High-Temperature Superconductor Coating for the Future Circular Collider Beam Pipe





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Future Circular Collider Study (FCC) Beam Screen Design



Cold bore diameter 44/47 mm

Nominal aperture:

- H: 29.6 mm
- V: 26.4 mm

Slit height: 2.9 mm \rightarrow 5 mm

Temperature: 50 ± 10 K

Beam screen wall

- 1.25 mm steel
- 0.3 mm copper
- 1 µm High-Temperature Superconductor
- Thin-film coating e.g amorphous carbon
 - \rightarrow Secondary electron emission yield



Future Circular Collider Study (FCC) Beam Screen Design



Basic Idea

- HTS carries the image currents to reduce power dissipated
- Provides lower electric impedance
- Reducing coupling impedance
 - Longitudinal and transversal direction
 - Small impedance for sufficient beam stability

Hybrid solution

- Stripes
- Overall symmetrical distributed
- Alternating HTSC and Cu
- Using Coated Conductors
- Overall covered with Carbon



Cu

YBCO

High-Temperature Superconductor



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Properties of HTS

- High critical values
- Anisotropic
- Ceramics
- **Brittle**
- Stoichiometry
- Type II Superconductor



Electroplat **Copper Stabilizer** Sputtering Silver Overlayer MOCVD (RE)BCO - HTS (epitaxial) 20 µm IBAD/Magnetron Sputtering **Buffer Stack** um Electropolishing -0.2 um Substrate 50 um ~1.8 µm 20 µm Configuration of SuperPower® 2G HTS Wire

	Parallel to c-axix	Parallel to ab-plane
$\lambda(nm)$	150	800
$J_C(Acm^{-2} 77K)$	10 ⁶	10 ⁷
$H_{c2}(\mathrm{T})$	110	240

W.Buckel, Supraleitung Grundlagen und Anwendung, Wiley-VCH Verlag GmbH & Co. KGaA

^(*) P.J.Ray, "Structural investigation of - Following staging as a function of temperature", M.Sc. Thesis, Niels Bohr Institute, Copenhagen, Denmark, 2015.

Coated Conductor

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Electrodynamics Surface Impedance

Normal Conductors

- Skin depth δ proportional to $\omega^{-1/2}$
- Surface Resistance proportional to $\omega^{1/2}$
- Surface resistance independent of temperature (low T)

ωμκ

$$Z_s = R_s + X_s = \frac{(1+i)}{\delta \cdot \kappa} \qquad \delta = \frac{1}{\delta \cdot \kappa}$$

Superconductors

- London Theory and Two-Fluid-Model
- London penetration depth λ_L independent of ω
- Surface resistance proportional to ω^2
- Surface resistance strongly dependent of temperature

$$R_{S} = \frac{1}{2} \kappa_{n} \mu_{o}^{2} \omega^{2} \lambda_{L}^{3} \qquad X_{S} = \omega \mu_{o} \lambda_{L} \qquad \lambda_{L} = \sqrt{\frac{m}{\mu_{o} n_{s} e^{2}}}$$





Surface Impedance
$$Z_s \equiv \frac{E_t}{H_t} \xrightarrow{\text{metal}} Z_{S_1} = \frac{(1+\iota)}{\kappa_1 \cdot \delta_1}$$

Boundary conditions

 $\vec{n} \vec{\varkappa} \left(\vec{E}_1 \vec{E}_2 \vec{E}_v \vec{B}_v \right) \oplus 0 \qquad \vec{n} \times \left(\vec{E}_1 - \vec{E}_2 \right) = 0$ $\vec{n} \vec{\varkappa} \left(\vec{H}_1 \vec{H}_2 \vec{H}_v \vec{F}_v \right) \oplus 0 \qquad \vec{n} \times \left(\vec{H}_1 - \vec{H}_2 \right) = 0$

Surface impedance for 2nd layer

$$Z_{S_{2}}(Z_{S_{1}}) = \frac{1+A_{2}}{1-A_{2}} \cdot \alpha_{2}$$

$$Z_{S_{2}} \qquad Z_{S_{1}}$$

$$A_{2} = \frac{[Z_{S_{1}}-\alpha_{2}]\cdot\exp\{-D_{2}\}}{[Z_{S_{1}}+\alpha_{2}]\cdot\exp\{+D_{2}\}} \cdot \alpha_{2}$$

$$\alpha_{2} = \frac{(1+i)}{2} \cdot \omega\delta_{2}\mu_{2}$$

$$D_{2} = (1+i) \cdot \frac{d_{2}}{\delta_{2}}$$

$$\delta_{2} = \sqrt{\frac{2}{\mu_{2}\kappa_{2}\omega}}$$







Boundary conditions

 $\vec{n} \times \left(\vec{E}_3 - \vec{E}_v\right) = 0 \qquad \qquad \vec{n} \times \left(\vec{E}_2 - \vec{E}_3\right) = 0$ $\vec{n} \times \left(\vec{H}_3 - \vec{H}_v\right) = 0 \qquad \qquad \vec{n} \times \left(\vec{H}_2 - \vec{H}_3\right) = 0$

Surface impedance for 3nd layer

$$Z_{S_3}(Z_{S_2}) = \frac{1+A_3}{1-A_3} \cdot \alpha_3 \qquad \qquad Z_{S_3} \quad Z_{S_2} \quad Z_{S_1}$$

$$A_3 = \frac{[Z_{S_2} - \alpha_3] \cdot \exp\{-D_3\}}{[Z_{S_2} + \alpha_3] \cdot \exp\{+D_3\}} \cdot \alpha_3$$

$$\alpha_3 = \frac{(1+i)}{2} \cdot \omega \delta_3 \mu_3 \qquad \qquad D_3 = (1+i) \cdot \frac{d_2}{\delta_2} \qquad \qquad \delta_3 = \sqrt{\frac{2}{\mu_3 \kappa_3 \omega}}$$





Surface Impedance
$$Z_s \equiv \frac{E_t}{H_t} \xrightarrow{\text{metal}} Z_{S_1} = \frac{(1+i)}{\kappa_1 \cdot \delta_1}$$

Boundary conditions

 $\vec{n} \times \left(\vec{E}_k - \vec{E}_{k+1}\right) = 0$ $\vec{n} \times \left(\vec{H}_k - \vec{H}_{k+1}\right) = 0$

Surface impedance for k-th layer

$$Z_{S_k}(Z_{S_{k-1}}) = \frac{1+A_k}{1-A_k} \cdot \alpha_k$$

$$A_k = \frac{[Z_{Sk-1} - \alpha_k] \cdot \exp\{-D_k\}}{[Z_{Sk-1} + \alpha_k] \cdot \exp\{+D_k\}} \cdot \alpha_k$$

$$\alpha_k = \frac{(1+i)}{2} \cdot \omega \delta_k \mu_k \qquad \qquad D_k = (1+i) \cdot \frac{d_k}{\delta_k}$$







Impedance Calculation **Beam Screen : 2D impedance simulation (BI2D)**



BI2D: 2D frequency domain solver

- Computational tool for longitudinal and transverse coupling impedance
- Finite-Element-Method (FEM) ٠
- Input: mesh of the geometry ٠

Two methods:

- Surface Impedance Boundary Condition (SIBC)
- Meshiper whole swart for Population frequency

$$f = \frac{1}{\kappa_n \pi \mu_0 \delta^2} \approx 469 \text{ Hz}$$

YBCO : penetration depth (frequency independent)

 $\lambda_L(50K) = 157 \text{ nm}$



- $\delta = 300 \ \mu m$
- 1 µm thickness

Space charge and resistive wall impedance computation in the frequency domain using the finite element method Uwe Niedermayer, Oliver Boine-Frankenheim, and Herbert De Gersem

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Impedance Calculation Transversal Impedance - Effect of the Radiation Slit



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YBCO







 10^{6} Cu YBCO -Wide tape 10^{5} position 2mm tapes dependent impedance 10^{4} 10³- $Re(Z_{\perp})[\Omega/m]$ 10² 10^{1} Low-frequency instability 10⁰ 10^{-1} 10^{-2} 10-3 10³ 10⁵ 106 10^{7} 10^{8} 10⁹ 10^{10} 10^{11} 10^{2} 10^{4} f [Hz]

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Cu YBCO

Conclusion & Outlook Impedance



Intended reduction of the impedance with HTS can be reached

- Effect of the radiation slit
 - Coating of the ante-chamber could reduce impedance further
 - What is the effect on the intended function?
- Hybrid solution
 - Percentage of superconducting surface yields a gradual reduction of the impedance
 - Using narrow stripes shows lower impedance at low frequencies
 - Suggested solution : Hybrid system with very narrow alternating stripes
 - Further simulations with different distributions should be done

Secondary electron emission yield

- Amorphous Carbon coating to reduce the electron cloud effect
- · Just a small raise in the resistance visible
- The question arises how the coating thickness relates to the SEY value
- Dust-particle-beam-interactions?
- Coating properties need extensive experimental validation

Conclusion & Outlook Experiments & Collaborations



There are two collaborations on HTS which have been established under the FCC by Sergio Calatroni - (Cern)



Conclusion & Outlook Magnetic field





Conclusion & Outlook Magnetic field



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Question:

What value has the distortion of the magnetic field?



- High synchrotron radiation load of protons @ 50 TeV
 - High synchrotron radiation load of protons @ 50 re
 Heat deposition of 31 W/m on beam screen edge (*)
 - HTS strongly dependent on temperature
 - Effect on the superconducting material

Assumption:

- not applicable for wide tapes
- Properties depend on field orientation
- Flux Line Lattice behaviour $T \approx \frac{T_c}{2}$
- stronger magnetic field
 - more vortices lead to less distortion
- Mechanical analysis needed
 - During magnetic quench
- Hysteresis
 - Pinning and impurities
 - Multipole field error
 - Destroying superconductivity by warming → unattainable
 - Has to be quantified

Superconductor surface resistance in the presence of a dc magnetic field: frequency and field intensity limits **Sergio Calatroni** – (submitted to PR-AB)



The presence of a magnetic field modifies the surface impedance of the superconductor

- Flux Line Lattice
 - Normal conducting cores
 - Ratio of normal conducting and superconducting surface area
 - Dependent on magnetic field strength

"Low frequency large field" – approximation

$$R_{sf} = \frac{R_n}{\sqrt{2}} \sqrt{\frac{B}{B_{c2}}} \left(\frac{\omega}{\omega_{dep}}\right)^{3/2}$$

- Resistivity scales with reduced magnetic field $\frac{B}{B_{c2}}$
 - At 50 K the critical field B_{c2} is in the range of 30-50 T
- $R_n \propto \sqrt{\omega} \rightarrow R_{sf}$ still dependent on frequency as ω^2



Conclusion & Outlook



Intended reduction of the impedance with HTS can be reached

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- The question arises how the coating thickness relates to the SEY value
- Dust-particle-beam-interactions?
- Coating properties need extensive experimental validation
- Need to measure RF performance as a function of T, B, J at the relevant frequencies for the FCC
- Thermal, mechanical and cryogenic aspects have to be studied

HTS coating? Lots of work Impedance, coating technology, ecloud, ...



Thank you for your attention!





Impedance Calculation Longitudinal Impedance - Effect of the Radiation Slit



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Various Configurations





Impedance Calculation (BI2D) Longitudinal Impedance





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YBCO

Impedance Calculation (BI2D) Longitudinal Impedance





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YBCO





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Cu YBCO





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YBCO



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YBCO

Cu

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YBCO

Cu

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— ҮВСО

Cu

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HTS Coating Magnetic Field – Vortex

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H-T phase diagram for cuprate superconductors shows several phases of vortex matters:

- ordered vortex-lattice at low fields $H_{c1} < H \ll H_{c2}$
- highly disordered vortex-solid at high fields and low temperatures (vortex-glass)
- vortex-liquid at high temperatures





N.Plakida, High-Temperature Cuprate Superconductors, Spinger Heidelberg

*HTS Coating normal dc-conductivity

Surface Resistance $R_S = \frac{1}{2}\kappa_n \mu_o^2 \omega^2 \lambda^3 \rightarrow \kappa_n = \frac{2R_S}{\mu_o^2 \omega^2 \lambda^3}$

London penetration depth for YBCO $\lambda \approx 150 nm$ Buckel, Werner: Supraleitung, Grundlagen und Anwendungen, 7.Auflage, Wiley-VHC Verlag GmbH & Co.KGaA, 2013

Two fluid model predicts temperature dependence

For
$$T = 77K$$
, $f = 1GHz$ and $R_S = 2.1 \cdot 10^{-6} \Omega$

$$\kappa_n(K) = \frac{2R_S}{\mu_o^2 \omega^2 \lambda (77K)^3} = 7.7 \cdot 10^6 \, S/m$$

Scaling the conductivity with TFM $\kappa_n(T) = \kappa_0 \left(\frac{T}{T_c}\right)^4$

Ratio of
$$\frac{\kappa_n(50K)}{\kappa_n(77K)} \rightarrow \kappa_n(50K) = 1.37 \cdot 10^6 S/m$$





$$\lambda(T) = \lambda(0) \left[1 - \left(\frac{T}{T_c}\right)^4 \right]^{-1/2} \lambda(50K) = 157 \, nm$$
$$\lambda(77K) = 206 \, nm$$

^{09.11.2016 |} TU Darmstadt | TEMF | Patrick Krkotić | EuroCirCol Meeting - Barcelona | 34 Springer-Verlag Berlin Heidelberg (1999)