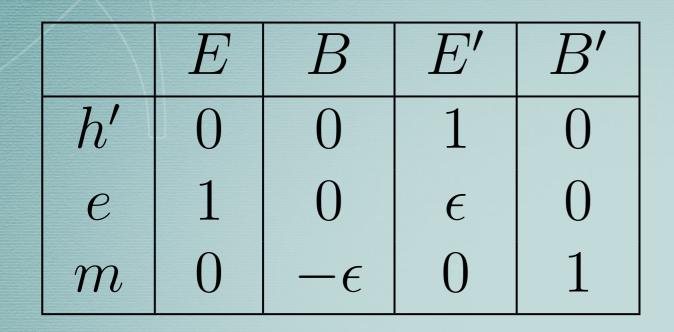




If dark photon mass is irrelevant

- Electron generates a field  $E, B, E' = \epsilon E, B' = \epsilon B$
- Monopole feels a field

 $E_{\text{eff}} = -\epsilon E + E' = 0$  $B_{\text{eff}} = -\epsilon B + B' = 0$ 



- Electron generates a field
- Monopole feels a field

$$B_{\text{eff}} = \epsilon B (e^{-m_{A'}r} - 1)$$

 As long as distance long enough that dark magnetic field is screened, then electron and monopole can interact  $B' = \epsilon B e^{-m_{A'}r}$ 

### Constraints

- If magnetic charges already exist around us
  - Astrophysical constraints/experimental constraints/...
- If they are not surrounding us
  - 1. No constraints if they are non-perturbative ('t Hooft Polyakov monopoles)
  - 2. Weak constraints if fundamental

#### Parker bound

$$\mathcal{F}_{\text{Parker}} = 10^{-16} \,\mathrm{cm}^{-2} \,\mathrm{s}^{-1} \,sr^{-1}$$

Magnetic fields accelerate monopoles

- Energy in monopoles comes from B field
- If too much energy is taken, then B field of Milky Way neutralized
- Usually applied to some cosmological abundance



 Obtain model independent bounds on millimagnetic charged particles



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- Perturbative production can be exponentially suppressed
- Need non-perturbative production



- Obtain model independent bounds on millimagnetic charged particles
- Perturbative production can be exponentially suppressed
- Need non-perturbative production
- Exponentially large number of initial states (photons)
- Extremely large electric and magnetic fields

# Schwinger pair production

 Production of electric (magnetic) particles in a strong electric (magnetic) field

$$\frac{P}{Vt} = \frac{e^2 E^2}{4\pi^3} e^{\frac{-\pi m^2}{eE}}$$

 $4\pi$ 

$$\frac{P}{Vt} = \frac{\epsilon^2 g^2 B^2}{8\pi^3} e^{\frac{-\pi m^2}{\epsilon g B} + \frac{g^2}{4}}$$

# Schwinger pair production

- We have monopoles connected by strings
- Easy to modify pair production to account for strings

$$\frac{P}{Vt} \sim e^{-\frac{\pi m^2}{eE - m_A^2}}$$

As long as string tension smaller than
ɛ Electric/Magnetic field then pair
production proceeds as before

# Schwinger pair production

- Unsuppressed production of milli-magnetic charged particles if there is a large magnetic field
- Largest magnetic fields in the universe are at magnetars
- Production of magnetic particles neutralizes magnetic field
  - Require that it is not neutralized over the lifetime of magnetar

# Magnetars

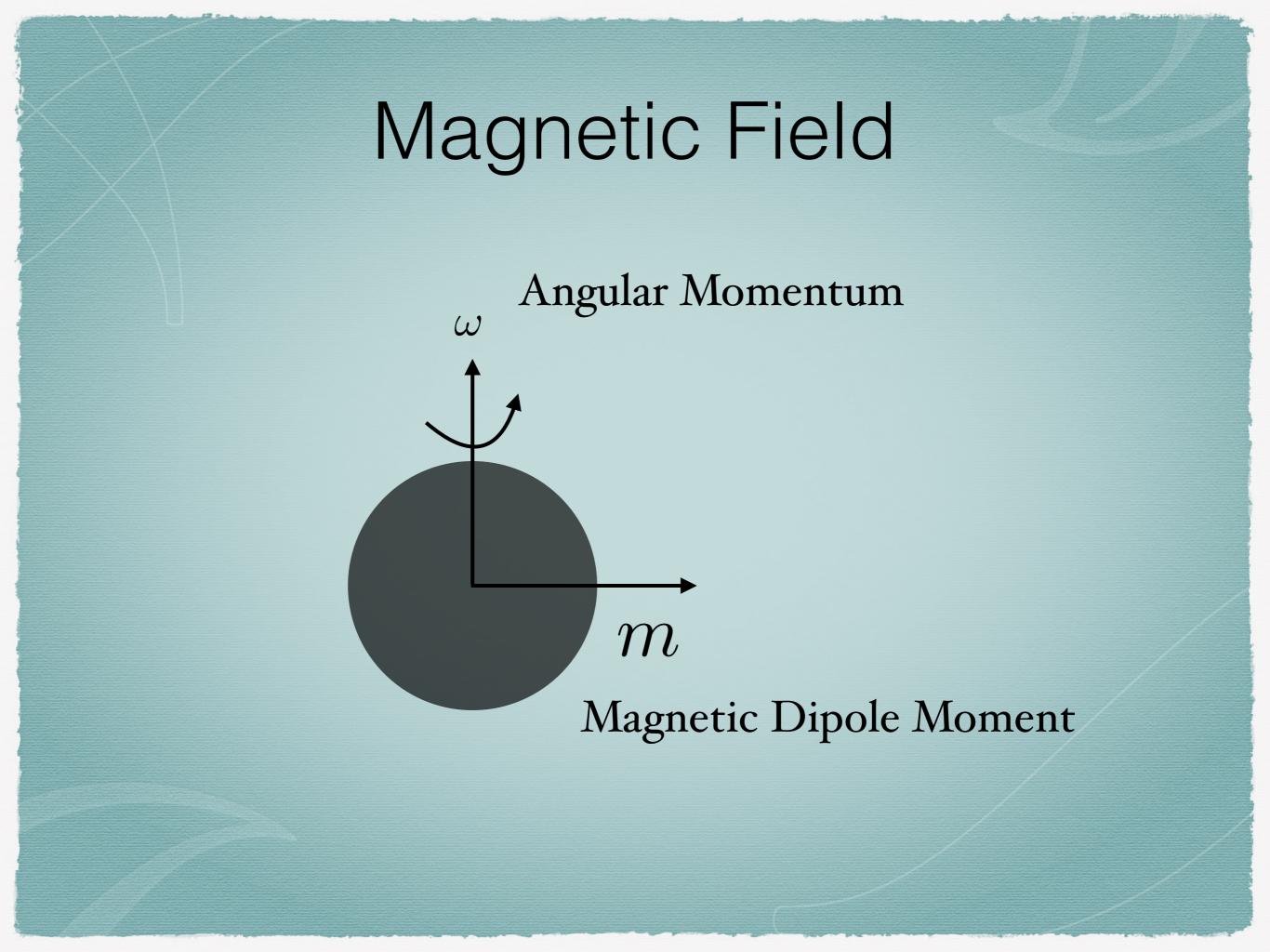
- Neutron stars with extremely large magnetic fields
  - Size ~ 10 km
  - B ~ 10<sup>13-16</sup> gauss ~ MeV<sup>2</sup>
  - Age ~ 10<sup>3-5</sup> years
  - Luminosity (persistent x-rays) ~ 10<sup>33-36</sup> ergs/s ~ B<sup>2</sup> V/t
  - ~ 20 observed, ~ kiloparsec away
  - Not much known about them

# Magnetars

- Anomalous x-ray pulsars
  - Emit soft x rays
  - Anomalous because not powered by standard means
- Soft gamma-ray repeaters
  - Peak luminosity larger than Eddington limit

# Magnetic Field

- Evidence for magnetic field is from soft gamma-ray bursts
  - Strong magnetic fields allow for super Eddington luminosity emissions
  - Fall off of burst depends on magnetic instabilities of the magnetar
- Crude estimate of the magnetic field can be made via loss of angular momentum



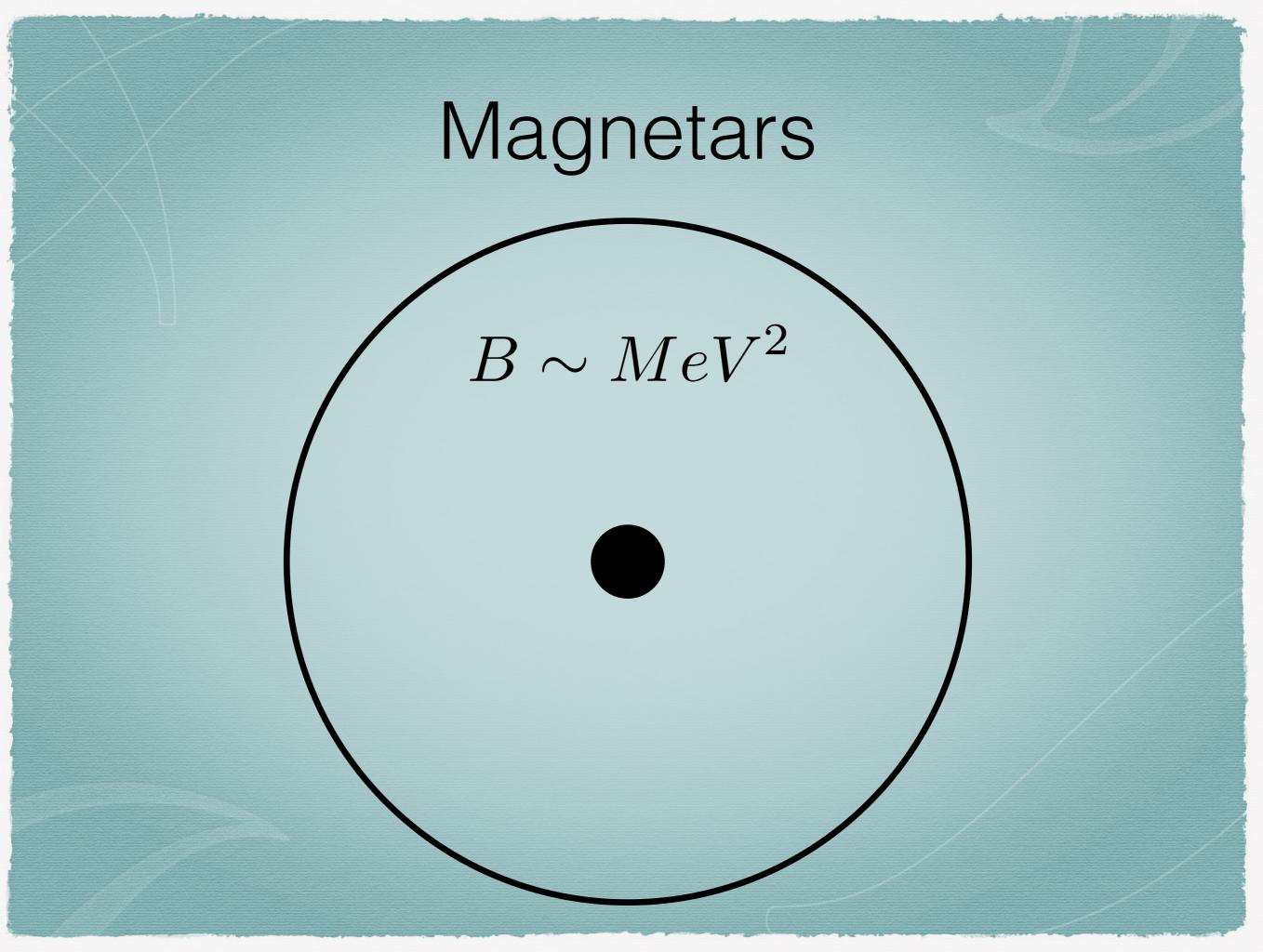
### Magnetic Field

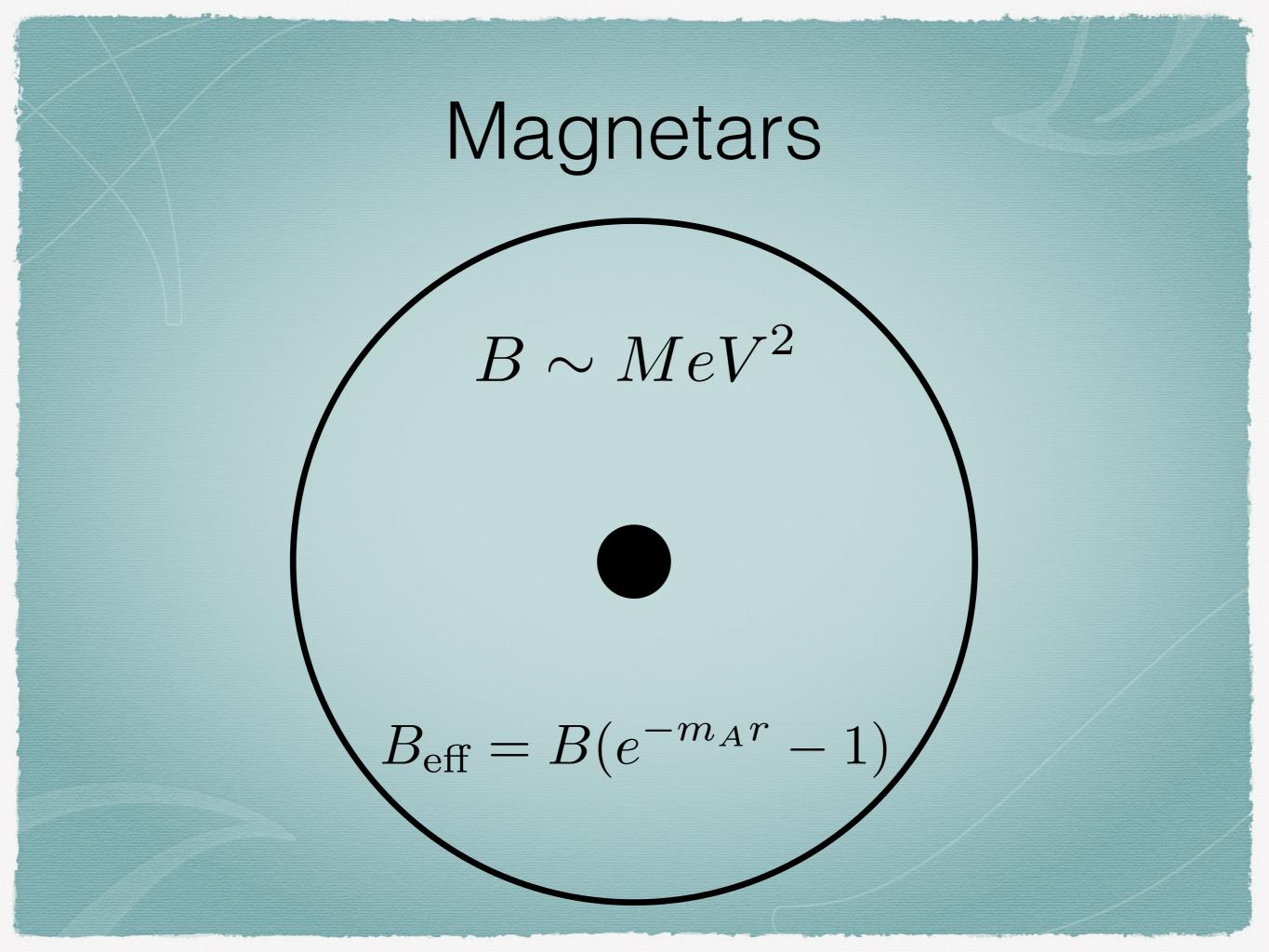
$$B \approx 3 \times 10^{19} \sqrt{\frac{P}{\text{second}}} \dot{P} \text{Gauss}$$

- Observed periods of a few seconds
- dP/dt ~ 10<sup>-11</sup>
- B ~ 10<sup>15</sup> gauss ~ MeV<sup>2</sup>
- More refined estimates change results by only O(1) factors

# Magnetars

- Given huge uncertainties, we will take the following values
  - Radius = 10 km
  - B = 10<sup>15</sup> gauss
  - Age =  $10^4$  years





#### Bound

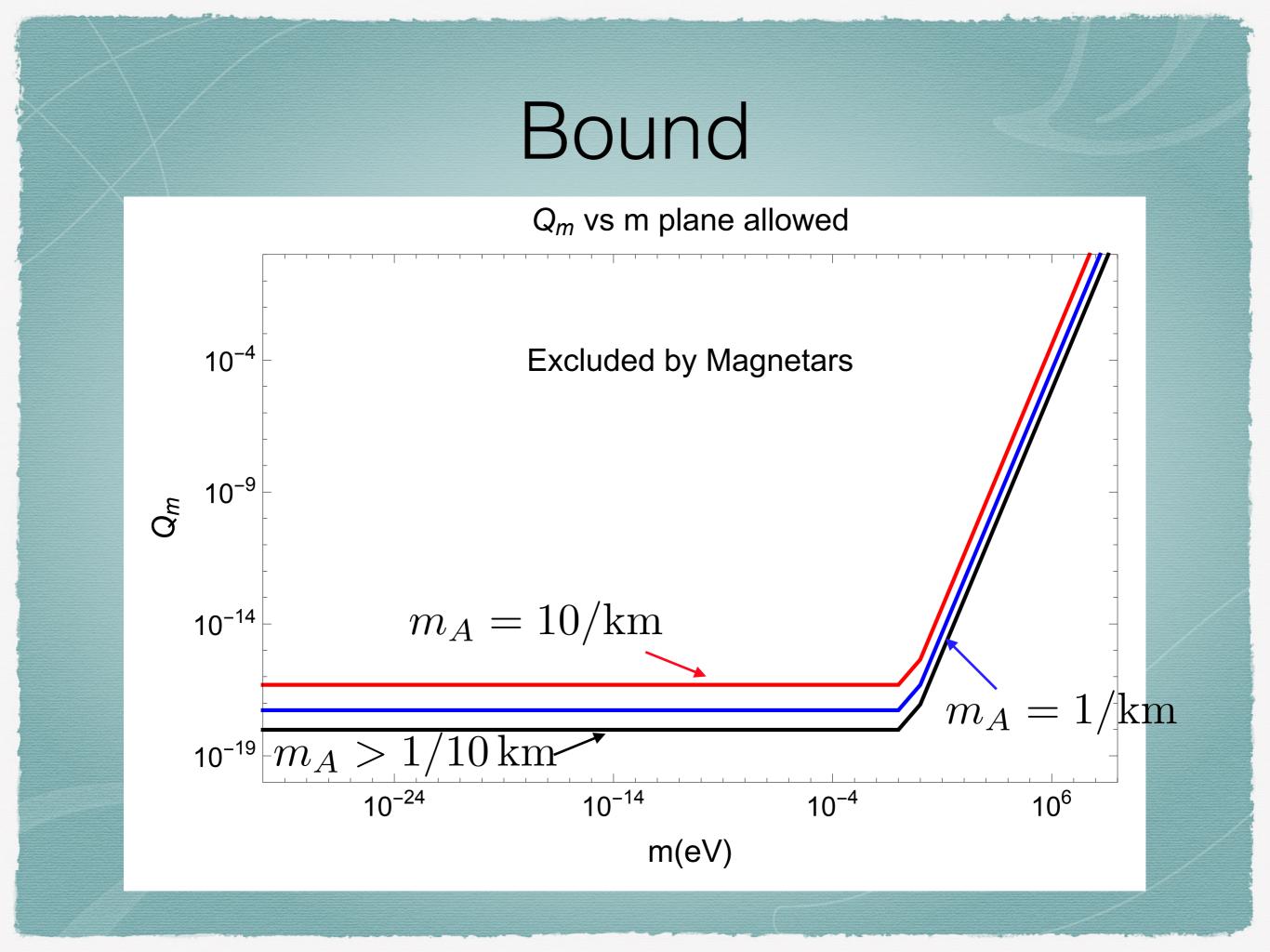
$$E_{\rm loss} = 2Q_m g B_{\rm eff} r$$

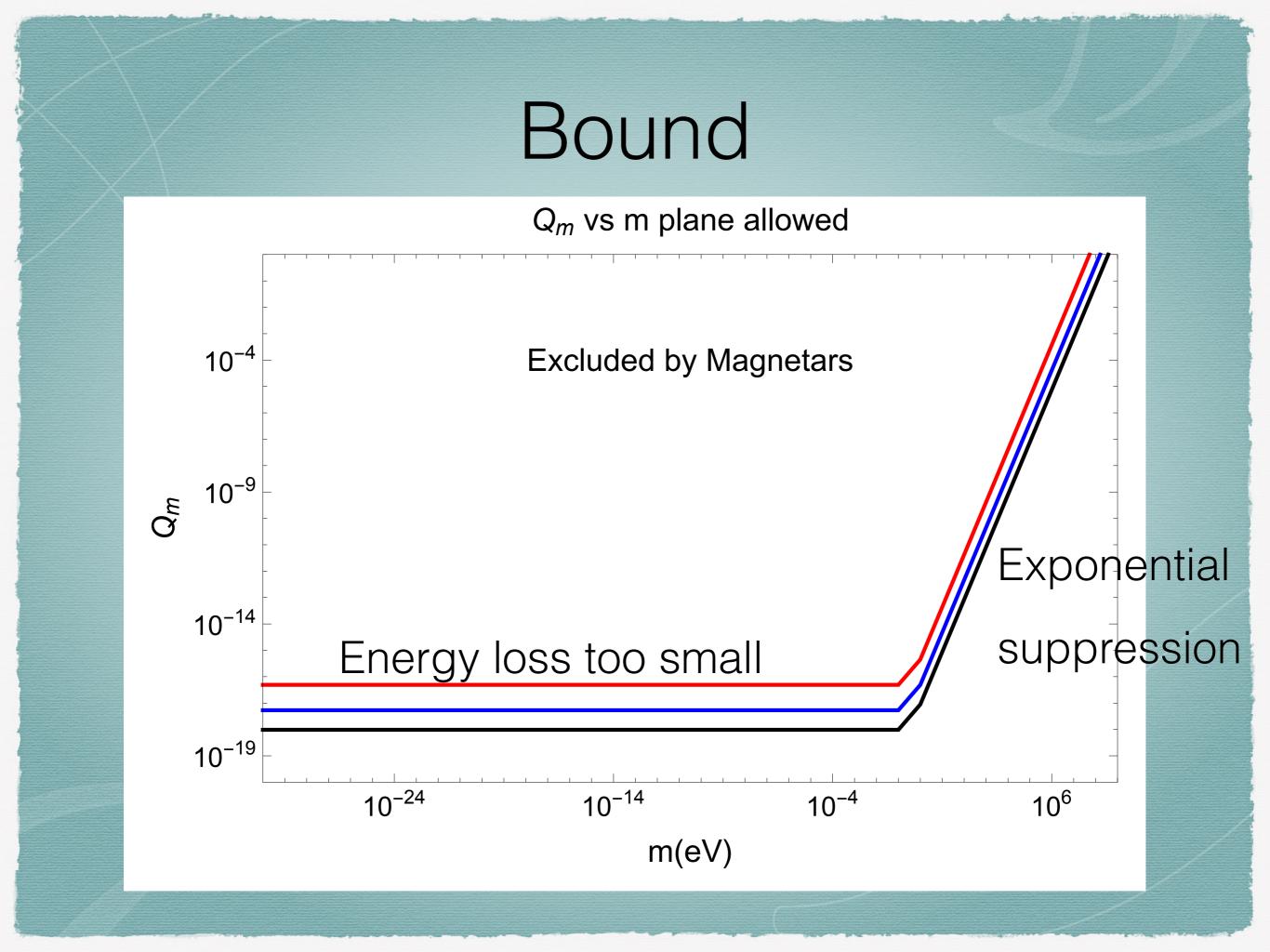
- Energy loss due to production of monopoles
- After pair production, monopoles carry away energy

### Bound

$$\frac{dE}{dt\,dV} = \frac{Q_m^2 g^2 B_{\text{eff}}^2}{4\pi^3} e^{\frac{-\pi m^2}{Q_m g B_{\text{eff}}}} E_{\text{loss}} \lesssim \frac{B^2}{2t_{\text{lifetime}}}$$

- Magnetic field is not neutralized by pair production during the lifetime of magnetar
- Equivalent to saying that total energy loss must be smaller than observed energy loss





### Conclusion

- Just like there can be milli-electric charged particles, there can be milli-magnetic charged particles
- Currently NO model independent bounds
  - Model dependent bounds are weak and often come from cosmology
- Magnetars have large magnetic fields that let one place new very strong model independent bounds