

What did we learn from the “QGP Brick” Problem?

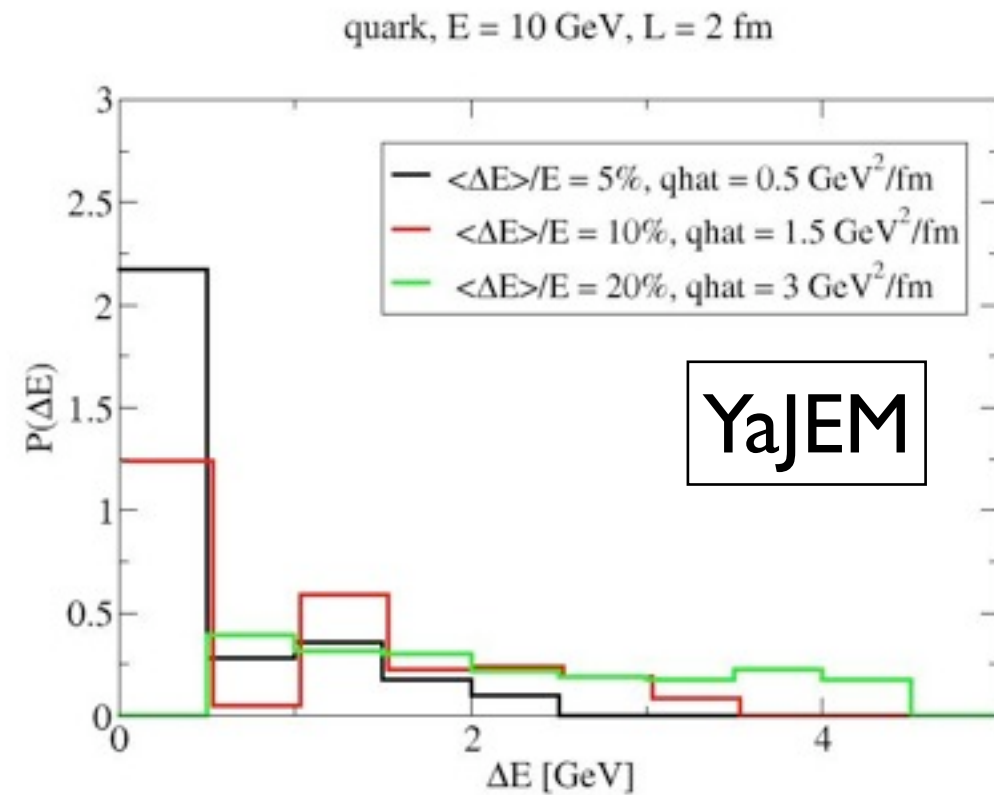
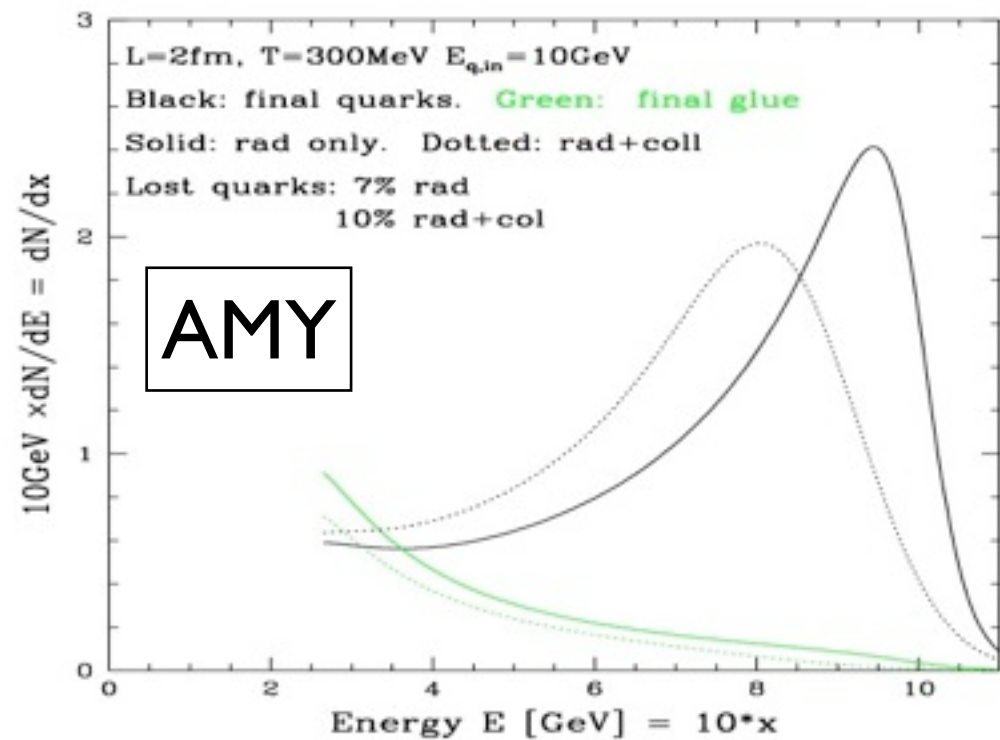
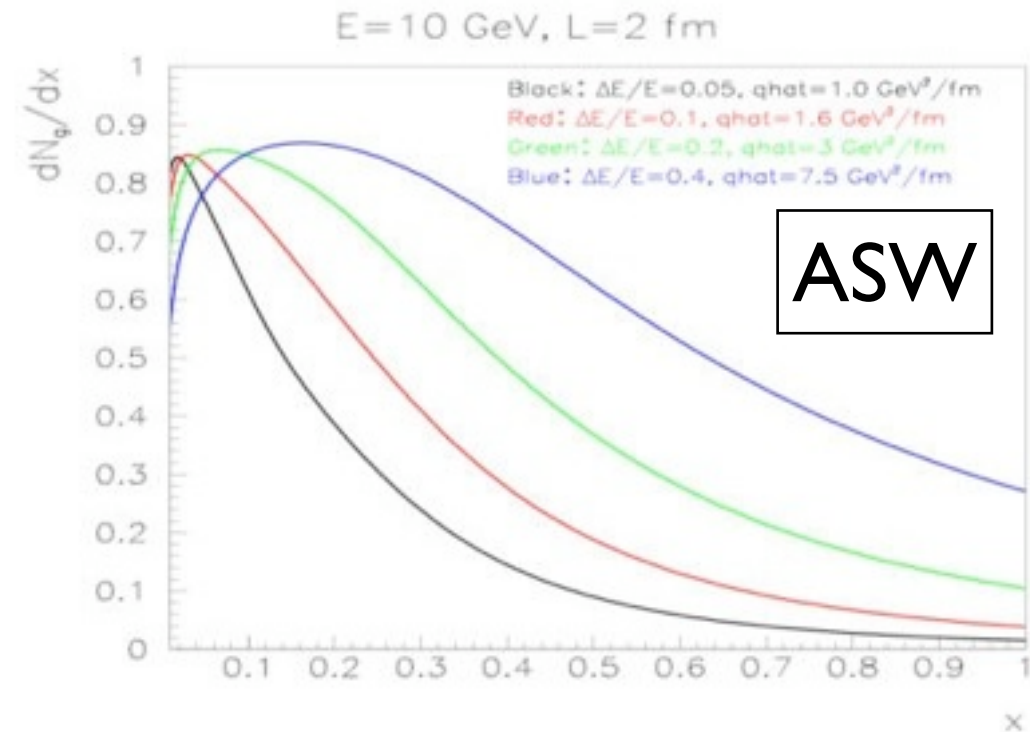
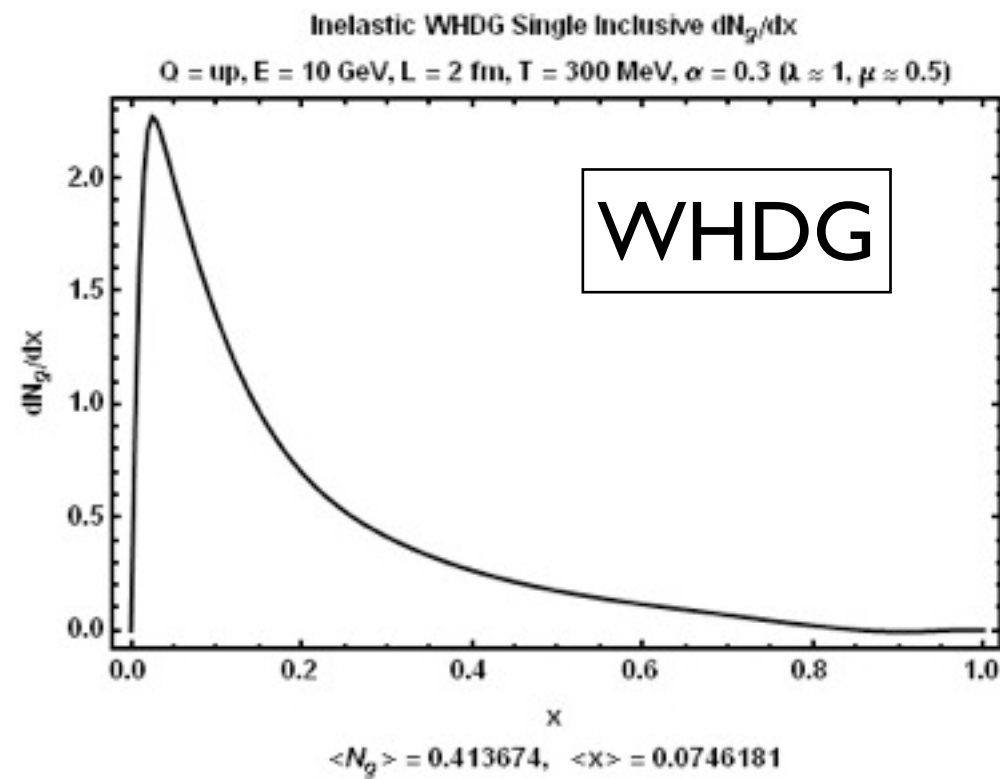
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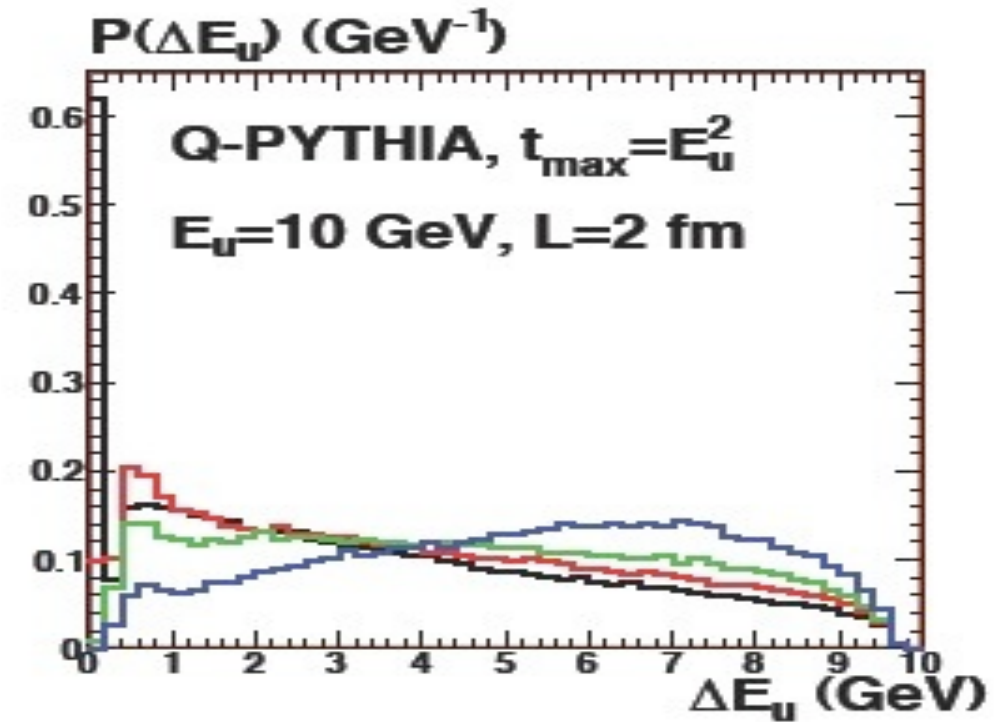
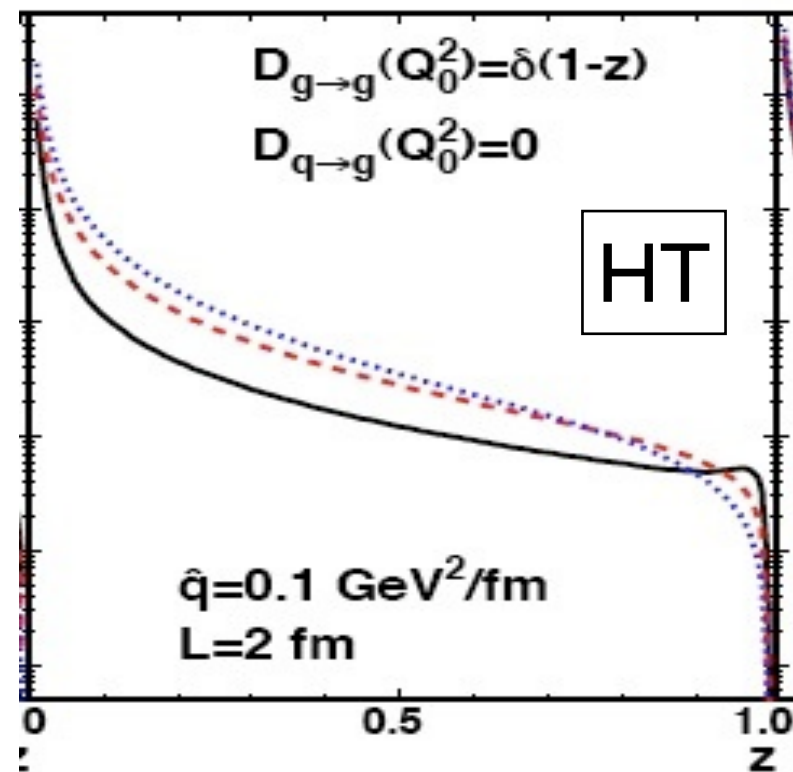
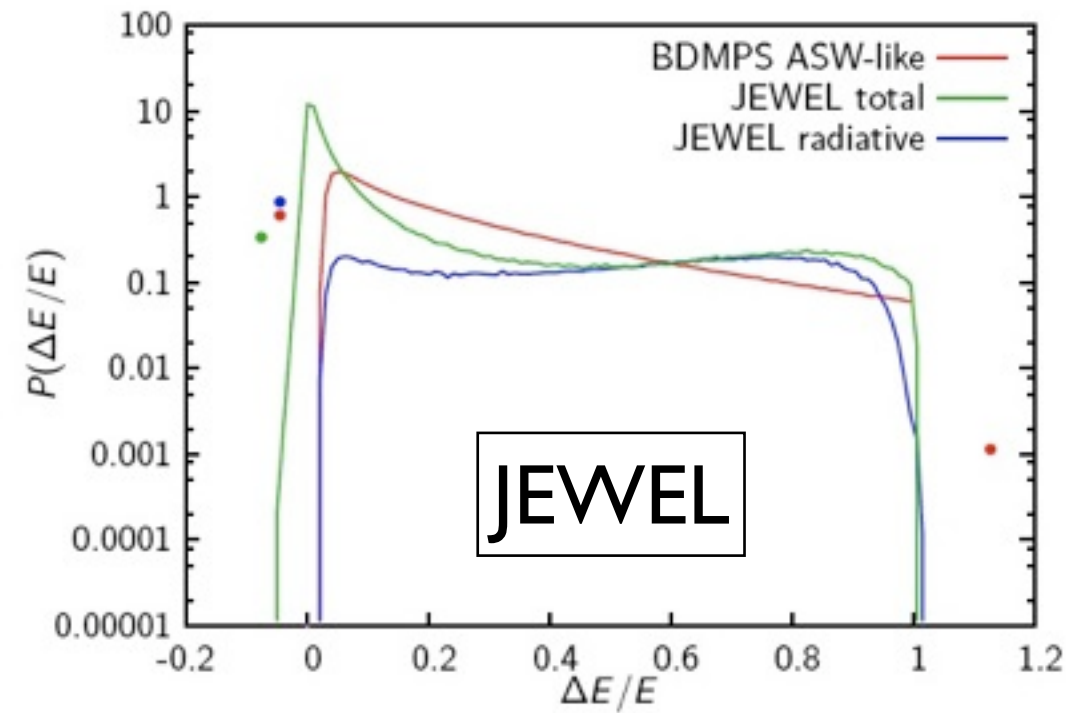
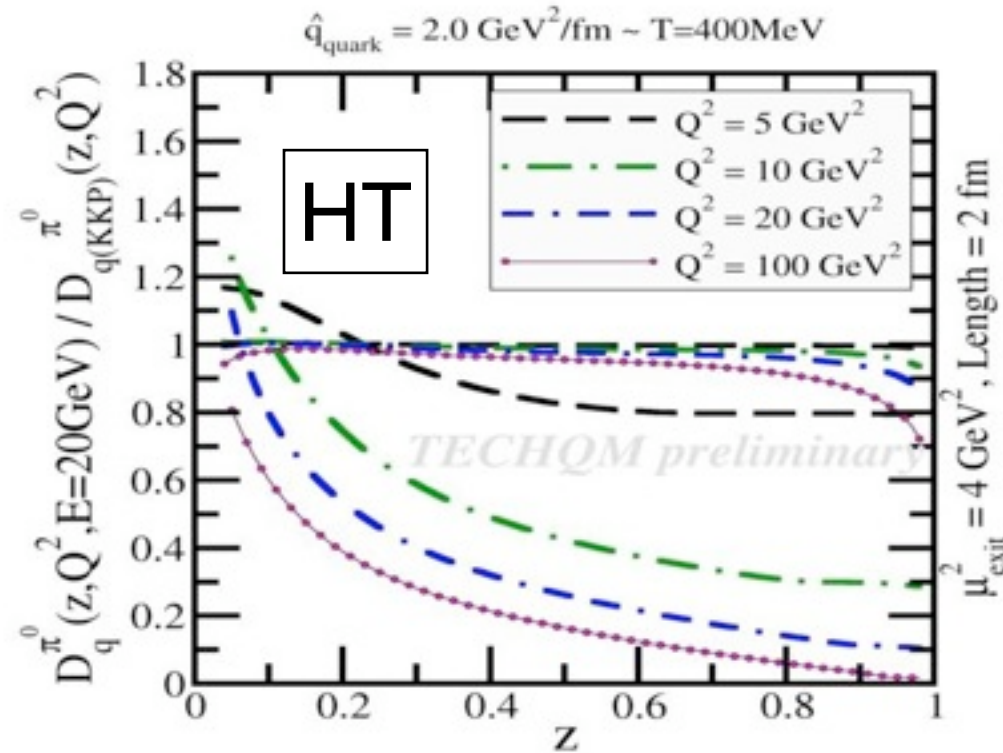
What did we learn from the “QGP Brick” Problem ?

- The community can collaborate
- More popular than anticipated
- Specifying a problem completely is not easy
- The “QGP Brick” challenge was useful
- The “QGP Brick” will remain a benchmark

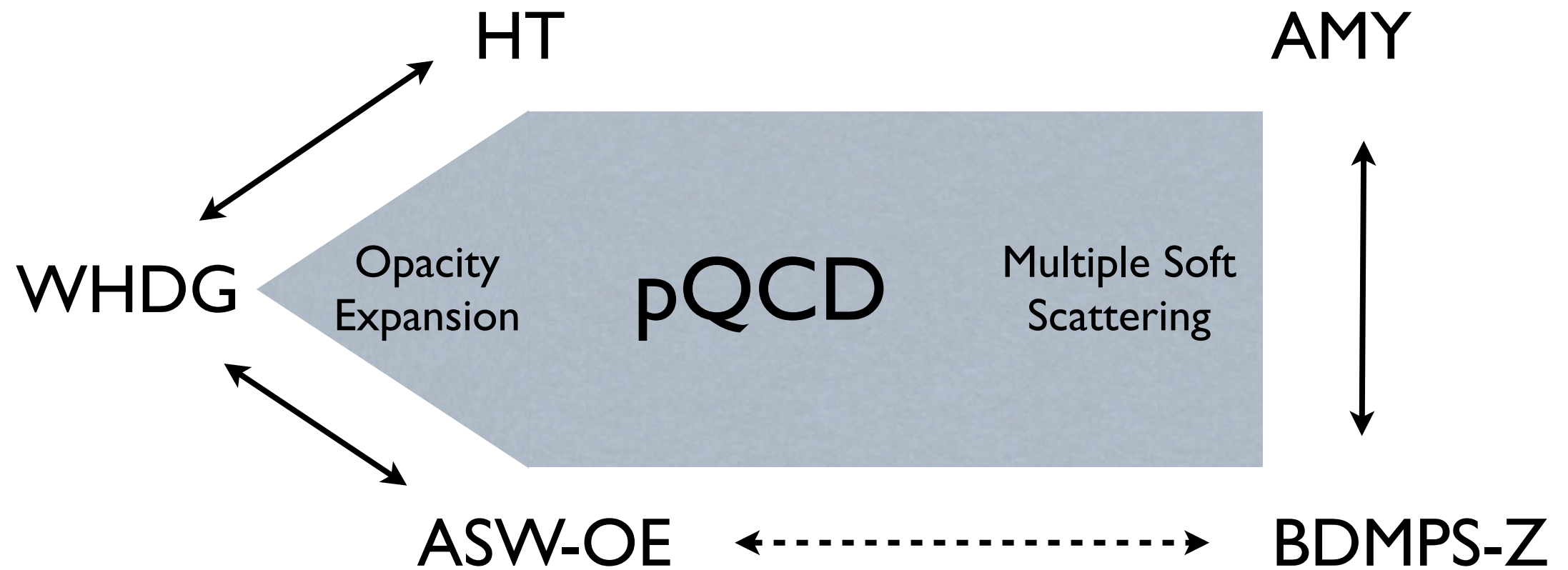
The Entrants - 1



The Entrants - 2



Making comparisons across the Jet Quenching Landscape



WHDG \leftrightarrow ASW-LOE

$$x \frac{dN_g^{\text{GLV}}}{dx} = \frac{C_R \alpha_s L}{\pi \lambda} \int \frac{d^2 \mathbf{q}}{\pi} \frac{\mu^2}{(\mathbf{q}^2 + \mu^2)^2} \int \frac{2d^2 \mathbf{k}}{\pi} \frac{\mathbf{k} \cdot \mathbf{q} (\mathbf{k} - \mathbf{q})^2}{\mathbf{k}^2 (\mathbf{k} - \mathbf{q})^4} \int dz \left[1 - \cos \left(\frac{(\mathbf{k} - \mathbf{q})^2}{2Ex} z \right) \right] \rho(z)$$

$$\rho(z) = \begin{cases} \frac{1}{L} \theta(L - z) \\ \frac{2}{L} \exp(-2z/L) \end{cases}$$

$$\omega \frac{dI^{\text{ASH-SH}}}{d\omega} = \frac{4\alpha_s C_R}{\pi} (n_0 L) \gamma \int_0^\infty \tilde{q} d\tilde{q} \left[\frac{\tilde{q}^2 - \sin \tilde{q}^2}{\tilde{q}^4} \right] \left(\frac{1}{\gamma + \tilde{q}^2} - \frac{1}{\sqrt{(\kappa^2 + \tilde{q}^2 + \gamma^2)^2 - 4\kappa^2 \tilde{q}^2}} \right)$$

$$\gamma = \tilde{\omega}_c / \omega, \tilde{\omega}_c = \frac{1}{2} \mu^2 L, \kappa = \sqrt{\omega L / 2}, \text{ and } n_0 L = L / \lambda$$

It all depends on what the meaning of “x” is....

$$x_+ = \frac{x_E}{2} \left(1 + \sqrt{1 - \left(\frac{k_\perp}{x_E E} \right)^2} \right)$$

WHDG

$$x_E = x_+ \left(1 + \left(\frac{k_\perp}{x_+ E_+} \right)^2 \right)$$

ASW

Identical in the k_\perp / ω limit, but ...

... but not when you consider the kinematic limit of k_{\perp} for given ω .

WHDG

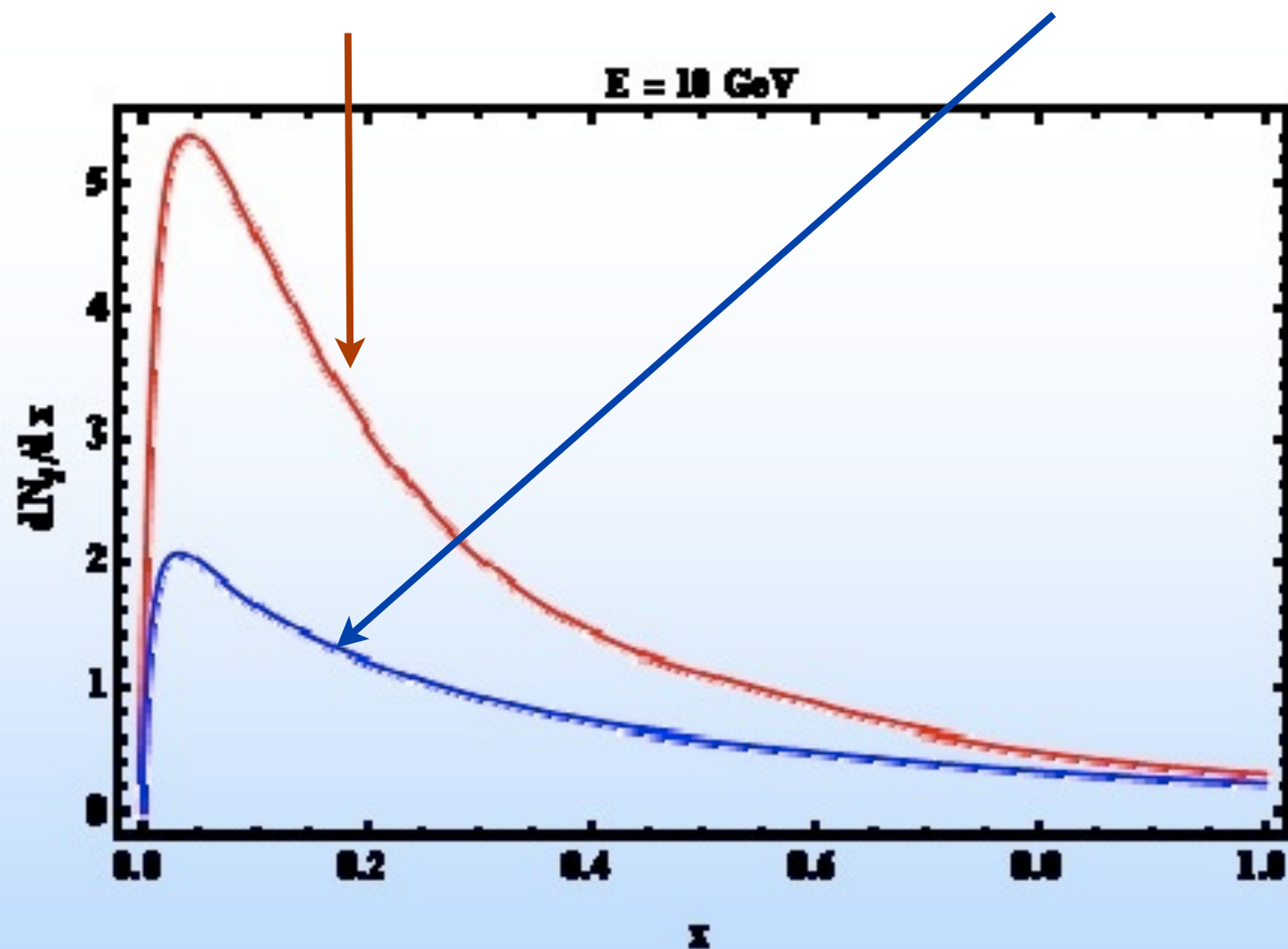
ASW

$$- k_T < x_+ E^+ = 2 x_+ E$$

$$- k_T < x_E E$$

$$\bullet k_{\max} = 2 x (1-x) E$$

$$\bullet k_{\max} = x E$$



Lesson:

Kinematic assumptions beyond the strict validity of the eikonal/collinear approximation can have drastic consequences even at high energies, because radiation always tries to exhaust the available phase space.

Vac-med interference vs. LPM

$$\omega \frac{dI^{(0)}}{d\omega dk_T} = \frac{\alpha_s C_R}{\pi^2} H(k_T) = \left| \text{diagram} \right|^2, \quad H(k_T) = \frac{1}{k_T^2}$$

Vacuum radiation

$$\left| \text{diagram 1} + \text{diagram 2} + \text{diagram 3} \right|^2 = |A(q_1)|^2 R(k; q_1) \equiv |A(q_1)|^2 \frac{q_1^2}{k^2 (k + q_1)^2}$$

WHDG
ASW-LOE
HT-LO

$$\left| \text{diagram 4} + \dots \right|^2$$

“True” LPM effect: Coherent action of multiple scatterings in the medium

$$\omega \frac{dI^{(1)}}{d\omega dk_T dq_1 dq_2} = \frac{\alpha_s C_R}{\pi^2} \left(|A(q_1)|^2 - V_{tot} \delta(q_1) \right) \left(|A(q_2)|^2 - V_{tot} \delta(q_2) \right)$$

BDMPS-Z AMY
GLV-HO ASW-HO

$$\times \left[\frac{(n_0 L)^2}{2} R(k + q_1, q_2) - n_0^2 \frac{1 - \cos LQ_1}{Q_1^2} R(k + q_1, q_2) + n_0^2 \frac{1 - \cos LQ_1}{Q_1^2} R(k; q_1 + q_2) \right]$$

Summary of insights: WHDG \leftrightarrow ASW-LOE

- Different definitions of the variables x_+ and x_E ;
- Importance of the kinematical region $k_\perp \sim \omega$, which violates the assumption of collinearity of the radiation process,
- Absence of exact energy and momentum conservation, both in the elementary process and in the convolution of successive radiation events (radiative cascade);
- Influence of different choices for the distribution of scattering centers (step function vs. exponential).
- Need for a consistent definition of q^\wedge for quantitative comparison with other models.

$$\hat{q} = \rho \int d^2 q_\perp q_\perp^2 \frac{d\sigma}{d^2 q_\perp}$$

Summary of insights: WHDG \leftrightarrow HT-LO

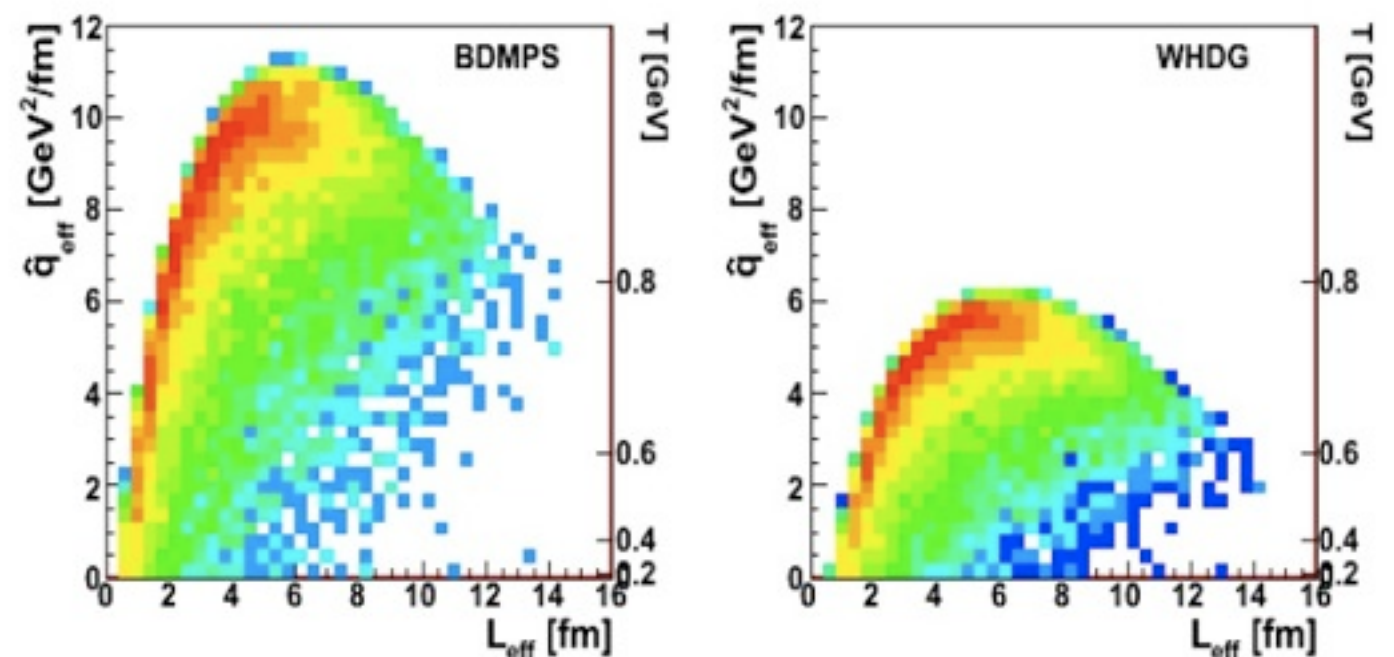
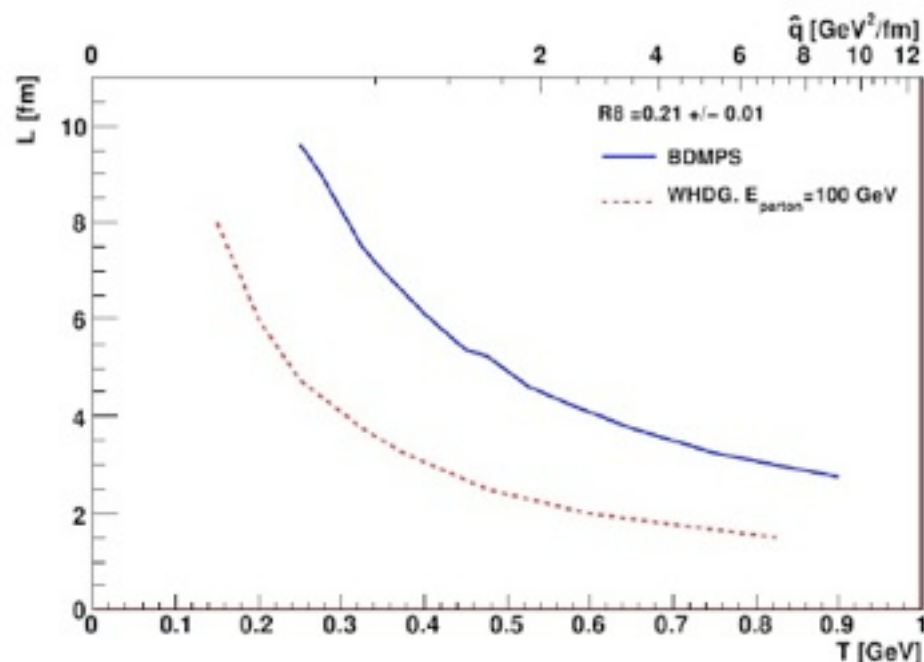
- HT implements energy-momentum conservation in the elementary process;
- HT assumes $k_{\perp} \gg q_{\perp}$, i.e. virtuality is dominated by primary hard scattering;
- HT encodes the running of $\alpha_s(Q^2)$. For a first quantitative comparison with WHDG one should set α_s constant, and then explore the quantitative importance of the running of α_s ;
- HT does not assume a specific model of the medium, but parametrizes the medium through a transport coefficient q^{\wedge} (and \hat{e} for elastic energy loss);
- HT exhibits sizable flavor change of leading parton.

Summary of insights: BDMPS-Z \leftrightarrow AMY

- AMY does not contain interference between vacuum and medium induced radiation;
- AMY implements exact energy and momentum conservation, both in the elementary process and in the radiative cascade;
- AMY treats the medium dynamically, not as collection of static scattering centers;
- Average energy loss is a bad approximation for true collisional energy loss;
- AMY and BDMPS both assume collinearity of the radiation; importance of large angle radiation needs to be studied;
- AMY exhibits sizable flavor change of leading parton.

Summary of insights: BDMPS-Z \leftrightarrow WHDG

- Assignment of $q^\wedge = \mu^2/\lambda$ in WHDG seems to underestimate the true value of q^\wedge by a factor 2-3;
- Correct definition: $q^\wedge = \langle k_\perp^2 \rangle / \lambda = \mu^2/\lambda_{\text{tr}}$.
- Results for WHDG and BDMPS-Z can be mapped into each other by a rescaling of T or L of the medium by factor ~ 2 ;
- Results of q^\wedge fit for a dynamical medium differ by factor ~ 2 .



Outlook (1)

- Origin of differences between 1st generation jet quenching formalisms is now well understood; they lie mostly outside the range of strict validity of the eikonal-collinear approximation.
- pQCD approach to jet quenching is alive and well.
- Reduction of uncertainty in q^{\wedge} from R_{AA} to \leq factor 2 seems possible with some effort.
- Most severe deficiencies are:
 - Energy-momentum conservation;
 - Vacuum radiation interference;
 - Consistent treatment of elastic & inelastic processes;
 - *Ad hoc* vacuum hadronization (?)

Outlook (2)

- A timely and comprehensive “TEC-HQM report” on the insights gained from the QGP Brick challenge would be a document of great value and with lasting impact.
 - “Timely” \approx 3 months (?)
 - “Comprehensive” \Leftrightarrow circulated outline sketch (?)
- “1st generation” jet quenching formalisms will remain the basis for MC schemes and detailed modeling of jet evolution.
- 1st generation jet quenching codes will also provide test cases for more sophisticated schemes.

Outlook (3)

- Some other questions:
 - Can I_{AA} vs. R_{AA} be used to check consistency of q^{\wedge} determination?
 - Can selection of jet virtuality (“jet mass”) be used to discriminate between VMI and LPM ?
 - Can we probe the validity of vacuum hadronization assumption ? It must fail somewhere! (recombination? heavy quark hadrons?)
 - Can we rule out that QCD jets become nonperturbative once they “see” the QGP? Can we rule out that pQCD does not apply to jets in a QCD medium? What kind of fragmentation pattern would a thin [i.e. $L \ll E/(dE/dx)$] “AdS/CFT Brick” produce?