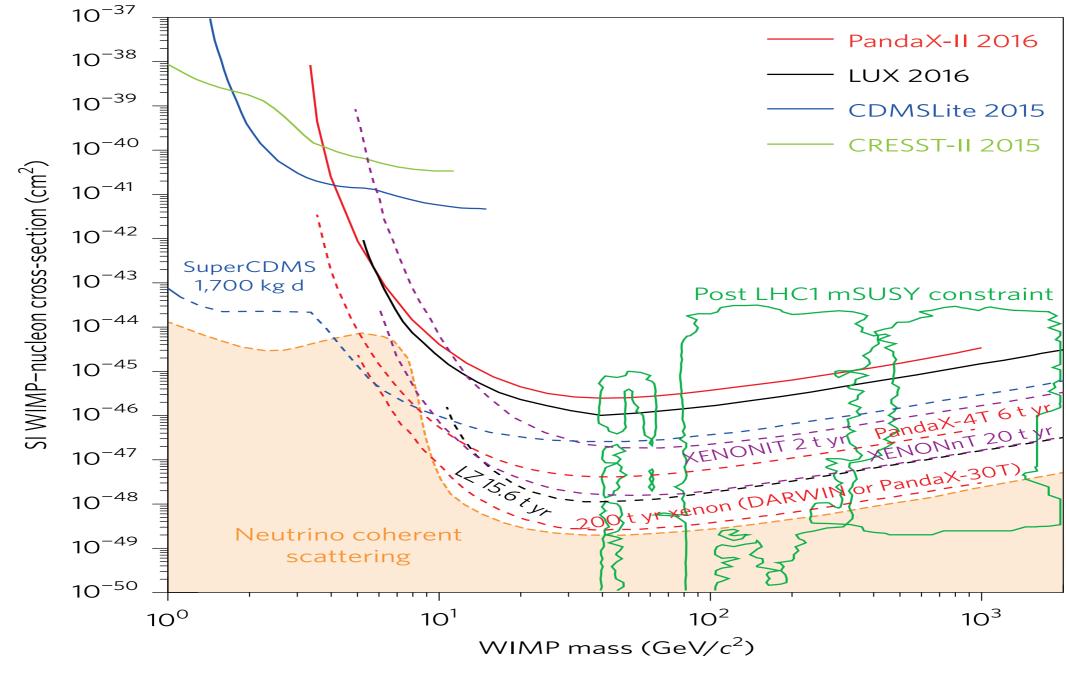
Discovering Dark Matter at High Nuclear Recoil

Adam Martin (<u>amarti41@nd.edu</u>) University of Notre Dame



CERN-CKC, June 6th, 2017

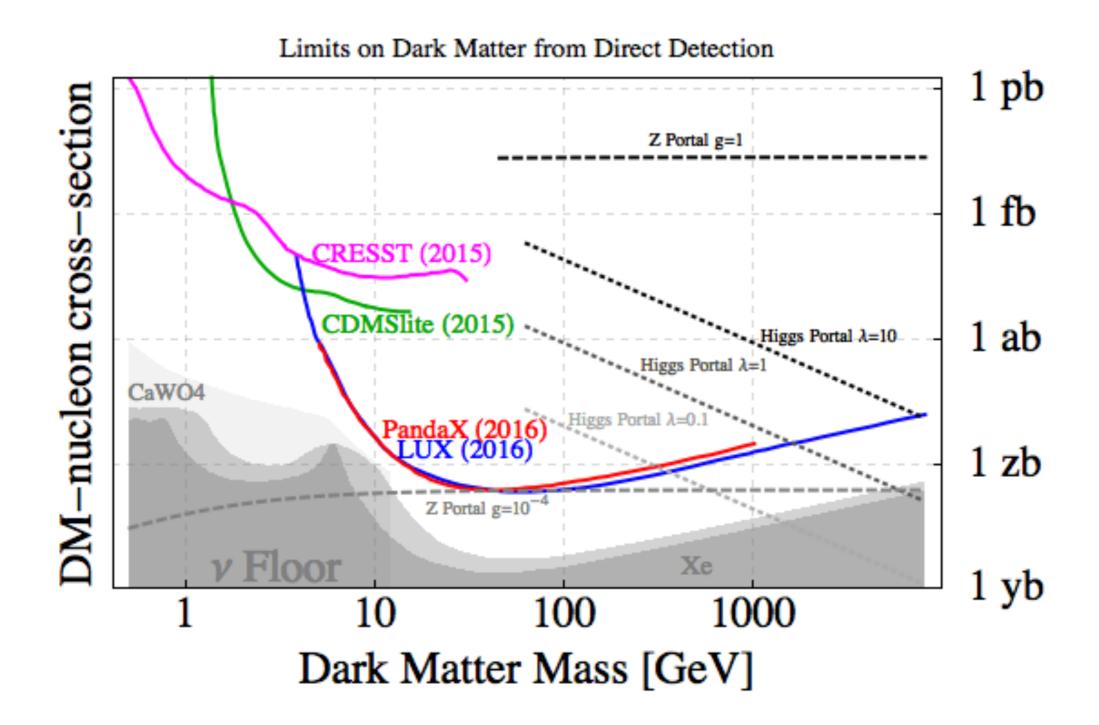




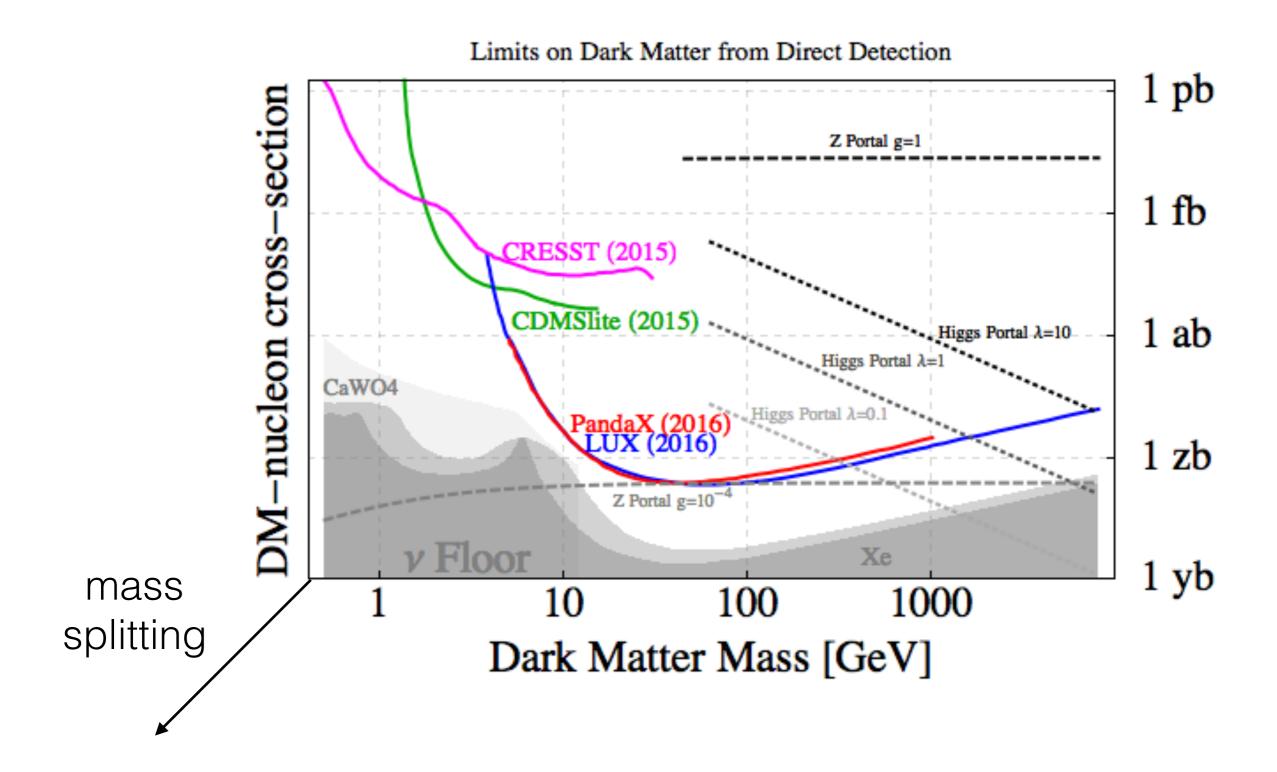
[Liu, Chen, Ji '17]

(not shown: latest Xe100 (1609.06154) and XENON1T (1705.06655) results)

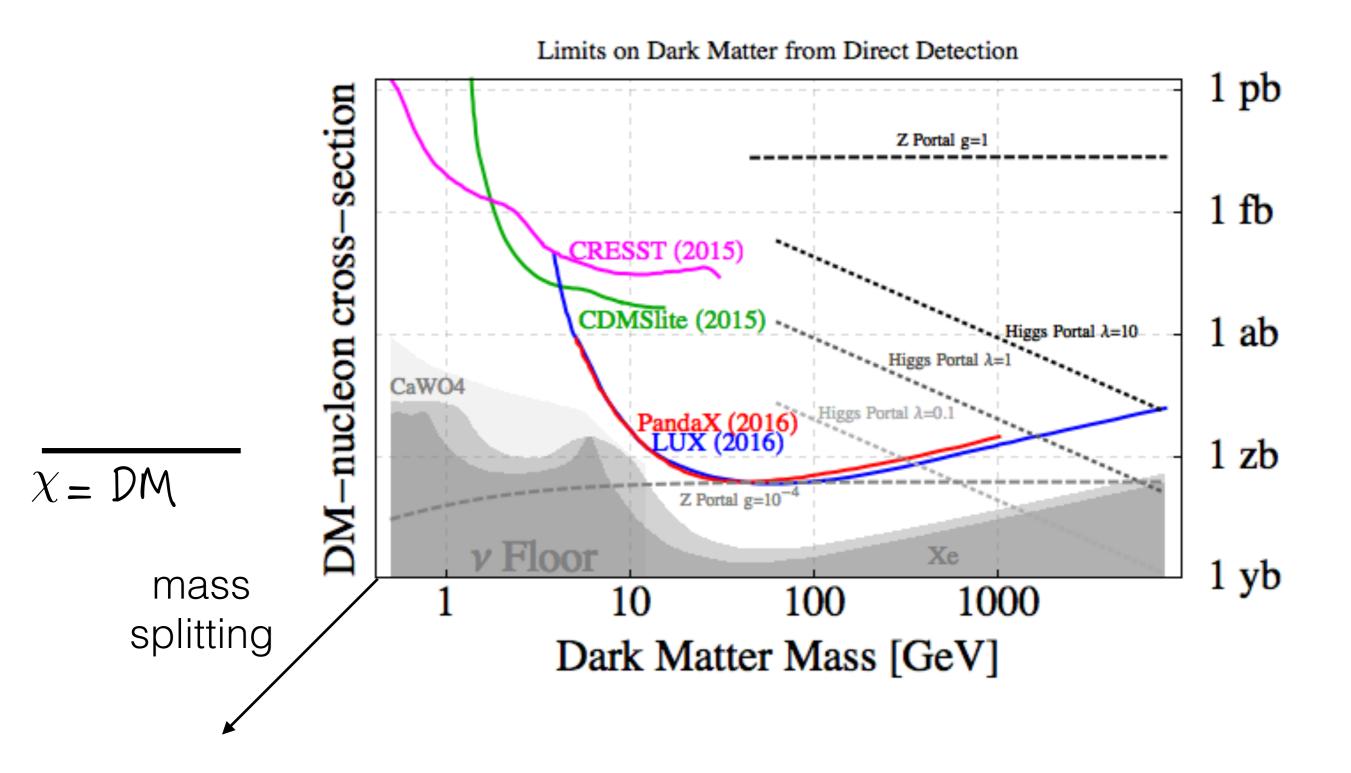
simplest WIMP DM is running out of room



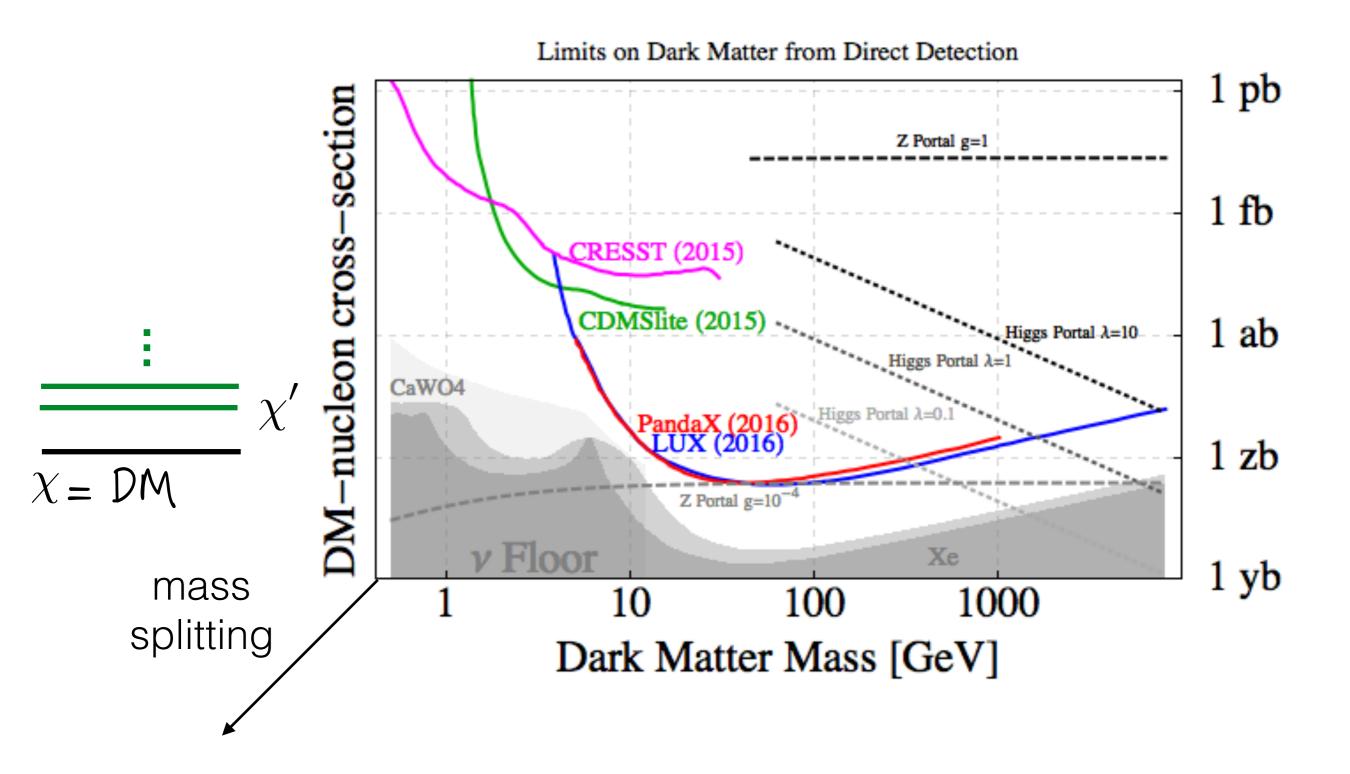
third direction: inelasticity

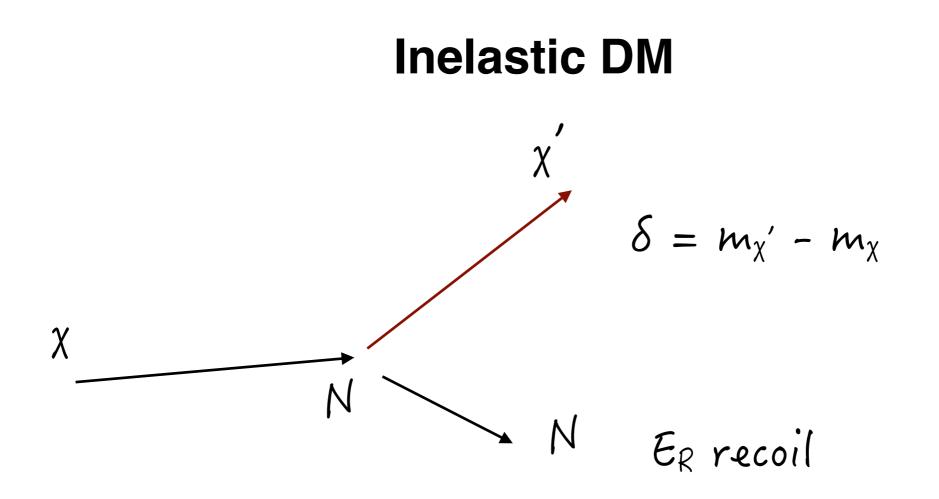


third direction: inelasticity

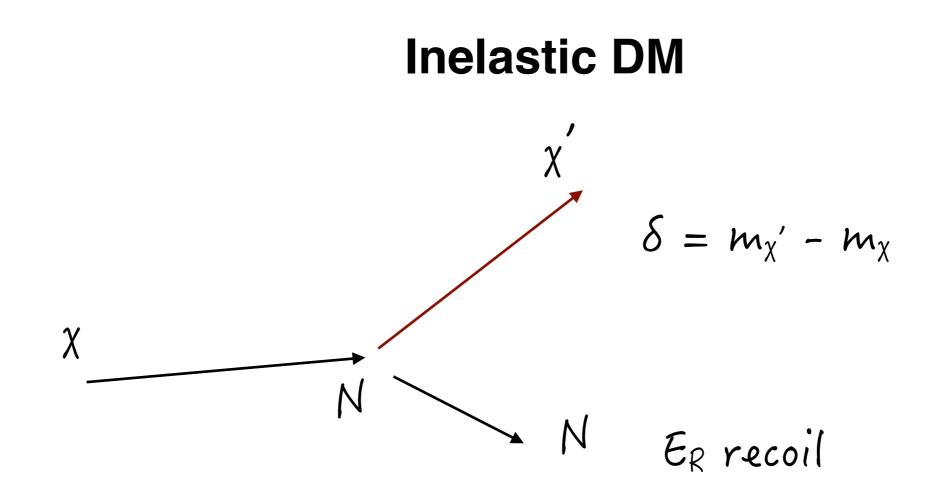


third direction: inelasticity





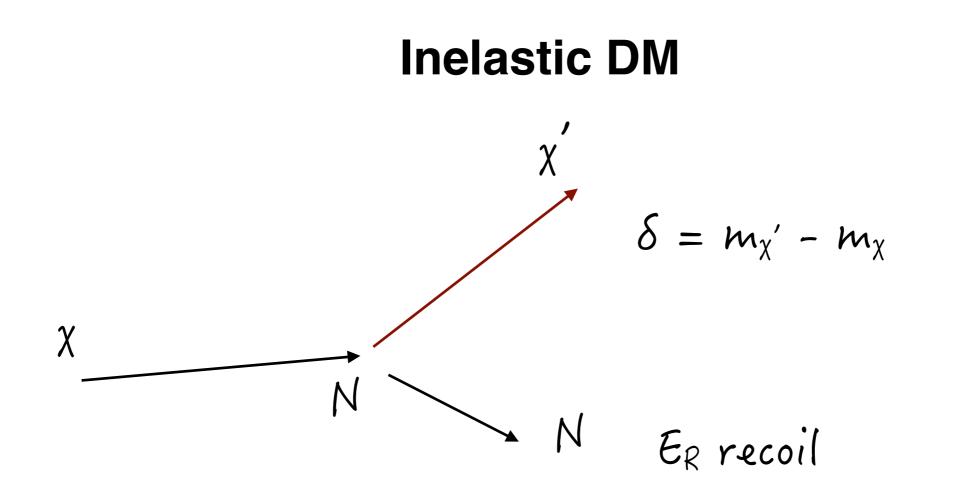
[L. J. Hall, T. Moroi and H. Murayama '97, Tucker-Smith, Weiner '01]



(target-dependent) minimum velocity required to scatter

$$\begin{split} KE_{\chi} \geq \delta \left(1 + \frac{m_{\chi}}{m_N} \right) \\ \sigma_{inelastic} = \sqrt{1 - \frac{2\,\delta}{\mu_{\chi N}\,v^2}} \sigma_{elastic} \end{split}$$

[L. J. Hall, T. Moroi and H. Murayama '97, Tucker-Smith, Weiner '01]

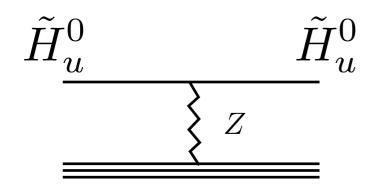


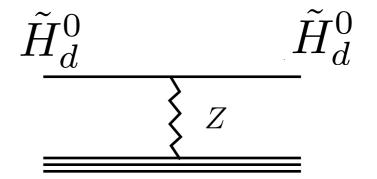
popularized to reconcile DAMA with CDMS (2001-) required δ ~ 100 keV for m_X ~ 100 GeV

forgetting DAMA, range of δ is wide open

for canonical DM velocity distribution, available KE ≤ 650 keV

(nearly) pure Higgsinos: $\mu \ll M_1, M_2$





(nearly) pure Higgsinos: $\mu \ll M_1, M_2$

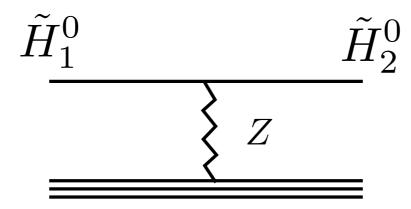


 $ilde{H}^0_u,\, ilde{H}^0_d
ightarrow ilde{H}^0_1,\, ilde{H}^0_2 \,\,$ once we turn on EWSB

(nearly) pure Higgsinos: $\mu \ll M_1, M_2$



 $\tilde{H}^0_u,\,\tilde{H}^0_d\to\tilde{H}^0_1,\,\tilde{H}^0_2$ once we turn on EWSB



Z-exchange inelastic

achieve right relic abundance for $m_{\tilde{H}} \sim 1.1 \text{ TeV}$

achieve right relic abundance for $m_{\tilde{H}} \sim 1.1 \mbox{ TeV}$

direct detection:

 $\sigma_{\tilde{H},inelastic} \sim (\text{velocity factor}) \times 10^{-39} \, \text{cm}^2 \times A^4$

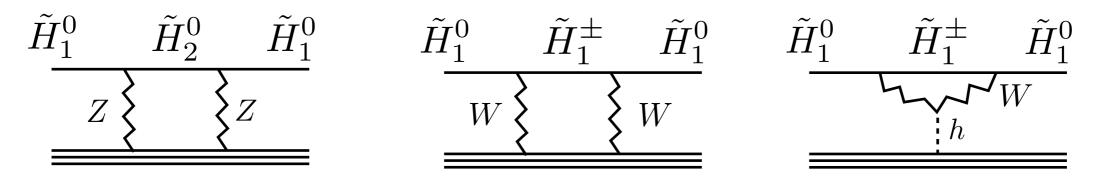
$$\tilde{H}_{2}^{0} \longrightarrow \delta_{\tilde{H}} \qquad \delta_{\tilde{H}} \simeq m_{Z}^{2} \left(\frac{\sin^{2} \theta_{W}}{M_{1}} + \frac{\cos^{2} \theta_{W}}{M_{2}} \right) + \mathcal{O}\left(\frac{1}{M_{1,2}^{2}} \right) = \left(\frac{1}{M_{1,2}^{2}} \right) \left(\frac{1}{M_{1,2}^{2}} + \frac{1}{M_{1,2}^{2}} \right) + \mathcal{O}\left(\frac{1}{M_{1,2}^{2}} \right) = \left(\frac{1}{M_{1,2}^{2}} + \frac{1}{M_{1,2}^{2}} \right) + \mathcal{O}\left(\frac{1}{M_{1,2}^{2}} \right) = \left(\frac{1}{M_{1,2}^{2}} + \frac{1}{M_{1,2}^{2}} \right) + \mathcal{O}\left(\frac{1}{M_{1,2}^{2}} \right) = \left(\frac{1}{M_{1,2}^{2}} + \frac{1}{M_{1,2}^{2}} \right) + \mathcal{O}\left(\frac{1}{M_{1,2}^{2}} \right) = \left(\frac{1}{M_{1,2}^{2}} + \frac{1}{M_{1,2}^{2}} + \frac{1}{M_{1,2}^{2}} \right) + \mathcal{O}\left(\frac{1}{M_{1,2}^{2}} + \frac{1}{M_{1,2}^{2}} + \frac{1}{M_{1,2}^{2}} \right) = \left(\frac{1}{M_{1,2}^{2}} + \frac{1}{$$

achieve right relic abundance for $m_{\tilde{H}} \sim 1.1 \text{ TeV}$

direct detection:

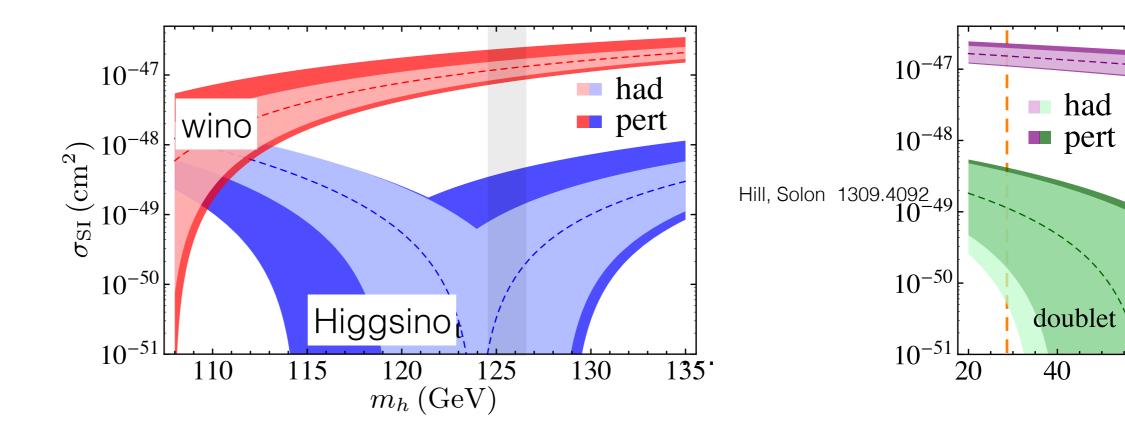
$$\sigma_{\tilde{H},inelastic} \sim (\text{velocity factor}) \times 10^{-39} \text{ cm}^2 \times A^4$$

elastic scattering at loop level: suppressed by m_n or E_R

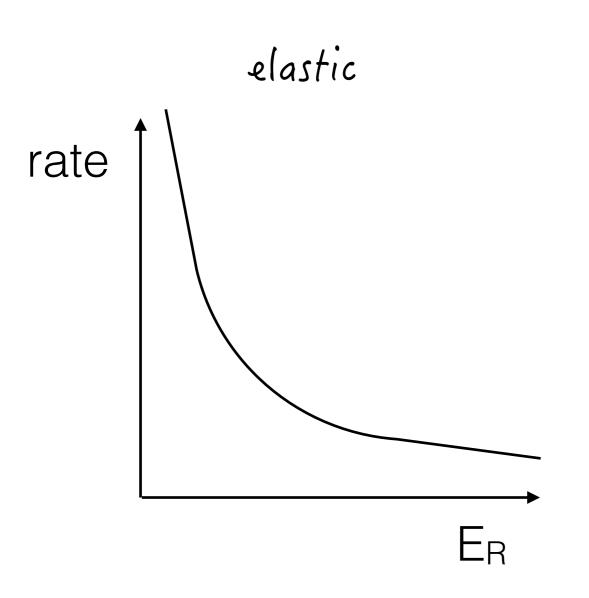


[[]Hisano et al '11, Hill+Solon '13]

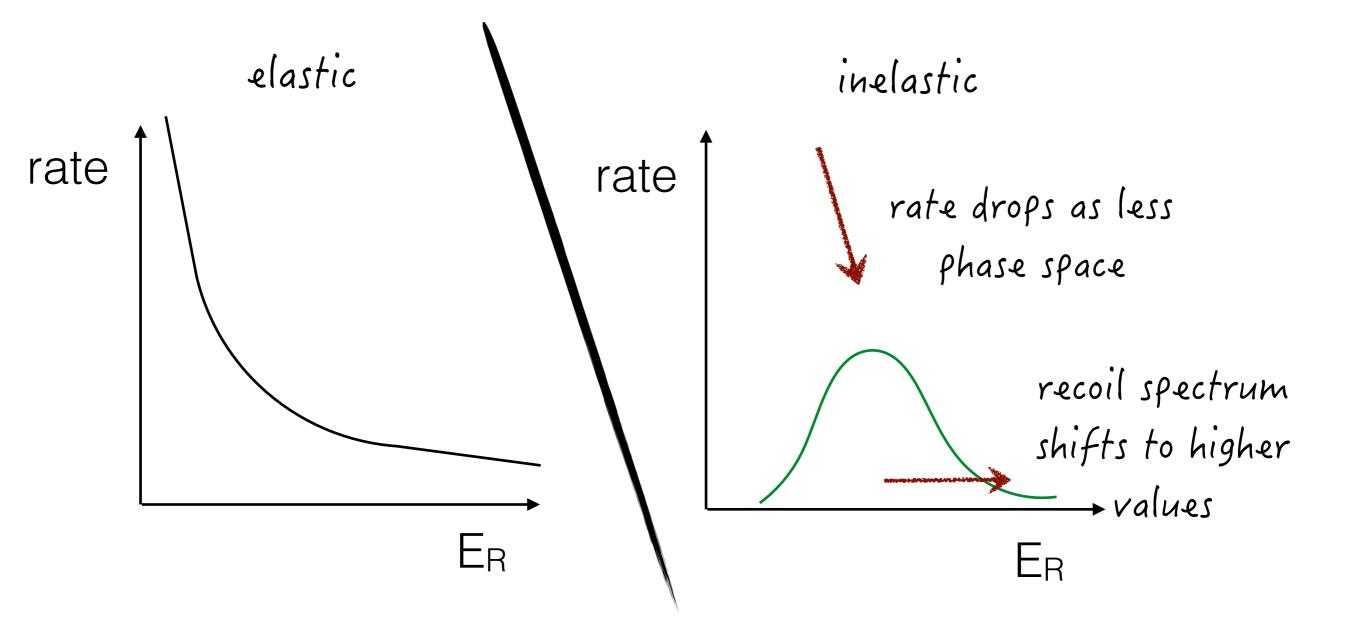
further suppressed by accidental cancellations



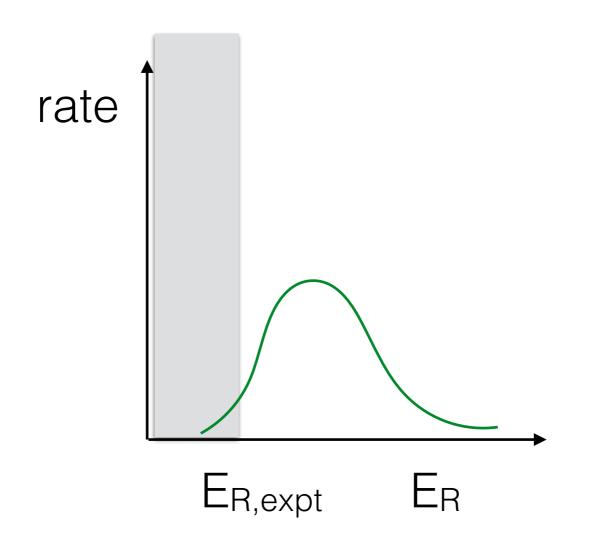
inelasticity changes nuclear recoil energy spectrum



inelasticity changes nuclear recoil energy spectrum

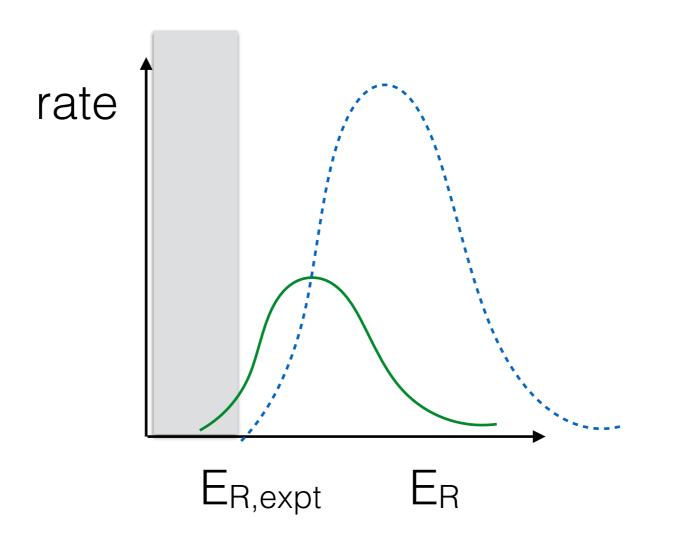


Experimental **signal windows** focused on low E_R

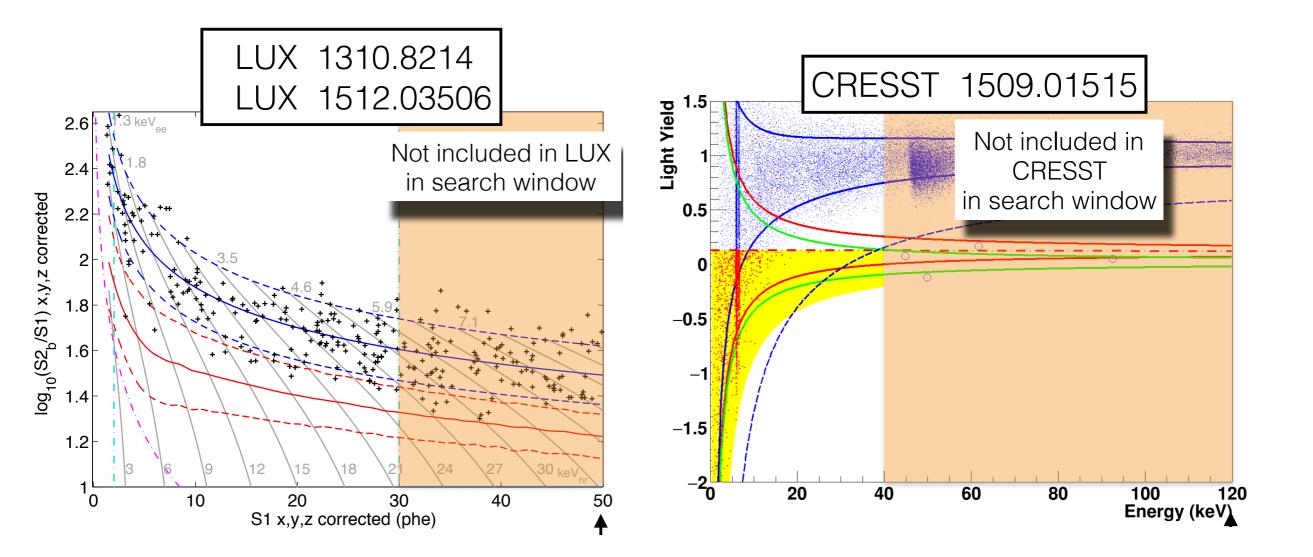


as a result, blind to sufficiently inelastic DM, even if $\sigma_{\!X^{\!N}}$ is large

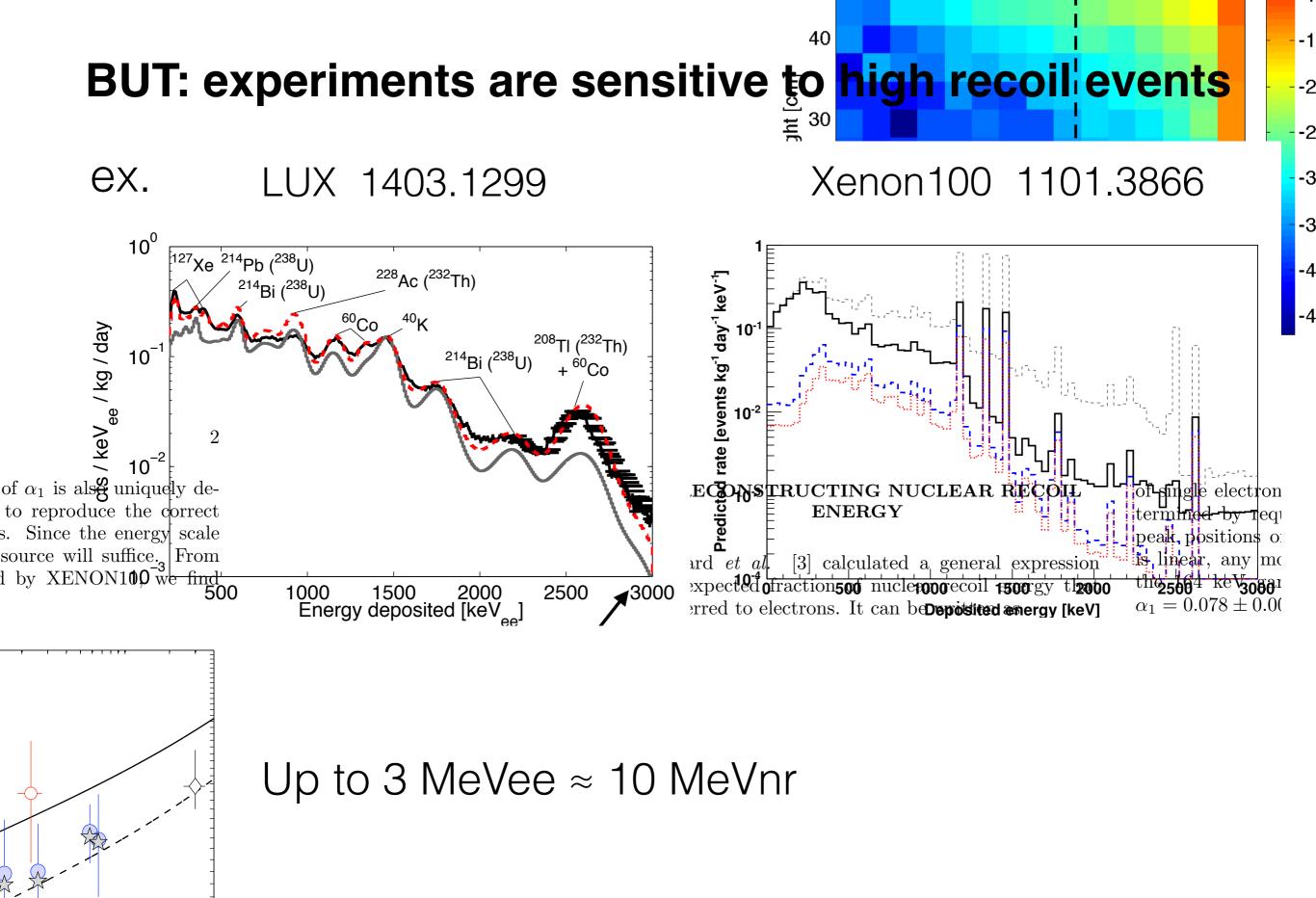
Experimental **signal windows** focused on low E_R

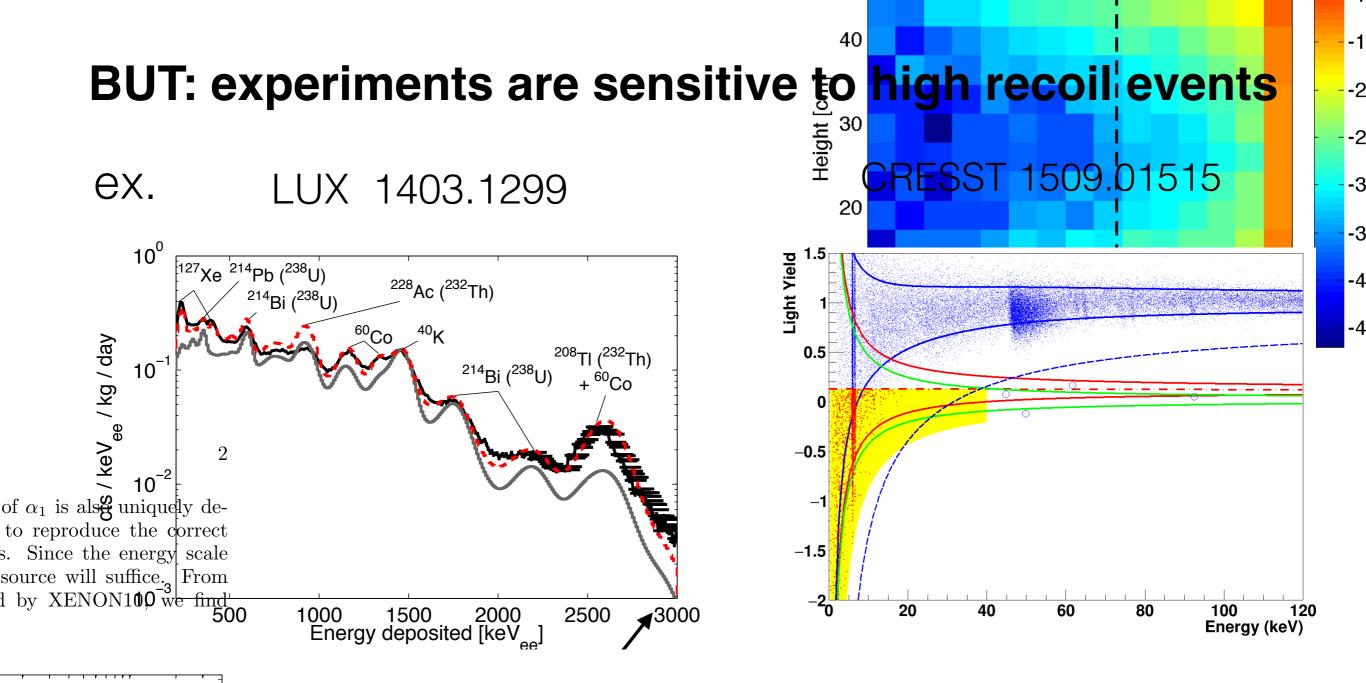


as a result, blind to sufficiently inelastic DM, even if $\sigma_{\!X^{\!N}}$ is large

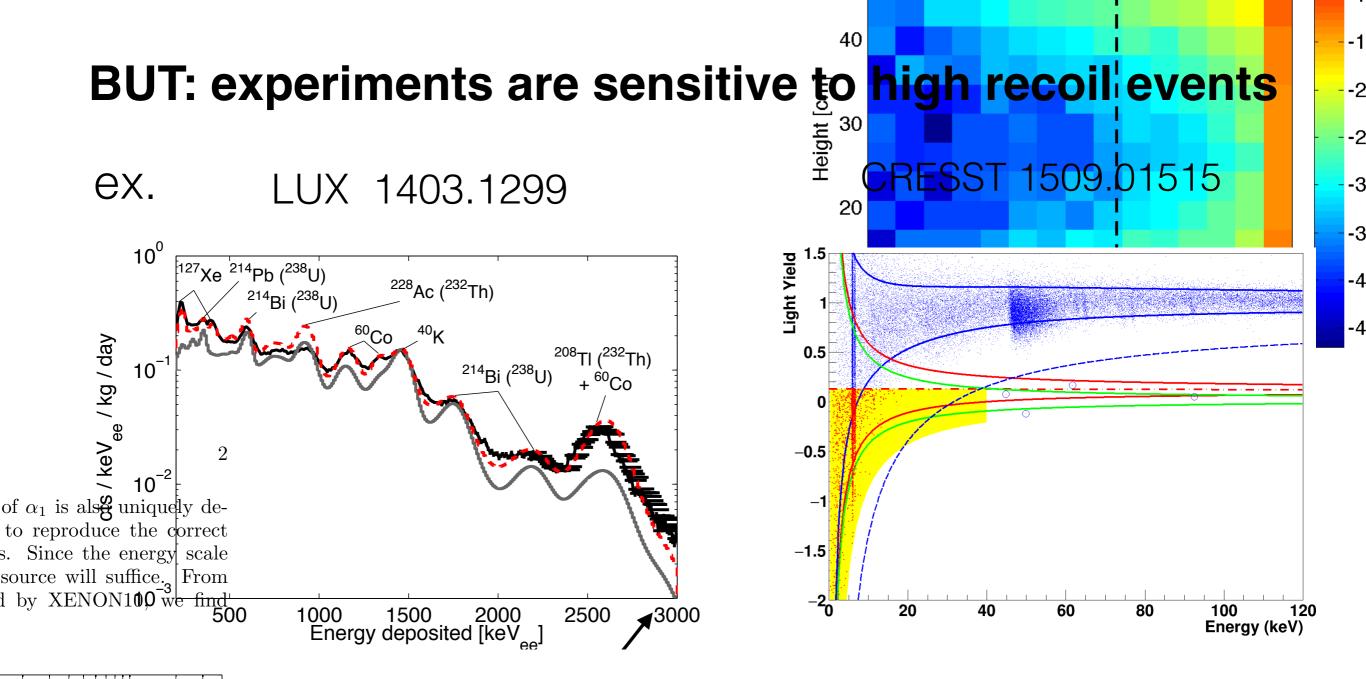


focus of current runs on low E_{R}

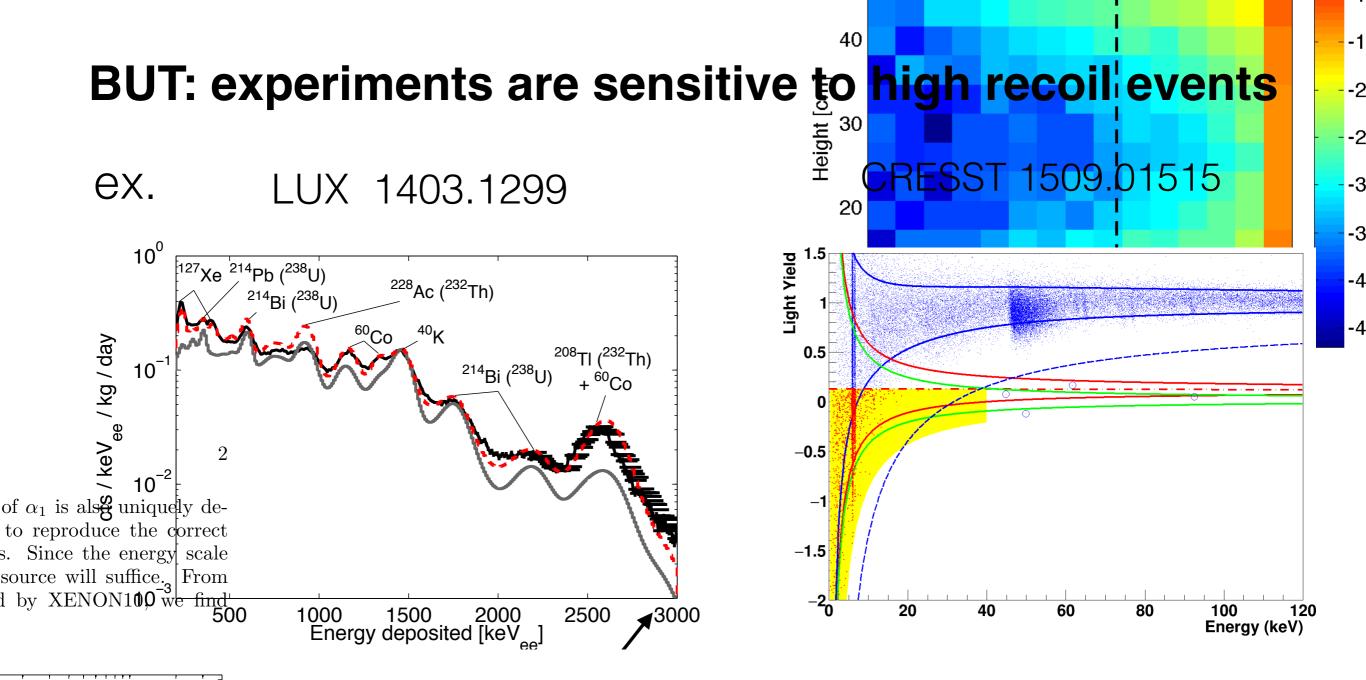




Up to 3 MeVee \approx 10 MeVnr

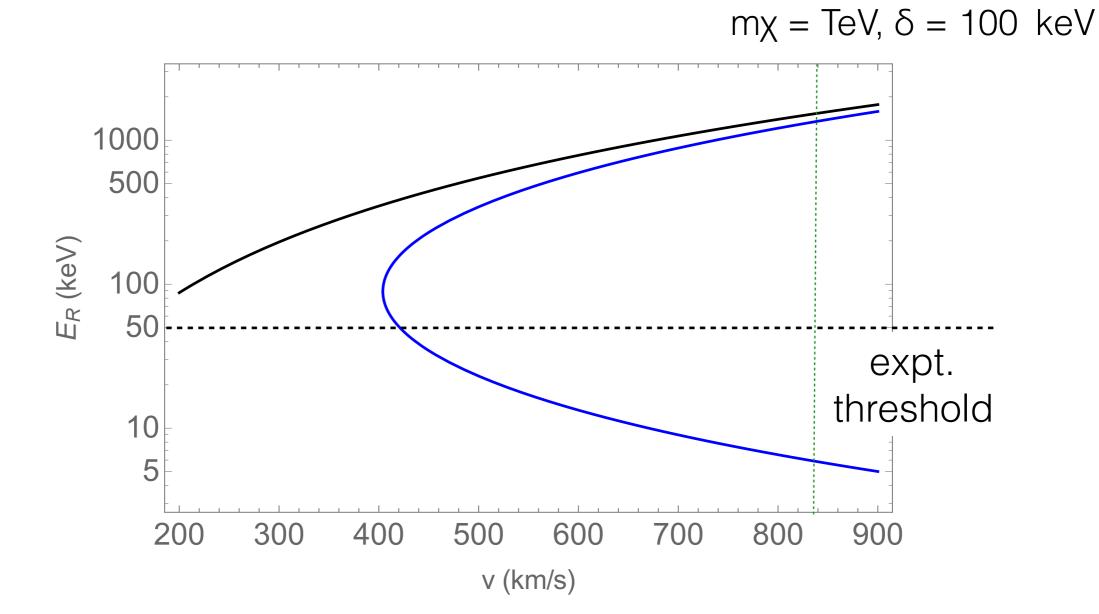


present: inelastic DM could be lurking in high nuclear recoil data of existing experiments: go and look! Up to 3 MeVee ≈ 10 MeVnr

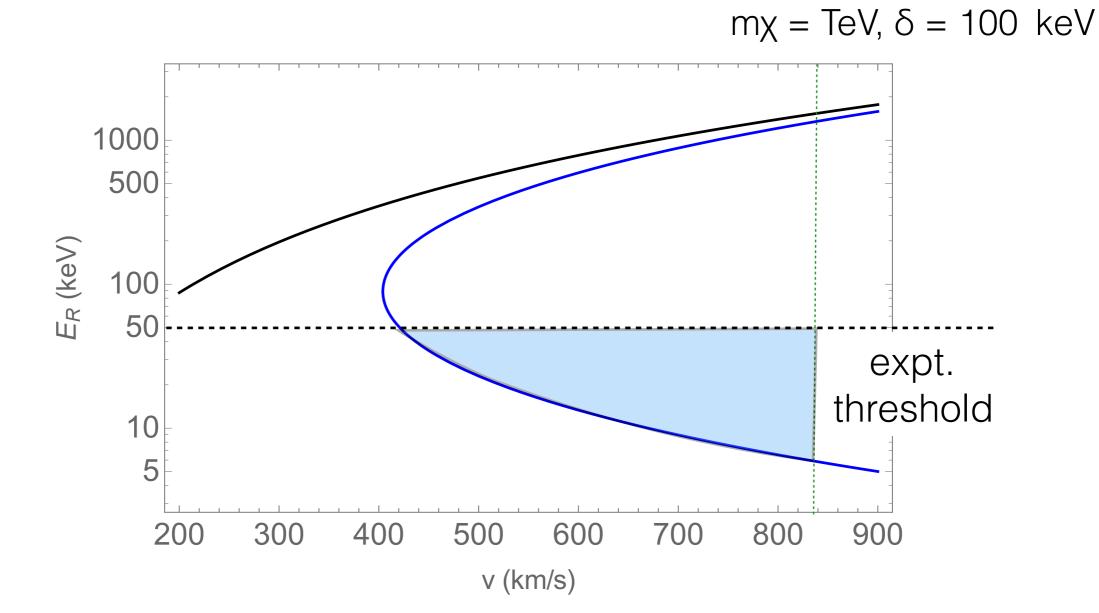


present: inelastic DM could be lurking in high nuclear recoil data of existing experiments: go and look! Up to 3 MeVee ≈ 10 MeVnr

uture: don't limit searches to low-recoil



$$E_R = \frac{\mu_{\chi N}}{m_N} \Big((\mu_{\chi N} v^2 \cos^2 \theta_{lab} - \delta) \pm (\mu_{\chi N} v^2 \cos^2 \theta_{lab})^{1/2} (\mu_{\chi N} v^2 \cos^2 \theta_{lab} - 2\delta)^{1/2} \Big)$$

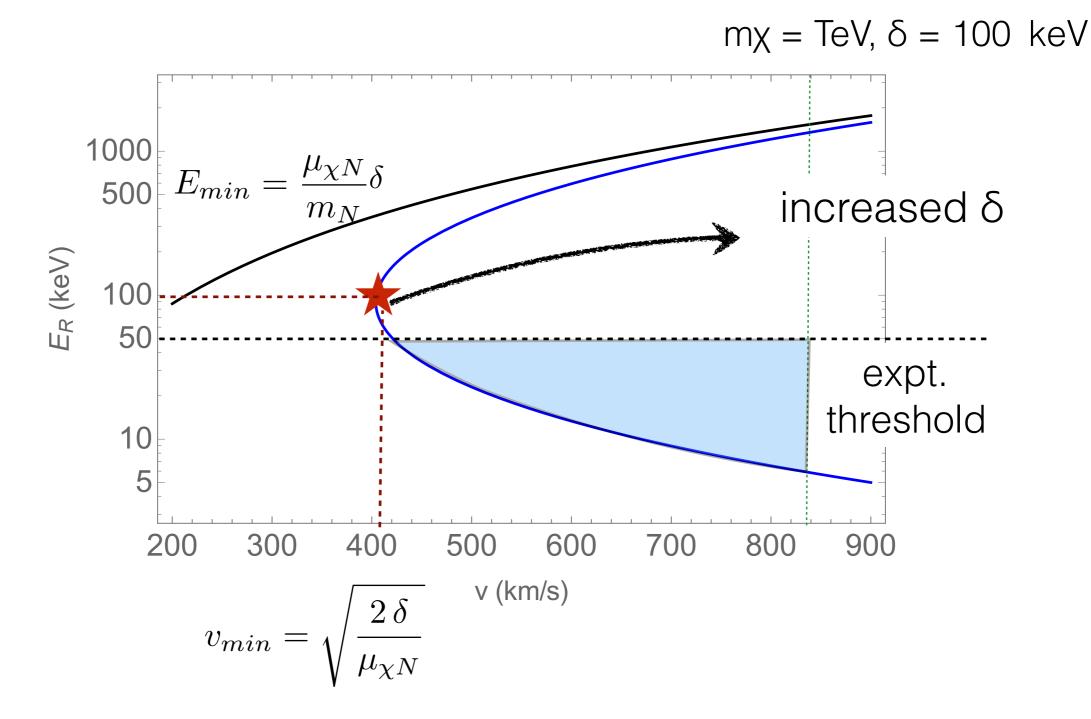


$$E_R = \frac{\mu_{\chi N}}{m_N} \left((\mu_{\chi N} v^2 \cos^2 \theta_{lab} - \delta) \pm (\mu_{\chi N} v^2 \cos^2 \theta_{lab})^{1/2} (\mu_{\chi N} v^2 \cos^2 \theta_{lab} - 2\delta)^{1/2} \right)$$

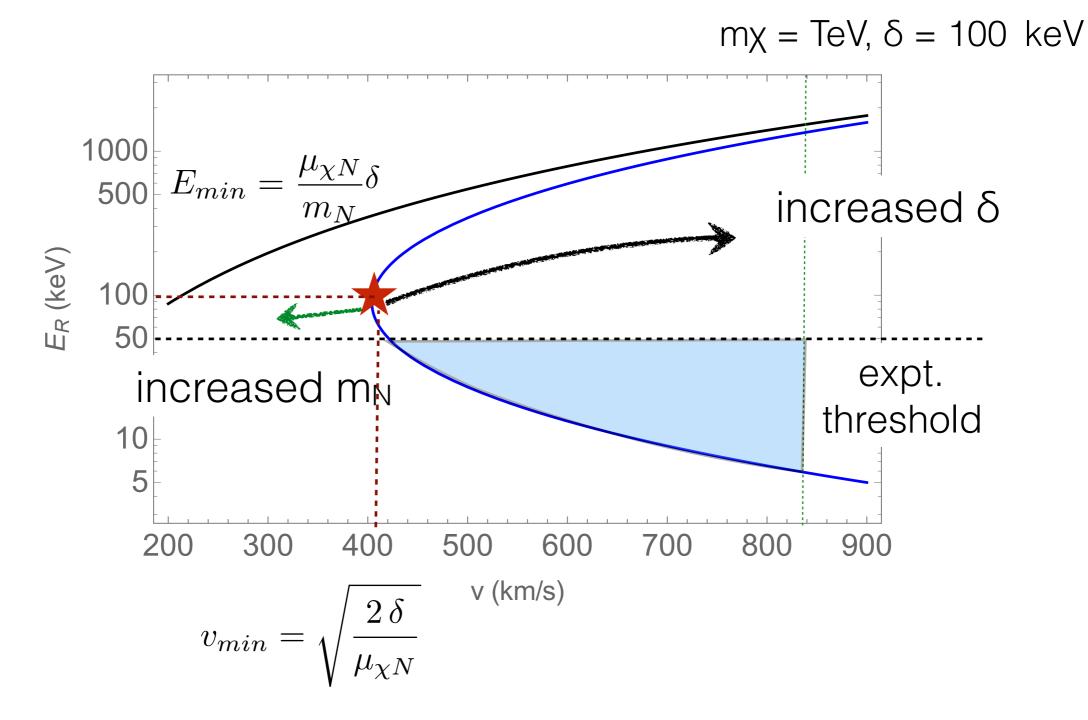
 E_R (keV)

 $m\chi = TeV, \delta = 100 keV$ $E_{min} = \frac{\mu_{\chi N}}{m_N} \delta$ expt. threshold v (km/s) $\left| {2\,\delta\over \mu_{\chi N}}
ight|$ $v_{min} = \sqrt{}$

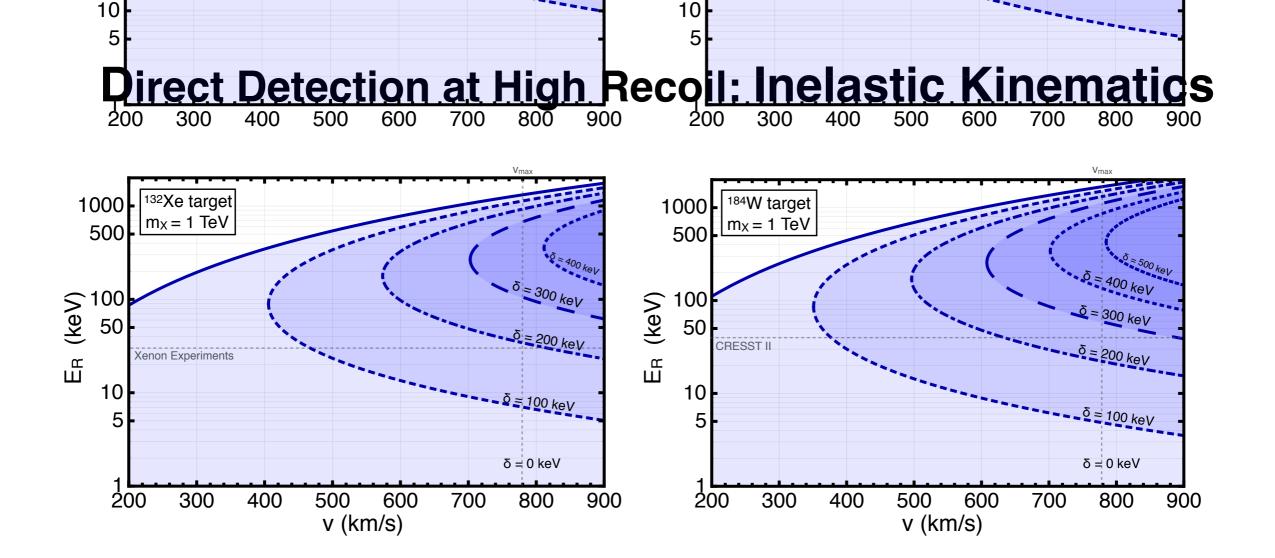
$$E_R = \frac{\mu_{\chi N}}{m_N} \Big((\mu_{\chi N} v^2 \cos^2 \theta_{lab} - \delta) \pm (\mu_{\chi N} v^2 \cos^2 \theta_{lab})^{1/2} (\mu_{\chi N} v^2 \cos^2 \theta_{lab} - 2\delta)^{1/2} \Big)$$

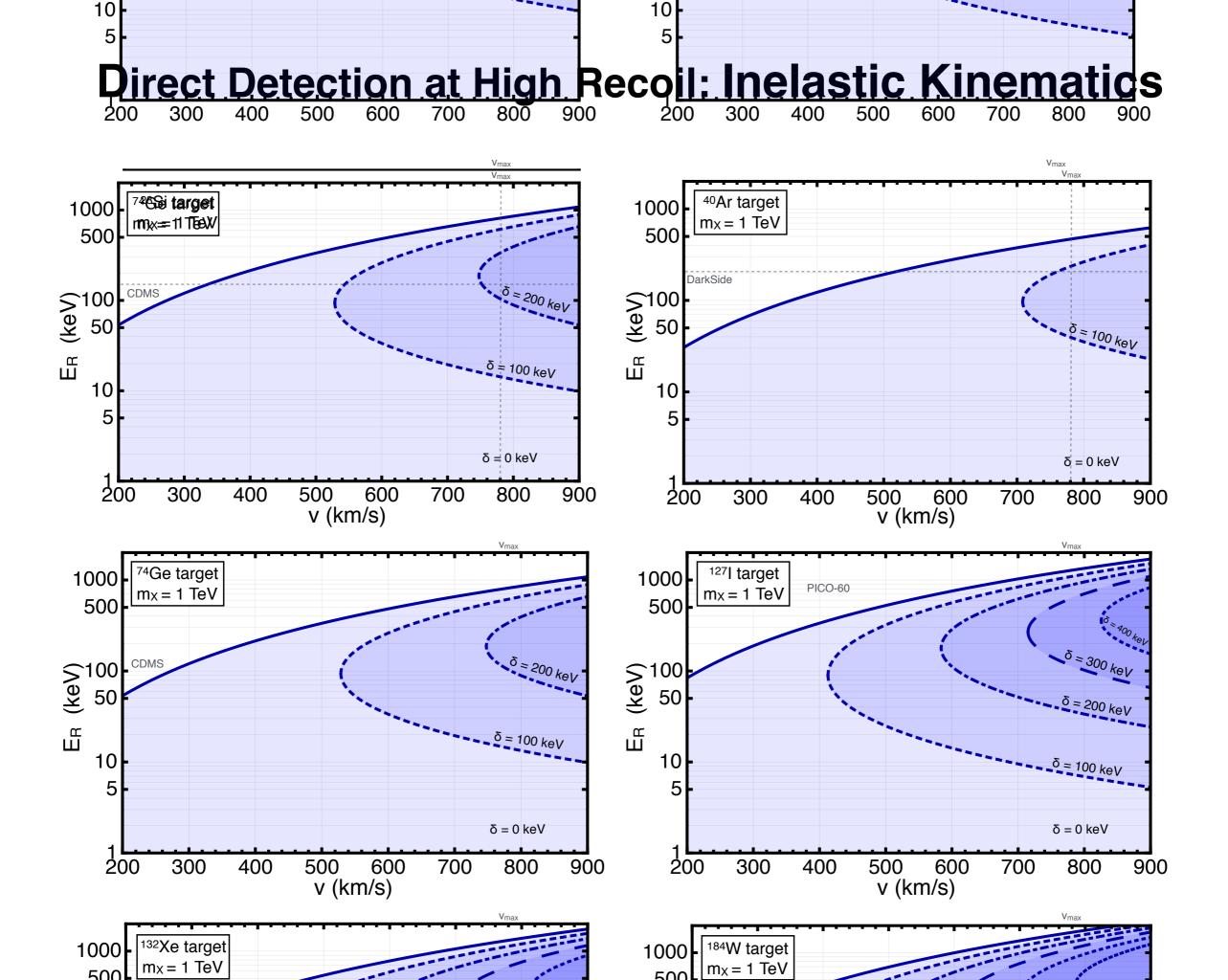


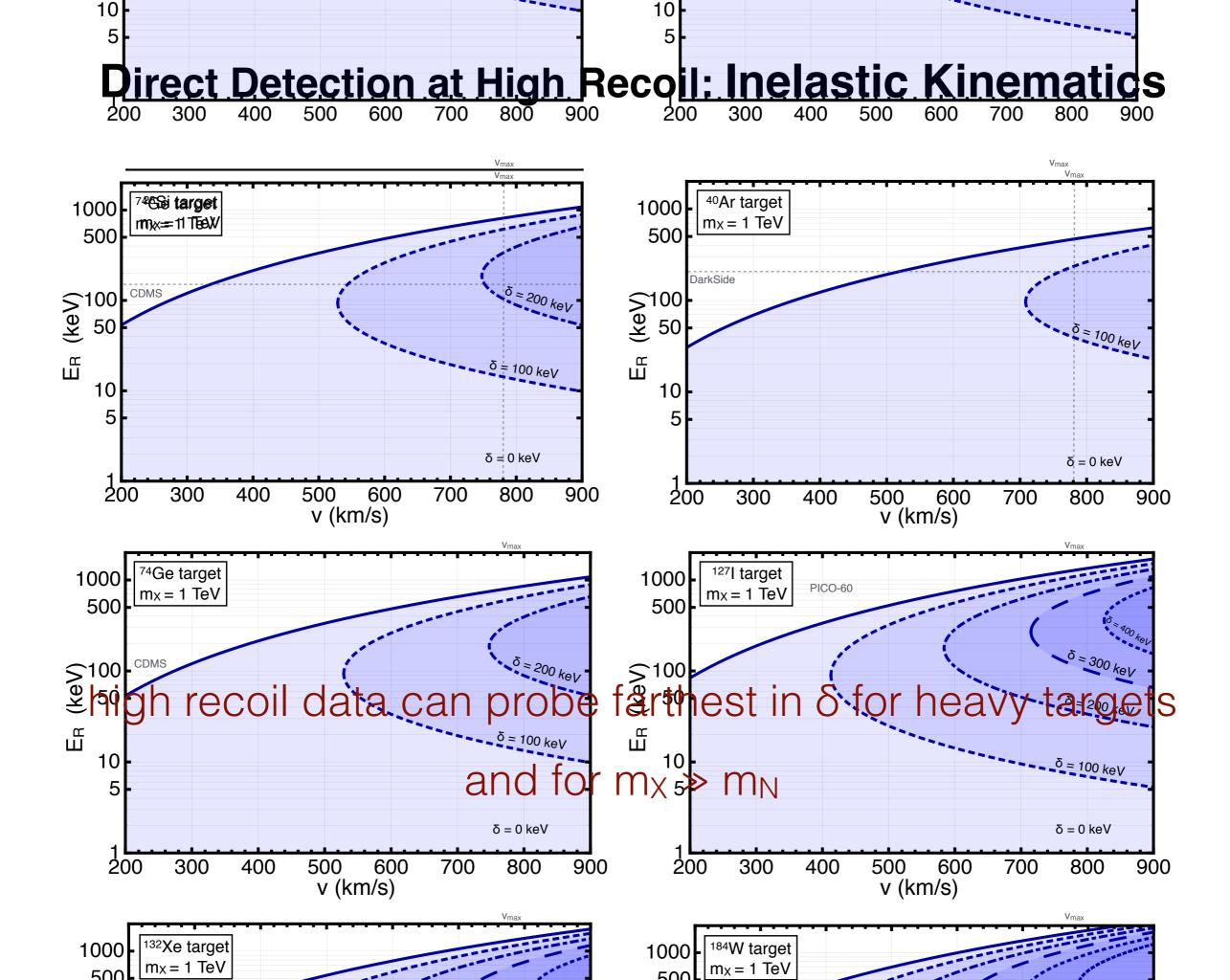
 $E_R = \frac{\mu_{\chi N}}{m_N} \Big((\mu_{\chi N} v^2 \cos^2 \theta_{lab} - \delta) \pm (\mu_{\chi N} v^2 \cos^2 \theta_{lab})^{1/2} (\mu_{\chi N} v^2 \cos^2 \theta_{lab} - 2\delta)^{1/2} \Big)$

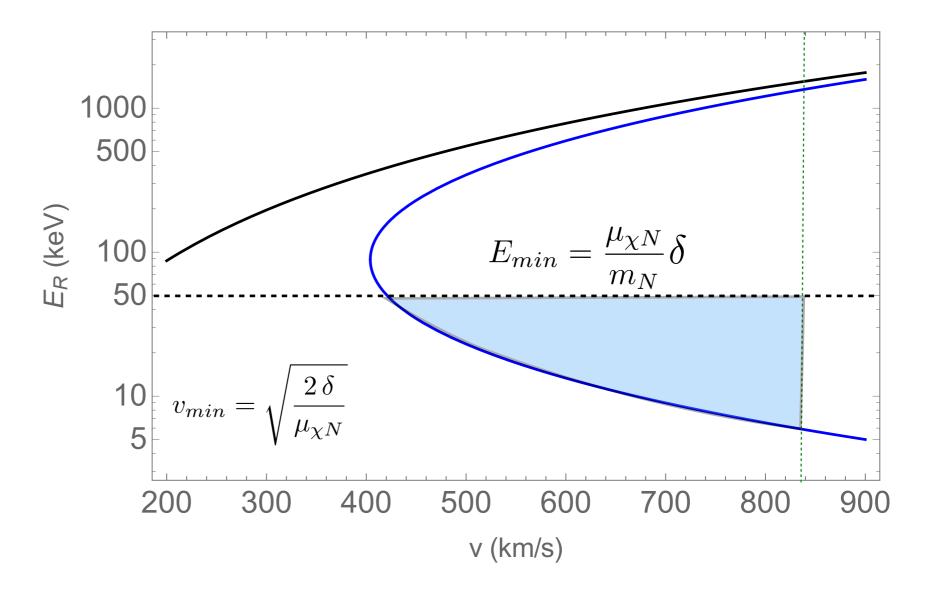


$$E_R = \frac{\mu_{\chi N}}{m_N} \left((\mu_{\chi N} v^2 \cos^2 \theta_{lab} - \delta) \pm (\mu_{\chi N} v^2 \cos^2 \theta_{lab})^{1/2} (\mu_{\chi N} v^2 \cos^2 \theta_{lab} - 2\delta)^{1/2} \right)$$

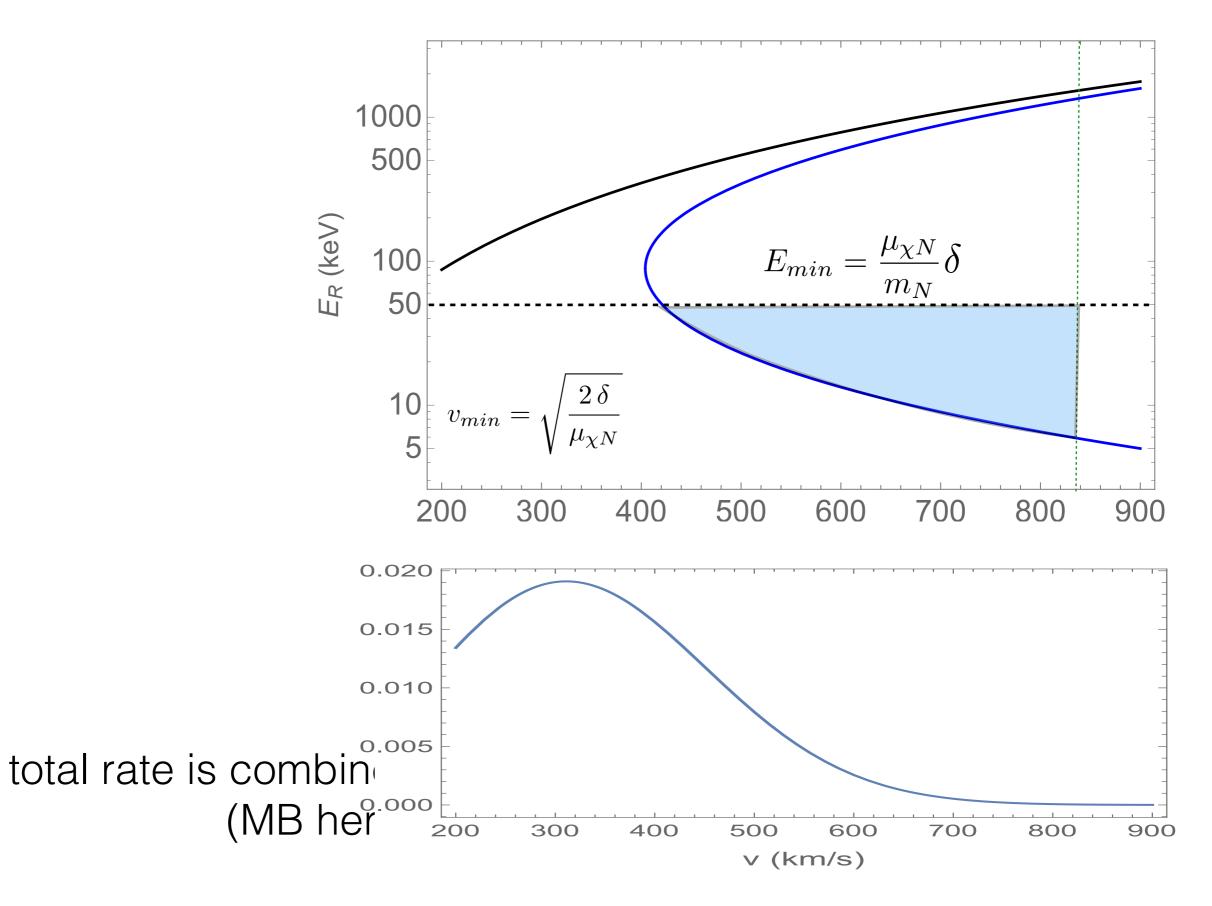




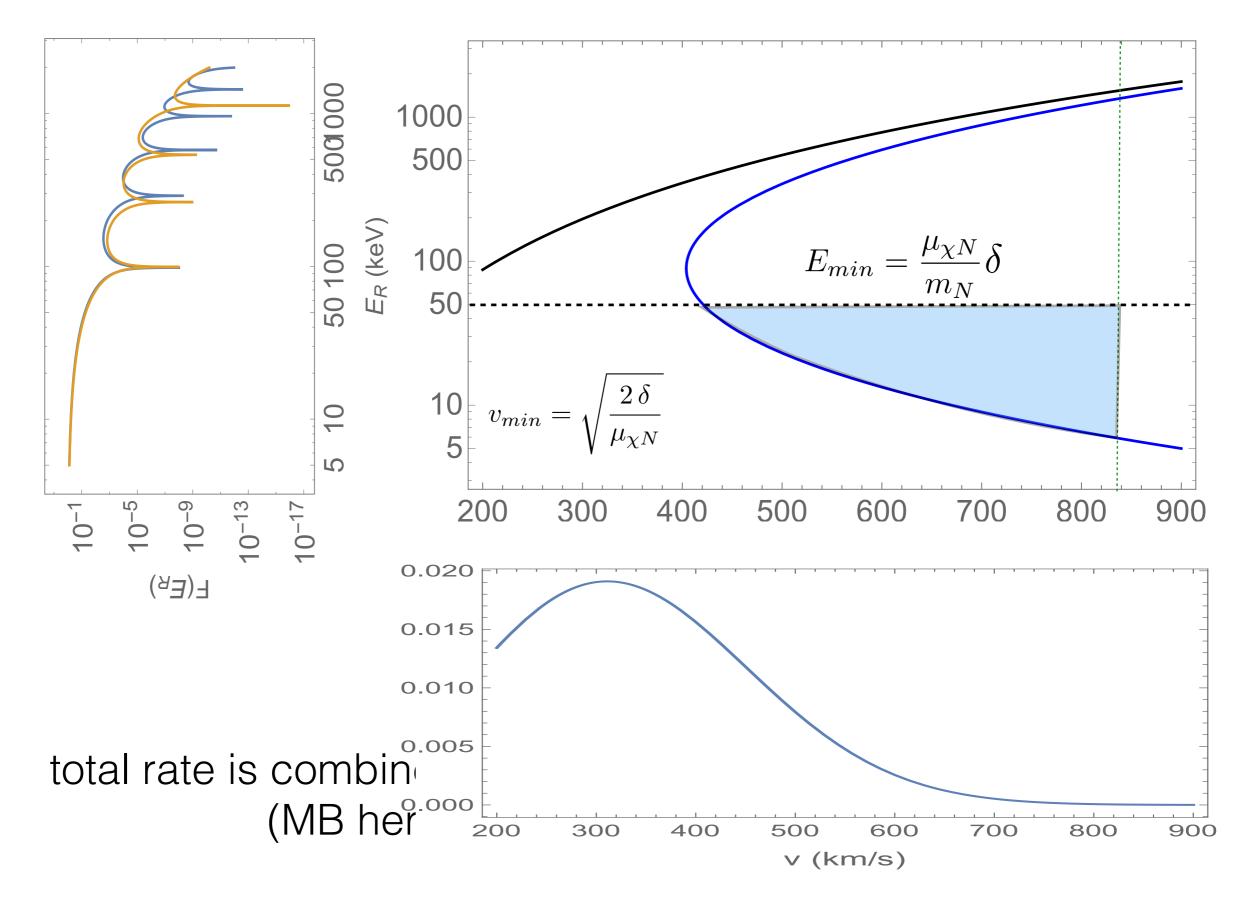




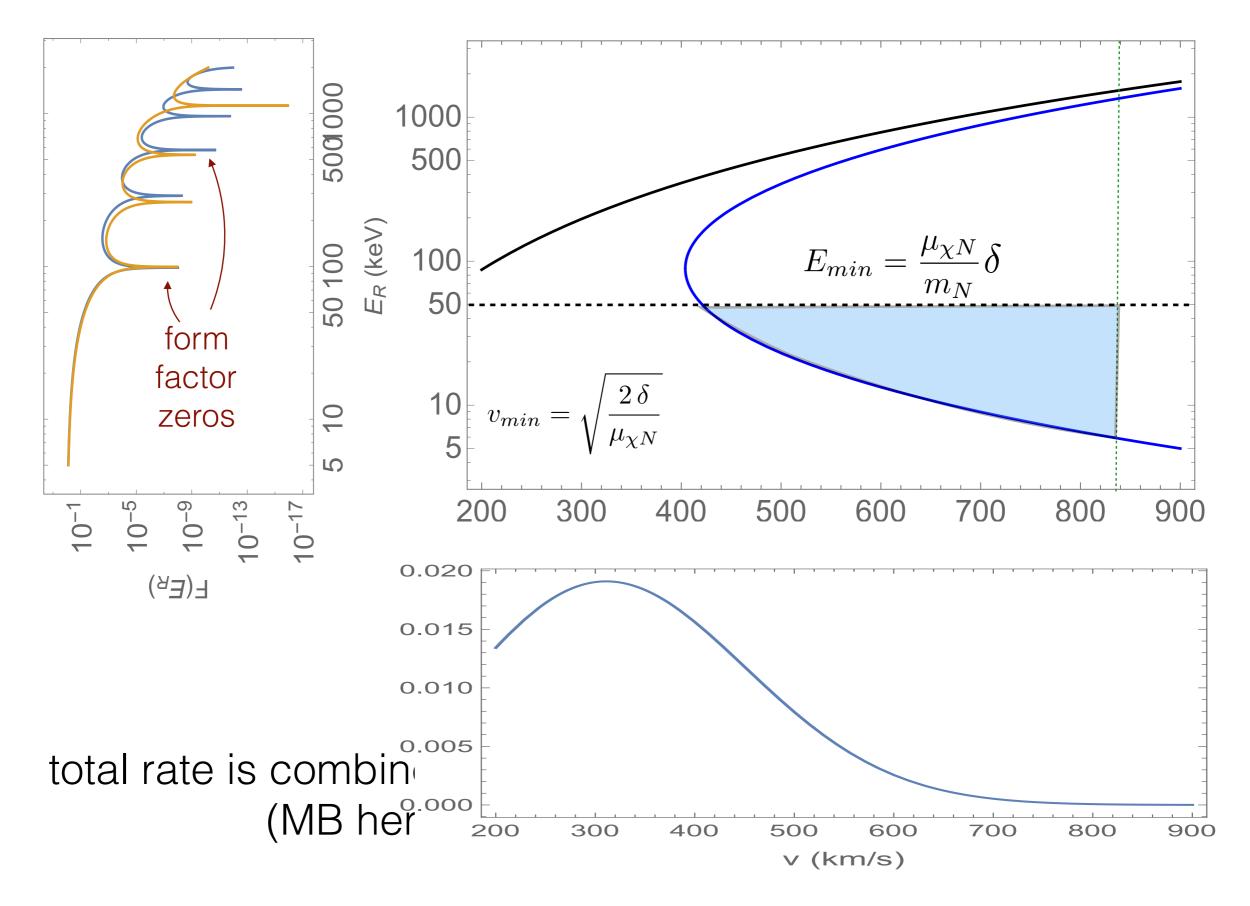
total rate is combines energy range, DM velocity spectrum (MB here), and nuclear form factors



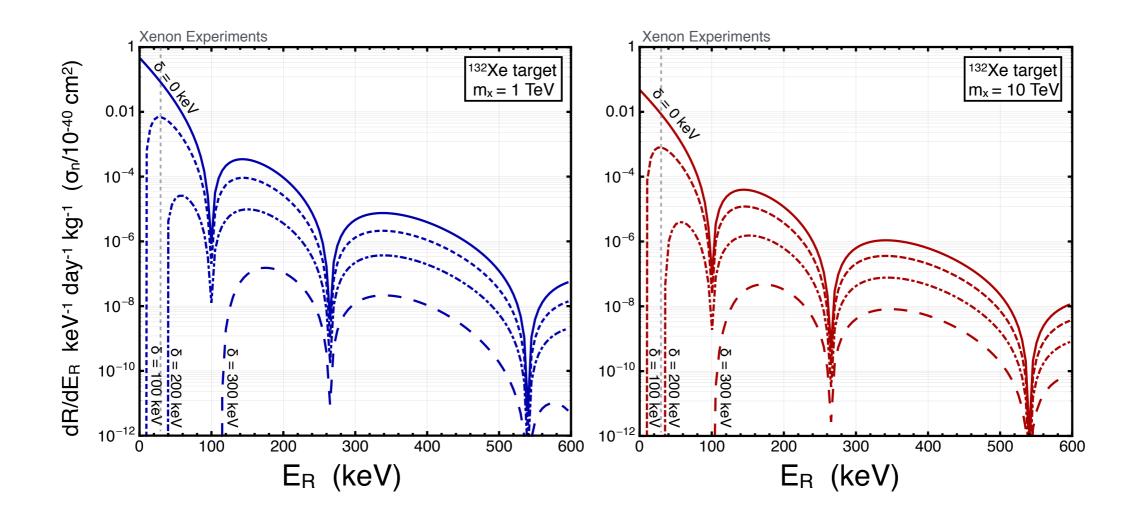
Direct Detection at High Recoil: Inelastic Kinematics



Direct Detection at High Recoil: Inelastic Kinematics

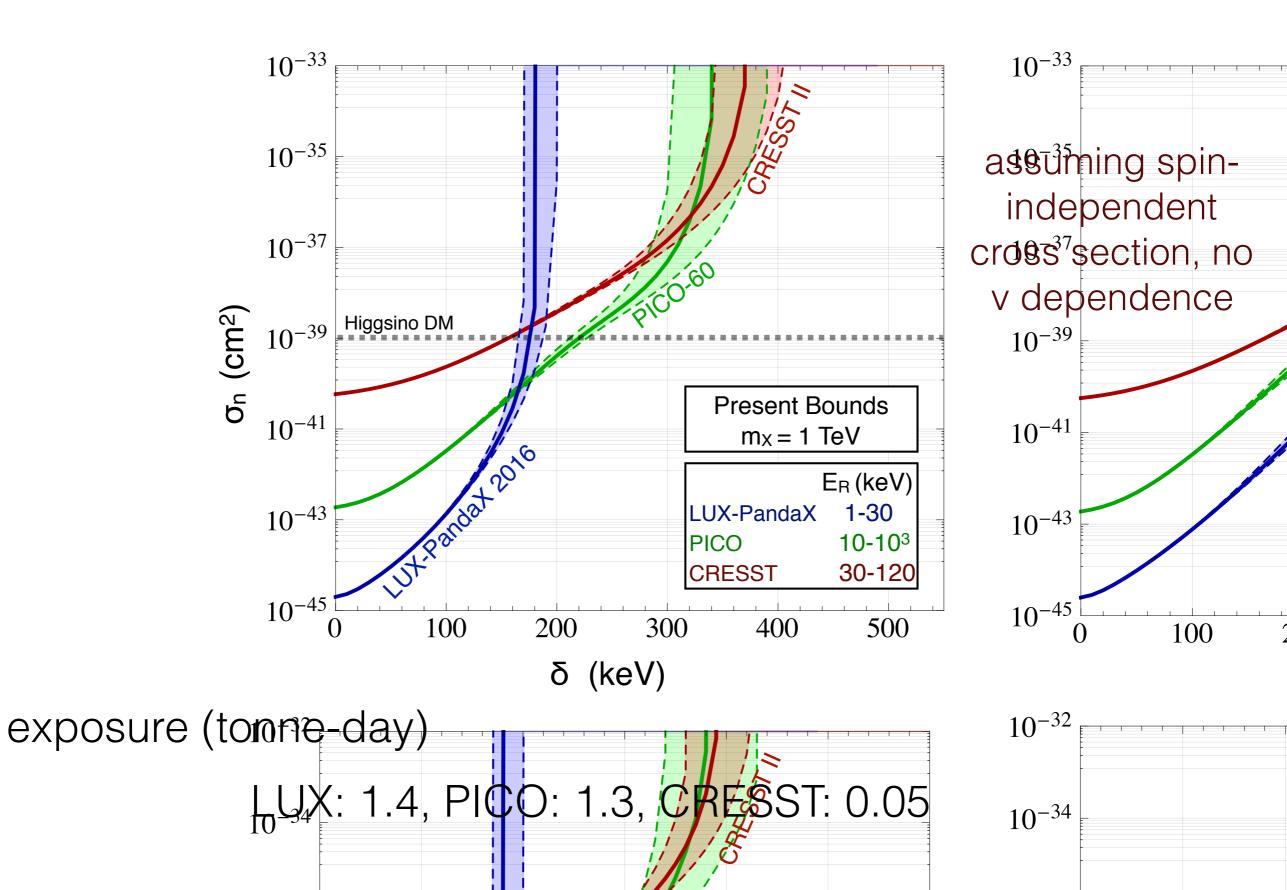


putting everything together

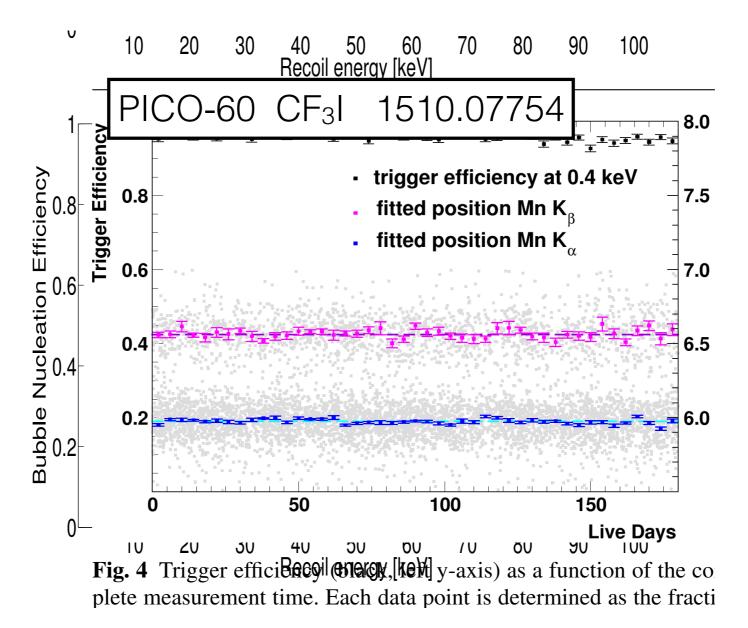


exact rate will be sensitive to the tail of the velocity distribution and large E_R part of form factors

bounds with current data, $m_X = 1$ TeV

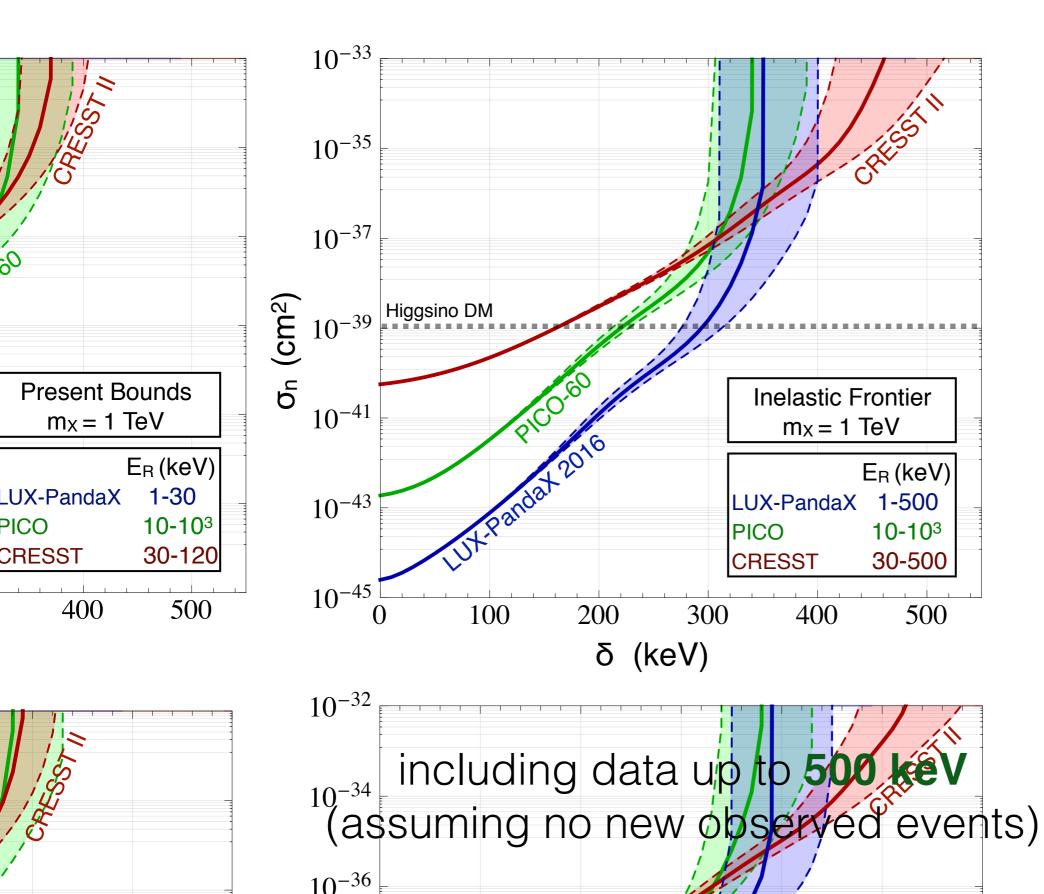


PICO?

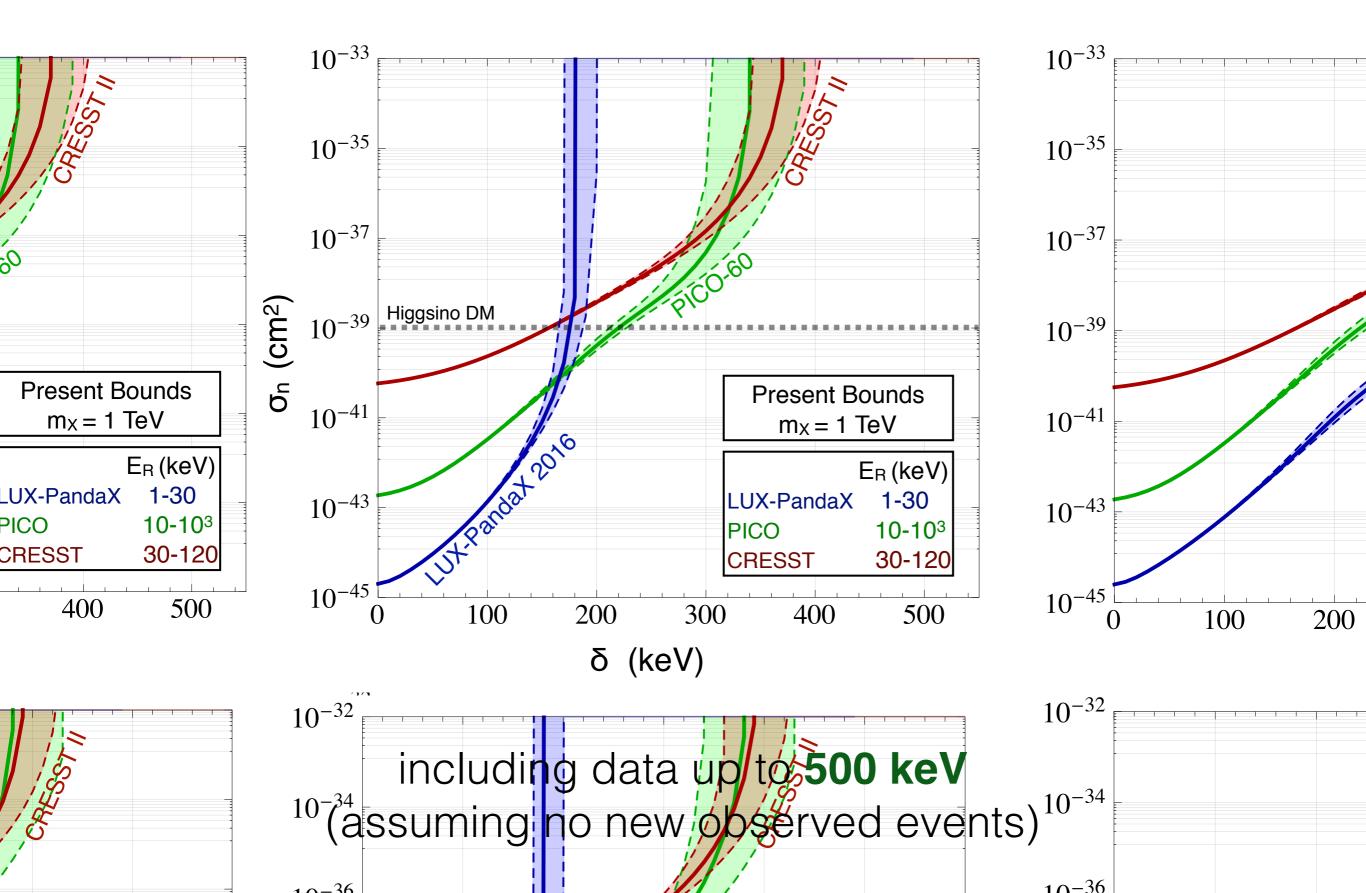


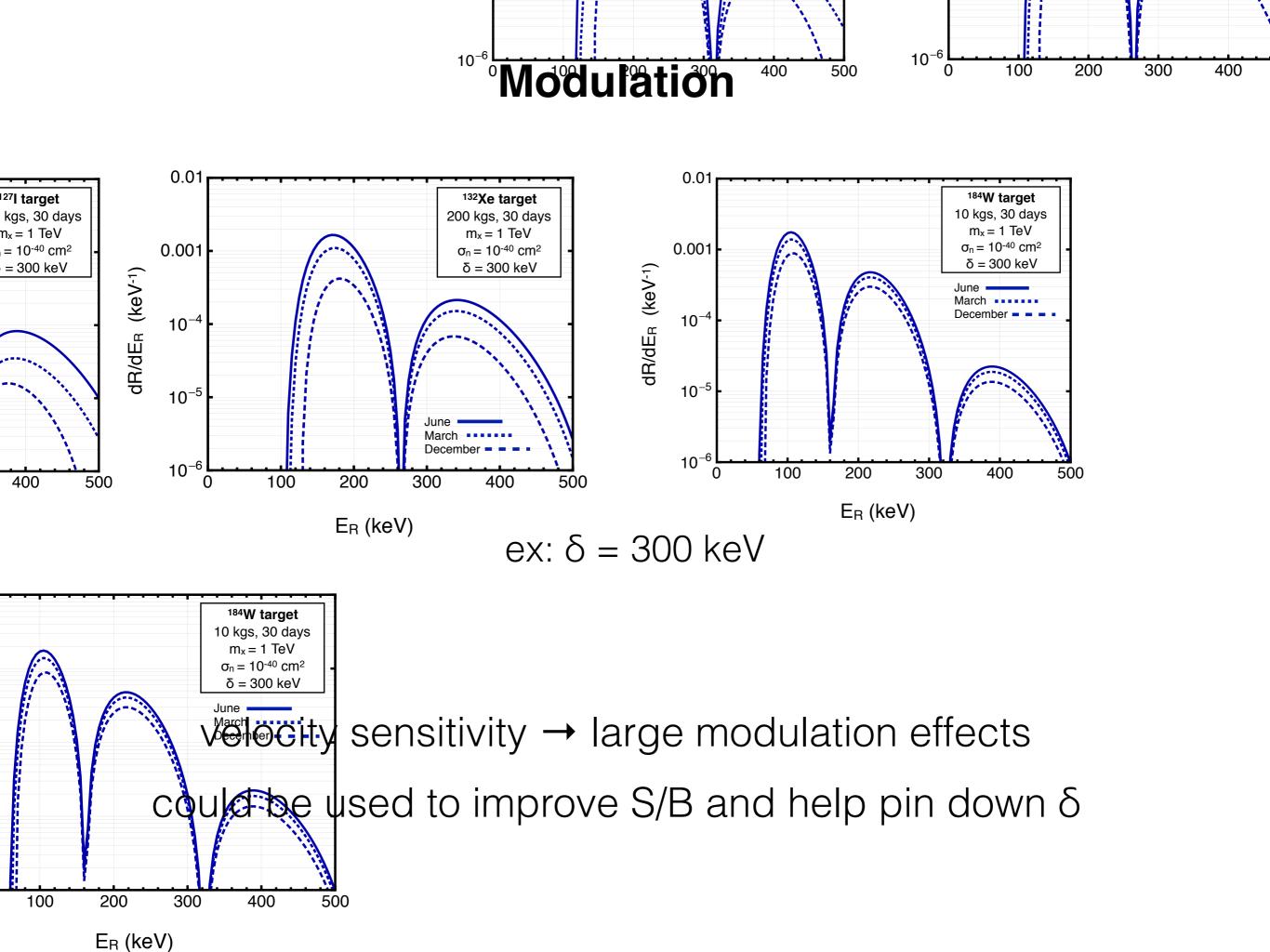
smaller fiducial volume than Xe experiments (& similar mass element).. but NO upper limit on search window

projected bounds including high recoil, current exposure: $m_X = 1$ TeV



projected bounds including high recoil, current exposure: $m_X = 1$ TeV





Dark photon + dark charged DM:

[Batell, Pospelov, Ritz '09]

$$\mathcal{L} = \mathcal{L}_{\rm SM} + |D_{\mu}\Phi|^{2} - V(\Phi) - \frac{1}{4}V_{\mu\nu}^{2} + \epsilon V_{\mu}\partial_{\nu}F^{\mu\nu} + \bar{\psi}(iD_{\mu}\gamma_{\mu} - m_{\psi})\psi + (\lambda_{\rm D}\Phi\psi^{T}C^{-1}\psi + \text{h.c.})$$

 $U(1)_D$ breaking splits Dirac DM (ψ) -> 2 Majorana (χ_1, χ_2)

Dark photon + dark charged DM:

[Batell, Pospelov, Ritz '09]

$$\mathcal{L} = \mathcal{L}_{SM} + |D_{\mu}\Phi|^{2} - V(\Phi) - \frac{1}{4}V_{\mu\nu}^{2} + \epsilon V_{\mu}\partial_{\nu}F^{\mu\nu} + \bar{\psi}(iD_{\mu}\gamma_{\mu} - m_{\psi})\psi + (\lambda_{D}\Phi\psi^{T}C^{-1}\psi + h.c.)$$

 $U(1)_D$ breaking splits Dirac DM (ψ) -> 2 Majorana (χ_1, χ_2)

annihilation $\chi_1 \chi_1 \rightarrow \gamma_D \gamma_D$ yields correct relic abundance for $\alpha_D \sim 0.04 \left(\frac{m_{\chi_1}}{\text{TeV}}\right)$

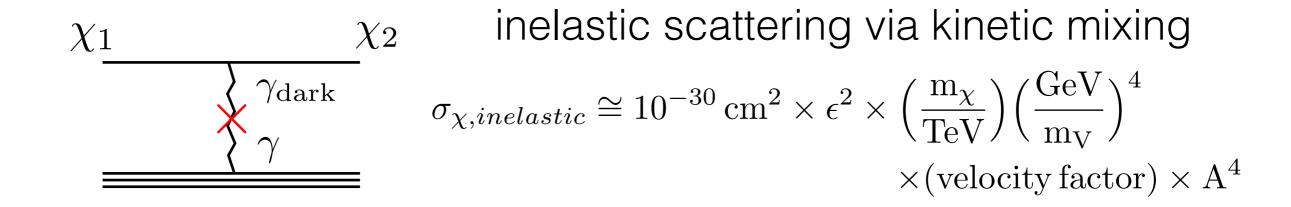
Dark photon + dark charged DM:

[Batell, Pospelov, Ritz '09]

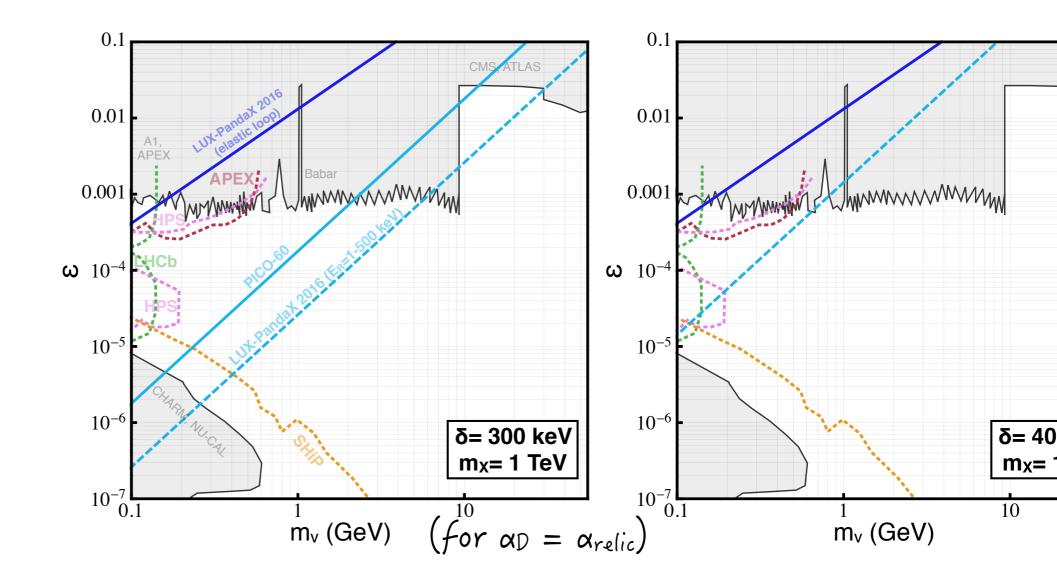
$$\mathcal{L} = \mathcal{L}_{SM} + |D_{\mu}\Phi|^{2} - V(\Phi) - \frac{1}{4}V_{\mu\nu}^{2} + \epsilon V_{\mu}\partial_{\nu}F^{\mu\nu} + \bar{\psi}(iD_{\mu}\gamma_{\mu} - m_{\psi})\psi + (\lambda_{D}\Phi\psi^{T}C^{-1}\psi + h.c.)$$

 $U(1)_D$ breaking splits Dirac DM (ψ) -> 2 Majorana (χ_1, χ_2)

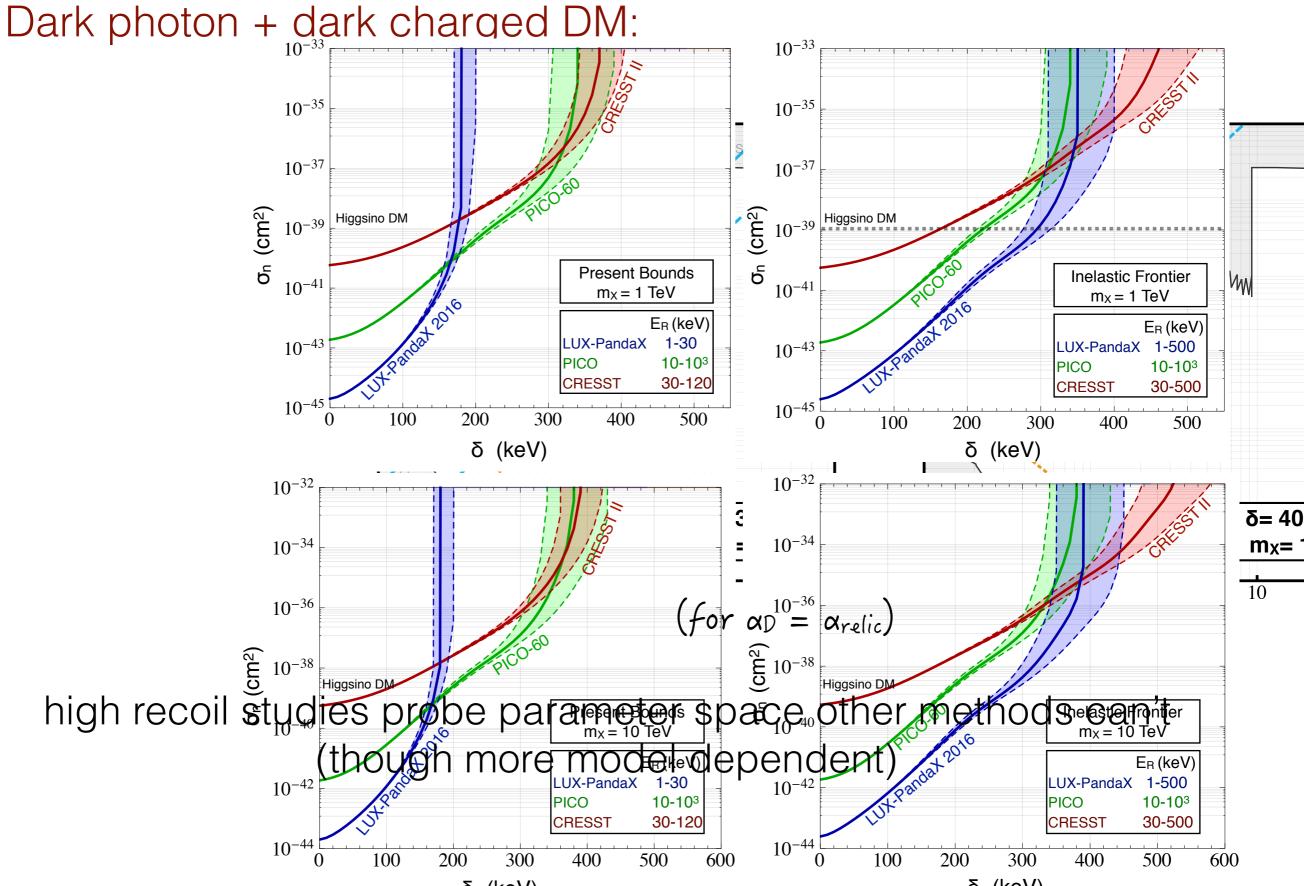
annihilation $\chi_1 \chi_1 \rightarrow \gamma_D \gamma_D$ yields correct relic abundance for $\alpha_D \sim 0.04 \left(\frac{m_{\chi_1}}{\text{TeV}}\right)$

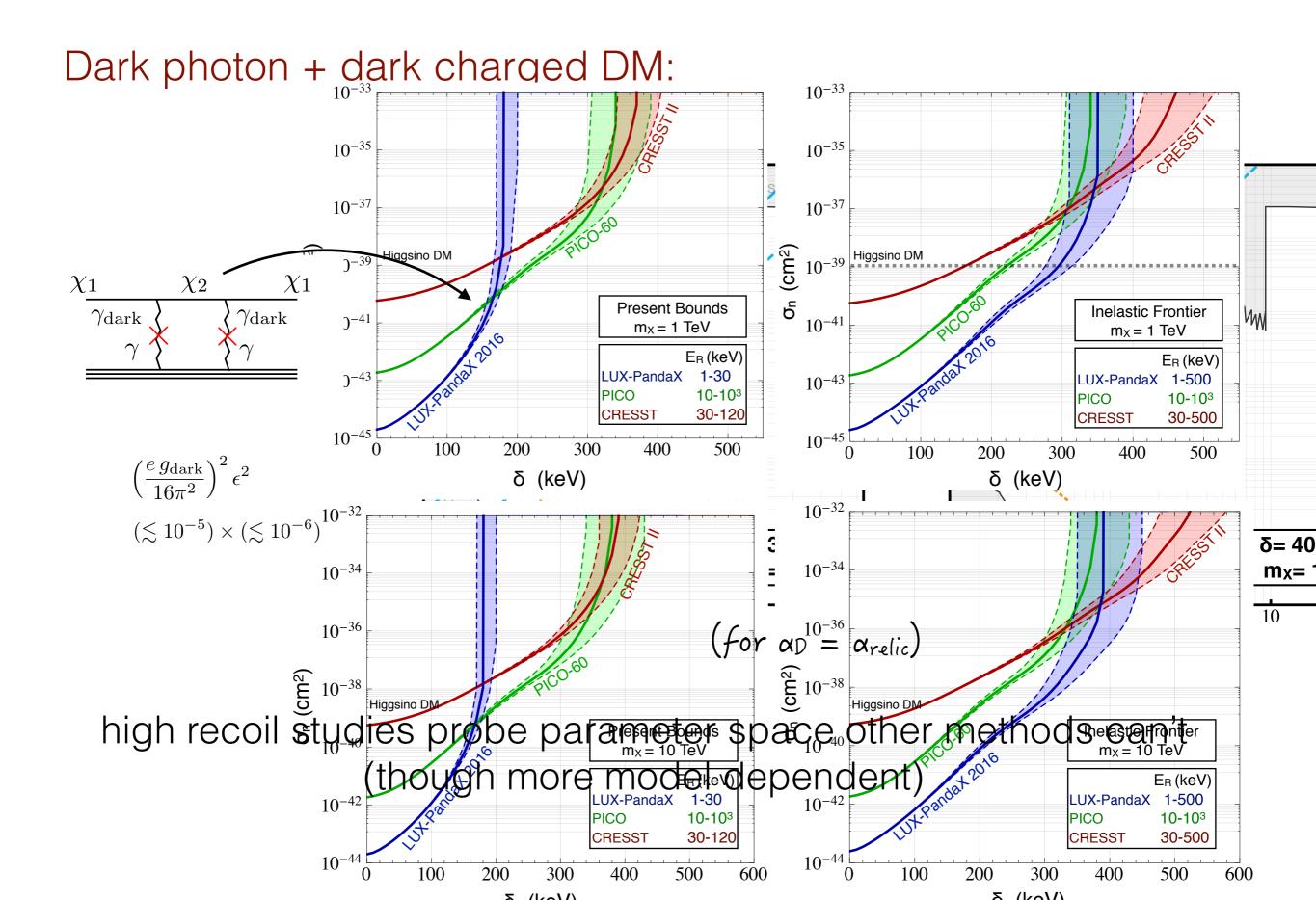


Dark photon + dark charged DM:



high recoil studies probe parameter space other methods can't (though more model dependent)





stic Dark Matter What other models could lie in high recoil data?

which again has two Majorana fermions χ_1, χ_2 nearby in mass interaction with the SM is through a magnetic dipole operator

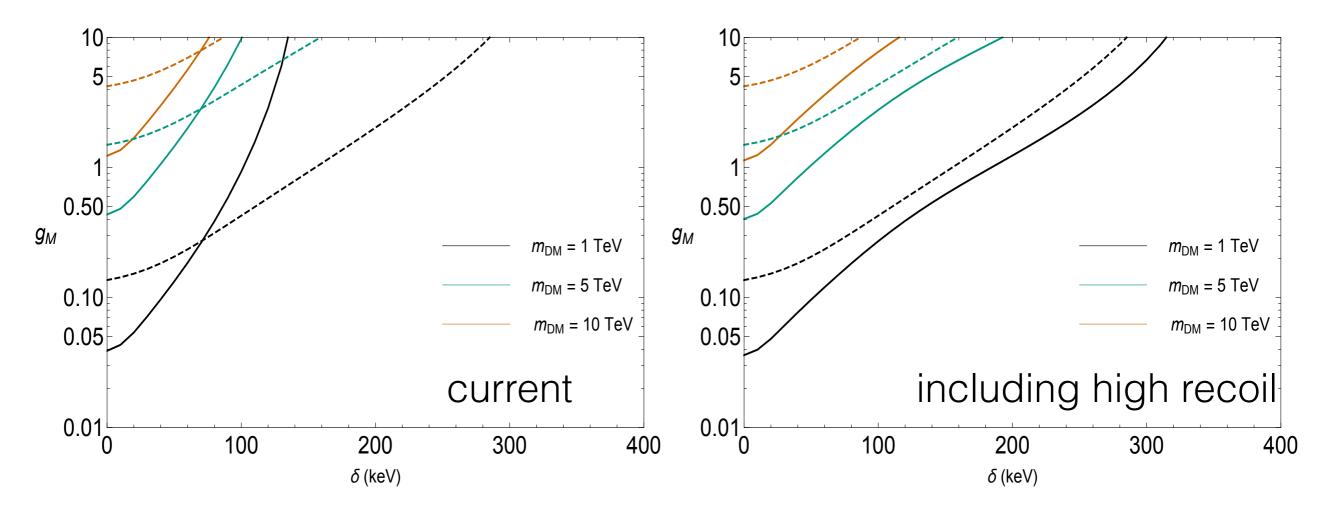
$$\left(\frac{g_M}{4}\right) \mathcal{L} \underbrace{\mathcal{S}}_{2m_{\chi}} \left(\underbrace{\frac{g_M e}{\chi_2}}_{2m_{\chi}}\right) \chi_2 \mathcal{F}_{2m_{\chi}} \mathcal{V}_1 F^{\mu\nu}$$
[Chang, Weiner, Yavin '10]

na in nature diversion policies and dipole exercises, and only turbative UV completion [46] of the theory generates this integrating out $\frac{1}{4}$ heavy charge fermion and scalar of mass have chosen to adopt the operator normalization inspired by p uld be expected if the DM was a composite of a new strongly ct $g_M \sim 1$.

Signal of MIDM is the inelastic collision of $(\Pr_{Te})^2$ with the SM through the the set of the charge and magnetic moment of the proton and TReliseapundance gy complete standing standards for M and M

as interaction is energy dependent, no analog of σ_n

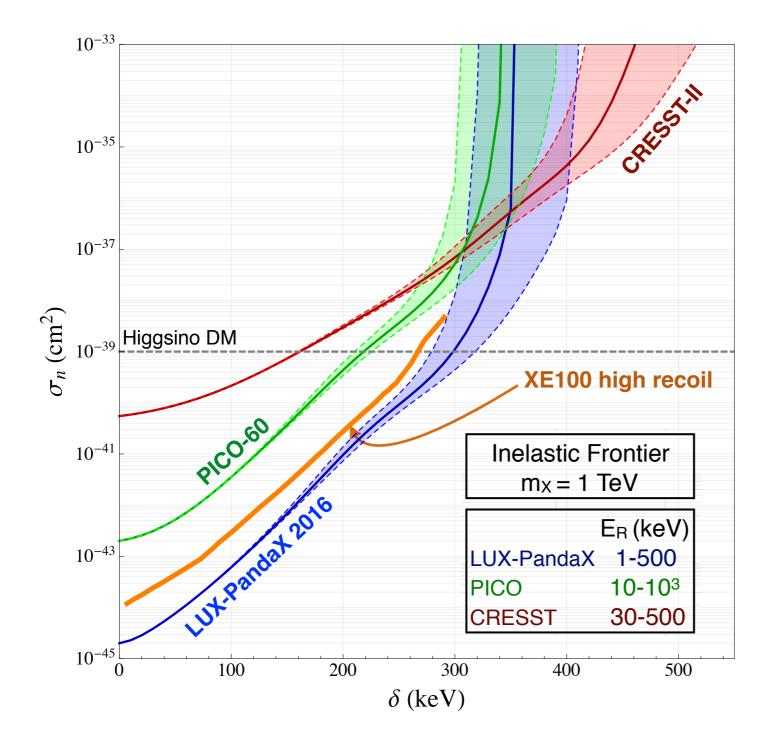
∴ bound g_M directly using nuclear response formalism of Fitzpatrick et al 1203.3542



dashed = PICO, **solid** = LUX/PandaX

strong bound from PICO due to large ¹²⁷I spin

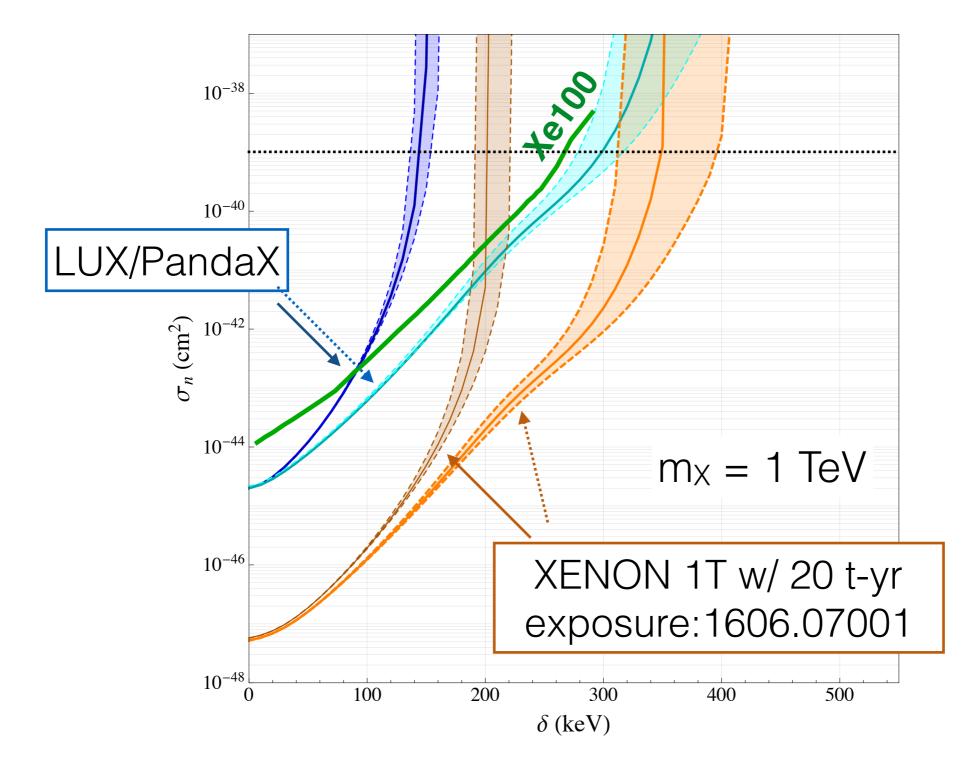
Recent progress at high recoil! : XE100 1705.02614



considered recoil energies out to 240 keVnr

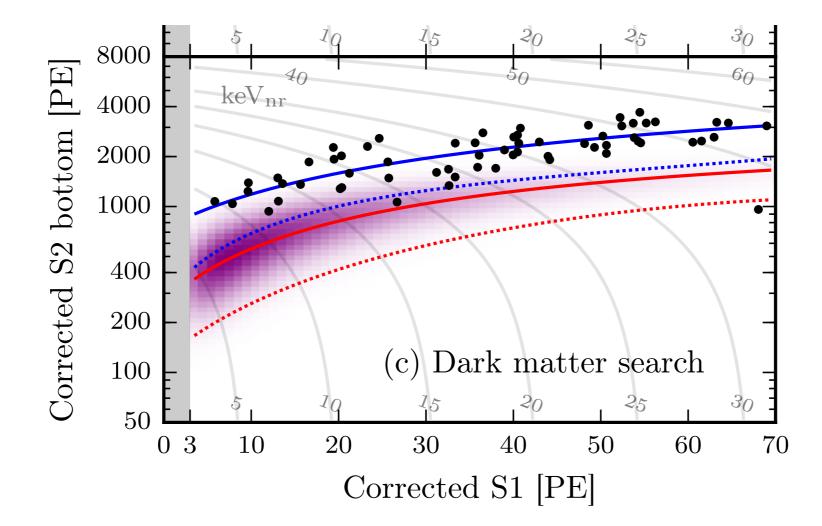
limit on 'weak strength' DM pushed out to $\delta \ge 260 \text{ keV}$

Projections for future experiments



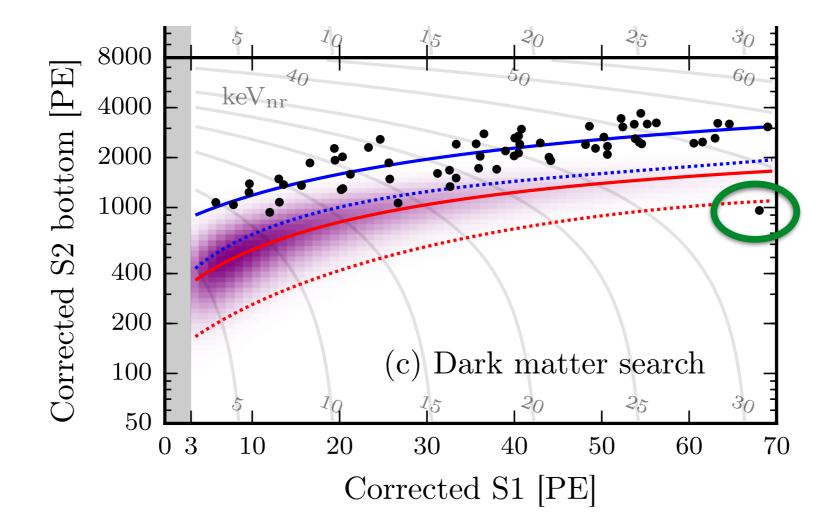
Perhaps the first signs of iDM?

from XENON1T 1705.06655



Perhaps the first signs of iDM?

from XENON1T 1705.06655



Conclusions

Explore the inelastic direction! Motivated models with sizable $\sigma_{X^{-N}}$ within **easy** reach

 current techniques work, just enlarge E_R signal regions; Xe100, LUX already looking into it

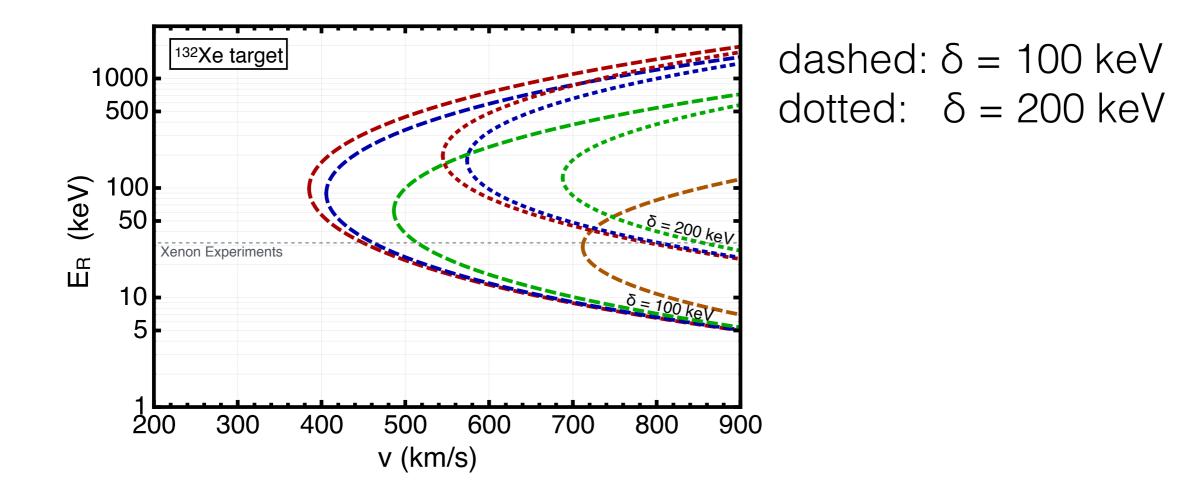
- most sensitive to heavy DM using heavy targets;
- sensitive to tails of DM velocity distribution, large modulation effects

what about $\delta \ge MeV$?



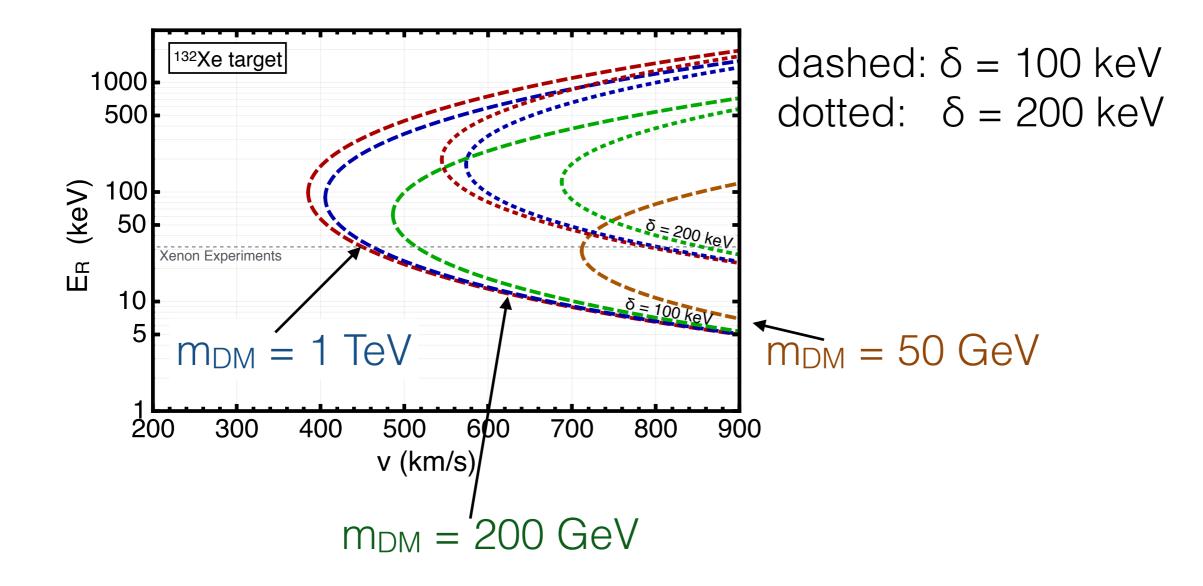
Inelastic Kinematics

shifts in kinematically allowed region as we change mDM



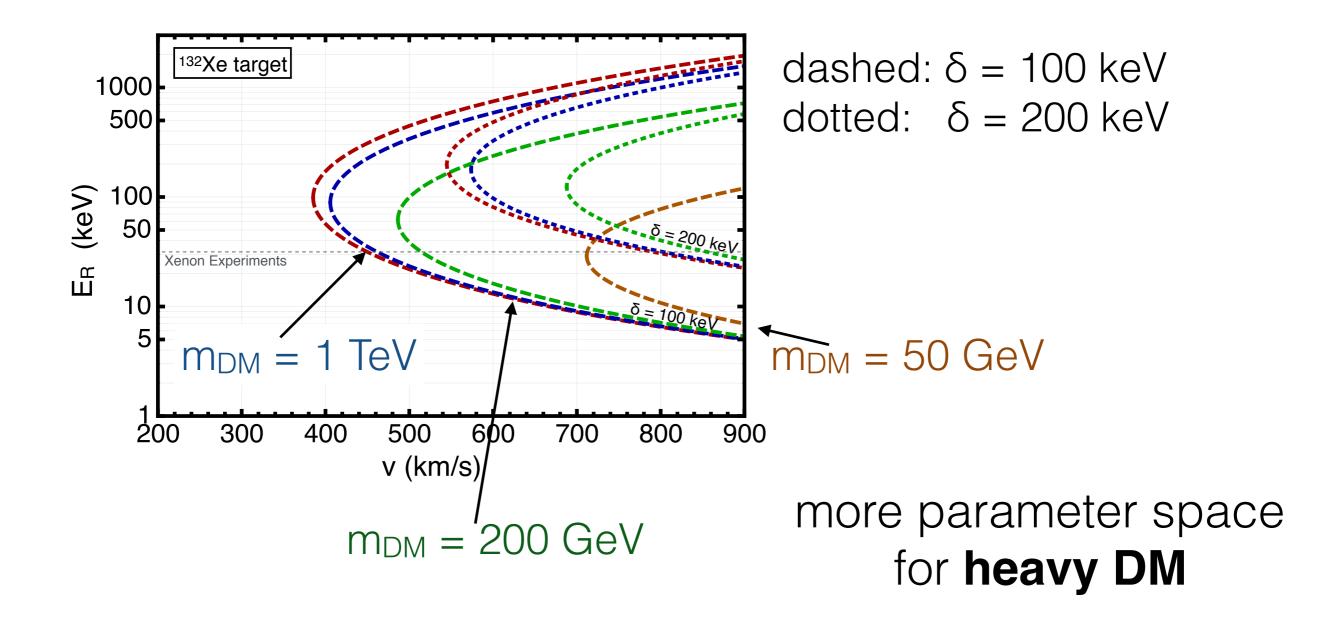
Inelastic Kinematics

shifts in kinematically allowed region as we change mDM

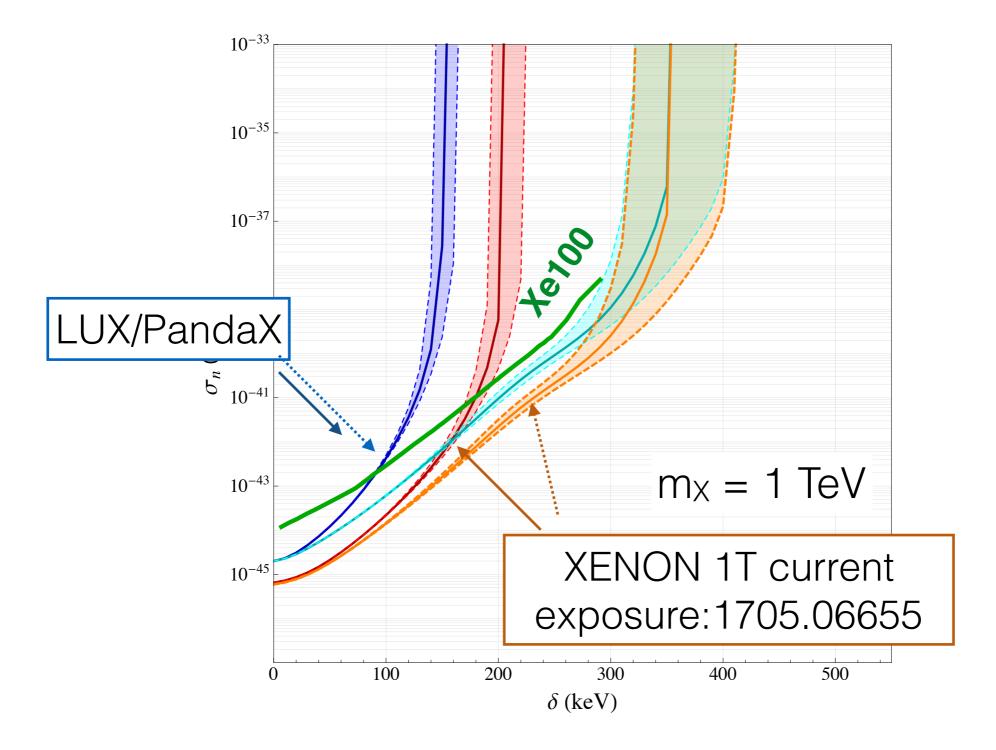


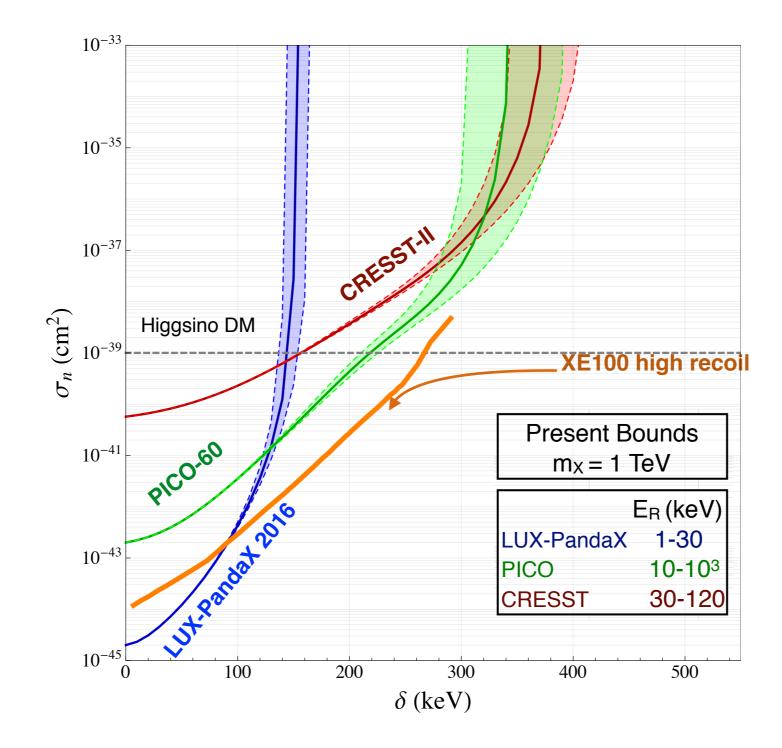
Inelastic Kinematics

shifts in kinematically allowed region as we change mDM

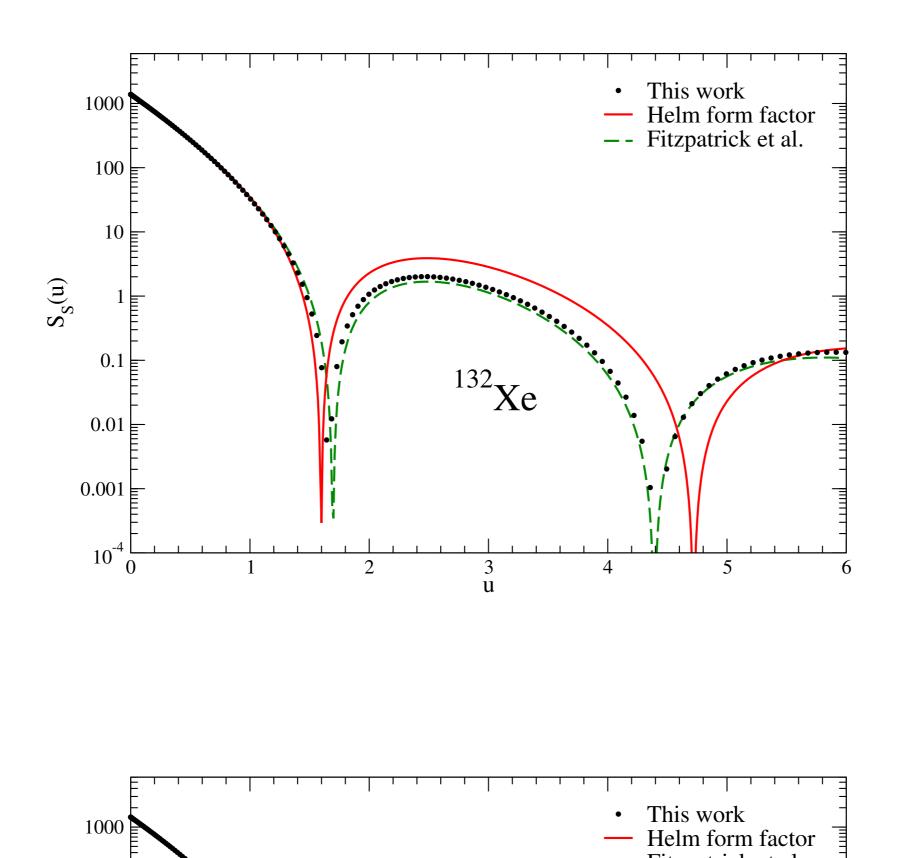


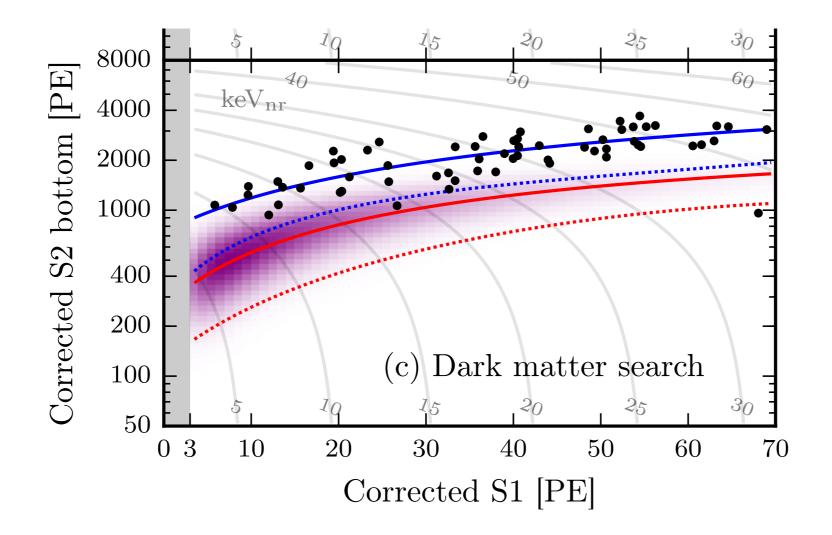
current XENON 1T comparison



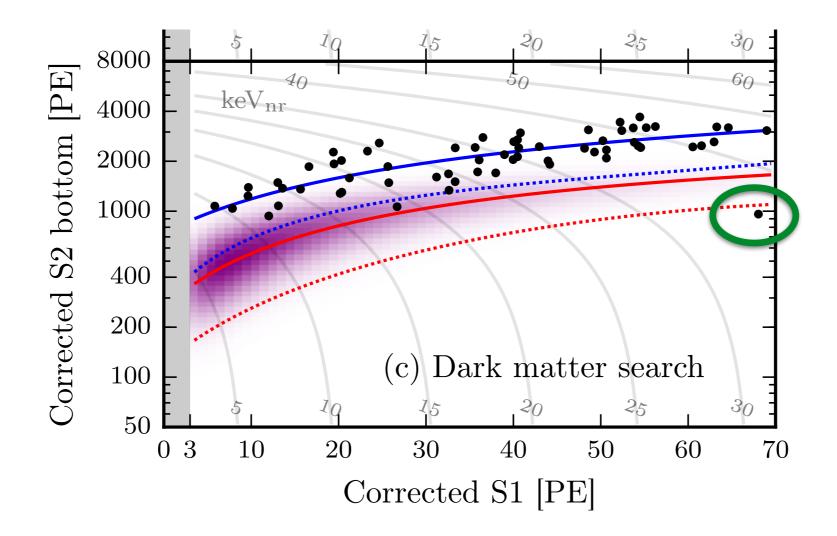


nuclear form factors





the lowest ER background ever achieved in a dark matter experiment. A single event far from the bulk distribution was observed at cS1 = 68.0 PE in the initial 4-day unblinding stage. This appears to be a *bona fide* event, though its location in $(cS1, cS2_b)$ (see Fig. 2c) is extreme for all our physical background models and WIMP signal models. One event at cS1 = 26.7 PE is at the -2.4σ ER



the lowest ER background ever achieved in a dark matter experiment. A single event far from the bulk distribution was observed at cS1 = 68.0 PE in the initial 4-day unblinding stage. This appears to be a *bona fide* event, though its location in $(cS1, cS2_b)$ (see Fig. 2c) is extreme for all our physical background models and WIMP signal models. One event at cS1 = 26.7 PE is at the -2.4σ ER