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# CT update: towards CT17 and studies of intrinsic charm

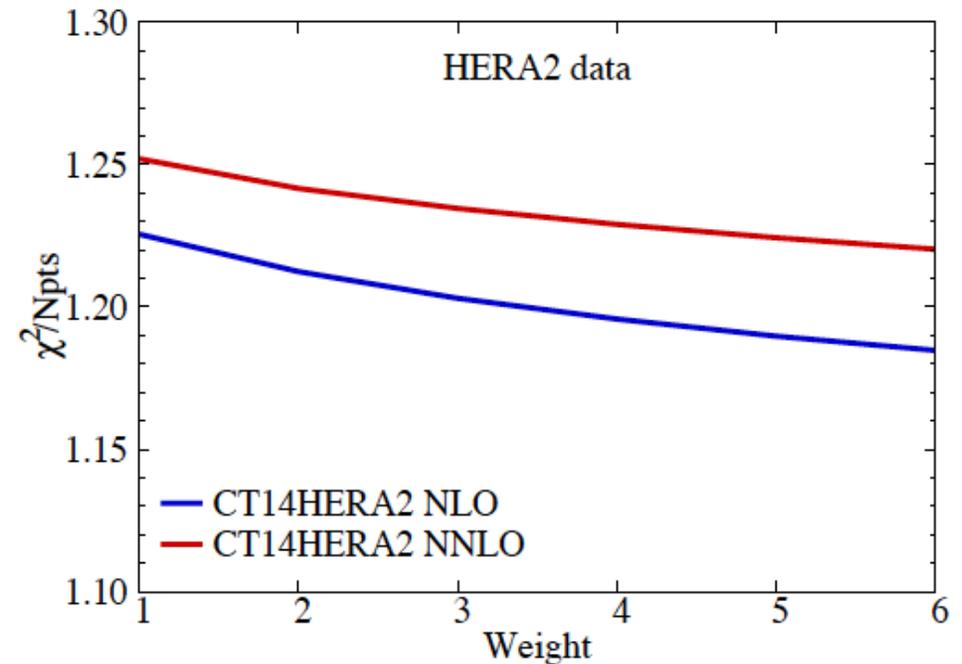
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MSU

for the CTEQ-TEA collaboration

# HERA2

- In CT17, HERA1 data will be replaced by HERA1+2 data
- Impact investigated in PRD95(2017)034003 (arXiv: 1609.07968)
- $\chi^2$  slightly better at NLO than at NNLO
- $\chi^2$  does not improve greatly as weight of the data-set is increased
- $\chi^2$  for rest of the data does not change much if HERA2 used instead of HERA1

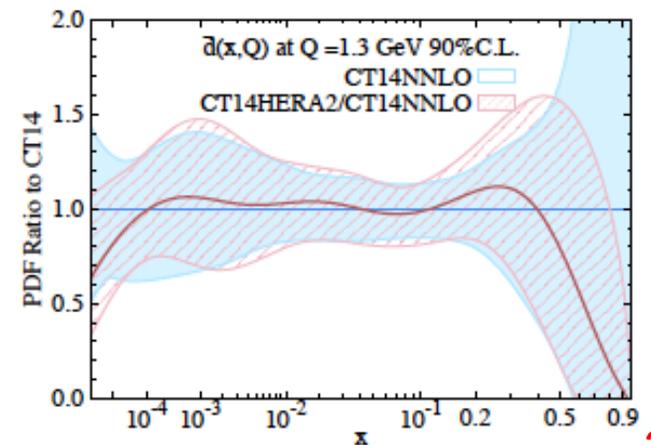
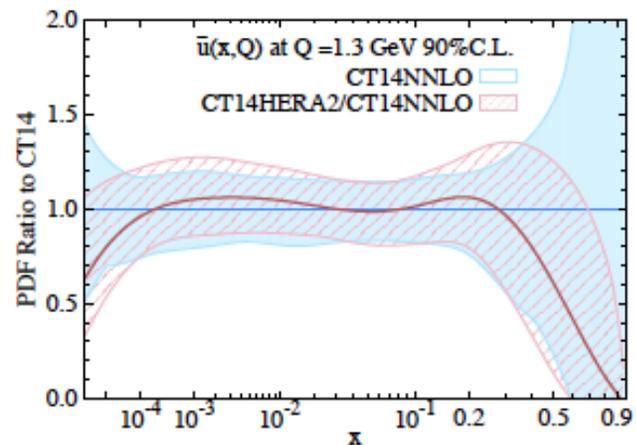
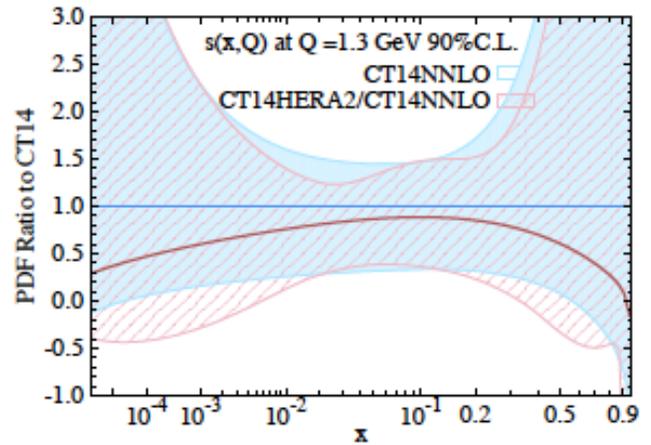
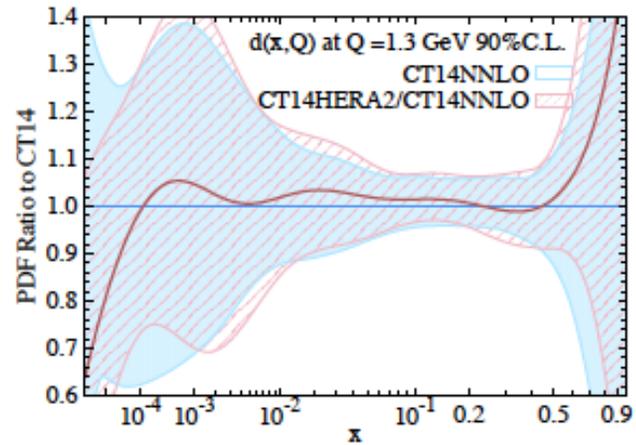
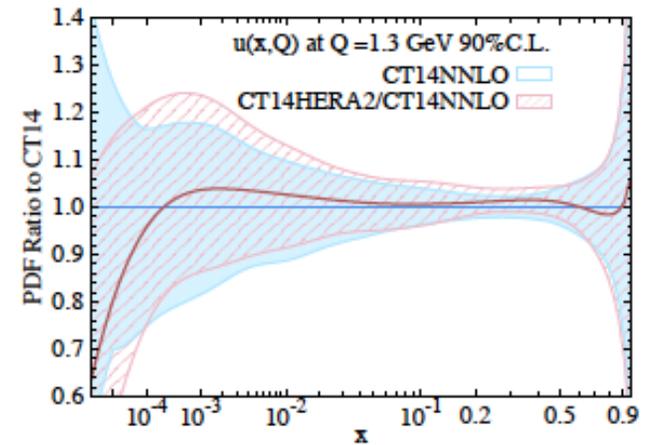
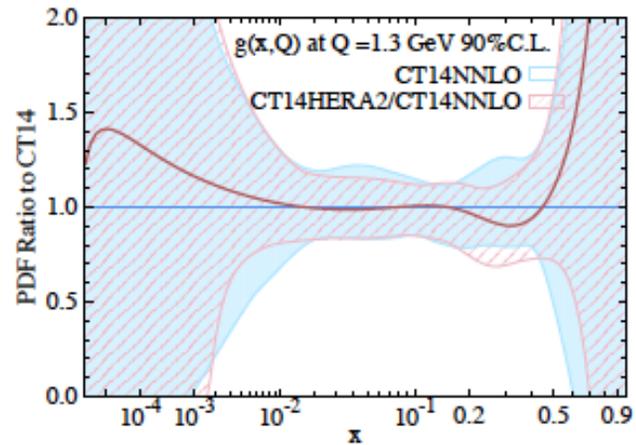


- Resultant changes in PDFs well within CT14 uncertainty bands

- In addition to replacing HERA1 with HERA1+2 data set in CT17, the NMC  $\mu p$  cross section will be dropped

◆ never fit well

- Strange quark distribution no longer tied to  $u\bar{u} + d\bar{d}$
- Consequences at low  $x$

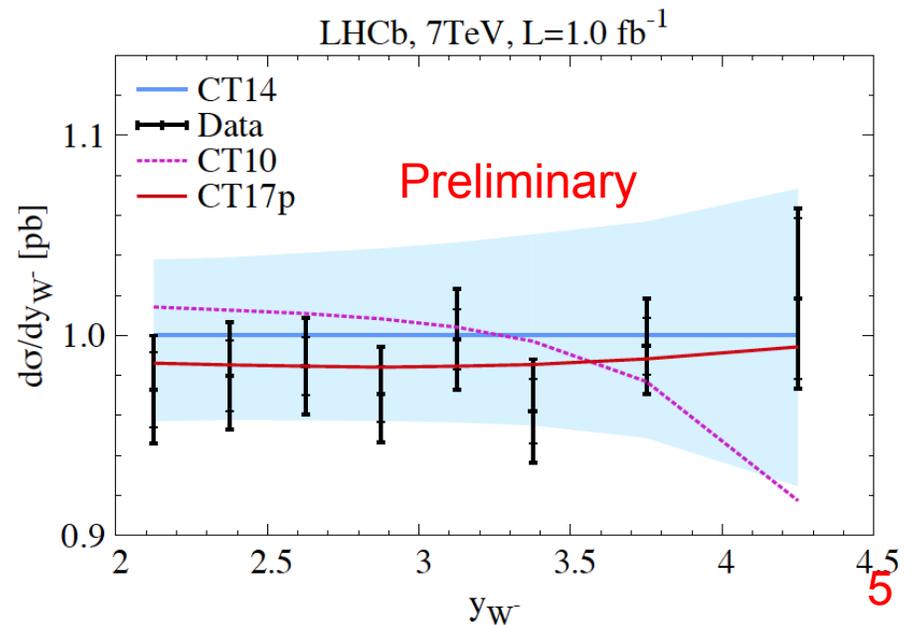
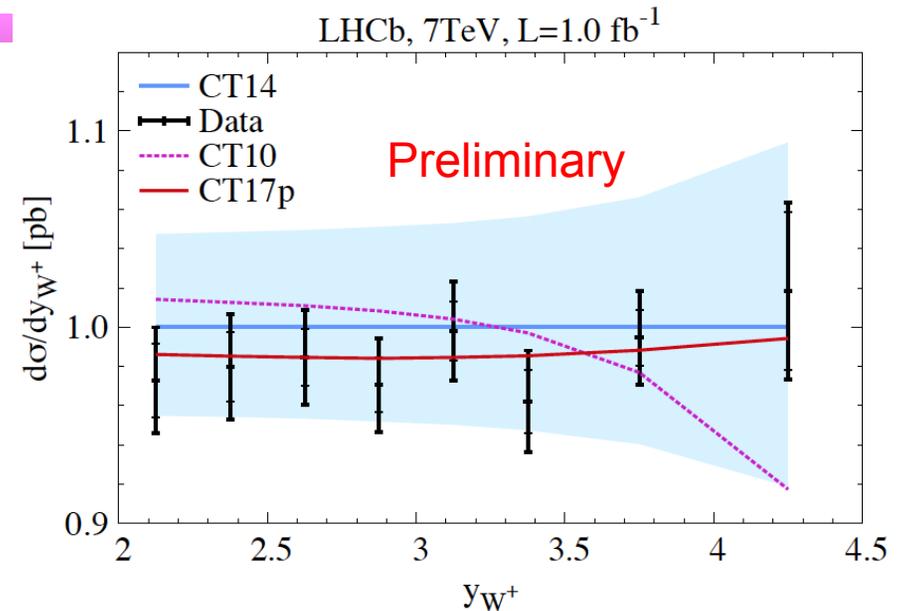
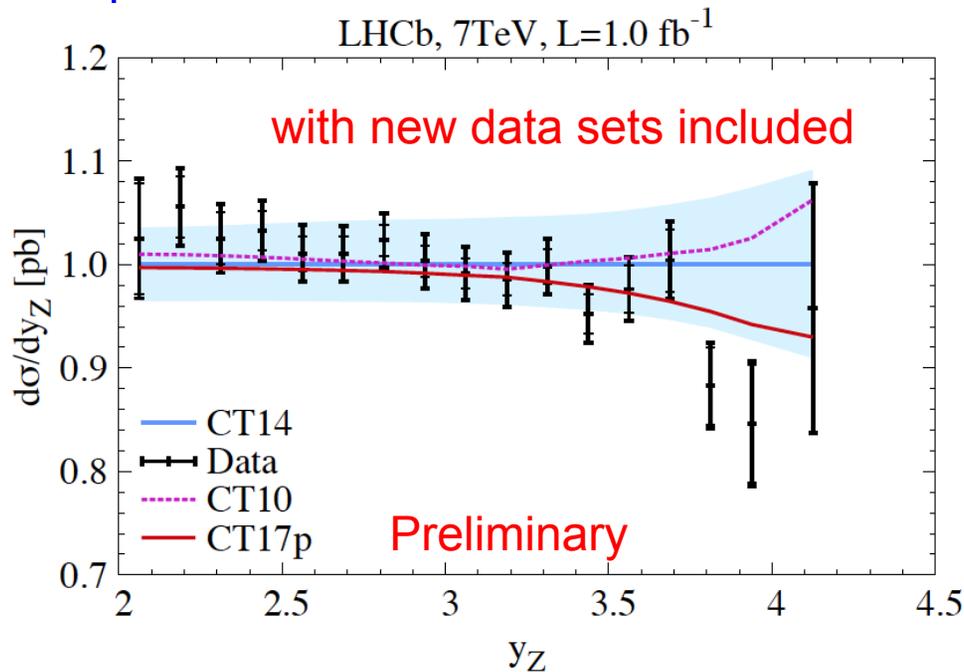


# Other data sets to be included in CT17

- Use as much relevant LHC data as possible using applgrid/fastNLO interfaces to data sets, with NNLO/NLO K-factors
- Implementing a parallelization of the global PDF fitting to allow for faster turn-around time
- Next few slides are a snapshot
- More details at DIS2017
- I will concentrate on impact of data in bold
- **LHCb W/Z rapidity at 7 and 8 TeV (applgrid)**
- LHCb Z  $y$  at 13 TeV (applgrid)
- ATLAS W/Z rapidity at 7 TeV (applgrid)
- ATLAS 8 TeV DY (applgrid)
- ATLAS 7 TeV W/Z  $p_T$  (applgrid)
- **ATLAS 7,8 Z  $p_T$ , as a function of mass (applgrid)**
- CMS W,Z  $p_T$ , as a function of  $y$ , at 8 TeV (applgrid)
- CMS W rapidity at 8 TeV (applgrid)
- CMS W/Z  $p_T$  at 8 TeV (applgrid)
- CMS inclusive jet cross section at 7 TeV with  $R=0.7$  (fastNLO)
- **ATLAS inclusive jet cross section at 7 TeV with  $R=0.6$  (applgrid)**
- ATLAS and CMS 7,8 TeV  $t\bar{t}$  differential distributions
- ATLAS low mass/high mass Drell-Yan at 7 TeV
- CMS low mass/high mass DY at 8 TeV

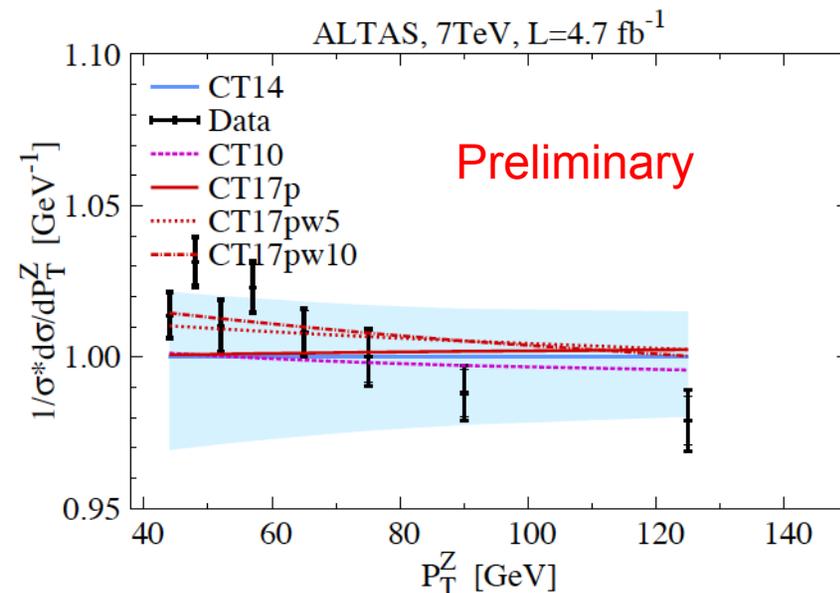
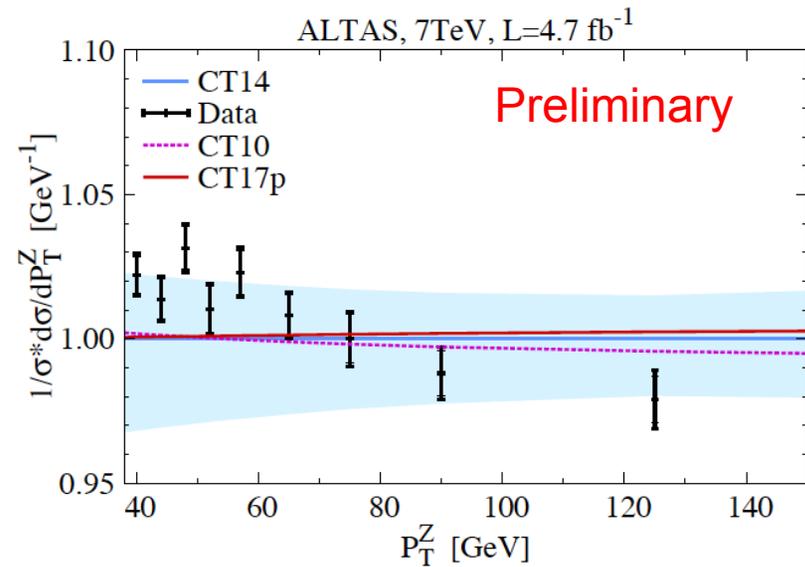
# LHCb W,Z rapidity data

- 7 TeV: W,Z muon rapidity
- 8 TeV Z(->ee) rapidity
- Unshifted data; outer error bars are with stat and syst errors added in quadrature; inner error bars are uncorrelated errors alone
- Major change is an increase in the strange quark distribution to that similar to CT10



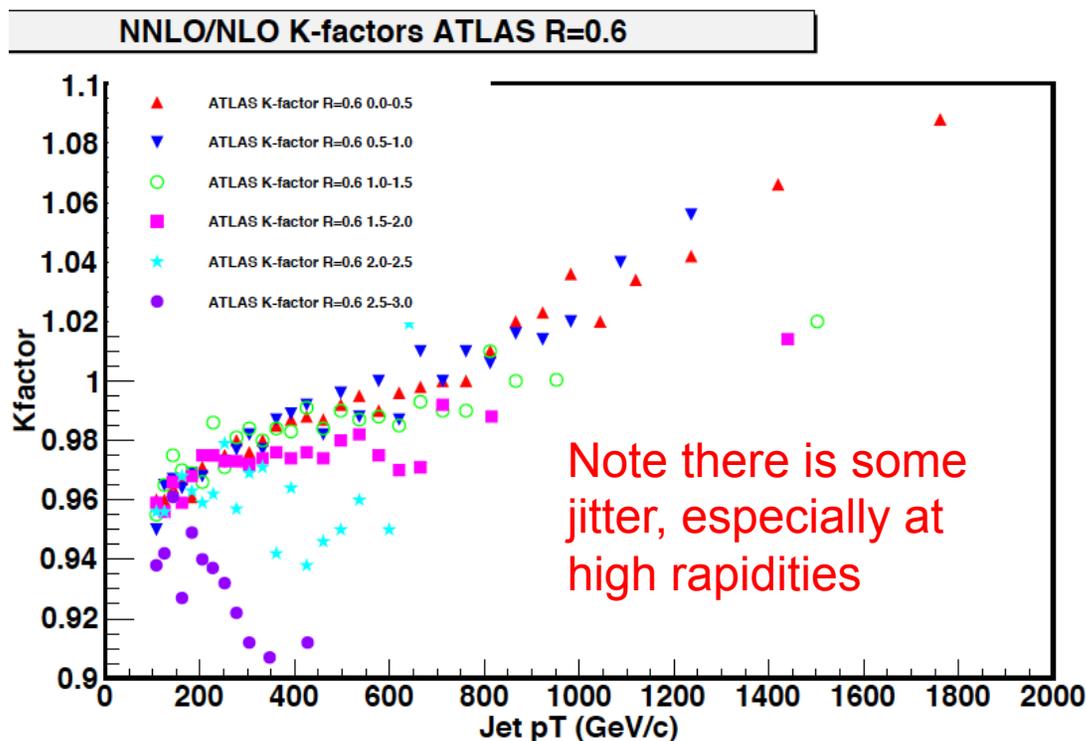
# ATLAS Z $p_T$

- Start with using normalized Z  $p_T$  data in range from 40 to 150 GeV/c at NLO, with  $\mu_R = \mu_F = m_T$ 
  - ◆ prefers softer gluon at high  $x$
  - ◆ adding additional weight improves agreement
- Also trying absolute cross section
- Will add in NNLO/NLO K-factors, extend to full range
  - ◆ NNLO/NLO K-factors for this process have jitter, need an additional uncorrelated error



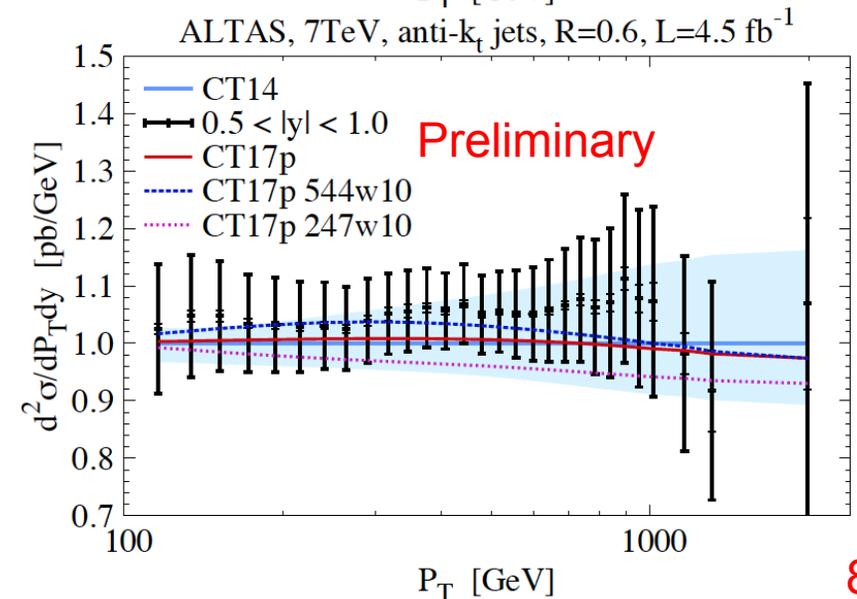
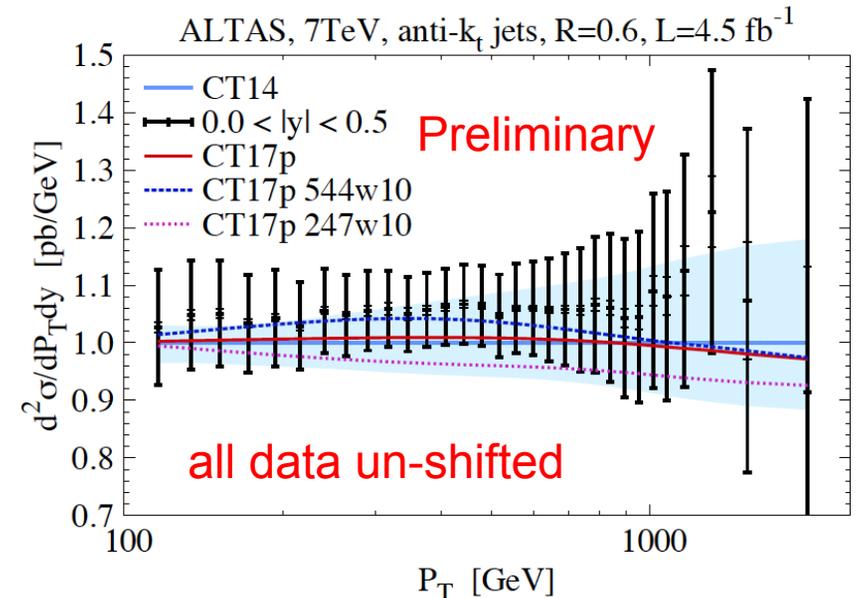
# Some jet issues

- For LHC jet cross sections, use largest jet sizes available ( $R=0.6$  ATLAS,  $R=0.7$  CMS) and a scale choice of  $p_{T}^{\text{jet}}$ 
  - ◆ traditional choice for PDF fitting
- With such choices, K-factors (NNLO/NLO) constant and close to unity
- But tension within ATLAS jet rapidity data sets, as well-known within this group
- Not quite fair to choose just one rapidity interval (say 0.0-0.5) if another rapidity interval (say 0.5-1.0) gives a different PDF result
- Currently under investigation; what is the envelope of PDFs resulting from fits to individual rapidity bins, i.e. how unanimous are the rapidity bin data

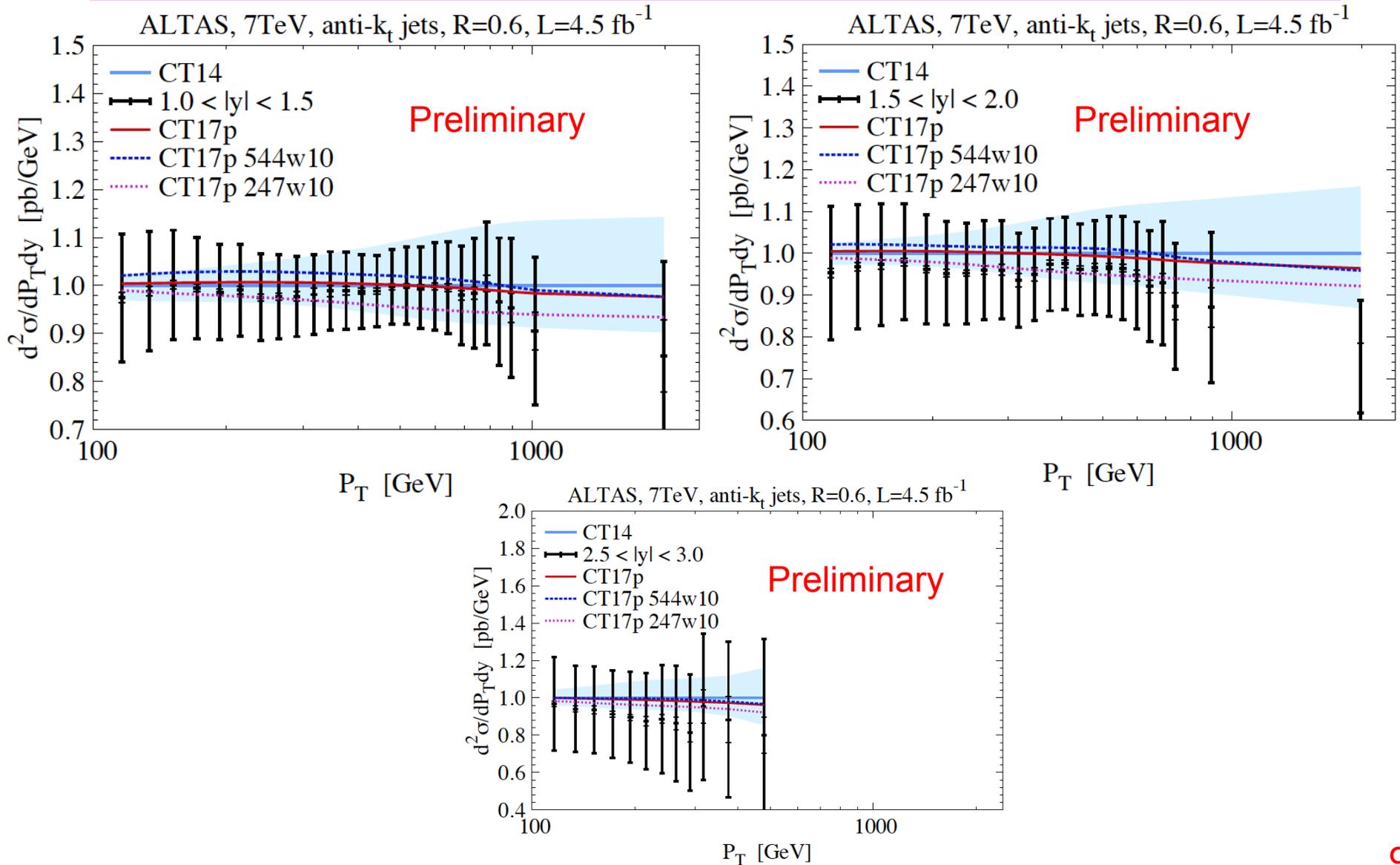


# ATLAS 7 TeV jet data

- Can not get a good fit simultaneously to all rapidity bins:  $\chi^2/\text{DOF} \sim 2$ , even with weight 10 for combined jet data
  - ◆ outer error bars are with stat and syst errors added in quadrature; inner error bars are uncorrelated errors alone
- Individual rapidity bins can be fit well
- In general, ATLAS jet data favors a larger gluon PDF around  $x=0.2$
- Some tension with ATLAS Z  $p_T$  data regarding high  $x$  gluon
  - ◆ jet data likes a harder gluon ( $x \sim 0.2$ )
  - ◆ Z  $p_T$  data (normalized) likes a softer gluon

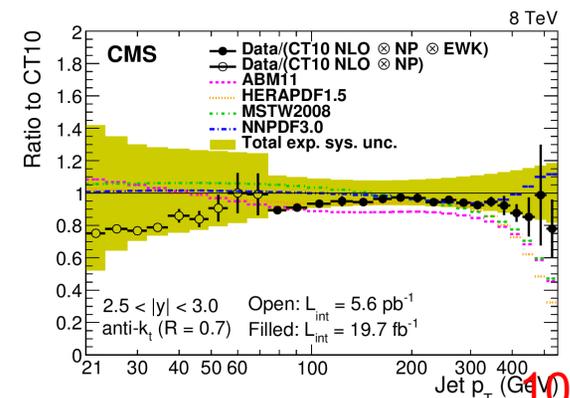
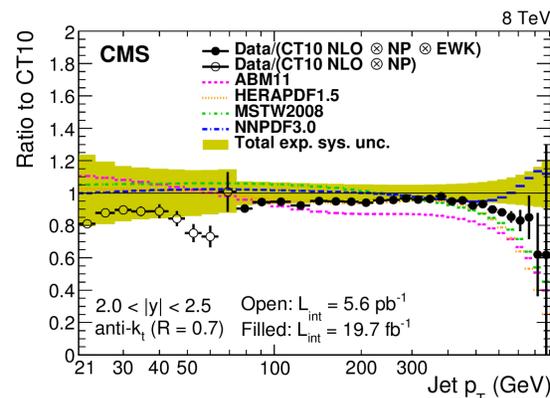
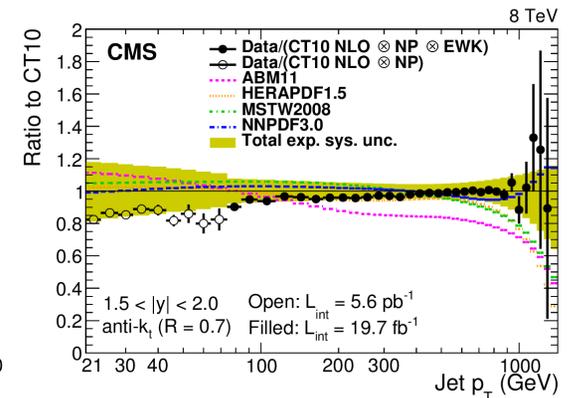
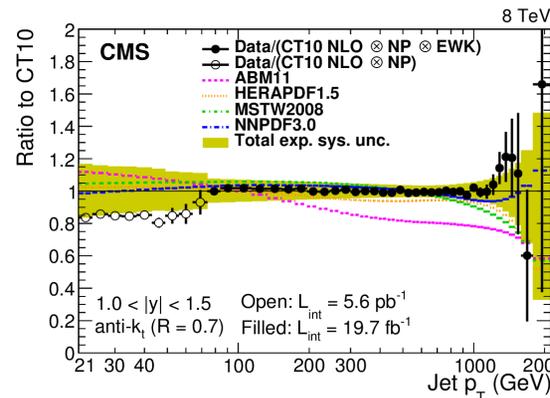
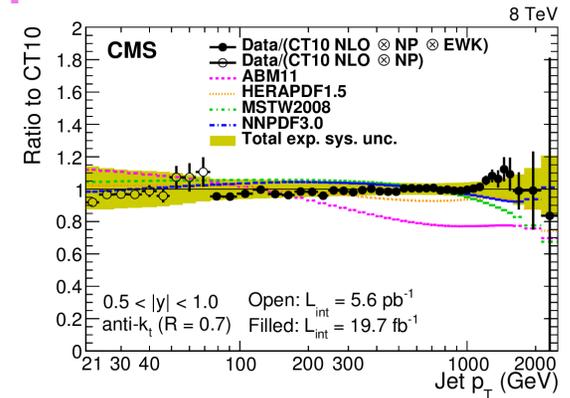
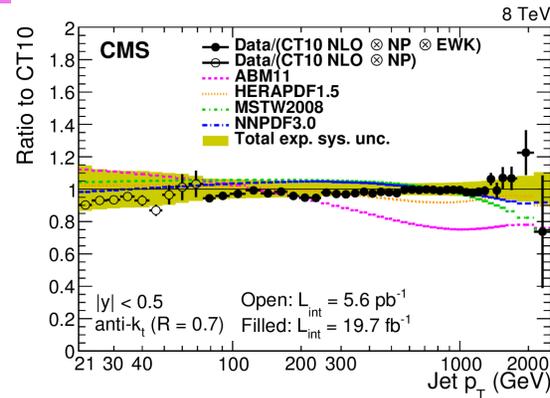


# ATLAS 7 TeV jet data



# CMS 8 TeV jet data

- CT10 has a harder gluon than CT14
- CMS data seems happy with that
- I'm happy with that
- ...but may point out a tension between the ATLAS and CMS jet data sets; if so, high  $x$  gluon uncertainty may not be reduced by these data sets
- Tension may need a better understanding
- Will use 8 TeV CMS data if available in near-term



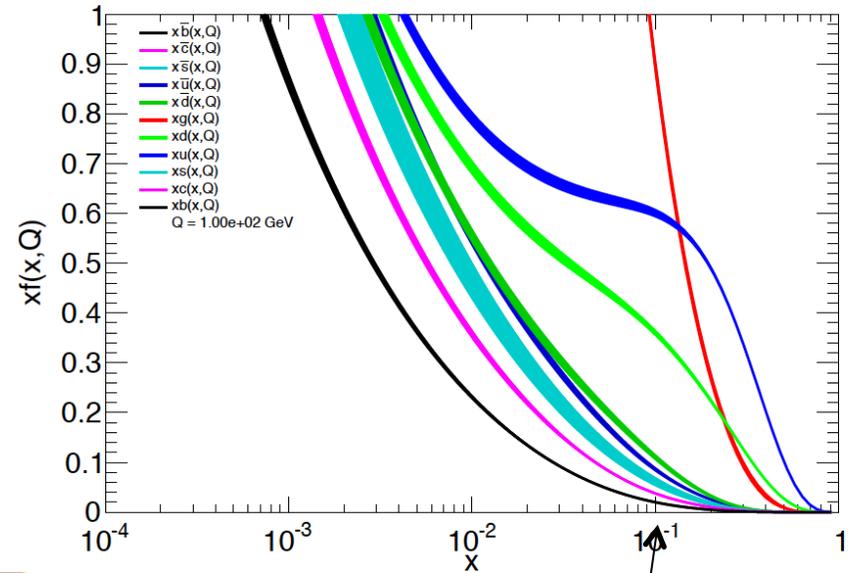
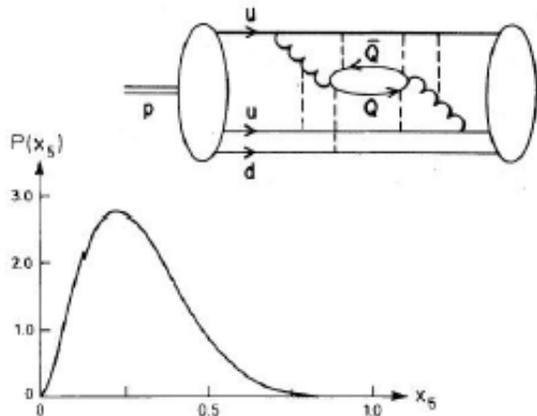
# Top distributions

- There are several distributions measured by ATLAS and CMS that have information on the high  $x$  gluon
  - ◆  $m_{tT}$ ,  $y_{tT}$ ,  $p_T^{t,T}$  directly
  - ◆  $y_{t,T}$ ,  $p_T^{tT}$  indirectly
- Only one distribution should be used, unless a correlation model can be developed
  - ◆ which one?
  - ◆ do they give the same answer? if not, do we understand why?
  - ◆ how do the constraints/trends from each distribution compare?
  - ◆ similar to what we were talking about with ATLAS jet data in different  $y$  bins
- We are currently doing exploratory studies at NLO using MCFM and DiffTop and at NNLO using DiffTop
  - ◆ starting with the  $p_T$  and rapidity of the top quark
- ATLAS and CMS have different trends; in this case, ATLAS favors harder gluon at high  $x$ , CMS weaker gluon
- In general, the ATLAS and CMS top results are in tension internally, and with each other (the latter more so in the case of normalized distributions where the experimental errors are smaller)
- This is similar to the tension that exists between the ATLAS and CMS jet data, although there the tension is in the opposite direction
- If tension, then gluon PDF uncertainty may not decrease and may even increase
- Study in progress

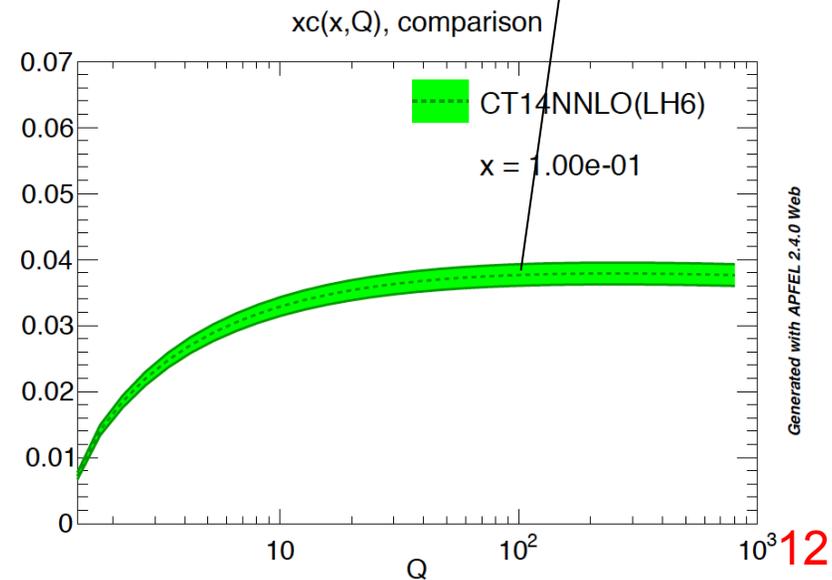
# Charm

- The charm quark distribution is generated perturbatively through gluon splitting
- So normally no charm below  $c\bar{c}$  threshold
- But what if there is an intrinsic charm present in the proton at low  $Q$
- This has been Stan Brodsky's dream for some time

BHPS PLB93B (1980) 451  
 Brodsky et al: arXiv:1504.06287



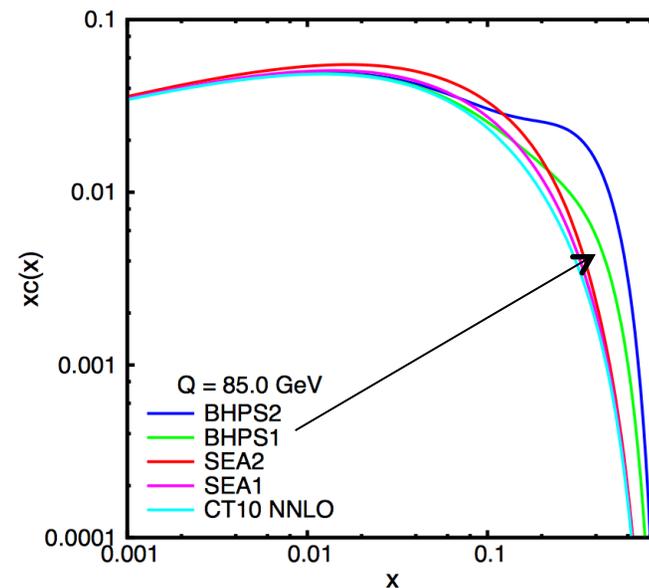
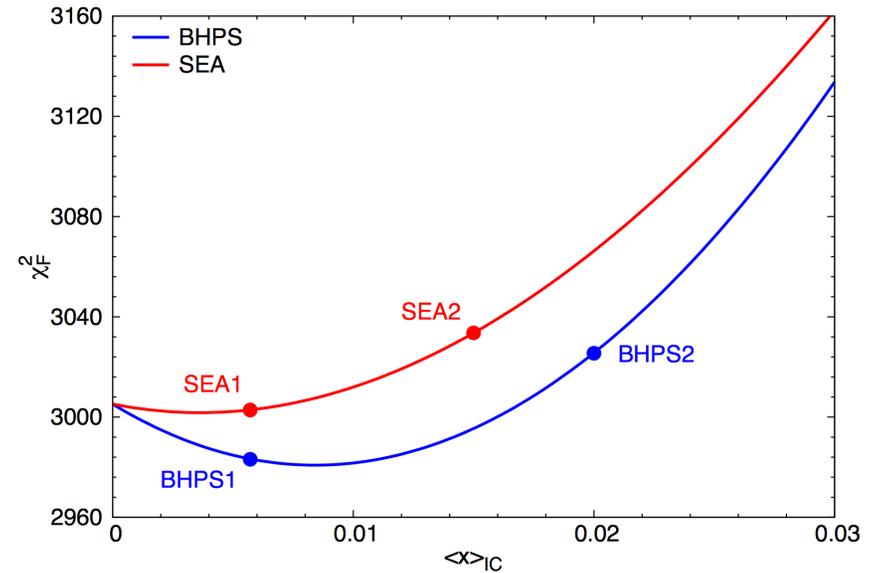
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# Intrinsic charm

- ...and has been studied by CTEQ in, for example, arXiv: 1309.0025 and in proceedings of DIS2014
  - ◆ these analyses carried out at NNLO
- Two types of models: Brodsky-like (valence-like) or Sea-quark like
- One Brodsky-like model, BHPS1 actually leads to a modest reduction in  $\chi^2$ , but as we said in the paper, it's interesting, but not enough to claim the discovery of intrinsic charm



# Update

- Return to study of intrinsic charm within CT14/CT14HERA2 framework, again with 2 models of intrinsic charm

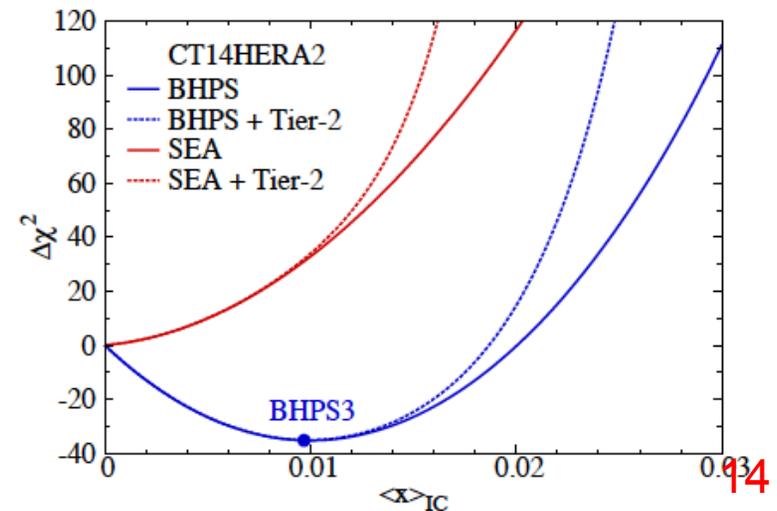
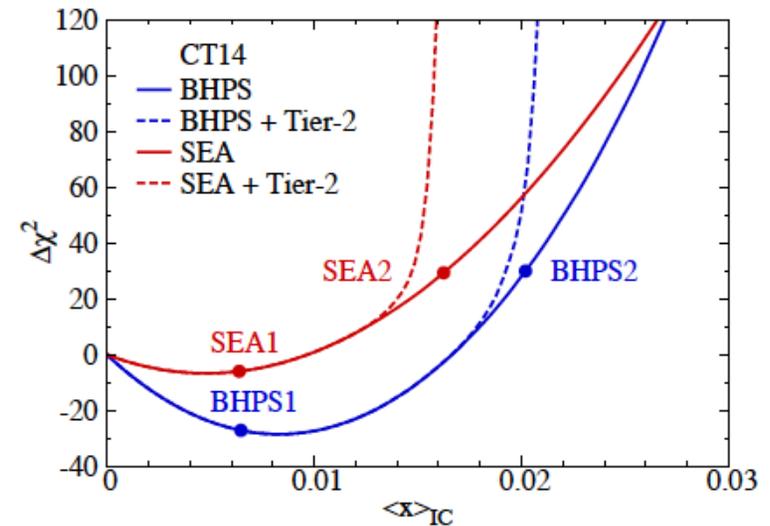
- ◆ Valence-like

$$\hat{c}(x) = A x^2 [6x(1+x)\ln x + (1-x)(1+10x+x^2)]$$

- ◆ Sea-like

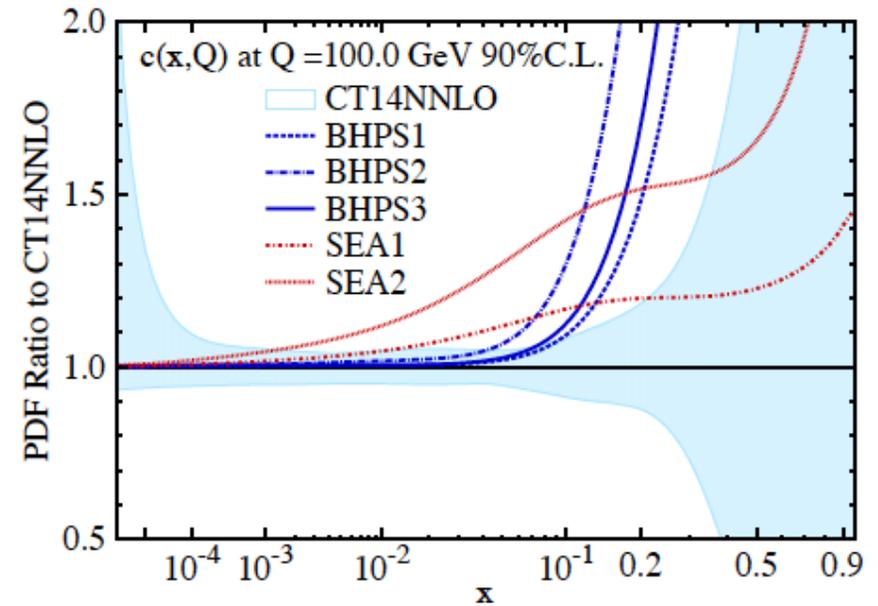
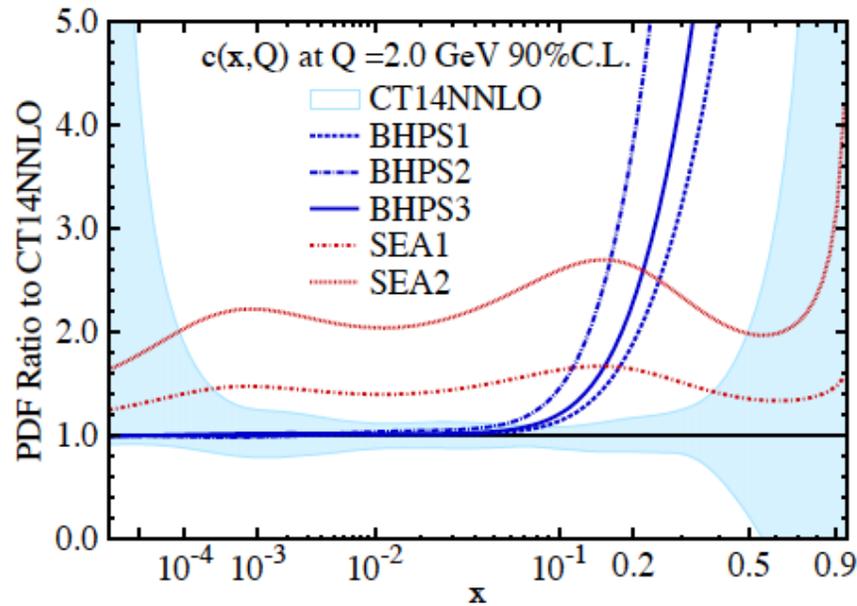
$$\hat{c}(x) = A (\bar{d}(x, Q_0) + \bar{u}(x, Q_0))$$

- See reduction in valence-like (BHPS) models, but with the drop in  $c_2$  interesting but less than our criterion (100)
- In addition, the decrease in  $c_2$  comes primarily from BCDMS data that does not have any particular sensitivity to charm

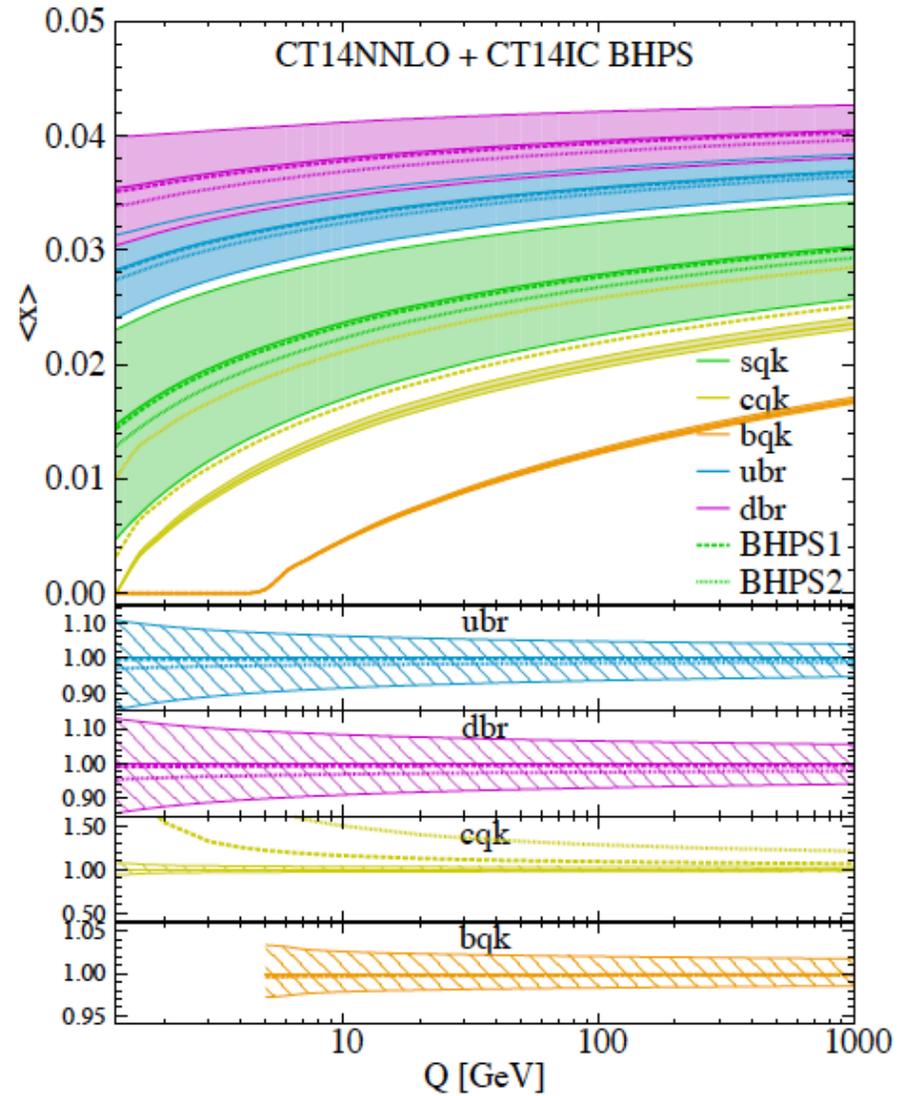
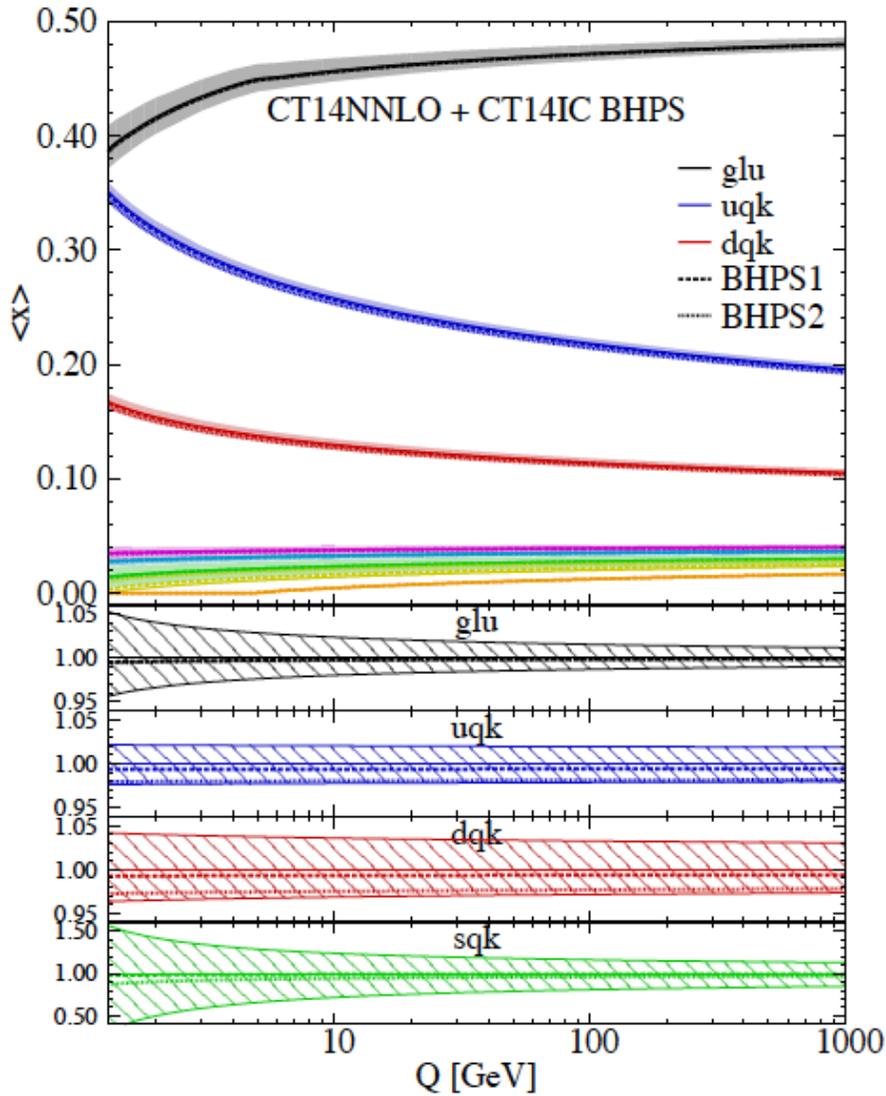


# Charm quark distributions

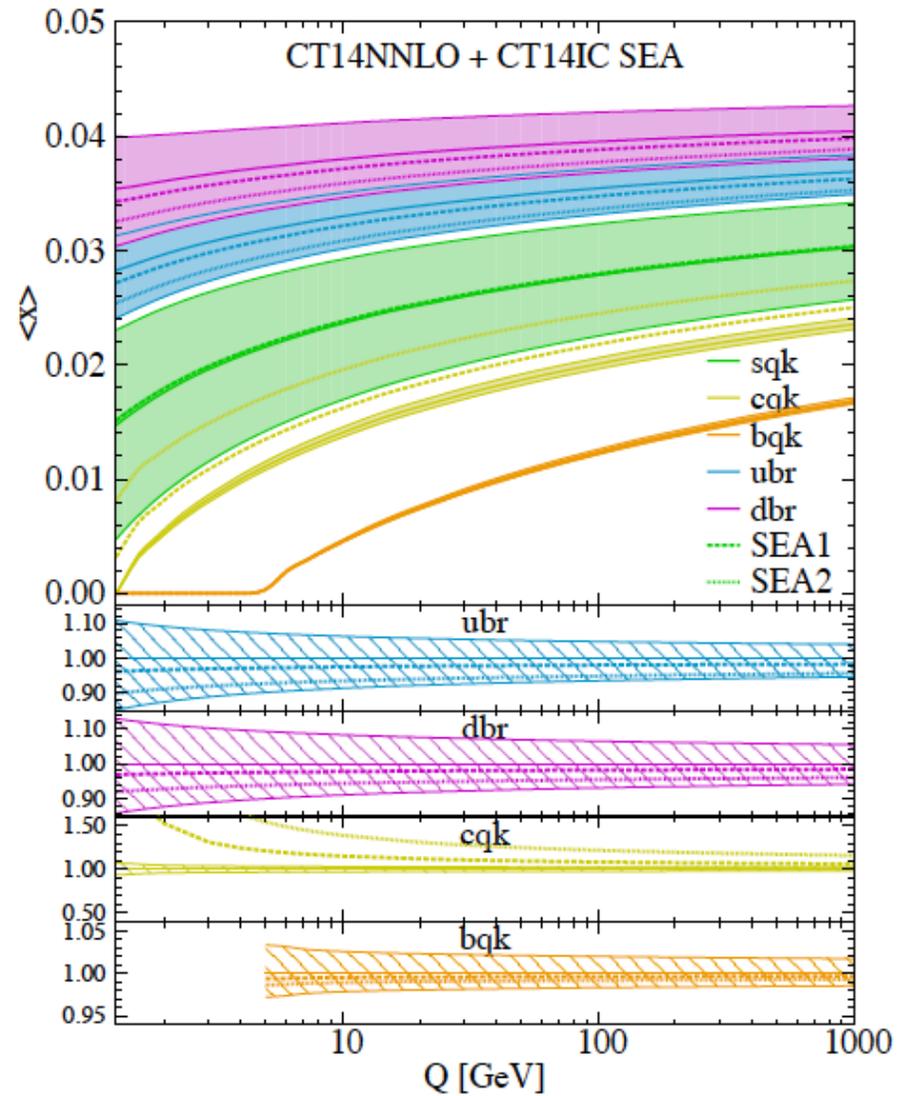
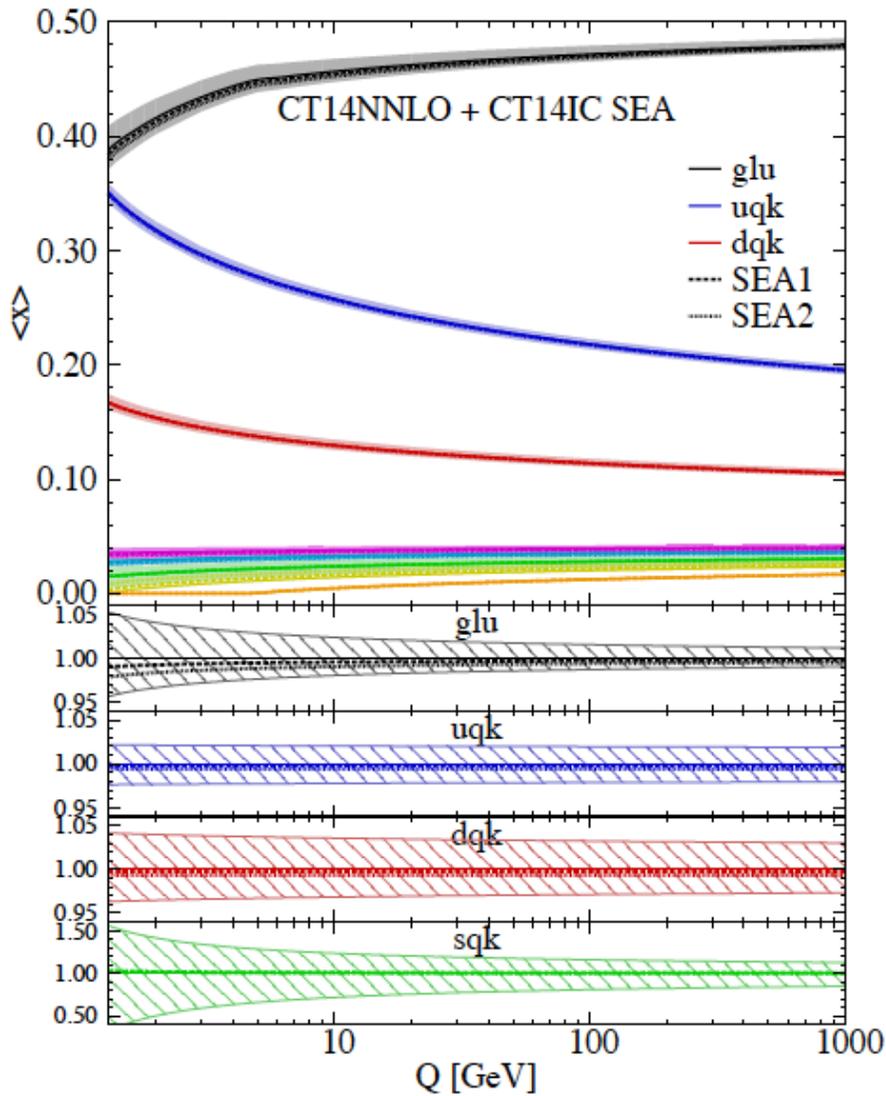
- BHPS solutions increase charm primarily at high  $x$ , as expected
- Sea-like solutions cause increase in charm at all  $x$  values



# Changes to PDFs within uncertainty bands

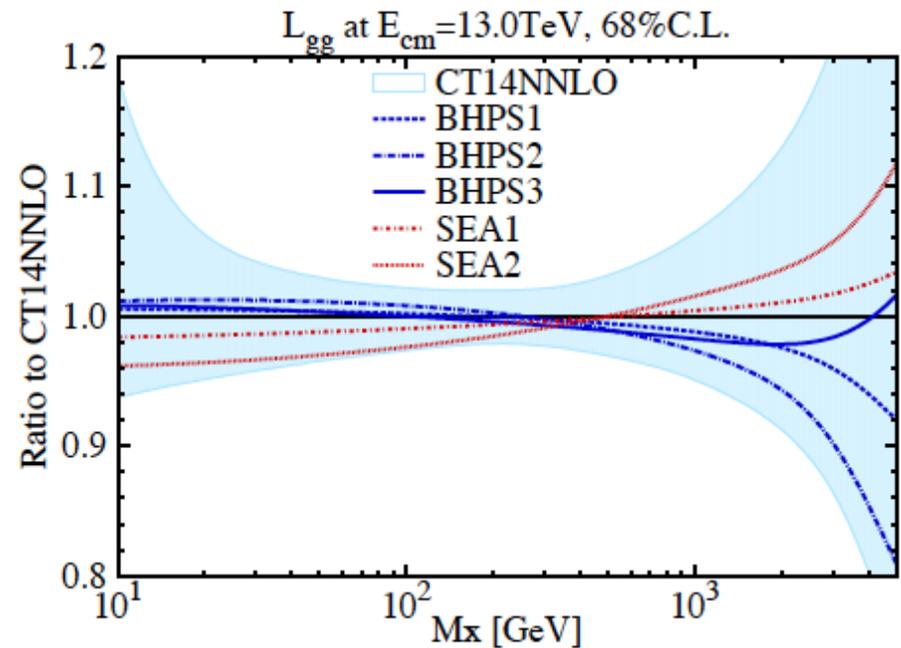
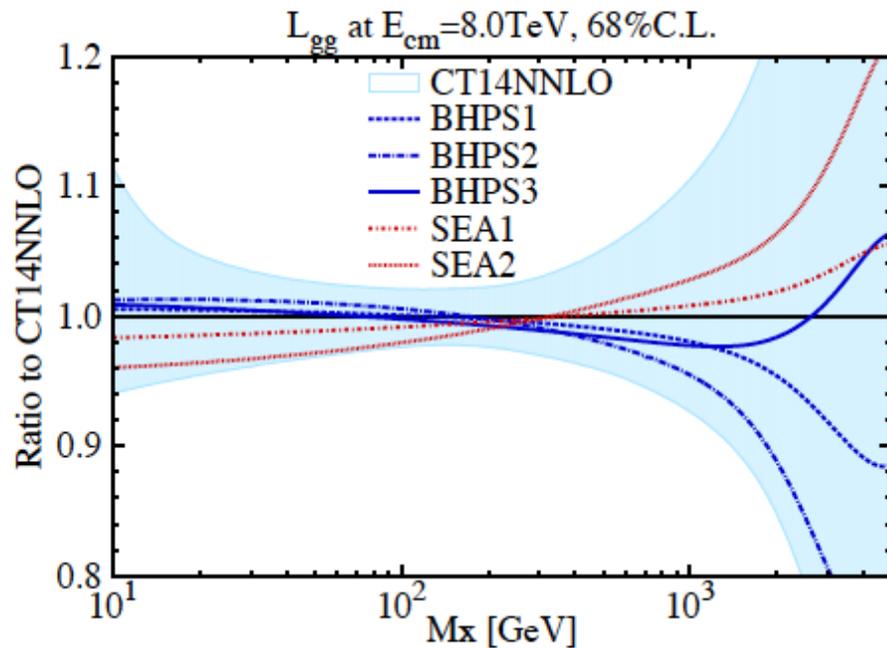


# Changes to PDFs within uncertainty bands



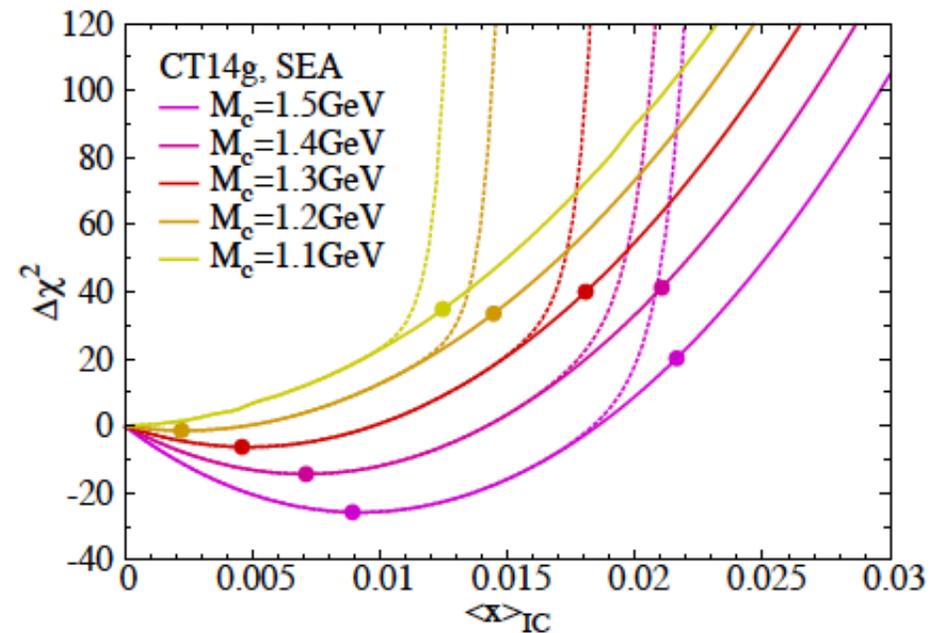
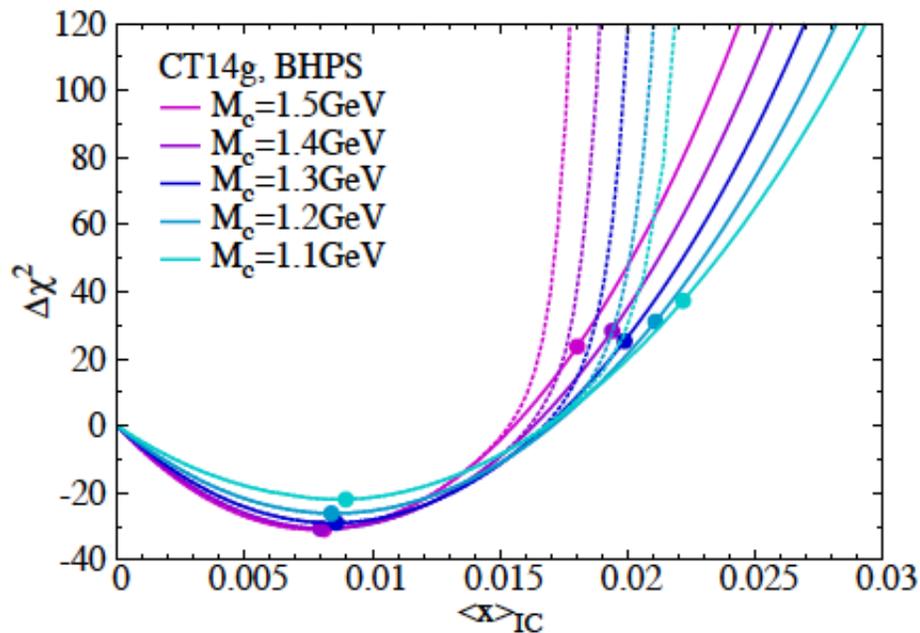
# gg PDF luminosity

- Impact of BHPS solutions on (gg) Higgs cross sections minimal
- Impact of sea-like solutions on Higgs cross sections more noticeable
- More impact at high mass, but within uncertainties



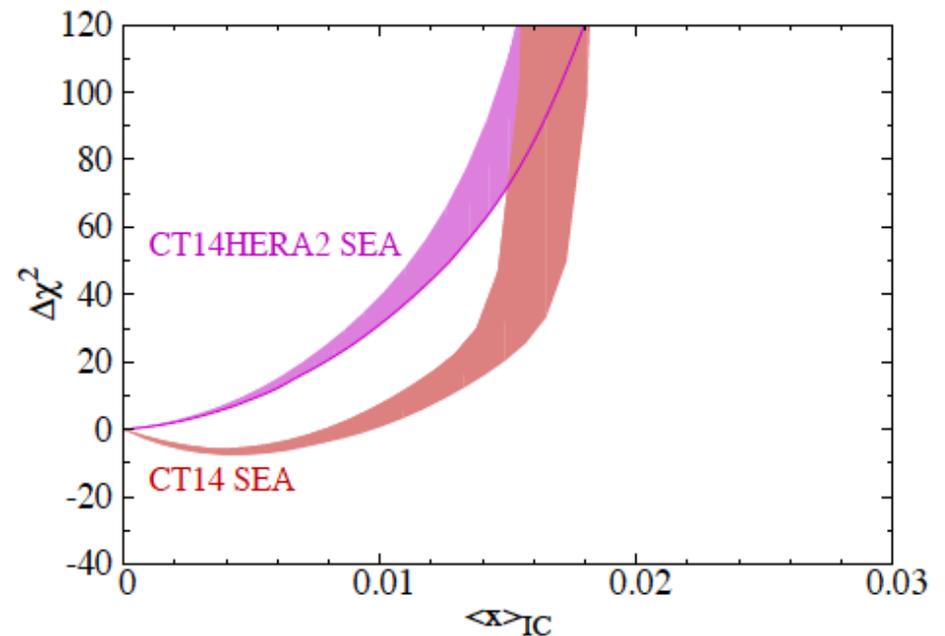
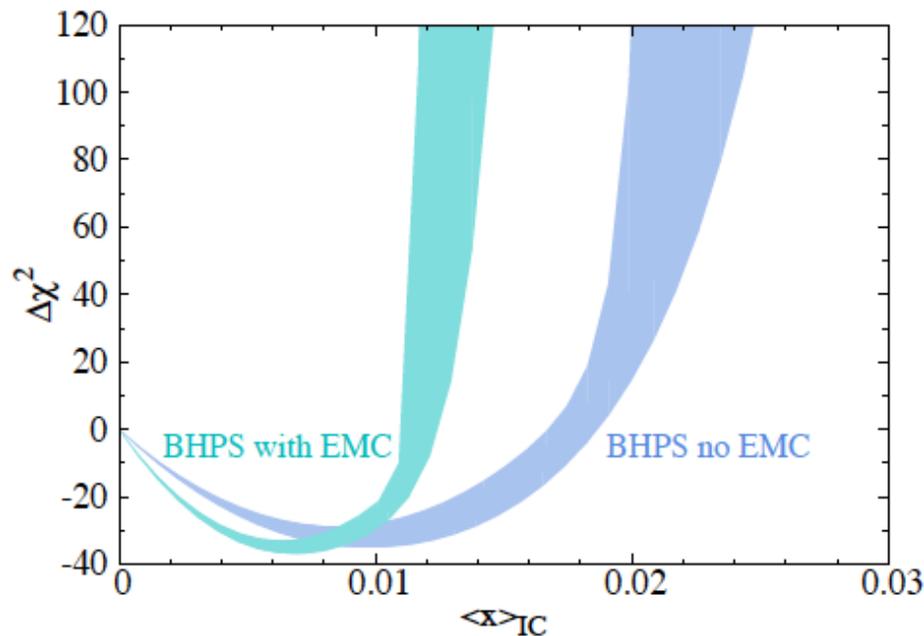
# Dependence on charm mass

- Variations of charm mass have little impact on amount of intrinsic charm in BHPS modes, stronger effects for SEA models
  - ◆ *Note: CT14g is a special CT14 fit with a more flexible gluon parameterization necessary for  $Q_0=1$  GeV, i.e. allowing the gluon to be negative*



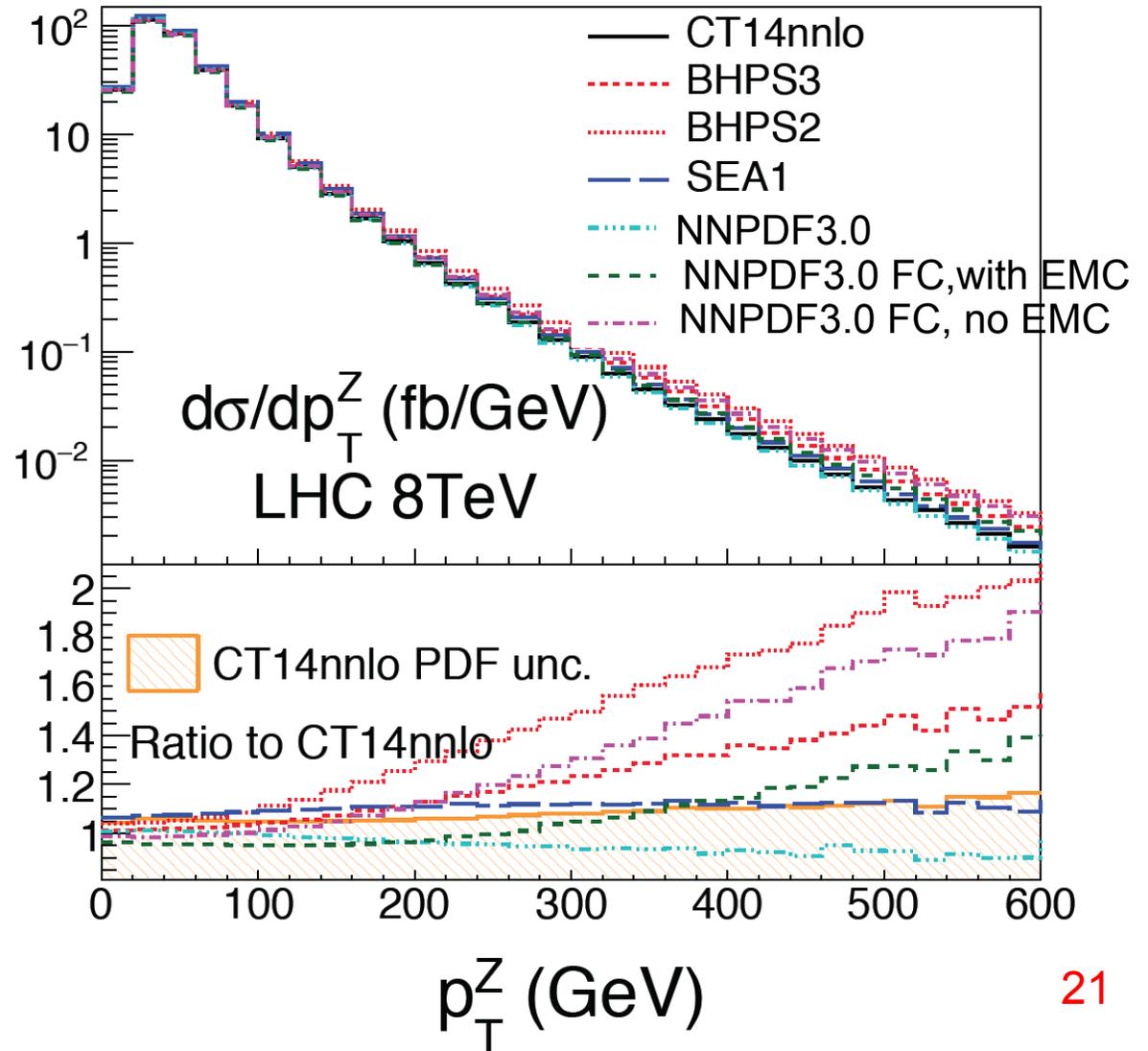
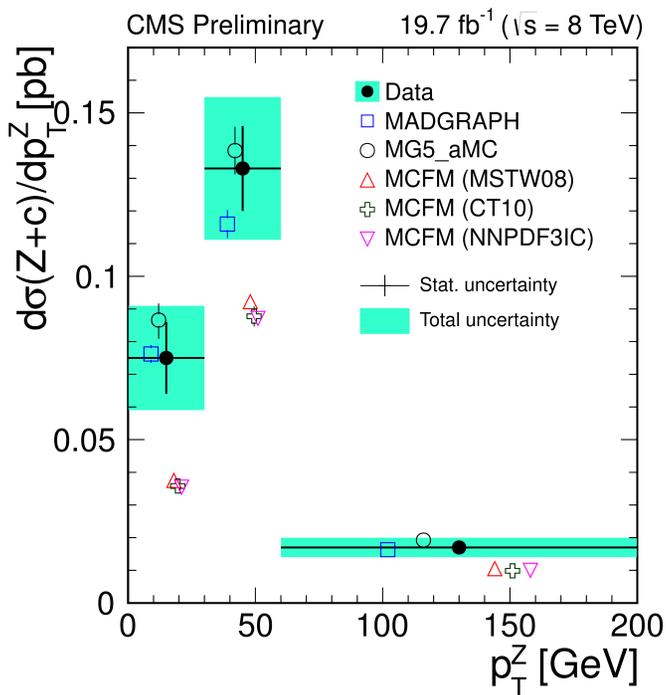
# Impact of EMC data

- Inclusion of EMC data in fit results in stronger limits on amount of intrinsic charm, move of minimum in  $\chi^2$  to slightly lower amounts of intrinsic charm
- EMC data is fit poorly, with or without IC,  $\chi^2$  between 2 and 3



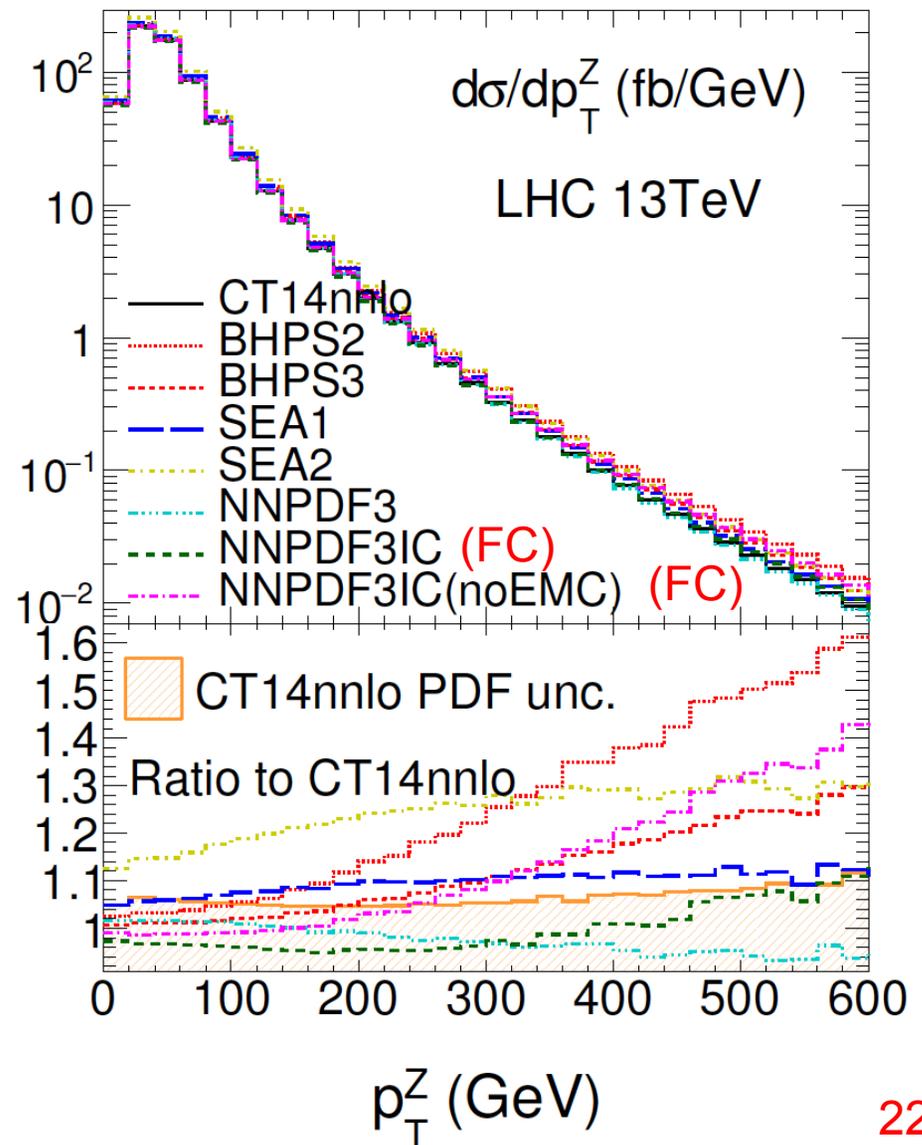
# Z+charm

- BHPS3, for example, would take a high integrated luminosity to confirm or rule out

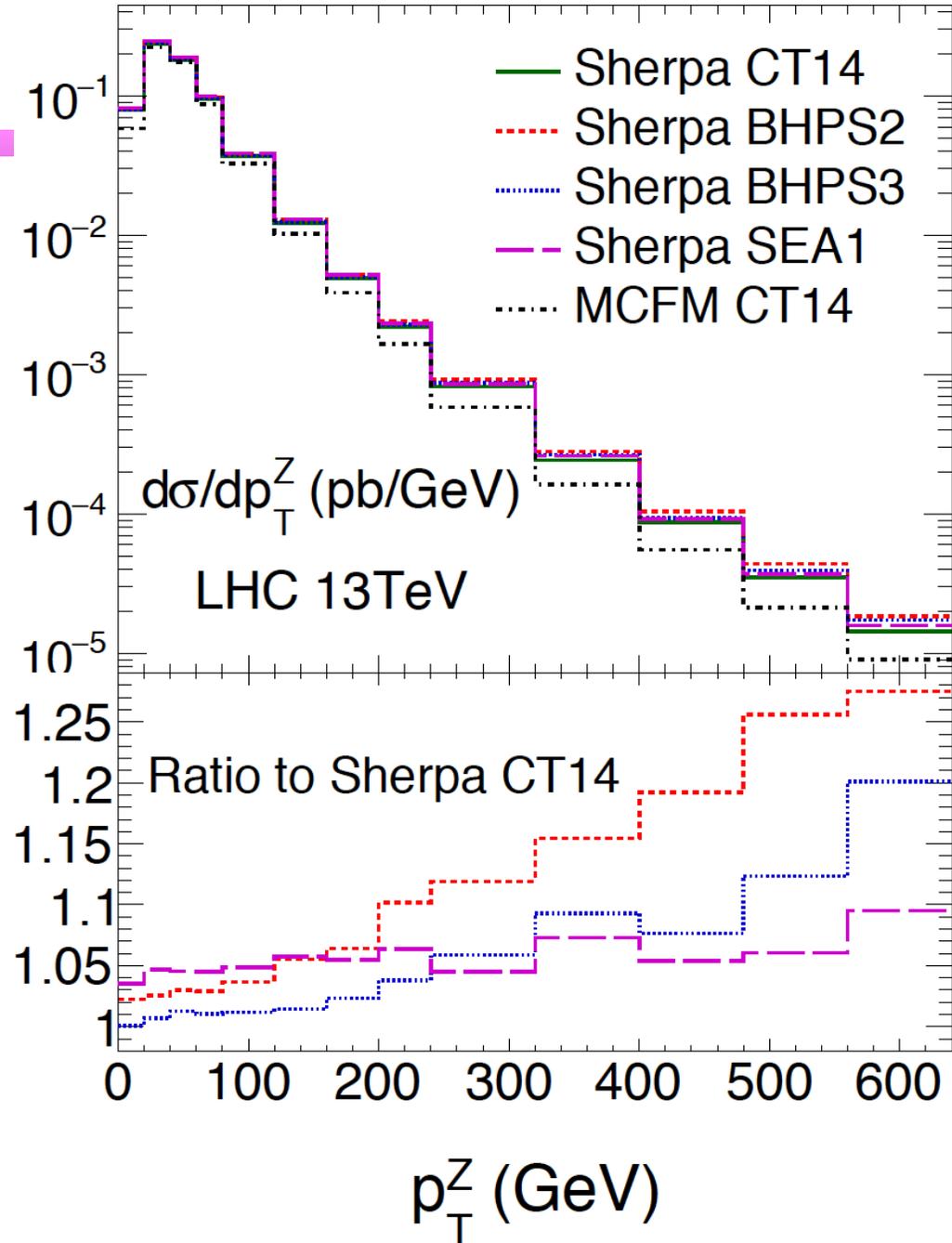


# Z+charm at 13 TeV

- Higher energy means lower  $x$  for fixed  $p_T$ , so in some sense confirming models becomes more difficult at 13 TeV, without a very large amount of integrated luminosity
- Except SEA2 again which results in a fairly sizeable DC shift to the cross section



- The effect of intrinsic charm is reduced due to the effects of radiative corrections, for example Sherpa with up to 3 jets by the matrix element in the final state



# Future studies/issues

- Add new processes known at NNLO
  - ◆ photon+jet
  - ◆ W+jet, Z+jet
- As more NNLO results are transferred by means of NNLO/NLO K-factors, can we agree on a standard PDF for the K-factors, i.e. PDF4LHC15?
- Scale uncertainties
  - ◆ perhaps most crucial for jet production
  - ◆ run fits with  $p_T^{\text{jet}}$  scale varied by a factor of 2
  - ◆ use  $p_T^{\text{lead-jet}}$  rather than  $p_T^{\text{jet}}$
  - ◆ add rapidity dependence to jet scales
- Conflicts between data sets
  - ◆ a better understanding of PDF uncertainties requires a better understanding of LHC data, and in particular conflicts between experiments, between data-sets, and even between observables
  - ◆ this trend will probably continue as systematic errors decrease and vanishingly small uncorrelated errors don't allow much wiggle room in PDF fits
- This is an area where PDF4LHC can contribute